Second Order Partial Differential Equations in Hilbert Spaces

Giuseppe Da Prato Scuola Normale Superiore di Pisa

Jerzy Zabczyk Polish Academy of Sciences, Warsaw



PUBLISHED BY THE PRESS SYNDICATE OF THE UNIVERSITY OF CAMBRIDGE The Pitt Building, Trumpington Street, Cambridge, United Kingdom

CAMBRIDGE UNIVERSITY PRESS The Edinburgh Building, Cambridge CB2 2RU, UK 40 West 20th Street, New York, NY 10011-4211, USA 477 Williamstown Road, Port Melbourne, VIC 3207, Australia Ruiz de Alarcón 13, 28014 Madrid, Spain Dock House, The Waterfront, Cape Town 8001, South Africa

http://www.cambridge.org

© Cambridge University Press 2002

This book is in copyright. Subject to statutory exception and to the provisions of relevant collective licensing agreements, no reproduction of any part may take place without the written permission of Cambridge University Press.

First published 2002

Printed in the United Kingdom at the University Press, Cambridge

Typeface Computer Modern 10/12pt System IAT_FX 2_{ε} [TB]

A catalogue record for this book is available from the British Library

Library of Congress Cataloguing in Publication data

Da Prato, Giuseppe.

Second order partial differential equations in Hilbert spaces / Giuseppe Da Prato & Jerzy Zabczyk.

p. cm. – (London Mathematical Society lecture note series; 293)
Includes bibliographical references and index.
ISBN 0 521 77729 1 (pbk.)
1. Differential equations, Partial. 2. Hilbert space. I. Zabczyk, Jerzy. II. Title.
III. Series.
QA374 .D27 2002
515'.353-dc21 2002022269

ISBN 0 521 77729 1 paperback

Contents

Pı	refac	e	x
Ι	TH FU	IEORY IN SPACES OF CONTINUOUS JNCTIONS	1
1	Gau	ussian measures	3
	1.1	Introduction and preliminaries	3
	1.2	Definition and first properties of Gaussian measures	$\overline{7}$
		1.2.1 Measures in metric spaces	7
		1.2.2 Gaussian measures	8
		1.2.3 Computation of some Gaussian integrals	11
		1.2.4 The reproducing kernel	12
	1.3	Absolute continuity of Gaussian measures	17
		1.3.1 Equivalence of product measures in \mathbb{R}^{∞}	18
		1.3.2 The Cameron-Martin formula	22
		1.3.3 The Feldman-Hajek theorem	24
	1.4	Brownian motion	27
2	Spa	ces of continuous functions	30
	2.1	Preliminary results	30
	2.2	Approximation of continuous functions	33
	2.3	Interpolation spaces	36
		2.3.1 Interpolation between $UC_b(H)$ and $UC_b^1(H)$	36
		2.3.2 Interpolatory estimates	39
		2.3.3 Additional interpolation results	42
3	$\mathbf{Th} \boldsymbol{\epsilon}$	e heat equation	44
	3.1	Preliminaries	44
	3.2	Strict solutions	48

	3.3	Regularity of generalized solutions
		3.3.2 O-derivatives of generalized solutions 57
	34	Comments on the Gross Laplacian 67
	3.5	The heat semigroup and its generator 69
	0.0	
4	Poi	sson's equation 76
	4.1	Existence and uniqueness results
	4.2	Regularity of solutions
	4.3	The equation $\Delta_Q u = g \dots \dots$
		4.3.1 The Liouville theorem
5	Elli	ptic equations with variable coefficients 90
	5.1	Small perturbations
	5.2	Large perturbations
6	Orn	stein-Uhlenbeck equations 99
	6.1	Existence and uniqueness of strict solutions 100
	6.2	Classical solutions
	6.3	The Ornstein-Uhlenbeck semigroup
		6.3.1 π -Convergence
		6.3.2 Properties of the π -semigroup (R_t)
		6.3.3 The infinitesimal generator
	6.4	Elliptic equations
		6.4.1 Schauder estimates
		6.4.2 The Liouville theorem $\ldots \ldots \ldots$
	6.5	Perturbation results for parabolic equations
	6.6	Perturbation results for elliptic equations
7	Car	and perchalia equations 197
1	7 1	Implicit function theorems 127
	7.1	Wiener processes and stochastic equations
	1.2	7.2.1 Infinite dimensional Wiener processes
		7.2.1 Infinite dimensional whener processes
	79	Dependence of the solutions to stochastic equations on initial
	1.5	deta 122
		uata \dots 130
		7.3.2 Convolution and evaluation maps
	74	Space and time regularity of the generalized solutions 120
	1.4 75	Fristoneo 140
	1.0	

vi

Contents

	7.6	Uniqueness	1
		7.6.1 Uniqueness for the heat equation	5
		7.6.2 Uniqueness in the general case	3
	7.7	Strong Feller property)
8	Para	abolic equations in open sets 150	3
	8.1	Introduction	3
	8.2	Regularity of the generalized solution	3
	8.3	Existence theorems	5
	8.4	Uniqueness of the solutions	3
II	Τł	IEORY IN SOBOLEV SPACES 18	5
•	τ9		_
9	L^2 a	Itâ Wienen Jacoma acitien	r S
	9.1	10-Wiener decomposition	5
		9.1.1 Real Hermite polynomials	5
		9.1.2 Chaos expansions	ן א
	0.2	9.1.5 The space $L^{-}(\Pi, \mu; \Pi)$) 1
	9.2	Sobolev spaces	ŧ
		9.2.1 The space $W = (\Pi, \mu)$	נ 7
		9.2.2 Some additional summability results $\dots \dots \dots$	2
		9.2.5 Compactness of the embedding $W^{-1}(\Pi,\mu) \subset L^{-1}(\Pi,\mu)$ 130 0.2.4 The space $W^{2,2}(\Pi,\mu)$ 200)
	9.3	The Malliavin derivative $\dots \dots \dots$	3
10	Orn	stein Uhlenbeck somigroups on $L^p(H,\mu)$ 20	5
10	10.1	Extension of (R_i) to $L^p(H,\mu)$ 200	, 3
	10.1	10.1.1 The adjoint of (R_t) in $L^2(H, \mu)$ 21	,
	10.2	The infinitesimal generator of (R_i) (R_i) (R_i) (R_i)	2
	10.2	10.2.1 Characterization of the domain of L_2	5
	10.3	The case when (R_t) is strong Feller	ź
		10.3.1 Additional regularity properties of (R_t)	I
		10.3.2 Hypercontractivity of (R_t)	1
	10.4	A representation formula for (R_t) in terms of the second quan-	
		tization operator	3
		10.4.1 The second quantization operator	3
		10.4.2 The adjoint of (R_t))
	10.5	Poincaré and log-Sobolev inequalities)
		10.5.1 The case when $M = 1$ and $Q = I \dots 23$	2

