# SUPERSOLUTIONS AND STABILIZATION OF THE SOLUTION OF A NONLINEAR PARABOLIC SYSTEM

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Abstract \_

Let us consider a nonlinear parabolic system of the following type:

$$(S) \frac{\frac{\partial u}{\partial t} - \operatorname{div} (|\nabla u|^{p-2} \nabla u) = \frac{\partial H}{\partial u}(x, u, v)}{\frac{\partial v}{\partial t} - \operatorname{div} (|\nabla v|^{p-2} \nabla v) = \frac{\partial H}{\partial u}(x, u, v)}$$

with Dirichlet boundary conditions and initial data.

In this paper, we construct sub-supersolutions of (S), and by use of them, we prove that, for  $t_n \to +\infty$ , the solution of (S) converges to some solution of the elliptic system associated with (S).

### 0. Introduction

This paper concerns the existence and asymptotic behaviour of bounded, non negative solutions of the following system of nonlinear equations:

$$(\$) \begin{cases} \frac{\partial u}{\partial t} - \Delta_p \, u = f(x, u, v) & \text{in } \Omega_x \mathbf{R}_+ \\ \frac{\partial v}{\partial t} - \Delta_q v = g(x, u, v) & \text{in } \Omega_x \mathbf{R}_+ \\ u(x, t) = v(x, t) = 0 & \text{in } \partial \Omega_x \mathbf{R}_+ \\ u(x, o) = \varphi_0(x), \, v(x, 0) = \psi_0(x) & \text{in } \Omega \end{cases}$$

where  $p \geq 2$ ,  $q \geq 2$ ,  $\Delta_p u = \text{div}(|\nabla u|^{p-2} \nabla u)$  and  $\Omega$  is a bounded regular open subset of  $\mathbb{R}^N$ .

For p = q = 2, Problem (\$) has been investigated by many authors [5, 6, 12]. (\$) is an example of a nonlinear parabolic system arising from non-Newtonian fluid mechanics. NAKAO [9] studies a similar system in which p = q > 2 and the right hand side is  $f, -\lambda f, \lambda$  constant. The case of a single equation of the type (\$) is studied in [2, 3, 8, 11]. The purpose of this paper is to extend the results of [3] to the system (\$).

First, using sub-supersolutions, we show that (5) has a solution. Moreover, supposing that there exist  $\lambda > 0$ ,  $\mu > 0$  and a function H(x, u, v) such that

 $f = \lambda \frac{\partial H}{\partial u}$  and  $g = \mu \frac{\partial H}{\partial v}$ , we prove that the solution of (§) converges to a solution of the Dirichlet problem for the elliptic system.

We obtain regularizing effects such that:

$$\begin{aligned} &\frac{\partial}{\partial t} \left( |\nabla u|^{p-2} \nabla u \right) \in L^2(t_0, +\infty; L^{p*}(\Omega)) \text{ and} \\ &\frac{\partial}{\partial t} \left( |\nabla v|^{q-2} \nabla v \right) \in L^2(t_0, +\infty; L^{q*}(\Omega)). \end{aligned}$$

Our method is closely related to the paper of LANGLAIS and PHILLIPS [7], and also to the paper of ELHACHIMI and DE THELIN [3] who study the stabilization of the solution of a single equation. Some examples are discussed in part IV, and include:

$$H(x, u, v) = k(x) uv + \lambda u^{\gamma_1 + 1} + \mu v^{\gamma_2 + 1}$$
  
$$H(x, u, v) = -u^m v^n + \lambda u^{\gamma_1 + 1} + \mu v^{\gamma_2 + 1}$$

Some numerical results related to the system (S) are given in [4].

All Theorems are written in the case p > 2, q > 2; obvious modifications (for Theorem 6) give the case p = 2 or q = 2.

#### 1. Preliminaries and sub-supersolutions

Throughout this paper,  $\Omega$  stands for a regular bounded open subset of  $\mathbb{R}^N$ . Left f and g be some functions from  $\mathbb{R}^{N+2}$  to  $\mathbb{R}$  such that:

(1.1) 
$$\begin{cases} f, g \in C^1(\overline{\Omega} \times \mathbb{R} \times \mathbb{R}) \\ \text{and for any } x \in \Omega, u \in \mathbb{R}_+, v \in \mathbb{R}_+ : f(x, 0, v) \ge 0, g(x, u, 0) \ge 0 \end{cases}$$

and

(1.2) 
$$\begin{cases} \text{For any } M > 0, N > 0, \text{ there exist } k_{M,N}^1 > 0, k_{M,N}^2 > 0\\ \text{such that :}\\ a)f(x, u, v) - f(x, w, v) \le k_{M,N}^1(u - w), \forall x \in \Omega, \\ \forall u, v, w : 0 \le w \le u \le M, v \in [0, N]\\ b)g(x, u, v) - g(x, u, w) \le k_{M,N}^2(v - w), \forall x \in \Omega\\ \forall u, v, w : 0 \le w \le v \le N, u \in [0, M]. \end{cases}$$

**Remark 1:** the condition 1.2. a) is satisfied if  $u \to f(x, u, v)$  is a non increasing function on  $\mathbb{R}_+$ .

