

57. Singular Hadamard's Variation of Domains and Eigenvalues of the Laplacian. II

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(Communicated by Kôzaku YOSIDA, M. J. A., May 12, 1981)

§ 1. This paper is a continuation of our previous note [2]. Let Ω be a bounded domain in \mathbf{R}^n with C^3 boundary γ and w be a fixed point in Ω . For any sufficiently small $\varepsilon > 0$, let B_ε be the ball defined by $B_\varepsilon = \{z \in \Omega; |z - w| < \varepsilon\}$. Let Ω_ε be the bounded domain defined by $\Omega_\varepsilon = \Omega \setminus \bar{B}_\varepsilon$. Then the boundary of Ω_ε consists of γ and ∂B_ε . Let $0 > \mu_1(\varepsilon) > \mu_2(\varepsilon) > \dots$ be the eigenvalues of the Laplacian with the Dirichlet condition on $\partial\Omega_\varepsilon$. We arrange them repeatedly according to their multiplicities. In [2], [3] we gave the asymptotic formulas for the j -th eigenvalue $\mu_j(\varepsilon)$ when $\varepsilon \searrow 0$ in case $n=2, 3$. In this note we treat the case $n=4$. We have the following

Theorem 1. *Assume $n=4$. Fix j . Suppose that the j -th eigenvalue μ_j of the Laplacian with the Dirichlet condition on γ is a simple eigenvalue, then*

$$(1.1) \quad \mu_j(\varepsilon) - \mu_j = -2S_4\varepsilon^2\varphi_j(w)^2 + O(\varepsilon^{5/2})$$

holds when ε tends to zero. Here φ_j denotes the eigenfunction of the Laplacian with the Dirichlet condition on γ satisfying

$$\int_{\Omega} \varphi_j(x)^2 dx = 1.$$

Here S_4 denotes the area of the unit sphere in \mathbf{R}^4 .

Our aim of this note is to offer a rough sketch of the proof of the above theorem. Calculation and technique which are used to prove Theorem 1 are more elaborate than in case $n=2$ and 3. L^p ($1 < p < \infty$) spaces are used in this note. We employed only L^2 spaces in case $n=2, 3$.

We review a generalization of the Schiffer-Spencer formula. See [6]. Also see [3]. In the following we assume $n=4$. Let $G(x, y)$ be the Green's function on Ω . Put

$$\omega_\varepsilon = \{x \in \Omega; G(x, w) < (2S_4\varepsilon^2)^{-1}\}$$

and $\beta_\varepsilon = \Omega \setminus \bar{\omega}_\varepsilon$. Let $G_\varepsilon(x, y)$ be the Green's function in ω_ε .

Variational formula for the Green's function [3]. Fix $x, y \in \Omega \setminus \{w\}$ such that $x \neq y$. Then

$$(1.2) \quad G_\varepsilon(x, y) = G(x, y) - 2S_4\varepsilon^2 G(x, w)G(y, w) + O(\varepsilon^3)$$

holds when ε tends to zero. The remainder term is not uniform with respect to x, y .

To prove Theorem 1 we use the iterated kernel $G_\varepsilon^{(2)}$ (resp. $G^{(2)}$) of

$G_\varepsilon(x, y)$ (resp. $G(x, y)$) and a variational formula for $G_\varepsilon^{(2)}(x, y)$. See § 2. It should be remarked that in case $n=2, 3$ only G_ε, G were used. See [3].

There are many related papers and topics. For example, see [1], [2], [4], [5] and [7]. Details and a further generalization of this note will be given elsewhere.

§ 2. Outline of proof of Theorem 1. Since $G(x, w) - (2S_\varepsilon|x - w|^2)^{-1}$ is bounded when x tends to w , we see there exists $C > 0$ such that

$$\omega_{\varepsilon + C\varepsilon^3} \subset \Omega_\varepsilon \subset \omega_{\varepsilon - C\varepsilon^3}$$

holds for any small $\varepsilon > 0$. Since the eigenvalues of the Laplacian with the Dirichlet condition depend monotonically on the domain, we can easily deduce Theorem 1 from the following

Proposition 1. Assume $n=4$. Let $0 > \tilde{\mu}_1(\varepsilon) > \dots \geq \tilde{\mu}_j(\varepsilon) \geq \dots$ be the eigenvalues of the Laplacian with the Dirichlet condition on $\partial\omega_\varepsilon$. We arrange them repeatedly according to their multiplicities. Fix j . If μ_j is a simple eigenvalue then

$$(2.1) \quad \tilde{\mu}_j(\varepsilon) - \mu_j = -2S_\varepsilon \varepsilon^2 \varphi_j(w)^2 + O(\varepsilon^{5/2})$$

holds when ε tends to zero.

