

Some Properties of Ferromagnetoelectric Nickel-Iodine Boracite, Ni₃B₇O₁₃I

ASCHER, Edgar, *et al.*

Abstract

Ferroelectricity and weak ferromagnetism have been found to set in simultaneously in Ni₃B₇O₁₃I at ~64°K. This is evidenced by dielectric hysteresis, spontaneous Faraday effect, quadratic magnetoelectric hysteresis, etc. The strong coupling between the mutually perpendicular spontaneous polarization--[001]--and spontaneous magnetization--[110]--is such that, when the former is reversed, the latter turns by 90°. The magnetic point group is most probably m'm2'. Dielectric constant, magnetic, and magnetoelectric susceptibilities and magnetic coercive field are shown as a function of temperature.

Reference

ASCHER, Edgar, *et al.* Some Properties of Ferromagnetoelectric Nickel-Iodine Boracite, Ni₃B₇O₁₃I. *Journal of applied physics*, 1966, vol. 37, no. 3, p. 1404-1405

DOI : 10.1063/1.1708493

Available at:

<http://archive-ouverte.unige.ch/unige:30949>

Disclaimer: layout of this document may differ from the published version.

Some Properties of Ferromagnetoelectric Nickel-Iodine Boracite, $\text{Ni}_3\text{B}_7\text{O}_{13}\text{I}$

E. ASCHER, H. RIEDER, H. SCHMID, AND H. STÖSSEL

Battelle Memorial Institute, Geneva, Switzerland

Ferroelectricity and weak ferromagnetism have been found to set on simultaneously in $\text{Ni}_3\text{B}_7\text{O}_{13}\text{I}$ at about 64°K . This is evidenced by dielectric hysteresis, spontaneous Faraday effect, quadratic magnetoelectric hysteresis, etc. The strong coupling between the mutually perpendicular spontaneous polarization— $[001]$ —and spontaneous magnetization— $[110]$ —is such that, when the former is reversed, the latter turns by 90° . The magnetic point group is most probably $m'm2'$. Dielectric constant, magnetic, and magnetoelectric susceptibilities and magnetic coercive field are shown as a function of temperature.

IN this paper we report on the discovery of ferro-magnetoelectricity—simultaneous ferroelectricity and (weak) ferromagnetism—in $\text{Ni}_3\text{B}_7\text{O}_{13}\text{I}$. This compound is essentially isostructural with the mineral boracite, $\text{Mg}_3\text{B}_7\text{O}_{13}\text{Cl}$, that has a piezoelectric high-temperature phase (T_d^5) and a pyroelectric low-temperature phase (C_{2v}^5).^{1,2}

Ni-Cl-boracite was shown to be ferroelectric³; the other paramagnetic boracites were expected to behave similarly. Since 3-dimensionally linked metal-halogen-metal chains occur in the boracite structure, it was

maximum of the susceptibility at $\sim 120^\circ\text{K}$); there is furthermore a sharp minimum at 64°K and a sharp maximum at 60°K (Fig. 1).⁴ The dielectric constant—measured on a (100) platelet⁵—shows a very small peak at 60°K (Fig. 1). Observation in polarized light (transmission approximately between 5300 and 7000 Å) reveals a spontaneous Faraday effect, and hence a spontaneous magnetization M^s , below 64°K . The simultaneous onset of ferroelectricity is evidenced by dielectric hysteresis (Fig. 2). The displaced loop reveals

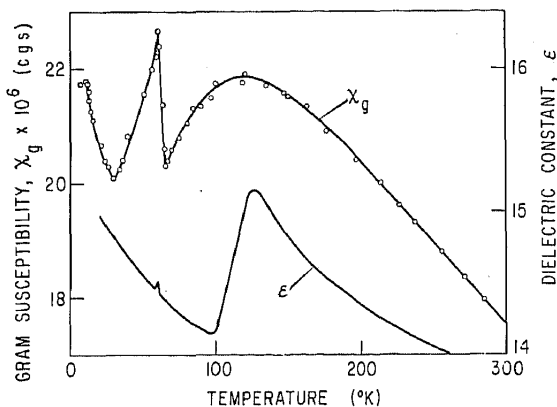


FIG. 1. Dielectric constant ϵ and gram susceptibility $\chi_g(H=2700 \text{ Oe})$ versus temperature.

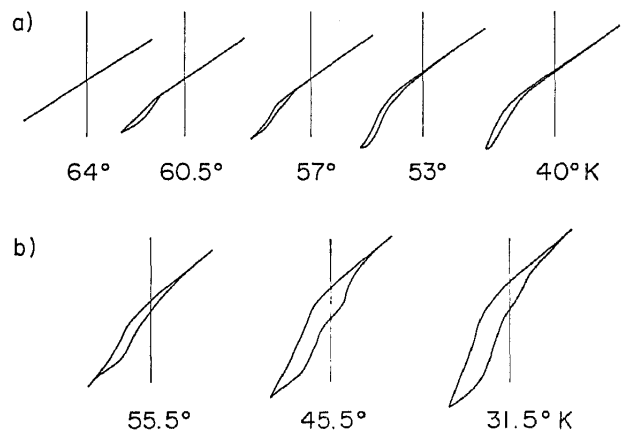


FIG. 2. Examples of ferroelectric hysteresis loops (a) 15 kV/cm; (b) 30 kV/cm, at 200 Hz.

