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Still Missing Dark Matter: KCWI High-resolution Stellar Kinematics of NGC1052-DF2

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Abstract

The velocity dispersion of the ultra diffuse galaxy NGC1052-DF2 was found to be $\sigma_{\text{gc}} = 7.8^{+5.2}_{-2.2} \text{ km s}^{-1}$, much lower than expected from the stellar mass–halo mass relation and nearly identical to the expected value from the stellar mass alone. This result was based on the radial velocities of 10 luminous globular clusters that were assumed to be associated with the galaxy. A more precise measurement is possible from high-resolution spectroscopy of the diffuse stellar light. Here we present an integrated spectrum of the diffuse light of NGC1052-DF2 obtained with the Keck Cosmic Web Imager (KCWI), with an instrumental resolution of $\sigma_{\text{instr}} \approx 12 \text{ km s}^{-1}$. The systemic velocity of the galaxy is $v_{\text{sys}} = 1805 \pm 1.1 \text{ km s}^{-1}$, in very good agreement with the average velocity of the globular clusters ($\langle v_{\text{gc}} \rangle = 1803 \pm 2 \text{ km s}^{-1}$). There is no evidence for rotation within the KCWI field of view. We find a stellar velocity dispersion of $\sigma_{\text{stars}} = 8.5^{+2.3}_{-3.1} \text{ km s}^{-1}$, consistent with the dispersion that was derived from the globular clusters. The implied dynamical mass within the half-light radius $r_{1/2} = 2.7 \text{ kpc}$ is $M_{\text{dyn}} = (1.3 \pm 0.8) \times 10^8 M_{\odot}$, similar to the stellar mass within that radius ($M_{\text{stars}} = (1.0 \pm 0.2) \times 10^8 M_{\odot}$). With this confirmation of the low velocity dispersion of NGC1052-DF2, the most urgent question is whether this “missing dark matter problem” is unique to this galaxy or applies more widely.

Key words: galaxies: individual (NGC1052-DF2) – galaxies: kinematics and dynamics

1. Introduction

A remarkable result of the past 20 years is the apparent regularity of galaxy formation as reflected in the existence of a well-defined relation between galaxy and halo mass with small scatter, the stellar mass–halo mass relation (Mandelbaum et al. 2006; Behroozi et al. 2010; Moster et al. 2010; Wechsler & Tinker 2018). The scatter in this relation constrains the possible evolutionary histories of galaxies and has been measured fairly well at high masses (Vale & Ostriker 2004; Gu et al. 2016). However, this scatter is relatively unconstrained at low masses, in the regime where galaxy formation is thought to be less efficient. Measuring or constraining the halo masses of low-mass galaxies therefore provides important information on the scatter in the stellar mass–halo mass relation and on the question of whether galaxy formation is less regulated, or even stochastic, at low masses.

There is rich literature on halo mass measurements of low-mass galaxies ($M_{*} \sim 10^8 M_{\odot}$ or lower) in the Local Group (e.g., Aaronson 1983; Geha et al. 2009; Martin et al. 2016; Torrealba et al. 2018). These studies use the velocities of individual stars to infer the kinematics, the mass density profile, and the halo mass. Other studies focus on gas-rich dwarf galaxies within a few megaparsecs, inferring the halo mass from H α and/or HI rotation curves (de Blok et al. 2001; van den Bosch & Swaters 2001; Oman et al. 2016). There seems to be considerable scatter in the halo mass in this stellar mass regime, although this may partly reflect uncertainties in inclination (see Oman et al. 2016) and variation in the inner density profiles of halos rather than in total halo masses.

Outside the Local Group, much less is known about the variation in the stellar mass–halo mass relation. The recently identified population of ultra diffuse galaxies (UDGs; van Dokkum et al. 2015) holds the promise of new constraints, as their large spatial extent and often significant globular cluster populations provide probes on spatial scales where dark matter should dominate the kinematics. Using the velocities of globular clusters (Beasley et al. 2016; Toloba et al. 2018) and stellar velocity dispersions (van Dokkum et al. 2016), UDGs are gradually adding to the sample of low-mass galaxies with constraints on their dark matter content beyond the Local Group.

Recently, a relatively nearby UDG at 20 Mpc, NGC1052-DF2, was inferred to have little or no dark matter, deviating by an unprecedented amount from the expected $M_{\text{halo}}/M_{\text{stars}}$ ratio. The constraints on the NGC1052-DF2 halo mass were derived by measuring the velocities of 10 globular clusters that were assumed to be associated with the galaxy itself. The velocity dispersion of the 10 clusters is $\sigma_{\text{gc}} = 7.8^{+5.2}_{-2.2} \text{ km s}^{-1}$ (van Dokkum et al. 2018a). Due to the small number of tracers, the results have a large random uncertainty (see also Martin et al. 2018), may suffer from small sample bias in the likelihood estimator (see Laporte et al. 2018), and are sensitive to systematic errors in individual measurements (as demonstrated by the cluster GC-98; see van Dokkum et al. 2018a).

A more precise way of constraining the kinematics is by measuring the stellar velocity dispersion of the galaxy. This is challenging because of the low surface brightness of NGC1052-DF2 and because a relatively high spectral resolution is required. The observed broadening of spectral features is

$\sigma_{\text{obs}}^2 = \sigma_{\text{instr}}^2 + \sigma_{\text{stars}}^2$; because of this quadratic behavior a spectral resolution $\sigma_{\text{instr}} \sim \sigma_{\text{stars}} \sim 10 \text{ km s}^{-1}$ is required. This is now possible with the Keck Cosmic Web Imager (KCWI), a new instrument on the Keck II telescope that is optimized for precision sky limited spectroscopy of low surface brightness phenomena at relatively high spectral resolution.

