

SMOOTH STEADY SOLUTIONS OF THE PLANAR VLASOV-POISSON SYSTEM WITH A MAGNETIC OBSTACLE

BY

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Abstract. The solar wind interacting with a magnetized obstacle is modeled with the Vlasov equation. The domain considered is a disk in the plane. Inflowing boundary conditions are given for the particle density. A magnetic field is prescribed, and the electric field is computed self consistently with potential zero on the boundary. Taking the boundary condition for the particle density to be sufficiently small, it is shown that there is a natural smooth steady solution. The speed of the inflowing plasma and the magnetic field are not size restricted.

1. Introduction. Let $R > 0$, $D = \{x \in \mathbb{R}^2 : |x| < R\}$, and $\partial D = \{x \in \mathbb{R}^2 : |x| = R\}$. Consider the problem

$$\left\{ \begin{array}{l} v \cdot \nabla_x f + (E_1 + v_2 B) \partial_{v_1} f + (E_2 - v_1 B) \partial_{v_2} f = 0 \text{ on } D \times \mathbb{R}^2 \\ f \text{ given if } x \in \partial D \text{ and } x \cdot v < 0 \\ \rho = \int f dv \\ \Delta U = -4\pi\rho \text{ in } D \\ U = 0 \text{ on } \partial D \\ E = -\nabla U \end{array} \right. \quad (1.1)$$

where B is a given function of x . Realistic modeling would require including an additional species of particles (with opposite charge), but this does not affect the methods used here, so it is omitted. It will be assumed that B has support contained in D . The boundary condition that $U = 0$ on ∂D says physically that the boundary is a perfect conductor.

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This is spurious, but in order to work on a bounded domain, some boundary condition must be given. This one gives a physically meaningful problem.

Let

$$S_+ = \{(x, v) \in \partial D \times \mathbb{R}^2 : x \cdot v > 0\},$$

$$S_- = \{(x, v) \in \partial D \times \mathbb{R}^2 : x \cdot v < 0\}$$

and

$$S_0 = \{(x, v) \in \partial D \times \mathbb{R}^2 : x \cdot v = 0\}.$$

THEOREM 1.1. Let $B \in C^1(\mathbb{R}^2)$ and assume there is $C_1 \in (0, R)$ such that

$$B(x) \neq 0 \Rightarrow |x| < C_1. \tag{1.2}$$

Let $F : S_- \cup S_0 \rightarrow [0, \infty)$ be C^1 and assume there are $C_2, C_3 \in (0, \infty)$ such that

$$F(x, v) \neq 0 \Rightarrow C_2 > |v| > C_3 \text{ and } v \cdot \frac{(-x)}{R} > C_3. \tag{1.3}$$

We also assume (1.5), which is stated below. Then there exists $C > 0$ such that for every $\varepsilon \in (0, C)$ there exists $f \in C^1(\overline{D} \times \mathbb{R}^2)$ and $U \in C^2(\overline{D})$ that satisfy (1.1) where the boundary condition for f is

$$f = \varepsilon F \text{ on } S_-. \tag{1.4}$$

Comments

1. In general, we do not expect uniqueness for the above problem. If there are characteristics which never intersect the boundary, then the value of f on these is not determined by (1.4); hence we expect uniqueness to fail. For the solution constructed in Theorem 1.1, f is zero along any characteristic that never intersects the boundary. Thus, all charge comes from upstream by way of condition (1.4).
2. Restricting ε limits the amount of charge the plasma carries. In the limit as $\varepsilon \rightarrow 0^+$, $E \rightarrow 0$ and the motion of the particles is determined by

$$\begin{cases} \frac{dX}{ds} = V \\ \frac{dV}{ds} = (V_2 B(X), -V_1 B(X)). \end{cases}$$

Condition (1.5) requires that when $\varepsilon = 0$ the particles exit the domain within a bounded time.

3. B and the v support of F are not size restricted. Thus, the inflowing plasma may move rapidly and encounter a large magnetic obstacle and hence change direction rapidly. Study of this problem was motivated by interest in the bow shock formed when the solar wind encounters the earth's magnetic field (see [14] for example).

4. The well-known theorems that ensure the existence of smooth solutions ([8], [9]) apply to the time dependent problem. The bounds they obtain increase without bound as t increases and hence cannot be used for the time independent case. Many papers on the steady problem (e.g. [1], [4]) concern solutions of the form

$$f(x, v) = G \left(\frac{1}{2}|v|^2 + U(x) \right).$$

However, this would require

$$f(x, v) = G \left(\frac{1}{2}|v|^2 \right) \text{ on } \partial D \times \mathbb{R}^2,$$

which conflicts with (1.4).

For excellent surveys of the mathematics of Vlasov equations see [3] and [10]. In the landmark papers [9] and [8] it is shown that solutions of the Poisson-Vlasov system in three space dimensions which start smooth remain smooth for all time. We mention [11], [12], [13] as other attempts to model a steady solar wind. Other papers that consider boundary value problems for the Vlasov equation are [5], [6], [7]. In particular, it was pointed out in [6] that regularity can be lost near characteristics that intersect the boundary in a tangential manner. Much effort is made in this work to ensure the solution is zero in a neighborhood of any such tangential characteristic. One factor in this is assumption (1.3). Another is the restriction of ε which in turn restricts the size of E . Note, also, that $B = 0$ near ∂D by (1.2).

The letter C will denote a generic positive constant which changes from line to line and may depend on B, F, R but not on ε or n . When a specific constant is chosen it will be given a subscript. For example, C_1, C_2, C_3 introduced in Theorem 1.1 will have the same values throughout the paper. The norm

$$\|E\|_\infty = \sup \{|E(x)| : x \in D\}$$

will be used. Also $\bar{D} = D \cup (\partial D)$ and

$$(v_1, v_2) \wedge B = (v_2 B, -v_1 B).$$

Consider the case $\varepsilon = 0$ first. For $(x, v) \in \mathbb{R}^2 \times \mathbb{R}^2$ define $(X^0(s, x, v), V^0(s, x, v))$ by

$$\begin{cases} \frac{dX^0}{ds} = V^0 & X^0(0, x, v) = x \\ \frac{dV^0}{ds} = V^0 \wedge B(X^0) & V^0(0, x, v) = v. \end{cases}$$

For $(x, v) \in S_-$ define

$$\omega^0(x, v) = \sup \{t > 0 : X(s, x, v) \in D \forall s \in (0, t)\}$$

and assume there exists $T^0 > 0$ such that

$$\omega^0(x, v) \leq T^0 \tag{1.5}$$

for all $(x, v) \in S_-$ with $F(x, v) \neq 0$.

