

Title: Sequential establishment of stripe patterns in an expanding cell population

One-sentence summary: A synthetic circuit implementing density-controlled motility autonomously produces a tunable number of stripes.

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Periodic stripe patterns are ubiquitous in living organisms, yet the underlying developmental processes are complex and difficult to disentangle. We describe a synthetic genetic circuit that couples cell density and motility. This system enabled programmed *Escherichia coli* cells to form periodic stripes of high- and low- cell densities sequentially and autonomously. Theoretical and experimental analyses reveal that the spatial structure arise from a recurrent aggregation process at the front of the continuously expanding cell population. The number of stripes formed could be tuned by modulating the basal expression of a single gene. The results establish motility control as a simple route to establishing recurrent structures without requiring an extrinsic pacemaker.

Living organisms display an amazing array of regular spatial patterns (1-4). Traditionally, elucidation of their developmental mechanisms is pursued through forward or reverse genetics (3, 5). However, essential components required for pattern formation and control are often buried in the overwhelmingly complex physiological context. Synthetic biology provides an engineering approach to examine strategies of pattern formation (6-8). Recently, efforts have been made to emulate patterning systems using pre-deposited positional cues (9, 10). These studies do not address spatial or temporal self-organization, which is key to morphogenesis. Indeed, coordinated cell movement in response to self-generated cues is important in embryonic development (11-13). Here, we investigate striking patterns that emerge in growing bacterial populations when cell motility is controlled by cell density, construed as one of the simplest examples of morphogenic cues (14).

We constructed a genetic circuit to suppress the motility of *Escherichia coli* cells at high cell densities (Fig. 1A). The circuit is composed of two modules: a density-sensing module and a motility-control module (Fig. 1B). The quorum-sensing system in *Vibrio fischeri* was adopted as the density-sensing module to signal the local cell density. This system synthesizes and excretes a small molecule acyl-homoserine lactone (AHL), which, at high extracellular levels (reflecting high cell density), accumulates intracellularly and activates a constitutively expressed regulator, LuxR (15). The motility-control module was devised to modify bacterial motility by regulating the transcription of *cheZ*. *CheZ* deletion causes cells to tumble incessantly, resulting in a non-motile phenotype in semi-solid agar (16). Reintroducing *cheZ* restores cell motility (17). In our circuit, coupling of the two modules was achieved via the lambda repressor (CI): the LuxR/AHL complex drove the expression of CI that in turn repressed *cheZ* transcription (Fig. 1B). Based on this design, we created an engineered strain CL3 (18).

Gene expression of the strain CL3 at various cell densities were quantitatively characterized using qRT-PCR. As designed, *cI* expression level (red) elevated over 40-fold as cell density

increased whilst *cheZ* expression (blue) decreased sharply to 5% of the high level (Fig. 1C). Cell motility was measured in semi-solid agar using a modified Continuous-Fluorescence-Photobleaching (CPB) method (18). It dropped abruptly at a density of $\sim 4 \times 10^8$ cells ml⁻¹ (Fig. 1D, purple symbols).

When a suspension of exponentially growing CL3 cells was inoculated at the centre of a 8.5-cm Petri dish containing 10-ml LB medium and 0.25% agar, a pattern of alternating white (high cell density) and dark (low cell density) stripes developed overnight (Fig. 1E, Movie S1) (18). The stripes formed sequentially, at a spacing of approximately 0.5 cm once every ~ 2 hours. These stripes were stable for days until the agar completely dried up. As a control, strain CL4 (carrying the density-sensing module but with the native *cheZ* regulation) behaved like wild type (16): from the position of the inoculum, two traveling waves successively moved radially outwards, followed by a uniform expansion of cell density without stripes (Fig. 1F and Movie S2). We further verified that each element in the designed genetic circuit is necessary for the stripe pattern formation (Fig. S1). We also surveyed the effect of nutrient content, agar thickness, humidity, starting cell number, volume, and growth phase, and found the pattern to be robust to most conditions tested (Fig. S2). Importantly, the stripe pattern persists even in *cheB cheR* double mutants incapable of chemotaxis [Fig. S2J, (19)] but still able to swim and tumble randomly, indicating that the stripes are not a result of chemotactic signaling (20). Also, no stripe pattern resulted if CL3 cells were spotted on agar pre-mixed with a low level of CL3 cells in otherwise identical conditions (Fig. S2K), indicating that the pattern is not a result of spontaneous amplification of density fluctuations, as would follow from the paradigm of Turing instability (21, 22).

We monitored the spatiotemporal cell density profiles from optical absorbance of the plates [Fig. S3 and (18)]. Fig. 1G displays the time course of light intensities along a diagonal that goes through the inoculum, with the color codes indicating the intensity as in Fig. 1E. It shows that once

formed, the pattern of low and high cell densities was frozen in space. Furthermore, new stripes of high cell density emerged at a regular time and distance intervals. The same phenomenon was observed in a rectangular geometry when bacterial growth was initiated along a thin line in the middle of the plate (18). Stripes developed and propagated towards the opposite edges for CL3 (Fig. S4A and Movie S3) but not for the wild type or CL4 cells. Fourier analysis revealed high spatial periodicity, with a wavelength of 0.5 cm (Fig. S4B).

To gain a quantitative understanding of the patterning process, we developed a mathematical model based on the characterized properties of the engineered circuit [Fig. 2 and (18)]. The stochastic swim-and-tumble motion of bacterial cells is described at the population level by a diffusion-like equation (Eq. [1]) for the cell density $\rho(\mathbf{x}, t)$. The synthesis, diffusion, and turnover of AHL are described by Eq. [2], while the consumption and diffusion of the nutrient $n(\mathbf{x}, t)$ are described by Eq. [3]. The AHL-dependent motility enters through a phenomenological diffusion coefficient $\mu(h)$ that drops abruptly where the local AHL concentration $h(\mathbf{x}, t)$ exceeds a threshold level K_h (Fig. 2A). Numerical simulations of Eqs. [1]-[3] (23) using realistic parameter values [Table S3 and (18)] generated patterns similar to the experimental observations in both one- and two- dimensional geometries (Fig. S5 and S20). The model also captures the more elaborate patterns formed by multiple seeding (Fig. S6).

The mathematical model reveals an intriguing sequence of events at the propagating front (Fig. 2B and Movie S4). Early on (T=150 min), cell density (blue line) grows and propagates laterally just as the Fisher wave front (2) since the AHL concentrations (red line) is uniformly below the threshold level K_h (grey dashed line) due to the low initial cell density. A while later (T=300 min), the AHL level exceeds the threshold in a region behind the front due to local cell proliferation: there, a “bud” (marked by B) shoots up in the density profile. Meanwhile, the front (marked by F), which is unaffected by what happens behind due to its low density, continues to

propagate laterally at the same speed. Later (T=450 min), the bud in the density profile grows into a “mound” (marked by M), separated from the frontal region by a “cleft” (marked by C). The continued cell proliferation in the region just behind the front drives up the AHL level and subsequently produces a new bud (B’ at T=600 min), while the lateral movement of the mound and cleft eventually stops (M and C at T=600 min). The new bud starts a new stripe of low motility cells which again expands into a mound (M’ at T=750 min) delimited by a cleft (C’ at T=750 min) from the propagating front, and the process repeats itself. Well behind the front, cell proliferation in both the mound and cleft eventually stops as the nutrient is depleted (height of green shade).

The development of a cleft that separates the propagating front from the dense zone of cells behind, and the budding within the frontal region result from an aggregation phenomenon driven by density-dependent motility discussed in recent theoretical studies (21, 24) and illustrated in Fig. 3A. Cells can diffuse freely in semi-solid agar when the cell density is low (top panel). As cells proliferate and the local AHL level exceeds the threshold, cell motility slows down as programmed (middle panel, with the non-motile cells in teal). These cells cannot move away but neighboring cells may continue to move into this high-density region and become non-motile (bottom panel), leading to a net cell-flow towards the high-density region.

Previously, it has been reported that wild-type *E. coli* cells may aggregate when grown in specific poor-nutrient media due to the secretion of potent chemo-attractants (25). To show that cell aggregation in our system is controlled by the synthetic circuit, we created two strains, the senders (CL6) that synthesize AHL but are non-motile, and the receivers (CL8), capable of receiving AHL and regulating motility but incapable of AHL synthesis. The receiver cells were uniformly mixed with semi-solid agar. Subsequently, a drop of sender cells was spotted at the center of the hardened cell-agar mixture. The non-motile sender cells could not move away from where they were spotted, but were expected to synthesize and excrete AHL as they proliferated at the center of the plate.

After 12-h incubation, a high-density stripe of receiver cells aggregated around the spotted sender cells (Fig. 3B upper panels, Movie S5). No aggregation was observed in a control with receiver cells incapable of reducing motility in response to high AHL concentration (CL10) (Fig. 3B lower panels), nor in other controls (Fig. S7). Thus, an effective aggregation phenomenon was mediated by density-dependent motility.

Detailed analysis of the mathematical model indicates that the aggregation effect alone is not sufficient to generate stripes. Budding in the frontal region requires the AHL level to exhibit a local maximum, which can happen only when the diffusion length of AHL molecules is sufficiently short (so that the AHL profile closely follows the cell density profile), and the cleft in the cell density profile just behind the front is sufficiently deep. The latter is determined in turn by the parameters, such as the maximum cell motility (18, 26). A key output of the model is summarized by the phase diagram (Fig. 4A), which predicts that in addition to the periodic stripe phase, the engineered strain may exhibit a qualitatively different behavior with no stripes, by passing through a transition region exhibiting a limited number of stripes.

A direct test of the occurrence of the phase transition predicted by the phase diagram (Fig. 4A) is to vary the maximum cell motility (D_p in Fig. 2A), thereby changing the number of stripes. Within our experimental design in which motility is set by *cheZ* expression, tuning of D_p can be implemented by adding a second *cI* gene, whose expression is titratable in an AHL-independent manner (Fig. 4B). This is implemented in strain CL5 (18). Fig. 4C shows that the basal *cI* expression level of CL5 cells could indeed be smoothly tuned by adjusting the dosage of an inducer, anhydrotetracycline (aTc), at a fixed cell density (red symbols). Corresponding suppression of *cheZ* expression at low cell density was also observed (blue symbols). For a fixed aTc level, the density dependence of *cheZ* expression remained (Fig. S8). When strain CL5 was spotted at the center of semi-solid agar, consistent with the model predictions, the number of stripes decreased gradually as

the aTc concentration increased from 0.1 to 3.0 ng ml⁻¹ (Fig. 4D). As a control, the pattern of strain CL3 did not change even when the aTc concentration was increased to 100 ng ml⁻¹ (Fig. S2G).

Natural occurrences of well-coordinated spatial and/or temporal patterns are abundant in developmental systems, and are believed to involve elaborate control mechanisms (3, 5). Similarly, stripe formation in various bacterial systems (27-29) has been attributed to complex effects involving chemotaxis, swarming, and differentiation (4). Using synthetic circuits, we demonstrate that precise and robust spatiotemporal patterns can be generated autonomously from a very simple interaction – motility control by density. A recurrent mechanism enables structures to form periodically and sequentially. Important features of the pattern such as the number of stripes can be manipulated by tuning components of the circuit, such as the basal expression level of a single gene.

The strategy of sequential stripe formation is likely not limited to cellular motility control. In the kinetic equations [1]-[3] that generated the stripe patterns, the cell density $\rho(x, t)$ is merely an example of space- and time- dependent variable. Another example could be the concentration of a diffusible molecule, which stimulates its own synthesis to mimic the effect of “cell growth”. Thus Eqs. [1]-[3] may be taken as a generalized reaction-diffusion system with the key feature that the mobility or transport of one of the regulators between adjacent cells is regulated by another mobile regulator. As regulated mobility of signaling molecules is a common strategy in metazoan development (30, 31), the lessons from this work may stimulate new insights and inspire new directions in the studies of developmental systems. For example, the sequential establishment of a periodic somite structure during vertebrate embryonic development is commonly assumed to be controlled by an enabling clock coupled to a sweeping morphogen gradient (32). Our results suggest that a spatially periodic structure can be formed autonomously without the need for a clock.

References and Notes

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23. We believe Eqs. [1]-[3] comprise the minimalistic, experimentally faithful model needed to generate the observed patterns. A simplified model in which the explicit description of nutrient dynamics was replaced by a growth saturation term (Eq. [S11]) was able to initiate but not maintain the patterns, since cell density eventually took on the saturation value everywhere in space (Fig. S17A). A related density-only model (21) generated stripe patterns by balancing cell aggregation with cell death [see also (18)]; it is not applicable to our experiments since cells stop growing but do not die on the relevant time scales. The density-only model also predicted patterns to form from small initial density fluctuation, in contrast to our observation (Fig. S18) and model output (18).
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26. Eqs. [1]-[3] can produce a variety of stripe patterns depending on the values of two key parameters, maximum motility D_p and AHL half life $\ln(2)/\beta$, as summarized in the phase diagram (Fig. 4A), see (18) for details. To the right of the solid red line, the system exhibits the periodic stripe phase with an infinite number of stripes. Moving to the left of the solid line, there is a regime where the system can still generate a limited number of stripes for some initial conditions. Far away from the solid line no stripes can form. The separation between the latter two regimes is not clear-cut and is indicated by the dashed green line. Additional effects such as chemotaxis (Eq. [S17]) have been investigated. They modify the detailed appearances of the patterns as well as the location of the phase boundaries (Fig. S21). Importantly, the general occurrence of the periodic stripes in this class of models does not require such effects.
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Supporting Online Material

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Methods

Model

References (33-63)

