

# Self-consistent optical constants of MgF<sub>2</sub>, LaF<sub>3</sub>, and CeF<sub>3</sub> films

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**Abstract:** Various fluorides are materials in nature that extend their transparency range to the shortest wavelengths in the far ultraviolet (FUV, 100nm< $\lambda$ <200 nm). These are relevant materials to prepare multilayer coatings in the FUV for demanding applications such as space instrumentation for astrophysics, solar physics and atmosphere physics, as well as free electron lasers, plasma diagnostics, synchrotron radiation, lithography, spectroscopy, etc. Multilayer design requires the optical constants of the coating materials. Multilayers optimally alternate two transparent materials with contrasting refractive indices. The optical constants of a low-index material, MgF<sub>2</sub>, and of two high-index materials, LaF<sub>3</sub> and CeF<sub>3</sub>, have been determined in a wide spectral range and are presented here. Thin films of MgF<sub>2</sub>, LaF<sub>3</sub>, and CeF<sub>3</sub> were deposited by boat evaporation onto substrates at 523 K. Transmittance, reflectance, and ellipsometry measurements were performed in ranges jointly covering the 30-950-nm spectral range. This range was extended with literature data and extrapolations to obtain self-consistent optical constants using the Kramers-Kronig (KK) analysis. An iterative, double KK analysis procedure (successive reflectance-phase and  $k$ - $n$  KK analyses) was carried out to obtain a self-consistent set of optical constants per material. With the final data sets, the experimental measurements were satisfactorily reproduced. Global self-consistency of the data sets was successfully evaluated through sum rules; additionally, local self-consistency at each photon energy range was also evaluated through a novel sum-rule method which involves window functions. The new sets of optical constant extend the data availability mainly to the FUV and beyond, particularly for CeF<sub>3</sub>, where few data had been reported.

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**OCIS codes:** (120.4530) Optical constants; (310.6860) Thin films, optical properties; (230.4170) Multilayers; (260.7210) Ultraviolet, vacuum; (260.7190) Ultraviolet; (260.7200) Ultraviolet, extreme.

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## Introduction

Fluorides are common materials in optical coatings. One attractive feature is their wide transparency range, that for various fluorides it extends down to wavelengths in the far ultraviolet (FUV,  $100\text{ nm} < \lambda < 200\text{ nm}$ ). Hence fluorides are first-choice materials to prepare FUV multilayer coatings which are demanded for various applications, such as space instrumentation for astrophysics, solar physics and atmosphere physics [1], as well as free electron lasers [2], lithography [3], plasma diagnostics, synchrotron radiation, spectroscopy, etc. A review of coatings for the FUV, including fluorides, was reported by Thielsch [4].

In order to design optical coatings, optical constants of the constituent materials are necessary. Multilayer coatings typically alternate two transparent materials with contrasting refractive indices. This research investigates the optical constants of MgF<sub>2</sub>, a material with a relatively small refractive index in the FUV, and of LaF<sub>3</sub> and CeF<sub>3</sub>, two materials with a relatively high refractive index in the same range.

Evaporation is a suitable technique to deposit fluoride films that are transparent in the FUV. One advantage of fluorides over oxides or nitrides is that usually their molecules do not dissociate upon evaporation, so that films are easier to grow with a good stoichiometry, and reactive deposition is not required. The optical constants of thin films depend on film-growth conditions, such as substrate temperature, the presence of ion assistance, etc. In the transparency range of the materials, transparency is typically favoured by a higher substrate temperature during thin-film deposition; this is most important in the short FUV close to the material cutoff in order to make the extinction coefficient  $k$  small enough to prepare efficient multilayer coatings. In the transparency range,  $k$  of the fluoride can vary with substrate temperature during deposition over orders of magnitudes. Therefore, optical-constant data should be available for films deposited at a temperature close to the one to be used in the application. Measuring small values of  $k$  is not straightforward since no photometric or ellipsometric variable is made to significantly vary with small values of  $k$ . On the other hand, at wavelengths beyond the fluoride cutoff, the dependence of  $k$  on substrate temperature is expected to be less dramatic. In this research, fluoride films have been deposited by evaporation with thermal sources onto substrates heated at a temperature of 523 K.

Literature has reported several sets of FUV optical-constant data on MgF<sub>2</sub> and LaF<sub>3</sub>. Regarding CeF<sub>3</sub>, almost no FUV data was found for thin films. A thesis dissertation systematically investigated the optical properties of most lanthanide trifluorides [5]; unfortunately, LaF<sub>3</sub> and CeF<sub>3</sub> were not included in that research. Here a summary is presented of the literature providing optical constants of films deposited by evaporation on substrates at a temperature relatively close to the one used in this research. Such data have been often reported in narrow spectral ranges and it typically lacks self-consistency.

A considerable number of papers have been devoted to the optical constants of MgF<sub>2</sub> films deposited by evaporation. Literature provides data on the reflective index  $n$  in a wide range from close to the cutoff ( $\sim 115\text{ nm}$ ) up to the visible [2, 3, 6–9], and an extension up to 2000 nm with a Sellmeier model [10]. Regarding the extinction coefficient  $k$  or the absorption

coefficient  $\alpha$ , there have been data reported between the cutoff up to  $\sim 600$  nm [2, 6–8, 11]. Some papers report on optical constants of  $\text{MgF}_2$  films deposited at considerably larger temperatures [4, 12, 13]. At wavelengths below  $\text{MgF}_2$  cutoff, no data was found for films deposited in conditions similar to the present ones and most reported data are for films deposited at room temperature (sometimes temperature was not specified, which might be interpreted as that there was no substrate heating) [14–18]. Anyway, as mentioned above, the dependence of  $k$  and probably also  $n$  on substrate temperature is expected to be less dramatic at wavelengths shorter than the cutoff. Most papers involve optical measurements, but there are also optical constants obtained from electron energy-loss spectroscopy measurements [19]. Several reports on optical constants of  $\text{MgF}_2$  deposited by other techniques can be also found in the literature [13, 20–22].

There is also a relatively large number of papers reporting optical constants of  $\text{LaF}_3$ . The following summary presents again the literature providing optical constants of  $\text{LaF}_3$  films deposited by evaporation on substrates at a temperature relatively close to the one used in this research. As with  $\text{MgF}_2$ , there are several papers that provide  $n$  in a wide range from close to the cutoff ( $\sim 120$  nm) up to the infrared [2, 3, 6, 7, 23–25], some of them using a Sellmeier model [10, 26]. Regarding  $k$ , again most data are for wavelengths in the UV down to the cutoff [2, 6, 7, 23, 27]. There are also papers with optical constants of  $\text{LaF}_3$  films deposited at much larger temperatures [4, 13]. At wavelengths below  $\text{LaF}_3$  cutoff, only one paper was found, which reports a film deposited by evaporation on a nonheated substrate [28]. Several reports on optical constants of  $\text{LaF}_3$  films deposited by other techniques can be also found in the literature [20, 22, 29].

In contrast, very limited data on optical constants of  $\text{CeF}_3$  films were found in the literature. Hass et al. [25] obtained the optical constants of  $\text{CeF}_3$  films grown with a substrate temperature of 573 K in the 220–2000-nm range. Smith and Baumeister obtained  $n$  in the 250–2000-nm range with a Sellmeier model [10]. With a deposition temperature much larger than here and unspecified between 673 and 973 K, Dujardin et al. [30] plotted the absorption coefficient in the 122–270-nm range of  $\text{CeF}_3$  films deposited by molecular beam epitaxy. Other than experimental data for films, Saini [31] theoretically calculated the refractive index of crystalline  $\text{CeF}_3$  in the 0–30-eV range.

