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# Some results for Laplace-type integral operator in quantum calculus

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**Abstract**

In the present article, we wish to discuss  $q$ -analogues of Laplace-type integrals on diverse types of  $q$ -special functions involving Fox's  $H_q$ -functions. Some of the discussed functions are the  $q$ -Bessel functions of the first kind, the  $q$ -Bessel functions of the second kind, the  $q$ -Bessel functions of the third kind, and the  $q$ -Struve functions as well. Also, we obtain some associated results related to  $q$ -analogues of the Laplace-type integral on hyperbolic sine (cosine) functions and some others of exponential order type as an application to the given theory.

**Keywords:**  $J_\nu(x; q)$  function;  $Y_\nu(x; q)$  function;  $K_\nu(x; q)$  function;  $H_\nu(x; q)$  function; Laplace-type integral

**1 Introduction and preliminaries**

Quantum calculus is a version of calculus where derivatives are differences and antiderivatives are sums, and no further limits are required. The quantum calculus or  $q$ -calculus, compared to the differential and integral calculus, has been very recently named. Hence some rules and definitions need to be recalled. For  $0 < q < 1$ , the  $q$ -calculus starts with the definition of the  $q$ -analogue of the differential and the  $q$ -analogue of derivatives as well. The  $q$ -analogue of the integer  $n$ , the factorial of  $n$ , and the binomial coefficient are respectively given as

$$[n]_q = \frac{1 - q^n}{1 - q}, \quad ([n]_q)! = \begin{cases} \prod_{k=1}^n [k]_q, & n \in \mathbb{N} \\ 1, & n = 0 \end{cases}, \quad \begin{bmatrix} n \\ k \end{bmatrix}_q = \prod_{j=1}^k \frac{1 - q^{n-k+j}}{1 - q^j}. \quad (1)$$

The  $q$ -analogue of  $(x + a)^n$  ( $n \in \mathbb{N}$ ) and its  $q$ -derivative are respectively given as

$$(x + a)_q^n = \prod_{j=0}^{n-1} (x + q^j a), \quad D_q(x + a)_q^n = [n]_q (x + a)_q^{n-1}, \quad (x + a)_q^0 = 1. \quad (2)$$

The  $q$ -Jackson integrals from 0 to  $a$  and from  $a$  to  $b$  are given as follows (see [1], see also [2]):

$$\int_0^a f(x) d_q x = (1 - q)a \sum_0^\infty f(aq^k) q^k \quad (3)$$



and

$$\int_b^a f(x) d_q x = \int_0^b f(x) d_q x - \int_0^a f(x) d_q x. \tag{4}$$

The improper  $q$ -Jackson integral is given as follows (see [1]):

$$\int_0^{\frac{\infty}{A}} f(x) d_q x = (1 - q) \sum_{n \in \mathbb{Z}} \frac{q^n}{A} f\left(\frac{q^n}{A}\right), \quad A \in \mathbb{C}.$$

The  $q$ -analogues of the gamma function are defined by

$$\Gamma_q(\alpha) = \int_0^{\frac{1}{1-q}} x^{\alpha-1} E_q(q(1-q)x) d_q x$$

and

$${}_q\Gamma(\alpha) = K(A; \alpha) \int_0^{\frac{\infty}{A(1-q)}} x^{\alpha-1} e_q(-(1-q)x) d_q x,$$

where  $\alpha > 0$  and, for every  $t \in \mathbb{R}$ ,

$$K(A; t) = A^{t-1} \frac{(-q/A; q)_{\infty}}{(-q^t/A; q)_{\infty}} \frac{(-A; q)_{\infty}}{(-Aq^{1-t}; q)_{\infty}}.$$

Here

$$(a; q)_n = \prod_{k=0}^{n-1} (1 - aq^k), \quad (a; q)_{\infty} = \lim_{n \rightarrow \infty} (a; q)_n.$$