2. Observations and Data Reduction

2.1. KCWI Spectroscopy

We obtained IFU spectroscopy of NGC1052-DF2 on 2018 October 11 with the KCWI (Morrissey et al. 2012, 2018) on Keck II. The highest resolution KCWI configuration was chosen where the spectra are still (nearly) sky limited. The medium slicer was used with the BH3 grating, for a field of view (FOV) of $16''.5 \times 20''.4$. The central wavelength was set to $\lambda_{\text{cen}} = 5080 \text{ \AA}$. The spectral resolution, as measured from arc lamps, ranges from 14 km s^{-1} at $\lambda = 4800 \text{ \AA}$ ($R \approx 9100$) to 11 km s^{-1} at $\lambda = 5300 \text{ \AA}$ ($R \approx 11,600$).

NGC1052-DF2 is larger than the KCWI FOV, which means that offset exposures have to be used to characterize the sky emission. In practice we alternated between two positions. In the first, “science” exposures were taken with the KCWI FOV placed just southwest of the center of NGC1052-DF2, covering the stellar component of the galaxy out to $0.7R_{\text{eff}}$ as well as GC-77, the second-brightest globular cluster associated with the galaxy. In the second pointing, “sky” exposures were taken with the FOV placed on a field $1'.3$ away, centered on the globular cluster GC-39. The globular cluster takes up only a small fraction of the KCWI area, and is masked in the sky analysis. The two pointings, along with stacked collapsed images of the science and sky exposures, are shown in Figure 1.

We obtained exposures of 1200 s at each position, for a total of 9600 s on the galaxy and 10,800 s on the offset field. The total science + sky time that is used in the analysis is 5.6 hr. Conditions were somewhat variable, with thin cirrus present during most of the observations.

2.2. Reduction

The KCWI Data Extraction and Reduction Pipeline (KDERP), with default settings, is used to perform basic reduction and calibration of the data (Morrissey et al. 2018). Each of the 17 science and sky frames is treated independently. A combination of “bars” exposures, arc lamps, and the science data is used to derive distortion corrections and wavelength calibration solutions. The transformations are used to convert the 2D image into a 3D data cube, consisting of the slice number, the position along the slice, and wavelength. These data cubes, dubbed “ocubed” files by the KDERP, are used in the subsequent steps.

The sky background in the science exposures is determined from the offset sky exposures. The sky frames cannot be used directly as the sky spectrum changes significantly over the 20 minute interval between successive exposures. Instead, we model the variation in the sky spectrum with a principal component analysis. The method is introduced and explained in detail in P. van Dokkum et al. (2019, in preparation). Briefly, 1D sky spectra are extracted from the nine offset exposures by averaging over the two spatial dimensions after masking GC-39 and serendipitous objects in the field. These nine spectra are analyzed with singular value decomposition using the

scikit-learn package, with six components. These eigenspectra, along with the average of the nine spectra and an approximate model for the galaxy spectrum, are then fitted to 1D extractions of each of the eight science exposures. The 1D sky model for each science exposure is subtracted from each spatial pixel in the science data cube. We note that the model does not take possible spatial variation in the sky into account and is insensitive to offsets that are not correlated with specific sky emission or absorption features (see Section 3).

Finally, 1D combined spectra for different spatial regions are created by extracting them from the individual science cubes and averaging them with optimal weighting. In the combination, step pixels that deviate $>3\sigma$ from the median are not included in the average. Ten spatial regions are extracted: the sum over the entire field and nine rectangular regions in a 3×3 grid (see Section 3). The summed spectrum is shown in Figure 2; selected redshifted absorption lines are marked.

3. Kinematics

3.1. Empirical Templates

The stellar kinematics of NGC1052-DF2 are measured by fitting template spectra to the extracted 1D spectrum. Key requirements are that the resolution of the template is well characterized and that template mismatch is minimized. This is not easily accomplished, given the high resolution of our data ($\approx 12 \text{ km s}^{-1}$ at $\lambda = 5000 \text{ \AA}$) and the low metallicity of the stellar population. We resolve these issues by using spatially integrated spectra of old, metal-poor Galactic globular clusters as templates, obtained with the same instrumental configuration.

The Galactic globular clusters M3 and M13 were observed on 2018 April 17. The metallicities of these clusters are $[\text{Fe}/\text{H}] \approx -1.5$ (Harris 1996), slightly lower than the expected metallicity of NGC1052-DF2 based on its stellar mass and slightly higher than that based on its velocity dispersion (Kirby et al. 2013; Gu et al. 2018). The total integration time was 600 s on each cluster, and 600 s on a nearby sky field. The data reduction and spectral extraction followed the same procedures as described in Section 2.2 for NGC1052-DF2; this ensures that any instrument-induced effects (such as small spatial variations in the wavelength calibration) are in common between the templates and the science data. The spectra are simple averages of the entire KCWI field, with individual stars contributing at most a few percent. We verified that the instrumental resolution, as measured from arc lamps, is consistent between the April globular clusters data and the October NGC1052-DF2 data.

3.2. Velocity Dispersion Measurement

The velocity dispersion was determined in the wavelength region $4830 \text{ \AA} < \lambda < 5235 \text{ \AA}$, using the M3 and M13 templates. The fit was performed with an implementation of the emcee Markov Chain Monte Carlo sampler (Foreman-Mackey et al. 2013), with the redshift and velocity dispersion as free parameters (see van Dokkum et al. 2016). Besides a multiplicative polynomial we fit for a third-order additive polynomial, to account for both sky subtraction errors and template mismatch. Varying the order of this polynomial does not change the results significantly. When fitting the full spectral range we also fit for any subtle wavelength calibration mismatch between the template and the data, parameterized as