10.5.2 A generalization 235				
10.6 Some additional regularity results when Q and A commute $\therefore 236$				
11 Perturbations of Ornstein-Uhlenbeck semigroups23811 1 D11 1 D				
11.1 Bounded perturbations				
11.2 Lipschitz perturbations				
11.2.1 Some additional results on the Ornstein-Uhlenbeck				
semigroup $\ldots \ldots 251$				
11.2.2 The semigroup (P_t) in $L^p(H,\nu)$				
11.2.3 The integration by parts formula $\ldots \ldots \ldots \ldots 260$				
11.2.4 Existence of a density $\ldots \ldots \ldots \ldots \ldots \ldots \ldots 263$				
12 Gradient systems 267				
12.1 General results				
12.1.1 Assumptions and setting of the problem				
12.1.2 The Sobolev space $W^{1,2}(H,\nu)$				
12.1.3 Symmetry of the operator N_0				
12.1.4 The <i>m</i> -dissipativity of N_1 on $L^1(H, \nu)$				
12.2 The <i>m</i> -dissipativity of N_2 on $L^2(H\nu)$ 277				
12.3 The case when U is convex 281				
12.3.1 Poincaré and log-Sobolev inequalities				
III ADDIICATIONS TO CONTROL THEORY 201				
III AFFEICATIONS TO CONTROL THEORY 291				
13 Second order Hamilton-Jacobi equations 293				
13.1 Assumptions and setting of the problem				
13.2 Hamilton-Jacobi equations with a Lipschitz Hamiltonian 300				
13.2.1 Stationary Hamilton-Jacobi equations				
13.3 Hamilton-Jacobi equation with a quadratic Hamiltonian \ldots 305				
13.3.1 Stationary equation $\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots 308$				
13.4 Solution of the control problem $\ldots \ldots \ldots \ldots \ldots \ldots \ldots 310$				
13.4.1 Finite horizon				
13.4.2 Infinite horizon $\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots 312$				
13.4.3 The limit as $\varepsilon \to 0$				
14 Hamilton-Jacobi inclusions				
14.1 Introduction				

	14.4	Excessive measures for Wiener processes	328
IV	A	PPENDICES	333
A	Inte A.1 A.2	erpolation spacesThe interpolation theoremInterpolation between a Banach space X and the domain ofa linear operator in X	335 335 336
В	Nul B.1 B.2 B.3	l controllability Definition of null controllability	338 338 339 340
С	Sem C.1 C.2 C.3	niconcave functions and Hamilton-Jacobi semigroups Continuity modulus Semiconcave and semiconvex functions The Hamilton-Jacobi semigroups	347 347 348 351
Bi	bliog	graphy	358
In	dex		376

ix

Chapter 1

Gaussian measures

This chapter is devoted to some basic results on Gaussian measures on separable Hilbert spaces, including the Cameron-Martin and Feldman-Hajek formulae. The greater part of the results are presented with complete proofs.

1.1 Introduction and preliminaries

We are given a real separable Hilbert space H (with norm $|\cdot|$ and inner product $\langle \cdot, \cdot \rangle$). The space of all linear bounded operators from H into H, equipped with the operator norm $\|\cdot\|$, will be denoted by L(H). If $T \in L(H)$, then T^* is the adjoint of T. Moreover, by $L^+(H)$ we shall denote the subset of L(H) consisting of all nonnegative symmetric operators. Finally, we shall denote by $\mathcal{B}(H)$ the σ -algebra of all Borel subsets of H.

Before introducing Gaussian measures we need some results about trace class and Hilbert-Schmidt operators.

A linear bounded operator $R \in L(H)$ is said to be of *trace class* if there exist two sequences (a_k) , (b_k) in H such that

$$Ry = \sum_{k=1}^{\infty} \langle y, a_k \rangle b_k, \quad y \in H,$$
(1.1.1)

and

$$\sum_{k=1}^{\infty} |a_k| \, |b_k| < +\infty. \tag{1.1.2}$$

Notice that if (1.1.2) holds then the series in (1.1.1) is norm convergent. Moreover, it is not difficult to show that R is compact. We shall denote by $L_1(H)$ the set of all operators of L(H) of trace class. $L_1(H)$, endowed with the usual linear operations, is a Banach space with the norm

$$||R||_{L_1(H)} = \inf\left\{\sum_{k=1}^{\infty} |a_k| \, |b_k| : \ Ry = \sum_{k=1}^{\infty} \langle y, a_k \rangle b_k, \ y \in H, \ (a_k), (b_k) \subset H\right\}.$$

We set $L_1^+(H) = L^+(H) \cap L_1(H)$. If an operator R is of trace class then its trace, Tr R, is defined by the formula

$$\operatorname{Tr} R = \sum_{j=1}^{\infty} \langle Re_j, e_j \rangle,$$

where (e_j) is an orthonormal and complete basis on H. Notice that, if R is given by (1.1.1), we have

Tr
$$R = \sum_{j=1}^{\infty} \langle a_j, b_j \rangle.$$

Thus the definition of the trace is independent on the choice of the basis and

$$|\text{Tr } R| \leq ||R||_{L_1(H)}.$$

Proposition 1.1.1 Let $S \in L_1(H)$ and $T \in L(H)$. Then

(i) $ST, TS \in L_1(H)$ and

$$||TS||_{L_1(H)} \le ||S||_{L_1(H)} ||T||, ||ST||_{L_1(H)} \le ||S||_{L_1(H)} ||T||.$$

(*ii*) $\operatorname{Tr}(ST) = \operatorname{Tr}(TS)$.

Proof. (i) Assume that $Sy = \sum_{k=1}^{\infty} \langle y, a_k \rangle b_k, y \in H$, where $\sum_{k=1}^{\infty} |a_k| |b_k| < +\infty$. Then

$$STy = \sum_{k=1}^{\infty} \langle y, T^*a_k \rangle b_k, \ y \in H,$$

and

$$\sum_{k=1}^{\infty} |T^*a_k| |b_k| \le ||T|| \sum_{k=1}^{\infty} |a_k| |b_k|.$$

Gaussian measures

It is therefore clear that $ST \in L_1(H)$ and $||ST||_{L_1(H)} \leq ||S||_{L_1(H)} ||T||$. Similarly we can prove that $||TS||_{L_1(H)} \leq ||S||_{L_1(H)} ||T||$.

(ii) From part (i) it follows that

$$\operatorname{Tr}(ST) = \sum_{k=1}^{\infty} \langle b_k, T^* a_k \rangle = \sum_{k=1}^{\infty} \langle Tb_k, a_k \rangle.$$

In the same way Tr $(TS) = \sum_{k=1}^{\infty} \langle a_k, Tb_k \rangle$, and the conclusion follows. \Box

We say that $R \in L(H)$ is of Hilbert-Schmidt class if there exists an orthonormal and complete basis (e_k) in H such that

$$\sum_{k,j=1}^{\infty} |\langle Se_k, e_j \rangle|^2 < +\infty.$$
(1.1.3)

If (1.1.3) holds then we have

$$\sum_{k=1}^{\infty} |Se_k|^2 = \sum_{k,j=1}^{\infty} |\langle Se_k, e_j \rangle|^2 = \sum_{k,j=1}^{\infty} |\langle e_k, S^*e_j \rangle|^2 = \sum_{j=1}^{\infty} |S^*e_j|^2.$$
(1.1.4)

Now if (f_k) is another complete orthonormal basis in H, we have

$$\sum_{m=1}^{\infty} |Sf_m|^2 = \sum_{m,n=1}^{\infty} |\langle Sf_m, e_n \rangle|^2 = \sum_{m,n=1}^{\infty} |\langle f_m, S^*e_n \rangle|^2 = \sum_{n=1}^{\infty} |S^*e_n|^2.$$

Thus, by (1.1.4) we see that the assertion (1.1.3) is independent of the choice of the complete orthonormal basis (e_k) . We shall denote by $L_2(H)$ the space of all Hilbert-Schmidt operators on H. $L_2(H)$, endowed with the norm

$$||S||_{L_2(H)}^2 = \sum_{k,j=1}^{\infty} |\langle Se_k, e_j \rangle|^2 = \sum_{k=1}^{\infty} |Se_k|^2,$$

is a Banach space.