The very useful identities used in this article are (cf. [2])

$$\Gamma_q(x) = \frac{(q; q)_{\infty}}{(q^x; q)_{\infty}} (1 - q)^{1-x} \quad \text{and} \quad (a; q)_t = \frac{(a; q)_{\infty}}{(aq^t; q)_{\infty}}, \quad t \in \mathbb{R}.$$

The  $q$ -hypergeometric functions are represented by

$${}_r\phi_s \left( \begin{matrix} a_1, a_2, \dots, a_r \\ \alpha_1, \alpha_2, \dots, \alpha_s \end{matrix} \middle| q, z \right) = \sum_0^{\infty} \frac{(a_1, a_2, \dots, a_r; q)_n}{(\alpha_1, \alpha_2, \dots, \alpha_s; q)_n} \frac{z^n}{(q; q)_n}$$

and

$${}_{m-k}\Phi_{m-1} \left( \begin{matrix} a_1, a_2, \dots, a_{m-k} \\ \alpha_1, \alpha_2, \dots, \alpha_{m-1} \end{matrix} \middle| q, z \right) = \sum_0^{\infty} \frac{(a_1, \dots, a_{m-k}; q)_n}{(\alpha_1, \dots, \alpha_{m-1}; q)_n} [(-1)^n q^{\binom{n}{2}}]^k \times \frac{z^n}{(q; q)_n},$$

where  $(a_1, a_2, \dots, a_p; q)_n = \prod_{k=0}^p (a_k; q)_n$ .

### 2 H-Function and related functions

The  $H$ -function, which is an extension of the hypergeometric functions  ${}_pF_q$ , introduced by Fox [3] (see also [4, 5]), has found various applications in a huge range of problems associated with reaction, reaction diffusion, communication, engineering, fractional differential equations, integral equations, theoretical physics, and statistical distribution theory as well. The  $H$ -functions have also been recognized to play a fundamental role in fractional calculus with its applications. Fox’s  $H$ -function, admitting to a standard notation, is presented as

$$H_{p,q}^{m,n}(\eta) = \frac{1}{2\pi i} \int_P J_{p,q}^{m,n}(w) \eta^w dw, \tag{5}$$

where  $P$  is a suitable complex path,  $\eta^w = \exp\{w(\log |\eta| + i \arg \eta)\}$ ,  $J_{p,q}^{m,n}(w) = \frac{A(w)B(w)}{C(s)D(w)}$ , and

$$\begin{aligned} A(w) &= \prod_1^m \Gamma(b_j - \beta_j w), & B(w) &= \prod_1^n \Gamma(1 - a_j + \alpha_j w), \\ C(w) &= \prod_{m+1}^q \Gamma(1 - b_j - \beta_j w), & D(w) &= \prod_{n+1}^p \Gamma(a_j + \alpha_j w), \end{aligned}$$

$0 \leq n \leq p$ ,  $1 \leq m \leq q$ ,  $\{a_j, b_j\} \in \mathbb{C}$ ,  $\{\alpha_j, \beta_j\} \in \mathbb{R}^+$ . Let  $\alpha_j$  and  $\beta_j$  be positive integers and  $0 \leq m \leq N$ ;  $0 \leq n \leq M$ . Then the  $q$ -analogue of Fox’s  $H$ -function is given as (see [6])

$$\begin{aligned} &H_{M,N}^{m,n} \left( x; q \left| \begin{matrix} (a_1, \alpha_1), (a_2, \alpha_2), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), (b_2, \beta_2), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \\ &= \frac{1}{2\pi i} \int_C \frac{\prod_{j=1}^m G(q^{b_j - \beta_j s}) \prod_{j=1}^n G(q^{1 - a_j + \alpha_j s}) \pi x^s}{\prod_{j=m+1}^N G(q^{1 - b_j + \beta_j s}) \prod_{j=n+1}^M G(q^{a_j - \alpha_j s}) G(q^{1-s}) \sin \pi s} d_q s, \end{aligned}$$

where  $G$  is defined in terms of the product

$$G(q^\alpha) = \prod_{k=0}^\infty (1 - q^{\alpha - k})^{-1} = \frac{1}{(q^\alpha; q)_\infty}. \tag{6}$$