We introduce various operators. For $\varepsilon > 0$, let $G_\varepsilon^{(2)}(x, y)$ be the kernel of the operator G_ε^2 defined by

$$(2.2) \quad (G_\varepsilon^2 g)(x) = \int_{\omega_\varepsilon} G_\varepsilon^{(2)}(x, y) g(y) dy \quad x \in \omega_\varepsilon.$$

Let G (resp. G^2) be the Green operator (resp. its iterated operator) given by

$$(2.3) \quad (Gf)(x) = \int_\Omega G(x, y) f(y) dy$$

(resp. $(G^2 f)(x) = \int_\Omega G^{(2)}(x, y) f(y) dy$).

To get Proposition 1 we compare the eigenvalues of G_ε^2 and G^2 .

To interpolate G_ε^2 and G^2 we introduce two operators H_ε with of \tilde{H}_ε given by the following :

$$(H_\varepsilon g)(x) = \int_{\omega_\varepsilon} h_\varepsilon(x, y) g(y) dy,$$

where

$$h_\varepsilon(x, y) = G^{(2)}(x, y) - 2S_\varepsilon \varepsilon^2 (G^{(2)}(x, w)G(y, w) + G(x, w)G^{(2)}(y, w))$$

for $x, y \in \omega_\varepsilon$.

$$(\tilde{H}_\varepsilon f)(x) = \int_\Omega \tilde{h}_\varepsilon(x, y) f(y) dy,$$

where

$$\tilde{h}_\varepsilon(x, y) = G^{(2)}(x, y) - 2S_\varepsilon \varepsilon^2 (G^{(2)}(x, w)G(y, w)\tilde{\Psi}_\varepsilon(y) + \tilde{\Psi}_\varepsilon(x)G(x, w)G^{(2)}(y, w))$$

for $x, y \in \Omega$. Here $\tilde{\Psi}_\varepsilon \in C^\infty(\mathbf{R}^4)$ is defined as $\tilde{\Psi}_\varepsilon(x) = 1$ on $|x - w| \geq \varepsilon/2$ and $\tilde{\Psi}_\varepsilon(x) = 0$ on $|x - w| \leq \varepsilon/4$.

We give the following

Theorem 2.

$$\|G_\varepsilon^2 - H_\varepsilon\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^{5/2}$$

for some constant C independent of $\varepsilon > 0$.

Our proof of Theorem 2 is rather involved, thus we put off its sketch until § 3.

We study \tilde{H}_ε . We have the following

Proposition 2. For any fixed $\varepsilon > 0$, \tilde{H}_ε is a compact selfadjoint operator in $L^2(\Omega)$.

It should be remarked that the perturbation family $\varepsilon \mapsto \tilde{H}_\varepsilon$ is not an analytic family of selfadjoint operators, thus some techniques are necessary to study eigenvalues of \tilde{H}_ε . Let λ_j be a simple eigenvalue of G . We construct an approximate eigenvalue of \tilde{H}_ε which tends to λ_j^2 when $\varepsilon \searrow 0$. We solve the following equation for $\tilde{\varphi}_\varepsilon$.

$$(2.4) \quad (G^2 - \lambda_j^2)\tilde{\varphi}_\varepsilon(x) = -2\lambda_j^2\varphi(w)\varphi(x)(G(\hat{\Psi}_\varepsilon \cdot \varphi))(w) \\ + G^{(2)}(x, w)(G(\hat{\Psi}_\varepsilon \cdot \varphi))(w) + \lambda_j^2 G(x, w)\hat{\Psi}_\varepsilon(x)\varphi(w).$$

Here φ denotes the eigenfunction of the Laplacian associated with λ_j and satisfying $\|\varphi\|_{L^2(\Omega)} = 1$. The above equation is solvable since the right-side of (2.4) is orthogonal to the kernel space spanned by φ and $G^2 - \lambda_j^2$ is the Fredholm operator. We have the following

Lemma 1. Put $r(\varepsilon) = 2S_4\varepsilon^2(G(\hat{\Psi}_\varepsilon \cdot \varphi))(w)$, then

$$(2.5) \quad (\tilde{H}_\varepsilon - (\lambda_j^2 - 2\lambda_j^2 r(\varepsilon)))(\varphi + 2S_4\varepsilon^2\tilde{\varphi}_\varepsilon^0) \\ = -4S_4^2\varepsilon^4(G^{(2)}(x, w)(G(\hat{\Psi}_\varepsilon \cdot \varphi))(w) + G(x, w)\hat{\Psi}_\varepsilon(x)(G^2\tilde{\varphi}_\varepsilon^0)(w) \\ - 2\lambda_j^2\varphi(w)\tilde{\varphi}_\varepsilon^0(x)(G(\hat{\Psi}_\varepsilon \cdot \varphi))(w)),$$

where $\tilde{\varphi}_\varepsilon^0$ is the unique solution of (2.4) orthogonal to φ .