anticipated that collective magnetic phenomena would occur at sufficiently low temperatures, and that there would be interactions with the electric polarization. Magnetic susceptibility measurements on powders of 3d-metal boracites indicate that these become anti-ferromagnetic at low temperature.⁴ $\text{Ni}_3\text{B}_7\text{O}_{13}\text{I}$ shows the highest magnetic ordering temperature (broad

a unidirectional anisotropy (due to growth condition) of the spontaneous polarization P^s . Optical observations and symmetry considerations show that the magnetic point group below 64°K is $m'm2'$, and that P^s is parallel to one of the cube edges, e.g. $[001]$; then M^s is either $[110]$ or $[\bar{1}\bar{1}0]$.⁶

(1) Electrical switching of P^s from $[001]$ to $[00\bar{1}]$ —by application of, e.g., 5 kV/cm at 56°K —results in a 90° change of M^s from $[110]$ to $[1\bar{1}0]$; a magnetic field of a few oersted in the $[100]$ direction prevents the new ferroelectric domain from splitting in several magnetic 180° domains. As a consequence of this cou-

¹ T. Ito, N. Morimoto and R. Sadanaga, *Acta Cryst.* **4**, 310 (1951).

² The low-temperature space group seems to be inadequate; see Ref. 3.

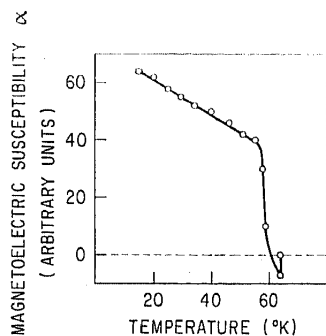
³ E. Ascher, H. Schmid, and D. Tar, *Solid State Commun.* **2**, 45 (1964).

⁴ H. Schmid, H. Rieder, and E. Ascher, *Solid State Commun.* **3**, 327 (1965).

⁵ H. Schmid, *J. Phys. Chem. Solids* **26**, 973 (1965).

⁶ Pseudocubic indices.

Fig. 3. Magnetoelectric susceptibility α versus temperature. Measured after cooling from above 65°K in $H=7.6$ kOe//[110] and $E=10$ kV/cm//[001].



pling, reversal of P^s in the plane of a (100) plate reverses the Faraday rotation for light propagating in the [100] direction from $+\rho$ to $-\rho$; reversal of P^s in the plane of a (110) plate changes the Faraday rotation for light propagating in the [110] direction from $\pm\rho$ to 0 (and vice versa).

(2) When an external magnetic field—e.g., 6 kOe at 56°K—is rotated from [110] to $[1\bar{1}0]$, the polarization P^s is reversed from [001] to $[00\bar{1}]$.

Denoting the direction of P^s by z , and that of M^s by

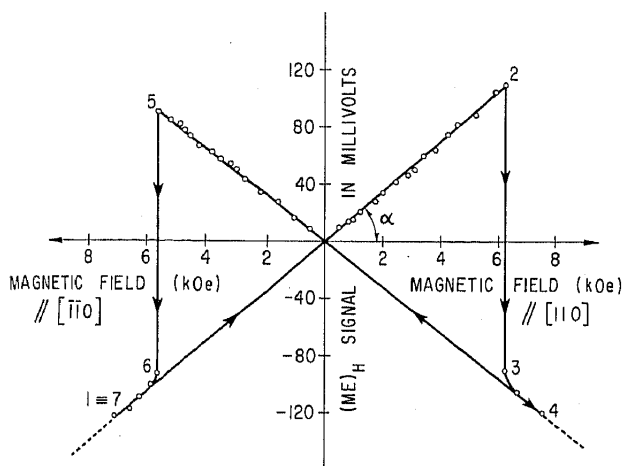
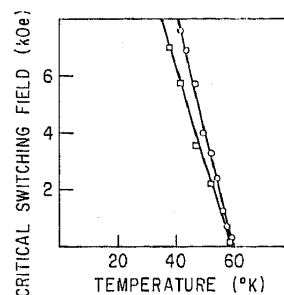


Fig. 4. Example of a quadratic magnetolectric hysteresis loop with H along $\pm[110]$ and P along $[001]$ at 46°K. After annealing as in Fig. 3.

y (point group $m_x'm_y'2_z'$), we obtain the following equations for the linear magnetolectric effects: $P_z = \alpha_{zy}H^y$, $M_y = \alpha_{zy}E^z$, $P_y = \alpha_{yz}H^z$, $M_z = \alpha_{yz}E^y$. We have measured the temperature dependence of α (as defined by the first equation) by a static method⁷ in a field annealed sample (Fig. 3). The value of α_{zy} at, e.g. 15°K, is 3.3×10^{-4} . The sign reversal of α_{zy} at about 60°K may be explained by the fact that, in about 1% of the sample, the magnetization assumes the direction of the applied field only close to the Curie point (this results in a change of polarization opposite to the magnetolectric effect). The presence of ferromagnetism gives rise to a quadratic magnetolectric

Fig. 5. Optically determined switching field versus temperature for the 180° magnetization reversal of two different ferromagnetic single domains of the same ferroelectric domain.



hysteresis loop (Fig. 4). The coercive field of these loops rises abruptly with falling temperature (Fig. 5) (as is typical for weak ferromagnets).

To summarize, it may be stated that $Ni_3B_7O_{13}I$ is a piezoelectric paramagnet above $\sim 120^\circ K$, a piezoelectric-antiferromagnet from 64° to $\sim 120^\circ K$, and a ferroelectric weak ferromagnet below $64^\circ K$.

ACKNOWLEDGMENTS

The authors wish to thank the Battelle Institute, Geneva, for support of this research. The help of H. Tippmann and F. Rademacher in some of the experiments is gratefully acknowledged.

⁷ G. T. Rado and V. J. Folen, Phys. Rev. Letters 7, 310 (1961).