

Supersymmetric Backgrounds From Generalized Calabi-Yau Manifolds

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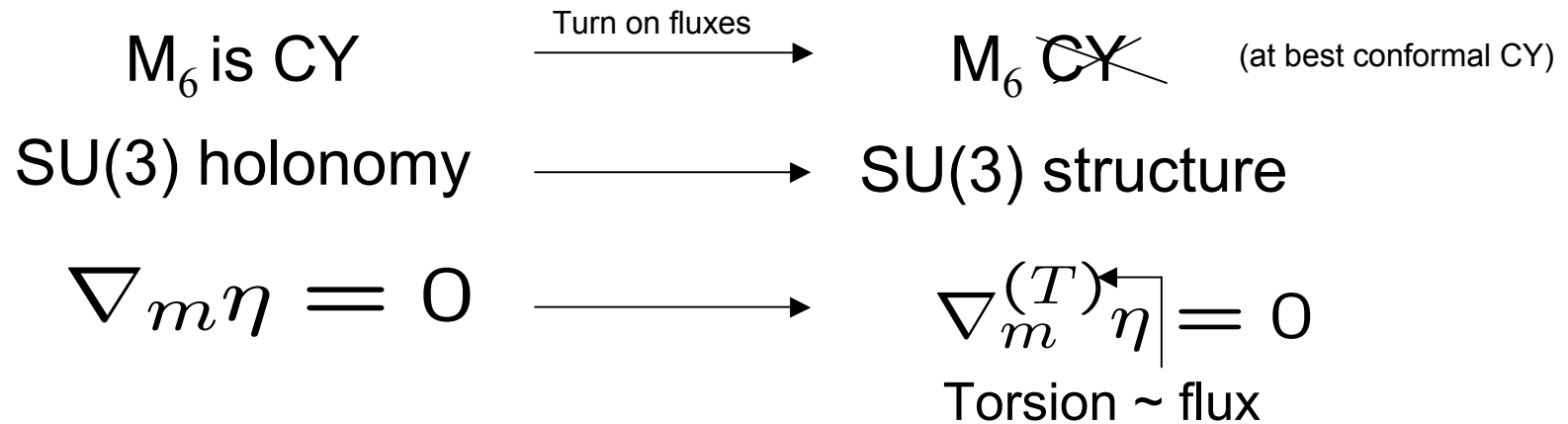
hep-th/0406137

Strings 2004 - Paris

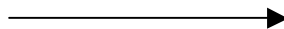
Fluxes \longleftrightarrow Geometry

Type II sugra on $M_{10} = M_4 \times M_6$

Minimal supersymmetry



CY



**Generalized
CY**

Outline

- Basic notions about $SU(3)$ structure
- Torsion versus fluxes
- Generalized complex geometry / Generalized CY
- Generalized CY from supersymmetry equations
- Conclusions / open questions

$$M_{10} = M_4 \times M_6$$

Basic notions about SU(3) structure

- No flux) M_6 is CY \rightarrow SU(3) holonomy

- 9 SU(3) invt. spinor η

$$\begin{cases} \eta^\dagger \gamma_{mn} \gamma \eta = i w_{mn} \\ \eta^\dagger \gamma_{mnp} (1 + \gamma) \eta = i \Omega_{mnp} \end{cases} \quad \begin{cases} w \wedge \Omega = 0 \\ i\Omega \wedge \bar{\Omega} = w^3 \end{cases}$$

- η is cov. constant: $r_m \eta = 0$

$$\begin{cases} dw = 0 \\ d\Omega = 0 \end{cases} \quad \epsilon_{10}^{1,2} = \theta^{1,2} \otimes \eta \Rightarrow \mathcal{N} = 2$$

- Turn on flux) back-reaction

- $M_{10} = M_4 \times_w M_6$ $[ds_{10}^2 = e^{2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu + ds_6^2(y)]$

- $\mathcal{N} = 2$! $\mathcal{N} = 1$ (relation between θ^1) θ^2)

- M_6 acquires torsion \rightarrow SU(3) structure

Gauntlett, Martelli, Pakis, Waldram '02

- 9 SU(3) invt. spinor η

$$\begin{cases} \eta^\dagger \gamma_{mn} \gamma \eta = i w_{mn} \\ \eta^\dagger \gamma_{mnp} (1 + \gamma) \eta = i \Omega_{mnp} \end{cases}$$

