

STATISTICS ON RIEMANNIAN MANIFOLDS: ASYMPTOTIC DISTRIBUTION AND CURVATURE

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ABSTRACT. In this article a nonsingular asymptotic distribution is derived for a broad class of underlying distributions on a Riemannian manifold in relation to its curvature. Also, the asymptotic dispersion is explicitly related to curvature. These results are applied and further strengthened for the planar shape space of k -ads.

1. INTRODUCTION

Statistical analysis of a probability measure Q on a differentiable manifold M has diverse applications in directional and axial statistics, morphometrics, medical diagnostics and machine vision ([1, 2, 3, 4, 6, 7, 8, 11, 15]). Most of this analysis focuses on nonintrinsic Fréchet means of Q . In this article we provide a distribution theory for the nonparametric analysis of intrinsic means which can be directly used for the one- and two-sample problems. To be precise, let (M, g) be a Riemannian manifold with metric tensor g and geodesic distance d_g . Define the **Fréchet function** F of Q as

$$(1.1) \quad F(p) = \int_M d_g^2(p, m)Q(dm), \quad p \in M.$$

Assume F to be finite. We consider probability measures Q whose support $\text{supp}(Q)$ are contained in **geodesic balls** $B(p, r) = \{m : d_g(p, m) < r\}$. If the Fréchet function, restricted to such a ball $B(p, r)$, has a unique minimizer μ_I in $B(p, r)$, we call it the **intrinsic mean of Q in $B(p, r)$** . The **sample intrinsic mean μ_{nI} in $B(p, r)$** is the intrinsic mean of $Q_n = \frac{1}{n} \sum_{j=1}^n \delta_{X_j}$ in $B(p, r)$, where X_1, X_2, \dots, X_n are independent and identically distributed (iid) observations from the underlying distribution Q . Crucial to nonparametric analysis is the asymptotic distribution of μ_{nI} . Our main goals are (i) to derive this asymptotic distribution, assuring its nonsingularity, under as broad a condition on $\text{supp}(Q)$ as possible, (ii) to explicitly compute the asymptotic dispersion, and (iii) to apply and refine the general theory to the particularly important **planar shape space** Σ_2^k of k landmarks introduced by Kendall [11].

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To indicate the role curvature plays in this endeavor, let $r_* = \min\{inj(M), \frac{\pi}{\sqrt{\bar{C}}}\}$, where \bar{C} is an upper bound of sectional curvatures of M if this upper bound is positive, and $\bar{C} = 0$ otherwise. Also, $inj(M) \equiv \inf\{d_g(p, C(p)) : p \in M\}$ is the **injectivity radius** of M , where $C(p)$ is the **cut locus** of p , i.e., the set of points of the form $\gamma(t_0)$, where γ is a geodesic and t_0 is the supremum of all $t > 0$ such that the geodesic from p to $\gamma(t)$ is distance minimizing. The **exponential map** exp_p is injective on $\{v \in T_p(M) : |v| < r\}$ if and only if $r \leq r_*$ (Do Carmo [5], p. 271). It follows that a geodesic ball $B(p, r)$ with $r \leq \frac{r_*}{2}$ is **strongly convex**, i.e., for every pair $q, q' \in B(p, r)$ there exists a unique geodesic connecting q, q' , entirely contained in $B(p, r)$, this geodesic being distance minimizing. By Proposition 2.1 and Theorem 2.2, if $supp(Q) \subseteq B(p, \frac{r_*}{2})$, then Q has a unique intrinsic mean μ_I in $B(p, \frac{r_*}{2})$. If in addition, $supp(Q) \subseteq B(\mu_I, \frac{r_*}{2})$, then the sample intrinsic mean has asymptotic normal distribution. Further, in the case of manifolds with constant sectional curvature, the asymptotic dispersion can be explicitly expressed in terms of curvature.

It may be noted that our results are not related to those of Pennek [16] who has a number of interesting results on distributions on manifolds, including one that provides an expansion of the density of the (analog of) normal distribution on the manifold in terms of its variance, for the case of small variance.

For background in differential geometry used here, we refer to Do Carmo [5] and Lee [14].

