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# Structural and magnetic properties and DFT analysis of ZnO:(Al,Er) nanoparticles

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The structural, optical and electrical properties of un-doped and (Al,Er) co-doped zinc oxide (ZnO) powders synthesized by hydrothermal method were investigated. The obtained samples were characterized by X-ray diffraction (XRD), transmission electron microscopy (TEM), energy-dispersive spectroscopy and magnetic measurements. XRD results reveal that the incorporation of Al and Er in ZnO matrix leads to the formation of a nanostructured hexagonal (würtzite) ZnO structure and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> secondary phase. High-resolution transmission electron microscopy image also shows the hexagonal shape of the ZnO nanoparticles. The magnetic behavior of the nanoparticles changes with concentration of dopant elements due to the competition between oxygen vacancies, secondary phase effect and exchange interaction between dopant elements.

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## 1. Introduction

Development in materials technology over the last few years realizes the usage of wide-band gap semiconductors for various optoelectronic device applications. The semiconductor ZnO is an attractive material in the research community. It has a wide range of properties that depend on doping, including a range of conductivity from metallic to insulating (including n-type and p-type conductivity), high transparency, piezoelectricity, wide-band gap semiconductivity, room-temperature ferromagnetism, and huge magneto-optic and chemical-sensing effects.<sup>1</sup>

ZnO can be doped with a wide variety of ions to meet the demands of several application domains.<sup>2–4</sup> The elements used for doped ZnO are part low valent metals (Al, In, Ga) belonging to group III of the periodic table and rare earth elements (lanthanides) (Er, Nd, Sm, Eu, Tm). Al has been the most used dopant element due to its small ionic radius and low material cost. Al doped ZnO nanopowders possess room temperature ferromagnetism.<sup>5</sup> In addition to that, the substitution of Zn<sup>2+</sup> ions with Al<sup>3+</sup> in ZnO lattice improves the electrical conductivity

through the increase of charge carriers where it is reported that the electron concentration increases from 10<sup>16</sup> to 10<sup>21</sup>/cm<sup>-3</sup>.<sup>6</sup> Dghoughi *et al.*<sup>7</sup> reported that 5 mol% Al doped ZnO film exhibited the preferential (002) orientation, and it also had a high transmittance, and minimum resistivity. Effects of the doping concentrations and annealing temperatures on the properties of the high oriented Al-doped ZnO films were deeply investigated.<sup>8</sup>

ZnO nanoparticles (NPs) doped with rare-earth ions are expected to be novel optical materials because of their sharp and intense optical emission spectra.<sup>9</sup> Er-doped ZnO has attracted attention for use in optical communication applications because of its 1.54  $\mu\text{m}$  emission.<sup>10</sup> Lamrani *et al.* studied the effects of the dopant erbium on the structural, morphological, luminescent and nonlinear optical properties.<sup>11</sup> They found that addition of erbium could effectively control the film surface morphology and its cathodoluminescent properties, and that the best crystallization and morphology were obtained at lower Er doping concentration. Chen *et al.*<sup>12</sup> reported that, when the dopant Er was added, the crystalline quality of the doped ZnO films decreased due to the oxygen deficiency. Zhang *et al.*<sup>13</sup> reported that the photoluminescence of Er-doped ZnO had an improvement in the visible emissions, compared with that of the pure ZnO, which was attributed to the defects or vacancies in the doped film.

The co-doping ZnO with erbium and aluminum ZnO:(Al,Er) could significantly improve its responses to visible light.<sup>14</sup> Various fabrication methods such as ball milling, hydrothermal, co-precipitation, spray pyrolysis, solution combustion, microwave method, sol-gel, vapor deposition *etc.* were employed to fabricate various ZnO nanostructures.<sup>15</sup> In this paper ZnO and (Al,Er) co-doped ZnO powders were fabricated by

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hydrothermal method. This technique is a promising alternative synthetic method because of the low process temperature, its suitability to any type of doping and very easy to control the particle size. The hydrothermal process has several advantages over other growth processes such as use of simple equipment, catalyst-free growth, low cost, ease of large scale production, eco-friendly and less hazardous.<sup>16</sup> This method has also been successfully employed to prepare nano scale ZnO and other luminescent materials. The influence of Al and Er dopants on the properties of ZnO were analyzed by performing various spectroscopic investigations using XRD, TEM, EDX and magnetic characterizations.

## 2. Experimental details

### 2.1. Synthesis

The pure and co-doped ZnO NPs with Er and Al were prepared by hydrothermal synthesis method to study their structural and magnetic properties. The chemicals used in the experiment are zinc nitrate  $Zn(NO_3)_2$ , erbium nitrate  $Er(NO_3)_3$  and aluminum nitrate  $Al(NO_3)_3$  (Table 1).

For un-doped ZnO NPs, 4 moles of zinc nitrate  $Zn(NO_3)_2$  were dissolved in 20 ml of ethylene glycol resulting in mixture of brown dispersion. After that, 8 moles of NaOH dissolved, separately, in 8 ml of water were added drop by drop to the above solution (500  $\mu$ l in one min). For co-doped ZnO NPs, a predetermined quantities of  $Er(NO_3)_3$  and  $Al(NO_3)_3$  were added to the above mixture. The mixture was maintained under agitation during 2 hours and then transferred in a container containing 45 ml of acid solvent. The hydrothermal treatment was carried out at 165 °C during 24 hours. The remaining solution was washed with distilled water and ethanol to separate impurities attached with nanoparticles. The as formed solution was kept in an oven at around 100 °C till the precipitate dried and finally grinded to obtain powder.

### 2.2. Characterization

The structural characterization of the samples were studied using X-ray diffraction (XRD, X'Pert Pro Super, Philips Co., the Netherlands) at room temperature, with monochromatic Cu-K $\alpha$  ( $\lambda_{Cu} = 1.54156$  Å), in a ( $\theta$ -2 $\theta$ ) Bragg-Brentano geometry. The scan range was carried out between 20° and 90° with a step increment of 0.02°. High resolution transmission electron microscopy analysis was done by using JEM-2010 transmission electron microscope operating at 200 kV equipped with energy

dispersive X-ray (EDX). Magnetic properties were measured by a Quantum Design superconducting quantum interference device (SQUID) magnetometer. We used a kind of plastic sample holder and cotton to fix the powder then we measure the sample holder and cotton alone as background, which will be removed after measuring our samples.

## 3. Results and discussion

### 3.1. Structure and microstructure analyses

The XRD pattern of the un-doped ZnO NPs is shown in Fig. 1. The diffraction peaks can be indexed with the hexagonal wurtzite type structure (space group  $P6_3mc$ ), which are in agreement with JCPDS card No. 75-0576. For the (Al,Er) co-doped ZnO samples, the characteristic peaks related to Er and Al impurities do not appear in XRD patterns in comparison with the standard XRD patterns of bulk ZnO. However, additional weak peaks of  $Al_2O_3$  appeared at about 43.35°, 50.91° and 61.38° (Fig. 2a-c), which shows that the phase segregation occurred in this system. The intensity of these peaks goes on increasing as Er content in the sample increased. Thus, it is understood that  $Al^{3+}$  and  $Er^{3+}$  ions were partially substituted into  $Zn^{2+}$  ions sites within the ZnO crystal due to their low solubility limit in the parent ZnO. In addition to that, all samples shows a preferred orientation along the (101) direction of the ZnO phase.

