## STRUCTURE FOR NONNEGATIVE SQUARE ROOTS OF UNBOUNDED NONNEGATIVE SELFADJOINT OPERATORS

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Abstract. It is well known that, for an unbounded nonnegative selfadjoint operator A on a Hilbert space, there is a unique nonnegative square root  $A^{1/2}$ , which is frequently associated with the structural damping in many practical vibration systems. In this paper we develop a general theory for the structure of  $A^{1/2}$ , which includes the expression of  $A^{1/2}$  and a program to find the domain of  $A^{1/2}$  explicitly from the domain of A. The relationship between  $A^{1/2}$  and related differential operators is determined for the selfadjoint differential operator A. Finally, the theoretical results given in this paper are applied to fourth-order "beam" operators and n-dimensional "wave" operators with sufficient complexity for applications to elastic vibration systems.

1. Introduction. Let H be a Hilbert space with the inner product  $(\cdot, \cdot)$  and the induced norm  $\|\cdot\|$ . Let A be an unbounded nonnegative selfadjoint operator on H and  $\mathcal{L}$  a closed linear operator on H. In the last two decades, great attention has been focused on the following elastic system:

$$\begin{cases} \ddot{y}(t) + \mathcal{L}\dot{y}(t) + Ay(t) = 0, \\ y(0) = y_0, \ \dot{y}(0) = y_1, \end{cases}$$
(1.1)

where a dot denotes  $\frac{d}{dt}$ , and  $y, y_0, y_1 \in H$ . The usual procedure for dealing with the system (1.1) is as follows.

Letting  $x_1 = A^{1/2}y, x_2 = y$ , the system (1.1) can be transformed into an equivalent first-order evolution system

$$\begin{cases} \frac{d}{dt}\tilde{x}(t) = \mathcal{A}\tilde{x}(t), \\ \tilde{x}(0) = Y_0, \end{cases}$$
(1.2)

where

$$ilde{x}(t) = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}, \qquad \mathcal{A} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -\mathcal{L} \end{pmatrix}, \qquad Y_0 = \begin{pmatrix} A^{1/2}y_0 \\ y_1 \end{pmatrix},$$

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and  $A^{1/2}$  is the nonnegative square root of A. The domain of  $\mathcal{A}$  is

$$D(A^{1/2}) \times D(A^{1/2})$$

if  $D(\mathcal{L}) \subset D(A^{1/2})$ , where  $D(\mathcal{L})$  is the domain of  $\mathcal{L}$ . For many real systems, however, the structure of  $D(A^{1/2})$  is not clear except that D(A) has special boundary conditions.

On the other hand,  $A^{1/2}\dot{y}$  corresponds to structural damping for elastic vibration systems in [1] and is related to frequency proportional damping in [4]. Thus, it is necessary to understand the relationship between  $A^{1/2}$  and differential operators in order that  $A^{1/2}\dot{y}$ admits the proper physical interpretation if A is a differential operator. D. L. Russell in [2] gave the relationship between  $A^{1/2}$  and differential operators for fourth-order "beam" operators with the "symmetric" boundary conditions but the structure of  $D(A^{1/2})$  is not dealt with. A. V. Balakrishnan in [3], [4] obtained  $D(A^{1/2})$  and the relationship between  $A^{1/2}$  and differential operators for two models of bending of uniform Bernoulli beams by the formula introduced in [5].

In this paper, we develop a general theory for the structure of the nonnegative square root of any unbounded nonnegative selfadjoint operator A. In Sec. 2, a program is given to find the domain of  $A^{1/2}$  and the expression of  $A^{1/2}$  is studied. Finally, in Secs. 3 and 4, we apply the theoretical results developed in the paper to fourth-order elastic "beam" operators and *n*-dimensional "wave" operators, respectively.

2. The main results. Throughout this paper we make the following assumptions.

Let H be a complex Hilbert space with the inner product  $(\cdot, \cdot)$  and the induced norm  $\|\cdot\|$ , and let A be an unbounded nonnegative selfadjoint operator on H with compact resolvent and domain D(A); let  $A^{1/2}$  be the nonnegative square root of A. Note that A is nonnegative if and only if

$$(Ax, x) \ge 0, \quad \forall x \in D(A).$$

The operator A is said to be positive if

$$(Ax, x) > 0, \quad \forall x \in D(A), \ x \neq 0.$$

Let  $H_1$  be another complex Hilbert space with the inner product  $(\cdot, \cdot)_1$  and the induced norm  $\|\cdot\|_1$  and  $B: D(B) \subset H \to H_1$  be a closed linear operator such that

$$D(B) \supset D(A); \tag{2.1}$$

$$(Ax, x) \ge \|Bx\|_1^2, \quad \forall x \in D(A).$$

$$(2.2)$$

Set

$$[x, y] = (x, Ay) - (Bx, By)_1, \quad \forall x \in D(B), \ \forall y \in D(A),$$
$$M_0 = \{x \mid x \in D(A), (x, Ax) = \|Bx\|_1^2\}.$$

We introduce several definitions in preparation for development of the structure of  $A^{1/2}$ . It should be noted that, throughout this paper, the definitions, given specially, all belong to the authors.

DEFINITION 2.1.  $(B, H_1)$  is said to be a pseudo-square-root of A if  $\overline{M}_0 = H$ , where  $\overline{M}_0$  is the closure of  $M_0$  in H.

REMARK 2.1. Any nonnegative selfadjoint operator A has at least one pseudo-squareroot. For example,  $(A^{1/2}, H)$  is a pseudo-square-root of A, where  $B = A^{1/2}$  and  $H_1 = H$ .

EXAMPLE 2.1. Let  $\Omega$  be a bounded domain of  $\mathbb{R}^n$  with smooth boundary  $\partial \Omega$  of  $\mathbb{C}^2$ and  $H = L^2(\Omega)$ . Consider the following Laplace operator:

$$D(A) = \left\{ u \mid u \in H^2(\Omega), \left( \frac{\partial u}{\partial n} + \alpha u \right) \Big|_{\partial \Omega} = 0 \right\},$$
  
 $Au = -\Delta u.$ 

where  $\alpha \ge 0$ ,  $\frac{\partial}{\partial n}$  is the normal derivative, and  $\Delta u = \sum_{j=1}^{n} \frac{\partial^2 u}{\partial x_j^2}$ . It is well known that A is an unbounded nonnegative selfadjoint operator on  $L^2(\Omega)$ .

Set  $H_1 = (L^2(\Omega))^n$ ;  $H_1$  is a product Hilbert space with the inner product

$$\langle \tilde{u}, \tilde{v} \rangle = \sum_{j=1}^{n} (u_j, v_j)_{L^2(\Omega)}, \quad \forall \tilde{u} = (u_1, u_2, \dots, u_n)^{\tau}, \ \tilde{v} = (v_1, v_2, \dots, v_n)^{\tau} \in (L^2(\Omega))^n.$$

Define  $B: L^2(\Omega) \to (L^2(\Omega))^n$  by

$$D(B) = H^{1}(\Omega), \qquad Bu = \nabla u = \left(\frac{\partial u}{\partial x_{1}}, \dots, \frac{\partial u}{\partial x_{n}}\right)^{\tau}$$

Conditions (2.1) and (2.2) can easily be verified by using Green's formula, and

$$H_0(\Omega) = \left\{ u \; \Big| \; u \in H^2(\Omega), u |_{\partial \Omega} = rac{\partial u}{\partial n} \Big|_{\partial \Omega} = 0 
ight\} \subset M_0.$$

Hence  $(\nabla, (L^2(\Omega))^n)$  is a pseudo-square-root of A.  $\Box$ 

DEFINITION 2.2. Let  $(B, H_1)$  be a pseudo-square-root of A. Let  $H_{\partial}^+$  be a Hilbert space with the inner product  $(\cdot, \cdot)_{\partial}^+$  and the induced norm  $\|\cdot\|_{\partial}^+$ . If there exist a bounded positive selfadjoint operator  $\Upsilon: H_{\partial}^+ \to H_{\partial}^+$  and a mapping  $\Gamma: D(B) \to H_{\partial}^+$  such that

$$[x,y] = (\Upsilon \Gamma x, \Gamma y)_{\partial}^{+}, \quad \forall x, y \in D(A),$$
(2.3)

then we say that  $(H^+_{\partial}, \Upsilon, \Gamma)$  is a positive boundary space of A corresponding to B.

