

SUBELLIPTIC ESTIMATES FOR THE $\bar{\partial}$ -NEUMANN PROBLEM IN C^2

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1. Introduction

In this paper we prove a conjecture of J. J. Kohn concerning precise subelliptic estimates for the local $\bar{\partial}$ -Neumann problem in C^2 . Let Ω be a bounded open set in C^2 with C^∞ boundary ω . If ω is pseudoconvex near a point $P \in \omega$, and P is of type m (the precise definitions are given in § 2), then Kohn proved that the subelliptic estimate

$$\|\phi\|_{(s)}^{(\Omega)} \leq C_s(\|\bar{\partial}\phi\|_{(0)}^{(\Omega)} + \|\theta\phi\|_{(0)}^{(\Omega)} + \|\phi\|_{(0)}^{(\Omega)})$$

holds for all $s < 1/(m+1)$ (see [8, (7.4)]). Here ϕ is a C^∞ one-form with compact support in $\Omega \cap U$ where U is some sufficiently small neighborhood of P , θ is the adjoint of $\bar{\partial}$, and ϕ is in the domain of θ .

In [7] and [8] Kohn suggested that (i) the subelliptic estimate in question holds with $s = 1/(m+1)$, and (ii) it cannot hold with $s > 1/(m+1)$. In Theorem 3.7 of this paper we shall prove the second conjecture. We do not know whether $s = 1/(m+1)$ is achieved. In proving Theorem 3.7 we make use of results obtained by Yu. V. Egorov [1], L. Hörmander [4], [5] and W. J. Sweeney [9], which enable us to reduce the problem to a similar question concerning a system of pseudo-differential operators on ω . We shall compute these pseudo-differential operators with great precision by utilizing some results of Kohn (see [7] and [8]) concerning the behavior of ω near a point of type m . Our notation and terminology are standard (see e.g. [3] and [4]).

2. The Levi invariants

We recall the basic definitions of [8]. Let Ω be a bounded open subset of C^2 with C^∞ boundary ω , and let $r(P)$ denote the distance of the point P from ω , and assume that $r < 0$ in Ω and $r > 0$ outside of Ω . A vector field L is said to be *holomorphic* in some open set $U \subset C^2$ if it can be written in the form

$$(2.1) \quad L = a^1 \frac{\partial}{\partial z_1} + a^2 \frac{\partial}{\partial z^2}, \quad a^i \in C^\infty(U),$$

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where $\partial/\partial z_j = \frac{1}{2}(\partial/\partial x_j - i(\partial/\partial y_j))$, $j = 1, 2$. A vector field L is said to be *tangential* if at each point of ω it is tangent to ω , that is, if $L(r) = 0$ at $r = 0$. As usual we define \bar{L} by

$$(2.2) \quad \bar{L} = \bar{a}^1 \frac{\partial}{\partial \bar{z}_1} + \bar{a}^2 \frac{\partial}{\partial \bar{z}_2} .$$

If T_1 and T_2 are two vector fields, we define the Lie bracket by $[T_1, T_2] = T_1 T_2 - T_2 T_1$. The Lie algebra generated by T_1 and T_2 over the C^∞ functions is the smallest module over the C^∞ functions closed under $[,]$, and is denoted by $\mathcal{L}\{T_1, T_2\}$. $\mathcal{L}\{T_1, T_2\}$ is filtered, that is,

$$\mathcal{L}\{T_1, T_2\} = \bigcup_{k=0}^{\infty} \mathcal{L}_k\{T_1, T_2\} ,$$

where $\mathcal{L}_0\{T_1, T_2\}$ is the module spanned by T_1 and T_2 , and $\mathcal{L}_{k+1}\{T_1, T_2\}$ is the module spanned by the elements of $\mathcal{L}_k\{T_1, T_2\}$ and the elements of the form $[A, T_i]$ with $A \in \mathcal{L}_k\{T_1, T_2\}$. Set

$$\mathcal{L} = \mathcal{L}\{L, \bar{L}\} , \quad \mathcal{L}_k = \mathcal{L}_k\{L, \bar{L}\} ,$$

where L is a holomorphic tangent vector in some neighborhood of a point $P \in \omega$, which is different from zero at P . We note that the \mathcal{L} and \mathcal{L}_k evaluated at P do not depend on the choice of L .

2.3. Definition. $P \in \omega$ is said to be of finite type if there exists $F \in \mathcal{L}$ such that $\langle (\partial r)_P, F_P \rangle \neq 0$. Here \langle , \rangle denotes contraction between cotangent vectors and tangent vectors, and the subscript P denotes evaluation at P . P of finite type is said to be of *type m* if m is the least integer such that there is an element in \mathcal{L}_m satisfying the above property.