We shall also use the following notations:

$$\begin{aligned} &\text{for } T > 0, \, Q_T = \Omega \times ]0, T[, \, S_T = \partial \Omega \times [0,T], \\ F(\nabla u) &= |\nabla u|^{p-2} \, \nabla u, G(\nabla v) = |\nabla v|^{q-2} \, \nabla v, \text{with } p > 2 \text{ and } q > 2. \\ &\Delta_p u = \operatorname{div}(F(\nabla u)), \, \Delta_q v = \operatorname{div}(G(\nabla v)). \end{aligned}$$

Let  $\varphi_0, \psi_0$  be given such that:

(1.3) 
$$\begin{cases} \varphi_0 \in W_0^{1,p}(\Omega) \cap L^{\infty}(\Omega), \, \varphi_0 \ge 0\\ \psi_0 \in W_0^{1,q}(\Omega) \cap L^{\infty}(\Omega), \, \psi_0 \ge 0, \end{cases}$$

We say that (u, v) is a solution of (\$) in  $Q_T$  (resp:  $(\hat{u}, \hat{v})$  is a supersolution of (\$) in  $Q_T$ ) iff

(1.4) 
$$\begin{cases} u(\operatorname{resp} \hat{u}) \in L^{\infty}(0,T; W^{1,p}(\Omega) \cap L^{\infty}(\Omega)) \\ v(\operatorname{resp} \hat{v}) \in L^{\infty}(0,T; W^{1,q}(\Omega) \cap L^{\infty}(\Omega)) \end{cases}$$

(1.5) 
$$\frac{\partial u}{\partial t} \left( \operatorname{resp} \frac{\partial \hat{u}}{\partial t} \right) \in L^2(Q_T), \frac{\partial v}{\partial t} \left( \operatorname{resp} \frac{\partial \hat{v}}{\partial t} \right) \in L^2(Q_T).$$

(1.6) 
$$\begin{cases} \frac{\partial u}{\partial t} - \Delta_p u - f(x, u, v) = 0\\ \frac{\partial v}{\partial t} - \Delta_q v - g(x, u, v) = 0\\ (\operatorname{resp} \cdots \ge 0 \quad \operatorname{in} Q_T) \end{cases} \quad \text{in } Q_T$$

(1.7) 
$$u = v = 0 \text{ (resp : } \hat{u} \ge 0, \, \hat{v} \ge 0 \text{) in } S_T.$$

(1.8) 
$$\begin{cases} u(.,0) = \varphi_0 (\operatorname{resp}: \hat{u}(.,0) \ge \varphi_0) \\ v(.,0) = \psi_0 (\operatorname{resp}: \hat{v}(.,0) \ge \psi_0) \end{cases} \text{ in } \Omega$$

Our method is based upon a comparison principale for the system (S); but the usual notion of supersolution does not work; so, following Hernandez [6], we set:

**Definition 1.**  $[(0,0), (\hat{u}, \hat{v})]$  is said to be a sub-supersolution of (S) in  $Q_T$  if it satisfies the following conditions:

(1.9) 
$$\hat{u} \in W^{1,p}(\Omega) \cap L^{\infty}(\Omega), \ \hat{v} \in W^{1,q}(\Omega) \cap L^{\infty}(\Omega)$$

$$(1.10) \qquad \forall x \in \Omega : 0 \le \varphi_0(x) \le \hat{u}(x) \le M_1, 0 \le \psi_0(x) \le \hat{v}(x) \le N_1$$

(1.11) 
$$\begin{cases} \forall x \in \Omega, \forall v \in [0, \hat{v}] : -f(x, 0, v) \le 0 \le -\Delta_p \hat{u} - f(x, \hat{u}, v) \\ \forall x \in \Omega, \forall u \in [0, \hat{u}] : -g(x, u, 0) \le 0 \le -\Delta_q \hat{v} - g(x, u, \hat{v}) \end{cases}$$

**Remark 2:** if we suppose that  $v \to f(x, u, v) \uparrow$  and  $u \to g(x, u, v) \uparrow$  any supersolution of (S) gives a sub-supersolution of (S).

Our first results are sufficient conditions for the existence of sub-supersolutions of (S).

It is well known (cf.[3]) that the problem:

$$\begin{cases} -\Delta_p u = k_0 + k_1 u^{\gamma}, & x \in \Omega \\ u = 0 & x \in \partial \Omega \end{cases}$$

has a supersolution if  $\gamma \in ]0, p-1[$ .