From (2.4) we easily get

$$(2.6) \quad \|\tilde{\varphi}_\varepsilon^0\|_{L^2(\Omega)} \leq C|\log \varepsilon|^{1/2}$$

for some constant C independent of $\varepsilon > 0$. By Lemma 1 and (2.6) we have the following

Lemma 2. $L^2(\Omega)$ -norm of the left-side of (2.5) does not exceed $C\varepsilon^4|\log \varepsilon|$. And $\|\varphi + 2S_4\varepsilon^2\tilde{\varphi}_\varepsilon^0\|_{L^2(\Omega)} \geq 1$.

From Lemma 2 we can deduce the existence result for an approximate eigenvalue. We get the following

Proposition 3. There exists at least one eigenvalue $\tilde{\lambda}_j^{(2)}(\varepsilon)$ of \tilde{H}_ε satisfying

$$\tilde{\lambda}_j^{(2)}(\varepsilon) = \lambda_j^2 - 2\lambda_j^2(2S_4)\varepsilon^2\varphi(w)^2 + O(\varepsilon^4|\log \varepsilon|).$$

We compare \tilde{H}_ε with H_ε . Let ψ_ε be the eigenfunction of \tilde{H}_ε with respect to $\tilde{\lambda}_j^{(2)}(\varepsilon)$. Assume $\|\psi_\varepsilon\|_{L^2(\Omega)} = 1$. We put $\psi_{\varepsilon,1} = \chi_\varepsilon\psi_\varepsilon$, $\psi_{\varepsilon,2} = (1 - \chi_\varepsilon)\psi_\varepsilon$ where χ_ε is the characteristic function of ω_ε . Then these equations are equivalent to the following equations:

$$(2.7) \quad (H_\varepsilon - \tilde{\lambda}_j^{(2)}(\varepsilon))\psi_{\varepsilon,1}(x) = \int_{\beta_\varepsilon} \tilde{h}_\varepsilon(x, y)\psi_{\varepsilon,2}(y)dy \quad x \in \omega_\varepsilon$$

$$(2.8) \quad \int_{\omega_\varepsilon} \tilde{h}_\varepsilon(x, y) \psi_{\varepsilon,1}(y) dy + \int_{\beta_\varepsilon} \tilde{h}_\varepsilon(x, y) \psi_{\varepsilon,2}(y) dy = \tilde{\lambda}_j^{(2)}(\varepsilon) \psi_{\varepsilon,2} \quad x \in \beta_\varepsilon$$

$$(2.9) \quad \|\psi_{\varepsilon,1}\|_{L^2(\omega_\varepsilon)}^2 + \|\psi_{\varepsilon,2}\|_{L^2(\beta_\varepsilon)}^2 = 1.$$

By (2.7) and arguments from which we deduce Theorem 2 we have

$$\|(H_\varepsilon - \tilde{\lambda}_j^{(2)}(\varepsilon))\psi_{\varepsilon,1}\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^{4/p} \|\psi_{\varepsilon,2}\|_{L^{p'}(\beta_\varepsilon)}$$

for any fixed p satisfying $1 < p < 2$. Here p' is the conjugate number of p . On the other hand we get from (2.8)

$$\|\psi_{\varepsilon,2}\|_{L^{p'}(\beta_\varepsilon)} \leq C\varepsilon^{4/p'}$$

for $2 < p' < \infty$, and

$$\|\psi_{\varepsilon,2}\|_{L^2(\beta_\varepsilon)} \leq C\varepsilon^2.$$

Summing up these facts we have the following

Proposition 4. *There exists a constant C independent of ε such that*

$$\|(H_\varepsilon - \tilde{\lambda}_j^{(2)}(\varepsilon))\psi_{\varepsilon,1}\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^4$$

and

$$\|\psi_{\varepsilon,1}\|_{L^2(\omega_\varepsilon)} \geq 1/2$$

hold.

Thus we get the following

Proposition 5. *There exists at least one eigenvalue $\lambda_*^{(2)}(\varepsilon)$ of H_ε satisfying $\lambda_*^{(2)}(\varepsilon) = \lambda_j^2 - 2\lambda_j^3(2S_4)\varepsilon^2\varphi(w)^2 + O(\varepsilon^4|\log \varepsilon|)$.*

Now we prove Proposition 1. Let λ_j be as above. Then the following is known. See [5], [3].

