

# Spectroscopic Studies of Electronically-Active Defects in Transition Metal Oxides for Advanced Si Devices

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**Abstract.** Based on spectroscopic studies, and guided by ab initio theory, the electron and hole traps in HfO<sub>2</sub> and other transition metal elemental oxides e.g., TiO<sub>2</sub>, are assigned to O-atom divacancies, clustered at internal grain boundaries. *Engineering solutions* for defect reduction are identified: i) deposition of ultra-thin, < 2 nm HfO<sub>2</sub> and phase separated Hf silicate dielectrics, in which grain boundary formation is suppressed by effectively eliminating inter-primitive unit cell  $\pi$ -bonding interactions, and ii) non-crystalline Zr/Hf Si oxynitrides without nanocrystalline grain boundaries.

**Keywords:** transition metal oxides, x-ray absorption spectroscopy, grain-boundary defects, engineering solutions

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## INTRODUCTION

Densities of electrically active defects in high-k dielectrics are typically about 10-100 time higher than in Si-SiO<sub>2</sub> devices. A crucial issue is to determine if these defects are intrinsic or if they are derived from processing, e.g., chemical impurities. Spectroscopic studies of high-k gate dielectrics are used to identify electron and hole trapping states in HfO<sub>2</sub> gate stack dielectrics [1]. An asymmetry in hole and electron trapping is a significant limitation for operation and reliability of Si integrated circuits. Based on spectroscopic studies, asymmetries are explained by different charge states of O-divacancies clustered at grain boundaries in nano-crystalline (na-) films. These films are prepared at 30-300°C by remote plasma processing, RPECVD, and reactive evaporation, RE, and processed at higher temperatures (700-900°C) for device integration. Spectroscopic studies have also been performed on na-TiO<sub>2</sub> and -ZrO<sub>2</sub>, and complex oxides, e.g., La<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>, with similar results.

## EXPERIMENTAL RESULTS

Soft x-ray photoelectron spectroscopy (SXPS), near edge x-ray absorption spectroscopy (NEXAS), visible and vacuum ultra-violet spectroscopic ellipsometry (vis-VUV SE), and vacuum ultra-violet photoemission spectroscopy (UPS) have been applied to na-Ti, -Zr and -Hf elemental dioxides. Figure 1(a) presents SXPS spectra for HfO<sub>2</sub> and TiO<sub>2</sub>, indicating

occupied defect states above the respective valence band edges. Figures 2(a) and (b) compared vis-VUV SE  $\epsilon_2$  spectra for TiO<sub>2</sub> and ZrO<sub>2</sub>, each displaying features ~1 to 3 eV below the respective d-state band edges [2]. A comparison between the energies of d-state features extracted from the second derivative of spectra of the TiO<sub>2</sub> O K<sub>1</sub> spectrum, with energies of d-state features in the TiO<sub>2</sub>  $\epsilon_2$  spectrum yields a linear scaling relationship with a slope of 1 providing the rationale for using d-state energies extracted from O K<sub>1</sub> edges of HfO<sub>2</sub> and ZrO<sub>2</sub>, in combination with defect state features in the XPS valence band and VUV SE  $\epsilon_2$  spectra, to locate the energies of defect states within the forbidden energy gaps of TiO<sub>2</sub>, ZrO<sub>2</sub> and HfO<sub>2</sub>. The correspondence between d-state features in O K<sub>1</sub> and  $\epsilon_2$  spectra is supported by hole localization in the respective final states: i) a core hole in the O 1s state, and ii) a localized hole in previously doubly occupied valence band edge O-atom 2p  $\pi$  state. The results of this analysis have been used to create electron energy level diagrams for HfO<sub>2</sub> and TiO<sub>2</sub> that are in excellent agreement with molecular orbital (MO) calculations applied to transition metal ion clusters [3]. Spectroscopic data for ZrO<sub>2</sub> yields essentially the same defect states as for HfO<sub>2</sub>, with relatively small energy shifts. Photoconductivity and cathode-luminescence data support the intrinsic nature of these defects [2].