EXAMPLE 2.2. Let  $H, A, H_1$ , and B be given by Example 2.1. By Green's formula, we have

$$[u, v] = -\int_{\partial\Omega} u \frac{\partial \overline{v}}{\partial n} d\sigma$$
  
=  $\alpha \int_{\partial\Omega} u \overline{v} \, d\sigma, \quad \forall u \in H^1(\Omega), \ \forall v \in D(A).$  (2.4)

Set  $H^+_{\partial} = L^2(\partial\Omega)$  and  $\Upsilon = \alpha I$ , where I is the identity mapping on  $L^2(\partial\Omega)$ . Define  $\Gamma: H^1(\Omega) \to L^2(\partial\Omega)$  by

$$\Gamma u = u|_{\partial\Omega}, \quad \forall u \in H^1(\Omega).$$

It is easy to verify from (2.4) that  $(L^2(\partial\Omega), \alpha I, \Gamma)$  is a positive boundary space of A corresponding to B.  $\Box$ 

The following proposition follows from a direct check.

**PROPOSITION 2.1.** Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Then

$$|[x,y]| \le 2||A^{1/2}x|| \, ||A^{1/2}x|| \, ||A^{1/2}y||, \quad \forall x,y \in D(A^{1/2});$$
(2.5)

$$|(\Upsilon\Gamma x, \Gamma y)_{\partial}^{+}| \le 2||A^{1/2}x|| \, ||A^{1/2}y||, \quad \forall x, y \in D(A).$$
(2.6)

It is well known that the spectrum of A consists of eigenvalues:

 $0 \leq \lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_n \leq \lambda_{n+1} \leq \cdots$ 

with  $\lim_{n\to\infty} \lambda_n = +\infty$ , where the multiple eigenvalues are listed according to their algebraic multiplicities. There exists an orthonormal basis  $\{x_n\}$  of H such that

$$Ax_n = \lambda_n x_n, \quad n = 1, 2, \dots$$

Let k be the nonnegative integer such that

$$\lambda_j = 0 \quad ext{if } j \leq k; \qquad \lambda_j > 0 \quad ext{if } j \geq k+1.$$

If A is positive, then k = 0.

PROPOSITION 2.2. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Set

$$g_n = \frac{1}{\sqrt{\lambda_n}} \Gamma x_n, \qquad \psi_n = \frac{1}{\sqrt{\lambda_n}} B x_n, \quad n = k+1, k+2, \dots$$
 (2.7)

Then

$$\left\{ \left( \begin{array}{c} \Upsilon^{1/2} g_n \\ \psi_n \end{array} \right) \ \middle| \ n \ge k+1 \right\}$$

is an orthonormal set in the product Hilbert space  $H^+_{\partial} \times H_1$ , where the inner product of  $H^+_{\partial} \times H_1$  is defined by

$$\left\langle \begin{pmatrix} f\\\psi \end{pmatrix}, \begin{pmatrix} g\\\phi \end{pmatrix} \right\rangle = (f,g) + (\psi,\phi)_1, \quad \forall f,g \in H^+_\partial, \; \forall \psi, \phi \in H_1.$$

*Proof.* For any  $n, m \ge k + 1$ , we have

$$\begin{aligned} (x_n, x_m) &= \frac{1}{\lambda_m} (x_n, Ax_m) = \frac{1}{\lambda_m} [x_n, x_m] + \frac{1}{\lambda_m} (Bx_n, Bx_m)_1 \\ &= \sqrt{\frac{\lambda_n}{\lambda_m}} [(\Upsilon g_n, g_m)_{\partial}^+ + (\psi_n, \psi_m)_1] \\ &= \sqrt{\frac{\lambda_n}{\lambda_m}} \left\langle \begin{pmatrix} \Upsilon^{1/2} g_n \\ \psi_n \end{pmatrix}, \begin{pmatrix} \Upsilon^{1/2} g_m \\ \psi_m \end{pmatrix} \right\rangle. \end{aligned}$$

The orthonormality of  $\{x_n\}$  in H yields the orthonormality of

$$\left\{ \left( \begin{array}{c} \Upsilon^{1/2} g_n \\ \psi_n \end{array} \right) \middle| n \ge k+1 \right\}$$

in  $H^+_{\partial} \times H_1$ .  $\square$ 

First, we consider the structure of  $D(A^{1/2})$ .

THEOREM 2.1. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Then

$$D(A^{1/2}) = \{x \mid x \in D(B), [x, y] = (\Upsilon \Gamma x, \Gamma y)_{\partial}^+, \forall y \in D(A)\}.$$

*Proof.* Let  $x \in D(A^{1/2})$ . Then there are  $\{z_n\}$  in D(A) such that in H

$$z_n \to x$$
 and  $A^{1/2}z_n \to A^{1/2}x$ , as  $n \to \infty$ .

By Proposition 2.1 and (2.2), we have that  $x \in D(B)$  and that

$$\begin{split} [x,y] &= \lim_{n \to \infty} [z_n,y] = \lim_{n \to \infty} (\Upsilon \Gamma z_n, \Gamma y)_{\partial}^+ \\ &= (\Upsilon \Gamma x, y)_{\partial}^+, \quad \forall y \in D(A). \end{split}$$

Conversely, suppose  $x \in D(B)$  such that

$$[x,y] = (\Upsilon \Gamma x, y)_{\partial}^{+}, \quad \forall y \in D(A).$$
(2.8)

Set

$$x = \sum_{j=1}^{\infty} \alpha_n x_n$$

Then

$$\begin{aligned} \alpha_n &= (x, x_n) = \frac{1}{\lambda_n} (x, A x_n) \\ &= \frac{1}{\lambda_n} [x, x_n] + \frac{1}{\lambda_n} (B x, B x_n)_1 \\ &= \frac{1}{\sqrt{\lambda_n}} \left[ x, \frac{1}{\sqrt{\lambda_n}} x_n \right] + \frac{1}{\sqrt{\lambda_n}} (B x, \psi_n)_1, \quad \forall n \ge k+1, \end{aligned}$$

or equivalently

$$\begin{split} \sqrt{\lambda_n} \alpha_n &= \left[ x, \frac{1}{\sqrt{\lambda_n}} x_n \right] + (Bx, \psi_n)_1 \\ &= (\Upsilon^{1/2} \Gamma x, \Upsilon^{1/2} g_n)_{\partial}^+ + (Bx, \psi_n)_1 \\ &= \left\langle \begin{pmatrix} \Upsilon^{1/2} \Gamma x \\ Bx \end{pmatrix}, \begin{pmatrix} \Upsilon^{1/2} g_n \\ \psi_n \end{pmatrix} \right\rangle, \quad \forall n \ge k+1. \end{split}$$
(2.9)