Proposition 1.1.2 Let $S, T \in L_2(H)$. Then $ST \in L_1(H)$ and

$$\|ST\|_{L_1(H)} \le \|S\|_{L_2(H)} \|T\|_{L_2(H)}.$$
(1.1.5)

Proof. Let (e_k) be a complete and orthonormal basis in H, then

$$Ty = \sum_{k=1}^{\infty} \langle Ty, e_k \rangle e_k = \sum_{k=1}^{\infty} \langle y, T^* e_k \rangle e_k,$$

$$STy = \sum_{k=1}^{\infty} \langle y, T^* e_k \rangle Se_k.$$

Consequently $ST \in L_1(H)$ and

$$\begin{split} \|ST\|_{L_1(H)} &\leq \sum_{k=1}^{\infty} |T^*e_k| \, |Se_k| \leq \left(\sum_{k=1}^{\infty} |T^*e_k|^2\right)^{1/2} \left(\sum_{k=1}^{\infty} |Se_k|^2\right)^{1/2} \\ &= \|T\|_{L_2(H)} \|S\|_{L_2(H)}. \end{split}$$

Therefore the conclusion follows. \Box

Warning. If S and T are bounded operators, and ST is of trace class then in general TS is not, as the following example, provided by S. Peszat [183], shows.

Define two linear operators S and T on the product space $H \times H$, by

$$S = \left(\begin{array}{cc} 0 & A \\ B & 0 \end{array}\right), \quad T = \left(\begin{array}{cc} I & 0 \\ 0 & 0 \end{array}\right).$$

Then

$$ST = \begin{pmatrix} 0 & 0 \\ B & 0 \end{pmatrix}, \quad TS = \begin{pmatrix} 0 & A \\ 0 & 0 \end{pmatrix},$$

and it is enough to take B of trace class and A not of trace class. \Box

We have also the following result, see e.g. A. Pietsch [187].

Proposition 1.1.3 Assume that S is a compact self-adjoint operator, and that (λ_k) are its eigenvalues (repeated according to their multiplicity).

$$(i) \ S \in L_1(H) \ if and only \ if \ \sum_{k=1}^{\infty} |\lambda_k| < +\infty. \ Moreover \ ||S||_{L_1(H)} = \sum_{k=1}^{\infty} |\lambda_k|,$$

and Tr $S = \sum_{k=1}^{\infty} \lambda_k.$
$$(ii) \ S \in L_2(H) \ if and only \ if \ \sum_{k=1}^{\infty} |\lambda_k|^2 < +\infty. \ Moreover$$
$$||S||_{L_2(H)} = \left(\sum_{k=1}^{\infty} |\lambda_k|^2\right)^{1/2}.$$

Gaussian measures

More generally let S be a compact operator on H. Denote by (λ_k) the sequence of all positive eigenvalues of the operator $(S^*S)^{1/2}$, repeated according to their multiplicity. Denote by $L_p(H)$, p > 0, the set of all operators S such that

$$\|S\|_{L_p(H)} = \left(\sum_{k=1}^{\infty} \lambda_k^p\right)^{1/p} < +\infty.$$
 (1.1.6)

Operators belonging to $L_1(H)$ and $L_2(H)$ are precisely the trace class and the Hilbert-Schmidt operators.

The following result holds, see N. Dunford and J. T. Schwartz [107].

Proposition 1.1.4 Let $S \in L_p(H)$, $T \in L_q(H)$ with p > 0, q > 0. Then $ST \in L_r(H)$ with $\frac{1}{r} = \frac{1}{p} + \frac{1}{q}$, and

$$||TS||_{L_r(H)} \le 2^{1/r} ||S||_{L_p(H)} ||T||_{L_q(H)}.$$
(1.1.7)

1.2 Definition and first properties of Gaussian measures

1.2.1 Measures in metric spaces

If E is a metric space, then $\mathcal{B}(E)$ will denote the Borel σ -algebra, that is the smallest σ -algebra of subsets of E which contains all closed (open) subsets of E.

Let metric spaces E_1, E_2 be equipped with σ -fields $\mathcal{E}_1, \mathcal{E}_2$ respectively. Measurable mappings $X : E_1 \to E_2$ will often be called *random variables*. If μ is a measure on (E_1, \mathcal{E}_1) , then its image by the transformation X will be denoted by $X \circ \mu$:

$$X \circ \mu(A) = \mu(X^{-1}(A)), \ A \in \mathcal{E}_2.$$

We call $X \circ \mu$ the *law* or the *distribution* of X, and we set $X \circ \mu = \mathcal{L}(X)$.

If ν and μ are two finite measures on (E, \mathcal{E}) such that $\Gamma \in \mathcal{E}$, $\mu(\Gamma) = 0$ implies $\nu(\Gamma) = 0$ then one writes $\nu \ll \mu$ and one says that ν is *absolutely* continuous with respect to μ . If there exist $A, B \in \mathcal{E}$ such that $A \cap B = \emptyset$, $\mu(A) = \nu(B) = 1$, one says that μ and ν are singular.

If $\nu \ll \mu$ then by the Radon-Nikodým theorem there exists $g \in L^1(E, \mathcal{E}, \mu)$ nonnegative such that

$$\nu(\Gamma) = \int_{\Gamma} g(x) \mu(dx), \ \ \Gamma \in \mathcal{E}.$$

The function g is denoted by $\frac{d\nu}{d\mu}$.

If $\nu \ll \mu$ and $\mu \ll \nu$ then one says that μ and ν are *equivalent* and writes $\mu \sim \nu$.

We have the following change of variable formula. If φ is a nonnegative measurable real function on E_2 , then

$$\int_{E_1} \varphi(X(x))\mu(dx) = \int_{E_2} \varphi(y) X \circ \mu(dy).$$
(1.2.1)

Let μ and ν be two measures on a separable Hilbert space H; if $T \circ \mu = T \circ \nu$ for any linear operator $T: H \to \mathbb{R}^n$, $n \in \mathbb{N}$, then $\mu = \nu$.

Random variables X_1, \ldots, X_n are said to be *independent* if

$$\mathcal{L}(X_1,\ldots,X_n) = \mathcal{L}(X_1) \times \cdots \times \mathcal{L}(X_n).$$

A family of random variables $(X_{\alpha})_{\alpha \in A}$ is said to be independent, if any finite subset of the family is independent.

Probability measures on a separable Hilbert space H will always be regarded as defined on $\mathcal{B}(H)$. If μ is a probability measure on H, then its Fourier transform is defined by

$$\hat{\mu}(\lambda) = \int_{H} e^{i \langle \lambda, x \rangle} \mu(dx), \ \lambda \in H;$$

 $\hat{\mu}$ is called the *characteristic function* of μ . One can show that if the characteristic functions of two measures are identical, then the measures are identical as well.

1.2.2 Gaussian measures

We first define Gaussian measures on \mathbb{R} . If $a \in \mathbb{R}$ we set

$$N_{a,0}(dx) = \delta_a(dx),$$

where δ_a is the Dirac measure at a. If moreover $\lambda > 0$ we set

$$N_{a,\lambda}(dx) = \frac{1}{\sqrt{2\pi\lambda}} e^{-\frac{(x-a)^2}{2\lambda}} dx.$$

The Fourier transform of $N_{a,\lambda}$ is given by

$$\widehat{N_{a,\lambda}}(h) = \int_{\mathbb{R}} e^{ihx} N_{a,\lambda}(dx) = e^{iah - \frac{1}{2}\lambda h^2}, \ h \in \mathbb{R}.$$

More generally we show now that in an arbitrary separable Hilbert space and for arbitrary $Q \in L_1^+(H)$ there exists a unique measure $N_{a,Q}$ such that

$$\widehat{N_{a,\lambda}}(h) = \int_{H} e^{i\langle h,x\rangle} N_{a,Q}(dx) = e^{i\langle h,x\rangle - \frac{1}{2}\langle Qh,h\rangle}, \ h \in H$$

Let in fact $Q \in L_1^+(H)$. Then there exist a complete orthonormal system (e_k) on H and a sequence of nonnegative numbers (λ_k) such that $Qe_k = \lambda_k e_k, \ k \in \mathbb{N}$. We set $x_h = \langle x, e_h \rangle, h \in \mathbb{N}$, and $P_n x = \sum_{k=1}^n x_k e_k, x \in H, n \in \mathbb{N}$. Let us introduce an isomorphism γ from H into ℓ^2 : (¹)

$$x \in H \to \gamma(x) = (x_k) \in \ell^2.$$

In the following we shall always identify H with ℓ^2 . In particular we shall write $P_n x = (x_1, ..., x_n), x \in \ell^2$.