- $r_m^{(T)} \eta = 0$

$$\begin{cases} r_m^{(T)} w = 0 \text{) } dw \neq 0 / T \\ r_m^{(T)} \Omega = 0 \text{) } d\Omega \neq 0 / T \end{cases}$$

Torsion versus fluxes

Torsion: $dw = \text{Im} (W_1 \Omega) + W_4 \mathcal{A}E w + W_3 \bar{\mathcal{A}}E w$
 SU(3) reps $\quad \quad \quad 1 \oplus 1 \quad \quad 3 \oplus \bar{3} \quad \quad 6 \oplus \bar{6}$

$d\Omega = W_1 w^2 + W_5 \mathcal{A}E \Omega + W_2 \bar{\mathcal{A}}E w$
 $\quad \quad \quad 1 \oplus 1 \quad \quad 3 \oplus \bar{3} \quad \quad 8 \oplus 8$

	$1 \odot 1$	$3 \odot \bar{3}$	$6 \odot \bar{6}$	$8 \odot 8$
Torsion	1 (W_1)	2 (W_4, W_5)	1 (W_3)	1 (W_2)
H_3	1	1	1	0
IIA: F_{2n}	2 (F_0, F_2, F_4)	2 (F_2, F_4)	0	1 (F_2, F_4)
IIB: F_{2n+1}	1 (F_3)	3 (F_1, F_3, F_5)	1 (F_3)	0

In IIB $W_2 = 0$ (~~integrability of complex structure~~)



In IIA $W_3 \gg H^{(6)}$ (~~symplectic geometry~~)

If also $W_1=0 \rightarrow$ IIB: $d\Omega = W_5 \mathcal{A}E \Omega$
 (true in all susy vacua)

M_6 is complex

IIA: $d w = W_4 \mathcal{A}E w + H^{(6)}$

M_6 is "twisted symplectic"

Is there a mathematical construction that ~~contains~~ extends complex and symplectic geometry?

Generalized complex geometry

Hitchin
Gualtieri

Huybrechts
Kapustin
Lindström
Zabzine

• Usual differential geometry \rightarrow T tangent bundle
sections vector fields X | T^* cotangent bundle
sections are 1-forms ζ

• Want differential geometry on $T \circledast T^*$ sections are $X + \zeta$

Natural metric I on $T \circledast T^*$: $(X + \zeta, X + \zeta) = i_X \zeta$ $I = \begin{pmatrix} 0 & 1_d \\ 1_d & 0 \end{pmatrix}$

On $T \circledast T^*$ define Generalized Almost Complex Structure (GACS) $J: T \circledast T^* \rightarrow T \circledast T^*$

$$\begin{aligned} J^2 &= -1_{2d} \\ J^t | J &= I \end{aligned}$$

	T	$T \circledast T^*$	project down
GACS	J s.t. $J^2 = -1_d$	J s.t. $J^2 = -1_{2d}$	$T \circledast T^* \rightarrow T$
Projectors	$\pi_\zeta = (1_d \ \S \ iJ)$	$\Pi_\zeta = (1_{2d} \ \S \ iJ)$	$X + \zeta \rightarrow X$
Integrability	$\pi_\cdot [\pi_\zeta X, \pi_\zeta Y] = 0$	$\Pi_\cdot [\Pi_\zeta (X + \zeta), \Pi_\zeta (Y + \zeta)]_C = 0$	$[\cdot , \cdot]_C \rightarrow [\cdot , \cdot]$
			$J \rightarrow J ?$

$$[X + \zeta, Y + \eta]_C = [X, Y] + L_X \eta - L_Y \zeta - \frac{1}{2} d(i_X \eta - i_Y \zeta)$$

when $\zeta = \eta = 0$) $[\cdot , \cdot]_C = [\cdot , \cdot]$

How is complex geometry embedded in generalized complex geometry?
Aside: GCG naturally incorporates B-field.