The increase of Er content from 1 to 3% causes notable increase in all ZnO peak intensities (Fig. 2d). However, the

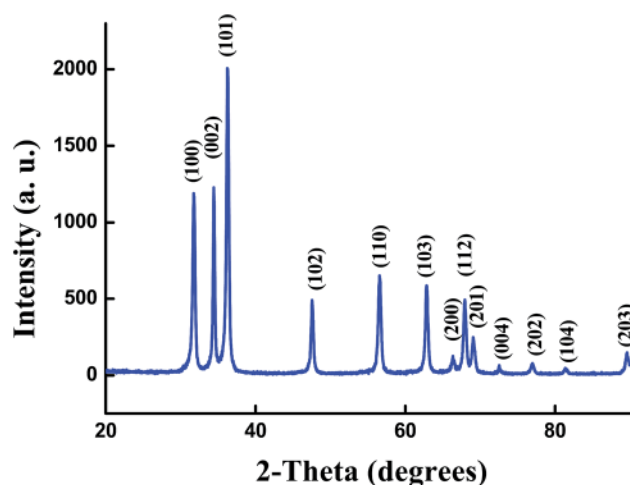


Fig. 1 X-ray diffraction pattern of un-doped ZnO NPs.

Table 1 Quantities of chemicals used for preparing un-doped and co-doped ZnO NPs

Composition	Amount of $Zn(NO_3)_2$ (297.49 g mol <sup>-1</sup> ) (purity: 99.97%)	NaOH (40 g mol <sup>-1</sup> )	Amount of $Er(OAc)_3 \cdot 4H_2O$ (416.45 g mol <sup>-1</sup> ) (purity: 99.9%)	Amount of $Al(NO_3)_3 \cdot 9H_2O$ (375.13 g mol <sup>-1</sup> ) (purity: 99.98%)
Pure ZnO	4 mmol in 20 ml (1.18 g)	8 mmol (320 mg) in 8 ml water	—	—
$Zn_{0.96}Al_{0.03}Er_{0.01}O$	4 mmol (1.18 g)	8 mmol (320 mg)	0.04 mmol (0.016 g)	0.12 mmol (0.045 g)
$Zn_{0.95}Al_{0.03}Er_{0.02}O$	4 mmol (1.18 g)	8 mmol (320 mg)	0.08 mmol (0.033 g)	0.12 mmol (0.045 g)
$Zn_{0.94}Al_{0.03}Er_{0.03}O$	4 mmol (1.18 g)	8 mmol (320 mg)	0.12 mmol (0.050 g)	0.12 mmol (0.045 g)



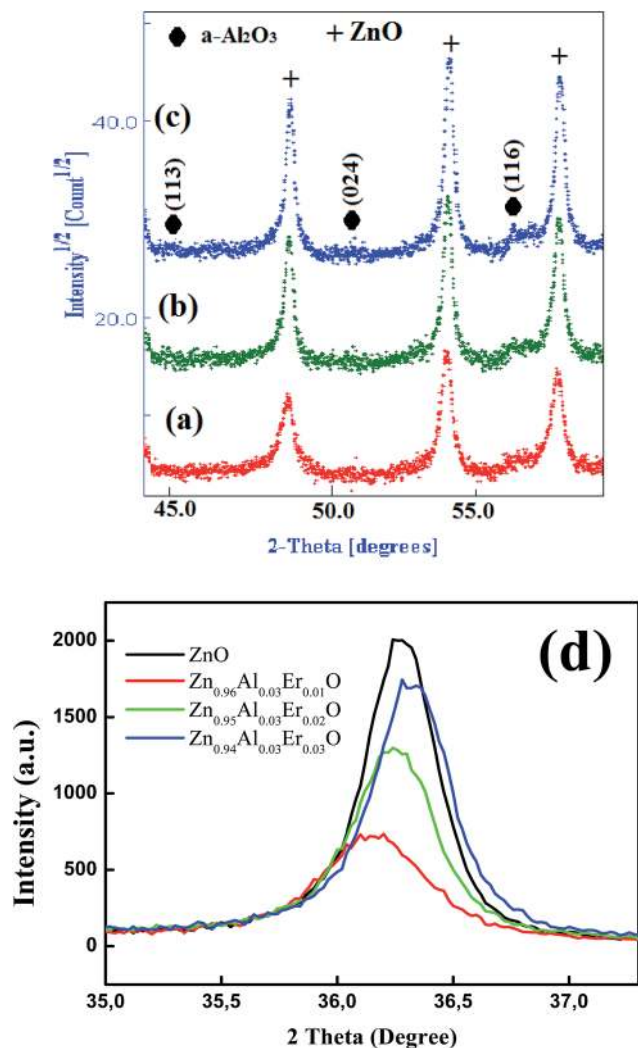


Fig. 2 Zoom of Al<sub>2</sub>O<sub>3</sub> diffraction peaks in (a) Zn<sub>0.96</sub>Al<sub>0.03</sub>Er<sub>0.01</sub>O, (b) Zn<sub>0.95</sub>Al<sub>0.03</sub>Er<sub>0.02</sub>O, (c) Zn<sub>0.94</sub>Al<sub>0.03</sub>Er<sub>0.03</sub>O and (d) main (101) ZnO diffraction peak evolution with dopant content.

FWHM decreases with increasing Er concentration. More intense diffraction peaks of co-doped samples point out that Er doping enhances crystallinity compared to ZnO:(3% Al,1% Er) sample. Interestingly, the XRD pattern clearly shows that there is a shift in the peak position with increasing (Al,Er) content, suggesting a change in the lattice parameters. To characterize in detail the crystalline structure, including both the structural (cell parameters, site occupancy, atomic positions, thermal parameters, *etc.*) and microstructural parameters (crystallite size, microstrain, preferred orientation, *etc.*), we have used Rietveld refinement method. In this method, the peak shape profile is fitted by Voigt or Pseudo-Voigt functions.<sup>17</sup> The Marquardt least-squares procedure was adopted for minimizing the difference between the observed and simulated powder diffraction patterns and the minimization was carried out by using the reliability parameters  $R_{wp}$  (weighted residual factor),  $R_B$  (Bragg factor) and  $R_{exp}$  (expected residual factor).<sup>17,18</sup> The goodness of fit (GoF) is obtained by comparing  $R_{wp}$  with  $R_{exp}$  ( $GoF = R_{wp}/R_{exp}$ ). Refinement continues till convergence is

reached with the value of the quality factor, GoF approaching 1, which confirms the goodness of refinement. The typical refinement plots of un-doped and (Al,Er) co-doped ZnO NPs are shown in Fig. 3 where the calculated patterns are shown in the same field as a solid line curve. The difference between the observed and calculated intensities is shown in the middle field. The best Rietveld refinement is obtained by the introduction of two phases: hexagonal wurtzite ZnO with the space group  $P6_3mc$  and trigonal Al<sub>2</sub>O<sub>3</sub> with the space group  $R\bar{3}c:H$ . The obtained structural and microstructural parameters are listed in Table 2.