Figs. S1 to S21

Tables S1 to S3

Movie S1 to S5

Figure Legends

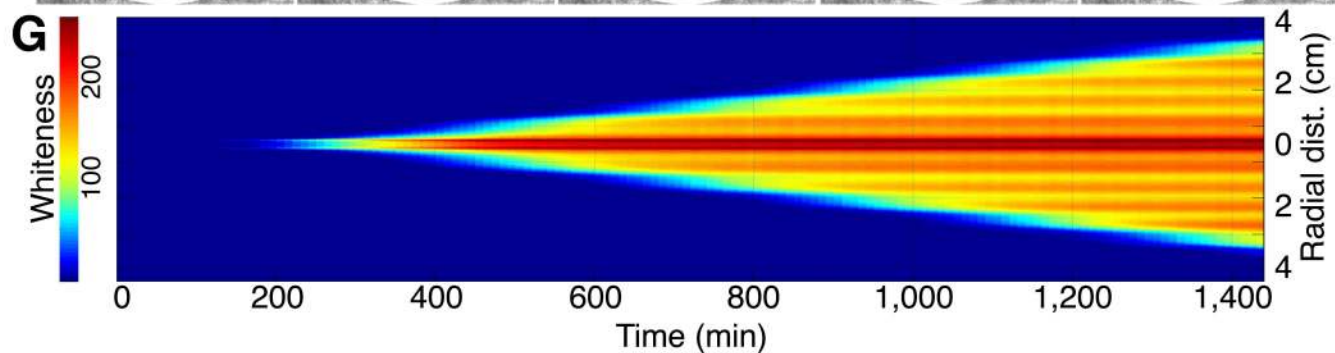
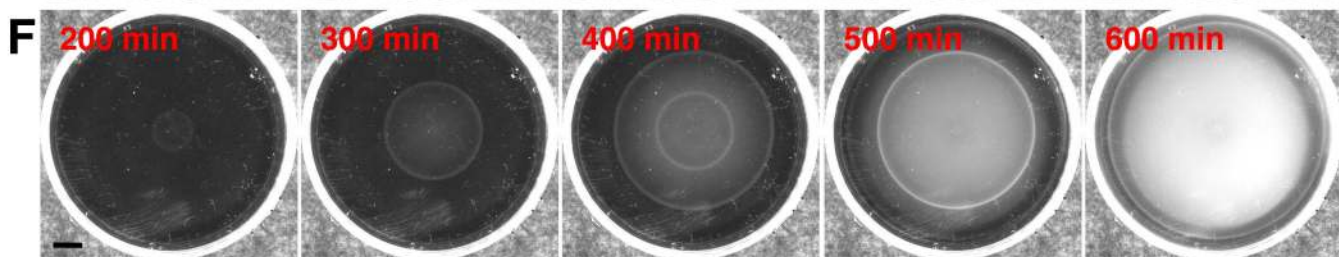
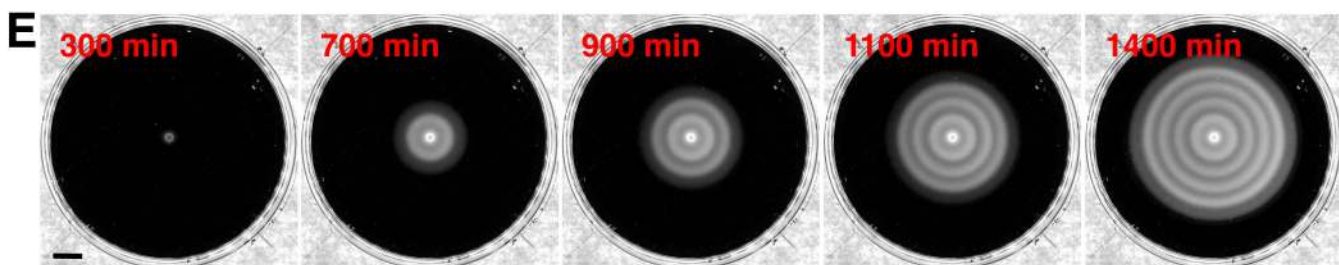
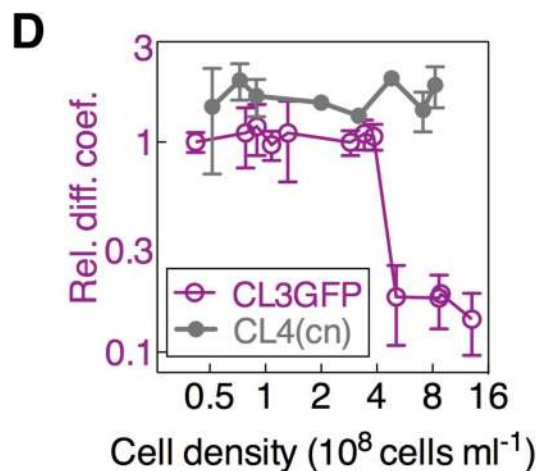
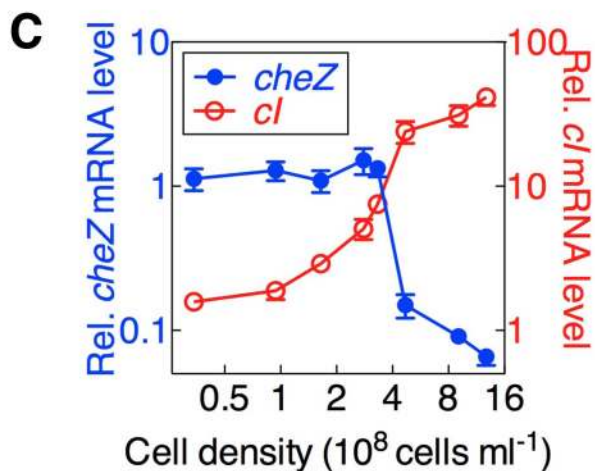
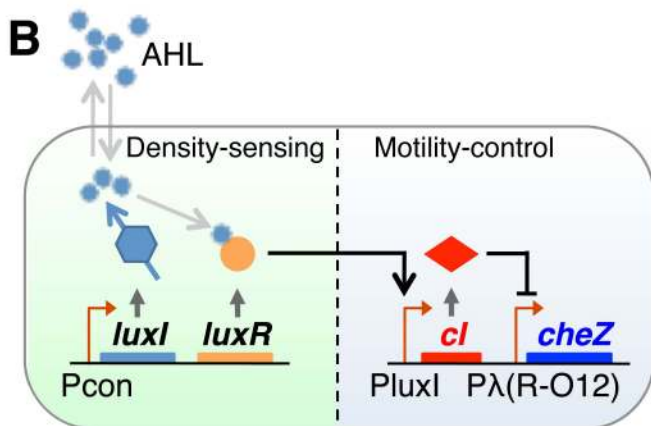
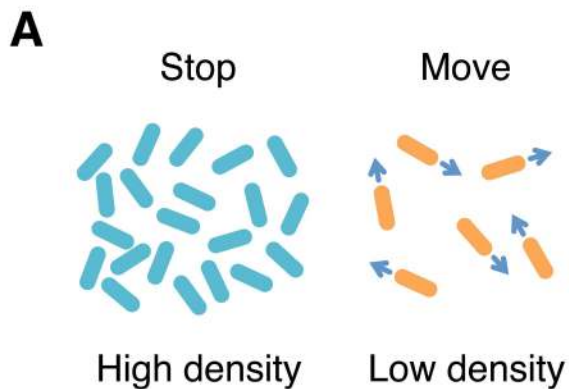
Fig. 1. Spatiotemporal patterns formed by engineered *E. coli* strains. (A) Illustration of the desired cell behavior. (B) Design of the genetic circuit; see text. (C) Relative *cheZ* and *ci* mRNA level in strain CL3 as a function of cell density in bulk culture by qRT-PCR. Data are normalized by the mean value of cells at 0.3×10^8 cells ml^{-1} . (D) The relative diffusion coefficient as a function of cell density. CL3GFP, strain CL3 carrying superfolder GFP. CL4 (Cn), strain CL4 as a control (carrying density-sensing module but with native *cheZ* regulation) (18). Relative diffusion coefficient values are normalized by the mean value of CL3GFP at 0.4×10^8 cells ml^{-1} . Results are representative data from three independent experiments. Error bars represent s.d. of three replicates. Time-lapsed photographs of typical patterns obtained for the engineered strain CL3 (E) and the control strain CL4 (F); see also Movie S1 and S2. (G) Spatiotemporal diagram of (E). All experiments were carried out at 37°C . Scale bars, 1 cm.

Fig. 2. Kinetic model of autonomous periodic stripe pattern formation. (A) The model comprises three key ingredients, cell-density (ρ), AHL concentrations (h), and nutrient levels (n), see text and (18). Crucial feature of this model is the AHL-dependent cell motility $\mu(h)$, in light of the data in Fig. 1D, is modeled by a steep Hill function (red line) with an abrupt transition between the values D_ρ and $D_{\rho,0}$ at $h \approx K_h$ (dashed vertical line). (B) Time-lapse plots of the simulated relative cell-density profiles $\rho(\mathbf{x},t)$ (blue), AHL concentrations $h(\mathbf{x},t)$ (red), and nutrient levels $n(\mathbf{x},t)$ (green shade); see also Movie S4. The AHL threshold of motility regulation K_h is shown as the horizontal dashed line; the vertical axis is rescaled to cell density (10^8 cells ml^{-1}) (18). The periodic pattern is seen as a recurrent process involving the formation of a bud (B), and the bud subsequent grows into

a stationary mound (M) and a cleft (C) structure forms in a region between the bud and the expanding front.

Fig. 3. Cell aggregation driven by density-dependent motility. (A) Illustration of aggregation via density-dependent motility; see text. (B) Experimental evidence of effective aggregation. Receiver cells were uniformly mixed with 0.25% agar. 5- μ l of sender cells (CL6, producing AHL but non-motile, GFP-expressing) were spotted at the center of receiver-cell-agar mixture, followed by 12-h incubation at 37°C. See (18) for details. Upper, strain CL8 (carrying the entire genetic circuit but lacking the AHL-producing gene) as receiver (Movie S5). Lower, strain CL10 (CL8 lacking *ci*) as receiver. Right panels, fluorescence photographs indicate the positions of the sender cells.

Fig. 4. Tunable stripe patterns. (A) Phase diagram; see (18) and (26). (B) Genetic circuit design for tuning stripe number. An aTc-inducible module was added to vary the expression of *cheZ*. (C) Relative *cheZ* and *ci* mRNA level of CL5 in bulk culture containing various concentrations of aTc (0.1-3 ng ml⁻¹). Cells were cultured to OD₆₀₀~0.05 and harvested for measurement. Data were normalized by the mean value of CL3 at 0.3 \times 10⁸ cells ml⁻¹. (D) Experimental patterns of CL3 (also see Fig. S2G) or CL5 inoculated at 0.25% agar containing various concentrations of aTc. Agar plates were incubated at 37°C for 40 h.

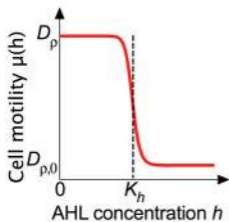


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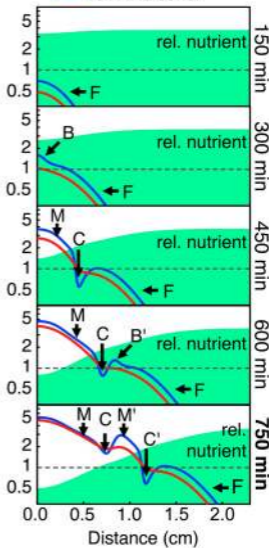
$$\frac{\partial \rho}{\partial t} = \nabla^2 [\mu(h)\rho] + \frac{\gamma n^2 \rho}{n^2 + K_n} \quad [1]$$

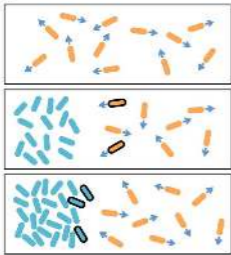
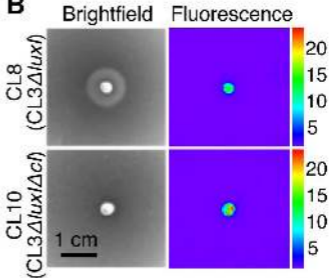
$$\frac{\partial h}{\partial t} = D_h \nabla^2 h + \alpha \rho - \beta h \quad [2]$$

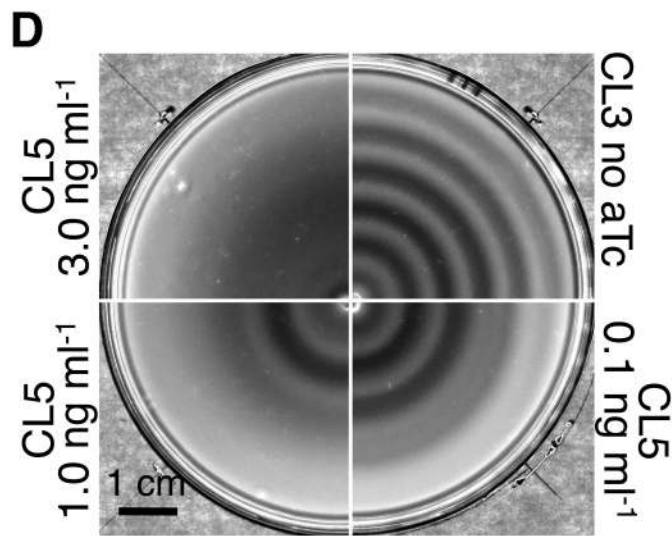
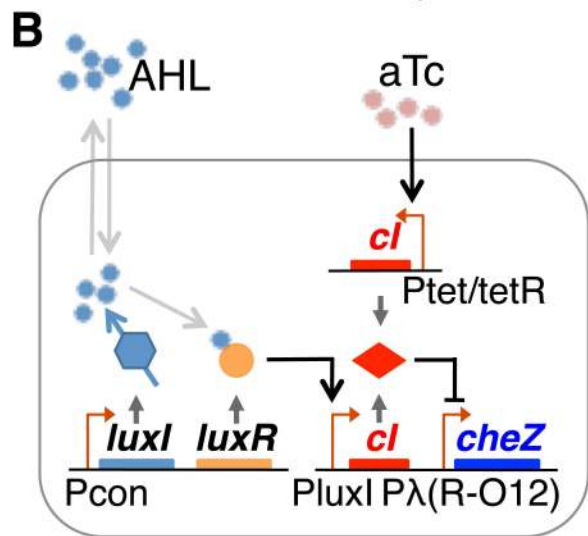
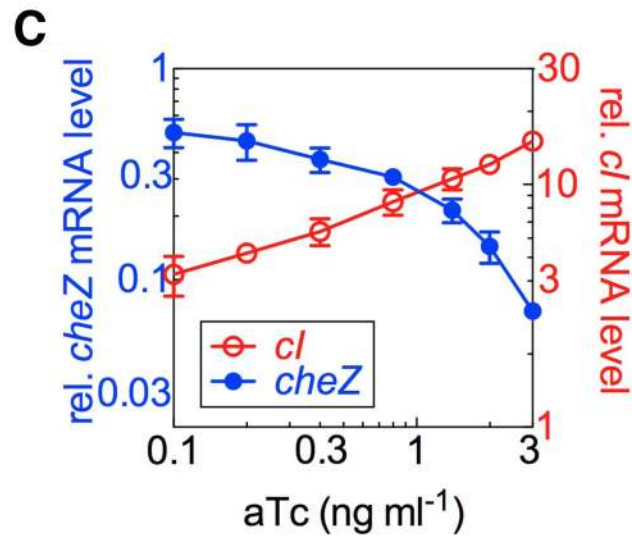
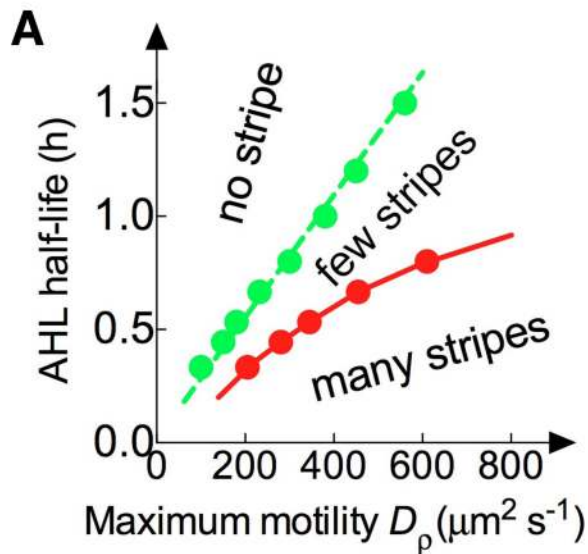
$$\frac{\partial n}{\partial t} = D_n \nabla^2 n - \frac{k_n \gamma n^2 \rho}{n^2 + K_n} \quad [3]$$

**B**

— rel. cell density
— rel. AHL conc.



A**B**



SUPPORTING ONLINE MATERIAL

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SUPPORTING TEXT

METHODS

1. Strain construction

Strains and plasmids used in this study are listed in Table S1. Oligonucleotides used are listed in Table S2.

Strain CL-M is a motile *E. coli* K12 AMB1655 strain, and strain CL-IM is a flagella-lacking mutant of *E. coli* K12 MG1655 strain. Both are kindly provided by Prof. Antoine Danchin (AMAbiotics, France). To rewire the motility regulation, full-length *cheZ* in CL-M was deleted “seamlessly” using a recombineering protocol (33) with the aid of pSIM6 (34). Briefly, loxP flanked chloramphenicol resistance cassette was PCR amplified from pEGFP-loxP-CmR-loxP (33) with primers *cheZ*-del-f and *cheZ*-del-r (Table S2), each composed of a 50-bp sequence at the 5'-end homologous to the *cheZ* gene region. The PCR products were treated with *DpnI* (NEB), gel purified using a QIAquick Gel Extraction Kit (Qiagen), and then electroporated into CL-M cells containing pSIM6 (34) that encodes lambda-Red proteins. Cm^r colonies were verified for replacement of *cheZ* with the loxP-CmR-loxP cassette using colony PCR with primers *cheZ*-del-conf-f and *cheZ*-del-conf-r, and followed by direct sequencing. Removal of the chloramphenicol gene insertion was mediated by p705Cre as described previously (33). pSIM6 and p705Cre were removed from the host by growing the mutants overnight at 37°C, since they both carried temperature-sensitive origin of replication, pSC101. Furthermore, we made seamless deletion of the *lac* operon (including *lacI*) using the same recombineering protocol, to avoid possible interference with the synthetic $P_{lac/ara-1}$ promoter used to drive the expression of *luxRI* in pLuxRI2 and its derivatives. The final mutant, designated as CL1, contains seamless deletion of *cheZ* and the *lac* operon. Similarly, strain CL1RB was constructed by replacing *cheR* *cheB* on the chromosome with a *bla* gene.