2. ASYMPTOTIC DISTRIBUTION AND CURVATURE

Let (M, g) be a Riemannian manifold. We continue to use the notation of Section 1. Let Q be a probability measure, $supp(Q) \subseteq B(p, \frac{r_*}{2})$ for some p . Then there is a unique (local) intrinsic mean μ_I in $B(p, \frac{r_*}{2})$ (Kendall [12]). This substantially extends Karchar's result on the existence and uniqueness of a local mean ([9]), but not his important result on the strict convexity of F . We are able to circumvent this difficulty in the case $supp(Q) \subseteq B(\mu_I, \frac{r_*}{2})$. Denote by μ_{nI} the intrinsic mean of Q_n in $B(p, \frac{r_*}{2})$. The inverse of the exponential map $\phi = exp_p^{-1}$ is a diffeomorphism on $B(p, \frac{r_*}{2})$ onto its image, say U , in $T_p(M)$. The image $\tilde{Q} = Q \circ \phi^{-1}$ of Q under ϕ is a probability measure in $T_p(M)$, and the image $\mu = \phi(\mu_I)$ of μ_I is the minimizer of

$$(2.1) \quad \tilde{F}(x) = \int_U d_g^2(\phi^{-1}x, \phi^{-1}y)\tilde{Q}(dy), \quad x \in U.$$

Similarly $\mu_n = \phi(\mu_{nI})$ is the corresponding minimizer when \tilde{Q} is replaced by $\tilde{Q}_n = \frac{1}{n} \sum_{j=1}^n \delta_{\phi(X_j)}$. As proved in [3], Theorem 2.1, a central limit theorem for the **M-estimator** μ_n may be derived and used to obtain the following result. The **normal coordinates** x, y used here are with respect to a chosen orthonormal basis in T_pM .

Proposition 2.1. *Suppose the support of Q is contained in the geodesic ball $B = B(p, \frac{r_*}{2})$. Let $\phi = exp_p^{-1} : B \rightarrow \phi(B)$. Define $h(x, y) = d_g^2(\phi^{-1}x, \phi^{-1}y)$; $x, y \in \phi(B)$. Let $((D_r h))_{r=1}^d$ and $((D_r D_s h))_{r,s=1}^d$ be the matrices of first and second order derivatives of $y \mapsto h(x, y)$. Let $\tilde{X}_j = \phi(X_j)$; $j = 1, \dots, n$, X_1, \dots, X_n being iid observations from Q . Define*

$$\Lambda = E((D_r D_s h(\tilde{X}_1, \mu))_{r,s=1}^d), \quad \Sigma = Cov((D_r h(\tilde{X}_1, \mu))_{r=1}^d).$$

If Λ is nonsingular, then

$$(2.2) \quad \sqrt{n}(\mu_n - \mu) \xrightarrow{\mathcal{L}} N(0, \Lambda^{-1}\Sigma\Lambda^{-1}).$$

The natural candidate for p in Proposition 2.1 is the intrinsic mean of Q in $B(p, \frac{r_*}{2})$, namely μ_I . Then we get expressions for Λ and Σ using an orthonormal basis in $T_{\mu_I}M$. Theorem 2.2 below gives a lower bound on Λ and an exact expression when M has constant sectional curvature. The lower bound gives a condition on the nonsingularity of Λ . The nonsingularity of Σ is a milder condition which holds, for example, when Q has a density with respect to the volume measure on M . In the statement of the theorem, the usual partial order $A \geq B$ between $d \times d$ symmetric matrices A, B , means that $A - B$ is nonnegative definite.

Theorem 2.2. *Assume $\text{supp}(Q) \subseteq B(p, \frac{r_*}{2})$. Let $\phi = \exp_{\mu_I}^{-1} : B(p, \frac{r_*}{2}) \rightarrow T_{\mu_I}M (\approx \mathbb{R}^d)$, and let C denote an upper bound of all sectional curvatures. Then in normal coordinates with respect to a chosen orthonormal basis in $T_{\mu_I}M$,*

$$(2.3) \quad D_r h(x, 0) = -2x^r, \quad 1 \leq r \leq d,$$

$$(2.4) \quad [D_r D_s h(x, 0)] \geq \left[2 \left(\frac{1 - f(|x|)}{|x|^2} \right) x^r x^s + f(|x|) \delta_{rs} \right]_{1 \leq r, s \leq d},$$

$$\text{where } |x| = \sqrt{(x^1)^2 + (x^2)^2 + \dots + (x^d)^2},$$

$$(2.5) \quad f(x) = \begin{cases} 1 & \text{if } C = 0, \\ \sqrt{C}x \frac{\cos(\sqrt{C}x)}{\sin(\sqrt{C}x)} & \text{if } C > 0, \\ \sqrt{-C}x \frac{\cosh(\sqrt{-C}x)}{\sinh(\sqrt{-C}x)} & \text{if } C < 0. \end{cases}$$