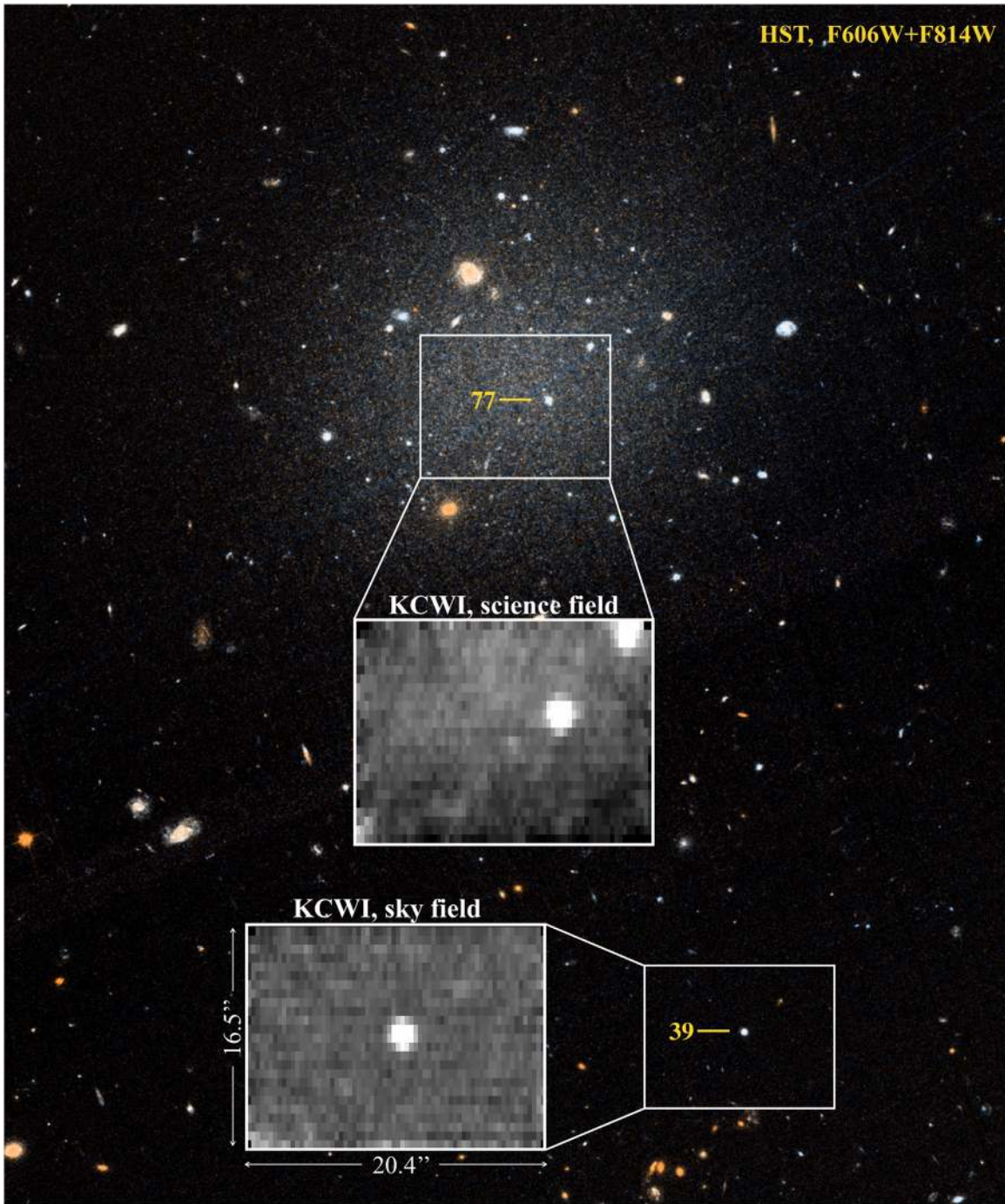


Figure 1. *HST*/ACS color image of NGC1052-DF2, created from the V_{606} and I_{814} bands. The white frames represent the two $20''.4 \times 16''.5$ KCWI pointings. The upper frame covers the diffuse light of the galaxy out to $0.7R_{\text{eff}}$. The lower frame was used for sky modeling. The insets show the collapsed summed KCWI data cubes.

an second-order polynomial with respect to the central wavelength. Although we find polynomial coefficients that are slightly different from zero, the resulting dispersion does not change when they are forced to zero. The best-fitting models are shown in Figure 2 by the red and blue lines. The errors describe the differences between the data and the models well; the reduced χ^2 values are 1.05 for the M3 template and 1.07 for the M13 template. The measured dispersion is the quadratic difference between the velocity dispersion of NGC1052-DF2 and that of the globular clusters. To obtain the stellar dispersion of NGC1052-DF2 we correct the

measured values:

$$\sigma_{\text{stars}}^2 = \sigma_{\text{meas}}^2 + \sigma_{\text{M3/M13}}^2, \quad (1)$$

with $\sigma_{\text{M3}} = 5.5 \pm 0.3 \text{ km s}^{-1}$ and $\sigma_{\text{M13}} = 7.1 \pm 0.4 \text{ km s}^{-1}$ (Harris 1996).

We find a stellar line-of-sight velocity dispersion of NGC1052-DF2 of $\sigma_{\text{stars}} = 7.9_{-3.4}^{+2.5} \text{ km s}^{-1}$ when fitting the M3 template and $\sigma_{\text{stars}} = 9.0_{-2.8}^{+2.0} \text{ km s}^{-1}$ when fitting the M13 template. These numbers are in excellent agreement. The mean is $\sigma_{\text{stars}} = 8.5_{-3.1}^{+2.3} \text{ km s}^{-1}$. We note that the lower bound of