**Theorem 1.** Assume that  $v \to f(x, u, v)$  and  $u \to g(x, u, v)$  are monotone non-increasing functions, and that there exist

$$\lambda_0 \ge 0, \ \mu_0 \ge 0, \ \lambda_1 > 0, \ \mu_1 > 0 \ and \ \gamma_1 \in ]0, \ p-1[, \ \lambda_2 \in ]0, \ q-1[$$

such that

(1.12) 
$$\begin{cases} f(x, u, 0) \leq \lambda_0 + \lambda_1 u^{\gamma_1}, \forall x \in \Omega, \forall u \in \mathbb{R}_+ \\ g(x, 0, v) \leq \mu_0 + \mu_1 v^{\gamma_2}, \forall x \in \Omega, \forall v \in \mathbb{R}_+ \end{cases}$$

Then (S) has a sub-supersolution.

Proof: By [3] and (1.12), the equations

$$egin{aligned} &-\Delta_p u = f(x,u,0) ext{ and } \ &-\Delta_g v = g(x,0,v) \end{aligned}$$

have supersolutions  $\bar{u}$  and  $\bar{v}$ . In fact, the monotonicity assumptions on f and g prove that  $[(0, 0), (\bar{u}, \bar{v})]$  is a sub-supersolution of (**S**).

**Theorem 2.** Let  $v \to f(x, u, v)$  be a non-decreasing function and  $u \to g(x, u, v)$  be a non-increasing function.

Assume that there exist constants

 $\mu_0 \ge 0, \, \mu_1 > 0, \, \gamma_2 \in ]0, q-1[$ 

and for any N:

 $\lambda_0 \geq 0, \, \lambda_1 > 0, \, \gamma_1 \in \left]0, p-1\right[$ 

such that:

(1.14) 
$$\begin{cases} f(x,u,N) \leq \lambda_0 + \lambda_1 u^{\gamma_1} \\ g(x,0,v) \leq \mu_0 + \mu_1 v^{\gamma_2} \end{cases}$$

Then (S) has a sub-supersolution.

**Proof:** By [3] there exits  $\bar{v}$  such that:

$$-\Delta_q \bar{v} \ge \mu_0 + \mu_1 \, \bar{v}^{\gamma_2}$$

Let  $N \leq \bar{v}$  and  $\bar{u}$  be such that:

$$-\Delta_p \bar{u} \le \lambda_0 + \lambda_1 \, \bar{u}^{\gamma_1}$$

Then:

$$-\Delta_p \bar{u} \ge f(x, u, N) \ge f(x, \bar{u}, \bar{v}) -\Delta_q \bar{v} \ge g(x, 0, \bar{v})$$

whence the result.

**Theorem 3.** Assume that there exist  $k_i > 0$ , ] = 0, 1 such that:

$$\begin{cases} f(x, u, v) = k_0 v + \varphi(x, u) \\ g(x, u, v) = k_1 u + \psi(x, v) \end{cases}$$

and that there exist  $\mu_1 \ge 0$ ,  $\lambda_1 \ge 0$ ,  $\mu_2 > 0$ ,  $\lambda_2 > 0$ ,  $\gamma_1 \in [1, p - 1]$  and  $\gamma_2 \in [1, q - 1]$  such that:

$$\varphi(x,u) \leq \mu_1 + \mu_2 \, u^{\gamma_1}, \, \psi(x,v) \leq \lambda_1 + \lambda_2 \, v^{\gamma_2}, \, \forall x \in \Omega, \, \forall u, v \in \mathbf{R}_+$$

Then (S) has a sub-supersolution.

Proof: By Remark 2, it is sufficient to show the existence of a supersolution. Let  $M_0 = ||\psi_0||_{L^{\infty}(\Omega)}$ ,  $N_0 = ||\varphi_0||_{L^{\infty}(\Omega)}$  and R be such that  $\Omega \subset B(0, R)$ We seek  $\bar{u}$  and  $\bar{v}$  of the following type:

$$ar{u}(x) = lpha r^{p*} + eta, \, ar{v}(x) = \gamma r^{q*} + \delta, ext{ where } r = |x|, \ lpha < 0, \, eta > 0, \, \gamma < 0, \, \delta > 0.$$

(1.7) and (1.8) are satisfied if:

(1.16) 
$$\begin{cases} \alpha R^{p*} + \beta = M_0 \\ \gamma R^{q+} + \delta = N_0 \end{cases}$$

We want

(1.17) 
$$\begin{cases} -\operatorname{div}(|\nabla \bar{u}|^{p-2} \nabla \bar{u}) = N |\alpha p^*|^{p-1} \ge k_0 \, \bar{v} + \mu_1 + \mu_2 \, \bar{u}^{\gamma_1} \\ -\operatorname{div}(|\nabla \bar{v}|^{q-2} \nabla \bar{v}) = N |\gamma q^*|^{q-1} \ge k_1 \, \bar{u} + \lambda_1 + \lambda_2 \, \bar{v}^{\gamma_2} \end{cases}$$