A subset I of H of the form $I = \{x \in H : (x_1, \dots, x_n) \in B\}$, where $B \in \mathcal{B}(\mathbb{R}^n)$, is said to be *cylindrical*. It is easy to see that the σ -algebra generated by all cylindrical subsets of H coincides with $\mathcal{B}(H)$.

Theorem 1.2.1 Let $a \in H$, $Q \in L_1^+(H)$. Then there exists a unique probability measure μ on $(H, \mathcal{B}(H))$ such that

$$\int_{H} e^{i\langle h,x\rangle} \mu(dx) = e^{i\langle a,h\rangle} e^{-\frac{1}{2}\langle Qh,h\rangle}, \ h \in H.$$
(1.2.2)

Moreover μ is the restriction to H (identified with ℓ^2) of the product measure

$$\bigotimes_{k=1}^{\infty} \mu_k = \bigotimes_{k=1}^{\infty} N_{a_k,\lambda_k},$$

defined on $(\mathbb{R}^{\infty}, \mathcal{B}(\mathbb{R}^{\infty}))$. (2)

We set $\mu = N_{a,Q}$, and call *a* the *mean* and *Q* the *covariance operator* of μ . Moreover $N_{0,Q}$ will be denoted by N_Q .

Proof of Theorem 1.2.1. Since a characteristic function uniquely determines the measure, we have only to prove existence.

Let us consider the sequence of Gaussian measures (μ_k) on \mathbb{R} defined as $\mu_k = N_{a_k,\lambda_k}, \ k \in \mathbb{N}$, and the product measure $\mu = \bigotimes_{k=1}^{\infty} \mu_k$ in \mathbb{R}^{∞} , see e.g

¹For any $p \ge 1$, we denote by ℓ^p the Banach space of all sequences (x_k) of real numbers such that $|x|_p := (\sum_{k=1}^{\infty} |x_k|^p)^{1/p} < +\infty$. ²We shall consider \mathbb{R}^{∞} as a metric space with the distance d(x, y) :=

²We shall consider \mathbb{R}^{∞} as a metric space with the distance $d(x,y) := \sum_{k=1}^{\infty} 2^{-k} \frac{|x_k - y_k|}{1 + |x_k - y_k|}, x, y \in \mathbb{R}^{\infty}$

P. R. Halmos [141, §38.B]. We want to prove that μ is concentrated on ℓ^2 , (that it is clearly a Borel subset of \mathbb{R}^{∞}). For this it is enough to show that

$$\int_{\ell^{\infty}} |x|_{\ell^{2}}^{2} \mu(dx) < +\infty.$$
 (1.2.3)

We have in fact, by the monotone convergence theorem,

$$\int_{\mathbb{R}^{\infty}} |x|_{\ell^{2}}^{2} \mu(dx) = \sum_{k=1}^{\infty} \int_{\mathbb{R}^{\infty}} x_{k}^{2} \, \mu(dx) = \sum_{k=1}^{\infty} \left(\int_{\mathbb{R}} (x_{k} - a_{k})^{2} \mu_{k}(dx) + a_{k}^{2} \right)$$
$$= \sum_{k=1}^{\infty} (\lambda_{k} + a_{k}^{2}) = \operatorname{Tr} Q + |a|^{2} < +\infty.$$

Now we consider the restriction of μ to ℓ^2 , which we still denote by μ . We have to prove that (1.2.2) holds. Setting $\nu_n = \prod_{k=1}^n \mu_k$, we have

$$\int_{\ell^2} e^{i\langle x,h\rangle} \mu(dx) = \lim_{n \to \infty} \int_{\ell^2} e^{i\langle P_n h, P_n x \rangle} \mu(dx)$$
$$= \lim_{n \to \infty} \int_{\mathbb{R}^n} e^{i\langle P_n h, P_n x \rangle} \nu_n(dx) = \lim_{n \to \infty} e^{i\langle P_n h, P_n a \rangle - \frac{1}{2}\langle QP_n h, P_n h \rangle}$$
$$= e^{i\langle h, a \rangle - \frac{1}{2}\langle Qh, h \rangle}. \square$$

If the law of a random variable is a Gaussian measure, then the random variable is called *Gaussian*. It easily follows from Theorem 1.2.1 that a random variable X with values in H is Gaussian if and only if for any $h \in H$ the real valued random variable $\langle h, X \rangle$ is Gaussian.

Remark 1.2.2 From the proof of Theorem 1.2.1 it follows that

$$\int_{H} |x|^2 N_{a,Q}(dx) = \text{Tr } Q + |a|^2.$$
(1.2.4)

Proposition 1.2.3 Let $T \in L(H)$, and $a \in H$, and let $\Gamma x = Tx + a$, $x \in H$. Then $\Gamma \circ N_{m,Q} = N_{Tm+a,TQT^*}$.

Proof. Notice that, by the change of variables formula (1.2.1), we have

$$\int_{H} e^{i\langle\lambda,y\rangle} \Gamma \circ N_{m,Q}(dy) = \int_{H} e^{i\langle\lambda,\Gamma x\rangle} N_{m,Q}(dy)$$
$$= \int_{H} e^{i\langle\lambda,Tx+a\rangle} N_{m,Q}(dy) = e^{i\langle\lambda,a\rangle} e^{i\langle T^*\lambda,m\rangle - \frac{1}{2}\langle QT^*\lambda,T^*\lambda\rangle}.$$

This shows the result. \Box

1.2.3 Computation of some Gaussian integrals

We are here given a Gaussian measure $N_{a,Q}$. We set

$$L^2(H, N_{a,Q}) = L^2(H, \mathcal{B}(H), N_{a,Q}).$$

The following identities can be easily proved, using (1.2.2).

Proposition 1.2.4 We have

$$\int_{H} x N_{a,Q}(dx) = a, \qquad (1.2.5)$$

$$\int_{H} \langle x - a, y \rangle \langle x - a, z \rangle N_{a,Q}(dx) = \langle Qy, z \rangle.$$
 (1.2.6)

$$\int_{H} |x-a|^2 N_{a,Q}(dx) = \text{Tr } Q.$$
 (1.2.7)

Proof. We prove as instance (1.2.6). We have

$$\int_{H} x N_{a,Q}(dx) = \lim_{n \to \infty} \int_{H} P_n x N_{a,Q}(dx).$$

But

$$\int_{H} P_n x N_{a,Q}(dx) = (2\pi)^{-n/2} \prod_{k=1}^n \int_{\mathbb{R}} x_k \lambda_k^{-1/2} e^{-\frac{(x_k - a_k)^2}{2\lambda_k}} dx_k = a_k,$$

and the conclusion follows. \Box

Proposition 1.2.5 For any $h \in H$, the exponential function E_h , defined as

$$E_h(x) = e^{\langle h, x \rangle}, \quad x \in H,$$

belongs to $L^p(H, N_{a,Q}), p \ge 1$, and

$$\int_{H} e^{\langle h, x \rangle} N_{a,Q}(dx) = e^{\langle a, h \rangle} e^{\frac{1}{2} \langle Qh, h \rangle}.$$
(1.2.8)

Moreover the subspace of $L^2(H, N_{a,Q})$ spanned by all E_h , $h \in H$, is dense on $L^2(H, N_{a,Q})$.