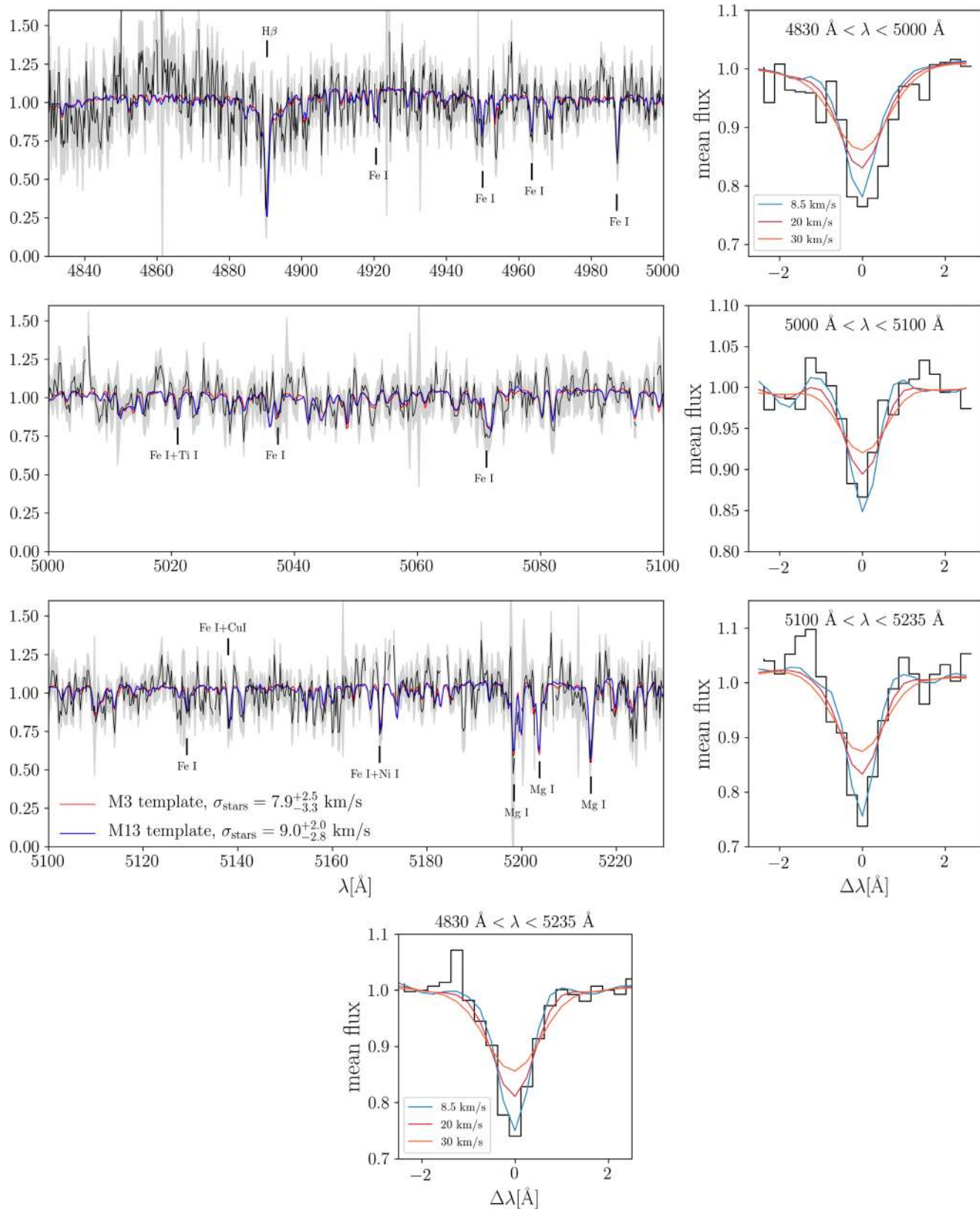


Figure 2. Main panels: integrated 2.67 hr KCWI spectrum of NGC1052-DF2 (black), with 1σ uncertainties (gray). The best fits of the two empirical M3 and M13 templates (see Section 3.1) that were used to determine the kinematics are shown in red and blue, respectively. The high resolution of KCWI allows us to detect a large number of absorption lines with high accuracy. Right: average of the 10 strongest absorption features in each spectral region, along with the best-fitting models and models with a higher dispersion. Bottom: average of the 20 strongest lines in the entire spectrum.

