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SPECTROSCOPIC CHARACTERIZATION OF HD 95086 b WITH THE GEMINI PLANET IMAGER

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# ABSTRACT

We present new H (1.5–1.8  $\mu$ m) photometric and K<sub>1</sub> (1.9–2.2  $\mu$ m) spectroscopic observations of the young exoplanet HD 95086 b obtained with the Gemini Planet Imager. The H-band magnitude has been significantly improved relative to previous measurements, whereas the low-resolution  $K_1$  ( $\lambda/\delta\lambda \approx 66$ ) spectrum is featureless within the measurement uncertainties and presents a monotonically increasing pseudo-continuum consistent with a cloudy atmosphere. By combining these new measurements with literature L' photometry, we compare the spectral energy distribution (SED) of the planet to other young planetary-mass companions, field brown dwarfs, and to the predictions of grids of model atmospheres. HD 95086 b is over a magnitude redder in  $K_1 - L'$  color than 2MASS J12073346–3932539 b and HR 8799 c and d, despite having a similar L' magnitude. Considering only the nearinfrared measurements, HD 95086 b is most analogous to the brown dwarfs 2MASS J2244316+204343 and 2MASS J21481633+4003594, both of which are thought to have dusty atmospheres. Morphologically, the SED of HD 95086 b is best fit by low temperature ( $T_{\rm eff} = 800-1300$  K), low surface gravity spectra from models which simulate high photospheric dust content. This range of effective temperatures is consistent with field L/T transition objects, but the spectral type of HD 95086 b is poorly constrained between early L and late T due to its unusual position the color-magnitude diagram, demonstrating the difficulty in spectral typing young, low surface gravity substellar objects. As one of the reddest such objects, HD 95086 b represents an important empirical benchmark against which our current understanding of the atmospheric properties of young extrasolar planets can be tested.

Key words: infrared: planetary systems – instrumentation: adaptive optics – planets and satellites: atmospheres – stars: individual (HD 95086)

# 1. INTRODUCTION

Directly imaged exoplanets form an important subset of the exoplanet population for which it is possible to characterize atmospheric properties from the optical and thermal-infrared emissions from the planet. Given the steep, monotonic decline in planet brightness with time (e.g., Burrows et al. 1997), the currently known, directly imaged exoplanets are companions to

young stars (<150 Myr), enabling investigations into their early evolution (Chauvin et al. 2004; Lafrenière et al. 2008; Marois et al. 2008; Lagrange et al. 2009; Delorme et al. 2013; Bailey et al. 2014; Kraus et al. 2014; Naud et al. 2014; Gauza et al. 2015; Macintosh et al. 2015). While these young exoplanets share many of the same atmospheric properties as isolated brown dwarfs recently identified in young moving groups (e.g., Liu et al. 2013; Gagné et al. 2015), differences in the elemental abundances of their atmospheres may provide evidence as to their formation mechanism (e.g., Öberg et al. 2011; Konopacky et al. 2013). Atmospheres of young, directly imaged planets also represent an important comparison to the transmission (e.g., Sing et al. 2011; Deming et al. 2013) and emission (e.g., Knutson et al. 2008; Bean et al. 2013) spectra of intensely irradiated planets transiting older stars.

Multi-wavelength infrared photometry of imaged planet atmospheres (e.g., Chauvin et al. 2004; Marois et al. 2008; Currie et al. 2013; Skemer et al. 2014) has defined a sequence that is under-luminous and redder than field brown dwarfs of the same spectral type. Low-resolution spectra of planetary mass companions (e.g., Janson et al. 2010; Barman et al. 2011b; Patience et al. 2012; Bonnefoy et al. 2014b) exhibit signatures of low surface gravity and spectral features indicating the effects of clouds and non-equilibrium chemistry (e.g., Hinz et al. 2010; Barman et al. 2011a; Skemer et al. 2012; Bonnefoy et al. 2014b). Altogether, these properties are difficult to reproduce with state-of-the-art atmospheric models, which have difficulty explaining the discrepancies seen with respect to field L- and T-dwarfs.

The Scorpius-Centaurus association (Sco-Cen) is the nearest OB association (de Zeeuw et al. 1999), and has undergone multiple and complex phases of massive star formation, resulting in a large sample of young (5-20 Myr) intermediate-mass  $(1.2-2.0 M_{\odot})$  stars and (Mamajek et al. 2002; Rizzuto et al. 2011; Pecaut et al. 2012; Song et al. 2012). Its well-constrained young age, proximity (90-150 pc), and the presence of near-infrared excess around a large number of its stars (Chen et al. 2012)—tracing leftover material from planet formation-make Sco-Cen ideal for giant planet searches with direct imaging. Thus far, five low-mass substellar companions ( $M < 25 M_{Jup}$ ) have been detected around Sco-Cen stars: 1RXS J160929.1-210524 b (Lafrenière et al. 2008), GSC 06214-00210 b (Ireland et al. 2011), HIP 78530 B (Lafrenière et al. 2011), HD 95086 b (Rameau et al. 2013a), and HD 106906 b (Bailey et al. 2014).

The star HD 95086 recently became a prime target to investigate the early phase of giant planet evolution. It is an A8, 1.7  $M_{\odot}$  member of the Sco-Cen subgroup, Lower Centaurus Crux (de Zeeuw et al. 1999; Madsen et al. 2002), with an age of  $17 \pm 4$  Myr (Meshkat et al. 2013) and at distance of 90.4  $\pm$ 3.4 pc (van Leeuwen 2007). HD 95086 also harbors a massive debris disk, first traced by a large infrared excess (Chen et al. 2012; Rizzuto et al. 2012) and then resolved with Herschel/ PACS (Moór et al. 2013). Modelling the spectral energy distribution (SED) and images of the system suggests a threecomponent debris structure, like that of HR 8799 (Su et al. 2015). With VLT/NaCo, Rameau et al. (2013a, 2013b) discovered and confirmed the presence of a co-moving giant planet, HD 95086 b, at a projected separation of 56 au. This planet may be responsible for sculpting the outer edge of the debris gap (Rameau et al. 2013a; Su et al. 2015). This discovery added another system to the relatively small population of young, nearby, intermediate-mass stars hosting both debris disks and directly imaged giant planets.

The observed L' (3.8  $\mu$ m) luminosity of HD 95086 b indicated a "hot-start" model-dependent mass of 5  $\pm$ 2  $M_{Jup}$  (Rameau et al. 2013a). "Warm-start" models predicted masses from 4 to 14  $M_{Jup}$  for a wide range of initial entropy values (8–13  $k_{\rm B}$ /baryon; Galicher et al. 2014). Follow-up observations with the Gemini Planet Imager at H (1.65  $\mu$ m) and  $K_1$  (2.06  $\mu$ m) enabled initial constraints of its atmospheric properties (Galicher et al. 2014); the very red colors of  $H - L' = 3.6 \pm 1.0$  mag and  $K_1 - L' = 2.7 \pm 0.70$  mag suggested the atmosphere of HD 95086 b contains a high amount of photospheric dust, placing the planet at the L/T transition. Empirical comparisons of the position of HD 95086 b on the color-magnitude diagram (CMD) showed that it is redder than dusty late field L-dwarfs (Stephens et al. 2009) and closer to, but still redder than, the young imaged planets HR 8799 cde (Marois et al. 2008, 2010) and 2M1207 b (Chauvin et al. 2004). Therefore, clouds may play an important role in the atmosphere of HD 95086 b. Finally, the predictions of the atmospheric properties of HD 95086 b from BT-SETTL (Allard et al. 2012) and LESIA (Baudino et al. 2015) atmospheric model fitting to its SED yielded  $T_{\text{eff}} = 600-1500 \text{ K}$  and  $\log g = 2.1-4.5 \text{ (cgs)}$ . Low surface gravity seemed to be favored, which, combined with the enhanced dust content, might inhibit the formation of methane expected in this temperature range due to nonequilibrium CO/CH<sub>4</sub> chemistry (Barman et al. 2011a; Zahnle & Marley 2014).

In this paper, we present the first  $K_1$  spectrum and revised H photometry of HD 95086 b, obtained with the Gemini Planet Imager (GPI; Macintosh et al. 2014), which we synthesize with existing L' photometry to constrain the SED of the planet (Section 2). In Section 3, we compare the SED of HD 95086 b to that of other substellar companions in Sco–Cen, other benchmark planetary-mass companions, and field brown dwarfs. We also fit the full SED to several grids of model atmospheric spectra in Section 4 in an attempt to further constrain the atmospheric and physical properties of the planet. We also took advantage of this study to derive color transformations between the 2MASS, MKO, and GPI photometric systems as a function of spectral type.

#### 2. OBSERVATIONS AND DATA REDUCTION

# 2.1. Observing Sequence

Integral field spectrograph observations of HD 95086 were made with GPI on 2015 April 8 UT at K<sub>1</sub> and 2016 February 29 UT at H under program IDs GS-2015A-Q-501 and GS-2015B-Q-500. The H-band filter within GPI is similar to the MKO H filter (Tokunaga et al. 2002), while the  $K_1$  filter spans the blue half of the K-band atmospheric window. The characterization of the  $H, K_1$ , and  $K_2$  GPI filters is described in the Appendix. The observations of HD 95086 were timed such that the star transited the meridian during the observing sequence, maximizing the field rotation for Angular Differential Imaging (ADI; Marois et al. 2006) observations. Wavelength calibration measurements using an argon arc lamp were taken immediately following each sequence. At  $K_1$ , 42 × 119.3 s exposures were obtained over the course of two hours, achieving a field rotation of 39°. Six sky measurements offset by 20" from the position of HD 95086 were also taken, three during the middle of the sequence and three at the end, to characterize the thermal emission from the sky, the telescope, and the instrument itself. Observing conditions were good, with an average DIMM seeing of 0."74, MASS coherence time of 1.28 ms, and wind speed of  $0.93 \text{ m s}^{-1}$ . Of the 42 exposures, four were rejected due to poor image quality, and one due to the instrument shutter failing to open. At H, 37  $\times$  59.6 s exposures were obtained over the course of an hour for a total rotation of 16°.9. The observing conditions were excellent, with an average

DIMM seeing of 0."94, MASS coherence time of 4.7 ms and wind speed of  $5.09 \text{ m s}^{-1}$ .