Strain CL2 was constructed via P1 transduction with a recipient strain (CL1) and a donor strain (EQ44) (35), by which the $P_{LTet-O1}$ -*tetR* construct (containing the *tetR* gene driven by the synthetic, TetR-repressible $P_{LTet-O1}$ promoter) and the closely linked *bla* gene was integrated into the *attB* site. Ampicillin-resistant colonies were verified for integration of $P_{LTet-O1}$ -*tetR* by PCR and sequencing.

Other strains listed in Table S1 were made by transforming appropriate plasmids into strain CL-M, CL-IM, CL1, CL2, or CL1RB.

2. Plasmid construction

pC*icheZ* (Fig. S9A) was constructed in two steps. First, *luxI* in pLD (9) was replaced by the *cheZ* structure gene amplified by PCR from the genomic DNA of strain CL-M. The ribosomal binding site (RBS) of *cheZ* was kept the same as that of *luxI* in pLD. Second, the resultant plasmid was linearized by inverse PCR with primers pLD-rot-f-*xba* and pLD-rot-r-*cla*, creating *XbaI* and *ClaI* sites, and ligation with PCR-amplified P_{luxI} -*cI* from pLD. pLuxRI/CI (Fig. S9B) was made by inserting a fragment containing P_{tet} -*cI* from restriction-digested pCRTcI into pLuxRI2 (36). pCRTcI was made by inserting PCR-amplified *cI* from pC*icheZ* with *ScaI* and *NotI* restriction sites into pCRT-1 (37).

To construct pLuxRI/GFP, superfolder-gfp (38) was PCR-amplified from BBa_I746916 (Biobrick) with primers gfp-r-sac and gfp-f-sac, composed of a 35-bp sequence of a synthetic constitutive promoter (Biobrick, BBa_J23100) at the 5' end. The fragment was then inserted into pLuxRI2 at the *SacI* site. The colonies with the correct orientation were confirmed by PCR with primers lac-del-conf-f and gfp-conf-r. To make pCIcheZΔcheZ, pCIcheZ was digested with *NheI* and *EcoRV*. The resultant fragment was blunted by T4 DNA polymerase (NEB), then self-ligated. pCIcheZΔcI was made by self-ligating inverse-PCR-amplified fragments from pCIcheZ to exclude the coding regions of *cI*. pLuxR was made by digesting pLuxRI2 with *BamHI* and *NotI* to delete *luxI*, blunting and self-ligating the resultant fragment. Similarly, pLuxI was made by digesting pLuxRI2 with *EcoRI* and *BamHI*. pCIcheZ16S was constructed by inserting PCR-amplified 16S rRNA gene with *Clal* and *XhoI* sites into pCIcheZ. All the constructs were confirmed by DNA sequencing.

Since CL1 and its derivatives are *lacI*-null and also arabinose is never added to the medium, *luxR* and *luxI* in pLuxRI2 and its derivatives may be regarded as constitutively expressed. For simplicity, we refer to the promoter of *luxR* and *luxI* as Pcon.

3. Culture conditions

Unless otherwise stated, the following conditions were used throughout this study. Plasmids were maintained with 50 μg ml⁻¹ Kanamycin and 25 μg ml⁻¹ Chloramphenicol. It had been shown that both AHL degradation rate (39) and *E. coli* cell motility (40) were influenced by pH change in media. To minimize the effects of pH, for bulk culture, we used Luria–Bertani (LB) medium containing 2.5 g Yeast extract, 5 g Bacto Tryptone, 5 g NaCl per liter, buffered by 0.1 M HEPES (pH 8.0). The pH variation was less than 0.3 (measured with PB-10 pH Meter, Sartorius). To prepare semi-solid agar, 10-ml of the above medium containing 0.25% agar (Difco, Bacto agar) was poured into a Petri dish (8.5-cm internal diameter, Sterilin), and allowed to harden at room temperature for 90 min. For the experiments with varying degree of *cI* expression, various concentrations of the inducer anhydrotetracycline (aTc, Clontech) were added to the same media before the agar hardened. When rich defined medium (Teknova, Hollister, CA) was utilized, 0.1 M HEPES (pH 8.0) was also added and 0.5% (w/v) glycerol was used as the carbon source. Unless otherwise stated, all other reagents were from Sigma. All experiments were carried out at 37 °C.

A single frozen glycerol stock was used throughout this study for each bacterial strain. For each experiment, cells were streaked onto LB agar with appropriate antibiotics from corresponding frozen glycerol stock. 3 to 5 colonies were randomly picked for the following experiments. Medium was pre-warmed prior to use. Before each experiment, cells underwent seed culture and preculture, and then went into experimental culture. Strains were first inoculated from fresh single colonies and cultured in 14-ml tubes (Falcon) with vigorous shaking (seed culture), followed by a round of preculture growth in the same medium (until the mid-exponential phase). After that, cells were subcultured (1:200 dilution) into fresh pre-warmed medium with appropriate antibiotics and aTc, when applicable. For gene expression experiments, samples were taken at different time points for measurements of OD₆₀₀ (measured with Cary Bio 300 UV-Vis Spectrophotometer, Varian) and RT-qPCR (see Method S4). For pattern formation

experiments, when OD₆₀₀ reached 0.1-0.2, a suspension of cells (2- μ l) was spotted at the center of a pre-warmed semi-solid agar plate. To measure relative diffusion coefficient in semi-solid agar, after preculture growth, cells were diluted 200-fold into the same pre-warmed media containing 0.25% agar. Cell-agar mixtures were poured into Petri dishes and allowed to harden at room temperature for 90 min, and then moved back into a 37 °C incubator. After various durations of incubation, plates were taken out for measurements of cell density (see Method S6) and diffusion coefficient (see Method S5) in triplicate. Similarly, for the aggregation test, receiver cells were diluted 10-fold into a fresh medium containing 0.25% agar when OD₆₀₀ reached 0.1-0.2. Sender cells were grown to the same OD value and concentrated 16-fold, unless otherwise stated. After the cell-agar mixture solidified, 5- μ l of sender cells was spotted at the center, followed by 12-h incubation at 37 °C.

4. Quantitative Real-Time RT-PCR

Approximately 10⁸ cells from each condition were immediately mixed with RNA Bacteria Protect Reagent (Qiagen). Total RNA was extracted using the RNeasy Mini kit (Qiagen) according to the manufacturer's protocol, followed by treatment with Amplification Grade DNase I (Invitrogen). RNA quality was checked by electrophoresis on agarose gel and the absence of DNA contamination was confirmed by PCR. RNA yield was estimated using a NanoDrop 1000 spectrophotometer (Thermo Scientific). To reverse transcribe RNA, about 0.4 μ g RNA, 200 ng random primer (Invitrogen), 0.5 mM dNTP Mix, and up to 13 μ l RNase-free distilled water was mixed and heated to 65 °C for 5 min and chilled in ice for 1 min. Then 10 μ M DTT, 10 unit SuperScriptTM III Reverse Transcriptase (Invitrogen) and 1 μ l RNaseOUTTM Recombinant RNase Inhibitor (Invitrogen) were added. The reaction was incubated for 5 min at 25 °C, 60 min at 50 °C, and 15 min at 80 °C. Reactions without reverse transcriptase were conducted as controls for the following qPCR reactions. The cDNA samples were diluted 1:40 with PCR grade water and stored at -20 °C until use. Power SYBR Green PCR Master Mix (ABI) was used for qPCR amplification of the amplified cDNA. 1 μ l of the 1:10 cDNA dilution, 200 nM of forward and reverse qPCR primers, 10 μ l of SYBR Green PCR Master Mix, and up to 20 μ l of PCR grade water were mixed in a well of a MicroAmp Fast Optical 96-well reaction plate (ABI). The non-template control (NTC), containing sterile water instead of cDNA template, was included during each qPCR experiment to check the purity of the reagents. Each reaction was performed in triplicate. The qPCR reactions were performed using StepOnePlusTM Real-time PCR system (ABI) with the following programme: 2 min at 50 °C, 10 min at 95 °C, and 40 cycles of denaturation (10 s, 95 °C), annealing and elongation (30 s, 60 °C). Data were acquired at the end of the elongation step. The specificities of accumulated products were verified by melting-curve analysis. A melting curve was run at the end of the 80 cycles to test for the presence of a unique PCR reaction. The qPCR products were checked by electrophoresis on 2% agarose gels. The primers for real-time PCR were designed with the Primer premier 5.0 software package (<http://www.premierbiosoft.com/>) and synthesized by Invitrogen (Table S2). DNA gyrase subunit B (*gyrB*), D-1-deoxyxylulose 5-phosphate synthase (*dxs*) and 16S ribosomal RNA (16S rRNA) were tested as reference genes. They were analyzed in various cell density conditions (OD₆₀₀=0.05-1.2) and aTc concentrations (0-3 ng ml⁻¹).

All of them showed comparative low variation in various conditions. Unless otherwise stated, 16S rRNA was used to normalize the abundance of various genes of interest. The relative fold change in mRNA quantity was calculated for each gene using the relative standard curve method (41). To construct the standard curves for 16S rRNA, *cl*, and *cheZ*, pCIcheZ16S (see Method S2) was 10-fold serially diluted, ranging from 10^{-5} to 10 ng μl^{-1} . For each RNA preparation, at least three independent real-time PCR measurements were performed.

5. Measurement of cell density in semi-solid agar

Following the same principle used in the regular spectrophotometer, we determined the cell density by measuring light transmittance of the cell-containing agar plate. Specifically, we positioned a parallel laser beam from 2 mW/650 nm laser diode (DA650-2-3, Huanic Co., Xi'an, China) perpendicularly to the agar plates and collected the light with a convex lens ($f = 75$ mm) with an amplified Si photodetector (PDA36A, Thorlabs, NJ, USA). To ensure spatial resolution, the laser beam was passed through a 300- μm pinhole. The Petri dish was placed on a motorized stage controlled by a computer running a customized LabVIEW program (National Instruments, Austin, TX, USA) to facilitate spatial and temporal measurements. The whole setup was kept in a warm room (37°C) throughout the experiment. See Fig. S10.

The 10-ml 0.25% agar in Petri dish usually forms an inverse tapered shape with the lowest depth at the center. To avoid such systematic error on the optical path length, we measured the location-dependent transmittances for known standard cell densities. Hence we established a series of standard curves for any given radial positions. The agar plates with standard cell densities were generated with cells collected at the mid-exponential phase ($\text{OD}_{600} = 0.1\text{-}0.2$). The cell growth was stopped by washing twice with nutrient-depleted LB (42), and concentrated to 9.6×10^9 cells ml^{-1} . Subsequently, serial dilutions were carried out. For each cell density, 15 ml of cells was vigorously mixed with an equal volume of pre-warmed nutrient-depleted medium containing 0.5% agar and poured into 3 Petri dishes with 10 ml each. All dishes were allowed to harden at room temperature for 90 min. The final cell densities ranged from 0.03 to 9.6×10^9 cells ml^{-1} . Dishes were moved into a warm room (37°C) and placed on the motorized stage, and light transmittance measurements were taken every 0.2-mm along the line across the center of the dish. The output laser intensities were digitized at 10 KHz (12 bit) using a DAQ device from National Instruments (NI USB-6009, National Instruments, Austin, TX, USA). The average of 1,000 data points was used to reduce the noise in the measurements at each position. The constant input laser intensity was ensured throughout the experiments with a regulated 5V power supply, verified using the photodetector. At each position along the radius, the ratio of the output intensities to the input intensities (measured using a blank agar plate), namely the transmittances (T), were plotted against the known cell densities ρ for 21 different seeding densities (an example curve is shown in Fig. S11A). In total we generated 376 position-dependent density-transmittance standard curves (using a step size of 0.2 mm for a scanning range of 75 mm) and fitted the results to an exponential function for each position x :

$$T(x) = A(x) \exp[B(x) \cdot \rho] \quad [\text{S.1}]$$

where $A(x)$ and $B(x)$ are the position-dependent fitting parameters (Fig. S11B). From these experimentally determined standard curves, we could invert them to compute the real cell density profile $\rho(x)$ from the measured transmittance profile $T(x)$. In subsequent experiments, we scanned the transmittances across the Petri dish for every experiment sample, and computed the spatial cell density profiles. The temporal density profiles were obtained by scanning the agar dish every 10 min with the help of a customized LabVIEW program.

6. Continuous fluorescence photobleaching (CPB) microscopy

A crucial feature needed to characterize the engineered strain is the density dependence of cell motility in semi-solid agar. Microscopy-based single-cell tracing methods are technically challenging to implement at high cell densities where the cells frequently run into each other, and in thick agar where the cells move in-or-out of focus in the z -direction. In addition, a large number of single-cell tracing experiments need to be carried out over long time periods to obtain good statistics. Since the swim-and-tumble motion of *E. coli* cells could be described macroscopically as a diffusive process (43), we aimed to directly quantify the density dependence of diffusion with a population of *E. coli* cells.

Here, we modified a continuous fluorescence photobleaching (CPB) method (44) to quantify the random motion of *E. coli* cells in semi-solid agar. This method utilized a fine laser beam to photobleach green fluorescent proteins (GFP) expressed by *E. coli* cells, and deduced the diffusion coefficient from the decay time of photobleaching. This is based on the fact that the diffusion of *E. coli* cells in and out of the photobleaching region affects the decay of fluorescent intensity. The faster the diffusion, the faster the unbleached GFP-carrying *E. coli* cells refreshed the bleaching region, which lead to slower decay of the total fluorescent intensity observed.

A fluorescence photobleaching microscopy was set up using a Nikon inverted microscope (Eclipse TE2000-E, Nikon Instruments, Kawasaki, Kanagawa, Japan) as illustrated in Fig. S12. An environmental chamber was built maintain the temperature inside at 37 °C throughout the experiments. A laser beam from a 473-nm solid-state laser (MBL-III-473/30mW, Changchun New Industries Co., Changchun, China) illuminated through the semi-solid agar from the top to the bottom, to excite and bleach fluorescent proteins in cells. The emitted fluorescence signal was collected with a 10 × objective (Nikon Plain Fluor 10×/NA0.3) and filtered through a narrow green bandpass filter (center wavelength=515 nm, FWHM=10 nm, Omega Optical, Brattleboro, VT, USA). The fluorescent image was recorded every second with a CMOS camera (DCC1545M, Thorlabs, NJ, USA).

In our experiments, CL3GFP, CL4 (CL-M expressing GFP), or CL14 (CL-IM expressing GFP) cells were uniformly mixed in 0.25% agar at low cell density and cultured at 37 °C (Method S3). At each time point, we measured the average fluorescent intensity in the photobleaching region for each fluorescence image, and obtained a time series of fluorescent intensity for each location of each sample, $i(t)$. To control for variations in cell densities and obtain specific GFP signals, we defined a GFP-specific

relative fluorescent intensity ($I(t)$) as

$$I(t) = \frac{i(t) - i_B(t)}{i(0) - i_B(0)} \quad [\text{S.2}]$$

where $i(t)$ is the average of the fluorescent intensity, and $i(0)$ is the value at $t=0$. Agar plates with similar cell densities but no GFP expression were measured in parallel as $i_B(t)$. The background signal $i_B(t)$ was constituted of autofluorescence from LB medium, agar and cells. The latter had minimal contribution. Examples of the continuous fluorescence photobleaching curve with various fitted values of cell diffusion coefficients are shown in Fig. S13A.

