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Hyperpolarized Long- T_1 Silicon Nanoparticles for Magnetic Resonance Imaging

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Silicon nanoparticles are experimentally investigated as a potential hyperpolarized, targetable MRI imaging agent. Nuclear T_1 times at room temperature for a variety of Si nanoparticles are found to be remarkably long $(10^2 \text{ to } 10^4 \text{ s})$ —roughly consistent with predictions of a core-shell diffusion model—allowing them to be transported, administered and imaged on practical time scales without significant loss of polarization. We also report surface functionalization of Si nanoparticles, comparable to approaches used in other biologically targeted nanoparticle systems.

The use of nanoparticles for biomedical applications has benefited from rapid progress in nanoscale synthesis of materials with specific optical and magnetic properties, as well as biofunctionalization of surfaces, allowing targeting^{1,2,3}, *in-vivo* tracking^{3,4,5}, and therapeutic action^{6,7}. For magnetic resonance imaging (MRI), superparamagnetic nanoparticles have extended susceptibilitybased contrast agents toward targeted imaging⁸, though achieving high spatial resolution with high contrast remains challenging. An alternative approach is direct MRI of hyperpolarized materials with little or no background signal. Hyperpolarized noble gases^{9,10,11} and ¹³C-enhanced biomolecules^{12,13} have demonstrated impressive image contrast, but are limited by short *in-vivo* enhancement times (~ 10 s for noble gases⁹, ~ 30 s for ¹³C biomolecules^{12,13}). It is known that bulk silicon can exhibit multi-hour nuclear spin relaxation (T_1) times at room temperature¹⁴ and that silicon nanoparticles can be hyperpolarized via dynamic nuclear polarization $(DNP)^{15}$.

Nuclear magnetic resonance (NMR) in silicon has been widely investigated for half a century¹⁴, and with renewed interest in the context of quantum computation 16 . The low natural abundance of spin-1/2²⁹Si nuclei (4.7%) embedded in a lattice of zero-spin²⁸Si nuclei isolates the active nuclear spins from one another and from the environment, leading to multi-hour spin relaxation (T_1) times, and decoherence (T_2) times of tens of seconds¹⁶. Moreover, the weak dipole-dipole coupling of the sparse ²⁹Si atoms, together with the isotropic crystal structure and the absence of nuclear electric quadrupole moment conspire to keep any induced nuclear polarization aligned with an external magnetic field, even as the nanoparticle tumbles in space. This is critical for tracking hyperpolarized nanoparticles in a fluid suspension using MRI, as the nuclear polarization direction can be fixed using a small (mT-scale) applied field.

Particle size determines regimes of application to

biomedicine¹⁷ as well as predicted NMR properties¹⁵. Here, we report NMR properties of Si particles spanning four orders of magnitude in mean diameter, from ~ 40 nm nanoparticles to millimeter-scale granules. We investigate particles made by ball milling high-resistivity $(30-100 \text{ k}\Omega\text{-cm}, \text{ residual p-type } (111), \text{ Silicon Quest})$ International), and low-resisitivity $(0.01 - 0.02 \ \Omega\text{-cm})$. boron-doped (p-type), (100) oriented, Virginia Semiconductor) commercial silicon wafers, followed by centrifugal segregation by size¹⁸. We also investigate chemically synthesized Si nanoparticles with mean diameters 40 nm (Meliorum), 60 nm (Meliorum), 140 nm (MTI) and 600 nm (NanoAmor), obtained commercially. Figure 1 shows representative scanning electron microscope (SEM) images of all measured particles, along with volume-weighted size distributions obtained by SEM image analysis. Dilute suspensions of silicon nanoparticles in ethanol were sonicated for ten minutes before being pipetted onto a vitreous carbon planchett which was mounted on a standard specimen holder with conducting carbon tape. For each sample, > 1000 particles were analvzed, sourced from ~ 50 images, with particles in contact excluded from the analysis. Particle agglomeration seen in dry Meliorum and MTI samples has been reported in similarly sized silica nanoparticles, but is significantly reduced after pegylation¹⁹. In these cases (Meliorum, MTI), individual measurement of the particle diameter from SEM images was used instead of software analysis.