Thus Proposition 2.2 yields

$$\sum_{n=k+1}^{\infty} |\alpha_n|^2 \lambda_n \leq \left\| \left( \frac{\Upsilon^{1/2} \Gamma x}{Bx} \right) \right\|_{H_{\partial}^+ \times H_1}^2 < +\infty,$$

so that  $x \in D(A^{1/2})$ .  $\Box$ 

When  $M_0 = D(A)$ , we can choose  $H_{\partial}^+ = \{0\}$ . Thus we have

COROLLARY 2.1. Let  $(B, H_1)$  be a pseudo-square-root of A. If  $M_0 = D(A)$ , then

$$D(A^{1/2}) = \{x \mid x \in D(B), [x, y] = 0, \ \forall y \in M_0\}.$$

EXAMPLE 2.3. Let A be a "string vibration" operator on  $L^2(0,1)$ , i.e.,  $Au = -u''(\cdot)$ ,  $u \in D(A)$ , such that

$$H_0^2(0,1) \subset D(A).$$
 (2.10)

Set  $H_1 = L^2(0,1)$ . The closed linear operator  $B: L^2(0,1) \to L^2(0,1)$  is defined by

$$D(B) = H^1(0, 1), \qquad Bu = iDu = iu'(\cdot).$$

By (2.10) and integration by parts it is easily checked that  $(iD, L^2(0, 1))$  is a pseudo-square-root of A.

(i) If A has the following boundary conditions,

$$D(A) = \{ u \mid u \in H^2(0,1), u(0) = u(1) = 0 \},$$
(2.11)

it is easily verified that  $M_0 = D(A)$ . For any  $u \in H^1(0,1)$ , it is clear that u satisfies

$$[u, v] = u(0)\overline{v'(0)} - u(1)\overline{v'(1)} = 0, \quad \forall v \in D(A)$$

if and only if u(0) = u(1) = 0. By Corollary 2.1, we have

$$D(A^{1/2}) = \{ u \mid u \in H^1(0,1), u(0) = u(1) = 0 \} = H^1_0(0,1)$$

(ii) If A has the boundary conditions

$$D(A) = \{ u \mid u \in H^2(0,1), u'(0) = u'(1) = 0 \},$$
(2.12)

then  $M_0 = D(A)$ . For any  $u \in H^1(0, 1)$ , it is clear that

$$[u,v] = u(0)\overline{v'(0)} - u(1)\overline{v'(1)} = 0, \quad \forall v \in D(A).$$

Therefore, Corollary 2.1 yields

$$D(A^{1/2}) = H^1(0,1).$$

(iii) If A has the boundary conditions

$$D(A) = \{ u \mid u \in H^2(0,1), u(0) = u'(1) = 0 \},$$
(2.13)

by a similar process, we can obtain

$$D(A^{1/2}) = \{ u \mid u \in H^1(0,1), u(0) = 0 \}.$$

(iv) For the boundary conditions

$$D(A) = \{ u \mid u \in H^2(0,1), u(0) = u'(0), u(1) = 0 \},$$
(2.14)

set  $H^+_{\partial} = \mathbb{C}$  and  $\Upsilon = 1$ . Define  $\Gamma : H^1(0,1) \to C$  by  $\Gamma u = u(0)$ . Then  $(C,1,\Gamma)$  is a positive boundary space of A corresponding to iD. Since

$$\begin{split} & [u,v] = -u(x)\overline{v'(x)}|_0^1 = u(0)\overline{v'(0)} - u(1)\overline{v'(1)} \\ & = (\Upsilon\Gamma u, \Gamma v)_\partial^+ - u(1)\overline{v'(1)}, \quad \forall u \in H^1(0,1), \; \forall v \in D(A), \end{split}$$

we have from Theorem 1.1 that

$$D(A^{1/2}) = \{ u \mid u \in H^1(0,1), \ u(1) = 0 \}.$$

Now we consider explicit representations of  $A^{1/2}$ .

PROPOSITION 2.3. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Let

$$\mathcal{M}_0 = \overline{\operatorname{span}} \left\{ \begin{pmatrix} \Upsilon^{1/2} g_n \\ \psi_n \end{pmatrix} \middle| n \ge k+1 \right\}$$

be a closed linear subspace of  $H_{\partial}^+ \times H_1$ . Then

$$\mathcal{M}_0 = \left\{ \left( egin{array}{c} \Upsilon^{1/2} \Gamma x \ Bx \end{array} 
ight| \, x \in D(A^{1/2}) 
ight\}.$$

*Proof.* By Theorem 2.1 and (2.2), we have

$$(\|\Upsilon^{1/2}\Gamma x\|_{\partial}^{+})^{2} = [x, x] = (x, Ax) - (Bx, Bx)_{1}$$
  
$$\leq 2\|A^{1/2}x\|^{2}, \quad \forall x \in D(A^{1/2}).$$
 (2.15)

We have that  $A^{1/2}x_n = 0$ ,  $\forall n, 1 \le n \le k$ , since  $Ax_n = 0$ , for any  $1 \le n \le k$ . Therefore, (2.2) yields

$$Bx_n = 0, \quad \forall n, \ 1 \le n \le k. \tag{2.16}$$

By (2.6), it follows that

$$\Gamma x_n = 0, \quad \forall n, \ 1 \le n \le k.$$
(2.17)

Suppose  $x \in D(A^{1/2})$ . Set  $x = \sum_{n=1}^{\infty} \alpha_n x_n$ . From (2.6) together with (2.7), we have

$$\Upsilon^{1/2}\Gamma x = \sum_{n=1}^{\infty} \alpha_n \Upsilon^{1/2}\Gamma x_n = \sum_{n=k+1}^{\infty} \alpha_n \Upsilon^{1/2}\Gamma x_n = \sum_{n=k+1}^{\infty} \sqrt{\lambda_n} \alpha_n \Upsilon^{1/2} g_n.$$
(2.18)

By (2.2) and (2.5), it follows that

$$Bx = \sum_{n=1}^{\infty} \alpha_n Bx_n = \sum_{n=k+1}^{\infty} \alpha_n Bx_n = \sum_{n=k+1}^{\infty} \sqrt{\lambda_n} \alpha_n \psi_n.$$
(2.19)

Therefore, (2.18) and (2.19) mean that

$$\begin{pmatrix} \Upsilon^{1/2}\Gamma x\\ Bx \end{pmatrix} = \sum_{n=k+1}^{\infty} \sqrt{\lambda_n} \alpha_n \begin{pmatrix} \Upsilon^{1/2}g_n\\ \psi_n \end{pmatrix} \in \mathcal{M}_0, \quad \forall x = \sum_{n=1}^{\infty} \alpha_n x_n \in D(A^{1/2}).$$
(2.20)

Conversely, suppose  $(g, \psi)^{\tau} \in \mathcal{M}_0$ , where  $g \in H^+_{\partial}$  and  $\psi \in H_1$ . By Proposition 2.2, we have

$$\begin{pmatrix} g \\ \psi \end{pmatrix} = \sum_{n=k+1}^{\infty} \beta_n \begin{pmatrix} \Upsilon^{1/2} g_n \\ \psi_n \end{pmatrix},$$

where  $\sum_{n=k+1}^{\infty} |\beta_n|^2 < +\infty$ . Setting  $x = \sum_{n=k+1}^{\infty} \frac{\beta_n}{\sqrt{\lambda_n}} x_n$ , it is easily checked that  $x \in D(A^{1/2})$  and  $(g, \psi)^{\tau} = (\Upsilon^{1/2} \Gamma x, Bx)^{\tau}$ .  $\Box$ 

PROPOSITION 2.4. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Define  $\mathcal{B}: H \to H_{\partial}^+ \times H_1$  by

$$D(\mathcal{B}) = D(A^{1/2}), \qquad \mathcal{B}x = \begin{pmatrix} \Upsilon^{1/2}\Gamma x \\ Bx \end{pmatrix}.$$
 (2.21)

Then

$$\mathcal{B}: H \to H_{\partial}^+ \times H_1$$

is a closed linear operator.

