

Engineering Notes

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Some Remarks on the Solution of the Lifting Line Equation

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Introduction

SO much has been said about Prandtl's lifting line equation that it seems futile to elaborate on it further. However, the author feels that a few salient features of this equation have not been remarked upon earlier. The equation is¹

$$\Gamma(y) = \pi UC(y)[\alpha(y) - \epsilon(y)] \quad (1)$$

where

$$\epsilon(y) = \frac{1}{4\pi U} \int_{-s}^s \frac{d\Gamma}{d\eta} \times \frac{d\eta}{(y-\eta)} \quad (2)$$

and

$$\Gamma(s) = \Gamma(-s) = 0 \quad (3)$$

Here Γ is the local circulation, α the local geometric angle of attack measured from the zero lift line, C the local chord, ϵ the local downwash induced by the trailing vortices, U the freestream velocity, s the wing semispan, and y the spanwise coordinate.

Equation (1) is usually solved using a sine series for Γ in the form

$$\Gamma(\theta) = \sum_{n=1}^{\infty} A_n \Gamma_n(\theta) = 2sU \sum_{n=1}^{\infty} A_n \sin n\theta \quad (4)$$

in a collocation method, and $\theta = \cos^{-1}(y/s)$.

Since Γ and α are related by a linear operator, it may be shown that

$$\alpha(\theta) = \sum_{n=1}^{\infty} A_n \alpha_n(\theta) = \frac{2s}{\pi C(y)} \sum_{n=1}^{\infty} A_n \sin n\theta + \epsilon(\theta) \quad (5)$$

and

$$\epsilon(\theta) = \sum_{n=1}^{\infty} A_n \epsilon_n(\theta) = \frac{1}{2} \sum_{n=1}^{\infty} A_n n \sin n\theta / \sin \theta \quad (6)$$

The lift and drag coefficients are, respectively,

$$C_L = \pi A R A_1 / 2 \quad (7)$$

$$C_D = \pi A R \sum_{n=1}^{\infty} n A_n^2 / 4 \quad (8)$$

where the aspect ratio $AR = 4s^2/S$, and S is the wing area.

We note that C_D contains only the square of the unknown constants A_n . In other words, the loadings (Γ_n , α_n) are orthogonal in the sense used by Graham.³ Orthogonality is usually associated with simplicity, and this is reflected in Eqs. (7) and (8).

Let us assume that a loading (Γ , α) is expressible by a finite number of terms m of Eqs. (4) and (5) to the desired accuracy. Then from the orthogonality of the loadings it may be easily shown that

$$\begin{aligned} \int_{-s}^s \Gamma(y) \epsilon_n(y) dy &= \int_{-s}^s \epsilon(y) \Gamma_n(y) dy \\ &= A_n \int_{-s}^s \Gamma_n(y) \epsilon_n(y) dy \\ &= A_n \frac{\pi}{2} s^2 U n \end{aligned} \quad (9)$$

for $n = 1, \dots, m$.

Hence if the circulation distribution is given, the unknown constants may be obtained from

$$A_n = 2 \int_{-s}^s \Gamma(y) \epsilon_n(y) dy / (\pi s^2 U n) \quad (10)$$

to yield $\alpha(y)$ from Eqs. (5) and (6). If the downwash distribution $\epsilon(y)$ is given, the constants may be obtained from

$$A_n = 2 \int_{-s}^s \epsilon(y) \Gamma_n(y) dy / (\pi s^2 U n) \quad (11)$$

to yield $\Gamma(y)$ from Eq. (4). However, more frequently the wing geometry is given in which case $\alpha(y)$ is specified. Using Eq. (5) to substitute for $\epsilon(y)$ in Eq. (11) results in m linear simultaneous equations

$$\begin{aligned} A_n \frac{\pi}{2} s^2 U n + (1/\pi U) \int_{-s}^s [\Gamma_n(y) \sum_{r=1}^m A_r \Gamma_r(y) / C(y)] dy \\ = \int_{-s}^s \alpha(y) \Gamma_n(y) dy \end{aligned} \quad (12)$$

which may be solved for the A_n . Results of Eqs. (10) and (11) are important since they allow the A_n to be evaluated explicitly. Equation (12) clearly shows that for elliptic planforms, i.e. $C(\theta) \sim \sin \theta$, the A_n can be obtained explicitly for any angle of attack distribution. In this form it is a generalization of Filotas' results.² It is further interesting to note that Eq. (12) is the same as would be obtained by an application of the Galerkin method⁴ to Eq. (1). The present derivation shows that the Galerkin approach is a very natural one and is the reason for its success in Ref. 4.

The advantage of using Eqs. (10-12) is that the solution does not depend only on the characteristics of the wing at a small number of isolated points as in a collocation procedure, but gives an approximate solution along the entire span. Hence discontinuities, flap deflections, etc., may be accounted for.

An Example—Rectangular Wing

Application of Eq. (12) to a wing of constant chord C , and constant incidence α , gives

$$\begin{aligned} \frac{\pi}{2} n A_n - \frac{16s}{\pi C} \sum_{r=1}^m \frac{n r A_r}{(r^2 - n^2)^2 - 2(r^2 + n^2) + 1} \\ = \pi \alpha ; \quad n = 1 \\ = 0 ; \quad n \neq 1 \end{aligned} \quad (13)$$

for $n = 1, \dots, m$;

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where Σ' denotes summation over only those terms for which $(n + r)$ is even. A one term approximation for Γ yields

$$A_1 = \pi\alpha / (\pi/2 + 16s/3\pi C) \quad (14)$$

A two term approximation shows

$$\begin{aligned} A_1 &= \pi\alpha / \left[\left(\pi/2 + \frac{16s}{3\pi C} \right) - \left(\frac{16s}{15\pi C} \right)^2 / \left(\frac{\pi}{2} - \frac{144s}{35\pi C} \right) \right] \\ A_2 &= 0 \\ A_3 &= \frac{16s\alpha}{15C} / \left[\left(\frac{\pi}{2} + \frac{16s}{3\pi C} \right) \left(\frac{\pi}{2} - \frac{144s}{35\pi C} \right) - \left(\frac{16s}{15\pi C} \right)^2 \right] \end{aligned} \quad (15)$$

The computation effort required in solving Eq. (13) is less compared to a collocation method where trigonometric functions must be evaluated.

Conclusions

The method outlined in this note is valid for any set of loadings (Γ_n, α_n) which are orthogonal in Graham's sense. It may be used for non-orthogonal loadings if they are first converted to an orthogonal set as suggested by Graham.³

References

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