2.4. Definition. Ω is said to be pseudo-convex near a point $P \in \omega$ if there is a neighborhood U of P such that

$$(2.5) \quad \langle \partial r, [\bar{L}, L] \rangle_{\omega \cap U} \geq 0 ,$$

where L is a nonzero tangential holomorphic vector field.

2.6. Definition. If Ω is pseudo-convex near a point $P \in \omega$, and P is of type m , we say that ω is pseudo-convex of order m at P .

3. The local $\bar{\partial}$ -Neumann problem in C^2

Let $H_{(s)}^{(\Omega)}$ and $H_{(s)}^{(\omega)}$ denote the Sobolev spaces on Ω and ω respectively (see e.g. [3]) with norms denoted by $\| \cdot \|_{(s)}^{(\Omega)}$ and $\| \cdot \|_{(s)}^{(\omega)}$ as usual. These spaces and norms are well defined for vector functions, in particular, for $(0, 1)$ -forms $\phi = \phi_1 d\bar{z}_1 + \phi_2 d\bar{z}_2$, $\phi_1, \phi_2 \in C^\infty(\Omega)$. On $(0, 1)$ -forms we have

$$(3.1) \quad \bar{\partial}\phi = (\partial\phi_2/\partial\bar{z}_1 - \partial\phi_1/\partial\bar{z}_2)d\bar{z}_1 \wedge dz_2 .$$

Let θ denote the formal adjoint of $\bar{\partial}$ operating on $(0, 1)$ -forms, that is,

$$(3.2) \quad (\bar{\partial}\phi, \psi)_{L^2(\Omega)} = (\phi, \theta\psi)_{L^2(\Omega)} ,$$

$\phi \in C_0^\infty(\Omega)$ and $\psi \in D_{(0,1)}(\Omega)$, where $D_{(0,1)}(\Omega)$ stands for C^∞ $(0, 1)$ -forms with compact support in Ω . More precisely we have

$$(3.3) \quad \theta(\phi_1 d\bar{z}_1 + \phi_2 d\bar{z}_2) = -\partial\phi_1/\partial z_1 - \partial\phi_2/\partial z_2 .$$

Now we can state the main result of [8].

3.4. Theorem. *Let $P \in \omega$ be a point of type m , and U be an open neighborhood of P such that $U \cap \omega$ is pseudoconvex. Then there exists a constant C_s for all $s, 0 < s < 1/(m + 1)$, such that*

$$(3.5) \quad \|\phi\|_{(s)}^{(a)} \leq C_s(\|\bar{\partial}\phi\|_{(0)}^{(a)} + \|\theta\phi\|_{(0)}^{(a)} + \|\phi\|_{(0)}^{(a)})$$

for all $\phi \in D_{(0,1)}(U \cap \bar{\Omega})$ satisfying $\langle \phi, \bar{\partial}r \rangle = 0$ on $\omega \cap U$.

We note that $\langle \psi, \bar{\partial}r \rangle = 0$ on $\omega \cap U$ is equivalent to

$$\langle \bar{\partial}\phi, \psi \rangle = \langle \phi, \theta\psi \rangle , \quad \phi \in D_{(0,1)}(U \cap \bar{\Omega}) .$$

When $m = 1$, (3.4) holds with $s = \frac{1}{2}$, and this is the best possible estimate (see [4], [6] and [10]). When $m > 1$, we do not have such a precise result. On the other hand, we have the following result.

3.7. Theorem. *Let $P \in \omega$ be a point of type m , and U a neighborhood of P . Then the estimate (3.5) does not hold with any $s > 1/(m + 1)$.*

The proof of Theorem 3.7 will be given in §§ 4, 5 and 6.

4. The $\bar{\partial}$ operator near a point of type m

Let $P \in \omega$ be a point of type m , and U a sufficiently small neighborhood of P . By an affine change of coordinates we construct coordinates z'_1, z'_2 in U such that

$$(4.1) \quad z'_1(P) = z'_2(P) = (\partial r/\partial z'_1)_P = (\partial r/\partial \bar{z}'_1)_P = (\partial r/\partial y'_2)_P = 0 ,$$

$$(\partial r/\partial x'_2)_P = 1 ,$$

where $z'_1 = x'_1 + iy'_1$ and $z'_2 = x'_2 + iy'_2$. Now r has the following Taylor series expansion

$$(4.2) \quad r(z') = \text{Re } h(z') + \psi(z') + O(|z'|^{m+2}) ,$$

where $\psi(z')$ is a polynomial of degree $m + 1$ such that each term contains $z'_i \bar{z}'_j$ as a factor and

$$(4.3) \quad h(z'_1, z'_2) = \sum_{s+t \leq m+1} \frac{1}{s! t!} \{(\partial/\partial z'_1)^s (\partial/\partial z'_2)^t r\} z'^s_1 z'^t_2 .$$