The reported literature lacks self-consistent optical-constant data for these materials, which may result in an inaccurate computation of the optical properties of multilayer coatings with these fluorides. This research is devoted to obtain self-consistent optical constants of thin films of the three fluorides deposited by evaporation onto substrates at 523 K in a broad spectral range. This paper is organized as follows. Section 2 describes sample preparation of thin films deposited by boat evaporation and measurement techniques. Section 3 presents the procedure used to calculate optical constants, which involves the use of two successive Kramers-Kronig (KK) analyses. Section 4 presents optical measurements on thin films of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$  through a combination of reflectance, transmittance, and ellipsometry measurements covering the spectral range from 30 to 950 nm. These measurements were extended to a much larger spectral range with literature data and extrapolations to perform the KK analyses, from which the optical constants were obtained. The self-consistency of the obtained optical constant sets is evaluated both with well-known sum rules that provide the global consistency evaluation [32] as well as with recent sum rules developed to evaluate consistency at the desired spectral range [33].

## 2. Experimental techniques

Thin films of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$  were prepared by evaporation from tungsten boats. Evaporation materials of 99.99% purity ( $\text{LaF}_3$  and  $\text{CeF}_3$ ) and VUV-grade ( $\text{MgF}_2$ ) were used. Deposition was performed in a 50-cm side cube chamber pumped with a turbomolecular system and a liquid- $\text{N}_2$  cooled, Ti sublimation pump. The distance between evaporation source and substrate was 30 cm.

Base pressure was  $\sim 1.5 \times 10^{-5}$  Pa; pressure increased during deposition up to  $4 \times 10^{-5}$  Pa ( $\text{MgF}_2$  and  $\text{LaF}_3$ ) or  $6 \times 10^{-4}$  Pa ( $\text{CeF}_3$ ). Average deposition rate was 0.3 nm/s. Substrate temperature was set at 523 K. A set of three samples were prepared in each run: one on a BK-7 glass substrate for FUV-extreme UV (EUV), and x-ray reflectance measurements ( $\text{MgF}_2$  and  $\text{CeF}_3$ ), one on a  $\text{MgF}_2$  crystal substrate for transmittance measurements, and one on a piece of a Si wafer for ellipsometry measurements and, only for  $\text{LaF}_3$ , also for reflectance and x-ray reflectance measurements. After deposition, samples were let to cool down at a rate not faster than 1 K/min. Film thickness was monitored with a quartz crystal; film thickness was measured a posteriori through stylus profilometry (except for  $\text{LaF}_3$ ), and from the fit to both grazing-incidence x-ray reflectometry and ellipsometry measurements. The available measurements are displayed in Table 1 and the average of these measurements was taken as the final thickness: 44, 48, and 40 nm for the sets of samples of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$ , respectively.

Table 1. Film thicknesses in nm obtained through different techniques

Material	Profilometry	Ellipsometry	x-ray reflectometry	Average
$\text{MgF}_2$	43	44	46	44
$\text{LaF}_3$	-	48	48	48
$\text{CeF}_3$	39	40	40	40

Sample reflectance and transmittance were measured in GOLD's (acronym for Grupo de Óptica de Láminas Delgadas) reflectometer system in the 30-190 nm range. The reflectometer has a grazing-incidence, toroidal-grating monochromator, in which the entrance and exit arms are  $146^\circ$  apart. The monochromator covers the 12.5-200 nm spectral range with two Pt-coated diffraction gratings that operate in the long (250 l/mm) or in the short (950 l/mm) spectral range. A windowless capillary discharge lamp was used in this work. The lamp is fed with various pure gases or gas mixtures with which it can generate many spectral lines to cover the spectral range of interest. The beam divergence was  $\sim 1.7$  mrad. The sample holder can fit samples up to an area of  $50.8 \times 50.8$  mm<sup>2</sup>. A channel electron multiplier with a CsI-coated photocathode was used as the detector. Reflectance and transmittance were obtained by alternately measuring the incident intensity and the intensity reflected or transmitted, respectively, by the sample. Transmittance was measured at normal incidence and reflectance was measured at  $5^\circ$  away from the normal. FUV-EUV reflectance and transmittance uncertainty are estimated to be  $\sim \pm 2\%$ .

In the near UV and longwards, Perkin-Elmer Lambda-9 and Lambda-900 double-beam spectrophotometers were used to measure regular transmittance and specular reflectance with the universal reflectance accessory, respectively, in the 190-2000-nm range. For reflectance and transmittance measurements, samples were situated at  $8^\circ$  and  $1.5^\circ$  from normal incidence, respectively.

Ellipsometry measurements were performed with a GES-5E Sopralab Spectroscopic Ellipsometer. The incidence angle at which measurements were performed was optimized around Si Brewster angle at the Si bandgap by confining the spectral distribution of  $\cos\Delta$  symmetrically around zero in order to maximize accuracy.

Grazing incidence x-ray reflectometry measurements were performed at Centro de Asistencia a la Investigación, Universidad Complutense de Madrid. The diffractometer was a PANalytical X'pert PRO MRD. The source was a Cu anode under 45 kV discharge. The Cu  $K\alpha$  ( $\lambda = 0.154$  nm) line was selected by means of a graphite monochromator. Measurements were performed at the grazing incidence angles from  $0.15^\circ$  to  $2.5^\circ$ , with a step of  $0.005^\circ$ . x-ray reflectometry was used to evaluate film thickness, surface roughness, and density. To fit these parameters, the optical constants of the materials were downloaded from CXRO's web [34] and reflectance versus incidence angle were calculated and compared with measurements. At these short wavelengths,  $k$  and  $\delta = 1-n$  can be considered proportional to

density. Film density, thickness, and roughness were let to vary until the best match was obtained.

### 3. Optical constants calculation procedure

The optical constants of MgF<sub>2</sub>, LaF<sub>3</sub>, and CeF<sub>3</sub> films were calculated through an iterative double KK analysis. The present procedure reproduces the scheme followed in a research performed to obtain the optical constants of SiO<sub>2</sub> and Ta<sub>2</sub>O<sub>5</sub> films [35]. The process was iterated twice for each material until a self-consistent set  $\{n, k\}$  was obtained that best fitted all ellipsometry and photometry measurements. One of the KK analyses was performed in reflectance and the other one in the extinction coefficient  $k$ :

$$\phi(E) = -\frac{E}{\pi} P \int_0^{\infty} \frac{\ln[R(E')]}{E'^2 - E^2} dE' \quad (1)$$

$$n(E) - 1 = \frac{2}{\pi} P \int_0^{\infty} \frac{E'k(E')}{E'^2 - E^2} dE' \quad (2)$$

where  $P$  stands for the Cauchy principal value. Equations (1) and (2) relate reflectance with its phase, and  $k$  with the real part of the refractive index  $n$ , respectively. Complex reflectance is given by  $r = R^{1/2} \exp(i\phi)$ , where  $R^{1/2}$  and  $\phi$  stand for its modulus and phase, respectively. Integration is performed over photon energy  $E$ . Each KK analysis requires data over the whole spectrum, for which experimental data are completed with extrapolations.