Nuclear T_1 times of the Si nanoparticles, segregated by size and packed dry in teflon NMR tubes, were measured at room temperature at a magnetic field of 2.9 T using a spin-echo Fourier transform method with a saturation recovery sequence. Following a train of sixteen hard $\pi/2$ pulses to null any initial polarization, the sample was left at field to polarize for a time τ_{pol} , followed by a CPMG sequence $(\pi/2)_X - [\tau - (\pi)_Y - \tau - \text{echo}]^n$ with $\tau = 1$ ms and n = 200. In Si and other nuclear-dipolecoupled materials echo sequences can yield anomalously

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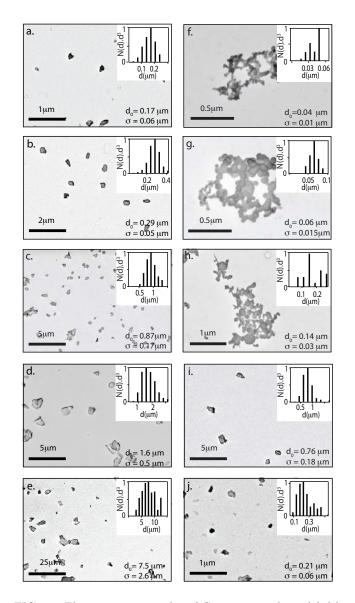


FIG. 1: Electron micrographs of Si nanoparticles. (a)-(e) ball milling high-resistivity silicon wafer, (f)-(g) wet synthesis (Meliorum), (h) plasma synthesis (MTI), (h) electrical explosion (NanoAmor), (j) ball milling low-resistivity wafer. Insets: Volume-weighted histograms of diameters following size segregation along with averages d_0 and standard deviations σ based on gaussian fits to distributions.

long decay tails²⁰. However, the Fourier amplitude of the echo train still provides a signal proportional to initial polarization²⁰. Values for T_1 are extracted from exponential fits, $A \propto 1 - e^{-\tau_{\rm pol}/T_1}$, to the Fourier amplitude, A, of the n = 200 echoes as a function of polarization time (see Fig. 2a, inset for an example).

Figure 2a shows T_1 as a function of (volume-weighted) average particle diameter for the various samples. The high-resistivity ball-milled samples follow a roughly linear dependence on size, $T_1 \propto d_0$, for $d_0 < \sim 10 \mu$ m, saturating at $T_1 \sim 5$ h for larger particles. The trend of in-

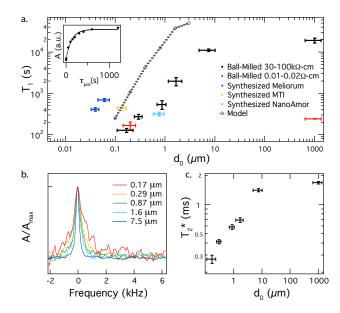


FIG. 2: (a) Nuclear spin relaxation (T_1) times at 2.9 T as a function of particle diameter d_0 for various Si particles. Vertical error bars are from exponential fits to relaxation data; horizontal error bars are σ of size distributions (see Fig. 1). Also shown is a core-shell diffusion model of T_1 dependence on particle diameter, from Ref. 15 (dashed curve). Inset: Fouriertransform NMR peak amplitude, A, as a function of polarization time τ_{pol} (see text) for the ball-milled high-resistivity particles with $d_0 = 170$ nm. T_1 values were measured using a saturation recovery spin echo pulse sequence described in the text. (b) Amplitude normalized NMR lineshape (A/A_{max}) versus frequency for ball-milled high resistivity samples at 4.7 T. (c) Inhomogeneous dephasing time, T_2^* , as a function of mean particle diameter for ball-milled high resistivity samples at 4.7 T.

creasing T_1 in larger particles is qualitatively consistent with a simple shell-core spin diffusion model¹⁵(dotted line), which predicts $T_1 \propto d_0^2$ for particles with no in-ternal defects or dopants. The low-resistivity ball-milled particles have $T_1 \sim 200$ s, independent of size. Smaller commercial particles formed by wet synthesis (Meliorum) and plasma synthesis (MTI) have T_1 times as long as 700 s. This is significantly longer than the predicted values, as this model assumes instantaneous spin relaxation at the surface of the particle. Larger commercial particles formed by electrical explosion (NanoAmor) have shorter T_1 than the comparably sized high-resistivity ball-milled particles. Powder x-ray diffraction measurements (not shown) indicate that ball milling induces partial polycrystallinity, consistent with previous studies²¹. We speculate that T_1 is reduced for the smallest ball-milled particles, compared to the synthesized particles, by coupling of nuclear spins to paramagnetic defects at the interface between crystallites. Electron Spin Resonance (ESR) measurements (see Supplementary Material S1) on ballmilled and synthesized particles show a single peak corresponding to a g-factor of g=2.006, characteristic of P_{b} -