According to (4.1) and (4.2)

$$(4.4) \quad (\partial h/\partial z'_1)_0 = 0, \quad (\partial h/\partial z'_2)_0 = (\partial r/\partial x_2)_0 = 1 .$$

Thus z'_1 and h are linearly independent in U (here we need U to be sufficiently small), and therefore we can introduce holomorphic coordinates $w_1 = u_1 + iv_1$, $w_2 = u_2 + iv_2$ defined by $w_1 = z'_1$ and $w_2 = h$. Then (4.2) becomes

$$(4.5) \quad r(w_1, w_2) = u_2 + \gamma(w_1, w_2) + O(|w|^{m+2}),$$

where

$$(4.6) \quad \gamma(w_1, w_2) = O(|w|^2)$$

is a polynomial of degree $m + 1$ which contains no pure terms, that is, holomorphic or antiholomorphic terms.

To derive a precise expression for the $\bar{\partial}$ operator we set

$$|\nabla r|\omega^1 = r_{w_2} dw_1 - r_{w_1} dw_2, \quad |\nabla r|\omega^2 = r_{w_1} dw_1 + r_{w_2} dw_2 = \bar{\partial} r,$$

where $r_{w_1} = \partial r/\partial w_1$, etc., so that ω_1 and ω_2 yield a basis of the (1,0)-forms in U . Let $\phi = \phi_1 \bar{w}^1 + \phi_2 \bar{w}^2$. From (3.1) it is easy to see that the $\bar{\partial}$ operator on (0,1)-forms ϕ has the following expression in terms of the basis \bar{w}^1 and \bar{w}^2 :

$$(4.7) \quad \bar{\partial} \phi = (-\bar{M}\phi_1 + \bar{L}\phi_2)\bar{w}^1 \wedge \bar{w}^2 + (\text{terms in which } \phi_1 \text{ and } \phi_2 \text{ remain undifferentiated}),$$

where

$$(4.8) \quad |\nabla r|L = r_{w_2} \frac{\partial}{\partial w_1} - r_{w_1} \frac{\partial}{\partial w_2},$$

$$(4.9) \quad |\nabla r|M = r_{w_1} \frac{\partial}{\partial w_1} + r_{w_2} \frac{\partial}{\partial w_2} .$$

Given $\phi = \phi_1 \bar{w}^1 + \phi_2 \bar{w}^2$ the $\bar{\partial}$ -Neumann boundary condition $\langle \phi, \bar{\partial} r \rangle = 0$ on ω is equivalent to the vanishing of ϕ_2 on ω . If $\phi = \phi_1 \bar{w}^1 + \phi_2 \bar{w}^2 \in C^\infty_0(U \cap \bar{\Omega})$ and $\phi_2 = 0$ on ω , then $\theta\phi$ is well defined and is given by the expression

$$(4.10) \quad \theta\phi = -(L\phi_1 + M\phi_2) + (\text{terms in which } \phi_1 \text{ and } \phi_2 \text{ remain undifferentiated}).$$

Thus in terms of the basis $\bar{\omega}^1, \bar{\omega}^2$ the principal part of the $\bar{\delta}$ -Neumann operator on $(0,1)$ -forms is given by

$$(4.11) \quad D_0 = \begin{pmatrix} -\bar{M} & \bar{L} \\ -L & -M \end{pmatrix}.$$

4.12. Lemma. *Let $P \in \omega$ be of type m . Then $\gamma(w_1, 0)$ is a homogeneous polynomial in w_1 of degree $m + 1$. More precisely*

$$(4.13) \quad \gamma(w_1, 0) = \sum_{s+t=m-1} \frac{1}{(s+1)!(t+1)!} L^s \bar{L}^t \langle \partial r, [L, \bar{L}] \rangle w_1^{s+1} \bar{w}_1^{t+1}.$$

Proof. See Kohn [8, Lemma 3.16].
Consider

$$(4.14) \quad r = u_2 + \gamma(u_1, v_1, u_2, v_2) + O(|u|^{m+2} + |v|^{m+2}) = 0,$$

where we set $u = (u_1, u_2)$ and $v = (v_1, v_2)$. Since $(\partial r / \partial u_2)_0 = 1$, we can solve (4.14) for $u_2 = u_2(u_1, v_1, v_2)$ in a neighborhood of 0.