As it can be seen from Table 2, the crystallite size decreases from 35.43 nm for un-doped ZnO to 24.83 nm for Zn<sub>0.96</sub>Al<sub>0.03</sub>Er<sub>0.01</sub>O, and then with increasing the Er doping concentration, it slightly increases reaching 28.97 nm for 3% Er. The co-doped ZnO NPs exhibit smaller crystallite size than the un-doped ZnO NPs, which may change the lattice parameters due to differences between ionic radii of Er<sup>3+</sup>, Al<sup>3+</sup> and Zn<sup>2+</sup>.<sup>19</sup> This mismatch may generate lattice strain (microstrain). This decrease of the crystallite size is further an evidence that the co-doping with (Al,Er) inhibits grain growth of ZnO, as it has been reported by other authors for ZnO:(Al,Er)<sup>20,21</sup> and ZnO:Al<sup>22</sup> systems. Senol *et al.*<sup>23</sup> have reported that the addition of 5% of Er to ZnO NPs lead to increasing the crystallite size from 29.35 to 63.82 nm. These authors suggested that the increase of Er concentrations in ZnO matrix will increase the nucleation of the particles, which enhances the grain growth of ZnO:Er NPs.<sup>23</sup>

The entering of Er<sup>3+</sup> and/or Al<sup>3+</sup> ions into the ZnO structure may cause the formation of more nucleation center. By increasing Er<sup>3+</sup> ion concentration, the intensity of the diffraction peaks increased and full width at half-maximum (FWHM) decreased. In fact, FWHM of the main pic (101) decreases from 0.58° to 0.40° for 1% and 3% Er, respectively. Similar results were obtained in other system such as Gd doped ZnO.<sup>24</sup> It is interesting that by increasing (Er,Al) doping concentration we saw that the crystallite size is increased. The obtained values are lower than those observed for un-doped sample. This may be attributed to the increase of the defects in the crystal structure, such as oxygen vacancies. Dislocation density ( $\delta$ ) that is a measure of the amount of defects in the crystal structure is defined as the length of the dislocation lines per unit volume and it can be found by:<sup>25</sup>

$$\delta = 1/d^2 \quad (1)$$

where  $d$  is the crystallite size. The results are given in Table 2.

It is well known that the lattice parameters of a semiconductor usually depend on the following factors: (i) free electron concentration acting *via* deformation potential of a conduction band minimum occupied by these electrons, (ii) concentration of foreign atoms and defects and their difference of ionic radii with respect to the substituted matrix ion, (iii) external strains (*e.g.*, those induced by substrate), and (iv) temperature. From Rietveld refinement, the obtained lattice constants of pure ZnO are  $a = 3.2495$  Å and  $c = 5.2085$  Å. These values deviate slightly from that of the ideal wurtzite crystal ( $a = 3.2417$  Å and  $c = 5.1876$  Å). This is probably due to lattice stability and ionicity. It has been reported that free charge is the



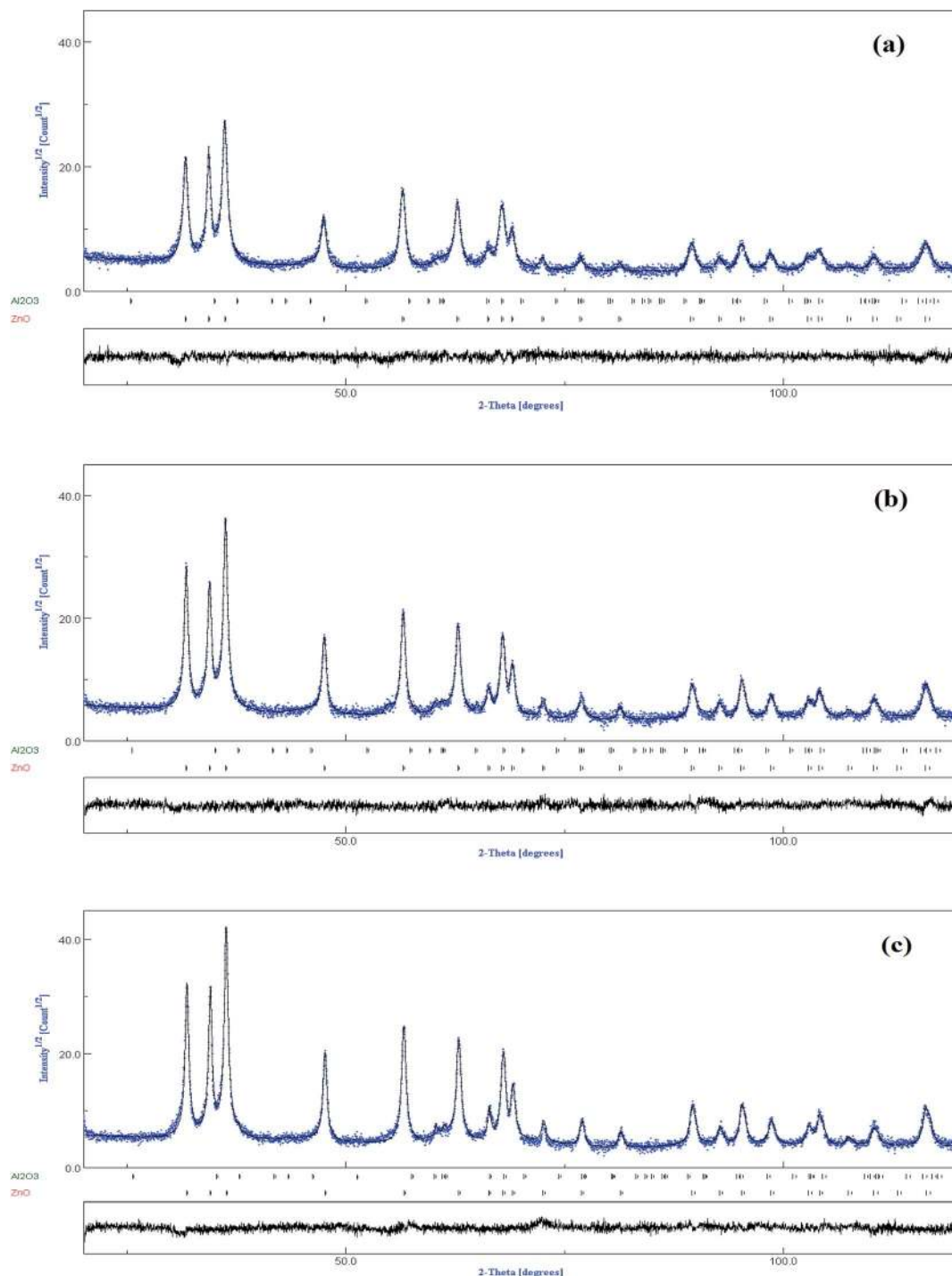


Fig. 3 Rietveld refinement of the XRD pattern of (a)  $\text{Zn}_{0.96}\text{Al}_{0.03}\text{Er}_{0.01}\text{O}$  (b)  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$  and (c)  $\text{Zn}_{0.94}\text{Al}_{0.03}\text{Er}_{0.03}\text{O}$  NPs. Experimental (dots) and calculated (full line) patterns are shown. Difference is given below.

dominant factor responsible for expanding the lattice proportional to the deformation potential of the conduction band minimum and inversely proportional to the carrier density and bulk modulus.<sup>26</sup> The point defects such as zinc antisites, oxygen vacancies, and extended defects, such as threading dislocations, also increase the lattice constant, albeit to a lesser extent in the hetero-epitaxial layers. For  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.01}\text{O}$  sample, the

lattice parameters  $a$  and  $c$  of the wurtzite phase increases to  $a = 3.2561 \text{ \AA}$  and  $c = 5.2109 \text{ \AA}$ . This increase may be attributed to the difference in the ionic radii of  $\text{Al}^{3+}$  ( $0.54 \text{ \AA}$ ) and  $\text{Er}^{3+}$  ( $0.89 \text{ \AA}$ ) as compared to  $\text{Zn}^{2+}$  ( $0.74 \text{ \AA}$ ) one. This result is justified by the shift of (101) main peak toward lower angle (Fig. 2d).

