# SYMMETRIES IN THE FOURTH PAINLEVÉ EQUATION AND OKAMOTO POLYNOMIALS 

MASATOSHI NOUMI AND YASUHIKO YAMADA


#### Abstract

The fourth Painlevé equation $P_{I V}$ is known to have symmetry of the affine Weyl group of type $A_{2}^{(1)}$ with respect to the Bäcklund transformations. We introduce a new representation of $P_{I V}$, called the symmetric form, by taking the three fundamental invariant divisors as the dependent variables. A complete description of the symmetry of $P_{I V}$ is given in terms of this representation. Through the symmetric form, it turns out that $P_{I V}$ is obtained as a similarity reduction of the 3-reduced modified KP hierarchy. It is proved in particular that the special polynomials for rational solutions $P_{I V}$, called Okamoto polynomials, are expressible in terms of the 3-reduced Schur functions.


It is known by K. Okamoto [7] that the fourth Painlevé equation has symmetries under the affine Weyl group of type $A_{2}^{(1)}$. In this paper we propose a new representation of the fourth Painleve equation in which the $A_{2}^{(1)}$-symmetries become clearly visible. By means of this representation, we clarify the internal relation between the fourth Painlevé equation and the modified KP hierarchy. We obtain in particular a complete description of the rational solutions of the fourth Painlevé equation in terms of Schur functions. This implies that the so-called Okamoto polynomials, which arise from the $\tau$-functions for rational solutions, are in fact expressible by the 3 reduced Schur functions. ${ }^{1}$

## §1. A symmetric form of the fourth Painlevé equation

The fourth Painlevé equation $\mathrm{P}_{\mathrm{IV}}$ is the following second order ordinary

[^0]differential equation
\[

$$
\begin{equation*}
y^{\prime \prime}=\frac{1}{2 y}\left(y^{\prime}\right)^{2}+\frac{3}{2} y^{3}+4 t y^{2}+2\left(t^{2}-a\right) y+\frac{b}{y} \tag{1.1}
\end{equation*}
$$

\]

for the unknown function $y=y(t)$, where ${ }^{\prime}=d / d t$ and $a, b \in \mathbb{C}$ are parameters. It is known by K. Okamoto [7] that equation 1.1 is represented as the following system for the two unknown functions $q=y$ and $p$ :

$$
\begin{align*}
& q^{\prime}=q(2 p-q-2 t)-2\left(v_{1}-v_{2}\right) \\
& p^{\prime}=p(2 q-p+2 t)+2\left(v_{2}-v_{3}\right) \tag{1.2}
\end{align*}
$$

This equation, called $\mathrm{H}_{\mathrm{IV}}$, is in fact a Hamiltonian system

$$
\begin{equation*}
q^{\prime}=\frac{\partial H}{\partial p}, \quad p^{\prime}=-\frac{\partial H}{\partial q} \tag{1.3}
\end{equation*}
$$

with polynomial Hamiltonian

$$
\begin{equation*}
H=q p^{2}-q^{2} p-2 t p q-2\left(v_{1}-v_{2}\right) p-2\left(v_{2}-v_{3}\right) q \tag{1.4}
\end{equation*}
$$

The parameters $\mathbf{v}=\left(v_{1}, v_{2}, v_{3}\right),\left(v_{1}+v_{2}+v_{3}=0\right)$ in 1.2 and $(a, b)$ in 1.1 are related through the formulas

$$
\begin{equation*}
a=1+3 v_{3}, \quad b=-2\left(v_{1}-v_{2}\right)^{2} \tag{1.5}
\end{equation*}
$$

The equivalence between 1.1 and 1.2 can be checked directly, but it requires a tedious calculation. (This calculation is fairly simplified by the "symmetric" representation which we will propose in this paper. See the proof of Theorem 1.1 below.)

It is clearly seen from 1.2 that, if $v_{1}-v_{2}=0$ or $v_{2}-v_{3}=0$, the Hamiltonian system $\mathrm{H}_{\mathrm{IV}}$ has classical solutions such that $q=0$ or $p=0$. In these cases, equation 1.2 is reduced to the Riccati equations $p^{\prime}=-p^{2}+$ $2 t p+2\left(v_{2}-v_{3}\right)$ and $q^{\prime}=-q^{2}-2 t-2\left(v_{1}-v_{2}\right)$ respectively, and they are furthermore linearized to Hermite-Weber equations. In this sense, the Hamiltonian system $\mathrm{H}_{\mathrm{IV}}, 1.2$, has invariant divisors $q=0$ and $p=0$ along the lines $v_{1}-v_{2}=0$ and $v_{2}-v_{3}=0$, respectively. It should be noted that equation 1.2 has one more typical invariant divisor $q-p+2 t=0$ along the line $v_{1}-v_{3}=1$. In fact equation 1.2 implies

$$
\begin{equation*}
(q-p+2 t)^{\prime}=-(q-p-2 t)(q+p)+2\left(1-v_{1}+v_{3}\right) \tag{1.6}
\end{equation*}
$$