Set  $\beta = \delta$ . Then, if  $\beta$  is sufficiently large, using the fact that  $\gamma_1 \in [1, p-1[$ and  $\gamma_2 \in [1, q-1[]$ , we can obtain

(1.18) 
$$\begin{cases} \frac{N(p^{\bullet})^{p-1}}{R^{p}} \frac{(\beta - M_{0})^{p-1}}{\mu_{1} + \mu_{2} \beta^{\lambda_{1}}} - 1 - \frac{k_{0}\beta}{\mu_{1} + \mu_{2} \beta^{\gamma_{1}}} \ge 0 \text{ and} \\ \frac{N(q^{\bullet})^{q-1}}{R^{q}} \frac{(\delta - N_{0})^{q-1}}{\lambda_{1} + \lambda_{2} \delta^{\gamma_{2}}} - 1 - \frac{k_{1}\delta}{\lambda_{1} + \lambda_{2} \delta^{\gamma_{2}}} \ge 0 \end{cases}$$

So (1.17) is satisfied and we check  $\alpha$  and  $\gamma$  with (1.16)

## 2. Existence results

Our main result is the following one:

Theorem 4. Let p > 2, q > 2 and  $\varphi_0 \in W_0^{1,p}(\Omega) \cap L^{\infty}(\Omega), \psi_0 \in W_0^{1,q}(\Omega) \cap L^{\infty}(\Omega), \varphi_0 \ge 0$ ,  $\psi_0 \ge 0$  be given.

Suppose that f and g verify (1.1) and (1.2) and that (S) has a sub-supersolution  $|(0, 0), (\bar{u}\bar{v})|$  in  $Q_T$ .

Then (§) has a unique solution (u, v) in  $Q_T$  satisfying:

$$\begin{cases} 0 \le u \le \bar{u} \\ 0 \le v \le \bar{v} \end{cases} \text{ in } Q_T$$

Proof:: By Theorem II.2[3], we can choose  $u_0 \in L^{\infty}(0, T; W^{1,p}(\Omega) \cap L^{\infty}(\Omega))$  $v_0 \in L^{\infty}(0, T; W^{1,q}(\Omega) \cap L^{\infty}(\Omega))$  satisfying  $0 \le u_0 \le \overline{u}$  and  $0 \le v_0 \le \overline{v}$ , such that:

$\int \frac{\partial u_0}{\partial t} - \Delta_p u_o = f(x, u_o, 0)$	in $Q_T$
$\left\langle u_0(x,t)=0\right\rangle$	in $S_T$
$\begin{cases} \frac{\partial u_0}{\partial t} - \Delta_p u_o = f(x, u_o, 0) \\ u_0(x, t) = 0 \\ v_0(x, 0) = \varphi_0(x) \end{cases}$	in $\Omega$
	in $Q_T$
$\begin{cases} \frac{\partial v_0}{\partial t} - \Delta_q v_0 = g(x, 0, v_0) \\ v_0(x, t) = 0 \end{cases}$	in $S_T$
$\left( egin{array}{c} v_0(x,0)=\psi_0(x) \end{array}  ight)$	in $\Omega$

By the existence theorem of Meike ([9], p 1024) we construct two sequences of functions,  $(u_n)$  and  $(v_n)$ , such that:

$$\begin{array}{c} (2.1)\\ (2.2)\\ (2.3)\\ (2.3)\end{array} \begin{cases} \frac{\partial u_{n+1}}{\partial t} - \Delta_p \, u_{n+1} = f(x, u_{n+1}, v_n) & \text{in } Q_T \\ u_{n+1}(x, t) = 0 & \text{in } S_T \\ u_{n+1}(x, 0) = \varphi_0(x) & \text{in } \Omega \end{cases}$$

and

$$\begin{array}{l} (2.4)\\ (2.5)\\ (2.5)\\ (2.6)\end{array} \begin{pmatrix} \frac{\partial v_{n+1}}{\partial t} - \Delta_q \, v_{n+1} = g(x, u_n, v_{n+1}) & \text{in } Q_T\\ v_{n+1}(x, t) = 0 & \text{in } S_T\\ u_{n+1}(x, 0) = \psi_0(x) & \text{in } \Omega \end{array}$$

We need several lemmas to complete the proof of Theorem 4:

Lemma 1. For any  $n \in \mathbb{N}$ , the relations  $0 \le u_n \le \overline{u}$ ,  $0 \le v_n \le \overline{v}$  imply that  $0 \le u_{n+1} \le \overline{u}$  and  $0 \le v_{n+1} \le \overline{v}$ 

Proof of lemma 1: By (1.10), (1.11) and the above assumptions, we have:

(2.7) 
$$\frac{\partial}{\partial t}(u_{n+1}-u) - (\Delta_p u_{n+1} - \Delta_p \bar{u}) \le f(x, u_{n+1}, v_n) - f(x, \bar{u}, v_n)$$

Multiplying (2.7) by  $(u_{n+1} - \bar{u})_+$ , the monoticity of  $\Delta_p$  implies:

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega}(u_{n+1}-\bar{u})_{+}^{2}\leq\lambda\int_{\Omega}(f(x,u_{n+1},v_{n})-f(x,\bar{u},v_{n}))(u_{n+1}-\bar{u})_{+}$$

By the Lipschitz condition (1.2), the initial condition and Gronwall's Lemma, we obtain :  $u_{n+1} \leq \bar{u}$ .