4.15. Lemma. *Let $u_2(u_1, v_1, v_2)$ be a solution of (4.14) in some neighborhood of 0. Then*

$$(4.16) \quad \frac{\partial^{l+k} u_2(0)}{\partial u_1^l \partial v_1^k} = 0 \quad \text{if } l + k \leq m.$$

Proof. According to Lemma 4.12

$$(4.17) \quad \frac{\partial^{l+k} \gamma(0)}{\partial u_1^l \partial v_1^k} = 0 \quad \text{if } l + k \leq m.$$

By the definition of $r, u_2(0) = 0$. Next, replacing u_2 by $u_2(u_1, v_1, v_2)$ in (4.17) we obtain

$$\frac{\partial u_2}{\partial u_1} + \frac{\partial \gamma}{\partial u_1} + \frac{\partial \gamma}{\partial u_2} \frac{\partial u_2}{\partial u_1} + O(|u_1|^{m+1} + |v|^{m+1}) = 0.$$

Since $\gamma = O(|u|^2 + |v|^2)$, this implies that

$$(4.18) \quad \frac{\partial u_2(0)}{\partial u_1} = 0, \text{ and similarly } \frac{\partial u_2(0)}{\partial v_1} = 0.$$

Now suppose that

$$(4.19) \quad \frac{\partial^{l+k} u_2(0)}{\partial u_1^l \partial v_1^k} = 0 \quad \text{if } l + k \leq p$$

for some $p < m$. Then for a fixed l and k satisfying $l + k = p + 1$ we have

$$\begin{aligned} & \frac{\partial^{p+1}u_2}{\partial u_1^l \partial v_1^k} + \frac{\partial^{p+1}\gamma}{\partial u_1^l \partial v_1^k} \\ & + \sum_{\sum_j (s_j + t_j + (q_j - 1)) \leq p+1} C_{\{s_j, t_j, q_j\}_j} \prod_j \left(\frac{\partial^{s_j + t_j} u_2(0)}{\partial u_1^{s_j} \partial v_1^{t_j}} \right)^{q_j} \\ & + O((|u_1| + |v|)^{m+1-p}) = 0 . \end{aligned}$$

In particular if we set $u_1 = v_1 = v_2 = 0$, the induction hypothesis (4.19) implies that

$$\frac{\partial^{p+1}u_2(0)}{\partial u_1^l \partial v_1^k} + \frac{\partial \gamma(0)}{\partial u_2} \frac{\partial^{p+1}u_2(0)}{\partial u_1^l \partial v_1^k} = \frac{\partial^{p+1}u_2(0)}{\partial u_1^l \partial v_1^k} = 0 .$$

This proves Lemma 4.15.

To utilize Lemma 4.15 we set $x_1 = u_1$, $x_2 = v_1$, $x_3 = v_2$ and $\rho = -r$. Then a simple computation yields

$$(4.20) \quad |\mathcal{V}\rho|L = -\frac{1}{2}\rho_{w_2} \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial x_2} \right) - \frac{1}{2}i\rho_{w_1} \frac{\partial}{\partial x_3} ,$$

$$(4.21) \quad |\mathcal{V}\rho|M = -|\mathcal{V}_w\rho|^2 \frac{\partial}{\partial \rho} - \frac{1}{2}\rho_{w_1} \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial x_2} \right) + \frac{1}{2}i\rho_{w_2} \frac{\partial}{\partial x_3} ,$$

where

$$(4.22) \quad |\mathcal{V}_w\rho|^2 = |\rho_{w_1}|^2 + |\rho_{w_2}|^2 .$$

5. Reduction to the boundary

In [4] L. Hörmander reduced the study of the estimate (3.5) from $U \cap \Omega$ to the study of similar estimates involving pseudo-differential operators on $U \cap \omega$, at least in the case $s = \frac{1}{2}$. This result was extended by W. J. Sweeney [10] to arbitrary s , $0 < s \leq 1$. To be able to state the result in our particular case we shall first compute the boundary system of pseudo-differential operators in question. From (4.11) we have

$$(5.1) \quad D_0^*D_0 = (L(-\bar{L}) + M(-\bar{M}))I_2 + \text{first order terms} ,$$

where I_2 stands for the two-by-two identity matrix. Let r^0 denote $d^{0*}d^0$, the principal symbol of $D_0^*D_0$. A somewhat messy calculation yields

