

Stiffness and damping characteristics of conical multirecess hybrid journal bearing for different load arrangements

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Abstract

This paper presents the theoretical analysis on the effect of two different load arrangements for a conical multirecess hybrid journal bearing. For this analysis constant flow valve restrictors are used under micropolar lubricant for different semi cone angles. The modified Reynolds equation using micropolar fluid flow on the conical surface has been solved with the use of FEM and Galerkin's technique applying proper boundary conditions in the iterative scheme. The study suggests that the configuration LA I can be a better alternative as compared to LA II for different semi-cone angles and external load in terms of bearing characteristics.

Keywords CFV restrictor · Conical bearing · Micropolar lubrication · FEM

List of symbols

a_b	Axial bearing land width, mm
A_p	Area of pocket, mm ²
A_b	Area of bearing, mm ²
A^e	Area of each element
c	Radial clearance, mm
C_{ij}	Fluid film damping coefficients (N/mm ²), ($i, j = x, z$)
D_m	Mean journal diameter of conical shaft, mm
e	Eccentricity, mm
F	Fluid film reaction
F_x, F_z	Fluid film reaction component, N
h	Fluid film thickness, mm
l_m, \bar{l}	Bearing characteristic length
L	Bearing length, mm
M	Journal mass, kg
p_s	Supply pressure, N/mm ²
Q	Lubricant flow (mm ³ /s)
R_j	Journal radius, mm
r	Radial coordinate
S_{ij}	Fluid film stiffness coefficients (N/mm), ($i, j = x, z$)
ω_j	Journal rotation speed, rad/s

ω_{th}	Threshold speed, rad/s
W_r	Radial load, N
X_j, Z_j	Journal center coordinate
x, y, z	Cartesian coordinate system

Greek symbol

α	Circumferential coordinate
β	Axial co-ordinate
γ	Semi cone angle, degree
μ	Dynamic viscosity (Pa·s)
λ	Aspect ratio, L/D _m
θ	Inter-recess angle
Ω	Speed parameter
ε	Eccentricity ratio e/c
\emptyset	Micro polar function
φ	Attitude angle

Non-dimensional number/parameter

\bar{a}_b	a_b/L
\bar{C}_{ij}	$C_{ij} \left(c^3 / \mu r_j^4 \right)$
\bar{F}_r	(F / p_s)
\bar{F}_x, \bar{F}_z	$(F_x, F_z / p_s r_j^2)$

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\bar{F}_0	Resultant fluid film force at steady state condition
\bar{h}_{min}	Minimum fluid film thickness
\bar{P}_{max}	Maximum fluid film pressure
\bar{M}_c	$M(c/p_s R_j^3)$
N^2	Coupling number
N_i, N_j	Shape function
\bar{P}_j	(p/p_s)
\bar{Q}	$Q(\mu/c^3 p_s) r$
\bar{Q}_c	Constant flow valve restrictor design parameter
\bar{S}_{ij}	$S_{ij}(c/p_s r_j^2)$
$\bar{\omega}_{th}$	$\sqrt{\bar{M}_c/\bar{F}_0}$
\bar{X}_j, \bar{Z}_j	$(X_j \cdot Y_j)/R_j$

Matrices

$\{\bar{F}_{ij}\}$	Fluidity matrix
$\{\bar{P}\}$	Nodal pressure vector
$\{\bar{Q}\}^e$	Nodal flow vector
$\{\bar{R}_{H_i}\}^e$	Column vector due to hydrodynamic term
$\{\bar{R}_{x_{ji}}, \bar{R}_{z_{ji}}\}$	Nodal R.H.S vector due to journal center velocity

1 Introduction

Conical journal bearings offer the advantages over cylindrical bearings. They can support radial as well as thrust load thereby can replace two bearings with one and they have adjustable radial clearance during assembly. However, these advantages are at the expense of corresponding requirement of accuracy level in machining of surfaces and accurate assembly of conical bearing. Many researchers have carried out the theoretical study on the conical bearing geometry. The multirecess externally pressurised journal bearing with micropolar lubrication has

been studied by Verma et al. [1]. Rana et al. have studied theoretically the effect of micropolar fluid parameters on the steady state as well as stiffness and damping characteristics of 4-pocket constant flow valve compensated hydrostatic conical journal bearing [2, 3]. Sharma et al. [4, 5] theoretically analyzed the effect of wear on the characteristics of a 4 pocket hydrostatic/hybrid conical journal bearing system for different cone angles with different restrictors for different load arrangements. The analysis of multirecessed hybrid conical journal bearing with various restrictors have been carried out by Rajput and Sharma [6, 7]. Dhawan and Verma [8] have studied non-circular hybrid journal bearings with micropolar fluid film lubricant. Sharma et al. [9] carried out the performance characteristics of multirecess hydrostatic/hybrid flexible journal bearing compensated with membrane type variable flow restrictor. Performance analysis of a hybrid conical journal bearing with capillary restrictors has been presented by Khakse et al. [10]. Guo et al. [11] studied the performance of hydrostatic deep/shallow pocket hybrid conical bearing compensated by capillary and seen that the stability and load carrying capacity improves under small eccentric operation. Zuo et al. [12] have analyzed the comparative performance of conical hydrostatic bearings compensated by variable slot and fixed slot. They have concluded that variable slot compensated bearing exhibits a better radial load carrying capacity as well as the stiffness coefficients. The stability of the rigid rotor supported by six recessed hybrid bearings with orifice compensation and influence of restrictor design parameters and eccentricity on it are studied by Chen et al. [13]. Rana et al. [14, 15] have analyzed the performance of four and six pockets conical hybrid journal bearing.

The performances of conical hybrid journal bearing with constant flow valve for micropolar lubricant for different load arrangements have not been exploited by the researchers. Hence, the efforts have been made to

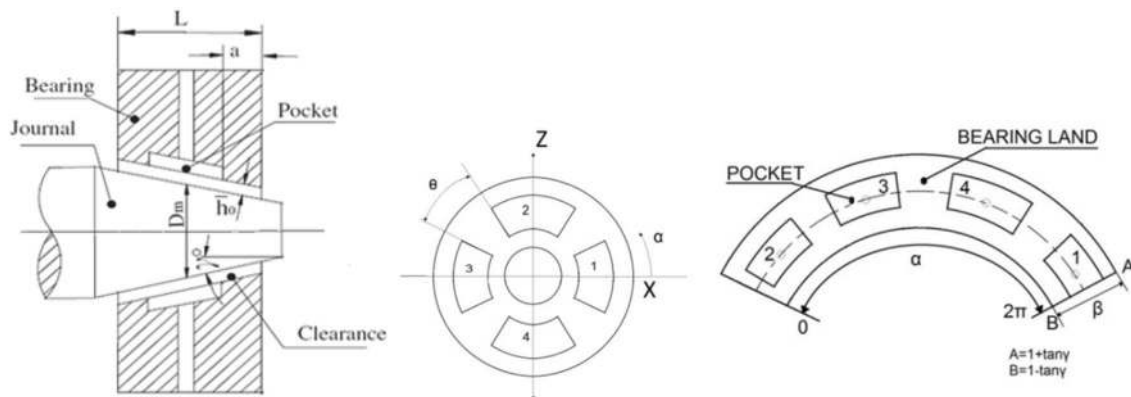


Fig. 1 Conical hybrid journal bearing [14]

highlight this issue in the present work to analyze the bearing performances. Based on the above fact, the CFV compensated conical hybrid journal bearing performance for different loading conditions LA I (i.e. load line in between the two recesses) and LA II (i.e. load line on a recess) under micropolar lubrication has been carried out. The data charts presented here for multirecess hybrid bearings may be useful for the bearing designers as a guide for bearing performance.