It is known by [6] that these three polynomials $q, p$ and $q-p+2 t$ generate essentially all the invariant divisors of the fourth Painlevé equation 1.2. Note that the three simple affine roots $1-v_{1}+v_{3}, v_{1}-v_{2}, v_{2}-v_{3}$ of type $A_{2}^{(1)}$ are already involved in these equations. We denote by

$$
\begin{equation*}
V=\left\{\mathbf{v}=\left(v_{1}, v_{2}, v_{3}\right) \in \mathbb{C}^{3} ; v_{1}+v_{2}+v_{3}=0\right\} \tag{1.7}
\end{equation*}
$$

the parameter space for the system 1.2.
We now propose to treat the three typical invariant divisors $q, p$ and $q-p+2 t$ equally so as to obtain a "symmetric" representation of the fourth Painlevé equation. We introduce the three dependent variables $f=$ $\left(f_{0}, f_{1}, f_{2}\right)$ as follows. Fixing a nonzero complex number $c \in \mathbb{C}^{\times}$, set

$$
\begin{equation*}
f_{0}=c(q-p+2 t), \quad f_{1}=-c q, \quad f_{2}=c p \tag{1.8}
\end{equation*}
$$

and rescale the independent variable as $x=-t / c$. Then we have

$$
\begin{align*}
f_{0}^{\prime} & =f_{0}\left(f_{2}-f_{1}\right)-2 c^{2}\left(1-v_{1}+v_{3}\right) \\
f_{1}^{\prime} & =f_{1}\left(f_{0}-f_{2}\right)-2 c^{2}\left(v_{1}-v_{2}\right)  \tag{1.9}\\
f_{2}^{\prime} & =f_{2}\left(f_{1}-f_{0}\right)-2 c^{2}\left(v_{2}-v_{3}\right)
\end{align*}
$$

where $^{\prime}=d / d x$. With the normalization $c=\sqrt{-3 / 2}$, we set

$$
\begin{equation*}
\alpha_{0}=3\left(1-v_{1}+v_{3}\right), \quad \alpha_{1}=3\left(v_{1}-v_{2}\right), \quad \alpha_{2}=3\left(v_{2}-v_{3}\right) \tag{1.10}
\end{equation*}
$$

Then we have
Theorem 1.1. The fourth Painlevé equation 1.1 (or 1.2) can be written in the following symmetric form:

$$
\begin{align*}
& f_{0}^{\prime}+f_{0}\left(f_{1}-f_{2}\right)=\alpha_{0}, \\
& f_{1}^{\prime}+f_{1}\left(f_{2}-f_{0}\right)=\alpha_{1},  \tag{1.11}\\
& f_{2}^{\prime}+f_{2}\left(f_{0}-f_{1}\right)=\alpha_{2},
\end{align*}
$$

with normalization

$$
\begin{equation*}
f_{0}+f_{1}+f_{2}=3 x \tag{1.12}
\end{equation*}
$$

where ${ }^{\prime}=d / d x$ and $\alpha_{0}, \alpha_{1}, \alpha_{2} \in \mathbb{C}$ are parameters with $\alpha_{0}+\alpha_{1}+\alpha_{2}=3$.

Proof. The equation 1.11 has been derived from the Hamiltonian system 1.2 ; it is clear that these two are equivalent. We will show the equivalence of 1.1 and 1.11 (with normalization 1.12). This gives in fact an easier way to establish the equivalence between 1.1 and 1.2. Taking a derivative of the second equation of 1.11 , we have

$$
f_{1}^{\prime \prime}+f_{1}^{\prime}\left(f_{2}-f_{0}\right)+f_{1}\left(f_{2}^{\prime}-f_{0}^{\prime}\right)=0
$$

Substituting the first and the third equations of 1.11 to this, we obtain

$$
f_{1}^{\prime \prime}+f_{1}^{\prime}\left(f_{2}-f_{0}\right)-2 f_{0} f_{2} f_{1}-\left(\alpha_{0}-\alpha_{2}\right) f_{1}+\left(f_{2}+f_{0}\right) f_{1}^{2}=0
$$

Then, by using the relations

$$
f_{2}-f_{0}=\frac{\alpha_{1}-f_{1}^{\prime}}{f_{1}}, \quad f_{2}+f_{0}=3 x-f_{1}, \quad 4 f_{0} f_{2}=\left(f_{2}+f_{0}\right)^{2}-\left(f_{2}-f_{0}\right)^{2}
$$

we have

$$
f_{1}^{\prime \prime}-\frac{1}{2} \frac{f_{1}^{\prime 2}}{f_{1}}-\frac{3}{2} f_{1}^{3}+6 x f_{1}^{2}+\left(-\frac{9}{2} x^{2}-\left(\alpha_{0}-\alpha_{2}\right)\right) f_{1}+\frac{\alpha_{1}^{2}}{2} \frac{1}{f_{1}}=0
$$

This is transformed into the equation 1.1 by the rescaling $f_{1}=-c y, x=$ $-t / c, c=\sqrt{-3 / 2}$ and the change of parameters 1.5, 1.10.