The need to use the two sorts of KK analyses arises in the initial lack of data for both of them. Hence, to perform the integration in reflectance (odd iterations), available measurements in part of the FUV range were affected by the contribution to reflectance of the film-substrate interface (due to the transparency of the fluoride film), which does not let us obtain  $n$  and  $k$  from  $R$  and  $\phi$ . Regarding integration in the extinction coefficient (even iterations),  $k$  data in the EUV range was not available due to the lack of transparent massive substrates. Each iteration of the present procedure obtains the necessary data to fill in the next iteration. Four iterative KK analyses resulted in self-consistent data sets which reasonably fitted the experimental photometric and ellipsometry measurements.

Experimental data included transmittance, reflectance, and ellipsometry measurements; the covered range was extended with extrapolations. In the following the data sources used to perform KK integrations are specified; moreover, data sources and the iterative procedure are summarized in Table 2:

- i) Between 0.041 nm ( $3 \cdot 10^4$  eV) and 30 nm, Henke's semiempirical data were used, which were downloaded from the Web site of Center for X-Ray Optics (CXRO) at Lawrence Berkeley National Laboratory [34]).  $k$  data was used in  $k$  integration and  $n$  and  $k$  were used to calculate the reflectance of an opaque film in reflectance ( $R$ ) integration.  $k$  was further extended to larger energies with a power function. The densities of the materials, which are required to complete these data, were measured and are presented in subsection 4.4.
- ii) Between 30 nm and  $\sim 114$  nm, there was no experimental  $k$  data due to the lack of transparent massive substrates in most of the range. In this spectral range reflectance was measured, and hence the first iteration was an  $R$ - $\phi$  KK analysis. Out of the  $R$ - $\phi$  data at near-normal incidence,  $n$  and  $k$  can be obtained using the well-known reflectance formula of a single interface:

$$r = \frac{(n + ik) - 1}{(n + ik) + 1} \quad (3)$$

where  $r = R^{1/2} \exp(i\phi)$ . For this formula to apply, the fluoride layer must be opaque, which is satisfied in this range. Eq. (3) is inverted to obtain the optical constants:

$$n + ik = \frac{1 + r}{1 - r} \quad (4)$$

Eq. (4) is a complex equation and hence both  $n$  and  $k$  can be calculated.  $k$  data obtained through Eq. (4) are used in the even iterations.

- iii) Between ~114 and 190 nm,  $k$  was calculated out of transmittance measurements. For the first KK analysis in  $k$  (iteration 2),  $k$  was calculated with the well-known Beer-Lambert law:

$$\frac{T_{fs}}{T_s} = \exp\left(-\frac{4\pi kz}{\lambda}\right) \quad (5)$$

where  $T_s$  and  $T_{fs}$  are the transmittance of the uncoated  $\text{MgF}_2$  substrate and of the substrate coated with the film, respectively;  $z$  and  $\lambda$  represent film thickness and wavelength. Light reflected both at the outer as well as at the inner interfaces is neglected in Eq. (5). After a preliminary set of optical constants was available, exact transmittance, including internal interferences and reflectance, could be calculated in iteration 3 [hence Eq. (5) was not used in this iteration].

Regarding reflectance, experimental values were used in the first R- $\phi$  KK analysis, whereas calculated reflectance (from optical constants obtained from iteration 2) was used in the second R- $\phi$  KK analysis. Calculated reflectance of an opaque film was used instead of experimental reflectance because the latter was affected by interferences with the substrate, as fluorides turn transparent in this spectral range. As a reminder, to obtain  $n$  and  $k$  from reflectance using Eq. (4), the contribution from the inner interfaces must be negligible.

- iv) Between 190 and 950 nm, ellipsometry measurements were fitted with one ( $\text{MgF}_2$  and  $\text{LaF}_3$ ) or five ( $\text{CeF}_3$ ) Lorentz oscillators.  $k$  or reflectance were calculated with the single or multiple Lorentz oscillators obtained in the fit and were used in all iterations.
- v) Above 950 nm, the Lorentz oscillator(s) fitted in the 190-950-nm range were extrapolated, which were used to calculate  $k$  or reflectance for the three materials.

The sequence of R- $\phi$  and  $n$ - $k$  KK integrals through four successive iterations is explained in the following. An initial  $\{R_1\}$  set was used for a first R- $\phi$  KK integration, from which an initial  $\{n_1, k_1\}$  set was obtained using Eq. (4). Then, for the first  $n$ - $k$  KK integration, a  $\{k_2\}$  data set was gathered, from which a second  $\{n_2, k_2\}$  set was obtained. Next, a new  $\{R_2\}$  set was gathered for a second R- $\phi$  KK integration, from which a third  $\{n_3, k_3\}$  set was obtained. Finally,  $\{k_4\}$  data set was gathered and used in a second  $n$ - $k$  KK integration, from which a fourth and final  $\{n_{final}, k_{final}\}$  set was obtained.  $\{k_2\}$  and  $\{k_{final}\}$  data sets included parts of  $\{k_1\}$  and  $\{k_3\}$  data obtained in the previous iteration, respectively, as displayed in Table 2.  $\{R_2\}$  included parts of calculated reflectance with  $\{n_2, k_2\}$ . The procedures and their data sources are summarized in Table 2.

The optical properties of fluoride films may somewhat depend on film thickness. When fluoride films become thicker, their packing density often reduces, which must result in some kind of a gradient in the optical constants. Film density of the present fluorides stand at ~95% of the bulk density (see Sub-section 4), which can be considered rather high if we compare it with the low packing density of fluoride films when deposited at room temperature. The thickness at which such gradient must be significant will depend on substrate temperature. For  $\text{MgF}_2$  films deposited by evaporation, Kolbe et al. [13] measured a packing density not far from 100% for a 300-nm thick film deposited at a temperature close to 523 K, so that no significant gradient is expected up to such a film thickness. For  $\text{LaF}_3$  films, the same research did not measure density but calculated the refractive index and absorption coefficient and found a gradient for 300-nm thick films. Thielsch et al. [24] found that  $\text{LaF}_3$  films deposited at 523 K grew with some gradient; they developed a model in which the inner 40-nm thick

sublayer was considered homogeneous. The present  $\text{LaF}_3$  film thickness is only slightly above that limit. Bischoff et al. [36] measured the packing density of  $\text{LaF}_3$  films deposited at 423 K and found it to be within ~90%-100% of the bulk value for films up to 50 nm; the limiting thickness for high packing density can be expected to be larger for films deposited at 523 K than at 423 K. According to the above, the present 48-nm thick  $\text{LaF}_3$  film can be considered mostly homogeneous, although packing density is expected to decrease with larger thickness. We found no data on homogeneity for  $\text{CeF}_3$  films, but its relatively high packing density suggests that the present film is mostly homogeneous. Summarizing, the present films for the three fluorides can be considered homogeneous in depth, but the use of the below optical constants on thicker films, at least for  $\text{LaF}_3$ , may require some gradient model of the layer and new optical measurements to characterize that gradient.

Table 2. Scheme of the data sources for the iterative KK analysis<sup>a</sup>.