Proof. Suppose that  $\{z_n\} \subset D(A^{1/2})$  such that in  $H, z_n \to x$ , as  $n \to \infty$ , such that in  $H_{\partial}^+, \Upsilon^{1/2}\Gamma z_n \to g$ , as  $n \to \infty$ , and such that in  $H_1, Bz_n \to \psi$ , as  $n \to \infty$ , where  $x \in H$ ,  $g \in H_{\partial}^+$ , and  $\psi \in H_1$ . From the closedness of B it follows that  $x \in D(B)$  and  $Bx = \psi$ . From (2.6) and (2.2), it follows that  $\Upsilon^{1/2}\Gamma x = g$ . Therefore  $\mathcal{B}$  is closed.  $\Box$ 

Now we consider the expression of  $A^{1/2}$ .

THEOREM 2.2. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B. Then there exists a bounded linear operator  $\mathcal{Q}: H_{\partial}^+ \times H_1 \to H$  such that

(i)  $\|Q\| \le 1$ ,  $Q : \mathcal{M}_0 \to H$  is isometric; (ii)  $A^{1/2}x = Q\mathcal{B}x$ , for any  $x \in D(A^{1/2})$ .

*Proof.* For any  $(g, \psi)^{\tau} \in \mathcal{M}_0$ , set

$$\begin{pmatrix} g \\ \psi \end{pmatrix} = \sum_{n=k+1}^{\infty} lpha_n \begin{pmatrix} \Upsilon^{1/2} g_n \\ \psi_n \end{pmatrix}.$$

Define a linear operator  $\mathcal{Q}: H_{\partial}^+ \times H_1 \to H$  as shown below: for  $(g, \psi)^{\tau} \in \mathcal{M}_0$ , set

$$\mathcal{Q}\left(\begin{array}{c}g\\\psi\end{array}\right) = \sum_{n=k+1}^{\infty} \alpha_n x_n;$$

in the orthogonal complement of  $\mathcal{M}_0$  in  $H^+_{\partial} \times H_1$ , define  $\mathcal{Q}$  by the zero operator.

It is easily checked that (i) holds. For any  $x \in D(A^{1/2})$ , by (2.20) we have

$$\mathcal{QB}x = \sum_{n=k+1}^{\infty} \sqrt{\lambda_n} \alpha_n x_n = A^{1/2} x.$$

From the proof of Theorem 2.2, we know that, if  $\Gamma x = 0$ ,  $\forall x \in D(A^{1/2})$ , then  $\mathcal{Q}$  can be defined from  $H_1$  to H. In fact, we have the following result.

THEOREM 2.3. Let  $(B, H_1)$  be a pseudo-square-root of A. Set

$$\mathbb{R}_0 = \overline{\operatorname{span}}\{\psi_n \mid n \ge k+1\}.$$

Then  $M_0 = D(A)$  if and only if there exists a bounded linear operator  $Q: H_1 \to H$  such that

- (i)  $||Q|| \leq 1, Q : \mathbb{R}_0 \to H$  is isometric;
- (ii)  $A^{1/2}x = QBx, \forall x \in D(A^{1/2}).$

*Proof.* The "only if" part follows from Theorem 2.2. The "if" part. By (ii) we have that  $\sqrt{\lambda_n}x_n = A^{1/2}x_n = QBx_n, \forall n \ge k+1$ , i.e.,

$$Q\psi_n = x_n, \quad n = k+1, k+2, \dots$$
 (2.22)

The isometry of  $Q : \mathbb{R}_0 \to H$  yields

$$(\psi_n, \psi_m)_1 = (x_n, x_m), \quad \forall n, m \ge k+1.$$
 (2.23)

From the orthogonality of  $\{x_n\}$  in H we obtain that

$$[x_n, x_m] = \lambda_m \left[ (x_n, x_m) - \sqrt{\frac{\lambda_n}{\lambda_m}} (\psi_n, \psi_m)_1 \right] = 0, \quad \forall n, m \ge k+1.$$
(2.24)

In addition, we have by a calculation that

$$[x_n, x] = \overline{[x, x_n]} = 0, \quad \forall x \in D(A), \ 1 \le n \le k.$$
(2.25)

For any  $x, y \in D(A)$ ,  $x = \sum_{n=1}^{\infty} \alpha_n x_n$ ,  $y = \sum_{n=1}^{\infty} \beta_n x_n$ , we have from (2.24) and (2.25) that

$$[x,y] = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \alpha_n \overline{\beta_m} [x_n, x_m] = 0.$$

Thus  $M_0 = D(A)$ .  $\Box$ 

EXAMPLE 2.4. Let A and B be given by Example 2.3.

(i) Let A have any one of the boundary conditions of (2.11), (3.12), and (2.13). Then  $M_0 = D(A)$ . By Theorem 2.3, there is a bounded linear operator  $Q: L^2(0,1) \to L^2(0,1)$  such that

$$A^{1/2}u = Q(iu'), \quad \forall u \in D(A^{1/2}).$$

Set

$$F(t,s) = \sum_{n=1}^{\infty} (Qx_n)(t)\overline{x_n(s)}, \quad 0 \le t, \ s \le 1.$$

Then

$$A^{1/2}u = i \int_0^1 F(t,\xi) u'(\xi) \, d\xi, \quad orall u \in D(A^{1/2}).$$

(ii) Let the boundary conditions of A be given as (2.14). Then  $(\mathbb{C}, 1, \Gamma)$  is a positive boundary space of A corresponding to B, where  $\Gamma u = u(0), \forall u \in H^1(0, 1)$ . By Theorem 2.2, there is a bounded linear operator  $\mathcal{Q} : \mathbb{C} \times L^2(0, 1) \to L^2(0, 1)$  such that

$$A^{1/2}u = \mathcal{Q}\left(egin{array}{c} u(0)\ iu'(\cdot) \end{array}
ight), \quad orall u \in D(A^{1/2}) = \{u \mid u \in H^1(0,1), u(1) = 0\}.$$

Let  $\{x_n(\cdot)\}\$  be the orthonormal basis of  $L^2(0,1)$  consisting of the eigenfunctions of A. Set

$${\mathcal Y}_n(t) = {\mathcal Q} \left(egin{array}{c} x_n(0) \ 0 \end{array}
ight) \in L^2(0,1), \quad n \geq 1, \; 0 \leq t \leq 1,$$

$$\mathcal{Z}_n(t) = \mathcal{Q}\begin{pmatrix} 0\\x_n \end{pmatrix} \in L^2(0,1), \quad n \ge 1, \ 0 \le t \le 1,$$
$$F_{\partial}(t,s) = \sum_{n=1}^{\infty} \mathcal{Y}_n(t) \overline{x_n(s)}, \qquad F(t,s) = \sum_{n=1}^{\infty} \mathcal{Z}_n(t) \overline{x_n(s)}, \quad 0 \le t, s \le 1$$