We remark that our equation 1.11 has the following rational solutions:

$$
\begin{array}{ll}
\text { (A) } & \left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)=(1,1,1 ; x, x, x)  \tag{1.13}\\
\text { (B) } & \left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)=(3,0,0 ; 3 x, 0,0) .
\end{array}
$$

From the work of Y. Murata [4], it follows that all the rational solutions of 1.11 are obtained from these two particular solutions by Bäcklund transformations. There are classical solutions obtained as Bäcklund transformations from the solutions of Riccati type along the three lines $\alpha_{0}=0, \alpha_{1}=0$, $\alpha_{2}=0$. Any other solutions are non-classical in the sense of H. Umemura [9] (see also [6], [7]). It should also be noted that our equation 1.9 reduces to the Kac-Moerbeke integrable system [3] in the degenerate limit $c \rightarrow 0$.

## §2. Bäcklund transformations and the affine Weyl group

We now discuss symmetries in the fourth Painlevé equation represented by 1.11 with the normalization of 1.12 . In what follows, we regard $\alpha_{0}, \alpha_{1}$, $\alpha_{2}$ as coordinate functions (with $\alpha_{0}+\alpha_{1}+\alpha_{2}=3$ ) of the parameter space $V$.

We consider the affine Weyl group $W=\left\langle s_{0}, s_{1}, s_{2}\right\rangle$ of type $A_{2}^{(1)}$ with fundamental relations

$$
\begin{equation*}
s_{i}^{2}=1, \quad s_{i} s_{i+1} s_{i}=s_{i+1} s_{i} s_{i+1} \quad(i=0,1,2) \tag{2.1}
\end{equation*}
$$

Here the subscripts are understood as elements of $\mathbb{Z} / 3 \mathbb{Z}$. This convention for subscripts will be applied to other variables $\alpha_{i}, f_{i}$, etc., as well. We denote by $\widetilde{W}=\left\langle s_{0}, s_{1}, s_{2}, \pi\right\rangle$ the extension of $W$ obtained by adjoining the following Dynkin diagram automorphism $\pi$ :

$$
\begin{equation*}
\pi^{3}=1, \quad \pi s_{i}=s_{i+1} \pi \quad(i=0,1,2) \tag{2.2}
\end{equation*}
$$

The affine Weyl group $\widetilde{W}$ acts naturally on the coordinate ring $\mathbb{C}[\alpha]$ of $V$ through the algebra automorphism $s_{0}, s_{1}, s_{2}$ and $\pi$ of $\mathbb{C}[\alpha]$ determined by

$$
\begin{equation*}
s_{i}\left(\alpha_{i}\right)=-\alpha_{i}, \quad s_{i}\left(\alpha_{j}\right)=\alpha_{j}+\alpha_{i} \quad(i \neq j), \quad \pi\left(\alpha_{j}\right)=\alpha_{j+1} \tag{2.3}
\end{equation*}
$$

for $i, j=0,1,2$. When we consider the action of $\widetilde{W}$ on the parameter space $V$, we will use the action such that $(w . \varphi)(\mathbf{v})=\varphi\left(w^{-1} \cdot \mathbf{v}\right)$ for any $\mathbf{v} \in V$ and $\varphi \in \mathbb{C}[\alpha]$. The action of $\widetilde{W}$ on $V$ is given as follows:

$$
\begin{array}{ll}
s_{0} \cdot \mathbf{v}=\left(v_{3}+1, v_{2}, v_{1}-1\right), & s_{1} \cdot \mathbf{v}=\left(v_{2}, v_{1}, v_{3}\right) \\
s_{2} \cdot \mathbf{v}=\left(v_{1}, v_{3}, v_{2}\right), & \pi \cdot \mathbf{v}=\left(v_{3}+\frac{2}{3}, v_{1}-\frac{1}{3}, v_{2}-\frac{1}{3}\right) \tag{2.4}
\end{array}
$$

for any $\mathbf{v}=\left(v_{1}, v_{2}, v_{3}\right) \in V$.
One advantage of our representation 1.11 is that the action of the affine Weyl group $\widetilde{W}$ on the fourth Painlevé equation can be described in a completely symmetric way on the dependent variables $f_{0}, f_{1}$ and $f_{2}$. The action of $\widetilde{W}$ on $\mathbb{C}[\alpha]$ extends in fact to the whole differential field $K=\mathbb{C}(\alpha ; f)$ as follows.