Iteration	Spectral range					Produced optical constants
	$\lambda < 30$ nm	30-114 nm	114-190 nm	190-950 nm	$\lambda > 950$ nm	
1. R- $\phi$ KK. $\{R_1\}$	R CXRO	R experimental		R from Lorentz oscillator(s) fitted from E	R from Lorentz oscillator(s) extrapolation	$\{n_1, k_1\}$
2. k-n KK. $\{k_2\}$	k CXRO	k from iteration 1	k from transmittance, neglecting R MgF <sub>2</sub> : T noisy → smooth connection at 180-190 nm	k from Lorentz oscillator(s) fitted from E	k from Lorentz oscillator(s) extrapolation	$\{n_2, k_2\}$
3. R- $\phi$ KK. $\{R_2\}$	R CXRO	R experimental	R calculated from iteration 2 with $\{n_2, k_2\}$	R from Lorentz oscillator(s) fitted from E	R from Lorentz oscillator(s) extrapolation	$\{n_3, k_3\}$
4. k-n KK. $\{k_{\text{final}}\}$	k CXRO	k from iteration 3	k from transmittance. Exact calculations including R and interferences MgF <sub>2</sub> : T noisy → smooth connection at 180-190 nm	k from Lorentz oscillator(s) fitted from E	k from Lorentz oscillator(s) extrapolation	$\{n_{\text{final}}, k_{\text{final}}\}$

<sup>a</sup>: R: reflectance. T: transmittance. E: ellipsometry

Other parameters, such as substrate material and deposition rate, might also somewhat influence the optical constants. The influence of substrate material was not investigated here because the procedure used to obtain the optical constants relies on measurements with various techniques over the spectrum that were performed on samples prepared with different sorts of substrates. Regarding deposition rate, all samples prepared to measure with the different techniques (reflectance, transmittance, ellipsometry) were deposited in the same run and hence their deposition rate was the same; no further effort was performed to repeat the characterization process for samples deposited at different rates.

Even though the three fluorides analyzed here are uniaxial crystals when bulk and hence are birefringent, thin films of these materials deposited at a temperature far below the material melting point are expected to grow in the form of nanocrystals with random orientations, which results in films with no effective birefringence.

## 4. Experimental results

### 4.1. Optical constants of MgF<sub>2</sub>

Figure 1(a) displays the transmittance of the MgF<sub>2</sub> film normalized to the transmittance of the bare MgF<sub>2</sub> substrate; half of the MgF<sub>2</sub> substrate was left uncoated for this purpose. Transmittance in the near UV and at longer wavelengths was indistinguishable from the



transmittance of the bare substrate. Even in the long FUV, absorption in the film was very small so that the transmittance ratio, i.e., normalized to substrate transmittance, was not far from unity and hence transmittance looks somewhat noisy. Figure 1(b) displays the reflectance of the  $\text{MgF}_2$  film deposited on a glass substrate. Figure 2 displays the ellipsometry measurements at  $78^\circ$  on the  $\text{MgF}_2$  film deposited on a Si wafer, and the fit performed with a single Lorentz oscillator. The oscillator was used in the iterative process described in section 3 to calculate the optical constants in the 190-950-nm spectral range and also at longer wavelengths. The Lorentz oscillator parameters of  $\text{MgF}_2$  and also of  $\text{LaF}_3$  and  $\text{CeF}_3$  are given in Table 3. Silicon wafer substrates were previously characterized: ellipsometry measurements were performed on a part of the wafer where no coating was deposited, from which the thickness of the native  $\text{SiO}_2$  was obtained; optical constants of Si [37] and  $\text{SiO}_2$  [38] were used for modelling the Si wafer and its native oxide.

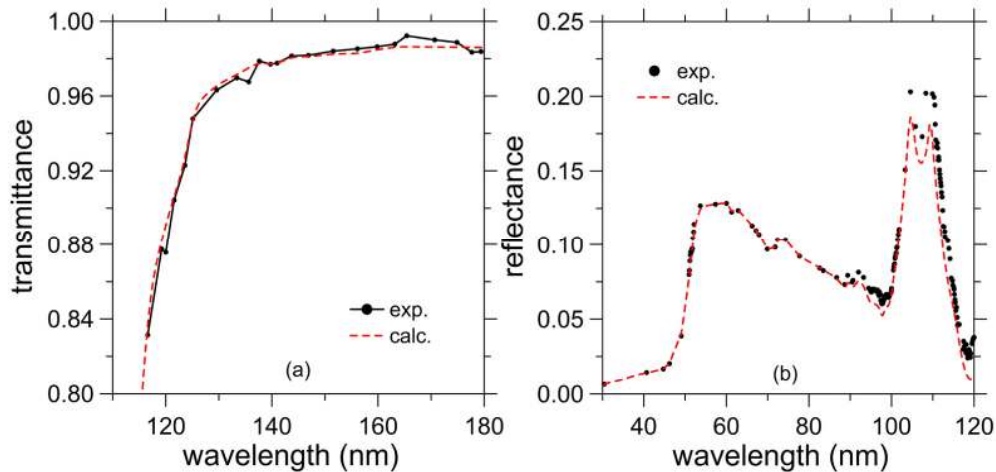


Fig. 1. Comparison between the experimental measurements and calculations with  $\text{MgF}_2$  optical constants obtained in this subsection. a: transmittance (normalized to the transmittance of the bare  $\text{MgF}_2$  substrate) of a 44-nm thick film. b: reflectance of a 44-nm thick  $\text{MgF}_2$  film on a glass substrate

The final self-consistent set of optical constants, referred to as  $\{n_{final}, k_{final}\}$ , obtained after iteration 4, is displayed in Fig. 3. With the final optical constants, normalized transmittance, reflectance, and ellipsometry parameters were calculated and are also plotted in Figs. 1 and 2. There is a satisfactory match between calculation and experimental values.

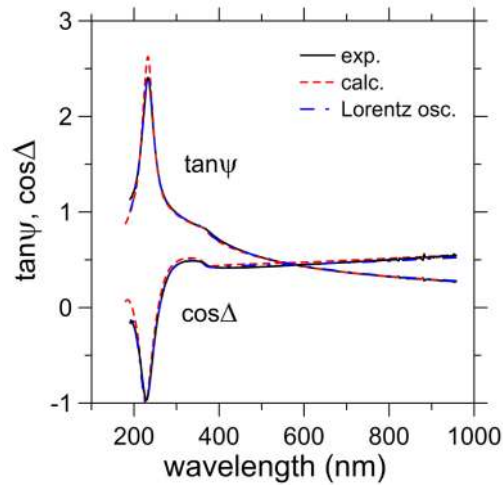


Fig. 2. Ellipsometry measurements at  $78^\circ$  on a 44-nm thick  $\text{MgF}_2$  film deposited on a Si wafer, along with the fit performed with a single Lorentz oscillator and calculations with optical constants obtained in this subsection

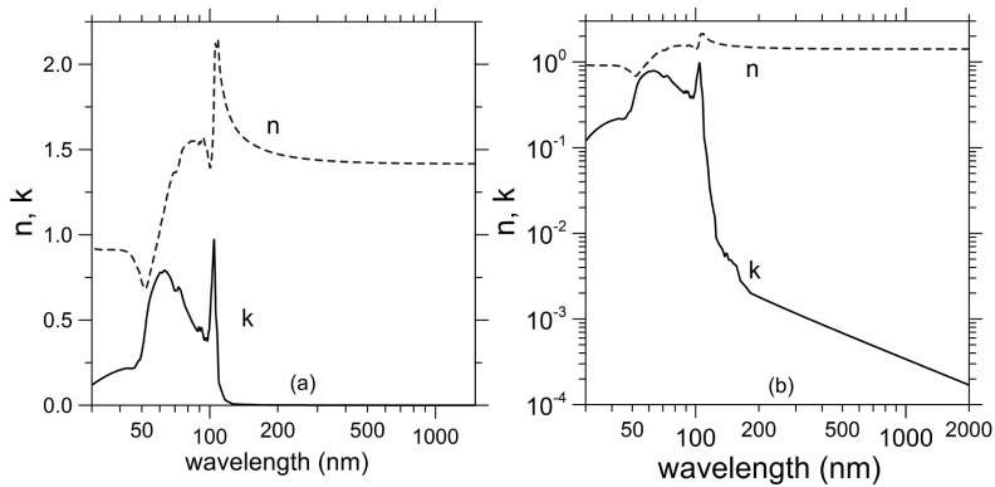


Fig. 3. Optical constants of  $\text{MgF}_2$  films deposited at 523 K (a: linear-axis; b: log-axis) versus wavelength in log scale.