1.1 Governing equation

The 4-pocket hybrid conical journal bearing configurations are shown in Figs. 1 and 2 with journal assuming equilibrium position under rotation with uniform angular velocity.

For conical surface using micropolar fluid the Reynolds equation in non-dimensional form is as follows [5–7]:

$$\frac{1}{\beta^2} \frac{\partial}{\partial \alpha} \left[\frac{\bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \alpha} \right] + \frac{\sin^2 \gamma}{\beta} \frac{\partial}{\partial \beta} \left[\frac{\beta \bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \beta} \right] = \frac{\Omega}{2} \frac{\partial \bar{h}}{\partial \alpha} + \frac{\partial \bar{h}}{\partial \bar{t}} \tag{1}$$

where

$$\beta = \frac{r \sin \gamma}{R_j}, \quad \bar{p} = \frac{p}{p_s}, \quad \bar{h} = \frac{h}{C}, \quad \bar{t} = t \left(\frac{C^2 p_s}{R_j^2 \mu} \right), \quad \bar{l} = \frac{c}{l_m}$$

$$\bar{\theta}(N, \bar{l}, \bar{h}) = 1 + \frac{12}{\bar{h}^2 \bar{l}^2} - \frac{6N}{\bar{h} \bar{l}} \coth \left(\frac{N \bar{h} \bar{l}}{2} \right)$$

1.2 Finite element formulation

The governing element equation is obtained with the help of Galerkin’s technique along with the orthogonality

conditions as well as approximate value of non-dimensional pressure. Therefore, Eq. (1) can be expressed as follows:

$$\iint_{A^e} \left\{ \frac{\sin^2 \gamma}{\beta} \frac{\partial}{\partial \beta} \left[\frac{\beta \bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \beta} \right] + \frac{1}{\beta^2} \frac{\partial}{\partial \alpha} \left[\frac{\bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \alpha} \right] - \frac{\Omega}{2} \frac{\partial \bar{h}}{\partial \alpha} - \frac{\partial \bar{h}}{\partial \bar{t}} \right\} N_j d\alpha d\beta = 0$$

$$\iint_{A^e} \left\{ \frac{\sin^2 \gamma}{\beta} \frac{\partial}{\partial \beta} \left[\frac{\beta \bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \beta} \right] + \frac{1}{\beta^2} \frac{\partial}{\partial \alpha} \left[\frac{\bar{\theta}(N, \bar{l}, \bar{h})}{12} \frac{\partial \bar{p}}{\partial \alpha} \right] - \frac{\Omega}{2} \frac{\partial \bar{h}}{\partial \alpha} + \bar{X}_j \cos \alpha \cos \gamma + \bar{Z}_j \sin \alpha \cos \gamma \right\} N_j d\alpha d\beta = 0 \tag{2}$$

This Eq. (2) is intergrated part wise and combined by multiplying with the term $\frac{\beta^2}{\sin \gamma}$, which transforms as,

$$\begin{aligned} & \iint_{A^e} \frac{\bar{\theta}(N, \bar{l}, \bar{h})}{12} \left[\beta^2 \sin \gamma \frac{\partial N_i}{\partial \beta} \frac{\partial N_j}{\partial \beta} + \frac{1}{\sin \gamma} \frac{\partial N_i}{\partial \alpha} \frac{\partial N_j}{\partial \alpha} \right] \bar{p}_j d\alpha d\beta \\ & + \int_{\Gamma^e} \frac{\bar{\theta}(N, \bar{l}, \bar{h})}{12} \left[\frac{1}{\sin \gamma} \frac{\partial \bar{p}}{\partial \alpha} l_1 + \beta^2 \sin \gamma \frac{\partial \bar{p}}{\partial \beta} l_2 \right] N_i d\Gamma^e \\ & - \frac{\Omega}{2} \int_{\Gamma^e} N_i \frac{\beta^2}{\sin \gamma} \bar{h} l_1 d\Gamma^e + \frac{\Omega}{2} \iint_{A^e} \bar{h} \frac{\beta^2}{\sin \gamma} \frac{\partial N_i}{\partial \alpha} d\alpha d\beta \\ & + \iint_{A^e} \bar{X}_j \frac{\beta^2}{\sin \gamma} \cos \alpha \cos \gamma N_i d\alpha d\beta \\ & + \iint_{A^e} \bar{Z}_j \frac{\beta^2}{\sin \gamma} \sin \alpha \cos \gamma N_i d\alpha d\beta = 0 \end{aligned} \tag{3}$$

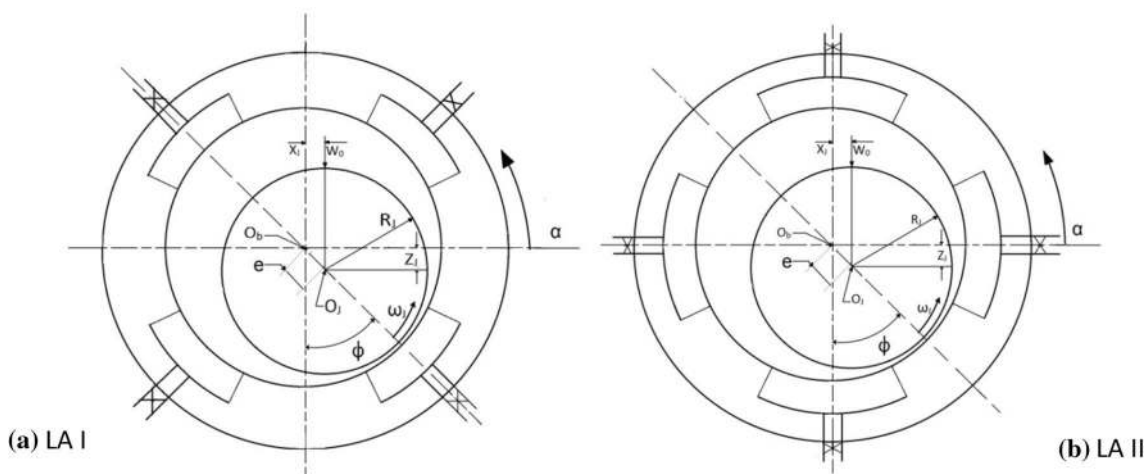


Fig. 2 LA I and LA II configurations of conical hybrid journal bearing

Now the above equation is represented in matrix form as:

$$[\bar{F}_{ij}]\{\bar{p}\} = \{\bar{Q}\}^e + \Omega\{\bar{R}_{H_i}\}^e + \bar{X}_j\{\bar{R}_{X_{ji}}\}^e + \bar{Z}_j\{\bar{R}_{Z_{ji}}\}^e \quad (4)$$