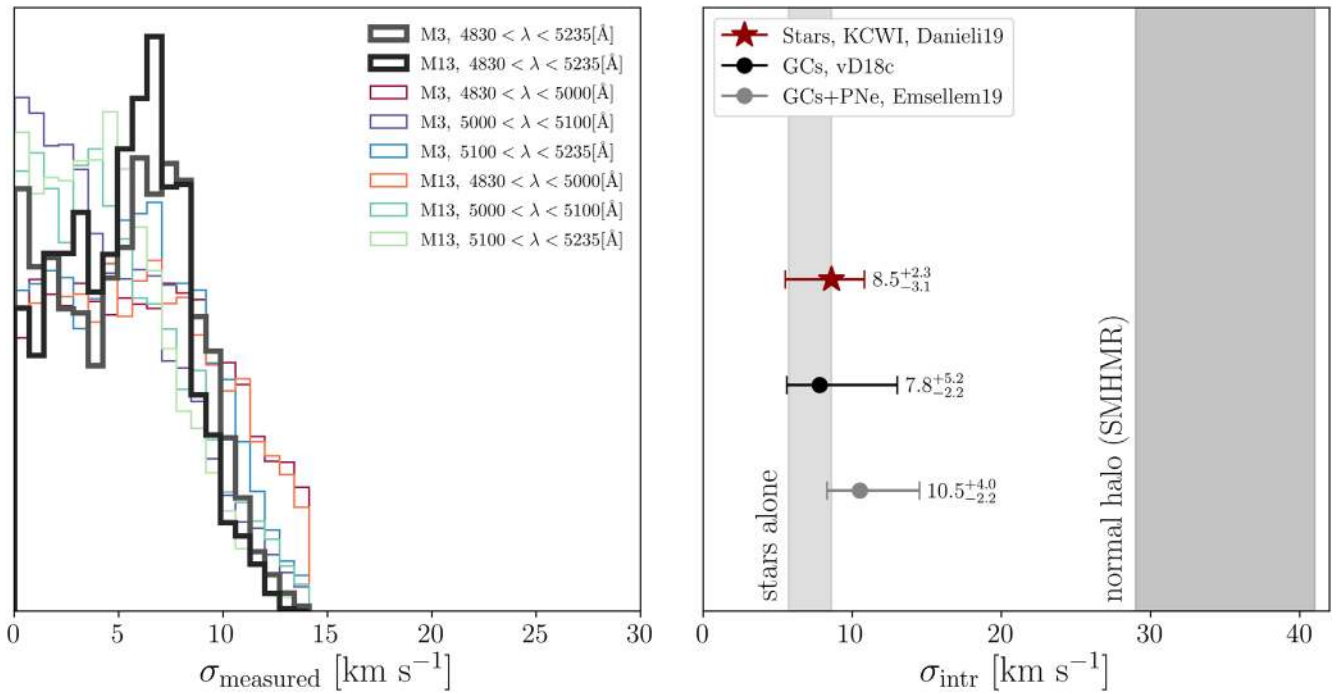


Figure 3. Left panel: posterior from eight MCMC runs of fitting the KCWI spectrum with two different templates and four different wavelength regimes. The upper limit is very well constrained in all MCMC runs with the full wavelength range giving the strongest constraints, as expected. Right panel: constraints on the intrinsic velocity dispersion of NGC1052-DF2. The stellar velocity dispersion measured in this study (dark red star) is consistent with the results obtained from 10 globular clusters in van Dokkum et al. (2018a). The stars alone contribute $\sigma_{\text{stars}} = 7.0^{+1.6}_{-1.3}$ km s⁻¹ (light gray band) and the expectation from the stellar mass–halo mass relation (Behroozi et al. 2013), assuming a standard NFW halo (Lokas & Mamon 2001), is $\sigma_{\text{SMHM}} = 35 \pm 6$ km s⁻¹.

5.4 km s⁻¹ is somewhat artificial, as it is partially determined by the internal dispersion of the globular clusters. The MC samples extend all the way to 0 km s⁻¹ (see Figure 3). The 95% confidence upper limit on the dispersion is 11.8 km s⁻¹. The central velocity dispersion for M13 is also somewhat uncertain; Kamann et al. (2014) find a higher value than Harris (1996), although they note that dispersions in this regime cannot be measured reliably given the instrumental resolution of PMAS ($\sigma_{\text{instr}} \approx 18$ km s⁻¹). An over- or underestimation of the intrinsic dispersion of M13 would result in a slightly different inferred dispersion for NGC1052-DF2. In that context, it is reassuring that the two independent measurements using the M3 and M13 templates are fully consistent with each other.

We performed two further tests of the stability of this result. First, the spectrum was split into three wavelength regions. The first region is dominated by H β , the second by relatively weak Fe lines, and the third by Mg (see Figure 2). The MCMC posteriors for these fits are shown in the left panel of Figure 3, and the corrected velocity dispersions are listed in Table 1. All inferred dispersions are consistent within 1.5 km s⁻¹. Next, we split the data into nine spatial bins and fitted those independently. In all cases, the best-fit dispersion is well within 1 σ of the value from fitting the full wavelength range. An additional test we performed was fitting the M3 spectrum with the M13 spectrum as a template. The measured dispersion of M3 is consistent with zero and the intrinsic dispersion is consistent with the value from the literature (Pryor & Meylan 1993; Harris 1996).

The robustness of our results is illustrated by the small panels in Figure 2, where we show the average observed absorption in the spectral regions corresponding to the strongest absorption lines in the templates. The 8.5 km s⁻¹ model is an excellent fit for all wavelength regions and also for the average of the 20 strongest lines in the entire spectrum.

Table 1
NGC1052-DF2 Stellar Dispersions

Template	λ (Å)	σ_{meas} (km s ⁻¹)	σ_{stars} (km s ⁻¹)
M3	[4830, 5235]	$5.8^{+3.1}_{-4.1}$	$7.9^{+2.5}_{-3.4}$
M13	[4830, 5235]	$5.5^{+2.5}_{-3.6}$	$9.0^{+2.0}_{-2.8}$
M3	[4830, 5000]	$6.4^{+4.7}_{-4.4}$	$8.5^{+4.0}_{-3.7}$
M3	[5000, 5100]	$4.1^{+4.1}_{-2.9}$	$6.9^{+3.2}_{-2.3}$
M3	[5100, 5235]	$5.7^{+3.7}_{-3.9}$	$7.9^{+3.1}_{-3.2}$
M13	[4830, 5000]	$6.2^{+4.7}_{-4.2}$	$9.4^{+3.8}_{-3.4}$
M13	[5000, 5100]	$4.4^{+4.2}_{-3.1}$	$8.4^{+3.2}_{-2.3}$
M13	[5100, 5235]	$4.3^{+3.3}_{-2.9}$	$8.3^{+2.5}_{-2.2}$

Measuring velocity dispersions in the 10–30 km s⁻¹ regime is well-suited to KCWI.

The inferred intrinsic stellar velocity dispersion is consistent with the constraints on the velocity dispersion derived using globular clusters in van Dokkum et al. (2018a) with $\sigma_{\text{gc}} = 7.8^{+5.2}_{-2.2}$ km s⁻¹ and in Emsellem et al. (2019) with an estimated value of $\sigma_{\text{gc}} = 10.5^{+4.0}_{-2.2}$ km s⁻¹. These various results are shown in the right panel of Figure 3.