#### 4.2. Optical constants of $\text{LaF}_3$

Figure 4(a) displays the transmittance of the  $\text{LaF}_3$  film normalized to the transmittance of the bare  $\text{MgF}_2$  substrate. Figure 4(b) displays the reflectance of the  $\text{LaF}_3$  film deposited on the Si wafer. Figure 5 displays the ellipsometry measurements at  $74^\circ$  on the  $\text{LaF}_3$  film deposited on Si wafer, and the fit performed with a single Lorentz oscillator, which precisely reproduces the experimental data. The Si wafer was previously characterized as described in sub-section 4.1. The final self-consistent set of optical constants, referred to as  $\{n_{final}, k_{final}\}$  obtained after iteration 4 is displayed in Fig. 6.

Table 3. Parameters of the fits to ellipsometry data with Lorentz oscillators<sup>a</sup> defined as

$$-A\lambda^2/(\lambda_0^2 - \lambda^2 + i\gamma\lambda)$$

Material	$\lambda_0$ (nm)	$\gamma$ (nm)	A
MgF <sub>2</sub>	80.5	1.87	1.73
LaF <sub>3</sub>	69.1	0.49	2.55
CeF <sub>3</sub> , LO1	73.7	2.69	2.39
CeF <sub>3</sub> , LO2	205.7	4.49	0.00081
CeF <sub>3</sub> , LO3	216.6	8.08	0.0060
CeF <sub>3</sub> , LO4	231.7	8.55	0.0061
CeF <sub>3</sub> , LO5	246.8	14.10	0.020

<sup>a</sup>: CeF<sub>3</sub> was fitted to the addition of 5 Lorentz oscillators

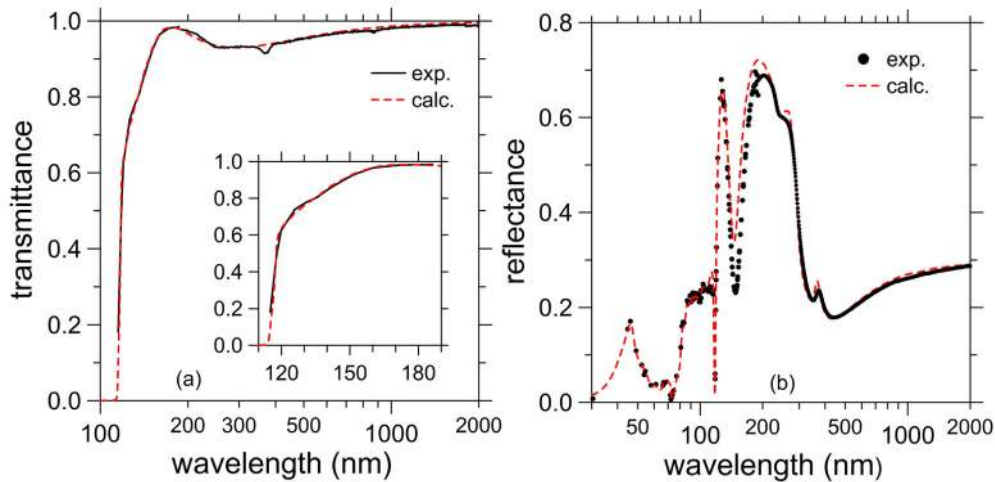


Fig. 4. Comparison between the experimental measurements and calculations with LaF<sub>3</sub> optical constants obtained in this subsection. a: transmittance (normalized to the transmittance of the bare MgF<sub>2</sub> substrate) of a 48-nm thick film; the inset highlights the 110-190 nm range. b: reflectance of a 48-nm thick film on a Si substrate. Wavelength is in log scale

With the final optical constants, normalized transmittance, reflectance, and ellipsometry parameters were calculated and are also plotted in Figs. 4 and 5, respectively. For reflectance calculations in the transparency range of LaF<sub>3</sub>, the Si wafer substrate was modelled with the optical constants mentioned above for ellipsometry fittings.

A small dip in experimental transmittance data was found at around ~360 nm, which was attributed to an artifact of the spectrophotometer, since the instrument changes a filter at that wavelength. Other than that feature, there is a good overall match between calculation and experimental values.

#### 4.3. Optical constants of CeF<sub>3</sub>

Figure 7(a) displays the transmittance of the CeF<sub>3</sub> film normalized to the transmittance of the bare MgF<sub>2</sub> substrate. Figure 7(b) displays the reflectance of the film deposited on a glass substrate. CeF<sub>3</sub> presents weak absorption peaks at various wavelengths (211 nm, 216 nm, 231 nm, and 246 nm) longer than the cutoff at ~120 nm. Similar weak absorption peaks have been reported in the literature [25, 30, 39], which were attributed to the 4f→5d transitions of Ce<sup>3+</sup> ions. Figure 8 displays the ellipsometry measurements at 72° on the film deposited on the Si

wafer and the fit performed with five Lorentz oscillators. The Si wafer was previously characterized as described above. For this material, a single oscillator did not provide a good fit of the experimental data because of the multi-peak absorption observed in the  $\sim 190$ - $250$  nm spectral range; the number of Lorentz oscillators in the fit was increased to five to reproduce this absorption. Four out of the five oscillators are included to fit the four small-absorption peaks that can be seen in the transmittance figure. The other oscillator plays the same role that the single oscillator for  $\text{MgF}_2$  and  $\text{LaF}_3$ : it provides the tail of decreasing  $k$  with increasing wavelength due to the strong absorption below the material cutoff. The final self-consistent set of optical constants, referred to as  $\{n_{final}, k_{final}\}$ , obtained after iteration 4, is displayed in Fig. 9.

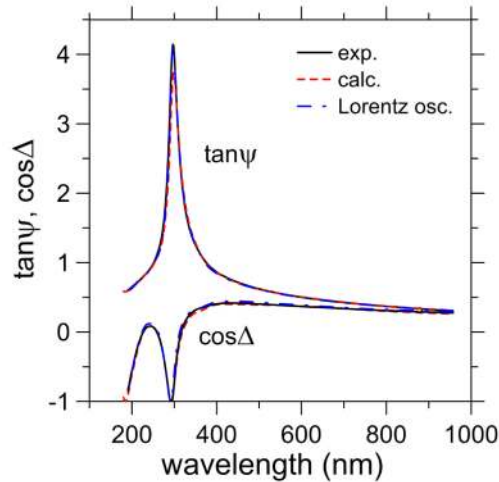


Fig. 5. Ellipsometry measurements at  $74^\circ$  on a 48-nm thick  $\text{LaF}_3$  film deposited on a Si wafer, along with the fit performed with a single Lorentz oscillator and calculations with optical constants obtained in this subsection