We used strain CL14, a flagella-lacking mutant, as an immobile reference ($D=0$) throughout the study. It provided us with a pure GFP photobleaching curve $I(t)=f(t)$ (Fig. S13A). The curve is best described by a double-exponential decay, with the half-lives of ~ 20 s and ~ 170 s for each component, respectively (Fig. S13A). This two-component photobleaching has been observed in a number of fluorescent proteins (45), especially for the GFP variant (superfolder GFP) we used (38). As we kept the laser power constant, the two decay constants remained the same throughout all our experiments. For the mobile strains, the decays were, as we expected, slower (Fig. S13A). Then we adapted a single-cell simulation-based process to derive the diffusion coefficients from experiments (Fig. S13A). For a given diffusion coefficient, we performed computer simulations to mimic the Brownian motion of many cells, and assigned the time-dependent fluorescent intensity for each cell based on the accumulated time that it was located in the bleaching zone and the two decay constants. Both the decay constants and the size of the bleaching zone used in the simulation were determined experimentally (Fig. S13A). With the trajectories of fluorescent intensities of many cells, we can compute the total fluorescent intensity in the photobleaching zone. We plotted the relation between diffusion coefficients and the averages of relative fluorescent intensities from 175 to 180 s (Fig. S13B). This was used as the standard curve to compute the diffusion coefficients from experimentally measured bleaching curves. However, the exact value of diffusion coefficient depends on evaluation of the sizes of the bleaching zone (R^2). This size is hard to determine due to the lacking of well-defined boundary, as laser beam exhibits a cross-sectional Gaussian intensity distribution. Therefore, we reported relative diffusion coefficients in Fig. 1D.

To measure the relative fluorescent intensity curves for different cell densities in semi-solid agar, cell-containing semi-solid agar plates were prepared as described in Method S3. Every 20-30 min, we took two plates out to measure the cell density (see Method S5), followed by the measurement of the fluorescent intensity curves for 200 s. Three positions were randomly measured for each plate. The exposure time was adjusted between 0.05 and 1 s to match the dynamic range of the camera. The boundary for the photobleaching region was computed using the segmentation algorithm developed by Otsu *et al* (46).

7. One-dimensional pattern formation

We made rectangular plastic plates of two sizes: 3 cm by 7 cm and 6 cm by 10 cm. 3- and 10-ml of 0.25% agar was added to the small and large size plates, respectively.

After the agar hardened, cells were horizontally inoculated as a line using a sterile coverslip (47) and incubated overnight at 37°C.

8. Image capturing

The pseudo darkfield images of agar plates were taken using a Canon EOS 450D digital camera with a Canon EF 50mm f/1.8 II lens and an exposure setting of f5.6/0.6s/ISO200. The agar plates were illuminated by a circular white LED array with the diameter of 30 cm and 10 cm bellow, similar to what was described by Budrene and Berg(48). The time-lapse photographs were acquired with the aid of a timer shutter release controller (TC-80N3a, YongNuo, HK, China). Fluorescence images of the agar plates were captured using a Xenogen IVIS 100 imaging system (Caliper LifeSciences, Hopkinton, MA, USA) with the exposure time of 30 s.

9. Fourier analysis of pattern formation

The time-lapse imaging for 1-D stripe formation was used for Fourier analysis. At each time point, a 1-D intensity profile was computed perpendicular to the stripes. These 1-D profiles were padded with mean intensities to 1024 pixels and applied with a Hamming window. A standard MATLAB (Mathworks, Natick, MA, USA) built-in Fast Fourier Transform (FFT) function was used to compute the mean square power spectra from the 1-D intensity profiles.

MODEL

1. Random walks and diffusion

The run-and-tumble motion of *E. coli* cells has been studied for decades (43). On the molecular level, *E. coli* cells switch between run-and-tumble states of motion by controlling the direction of flagella rotation. When the flagella rotate in a counterclockwise fashion, the bacterium swims forward, otherwise it tumbles. The typical swimming speed of wild type *E. coli* cells in liquid culture is about 10~20 $\mu\text{m s}^{-1}$. Cells tumble once every 0.5~1 s (43). Macroscopically, the continuous run-and-tumble motion of *E. coli* cells can be described by random walks with step size of 10~20 μm , changing direction completely every 0.5~1 s. Such random walks can be described by a diffusion equation, particularly at a population level. *E. coli* K12 cell motion in liquid culture are described with a diffusion coefficient ranging from 200 to 1,000 $\mu\text{m}^2 \text{s}^{-1}$ (49-51).

2. Growth and expansion of the wild type *E. coli* population

For wild type (WT) *E. coli* grown in the semi-solid agar plate (0.1-0.3%), which allows the cells to swim, the spatiotemporal dynamics is governed by two ingredients: cell motility and cell growth. Modeling cell motility by diffusion and cell growth by logistic growth (52) in the simplest case, we have

$$\frac{\partial \rho}{\partial t} = D_\rho \nabla^2 \rho + \gamma \rho \cdot (1 - \rho / \rho_s) \quad [\text{S.3}]$$

for the dynamics of cell density ρ , with D_ρ being the diffusion coefficient, γ being the growth rate, and ρ_s being the saturated cell density (also known as the ‘‘carrying capacity’’). [Note that while the term $-\rho/\rho_s$ in the above equation may be interpreted mathematically as due to cell death (21), bacterial death is negligible over the time scale of our experiments and this term is used here as a mathematical convenient way to describe the effect of nutrient depletion. All the results presented in this work are generated by explicitly including the nutrient field as described below.]

Eq. [S.3] is known as the Fisher’s equation²¹. It admits a propagating front solution for localized initial cell populations. In one dimension, the propagation speed is given by $2\sqrt{D_\rho \gamma}$ (52, 53). In two dimensions, the equation for a radial symmetric profile can be written as

$$\frac{\partial \rho}{\partial t} = D_\rho \left(\frac{\partial^2 \rho}{\partial r^2} + \frac{1}{r} \frac{\partial \rho}{\partial r} \right) + \gamma \rho \cdot (1 - \rho / \rho_s) \quad [\text{S.4}]$$

in polar coordinate with radial distance r . The front propagation speed of the traveling wave approaches that of the 1D case when r is large (53).

Using typical parameters relevant to the WT strain under our experimental conditions, $D_\rho = 200\text{--}1,000 \mu\text{m}^2 \text{ s}^{-1}$, (49-51) $\rho_s = 15 \times 10^8 \text{ cells ml}^{-1}$ (corresponding to $\text{OD}_{600} \sim 2.5$), $\gamma = 1.4 \text{ h}^{-1}$ (doubling time of $\sim 30 \text{ min}$) as summarized in Table S3, the propagating speed is not larger than 2 mm h^{-1} (see Fig. S15A), which is significantly below the observed speed of 5 mm h^{-1} (i.e. $1.4 \mu\text{m s}^{-1}$) (Fig. S15C). One likely reason for this discrepancy is bacterial chemotaxis, which is believed to be the cause for the bright rings leading the propagating front (known as the Adler’s ring) (54); see Fig. S15C. There are a number of models to describe the chemotaxis effect. Here we follow the model by Woodward *et al* (55), in which a drift term (in the direction of the gradient of the attractant a) is added to Fisher’s equation, and growth is coupled directly to the nutrient concentration (n), with a and n both consumed by the cells:

$$\frac{\partial \rho}{\partial t} = D_\rho \nabla^2 \rho - \nabla \cdot \left(\frac{D_\rho K_{C_1} \rho}{(K_{C_2} + a)^2} \nabla a \right) + \frac{\gamma n^2 \rho}{n^2 + K_n^2} \quad [\text{S.5}]$$

$$\frac{\partial a}{\partial t} = D_a \nabla^2 a - \frac{k_a \gamma a \rho}{a + K_a} \quad [\text{S.6}]$$

$$\frac{\partial n}{\partial t} = D_n \nabla^2 n - \frac{k_n \gamma n^2 \rho}{n^2 + K_n^2}. \quad [\text{S.7}]$$

Here, the strength of the chemotaxis is taken to be proportional to the cell diffusion coefficient D_ρ . K_{C_1} and K_{C_2} are two additional parameters associated with chemotaxis. The consumption rates of the chemoattractant a and nutrient n are assumed to be proportional to the growth rate $\gamma \rho$ with proportionality constants k_a and k_n respectively. The parameters K_a and K_n represent the saturation levels of consumption of a and n , respectively.

Eqs.[S.5] -[S.7] again generate a propagating front in one- and two- dimensions, but this time preceded by a ring (Fig. S15B) analogous to the Adler ring observed in experiment (Fig. S15C). The values of the chemotaxis parameters K_{c_1} and K_{c_2} were adjusted to reproduce the observed propagation speed (Fig. S15B). Of the remaining parameters, the values of D_a and D_n were chosen based on the diffusion coefficients of similar small molecules ($\sim 800 \mu\text{m}^2 \text{s}^{-1}$)(56, 57), while k_a and k_n can be scaled out. K_a is unknown, but within a reasonable range does not affect the simulation results much. K_n can be estimated by the cell growth curve measured in the agar plate. From our experimental data (Fig. S16), the cell growth curve for strain (CL4) shows nearly exponential growth until the cell density reaches $10^9 \text{ cells ml}^{-1}$. The cell density for the stationary phase is $15 \times 10^8 \text{ cells ml}^{-1}$ ($\text{OD}_{600} \sim 2.5$), which can be considered as the initial nutrient level in terms of cell density. So K_n is about two thirds of the initial nutrient concentration.

3. Density-dependent motility

To understand how the uniform propagating front solution exhibited by wild type *E. coli* cells break into spatially nontrivial structures generated by the engineered strain, we start by including the effect of density-dependent cell motility of the engineered strain into Fisher's equation (Eq. [S.3]). We first consider an isotropic random walk, *i.e.* without chemotaxis. The diffusion of the bacterial population with diffusion coefficient D_ρ may be characterized in the following microscopic model: Over some small time interval Δt , most cells are localized within a spatial region of linear dimension Δx and only a small fraction $D_\rho \Delta t / (\Delta x)^2$ of randomly chosen cells leave the region (in random directions) to neighboring regions. Density-dependence may be incorporated by multiplying D_ρ by a density-dependent function $\varepsilon(\rho)$. Over a region where density varies in space, the macroscopic dynamics of the population is generally not described by diffusion (53), and the following equation of motion can be derived:

$$\frac{\partial \rho}{\partial t} = D_\rho \nabla^2 (\varepsilon(\rho) \rho) + \gamma \rho \cdot (1 - \rho / \rho_s). \quad [\text{S.8}]$$

More generally, the spatial coupling term proportional to D_ρ is given by a mixture of two terms, $\theta \nabla^2 (\varepsilon(\rho) \rho) + (1 - \theta) \nabla \cdot (\rho \nabla \varepsilon(\rho))$, with the value of the parameter θ being dependent on the details of the microscopic dynamics; *e.g.*, $\theta=1$ for discrete-time (Ito) dynamics and $\theta=1/2$ for continuous-time (Stratanovich) dynamics (58). Detailed derivation of a related chemotaxis model found a mixture of two terms (24). As we will see, the important term for pattern formation is $\nabla^2 (\varepsilon(\rho) \rho)$; to minimize the number of parameters, we will take $\theta=1$ and only keep track of the effect of the $\nabla^2 (\varepsilon(\rho) \rho)$ term in this study.

Eq. [S.8] is linearly unstable. To see this behavior, consider a small density perturbation of amplitude $\delta \hat{\rho}$ and wavenumber \vec{k} around a time dependent and spatially homogenous solution $\rho_0(t)$ that follows logistic growth. If $\rho(\vec{r}, t) = \rho_0(t) + \int \delta \hat{\rho}(\vec{k}) e^{i\vec{k} \cdot \vec{r}} \cdot e^{-\lambda t} dk$, then

$$\lambda = D_\rho k^2 [\varepsilon'(\rho) \rho + \varepsilon(\rho)] \quad [\text{S.9}]$$

So if $\varepsilon'(\rho)\rho + \varepsilon(\rho) < 0$, then Eq. [S.8] would become unstable for all wave numbers, diverging at large k . This means Eq. [S.8] is numerically unstable. To avoid the high wavelength instabilities, we introduce the effect of AHL in accordance with our genetic circuit (see next section). In the recent work by Cates *et al* (21) an ad-hoc k^4 -term was added to Eq. [S.8] to achieve a similar effect.

4. AHL-mediated signaling

In the designed genetic circuit (Fig. 1B), a small molecule AHL was used to signal cell density. Increased cell density (implying increased AHL level) reduced the cell's motility abruptly as shown in Fig. 1D. To describe this effect mathematically, we introduce a motility function $\mu(h)$, which depends explicitly on the AHL level $h(\mathbf{x}, t)$. At constant cell density ρ , macroscopic dynamics is described by diffusion with diffusion coefficient $\mu(h(\rho))$. Assuming that h and ρ are linearly related, then the measured density-dependent diffusion coefficient (Fig. 1D) suggests an abrupt form of $\mu(h)$, which was modeled with a Hill function,

$$\mu(h) = \frac{D_\rho + D_{\rho,0} (h / K_h)^m}{1 + (h / K_h)^m} \quad [\text{S.10}]$$

which has a maximum value of D_ρ switching to a minimum value of $D_{\rho,0} \ll D_\rho$ at $h \approx K_h$ as shown in Fig. 2A.

Including the nonlinear diffusion term as in Eq. [S.8] with $D_\rho \varepsilon(\rho)$ replaced by $\mu(h)$ and further including the dynamics of h , we have

$$\frac{\partial \rho}{\partial t} = \nabla^2 [\mu(h)\rho] + \gamma\rho \cdot (1 - \rho / \rho_s) \quad [\text{S.11}]$$

$$\frac{\partial h}{\partial t} = D_h \nabla^2 h + \alpha\rho - \beta h \quad [\text{S.12}]$$

where α and β are the synthesis and degradation rate of AHL respectively, and D_ρ is the diffusion coefficient of AHL. AHL diffusion is expected to be in the range of small molecule diffusion, $200 \sim 1,000 \mu\text{m}^2 \text{s}^{-1}$ (56, 57). The range of AHL molecule half-life in a slightly alkaline environment (pH7~9) was estimated to be 10~1,000 min (39, 59). The value of α in our experiment is not known and depends on the expression level of AHL-synthesizing LuxI. However, it is only the combination α / K_h , the rate to synthesize enough AHL to reach the level of K_h , that affects the outcome. This ratio can be obtained from the density-dependent diffusion coefficient measured in Fig. S1b, for homogenous distribution of cells of different densities mixed in semi-solid agar plate. The data may be fitted by the form of $\mu(h)$ in Eq. [S.10] with $h = \alpha\rho / \beta$. We obtain $\beta K_h / \alpha \approx 4 \times 10^8 \text{ cells ml}^{-1}$ together with a lower bound of 10 on the Hill coefficient m . In the following calculation, we will use $m=20$.

Eqs. [S.11] and [S.12] were able to initiate patterns. However over time, cell density increased towards ρ_s everywhere in space (see Fig. S17A) and no stable pattern could be maintained. Mathematically, patterns could in principle be stabilized at $\rho > \rho_s$ provided that K_h is large enough so that the instability arising from nonlinear diffusion

(occurring at $\rho \sim \beta K_h / \alpha$) balances negative contribution from the logistic growth term (Fig. S17B). This type of patterns could be initiated from small density variation in the initial conditions due to linear instability as discussed by Cates *et al* (21). However, we believe that they are not applicable to our experimental system for a number of reasons: (i) cell death simulated by the negative growth term does not play a role in our system; and (ii) according to the data $\beta K_h / \alpha \approx 0.25 \rho_s$, so the instability zone is below ρ_s should one wish to cast the numbers within the logistic description (the exact value of ρ_s could be affected by the initial nutrient used; in an experimental relevant range, ρ_s does not affect the pattern, see Fig. S2H). Under these conditions, our model admits no nontrivial stable states. This is also verified experimentally: starting from initial semi-uniform cell distribution (small fluctuation of densities) on the semi-solid agar plate, no periodic patterns can be formed (Fig. S18).