3.3. Systemic Velocity and Stellar Velocity Field

Besides the velocity dispersion we also obtain a measurement of the mean systemic velocity. The best-fit values for the M3 and M13 templates are $v_{\text{stars}} = 1805.2^{+1.1}_{-1.1}$ km s⁻¹ and $v_{\text{stars}} = 1804.7^{+1.0}_{-1.1}$ km s⁻¹, respectively. The two values are consistent with each other and also with the mean velocity of the 10 globular clusters as measured in van Dokkum et al. (2018b): $\langle v_{\text{gc}} \rangle = 1803^{+2}_{-2}$ km s⁻¹.

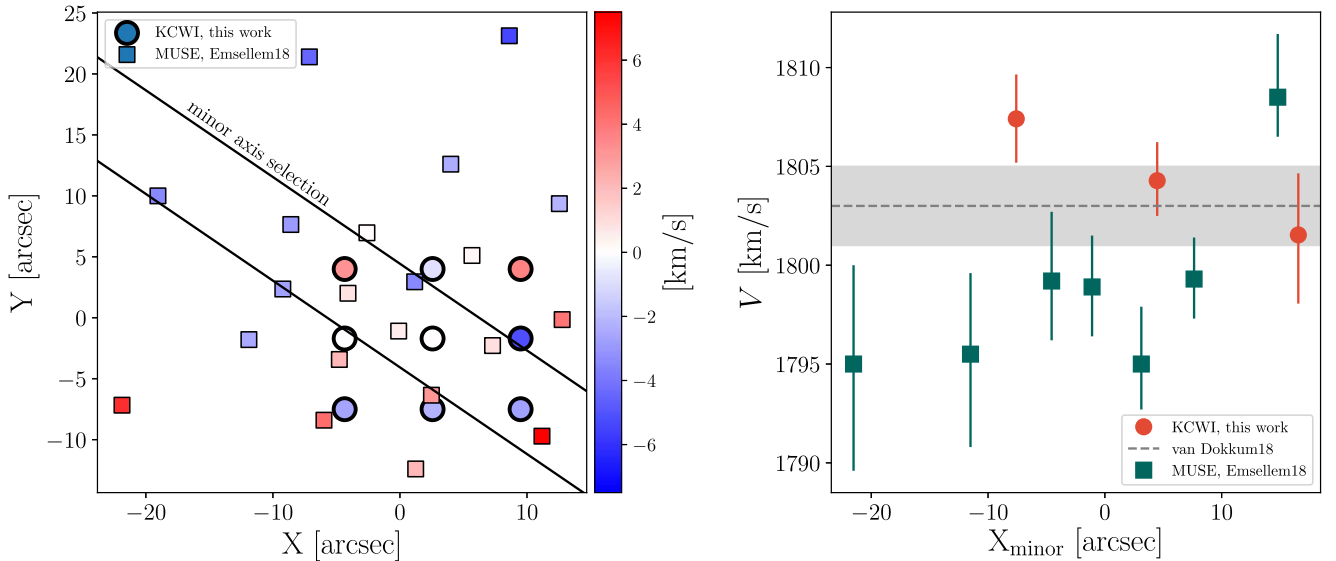


Figure 4. Left panel: velocity field measured from stars in different spatial positions of the galaxy, as measured by *MUSE* (squares; Emsellem et al. 2019) and KCWI (circles). The velocities are relative to the mean velocity in each data set. Right panel: absolute velocities in a band (black diagonal lines) roughly corresponding to the minor axis. There is an offset between the *MUSE* and KCWI velocities, and no clear gradient in our data.

Next, we examine the systemic velocities in the nine spatial bins described in Section 3.2. The rms scatter among the nine velocities is 2.8 km s^{-1} . This is very similar to the mean velocity uncertainty (2.2 km s^{-1}), leaving little room for velocity gradients of the same order as the velocity dispersion. In the left panel of Figure 4 we show the stellar velocity field. We find no clear gradient in the velocities measured from our data within our FOV. This is in contrast to results from Emsellem et al. (2019), who report a gradient of $2.8 \pm 0.9 \text{ km s}^{-1}$ per $10''$ along the minor axis. In the right panel of Figure 4 we compare the absolute velocities of several spatial bins along the minor axis directly, as were obtained from our data (orange circles) and from the *MUSE* data (green squares). We shift the *MUSE* velocities by 5.4 km s^{-1} to account for the different applied redshift-velocity transformation ($c z$ in our analysis versus $c \ln(1+z)$ in Emsellem et al. 2019). If there is a slight trend in the KCWI data, it is in the opposite direction from that seen in the *MUSE* data. Given the large offset in the absolute velocities between the *MUSE* results on the one hand and the KCWI, LRIS, and DEIMOS results on the other, and the lack of a clear trend in our data, we do not confirm the presence of the gradient claimed in Emsellem et al. (2019).⁹ We note that the *MUSE* velocities are consistent with ours in the lower (Southern) part of the *MUSE* data cube.

4. Discussion

In this Letter we have presented stellar kinematics measurements for the galaxy NGC1052-DF2, using high-resolution ($\sim 12 \text{ km s}^{-1}$) integral-field spectroscopy with the KCWI on Keck II. We measure a systemic velocity of $\langle v_{\text{stars}} \rangle = 1804.9_{-1.1}^{+1.0} \text{ km s}^{-1}$, and confirm that the 10 star clusters that were previously used to constrain the kinematics of NGC1052-DF2 (van Dokkum et al. 2018a; Wasserman et al. 2018) are indeed associated with the diffuse stellar light of the galaxy.

⁹ These authors have also measured the stellar velocity dispersion, but owing to the complexities of measuring well below the instrumental resolution, the final values are not yet known at the time of writing and we therefore cannot compare them directly to ours (E. Emsellem 2019, private communication).

Thanks to the exceptionally high resolution of KCWI, we are providing a robust measurement of the stellar velocity dispersion of NGC1052-DF2. We measure $\sigma_{\text{stars}} = 8.5_{-3.1}^{+2.3} \text{ km s}^{-1}$ within the effective radius, consistent with the revised dispersion of $\sigma_{\text{gc}} = 7.8_{-2.2}^{+5.2} \text{ km s}^{-1}$ measured from the 10 globular clusters using the maximum likelihood method in van Dokkum et al. (2018a).