#### 2.2. Initial Reduction Steps

The observations of HD 95086 were reduced using the standard primitives contained within the GPI Data Reduction Pipeline<sup>25</sup> (DRP; Perrin et al. 2014) v1.3.0. The reduction process consisted of seven steps: (1) dark current subtraction; (2) bad pixel identification and interpolation; (3) compensation for instrument flexure between reference arcs taken several months previous, and the arcs taken during the observing sequence (Wolff et al. 2014); (4) microspectra extraction to create an  $(x, y, \lambda)$  data cube (Maire et al. 2014); (5) division by a flat field to correct for lenslet throughput; (6) interpolation to a common wavelength scale; and (7) correction for geometric distortion of the image, as measured with a calibrated pinhole mask in the laboratory (Konopacky et al. 2014). The interpolation step oversamples the final spectra, which have a spectral resolving power of  $\lambda/\delta\lambda \approx 66$  at  $K_1$ . This process was repeated for the sky measurements taken during the observing sequence at  $K_1$ , an average of which was subtracted from the reduced science data cubes. Images were then registered to a common center using the barycenter of the four satellite spots, which are attenuated replicas of the central point-spread function (PSF) generated by a pupil plane diffraction grating, and for which the positions are measured by the GPI DRP (Wang et al. 2014).

#### 2.3. Speckle Noise Minimization

The reduced data cubes from each observing sequence were further processed through two independent pipelines using two speckle subtraction algorithms: the Karhunen–Loève Image Projection algorithm (KLIP; Soummer et al. 2012), and the Locally Optimized Combination of Images algorithm (LOCI; Lafrenière et al. 2007). The speckle noise was minimized in each wavelength channel to create a residual spectral cube.

The KLIP-based pipeline processed the reduced data using pyKLIP<sup>26</sup> (Wang et al. 2015), an open-source Python implementation of the KLIP algorithm. First, the individual wavelength slices from each reduced data cube were high-pass filtered in the Fourier domain to remove the low spatial frequency seeing halo. For the  $K_1$  observations, the pipeline only utilized ADI due to the relatively small wavelength range of the  $K_1$  filter and in an attempt to avoid biases inherent to Spectral Differential Imaging (Marois et al. 2000). PSF subtraction was performed over six annuli of width 13 pixels (0.184), starting from an inner separation of 11 pixels (0.156). Each annulus was subdivided into four 90° arcs. A PSF constructed from the projection of the individual images onto the first five Karhunen-Loève modes was subtracted from each region. The PSF-subtracted images were de-rotated and averaged together. HD 95086 b was detected at a signal-tonoise ratio (S/N) of ~5 and ~7 in the collapsed H and  $K_1$ broadband images, which are shown in Figure 1.

For the LOCI-based pipeline, the low spatial frequency signal was removed from each image by subtracting a medianfiltered version of the image, using a 15 pixel wide box. The speckle subtraction was then performed using LOCI with subtraction annuli of width five pixels, optimization regions of  $N_a = 300$  full width at half maximum (FWHM,  $\simeq 3.5/3.9$  px at  $H/K_1$ ) with a geometry factor of g = 1, and a separation criteria of  $N_{\delta} = 0.75$  FWHM. In each case, the parameters are as defined in Lafrenière et al. (2007). Each residual image was then derotated, mean combined with a trimmed mean (10%), and aligned north up to generate each final residual slice. The final residual images within each filter are shown in Figure 1. At  $K_1$ , HD 95086 b is detected at an S/N of 3–5 in individual channels and at an S/N of ~7.5 in the broadband image. At H, HD 95086 b was recovered in the broadband image with a similar S/N. The astrometry of HD 95086 b will be discussed in a seperate study (Rameau et al. 2016).

# 2.4. Spectrophotometry Extraction

The  $K_1$  spectrum of HD 95086 b was extracted from the final PSF-subtracted image of each pipeline using two different techniques. For the KLIP-based pipeline, the spectrum was extracted from each wavelength slice using a  $1.5\lambda/D$  aperture, with a center based on the position of the planet in the wavelength-collapsed image. Uncertainties were estimated by measuring the noise within an annulus of width  $1.5\lambda/D$  and radius equivalent to the star to planet separation. In order to account for self-subtraction within the KLIP algorithm, the throughput was calculated by injecting a copy of the average of the four satellite spots, scaled to a similar brightness to HD 90586 b, into the data cubes prior to PSF subtraction. This copy was injected at the same separation as HD 95086 b. The throughput was then estimated at each wavelength by comparing the flux of the injected source to the flux measured in the PSF-subtracted image. This process was repeated 25 times, with the position angle of the injected source rotated for each trial. For each wavelength, the average throughput was calculated from these trials, and the standard deviation was used as the uncertainty on the throughput. To correct for throughput, the flux of the planet at a given wavelength was scaled by the average throughput, and the uncertainties on the flux and throughput were combined in quadrature. The flux from the central star was estimated by performing aperture photometry on an average of the four satellite spots within each wavelength slice prior to PSF subtraction, multiplied by the satellite spot to star flux ratio measured in the laboratory to be  $\Delta m = 8.92 \pm 0.06 \text{ mag}$  (Maire et al. 2014). The contrast between star and planet was then calculated as the ratio of their fluxes.

The LOCI-based pipeline first derived the position and contrast of HD 95086 b from the broadband image-created by averaging the 37 wavelength channels of the final PSFsubtracted data cube-using the fake planet injection technique (e.g., Lagrange et al. 2010). A PSF was generated from the average of the four satellite spots prior to PSF subtraction, integrated over the band and averaged over the temporal sequence, which was scaled and subtracted from the data at the estimated flux and position of the planet before performing the LOCI-based PSF subtraction. The flux and the position of the planet were iterated using an amoeba-simplex optimization algorithm (Nelder & Mead 1965) until the sum of the squared pixel intensities in an arc of size  $2 \times 4$  FWHM was minimized. Uncertainties on the position and contrast were estimated from the root mean square of fake planets injected with the measured contrasts and separations of HD 95086 b at ten different position angles uniformly distributed between 90° and 270°

<sup>&</sup>lt;sup>25</sup> http://docs.planetimager.org/pipeline

<sup>&</sup>lt;sup>26</sup> https://bitbucket.org/pyKLIP/pyklip



**Figure 1.** S/N map of the HD 95086 observations taken at *H* (left column) and  $K_1$  (right column). Each dataset was processed with both the LOCI-based (top row) and pyKLIP (bottom row) pipelines. The S/N was calculated using a  $1.5\lambda/D$  aperture. HD 95086 b was detected at a S/N of ~5–7 at *H* and ~7 at  $K_1$ . The central black mask conceals the region within which PSF subtraction is not performed, and the position of the central star position is designated by the star symbol.

away from the planet. These injected planets were then characterized using the same strategy as for HD 95086 b (e.g., Lagrange et al. 2010; Marois et al. 2010; Macintosh et al. 2015). From the  $K_1$  observations, a revised contrast of  $\Delta K_1 = 12.20 \pm 0.20$  was measured for HD 95086 b, the uncertainty on which was calculated by combining in quadrature the star-to-spot ratio uncertainty (0.07 mag), the PSF variability (0.07 mag), and the measurement error (0.18 mag). The *H*-band contrast was similarly measured as  $\Delta H = 13.70 \pm 0.25$ , the uncertainty on which combines the star-to-spot ratio uncertainty (0.06 mag), PSF variability (0.05 mag), and measurement error (0.24 mag).

Because the S/N of HD 95086 b is lower within the individual slices of the PSF-subtracted  $K_1$  data cube, the contrast at each wavelength—and their associated errors—were extracted using the same technique as described above, except that the position of the planet was fixed to that measured within the broadband image. Because LOCI biases due to oversubtraction can be difficult to calibrate, we also extracted the spectrum using the cADI algorithm (Marois et al. 2006). While

the contrasts measured using LOCI and cADI were very similar, HD 95086 b was recovered in fewer wavelength channels in the cADI reduction. Although the  $K_1$  spectrum of HD 95086 b extracted from the pyKLIP and LOCI-based pipelines agreed within the uncertainties (Figure 2), the spectrum extracted from the LOCI-based pipeline was used for the remainder of this study due to the smaller point-to-point scatter. The adopted contrasts measured within the  $K_1$  bandpass between the star and planet are given in Table 1.

The  $K_1$  spectrum of HD 95086 b was then calculated by multiplying the contrasts by the predicted flux of HD 95086 at each wavelength. As no flux-calibrated  $K_1$  spectrum of HD 95086 exists, a model stellar atmosphere from the BT-NEXTGEN grid<sup>27</sup> was used (Allard et al. 2012). The effective temperature, surface gravity, and metallicity of HD 95086 were fit using optical broadband (Høg et al. 2000) and intermediate-band (Rufener 1988) photometry, and broadband near-infrared (Skrutskie et al. 2006) and thermal-infrared (Cutri 2014)

<sup>&</sup>lt;sup>27</sup> https://phoenix.ens-lyon.fr/Grids/BT-NextGen/SPECTRA/



**Figure 2.** GPI  $K_1$  spectrum of HD 95086 b extracted from the LOCI-based (filled black circles) and pyKLIP (open blue circles) pipelines. For clarity the pyKLIP spectrum has been slightly offset in the wavelength direction. The first eight channels (1.886–1.946  $\mu$ m) were excluded due to the poor quality of the extracted spectrum at those wavelengths. The covariance between neighboring spectral channels in the adopted spectrum is described in Section 2.5.

photometry. By performing linear interpolation of the predicted flux for each filter between the model grid points, the best fit  $T_{\rm eff} = 7633 \pm 24$ atmosphere was with Κ, found  $\log g = 3.93 \pm 0.05$ ,  $[M/H] = -0.35 \pm 0.01$ , and a dilution factor of  $R^2/d^2 = 10^{-18.8 \pm 0.01}$ , where R and d are the radius of and distance to HD 95086. These values are consistent with previous literature estimates (e.g., Allende Prieto & Lambert 1999; Wright et al. 2003; Bonneau et al. 2006; McDonald et al. 2012). The apparent  $K_1$  spectrum of HD 95086 was estimated from the best fitting interpolated model and is given in Table 1. The GPI H and  $K_1$ , and NaCo L' apparent magnitudes of HD 95086 were also estimated from the same model atmosphere using the filter transmission profiles discussed in the Appendix and are given in Table 2. The final  $K_1$  spectrum of HD 95086 b is plotted in Figure 2 and is given in Table 1, in addition to the apparent fluxes calculated from the revised H-band contrast and the L' contrast from Galicher et al. (2014).