The hypothesis  $f(x, 0, v_n) \ge 0$ , gives  $u_{n+1} \ge 0$ ; similarly, we get  $0 \le v_{n+1} \le \overline{v}$ .

**Lemma 2.** There exists  $C = C(M_1, N_1, T)$  such that:

(2.8) 
$$||u_{n+1}||_{L^{\infty}(Q_T)} \leq C$$

(2.9) 
$$||u_{n+1}||_{L^{\infty}(0,T;W_{\theta}^{1,p})} \leq C$$

$$(2.10) ||\frac{\partial u_{n+1}}{\partial t}||_{L^2(Q_T)} \le C$$

The same estimates hold for  $v_{n+1}$  with p replaced by q.

Proof of Lemma 2: By lemma 1, for any  $n \in N, u_n$  and  $v_n$  are bounded; whence (2.8). The properties of the functions f and g, then imply that  $f(x, u_{n+1}, v_n)$  is bounded.

We therefore obtain:

$$\int_{\Omega} f(x, u_{n+1}, v_n) \frac{\partial u_{n+1}}{\partial t} \leq \frac{1}{2} \int_{\Omega} (f(x, u_{n+1}, v_n))^2 + \frac{1}{2} \int_{\Omega} \left(\frac{\partial u_{n+1}}{\partial t}\right)^2 \leq C_0 + \frac{1}{2} \int_{\Omega} \left(\frac{\partial u_{n+1}}{\partial t}\right)^2 dx$$

Multiplying (2.1) by  $\frac{\partial u_{n+1}}{\partial t}$ , we get:

$$\frac{1}{2}\int_0^T\int_\Omega \left(\frac{\partial u_{n+1}}{\partial t}\right)^2\,dx\,dt+\frac{1}{p}\int_\Omega |\nabla u_{n+1}(\cdot,T)|^p\,dx\leq C_0T+\frac{1}{p}\int_\Omega |\nabla \varphi_0|^p\,dx$$

It is the same for  $v_{n+1}$ .

Proof of theorem 4: By (2.8), (2.9), (2.10), there is a subsequence  $(u_n, v_n)$  with the following properties:

 $u_n$  converges to u in the weak \* sense in  $L^{\infty}(0, T; W_0^{1,p}(\Omega) \cap L^{\infty}(\Omega))$  and  $u_n$  converges weakly in  $L^p(0, T; W_0^{1,p}(\Omega)); u_n$  is such that  $\frac{\partial u_n}{\partial t}$  converges to  $\frac{\partial u}{\partial t}$  in weak  $L^2(Q_T)$ ; the same holds also for  $v_n$  with p replaced by q.

By standard monotonicity argument [8],  $\Delta_p u_{n+1}$ , converges to  $\Delta_p u$  in weak  $L^{p*}(0,T; W^{-1,p*}(\Omega))$ ,  $\Delta_q v_{n+1}$  converges to  $\Delta_q v$  in weak  $L^{q*}(0,T; W^{-1,q*}(\Omega))$ .  $u_n$  converges almost everywhere to u and  $v_n$  converges almost everywhere to v.

By Lebesgue's theorem:

$$f(\cdot, u_{n+1}, v_n)$$
 converges to  $f(\cdot, u, v)$   
 $g(\cdot, u_n, v_{n+1})$  converges to  $g(\cdot, u, v)$ 

whence (u, v) is a solution of (S) in  $Q_T$ .

Applying lemma 1, we have  $0 \le u \le \overline{u}, 0 \le v \le \overline{v}$ .

**Remark 3:** Uniqueness follows from the Lipschitz condition on f and g.

#### 3. Asymptotic behaviour

Hereafter, we assume that there exist positive constants  $\lambda > 0$  and  $\mu > 0$ and a function H from  $\mathbb{R}^{N+2}$  to  $\mathbb{R}$  such that

$$(1.1') - (1.2') \begin{cases} f = \lambda \frac{\partial H}{\partial u}, & g = \mu \frac{\partial H}{\partial v} \\ f \text{ and } g \text{ satisfy (1.1) and (1.2)} \end{cases}$$

For a solution (u, v) of (S), we define the  $\omega$ -limit set by:

 $\omega(\varphi_0,\psi_0) = \{w = (w_1,w_2) : w_1 \in W^{1,p}_0(\Omega) \cap L^{\infty}(\Omega), w_2 \in w^{1,q}_0(\Omega) \cap L^{\infty}(\Omega) |$ 

$$\exists t_n \to +\infty : u(\cdot, t_n) \to w_1 \text{ in } W_0^{1,p}(\Omega)$$
$$v(\cdot, t_n) \to w_2 \text{ in } W_0^{1,q}(\Omega) \}$$