$$e^B (X + \zeta) = X + \zeta + i_X B \rightarrow [e^B (X + \zeta), e^B (Y + \eta)]_C = e^B [X + \zeta, Y + \eta]_C$$

automorphisms of $[\cdot , \cdot]_C$

_ A GACS has the form

$$J = \begin{pmatrix} J & P \\ L & K \end{pmatrix}$$

J: TM ! TM	J_{mn}^m
P: T*M ! TM	P^{mn}
L: TM ! T*M	L_{mn}
K: T*M ! T*M	K_m^n

Demanding

$$J^2 = -1_{2d}$$

$$J^t J = I$$

Get conditions

$$K = -J^t$$

$$P^t = -P$$

$$L^t = -L$$

$$\rightarrow J = \begin{pmatrix} J & P \\ L & -J^t \end{pmatrix}$$

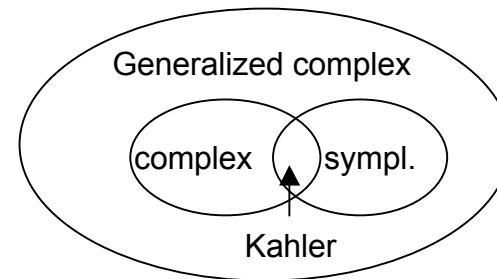
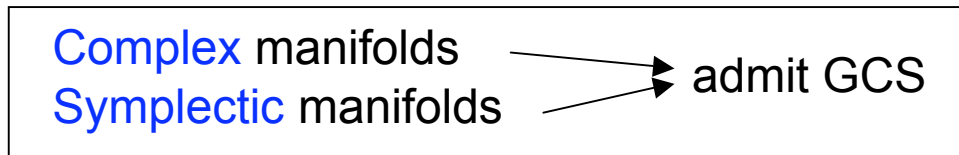
_ It is easy to guess how ACS _ GACS :

$$J_1 = \begin{pmatrix} J & 0 \\ 0 & -J^t \end{pmatrix}$$

Integrability of GACS J_1) is J integrable (it is a CS in T)) M is **complex**

_ But GACS have more...

Consider $J_2 = \begin{pmatrix} 0 & -w^{-1} \\ w & 0 \end{pmatrix}$ Integrability of J_2) $dw=0$. If w is non-degenerate) M is **symplectic**



Complex: locally equivalent to $C^{d/2}$

Symplectic: locally equivalent to (R^d, w) ; $w = dx^1 \wedge dx^2 + \dots + dx^{d-1} \wedge dx^d$

Generalized complex: locally equivalent to $C^k _ (R^{d-2k}, w)$ k : rank. $k=0$ for **symplectic**
 $k=d/2$ for **complex**

GACS ↔ Pure spinors of Clifford (d,d)

	T	T ⊗ T*
Algebra	Clifford 6 $\{\gamma^m, \gamma^n\} = g^{mn}$	Clifford (6,6) $\{\gamma^m, \gamma^n\} = \{\gamma_m, \gamma_n\} = 0, \{\gamma^m, \gamma_n\} = \delta^m_n$
Representation in terms of forms	$\gamma^i = dz^i \wedge \epsilon$ and $\gamma_- = g^{-j} \iota_{\partial_j}$	$\gamma^m = dx^m \wedge \epsilon$ and $\gamma_m = \iota_{\partial_m}$
Clifford vacuum	Ω	Ω
Pure spinor	$\gamma^i \Omega = 0$	$\gamma^i \Omega = \gamma_- \Omega = 0$
Basis	$\Omega : (3,0)$ $\gamma_- \Omega : (2,0)$ $\gamma^j \Omega : (1,0)$ $\overline{\gamma}^{jk} \Omega : (0,0)$	$\Omega : (3,0)$ $\Omega : (3,0)$ $\gamma_- \Omega : (3,1)$ $\gamma_i \Omega : (2,0)$ $\gamma^j \Omega : (3,2)$ $\gamma_{ij} \Omega : (1,0)$ $\gamma^{ijk} \Omega : (3,3)$ $\gamma_{ijk} \Omega : (0,0)$
	Spinors \$ (p,0) forms	Spinors \$ (p,q) forms

$$e^{iw} = 1 + iw - \frac{1}{2}w^2 - i\frac{1}{3!}w^3$$

$$(\gamma_m + i W_{mn} \gamma^n) e^{iw} = 0$$

- On a manifold with SU(3) structure in T → pure spinors of Clifford(6,6) $\Omega; e^{iw}$
- 1-1 correspondence between pure spinors and GACS J

$$\begin{array}{ccc} \Omega & & e^{iw} \\ \downarrow & & \downarrow \\ J_1 = \begin{pmatrix} J & 0 \\ 0 & -J^t \end{pmatrix} & & J_2 = \begin{pmatrix} 0 & -w^{-1} \\ w & 0 \end{pmatrix} \end{array}$$