There is equality in (2.4) when M has constant sectional curvature C , and in this case Λ has the expression

$$(2.6) \quad \Lambda_{rs} = 2E \left(\frac{1 - f(|\tilde{X}_1|)}{|\tilde{X}_1|^2} \right) \tilde{X}_1^r \tilde{X}_1^s + (f(|\tilde{X}_1|)) \delta_{rs}, \quad 1 \leq r, s \leq d,$$

Λ being positive definite if Q has support in $B(\mu_I, \frac{r_*}{2})$.

Proof. Let $\gamma(s)$ be a geodesic, $\gamma(0) = \mu_I$. Define $c(s, t) = \exp_m(t \exp_m^{-1} \gamma(s))$, $s \in [0, \epsilon]$, $t \in [0, 1]$, as a smooth variation of γ through geodesics lying entirely in $B(p, \frac{r_*}{2})$. Let $T = \frac{\partial}{\partial t} c(s, t)$, $S = \frac{\partial}{\partial s} c(s, t)$. Since $c(s, 0) = m$, $S(s, 0) = 0$, and since $c(s, 1) = \gamma(s)$, $S(s, 1) = \dot{\gamma}(s)$. Also $\langle T, T \rangle = d_g^2(\gamma(s), m)$ is independent of t , and the covariant derivative $D_t T$ vanishes because $t \mapsto c(s, t)$ is a geodesic (for each s). Then

$$(2.7) \quad d_g^2(\gamma(s), m) = \langle T(s, t), T(s, t) \rangle = \int_0^1 \langle T(s, t), T(s, t) \rangle dt.$$

Hence $d_g^2(\gamma(s), m)$ is C^∞ smooth, and using the symmetry of the connection on a parametrized surface (see Lemma 3.4, p. 68 in Do Carmo [5]), we get

$$(2.8) \quad \begin{aligned} \frac{d}{ds} d_g^2(\gamma(s), m) &= 2 \int_0^1 \langle D_s T, T \rangle dt = 2 \int_0^1 \frac{d}{dt} \langle T, S \rangle dt \\ &= 2 \langle T(s, 1), S(s, 1) \rangle = -2 \langle \exp_{\gamma(s)}^{-1} m, \dot{\gamma}(s) \rangle. \end{aligned}$$

Substituting $s = 0$ in (2.8), we get expressions for $D_r h(x, 0)$ as in (2.3). Also

$$(2.9) \quad \frac{d^2}{ds^2} d_g^2(\gamma(s), m) = 2\langle D_s T(s, 1), S(s, 1) \rangle$$

$$(2.10) \quad = 2\langle D_t S(s, 1), S(s, 1) \rangle = 2\langle D_t J_s(1), J_s(1) \rangle,$$

where $J_s(t) = S(s, t)$. Note that J_s is a Jacobi field along $c(s, \cdot)$ with $J_s(0) = 0$, $J_s(1) = \dot{\gamma}(s)$. Let J_s^\perp and J_s^- be the normal and tangential components of J_s . The relations (2.13)-(2.15) below may be obtained from Jost [10], p. 197, and Lee [14], Lemma 10.8. For the sake of exposition, we indicate the arguments here. Let η be a unit speed geodesic in M and J a normal Jacobi field along η , $J(0) = 0$. Define

$$(2.11) \quad u(t) = \begin{cases} t & \text{if } C = 0, \\ \frac{\sin(\sqrt{C}t)}{\sqrt{C}} & \text{if } C > 0, \\ \frac{\sinh(\sqrt{-C}t)}{\sqrt{-C}} & \text{if } C < 0. \end{cases}$$