Increasing Er content (2%, 3%) is observed to reduce the lattice parameters as indicated in Table 2. However, the



Table 2 Structural parameters deduced from Rietveld refinement<sup>a</sup>

	Phases	Lattice parameters (Å) (±0.001)		$d$ (nm) (±2)	$\langle\sigma^2\rangle^{1/2}$ (%) (±0.02)	Amount (%) (±0.1)	Refinement parameters		Dislocation density ( $\times 10^{15}/\text{m}^2$ )	$R$
		$a$	$c$				$R_B$ (%)	GoF		
Pure ZnO	ZnO	3.2495	5.2085	35.431	0.115	100	6.879	1.100	0.796	1.1274
$\text{Zn}_{0.96}\text{Al}_{0.03}\text{Er}_{0.01}\text{O}$	ZnO	3.2561	5.2109	24.832	0.117	95.214	10.053	0.993	1.622	1.0203
	$\text{Al}_2\text{O}_3$	4.7429	12.9458	21.221	0.205	4.786			2.220	—
$\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$	ZnO	3.2502	5.2104	27.274	0.114	93.772	9.795	0.986	1.344	1.0185
	$\text{Al}_2\text{O}_3$	4.7435	12.9481	20.844	0.196	6.228			2.302	—
$\text{Zn}_{0.94}\text{Al}_{0.03}\text{Er}_{0.03}\text{O}$	ZnO	3.2498	5.2089	28.970	0.109	91.285	7.855	1.012	1.191	1.0188
	$\text{Al}_2\text{O}_3$	4.7438	12.9502	22.783	0.127	8.715			1.926	—

<sup>a</sup>  $\langle\sigma^2\rangle^{1/2}$ : root mean square microstrain (r.m.s.);  $d$ : crystallite size.

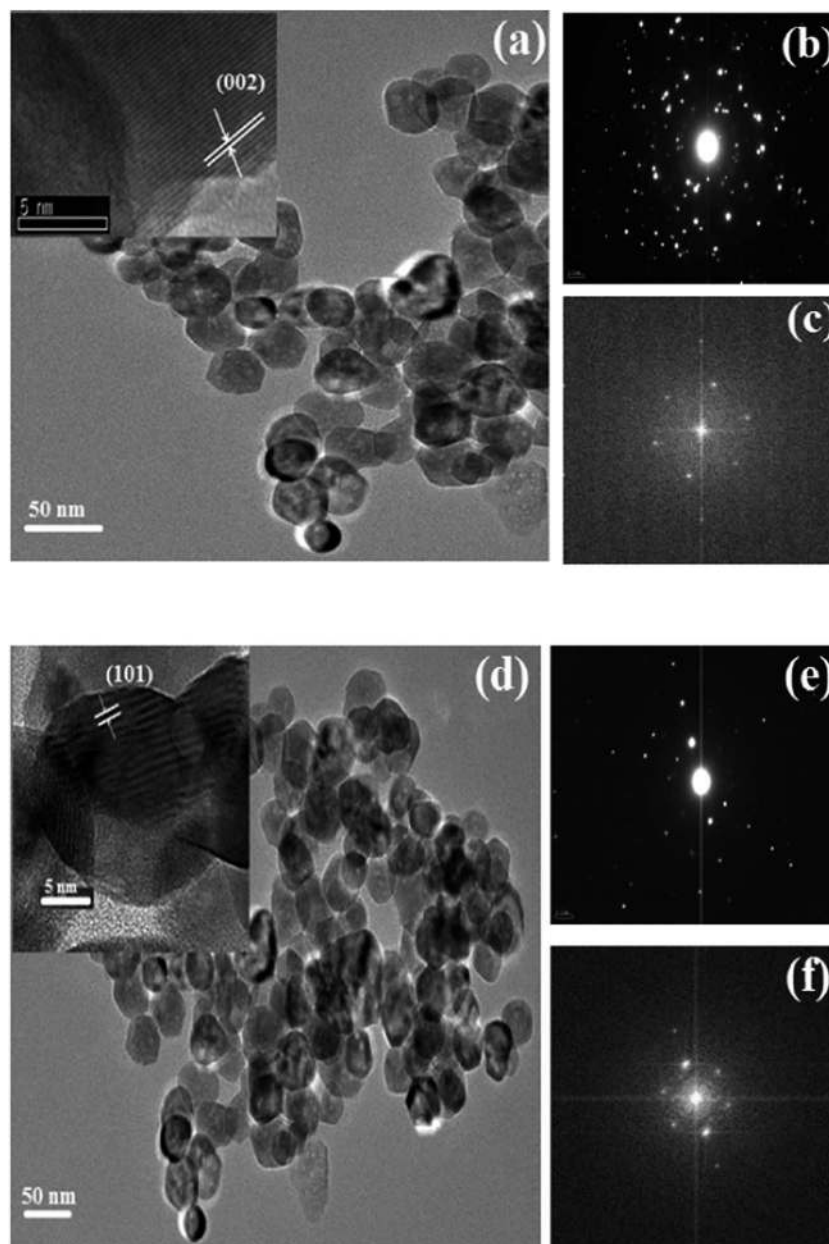


Fig. 4 TEM micrographs of un-doped ZnO (a–c), and of the  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$  nanoparticles (d–f).



obtained values are higher than the pure ZnO sample. This result may be explained by the diffusion of Er ions in Al<sub>2</sub>O<sub>3</sub> phase. In fact the lattice parameters of this later, increase with increasing Er content from 2 to 3% as indicated in Table 2.

The change in ZnO crystal parameters indicates the distortion or compression of the crystal lattice when Al<sup>3+</sup> and Er<sup>3+</sup> ions enter into the ZnO lattice. The crystal lattice distortion degree is calculated by following relation:<sup>27</sup>

$$R = \frac{2a(2/3)^{1/2}}{c} \quad (2)$$

where  $a$  and  $c$  are calculated lattice constants of ZnO phase from the XRD data. The calculated values are listed in Table 2.

### 3.2. Morphological observations

The crystallinity, morphology, and size of un-doped and co-doped ZnO nanoparticles were studied using transmission electron microscopy (TEM), as shown in Fig. 4a and d. The ZnO nanoparticles are not monodisperse and the nanoparticles agglomerate each other, which is very common for ZnO nanoparticles using hydrothermal method<sup>28</sup> in consistence with XRD analysis.