Experiments by our group [4], yield essentially the same asymmetric trapping as in Ref. 1. These traps are present in films prepared by RE and RPECVD, and

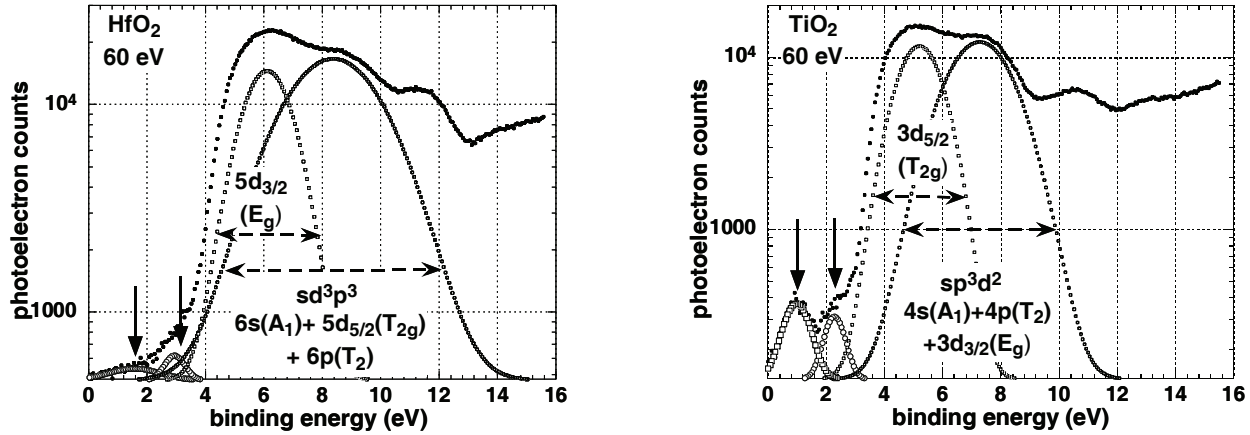


FIGURE 1. SXPS valence band spectra for na-HfO<sub>2</sub> and na-TiO<sub>2</sub>, prepared by reactive evaporation.

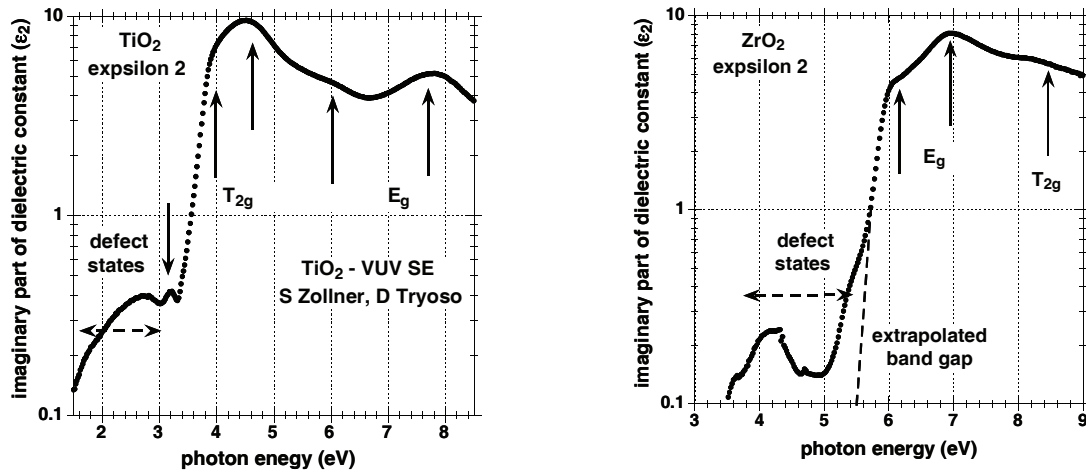


FIGURE 2.  $\epsilon_2$  spectra from vis-VUV-SE for (a) na-TiO<sub>2</sub> and (b) na-ZrO<sub>2</sub> films prepared by reactive evaporation.

atomic layer deposition, ALD, as well. Analysis of capacitance-voltage (C-V) and current voltage (J-V) indicate: i) relatively shallow electron traps,  $\sim 0.5$  eV below the conduction band edges of HfO<sub>2</sub> and ZrO<sub>2</sub>, as well as ii) deep hole traps below the valence band edge of Si;  $\sim 3$  eV above the HfO<sub>2</sub> valence band.