Theorem 2.1. The fourth Painlevé equation 1.11 is invariant under the following transformations $s_{0}, s_{1}, s_{2}$ and $\pi$ :

$$
\begin{array}{llll}
s_{0}\left(f_{0}\right)=f_{0}, & s_{1}\left(f_{1}\right)=f_{1}, & s_{2}\left(f_{2}\right)=f_{2}, & \pi\left(f_{0}\right)=f_{1} \\
s_{0}\left(f_{1}\right)=f_{1}-\frac{\alpha_{0}}{f_{0}}, & s_{1}\left(f_{2}\right)=f_{2}-\frac{\alpha_{1}}{f_{1}}, & s_{2}\left(f_{0}\right)=f_{0}-\frac{\alpha_{2}}{f_{2}}, & \pi\left(f_{1}\right)=f_{2} \\
s_{0}\left(f_{2}\right)=f_{2}+\frac{\alpha_{0}}{f_{0}}, & s_{1}\left(f_{0}\right)=f_{0}+\frac{\alpha_{1}}{f_{1}}, & s_{2}\left(f_{1}\right)=f_{1}+\frac{\alpha_{2}}{f_{2}}, & \pi\left(f_{2}\right)=f_{0}  \tag{2.5}\\
s_{0}\left(\alpha_{0}\right)=-\alpha_{0}, & s_{1}\left(\alpha_{1}\right)=-\alpha_{1}, & s_{2}\left(\alpha_{2}\right)=-\alpha_{2}, & \pi\left(\alpha_{0}\right)=\alpha_{1} \\
s_{0}\left(\alpha_{1}\right)=\alpha_{1}+\alpha_{0}, & s_{1}\left(\alpha_{2}\right)=\alpha_{2}+\alpha_{1}, & s_{2}\left(\alpha_{0}\right)=\alpha_{0}+\alpha_{2}, & \pi\left(\alpha_{1}\right)=\alpha_{2} \\
s_{0}\left(\alpha_{2}\right)=\alpha_{2}+\alpha_{0}, & s_{1}\left(\alpha_{0}\right)=\alpha_{0}+\alpha_{1}, & s_{2}\left(\alpha_{1}\right)=\alpha_{1}+\alpha_{2}, & \pi\left(\alpha_{2}\right)=\alpha_{0}
\end{array}
$$

Furthermore, these transformations define a representation of the affine Weyl group $\widetilde{W}=\left\langle s_{0}, s_{1}, s_{2}, \pi\right\rangle$. Namely, $\widetilde{W}$ acts on the differential field $K=\mathbb{C}(\alpha ; f)$ as a group of differential automorphisms.

Theorem 2.1 is proved by straightforward computations. The transformations described above will be called the Bäcklund transformations of the fourth Painlevé equation 1.11. Note that the independent variable $x=\left(f_{0}+f_{1}+f_{2}\right) / 3$ is fixed under the action of $\widetilde{W}$.

Note that, for any $w \in W$, one obtains three linear functions $\beta_{0}=$ $w\left(\alpha_{0}\right), \beta_{1}=w\left(\alpha_{1}\right), \beta_{2}=w\left(\alpha_{2}\right)$ in $\alpha_{0}, \alpha_{1}, \alpha_{2}$. Theorem 2.1 then implies that, one can specify certain rational functions $g_{0}=w\left(f_{0}\right), g_{1}=w\left(f_{1}\right), g_{2}=$ $w\left(f_{2}\right)$ in $f_{0}, f_{1}, f_{2}, \alpha_{0}, \alpha_{1}, \alpha_{2}$ such that

$$
\begin{equation*}
g_{i}^{\prime}+g_{i}\left(g_{i+1}-g_{i+2}\right)=\beta_{i} \quad(i=0,1,2) \tag{2.6}
\end{equation*}
$$

Namely, if $\left(f_{0}, f_{1}, f_{2}\right)$ is a (generic) solution of 1.11 with parameters $\left(\alpha_{0}, \alpha_{1}\right.$, $\left.\alpha_{2}\right)$, then $\left(g_{0}, g_{1}, g_{2}\right)$ is again a solution of the same system with parameters $\left(\beta_{0}, \beta_{1}, \beta_{2}\right)$. We give an example below to show how the dependent variables $f_{0}, f_{1}, f_{2}$ are transformed under the action of the affine Weyl group.