#### 2.5. Spectral Covariances

To properly assess the uncertainties on any fit to the  $K_1$  spectrum of HD 95086 b, the covariances between neighboring wavelength channels caused by both residual speckle noise in the final PSF-subtracted image, and the oversampling of the microspectra in the extraction process, must be taken into account. Following the method described in Greco & Brandt (2016), the correlation  $\psi_{ij}$  between pixel values at wavelengths  $\lambda_i$  and  $\lambda_j$  within an annulus of width 1.5  $\lambda/D$  at a given separation was estimated as

$$\psi_{ij} = \frac{\langle I_i I_j \rangle}{\sqrt{\langle I_i^2 \rangle \langle I_j^2 \rangle}} \tag{1}$$

where  $\langle I_i \rangle$  is the average intensity within the annulus at wavelength  $\lambda_i$ . This was repeated for all pairs of wavelengths

 Table 1

 Contrast between and Spectral Energy Distribution of HD 95086 and HD 95086 b, with Corresponding 68% Confidence Intervals

$\overline{\lambda_{\text{eff}}}$	Contrast	$F_{\lambda}(A) \times 10^{-13}$	$F_{\lambda}(b) \times 10^{-18}$
(µm)	$(10^{-6})$	(W m <sup>-</sup>	$^{2} \mu m^{-1}$ )
1.6330	3. $31^{+0.86}_{-0.68}$	$21.79\pm0.23$	7. $22^{+1.87}_{-1.49}$
÷	:	:	:
1.9548	$4.78\pm4.30$	$11.54\pm0.13$	$5.51\pm4.96$
1.9634	$4.65\pm3.53$	$11.57\pm0.13$	$5.38\pm4.08$
1.9719	$6.99\pm3.22$	$11.54\pm0.12$	$8.07\pm3.72$
1.9805	$5.66\pm2.77$	$11.44 \pm 0.12$	$6.47\pm3.17$
1.9891	$8.32\pm2.61$	$11.29\pm0.12$	$9.38\pm2.95$
1.9977	$8.63 \pm 2.44$	$11.12\pm0.12$	$9.59\pm2.71$
2.0063	$8.45\pm2.60$	$10.95\pm0.12$	$9.25\pm2.85$
2.0149	$9.99 \pm 3.10$	$10.77 \pm 0.11$	$10.77 \pm 3.34$
2.0234	$9.84 \pm 3.35$	$10.60 \pm 0.11$	$10.43\pm3.55$
2.0320	$10.78\pm2.66$	$10.43\pm0.11$	$11.24 \pm 2.78$
2.0406	$10.72 \pm 2.40$	$10.27\pm0.11$	$11.01 \pm 2.46$
2.0492	$9.51 \pm 2.30$	$10.11\pm0.11$	$9.61\pm2.33$
2.0578	$9.98\pm2.59$	$9.95 \pm 0.11$	$9.93\pm2.58$
2.0664	$10.73\pm2.70$	$9.79\pm0.10$	$10.51\pm2.65$
2.0749	$9.55\pm2.46$	$9.64 \pm 0.10$	$9.21 \pm 2.37$
2.0835	$9.92\pm2.43$	$9.50\pm0.10$	$9.42\pm2.31$
2.0921	$10.83\pm2.63$	$9.35 \pm 0.10$	$10.13\pm2.46$
2.1007	$10.34 \pm 2.71$	$9.20\pm0.10$	$9.51 \pm 2.49$
2.1093	$13.58\pm2.73$	$9.06\pm0.10$	$12.30\pm2.47$
2.1179	$12.81\pm2.84$	$8.91\pm0.10$	$11.41\pm2.53$
2.1264	$10.88 \pm 2.79$	$8.77\pm0.09$	$9.53\pm2.44$
2.1350	$13.88 \pm 2.79$	$8.61\pm0.09$	$11.95\pm2.40$
2.1436	$14.69\pm3.45$	$8.41\pm0.09$	$12.36\pm2.90$
2.1522	$14.53\pm3.69$	$8.17\pm0.09$	$11.86\pm3.01$
2.1608	$18.99\pm4.19$	$7.92 \pm 0.09$	$15.04\pm3.32$
2.1694	$18.36\pm5.17$	$7.78\pm0.09$	$14.29\pm4.02$
2.1779	$16.76\pm5.82$	$7.77\pm0.09$	$13.02\pm4.52$
2.1865	$17.42\pm 6.58$	$7.78\pm0.08$	$13.56\pm5.13$
2.1951	$17.02\pm9.91$	$7.76\pm0.08$	$13.20\pm7.68$
:	÷	:	
3.7697	$161.42\substack{+30.88\\-25.91}$	$0.99\pm0.01$	$15.95^{+3.05}_{-2.56}$

and at separations of  $\rho = 400$ , 500, 620 (the separation of HD 95086 b), 700, 800, 900, and 1000 mas. The companion was masked in the 620 mas annulus. Greco & Brandt (2016) parametrize the correlation  $\psi_{ii}$  into three terms as

$$\psi_{ij} \approx A_{\rho} \exp\left[-\frac{1}{2}\left(\frac{\rho}{\sigma_{\rho}}\frac{\lambda_{i}-\lambda_{j}}{\lambda_{c}}\right)^{2}\right] + A_{\lambda} \exp\left[-\frac{1}{2}\left(\frac{1}{\sigma_{\lambda}}\frac{\lambda_{i}-\lambda_{j}}{\lambda_{c}}\right)^{2}\right] + A_{\delta}\delta_{ij}$$
(2)

where the symbols are as in Greco & Brandt (2016). The first two terms are correlated; the first models the contribution of speckle noise, while the second models other spectral correlations, such as correlations induced by the interpolation of the micro-spectra in the reduction process. The third term models uncorrelated noise and does not contribute to the offdiagonal elements of the correlation matrix.

The measurements of  $\psi_{ij}$  at the seven angular separations within the final PSF-subtracted image of HD 95086 were used to fit the parameters of the model in Equation (2). The amplitudes of the three noise terms  $(A_{\rho}, A_{\lambda}, A_{\delta})$  were allowed to vary as a function of angular separation, while the two

 Table 2

 Properties of the HD 95086 System

Property	Value	•	Unit
Parallax	11.06 ± 0	).41 <sup>a</sup>	mas
Distance	$90.42 \pm 3$	3.35 <sup>a</sup>	pc
$\mu_{lpha}$	$-41.41~\pm$	0.42 <sup>a</sup>	mas $yr^{-1}$
$\mu_{\delta}$	$12.47 \pm 0$	0.36 <sup>a</sup>	mas yr <sup>-1</sup>
Age	$17 \pm 4$	4 <sup>b</sup>	Myr
	HD 95086	HD 95086 b	
GPI $\Delta H$		$13.70\pm0.25$	mag
GPI $\Delta K_1$	•••	$12.20\pm0.20$	mag
NaCo $\Delta L'$	•••	$9.48 \pm 0.19^{\circ}$	mag
2MASS H	$6.867 \pm 0.047^{d}$		mag
GPI H	$6.807 \pm 0.025^{\rm e}$	$\textbf{20.51} \pm \textbf{0.25}$	mag
2MASS $K_S$	$6.789 \pm 0.021^{d}$		mag
GPI $K_1$	$6.785 \pm 0.025^{e}$	$18.99\pm0.20$	mag
WISE W1	$6.717 \pm 0.067^{f}$		mag
NaCo L'	$6.787 \pm 0.025^{e}$	$16.27\pm0.19$	mag
$H - K_1$	$0.022 \pm 0.035$	$1.52\pm0.32$	mag
H - L'	$0.019\pm0.035$	$\textbf{4.24} \pm \textbf{0.32}$	mag
$K_1 - L'$	$0.035\pm0.035$	$2.72\pm0.28$	mag
$M_H$	$2.025 \pm 0.084$	$\textbf{15.73} \pm \textbf{0.26}$	mag
$M_{K_1}$	$2.004 \pm 0.084$	$14.20\pm0.22$	mag
$M_{L'}$	$2.006 \pm 0.084$	$11.49\pm0.21$	mag
SpT	$A8 \pm 1^{b}$	L1-T3	
$T_{\rm eff}$	$7609 \pm 128^{g}$	800-1300	K
	$7633 \pm 24^{e}$		K
log g	~4.15 <sup>h</sup>	$\lesssim 4.5$	dex
	$3.93 \pm 0.05^{e}$		dex
$\log L/L_{\odot}$	$0.836 \pm 0.035^{b}$	$-4.96 \pm 0.10^{i}$	dex
Mass	~1.7 <sup>b</sup>		$M_{\odot}$
		$4.4 \pm 0.8^{j}$	$M_{ m Jup}$

Notes.

<sup>a</sup> van Leeuwen (2007).

<sup>b</sup> Meshkat et al. (2013).

<sup>c</sup> Galicher et al. (2014).

<sup>d</sup> Skrutskie et al. (2006).

<sup>e</sup> Estimated from best fit BT-NEXTGEN model atmosphere (Section 2.4).

<sup>f</sup> Cutri (2014).

<sup>g</sup> Based on an average of the measurements from Allende Prieto & Lambert (1999), Wright et al. (2003), Bonneau et al. (2006), McDonald et al. (2012). <sup>h</sup> Using Dartmouth isochrones from Dotter et al. (2008).

<sup>i</sup> Using *K*-band bolometric correction.

<sup>j</sup> Based on Baraffe et al. (2003) evolutionary models.

correlation lengths ( $\sigma_{\rho}$ ,  $\sigma_{\lambda}$ ) were fixed over the entire field-ofview. As the sum of the three amplitudes was fixed at unity, there were a total of sixteen free parameters to fit. Due to the high dimensionality of the problem, a parallel-tempered Markov Chain Monte Carlo algorithm was used to properly sample the posterior distributions of each free parameter and to find the global minimum (Foreman-Mackey et al. 2013). The best fit parameters at the separation of HD 95086 b were found to be:  $A_{\rho} = 0.077$ ,  $A_{\lambda} = 0.851$ ,  $A_{\delta} = 0.073$ ,  $\sigma_{\rho} = 0.375 \lambda/D$ ,  $\sigma_{\lambda} = 0.004 \lambda/D$ . From this model of the correlation, the offdiagonal elements of the covariance matrix *C* were then calculated using

$$\psi_{ij} \equiv \frac{C_{ij}}{\sqrt{C_{ii}C_{jj}}} \tag{3}$$

where the diagonal elements of *C* were the square of the stated uncertainties on the spectrum of the planet given in Table 1. Based on this analysis, the spectrum was shown to be highly correlated for flux measurements separated by up to two wavelength channels (each channel is 0.0086  $\mu$ m wide), with a significantly lower correlation for fluxes separated by more than two wavelength channels. Both the parametrized model and a visual inspection of the final images showed that the data were not dominated by speckles, such that the correlation can be ascribed almost entirely to the wavelength interpolation carried out as a part of the standard reduction procedure for GPI data.

# 3. RESULTS AND EMPIRICAL COMPARISONS

#### 3.1. Photometry

We measured contrasts of  $\Delta H = 13.70 \pm 0.25$  and  $\Delta K_1 = 12.20 \pm 0.20$  mag for HD 95086b. These new measurements are in agreement with the contrasts given in Galicher et al. (2014), but the uncertainties have been significantly reduced due to an improvement of the S/N of the detections. These contrasts, and the *L'* contrast given in Galicher et al. (2014), were combined with the predicted flux of HD 95086 in each filter reported in Section 2.4 to calculate the apparent magnitudes of HD 95086 b. Absolute magnitudes were then calculated using these apparent magnitudes, and the *Hipparcos* parallax of HD 95086. A summary of all available photometry of HD 95086 b is given in Table 2.