The contour  $C$  is parallel to  $\text{Re}(ws) = 0$ , such that all poles of  $G(q^{b_j - \beta_j s})$ ,  $1 \leq j \leq m$ , are its right and those of  $G(q^{1 - a_j + \alpha_j s})$ ,  $1 \leq j \leq n$ , are the left of  $C$ . The above integral converges if  $\text{Re}(s \log x - \log \sin \pi s) < 0$ , for huge values of  $|s|$  on  $C$ . Hence,

$$|\arg(x) - w_2 w_1^{-1} \log |x|| < \pi, \quad |q| < 1, \quad \log q = -w = -w_1 - iw_2,$$

where  $w_1$  and  $w_2$  are real numbers.

Indeed, for  $\alpha_i = \beta_j = 1$ , for all  $i, j$ , we write the  $q$ -analogue of Meijer’s  $G$ -function as

$$\begin{aligned} &G_{M,N}^{m,n} \left( x; q \left| \begin{matrix} a_1, a_1, \dots, a_M \\ b_1, b_2, \dots, b_N \end{matrix} \right. \right) \\ &= \frac{1}{2\pi i} \int_C \frac{\prod_{j=1}^m G(q^{b_j - s}) \prod_{j=1}^n G(q^{1 - a_j + s}) \pi x^2}{\prod_{j=m+1}^N G(q^{1 - b_j + s}) \prod_{j=n+1}^M G(q^{a_j - s}) G(q^{1-s}) \sin \pi s} d_q s, \end{aligned} \tag{7}$$

where  $0 \leq m \leq N$ ;  $0 \leq n \leq M$  and  $\text{Re}(s \log x - \log \sin \pi s) < 0$ .

Additionally, the  $q$ -analogues of the Bessel function  $J_\nu(x)$  of the first kind, the Bessel function of  $Y_\nu(x)$ , the Bessel function of the third kind  $K_\nu(x)$ , and Struve's function  $H_\nu(x)$  are, respectively, defined in terms of Fox's  $H_q$ -function by [7] as follows:

$$J_\nu(x; q) = \{G(a)\}^2 H_{0,3}^{1,0} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (\frac{\nu}{2}, 1), (-\frac{\nu}{2}, 1) \end{matrix} (1, 1) \right. \right), \tag{8}$$

$$Y_\nu(x; q) = \{G(a)\}^2 \times H_{1,4}^{2,0} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (-\frac{\nu-1}{2}, 1) \end{matrix} (\frac{\nu}{2}, 1), (-\frac{\nu}{2}, 1), (-\frac{\nu-1}{2}, 1) (1, 1) \right. \right), \tag{9}$$

$$K_\nu(x; q) = (1-q) H_{0,3}^{2,0} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (\frac{\nu}{2}, 1), (-\frac{\nu}{2}, 1) \end{matrix} (1, 1) \right. \right), \tag{10}$$

$$H_\nu(x; q) = \left( \frac{1-q}{2} \right)^{1-\alpha} \times H_{1,4}^{3,1} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (\frac{1+\alpha}{2}, 1) \end{matrix} (\frac{\nu}{2}, 1), (-\frac{\nu}{2}, 1), (\frac{\nu+\alpha}{2}, 1) (1, 1) \right. \right). \tag{11}$$