## DISCUSSION

The nature of the intrinsic defects in na-films has been identified by comparing MO calculations for octahedrally coordinated Ti<sup>4+</sup> atoms in TiO<sub>2</sub> with Ti<sup>3+</sup> in Ti(H<sub>2</sub>O)<sub>6</sub><sup>3+</sup> clusters and Ti<sub>2</sub>O<sub>3</sub>, as well [3]. The defect electronic structure is *consistent with* the trap being a divacancy. Initial comparisons were based on crystal field splittings, but these have been extended to include E<sub>g</sub> and T<sub>2g</sub> state degeneracy removal by Jahn-Teller (J-T) distortions. The inherent properties of the divacancy defects are: i) electronic states of Ti<sup>3+</sup> that

fall within the energy gap of the Ti<sup>4+</sup>-O states in TiO<sub>2</sub>, and J-T degeneracy removal for ii) partially-occupied T<sub>2g</sub> states at the valence band edge, and iii) unoccupied E<sub>g</sub> states at the conduction band edge. The T<sub>2g</sub> states of Ti<sup>3+</sup> are partially occupied and act as traps for substrate hole injection, whilst E<sub>g</sub> states are empty, and act as traps for transport via substrate injection of electrons [1], and sub-band-gap photoconductivity [2].

Similar energy level schemes apply to i) Zr(Hf)<sup>3+</sup> defects at grain-boundary clustered divacancies, and ii) Zr(Hf)<sup>4+</sup> states in bulk Zr(Hf)O<sub>2</sub> where the symmetries of the defect d-states are reversed: occupied states near the valence band are J-T term-split E<sub>g</sub> states, and those near the conduction band, J-T term split T<sub>g</sub> states.

*Engineering solutions* for elimination of these intrinsic defects have been identified. Two involve suppression of defect states in ultra-thin HfO<sub>2</sub> films, and chemically-phase separated high HfO<sub>2</sub> content silicates. The third uses the deposition of non-crystalline Si<sub>3</sub>N<sub>4</sub>-rich Ti(Zr,Hf) Si oxynitrides: e.g.,

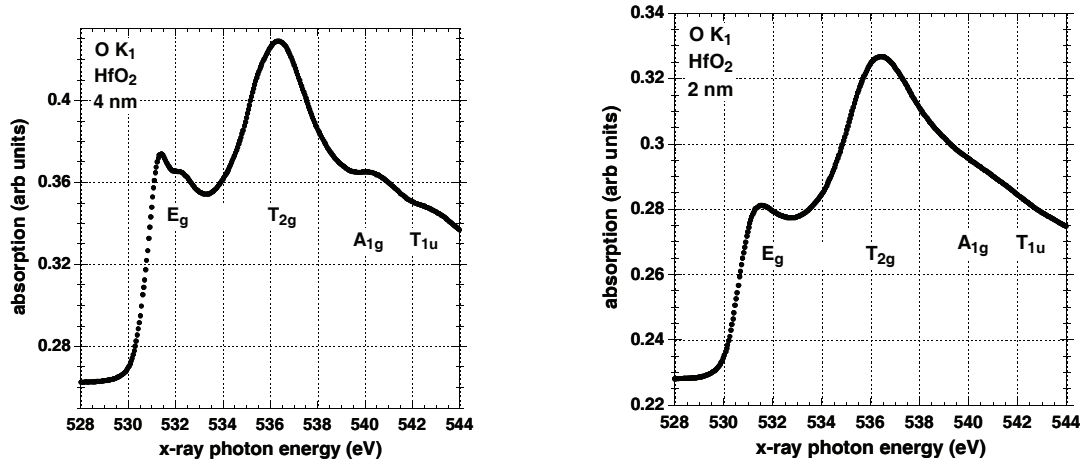


FIGURE 3. O  $K_1$  edge spectra for na-HfO<sub>2</sub>: (a) 4 nm thick film, and (b) 2 nm thick film.