where

$$[\bar{F}_{ij}] = \iint_{A^e} \frac{\bar{\theta}(N_i, \bar{l}, \bar{h})}{12} \left[\frac{1}{\sin\gamma} \frac{\partial N_i}{\partial \alpha} \frac{\partial N_j}{\partial \alpha} + \beta^2 \sin\gamma \frac{\partial N_i}{\partial \beta} \frac{\partial N_j}{\partial \beta} \right] d\alpha d\beta$$

$$\{\bar{Q}_i\}^e = \int_{\Gamma^e} \frac{\bar{\theta}(N_i, \bar{l}, \bar{h})}{12} \left[\frac{1}{\sin\gamma} \frac{\partial \bar{p}}{\partial \alpha} l_1 + \beta^2 \sin\gamma \frac{\partial \bar{p}}{\partial \beta} l_2 \right] N_i d\Gamma^e - \frac{\Omega}{2} \int_{\Gamma^e} \frac{\beta^2}{\sin\gamma} \bar{h} l_1 N_i d\Gamma^e$$

$$\bar{R}_{H_i}^e = \iint_{A^e} \frac{\bar{h}}{2} \frac{\partial N_i}{\partial \alpha} \frac{\beta^2}{\sin\gamma} d\alpha d\beta$$

$$\bar{R}_{X_{ji}}^e = \iint_{A^e} N_i \cos\alpha \cos\gamma \frac{\beta^2}{\sin\gamma} d\alpha d\beta$$

$$\bar{R}_{Z_{ji}}^e = \iint_{A^e} N_i \sin\alpha \cos\gamma \frac{\beta^2}{\sin\gamma} d\alpha d\beta$$

1.3 Stiffness and damping coefficients

The stiffness coefficients for fluid film in bearing are expressed as

$$\begin{bmatrix} \bar{S}_{xx} & \bar{S}_{xz} \\ \bar{S}_{zx} & \bar{S}_{zz} \end{bmatrix} = - \begin{bmatrix} \partial \bar{F}_x / \partial \bar{X}_j & \partial \bar{F}_x / \partial \bar{Z}_j \\ \partial \bar{F}_z / \partial \bar{X}_j & \partial \bar{F}_z / \partial \bar{Z}_j \end{bmatrix}$$

and the damping coefficients are expressed as

$$\begin{bmatrix} \bar{C}_{xx} & \bar{C}_{xz} \\ \bar{C}_{zx} & \bar{C}_{zz} \end{bmatrix} = - \begin{bmatrix} \partial \bar{F}_x / \partial \bar{X}_j & \partial \bar{F}_x / \partial \bar{Z}_j \\ \partial \bar{F}_z / \partial \bar{X}_j & \partial \bar{F}_z / \partial \bar{Z}_j \end{bmatrix}$$

1.4 Stability parameters

Equation for non-dimensional critical mass, \bar{M}_c

$$\bar{M}_c = \frac{\bar{G}_1}{\bar{G}_2 - \bar{G}_3}$$

$$\bar{G}_1 = [\bar{C}_{xx} \bar{C}_{zz} - \bar{C}_{zx} \bar{C}_{xz}]$$

$$\bar{G}_2 = \frac{[\bar{S}_{xx} \bar{S}_{zz} - \bar{S}_{zx} \bar{S}_{xz}][\bar{C}_{xx} + \bar{C}_{zz}]}{[\bar{S}_{xx} \bar{C}_{zz} + \bar{S}_{zz} \bar{C}_{xx} - \bar{S}_{xz} \bar{C}_{zx} - \bar{S}_{zx} \bar{C}_{xz}]}$$

$$\bar{G}_3 = \frac{[\bar{S}_{xx} \bar{C}_{xx} + \bar{S}_{xz} \bar{C}_{xz} + \bar{S}_{zx} \bar{C}_{zx} + \bar{S}_{zz} \bar{C}_{zz}]}{\bar{C}_{xx} + \bar{C}_{zz}}$$

Equation for threshold speed, $\bar{\omega}_{th}$

$$\bar{\omega}_{th} = \sqrt{\frac{\bar{M}_c}{\bar{F}_0}} \quad (5)$$

1.5 Fluid film thickness

For conical journal bearing, the fluid film thickness can be expressed in non-dimensional form as,

$$\bar{h} = (1 - \bar{X}_j \cos\alpha - \bar{Z}_j \sin\alpha) \cos\gamma \quad (6)$$

1.6 Restrictor flow equation

The flow control device used here is CFV restrictor. The lubricant supply passing through it must be a fixed quantity. The lubricant flow \bar{Q}_R can be written as

$$\bar{Q}_R = constant = \bar{Q}_c \quad (7)$$

here \bar{Q}_R and \bar{Q}_c are flow through restrictor and pocket, respectively.

1.7 Boundary conditions

1. Lubricant flow through restrictor equals input flow of the bearing ($\bar{Q}_R = \bar{Q}_c$).
2. All the nodes situated on a pocket/recess have the same pressure.
3. The nodes situated at the edges of bearing have atmospheric pressure i.e. zero gauge pressure ($\bar{p}|_{\beta=\mp 1.0} = 0.0$).
4. It is assumed that Reynolds boundary condition is applicable at the trailing portion of positive pressure region i.e.

$$\bar{p} = \frac{\partial \bar{p}}{\partial \alpha} = 0.0$$

In the present study the Reynolds boundary condition is used due to absence of rupture and reformation of the lubricant film.

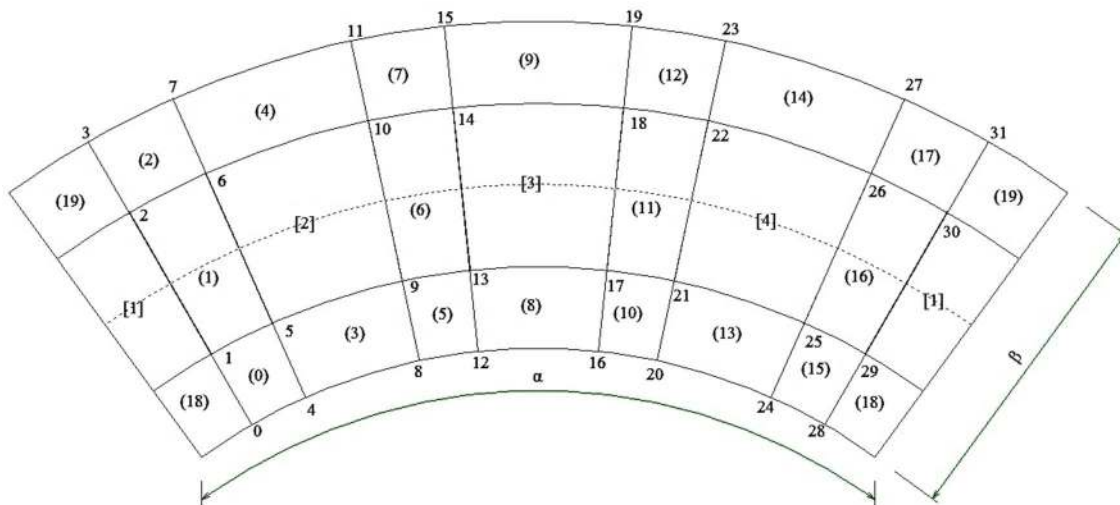


Fig. 3 Discretized 4-pocket conical journal bearing [14]

1.8 Solution procedure

The discretization of flow field of micropolar fluid in the lubrication zone of the conical journal bearing system using 4 noded isoparametric elements and Langragian interpolation function is done. The pressure is distributed in the element of discretized zone and can be expressed for 4 pockets by:

$$\bar{p} = \sum_{j=1}^4 N_j \bar{p}_j$$

The 2-dimensional mesh of 4 pockets conical journal bearing with nodes as well as elements are shown in Fig. 3.