LEMMA 1.1. There exists $C_4 > 0$ such that $X^0(\omega^0(x, v), x, v) \in \partial D$ and

$$V^0(\omega^0(x, v), x, v) \cdot n \geq C_4, \tag{1.6}$$

where $n = R^{-1}X^0(\omega^0(x, v), x, v)$, for all $(x, v) \in S_-$ with $F(x, v) \neq 0$.

Proof. For brevity we drop the dependence on (x, v) . It follows by a standard continuity argument (and since ω^0 is finite) that $X^0(\omega^0) \in \partial D$. Note, also, that for all $s \in (0, \omega^0)$,

$$|V^0(s)| = |v|.$$

Consider the case that $B(X^0(s)) \neq 0$ for some $s \in (0, \omega^0)$. Define

$$t_1 = \sup \{t \in (s, \omega^0) : |X^0| \leq C_1 \text{ on } (s, t)\}.$$

Then $|X^0| > C_1$ on (t_1, ω^0) so $B(X^0) = 0$ and V^0 is constant on $[t_1, \omega^0]$. Using some geometry it follows that

$$\sin \theta \leq C_1/R$$

where θ is the angle between $V(\omega^0)$ and n . Hence

$$V^0(\tau^0) \cdot n = |v| \cos \theta \geq C_3 \cos \left(\sin^{-1} \left(\frac{C_1}{R} \right) \right) = C.$$

Now consider the case that $B(X^0(s)) = 0$ for all $s \in (0, \omega^0)$. Then $V^0(s) = v$ for all $s \in [0, \omega^0]$ and

$$V^0(\tau^0) \cdot n = v \cdot \left(\frac{-x}{R} \right) > C_3,$$

completing the proof. □

2. The linear Vlasov equation. In this section, we consider $E \in C^1(\overline{D})$ given. A solution of the linear Vlasov equation will be defined that satisfies

$$f = \varepsilon F \quad \text{on } S_- \cup S_0.$$

The goal of this section is to show that there is a constant $C > 0$ such that for $\|E\|_\infty < C$ we have $f \in C^1(\overline{D} \times \mathbb{R}^2)$ and

$$\|f\|_\infty + \|\nabla f\|_\infty \leq C\varepsilon.$$

Consider $(x, v) \in (D \times \mathbb{R}^2) \cup S_- \cup S_+$. Define $(X(s, x, v), V(s, x, v))$ by

$$\begin{cases} \frac{dX}{ds} = V & X(0, x, v) = x \\ \frac{dV}{ds} = E(X) + V \wedge B(X) & V(0, x, v) = v. \end{cases} \tag{2.1}$$

Next we define f . On $S_- \cup S_0$ define

$$f = \varepsilon F. \tag{2.2}$$

Consider $(x, v) \in (D \times \mathbb{R}^2) \cup S_+$. Define

$$\alpha(x, v) = \inf \{t < 0 : X(s, x, v) \in D \quad \forall s \in (t, 0)\} \tag{2.3}$$

and then

$$f(x, v) = \varepsilon F(X(\alpha(x, v), x, v), V(\alpha(x, v), x, v)) \tag{2.4}$$

if $\alpha(x, v)$ is finite and

$$f(x, v) = 0 \tag{2.5}$$

otherwise.

When $f(x, v) \neq 0$ we may bound $\alpha(x, v)$ by using T^0 :

LEMMA 2.1. There exists $C_5 > 0$ such that for $\|E\|_\infty < C_5$ the following holds: Consider $(x, v) \in S_-$ with $F(x, v) \neq 0$ and define

$$\omega(x, v) = \sup \{t > 0 : X(s, x, v) \in D \ \forall s \in (0, t)\}.$$

Then

$$\omega(x, v) \leq 2T^0$$

and

$$V(\omega(x, v), x, v) \cdot n \geq \frac{1}{2}C_4$$

where

$$n = R^{-1}X(\omega(x, v), x, v).$$

Proof. Dropping the dependence on (x, v) we have

$$\begin{aligned} \left| \frac{dV}{ds} - \frac{dV^0}{ds} \right| &= |E(X) + V \wedge B(X) - V^0 \wedge B(X^0)| \\ &\leq \|E\|_\infty + |V - V^0||B(X)| + |V^0| |B(X) - B(X^0)| \\ &\leq \|E\|_\infty + C(|X - X^0| + |V - V^0|). \end{aligned}$$

Hence, for $s \in (0, \min(\omega, 2T^0))$ we have

$$|X(s) - X^0(s)| + |V(s) - V^0(s)| \leq 2T^0\|E\|_\infty + C \int_0^s (|X - X^0| + |V - V^0|)du,$$

and Gronwall's inequality yields

$$\begin{aligned} |X(s) - X^0(s)| + |V(s) - V^0(s)| &\leq 2T\|E\|_\infty e^{Cs} \\ &\leq 2T^0\|E\|_\infty e^{C2T^0} = C\|E\|_\infty. \end{aligned} \tag{2.6}$$

Now we sketch the rest of the proof. Let

$$\mathcal{E}^0 = V^0 \cdot (|X^0|^{-1}X^0) \text{ and } \mathcal{E} = V \cdot (|X|^{-1}X).$$

For illustration consider the case when $\omega < \omega^0$ (and hence $\omega < T^0$). Since we may restrict $\|E\|_\infty$, (2.6) allows us to restrict $|X(\omega) - X^0(\omega)| + |V(\omega) - V^0(\omega)|$, and hence

$$R - |X^0(\omega)| = |X(\omega)| - |X^0(\omega)|$$

may be made as small as desired. As in the proof of Lemma 1.1 we may obtain

$$\mathcal{E}^0 \geq C$$

on $[\omega, \omega^0]$. It follows that $\omega - \omega^0$ may be made arbitrarily small and that

$$\mathcal{E}(\omega) \geq \mathcal{E}^0(\omega^0) - |\mathcal{E}^0(\omega) - \mathcal{E}^0(\omega^0)| - |\mathcal{E}^0(\omega) - \mathcal{E}(\omega)| \geq C.$$

The case when $\omega \geq \omega^0$ may be handled similarly, so the proof is complete. □

Comment

It follows from Lemma 2.1 that if $(x, v) \in (D \times \mathbb{R}^2) \cup S_+$ and $f(x, v) \neq 0$, then

$$\alpha(x, v) \geq -2T^0.$$

LEMMA 2.2. There is $C_6 \in (0, C_5]$ such that if $\|E\|_\infty < C_6$, then $f \in C^1(\bar{D} \times \mathbb{R}^2)$.