The location of HD 95086 b on a  $K_1 - L'$  versus  $H - K_1$ color-color diagram (CCD) is shown in Figure 3, alongside Sco-Cen substellar companions, benchmark planetary-mass companions, and isolated young and/or dusty brown dwarfs (listed in Table 7), and field M-dwarfs and brown dwarfs within both the SpeX Prism Spectral Library<sup>28</sup> and those listed in Dupuy & Liu (2012). As no spectral information exists within either the H- or L-band for HD 95086 b, a color transformation could not be determined between either GPI H and MKO H or NaCo L' and MKO L'. Based on the empirical color transformations given in Table 6, the H-band color transformation is estimated to be small (<0.1 mag), despite the spectral type of HD 95086 b being relatively unconstrained (see Section 3.2). No empirical estimate for the *L*-band color transformation was made given the paucity of thermal-infrared spectra of brown dwarfs, although it is likely to be small relative to the size of the measurement uncertainties.

Photometry for the comparison objects was collated from a variety of sources: Lafrenière et al. (2008) and Lachapelle et al. (2015) for 1RXS J160929.1-210524 b (1RXS J1609 b); Bailey et al. (2014) for HD 106906 b; Ireland et al. (2011) and Lachapelle et al. (2015) for GSC 06214-00210 b (GSC 6214 b); Bailey et al. (2013) and Lachapelle et al. (2015) for HIP 78530 B (Lafrenière et al. 2011); Chauvin et al. (2004) for 2MASS J120734-393253 b (2M1207 b); Marois et al. (2008) and Esposito et al. (2013) for HR 8799 bcd; Skrutskie et al. (2006) and Cutri (2014) for 2MASS J11193254-1137466 (2M 1119-11; Kellogg et al. 2015), 2MASS J21481628+4003593 (2M2148+40; Looper et al. 2008), and 2MASS J22443167 +2043433 (2M2244+20; Dahn et al. 2002); and Liu et al. (2013) for PSO J318.5338-22.8603 (PSO 318.5-22).

<sup>&</sup>lt;sup>28</sup> http://pono.ucsd.edu/~adam/browndwarfs/spexprism



**Figure 3.** MKO  $H - \text{GPI} K_1$  vs. GPI  $K_1 - \text{MKO} L'$  color-color diagram showing the position of HD 95086 b (red filled star) with respect to Sco-Cen substellar companions, benchmark planetary-mass companions, isolated young and/or dusty brown dwarfs, and field M-, L-, and T-dwarfs. HR 8799 bcd are plotted as blue filled circles, and are distinguished by a small letter within the symbol. Sources for the photometric measurements of the comparison objects are given in Section 3.1. The colors of HD 95086 b were calculated using the GPI *H*-band and NaCo *L'* photometry, as discussed in Section 3.1. For comparison, the 17 Myr isochrone was computed using the AMES-COND (dotted line), AMES-DUSTY (solid line), and BT-SETTL (dashed line) atmospheric models, in conjunction with the Baraffe et al. (2003) evolutionary models.

The MKO K-band  $(K, K_S, K')$  or 2MASS  $K_S$  magnitudes for each comparison object were converted into GPI  $K_1$  magnitudes using their K-band spectra (see Section 3.2.2) using the procedure described in the Appendix. For the M-dwarfs and brown dwarfs within the SpeX library, the spectra were flux calibrated using either MKO or 2MASS H-band absolute magnitudes from Dupuy & Liu (2012). The GPI  $K_1$  magnitude was then estimated using this flux calibrated spectrum, and the  $K_1$  filter profile, as described in the Appendix. For brown dwarfs within Dupuy & Liu (2012) without a spectrum within the SpeX library, the empirical color transformation given in Table 6 was used. MKO L' photometry for these objects was obtained from a variety of sources (Dupuy & Liu 2012, and references therein). If no MKO L' measurement was available, it was instead estimated from the WISE W1 magnitude, and an empirical MKO L' - WISE W1 color transformation given in the Appendix (Table 6).

HD 95086 b was already known to have unusually red H - L' colors (Meshkat et al. 2013; Galicher et al. 2014). Figure 3 further demonstrates the peculiar nature of this object. It has a  $K_1 - L'$  consistent with some late T-dwarfs, but an  $H - K_1$  that is redder than known field objects. Relative to other young substellar objects, HD 95086 b has a  $K_1 - L'$  and  $H - K_1$  color that is approximately a magnitude redder than that of 2M1207 b and HR 8799 d. Of the four isolated brown dwarfs identified as being young or possessing unusually dusty atmospheres, HD 95086 b is most analogous to PSO 318.5-22, one of the reddest field brown dwarfs known (Liu et al. 2013), although HD 95086 b is half a magnitude redder in both  $K_1 - L'$  and  $H - K_1$ . When compared with the location of the



**Figure 4.** GPI  $M_{K_1}$  vs. MKO H – GPI  $K_1$  color–magnitude diagram showing the position of HD 95086 b relative to substellar companions within Sco–Cen, benchmark planetary-mass companions, and field M-, L-, and T-dwarfs. Symbols and lines have the same meaning as in Figure 3.

17 Myr isochrones calculated from three different types of the atmospheric models, HD 95086 b is closer to the predicted colors of the models, which simulate a large photospheric dust content (AMES-DUSTY, Chabrier et al. 2000; Allard et al. 2001). The location of these isochrones on the CCD does not vary significantly within the uncertainty of the age of HD 95086.

The updated photometry of HD 95086 b was also used to place it on two CMDs:  $M_{K_1}$  versus  $H - K_1$  in Figure 4, and  $M_{L'}$ versus  $K_1 - L'$  in Figure 5. In both diagrams, the Sco-Cen substellar companions are both brighter and bluer than HD 95086 b, consistent with the planet being cooler, less massive, and with a dustier photosphere. HD 95086 b has a  $K_1$ luminosity consistent with mid-T dwarfs, and an L' luminosity consistent with L/T transition objects. In both instances, the planet is significantly redder than the field population and lies on an apparent extension to the L-dwarf sequence. This extension is populated with lower surface gravity objects that have maintained cloudy atmospheres at temperatures for which higher-surface gravity field brown dwarfs have transitioned to clear atmospheres (Kirkpatrick 2005). As previously recognized (Galicher et al. 2014), HD 95086 b has a similar absolute L' magnitude to HR 8799 cd and 2M1207 b, suggesting similar effective temperatures, but it is approximately a magnitude fainter in  $K_1$ , with a similar luminosity to HR 8799 b.

Finally, the very red  $K_1 - L'$  color of HD 95086 b could also be explained by the presence of circumplanetary material, either a dust shell or a disc, as is the case for GSC 6214 b (Bailey et al. 2013). Contrary to GSC 6214 b, which has strong evidence of accretion (Bowler et al. 2014; Zhou et al. 2014; Lachapelle et al. 2015), no additional constraints beyond the red color support this hypothesis for HD 95086 b. Improvements in atmospheric modelling and the treatment of dust were able to explain the underluminosity of 2M1207 b (Barman et al. 2011b), which had previously been suggested was due to the presence of circumstellar material (Mohanty et al. 2007).



**Figure 5.** MKO L' vs. GPI  $K_1$  – MKO L' color–magnitude diagram showing the position of HD 95086 b relative to substellar companions within Sco–Cen, benchmark planetary-mass companions, and field M-, L-, and T-dwarfs. Symbols and lines have the same meaning as in Figure 3. The extremely red  $K_1 - L'$  color is more consistent with field late T-dwarfs, while the L' absolute magnitude suggests an earlier spectral type.

Future narrow-band observations sensitive to, for example,  $H\alpha$  or  $Br\gamma$  emission are required in order to further test this hypothesis.

# 3.2. Spectroscopy and the Full SED

The  $K_1$  spectrum of HD 95086 b shown in Figure 2 exhibits a relatively featureless slope rising toward longer wavelengths, as is seen for other young and/or dusty brown dwarfs (e.g., Stephens et al. 2009; Patience et al. 2010; Allers & Liu 2013; Liu et al. 2013). There is no evidence of strong molecular absorption due to CH<sub>4</sub> and H<sub>2</sub>O beyond 2.08  $\mu$ m (Kirkpatrick 2005), as is seen within the cloud-free atmospheres of mid- to late-T dwarfs (e.g., Naud et al. 2014; Gagné et al. 2015). The spectrum of HD 95086 b does show a slight decrease in flux between 2.03 and 2.12  $\mu$ m. However, given the size of the uncertainties on the spectrum relative to the size of the decrease, it is unclear whether this is an absorption feature within the spectrum or the result of the low S/N of the spectrum.

#### 3.2.1. Spectral Type and Comparison to Field Objects

To assess the spectral type of HD 90586 b, the literature photometry and the new  $K_1$  spectrum were compared with stars and brown dwarfs within the SpeX Prism Spectral Library and the IRTF Spectral Library (Cushing et al. 2005; Rayner et al. 2009). As most objects within both spectral libraries had no *L*-band spectroscopy from which a synthetic magnitude could be calculated, MKO L' photometric measurements were obtained from a number of literature sources (Stephens et al. 2001; Leggett et al. 2003, 2010; Golimowski et al. 2004; Dupuy & Liu 2012; Faherty et al. 2012). For those objects without literature L' photometry, it was instead estimated using the W1 magnitude from the AllWISE catalog (Cutri 2014), and



**Figure 6.** The minimum  $\chi^2$  for each object within the SpeX Prism Spectral Library and the IRTF Spectral Library considering only the near-infrared SED (*H* & *K*<sub>1</sub>, top panel) and full SED (*H*, *K*<sub>1</sub> & *L'*, bottom panel) of HD 95086 b. Those objects with literature MKO *L'* photometry are highlighted as filled red squares. The *L'* photometry for the remaining objects was estimated from their *WISE* W1 magnitudes, and the empirical L' - W1 color transformation given in Table 6. For both families of fits, the average and standard deviation of the  $\chi^2$  values were computed in bins two spectral subtypes wide (open symbols). From these fits, the spectral type of HD 95086 b was found to be between L7 ± 6 and L9<sup>+4</sup><sub>-6</sub> using the near-infrared and full SEDs, respectively. The lower and upper bounds are defined as the spectral type bin at which the average of the  $\chi^2$  of that bin is significantly higher than that of the minimum average  $\chi^2$ . Two of the best fitting field brown dwarfs, 2M2244+20 and 2M2148+40, are indicated.

the empirical spectral type versus L' - W1 color relation given in Table 6. As the spectra within the SpeX library were not flux calibrated, they were scaled to best fit the 2MASS *J*, *H*, and *K*<sub>S</sub> apparent magnitudes for each object.