Example. For $w=s_{1} s_{0}$, the Bäcklund transformation $w\left(f_{1}\right)=$ $s_{1} s_{0}\left(f_{1}\right)$ is computed as follows:

$$
\begin{equation*}
f_{1} \xrightarrow{s_{0}} \frac{f_{0} f_{1}-\alpha_{0}}{f_{0}} \xrightarrow{s_{1}} \frac{f_{1}\left(f_{0} f_{1}-\alpha_{0}\right)}{f_{0} f_{1}+\alpha_{1}} . \tag{2.7}
\end{equation*}
$$

Similarly we have

$$
\begin{align*}
& \beta_{0}=w\left(\alpha_{0}\right)=\alpha_{2}-3, \quad \beta_{1}=w\left(\alpha_{1}\right)=\alpha_{0}, \quad \beta_{2}=w\left(\alpha_{2}\right)=\alpha_{1}+3 \\
& g_{0}=w\left(f_{0}\right)=\frac{f_{0} f_{1}+\alpha_{1}}{f_{1}}, \quad g_{1}=w\left(f_{1}\right)=\frac{f_{1}\left(f_{0} f_{1}-\alpha_{0}\right)}{f_{0} f_{1}+\alpha_{1}}  \tag{2.8}\\
& g_{2}=w\left(f_{2}\right)=\frac{\left(f_{0} f_{1}+\alpha_{1}\right)\left(f_{1} f_{2}-\alpha_{1}\right)+\left(3-\alpha_{2}\right) f_{1}^{2}}{f_{1}\left(f_{0} f_{1}+\alpha_{1}\right)}
\end{align*}
$$

If we specialize these formula to the particular solution

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)=(1,1,1 ; x, x, x) \tag{2.9}
\end{equation*}
$$

we obtain another rational solution

$$
\begin{align*}
& \left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)  \tag{2.10}\\
& \quad=\left(-2,1,4 ; \frac{x^{2}+1}{x}, \frac{x\left(x^{2}-1\right)}{\left(x^{2}+1\right)}, \frac{x^{4}+2 x^{2}-1}{x\left(x^{2}+1\right)}\right) .
\end{align*}
$$

A complete description of rational functions in $x$ arising in this way will be given later in this paper.

Remark. The Bäcklund transformation $s_{0}\left(f_{1}\right)=f_{1}-\frac{\alpha_{0}}{f_{0}}$, for example, becomes singular when applied to a particular solution such that $f_{0}=$ 0 . This sort of problem, however, is only apparent since such a solution arises only under the condition $\alpha_{0}=0$ as one sees immediately from 1.11. When $\alpha_{0}=0$, it is natural to understand that the Bäcklund transformation $s_{0}$ becomes the identity transformation. In general, each $g_{i}=w\left(f_{i}\right)$ is a rational function in $(\alpha ; f)$ and its denominator possibly becomes identically zero when one specializes $(\alpha ; f)$ to certain particular solutions. Such a phenomenon occurs however only when some of the parameters $\alpha_{0}, \alpha_{1}$, $\alpha_{2}$ are in $3 \mathbb{Z}$. In such cases, critical factors in the denominator of $g_{i}=$ $w\left(f_{i}\right)$ can actually be eliminated by specializing the parameters $\left(\alpha_{0}, \alpha_{1}, \alpha_{2}\right)$ in advance. With this regularization, our Bäcklund transformations $w\left(f_{i}\right)$ make sense for any particular solution.

## §3. $\tau$-Functions

In this section, we show that our equation 1.11 for $f=\left(f_{0}, f_{1}, f_{2}\right)$ can be bilinearized by introducing a triple of $\tau$-functions $\tau=\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$. We also study the Bäcklund transformations on the level of $\tau$-functions.

We introduce the $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$ to be the dependent variables satisfying the following equations:

$$
\begin{align*}
& f_{0}=\left(\log \frac{\tau_{1}}{\tau_{2}}\right)^{\prime}+x=\frac{\tau_{1}^{\prime}}{\tau_{1}}-\frac{\tau_{2}^{\prime}}{\tau_{2}}+x \\
& f_{1}=\left(\log \frac{\tau_{2}}{\tau_{0}}\right)^{\prime}+x=\frac{\tau_{2}^{\prime}}{\tau_{2}}-\frac{\tau_{0}^{\prime}}{\tau_{0}}+x  \tag{3.1}\\
& f_{2}=\left(\log \frac{\tau_{0}}{\tau_{1}}\right)^{\prime}+x=\frac{\tau_{0}^{\prime}}{\tau_{0}}-\frac{\tau_{1}^{\prime}}{\tau_{1}}+x
\end{align*}
$$

We fix the freedom of overall multiplication by a function in defining $\tau_{0}, \tau_{1}$, $\tau_{2}$, by imposing the equation

$$
\begin{equation*}
2\left(\log \tau_{0} \tau_{1} \tau_{2}\right)^{\prime \prime}+\left(f_{0}-x\right)^{2}+\left(f_{1}-x\right)^{2}+\left(f_{2}-x\right)^{2}=0 \tag{3.2}
\end{equation*}
$$