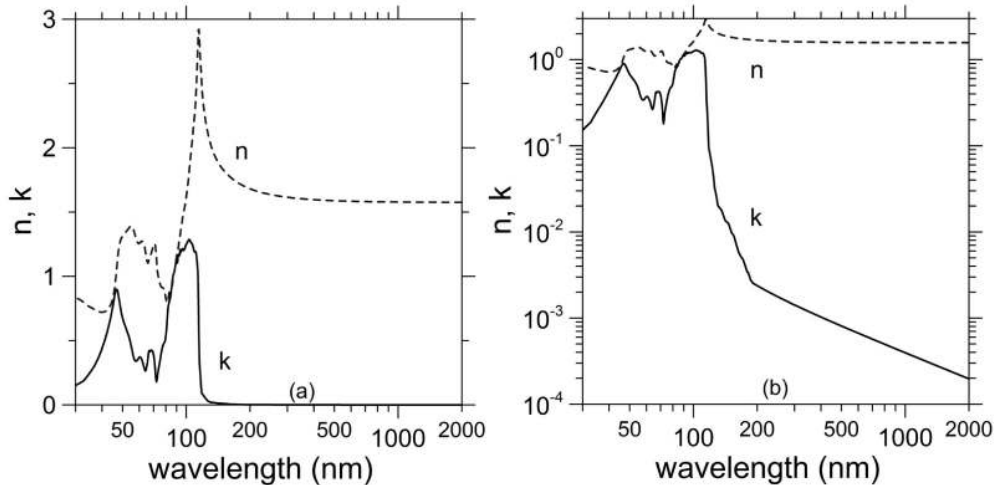


Fig. 6. Optical constants of  $\text{LaF}_3$  films deposited at 523 K (a: linear axis; b: log-axis) versus wavelength in log scale.

With the final optical constants, normalized transmittance, reflectance, and ellipsometry parameters were calculated and are also plotted in Figs. 7 and 8, respectively. There is an excellent match between calculation and experimental values. As with  $\text{LaF}_3$ , a small dip was

found in experimental transmittance data at around  $\sim 360$  nm, which was again attributed to a filter change in the spectrophotometer.

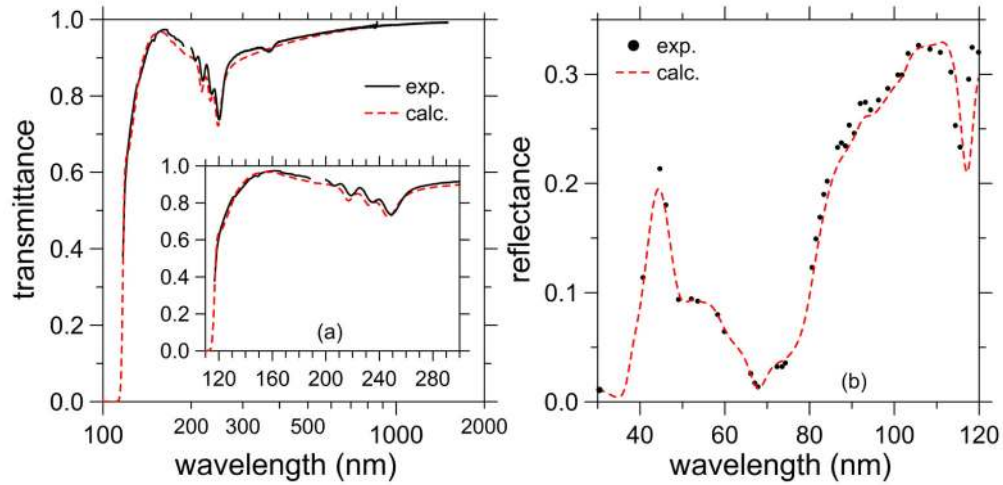


Fig. 7. Comparison between the experimental measurements and calculations with  $\text{CeF}_3$  optical constants obtained in this subsection. a: transmittance (normalized to the transmittance of the bare  $\text{MgF}_2$  substrate) of a 40-nm thick film vs. wavelength in log scale; the inset highlights the 110-300 nm range. b: reflectance of a 40-nm thick film on a glass substrate.

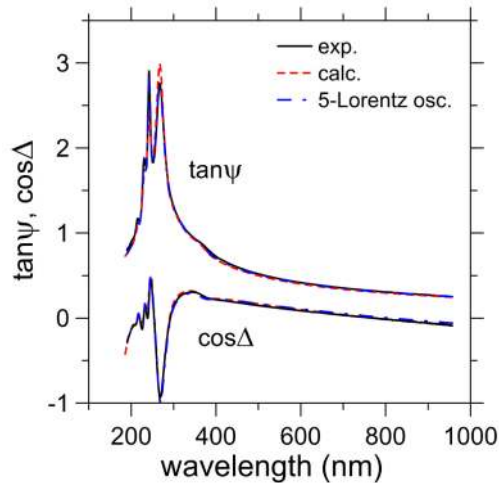


Fig. 8. Ellipsometry measurements at  $72^\circ$  on a 40-nm thick  $\text{CeF}_3$  film deposited on a Si wafer, along with the fit performed with 5 Lorentz oscillators and calculations with the optical constants obtained in this subsection.

#### 4.4. Optical constant consistency

The consistency of the optical constants was evaluated with sum rules, named as  $f$ - and inertial sum rules. Conventional sum rules applied to a set of optical constants provides a global evaluation of their consistency, but give no information of the possible lack of consistency at a specific spectral range. Novel sum rules have been recently proposed [33] that enable obtain information of optical-constant consistency at the desired spectral range. The consistency of the present optical constant sets of fluoride films are checked below both globally and locally.  $f$ -sum rule is expressed by:

$$n_{\text{eff}} = \frac{4\epsilon_0 m}{\pi N e^2 \hbar^2} \int_0^{\infty} E' k(E') dE' \quad (6)$$

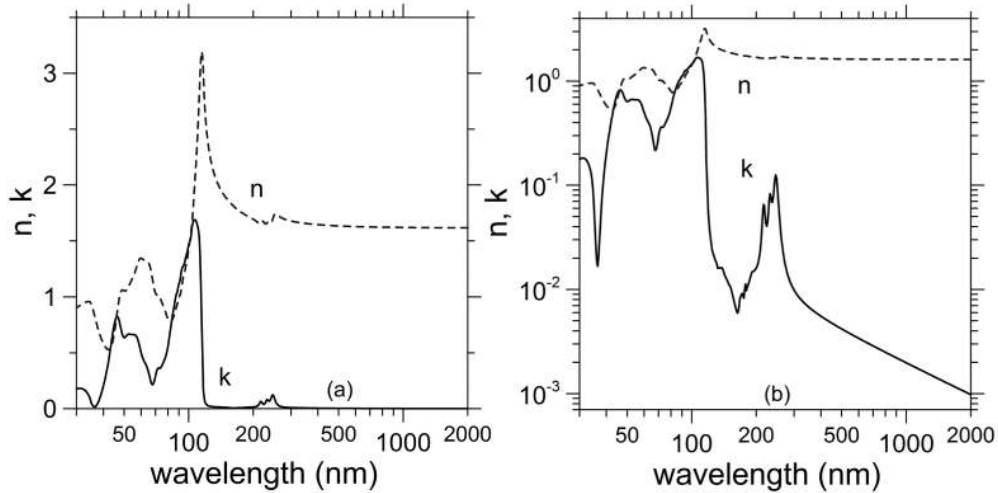


Fig. 9. Optical constants of  $\text{CeF}_3$  films deposited at 523 K (a: linear axis; b: log-axis) versus wavelength in log scale.