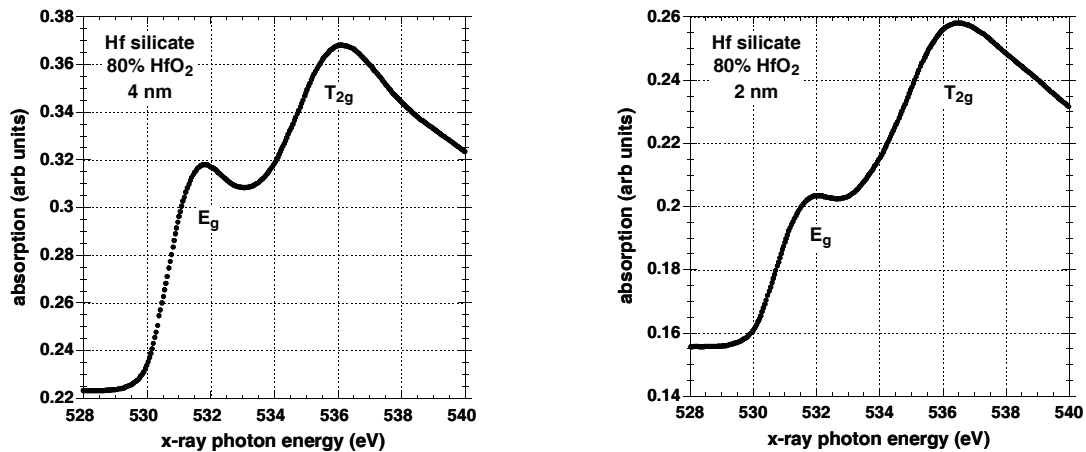


FIGURE 4. O  $K_1$  edge spectra of phase separated Hf silicates (80% HfO<sub>2</sub>). (a) 4 nm thick film, and (b) 2 nm thick film.

[Si<sub>3</sub>N<sub>4</sub>] ~33-40%, with approximately equal [SiO<sub>2</sub>] and [Ti(Zr,Hf)O<sub>2</sub>] concentrations of 33 to 30%.

In i) ultra-thin [ $<2$  nm] nc-HfO<sub>2</sub>, and ii) high HfO<sub>2</sub> (~80%) phase separated Hf silicates, J-T splitting of band edge  $E_g$  ( $5d_{3/2}$ ) anti-bonding  $\pi$ -states is suppressed by decoupling  $\pi$ -bonding interactions between strings of metal and oxide atoms in neighboring *primitive unit cells* (Figs. 3(a) and (b) and 4(a) and (b)). The ~2 nm scale of order is too small for grain boundaries; however, broadening of band edge states occurs, and may effect transport properties.

Studies of O  $K_1$  and N  $K_1$  edges of Zr Si oxynitrides indicate 4-fold coordinated Zr with O nearest-neighbors, and Si, O and N bonding similar to Si<sub>2</sub>ON<sub>2</sub>. Si<sub>3</sub>N<sub>4</sub> rich Zr/Hf Si oxynitrides are stable against phase separation up to 1100°C [5].

Electrical measurements show that ultra-thin HfO<sub>2</sub> annealed in N<sub>2</sub> or NH<sub>3</sub> at 700°C, and Zr/Hf Si

oxynitrides, have low defect densities, and additionally can be scaled to EOT values of 0.7-0.8 nm, thereby extending device scaling below the 0.45 nm node [6].

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