Proof. We'll take

$$\|E\|_\infty < C_6 = \min \left(C_5, \frac{1}{4}C_3(2T_0)^{-1}, \left(\frac{1}{4}C_3 \right)^2 R^{-1} \right).$$

Consider $(x, v) \in (D \times \mathbb{R}^2) \cup S_+$. If $f(x, v) \neq 0$, then $\alpha(x, v) \geq -2T^0$ so

$$|V(\alpha(x, v), x, v) - v| \leq \|E\|_\infty |\alpha| < \frac{1}{4}C_3.$$

But

$$f(x, v) = \varepsilon F(X(\alpha(x, v), x, v), V(\alpha(x, v), x, v)) \neq 0,$$

so

$$|V(\alpha(x, v), x, v)| \geq C_3$$

and

$$|v| > \frac{3}{4}C_3.$$

So

$$|v| \leq \frac{3}{4}C_3 \Rightarrow f(x, v) = 0. \tag{2.7}$$

Suppose $\alpha(x, v) \geq -2T^0$ and

$$X(\alpha(x, v), x, v) \cdot V(\alpha(x, v), x, v) = 0.$$

We claim that there is a neighborhood of (x, v) on which

$$f(y, w) = 0. \tag{2.8}$$

To show this, note that (writing $X(\alpha) = X(\alpha(x, v), x, v)$, etc.) $|X(\alpha)|^2 = R^2$,

$$\left. \frac{d}{ds} |X(s)|^2 \right|_{s=\alpha} = 0,$$

and $|X(s)|^2 \leq R^2$ for all $s \in [\alpha, 0]$, so

$$\begin{aligned} 0 &\geq \left. \frac{d^2}{ds^2} |X(s)|^2 \right|_{s=\alpha} = 2|V(\alpha)|^2 + 2X(\alpha) \cdot E(X(\alpha)) \\ &\geq 2(|V(\alpha)|^2 - R\|E\|_\infty) \end{aligned}$$

and

$$|V(\alpha)| \leq \sqrt{R\|E\|_\infty} \leq \frac{1}{4}C_3.$$

Hence

$$|v| \leq |V(\alpha)| + \|E\|_\infty |\alpha| \leq \frac{1}{4}C_3 + \|E\|_\infty 2T^0 < \frac{1}{2}C_3.$$

Now (2.8) follows from (2.7).

Suppose $\alpha(x, v) \in [-\infty, -2T^0)$. Then there is a neighborhood of (x, v) on which $\alpha(y, w) \in [-\infty, -2T^0)$ and hence

$$f(y, w) = 0. \tag{2.9}$$

Suppose $\alpha(x, v) \geq -2T^0$ and

$$X(\alpha(x, v), x, v) \cdot V(\alpha(x, v), x, v) \neq 0.$$

Then there is a neighborhood of (x, v) on which

$$\alpha(y, w) \geq -3T^0 \tag{2.10}$$

and

$$f(y, w) = \varepsilon F(X(\alpha(y, w), y, w), V(\alpha(y, w), y, w)). \tag{2.11}$$

On this neighborhood f is the composition of C^1 functions.

Finally, for $(x, v) \in S_0$,

$$f(y, w) = 0$$

on a neighborhood of (x, v) , and if $(x, v) \in S_-$ there is a neighborhood of (x, v) on which f is the composition of C^1 functions. Thus the proof is complete. \square

LEMMA 2.3. There exists $C_7 > 0$ such that if $\|E\|_\infty < C_6$ we have

$$f(x, v) \neq 0 \Rightarrow |v| \leq C_7$$

and

$$\int f(x, v) dv \leq C_7 \varepsilon. \tag{2.12}$$

Proof. Assume $f(x, v) \neq 0$. If $(x, v) \in S_- \cup S_0$, then $f(x, v) = \varepsilon F(x, v)$ so $|v| \leq C_2$. Suppose $(x, v) \in (D \times \mathbb{R}^2) \cup S_+$. By (2.7) $|V(s)| \neq 0$ for all $s \in [\alpha, 0]$, so

$$\begin{aligned} |v| &= |V(\alpha)| + \int_\alpha^0 \frac{V(s)}{|V(s)|} \cdot E(X(s)) ds \\ &\leq C_2 + |\alpha| \|E\|_\infty \leq C_2 + 2T^0 C_6. \end{aligned}$$

Also, (2.12) follows since $\|f\|_\infty = \|\varepsilon F\|_\infty$, completing the proof. \square

LEMMA 2.4. There exists $C_8 > 0$ such that

$$\|\nabla f\|_\infty \leq C_8 e^{C_8 \|\nabla E\|_\infty} \varepsilon$$

if $\|E\|_\infty < C_6$.

Proof. Let $(x, v) \in S_-$. Then $f(X(s), V(s)) = \varepsilon F(x, v)$, and since f is C^1 it follows that

$$v \cdot \nabla_x f + (E + v \wedge B) \cdot \nabla_v f = 0.$$

By using a regularization argument it follows that

$$\begin{aligned} \partial_{x_i} f(X(s), V(s)) &= \varepsilon \partial_{x_i} F(x, v) \\ &\quad - \int_0^s (\partial_{x_i} E(X) + V \wedge \partial_{x_i} B(X)) \cdot \nabla_v f(X, V) du \end{aligned}$$

for $i = 1, 2$,

$$\begin{aligned} \partial_{v_1} f(X(s), V(s)) &= \varepsilon \partial_{v_1} F(x, v) \\ &\quad - \int_0^s (\partial_{x_1} f(X, V) - B(X) \partial_{v_2} f(X, V)) du \end{aligned}$$

and

$$\begin{aligned} \partial_{v_2} f(X(s), V(s)) &= \varepsilon \partial_{v_2} F(x, v) \\ &\quad - \int_0^s (\partial_{x_2} f(X, V) + B(X) \partial_{v_1} f(X, V)) du. \end{aligned}$$

Hence

$$\begin{aligned} |\nabla_x f(X(s), V(s))| &\leq \|\nabla_x F\|_\infty \varepsilon \\ &\quad + (\sqrt{2} \|E\|_\infty + C_7 \|\nabla B\|_\infty) \int_0^s |\nabla_v f(X, V)| du \end{aligned}$$

and

$$\begin{aligned} |\nabla_v f(X(s), V(s))| &\leq \|\nabla_v F\|_\infty \varepsilon \\ &\quad + \int_0^s (|\nabla_x f(X, V)| + \|B\|_\infty |\nabla_v f(X, V)|) du. \end{aligned}$$

Hence

$$|\nabla f(X(s), V(s))| \leq C\varepsilon + C(1 + \|\nabla E\|_\infty) \int_0^s |\nabla f(X, V)| du,$$

and by Gronwall's inequality

$$\begin{aligned} |\nabla f(X(s), V(s))| &\leq C\varepsilon e^{C(1+\|\nabla E\|_\infty)s} \\ &\leq C\varepsilon e^{C(1+\|\nabla E\|_\infty)2T^0} = C e^{C\|\nabla E\|_\infty} \varepsilon. \end{aligned}$$

The lemma now follows. □

Next define

$$\rho = \int f dv,$$

\tilde{U} by

$$\begin{cases} \Delta \tilde{U} &= -4\pi\rho \quad \text{on } \bar{D} \\ \tilde{U} &= 0 \quad \text{on } \partial D, \end{cases}$$

and

$$\tilde{E} = -\nabla \tilde{U}.$$

LEMMA 2.5. There exist $C_9 > 1$ such that

$$\|\tilde{E}\|_\infty \leq C_9 \|\rho\|_\infty$$

and

$$\|\nabla \tilde{E}\|_\infty \leq C_9(1 + \|\rho\|_\infty) \ln(C_9(1 + \|\rho\|_\infty + \|\nabla \rho\|_\infty)).$$

The proof is deferred to the appendix.