Then we have that

$$A^{1/2}u = \mathcal{Q}\begin{pmatrix} u(0) \\ 0 \end{pmatrix} + \mathcal{Q}\begin{pmatrix} 0 \\ iu'(\cdot) \end{pmatrix}$$
  
=  $\sum_{n=1}^{\infty} (u, x_n)_{L^2(0,1)} \mathcal{Y}_n(t) + \sum_{n=1}^{\infty} (iu', x_n)_{L^2(0,1)} \mathcal{Z}_n(t)$   
=  $\int_0^1 F_{\partial}(t, \xi) u(\xi) d\xi + i \int_0^1 F(t, \xi) u'(\xi) d\xi, \quad \forall u \in D(A^{1/2}).$   $\Box$  (2.26)

Finally, we consider the relationship between  $A^{1/2}$  and B from a different point of view. For the elastic "beam" operator A, i.e.,  $Au = u^{(4)}(\cdot)$ ,  $u \in D(A) \subset H^4(0,1)$ , and the assumption that the boundary conditions of A are "symmetric" (see Section 3, (3.3)), D. L. Russell [2] proved that there is a bounded linear operator  $P : L^2(0,1) \to L^2(0,1)$  such that

$$PA^{1/2}u = -u''(\cdot), \quad \forall u \in D(A^{1/2}).$$

Here we have the following general result.

THEOREM 2.4. Let  $(B, H_1)$  be a pseudo-square-root of A. Then there is a bounded linear operator  $P: H \to H_1$  with  $||P|| \leq 1$  such that

$$PA^{1/2}x = Bx, \quad \forall x \in D(A^{1/2}).$$
 (2.27)

*Proof.* Define  $P: H \to H_1$  by

$$Px = \sum_{n=k+1}^{\infty} \alpha_n \psi_n, \quad \forall x = \sum_{n=1}^{\infty} \alpha_n x_n \in H.$$
(2.28)

Let  $(H^+_{\partial}, \Upsilon, \Gamma)$  be a positive boundary space of A corresponding to B. Set

$$\begin{pmatrix} g\\ \psi \end{pmatrix} = \sum_{n=k+1}^{\infty} \alpha_n \begin{pmatrix} \Upsilon^{1/2} g_n\\ \psi_n \end{pmatrix}, \quad \forall x = \sum_{n=1}^{\infty} \alpha_n x_n \in H.$$

Since  $\sum_{n=1}^{\infty} |\alpha_n|^2 < +\infty$  we know that  $(g, \psi)^{\tau} \in H_{\partial}^+ \times H_1$ . Hence

$$Px = \psi \in H_1, \quad \forall x \in H,$$

and

$$\|Px\|_{1}^{2} \leq \|g\|_{\partial_{+}}^{2} + \|\psi\|_{1}^{2} = \sum_{n=k+1}^{\infty} |\alpha_{n}|^{2} \leq \|x\|^{2}, \quad \forall x = \sum_{n=1}^{\infty} \alpha_{n} x_{n} \in H,$$

that is,  $||P|| \leq 1$ . Furthermore, for any  $x = \sum_{n=1}^{\infty} \alpha_n x_n \in D(A^{1/2})$ , we have

$$PA^{1/2}x = \sum_{n=k+1}^{\infty} \sqrt{\lambda_n} \alpha_n P x_n = Bx.$$

By Theorem 2.4, if  $P: H \to \mathbb{R}_0$  has a bounded inverse  $P^{-1}: \mathbb{R}_0 \to H$ , then  $A^{1/2}$  can be written as  $A^{1/2} = P^{-1}B$ , where  $P^{-1}: \mathbb{R}_0 \to H$  is not isometric in general. In particular, if  $\{\psi_n | n \ge k+1\}$  is an orthogonal set in  $H_1$  (which is not true in general), then the above assertion holds. In the following we give a necessary condition on the orthogonality of  $\{\psi_n | n \ge k+1\}$  if  $\dim(H^+_{\partial}) < +\infty$  by which an error in [2] can be corrected (see Sec. 3).

PROPOSITION 2.5. Let  $(B, H_1)$  be a pseudo-square-root of A and  $(H_{\partial}^+, \Upsilon, \Gamma)$  a positive boundary space of A corresponding to B with  $\dim(H_{\partial}^+) < \infty$ . If  $\{\psi_n | n \ge k+1\}$  is an orthogonal set in  $H_1$ , then there is a positive integer  $n_0$  such that

$$[x_n, x_n] = 0, \quad \forall n \ge n_0. \tag{2.29}$$

*Proof.* Suppose dim $(H_{\partial}^+) = m$ . By an argument similar to the proof of Proposition 2.2, we have

$$(x_n, x_j) = \sqrt{\frac{\lambda_n}{\lambda_j}} \left[ (\Upsilon^{1/2} g_n, \Upsilon^{1/2} g_j)_{\partial}^+ + (\psi_n, \psi_j)_1 \right], \quad \forall n, j \ge k+1.$$

$$(2.30)$$

Thus  $\{\Upsilon^{1/2}g_n|n \ge k+1\}$  is an orthogonal set in  $H^+_{\partial}$ . Since dim $(H^+_{\partial}) = m$ , there are at most *m* nonzero elements in  $\{\Upsilon^{1/2}g_n|n \ge k+1\}$ . Therefore, there exists a positive integer  $n_0 \ge k+1$  such that

$$\Upsilon^{1/2}g_n=0,\quad orall n\geq n_0.$$

From (2.24), we have

$$[x_n, x_n] = \lambda_n [(x_n, x_n) - (\psi_n, \psi_n)_1] = 0, \quad \forall n \ge n_0. \quad \Box$$

3. Applications to elastic "beam" operators. Let  $H = L^2(0, 1)$ . In this section, we shall consider elastic "beam" operators defined by

$$Au = u^{(4)}(\cdot), \quad u \in D(A) \subset H^4(0,1).$$
 (3.1)

Since

$$(Au, u)_{L^{2}(0,1)} = u'''\overline{u}|_{0}^{1} - u''\overline{u'}|_{0}^{1} + \int_{0}^{1} |u''|^{2} dt, \quad \forall u \in D(A),$$

it is supposed that

$$u'''\overline{u}|_{0}^{1} - u''\overline{u'}|_{0}^{1} \ge 0, \quad \forall u \in D(A).$$
(3.2)

Set  $H_1 = L^2(0,1)$ . Define  $B: L^2(0,1) \to L^2(0,1)$  by

$$D(B) = H^2(0,1), \qquad Bu = -D^2u = -u''(\cdot).$$

The following proposition can be easily verified.

PROPOSITION 3.1. Let A be a nonnegative selfadjoint operator that satisfies (3.1) and (3.2). Then  $(-D^2, L^2(0, 1))$  is a pseudo-square-root of A.

First, we introduce a main result given in [2]. Consider the real Hilbert space  $L^2(0, 1)$ . Set

$$B(u,v) = u'''(\cdot)v(\cdot) - u''(\cdot)v' + u'(\cdot)v''(\cdot) - u(\cdot)v'''(\cdot), \quad \forall u, v \in D(A);$$
  
$$\psi_n(\cdot) = -(\sqrt{\lambda_n})^{-1}D^2x_n(\cdot), \quad n \ge k+1,$$

where  $\lambda_n, x_n$ , and k are the same as those in Sec. 2. D. L. Russell in [2] has shown the following result.