For each comparison object, two fits were attempted: a fit to only the *H*-band photometry and  $K_1$  spectrum (Figure 6, lower curve), and a fit also including the *L'* photometry (Figure 6, upper curve). Synthetic *H*-band photometry was calculated for each object using the procedure described in the Appendix, and a synthetic  $K_1$  spectrum at the resolution of GPI was created by degrading the resolution of the library spectrum by convolution with a Gaussian kernel. The goodness of fit  $\chi^2$  was then calculated for each comparison object as

$$\chi^2 = R^{\mathrm{T}} C^{-1} R + \sum_i \left( \frac{f_i - \alpha s_i}{\sigma_i} \right)^2. \tag{4}$$

The first term computes the  $\chi^2$  of the fit to the  $K_1$  spectrum alone, accounting for the covariances between neighboring wavelength channels.  $R = F - \alpha S$ , the residual vector, is the difference between the  $K_1$  spectrum of HD 95086 b  $F = \{f_1 \dots f_n\}$ , and the  $K_1$  spectrum of the comparison object  $S = \{s_1 \dots s_n\}$  multiplied by a scaling factor  $\alpha$ , where n = 29the number of wavelength channels within the  $K_1$  spectrum. *C* is the covariance matrix calculated from the parametrized noise model described in Section 2.5. The second term incorporates the *H* and *L'* photometry of HD 95086 b in the standard  $\chi^2$ formalism:  $f_i$  and  $\sigma_i$  are the flux and corresponding uncertainty of HD 95086 b in the ith filter,  $\alpha$  is the same multiplicative factor as in the first term, and  $s_i$  is the flux of the comparison object in the same filter. For each comparison object, the value of the multiplicative factor  $\alpha$  which minimizes  $\chi^2$  was found using a truncated-Newton minimization algorithm. The resulting minimum  $\chi^2$  are plotted as a function of spectral type in Figure 6.

Considering only the fit to the near-infrared SED of HD 95086 b,  $\chi^2$  reaches a minimum at spectral type of L7, although this minimum is poorly constrained between L1 and T3 (Figure 6, top panel). Later T-dwarfs are strongly excluded due to the decline in flux seen in the *K*-band beyond 2.08  $\mu$ m caused by increasing CH<sub>4</sub> and H<sub>2</sub>O absorption within the atmospheres of these cool objects (Kirkpatrick 2005), compared with the monotonically rising spectrum of HD 95086 b. The addition of the *L'* photometry to the fit results in a significant increase in the value of  $\chi^2$  at all spectral types (Figure 6, bottom panel), consistent with the unusual location of HD 95086 b on the CCD and CMDs (Figures 3, 4, 5). Fitting to the full SED of HD 95086 b shifted the minimum  $\chi^2$  to the later spectral type of L9, with the minimum similarly poorly constrained between L3 and T3.

The best fitting object was the very red L6.5 brown dwarf 2M2244+20 (Dahn et al. 2002), which possesses an extremely cloudy atmosphere (Leggett et al. 2007), may have a low surface gravity (McLean et al. 2003; Looper et al. 2008), and was found to be the best match to the near-infrared SED of HR 8799 b (Bowler et al. 2010). Another good fit was the L6 brown dwarf 2M2148+40 (Looper et al. 2008), which has been identified as being old but with extremely red near-infrared colors indicative of thick clouds (Cruz et al. 2009). The spectral type of HD 95086 b was also estimated using the empirical absolute magnitude-spectral type relations given in Dupuy & Liu (2012) for field brown dwarfs. Using the absolute magnitudes of HD 95086 b given in Table 2, the spectral type was estimated as T7  $\pm$  0.5, T3  $\pm$  1.5, and L9  $\pm$  2 using  $M_{H}$ ,  $M_{K_{\rm I}}$ , and  $M_{L'}$ , respectively. These estimates further emphasize the peculiar morphology of the SED of HD 95086 b and how discrepant it is from that of old field brown dwarfs.

# 3.2.2. Comparison to Young Substellar Companions and Isolated Young/Dusty Brown Dwarfs

The complete SED of HD 95086 b was also compared to that of substellar companions within Sco-Cen, and around other young stars, and of isolated young and/or dusty brown dwarfs in Figure 7. For each comparison object, the SED was normalized to the integrated flux in the  $K_1$  filter bandpass. The  $K_1$  spectrum of HIP 78530 B (Lafrenière et al. 2011) and HIP 106906 b (Bailey et al. 2014) are flatter than that of HD 95086 b, which may indicate a lower effective temperature and a spectral type later than L2.5 for HD 95086 b. The  $K_1$  spectrum of HD 95086 b is most similar to that of GSC 6214 b (Bowler et al. 2011), 1RXS 1609 b (Lafrenière et al. 2010), 2M1207 b (Patience et al. 2010), and HR 8799 d (Ingraham et al. 2014). HR 8799 b (Barman et al. 2011a) and c (Ingraham et al. 2014) have a somewhat steeper slope in  $K_1$  than HD 95086 b. The four young and/or dusty brown dwarfs have very similar  $K_1$ spectra, consistent with their similar effective temperatures and spectral types. These comparison objects span a range of effective temperatures (750-2700 K), and spectral types (from mid-M to L/T transition objects). With the exception of HIP 78530 B, which exhibits an almost flat  $K_1$  spectrum due to its



**Figure 7.** Spectral energy distribution of HD 95086 b (gray points) compared with several substellar companions within Sco–Cen (top four), other young planetary-mass companions (middle four), and isolated young and/or dusty brown dwarfs (bottom four). The spectral energy distribution of each object has been normalized by the flux within the  $K_1$  bandpass. For clarity, dotted lines connect the *L'* flux of HD 95086 b and the comparison object. The discrepancy between the *H*-band photometry and *H*-band spectrum for HR 8799 b is because the *H*-band spectrum was normalized by Barman et al. (2011a) using H = 18.06 (Metchev et al. 2009), whereas the plotted *H*-band photometry uses H = 17.88 (Esposito et al. 2013).

Model		Parameter R	ange/Spacing		HD 95086 b					
	T <sub>eff</sub> (K)	$\Delta T_{\rm eff}$ (K)	log g (dex)	$\Delta \log g$ (dex)	T <sub>eff</sub> (K)	log g (dex)	R (R <sub>Jup</sub> )	$\log L/L_{\odot}$ (dex)	$\chi^2$	$\chi^2_{\nu}$
			Cloud	y/Dusty Atmo	spheres					
AMES-DUSTY	500-1500	100	3.50-6.00	0.50	1200 1250	5.00 3.50	0.83	-4.86 	14.35 13.51	0.55
BT-DUSTY	1000-1500	100 5	4.50-5.50	0.50	1000	4.50 4.50	1.24	-4.86 	18.42 18.42	0.71
DRIFT-PHOENIX	1000-1500	100 5	3.00-6.00	0.50	1300 1270	3.50 3.50	0.72	-4.88 	11.53 10.99	0.44
Madhusudhan (A60-F)	700–1500	100 5	3.75-4.25	0.25 0.01	800 780	3.75 3.75	1.87 2.00	-4.90 	16.26 15.31	0.63 0.59
			C	lear Atmosphe	res					
AMES-Cond	500-1500	100 5	2.50-6.00	0.50 0.01	1000 965	2.50 2.68	0.81 0.88	-5.23 	47.92 47.82	1.84 1.84
BT-COND	800-1500	100 5	4.00-5.50	0.50 0.01	1000 975	4.00 4.00	1.00 1.07	-5.05 	49.95 49.93	1.92 1.92
BT-Settl	500-1500	50 5	3.00-5.50	0.50 0.01	1150 1150	3.50 3.51	0.69 0.68	-5.17 	20.80 20.80	0.80 0.80

Table 3 Model Grid Properties and Best Fit

high effective temperature (2700 K), there is very little diversity in the shape of the  $K_1$  spectra, despite the large range of effective temperatures. This suggests that the shape of the  $K_1$ spectrum on its own is a poor diagnostic of the effective temperatures of young planetary-mass objects when at a resolution and S/N similar to the HD 95086 b spectrum.

# 4. MODEL ATMOSPHERE FITTING

Combining the revised *H*-band photometry and the new  $K_1$ spectrum obtained with GPI with literature L' photometry (Table 2; Galicher et al. 2014) further constrains the nearinfrared SED of HD 95086 b. The SED was fit to seven publicly available grids of model atmospheres: AMES-COND<sup>2</sup> (Allard et al. 2001; Baraffe et al. 2003), AMES-DUSTY<sup>30</sup> (Chabrier et al. 2000; Allard et al. 2001), BT-COND<sup>31</sup>, BT-DUSTY<sup>32</sup> and BT-SETTL<sup>33</sup>(e.g., Allard et al. 2012), DRIFT-PHOENIX<sup>34</sup> (e.g., Helling et al. 2008), and the various models published by A. Burrows and collaborators<sup>35</sup> (Burrows et al. 2006; Hubeny & Burrows 2007; Madhusudhan et al. 2011). The range of effective temperatures  $(T_{eff})$  and surface gravities  $(\log g)$ , and the grid spacing for each parameter, is given for each model grid in Table 3. Due to the irregular sampling of some grids at non-solar metallicities, a solar metallicity for HD 95086 b was assumed.

#### 4.1. AMES-COND and AMES-DUSTY

The AMES family of models (Allard et al. 2001) were created by combining the PHOENIX plane-parallel radiative transfer atmosphere model (Hauschildt 1992) with the NASA AMES molecular H<sub>2</sub>O and TiO line lists (Partridge &

<sup>35</sup> http://www.astro.princeton.edu/~burrows

Schwenke 1997) using the solar abundances of Grevesse et al. (1993). The AMES-DUSTY grid represents a limiting case, including treatments for both the condensation of dust within the photosphere and for dust opacities. Dust grain sizes were drawn from the interstellar grain size distribution (Mathis et al. 1977), and opacities were calculated assuming Mie scattering for spherical grains. Gravitational sedimentation was ignored, leading to a significant amount of photospheric dust. The effect of the dust was to increase the gas temperature, which in turn decreased the strength of molecular absorption features in the optical and near-infrared. On the other extreme, the AMES-COND grid included the treatment for dust condensation but ignored dust opacities entirely in order to simulate the immediate sedimentation of dust to the lower atmosphere, leading to a dust-free photosphere (Allard et al. 2001).

# 4.2. BT-COND, BT-DUSTY and BT-SETTL

As with the AMES models, the BT family of models (e.g., Allard et al. 2012) also used the PHOENIX atmosphere model but with updated molecular line lists, most notably for H<sub>2</sub>O (Barber et al. 2006), and revised solar abundances of Asplund et al. (2009) for BT-Dusty and Caffau et al. (2011) for BT-COND and BT-SETTL. The BT-COND and BT-DUSTY grids used the same treatment for dust as in their counterpart within the AMES grid, while the BT-SETTL grid used a more detailed cloud model to define the number density and size distribution of condensates within the atmosphere (Allard et al. 2012). The BT-SETTL models are able to reproduce the L/T transition observed for field brown dwarfs, where the atmosphere transitions from fully cloudy to fully clear (e.g., Marley et al. 2002), due in part to the formation of methane at lower temperatures.