Then $u''(t) = -Cu(t)$ and

$$(2.12) \quad (|J|'u - |J|u')'(t) = (|J|'' + C|J|)u(t).$$

By exact differentiation and Schwartz inequality, it is easy to show that $|J|'' + C|J| \geq 0$, hence $(|J|'u - |J|u')'(t) \geq 0$ whenever $u(t) \geq 0$. This implies that $|J|'u - |J|u' \geq 0$ if $t \leq t_0$, where u is positive on $(0, t_0)$. Also $|J|' = \frac{\langle J', J \rangle}{|J|}$. Therefore $\langle J(t), D_t J(t) \rangle \geq \frac{u'(t)}{u(t)} |J(t)|^2 \forall t < t_0$. If we drop the unit speed assumption on η , we get

$$(2.13) \quad \langle J(1), D_t J(1) \rangle \geq |\dot{\eta}| \frac{u'(|\dot{\eta}|)}{u(|\dot{\eta}|)} |J(1)|^2 \text{ if } |\dot{\eta}| < t_0.$$

Here $t_0 = \infty$ if $C \leq 0$ and equals $\frac{\pi}{\sqrt{C}}$ if $C > 0$. When M has constant sectional curvature C , $J(t) = u(t)E(t)$, where E is a parallel normal vector field along η . Hence

$$(2.14) \quad \langle J(t), D_t J(t) \rangle = u(t)u'(t)|E(t)|^2 = \frac{u'(t)}{u(t)} |J(t)|^2.$$

If we drop the unit speed assumption, we get

$$(2.15) \quad \langle J(t), D_t J(t) \rangle = |\dot{\eta}| \frac{u'(|\dot{\eta}|t)}{u(|\dot{\eta}|t)} |J(t)|^2.$$

Since J_s^\perp is a normal Jacobi field along the geodesic $c(s, \cdot)$, from (2.13) and (2.15) it follows that

$$(2.16) \quad \langle J_s^\perp(1), D_t J_s^\perp(1) \rangle \geq f(d(\gamma(s), m)) |J_s^\perp(1)|^2$$

with equality in (2.16) when M has constant sectional curvature C , f being defined in (2.5).

Next suppose J is a Jacobi field along a geodesic η , $J(0) = 0$ and let $J^-(t)$ be its tangential component. Then $J^-(t) = \lambda t \dot{\eta}(t)$ where $\lambda t = \frac{\langle J(t), \dot{\eta}(t) \rangle}{|\dot{\eta}|^2}$, λ being independent of t . Hence

$$(2.17) \quad \begin{aligned} (D_t J)^-(t) &= \frac{\langle D_t J(t), \dot{\eta}(t) \rangle}{|\dot{\eta}|^2} \dot{\eta}(t) \\ &= \frac{d}{dt} \left(\frac{\langle J(t), \dot{\eta}(t) \rangle}{|\dot{\eta}|^2} \right) \dot{\eta}(t) = \lambda \dot{\eta}(t) = D_t(J^-)(t) \end{aligned}$$

and

$$\begin{aligned} D_t|J^-|^2(1) &= 2\lambda^2|\dot{\eta}|^2 = 2\frac{\langle J(1), \dot{\eta}(1) \rangle^2}{|\dot{\eta}(1)|^2} \\ &= D_t\langle J, J^- \rangle(1) = \langle D_tJ(1), J^-(1) \rangle + |J^-(1)|^2 \end{aligned}$$

which implies

$$(2.18) \quad \langle D_tJ(1), J^-(1) \rangle = 2\frac{\langle J(1), \dot{\eta}(1) \rangle^2}{|\dot{\eta}(1)|^2} - |J^-(1)|^2 = \frac{\langle J(1), \dot{\eta}(1) \rangle^2}{|\dot{\eta}(1)|^2}.$$

Apply (2.17) and (2.18) to the Jacobi field J_s to get

$$(2.19) \quad D_t(J_s^-)(1) = (D_tJ_s)^-(1) = J_s^-(1) = \frac{\langle J_s(1), T(s, 1) \rangle}{|T(s, 1)|^2}T(s, 1),$$

$$(2.20) \quad \langle D_tJ_s(1), J_s^-(1) \rangle = \frac{\langle J_s(1), T(s, 1) \rangle^2}{|T(s, 1)|^2}.$$