To be more precise, we first introduce a variable $g$ (determined from $f_{0}, f_{1}$, $f_{2}$ up to an additive constant) as an integral of the equation

$$
\begin{equation*}
2 g^{\prime}+\left(f_{0}-x\right)^{2}+\left(f_{1}-x\right)^{2}+\left(f_{2}-x\right)^{2}=0 . \tag{3.3}
\end{equation*}
$$

Then we require that the $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$ should satisfy

$$
\begin{equation*}
g=\left(\log \tau_{0} \tau_{1} \tau_{2}\right)^{\prime}=\frac{\tau_{0}^{\prime}}{\tau_{0}}+\frac{\tau_{1}^{\prime}}{\tau_{1}}+\frac{\tau_{1}^{\prime}}{\tau_{1}} \tag{3.4}
\end{equation*}
$$

Note that, under the conditions 3.1 and 3.4 , the $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$ are determined by the equations

$$
\begin{align*}
& \left(\log \tau_{0}\right)^{\prime}=\frac{\tau_{0}^{\prime}}{\tau_{0}}=\frac{1}{3}\left(g-f_{1}+f_{2}\right) \\
& \left(\log \tau_{1}\right)^{\prime}=\frac{\tau_{1}^{\prime}}{\tau_{1}}=\frac{1}{3}\left(g-f_{2}+f_{0}\right)  \tag{3.5}\\
& \left(\log \tau_{2}\right)^{\prime}=\frac{\tau_{2}^{\prime}}{\tau_{2}}=\frac{1}{3}\left(g-f_{0}+f_{1}\right)
\end{align*}
$$

up to multiplicative constants, respectively. We remark that the integration constant in $g$ has the effect of multiplying each $\tau_{i}$ by the exponential of a linear function in $x$.

In order to describe the differential equations to be satisfied by the $\tau$ functions, we recall the definition of Hirota's bilinear equations. Let $P\left(\partial_{x}\right)$
$\left(\partial_{x}=d / d x\right)$ be a linear differential operator in the $x$-variable with constant coefficients. Then Hirota's bilinear operator $P\left(D_{x}\right)$ is defined by

$$
\begin{equation*}
P\left(D_{x}\right) F(x) \cdot G(x)=\left.P\left(\partial_{y}\right) F(x+y) G(x-y)\right|_{y=0} \tag{3.6}
\end{equation*}
$$

for a given pair of functions $F(x), G(x)$.
Theorem 3.1. The fourth Painlevé equation 1.11 for $f_{0}, f_{1}, f_{2}$, together with the integral $g$ of 3.3, is equivalent to the following system of Hirota bilinear equations for the triple of $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$ :

$$
\begin{align*}
& \left(D_{x}^{2}-x D_{x}-\frac{\alpha_{0}-\alpha_{1}}{3}\right) \tau_{0} \cdot \tau_{1}=0 \\
& \left(D_{x}^{2}-x D_{x}-\frac{\alpha_{1}-\alpha_{2}}{3}\right) \tau_{1} \cdot \tau_{2}=0  \tag{3.7}\\
& \left(D_{x}^{2}-x D_{x}-\frac{\alpha_{2}-\alpha_{0}}{3}\right) \tau_{2} \cdot \tau_{0}=0
\end{align*}
$$

Proof. Note first that, in terms of the logarithms $F_{i}=\log \tau_{i}(i=0,1,2)$ of $\tau$-functions, the dependent variables $f_{0}, f_{1}, f_{2}$ are expressed as follows:

$$
\begin{align*}
& f_{0}=F_{1}^{\prime}-F_{2}^{\prime}+x, \quad f_{1}=F_{2}^{\prime}-F_{0}^{\prime}+x, \quad f_{2}=F_{0}^{\prime}-F_{1}^{\prime}+x  \tag{3.8}\\
& g=F_{0}^{\prime}+F_{1}^{\prime}+F_{2}^{\prime}
\end{align*}
$$

The three equations of Theorem are rewritten into the following equations for $F_{0}, F_{1}, F_{2}$ :

$$
\begin{align*}
& F_{0}^{\prime \prime}+F_{1}^{\prime \prime}+\left(F_{0}^{\prime}-F_{1}^{\prime}\right)^{2}-x\left(F_{0}^{\prime}-F_{1}^{\prime}\right)-\frac{\alpha_{0}-\alpha_{1}}{3}=0 \\
& F_{1}^{\prime \prime}+F_{2}^{\prime \prime}+\left(F_{1}^{\prime}-F_{2}^{\prime}\right)^{2}-x\left(F_{1}^{\prime}-F_{2}^{\prime}\right)-\frac{\alpha_{1}-\alpha_{2}}{3}=0  \tag{3.9}\\
& F_{2}^{\prime \prime}+F_{0}^{\prime \prime}+\left(F_{2}^{\prime}-F_{0}^{\prime}\right)^{2}-x\left(F_{2}^{\prime}-F_{0}^{\prime}\right)-\frac{\alpha_{2}-\alpha_{0}}{3}=0
\end{align*}
$$