The governing Reynolds equation, which governs the micropolar fluid, is solved taking into consideration of restrictor flow and suitable boundary conditions by finite element technique to obtain the fluid film pressures. The iterations are repeated till solutions are converged for calculation of fluid film pressure.

The computations for the steady state characteristics of hybrid journal bearing system requires input data file. The INPDAT module reads input data for 2-D meshing. The fluid film thickness for different nodal points and tentative journal center positions is calculated by the subroutine FLDTHK. The fluidity matrix is created and generalized according to connectivity of various elements through subroutine FEQN. The boundary condition of the system is specified by the CALL_BOUNDARY subroutine. The governing Reynolds equation for conical bearing is solved using Gaussian elimination procedure for the calculation of fluid film pressures at different nodes through subroutine CALL_SOLVER. The load carrying limits are

calculated as well. The journal center and its equilibrium position are ascertained following iterations shown in solution scheme by JCE with the use of \bar{F}_x and \bar{F}_z components. Finally, the SPC subroutine calculates the steady state parameters of the conical bearing that can be seen in Fig. 12 ("Appendix").

1.9 Discussion of results

The mathematical model and the solution scheme presented in the previous section are used to calculate the bearing performances and stability parameters of CFV compensated 4-pocket conical hybrid journal bearing under micropolar lubrication. The hybrid journal bearing must sustain the operating load at maximum speed and the partial load at zero speed respectively. The maximum speed considered in this analysis is $\Omega = 1.0$ for all data charts. The values are computed for the following non-dimensional data of the bearing and geometric parameters:

$$\Omega = 1.0, \lambda = 1.0, \bar{a}_b = 0.25, \bar{W}_r = 0.5, \bar{Q}_c = 0.1 - 0.3, \theta = 18^\circ,$$

$$A_p/A_b = 0.333, N^2 = 0.3 - 0.9, \bar{l} = 10 - 30, \gamma = 10^\circ - 40^\circ$$

To build confidence in the results, the validity of the formulation of Reynolds equation and its solution scheme developed is established by employing the results available in literature for different cases. The results obtained are compared with the data of Sharma and Rajput [6] and are shown in Table 1. The compared results are seen to be in close agreement with the published data. The various performance characteristics for Newtonian and micropolar fluid parameters for load arrangements I & II are presented in Table 2.

Table 1 Validation for conical bearing with capillary: $\lambda = 1.0$, $\Omega = 1.0, \bar{W}_r = 0.5, C_{s2} = 0.5, A_p/A_b = 0.333, \gamma = 10^0$

L_M	N^2	Comparison	\bar{P}_{max}	\bar{h}_{min}	\bar{Q}	φ
10	0.4	Present result	0.6747	0.8953	0.8464	61.77
	0.4	Sharma and Rajput [6]	0.6796	0.8884	0.8522	61.00
	0.5	Present result	0.6898	0.8981	0.8161	62.10
	0.5	Sharma and Rajput [6]	0.6947	0.8914	0.8218	61.00
20	0.4	Present result	0.6436	0.8898	0.9096	60.40
	0.4	Sharma and Rajput [6]	0.6485	0.8825	0.9156	59.00
	0.5	Present result	0.6503	0.8910	0.8960	60.52
	0.5	Sharma and Rajput [6]	0.6553	0.8840	0.9020	59.00
30	0.4	Present result	0.6302	0.8870	0.9366	60.02
	0.4	Sharma and Rajput [6]	0.6352	0.8795	0.9426	59.00
	0.5	Present result	0.6344	0.8879	0.9281	60.07
	0.5	Sharma and Rajput [6]	0.6394	0.8805	0.9341	59.00

Table 2 Comparisons of Newtonian and micropolar fluid in terms of different bearing characteristics at $\Omega = 1, \bar{Q}_c = 0.2, \bar{W}_r = 0.5, \bar{a} = 0.25, \lambda = 1, \gamma = 20^\circ$

	Newtonian		$N^2 = 0.3, \bar{J} = 10$		$N^2 = 0.9, \bar{J} = 10$		$N^2 = 0.3, \bar{J} = 30$		$N^2 = 0.9, \bar{J} = 30$	
	LA I	LA II	LA I	LA II	LA I	LA II	LA I	LA II	LA I	LA II
\bar{h}_{min}	0.745	0.709	0.778	0.733	0.803	0.766	0.764	0.721	0.774	0.731
\bar{p}_{max}	0.493	0.560	0.566	0.632	0.693	0.758	0.524	0.591	0.555	0.621
\bar{S}_{11}	7.004	4.274	8.297	5.208	11.102	7.114	7.609	4.682	8.241	5.099
\bar{S}_{22}	3.280	7.823	4.090	9.287	5.833	12.42	3.64	8.507	4.013	9.213
\bar{C}_{11}	4.944	4.067	6.118	4.913	8.329	6.548	5.428	4.415	5.909	4.765
\bar{C}_{22}	3.342	2.394	4.078	3.092	5.517	4.398	3.645	2.684	3.951	2.971

1.9.1 Influence on minimum fluid film thickness (\bar{h}_{min}) and maximum fluid film pressure (\bar{p}_{max})

Here the numerical analysis has been performed for the recessed hybrid conical bearing, having symmetrically placed 4-pockets, for two possible loading arrangements namely LA I and LA II. In LA I, the vertical downward load line passes between the two recesses i.e. through land-width of bearing while in LA II arrangement, the load line passes through the recess in the bearing and is the conventional arrangement used in the bearing design.

The plot of minimum fluid film thickness \bar{h}_{min} with \bar{W}_r is shown in Fig. 4. The value of \bar{h}_{min} decreases with an increase in \bar{W}_r for different angles γ and for different loading conditions LA I and LA II. It is further noticed that \bar{h}_{min} is minimum for LA II than LA I for $\gamma = 20^\circ$ and $\gamma = 40^\circ$ for the same operating and geometric conditions considered. As the \bar{h}_{min} for LA I is more as compared to LA II, so it can accommodate a significant amount of surface worn and surface roughness.

The plot of \bar{p}_{max} versus \bar{W}_r at $\Omega = 1$ for LA I & LA II are shown in Fig. 5. The value of \bar{p}_{max} increases with an increase in \bar{W}_r and the trend is same for both the loading conditions. Also, as the maximum fluid film pressure is reciprocal to minimum fluid film thickness, the \bar{p}_{max} is seen to be higher for LA II when compared to LA I for the same value of \bar{W}_r considered in the range of 0.1 to 0.5. Whereas, LA I could have been a better alternative to reduce the \bar{p}_{max} for the same \bar{W}_r .