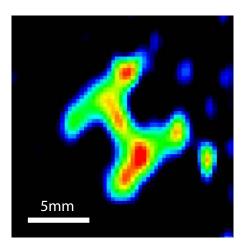


FIG. 3: Magnetic resonance imaging of hyperpolarized Si nanoparticles. An H-shaped cavity filled with high-resistivity Si particles ($d_0 = 1.6 \ \mu m$) pre-polarized at low temperature (T = 4.2 K) and high magnetic field (B = 5 T) for 60 h and warmed and transferred to a 4.7 T imager. See text for imaging details. No Si image could be obtained without hyperpolarization.

type defect centers²². For a fixed volume of particles, a larger ESR signal is seen in the smaller samples, suggesting that the surface, rather than the bulk, is the source of paramagnetic defects. A detailed study relating ESR and NMR properties will be reported later.

In addition to noting the remarkably long T_1 times for all Si measured particles relative to previously reported hyperpolarized MRI imaging agents^{9,11,12,13}, we also note that T_1 in the Si system can be tuned by size and doping, allowing optimization for specific applications, from virtual colonoscopy to molecular imaging.

Averaged NMR spectra for ball-milled high-resistivity samples of various particle sizes are shown in Fig. 2b. Each spectrum is taken from a series of summed free induction decay traces following polarization for a time $3T_1$ at a field of 4.7 T corresponding to 39.7 MHz, taken on a Bruker DMX-200 NMR console. In post processing, these spectra have been individually phase adjusted. Whereas T_1 changes by two orders of magnitude over the range of measured particle sizes, T_2^* changes only by factor of ~6 over the same range. It is expected that the rapid tumbling of particles in a liquid suspension will increase T_2^* considerably, restoring resolution. For biologically targeted imaging, the bound particles may not tumble, but may move enough for spectral narrowing to take place. This will be addressed in future work.

To demonstrate MRI of hyperpolarized Si nanoparticles, we filled a small cavity (in the shape of the letter H) with high-resistivity ball-milled particles ($d_0 = 1.6\mu$ m). The sample was left to polarize at lowtemperature (4.2 K) and high field (5 T) for 60 h, which enhanced the ²⁹Si nuclear spin polarization by approximately a factor of 15 over room temperature. We then removed the sample from the polarizing cryostat and transferred it to the 4.7 T Bruker DMX-200 imager with a micro-imaging gradient set, requiring ~ 1 minute for the transfer, and imaged the phantom using a small tip angle gradient echo sequence²³. Imaging parameters were: tip angle $\theta = 9^{\circ}$, echo time $\tau = 1.2$ ms, field of view = 15 mm, sample thickness = 2.5 cm, single pass (no averaging), acquisition time = 11 s. The resulting image is shown in Fig. 3. We note that much higher image resolution will be possible with DNP polarization^{9,10,11,12,13}.

To examine the applicability of Si nanoparticles to targeted MRI, we prepared the Si nanoparticle surface for attachment to biological-targeting ligands. Ballmilled high resistivity nanoparticles $(d_0 = 200 \text{ nm})$ were aminated using either (3-Aminopropyl)triethoxysilane (APTES, Sigma, 99%) or a 1:2 mixture by volume of APTES with bis-(triethoxysilvl)ethane (BTEOSE, Aldrich, 96%) or (3-trihydroxysilyl)propyl methylphosphonate (THPMP, Aldrich, 42 wt% in H₂O) (see Fig. $(4 a)^{24}$. The surface oxide was first etched with a dilute solution of hydrofluoric acid (8% in ethanol) followed by resuspension of the particles in ethanol. Approximately 100 mg of silicon nanoparticles were added to 45 mL of acidified 70% ethanol (0.04% v/v, adjusted to pH 3.5 with HCl) and the solution was placed in an ultrasonic bath for five minutes. Saline (0.15 M) was then added and the solution was shaken for 18-24 hours. Silanes were removed from the nanoparticle solution by washing and resuspending three times in methanol buffer, with the final resuspension performed with 10 mL of methanol buffer. Successful amination was assessed using fluorescence spectroscopy (Fig. 4b) using an excitation at 390 nm and emission at 465 nm (SpectraMax Plus, Molecular Devices). The high level of fluorescence observed for aminated particles results from the covalent bonding of surface amino groups with fluorescamine, showing these functional groups were accessible for further reaction.