Using (2.16), (2.19) and (2.20), (2.10) becomes

$$\begin{aligned} \frac{d^2}{ds^2}d_g^2(\gamma(s), m) &= 2\langle D_tJ_s(1), J_s(1) \rangle \\ &= 2\langle D_tJ_s(1), J_s^-(1) \rangle + 2\langle D_tJ_s(1), J_s^\perp(1) \rangle \\ &= 2\langle D_tJ_s(1), J_s^-(1) \rangle + 2\langle D_t(J_s^\perp)(1), J_s^\perp(1) \rangle \\ (2.21) \quad &\geq 2\frac{\langle J_s(1), T(s, 1) \rangle^2}{|T(s, 1)|^2} + 2f(|T(s, 1)|)|J_s^\perp(1)|^2 \\ &= 2\frac{\langle J_s(1), T(s, 1) \rangle^2}{|T(s, 1)|^2} + 2f(|T(s, 1)|)|J_s(1)|^2 \\ &\quad - 2f(|T(s, 1)|)\frac{\langle J_s(1), T(s, 1) \rangle^2}{|T(s, 1)|^2} \\ (2.22) \quad &= 2f(d_g(\gamma(s), m))|\dot{\gamma}(s)|^2 + 2(1 - f(d_g(\gamma(s), m)))\frac{\langle \dot{\gamma}(s), \exp_{\gamma(s)}^{-1}m \rangle^2}{d_g^2(\gamma(s), m)} \end{aligned}$$

with equality in (2.21) when M has constant sectional curvature C . Substituting $s = 0$ in (2.22), we get a lower bound for $[D_rD_s h(x, 0)]$ as in (2.4) and an exact expression for $D_rD_s h(x, 0)$ when M has constant sectional curvature. To see this, let $\dot{\gamma}(0) = v$. Then writing $m = \phi^{-1}(x)$, $\gamma(s) = \phi^{-1}(sv)$, one has

$$\begin{aligned} \frac{d^2}{ds^2}d_g^2(\gamma(s), m)|_{s=0} &= \frac{d^2}{ds^2}d_g^2(\phi^{-1}(x), \phi^{-1}(sv))|_{s=0} \\ (2.23) \quad &= \frac{d^2}{ds^2}h(x, sv)|_{s=0} = \sum_{r,s=1}^d v_r v_s D_r D_s h(x, 0). \end{aligned}$$

Since $d^2(\gamma(s), m)$ is twice continuously differentiable and Q has compact support, using the Lebesgue DCT, we get

$$(2.24) \quad \frac{d^2}{ds^2}F(\gamma(s))|_{s=0} = \int \frac{d^2}{ds^2}d^2(\gamma(s), m)|_{s=0}Q(dm).$$

Then (2.6) follows from (2.22). If $supp(Q) \subseteq B(\mu_I, \frac{r_*}{2})$, then the expression in (2.22) is strictly positive at $s = 0$ for all $m \in supp(Q)$, hence Λ is positive definite. This completes the proof. \square

Under the assumptions of Theorem 2.2, it follows that $E(\tilde{X}_1) = 0$ and $\Sigma = 4E(\tilde{X}_1\tilde{X}_1')$. This is also stated in Theorem 2.1 in Bhattacharya [2].

Remark 2.1. It may be noted that the spaces S^d , $\mathbb{R}P^d$ have constant positive curvature. One may also endow the projective shape space with a metric which makes it a space of constant positive curvature, since it is diffeomorphic to a product of real projective spaces ([15]). In the next section we turn to Σ_2^k , whose sectional curvatures range from 1 to 4.

3. APPLICATION TO THE PLANAR SHAPE SPACE Σ_2^k

Consider the planar shape space Σ_2^k of **k-ads** in \mathbb{R}^2 . An element of Σ_2^k is a set of k landmarks, or points in the plane (not all equal), modulo translation, rotation and scaling. Let S_2^k be the **pre-shape sphere** which is the space of column vectors in \mathbb{C}^k with mean 0 and norm 1. Its tangent space is

$$T_z S_2^k = \{v \in \mathbb{C}^k : v' \mathbf{1}_k = \operatorname{Re}(z' \bar{v}) = 0\}.$$

Here $\mathbf{1}_k$ denotes the column vector of ones of size k . To apply Proposition 2.1 to carry out nonparametric inference on Σ_2^k , we need to identify the exponential and inverse exponential maps on Σ_2^k . For that we consider their lifts to S_2^k as in Section 4 in Le [13], and Kendall [11]. The map $\pi : S_2^k \rightarrow \Sigma_2^k$,

$$z \mapsto \pi(z) = [z] = \{\lambda z : \lambda \in \mathbb{C}, |\lambda| = 1\},$$

is a Riemannian submersion. So the tangent space $T_{[z]}\Sigma_2^k$ is isometric with a subspace of $T_z S_2^k$ called the **horizontal subspace** H_z which is