Let  $\mathcal{E}$  be the set of non negative solutions  $w = (w_1, w_2)$  of the elliptic problem:

$$\begin{cases} -\Delta_p w_1 = \lambda \frac{\partial H}{\partial u}(x, w_1, w_2) & \text{in } \Omega \\ -\Delta_q w_2 = \mu \frac{\partial H}{\partial v}(x, w_1, w_2) & \text{in } \Omega \\ w_1 = w_2 = 0 \end{cases}$$

Our main result is the following:

**Theorem 5.** Let p > 2, q > 2 and  $\varphi_0 \in W_0^{1,p}(\Omega) \cap L^{\infty}(\Omega)$ ,  $\psi_0 \in W_0^{1,q}(\Omega) \cap L^{\infty}(\Omega)$ ,  $\varphi_0 \ge 0$ ,  $\psi_0 \ge 0$ . Suppose that H satisfies (1.1'), (1.2') and that (S) has a sub-supersolution. Then  $\omega(\varphi_0, \psi_0) \neq \phi$  and  $\omega(\varphi_0, \psi_0) \subset \mathcal{E}$ .

To prove this Theorem, we need the following lemmas:

**Lemma 3.** Under the assumptions of Theorem 5, there exists a constant  $C = C(M_1, N_1)$  such that for any T > 0:

- (3.1)  $||u||_{L^{\infty}(Q_T)} \leq C < +\infty, ||v||_{L^{\infty}(Q_T)} \leq C < +\infty$
- (3.2)  $||u||_{L^{\infty}(0,T;W_0^{1,p}(\Omega))} \le C < +\infty, ||v||_{L^{\infty}(0,T;W_0^{1,q}(\Omega))} \le C < +\infty$
- (3.3)  $\left\|\frac{\partial u}{\partial t}\right\|_{L^{2}(\mathcal{O}_{T})} \leq C < +\infty, \left\|\frac{\partial u}{\partial t}\right\|_{L^{2}(\mathcal{O}_{T})} \leq C < +\infty$

Proof of lemma 3: By Theorem 4 we have (3.1).

Multiplying the first equation (1.6) by  $\frac{1}{\lambda} \frac{\partial u}{\partial t}$  and the second equation by  $\frac{1}{\mu} \frac{\partial v}{\partial t}$ , we obtain:

$$(3.4) \qquad \frac{1}{\lambda} \int_{Q_T} \left(\frac{\partial u}{\partial t}\right)^2 dx \, dt + \frac{1}{\mu} \int_{Q_T} \left(\frac{\partial v}{\partial t}\right)^2 dx \, dt \\ + \frac{1}{\lambda p} \int_{\Omega} |\nabla u(\cdot, T)|^p \, dx + \frac{1}{\mu q} \int_{\Omega} |\nabla v(\cdot, T)|^q \, dx \\ = \int_{Q_T} \left(\frac{\partial H}{\partial u}(\cdot, u, v) \frac{\partial u}{\partial t} + \frac{\partial H}{\partial v}(\cdot, u, v) \frac{\partial v}{\partial t}\right) \, dx \, dt \\ + \frac{1}{\lambda p} \int_{\Omega} |\nabla \varphi_0|^p \, dx + \frac{1}{\mu q} \int_{\Omega} |\nabla \psi_0|^q \, dx \\ = \int_{\Omega} (H(\cdot, u(T), v(T)) - H(\cdot, \varphi_0, \psi_0)) \, dx \\ + \frac{1}{\lambda p} \int_{\Omega} |\nabla \varphi_0|^p \, dx + \frac{1}{\mu q} \int_{\Omega} |\nabla \psi_0|^q$$

H is continuous and (u, v) is bounded; we then obtain:

$$\frac{1}{\lambda} \int_{Q_T} \left(\frac{\partial u}{\partial t}\right)^2 + \frac{1}{\mu} \int_{Q_T} \left(\frac{\partial v}{\partial t}\right)^2 + \frac{1}{\lambda p} \int_{\Omega} |\nabla u(\cdot, T)|^p \, dx + \frac{1}{\mu q} \int_{\Omega} |\nabla v(\cdot, T)|^q \, dx$$

whence (3.2) and (3.3).