Let A be a nonnegative selfadjoint operator on  $L^2(0,1)$  with the compact resolvent such that (3.1) and (3.2) hold. Suppose that

$$B(u,v) = 0, \quad \forall u, v \in D(A), \tag{3.3}$$

at x = 0 and at x = 1. Then

(i) there is a bounded linear operator  $P: L^2(0,1) \to L^2(0,1)$  such that

$$PA^{1/2}u = -D^2u, \quad \forall u \in D(A^{1/2}).$$
 (3.4)

(ii) Set  $\mathbb{R}_0 = \overline{\operatorname{span}}\{\psi_n | n \ge k+1\}$ . Then P has a bounded inverse on  $\mathbb{R}_0$ , which extends to a bounded linear operator Q on  $L^2(0,1)$ , so that likewise

$$A^{1/2}u = Q(-D^2)u, \quad \forall u \in D(A^{1/2}).$$

From the proof of the above assertion (ii) in [2, pp. 761–765] we know that the assertion (ii) is based on the assertion that  $\{\psi_n | n \ge k + 1\}$  is an orthogonal set in  $L^2(0, 1)$ . Unfortunately, Proposition 3.2 below shows that in general  $\{\psi_n | n \ge k + 1\}$  is not an orthogonal set in  $L^2(0, 1)$  under the assumption (3.3) only. Therefore, the proof of the assertion (ii) given in [2] failed. By Theorem 2.2 and Theorem 2.3, the assertion (ii) can be revised (see Example 3.1 below).

EXAMPLE 3.1. Let us consider a cantilever with elastic forces applying at the free end, x = 1. The boundary conditions become

$$u(0) = u'(0) = 0, \quad u'''(1) - \alpha u(1) = 0, \quad u''(1) + \beta u'(1) = 0,$$
 (3.5)

where  $\alpha > 0, \beta > 0$ .

Set  $H^+_{\partial} = \mathbb{C}^2$ . Define  $\Upsilon : \mathbb{C}^2 \to \mathbb{C}^2$  by

$$\Upsilon = \begin{pmatrix} \alpha & 0 \\ 0 & \beta \end{pmatrix}.$$

Define  $\Gamma: H^2(0,1) \to \mathbb{C}^2$  by

$$\Gamma u = \begin{pmatrix} u(1) \\ u'(1) \end{pmatrix}, \quad \forall u \in H^2(0,1).$$

It can easily be verified that

$$[u,v] = \alpha u(1)\overline{v(1)} + \beta u'(1)\overline{v'(1)}, \quad \forall u,v \in D(A).$$
(3.6)

Then  $(\mathbb{C}^2, \Upsilon, \Gamma)$  is a positive boundary space of A corresponding to  $-D^2$ . Since

$$\begin{split} [u,v] &= \alpha u(1)\overline{v(1)} + \beta u'(1)\overline{v'(1)} - u(0)\overline{v''(0)} + u'(0)\overline{v''(0)} \\ &= (\Upsilon\Gamma u, \Gamma v)_{\mathbb{C}^2} - u(0)\overline{v''(0)} + u'(0)\overline{v''(0)}, \quad \forall u \in H^2(0,1), \ v \in D(A), \end{split}$$

from Theorem 2.1, we have

$$D(A^{1/2}) = \{ u \mid u \in H^2(0,1), [u,v] = (\Upsilon \Gamma u, \Gamma v)_{\mathbb{C}^2}, \ \forall v \in D(A) \}$$
$$= \{ u \mid u \in H^2(0,1), u(0) = u'(0) = 0 \}.$$

By Theorem 2.2, there is a bounded linear operator  $\mathcal{Q}: \mathbb{C}^2 \times L^2(0,1) \to L^2(0,1)$  such that

$$A^{1/2}u = \mathcal{Q}\begin{pmatrix} \sqrt{\alpha}u(1)\\ \sqrt{\beta}u'(1)\\ -u''(\cdot) \end{pmatrix}, \quad \forall u \in D(A^{1/2}).$$

By an argument similar to Example 2.4, we can find  $F_{\partial_1}(t,\xi)$ ,  $F_{\partial_2}(t,\xi)$ , and  $F(t,\xi)$ , for  $0 \le t, \xi \le 1$ , such that

$$\begin{aligned} A^{1/2}u &= \sqrt{\alpha} \int_0^1 F_{\partial_1}(t,\xi) u(\xi) d\xi + \sqrt{\beta} \int_0^1 F_{\partial_2}(t,\xi) u'(\xi) d\xi - \int_0^t F(t,\xi) u''(\xi) d\xi, \\ \forall u \in D(A^{1/2}). \quad \Box \end{aligned}$$

PROPOSITION 3.2. Let A be the "beam" operator on  $L^2(0,1)$  with the boundary conditions given by (3.5). Then (3.3) holds but  $\{\psi_n | n \ge k+1\}$  is not an orthogonal set in  $L^2(0,1)$ , where  $\psi_n = -(\sqrt{\lambda_n})^{-1} x''_n(\cdot)$ ,  $\lambda_n$  is the eigenvalue of A, and  $x_n(\cdot)$  is the eigenfunction of A corresponding to  $\lambda_n$ , for all  $n \ge 1$ .

*Proof.* It is easily checked that (3.3) holds.

Suppose that  $\{\psi_n | n \ge k + 1\}$  is an orthogonal set in  $L^2(0, 1)$ . By Proposition 2.5, there is a positive integer  $n_0$  such that

$$[x_n, x_n] = \alpha |x_n(1)|^2 + \beta |x_n'(1)|^2 = 0, \quad \forall n \ge n_0.$$

Therefore,  $x_n$  is a solution of the following boundary value problem:

$$\begin{cases} x_n^{(4)}(t) = \lambda_n x_n(t), & 0 < t < 1, \\ x_n(0) = x'_n(0) = x_n(1) = x'_n(1) = x''_n(1) = x''_n(1) = 0, \end{cases}$$
(3.7)

for all  $n \ge n_0$ . It is obvious that the problem (3.7) has the unique zero solution, for all  $n \ge n_0$ . Thus

$$x_n = 0, \quad \forall n \ge n_0.$$

This is a contradiction.  $\Box$ 

REMARK 3.1. It should be noted that there is a modified inner product, as shown in Proposition 2.2, relative to which  $\psi_n$  may be considered orthogonal. Thus, the statement made in [2] is correct if restricted to "SDE" boundary conditions (strictly distributed energy [2]). Without the SDE assumption, the potential energy form associated with the operator A includes some boundary terms.

Finally, we conclude this section by giving an example derived from the pointwise control of a flexible manipulator arm [6].