To circumvent the saturation problem (Fig. S17A), we explicitly include the nutrient field $n(\mathbf{x}, t)$ that fuels cell growth and, at the same time is consumed by cell growth. We have

$$\frac{\partial \rho}{\partial t} = \nabla^2 [\mu(h)\rho] + \frac{\gamma n^2 \rho}{n^2 + K_n^2} \quad [\text{S.13}]$$

$$\frac{\partial n}{\partial t} = D_n \nabla^2 n - \frac{k_n \gamma n^2 \rho}{n^2 + K_n^2} \quad [\text{S.14}]$$

where the parameters and their values are as defined in Eqs. [S.5] and [S.7] and summarized in Table S3. [In the above equations, a Hill coefficient of 2 was used following the model by Woodward *et al* (55). However, the results on pattern formation can also be obtained using Hill coefficient of 1.]

The system defined by Eqs. [S.12] -[S.14] is capable of generating stable stripes (60). The spatiotemporal dynamics of the stripe formation process in 1D space, shown in Fig. S5, was generated using the parameter values in Table S3, a Gaussian initial cell density profile:

$$\rho(x, 0) = 2 \exp(-x^2 / 4mm^2) \times 10^8 \text{ cells ml}^{-1} \text{ for } 0mm \leq x \leq 3mm \quad [\text{S.15}]$$

and $15 \times 10^8 \text{ cells ml}^{-1}$ as the initial nutrient level. We have also performed simulations in 2D. The 2D space model gives the stripe pattern (Fig. S19A) similar to the 1D case, except for minor differences in the amplitude of the periodic structure.

Depending on the detailed microscopic dynamics, in principle there could be also a density-dependent drift term as shown in Eq. [S.16] as discussed above and by Tailleur and Cates (24) (who explicitly coarse-grained the dynamics of a related process where the bacteria's run length is density dependent).

$$\frac{\partial \rho}{\partial t} = \theta \nabla^2 [\mu(h)\rho] + (1-\theta) \nabla \cdot (\mu(h) \nabla \rho) + \frac{\gamma n^2 \rho}{n^2 + K_n^2} \quad [\text{S.16}]$$

We examined its effect (using $\theta = 1/2$), together with Eqs. [S.12] and [S.14] in 1D. The stripes pattern can still be generated as shown in Fig. S20A, although the details appear differently from Fig. S5.

Finally, the engineered strain is also expected to exhibit chemotaxis as the wild type. As discussed in the previous section, the effect of chemotaxis can be modeled as a drift term towards some chemoattractant gradient. This is implemented by replacing Eq. [S.13] by

$$\frac{\partial \rho}{\partial t} = \nabla^2 (\mu(h)\rho) - \nabla \cdot \left(\frac{D_\rho K_{c_1} \rho}{(K_{c_2} + a)^2} \nabla a \right) + \frac{\gamma n^2 \rho}{n^2 + K_n^2} \quad [\text{S.17}]$$

and further supplementing Eqs. [S.6], [S.12] and [S.14]. Simulation shows that the stripe pattern formed is robust to the presence of chemotaxis; see Fig. S20B.

5. Phase diagram for stripe formation

Eqs. [S.12] -[S.14] do not generate periodic stripes for arbitrary combination of parameters. As will be explained below, the behavior of this system is controlled mostly by the two dimensionless combinations of parameters, D_ρ / D_h and γ / β . By performing the 1D simulation with varying parameters and a smaller initial cell density:

$$\rho(x, 0) = 2 \exp(-x^2 / 0.1 \text{mm}^2) \times 10^5 \text{ cells ml}^{-1} \quad \text{for } 0 \text{mm} \leq x \leq 0.2 \text{mm} \quad [\text{S.18}]$$

we found that the periodic stripe pattern occurs to the right of the red line in the phase diagram (Fig. 4A). Immediately to the left of the red line, the system can still generate a finite number of stripes if the initial cell density is sufficiently large (e.g. Eq.[S.18]). But in the long-time limit, it eventually turns into a propagating front. No stripes formed far away from the red line.

The separation between “few stripes” and “no stripe”, as indicated by the dashed green line, is not crisp and depends on the form of the initial density profile used (here we used the initial condition as Eq.[S.18]). However, we believe the transition across the red line is a real phase transition. Since a well-defined period can be assigned in one phase while no stripes formed in the long-time limit of the other phase, this transition is expected to be a Hopf bifurcation (61).

Similar phase boundary can be found by including the chemotaxis effect Eq. [S.17] with Eqs. [S.6], [S.12] and [S.14]; see Fig. S21. Over the range of the experimental relevant parameters, chemotaxis could either expand or reduce the parameter space for stripe formation.

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60. Strictly speaking, the pattern formed is not stable at a very long time scale because the diffusion coefficient $D_{\rho,0}$ at high cell densities is not strictly zero. However, practical time scales reachable by simulations, this effect is not noticeable. For the experiments, we expect the cell's motility to cease completely once the nutrient is depleted. Hence motion over the very long time scale is not experimentally relevant.
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SUPPLEMENTARY FIGURES

Figure S1

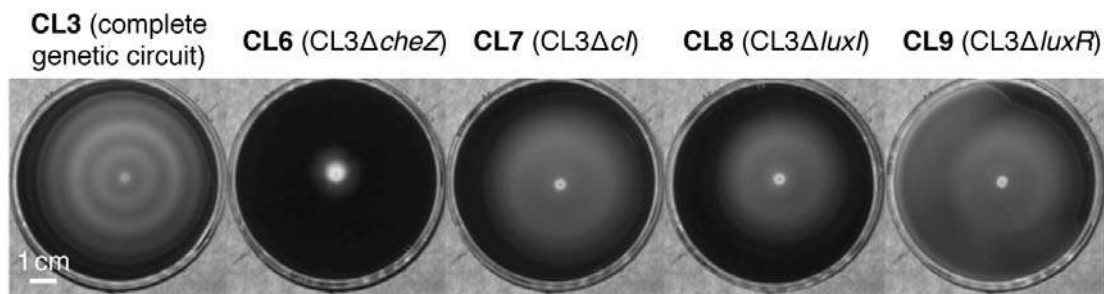


Fig. S1. The entire genetic circuit is required for stripe pattern formation. After seed culture and preculture, strains containing a complete (CL3) or partial genetic circuit (CL6, CL7, CL8, and CL9) were diluted 200-fold into fresh medium and cultured until the mid-exponential phase ($OD_{600}=0.1-0.2$) separately. 2- μ l cell suspension was then inoculated at the center of 0.23% agar plates, followed by 20-h incubation at 37 °C. Photographs were taken against a black background illuminated by an annular LED array from below. The property of each strain is indicated in the brackets.

Figure S2

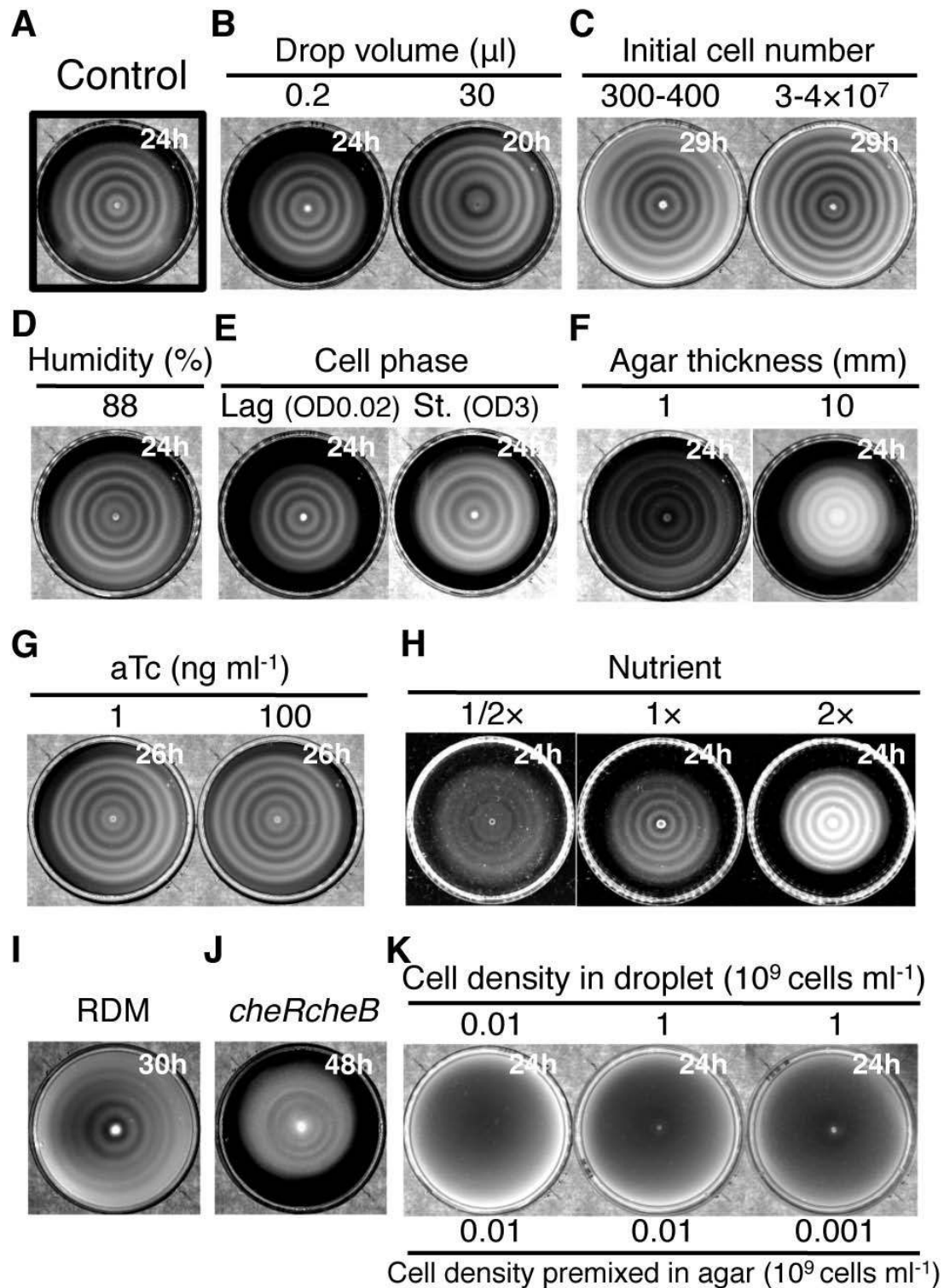


Figure S2. Pattern formation under various conditions. (A) The common experimental conditions applied in this study were used as the control, particularly, the initial drop volume was $2 \mu\text{l}$; initial cell number was $3-4 \times 10^5$; initial cell phase was

mid-exponential phase ($OD_{600}=0.1-0.2$); agar thickness was 2 mm; and humidity was 30%. **(B)** Initial drop volume from 0.2 to 30 μ l with the same cell density. Other conditions were kept unchanged. **(C)** Initial cell number, 1,000-fold dilution with experimental medium, or 100-fold concentration by centrifugation. 2- μ l of cell suspension was inoculated. **(D)** Initial cell phase, after seed-culture, subculture, cells were 200-fold diluted into fresh experimental medium and grown until the lag-phase ($OD_{600}=0.02$), or stationary phase ($OD_{600}=3$). **(E)** Thickness of agar layer, LB medium supplemented with 0.1 M HEPES (pH 8.0) containing 0.25% agar in various volumes, ranging from 5- to 50-ml, were poured into Petri dishes (8.5-cm internal diameter) and allowed to harden at room temperature for 90 min. 2- μ l of cell suspension was inoculated. **(F)** Effect of humidity. The conditions were the same as the control except for 88% humidity. **(G)** Effect of aTc. After seed-culture, cells were subcultured into experimental medium supplemented with various concentrations of aTc, grown until the mid-exponential phase. 2- μ l of cell suspension was spotted at the center of semi-solid agar containing the medium with the same concentration of aTc as subculture. **(H)** Effect of the initial nutrient level. Keep NaCl and HEPES concentration unchanged, other ingredients of LB medium were diluted or concentrated two times. **(I)** Rich defined medium (RDM). The culture conditions were the same as the control except for the medium used (see Method S3). **(J)** *cheR cheB* double mutant. After seed-culture, subculture, CL15 cells (see Method S1 and Table S1) were inoculated onto LB medium supplemented with 0.1 M HEPES (pH 8.0) containing 0.13% agar. **(K)** Initial cell densities in droplet and uniform mixture in agar. Cells were cultured until the mid-exponential phase, diluted or concentrated to the desired cell densities and mixed with pre-warmed 0.25% agar (as described in Method S3). Subsequently, cells in various densities were spotted in the center of the hardened cell-agar mixture. The times in the figures indicate the duration of the agar plate incubated at 37 °C after cells were spotted. Pattern formation is robust in all the above conditions; however, conditions that significantly affect cell growth (e.g., temperature and nutrient) can affect the patterns in complex ways, presumably due to the cell's innate control on motility. Moreover, given larger space and longer time to develop patterns, “loss-of-pattern” mutations were observed.

Figure S3

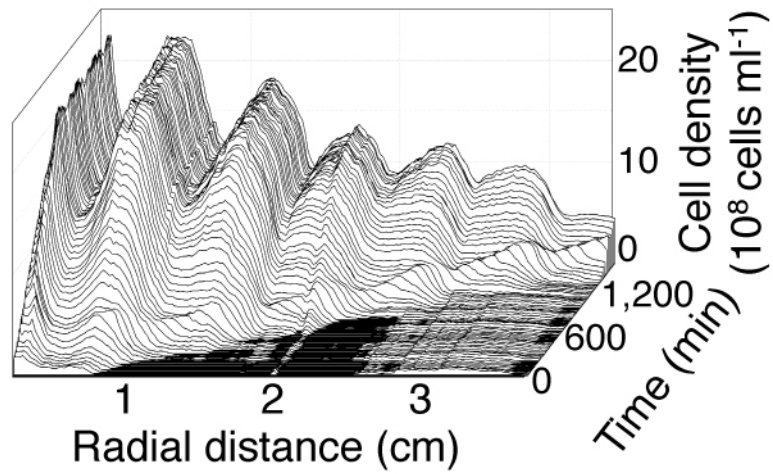


Fig. S3. Spatial-temporal cell density profile. To measure the absolute cell density profile across the center of the Petri dish as a function of time, the Petri dish was placed on a motorized stage, and the cell density profile was obtained by periodically scanning along a line through the center of the agar plate with a fine laser beam during the pattern formation process, and comparing the transmitted light with the standard; see Method S5 for details.

Figure S4

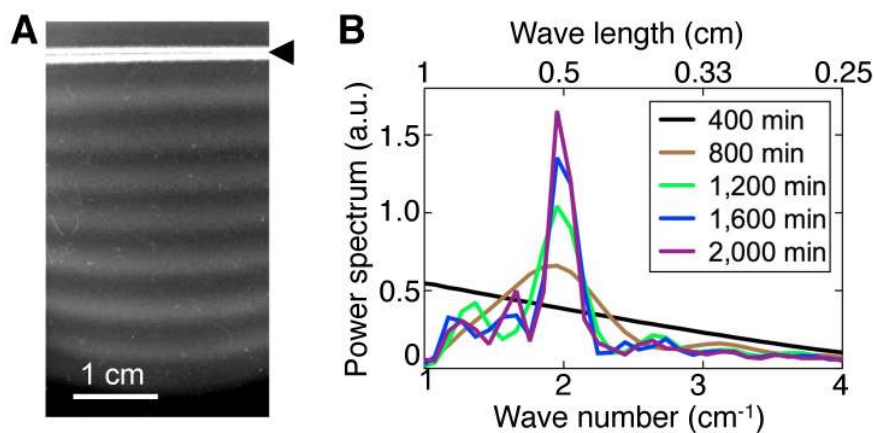


Fig. S4. One-dimensional pattern formed by strain CL3. **(A)** One-dimensional pattern observed 2,000 min after seeding. Arrow points to the position of inoculation (Method S7 and Movie S3). **(B)** Fourier analysis of one-dimensional band pattern formation. To evaluate the periodicity of the band pattern, the light intensity of the band pattern (**A** and Movie S3) was transformed from the real-space to the wave-number domain (see Method S9). The Fourier power spectra at various time points are shown. A peak around the wave number of $\sim 2 \text{ cm}^{-1}$ appeared at time 800 min, getting narrower and taller over time, indicating an expanding band pattern with a constant wavelength of $\sim 0.5 \text{ cm}$.