3. Iteration. Define $f_0 = 0$, $\rho_0 = 0$, and $E_0 = 0$. For $E_n \in C^1(\overline{D})$ known with $\|E_n\|_\infty < C_6$, define f_{n+1} , ρ_{n+1} , and E_{n+1} as follows. Let $E = E_n$. Then take f, ρ , and \tilde{E} as in the previous section. Now define $f_{n+1} = f$, $\rho_{n+1} = \rho$, and $E_{n+1} = \tilde{E}$. By Lemmas 2.5 and 2.3

$$\|E_{n+1}\|_\infty \leq C_9 \|\rho_{n+1}\|_\infty$$

and

$$\|\rho_{n+1}\|_\infty \leq C_7 \varepsilon.$$

We require

$$\varepsilon < \frac{C_6}{C_9 C_7}$$

so that

$$\|E_{n+1}\|_\infty < C_6.$$

By induction it follows that f_n, ρ_n , and E_n are defined and satisfy

$$\|E_n\|_\infty < C_6$$

and

$$\|\rho_n\|_\infty \leq C_7 \varepsilon < C_6 / C_9 \tag{3.1}$$

for all n . Also

$$\|E_n\|_\infty < C_9 C_7 \varepsilon \tag{3.2}$$

for all n .

LEMMA 3.1. There exist $C_{10} > 0$ and $C_{11} > 0$ such that if $\varepsilon < C_{10}$, then

$$\|\nabla E_n\|_\infty < C_{11}$$

for all n .

Proof. By Lemmas 2.5, 2.3, 2.4 and by (3.1) we have

$$\begin{aligned} \|\nabla E_{n+1}\|_\infty &\leq C_9(1 + \|\rho_{n+1}\|_\infty) \ln(C_9(1 + \|\rho_{n+1}\|_\infty + \|\nabla \rho_{n+1}\|_\infty)) \\ &\leq (C_9 + C_6) \ln(C_9 + C_6 + C_9 \|\nabla \rho_{n+1}\|_\infty) \\ &\leq (C_9 + C_6) \ln(C_9 + C_6 + C_9 \pi C_7^2 C_8 e^{C_8 \|\nabla E_n\|_\infty} \varepsilon). \end{aligned}$$

So there is $C_{12} > 0$ such that

$$\|\nabla E_{n+1}\| \leq C_{12} \ln(C_{12}(1 + e^{C_{12} \|\nabla E_n\|_\infty} \varepsilon)).$$

Take

$$C_{11} = C_{12} \ln(2C_{12}),$$

$$C_{10} = \min(C_6 C_7^{-1} C_9^{-1}, e^{-C_{12} C_{11}}),$$

and

$$\varepsilon < C_{10}.$$

If $\|\nabla E_n\|_\infty < C_{11}$, then

$$\|\nabla E_{n+1}\|_\infty \leq C_{12} \ln(C_{12}(1 + e^{C_{12} C_{11}} \varepsilon))$$

$$< C_{12} \ln(C_{12}(1 + 1)) = C_{11}.$$

The lemma now follows by induction. □

Before addressing the convergence of this iteration, a preliminary lemma is needed.

LEMMA 3.2. Assume $E \in C^1(\overline{D})$ and $\|E\|_\infty < C_6$. For any $\delta > 0$ there is $r \in [C_1, R)$ such that if $|x| \in [r, R)$ and

$$v \cdot \frac{(-x)}{|x|} \geq \delta,$$

then

$$\alpha(x, v) \geq -\frac{R - |x|}{\frac{1}{2}\delta} \tag{3.3}$$

and

$$V(s, x, v) \cdot \frac{(-X(s, x, v))}{|X(s, x, v)|} \geq \frac{1}{2}\delta \tag{3.4}$$

for all $s \in [\alpha(x, v), 0]$.

Proof. Let $\alpha_\delta = \max\left(\alpha, \frac{-\delta}{2C_6}\right)$. For $s \in [\alpha_\delta, 0]$,

$$\begin{aligned} V(s) \cdot \frac{(-X(s))}{|X(s)|} &= v \cdot \frac{(-x)}{|x|} \\ &+ \int_s^0 \left(E(X) \cdot \frac{X}{|X|} + \frac{|X|^2 |V|^2 - (X \cdot V)^2}{|X|^3} \right) du \end{aligned} \tag{3.5}$$

$$\geq \delta - \int_s^0 C_6 du = \delta + C_6 s \geq \frac{1}{2}\delta$$

and

$$|X(s)| = |x| + \int_s^0 V \cdot \frac{(-X)}{|X|} du \geq |x| - \frac{1}{2}\delta s. \tag{3.6}$$

Take

$$r = \max\left(C_1, R - \frac{\delta^2}{8C_6}\right).$$

If $\alpha \leq \frac{-\delta}{2C_6}$, then $\alpha_\delta = \frac{-\delta}{2C_6}$ and

$$R \geq |X(\alpha_\delta)| \geq |x| - \frac{1}{2}\delta\alpha_\delta \geq r + \frac{\delta^2}{4C_6} \geq R + \frac{\delta^2}{8C_6},$$

a contradiction. Hence $\alpha > \frac{-\delta}{2C_6}$, and (3.4) follows from (3.5). Also by (3.6) we have

$$R = |X(\alpha)| \geq |x| - \frac{1}{2}\delta\alpha,$$

and (3.3) follows, completing the proof. □

Define $(X_{n+1}(s, x, v), V_{n+1}(s, x, v))$ by

$$\begin{cases} \frac{dX_{n+1}}{ds} = V_{n+1} & X_{n+1}(0, x, v) = x \\ \frac{dV_{n+1}}{ds} = E_n(X_{n+1}) + V_{n+1} \wedge B(X_{n+1}) & V_{n+1}(0, x, v) = v \end{cases}$$

and

$$\alpha_n(x, v) = \inf \{t < 0 : X(s, x, v) \in D \ \forall s \in (t, 0)\},$$

as in (2.1) and (2.3).