From HRTEM images (inset in Fig. 4a and d), the calculated interplanar  $d$ -spacing matches closely with the (002) (0.260 nm) and (101) (0.281 nm) plane  $d$ -spacing of hexagonal ZnO.<sup>29</sup> In Fig. 4b and e, the selected area electron diffraction (SAED) patterns are presented for the un-doped ZnO and Zn<sub>0.95</sub>Al<sub>0.03</sub>Er<sub>0.02</sub>O samples, respectively, and the crystalline nature and preferential orientation of the patterns are found to match well with the würtzite structure of the XRD results. The Fast Fourier Transformation (FFT) images shown in Fig. 4c and f shows that it is a typical hexagonal structure consistent with würtzite ZnO.

The TEM images are used to prepare the histogram distribution of the particle size for the un-doped and co-doped samples. As observed Fig. 5, the histogram distribution is well modeled by a lognormal distribution. The mean particle size that exhibit hexagonal shape, estimated from the TEM micrograph using standard software (IMAGE J), are found to be 28 and 36 nm for pure ZnO and Zn<sub>0.95</sub>Al<sub>0.03</sub>Er<sub>0.02</sub>O, respectively, which is in good agreement with the mean crystalline size value determined from the XRD data analysis. The particles with such small size have very large specific surface area so that strong local aggregation can be found in the images. Few bigger particles are also seen in the micrograph, which may be due to the aggregation or overlapping of smaller particles.

Fig. 6 depicts the energy-dispersive spectra (EDX) of pure ZnO and 2% Er-doped ZnO. The spectra reveal the existence of Zn, O, Al, Er and Cu elements in co-doped ZnO nanoparticles. Among them, Cu (8.17 keV) signal come from supporting Cu grid. The EDX results (insert table) show that the chemical composition vary slightly, almost close to the nominal composition. We note that for small concentrations, EDX is not very precise and the detection limit is dependent on the matrix and its surface. Furthermore, the greatest source of error – or at least uncertainty – in qualitative analysis can be found for constituents present at minor ( $1 < C < 10$  wt%) and trace ( $C < 1$  wt%) levels<sup>30</sup> which is the case of Er and Al elements.

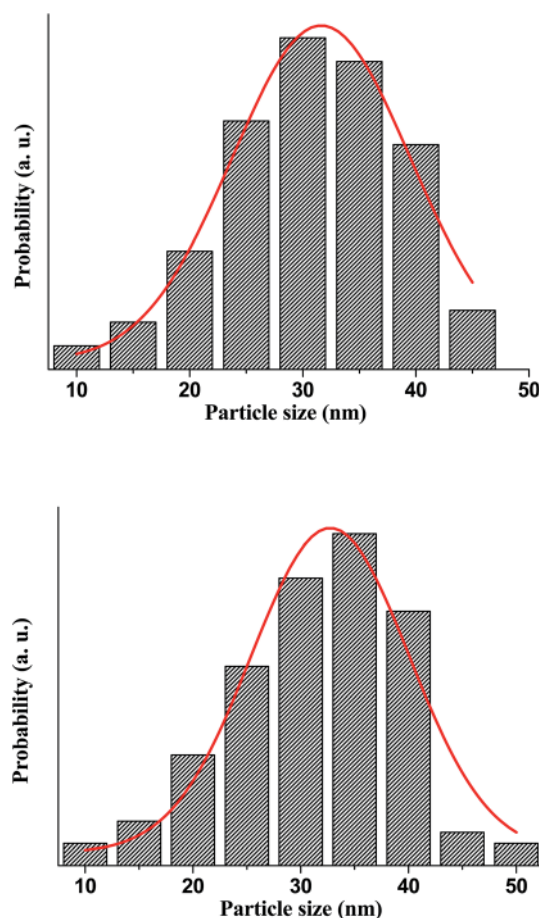


Fig. 5 The histogram distribution of the particles size, where the solid line represents the lognormal function.

### 3.3. Magnetic properties

Fig. 7 shows the magnetization curve measured at room temperature for un-doped and doped ZnO NPs. When ZnO is un-doped, there is no hysteresis loop and the sample is paramagnetic at room temperature. Since bulk ZnO is diamagnetic, the origin of this behavior at nanoscale may be usually linked to defects.<sup>31</sup> Extensive researches have indicated that defects like Zn and O vacancies ( $V_{Zn}$ ,  $V_O$ ), oxygen interstitial ( $O_i$ ) and cation interstitial (*i.e.*  $Zn_i$ ) in wide gap oxide semiconductors are possibly the origin of the magnetism.<sup>30</sup> We note that the oxygen deficiency exhibits a local magnetic moment.<sup>31,32</sup> Therefore, it was reported that ferromagnetism can appear in un-doped ZnO thin films when Zn vacancies are introduced.<sup>33</sup> The magnetization of very thin ZnO films was found to be much larger than that of the thicker films, suggesting that the vacancies must be located mostly at the surface. The origin of this magnetism has been shown theoretically<sup>34</sup> to be due to the unpaired 2p electrons at oxygen sites around the Zn vacancy. In addition, Zn vacancies prefer to occupy the surface sites, and the formation energy of Zn vacancy is higher than that of the oxygen vacancies.

For Zn<sub>0.96</sub>Al<sub>0.03</sub>Er<sub>0.01</sub>O sample, the data show the existence of two components, one is diamagnetic and the other is