1.9.2 Influence on stiffness coefficients (S_{ij}) and damping coefficients (C_{ij})

The plots of direct stiffnesses coefficient versus \bar{W}_r are given in Figs. 6 and 7. The stiffness coefficients (\bar{S}_{11}) and (\bar{S}_{22}) remains almost same with an increase in \bar{W}_r for LA I and LA II as well as for different semi cone angles. The value of \bar{S}_{11} is seen to be higher for LA I position as shown in Fig. 6. Whereas, the value of stiffness coefficient (\bar{S}_{22}) is higher for LA II position in comparison to LA I. The

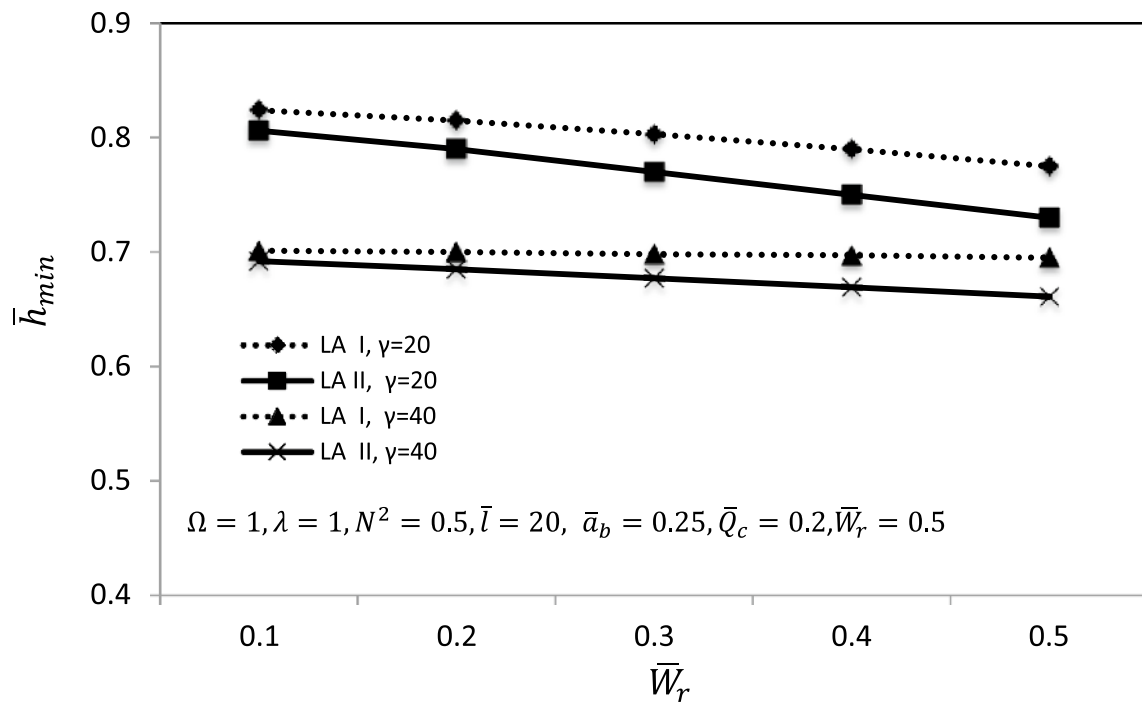


Fig. 4 Minimum fluid film thickness (\bar{h}_{min}) versus radial load (\bar{W}_r)

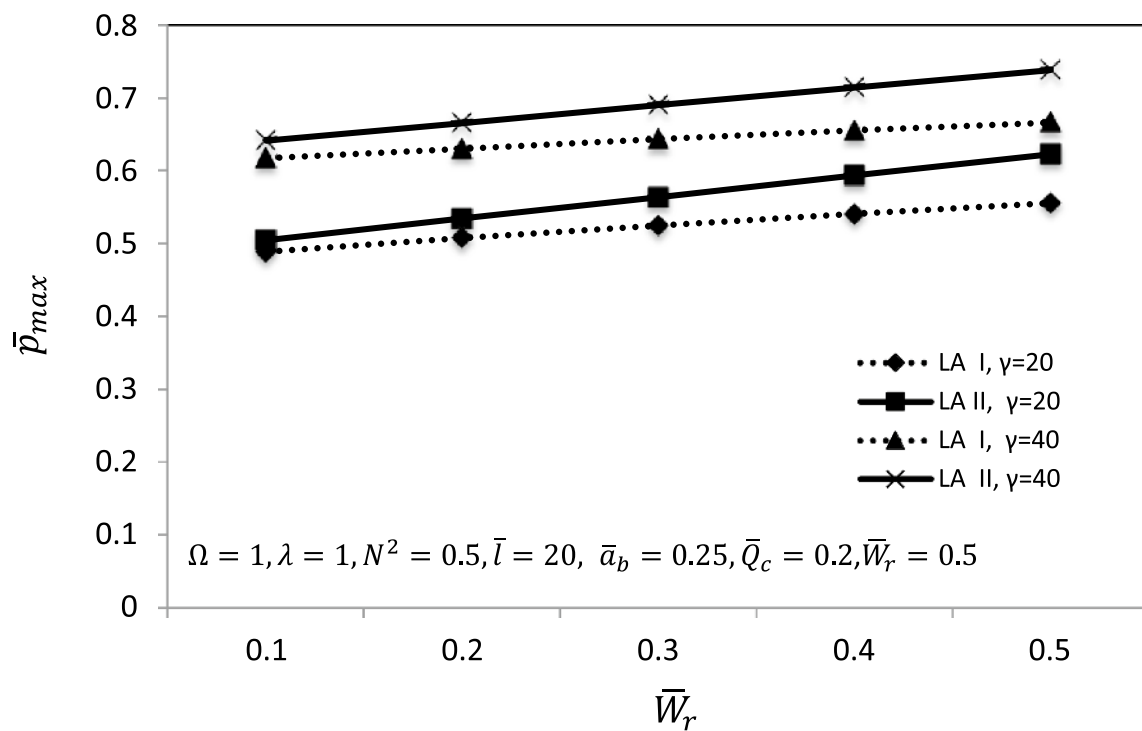


Fig. 5 Maximum fluid pressure (\bar{p}_{max}) versus radial load (\bar{W}_r)