Figure S5

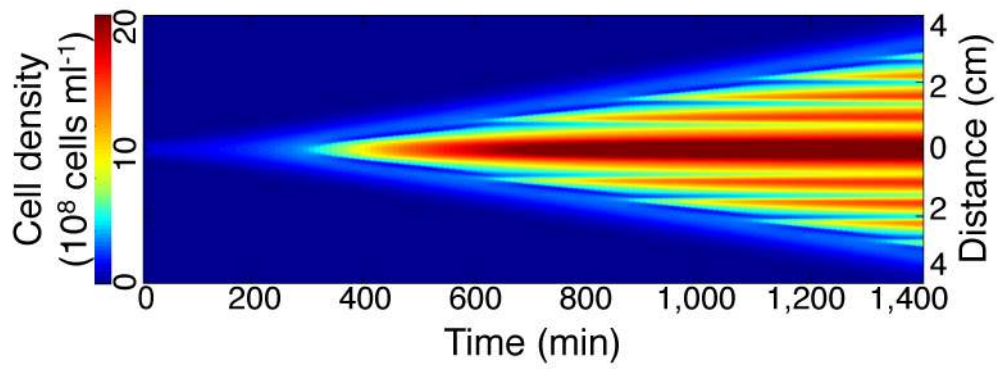


Fig. S5. Spatiotemporal diagram of cell density profile obtained from 1D simulation of Eqs. [1]-[3]. The simulation was done using the parameters in Table S3, starting with the initial condition as Eq. [S.15].

Figure S6

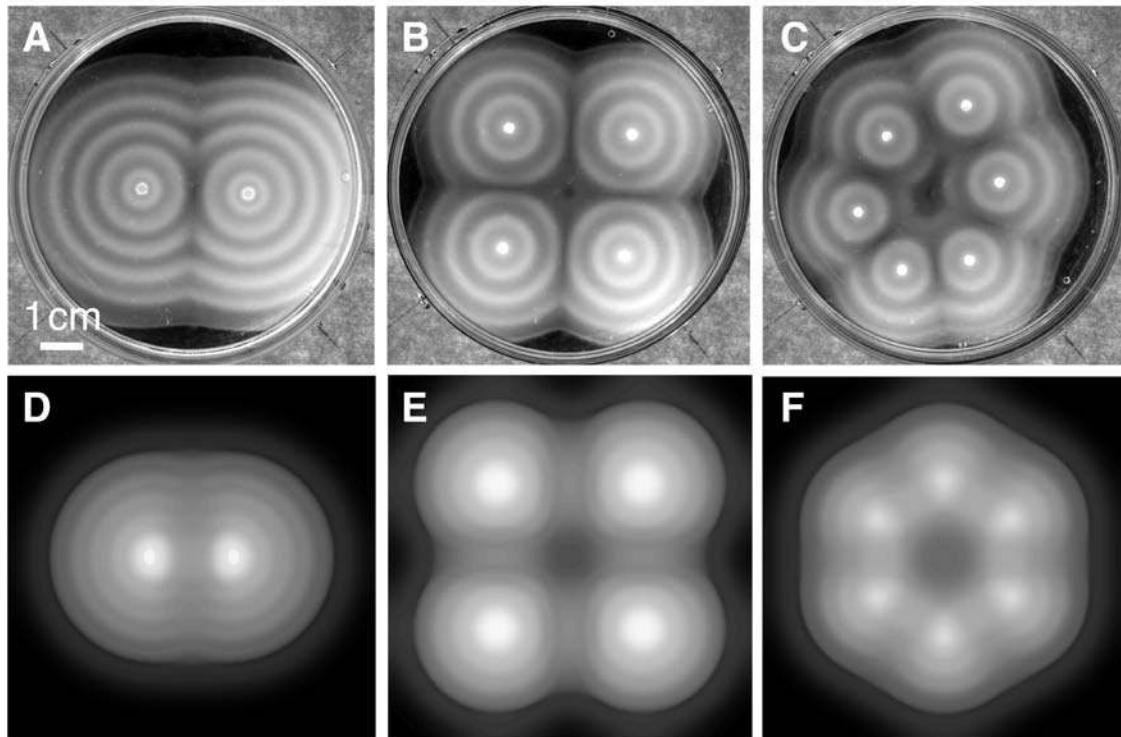


Fig. S6. Formation of intricate patterns. (A), (B), (C) Experimental results showing various patterns formed based on the placement of the inoculums. (D), (E), (F) Simulation results corresponding to (A), (B), and (C), respectively.

Figure S7

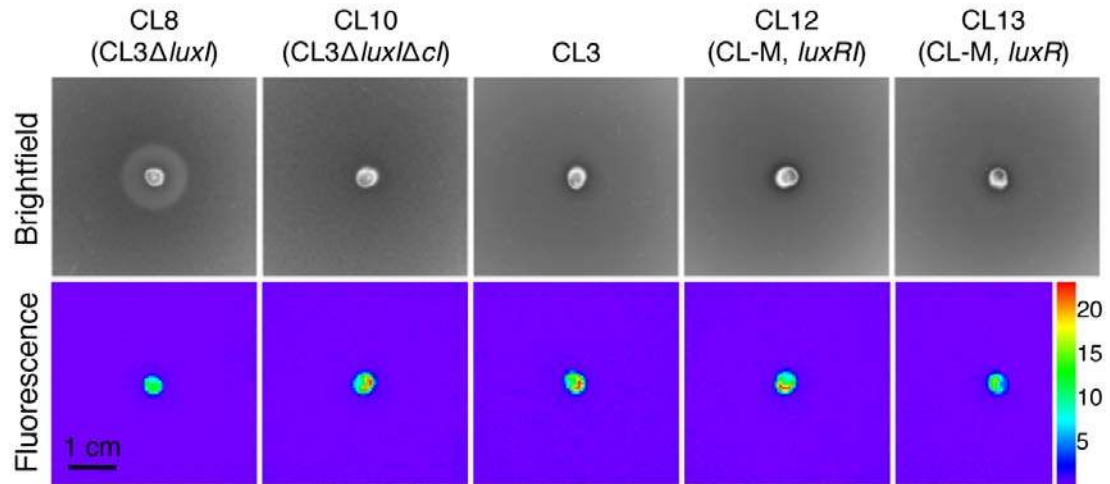


Fig. S7. Evidence of effective aggregation. To test the hypothesis of effective aggregation, various combinations of receiver and sender cells were conducted. All sender cells (CL6, non-motile but synthesizing AHL) were located at the position of inoculum as indicated by green fluorescence. Only those receiver cells with an entire genetic circuit but lacking the AHL-producing gene (*luxI*) formed a distinct cell aggregation around the spotted sender cells (CL8). Aggregation could not occur under other conditions, such as receiver cells lacking *cI*, which failed to regulate *cheZ* (CL10); receiver cells with AHL-producing ability, which might affect the formation of the AHL gradient (CL3); receiver cells with native *cheZ* and AHL-producing ability (CL12); receiver cells with native *cheZ* but incapable of AHL synthesis (CL13). Upper panels represent bright field images. Lower panels represent corresponding fluorescence images.

Figure S8

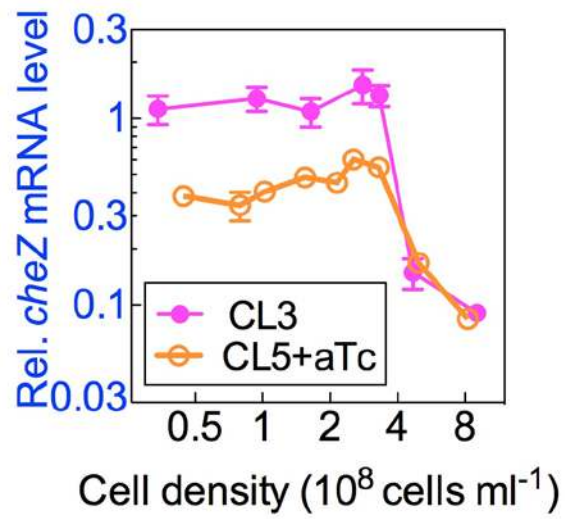


Fig. S8. Relative *cheZ* mRNA levels of CL5 as a function of cell density in bulk culture supplemented with 0.4 ng ml^{-1} aTc. CL3 was included as a control. Data are normalized by the mean value of CL3 at 0.3×10^8 cells ml^{-1} .

Figure S9

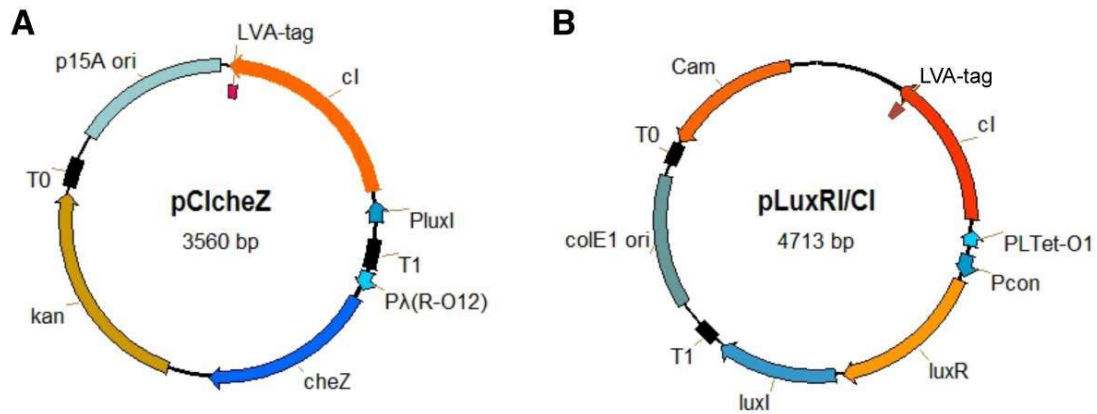


Fig. S9. Plasmids for the tunable stripe pattern formation. **(A)** pCIcheZ. *cheZ* and *cl* are under the control of CI repressible promoter $P_{\lambda(R-012)}$ and cell density-inducible promoter P_{luxI} , respectively. They are oriented back to back. **(B)** pLuxRI/CI. *luxR* and *luxI* are constitutively expressed (see Method S2 for details). The *cl* gene is driven by a TetR-repressed promoter, $P_{LTet-O1}$. The LVA-tag at the C-terminal of the CI coding region will facilitate the turnover of CI in the cells (62). T0 and T1 indicate transcriptional terminators.

Figure S10

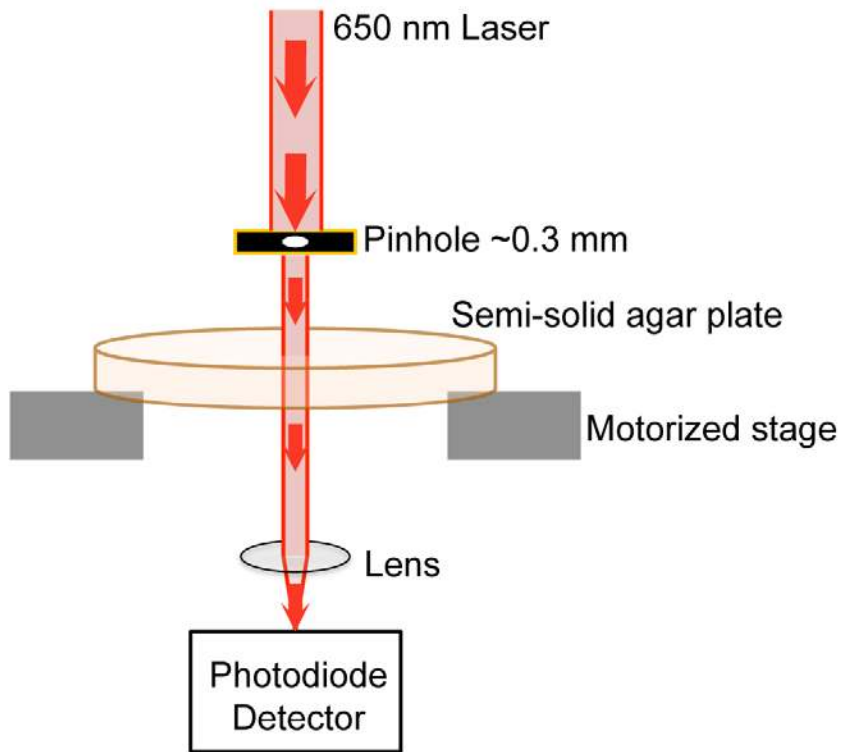


Fig. S10. The setup of the customized device for the real-time measurement of the spatiotemporal cell density profile in semi-solid agar dishes. A parallel 650-nm laser beam is guided through a 300- μm pinhole, passed through the sample in semi-solid agar on the Petri dish, and collected via a convex lens to a photodiode detector. The light intensity is digitized with a DAQ device and stored on a personal computer. The spatiotemporal scanning is realized with a motorized stage controlled with the PC. The whole apparatus was placed in a warm room (37 °C) throughout the experiment.

Figure S11

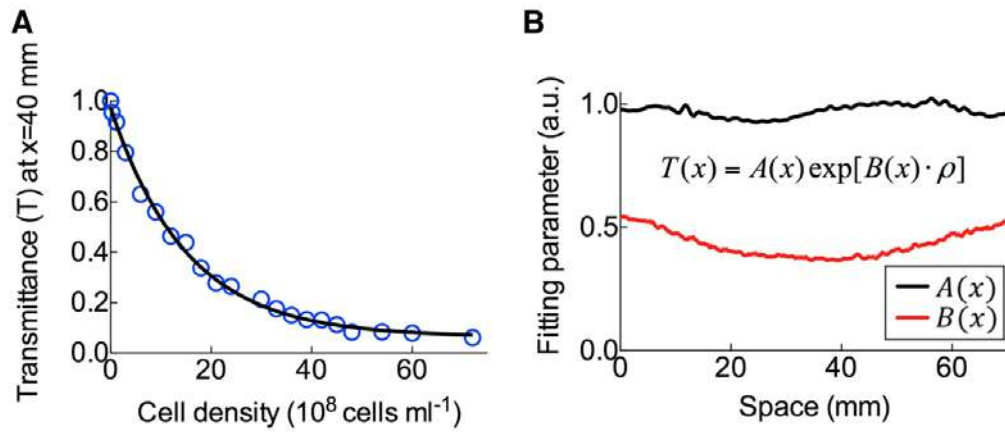


Fig. S11. The calibration curve of light transmittance vs. real cell density in semi-solid agar plates. **(A)** An example of the fitted curve of the transmittance as the function of cell density at a specific position. **(B)** The spatial distribution across the center of the dish, of the two fitting parameters for the standard curve, $A(x)$ (black line) and $B(x)$ (red line). This is derived from 376 standard curves (scanning range = 75 mm; step size = 0.2 mm).