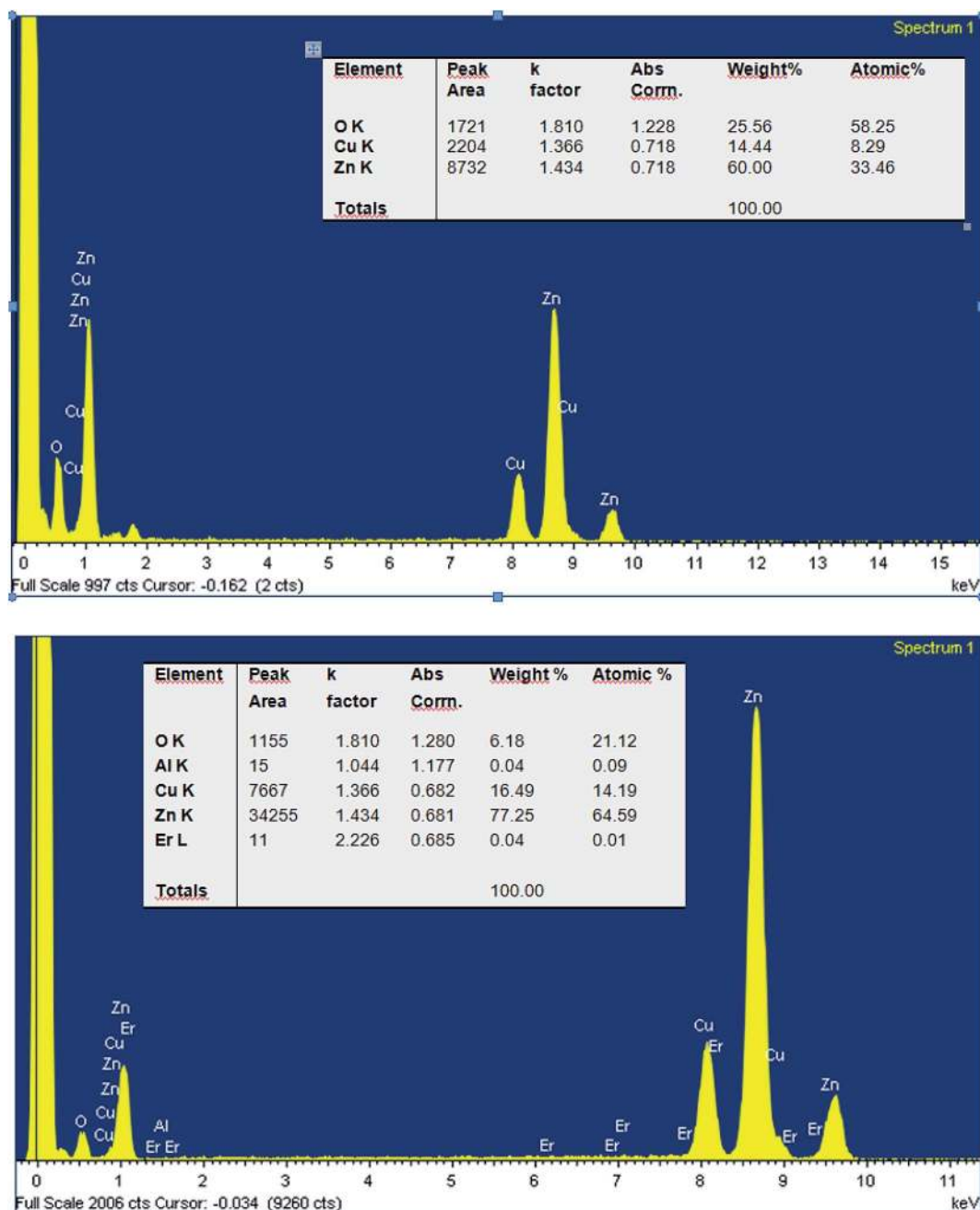


Fig. 6 EDX spectra of the pure ZnO (top) and Zn<sub>0.95</sub>Al<sub>0.03</sub>Er<sub>0.02</sub>O nanoparticles (bottom).

ferromagnetic. After removing the diamagnetic contribution, the magnetization curve exhibits ferromagnetic behavior with a coercive field of 53.51 Oe (Table 3). A linear trend is determined in the high magnetic field region for Zn<sub>0.95</sub>Al<sub>0.03</sub>Er<sub>0.02</sub>O sample which has been assigned to the paramagnetic behavior. However, in the low-field region, a ferromagnetic contribution is determined. After subtracting the paramagnetic contribution, a coercive field ( $H_c$ ) of 17.32 Oe has been determined. The ferromagnetic contribution becomes negligible as the Er content is increased to 3% as can be observed in Fig. 7 where the behavior is purely paramagnetic. A similar behavior has been determined for Er-doped SnO<sub>2</sub> nanoparticles.<sup>35</sup>

There are still debates on the origin of the ferromagnetism in non magnetic elements-doped oxide semiconductors. Similar to that of transition metal doped oxide semiconductors, both element doping, defects and secondary phase are believed to contribute to the ferromagnetism dependent on the materials and fabrication parameters.<sup>28</sup> Therefore, in Eu-doped ZnO thin films, the ferromagnetism was believed to originate from both the native and the substitutional defects through bounded magnetic polaron (BMP) model.<sup>36</sup> The overlapping of BMP can lead to the ferromagnetism, which is formed by the coupling between the localised carriers from the defects and magnetic dopants within the same region. BMP theory is proposed to



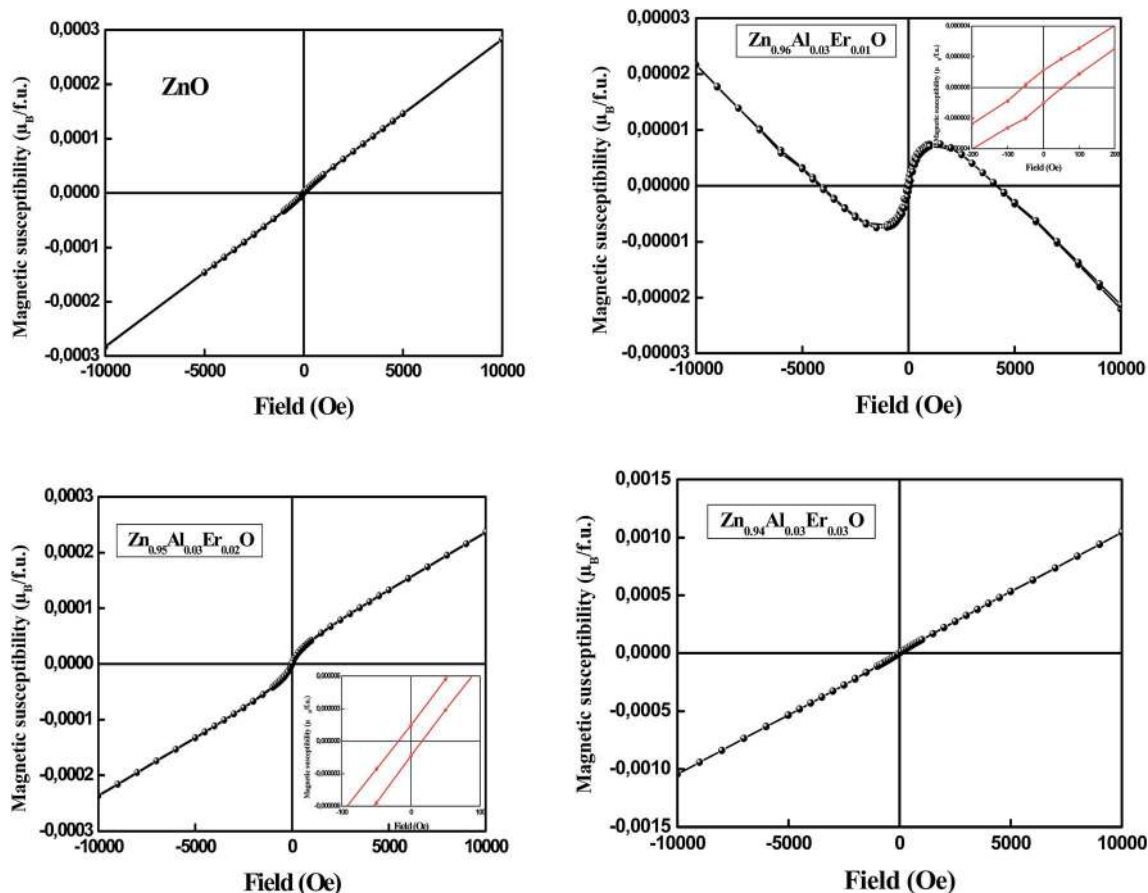


Fig. 7  $M$ - $H$  of magnetic susceptibility of pure ZnO and (Al,Er) co-doped ZnO NPs.

explain the ferromagnetism in Er-doped ZnO thin films as well.<sup>37,38</sup> The BMP in Er-ZnO is formed by the magnetic exchange interaction between Zn or O vacancies and Er<sup>3+</sup> ions, which leads to the alignment of Er<sup>3+</sup> spins. Moreover, the synthesis method or technique seems to play a major role in the magnetism of these ODS systems. ZnO:Er films synthesized by magnetron sputtering technique are reported to exhibit a coexistence of ferromagnetic and paramagnetic phases at room temperature.<sup>37</sup> According to the authors, the strongest ferromagnetic signal has been observed for the 4.0% Er doped sample and that ferromagnetic contribution decreases with the increase of the Er content.