In [8] (see also [9]), some  $q$ -analogues of the natural exponential functions, sine functions, cosine functions, hyperbolic sine functions, and hyperbolic cosine functions are, respectively, given in terms of Fox's  $H$ -function as follows:

$$e_q(-x) = G(q) H_{0,2}^{1,0} \left( x(1-q); q \left| \begin{matrix} (0, 1) \end{matrix} (1, 1) \right. \right), \tag{12}$$

$$\sin_q(x) = \sqrt{\pi}(1-q)^{-\frac{1}{2}} \{G(q)\}^2 \times H_{0,3}^{1,0} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (\frac{1}{2}, 1) \end{matrix} (0, 1) (1, 1) \right. \right), \tag{13}$$

$$\cos_q(x) = \sqrt{\pi}(1-q)^{-\frac{1}{2}} \{G(q)\}^2 \times H_{0,3}^{1,0} \left( \frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (0, 1) \end{matrix} (\frac{1}{2}, 1) (1, 1) \right. \right), \tag{14}$$

$$\sinh_q(x) = \frac{\sqrt{\pi}}{i}(1-q)^{-\frac{1}{2}} \{G(q)\}^2 \times H_{0,3}^{1,0} \left( -\frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (\frac{1}{2}, 1) \end{matrix} (0, 1) (1, 1) \right. \right), \tag{15}$$

$$\cosh_q(x) = \sqrt{\pi}(1-q)^{-\frac{1}{2}} \{G(q)\}^2 \times H_{0,3}^{1,0} \left( -\frac{x^2(1-q)^2}{4}; q \left| \begin{matrix} (0, 1) \end{matrix} (\frac{1}{2}, 1) (1, 1) \right. \right). \tag{16}$$

On the other hand, some impressive integral transforms also have the corresponding  $q$ -analogues in the concept of  $q$ -calculus; they include the  $q$ -Laplace transforms [10], the  $q$ -Sumudu transforms [9, 11–13], the  $q$ -Wavelet transform [14], the  $q$ -Mellin transform [15],  $q$ - $E_{2,1}$ -transform [16],  $q$ -Mangontarum transforms [17, 18],  $q$ -natural transforms [19], and

so on. Recently, a number of authors have studied various image formulas for these  $q$ -integral transforms, associated with a variety of special functions. In this sequel, we aim to investigate the  $q$ -analogues of Laplace-type integrals on diverse types of  $q$ -special functions involving Fox's  $H_q$ -function.

### 3 $q$ -Laplace-type transforms for $H_q$ -function

A Laplace-type integral was introduced in [20, 21]. The  $q$ -analogues of the Laplace-type integral of the first kind were defined later by [22] as follows:

$$\begin{aligned} {}_qL_2(f(\xi); y) &= \frac{1}{1 - q^2} \int_0^{y^{-1}} \xi E_{q^2}(q^2 y^2 \xi^2) f(\xi) d\xi \\ &= \frac{(q^2; q^2)_\infty}{[2]_q y^2} \sum_{i=0}^\infty \frac{q^{2i}}{(q^2; q^2)_i} f(q^i y^{-1}), \end{aligned} \tag{17}$$

whereas the  $q$ -analogues of the Laplace-type integral of the second kind were defined by

$$\begin{aligned} {}_q\ell_2(f(\xi); y) &= \frac{1}{1 - q^2} \int_0^\infty \xi e_{q^2}(y^2 \xi^2) d_q \xi \\ &= \frac{1}{[2]_q (-y^2; q^2)_\infty} \sum_{i \in \mathbb{Z}} q^{2i} f(q^i) (-y^2; q^2)_i. \end{aligned} \tag{18}$$

For the sake of convenience, we establish some formulas for the  ${}_qL_2$  operator. A similar argument can give certain corresponding results for the operator  ${}_q\ell_2$ .

**Theorem 1** *Let  $\beta$  be a positive real number. Then*

$${}_qL_2(\xi^{2\beta-2}; y) = \frac{(q^2; q^2)_\infty}{[2]_q y^2 (q^\beta; q^2)_\infty}.$$

*Proof* By using (17), we have

$$\begin{aligned} {}_qL_2(\xi^{2\beta-2}; y) &= \frac{(q^2; q^2)_\infty}{[2]_q y^{2\beta}} \sum_{i=0}^\infty \frac{q^{2i}}{(q^2; q^2)_i} (q^i y^{-1})^{2\beta-2} \\ &= \frac{(q^2; q^2)_\infty}{[2]_q y^{2\beta}} \sum_{i=0}^\infty \frac{q^{2\beta i} y^{2\beta-2}}{(q^2; q^2)_i}. \end{aligned}$$