Figure S12

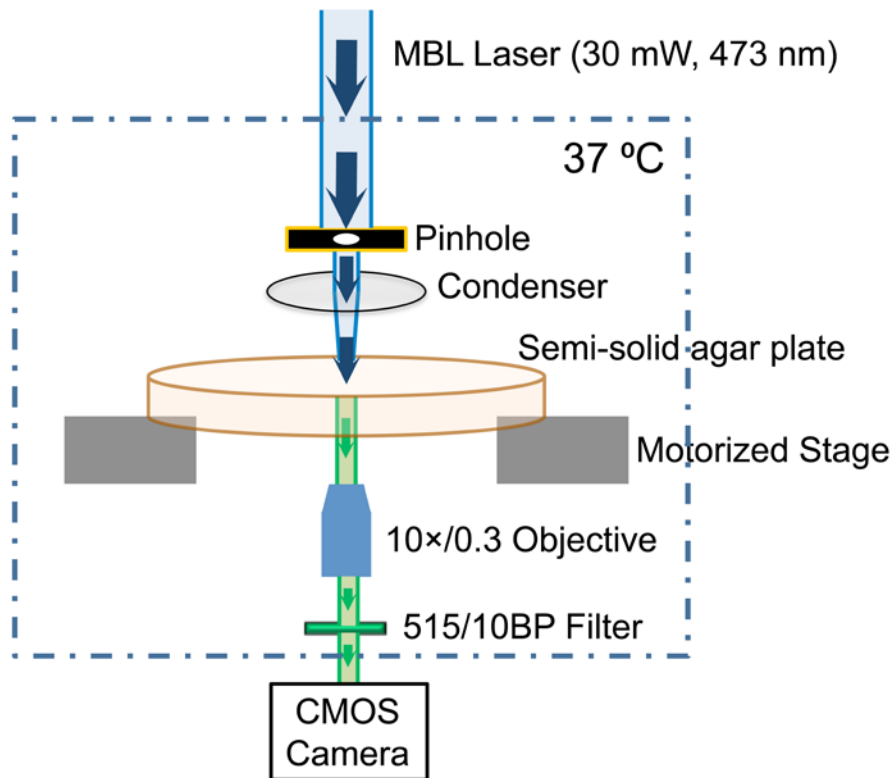


Fig. S12. A schematic diagram for the continuous fluorescence photobleaching (CPB) microscopy system. A parallel 473-nm laser beam is guided through a pinhole and a condenser to excite the GFP molecules in the cells in a cylindrical volume. The emitted fluorescence is collected via an objective lens, passed through a narrow band green filter, and detected with a CMOS camera. Throughout the measurement, the semi-solid agar plate was placed on a motorized stage within an environmental chamber, which is kept constant at 37 °C.

Figure S13

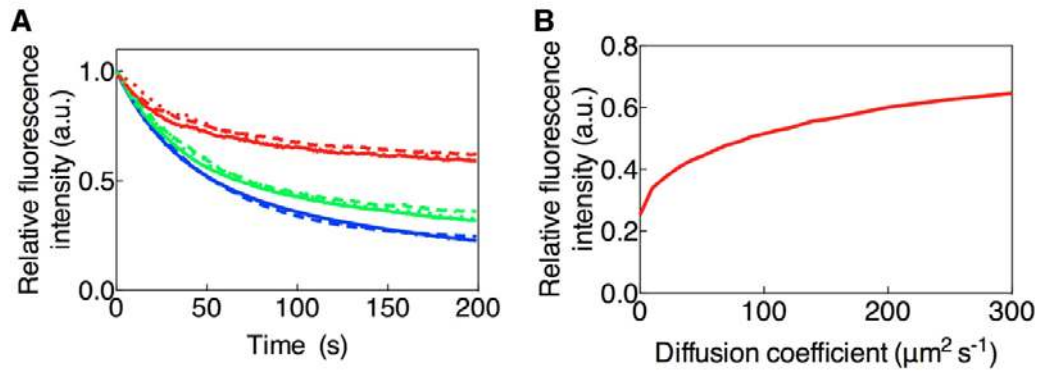


Fig. S13. Measurement of cell diffusion coefficient via continuous fluorescence photobleaching (CPB). **(A)** The representative CPB curves of different diffusion coefficients. Blue dots are experimental data of immobile strain CL14. Red and green dots are representative data of CL3GFP with fast and slow diffusion coefficients, respectively. Blue lines are fitting curves of blue dots ($D=0 \mu\text{m}^2 \text{s}^{-1}$) using double component decay function $f(t) = 0.48e^{0.03t} + 0.52e^{0.004t}$ (solid) and single decay function $f'(t) = 0.75e^{0.018t} + 0.25$ (dashed). Red and green lines are the simulation results of $D=110 \mu\text{m}^2 \text{s}^{-1}$ and $D=5 \mu\text{m}^2 \text{s}^{-1}$, respectively, with the corresponding blue lines as $D=0$ standard. **(B)** The relationship between diffusion coefficient and relative fluorescence intensity (mean value from 175 to 180 s) by using $f(t)$ as the $D=0$ standard.

Figure S14

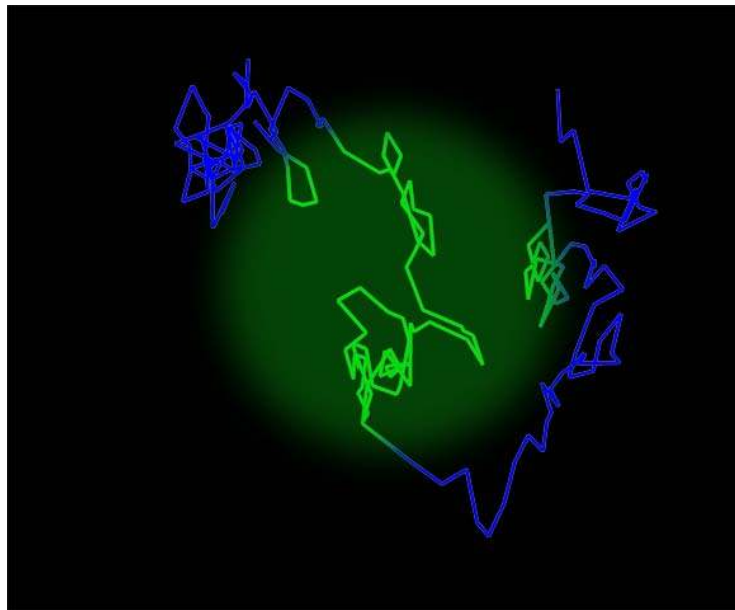


Fig. S14. Illustration of the trajectory of one cell in and out of the photobleaching region. The line represents a simulated 2D random walk trajectory of an *E. coli* cell. The GFP molecules would undergo photobleaching only when the cell locates in the laser-illuminated region (green). The accumulated photobleaching time (τ) was derived from the portions of the trajectory in the green region.

Figure S15

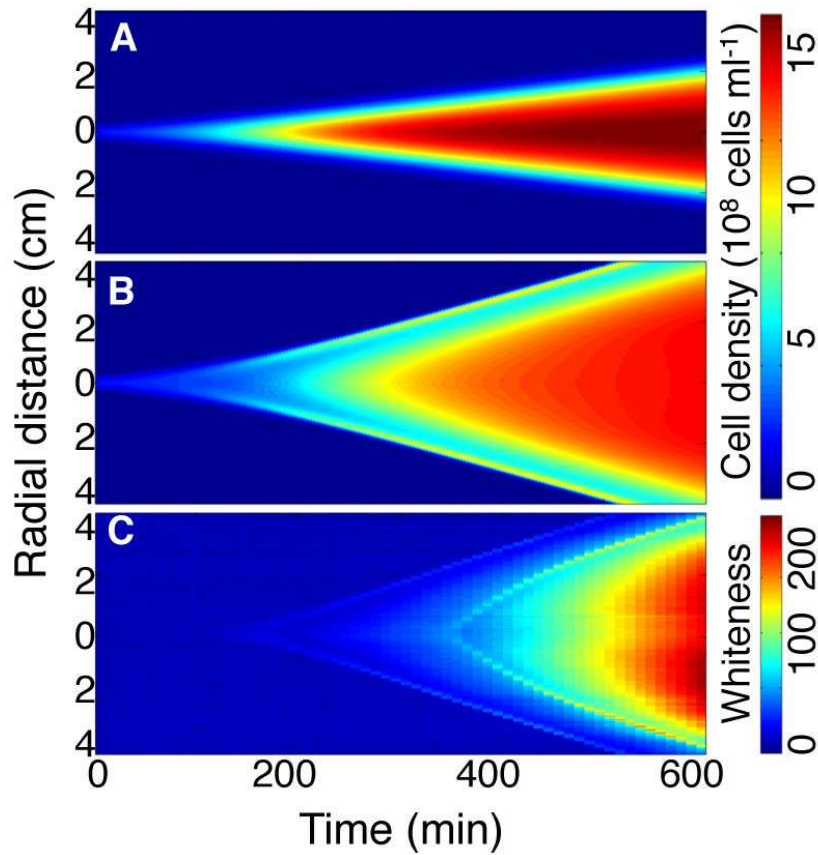


Fig. S15. The simulated spatiotemporal density profile in 1D. **(A)** without chemotaxis (Eqs.[S.5] -[S.7] with $D_\rho = 450 \mu\text{m}^2 \text{s}^{-1}$, $\gamma \sim 1.4\text{h}^{-1}$, $K_{C_1} = 0$, $K_n = 1 \times 10^9 \text{ cells ml}^{-1}$, $k_n = 1$, and $n(t=0) = 1.5 \times 10^9 \text{ cells ml}^{-1}$). **(B)** with chemotaxis ($K_{C_1} = 2 \times 10^9 \text{ cells ml}^{-1}$, $K_{C_2} = 1.5 \times 10^8 \text{ cells ml}^{-1}$, $K_a = 0.1 \times 10^9 \text{ cells ml}^{-1}$, $k_a = 1$ and $a(t=0) = 0.35 \times 10^9 \text{ cells ml}^{-1}$). The drift term changes the Fisher's wave front shape and increases its propagation speed. **(C)** The experimental spatiotemporal diagram of CL4 (harboring quorum-sensing module but with native *cheZ* regulation), showing the light intensity along a vertical line through the center of the Petri dish as a function of time.

Figure S16

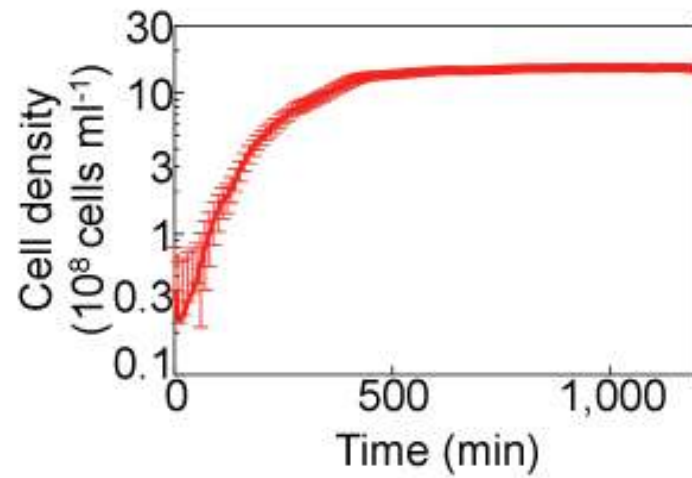


Fig. S16. Growth curve of CL3 in a Petri dish containing 10-ml 0.25% LB agar. After seed culture and preculture growth, CL3 cells were diluted 200-fold into the pre-warmed LB media containing 0.25% agar. Cell-agar mixtures were poured into Petri dishes and allowed to harden at room temperature for 90 min, and then moved back into a 37°C incubator. The cell densities in the semi-solid agar dish were automatically measured every 10 min (see Method S5).

Figure S17

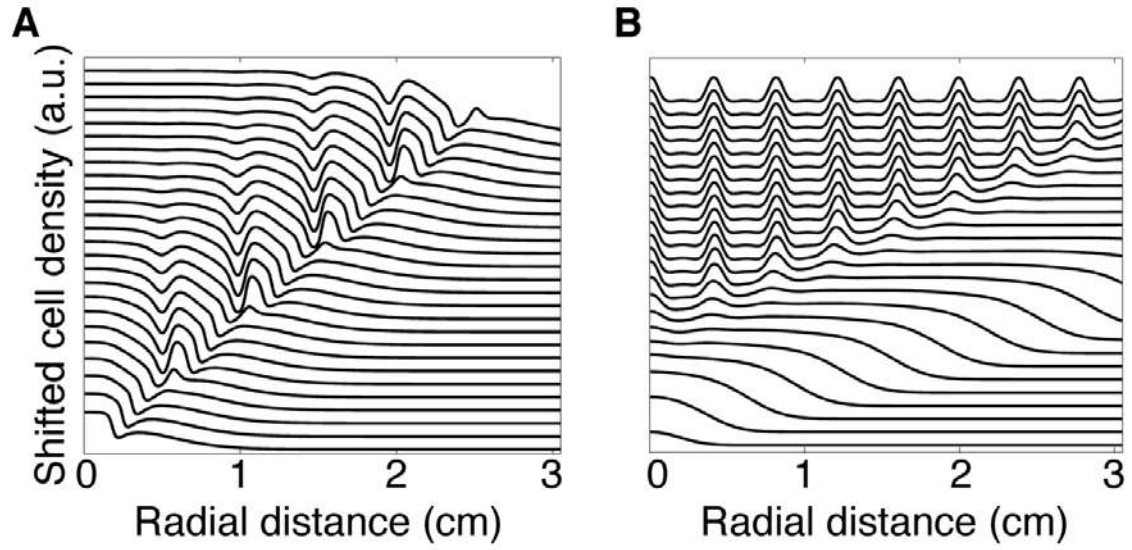


Fig. S17. The simulated cell density profile of Eqs. [S.11] and [S.12]. The y-axis represents the cell density, (shifted every 30 min). **(A)** simulated pattern with $D_\rho = 500 \mu\text{m}^2 \text{s}^{-1}$, $\gamma \sim 0.7 \text{ h}^{-1}$, $\rho_s = 1.5 \times 10^9 \text{ cells ml}^{-1}$, $\alpha = \beta : 2 \text{ h}^{-1}$, and $K_h = 4 \times 10^8 \text{ cells ml}^{-1}$. **(B)** simulated pattern with $D_\rho = 100 \mu\text{m}^2 \text{s}^{-1}$, and $K_h = 1.5 \times 10^9 \text{ cells ml}^{-1}$.

Figure S18

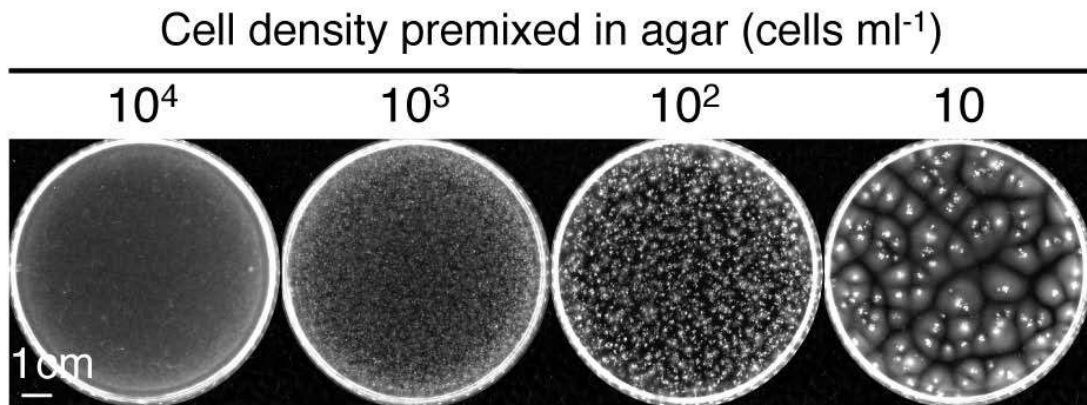


Fig. S18. Patterns formed by mixing agar with cells in various initial cell densities. Cells were cultured until the mid-exponential phase, diluted or concentrated to desired cell densities and mixed with pre-warmed 0.25% agar (see Method S3 for details), followed by 20-h incubation at 37 °C. Random patches of cells were formed with no periodic patterns. Scale bar, 1 cm.