The increase of saturation magnetic susceptibility between 1 and 2% Er can be explained by the effect of Er ion since the

magnetic moment attributing paramagnetism only to Er was observed to be 3.612  $\mu_B$  per f.u. for 2% Er (Table 3). On the other hand, the origin of room-temperature magnetism could be derived from secondary phase. In fact, Li *et al.* have confirmed recently by the first-principle calculations, that Al-vacancy can introduce magnetic moment for 3  $\mu_B$  in Al<sub>2</sub>O<sub>3</sub> crystal and form stable V<sub>Al</sub>-V<sub>Al</sub> ferromagnetic coupling at room temperature.<sup>39,40</sup> Ma *et al.* observed RT FM in Al doped ZnO after vacuum annealing and ascribed it to the charge transfer between Al and ZnO, which can cause the electronic structure alternation of Zn and Al.<sup>41</sup> so, when analyzing the magnetic properties of these nano-oxides, it is particularly important to ensure that there is no contamination due to the use of iron based tools in any step of the sample handling.<sup>42,43</sup>

Table 3 Magnetic parameters of (Al,Er) co-doped ZnO NPs

Sample molecular formula	Molecular weight	Mass (mg)	Mass susceptibility (cgs)	Mole susceptibility $\chi_m$	Effective moment $P_{\text{eff}}$ ( $\mu_B$ per f.u.)	Moment of Er ( $\mu_B$ per f.u.) attributing paramagnetism only to Er	Saturation magnetic susceptibility ( $\mu_B$ per f.u.)
ZnO	81.3894	34.29	$+1.89 \times 10^{-6}$	$1.53826 \times 10^{-4}$	0.60579	N/A	
Zn <sub>0.96</sub> Al <sub>0.03</sub> Er <sub>0.01</sub> O	81.2558	24.83	$-0.264 \times 10^{-6}$	$-2.14515 \times 10^{-5}$	—	N/A	$1.72 \times 10^{-5}$
Zn <sub>0.95</sub> Al <sub>0.03</sub> Er <sub>0.02</sub> O	82.2745	22.35	$+1.33 \times 10^{-6}$	$1.09425 \times 10^{-4}$	0.51094	3.612	$6.67 \times 10^{-4}$
Zn <sub>0.94</sub> Al <sub>0.03</sub> Er <sub>0.03</sub> O	83.2932	38.18	$+7.07 \times 10^{-6}$	$5.88883 \times 10^{-4}$	1.18529	6.843	$3.32 \times 10^{-4}$





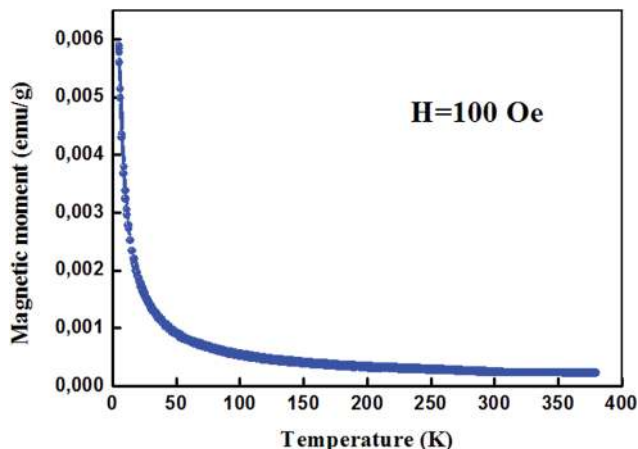


Fig. 8 Magnetization (magnetic moment per unit mass) for  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$  NPs.

Increasing Er content to 3% lead to a drop in estimated magnetic susceptibility to about  $3.32 \times 10^{-4} \mu_{\text{B}}$  per f.u. (Table 3) as compared to 2% Er doped sample ( $6.67 \times 10^{-4} \mu_{\text{B}}$  per f.u.), could be due to an increasing occurrence of antiferromagnetic coupling or super exchange interaction between neighboring ions at shorter separation distances, which was similar to the doped oxide semiconductor films like Cu-doped ZnO and Er-doped ZnO.<sup>38,44,45</sup> Similarly,  $\text{SnO}_2:\text{Er}$  nanoparticles prepared by sol-gel method show only the occurrence of a paramagnetism with the presence of a weak Er–Er antiferromagnetic interactions, which become stronger as the Er concentration increases.<sup>46</sup> In fact, a strong paramagnetic-like behavior coexisting with a ferromagnetic order at room temperature has been determined for samples with Er content up to 5.0%. Above this concentration, only a paramagnetic behavior has been determined.

For proper investigation of the magnetic properties of  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$  nanoparticles, the temperature dependence of the magnetic moment has been estimated at a constant magnetic field of 100 Oe. The behavior of the curve shown in Fig. 8 is a Curie type paramagnet, which is a type of magnetism resulting from the presence of atoms with unpaired electrons. Curie-type paramagnetism has a particular temperature dependence:

$$\chi_{\text{m}}(T) = C_{\text{m}}/T \quad (3)$$

where,  $\chi_{\text{m}}$  is molar susceptibility in (emu/mass/molecular weight/field);  $C_{\text{m}}$  is the molar Curie constant in (emu K mol<sup>-1</sup>). The plot of the inverse magnetic susceptibility *versus* temperature is very useful for characterizing Curie paramagnets (Fig. 9). The curve is fitted with a Curie–Weiss linear relation for the high temperature region. The slope of the curve is  $1/C_{\text{m}}$  and the Curie constant is given as:

$$C_{\text{m}} = bP_{\text{eff}}^2N \quad (4)$$

where  $P_{\text{eff}}$  is known as the effective magnetic moment,  $b$  is the universal constant and  $N$  is the concentration of magnetic

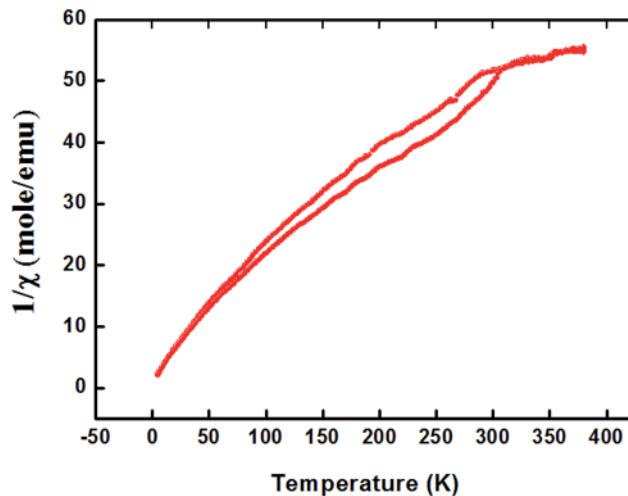


Fig. 9 The inverse susceptibility as a function of temperature for  $\text{Zn}_{0.95}\text{Al}_{0.03}\text{Er}_{0.02}\text{O}$  sample.

atoms with that moment. The calculated values are listed in Table 3. Extrapolation of the inverse susceptibility data in the high temperature region gave negative Curie–Weiss temperatures of  $-73.46$  K. The negative value of temperature indicates the antiferromagnetic interactions between doping ions. Such antiferromagnetic interactions were also reported earlier by several researchers in Co and Mn doped ZnO systems.<sup>47</sup> The calculated effective magnetic moment of co-doped ZnO NPs are listed in Table 3. The results show that the magnetic moment of Er is smaller indicating that not all Er atoms contribute to the ferromagnetic behaviour.