$$\begin{aligned} (5.2) \quad r^0(x, \xi, \tau) &= (|L(x, \xi)|^2 + |M(x, \xi, \tau)|^2)I_2 \\ &= \frac{1}{4}\{|\mathcal{V}_w\rho|^2\tau^2 + [\text{Re}(\rho_{w_1}(\xi_1 - i\xi_2)) \\ &\quad - (\text{Im} \rho_{w_2})\xi_3]\tau + \frac{1}{4}|\xi|^2\}I_2 |\mathcal{V}\rho|^{-2} , \end{aligned}$$

where τ stands for the symbol of $\partial/i\partial\rho$, and ρ is assumed to be zero. The equation $r^0(x, \xi, D_\rho)U(\rho) = 0$ has a unique exponentially decreasing solution on R_+ such that $U(0) = u$, which is given by

$$(5.3) \quad \begin{pmatrix} u_1 & e^{m\rho} \\ u_2 & e^{m\rho} \end{pmatrix}$$

where

$$(5.4) \quad m = \frac{1}{2}|\nabla_m \rho|^{-2} \{ -i[\operatorname{Re}(\rho_{w_1}(\xi_1 - i\xi_2)) - (\operatorname{Im} \rho_{w_2})\xi_3] - (|\nabla_{w\rho}|^2|\xi|^2 - [\operatorname{Re}(\rho_{w_1}(\xi_1 - i\xi_2)) - (\operatorname{Im} \rho_{w_2})\xi_3]^2)^{\frac{1}{2}} \} .$$

Following Hörmander (see [4, Theorem 2.3.1]) we define pseudo-differential operators P_1 and P_2 on $U \cap \omega$ with principal symbols $p_1^0(x, \xi)$ and $p_2^0(x, \xi)$, respectively, given by the first column of

$$(5.5) \quad d^0(x, \rho = 0, \xi, D_\rho) \begin{pmatrix} e^{m\rho} \\ 0 \end{pmatrix}$$

evaluated at $\rho = 0$. More explicitly we have

$$(5.6) \quad p_1^0(x, \xi) = \frac{1}{2} \operatorname{Im}(\rho_{w_1}(\xi_1 - i\xi_2)) - \frac{1}{2}(\operatorname{Re} \rho_{w_2})\xi_3 - \frac{1}{2}\{|\nabla_{w\rho}|^2|\xi|^2 - [\operatorname{Re}(\rho_{w_1}(\xi_1 - i\xi_2)) - \operatorname{Im}(\rho_{w_2})\xi_3]^2\}^{\frac{1}{2}} ,$$

$$(5.7) \quad p_2^0(x, \xi) = -\frac{1}{2}i\rho_{w_2}(\xi_1 - i\xi_2) + \frac{1}{2}\rho_{w_1}\xi_3 .$$

5.8. Proposition. *Let $0 < s \leq 1$. Then (3.5) implies the following estimate*

$$(5.9) \quad \|\phi\|_{(s)}^{(\omega)} \leq C_s(\|P_1\phi\|_{(0)}^{(\omega)} + \|P_2\phi\|_{(0)}^{(\omega)} + \|\phi\|_{(0)}^{(\omega)})$$

for all $\phi \in C_0^\infty(U \cap \omega)$.

Proof. Recall that the $\bar{\delta}$ -Neumann boundary condition is equivalent to $\phi_2 = 0$ on ω . Then Proposition 5.8 is a special case of the results of Hörmander (see [4, Theorems 2.3.1 and 2.3.2]) and of Sweeney (see [10, Propositions 5.7 and 5.8]).

6. Proof of Theorem 3.7

First we localize the estimate (5.9).

6.1. Proposition. *Let $0 < s \leq 1$, and set $\delta = 1 - s$. Suppose that the estimate (5.9) holds with*

$$(6.2) \quad \frac{k}{k+1} \leq \delta < \frac{k+1}{k+2},$$

where k is a positive integer. Then for every $(x, \xi) \in T^*(\omega)$, $|\xi| = 1$, there exists a constant C such that

$$(6.3) \quad \int_{\mathbf{R}_3} |\phi(y)|^2 dy \leq C \left\{ \sum_{j=1,2} \int_{\mathbf{R}_3} \left| \sum_{|\alpha+\beta| \leq k} \frac{1}{\alpha! \beta!} \frac{\partial^{\alpha+\beta} p_j^0(x, \xi)}{\partial \xi^\alpha \partial x^\beta} y^\beta (D^\alpha \phi)(y) \lambda^{\delta - |\alpha| \delta - (1-\delta)|\beta|} \right|^2 dy + \lambda^{2\delta - 2(k+1)(1-\delta)} \sum_{|\alpha+\beta| \leq k+1} \int_{\mathbf{R}_3} |y^\beta (D^\alpha \phi)(y)|^2 dy \lambda^{-2|\alpha|(2\delta-1)} \right\},$$

for all $\lambda \geq 1$ and $\phi \in C_0^\infty(\mathbf{R}_3)$.