# 4.3. Drift-Phoenix

The DRIFT-PHOENIX grids (Woitke & Helling 2003, 2004; Helling & Woitke 2006; Helling et al. 2008) also used the PHOENIX model with the same solar abundances as the AMES

<sup>&</sup>lt;sup>29</sup> https://phoenix.ens-lyon.fr/Grids/AMES-Cond/SPECTRA/

<sup>&</sup>lt;sup>30</sup> https://phoenix.ens-lyon.fr/Grids/AMES-Dusty/SPECTRA/

<sup>&</sup>lt;sup>31</sup> https://phoenix.ens-lyon.fr/Grids/BT-Cond/

<sup>&</sup>lt;sup>32</sup> https://phoenix.ens-lyon.fr/Grids/BT-Dusty/AGSS2009

<sup>&</sup>lt;sup>33</sup> https://phoenix.ens-lyon.fr/Grids/BT-Settl/CIFIST2011

<sup>&</sup>lt;sup>34</sup> http://svo2.cab.inta-csic.es/theory/main

models but is coupled with the non-equilibrium cloud model DRIFT, which includes treatments for the formation, growth, evaporation, settling, and advection of grains (Helling et al. 2008). Contrary to the other models, they used a kinetic approach to describe the formation of grains and a top-down approach to simulate the motion of the grains within the atmosphere. The DRIFT-PHOENIX models have been successful in reproducing the near-infrared SEDs of young substellar companions over a wide range of effective temperatures (e.g., Patience et al. 2012; Bonnefoy et al. 2014a; Lachapelle et al. 2015).

#### 4.4. Madhusudhan et al.

The models of A. Burrows and collaborators (Burrows et al. 2006; Hubeny & Burrows 2007; Madhusudhan et al. 2011) use a variant of the TLUSTY model stellar atmosphere code (Hubeny 1988; Hubeny & Lanz 1995) where the opacity as a function of wavelength is pre-computed using a chemical equilibrium code to account for the sedimentation of particles in a gravitational field (Hubeny et al. 2003). The Burrows et al. (2006) models are able to reproduce the overall shape of the L/T transition, although no single combination of surface gravity and cloud particle size can fit all of the observations. Using an updated version of the atmospheric model, Hubeny & Burrows (2007) explored the effects of non-equilibrium chemistry, although these effects are primarily limited to the *M* (4.8  $\mu$ m) and *N* (10.5  $\mu$ m) bands (e.g., Saumon et al. 2000).

Although these models were able to reproduce the overall shape of the L through T sequence, they were a poor fit to the observations of the HR 8799 planets. By invoking a significantly thicker cloud layer and adjusting the modal dust particle size, Madhusudhan et al. (2011) were able to reproduce the available photometry of the three outermost planets of the HR 8799 system, which are displaced from field brown dwarf sequence on the CMD. Madhusudhan et al. (2011) considered four models of the cloud structure, modal particle sizes between 1–100  $\mu$ m, and spherical dust grains composed of either fosterite or iron. While the SED of HD 95086 b was fit to each of the nine grids from Madhusudhan et al. (2011), and the 21 grids from Burrows et al. (2006) and Hubeny & Burrows (2007), only the results of the grid containing the best fit are discussed below (Model A, 60  $\mu$ m modal particle size, fosterite grains; Madhusudhan et al. 2011).

# 4.5. Fitting Procedure

For the individual spectra within each model grid the fitting procedure was the same. Synthetic H and L' photometry was computed for each model spectrum by folding the spectrum through the GPI H and NaCo L' filter transmission curves. Synthetic  $K_1$  spectra were also computed by degrading the resolution of the model spectra to that of the GPI K1 data  $(\lambda/\delta\lambda \approx 66)$  and interpolating the resulting smoothed spectra to the same wavelength values. The goodness of fit  $(\chi^2)$  for each model to the measurements of HD 95086 was calculated using the method described in Section 3.2.1, where  $\alpha = R^2/d^2$ was allowed to vary such that the radius R was in the range 0.6–2.0  $R_{Jup}$  as in Galicher et al. (2014), given the distance d to HD 95086 measured by Hipparcos (van Leeuwen 2007). This process was repeated on a finer version of the grid, obtained by linearly interpolating between the logarithm of the flux at each grid point, where the spacing in  $T_{\rm eff}$  and log g were reduced to 5 K and 0.01 dex in order to evaluate the confidence interval on each parameter. The 68%, 95%, and 99% confidence regions were calculated by integrating the likelihood  $\mathcal{L}$  over the full parameter range, where  $\mathcal{L} = e^{-\chi^2/2}$ .

#### 4.6. Effective Temperature and Surface Gravity

The effective temperature and surface gravity of the best fit spectrum from each of the seven model grids, and their interpolated versions, are given in Table 3. For each fit, the reduced  $\chi^2$  was calculated assuming 28 degrees of freedom. The best-fit spectra within the original non-interpolated grids are plotted in Figure 8. The  $\Delta\chi^2 \equiv \chi^2 - \chi^2_{\rm min}$  surfaces for each grid are plotted as a function of temperature and surface gravity in Figure 9. Of the Madhusudhan et al. (2011) models, the best fit was found in the Model A grid, the thickest (in vertical extent) of the four cloud structures simulated, and thicker than the clouds required to explain the SEDs of the HR 8799 planets (Madhusudhan et al. 2011). As such, only the  $\Delta\chi^2$  surface from this particular grid is shown in Figure 9.

The SED of HD 95086 b was best fit by an atmosphere within the DRIFT-PHOENIX grid which simulates a significant photospheric dust content. Within this grid, two distinct minima in the  $\Delta \chi^2$  surface were found (Figure 9). One was found at high surface gravities ( $\log g = 6.0$ ) at a temperature of 1100 K, with a  $\chi^2 = 10.83$  and a radius of 0.95  $R_{Jup}$ ; and the other at a lower surface gravity ( $\log g = 3.5$ ) at 1300 K, with a  $\chi^2 = 11.53$  and a smaller radius of 0.72  $R_{Jup}$ . As young planetary-mass objects are expected to have relatively low surface gravities, only the lower surface gravity fit is reported in Table 3. The SED was also fit well by an AMES-DUSTY  $T_{\rm eff} = 1200$  K,  $\log g = 5.0$  model atmosphere, with a radius of 0.83  $R_{Jup}$  ( $\chi^2 = 14.35$ ). Unfortunately, the AMES-DUSTY grid<sup>36</sup> contains several erroneous models at 1200 K: the  $\log g = 5.5$  model is incomplete, and the  $\log g = 3.5-4.5$ models are all identical. For the purposes of this study, we assume that these three identical models have  $\log g = 4.5$ . In the interpolated grid, where these gaps are interpolated over, the best fit is at a similar  $T_{\rm eff} = 1250$  K, but at a lower  $\log g = 3.50$  ( $\chi^2 = 13.51$ ). The two remaining grids which simulate a dusty atmosphere-the Madhusudhan et al. and BT-DUSTY models-both predict a low surface gravity, although the goodness of fit statistic for each are worse at  $\chi^2 = 16.26$ and  $\chi^2 = 18.42$ . In each case the best fit model is railed at the grid boundary, suggesting that an expanded parameter search is necessary to locate the true minimum  $\chi^2$ .

The fits to the clear atmospheres are worse, with a goodness of fit ranging between  $\chi^2 = 20.80$  for the BT-SETTL grid and  $\chi^2 = 49.95$  for the BT-COND grid. The best-fitting temperatures for HD 95086 b within these grids are cooler than the predicted temperature of the L–T transition, where brown dwarfs transition from having dusty to clear atmospheres with decreasing temperature. The position of HD 95086 b on the CMD (Figures 4 and 5) in the extension to the L-dwarf sequence, is inconsistent with the predicted luminosity and colors of objects with fully clear atmospheres. As the current observational evidence is consistent with a dusty atmosphere, we use only the fits to the dusty atmosphere models to estimate a  $T_{\rm iff}$  and log g for HD 95086 b. Unfortunately, the variation in which part of parameter space is sampled by each grid means

<sup>&</sup>lt;sup>36</sup> https://phoenix.ens-lyon.fr/Grids/AMES-Dusty/SPECTRA/



Figure 8. Photometry and spectroscopy of HD 95086 b (red open circles) and best-fit model atmosphere spectra. Synthetic H and L' photometry for the model spectra are plotted as open squares, with the vertical error bar corresponding to the effective width of the filter. The top panel contains the best fit from each model grid for ease of comparison, while the bottom panels contain each fit individually.

that these two parameters are still relatively unconstrained. In terms of  $T_{\rm iff}$ , best fit models were found between 800 and 1300 K, whereas log g can only be constrained to be  $\lesssim 4.5$  as the best fit is railed at the lowest surface gravity within three of the four model grids.

#### 4.7. Bolometric Luminosity and Mass

Using the revised *H* and *K*<sub>1</sub> magnitudes presented in this study, the bolometric luminosity of HD 95086 b was estimated. The GPI magnitudes were converted into MKO *H* and *K* magnitudes using the empirical filter transformations in Table 6, assuming a spectral type of L7 ± 6 (Section 3.2). The bolometric corrections for each band were then estimated from the spectral type as  $BC(H) = 2.53^{+0.13}_{-0.08}$  from the fit to field objects of Liu et al. (2010), and  $BC(K) = 3.19^{+0.13}_{-0.38}$ from the fit to young low-surface gravity objects of Filippazzo et al. (2015), corresponding to bolometric magnitudes of  $M_{bol}(H) = 18.30 \pm 0.28$  and  $M_{bol}(K) = 17.15^{+0.24}_{-0.24}$ .

Assuming a solar bolometric magnitude of  $M_{\text{bol},\odot} = 4.74$ , these bolometric magnitudes correspond to luminosities of  $\log L/L_{\odot} = -5.42 \pm 0.11$  using the *H*-band correction and  $\log L/L_{\odot} = -4.96 \pm 0.10$  using the K-band correction. While the K-band bolometric correction was based on young lowsurface gravity objects, the H-band correction was based on field brown dwarfs, which have significantly different SEDs to that of young substellar objects. For instance, Filippazzo et al. (2015) measured a one magnitude offset in the value of BC(J)between field and young objects. Given the unusual nature of the SED of HD 95086 b, the bolometric magnitudes and the associated derived luminosities may be similarly biased. The bolometric luminosity was also estimated by integrating the best fit model atmosphere from each of the model grids, ranging from  $\log L/L_{\odot} = -5.23$  to -4.86 (Table 3), with cloudy atmosphere models predicting higher luminosities (-4.90 to -4.86) than clear atmosphere models (-5.23)to -5.05).