where  $e$  and  $m$  are the electron charge and mass, respectively,  $\epsilon_0$  is the permittivity of vacuum,  $\hbar$  is the reduced Planck's constant, and  $N$  is the electron density.  $n_{\text{eff}}$  represents the total number of electrons in the molecule, which should be 30, 84, and 85, for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , respectively. When the relativistic correction on scattering factors is taken into account,  $n_{\text{eff}}$  is somewhat modified: the theoretical effective number of electrons is reduced to 29.98 ( $\text{MgF}_2$ ), 83.53 ( $\text{LaF}_3$ ) and 84.51 ( $\text{CeF}_3$ ) [40]. The density of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$  film samples was measured through grazing incidence x-ray reflectance measurements at  $\text{Cu K}\alpha$  line ( $\lambda = 0.154$  nm), which were fitted with IMD software [41]. Density values of  $2.99$   $\text{g/cm}^3$ ,  $5.79$   $\text{g/cm}^3$ , and  $5.87$   $\text{g/cm}^3$  were obtained for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , which are 94.6%, 97.5%, and 95.3% of their bulk density, respectively. The integral of Eq. (6) with the present  $k$  data results in 30.62, 83.01, and 86.21 for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , respectively. The deviations are 2.1%,  $-0.6\%$  and  $2.0\%$  for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , respectively, which can be considered acceptable numbers.

The other global consistency check is obtained with the inertial sum rule:

$$\int_0^{\infty} [n(E) - 1] dE = 0 \quad (7)$$

which is evaluated with the following normalization [42]:

$$\zeta = \frac{\int_0^{\infty} [n(E) - 1] dE}{\int_0^{\infty} |n(E) - 1| dE} \quad (8)$$

A satisfactory range for  $\zeta$  is usually established within  $\pm 0.005$  [43]. Values of  $\zeta = 1.8 \times 10^{-5}$ ,  $-4.8 \times 10^{-5}$ , and  $-0.7 \times 10^{-5}$  were obtained for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , respectively, well within the above limit.

Let us now evaluate the local consistency of the obtained optical constants [33]. The method that is applied enables enhancing the weight of a desired spectral range within the

sum-rule integral. The procedure consists in multiplying the complex refractive index with a complex window function constructed to have more weight at the desired spectral range. The window functions proposed in Ref [33] were used:

$$H_1(E) = \frac{1}{\pi} [L(E; E_2, c) - L(E; E_1, c)] \quad (9)$$

and

$$H_2(E) = \frac{1}{\pi} \left[ L(E; E_2, c_2) + L(E; E_1, c_1) - 2L\left(E; \sqrt{E_1 E_2}, \frac{c_1 + c_2}{2}\right) \right] \quad (10)$$

with:

$$L(E; E_j, c) \equiv \ln\left(E_j^2 - E^2 - icE\right) \quad j = 1, 2 \quad (11)$$

where  $\ln$  stands for natural logarithm, which is understood as the logarithm principal value, i.e., its imaginary part lies in the interval  $(-\pi, \pi]$ .  $[E_1, E_2]$  is the photon-energy spectral range where the above window functions take larger values, in order to increase the relative weight of the optical constants in this spectral range in the sum-rule integration. The shape of  $H_1$  and  $H_2$  can be seen in Ref [33].

The following sum rules were used here [33]:

$$\int_0^{\infty} E'^{-2} \operatorname{Re}\{H_2(E') [N(E') - 1]\} dE' = 0 \quad (12)$$

$$\int_0^{\infty} E'^{-1} \operatorname{Im}\{H_2(E') [N(E') - 1]\} dE' = 0 \quad (13)$$

$$\int_0^{\infty} \operatorname{Re}\{H_2(E') [N(E') - 1]\} dE' = 0 \quad (14)$$

$$\int_0^{\infty} E' \operatorname{Im}\{H_2(E') [N(E') - 1]\} dE' = 0 \quad (15)$$

$$\int_0^{\infty} E'^2 \operatorname{Re}\{H_2(E') [N(E') - 1]\} dE' = 0 \quad (16)$$

$$n_{\text{eff}} = \frac{4m\epsilon_0}{\hbar^2 N e^2 (E_1^2 - E_2^2)} \int_0^{\infty} E'^3 \operatorname{Im}\{H_1(E') [N(E') - 1]\} dE' \quad (17)$$

The last sum rule resembles  $f$ -sum rule and the other five sum rules show similarities with the inertial sum rule. In spite of the resemblance, the novel sum rules involve a larger contribution to the integral of a desired spectral range, which is achieved through the election of  $E_1$  and  $E_2$ .

In order to obtain the consistency at each spectral range, the window function is made to continuously scan the spectrum. Hence the window function is centred at a variable energy  $E_w$ , which is scanned over the spectrum. In the below calculations,  $E_w$  and the window limits  $E_1$  and  $E_2$  were set at  $E_w = (E_1 E_2)^{0.5}$  ( $E_w$  is the geometric average of  $E_1$  and  $E_2$ ) and  $E_2/E_1 = 3$ .

Figure 10 plots  $\zeta$  vs  $E_w$  for the above sets of optical constants of  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$ , respectively, using the five sorts of inertial-like sum-rules (Eq. (12) to Eq. (16)) with window

function  $H_2$  and window parameters  $c_{1(2)} = E_{1(2)}/10$ . A straightforward generalization of the evaluation parameter given by Eq. (8) has been used [33]. Each point in Fig. 10 represents  $\zeta$  calculated with one of the above sum rules where the window function has been centred at the corresponding energy  $E_w$ . Parameter  $\zeta$  comfortably stands within the  $\pm 0.005$  limit in almost all the plotted spectral range investigated for all sum rules and materials. This indicates that the optical constants are locally consistent at each photon energy range in this wide spectrum. There is a trend to diverge at energies larger than the plotted ones, which starts at smaller energies for sum rules with a larger energy power in the integrand.