In Figure 5 we show the distribution of Local Group galaxies in the plane of velocity dispersion versus stellar mass, using the compilation of McConnachie (2012). Several galaxies are displayed as model images, created with the ArtPop code (Danieli et al. 2018). This figure graphically illustrates the unusual nature of NGC1052-DF2: the galaxy combines a relatively high stellar mass with a large size and a very low velocity dispersion. The dashed line is an indicative relation between stellar mass and velocity dispersion in the absence of dark matter, for the radial regime where the dispersion profile is approximately isothermal: $\sigma \sim 5 \times 10^{-4} (M_*/M_\odot)^{0.5} \text{ km s}^{-1}$. Typical dwarf galaxies fall above the line, as they are dark matter dominated, but NGC1052-DF2 is on the line within the errors.

We quantify this by using the newly measured stellar velocity dispersion along with the projected circularized half-light radius of $R_{\text{e,c}} = 2.0 \text{ kpc}$ (Cohen et al. 2018) to determine the dynamical mass of NGC1052-DF2 within the 3D half-light radius $r_{1/2} \approx 4/3 R_{\text{e,c}}$ (Wolf et al. 2010). We find $M(r < r_{1/2}) = 1.3 \pm 0.8 \times 10^8 M_\odot$. The stellar mass within the half-light radius is $M_{\text{stars}}(r < r_{1/2}) = 1.0_{-0.2}^{+0.2} \times 10^8 M_\odot$ (see van Dokkum et al. 2018b), and we infer that the dynamical mass is consistent with the mass in stars alone. We refer to Wasserman et al. (2018) for quantitative constraints on the halo mass as derived from the globular clusters; our measurement confirms the central assumption in Wasserman et al. (2018) that the globular clusters indeed trace the potential of NGC1052-DF2. We note that if NGC1052-DF2 is a thin rotating disk seen close to face-on, its axis ratio of 0.85 implies inclination-corrected velocities that are (at most) a factor of 1.9 higher than the observed ones (van Dokkum et al. 2018b). This scenario is unlikely given the lack of detected rotation along the major axis or in the globular clusters, and the discovery of a second galaxy

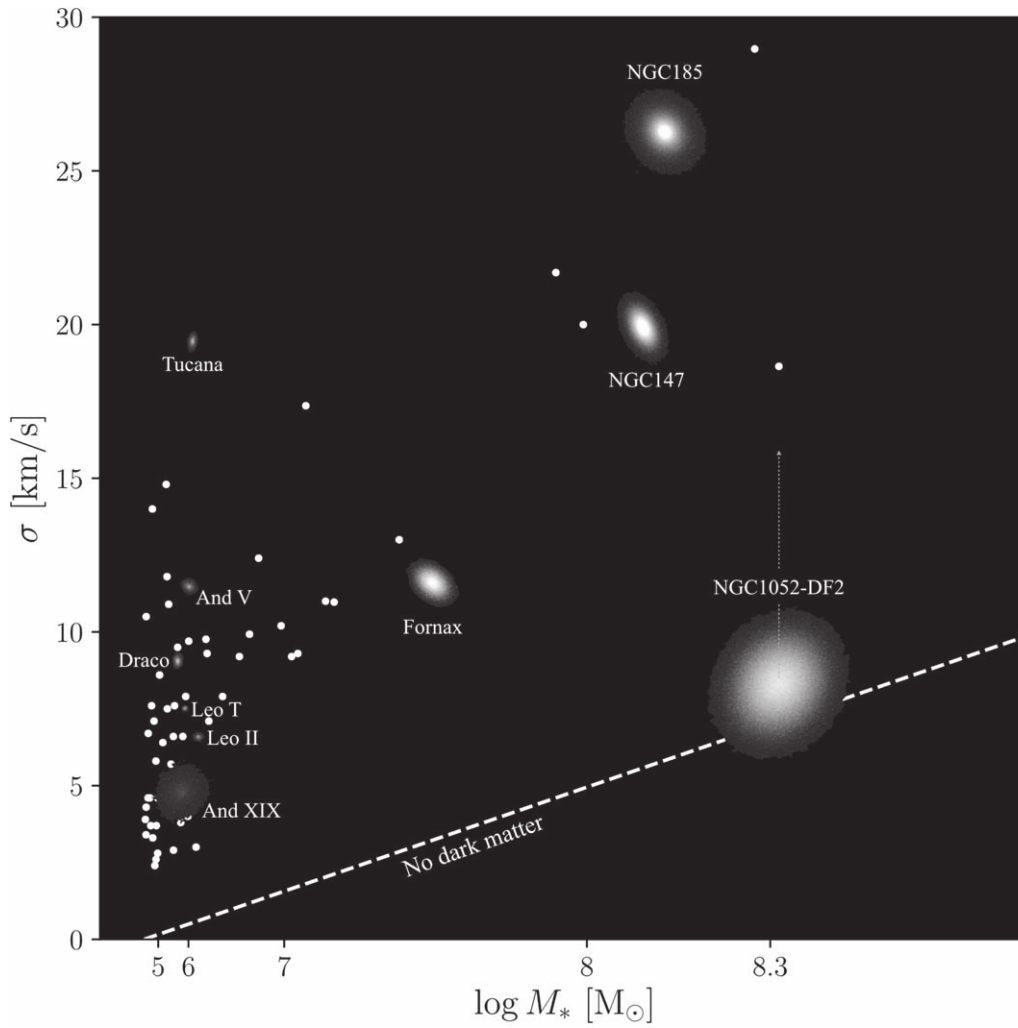


Figure 5. Local Group dwarf galaxies (white dots) and NGC1052-DF2 in the velocity dispersion–stellar mass plane. NGC1052-DF2 and several illustrative Local Group galaxies were modeled and placed at 20 Mpc (using the ArtPop code, Danieli et al. 2018). Its large size and relatively high stellar mass, along with its low velocity dispersion, distinguish it from Local Group galaxies. The dashed line is the approximate relation between stellar mass and velocity dispersion in the absence of dark matter.

missing dark matter, NGC1052-DF4, in the same group (van Dokkum et al. 2019). Nevertheless, for consistency with the other data points, we show this inclination correction with a dotted line in Figure 5.