Proof. See Egorov [1, Theorem 1] and Hörmander [5, Theorems 6.1 and 6.3].

Let $x = 0$ be a point of type m . We shall show that the estimate (6.3) cannot hold at the point $(x_0, \xi^0) = (0, 0, 0, 0, 0, 1)$ when $k < m$. According to Proposition 6.1 this proves that the estimate (5.9) does not hold with $s > 1/(m+1)$, which proves Theorem 3.7. Since $\rho_{w_2}(0) = -\frac{1}{2}$, according to (5.6) and (5.7) we have

$$(6.4) \quad p_1^0(x_0, \xi^0) = p_2^0(x_0, \xi^0) = 0,$$

and therefore we can assume that $k \geq 1$. Furthermore Lemmas 4.12 and 4.15 imply that

$$(6.5) \quad \rho_{w_1}(x) = x_3 h(x) + O(|x|^m).$$

Assume that the estimate (6.3) holds for some δ such that $k/(k+1) \leq \delta < (k+1)/(k+2)$ with $k < m$. We substitute

$$(6.6) \quad \phi(y) = \psi(y_1, y_2, y_3 \lambda^{2\delta-1+\epsilon}), \quad \psi \in C_0^\infty(\mathbf{R}_3)$$

into (6.3) with some $\epsilon > 0$ such that

$$(6.7) \quad (k+1)\epsilon + \delta < (k+1)(1-\delta) \leq 1.$$

According to the right hand side of (6.2) we have $\delta < (k+1)(1-\delta)$ so that such an ϵ can always be found. We change coordinates $y_1 = y'_1$, $y_2 = y'_2$, $y_3 \lambda^{2\delta-1+\epsilon} = y'_3$, divide both sides of (6.3) by $\lambda^{-2\delta+1-\epsilon}$, and let $\lambda \rightarrow \infty$. Then the left hand side of (6.3) becomes

$$(6.8) \quad \int_{\mathbf{R}_3} |\psi(y')|^2 dy'.$$

Next we compute the right hand side of (6.3).

- 1) Terms involving $p_1^0(x, \xi)$ and its derivatives at (x_0, ξ^0) .
- (i) Set $q_1(x, \xi) = \text{Im}(\rho_{w_1}(\xi_1 - i\xi_2))$. Then

$$\begin{aligned}
 (6.9) \quad & \sum_{\substack{|\beta| \leq k-1 \\ j=1,2}} \frac{1}{\beta!} \frac{\partial^{\beta+1} q_1(x_0, \xi^0)}{\partial x^\beta \partial \xi_j} y^\beta (D_j \phi)(y) \lambda^{-(1-\delta)|\beta|} \\
 &= \sum_{\substack{|\beta| \leq k-1 \\ j=1,2}} \frac{1}{\beta!} \frac{\partial^{\beta+1} q_1(x_0, \xi^0)}{\partial x^\beta \partial \xi_j} y^\beta (D_j \psi)(y') \lambda^{-(1-\delta)|\beta| + \beta_3(-2\delta+1-\varepsilon)} \\
 &= O(\lambda^{-(1-\delta)}),
 \end{aligned}$$

since $\beta_3 > 0$ by (6.5).

- (ii) Set $p_1 = \frac{1}{2}q_1 - \frac{1}{2}q_2$. Then

$$\begin{aligned}
 (6.10) \quad & \sum_{\substack{|\alpha+\beta| \leq k \\ \alpha_1+\alpha_2 \neq 0}} \frac{1}{\alpha! \beta!} \frac{\partial^{\alpha+\beta} q_2(x_0, \xi^0)}{\partial \xi^\alpha \partial x^\beta} y^\beta (D^\alpha \phi)(y) \lambda^{\delta - |\alpha|\delta - (1-\delta)|\beta|} \\
 &= \sum_{\substack{|\alpha+\beta| \leq k \\ \alpha_1+\alpha_2 \neq 0}} \frac{1}{\alpha! \beta!} \frac{\partial^{\alpha+\beta} q_2(x_0, \xi^0)}{\partial \xi^\alpha \partial x^\beta} y^\beta (D^\alpha \psi)(y') \\
 &\quad \cdot \lambda^{-(\alpha_1 + \alpha_2 - 1)\delta - \alpha_3(1-\delta-\varepsilon) - (1-\delta)(\beta_1 + \beta_2) - \beta_3(\delta + \varepsilon)} \\
 &= \sum_{j=1}^2 \frac{\partial q_2(x_0, \xi^0)}{\partial \xi_j} (D_j \psi)(y') + o(1) = o(1),
 \end{aligned}$$

because

$$\frac{\partial q_2(x_0, \xi^0)}{\partial \xi_j} = \left(\frac{\partial}{\partial \xi_j} \left(-\frac{1}{2} + \frac{1}{2} \sqrt{\xi_j^2 + 1} \right) \right)_{\xi_j=0} = 0, \quad j = 1, 2.$$