Proof. We have

$$\int_{H} e^{\langle P_n h, P_n x \rangle} N_{a,Q}(dx) = e^{\langle P_n a, P_n h \rangle} e^{\frac{1}{2} \langle Q P_n h, P_n h \rangle}.$$

Letting n tend to 0 this gives (1.2.8).

Let us prove the last statement. Let $\varphi \in L^2(H, N_{a,Q})$ be such that

$$\int_{H} e^{\langle h, x \rangle} \varphi(x) N_{a,Q}(dx) = 0, \ h \in H.$$

Denote by φ^+ and φ^- the positive and negative parts of φ . Then

$$\int_{H} e^{\langle h, x \rangle} \varphi^{+}(x) N_{a,Q}(dx) = \int_{H} e^{\langle h, x \rangle} \varphi^{-}(x) N_{a,Q}(dx), \quad h \in H.$$

Let us define two measures

$$\mu(dx) = \varphi^+(x) N_{a,Q}(dx), \quad \nu(dx) = \varphi^-(x) N_{a,Q}(dx).$$

Then μ and ν are finite measures such that

$$\int_{H} e^{\langle h, x \rangle} \mu(dx) = \int_{H} e^{\langle h, x \rangle} \nu(dx), \ h \in H.$$

Let T be any linear transformation from H into \mathbb{R}^n , $n \in \mathbb{N}$. Then for any $\lambda \in \mathbb{R}^n$

$$\int_{\mathbb{R}^n} e^{\langle \lambda, z \rangle} T \circ \mu(dz) = \int_H e^{\langle \lambda, Tx \rangle} \mu(dx) = \int_H e^{\langle T^*\lambda, \rangle \rangle} \mu(dx)$$
$$= \int_H e^{\langle T^*\lambda, x \rangle} \nu(dx) = \int_{\mathbb{R}^n} e^{\langle \lambda, z \rangle} T \circ \nu(dz)$$

By a well known finite dimensional result $T \circ \mu = T \circ \nu$. Consequently measures μ and ν are identical and so $\varphi = 0$. \Box

1.2.4 The reproducing kernel

Here we are given an operator $Q \in L_1^+(H)$. We denote as before by (e_k) a complete orthonormal system in H and by (λ_k) a sequence of positive numbers such that $Qe_k = \lambda_k e_k, k \in \mathbb{N}$.

The subspace $Q^{1/2}(H)$ is called the *reproducing kernel* of the measure N_Q . If Ker $Q = \{0\}$, $Q^{1/2}(H)$ is dense on H. In fact, if $x_0 \in H$ is such that $\langle Q^{1/2}h, x_0 \rangle = 0$ for all $h \in H$, we have $Q^{1/2}x_0 = 0$ and so $Qx_0 = 0$, which yields $x_0 = 0$.

Let Ker $Q = \{0\}$. We are now going to introduce an isomorphism Wfrom H into $L^2(H, N_Q)$ that will play an important rôle in the following. The isomorphism W is defined by

$$f \in Q^{1/2}(H) \to W_f \in L^2(H, N_Q), \ W_f(x) = \langle Q^{-1/2}f, x \rangle, \ x \in H.$$

By (1.2.7) it follows that

$$\int_{H} W_f(x) W_g(x) N_Q(dx) = \langle f, g \rangle, \ f, g \in H$$

Thus W is an isometry and it can be uniquely extended to all of H. It will be denoted by the same symbol. For any $f \in H$, W_f is a real Gaussian random variable $N_{|f|^2}$.

More generally, for arbitrary elements $f_1, ..., f_n, (W_{f_1}, ..., W_{f_n})$ is a Gaussian vector with mean 0 and covariance matrix $(\langle f_i, f_j \rangle)$. If Ker $Q \neq \{0\}$ then the transformation $f \to W_f$ can be defined in exactly the same way but only for $f \in H_0 = \overline{Q^{1/2}(H)}$. We will write in some cases $\langle Q^{-1/2}y, f \rangle$ instead of $W_f(y)$.

The proof of the following proposition is left as an exercise to the reader.

Proposition 1.2.6 For any orthonormal sequence (f_n) in H, the family

1, W_{f_n} , $W_{f_k}W_{f_l}$, $2^{-1/2} \left(W_{f_m}^2 - 1\right)$, $m, n, k, l \in \mathbb{N}, k \neq l$,

is orthonormal in $L^2(H, N_Q)$.

Next we consider the function $f \to e^{W_f}$.

Proposition 1.2.7 The transformation $f \to e^{W_f}$ acts continuously from H into $L^2(H, N_Q)$, and

$$\int_{H} e^{W_{f}(x)} N_{Q}(dx) = e^{\frac{1}{2}|f|^{2}},$$

$$\int_{H} e^{i \ \lambda W_{f}(x)} N_{Q}(dx) = e^{-\frac{1}{2}\lambda^{2}|f|^{2}}, \ \lambda \in \mathbb{R}.$$
(1.2.9)

Proof. Since W_f is Gaussian with law $N_{0,|f|^2}$, (1.2.9) follows. Moreover, taking into account (1.2.8) it follows that

$$\int_{H} \left[e^{W_f} - e^{W_g} \right]^2 dN_Q = \int_{H} \left[e^{2W_f} - 2e^{W_{f+g}} + e^{2W_g} \right] dN_Q$$
$$= e^{2|f|^2} - 2e^{\frac{1}{2}|f+g|^2} + e^{2|g|^2} = \left[e^{|f|^2} - e^{|g|^2} \right]^2 + 2e^{|f|^2 + |g|^2} \left[1 - e^{-\frac{1}{2}|f-g|^2} \right],$$

which shows that W_f is locally uniformly continuous on H. \Box

Let us define the determinant of 1 + S where S is a compact self-adjoint operator in $L_1(H)$:

det
$$(1+S) = \prod_{k=1}^{\infty} (1+s_k),$$

where (s_k) is the sequence of eigenvalues of S (repeated according to their multiplicity).

Proposition 1.2.8 Assume that M is a symmetric operator such that $Q^{1/2}MQ^{1/2} < 1$, (³) and let $b \in H$. Then

$$\int_{H} \exp\left\{\frac{1}{2}\langle My, y \rangle + \langle b, y \rangle\right\} N_Q(dy)$$

= $\left[\det(1 - Q^{1/2}MQ^{1/2})\right]^{-1/2} \exp\left\{\frac{1}{2}|(1 - Q^{1/2}MQ^{1/2})^{-1/2}Q^{1/2}b|^2\right\}.$
(1.2.10)

Proof. Let (g_n) be an orthonormal basis for the operator $Q^{1/2}MQ^{1/2}$, and let (γ_n) be the sequence of the corresponding eigenvalues.

Claim 1. We have

$$\langle b, x \rangle = \sum_{k=1}^{\infty} \langle Q^{1/2}b, g_n \rangle W_{g_n}(x), N_Q$$
-a.e.

Claim 2. We have

$$\langle Mx, x \rangle = \sum_{n=1}^{\infty} \gamma_n |W_{g_n}(x)|^2, N_Q$$
-a.e,

the series being convergent in $L^1(H, N_Q)$.