$$H_z = \{v \in \mathbb{C}^k : z' \bar{v} = 0, v' \mathbf{1}_k = 0\}.$$

Denote the corresponding isometric mapping by $\chi_{[z]} : T_{[z]}\Sigma_2^k \rightarrow H_z$. Then $\exp_{[z]} = \pi \circ \exp_z \circ \chi_{[z]}$, and

$$(3.1) \quad \chi_{[z]} \circ \exp_{[z]}^{-1} : \Sigma_2^k \setminus C([z]) \rightarrow H_z, [w] \mapsto \frac{r}{\sin r} \{-z \cos r + e^{i\theta} w\},$$

$$(3.2) \quad r = d_g([z], [w]) = \arccos(|z' \bar{w}|) \in [0, \frac{\pi}{2}), e^{i\theta} = \frac{z' \bar{w}}{|z' \bar{w}|}.$$

In (3.1), $C([z])$ is the cut-locus of $[z]$, which is

$$C([z]) = \{[x] \in \Sigma_2^k : d_g([x], [z]) = \frac{\pi}{2}\} = \{[x] : z' \bar{x} = 0\}.$$

Σ_2^k has all sectional curvatures bounded between 1 and 4, and its injectivity radius is $\frac{\pi}{2}$. From a result due to Kendall [12], Q has an intrinsic mean if its support is contained in a geodesic ball of radius $\frac{\pi}{4}$. Suppose $\operatorname{supp}(Q) \subseteq B(p, \frac{\pi}{4})$ and let $\mu_I = [\mu]$ be the intrinsic mean of Q in the support, with μ being one of its pre-shapes. The following theorem gives the expression for Λ in Theorem 2.2 and derives a sufficient condition for its nonsingularity.

Theorem 3.1. *Let $\phi : B(p, \frac{\pi}{4}) \rightarrow \mathbb{C}^{k-2} (\approx \mathbb{R}^{2k-4})$ be the coordinates of $\chi_{\mu_I} \circ \exp_{\mu_I}^{-1} : B(p, \frac{\pi}{4}) \rightarrow H_\mu$ with respect to some orthonormal basis $\{v_1, \dots, v_{k-2}, iv_1, \dots, iv_{k-2}\}$ for H_μ . Define $h(x, y) = d_g^2(\phi^{-1}x, \phi^{-1}y)$. Let $((D_r h))_{r=1}^{2k-4}$ and $((D_r D_s h))_{r,s=1}^{2k-4}$ be the matrix of first and second order derivatives of $y \mapsto h(x, y)$. Let $\tilde{X}_j = \phi(X_j) = (\tilde{X}_j^1, \dots, \tilde{X}_j^{k-2})'$; $j = 1, \dots, n$, X_1, \dots, X_n being iid observations from Q . Define*

$\Lambda = E((D_r D_s h(\tilde{X}_1, 0))_{r,s=1}^{2k-4})$. Then Λ is positive definite if the support of Q is contained in $B(\mu_I, R)$, where R is the unique solution of $\tan(x) = 2x$, $x \in (0, \frac{\pi}{2})$.

Proof. For a geodesic γ starting at μ_I , write $\gamma = \pi \circ \tilde{\gamma}$, where $\tilde{\gamma}$ is a geodesic in S_2^k starting at μ . From the proof of Theorem 2.2, for $m = [z] \in B(p, \frac{\pi}{4})$,

$$(3.3) \quad \frac{d}{ds} d_g^2(\gamma(s), m) = 2\langle T(s, 1), \dot{\gamma}(s) \rangle = 2\langle \tilde{T}(s, 1), \dot{\tilde{\gamma}}(s) \rangle,$$

$$(3.4) \quad \frac{d^2}{ds^2} d_g^2(\gamma(s), m) = 2\langle D_s T(s, 1), \dot{\gamma}(s) \rangle = 2\langle D_s \tilde{T}(s, 1), \dot{\tilde{\gamma}}(s) \rangle,$$

where $\tilde{T}(s, 1) = \chi_{\gamma(s)}(T(s, 1))$. From (3.1), this has the expression

$$(3.5) \quad \tilde{T}(s, 1) = -\frac{\rho(s)}{\sin(\rho(s))} [-\cos(\rho(s))\tilde{\gamma}(s) + e^{i\theta(s)}z],$$

$$\text{where } e^{i\theta(s)} = \frac{\bar{z}'\tilde{\gamma}(s)}{\cos(\rho(s))}, \quad \rho(s) = d_g(\gamma(s), m).$$