- (iii) Finally set $\xi_1 = \xi_2 = 0$ in q_2 . Then

$$\begin{aligned}
 (6.11) \quad q_2(x, 0, 0, \xi_3) &= (|\rho_{w_1}|^2 + (\text{Re } \rho_{w_2})^2)^{\frac{1}{2}} + \text{Re } \rho_{w_2} \xi_3 \\
 &= -\text{Re } \rho_{w_2} \{1 - (1 + |\rho_{w_1}|^2 / (\text{Re } \rho_{w_2})^2)^{\frac{1}{2}}\} \xi_3 \\
 &= \left(-(\text{Re } \rho_{w_2}) \sum_{j=1}^m a_j \frac{|\rho_{w_1}|^{2j}}{(\text{Re } \rho_{w_2})^{2j}} + O(|x|^{2m+2}) \right) \xi_3 \\
 &= (x_3^2 H(x) + O(|x|^{2m+2})) \xi_3,
 \end{aligned}$$

where $a_j, j = 1, \dots, m$, are the coefficients in the Taylor series expansion $1 + \sum_{j=1}^m a_j x^j + O(|x|^{m+1})$ of $\sqrt{1+x}$ about $x = 0$, and $H(x)$ is a C^∞ function near $x = 0$. Thus

$$\begin{aligned}
 (6.12) \quad & \sum_{|\beta+\alpha_3| \leq k} \frac{1}{\beta!} \frac{\partial^{\beta+\alpha_3} q_2(x_0, \xi^0)}{\partial \xi_3^{\alpha_3} \partial x^\beta} y^\beta (D^{\alpha_3} \phi)(y) \lambda^{\delta - \alpha_3 \delta - (1-\delta)|\beta|} \\
 &= \sum_{|\beta| \leq k} \frac{1}{\beta!} \frac{\partial^\beta q_2(x_0, \xi^0)}{\partial x^\beta} y^\beta \psi(y') \lambda^{\delta + 2(1-2\delta-\varepsilon) - (1-\delta)|\beta| + (\beta_3-2)(1-2\delta-\varepsilon)}
 \end{aligned}$$

$$\begin{aligned}
 & + \sum_{|\beta| \leq k-1} \frac{1}{\beta!} \frac{\partial^{\beta+1} q_2(x_0, \xi^0)}{\partial \xi_3 \partial x^\beta} y'^\beta (D_3 \psi)(y') \lambda^{-(1-\delta)|\beta| + (\beta_3-1)(1-2\delta-\epsilon)} \\
 & = O(\lambda^{-(\delta+2\epsilon)}),
 \end{aligned}$$

since $\beta_3 \geq 2$ according to (6.11).

2) As for $p_2^0(x, \xi)$ we have

$$\begin{aligned}
 & - \frac{1}{4}(\partial \psi / \partial y_1 - i \partial \psi / \partial y_2) \\
 & - \frac{1}{2} i \sum_{0 < |\beta| \leq k-1} \frac{1}{\beta!} \frac{\partial^\beta \rho_{w_2}(x_0)}{\partial x^\beta} y'^\beta ((D_1 \psi)(y') - i(D_2 \psi)(y')) \\
 & \qquad \qquad \qquad \cdot \lambda^{-(1-\delta)|\beta| + \beta_3(1-2\delta-\epsilon)} \\
 (6.13) \quad & + \frac{1}{2} \sum_{\substack{0 < |\beta| \leq k \\ \beta_3 \geq 1}} \frac{1}{\beta!} \frac{\partial^\beta \rho_{w_1}(x_0)}{\partial x^\beta} y'^\beta \psi(y') \lambda^{\delta - (1-\delta)|\beta| + \beta_3(1-2\delta-\epsilon)} \\
 & + \frac{1}{2} \sum_{\substack{0 < |\beta| \leq k-1 \\ \beta_3 \geq 1}} \frac{1}{\beta!} \frac{\partial^\beta \rho_{w_1}(x_0)}{\partial x^\beta} y'^\beta (D_3 \psi)(y') \lambda^{-(1-\delta)|\beta| + (\beta_3-1)(1-2\delta-\epsilon)} \\
 & = - \frac{1}{4}(\partial \psi / \partial y'_1 - i \partial \psi / \partial y'_2) + O(\lambda^{-\epsilon}),
 \end{aligned}$$

where we have used (6.5).