In addition to chemical assays, the accumulation of amines was indirectly monitored by measuring the particles' surface charge in solution, known as the zeta potential²⁵ (Fig. 4c). The surface of the unmodified silicon nanoparticles is composed of hydroxyl groups from the silicon dioxide and thus shows a negative zeta potential. Particles treated with APTES have surfaces coated with propylamines, which become protonated and positively charged in acidic solutions and show a positive zeta potential²⁵.

Aminated particles were coated with polyethylene glycol (PEG) polymers to confer stability and biocompatibility. PEG coating of silica and iron-oxide nanoparticles has been shown to be non-toxic²⁶ and to reduce the rate of clearance by organs such as the liver or kidneys, thus increasing the particle's circulation time *in-vivo*²⁶. Pegylation was performed with either α -methyl-PEGsuccinimidyl α -methylbutanoate (mPEG-SMB) (Nektar)

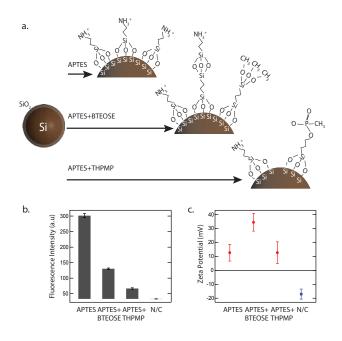


FIG. 4: Biological surface modification of silicon nanoparticles. (a) Silicon particles ($d_0 = 0.2 \ \mu m$) were aminated using either (3-Aminopropyl)triethoxysilane (APTES) alone or as a 1:2 mixture by volume of APTES with bis-(triethoxysilyl)ethane (BTEOSE) or (3-trihydroxysilyl)propyl methylphosphonate (THPMP in H₂O). (b) Fluorescence spectroscopy confirmed the success of the amination reaction. No fluorescence was evident with the negative control (N/C). (c) A change in the sign of the surface charge, or zeta potential of the particles was evident after amination with the three amine groups when compared to the negative control.

or maleimide-PEG-N-hydroxysuccinimide (MAL-PEG-NHS) (Nektar). Both SMB and NHS are reactive with amines on the particle surface. 10 mg of PEG was mixed in 500 μ L of methanol buffer and heated briefly at 50 °C to dissolve. Approximately 0.1 mg of aminated particles

(100 μ L in solution) were added to this solution and it was placed in an ultrasonic bath for 13 h. To remove the unreacted PEG, samples were centrifuged and resuspended twice in methanol and finally in a phosphatebuffered saline solution (PBS, 0.1 M Na₂HPO₄, 0.015 M NaCl buffer).

The stability of nanoparticles in solution was assessed using both dynamic light scattering (DLS) (Nano ZS90, Malvern) as a measure of the particles' hydrodynamic radius, and visual determination of flocculation and sedimentation. The particles treated with mPEG-SMB and NHS-PEG-MAL were both stable in phosphate-buffered saline (PBS) for a period of two days, with no significant change in the particles' hydrodynamic radius (see Supplementary Information S2). As a control, mPEG-Amine polymer, which does not contain amine-reactive groups, was used. The aminated particles treated with mPEG-Amine aggregated after centrifugation and resuspension in PBS. These results are consistent with other reports of the successful pegylation of SiO₂ nanoparticles^{19,27}.

In conclusion, we demonstrate that Si nanoparticles show promise as biologically targeted MRI imaging agents based on their exceptional NMR properties, including their receptivity to hyperpolarization and long nuclear relaxation (T_1) times, in the range of minutes to hours. We investigated Si T_1 times as a function of nanoparticle size, dopant concentration and synthesis method. Furthermore, we have demonstrated techniques for fabricating, size-separating and coating Si nanoparticles to satisfy a broad spectrum of design criteria. Future developments in the chemical synthesis of larger, monodisperse single-crystal silicon nanoparticles may provide even longer T_1 times. We note that Si nanoparticles may be combined with other material components to provide MRI tracking of the delivery of drugs²⁸ or as a therapeutic agent that allows simultaneous MRI tracking. The addition of APTES and PEG to the surface of these nanoparticles is a critical step for further surface functionalization and, ultimately, biological targeting.

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