That is,

$${}_qL_2(\xi^{2\beta-2}; y) = \frac{(q^2; q^2)_\infty}{[2]_q y^2} \sum_{i=0}^\infty \frac{q^{2\beta i}}{(q^2; q^2)_i}. \tag{19}$$

By the fact that

$$e_q(z) = \sum_{i=0}^\infty \frac{z^i}{(q; q)_i},$$

we have

$$\begin{aligned} {}_qL_2(\xi^{2\beta-2}; y) &= \frac{(q^2; q^2)_\infty}{[2]_q y^2} e_{q^2}(q^\beta) \\ &= \frac{(q^2; q^2)_\infty}{[2]_q y^2} \frac{1}{(q^{2\beta}; q^2)_\infty}. \end{aligned}$$

This completes the establishment of the belief. □

**Theorem 2** *Let  $\lambda$  be a complex number. Then*

$$\begin{aligned} &{}_qL_2 \left( x^{2\lambda} H_{M,N}^{m,n} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \right) (y) \\ &= \frac{(q^2; q^2)_\infty}{y^{2\lambda+2} [2]_q} H_{M+1,N}^{m,n+1} \left( \frac{\gamma}{y^{2k}}, q^2 \left| \begin{matrix} (-\lambda, k), (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right), \end{aligned}$$

where  $0 \leq n \leq m$  and  $0 \leq m \leq N$  and  $\lambda$  is an arbitrary complex number.

*Proof* Let  $\lambda$  be a complex number. Then by (17) we obtain

$$\begin{aligned} &{}_qL_2 \left( x^{2\lambda} H_{M,N}^{m,n} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \right) (y) \\ &= \frac{1}{2\pi i} \int_c \frac{\prod_{j=1}^m G(q^{2b_j-2\beta_j z}) \prod_{j=1}^n G(q^{2-2a_j+2\alpha_j z}) \pi \gamma^z}{\prod_{j=m+1}^N G(q^{2-2b_j+2\beta_j z}) \prod_{j=n+1}^M G(q^{2a_j-2\alpha_j z}) G(q^{2-2z}) \sin \pi z} \\ &\quad \times {}_qL_2(x^{2\lambda+2kz})(y) d_q z. \end{aligned} \tag{20}$$

Let  $\beta = \lambda + kz + 1$ , then by Theorem 1 we have

$${}_qL_2(x^{2(\lambda+kz)})(y) = {}_qL_2(x^{2B-2})(y) = \frac{(q^2; q^2)_\infty}{[2]_q y^2 (q^{2(\lambda+kz+1)}; q^2)_\infty}. \tag{21}$$

By invoking (21) in (20), we get

$$\begin{aligned} &{}_qL_2 \left( x^{2\lambda} H_{M,N}^{m,n} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \right) (y) \\ &= \frac{1}{2\pi i} \int_c \frac{\prod_{j=1}^m G(q^{2b_j-2\beta_j z}) \prod_{j=1}^n G(q^{2-2a_j+2\alpha_j z}) \pi \gamma^z}{\prod_{j=m+1}^N G(q^{2-2b_j+2\beta_j z}) \prod_{j=n+1}^M G(q^{2a_j-2\alpha_j z}) G(q^{2-2z}) \sin \pi z} \\ &\quad \times \frac{(q^2; q^2)_\infty}{[2]_q y^2 (q^{2(\lambda+kz+1)}; q^2)_\infty} d_q z. \end{aligned} \tag{22}$$