**Lemma 4.** Let  $t_0 \in ]0,1[$ . Under the assumptions of Theorem 5, there exists  $C = C(t_0) > 0$  such that for any  $T \ge t_0$ :

(3.5) 
$$\left\|\frac{\partial u}{\partial t}\right\|_{L^{\infty}(t_{0},+\infty;L^{2}(\Omega))} \leq C, \left\|\frac{\partial v}{\partial t}\right\|_{L^{\infty}(t_{0},+\infty;L^{2}(\Omega))} \leq C$$

(3.6) 
$$\left\|\frac{\partial}{\partial t}F(\nabla u)\right\|_{L^{2}(t_{0},T;L^{p^{*}}(\Omega))} \leq C, \left\|\frac{\partial}{\partial t}G(\nabla v)\right\|_{L^{2}(t_{0},T;L^{q^{*}}(\Omega))} \leq C$$

Proof of Lemma 4: Let  $E_p(\nabla u) = |\nabla u|^{\frac{p-2}{2}} \nabla u$ . Calculations, [cf.3], give:

(3.7)  $|\frac{\partial}{\partial t} E_p(\nabla u)|^2 \le \frac{p+2}{4} \frac{\partial}{\partial t} F(\nabla u) \cdot \frac{\partial}{\partial t} \nabla u$ 

$$(3.8) \qquad \qquad |\frac{\partial}{\partial t} F(\nabla u)| \le \left(\frac{4p}{p+2}\right)^{1/2} |\nabla u|^{\frac{p-2}{2}} \left|\frac{\partial}{\partial t} E_p(\nabla u)\right|$$

Similary for  $E_q(\nabla v) = |\nabla v|^{\frac{q-2}{2}} \nabla v$ .

By formal derivation of the first equation of (1.6), we get

(3.9) 
$$\frac{\partial^2 u}{\partial t^2} - \operatorname{div}\left(\frac{\partial}{\partial t}F(\nabla u)\right) = \lambda \frac{\partial}{\partial t}\left(\frac{\partial H}{\partial u}\right) = \\ = \lambda \frac{\partial^2 H}{\partial u^2} \cdot \frac{\partial u}{\partial t} + \lambda \frac{\partial^2 H}{\partial u \partial v} \cdot \frac{\partial v}{\partial t}$$

Multiplying (3.9) by  $\frac{\partial u}{\partial t}$ , we get with (3.7) and (1.1') (3.10)  $\frac{1}{2}\frac{\partial}{\partial t}\int_{\Omega}\left(\frac{\partial u}{\partial t}\right)^{2} + \frac{4}{p+2}\int_{\Omega}\left|\frac{\partial}{\partial t}E_{p}(\nabla u)\right|^{2}dx \leq k_{0}\int_{\Omega}\left(\frac{\partial u}{\partial t}\right)^{2} + k_{1}\int_{\Omega}\left(\frac{\partial v}{\partial t}\right)^{2}$ Similarly, we have:

Similarly, we have: (3.11)

$$\frac{1}{2}\frac{\partial}{\partial t}\int_{\Omega}\left(\frac{\partial v}{\partial t}\right)^{2} + \frac{4}{q+2}\int_{\Omega}\left|\frac{\partial}{\partial t}E_{q}(\nabla v)\right|^{2}\,dx \leq k_{0}^{'}\int_{\Omega}\left(\frac{\partial u}{\partial t}\right)^{2} + k_{1}^{'}\int_{\Omega}\left(\frac{\partial v}{\partial t}\right)^{2}$$

By (3.3), there exists  $t_1 \in ]0, t_0[$  such that:

(3.12) 
$$\left\|\frac{\partial u}{\partial t}(\cdot,t_1)\right\|_{L^2(\Omega)}^2 + \left\|\frac{\partial v}{\partial t}(\cdot,t_1)\right\|_{L^2(\Omega)}^2 = \int_0^{t_0} \left(\left\|\frac{\partial u}{\partial t}(\cdot,t)\right\|_{L^2(\Omega)}^2 + \left\|\frac{\partial v}{\partial t}(\cdot,t)\right\|_{L^2(\Omega)}^2\right) dt \le C < +\infty$$

and integrating (3.10) + (3.11) on  $(t_1, T)$ , we obtain with (3.3):

(3.13)

$$\begin{split} &\frac{1}{2} \left\| \frac{\partial u}{\partial t}(\cdot,T) \right\|_{L^{2}(\Omega)}^{2} + \frac{1}{2} \left\| \frac{\partial v}{\partial t}(\cdot,T) \right\|_{L^{2}(\Omega)}^{2} \\ &+ \frac{4}{p+2} \int_{t_{0}}^{T} \int_{\Omega} \left| \frac{\partial}{\partial t} E_{p}(\nabla u) \right|^{2} dx \, dt + \frac{4}{q+2} \int_{t_{0}}^{T} \int_{\Omega} \left| \frac{\partial}{\partial t} E_{q}(\nabla v) \right|^{2} dx \, dt \\ &\leq K \left[ \int_{t_{1}}^{T} \int_{\Omega} \left( \left( \frac{\partial u}{\partial t} \right)^{2} + \left( \frac{\partial v}{\partial t} \right)^{2} \right) \right] dx \, dt + \frac{1}{2} \left\| \frac{\partial u}{\partial t}(\cdot,t_{1}) \right\|_{L^{2}(\Omega)}^{2} \\ &+ \frac{1}{2} \left\| \frac{\partial v}{\partial t}(\cdot,t_{1}) \right\|_{L^{2}(\Omega)}^{2} \leq C < +\infty \end{split}$$