Lemma 3. *Let V be a sufficiently small fixed neighbourhood of λ_j . Then there exists $\varepsilon_0 > 0$ such that only one simple eigenvalue $\lambda_j(\varepsilon)$ of G_ε is in V for any fixed ε satisfying $\varepsilon_0 > \varepsilon > 0$. Moreover $\lim_{\varepsilon \rightarrow 0} \lambda_j(\varepsilon) = \lambda_j$.*

From Theorem 3, Proposition 4 and Lemma 3 we get

$$\lambda_*^{(2)}(\varepsilon) - (\lambda_j(\varepsilon))^2 = O(\varepsilon^{5/2})$$

and

$$\lambda_j(\varepsilon)^2 = \lambda_j^2 - 2\lambda_j^3(2S_4)\varepsilon^2\varphi_j(w)^2 + O(\varepsilon^{5/2}).$$

Then we have Proposition 1.

§ 3. Rough sketch of a proof of Theorem 2. We need two Lemmas.

Lemma 4. *If u_ε satisfies*

$$\begin{cases} \Delta u_\varepsilon = 0 & \text{in } \omega_\varepsilon \\ u_\varepsilon|_{\partial\omega} = 0, & |u_\varepsilon|_{\partial\beta_\varepsilon} \leq H(\varepsilon) \end{cases}$$

then $|u| \leq CH(\varepsilon)\varepsilon^2|x-w|^{-2}$ for $x \in \omega_\varepsilon$. C is a constant independent of ε .

Lemma 5. *If u_ε satisfies*

$$\begin{cases} \Delta^2 u_\varepsilon = 0 & \text{in } \omega_\varepsilon \\ u_\varepsilon|_{\partial\omega} = \Delta u_\varepsilon|_{\partial\omega} = 0 \\ |u_\varepsilon|_{\partial\beta_\varepsilon} \leq M(\varepsilon), & |\Delta u_\varepsilon|_{\partial\beta_\varepsilon} \leq N(\varepsilon) \end{cases}$$

then

$$\|u_\varepsilon\|_{L^{8/3}(\omega_\varepsilon)} \leq C(N(\varepsilon)\varepsilon^2 + M(\varepsilon)\varepsilon^{3/2})$$

holds for some constant C independent of ε .

Sketch of proof of Theorem 2. Fix $f \in C_0^\infty(\omega_\varepsilon)$. And we put $u_\varepsilon = (G_\varepsilon^2 - H_\varepsilon)f$. Then we have $\Delta^2 u_\varepsilon = 0$ in ω_ε and $u_\varepsilon|_{\partial\Omega} = \Delta u_\varepsilon|_{\partial\Omega} = 0$. To estimate $L^{8/3}(\omega_\varepsilon)$ -norm of u_ε , we need bounds for $M(\varepsilon)$ and $N(\varepsilon)$ in Lemma 5.

Since we have

$$|u_\varepsilon|_{\partial\beta_\varepsilon} \leq \int_{\omega_\varepsilon} |G^{(2)}(x, y) - G^{(2)}(y, w)|_{x \in \partial\beta_\varepsilon} |f(y)| dy + CG^{(2)}(x, w)|_{x \in \partial\beta_\varepsilon} \varepsilon^2 |(Gf)(w)|$$

and

$$|\Delta u_\varepsilon|_{\partial\beta_\varepsilon} \leq \int_{\omega_\varepsilon} |G(x, y) - G(y, w)|_{x \in \partial\beta_\varepsilon} |f(y)| dy,$$

we can take $M(\varepsilon)$, $N(\varepsilon)$ as

$$M(\varepsilon) = \tilde{C} |\log \varepsilon| \varepsilon^2 \|f\|_{L^{8/3}(\omega_\varepsilon)} + \tilde{C} \varepsilon \|f\|_{L^2(\omega_\varepsilon)}$$

$$N(\varepsilon) = \tilde{C} \varepsilon^{1/2} \|f\|_{L^{8/3}(\omega_\varepsilon)}$$

for a constant \tilde{C} independent of ε . Therefore we have

$$\|G_\varepsilon^2 - H_\varepsilon\|_{L^{8/3}(\omega_\varepsilon)} \leq C\varepsilon^{5/2}.$$

Since we have

$$\|(G_\varepsilon^2 - H_\varepsilon)^*\|_{L^{8/5}(\omega_\varepsilon)} = \|G_\varepsilon^2 - H_\varepsilon\|_{L^{8/3}(\omega_\varepsilon)}$$

and

$$(G_\varepsilon^2 - H)^*|_{C_0^\infty(\omega_\varepsilon)} = G_\varepsilon^2 - H_\varepsilon,$$

we get Theorem 2.

Errata in [1]. The right-side of the formula in Theorem 1 in [1] should be replaced by

$$G(x, y) - (m - n - 2)S_{m-n} \varepsilon^{m-n-2} \int_N G(x, w)G(y, w)dw + O(\varepsilon^{m-n}).$$

References

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