By inserting the identity

$$G(q^{2\lambda+2kz+2}) = \frac{1}{(q^{2\lambda+2kz+2}; q^2)_\infty}$$

in (22) yields

$$\begin{aligned} & {}_qL_2 \left( x^{2\lambda} H_{M,N}^{m,n} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \right) (y) \\ &= \frac{(q^2; q^2)_\infty}{2\pi i y^{2\lambda+2} [2]_q} \int_c \frac{\prod_{j=1}^m G(q^{2b_j-2\beta_j z}) \prod_{j=1}^n G(q^{2-2a_j+2\alpha_j z})}{\prod_{j=m+1}^N G(q^{2-2b_j+2\beta_j z}) \prod_{j=n+1}^M G(q^{2a_j-2\alpha_j z})} \\ & \quad \times \frac{G(q^{1+\lambda+kz})}{G(q^{2(1-z)}) \sin \pi z} \pi \left( \frac{\gamma}{y^2 k} \right)^z d_q z. \end{aligned}$$

Now, on account of the definition of  $H_q$ -function, we may establish that

$$\begin{aligned} & {}_qL_2 \left( x^{2\lambda} H_{M,N}^{m,n} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (a_1, \alpha_1), \dots, (a_M, \alpha_M) \\ (b_1, \beta_1), \dots, (b_N, \beta_N) \end{matrix} \right. \right) \right) (y) \\ &= \frac{(q^2; q^2)_\infty}{y^{2\lambda+2} [2]_q} H_{N,M+1}^{n+1,m} \left( \gamma x^{2k}; q^2 \left| \begin{matrix} (1-b_1, \beta_1), \dots, (1-b_N, \beta_N) \\ (1+\lambda, k), (1-a_1, \alpha_1), \dots, (1-a_M, \alpha_M) \end{matrix} \right. \right), \end{aligned}$$

provided  $k < 0$ .

The proof is completed. □

#### 4 Applications to trigonometric and hyperbolic functions

In this part, we shall give certain natural relevance to the leading results.

**Theorem 3** *Let  $e_q$  be defined in terms of (12). Then*

$${}_qL_2(e_{q^2}(-x))(y) = \frac{G(q^2)(q^2; q^2)_\infty}{[2]_q y^2} H_{1,2}^{1,1} \left( \frac{1-q^2}{y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (0, 1), (1, 1) \end{matrix} \right. \right).$$

*Proof* By setting  $\lambda = 0$ ,  $\gamma = 1 - q^2$ , and  $k = 1$ , Theorem 3 immediately follows from Theorem 2. □

The demonstration of this theorem is finished.

**Theorem 4** *Let  $\sin_q$  be defined in terms of (13). Then we have*

$$\begin{aligned} & {}_qL_2(\sin_{q^2}(x))(y) = \frac{\sqrt{\pi}(1-q^2)^{-\frac{1}{2}} \{G(q^2)\}^2}{[2]_q y^4} (q^2; q^2)_\infty \\ & \quad \times H_{1,3}^{1,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (\frac{1}{2}, 1)(0, 1), (1, 1) \end{matrix} \right. \right). \end{aligned}$$

*Proof* The proof of this theorem indeed follows from substituting the values  $\lambda = 0$ ,  $k = 1$ , and  $\gamma = \frac{(1-q^2)^2}{4}$  and from multiplying by  $\sqrt{\pi}(1-q^2)^{-\frac{1}{2}} \{G(q^2)\}^2$ .

Hence, the proof is completed. □

**Theorem 5** Let  $\cos_q$  be defined in terms of (14). Then

$$\begin{aligned}
 {}_qL_2(\cos_{q^2}(x))(y) &= \frac{\sqrt{\pi}(1-q^2)^{-\frac{1}{2}} \{G(q^2)\}^2}{[2]_q y^4} (q^2; q^2)_\infty \\
 &\quad \times H_{1,3}^{1,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (0, 1), (\frac{1}{2}, 1), (1, 1) \end{matrix} \right. \right).
 \end{aligned}$$

*Proof* Proof follows from Theorem 2 for  $\lambda = 0, k = 1, \gamma = \frac{(1-q^2)^2}{4}$ .  
 The proof is completed. □