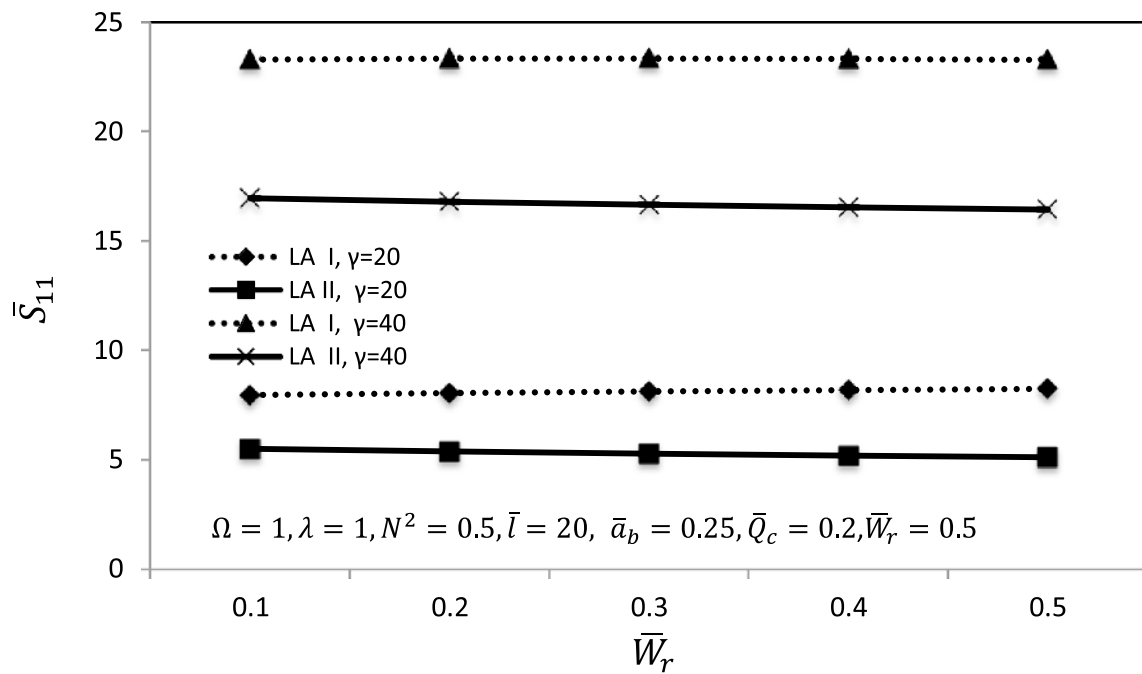


Fig. 6 Direct stiffness (\bar{S}_{11}) versus radial load (\bar{W}_r)

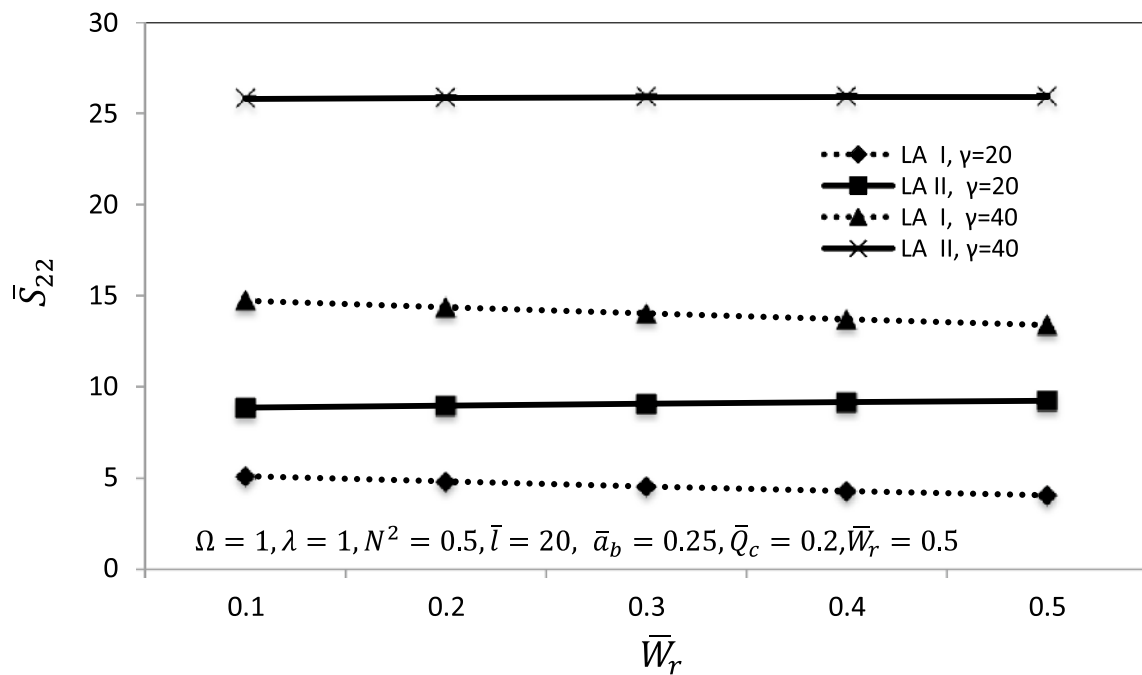


Fig. 7 Direct stiffness (\bar{S}_{22}) versus radial load (\bar{W}_r)

variation of stiffness coefficients \bar{S}_{11} and \bar{S}_{22} compliment each other for both the loading arrangements LA I and LA II respectively.

The direct damping coefficients \bar{C}_{11} and \bar{C}_{22} versus \bar{W}_r are plotted in Figs. 8 and 9. Both the damping coefficients \bar{C}_{11} and \bar{C}_{22} vary marginally with an increase in \bar{W}_r for LA I and LA II for different γ . However, \bar{C}_{11} and \bar{C}_{22} are found to

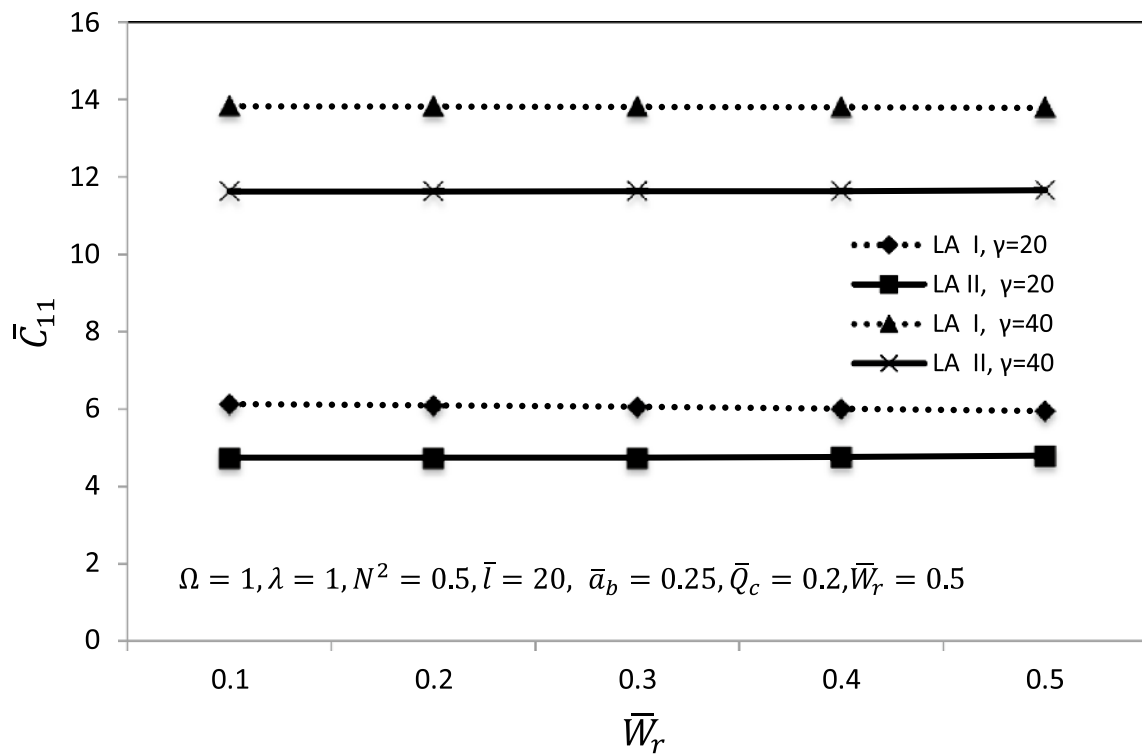


Fig. 8 Direct damping (\bar{C}_{11}) versus radial load (\bar{W}_r)