Taking the sum of these three equations, we have (3.10) $2\left(F_{0}^{\prime \prime}+F_{1}^{\prime \prime}+F_{2}^{\prime \prime}\right)+\left(F_{1}^{\prime}-F_{2}^{\prime}\right)^{2}+\left(F_{2}^{\prime}-F_{0}^{\prime}\right)^{2}+\left(F_{0}^{\prime}-F_{1}^{\prime}\right)^{2}=0$,
which corresponds to the equation 3.3 for $g$. By subtracting the third equation of 3.9 from the first, we have

$$
\begin{equation*}
F_{1}^{\prime \prime}-F_{2}^{\prime \prime}-\left(F_{1}^{\prime}-F_{2}^{\prime}+x\right)\left(2 F_{0}^{\prime}-F_{1}^{\prime}-F_{2}^{\prime}\right)-\alpha_{0}+1=0 \tag{3.11}
\end{equation*}
$$

which corresponds to the differential equation for $f_{0}$. Similarly we have the equations for $f_{1}$ and $f_{2}$ from 3.9. It is also clear that the equations 3.9 are recovered from 3.10 and the three equations which correspond to 1.11 .

Remark. Consider the differential field $K(g)=\mathbb{C}(\alpha ; f)(g)$ obtained from $K=\mathbb{C}(\alpha ; f)$ by adjoining a variable $g$ on which the derivation ' acts by the formula 3.3. Then Theorem 3.1 implies that this differential field is isomorphic to the differential field $\mathbb{C}(\alpha)\left(x, F_{0}^{\prime}, F_{1}^{\prime}, F_{2}^{\prime}\right)$ defined by the relations 3.9. Note that, by 3.9 and 3.10 , each second derivative $F_{i}^{\prime \prime}$ $(i=0,1,2)$ can be expressed in terms of $x$ and $F_{0}^{\prime}, F_{1}^{\prime}, F_{2}^{\prime}$ :

$$
\begin{equation*}
F_{i}^{\prime \prime}+x\left(F_{i+1}^{\prime}-F_{i+2}^{\prime}\right)+\left(F_{i}^{\prime}-F_{i+1}^{\prime}\right)\left(F_{i}^{\prime}-F_{i+2}^{\prime}\right)+\frac{\alpha_{i+1}-\alpha_{i+2}}{3}=0 \tag{3.12}
\end{equation*}
$$

for $i=0,1,2$. This system is also equivalent to the equation 3.7 for the triple $\tau_{0}, \tau_{1}, \tau_{2}$ of $\tau$-functions. Note that the differential field of our $\tau$-functions is defined as $\mathbb{C}(\alpha)\left(x, \tau_{0}, \tau_{1}, \tau_{2}, \tau_{0}^{\prime}, \tau_{1}^{\prime}, \tau_{2}^{\prime}\right)$ by 3.12 , regarded as equations for $\tau$-functions.

One important fact is that the action of the affine Weyl group on the $f$-variables lifts to the level of $\tau$-functions.

TheOrem 3.2. The $\tau$-functions $\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$ allow an action of the affine Weyl group $\widetilde{W}$ which is compatible with the action of $\widetilde{W}$ on $f_{0}, f_{1}, f_{2}$ of Theorem 2.1. Their Bäcklund transformations are again expressed by Hirota's bilinear operators as follows:

$$
\begin{align*}
& s_{0}\left(\tau_{0}\right)=\frac{1}{\tau_{0}}\left(D_{x}+x\right) \tau_{1} \cdot \tau_{2}=\frac{1}{\tau_{0}}\left(\tau_{1}^{\prime} \tau_{2}-\tau_{1} \tau_{2}^{\prime}+x \tau_{1} \tau_{2}\right), \\
& s_{1}\left(\tau_{1}\right)=\frac{1}{\tau_{1}}\left(D_{x}+x\right) \tau_{2} \cdot \tau_{0}=\frac{1}{\tau_{1}}\left(\tau_{2}^{\prime} \tau_{0}-\tau_{2} \tau_{0}^{\prime}+x \tau_{2} \tau_{0}\right),  \tag{3.13}\\
& s_{2}\left(\tau_{2}\right)=\frac{1}{\tau_{2}}\left(D_{x}+x\right) \tau_{0} \cdot \tau_{1}=\frac{1}{\tau_{2}}\left(\tau_{0}^{\prime} \tau_{1}-\tau_{0} \tau_{1}^{\prime}+x \tau_{0} \tau_{1}\right), \\
& s_{i}\left(\tau_{j}\right)=\tau_{j} \quad(i \neq j), \quad \pi\left(\tau_{j}\right)=\tau_{j+1} \quad(i, j=0,1,2),
\end{align*}
$$