**Theorem 6** Let  $\sinh_q$  be defined in terms of (15). Then

$$\begin{aligned}
 {}_qL_2(\sinh_{q^2}(x))(y) &= \frac{\sqrt{\pi}(1-q^2)^{-\frac{1}{2}} \{G(q^2)\}^2}{i[2]_q y^4} (q^2; q^2)_\infty \\
 &\quad \times H_{1,3}^{1,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (\frac{1}{2}, 1), (0, 1), (1, 1) \end{matrix} \right. \right).
 \end{aligned}$$

*Proof* By using the special case,  $\lambda = 0, k = 1, \gamma = \frac{(1-q^2)^2}{4}$ .  
 The proof is completed. □

**Theorem 7** Let  $\cosh_q$  be defined in terms of (16). Then

$$\begin{aligned}
 {}_qL_2(\cosh_{q^2}(x))(y) &= \frac{\sqrt{\pi}(1-q^2)^{-\frac{1}{2}} \{G(q^2)\}^2}{[2]_q y^4} (q^2; q^2)_\infty \\
 &\quad \times H_{1,3}^{1,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (0, 1), (\frac{1}{2}, 1), (1, 1) \end{matrix} \right. \right).
 \end{aligned}$$

*Proof* The validation of this theorem is identical to that of the previous theorem. □

**Theorem 8** Let the Bessel function be defined in terms of (8). Then

$$\begin{aligned}
 {}_qL_2(J_\nu(x; q^2))(y) &= \frac{\{G(q^2)\}^2}{[2]_q y^4} (q^2; q^2)_\infty \\
 &\quad \times H_{1,3}^{1,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (\frac{\nu}{2}, 1), (\frac{-\nu}{2}, 1), (1, 1) \end{matrix} \right. \right).
 \end{aligned}$$

*Proof* By setting  $\lambda = 0, k = 1, \gamma = \frac{1-q^2}{4}$  and multiplying by  $\{G(q^2)\}^2$ , the result follows. □

**Theorem 9** Let the  $q$ -Bessel function of the second kind be defined in terms of (9)–(11).  
 Then

$$\begin{aligned}
 {}_qL_2(Y_\nu(x; q^2))(y) &= \frac{\{G(q^2)\}^2}{[2]_q y^4} (q^2; q^2)_\infty \\
 &\quad \times H_{2,4}^{2,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1), (\frac{-\nu}{2}, 1) \\ (\frac{\nu}{2}, 1), (\frac{-\nu}{2}, 1), (\frac{-\nu-1}{2}, 1), (1, 1) \end{matrix} \right. \right),
 \end{aligned}$$



$$\begin{aligned}
{}_qL_2(K_\nu(x; q^2))(y) &= \frac{(1-q^2)}{[2]_q y^4} (q^2; q^2)_\infty \\
&\quad \times H_{1,3}^{2,1} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1) \\ (\frac{\nu}{2}, 1), (\frac{-\nu}{2}, 1), (1, 1) \end{matrix} \right. \right), \\
{}_qL_2(H_\nu(x; q^2))(y) &= \frac{(1-q^2)^{1-\alpha}}{2^{1-\alpha} [2]_q y^4} (q^2; q^2)_\infty \\
&\quad \times H_{2,4}^{3,2} \left( \frac{(1-q^2)^2}{4y^2}; q^2 \left| \begin{matrix} (0, 1), (\frac{1-\alpha}{2}, 1) \\ (\frac{\nu}{2}, 1), (\frac{-\nu}{2}, 1), (\frac{1+\alpha}{2}, 1), (1, 1) \end{matrix} \right. \right).
\end{aligned}$$

**Proof** Proof of this theorem follows from (9)–(11) and the technique quite similar to that of Theorems 3–8. We omit the details.  $\square$

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#### Competing interests

The authors declare that they have no competing interests.

#### Authors' contributions

The authors contributed equally and significantly in writing this paper. All authors read and approved the final manuscript.

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