We shall only prove the more difficult second claim.
Let
$$P_N = \sum_{k=1}^{N} e_k \otimes e_k$$
. (⁴) Then for any $x \in H$ we have
 $\langle MP_N x, P_N x \rangle = \langle (Q^{1/2} M Q^{1/2}) Q^{-1/2} P_N x, Q^{-1/2} P_N x \rangle$
 $= \sum_{n=1}^{\infty} \langle (Q^{1/2} M Q^{1/2}) Q^{-1/2} P_N x, g_n \rangle \langle Q^{-1/2} P_N x, g_n \rangle$
 $= \sum_{n=1}^{\infty} \gamma_n |\langle Q^{-1/2} P_N x, g_n \rangle|^2.$

Consequently, for each fixed x

$$\langle MP_N x, P_N x \rangle = \sum_{n=1}^{\infty} \gamma_n |W_{P_N g_n}|^2, \ N \in \mathbb{N}.$$

³This means that $\langle Q^{1/2}MQ^{1/2}x, x \rangle < |x|^2$ for any $x \in H$ different from 0.

⁴We rember that (e_k) is the sequence of eigenvectors of Q.

Gaussian measures

Moreover for each $L \in \mathbb{N}$

$$\int_{H} \left| \langle MP_N x, P_N x \rangle - \sum_{n=1}^{L} \gamma_n |W_{P_N g_n}|^2 \right| N_Q(dx)$$
$$\leq \sum_{n=L+1}^{\infty} |\gamma_n| \int_{H} |W_{P_N g_n}|^2 N_Q(dx)$$
$$= \sum_{n=L+1}^{\infty} |\gamma_n| |P_N g_n|^2 \leq \sum_{n=L+1}^{\infty} |\gamma_n|.$$

As $N \to \infty$ then $P_N x \to x$ and $W_{P_N g_n} \to W_{g_n}$ in $L^2(H, N_Q)$. Passing to subsequences if needed, and using the Fatou lemma, we see that

$$\int_{H} \left| \langle Mx, x \rangle - \sum_{n=1}^{L} \gamma_n |W_{g_n}|^2 \right| N_Q(dx) \le \sum_{n=L+1}^{\infty} |\gamma_n|.$$

Therefore the claim is proved.

By the claims it follows that

$$\exp\left\{\frac{1}{2}\langle Mx,x\rangle + \langle b,x\rangle\right\}$$
$$= \lim_{L \to \infty} \exp\left\{\sum_{n=1}^{L} \frac{1}{2}\gamma_n |W_{g_n}(x)|^2 + \langle Q^{1/2}b,g_n\rangle Wg_n(x)\right\},$$

with a.e. convergence with respect to N_Q for a suitable subsequence. Using the fact that (Wg_n) are independent Gaussian random variables, we obtain, by a direct calculation, for $p \ge 1$,

$$\int_{H} \exp\left\{p\sum_{n=1}^{L} \frac{1}{2}\gamma_{n}|W_{g_{n}}(x)|^{2} + p\langle Q^{1/2}b, g_{n}\rangle Wg_{n}(x)\right\} N_{Q}(dx)$$
$$=\left[\prod_{n=1}^{L} (1-p\gamma_{n})\right]^{-1/2} \exp\left\{\frac{1}{2}\sum_{n=1}^{\infty} \frac{|\langle Q^{1/2}b, g_{n}\rangle|^{2}}{1-p\gamma_{n}}\right\}.$$

Since $\gamma_n < 1$, and $\sum_{n=1}^{\infty} |\gamma_n| < \infty$, there exists p > 1 such that $p\gamma_n < 1$, for all $n \in \mathbb{N}$. Therefore

$$\lim_{L \to \infty} \prod_{n=1}^{L} (1 - p\gamma_n)^{-1/2} \exp\left\{\frac{1}{2} \frac{|\langle Q^{1/2}b, g_n \rangle|^2}{1 - p\gamma_n}\right\}$$
$$= \left[\prod_{n=1}^{\infty} (1 - p\gamma_n)\right]^{-1/2} \exp\left\{\frac{1}{2} \sum_{n=1}^{\infty} \frac{|\langle Q^{1/2}b, g_n \rangle|^2}{1 - p\gamma_n}\right\}.$$

So the sequence $\left(\exp\left\{\sum_{n=1}^{L}\left[\frac{1}{2}\gamma_{n}|W_{g_{n}}(x)|^{2}+\langle Q^{1/2}b,g_{n}\rangle W_{g_{n}}(x)\right]\right\}\right)$ is uniformly integrable. Consequently, passing to the limit, we find

$$\begin{split} &\int_{H} \exp\left\{1/2 \langle My, y \rangle + \langle b, y \rangle\right\} N_Q(dy) \\ &= \lim_{L \to \infty} \int_{H} \exp\left\{\sum_{n=1}^{L} \left[1/2 \gamma_n |W_{g_n}(x)|^2 + \langle Q^{1/2}b, g_n \rangle W_{g_n}(x)\right]\right\} N_Q(dx) \\ &= \lim_{L \to \infty} \prod_{n=1}^{L} (1 - \gamma_n)^{-1/2} \exp\left\{\frac{1}{2} \frac{|\langle Q^{1/2}b, g_n \rangle|^2}{1 - \gamma_n}\right\} \\ &= \prod_{n=1}^{\infty} (1 - \gamma_n)^{-1/2} \exp\left\{\frac{1}{2} \frac{|\langle Q^{1/2}b, g_n \rangle|^2}{1 - \gamma_n}\right\} \\ &= \left(\det(1 - Q^{1/2}MQ^{1/2})\right)^{-1/2} \exp\left\{\frac{1}{2} |(1 - Q^{1/2}MQ^{1/2})^{-1/2}Q^{1/2}b|^2\right\}. \Box$$

Remark 1.2.9 It follows from the proof of the proposition that

$$\langle Mx, x \rangle = \sum_{k=1}^{\infty} \gamma_n W_{g_n}^2(x) = \sqrt{2} \sum_{k=1}^{\infty} \gamma_n \left[2^{-1/2} (W_{g_n}^2(x) - 1) \right] + \sum_{k=1}^{\infty} \gamma_n,$$

and so, by Proposition 1.2.6, we have

$$\begin{split} \int_{H} [\langle Mx, x \rangle]^2 N_Q(dx) &= 2 \sum_{k=1}^{\infty} \gamma_n^2 + \left(\sum_{k=1}^{\infty} \gamma_n \right)^2 \\ &= 2 \|Q^{1/2} M Q^{1/2}\|_{L_2(H)}^2 + (\text{Tr } Q^{1/2} M Q^{1/2})^2 \\ &< +\infty. \end{split}$$

Proposition 1.2.10 Let $T \in L_1(H)$. Then there exists the limit

$$\langle TQ^{-1/2}y, Q^{-1/2}y \rangle := \lim_{n \to \infty} \langle TQ^{-1/2}P_n y, Q^{-1/2}P_n y \rangle, N_Q\text{-}a.e.,$$

where $P_n = \sum_{k=1}^n e_k \otimes e_k$.

Moreover we have the following expansion in $L^2(H, N_Q)$:

$$\langle TQ^{-1/2}y, Q^{-1/2}y \rangle = \sum_{n=1}^{\infty} \langle Tg_n, g_n \rangle + \sum_{m \neq n=1}^{\infty} \langle Tg_n, g_m \rangle W_{g_n} W_{g_m}$$

$$\times \sqrt{2} \sum_{n=1}^{\infty} \langle Tg_n, g_n \rangle \left[2^{-1/2} \left(W_{g_n}^2 - 1 \right) \right]. \quad (1.2.11)$$

The proof of the following result is similar to that of Claim 2 in the proof of Proposition 1.2.8 and it is left to the reader.