Figure 9.  $\Delta \chi^2 \equiv \chi^2 - \chi^2_{min}$  surfaces for the seven model grids (left column), and the six interpolated grids (right column) plotted as a function of effective temperature and surface gravity. For each grid, the best fitting model is indicated by the circle symbol. For the DRIFT-PHOENX grid, the best fit low surface gravity model is also indicated. For the interpolated grids, the black contours indicate the 68% (solid), 95% (dashed), and 99% (dotted) confidence regions. White contours define radii of 0.6, 1.0, and 2.0  $R_{Jup}$ , based on the dilution factor required to minimize  $\chi^2$  for each spectrum. Positive hatching corresponds to regions of phase space within each grid for which no model spectra exist. Negative hatching indicate models with a dilution factor corresponding to a radius of  $R < 0.6 R_{Jup}$  and  $R > 2.0 R_{Jup}$ , which are not physical for HD 95086 b (Mordasini et al. 2012). Only the results of the fit to the spectra within the Model A, 60  $\mu$ m modal grain size, fosterite grain grid are shown for the Madhusudhan et al. (2011) grid.

As the age of HD 95086 b  $(17 \pm 4 \text{ Myr})$  is well constrained due to the membership of the Lower Centaurus Crux subgroup of the Sco–Cen association, the luminosity can be used to obtain a model-dependent estimate of the mass. The luminosity of HD 95086 b is plotted relative to three evolutionary models in Figure 10. At the two extremes are the "hot-start" Baraffe et al. (2003) models, in which gas accreting onto the forming planet is unable to efficiently cool, resulting in a high initial luminosity and entropy, and the "cold-start" models of Marley et al. (2007), in which the accreting gas has radiated away most of its heat, resulting in a low initial luminosity and entropy for the planet. Intermediate to these are the models of Mordasini (2013), which predict a strong positive correlation between the luminosity and the mass of the solid core of the planet. In order to assess the range of probable masses for HD 95086 b, only the evolutionary models with the smallest core mass of  $22 M_{\oplus}$  were used from the Mordasini (2013) grid. Higher core masses result in a higher initial luminosity for a given final planet mass, converging with the luminosity predictions of the Baraffe et al. (2003) models for core masses >100  $M_{\oplus}$ .

The mass, radius, and effective temperature derived from fitting both the *H*- and *K*-band luminosity to the Baraffe et al.



**Figure 10.** Luminosity as a function of age for substellar objects from the "hotstart" evolutionary models of Baraffe et al. (2003) (dotted curves), and the "cold-start" evolutionary models of Marley et al. (2007) (dotted–dashed curves) and Mordasini (2013) (solid gray curves). For clarity, only the two most massive models from the Marley et al. (2007) grid are plotted due to the degeneracy between mass and luminosity at the age of HD 95086 b. Two estimates of the luminosity of HD 95086 b, derived from the two bolometric corrections, are plotted (blue square: *H*-band, red circle: *K*-band).

(2003) and Mordasini (2013) evolutionary models are listed in Table 4. No attempt was made to use the Marley et al. (2007) models as the two luminosity estimates were significantly higher than that for the most massive object within the grid, as shown in Figure 10. Based on the Baraffe et al. (2003) models, the mass of HD 95086 b was estimated to be between  $2.7^{+0.6}_{-0.5}$  M<sub>Jup</sub>, using the H-band luminosity, and  $4.4 \pm 0.8$  M<sub>Jup</sub>, using the K-band luminosity. The mass estimated from the Mordasini (2013) models using the H-band luminosity  $(3.1^{+1.0}_{-0.7} M_{Jup})$  was consistent with the two previous estimates, however, as the K-band luminosity was significantly higher than that predicted for the most massive object within the grid, only a lower limit on the mass of  $M > 9.1 M_{Jup}$  could be estimated. The quoted uncertainties on the model-dependent masses incorporate the uncertainty on the age of HD 95086 b, and the uncertainty on the estimated bolometric luminosity derived from the H- and K-band bolometric corrections.

# 5. CONCLUDING REMARKS

We have presented the first spectroscopic measurement of the young, low-mass, and extremely red exoplanet HD 95086 b. Obtained with the Gemini Planet Imager and spanning the blue half of the *K*-band ( $K_1$ , 1.9–2.2  $\mu$ m), the spectrum was recovered at an S/N of 3–5 within the individual wavelength channels, and shows a monotonic increase in flux toward longer wavelengths within the measurement uncertainties. Combining the new  $K_1$  spectrum and the revised *H* photometry with literature NaCo *L'* photometry (Galicher et al. 2014), the SED of the planet between 1.5 and 4.0  $\mu$ m was constructed and compared with other young substellar objects, and older field brown dwarfs. When placed on the MKO  $M_{L'}$  versus GPI  $K_1$  – MKO *L'* CMD (Figure 5), HD 95086 b occupies a region of the

 Table 4

 Physical Properties of HD 95086 b Estimated from the Bolometric Luminosity

Model	Band	М	R	$T_{\rm eff}$
		$(M_{\rm Jup})$	$(R_{Jup})$	(K)
Baraffe	Н	2. $7^{+0.6}_{-0.5}$	$1.32^{+0.02}_{-0.01}$	$703_{-43}^{+47}$
	K	$4.4\substack{+0.8\\-0.8}$	$1.34\substack{+0.02\\-0.02}$	$910_{-47}^{+54}$
Mordasini	Н	3. $1^{+1.0}_{-0.7}$	1. $30^{+0.01}_{-0.01}$	$702_{-46}^{+51}$
	K	>9.1		

Table 5GPI Filter Relative Spectral Response

1	Ч	K	-1	<i>K</i> <sub>2</sub>		
$\lambda(\mu m)$	$S(\lambda)\lambda$	$\lambda$ (µm)	$S(\lambda)\lambda$	$\lambda$ (µm)	$S(\lambda)\lambda$	
1.4946	0.3327	1.8861	0.0376	2.1074	0.2099	
1.5030	0.5142	1.8947	0.0539	2.1154	0.3846	
1.5114	0.6807	1.9033	0.0745	2.1235	0.6079	
1.5198	0.8029	1.9118	0.1082	2.1315	0.8342	
1.5283	0.8679	1.9204	0.1550	2.1395	0.9576	
1.5367	0.9134	1.9290	0.2197	2.1476	1.0000	
1.5451	0.9513	1.9376	0.3010	2.1556	0.9943	
1.5535	0.9749	1.9462	0.3698	2.1636	0.9611	
1.5619	0.9783	1.9548	0.4358	2.1716	0.9211	
1.5703	0.9748	1.9634	0.5021	2.1797	0.8799	
1.5787	0.9716	1.9719	0.5988	2.1877	0.8501	
1.5871	0.9743	1.9805	0.7119	2.1957	0.8153	
1.5955	0.9754	1.9891	0.7655	2.2037	0.7885	
1.6040	0.9788	1.9977	0.7436	2.2118	0.7550	
1.6124	0.9818	2.0063	0.6658	2.2198	0.7103	
1.6208	0.9823	2.0149	0.6268	2.2278	0.6745	
1.6292	0.9825	2.0234	0.6942	2.2359	0.6473	
1.6376	0.9843	2.0320	0.8139	2.2439	0.6027	
1.6460	0.9844	2.0406	0.8905	2.2519	0.5747	
1.6544	0.9823	2.0492	0.8909	2.2599	0.5509	
1.6628	0.9831	2.0578	0.8625	2.2680	0.5300	
1.6712	0.9828	2.0664	0.8599	2.2760	0.5258	
1.6796	0.9798	2.0749	0.8955	2.2840	0.5178	
1.6881	0.9738	2.0835	0.9502	2.2920	0.4992	
1.6965	0.9688	2.0921	0.9905	2.3001	0.4973	
1.7049	0.9701	2.1007	1.0000	2.3081	0.4991	
1.7133	0.9796	2.1093	0.9896	2.3161	0.4899	
1.7217	0.9927	2.1179	0.9688	2.3241	0.4854	
1.7301	0.9997	2.1264	0.9526	2.3322	0.4728	
1.7385	1.0000	2.1350	0.9347	2.3402	0.4586	
1.7469	0.9951	2.1436	0.9100	2.3482	0.4435	
1.7553	0.9794	2.1522	0.8776	2.3563	0.4308	
1.7638	0.9373	2.1608	0.8242	2.3643	0.4158	
1.7722	0.8557	2.1694	0.7436	2.3723	0.3959	
1.7806	0.7153	2.1779	0.6285	2.3803	0.3586	
1.7890	0.5268	2.1865	0.4750	2.3884	0.3013	
1.7974	0.3366	2.1951	0.3185	2.3964	0.2272	

diagram devoid of any comparable objects. While similar in L' luminosity to field late L-dwarfs, HD 95086 b is over a magnitude redder, and while having a similar  $K_1 - L'$  color to field late T-dwarfs, it is over two magnitudes brighter. Comparing to young substellar objects, HD 95086 b is most analogous to 2M1207 b, with a similar L' luminosity, but almost a magnitude redder in  $K_1 - L'$ . 2M1207 b was itself noted as being extremely red (Chauvin et al. 2004), which was explained by invoking thick clouds and non-equilibrium chemistry (Barman et al. 2011b; Skemer et al. 2014). The unusually red color of HD 95086 b may also be explained by

Table 6Empirical Filter Transformations

Color	$c_0$	$c_1$	<i>c</i> <sub>2</sub>	<i>c</i> <sub>3</sub>	$c_4$	C5	rms
		$\times 10^{-2}$	$\times 10^{-3}$	$\times 10^{-4}$	$ imes 10^{-6}$	$\times 10^{-7}$	
2MASS H – GPI H	0	-0.32073	0.22029	-0.13696	0.44056	-0.04090	0.003
MKO $H - \text{GPI } H$	0	0.52976	-0.45040	0.35256	-1.58498	0.26114	0.004
2MASS $K_{\rm S}$ – GPI $K_1$	0	-4.06008	5.24441	-5.47733	26.8625	-4.16143	0.031
2MASS $K_{\rm S}$ – GPI $K_2$	0	2.55246	-3.00911	2.23060	-7.68838	0.55557	0.027
MKO $K$ – GPI $K_1$	0	-5.81313	8.81710	-8.88431	41.3378	-6.23493	0.040
MKO $K$ – GPI $K_2$	0	0.74663	0.70976	-1.31673	7.34113	-1.59507	0.019
MKO $K'$ – GPI $K_1$	0	-2.53922	3.55111	-3.81046	18.9157	-2.91471	0.023
MKO $K'$ – GPI $K_2$	0	3.98467	-4.42024	3.59249	-14.2935	1.59708	0.033
MKO $K_{\rm S}$ – GPI $K_1$	0	-3.59593	4.61246	-4.85141	23.9600	-3.73692	0.027
MKO $K_{\rm S}$ – GPI $K_2$	0	2.90088	-3.28302	2.47902	-8.95925	0.73357	0.030
MKO $L' - WISE W1$	0	0.99889	-15.2677	15.7383	-65.6107	9.14685	0.113

Note. Polynomial coefficients are defined as  $y = \sum_i c_i x^i$ , where y is the color transform, x is the spectral type. Spectral types are defined such that M0 = 0, L0 = 10, T0 = 20.