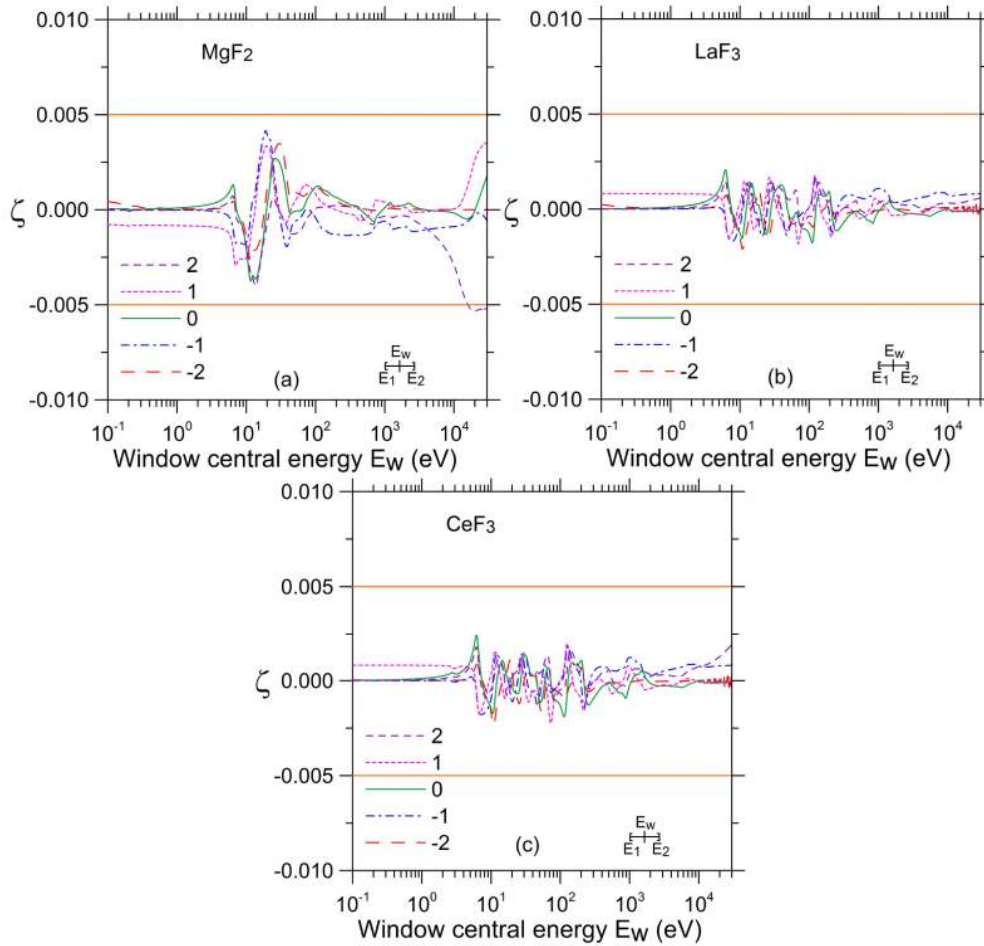


Fig. 10. The evaluation parameter  $\zeta$  versus the central energy  $E_w$  for sum rules represented through Eqs. (12) to (16) calculated with  $H_2$  window function and with the optical constants of  $MgF_2$  (a),  $LaF_3$  (b), and  $CeF_3$  (c). The five inertial-like sum rules are identified in the legend with the power of photon energy in the integral. Window function parameters at each  $E_w$  are given by:  $E_w = (E_1 E_2)^{0.5}$ ,  $E_2/E_1 = 3$ , and  $c_{1(2)} = E_{1(2)}/10$ . The suggested acceptable limits at  $\pm 0.005$  are also plotted

Figure 11 plots  $n_{eff}$  for the sum rule given by Eq. (17) with  $H_1$  window function applied to the optical constants of the three fluorides with the same value of  $E_w$  and with  $c = E_1/5$ . The sum of electrons in the molecules, once corrected for relativistic effects, was mentioned to be 29.98 ( $MgF_2$ ), 83.53 ( $LaF_3$ ) and 84.51 ( $CeF_3$ ); Fig. 11 shows that the deviation from these numbers over the plotted spectral range reaches a maximum of 2.1% and  $-0.6\%$  for  $MgF_2$  and  $LaF_3$ , respectively, which equal what was obtained with the application of the global  $f$ -sum rule; the deviation stands between 2.0% and 2.3% for  $CeF_3$ , the top value only slightly larger



than for the global sum rule. These numbers suggest that the optical constants are consistent at each part of the spectrum. The curves start diverging at large energies, what was also observed with the inertial-like sum rules. Sum rule given by Eq. (17) is strongly biased towards high energies due to the  $E^3$  term within the integrand, which may be the reason why  $n_{\text{eff}}$  starts diverging in this sum rule at a smaller energy than for the inertial-like sum rules. This feature has not been investigated in depth but it suggests that experimental values of the optical constants might be required at larger energies. Nevertheless, since the conjectural lack of self-consistency obtained is found far away from the spectral range covered in this research, it seems not to be a challenge for the consistency of the optical constants in the present range.

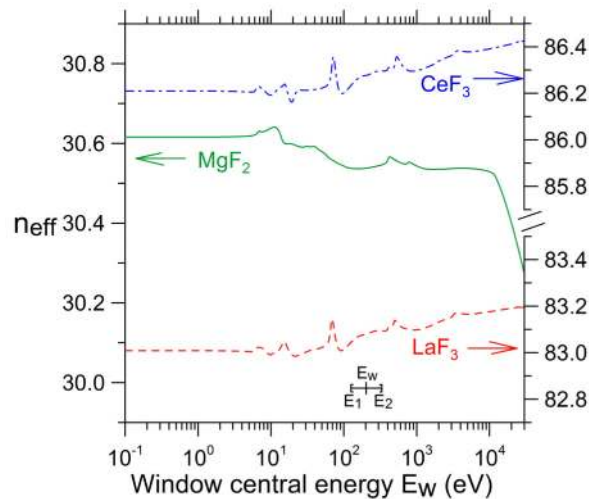


Fig. 11.  $n_{\text{eff}}$  versus the central energy  $E_w$  for sum rule represented through Eq. (17) calculated with  $H_1$  window function and with the optical constants of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$ . Window function parameters at each  $E_w$  are given by:  $E_w = (E_1 E_2)^{0.5}$ ,  $E_2/E_1 = 3$ , and  $c = E_1/5$

Summarizing, both classical  $f$ - and inertial-sum rules, along with new sum rules involving window functions, resulted in satisfactory evaluation parameters, which suggests a good consistency of the sets of optical constants obtained for  $\text{MgF}_2$ ,  $\text{LaF}_3$  and  $\text{CeF}_3$  films. The optical constants here presented are available upon request at the following e-mail address: [j.larruquert@csic.es](mailto:j.larruquert@csic.es).

## Conclusions

Self-consistent optical constants of thin films of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$  deposited by evaporation onto substrates at 523 K have been determined in the spectral range of 30-950 nm. The optical constants obtained for the three fluorides can be useful for the design of interference optical coatings. The optical constants were obtained from a combination of transmittance (FUV range), reflectance (EUV and FUV ranges), and ellipsometry measurements (from the near UV to the near IR ranges), along with extrapolations, from which data sets of  $k$  and reflectance were gathered over the whole spectrum. A consistent set of optical constants for each fluoride was obtained using KK analyses relating both reflectance with its phase and  $k$  with  $n$ ; the KK analyses were applied iteratively until the obtained sets of optical constants were successful to satisfactorily reproduce all experimental measurements.

The consistency of the optical constants of the three fluoride films was evaluated with the use of sum rules. The global consistency of optical-constant data over the whole spectrum was found to be satisfactory through the use of standard  $f$ - and inertial sum rules. Furthermore, the local consistency of the optical-constant data sets at each photon energy

range was also found satisfactory through the use of sum rules with window functions. To the best of our knowledge, these are the first self-consistent sets of optical constants of  $\text{MgF}_2$ ,  $\text{LaF}_3$ , and  $\text{CeF}_3$  thin films deposited onto substrates heated at a temperature close to 523K that involve experimental data from the EUV to the NIR. For  $\text{CeF}_3$  films, no FUV data were found for films deposited at a close temperature (no  $n$  data in the FUV at all), so that the present set provides a broad range of new and self-consistent data.

### Funding

Spanish Programa Estatal de Investigación Científica y Técnica de Excelencia, Secretaría de Estado de Investigación, Desarrollo e Innovación (AYA2013-42590-P and ESP2016-76591-P).

### Acknowledgments

We gratefully acknowledge J. Campos, A. Pons, and J. L. Bris for measurements with the spectrophotometers and I. Carabias and Centro de Asistencia a la Investigación, Universidad Complutense de Madrid, for performing the grazing incidence x-ray reflectometry measurements. Portions of this work were presented at the SPIE conference *Advances in Optical Thin Films V* in 2015; presentation title: “Multilayers and optical constants of various fluorides in the far UV”.