### 3.4. First principle calculations

**3.4.1 Methodology.** The total energies and electronic structures of Er and Al doped ZnO systems were calculated by Vienna *Ab initio* Simulation Package (VASP) code using the generalized gradient approximation of Perdew and Wang<sup>48–50</sup> and the project augmented wave (PAW) method.<sup>51</sup> A cutoff energy of 500 eV and  $k$ -mesh of  $3 \times 3 \times 2$  were used for the relaxations of the supercells with 108 atoms. The criteria for self-consistent calculations of energies and forces are less than 0.01 meV and 0.01 eV Å<sup>-1</sup>, respectively. The optimized lattice parameters of wurtzite ZnO are  $a = 0.3241$  nm and  $c = 0.5199$  nm, which are well consistent with the experimental values ( $a = 0.325$  nm and  $c = 0.521$  nm),<sup>52</sup> and other theoretical calculations.<sup>53</sup>

The Er and Al co-dopings were simulated by substituting four Zn atoms in the supercell with five different configurations. The concentration of dopants is 3.704 at%. The concentration of dopants is hard to decrease to close to the experimental value due to huge computation consume. The ferromagnetic (FM) and antiferromagnetic (AFM) orderings of the doped systems are simulated by altering the magnetic moment of the Er. The GGA + U (ref. 54) method was used to describe the exchange and correlation potential. The on-site Coulomb interaction  $U_{\text{eff}} = 5.5$  eV and 5.0 eV is employed to describe the correlation effects in localized f orbital of Er and d orbital of Zn.



**3.4.2 Formation energy and phase stability.** Formation energy of the Er and Al doped ZnO is evaluated *via* the definition below,

$$E^f = \frac{1}{N} \left\{ \left[ E_{\text{ZnO}}^{\text{doped}} - E_{\text{ZnO}}^{\text{undoped}} \right] + \left[ N E_{\text{Zn}} - N_{\text{Er}} E_{\text{Er}} - N_{\text{Al}} E_{\text{Al}} \right] \right\} \quad (5)$$

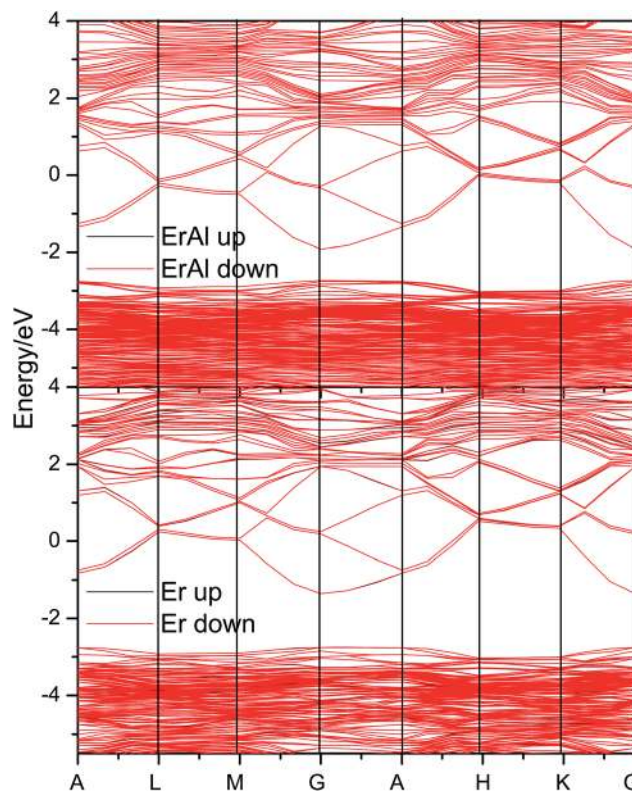
where  $E_X$  is the total energy of the X system under ground states. The  $N$ ,  $N_{\text{Er}}$  and  $N_{\text{Al}}$  are the numbers of Zn and dopants Er and Al atoms, respectively. The magnetic ordering of doped systems can be evaluated by comparing the formation energies of FM and AFM configurations.

Five different configurations of Er and Al co-doped ZnO systems with different distance between Er and Al are considered as shown in our previous study.<sup>55</sup> The distance between Er and Al changes from 3.2 to 9.2 Å in these simulated cells. Table 4 shows the formation energies of Er and Al co-doped systems with different configurations 'a' to 'e'. All co-doped configurations show negative formations energies indicating the Er and Al can be thermodynamically doped into ZnO. The AFM and FM magnetic orderings almost own same formation energies indicating the similar stabilities. Therefore, the total magnetic moments of these doped systems are almost the zero. For the Er doped systems, four different configurations are considered (a is equivalent to b configuration). The Er-d system is the most stable configuration among the studied systems. Its magnetic moment is 0.034  $\mu_B$  per cell, and therefore it shows weak FM property. However, there are some discrepancies with our experimental findings, which may be caused by the influence of defects and the segregations of Er element and/or the secondary phase.

In order to study the influence of dopants on the electronic structures and magnetic properties, the band structures of the most stable Er and Al co-doped, and Er doped systems are evaluated and shown in Fig. 10. These two doped systems show similar band structure, except the different band gap. The Al pulls about 0.6 eV closer the conduction band to the valence band comparing the Er solely doped system. Therefore, the further doping of Al can adjust the band gap in Er doped systems. Furthermore, the Er and Al co-doped systems almost have no spin-polarization, *i.e.*, the spin-up and spin-down band overlaps well. There are small difference between spin-up and spin-down band, which causing small magnetic moments.

**Table 4** The formation energies (eV) of Er and Al co-doped, and Er doped ZnO systems with different configurations 'a' to 'e'

Doped systems	AFM		FM	
	ErAl	Er	ErAl	Er
a	-2.5306	-3.170	-2.5305	-3.170
b	-2.5738	—	-2.5740	—
c	-2.6526	-3.266	-2.6527	-3.266
d	-2.6644	-3.388	-2.6643	-3.388
e	-2.7004	-3.379	-2.7004	-3.378



**Fig. 10** The electronic band structures of the most stable Er and Al co-doped, and Er doped ZnO systems.

## 4. Conclusion

In this study un-doped and (Al,Er) co-doped ZnO NPs have been successfully synthesized by hydrothermal method. The XRD results show that all doped samples have hexagonal structure with the preferential orientation along the (101) and low amount of  $\alpha$ - $\text{Al}_2\text{O}_3$  secondary phase. The ZnO crystallites are in the nanometer range. The TEM results display that the incorporation of Al and Er ions in crystal structure of ZnO can't change the morphology of ZnO NPs. The magnetic properties are found to change with varying dopant content. The *ab initio* calculations reveals that Er and Al co-doped systems are thermodynamically stable. The Al can adjust the band gap of the Er doped system and hardly change the distributions of electronic structures.

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