Proposition 1.2.11 Assume that M is a symmetric trace-class operator such that M < 1, (5) and $b \in H$. Then

$$\int_{H} \exp\left\{ 1/2 \left\langle MQ^{-1/2}y, Q^{-1/2}y \right\rangle + \left\langle b, Q^{-1/2}y \right\rangle \right\} N_{Q}(dy)$$
$$= (\det(1-M))^{-1/2} e^{\frac{1}{2}|(1-M)^{-1/2}b|^{2}}. \quad (1.2.12)$$

1.3 Absolute continuity of Gaussian measures

We consider here two Gaussian measures μ, ν . We want to prove the Feldman-Hajek theorem , that is they are either singular or equivalent.

⁵That is $\langle Mx, x \rangle < |x|^2$ for all $x \neq 0$.

In §1.3.1 we recall some results on equivalence of measures on \mathbb{R}^{∞} including the Kakutani theorem. In §1.3.2 we consider the case when $\mu = N_Q$ and $\nu = N_{a,Q}$ with $Q \in L_1^+(H)$ and $a \in H$, proving the Cameron-Martin formula. Finally in §1.3.3 we consider the more difficult case when $\mu = N_Q$ and $\nu = N_R$ with $Q, R \in L_1^+(H)$.

1.3.1 Equivalence of product measures in \mathbb{R}^{∞}

It is convenient to introduce the notion of *Hellinger* integral.

Let μ, ν be probability measures on a measurable space (E, \mathcal{E}) . Then $\lambda = \frac{1}{2}(\mu + \nu)$ is also a probability measure on (E, \mathcal{E}) and we have obviously

$$\mu << \lambda, \quad \nu << \lambda$$

We define the *Hellinger integral* by

$$H(\mu,\nu) = \int_E \left[\frac{d\mu}{d\lambda}(x)\frac{d\nu}{d\lambda}(x)\right]^{1/2}\lambda(dx).$$

Instead of $\frac{1}{2}(\mu + \nu)$ one could choose as λ any measure equivalent to $\frac{1}{2}(\mu + \nu)$ without changing the value of $H(\mu, \nu)$.

By using Hölder's inequality we see that

$$|H(\mu,\nu)|^2 \le \int_E \frac{d\mu}{d\lambda}(x)\lambda(dx) \int_E \frac{d\nu}{d\lambda}(x)\lambda(dx) = 1,$$

so that $0 \leq H(\mu, \nu) \leq 1$.

Exercise 1.3.1 (a) Let $\mu = N_q$ and $\nu = N_{a,q}$, where $a \in \mathbb{R}$ and q > 0. Show that we have

$$H(\mu,\nu) = e^{-\frac{a^2}{4q}}.$$
 (1.3.1)

(b) Let $\mu = N_q$ and $\nu = N_{\rho}$, where $q, \rho > 0$. Show that we have

$$H(\mu,\nu) = \left[\frac{4q\rho}{(q+\rho)^2}\right]^{1/4}.$$
 (1.3.2)

Proposition 1.3.2 Assume that $H(\mu, \nu) = 0$. Then the measures μ and ν are singular.

Gaussian measures

Proof. Set $\alpha = \frac{d\mu}{d\lambda}$, $\beta = \frac{d\nu}{d\lambda}$. Since $H(\mu, \nu) = \int_{\Omega} \sqrt{\alpha\beta} \ d\lambda = 0$, we have $\alpha\beta = 0$, λ -a.e. Consequently, setting

$$A = \left\{ \omega \in \Omega : \ \alpha(\omega) = 0 \right\}, \quad B = \left\{ \omega \in \Omega : \ \beta(\omega) = 0 \right\},$$

we have $\lambda(A \cup B) = 1$. This means that $\lambda(C) = 0$ where $C = \Omega \setminus (A \cup B)$, and hence $\mu(C) = \nu(C) = 0$. Then, as

$$\mu(A) = \int_A \alpha \ d\lambda = 0, \ \nu(B) = \int_B \beta \ d\lambda = 0,$$

we have that μ and ν are singular since

$$\mu(A\cup C)=\nu(B)=0, \ (A\cup C)\cap B=\emptyset. \ \Box$$

Proposition 1.3.3 Let $\mathcal{G} \subset \mathcal{E}$ be a σ -algebra, and let $\mu_{\mathcal{G}}$ and $\nu_{\mathcal{G}}$ be the restrictions of μ and ν to (E, \mathcal{G}) . Then we have $H(\mu, \nu) \leq H(\mu_{\mathcal{G}}, \nu_{\mathcal{G}})$.

Proof. Let $\lambda_{\mathcal{G}}$ be the restriction of λ to (E, \mathcal{G}) . It is easy to check that

$$\frac{d\mu_{\mathcal{G}}}{d\lambda_{\mathcal{G}}} = E_{\lambda} \left(\frac{d\mu}{d\lambda} \Big| \mathcal{G} \right) \quad \frac{d\nu_{\mathcal{G}}}{d\lambda_{\mathcal{G}}} = E_{\lambda} \left(\frac{d\nu}{d\lambda} \Big| \mathcal{G} \right), \ \lambda\text{-a.e.}^{(6)}$$

Consequently we have $(^7)$

$$H(\mu_{\mathcal{G}},\nu_{\mathcal{G}}) = \int_{E} \left[\mathbb{E}_{\lambda} \left(\frac{d\mu}{d\lambda} \Big| \mathcal{G} \right) \mathbb{E}_{\lambda} \left(\frac{d\nu}{d\lambda} \Big| \mathcal{G} \right) \right]^{1/2} d\lambda.$$

Since λ -a.e.

$$\frac{\left[\frac{d\mu}{d\lambda} \frac{d\nu}{d\lambda}\right]^{1/2}}{\left[\mathbb{E}_{\lambda}\left(\frac{d\mu}{d\lambda}|\mathcal{G}\right) \mathbb{E}_{\lambda}\left(\frac{d\nu}{d\lambda}|\mathcal{G}\right)\right]^{1/2}} \leq \frac{1}{2}\left(\frac{\frac{d\mu}{d\lambda}}{\mathbb{E}_{\lambda}\left(\frac{d\mu}{d\lambda}|\mathcal{G}\right)} + \frac{\frac{d\nu}{d\lambda}}{\mathbb{E}_{\lambda}\left(\frac{d\nu}{d\lambda}|\mathcal{G}\right)}\right),$$

taking conditional expectations of both sides one finds, λ -a.e.,

$$\left[\mathbb{E}_{\lambda}\left(\frac{d\mu}{d\lambda}\middle|\mathcal{G}\right)\ \mathbb{E}_{\lambda}\left(\frac{d\nu}{d\lambda}\middle|\mathcal{G}\right)\right]^{1/2} \ge \mathbb{E}_{\lambda}\left(\left(\frac{d\mu}{d\lambda}\right)^{1/2}\ \left(\frac{d\nu}{d\lambda}\right)^{1/2}\middle|\mathcal{G}\right). \quad (1.3.3)$$

 ${}^{6}E_{\lambda}(\eta|\mathcal{G})$ is the conditional expectation of the random variable η with respect to \mathcal{G} and measure λ .

⁷For positive numbers $a, b, c, d, \sqrt{\frac{ab}{cd}} \leq \frac{1}{2} \left(\frac{a}{c} + \frac{b}{d}\right)$.