Now consider $(x, v) \in (D \times \mathbb{R}^2) \cup S_+$ with either $f_n(x, v) \neq 0$ or $f_{n+1}(x, v) \neq 0$. We'll consider the case $f_n(x, v) \neq 0$; the other case may be handled similarly. Then

$$\alpha_n(x, v) \geq -2T^0,$$

$$|V_n| \leq C_7,$$

and

$$V_n \cdot \frac{(-X_n)}{|X_n|} \Big|_{(\alpha_n(x,v), x, v)} \geq C_3. \tag{3.7}$$

For $s \in [\max(\alpha_n, \alpha_{n+1}), 0]$ let

$$d_n(s, x, v) = d_n(s) = |X_{n+1}(s) - X_n(s)| + |V_{n+1}(s) - V_n(s)|.$$

Then

$$\begin{aligned} d_n(s) &= \left| \int_s^0 (V_{n+1} - V_n) du \right| \\ &\quad + \left| \int_s^0 (E_n(X_{n+1}) + V_{n+1} \wedge B(X_{n+1}) - E_{n-1}(X_n) - V_n \wedge B(X_n)) du \right| \\ &\leq \int_s^0 (d_n + \|\nabla E_n\|_\infty d_n + \|E_n - E_{n-1}\|_\infty + \|B\|_\infty d_n + |V_n| \|\nabla B\|_\infty d_n) du \\ &\leq \int_s^0 (Cd_n + \|E_n - E_{n-1}\|_\infty) du. \end{aligned}$$

Since

$$s \geq \max(\alpha_n, \alpha_{n+1}) \geq \alpha_n \geq -2T^0$$

we have

$$d_n(s) \leq 2T^0 \|E_n - E_{n-1}\|_\infty + C \int_s^0 d_n du,$$

and Gronwall’s inequality yields

$$\begin{aligned}
 d_n(s) &\leq 2T^0 \|E_n - E_{n-1}\|_\infty e^{C|s|} \\
 &\leq 2T^0 \|E_n - E_{n-1}\|_\infty e^{C2T^0} = C_{13} \|E_n - E_{n-1}\|_\infty.
 \end{aligned}
 \tag{3.8}$$

By (3.2) we have

$$d_n(s) < C_{13} 2C_9 C_7 \varepsilon. \tag{3.9}$$

We further restrict ε so that

$$C_{13} 2C_9 C_7 \varepsilon < \frac{1}{2} C_3. \tag{3.10}$$

Now we choose r by using Lemma 3.2 with $\delta = \frac{1}{2} C_3$. Further restrict ε again so that

$$C_{13} 2C_9 C_7 \varepsilon < R - r. \tag{3.11}$$

Define

$$D_n(x, v) = D_n = \sup \{d_n(s) : s \in [\max(\alpha_n, \alpha_{n+1}), 0]\}$$

and claim (with an additional restriction on ε) that

$$|\alpha_{n+1} - \alpha_n| \leq 4C_3^{-1} D_n \tag{3.12}$$

and there is $C_{14} > 0$ such that

$$|f_{n+1}(x, v) - f_n(x, v)| \leq C_{14} \varepsilon D_n. \tag{3.13}$$

To establish (3.12) consider the case that $\alpha_{n+1} \leq \alpha_n$ first. By (3.9) and (3.11) we have

$$\begin{aligned}
 |X_{n+1}(\alpha_n)| &\geq |X_n(\alpha_n)| - D_n = R - D_n \\
 &\geq R - (R - r) = r.
 \end{aligned}$$

Also, by (3.9)

$$\begin{aligned}
 &\left| V_{n+1} \cdot \frac{X_{n+1}}{|X_{n+1}|} - V_n \cdot \frac{X_n}{|X_n|} \right|_{\alpha_n} \\
 &\leq |V_{n+1} - V_n| + |V_n| \frac{|(X_{n+1} - X_n)|X_n| + X_n(|X_n| - |X_{n+1}|)}{|X_{n+1}||X_n|} \\
 &\leq D_n + \frac{C_7}{C_1} 2D_n < \left(1 + \frac{2C_7}{C_1}\right) C_{13} 2C_9 C_7 \varepsilon.
 \end{aligned}$$

We restrict ε so that

$$\left(1 + \frac{2C_7}{C_1}\right) C_{13} 2C_9 C_7 \varepsilon < \frac{1}{2} C_3.$$

Now using (3.7) we have

$$V_{n+1} \cdot \frac{(-X_{n+1})}{|X_{n+1}|} \Big|_{\alpha_n} \geq V_n \cdot \frac{(-X_n)}{|X_n|} \Big|_{\alpha_n} - \frac{1}{2} C_3 \geq \frac{1}{2} C_3.$$

Now by Lemma 3.2 we have

$$\begin{aligned} \alpha_{n+1} &\geq \alpha_n - \frac{R - |X_{n+1}(\alpha_n)|}{\frac{1}{2}\delta} = \alpha_n - \frac{|X_n(\alpha_n)| - |X_{n+1}(\alpha_n)|}{\frac{1}{4}C_3} \\ &\geq \alpha_n - \frac{4D_n}{C_3}, \end{aligned}$$

and (3.12) follows when $\alpha_{n+1} \leq \alpha_n$.

To establish (3.13) note first that by (3.9) we have

$$D_n \leq C,$$

so (3.12) implies

$$\alpha_{n+1} \geq \alpha_n - 4C_3^{-1}D_n \geq -2T^0 - 4C_3^{-1}C = -C,$$

so for all $s \in [\alpha_{n+1}, 0]$,

$$|V_{n+1}(s)| \leq |v| + \|E\|_\infty |s| \leq C_7 + C_6C = C.$$

Hence,

$$\begin{aligned} &|f_{n+1}(x, v) - f_n(x, v)| \\ &= \varepsilon |F(X_{n+1}(\alpha_{n+1}), V_{n+1}(\alpha_{n+1})) - F(X_n(\alpha_n), V_n(\alpha_n))| \\ &\leq \varepsilon \|\nabla F\|_\infty (|X_{n+1}(\alpha_{n+1}) - X_n(\alpha_n)| + |V_{n+1}(\alpha_{n+1}) - V_n(\alpha_n)|) \\ &\leq C\varepsilon \left(\sup_{[\alpha_{n+1}, 0]} |V_{n+1}| |\alpha_{n+1} - \alpha_n| + |X_{n+1}(\alpha_n) - X_n(\alpha_n)| \right. \\ &\quad \left. + C_6|\alpha_{n+1} - \alpha_n| + |V_{n+1}(\alpha_n) - V_n(\alpha_n)| \right) \\ &\leq C\varepsilon (|\alpha_{n+1} - \alpha_n| + D_n) \leq C\varepsilon D_n, \end{aligned}$$

which is (3.13).

Establishing (3.12) and (3.13) in the case that $\alpha_{n+1} > \alpha_n$ is highly similar and is omitted. Also, (3.12) and (3.13) may be shown to hold in the case that $f_{n+1}(x, v) \neq 0$ with minor changes to the above estimates.