Our study confirms that NGC1052-DF2 has far less dark matter than expected, and perhaps no dark matter at all. Future studies can examine what physical processes and formation schemes can result in this deficiency of dark matter on kiloparsec scales. This is particularly challenging given that other similar-looking UDGs appear to have normal (or even “overmassive”) halos (see, e.g., Toloba et al. 2018). It is now critical to determine whether NGC1052-DF2 is a unique galaxy or whether this “missing dark matter problem” is relatively common. If it is, it implies that the scatter in stellar mass at low halo masses is extremely large (see, e.g., Behroozi et al. 2018).

Another essential question is whether other properties of NGC1052-DF2, such as its nature as a UDG and its intriguing population of globular clusters, are related to its dark matter deficiency. Finding a closer-by system with a low velocity dispersion would allow us to constrain its properties (even) more accurately, and place strong constraints on dark matter and halo models.

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References

- Aaronsen, M. 1983, *ApJL*, 266, L11
 Beasley, M. A., Romanowsky, A. J., Pota, V., et al. 2016, *ApJL*, 819, L20
 Behroozi, P., Wechsler, R., Hearin, A., & Conroy, C. 2018, arXiv:1806.07893
 Behroozi, P. S., Conroy, C., & Wechsler, R. H. 2010, *ApJ*, 717, 379
 Behroozi, P. S., Wechsler, R. H., & Conroy, C. 2013, *ApJ*, 770, 57
 Cohen, Y., van Dokkum, P., Danieli, S., et al. 2018, *ApJ*, 868, 96

- Danieli, S., van Dokkum, P., & Conroy, C. 2018, *ApJ*, 856, 69
- de Blok, W. J. G., McGaugh, S. S., & Rubin, V. C. 2001, *AJ*, 122, 2396
- Emsellem, E., van der Burg, R. F. J., Fensch, J., et al. 2019, *A&A*, submitted (arXiv:1812.07345)
- Foreman-Mackey, D., Hogg, D. W., Lang, D., & Goodman, J. 2013, *PASP*, 125, 306
- Geha, M., Willman, B., Simon, J. D., et al. 2009, *ApJ*, 692, 1464
- Gu, M., Conroy, C., & Behroozi, P. 2016, *ApJ*, 833, 2
- Gu, M., Conroy, C., Law, D., et al. 2018, *ApJ*, 859, 37
- Harris, W. E. 1996, *AJ*, 112, 1487
- Kamann, S., Wisotzki, L., Roth, M. M., et al. 2014, *A&A*, 566, A58
- Kirby, E. N., Cohen, J. G., Guhathakurta, P., et al. 2013, *ApJ*, 779, 102
- Laporte, C. F. P., Agnello, A., & Navarro, J. F. 2018, *MNRAS*, 484, 245
- Łokas, E. L., & Mamon, G. A. 2001, *MNRAS*, 321, 155
- Mandelbaum, R., Seljak, U., Kauffmann, G., Hirata, C. M., & Brinkmann, J. 2006, *MNRAS*, 368, 715
- Martin, N. F., Collins, M. L. M., Longeard, N., & Tollerud, E. 2018, *ApJL*, 859, L5
- Martin, N. F., Geha, M., Ibata, R. A., et al. 2016, *MNRAS*, 458, L59
- McConnachie, A. W. 2012, *AJ*, 144, 4
- Morrissey, P., Matuszewski, M., Martin, C., et al. 2012, *Proc. SPIE*, 8446, 844613
- Morrissey, P., Matuszewski, M., Martin, D. C., et al. 2018, *ApJ*, 864, 93
- Moster, B. P., Somerville, R. S., Maulbetsch, C., et al. 2010, *ApJ*, 710, 903
- Oman, K. A., Navarro, J. F., Sales, L. V., et al. 2016, *MNRAS*, 460, 3610
- Pryor, C., & Meylan, G. 1993, in *ASP Conf. Ser. 50, Structure and Dynamics of Globular Clusters*, ed. S. G. Djorgovski & G. Meylan (San Francisco, CA: ASP), 357
- Toloba, E., Lim, S., Peng, E., et al. 2018, *ApJL*, 856, L31
- Torrealba, G., Belokurov, V., Koposov, S. E., et al. 2018, arXiv:1811.04082
- Vale, A., & Ostriker, J. P. 2004, *MNRAS*, 353, 189
- van den Bosch, F. C., & Swaters, R. A. 2001, *MNRAS*, 325, 1017
- van Dokkum, P., Abraham, R., Brodie, J., et al. 2016, *ApJL*, 828, L6
- van Dokkum, P., Cohen, Y., Danieli, S., et al. 2018a, *RNAAS*, 2, 54
- van Dokkum, P., Danieli, S., Abraham, R., Conroy, C., & Romanowsky, A. J. 2019, arXiv:1901.05973
- van Dokkum, P., Danieli, S., Cohen, Y., et al. 2018b, *Natur*, 555, 629
- van Dokkum, P. G., Abraham, R., Merritt, A., et al. 2015, *ApJL*, 798, L45
- Wasserman, A., Romanowsky, A. J., Brodie, J., et al. 2018, *ApJL*, 863, L15
- Wechsler, R. H., & Tinker, J. L. 2018, *ARA&A*, 56, 435
- Wolf, J., Martinez, G. D., Bullock, J. S., et al. 2010, *MNRAS*, 406, 1220