Figure S19

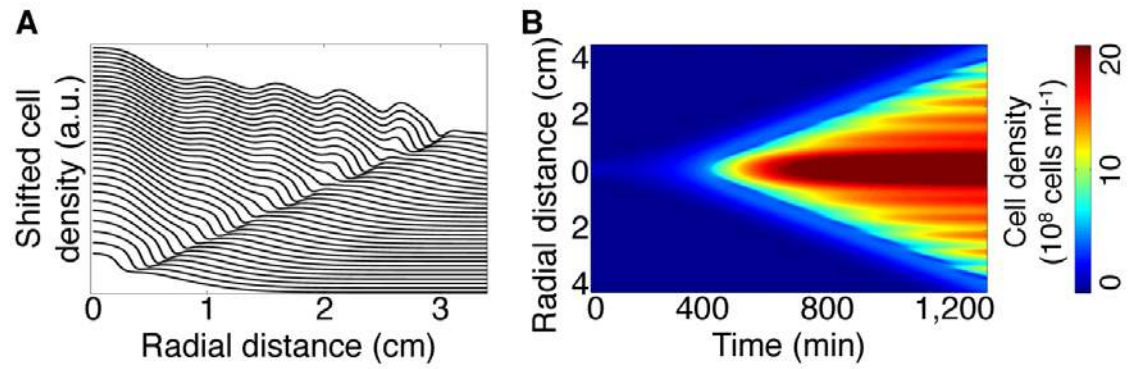


Fig. S19. The 2D simulation of Eqs.[S.12]-[S.14] with $D_\rho = 500 \mu\text{m}^2 \text{s}^{-1}$. **(A)** The y-axis represent the cell density profile (shifted every 30 min). **(B)** The simulated spatiotemporal density profile.

Figure S20

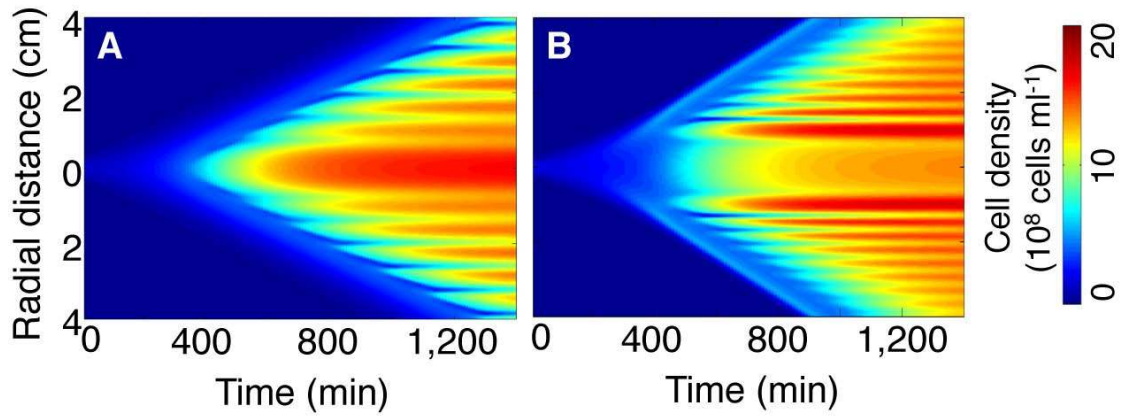


Fig. S20. The simulated spatiotemporal density profiles. **(A)** Eqs.[S.12], [S.14] and [S.16] with $D_\rho = 500 \mu\text{m}^2 \text{s}^{-1}$ and $\alpha = \beta : 2 \text{ h}^{-1}$. **(B)** Eqs. [S.6], [S.12] and [S.14] and [S.17] with $D_\rho = 450 \mu\text{m}^2 \text{s}^{-1}$, $\alpha = \beta : 1.4 \text{ h}^{-1}$, $K_{C_1} = 2 \times 10^9 \text{ cells ml}^{-1}$, $K_{C_2} = 1.5 \times 10^8 \text{ cells ml}^{-1}$, $K_a = 0.1 \times 10^9 \text{ cells ml}^{-1}$, $k_a = 1$ and $a(t=0) = 0.35 \times 10^9 \text{ cells ml}^{-1}$.

Figure S21

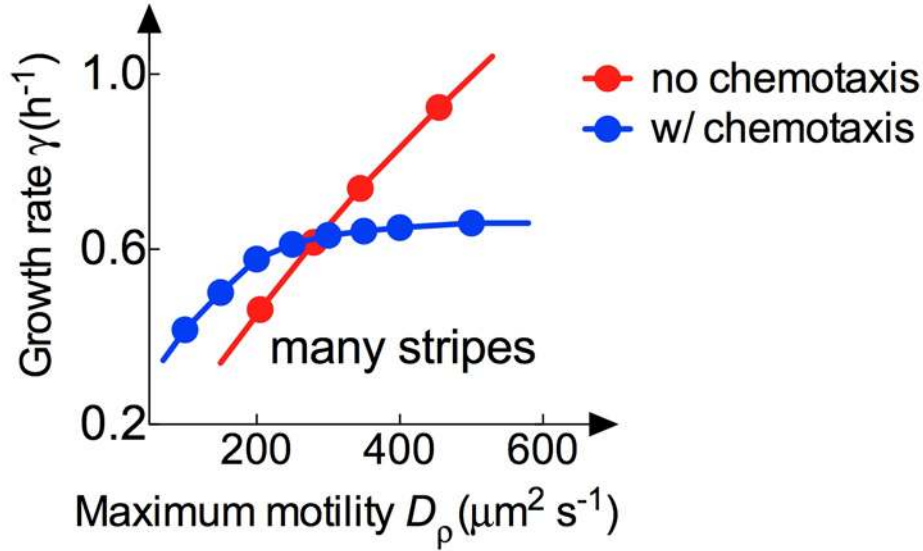


Fig. S21. The phase boundary of the stripe pattern formation. The behavior of this system is controlled mostly by the two dimensionless combinations of parameters, D_ρ/D_h and γ/β . Since the AHL diffusion coefficient and half-life is not easy to be modified experimentally, we fixed the D_h ($400 \mu\text{m}^2 \text{s}^{-1}$) and β (1.4h^{-1}) and simulate the phase boundary by varying the maximum cell motility D_ρ and the growth rate γ with the initial condition described with Eq.[S.18]. The red dots represent the phase boundary without the chemotaxis term (same as in Fig. 4A). The system can only form many stripes below this boundary. With the model including chemotaxis (Eqs. [S.6], [S.12], [S.14] and [S.17] with $K_{C_1} = 2 \times 10^9 \text{ cells ml}^{-1}$, $K_{C_2} = 1.5 \times 10^8 \text{ cells ml}^{-1}$, $K_a = 0.1 \times 10^9 \text{ cells ml}^{-1}$, $k_a = 1$ and $a(t=0) = 0.35 \times 10^9 \text{ cells ml}^{-1}$), the phase boundary shifted a little. The detailed position of this boundary can also be affected by the other parameters, such as K_{C_1} , K_{C_2} , K_a , etc.

SUPPLEMENTARY TABLES

Table S1 Strains and plasmids used in this study

Strains or plasmids	Genotype or description	Reference or source
Strains		
CL-M	A motile wild type <i>E. coli</i> K12 strain AMB 1655	Prof. Danchin (AMAbiotics, France)
CL-IM	A immotile <i>E. coli</i> K12 strain MG1655	Prof. Danchin (AMAbiotics, France)
CL1	$\Delta cheZ, \Delta lac$	CL-M, this study
CL1RB	$\Delta cheZ, \Delta lac, \Delta cheR, \Delta cheB$	CL1, this study
CL2	$\Delta cheZ, \Delta lac, bla:P_{Ltet-O1}-tetR$ at the <i>attB</i> site	CL1, this study
CL3	pCIcheZ, pLuxRI2	CL1, this study
CL3GFP	pCIcheZ, pLuxRI/GFP	CL1, this study
CL4	pLuxRI/GFP	CL-M, this study
CL5	pCIcheZ, pLuxRI/CI	CL2, this study
CL6	pCIcheZ $\Delta cheZ$, pLuxRI/GFP	CL1, this study
CL7	pCIcheZ ΔcI , pLuxRI2	CL1, this study
CL8	pCIcheZ, pLuxR	CL1, this study
CL9	pCIcheZ, pLuxI	CL1, this study
CL10	pCIcheZ ΔcI , pLuxR	CL1, this study
CL11	pCIcheZ $\Delta cheZ$, pLuxR/GFP	CL1, this study
CL12	pCIcheZ $\Delta cheZ$, pLuxRI2	CL-M, this study
CL13	pCIcheZ $\Delta cheZ$, pLuxR	CL-M, this study
CL14	pLuxRI/GFP	CL-IM, this study
CL15	pCIcheZ, pLuxRI2	CL1RB, this study
EQ44	$\Delta lacY, bla:P_{Ltet-O1}-tetR$ at the <i>attB</i> site	(35)
Plasmids		
pSim6	Ap ^r , pSC101 <i>ori</i> , λ Red	(34)
pEGFP-loxP-CmR-loxP	Ap ^r , Cm ^r , pUC <i>ori</i> , <i>loxP-Cm^r-loxP</i>	(33)
p705cre	Km ^r , pSC101 <i>ori</i> , <i>cre</i>	(63)
pLuxRI2	Cm ^r , ColE1 <i>ori</i> , Pcon- <i>luxR-luxI</i>	(36)
pLD	Km ^r , p15A <i>ori</i> , $P_{\lambda(0-12)}$ - <i>lacI</i> , P_{luxI} - <i>cI</i>	(9)
pCIcheZ	Km ^r , p15A <i>ori</i> , $P_{\lambda(0-12)}$ - <i>cheZ</i> , P_{luxI} - <i>cI</i>	This study
pLuxRI/CI	pLuxRI2 carrying P_{tet} - <i>cI</i>	This study

pLuxRI/GFP	pLuxRI2 carrying Pcon-superfolder gfp	This study
pLuxR/GFP	pLuxR carrying Pcon-superfolder gfp	This study
pCIcheZΔcheZ	pCIcheZ lacking <i>cheZ</i> -coding region	This study
pCIcheZΔcI	pCIcheZ lacking <i>cI</i> -coding region	This study
pLuxR	pLuxRI2 lacking <i>luxI</i> -coding region	This study
pLuxI	pLuxRI2 lacking <i>luxR</i> -coding region	This study
pCIcheZ16S	pCIcheZ carrying a copy of 16S rRNA gene	This study

Table S2 Oligonucleotides used in this study

Name	Sequence	Use
cheZ-del-f	ggaaaaactcaacaaaactttgagaaactgggcatgtga ggatgcgactctcgagccgatcatattcaataa	<i>cheZ</i> deletion
cheZ-del-r	ttatcagaccgctgatatgacgtgggtcacgccacatcag gcaatacaaaagacctcgaggactagtgaacctc	<i>cheZ</i> deletion
cheZ-del-conf-f	aactgggcatgtgaggat	Confirmation
cheZ-del-conf-r	gccacatcaggcaataca	Confirmation
lacIZYA-del-f	gcggtatggcatgatagcggcgggaagagagtcaattca gggtgggtaactctcgagccgatcatattcaataac	<i>lac</i> operon deletion
lacIZYA-del-r	ttcatgccggatcgcgctaattgatagctgctgaactgtag gcctgataagacctcgaggactagtgaacctc	<i>lac</i> operon deletion
lac-del-conf-f	ggcagaatgcttaataa	Confirmation
lac-del-conf-r	ttctgcgctttgctcat	Confirmation
cheRB-bla-f	tgcagttacaattgcgccagtggtatcctgaagtattga gaaggcgtatgagtattcaacattccgtgt	<i>cheRcheB</i> deletion
cheRB-bla-r	atccacaacaaaaatttaagtctttatccgccatttcaca ctcctgattaccaatgcttaacagtgagg	<i>cheRcheB</i> deletion
bla-conf-f	aatagtgtatcgggcgacc	Confirmation
cheRB-conf-r	tgtggcggtatttacctt	Confirmation
cheZ-f	atgatgcaacctcaatcaaa	Cloning <i>cheZ</i> into pLD
cheZ-r-Nhe	ctagctagctcaaatccaagactatccaac	Cloning <i>cheZ</i> into pLD
pLD-rot-f-xba	gctctagagttttccataggctccgc	Linearizing pLD <i>cheZ</i> flanking P_{luxI} - <i>cI</i>
pLD-rot-r-cla	ccatcgatgagtaggacaatccgccg	Linearizing pLD <i>cheZ</i> flanking P_{luxI} - <i>cI</i>
cI-f-cla	ccatcgatgctttacgttcacactattgtatcg	Cloning P_{luxI} - <i>cI</i> into linearized pLD <i>cheZ</i>
cI-r-xba	gctctagaagaacacctgcccgatcattaa	Cloning P_{luxI} - <i>cI</i> into linearized pLD <i>cheZ</i>
cI-f-sca	aaaagtactgagaatttcacacaggggtac	Cloning <i>cI</i> into pCRT-1
cI-r-not	aaagcgccgctagaagaacacctgcccga	Cloning <i>cI</i> into pCRT-1
gfp-f-sac	<u>cgagctctgacggctagctcagctctaggtacagtgcta</u> <u>gcttgggctagctactagagaagag</u>	Cloning superfolder-gfp into pLuxRI2, the synthetic promoter is underlined
gfp-r-sac	cgagctcttatcatcattgtacagttcatcc	Cloning superfolder-gfp into pLuxRI2
gfp-conf-r	acgcacggaaaacttatgac	confirmation
delcI-f-cla	acaatagtgtgaacgtaaacatc	Deleting <i>cI</i> from pCIcheZ
delcI-r-cla	gcatcgatcggaaggtgttctctag	Deleting <i>cI</i> from pCIcheZ

16S-f-cla	ccatcgatattgaagagttgatcatggc	Cloning 16S rRNA into pCIcheZ
16S-r-xho	ccgctcgagaggaggtgatccaaccg	Cloning 16S rRNA into pCIcheZ
RT-cheZ-f	ctggaaatcatgatggcg	RT qPCR primer for <i>cheZ</i>
RT-cheZ-r	tgctggtatcgacctgagg	RT qPCR primer for <i>cheZ</i>
RT-cI-f	ggttagtatgcagccgtcaact	RT qPCR primer for <i>cI</i>
RT-cI-r	caacctcaagccagaatgc	RT qPCR primer for <i>cI</i>
RT-rrn-r	gctacaatggcgcatacaaa	RT qPCR primer for 16S rRNA (41)
RT-rrn-f	ttcatggagtcgagttgcag	RT qPCR primer for 16S rRNA (41)

Table S3 The parameters most used in the simulation

Parameters	Value	Comments
D_ρ	$450 \mu\text{m}^2 \text{s}^{-1}$	The normal value (200-1,000 $\mu\text{m}^2 \text{s}^{-1}$) (49-51)
$D_{\rho,0}$	$10 \mu\text{m}^2 \text{s}^{-1}$	Fig. 1D shows it is almost zero
D_h	$400 \mu\text{m}^2 \text{s}^{-1}$	As small molecular diffusion (100-1,000 $\mu\text{m}^2 \text{s}^{-1}$) (56, 57)
D_n	$800 \mu\text{m}^2 \text{s}^{-1}$	As small molecular diffusion (100-1,000 $\mu\text{m}^2 \text{s}^{-1}$) (56, 57)
γ	0.7h^{-1}	Measured shown in Fig. S16
β	1.04h^{-1}	AHL half-life ranging 10~1,000 min (39, 59)
m	20	Fig. 1D shows an abrupt fall of the cell motility
$n (t=0)$	$15 \times 10^8 \text{ cells ml}^{-1}$	The saturated density cell as shown in Fig. S16
k_n	1	Rescaled with n to the unit of ρ
K_n	$10^9 \text{ cells ml}^{-1}$	Estimated from Fig. S16
K_h	$4 \times 10^8 \text{ cells ml}^{-1}$	Estimated from Fig. 1D