Combining (3.8) and (3.13) yields

$$|f_{n+1}(x, v) - f_n(x, v)| \leq C\varepsilon \|E_n - E_{n-1}\|_\infty \tag{3.14}$$

if $f_{n+1}(x, v) \neq 0$ or $f_n(x, v) \neq 0$. But if $f_{n+1}(x, v) = f_n(x, v) = 0$, (3.14) still holds, so

$$\|f_{n+1} - f_n\|_\infty \leq C\varepsilon \|E_n - E_{n-1}\|_\infty.$$

But by Lemmas 2.3 and 2.5 we have

$$\begin{aligned} \|E_n - E_{n-1}\|_\infty &\leq C_9 \|\rho_n - \rho_{n-1}\|_\infty \\ &\leq C_9 \pi C_7^2 \|f_n - f_{n-1}\|_\infty \end{aligned}$$

so

$$\|f_{n+1} - f_n\|_\infty \leq C_{15}\varepsilon\|f_n - f_{n-1}\|_\infty.$$

Requiring $\varepsilon < C_{15}^{-1}$ implies that f_n converges uniformly on \bar{D} to some continuous function f . Let

$$\begin{cases} \rho = \int f \, dv, \\ \Delta U = -4\pi\rho \\ U|_{\partial D} = 0, \\ E = -\nabla U. \end{cases}$$

Then $\rho_n \rightarrow \rho$ and $E_n \rightarrow E$ uniformly. By (3.8) and (3.12) it also follows that α_n converges uniformly; call the limit α . From (3.8) it follows that $D_n \rightarrow 0$ uniformly, and hence X_n and V_n converge uniformly on $\{(s, x, v) : (x, v) \in (D \times \mathbb{R}^2) \cup S_+ \text{ and } s \in (\alpha(x, v), 0]\}$, say to X and V . Furthermore,

$$\begin{cases} X(s, x, v) = x + \int_0^s V(u, x, v) \, du \\ V(s, x, v) = v + \int_0^s (E(X) + V \wedge B(X)) \, du. \end{cases} \tag{3.15}$$

By Lemmas 3.1 and 2.4 we have

$$\|\nabla f_n\|_\infty \leq C_8 e^{C_8 \|\nabla E_n\|_\infty} \varepsilon \leq C_8 e^{C_8 C_{11} \varepsilon} \varepsilon \leq C,$$

so

$$\begin{aligned} |f(x, v) - f(y, w)| &\leq |f(x, v) - f_n(x, v)| \\ &\quad + |f_n(x, v) - f_n(y, w)| + |f_n(y, w) - f(y, w)| \\ &\leq 2\|f_n - f\|_\infty + C|(x, v) - (y, w)| \end{aligned}$$

for all n . Hence

$$|f(x, v) - f(y, w)| \leq C|(x, v) - (y, w)|.$$

Also, ρ is Lipschitz continuous, and it follows from Theorem 4.13 of [2] that E is C^1 . Now by (3.15) it follows that X and V are C^1 . Then it further follows that α and hence

$$f(x, v) = \varepsilon F(X(\alpha(x, v), x, v), V(\alpha(x, v), x, v)) \tag{3.16}$$

are C^1 . Finally, that f satisfies the Vlasov equation follows from (3.16), and the proof of Theorem 1.1 is complete.

Appendix.

Proof of Lemma 2.5. For $x \in D$ and $\varepsilon \geq 0$ let

$$P^\varepsilon(x) = \int_D \rho(y) G^\varepsilon(x - y) \, dy$$

where

$$G^\varepsilon(z) = \ln(\sqrt{\varepsilon^2 + |z|^2}).$$

Let $\delta \in (0, R]$ and

$$B_\delta = \{y \in \mathbb{R}^2 : |x - y| < \delta\}.$$

Also let $e^{(1)} = (1, 0)$ and $e^{(2)} = (0, 1)$. Then for $\varepsilon > 0$,

$$\begin{aligned} \partial_{x_k x_\ell} P^\varepsilon(x) &= \int_D \rho(y) \partial_{x_k x_\ell} G^\varepsilon(x - y) dy \\ &= \int_{D \setminus B_\delta} \rho(y) \partial_{x_k x_\ell} G^\varepsilon(x - y) dy \\ &\quad - \int_{D \cap B_\delta} \rho(y) \nabla_y \cdot (\partial_{x_\ell} G^\varepsilon(x - y) e^{(k)}) dy \\ &= \int_{D \setminus B_\delta} \rho(y) \partial_{x_k x_\ell} G^\varepsilon(x - y) dy \\ &\quad + \int_{D \cap B_\delta} \nabla \rho(y) \cdot e^{(k)} \partial_{x_\ell} G^\varepsilon(x - y) dy \\ &\quad - \int_{(\partial B_\delta) \cap D} \rho(y) \partial_{x_\ell} G^\varepsilon(x - y) e^{(k)} \cdot ndS_y \\ &\quad - \int_{(\partial D) \cap B_\delta} \rho(y) \partial_{x_\ell} G^\varepsilon(x - y) e^{(k)} \cdot ndS_y \\ &= I + II + III + IV. \end{aligned}$$

Note that

$$|\partial_{x_k} G^\varepsilon(z)| \leq \frac{1}{\sqrt{\varepsilon^2 + |z|^2}}$$

and

$$|\partial_{x_k x_\ell} G^\varepsilon(z)| \leq \frac{3}{\varepsilon^2 + |z|^2},$$

so

$$\begin{aligned}
 |I| &\leq \|\rho\|_\infty \int_{D \setminus B_\delta} \frac{3}{\varepsilon^2 + |x - y|^2} dy \\
 &\leq 3\|\rho\|_\infty \int_\delta^{2R} \frac{2\pi r dr}{r^2} = 6\pi\|\rho\|_\infty \ln\left(\frac{2R}{\delta}\right),
 \end{aligned}$$

$$\begin{aligned}
 |II| &\leq \|\nabla\rho\|_\infty \int_{D \cap B_\delta} \frac{dy}{\sqrt{\varepsilon^2 + |x - y|^2}} \\
 &\leq \|\nabla\rho\|_\infty \int_0^\delta \frac{2\pi r dr}{r} = 2\pi\|\nabla\rho\|_\infty \delta,
 \end{aligned}$$

$$\begin{aligned}
 |III| &\leq \|\rho\|_\infty \int_{(\partial B_\delta) \cap D} \frac{1}{\sqrt{\varepsilon^2 + |x - y|^2}} dS_y \\
 &\leq \|\rho\|_\infty \int_0^{2\pi} \frac{\delta d\theta}{\sqrt{\varepsilon^2 + \delta^2}} \leq 2\pi\|\rho\|_\infty.
 \end{aligned}$$