Integrating with respect to λ both sides of (1.3.3), the required result follows. \Box

Now let us consider two sequences of measures (μ_k) and (ν_k) on $(\mathbb{R}, \mathcal{B}(\mathbb{R}))$ such that $\nu_k \sim \mu_k$ for all $k \in \mathbb{N}$. We set $\lambda_k = \frac{1}{2}(\mu_k + \nu_k)$, and we consider the Hellinger integral

$$H(\mu_k,\nu_k) = \int_{\mathbb{R}} \left[\frac{d\mu_k}{d\lambda_k}(x) \frac{d\nu_k}{d\lambda_k}(x) \right]^{1/2} \lambda_k(dx), \ k \in \mathbb{N}.$$

Remark 1.3.4 Since (μ_k) and (ν_k) are equivalent, we have

$$\frac{d\mu_k}{d\lambda_k} \frac{d\nu_k}{d\lambda_k} = \frac{d\mu_k}{d\lambda_k} \frac{d\nu_k}{d\mu_k} \frac{d\mu_k}{d\lambda_k} = \frac{d\nu_k}{d\mu_k} \left(\frac{d\mu_k}{d\lambda_k}\right)^2.$$

Thus

$$H(\mu_k, \nu_k) = \int_{\mathbb{R}} \left[\frac{d\nu_k}{d\mu_k}(x) \right]^{1/2} \mu_k(dx).$$
 (1.3.4)

We also consider the product measures on \mathbb{R}^∞

$$\mu = \prod_{k=1}^{\infty} \mu_k, \ \nu = \prod_{k=1}^{\infty} \nu_k,$$

and the corresponding Hellinger integral $H(\mu, \nu)$. As is easily checked we have

$$H(\mu,\nu) = \prod_{k=1}^{\infty} H(\mu_k,\nu_k).$$

Proposition 1.3.5 (Kakutani) If $H(\mu, \nu) > 0$ then μ and ν are equivalent. Moreover

$$f(x) := \frac{d\nu}{d\mu}(x) = \prod_{k=1}^{\infty} \frac{d\nu_k}{d\mu_k}(x_k), \ x \in \mathbb{R}^{\infty}, \ \mu\text{-a.e.}$$
(1.3.5)

Proof. We set

$$f_n(x) = \prod_{k=1}^n \frac{d\nu_k}{d\mu_k}(x_k), \ x \in \mathbb{R}^\infty, \ n \in \mathbb{N}.$$

We are going to prove that the sequence (f_n) is convergent on $L^1(\mathbb{R}^\infty, \mathcal{B}(\mathbb{R}^\infty), \mu)$. Let $m, n \in \mathbb{N}$, then we have

$$\begin{split} & \int_{\mathbb{R}^{\infty}} \left| f_{n+m}^{1/2}(x) - f_{n}^{1/2}(x) \right|^{2} \mu(dx) \\ &= \int_{\mathbb{R}^{\infty}} \prod_{k=1}^{n} \frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \left| \prod_{k=n+1}^{n+m} \left(\frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \right)^{1/2} - 1 \right|^{2} \mu(dx) \\ &= \prod_{k=1}^{n} \int_{\mathbb{R}^{\infty}} \frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \mu(dx) \int_{\mathbb{R}^{\infty}} \left| \prod_{k=n+1}^{n+m} \left(\frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \right)^{1/2} - 1 \right|^{2} \mu(dx). \end{split}$$

Consequently

$$\begin{split} &\int_{\mathbb{R}^{\infty}} |f_{n+p}^{1/2}(x) - f_{n}^{1/2}(x)|^{2} \mu(dx) \\ &= \int_{\mathbb{R}^{\infty}} \left[\prod_{k=n+1}^{n+p} \frac{d\nu_{k}}{d\mu_{k}}(x_{k}) - 2 \prod_{k=n+1}^{n+p} \left(\frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \right)^{1/2} + 1 \right] \mu(dx) \\ &= 2 \left(1 - \prod_{k=n+1}^{n+p} \int_{\mathbb{R}} \left(\frac{d\nu_{k}}{d\mu_{k}}(x_{k}) \right)^{1/2} \mu_{k}(dx_{k}) \right) \\ &= 2 \left(1 - \prod_{k=n+1}^{n+p} H(\mu_{k}, \nu_{k}) \right). \end{split}$$
(1.3.6)

On the other hand we know by assumption that

$$H(\mu, \nu) = \prod_{k=1}^{\infty} H(\mu_k, \nu_k) > 0,$$

or, equivalently, that

$$-\log H(\mu,\nu) = -\sum_{k=1}^{\infty} \log[H(\mu_k,\nu_k)] < +\infty.$$

Consequently, for any $\varepsilon > 0$ there exists $n_{\varepsilon} \in \mathbb{N}$ such that if $n > n_{\varepsilon}$ and $p \in \mathbb{N}$, we have

$$-\sum_{k=n+1}^{n+p}\log[H(\mu_k,\nu_k)]<\varepsilon.$$

By (1.3.6) if $n > n_{\varepsilon}$ we have

$$\int_{\mathbb{R}^{\infty}} |\sqrt{f_{n+p}} - \sqrt{f_n}|^2 d\mu \le 2(1 - e^{-\varepsilon}).$$

Thus the sequence $(f_n^{1/2})$ is convergent on $L^2(\mathbb{R}^{\infty}, \mathcal{B}(\mathbb{R}^{\infty}), \mu)$ to some function $f^{1/2}$. Therefore $f_n \to f$ in $L^1(\mathbb{R}^{\infty}, \mathcal{B}(\mathbb{R}^{\infty}), \mu)$.

Finally, we prove that $\nu \ll \mu$ and $f = \frac{d\nu}{d\mu}$. Let φ be a continuous bounded Borel function on \mathbb{R}^{∞} , and set $\varphi_n(x) = \varphi(P_n(x)), x \in \mathbb{R}^{\infty}$, where $P_n x = \{x_1, \ldots, x_n, 0, 0, \ldots\}$. Then we have

$$\int_{\mathbb{R}^{\infty}} \varphi(P_n x) \nu(dx) = \int_{\mathbb{R}^n} \varphi(P_n x) \nu_1(dx_1) \dots \nu_n(dx_n)$$
$$= \int_{\mathbb{R}^n} \varphi(P_n x) \frac{d\nu_1}{d\mu_1}(x_1) \dots \frac{d\nu_n}{d\mu_n}(x_n) \mu_1(dx_1) \dots \mu_n(dx_n)$$
$$= \int_{\mathbb{R}^{\infty}} \varphi(P_n x) f_n(x) \mu(dx).$$

Letting n tend to infinity, we find

$$\int_{\mathbb{R}^{\infty}} \varphi(x) \nu(dx) = \int_{\mathbb{R}^{\infty}} \varphi(x) f(x) \mu(dx),$$

so that $\nu \ll \mu$. Finally, by exchanging the rôles of μ and ν , we find $\mu \ll \nu$.

1.3.2 The Cameron-Martin formula

We consider here the measures $\mu = N_{a,Q}$ and $\nu = N_Q$, and for any $a \in Q^{1/2}(H)$ we set

$$\rho_a(x) = \exp\left\{-\frac{1}{2}|Q^{-1/2}a|^2 + \langle Q^{-1/2}a, Q^{-1/2}x\rangle\right\}, \ x \in H.$$
(1.3.7)

Let us recall, see §1.2.4, that $W_f(x) = \langle f, Q^{-1/2}x \rangle$ was defined for all $f \in \overline{Q^{1/2}(H)}$. Since $Q^{-1/2}a \in Q^{1/2}(H)$ the definition (1.3.7) is meaningful.