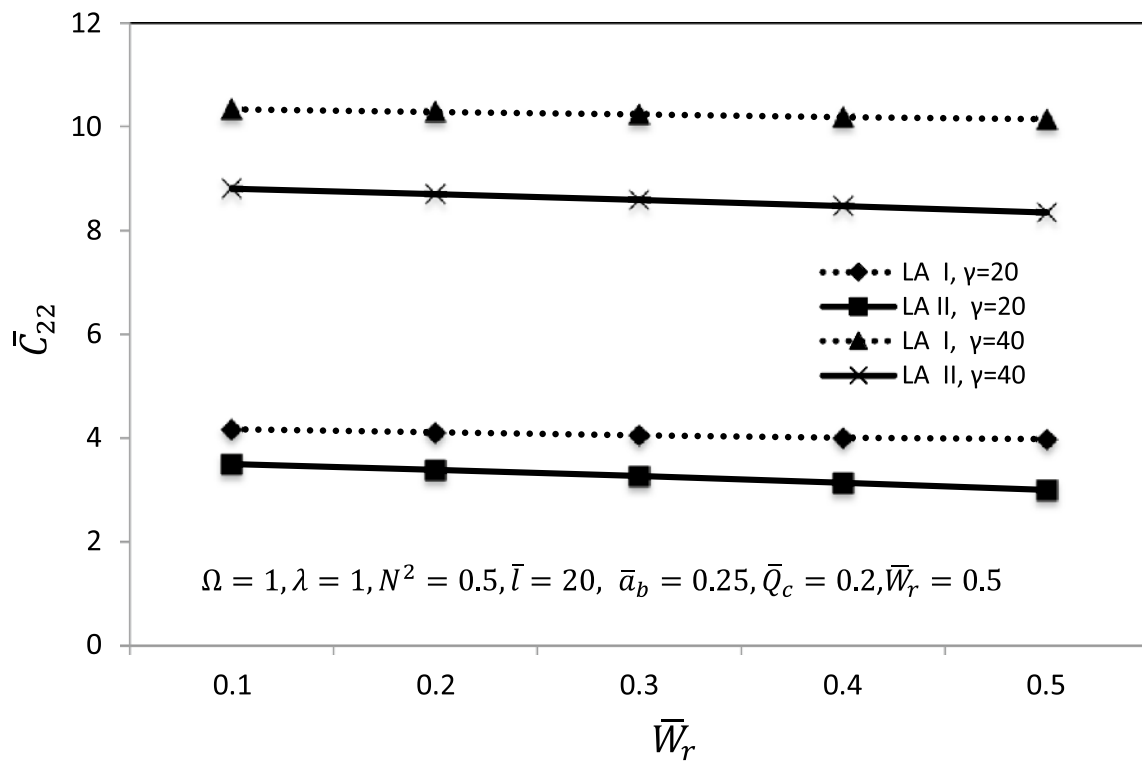


Fig. 9 Direct damping (\bar{C}_{22}) versus radial load (\bar{W}_r)

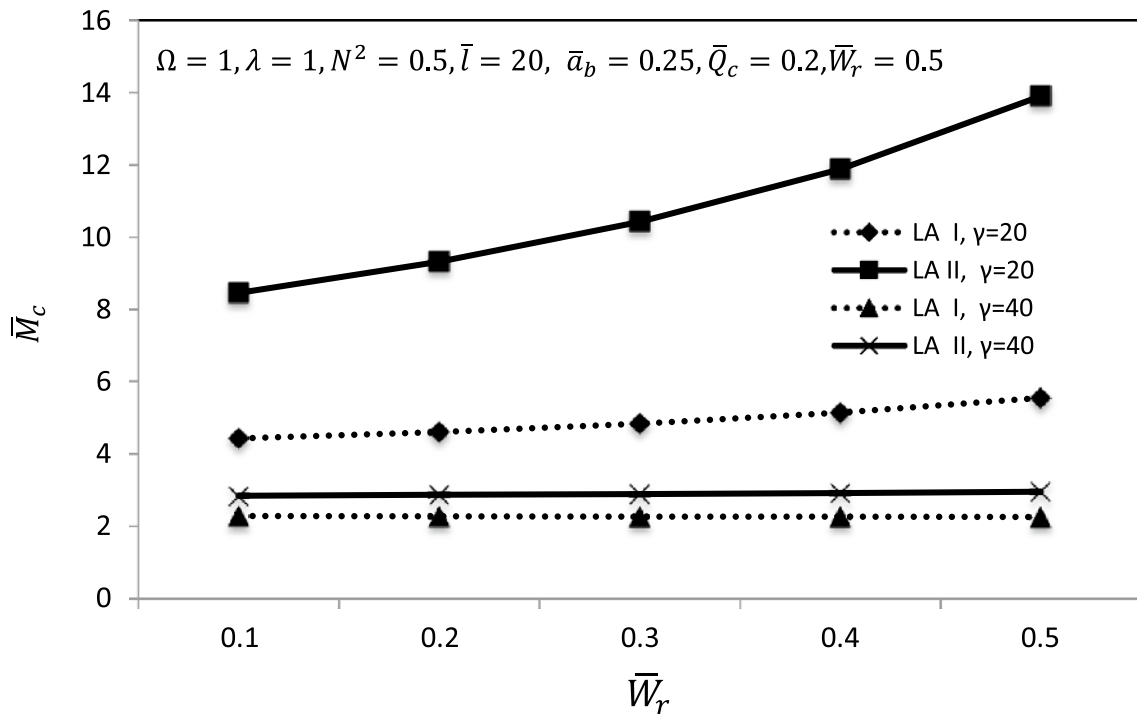


Fig. 10 Critical mass parameter (\bar{M}_c) versus radial load (\bar{W}_r)