To estimate IV consider $x \neq 0$ and let

$$n_0 = \frac{x}{|x|} = (\cos \theta_0, \sin \theta_0).$$

Write

$$\begin{aligned}
 IV &= \int_{(\partial D) \cap B_\delta} \left((\rho(x) - \rho(y))e^{(k)} \cdot n - \rho(x)e^{(k)} \cdot (n - n_0) \right. \\
 &\quad \left. - \rho(x)e^{(k)} \cdot n_0 \right) \partial_{x_\ell} G^\varepsilon(x - y) dS_y = IV_1 + IV_2 + IV_3.
 \end{aligned}$$

Then

$$|IV_1| \leq \|\nabla\rho\|_\infty \int_{(\partial D) \cap B_\delta} |x - y| \frac{dS_y}{\sqrt{\varepsilon^2 + |x - y|^2}} \leq C\|\nabla\rho\|_\infty \delta.$$

For IV_2 note that

$$\begin{aligned}
 |x - y|^2 &= |x|^2 + R^2 - 2x \cdot y = |x|^2 + R^2 - 2|x|Rn_0 \cdot n \\
 &= |x|^2 + R^2 - |x|R(2 - |n - n_0|^2) \\
 &= (R - |x|)^2 + |x|R|n - n_0|^2,
 \end{aligned}$$

so for $|x| \leq \frac{1}{2}R$,

$$|x - y|^2 \geq (R - \frac{1}{2}R)^2 = \frac{1}{4}R^2 \geq \frac{1}{16}R^2 |n - n_0|^2$$

and for $|x| \geq \frac{1}{2}R$,

$$|x - y|^2 \geq \left(\frac{1}{2}R\right) R |n - n_0|^2.$$

Hence

$$\begin{aligned} |IV_2| &\leq \|\rho\|_\infty \int_{(\partial D) \cap B_\delta} \frac{|n - n_0|}{\sqrt{\varepsilon^2 + |x - y|^2}} dS_y \\ &\leq \|\rho\|_\infty \int_{\partial D} \frac{|n - n_0|}{\frac{1}{4}R|n - n_0|} dS_y = 8\pi\|\rho\|_\infty. \end{aligned}$$

For $|IV_3|$ note that

$$IV_3 = -\rho(x)e^{(k)} \cdot n_0 \left(\int_{(\partial D) \cap B_\delta} \frac{x - y}{\varepsilon^2 + |x - y|^2} dS_y \right)_\ell.$$

It suffices to bound this integral when $x = (|x|, 0)$, in which case

$$\int_{(\partial D) \cap B_\delta} \frac{x - y}{\varepsilon^2 + |x - y|^2} dS_y = (\mathcal{I}, 0)$$

where

$$\mathcal{I} = \int_{(\partial D) \cap B_\delta} \frac{|x| - y_1}{\varepsilon^2 + |x - y|^2} dS_y.$$

For $|x| \leq \frac{1}{2}R$,

$$\frac{||x| - y_1|}{\varepsilon^2 + |x - y|^2} \leq \frac{\frac{1}{2}R + R}{(\frac{1}{2}R)^2} = \frac{6}{R}$$

so

$$|\mathcal{I}| \leq \int_{\partial D} \frac{6}{R} dS_y = 12\pi.$$

Consider $|x| \geq \frac{1}{2}R$. Writing $y = R(\cos \theta, \sin \theta)$ with $|\theta| \leq \pi$ we have

$$\begin{aligned} \left| \frac{|x| - y_1}{\varepsilon^2 + |x - y|^2} \right| &= \left| \frac{|x| - R \cos \theta}{\varepsilon^2 + |x|^2 + R^2 - 2|x|R \cos \theta} \right| \\ &= \frac{1}{2|x|} \left| 1 - \frac{\varepsilon^2 + R^2 - |x|^2}{\varepsilon^2 + (R - |x|)^2 + 2|x|R(1 - \cos \theta)} \right| \\ &\leq \frac{1}{2(\frac{1}{2}R)} \left(1 + \frac{\varepsilon^2 + R^2 - |x|^2}{\varepsilon^2 + (R - |x|)^2 + 2(\frac{1}{2}R)RC\theta^2} \right), \end{aligned}$$

and hence for $\varepsilon \leq 2R$,

$$\begin{aligned}
 |\mathcal{I}| &\leq \int_{-\pi}^{\pi} R^{-1} \left(1 + \frac{\varepsilon^2 + R^2 - |x|^2}{\varepsilon^2 + (R - |x|)^2 + CR^2\theta^2} \right) R d\theta \\
 &\leq 2\pi + \int_{-\infty}^{\infty} \frac{\varepsilon^2 + R^2 - |x|^2}{\varepsilon^2 + (R - |x|)^2 + CR^2\theta^2} d\theta \\
 &= 2\pi + (\varepsilon^2 + R^2 - |x|^2) \frac{\pi}{\sqrt{CR^2} \sqrt{\varepsilon^2 + (R - |x|)^2}} \\
 &= 2\pi + C \frac{\varepsilon^2 + R^2 - |x|^2}{\sqrt{\varepsilon^2 + (R - |x|)^2}} \leq 2\pi + C \frac{\varepsilon^2 + 2R(R - |x|)}{\frac{1}{\sqrt{2}}(\varepsilon + R - |x|)} \\
 &= 2\pi + \sqrt{2} C \frac{2R\varepsilon + 2R(R - |x|)}{\varepsilon + R - |x|} = C.
 \end{aligned}$$

Hence

$$|IV_3| \leq C\|\rho\|_{\infty}$$

and

$$|IV| \leq C\|\rho\|_{\infty} + C\|\nabla\rho\|_{\infty}\delta \tag{A.1}$$

for $x \neq 0$. By continuity (A.1) holds for $x = 0$ also.

Collecting terms we have

$$|\partial_{x_k x_\ell} P^\varepsilon(x)| \leq C \left(\|\rho\|_{\infty} \ln \left(\frac{2R}{\delta} \right) + \|\nabla\rho\|_{\infty}\delta \right)$$

for all $\delta \in (0, R]$. If $\|\rho\|_{\infty} < \|\nabla\rho\|_{\infty}R$ take

$$\delta = \|\rho\|_{\infty} / \|\nabla\rho\|_{\infty},$$

which yields

$$|\partial_{x_k x_\ell} P^\varepsilon(x)| \leq C\|\rho\|_{\infty} \left(1 + \ln \left(\frac{2R\|\nabla\rho\|_{\infty}}{\|\rho\|_{\infty}} \right) \right).$$

If $\|\rho\|_{\infty} \geq \|\nabla\rho\|_{\infty}R$ take $\delta = R$, which yields

$$|\partial_{x_k x_\ell} P^\varepsilon(x)| \leq C\|\rho\|_{\infty} + CR\|\nabla\rho\|_{\infty} \leq C\|\rho\|_{\infty}.$$

In both cases it follows that

$$|\partial_{x_k x_\ell} P^\varepsilon(x)| \leq C(1 + \|\rho\|_{\infty})(1 + \ln(1 + \|\nabla\rho\|_{\infty})). \tag{A.2}$$

Note also that

$$\begin{aligned}
 |\partial_{x_k} P^\varepsilon(x)| &= \left| \int_D \rho(y) \partial_{x_k} G_\varepsilon(x - y) dy \right| \\
 &\leq \|\rho\|_{\infty} \int_D \frac{dy}{\sqrt{\varepsilon^2 + |x - y|^2}} \leq \|\rho\|_{\infty} \int_0^{2R} \frac{2\pi r dr}{r} \\
 &\leq 4\pi R\|\rho\|_{\infty}.
 \end{aligned} \tag{A.3}$$

Letting $\varepsilon \rightarrow 0$ shows that (A.2) and (A.3) hold for $\varepsilon = 0$ also.