The inner product in (3.3) and (3.4) is the Riemannian metric on TS_2^k which is $\langle v, w \rangle = \text{Re}(v'\bar{w})$. Observe that $D_s \tilde{T}(s, 1)$ is $\frac{d}{ds} \tilde{T}(s, 1)$ projected onto $H_{\tilde{\gamma}(s)}$. Since $\langle \mu, \dot{\tilde{\gamma}}(0) \rangle = 0$,

$$(3.6) \quad \frac{d^2}{ds^2} d_g^2(\gamma(s), m)|_{s=0} = 2\langle \frac{d}{ds} \tilde{T}(s, 1)|_{s=0}, \dot{\tilde{\gamma}}(0) \rangle.$$

From (3.5) we have

$$(3.7) \quad \begin{aligned} \frac{d}{ds} \tilde{T}(s, 1)|_{s=0} &= \left(\frac{d}{ds} \left(\frac{\rho(s)\cos(\rho(s))}{\sin(\rho(s))} \right) \Big|_{s=0} \right) \mu + \left(\frac{\rho(s)\cos(\rho(s))}{\sin(\rho(s))} \Big|_{s=0} \right) \dot{\tilde{\gamma}}(0) \\ &\quad - \left(\frac{d}{ds} \left(\frac{\rho(s)}{\sin(\rho(s))\cos(\rho(s))} \right) \Big|_{s=0} \right) (\bar{z}'\mu)z \\ &\quad - \left(\frac{\rho(s)}{\sin(\rho(s))\cos(\rho(s))} \Big|_{s=0} \right) (\bar{z}'\dot{\tilde{\gamma}}(0))z, \end{aligned}$$

and along with (3.3), we get

$$(3.8) \quad \frac{d}{ds} \rho(s)|_{s=0} = \frac{-1}{\sin(r)} \langle \dot{\tilde{\gamma}}(0), \frac{\bar{z}'\mu}{\cos(r)}z \rangle \quad (r := d_g(m, \mu_I)).$$

Hence

$$(3.9) \quad \begin{aligned} \langle \frac{d}{ds} \tilde{T}(s, 1)|_{s=0}, \dot{\tilde{\gamma}}(0) \rangle &= r \frac{\cos(r)}{\sin(r)} \|\dot{\tilde{\gamma}}(0)\|^2 - \left(\frac{1}{\sin^2 r} - r \frac{\cos(r)}{\sin^3(r)} \right) (\text{Re}x)^2 \\ &\quad + \frac{r}{\sin(r)\cos(r)} (\text{Im}x)^2, \end{aligned}$$

where

$$(3.10) \quad x = e^{i\theta} z' \overline{\dot{\tilde{\gamma}}(0)}, \quad e^{i\theta} = \frac{\bar{z}'\mu}{\cos r}.$$

The value of x in (3.10), and hence the expression in (3.9), depends on z only through $m = [z]$. Also if $\gamma = \pi(\gamma_1) = \pi(\gamma_2)$, γ_1 and γ_2 being two geodesics on S_2^k starting at μ_1 and μ_2 respectively, with $[\mu_1] = [\mu_2] = [\mu]$, then $\gamma_1(t) = \lambda\gamma_2(t)$, where $\mu_2 = \lambda\mu_1$, $\lambda \in \mathbb{C}$. Now it is easy to check that the expression in (3.9) depends

on μ only through $[\mu] = \mu_I$. Note that $|x|^2 < 1 - \cos^2 r$. So when $|\dot{\gamma}(0)| = 1$, (3.9) is

$$(3.11) \quad \begin{aligned} & r \frac{\cos(r)}{\sin(r)} - \left(\frac{1}{\sin^2 r} - r \frac{\cos(r)}{\sin^3 r} \right) (\operatorname{Re}x)^2 + \frac{r}{\sin r \cos r} (\operatorname{Im}x)^2 \\ & > r \frac{\cos(r)}{\sin(r)} - \left(\frac{1}{\sin^2 r} - r \frac{\cos(r)}{\sin^3 r} \right) \sin^2 r = \frac{2r - \tan r}{\tan r}, \end{aligned}$$

which is > 0 if $r \leq R$ where $\tan(R) = 2R$, $R \in (0, \frac{\pi}{2})$. Thus if $\operatorname{supp}(Q) \subseteq B(\mu_I, R)$, then $\frac{d^2}{ds^2} d^2(\gamma(s), m)|_{s=0} > 0$, and hence Λ is positive definite. \square

Remark 3.1. It can be shown that $R \in (\frac{\pi}{3}, \frac{2\pi}{5})$. It is approximately 0.37101π .