Table 7
Properties of Young Low-Mass Companions and Isolated Young and/or Dusty Brown Dwarfs used as Comparison Objects

Name	Age (Myr)	Spectral Type	T <sub>eff</sub> (K)	$\log g$ (dex)	$\log L/L_{\odot}$ (dex)	Mass (M <sub>Jup</sub> )	References
HIP 78530 B	5	$\mathrm{M7}_{\!eta}\pm0.5$	$2700\pm100$	$4.5\pm1.0$	$-2.53 \pm 0.09$	$22 \pm 1$	(1), (2)
GSC 06214-0010 B	5	$M9_{\gamma} \pm 0.5$	$2300\pm100$	$3\pm0.5$	$-3.01 \pm 0.09$	$15 \pm 1$	(2), (3), (4)
1 RXS J160929.1-210524b	5	$L4_{\gamma} \pm 1$	$1700\pm100$	$3.5\pm0.5$	$-3.5\pm0.2$	$8 \pm 1$	(2), (5)
2MASS J120734-393253b	8	$>$ L5 $_{\gamma}$	$1600\pm100$	$4.0\pm0.5$	$-4.72\pm0.14$	$8\pm2$	(6), (7), (8), (9)
2MASS J11193254-1137466 PSO J318.5338-22.8603	$10 \pm 3 \\ 12^{+8}_{-4}$	L7 $L7_{\gamma} \pm 1$	$1160^{+30}_{-40}$	$3.86^{+0.10}_{-0.08}$	$$ $-4.42\pm0.06$	5-6 $6.5^{+1.3}_{-1.0}$	(10) (11)
HIP 106906b	$13\pm2$	$L2.5 \pm 1$	$1800\pm100$		$-3.64\pm0.08$	$11 \pm 2$	(12)
HR 8799b	30	L–T	750 - 1100	$4.0\pm0.5$	$-5.1\pm0.1$	~5	(13), (14), (15)
HR 8799c	30	L–T	$1100\pm100$	~4.0	$-4.7\pm0.1$	~7	(14), (15)
HR 8799d	30	L–T	$1100\pm100$	~4.0	$-4.7\pm0.1$	~7	(14), (15)
HR 8799 e	30	L–T	$1100 \pm 100$	~4.0	$-4.7\pm0.2$	~7	(16)
2MASS J21481628+4003593		L6 <sup>a</sup>	1500	3.0	$-4.07 \pm 0.12$		(17), (18)
2MASS J22443167+2043433	•••	L6.5 <sup>a</sup>		•••			(19)

#### Note.

<sup>a</sup> Optical spectral type.

**References.** (1) Lafrenière et al. (2011), (2) Lachapelle et al. (2015), (3) Ireland et al. (2011), (4) Bowler et al. (2011), (5) Lafrenière et al. (2010), (6) Chauvin et al. (2004), (7) Patience et al. (2010), (8) Faherty et al. (2013), (9) Mohanty et al. (2007), (10) Kellogg et al. (2015), (11) Liu et al. (2013), (12) Bailey et al. (2014), (13) Marois et al. (2008), (14) Barman et al. (2011a), (15) Marley et al. (2012), (16) Marois et al. (2010), (17) Looper et al. (2008), (18) Witte et al. (2011), (19) Dahn et al. (2002).

the presence of circumplanetary material, however further observations sensitive to emission from accretion are require to test this hypothesis.

Comparing the near-infrared portion of the SED of HD 95086 b to field brown dwarfs within the SpeX Prism and IRTF Spectral Libraries, a best fit was found at L7, although the relatively sparse coverage of the SED for HD 95086 b led to a broad minimum spanning between L1 and T3 (Figure 6). The best fitting objects within the library were the brown dwarfs 2M2244+20 and 2M2148+40, which have been identified as possessing thick clouds, and in the case of 2M2244+20, may have low surface gravity. The full SED of HD 95086 b was also compared with the predictions of grids of model atmospheres spanning a range of temperatures, surface gravities, and cloud properties (Figures 8 and 9, and Table 3). Morphologically, the SED was best fit by the models that incorporate a high photospheric dust content and was poorly fit by those with clear atmospheres. Considering the best fits

within the dusty atmosphere grids, the effective temperature of HD 95086 b was constrained to between 800–1300 K, with a surface gravity of  $\log g \leq 4.5$ . A higher surface gravity field brown dwarf at the same effective temperature would have already transitioned to a clear photosphere, within which dust particulates would have rained out to the lower atmosphere. This places HD 95086 b in the extension to the L-dwarf sequence occupied by young low surface gravity substellar companions that have retained their cloudy atmospheres despite low effective temperatures (e.g., Skemer et al. 2012).

HD 95086 b occupies a region of the CMD within which no comparable object exists, with an SED consistent with an atmosphere dominated by thick clouds. Future observations of HD 95086 b to perform near-infrared spectroscopy at *H* and  $K_2$  (covering the red half of the *K*-band atmospheric window) and to obtain thermal-infrared photometry at 3.3  $\mu$ m, will help distinguish between the predictions of cloudy and cloud-free models. Given its extremely red color, and the diagnostic

power of thermal-infrared measurements, HD 95086 b also represents an ideal target for a future thermal-infrared integral field spectrograph in the Southern hemisphere, similar to those already being commissioned in the North (Skemer et al. 2015).

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#### APPENDIX

#### A.1. GPI Filter Characterization

As the photometric measurements of HD 95086 b used in this study were made using non-standard filters (GPI H,  $K_1$ , and NaCo L'), the relative spectral response of each filter was obtained in order to extract synthetic photometry from model atmospheres, and a color transformation was calculated to convert between these filters, and the more standard MKO, 2MASS, and *WISE* photometric systems. For the GPI filters, the transmission curve and instrument throughput were measured using coronagraphic H,  $K_1$ , and  $K_2$  observations of the white dwarf companion to HD 8049 (Zurlo et al. 2013), and are plotted in Figure 11 and tabulated in Table 5. For the NaCo L' filter, the transmission curve given in the instrument manual was multiplied by a SkyCalc<sup>37</sup> atmospheric transmission model for Cerro Paranal (Noll et al. 2012; Jones et al. 2013),



**Figure 11.** The normalized relative spectral response  $S(\lambda)\lambda$  of the *H*,  $K_1$ , and  $K_2$  GPI filters (blue solid, green solid, and red dashed curves), measured from coronagraphic observations of the white dwarf companion to HD 8049 (Zurlo et al. 2013).  $S(\lambda)\lambda$  for the 2MASS *H* and  $K_5$  filters are also plotted (black dotted–dashed curves; Cohen et al. 2003). As in Tokunaga & Vacca (2005),  $S(\lambda)$  includes the wavelength-dependent terrestrial atmospheric absorption,  $T(\lambda)$ , an example of which is plotted for Cerro Pachón, assuming an airmass of 1.5 and a vertical water vapor column of 4.3 mm (thin gray curve; Lord 1992).

computed at the average airmass of the NaCo observations of HD 95086 b, assuming a seasonal average for the precipitable water vapor. The instrument throughput of NaCo was assumed to be constant over the L' band-pass. For the 2MASS and *WISE*, the relative spectral response of the filters were obtained from Cohen et al. (2003) and Wright et al. (2010), respectively. The zero point of each filter was estimated by integrating the product of the flux-calibrated spectrum of Vega  $F_{\lambda}(\lambda)$  (Bohlin 2014)<sup>38</sup> and the relative spectral response  $S(\lambda)\lambda$  of the filter, divided by the integral of the relative spectral response, i.e.,:

$$F_{\lambda} = \frac{\int F_{\lambda}(\lambda) S(\lambda) \lambda d\lambda}{\int S(\lambda) \lambda d\lambda}$$
(5)

where  $F_{\lambda}$  is the flux density of Vega for the filter, which is defined to have a magnitude of zero. The adopted zero points were: GPI *H*: 1.151 × 10<sup>-9</sup> Wm<sup>-2</sup>  $\mu$ m<sup>-1</sup>,  $K_1$ : 5.040 × 10<sup>-10</sup> Wm<sup>-2</sup>  $\mu$ m<sup>-1</sup>,  $K_2$ : 3.797 × 10<sup>-10</sup> Wm<sup>-2</sup>  $\mu$ m<sup>-1</sup>, and NaCo *L'*: 5.127 × 10<sup>-11</sup> Wm<sup>-2</sup>  $\mu$ m<sup>-1</sup>. An uncertainty of 2% on the zero point for each filter was assumed. The zero points for the 2MASS and *WISE* filters were consistent with the values reported in Cohen et al. (2003) and Wright et al. (2010).

Empirical color transformations were computed between; the 2MASS (Cohen et al. 2003) and MKO (Tokunaga et al. 2002) photometric systems, and the GPI H,  $K_1$ , and  $K_2$  filters; and between the MKO L' and WISE W1 filters. Empirical spectra of stars and brown dwarfs were obtained from the IRTF Spectral Library and the SpeX Prism Spectral Library. For each pair of

<sup>&</sup>lt;sup>37</sup> https://www.eso.org/observing/etc/skycalc

<sup>&</sup>lt;sup>38</sup> ftp://ftp.stsci.edu/cdbs/current\_calspec/alpha\_lyr\_stis\_008.fits



Figure 12. Empirical color transformation between the 2MASS  $K_S$  and GPI  $K_1$ filters, derived from objects within the IRTF Spectral Library (black squares), and the SpeX Prism Spectral Libraries (gray circles). The transformation is expressed as a fifth order polynomial over the MO-T8 spectral type range (red curve), the coefficients for which are given in Table 6.

filters for which a color transformation was required, the flux of the object in both filters was calculated (e.g., Bessell & Murphy 2012), converted into a magnitude using the zero points described previously, and subtracted to compute the color. For each filter, Vega was assumed to have a magnitude of zero. To assess the uncertainty on the color transformation for a given object, the process was repeated  $10^{\circ}$  times, each time drawing randomly from the uncertainty on the zero points of each filter, and the uncertainty on the flux of the object. This process was repeated for each object in both libraries, leading to an empirical color transformation as a function of spectral type, an example of which is shown in Figure 12. For each object, the optical spectral type was preferentially used, although in some cases only an infrared estimate was available. A fifth order polynomial was fit to the color transformation for each filter pair over the M0-T8 spectral type range. As the color transformation was found to be negligible for stars of spectral type M0 and earlier, the first term of the polynomial was fixed to zero, such that the fit passed through zero for MO spectral type. An example fit is shown in Figure 12, and the polynomial coefficients for each filter pair are given in Table 6.

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