Next let

$$H(x) = \int_D \rho(y) \ln \left(\frac{|y|}{R} \left| x - \frac{R^2 y}{|y|^2} \right| \right) dy$$

for $x \in \bar{D}$. Then

$$\begin{cases} \Delta(P^0 - H) = 2\pi\rho & \text{in } D \\ P^0 - H = 0 & \text{on } \partial D. \end{cases}$$

Letting $z = R^2 y / |y|^2$ we have $y = R^2 z / |z|^2$ and

$$dy = R^4 |z|^{-4} dz,$$

so

$$H(x) = \int_{|z|>R} \rho \left(\frac{R^2 z}{|z|^2} \right) \ln \left(\frac{R}{|z|} |x - z| \right) R^4 |z|^{-4} dz.$$

Let

$$\sigma(z) = R^4 |z|^{-4} \rho \left(\frac{R^2 z}{|z|^2} \right)$$

and

$$C_H = \int_{|z|>R} \sigma(z) \ln \left(\frac{R}{|z|} \right) dz;$$

then

$$H(x) = C_H + \int_{|z|>R} \sigma(z) \ln(|x - z|) dz.$$

Note that

$$\|\sigma\|_\infty \leq \|\rho\|_\infty$$

and

$$\begin{aligned} |\partial_{z_k} \sigma(z)| &\leq R^4 (4|z|^{-5} \|\rho\|_\infty + 6R^2 |z|^{-6} \|\nabla \rho\|_\infty) \\ &\leq 4R^{-1} \|\rho\|_\infty + 6\|\nabla \rho\|_\infty. \end{aligned}$$

Proceeding as for P^0 yields

$$\begin{aligned} |\partial_{x_k x_\ell} H(x)| &\leq C(1 + \|\sigma\|_\infty)(1 + \ln(1 + \|\nabla \sigma\|_\infty)) \\ &\leq C(1 + \|\rho\|_\infty)(1 + \ln(1 + C\|\rho\|_\infty + C\|\nabla \rho\|_\infty)) \end{aligned}$$

and

$$|\partial_{x_k} H(x)| \leq C\|\rho\|_\infty.$$

Hence

$$|\partial_{x_k x_\ell} (P^0 - H)(x)| \leq C(1 + \|\rho\|_\infty)(1 + \ln(1 + C\|\rho\|_\infty + C\|\nabla \rho\|_\infty)) \tag{A.4}$$

and

$$|\partial_{x_k} (P^0 - H)(x)| \leq C\|\rho\|_\infty \tag{A.5}$$

for $x \in D$. By Theorem 6.14 of [2], $P^0 - H \in C^2(\bar{D})$, so (A.4) and (A.5) hold on \bar{D} . Lemma 2.5 now follows. □

REFERENCES

- [1] J. Batt, W. Faltenbacher, and E. Horst, *Stationary spherically symmetric models in stellar dynamics*, Arch. Rational Mech. Anal. **93** (1986), no. 2, 159–183, DOI 10.1007/BF00279958. MR823117 (87i:85001)
- [2] David Gilbarg and Neil S. Trudinger, *Elliptic partial differential equations of second order*, 2nd ed., Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], vol. 224, Springer-Verlag, Berlin, 1983. MR737190 (86c:35035)
- [3] Robert T. Glassey, *The Cauchy problem in kinetic theory*, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1996. MR1379589 (97i:82070)
- [4] Yan Guo and Gerhard Rein, *Isotropic steady states in galactic dynamics*, Comm. Math. Phys. **219** (2001), no. 3, 607–629, DOI 10.1007/s002200100434. MR1838751 (2002g:85001)
- [5] Yan Guo, *Singular solutions of the Vlasov-Maxwell system on a half line*, Arch. Rational Mech. Anal. **131** (1995), no. 3, 241–304, DOI 10.1007/BF00382888. MR1354697 (96h:35228)
- [6] Yan Guo, *Regularity for the Vlasov equations in a half-space*, Indiana Univ. Math. J. **43** (1994), no. 1, 255–320, DOI 10.1512/iumj.1994.43.43013. MR1275462 (95d:35178)
- [7] Hyung Ju Hwang, *Regularity for the Vlasov-Poisson system in a convex domain*, SIAM J. Math. Anal. **36** (2004), no. 1, 121–171 (electronic), DOI 10.1137/S0036141003422278. MR2083855 (2005f:35036)
- [8] P.-L. Lions and B. Perthame, *Propagation of moments and regularity for the 3-dimensional Vlasov-Poisson system* (English, with French summary), Invent. Math. **105** (1991), no. 2, 415–430, DOI 10.1007/BF01232273. MR1115549 (92e:35160)
- [9] K. Pfaffelmoser, *Global classical solutions of the Vlasov-Poisson system in three dimensions for general initial data*, J. Differential Equations **95** (1992), no. 2, 281–303, DOI 10.1016/0022-0396(92)90033-J. MR1165424 (93d:35170)
- [10] Gerhard Rein, *Collisionless kinetic equations from astrophysics—the Vlasov-Poisson system*, Handbook of differential equations: evolutionary equations. Vol. III, Handb. Differ. Equ., Elsevier/North-Holland, Amsterdam, 2007, pp. 383–476, DOI 10.1016/S1874-5717(07)80008-9. MR2549372 (2011b:85005)
- [11] Jack Schaeffer, *Steady states for a one-dimensional model of the solar wind*, Quart. Appl. Math. **59** (2001), no. 3, 507–528. MR1848532 (2002j:82114)
- [12] Jack Schaeffer, *Slow decay for a linearized model of the solar wind*, Quart. Appl. Math. **70** (2012), no. 1, 181–198, DOI 10.1090/S0033-569X-2011-01252-2. MR2920623
- [13] Jack Schaeffer, *Steady states of the Vlasov-Maxwell system*, Quart. Appl. Math. **63** (2005), no. 4, 619–643. MR2187923 (2006k:82148)
- [14] D. Tidman and N. Krall, *Shock waves in collisionless plasmas*, Wiley-Interscience, 1971.