3) Finally, the remainder yields

$$\begin{aligned}
 (6.14) \quad & \sum_{|\alpha + \beta| \leq k+1} y'^\beta (D^\alpha \psi)(y') \lambda^{\delta - (k+1)(1-\delta) - (\alpha_1 + \alpha_2)(2\delta-1) + \beta_3(1-2\delta-\epsilon) + \alpha_3 \epsilon} \\
 & = O(\lambda^{(k+1)\epsilon + \delta - (k+1)(1-\delta)}) = o(1),
 \end{aligned}$$

where we have used (6.7). Thus (6.8), (6.9), (6.10), (6.12), (6.13) and (6.14) yield

$$(6.15) \quad \int_{R_3} |\psi(y)|^2 dy \leq \frac{1}{4} C \int_{R_3} \left| \frac{\partial \psi(y)}{\partial y_1} - i \frac{\partial \psi(y)}{\partial y_2} \right|^2 dy,$$

where $\psi \in C_0^\infty(R_3)$. This is impossible. To see that set $\psi(y) = f(\epsilon y)$, $f \in C_0^\infty(R_3)$, and let $\epsilon \rightarrow 0$. Then the left hand side of (6.15) is $O(\epsilon^{-3})$, while the right hand side is only $O(\epsilon^{-2})$. Hence Theorem 3.7 is proved..

7. Remarks on the estimate (3.5)

In [8] Kohn proved that if $P \in \omega$ is of type m , and ω is pseudo-convex at P , then m must be odd. This result also follows by applying Propositions 2.4 of [9] to the symbol (5.7). Furthermore Kohn conjectured that under the hypothesis of Theorem 3.4 the estimate (3.5) holds with $s = 1/(m + 1)$.

7.1. Proposition. *Let $P \in \omega$ be a point of type m , and suppose that the estimate (3.5) holds with $s = 1/(m + 1)$. Then m is necessarily odd.*

Proof. It suffices to show that if the estimate (6.3) holds with $k = m$, $\delta = m/(m + 1)$ and $(x_0, \xi^0) = (0, 0, 0, 0, 0, 1)$, then m is odd. We shall follow the arguments of § 6 and indicate the necessary changes. Thus we substitute

$$(7.2) \quad \phi(y) = \psi(y_1, y_2, y_3 \lambda^{2\delta-1+\epsilon}) , \quad \psi \in C_0^\infty(\mathbf{R}_3)$$

into (6.3), where

$$(7.3) \quad (m + 1)\epsilon + \delta < (m + 1)(1 - \delta) = 1 .$$

The left hand side of (6.3) again becomes (6.8). (6.9), (6.10) and (6.12) go through unchanged. (6.13) becomes

$$(7.4) \quad -\frac{1}{4}(\partial\psi/\partial y'_1 - i\partial\psi/\partial y'_2) + \frac{1}{2}\gamma_{w_1}(w_1, 0)\psi + O(\lambda^{-\epsilon}) ,$$

and there is no change in (6.14). Thus the hypothesis of Proposition 7.1 implies the following estimate

$$(7.5) \quad \int_{\mathbf{R}_3} |\psi(y)|^2 dy \leq C \int_{\mathbf{R}_3} \left| \frac{\partial\psi}{\partial y_1} - i \frac{\partial\psi}{\partial y_2} - 2\gamma_{w_1}(w_1, 0)\psi(y) \right|^2 dy ,$$

where $\psi \in C_0^\infty(\mathbf{R}_3)$. Set

$$\psi(y_1, y_2, y_3) = \overline{f(y_1, y_2)}g(y_3)e^{2\gamma(w_1, 0)} .$$

Then (7.5) yields

$$(7.6) \quad \int_{\mathbf{R}_2} |f(y)|^2 e^{4\gamma(w_1, 0)} dy \leq C \int_{\mathbf{R}_3} \left| \frac{\partial f}{\partial \bar{w}_1} \right|^2 e^{4\gamma(w_1, 0)} dy ,$$

for all $f \in C_0^\infty(\mathbf{R}_2)$. According to Theorem 2 of [2], (7.6) implies

$$(7.7) \quad \frac{\partial^2 \gamma(w_1, 0)}{\partial w_1 \partial \bar{w}_1} \geq 0 .$$

(Compare Kohn [8, formula (3.10)]. Egorov's proof of (7.7) is based on one of Hörmander's arguments in [4]; see [4, Lemma 1.2.4, especially (1.2.16)].) Now (7.7) clearly implies Proposition 7.1.

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