By (3.8) and Holder's inequality, we obtain with (3.13) and (3.2): (3.14)

$$\begin{split} & \left\|\frac{\partial u}{\partial t}(\cdot,T)\right\|_{L^{2}(\Omega)}^{2} + \left\|\frac{\partial v}{\partial t}(\cdot,T)\right\|_{L^{2}(\Omega)}^{2} + \left\|\frac{\partial}{\partial t}F(\nabla u)\right\|_{L^{2}(t_{0},T;L^{p^{*}}(\Omega))}^{2} \\ & + \left\|\frac{\partial}{\partial t}G(\nabla v)\right\|_{L^{2}(t_{0},T;L^{p^{*}}(\Omega))}^{2} \leq C \end{split}$$

(3.14) gives the estimates (3.5) and (3.6).

This formal proof of (3.14) can be made rigorous by means of the finite dimensional problems associated with (S).

The details are in [4, p 35] and are omitted.

Theorem 6. Let  $p > 2, q > 2, \varphi_0 \in W_0^{1,p}(\Omega) \cap L^{\infty}(\Omega), \psi_0 \in W_0^{1,q}(\Omega) \cap$  $L^{\infty}(\Omega)$ , and  $\varphi_0 \geq 0$ ,  $\psi_0 \geq 0$ .

Suppose that H satisfies (1.1'), (1.2') and that (S) has a sub-supersolution  $[(0,0), (\bar{u}, \bar{v})]$ . Then, for any  $t_0 \in ]0, 1[$ , the solution (u, v) of (S) satisifies the following regularizing estimates:

- $u \in L^{\infty}\left(t_0, +\infty; B^{1+1/(p-1)^2, p}_{\infty}(\Omega)\right)$ (3.15)
- $\frac{\partial u}{\partial t} \in L^{2}(t_{0}, +\infty; L^{2}(\Omega)) \cap L^{\infty}(t_{0}, +\infty; L^{2}(\Omega))$  $\frac{\partial u}{\partial t} F(\nabla u) \in L^{2}(t_{0}, +\infty; L^{p*}(\Omega)),$ (3.16)

(3.17)

where  $B^{1+1/(p-1)^2,p}_{\infty}(\Omega)$  is a BESOV space defined by the real interpolation method (cf./1], /13/).

The same estimates hold for v provided p and F are replaced by q and Grespectively.

Proof of Theorem 6: By (3.3), (3.5) and (3.6) we have (3.16) and (3.17), whence  $\frac{\partial u}{\partial t} \in L^{\infty}(t_0, +\infty; L^{p*}(\Omega))$  and  $\frac{\partial v}{\partial t} \in L^{\infty}(t_0, +\infty; L^{q*}(\Omega))$ .

By SIMON'S regularity results [13], we have:

$$\left\| u(\cdot,t) \right\|_{B^{1+1/p-1)^2,p}_{\infty}(\Omega)} \leq C \left\| \frac{\partial H}{\partial u}(\cdot,u,v) - \frac{\partial u}{\partial t}(\cdot,t) \right\|_{L^{p_*}(\Omega)} + C'$$

whence (3.15). The proof is the same for v.

Proof of theorem 5:

a)  $\omega(\varphi_0, \psi_0) \neq \phi$  because  $B^{1+1/(r-1)^2, r}_{\infty}(\Omega)$  is Compactly inbedded in  $W^{1, r}(\Omega)$ for r = p and q[1]. By Theorem 6, letting

$$w_1 = \lim_{n \to +\infty} u(\cdot, t_n), w_2 = \lim_{n \to +\infty} v(\cdot, t_n)$$
 and

 $w = (w_1, w_2) \in \omega(\varphi_0, \psi_0)$ , we get that  $w = (w, w_2) \in \mathcal{E}$ . The proof is analogous to EL HACHIMI and DE THELIN [3] and is omitted

## 4. Examples

Example 1. Let  $H(x, u, v) = K(x)uv + \lambda u^{\gamma_1+1} + \mu v^{\gamma_2+1}$ , where  $K \in$  $C(\Omega), K(x) > 0, \lambda > 0, \mu > 0, \gamma_1 \in [1, p-1] \text{ and } \gamma_2 \in [1, q-1].$ 

Then we can apply Theorem 3, 4, 5, 6.

**Example 2.** Let  $H(x, u, v) = -u^m v^n + \lambda^{\gamma_1+1} + \mu v^{\gamma_2+1}$ , where  $m \ge 1, n \ge 1$  $1, \lambda > 0, \mu > 0, \gamma_1 \in [1, p-1]$  and  $\gamma_2 \in [1, q-1]$ . Then we can apply theorems 1, 4, 5, 6.

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