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EXAMPLE 3.2. Let  $H = \mathbb{C}^3 \times L^2(0,1)$  with the inner product

$$(\Phi_1, \Phi_2) = \alpha_1 \overline{\alpha_2} + \mu_1 \overline{\mu_2} + \xi_1 \overline{\xi_2} + \int_0^1 \varphi_1 \overline{\varphi_2} dt$$

where  $\Phi_j = (\alpha_j, \mu_j, \xi_j, \varphi_j)^{\tau} \in \mathbb{C}^3 \times L^2(0, 1)$ . Define the linear operator A by

$$A\tilde{\varphi} = \begin{pmatrix} -\varphi''(0) \\ -\varphi'''(1) \\ \varphi''(1) \\ \varphi^{(4)}(\cdot) \end{pmatrix}, \qquad \tilde{\varphi} = \begin{pmatrix} \varphi'(0) \\ \varphi(1) \\ \varphi'(1) \\ \varphi(\cdot) \end{pmatrix},$$
$$D(A) = \{\tilde{\varphi} = (\varphi'(0), \varphi(1), \varphi'(1), \varphi(\cdot))^{\intercal} \mid \varphi \in H^4(0, 1), \varphi(0) = 0\}.$$

It is easily checked that A is a nonnegative selfadjoint operator on  $\mathbb{C}^3 \times L^2(0,1)$ . Set  $H_1 = L^2(0,1)$ . Define  $B : \mathbb{C}^3 \times L^2(0,1) \to L^2(0,1)$  by

$$B\tilde{\varphi} = -D^2\varphi = -\varphi''(\cdot), \qquad \tilde{\varphi} = \begin{pmatrix} \varphi'(0) \\ \varphi(1) \\ \varphi'(1) \\ \varphi(\cdot) \end{pmatrix},$$
$$D(B) = \{\tilde{\varphi} = (\varphi'(0), \varphi(1), \varphi'(1), \varphi(\cdot))^{\tau} \mid \varphi \in H^2(0, 1)\}.$$

It is easily checked that  $(B, L^2(0, 1))$  is a pseudo-square-root of A. Since

$$[\tilde{\psi}, \tilde{\varphi}] = 0, \quad \forall \tilde{\psi}, \tilde{\varphi} \in D(A),$$

then  $M_0 = D(A)$ .Since

$$[\tilde{\psi}, \tilde{\varphi}] = -\psi(0)\overline{\varphi'''(0)}, \quad \forall \tilde{\psi} \in D(B), \tilde{\varphi} \in D(A).$$

by Corollary 2.1, we have

$$D(A^{1/2}) = \{ \tilde{\psi} \mid \tilde{\psi} \in D(B), [\tilde{\psi}, \tilde{\varphi}] = 0, \forall \tilde{\varphi} \in D(A) \}$$
  
=  $\{ \tilde{\psi} = (\psi'(0), \psi(1), \psi'(1), \psi(\cdot))^{\tau} \mid \psi \in H^2(0, 1), \psi(0) = 0 \}.$ 

Furthermore, by Theorem 2.3, there is a bounded linear operator  $Q: L^2(0,1) \to \mathbb{C}^3 \times L^2(0,1)$  such that

$$A^{1/2}\tilde{\psi} = -Q(\psi''), \quad \forall \tilde{\psi} = (\psi'(0), \psi(1), \psi'(1), \psi(\cdot))^{\tau} \in D(A^{1/2}).$$

Since  $Q: L^2(0,1) \to \mathbb{C}^3 \times L^2(0,1)$  is bounded, Q has the form

$$Qu = egin{pmatrix} g_1(u) \ g_2(u) \ g_3(u) \ Q_0 u \end{pmatrix}, \quad orall u \in L^2(0,1),$$

where  $g_j(\cdot)$  are bounded linear functionals, for j = 1, 2, 3, and  $Q_0 : L^2(0, 1) \to L^2(0, 1)$  is a bounded linear operator. By the Riesz theorem and an argument similar to Example 2.4, we obtain  $f_j \in L^2(0, 1)$  and F(t, s), j = 1, 2, 3,  $0 \le t, s \le 1$ , such that

$$A^{1/2}\tilde{\varphi} = \begin{pmatrix} -\int_0^1 f_1(\xi)\varphi''(\xi)d\xi \\ -\int_0^1 f_2(\xi)\varphi''(\xi)d\xi \\ -\int_0^1 f_3(\xi)\varphi''(\xi)d\xi \\ -\int_0^1 F(\cdot,\xi)\varphi''(\xi)d\xi \end{pmatrix}, \quad \forall \tilde{\varphi} = (\varphi'(0),\varphi(1),\varphi'(1),\varphi(\cdot))^{\tau} \in D(A^{1/2}). \quad \Box$$

4. Applications to *n*-dimensional "wave" operators. In this section we shall apply the results given in previous sections to a variety of boundary value problems of *n*-dimensional "wave" operators. To our knowledge, there presently is no available method for readily calculating the square root of high-dimensional Laplace operators. Let  $\Omega \subset \mathbb{R}^n$  be a bounded domain of  $\mathbb{R}^n$  with smooth boundary  $\partial\Omega$  of class  $C^2$  and  $\frac{\partial}{\partial n}$ the normal derivative.

EXAMPLE 4.1. Let  $\partial \Omega = \Gamma_1 \cup \Gamma_2 \cup \Gamma_3$ . Consider the Laplace operator on  $L^2(\Omega)$ 

$$D(A) = \left\{ u \left| u \in H^{2}(\Omega), u |_{\Gamma_{1}} = 0, \frac{\partial u}{\partial n} \right|_{\Gamma_{2}} = 0, \left( \frac{\partial u}{\partial n} + \alpha u \right) \Big|_{\Gamma_{3}} = 0 \right\},$$

$$Au = -\Delta u, \qquad (4.1)$$

where  $\alpha > 0$ .

Set  $H_1 = (L^2(\Omega))^n$ . Define the closed linear operator  $B: L^2(\Omega) \to (L^2(\Omega))^n$  by

$$D(B) = H^{1}(\Omega), \qquad Bu = \nabla u = \left(\frac{\partial u}{\partial x_{1}}, \frac{\partial u}{\partial x_{2}}, \dots, \frac{\partial u}{\partial x_{n}}\right)^{\tau}$$

It is easily checked that  $(\nabla, (L^2(\Omega))^n)$  is a pseudo-square-root of A. By Green's formula, we have

$$[u,v] = -\int_{\partial\Omega} u \frac{\partial \overline{v}}{\partial n} d\sigma$$
  
=  $-\int_{\Gamma_1} u \frac{\partial \overline{v}}{\partial n} d\sigma + \alpha \int_{\Gamma_3} u \overline{v} d\sigma, \quad \forall u \in H^1(\Omega), \ \forall v \in D(A).$  (4.2)

Set  $H^+_{\partial} = L^2(\Gamma_3)$  and  $\Upsilon = \alpha I$ , where I is the identity mapping on  $L^2(\Gamma_3)$ . Define  $\Gamma : H^1(\Omega) \to L^2(\Gamma_3)$  by

$$\Gamma u = u|_{\Gamma_3}, \quad \forall u \in H^1(\Omega).$$

By (4.2), we have

$$[u,v] = \alpha \int_{\Gamma_3} u\overline{v} \, d\sigma = (\Upsilon \Gamma u, \Gamma v)_{L^2(\Gamma_3)}, \quad \forall u,v \in D(A).$$

Thus  $(L^2(\Gamma_3), \alpha I, \Gamma)$  is a positive boundary space of A corresponding to  $\nabla$ . From Theorem 2.1 and (4.2), we have

$$D(A^{1/2}) = \{ u \mid u \in H^1(\Omega), u |_{\Gamma_1} = 0 \}.$$
(4.3)

By Theorem 2.2, there is a bounded linear operator  $\mathcal{Q}: L^2(\Gamma_3) \times (L^2(\Omega))^n \to L^2(\Omega)$ such that

$$A^{1/2}u = \mathcal{Q}\left(\frac{\sqrt{\alpha}u|_{\Gamma_3}}{\nabla u}\right), \quad \forall u \in D(A^{1/2}).$$
(4.4)