Remark 3.2. The nonsingularity of Σ defined in Theorem 2.2 is a mild condition which holds in particular if Q has a density (component) with respect to the volume measure on Σ_2^k .

From Proposition 2.1, Theorem 3.1 and Remark 3.2, we conclude that if $\operatorname{supp}(Q) \subseteq B(\mu_I, R)$ and if Σ is nonsingular (e.g., if Q is absolutely continuous), then the sample mean from an iid sample has an asymptotically normal distribution with nonsingular dispersion. To get the expressions for Σ and Λ , note that the coordinate ϕ in Theorem 3.1 has the form

$$\phi(m) = (\tilde{m}^1, \dots, \tilde{m}^{k-2})', \quad \tilde{m}^j = \frac{r}{\sin r} e^{i\theta} \bar{v}_j' z.$$

Hence

$$(3.12) \quad \begin{aligned} \Sigma_{(2k-4) \times (2k-4)} &= \begin{bmatrix} \Sigma_{11} & \Sigma_{12} \\ \Sigma'_{12} & \Sigma_{22} \end{bmatrix}, \\ (\Sigma_{11})_{ij} &= 4E(\operatorname{Re}(\tilde{X}_1^i) \operatorname{Re}(\tilde{X}_1^j)), \quad (\Sigma_{12})_{ij} = 4E(\operatorname{Re}(\tilde{X}_1^i) \operatorname{Im}(\tilde{X}_1^j)), \\ (\Sigma_{22})_{ij} &= 4E(\operatorname{Im}(\tilde{X}_1^i) \operatorname{Im}(\tilde{X}_1^j)), \quad 1 \leq i, j \leq k-2, \end{aligned}$$

and

$$(3.13) \quad \Lambda_{(2k-4) \times (2k-4)} = \begin{bmatrix} \Lambda_{11} & \Lambda_{12} \\ \Lambda'_{12} & \Lambda_{22} \end{bmatrix},$$

where if $\dot{\gamma}(0) = \sum_{j=1}^{k-2} x^j v_j + \sum_{j=1}^{k-2} y^j (iv_j)$, $x = [x^1 \dots x^{k-2}]'$, $y = [y^1 \dots y^{k-2}]'$, then

$$E \left(\frac{d^2}{ds^2} d_g^2(\gamma(s), X_1) \right) |_{s=0} = x' \Lambda_{11} x + y' \Lambda_{22} y + 2x' \Lambda_{12} y.$$

This gives for $1 \leq r, s \leq k-2$,

$$\begin{aligned} (\Lambda_{11})_{rs} &= 2E \left[d_1 \cot(d_1) \delta_{rs} - \frac{(1 - d_1 \cot(d_1))}{d_1^2} (\operatorname{Re} \tilde{X}_1^r) (\operatorname{Re} \tilde{X}_1^s) \right. \\ &\quad \left. + \frac{\tan(d_1)}{d_1} (\operatorname{Im} \tilde{X}_1^r) (\operatorname{Im} \tilde{X}_1^s) \right], \\ (\Lambda_{22})_{rs} &= 2E \left[d_1 \cot(d_1) \delta_{rs} - \frac{(1 - d_1 \cot(d_1))}{d_1^2} (\operatorname{Im} \tilde{X}_1^r) (\operatorname{Im} \tilde{X}_1^s) \right. \\ &\quad \left. + \frac{\tan(d_1)}{d_1} (\operatorname{Re} \tilde{X}_1^r) (\operatorname{Re} \tilde{X}_1^s) \right], \\ (\Lambda_{12})_{rs} &= -2E \left[\frac{(1 - d_1 \cot(d_1))}{d_1^2} (\operatorname{Re} \tilde{X}_1^r) (\operatorname{Im} \tilde{X}_1^s) + \frac{\tan(d_1)}{d_1} (\operatorname{Im} \tilde{X}_1^r) (\operatorname{Re} \tilde{X}_1^s) \right], \end{aligned}$$

where $d_1 = d_g(X_1, \mu_I)$.

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