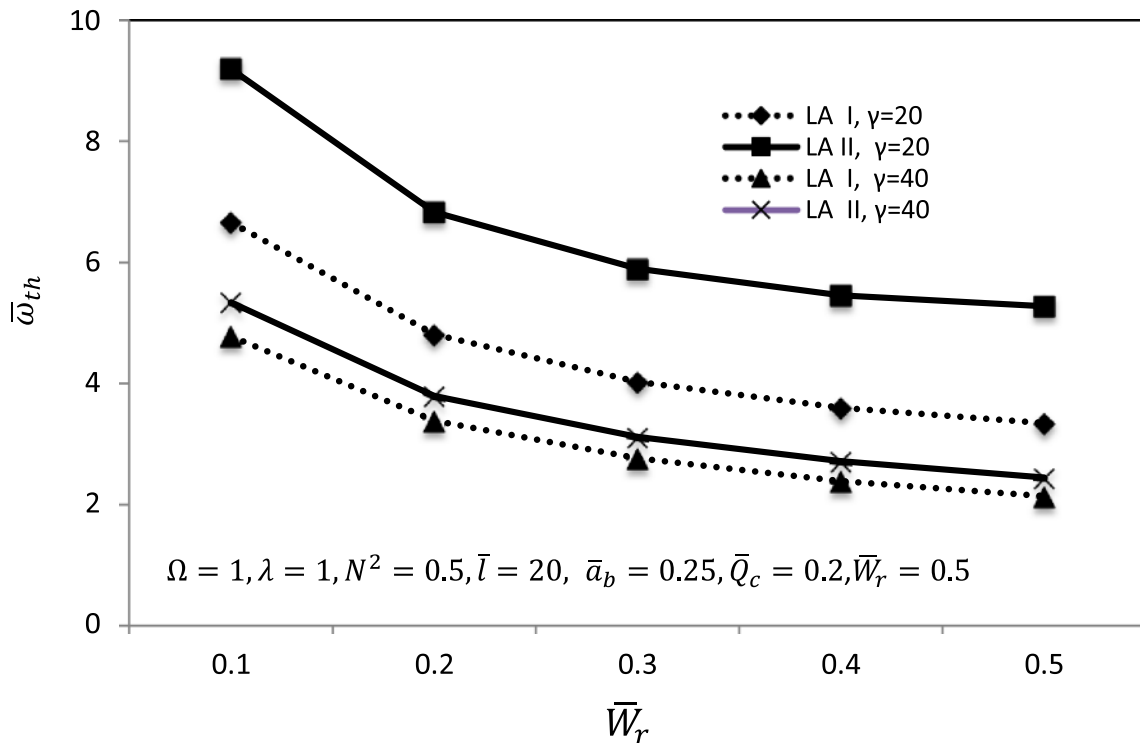


Fig. 11 Threshold frequency ($\bar{\omega}_{th}$) versus radial load (\bar{W}_r)

be increased for LA I than LA II position against \bar{W}_r . It can further be seen that for same amount of \bar{Q}_c and \bar{W}_r , the value of \bar{h}_{min} is quite high for LA I than LA II and can play a significant role in the calculation of stiffness and damping coefficients.

1.9.3 Influence on stability parameters

The plots of critical mass parameter and threshold speed at $\gamma=20^\circ$ and $\gamma=40^\circ$ for both the loading conditions are presented in Figs. 10 and 11. It is found that at $\gamma=40^\circ$ both the stability parameters show the similar trend and are seen to be very close for LA I & LA II. Whereas, at $\gamma=20^\circ$ the stability parameters are found to be higher when compared to $\gamma=40^\circ$. Hence, a combination of $\gamma=40^\circ$ with LA I can be considered from stiffness and damping coefficients point of view and $\gamma=20^\circ$ with LA II may be chosen from stability point of view.

2 Conclusions

The following conclusions can be drawn from the performance data obtained through present work that may be useful for the design of CFV compensated 4-pocket conical hybrid journal bearings:

1. The \bar{h}_{min} is found to be higher for LA I for a given range of the semi cone angle γ . Therefore, LA I configuration can be more useful than conventional configuration LA II from fluid film thickness point of view.
2. The stiffness coefficients (\bar{S}_{11}) and (\bar{S}_{22}) compliment each other for both the loading arrangements.
3. The damping coefficients (\bar{C}_{11} , \bar{C}_{22}) are found to be higher for LA I load arrangement as compared to LA II.
4. The stability parameters for $\gamma=40^\circ$ with LA I arrangement may be useful from stiffness and damping coefficients point of view and $\gamma=20^\circ$ with LA II may be chosen from stability point of view.

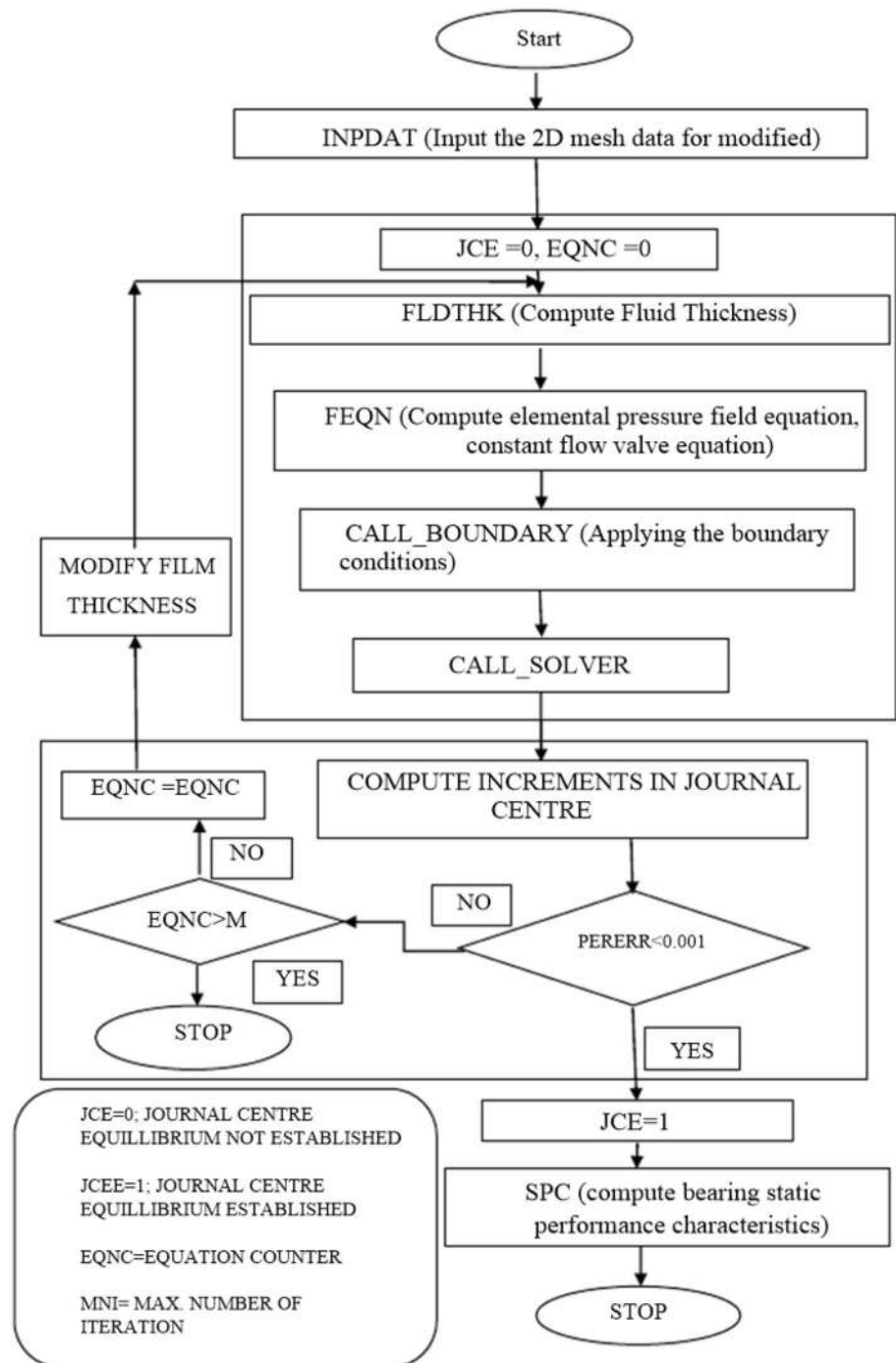
Compliance with ethical standards

Conflict of interest The authors declare that they have no conflict of interest.

Appendix

See Fig. 12.

Fig. 12 Solution scheme



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