Suppose that  $\{x_j(\zeta)\}$  is the orthonormal basis of  $L^2(\Omega)$  consisting of the eigenfunctions of A and  $e_j = (0, \ldots, 1, \ldots, 0)^{\tau} \in \mathbb{R}^{n+1}, j = 0, 1, \ldots, n$ . Setting

$$y_j(\zeta) = \mathcal{Q}(x_j|_{\Gamma_3}e_0) \in L^2(\Omega), \quad Z_{j,m}(\zeta) = \mathcal{Q}(x_je_m) \in L^2(\Omega), \quad j \ge 1, \ 1 \le m \le n, \ \zeta \in \Omega,$$

$$F_{\partial}(\zeta,\xi) = \sum_{j=1}^{\infty} y_j(\zeta) \overline{x_j(\xi)}, \qquad F_m(\zeta,\xi) = \sum_{j=1}^{\infty} Z_{j,m}(\zeta) x_j(\xi), \quad \forall \zeta, \ \xi \in \Omega, \ 1 \le m \le n.$$

We finally obtain

$$A^{1/2}u = \sqrt{\alpha} \int_{\Omega} F_{\partial}(\zeta,\xi)u(\xi) \,d\xi + \sum_{m=1}^{n} \int_{\Omega} F_{m}(\zeta,\xi) \frac{\partial}{\partial x_{m}}u(\xi) \,d\xi, \quad \forall u \in D(A^{1/2}). \quad \Box$$

EXAMPLE 4.2. Let A be the Laplace operator with the domain

$$D(A) = \left\{ u \mid u \in H^2(\Omega), \int_{\partial \Omega} u \, d\sigma = 0, \frac{\partial u}{\partial n} = \frac{1}{|\partial \Omega|} \int_{\partial \Omega} \Delta u \, d\sigma \right\},$$

where  $|\partial \Omega|$  is the measure of  $\partial \Omega$  in  $\mathbb{R}^{n-1}$ .

Let  $(\nabla, (L^2(\Omega))^n)$  be given by Example 4.1. Then  $(\nabla, (L^2(\Omega))^n)$  is a pseudo-square-root of A. Since

$$[u,v] = -\int_{\partial} u \frac{\partial \overline{v}}{\partial n} d\sigma = \frac{-1}{|\partial \Omega|} \int_{\partial \Omega} \Delta \overline{v} \, d\sigma \int_{\partial} u \, d\sigma = 0, \quad \forall u, v \in D(A),$$

then  $M_0 = D(A)$ . By Corollary 2.1, we have that

$$D(A^{1/2}) = \{ u \mid u \in H^{1}(\Omega), [u, v] = 0, \forall v \in D(A) \}$$
  
=  $\{ u \mid u \in H^{1}(\Omega), \int_{\partial \Omega} u \, d\sigma = 0 \}.$  (4.5)

By Theorem 2.3, there exists a bounded linear operator  $Q: (L^2(\Omega))^n \to L^2(\Omega)$  such that  $A^{1/2}u = Q(\nabla u), \quad u \in D(A^{1/2}).$ 

By an argument similar to Example 4.1, there exist 
$$F_m(\zeta,\xi)$$
, for  $m = 1, 2, ..., n$  such that

$$A^{1/2}u = \sum_{m=1}^{n} \int_{\Omega} F_m(\zeta,\xi) \frac{\partial u(\xi)}{\partial x_m} d\xi, \quad \forall u \in D(A^{1/2}). \quad \Box$$
(4.6)

EXAMPLE 4.3. Let A be the Laplace operator with the domain

$$D(A) = \{ u \mid u \in H^2(\Omega), \int_{\partial\Omega} \frac{\partial u}{\partial n} d\sigma = 0, u|_{\partial\Omega} = ext{constant} \}.$$

Then  $(\nabla, (L^2(\Omega))^n)$  is a pseudo-square-root of A.

Since

$$[u,v] = -\int_{\partial\Omega} u \frac{\partial \overline{v}}{\partial n} d\sigma = -u|_{\partial\Omega} \int_{\partial\Omega} \frac{\partial \overline{v}}{\partial n} d\sigma = 0, \quad \forall u,v \in D(A),$$

we know that  $M_0 = D(A)$ . Suppose  $u \in H^1(\Omega)$  such that

$$[u,v] = -\int_{\partial\Omega} u \frac{\partial\overline{v}}{\partial n} d\sigma = 0, \quad \forall v \in D(A).$$

It is easily checked by the trace operator theorem that

$$\int_{\partial\Omega} ug \, d\sigma = 0, \quad \forall g \in L^2(\partial\Omega), \quad \int_{\partial\Omega} g \, d\sigma = 0.$$

Thus  $u|_{\partial\Omega} = \text{constant}$ . Therefore, by Corollary 2.1, we have

$$D(A^{1/2}) = \{ u \mid u \in H^1(\Omega), [u, v] = 0, \forall v \in D(A) \}$$
$$= \{ u \mid u \in H^1(\Omega), u|_{\partial\Omega} = \text{constant} \}.$$

In addition, by Theorem 2.3,  $A^{1/2}$  has the form of (4.6).  $\Box$ 

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EXAMPLE 4.4. Let

$$a(\zeta) = (a_{ij}(\zeta))_{n \times n}, \quad a_{ij} \in C^1(\overline{\Omega}), \ 1 \le i, j \le n,$$

be a symmetric matrix on  $\mathbb{C}^n$  and suppose that there is  $\delta > 0$  such that

$$\sum_{j=1}^{n}\sum_{i=1}^{n}a_{ij}(\zeta)\mu_{i}\overline{\mu_{j}}\geq\delta|\mu|^{2},\quad\forall\mu=(\mu_{1},\mu_{2},\ldots,\mu_{n})^{\tau}\in\mathbb{C}^{n},\;\zeta\in\overline{\Omega}.$$

Define A by

$$D(A) = \{ u \mid u \in H^2(\Omega), u|_{\partial\Omega} = 0 \},\$$
$$Au = -\sum_{j=1}^n \frac{\partial}{\partial x_j} \left( \sum_{i=1}^n a_{ij}(\zeta) \frac{\partial u}{\partial x_i} \right).$$

It is easily checked that A is a nonnegative selfadjoint operator on  $L^2(\Omega)$ .

Set  $H_1 = (L^2(\Omega))^n$ . Define  $B : L^2(\Omega) \to (L^2(\Omega))^n$  by

$$D(B) = H^1(\Omega), \qquad Bu = (a(\zeta))^{1/2} \nabla u$$

It is easily checked that  $((a(\zeta))^{1/2}\nabla, (L^2(\Omega))^n)$  is a pseudo-square-root of A. By Green's formula, we have

$$[u,v] = -\int_{\partial\Omega} u \sum_{j=1}^{n} \left( \sum_{i=1}^{n} \overline{a_{ij}} \frac{\partial \overline{v}}{\partial x_i} \right) n_j \, d\sigma = 0, \quad \forall u, v \in D(A),$$

where  $n = (n_1, n_2, ..., n_n)$  is the unit normal direction of  $\partial \Omega$ . Thus  $M_0 = D(A)$ . Hence

$$D(A^{1/2}) = \{ u \mid u \in H^1(\Omega), [u, v] = 0, \ \forall v \in D(A) \}$$
$$= \{ u \mid u \in H^1(\Omega), u|_{\partial\Omega} = 0 \}.$$

By Theorem 2.3, there is a bounded linear operator  $Q: (L^2(\Omega))^n \to L^2(\Omega)$  such that

$$A^{1/2}u = Q((a(\zeta))^{1/2}\nabla u), \quad \forall u \in D(A^{1/2}).$$

By an argument similar to Example 4.2,  $A^{1/2}$  also has a form of (4.6).

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