Integrability of J, η and ξ s.t

$$d\varphi = (\eta + \xi \wedge) \varphi$$

has one GCS



Hitchin

Generalized Calabi-Yau → manifold that has one closed pure spinor ($d\varphi = 0$)

Supersymmetry equations written in terms of pure spinors

$$\delta \psi_m: \begin{array}{l} IIA : D_m^H \eta_+ + e^\phi (\bar{\Omega} F^A)_m \eta_+ + e^\phi \left[(F_m^A e^{i\omega})_n + (F^A e^{i\omega})_0 g_{mn} + (F^A e^{i\omega})_{mn} \right] \gamma^n \eta_- = 0 \\ \quad \quad \quad \updownarrow \\ IIB : D_m^H \eta_+ + e^\phi (e^{-i\omega} F^B)_m \eta_+ + e^\phi \left[(F_m^B \Omega)_n + (F^B \Omega)_0 g_{mn} + (F^B \Omega)_{mn} \right] \gamma^n \eta_- = 0 \end{array}$$

NSNS sector: $D_m^H \eta \equiv \left(\nabla_m + \frac{1}{8} H_{mnp} \gamma^{np} \right) \eta$ We are not in CY) $0 \neq \nabla_m \eta \sim W \eta$

$$D_m^H \eta_+ = (W_4 + W_5 + iH^{(3)})_m \eta_+ + \left[(W_1 + iH^{(1)}) g_{mn} + (W_3 + iH^{(6)} + W_2)_{mn} \right] \gamma^n \eta_-$$

RR fluxes: $F^A = F_0 + F_2 + F_4$ $F^B = F_1 + F_3 + F_5$

$$(F\Omega)_{mn} \equiv (F \lrcorner \Omega)_{mn} = \left(F_{i_1 \dots i_k} \Omega_{abc} \gamma^{i_1 \dots i_k} \gamma^{abc} \right)_{mn}$$

IIA \$ IIB

F^A \$ F^B

$e^{i\omega}$ \$ Ω

Take coefficient of term with 2 γ 's

$$\text{Ex: } (F_1 \Omega)_{mn} = \frac{1}{2} F^a \Omega_{amn}$$

exchange of two pure spinors -- action of mirror symmetry

Using $\eta_+ \otimes \eta_+^\dagger = e^{i\psi}$
 $\eta_+ \otimes \eta_-^\dagger = i \not{\Sigma}$ we can go further and find equations for pure spinors

IIA

$$e^{-f} d(e^f e^{i\omega}) = H \bullet e^{i\omega}$$

$$e^{-g} d(e^g \Omega) = H \bullet \Omega + F_A e^{i\omega}$$



$e^f e^{i\omega}$ is "twisted" closed

M is twisted symplectic

Sugra "twisting":

$$\left[d - H_{mnp} \left(dx^m \wedge dx^n \wedge dx^p - \frac{1}{3} \iota^m \iota^n \iota^p \right) \right] (e^f \varphi) = 0$$

Susy vacua are all twisted generalized Calabi-Yau's !

M twisted symplectic

$$d^H \Omega \quad \uparrow \quad \uparrow F_A$$

Caveat: Hitchin considered twisting by H

Hitchin's twisting

$$[d + H \wedge] \varphi = 0$$

Associated GACS integrable: with $[\Pi_S, \Pi_S]_C = 0$

$$[X + \zeta, Y + \eta]_H = F_B [X + \zeta, Y + \eta]_C + \iota_X \iota_Y H$$

Twisting by H in supergravity equations does not work a la Hitchin

Strominger '86
(Maldacena - Nuñez)

IIB

$$e^{-f} d(e^f e^{i\omega}) = H \bullet e^{i\omega} + F_B e^{i\omega}$$

$$e^{-g} d(e^g \Omega) = H \bullet \Omega$$



$e^g \Omega$ is "twisted" closed

M is complex

(no (2,2) piece in $H \bullet \Omega$)

$\varphi = e^{i\omega}$ in IIA

$\varphi = \Omega$ in IIB

$$e^f \sim e^g \sim e^\phi$$

mirror symmetry

Conclusions / Open questions

- Type II supersymmetric vacua are twisted generalized Calabi-Yau's
- What is the meaning of supergravity twisting by H?
- How to compactify on these manifolds? (evade no-go theorems)
- Moduli spaces?

Generalized complex geometry is a nice tool for a systematic description of flux backgrounds.

But maybe strings see $SO(d,d)$ structure of $T \circledast T^*$ and we will learn more...