while $s_{0}, s_{1}, s_{2}$ and $\pi$ act on $\alpha_{0}, \alpha_{1}, \alpha_{2}$ in the same way as in Theorem 2.1.
Proof. We first extend the action of $\widetilde{W}$ on $\mathbb{C}(\alpha ; f)$ to $\mathbb{C}(\alpha ; f)(g)$, or equivalently to $\mathbb{C}(\alpha)\left(x, F_{0}^{\prime}, F_{1}^{\prime}, F_{2}^{\prime}\right)$. From 3.3 we have

$$
\begin{equation*}
s_{0}\left(g^{\prime}\right)=g^{\prime}+\left(f_{1}-f_{2}\right) \frac{\alpha_{0}}{f_{0}}-\left(\frac{\alpha_{0}}{f_{0}}\right)^{2}=g^{\prime}-\alpha_{0} \frac{f_{0}^{\prime}}{f_{0}^{2}} \tag{3.14}
\end{equation*}
$$

by 1.11. Hence we have

$$
\begin{equation*}
s_{i}\left(g^{\prime}\right)=g^{\prime}-\alpha_{i} \frac{f_{i}^{\prime}}{f_{i}^{2}} \quad(i=0,1,2), \quad \pi\left(g^{\prime}\right)=g^{\prime} \tag{3.15}
\end{equation*}
$$

In view of these, we define the action of $\widetilde{W}$ on $g$ by

$$
\begin{equation*}
s_{i}(g)=g+\frac{\alpha_{i}}{f_{i}} \quad(i=0,1,2), \quad \pi(g)=g \tag{3.16}
\end{equation*}
$$

One can check that 3.16 gives rise in fact to a representation of $\widetilde{W}$ on $\mathbb{C}(\alpha ; f)(g)$. On the variables $F_{0}^{\prime}, F_{1}^{\prime}, F_{2}^{\prime}$, equation 3.5 together with 3.16 immediately implies

$$
\begin{equation*}
s_{i}\left(F_{j}^{\prime}\right)=F_{j}^{\prime} \quad(i \neq j), \quad \pi\left(F_{j}^{\prime}\right)=F_{j+1}^{\prime} \quad(i, j=0,1,2) \tag{3.17}
\end{equation*}
$$

These formulas justify the definitions of 3.13 other than those for $s_{i}\left(\tau_{i}\right)$ ( $i=0,1,2$ ). As to $s_{0}\left(\tau_{0}\right)$, we compute

$$
\begin{equation*}
s_{0}\left(F_{0}^{\prime}\right)=F_{0}^{\prime}+\frac{\alpha_{0}}{f_{0}}=F_{0}^{\prime}+\frac{f_{0}^{\prime}}{f_{0}}+f_{1}-f_{2}=-F_{0}^{\prime}+F_{1}^{\prime}+F_{2}^{\prime}+\frac{f_{0}^{\prime}}{f_{0}} \tag{3.18}
\end{equation*}
$$

This leads to the definition

$$
\begin{equation*}
s_{0}\left(\tau_{0}\right)=\frac{\tau_{1} \tau_{2}}{\tau_{0}} f_{0}=\frac{\tau_{1} \tau_{2}}{\tau_{0}}\left(\frac{\tau_{1}^{\prime}}{\tau_{1}}-\frac{\tau_{2}^{\prime}}{\tau_{2}}+x\right)=\frac{1}{\tau_{0}}\left(D_{x}+x\right) \tau_{1} \cdot \tau_{2} \tag{3.19}
\end{equation*}
$$

One can check by straightforward computations that the definition 3.13 thus obtained defines an action of $\widetilde{W}$ on the differential field $\mathbb{C}(\alpha)\left(x, \tau_{0}, \tau_{1}, \tau_{2}, \tau_{0}^{\prime}\right.$, $\left.\tau_{1}^{\prime}, \tau_{2}^{\prime}\right)$ as a group of differential automorphisms.

We remark that the Bäcklund transformations $s_{i}\left(\tau_{i}\right)$ of Theorem 3.2 possibly become zero for solutions reducible to Riccati equations, while they can be applied repeatedly as long as the $\tau$-functions remain nonzero. If $\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$ is a generic solution, we obtain the Bäcklund transformations $\left(w\left(\tau_{0}\right), w\left(\tau_{1}\right), w\left(\tau_{2}\right)\right)$ for any $w \in \widetilde{W}$, by Theorem 3.2.

From the formula 3.19 in the proof of Theorem 3.2, we have
Corollary 3.3. In terms of the $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$, the dependent variables $f_{0}, f_{1}, f_{2}$ of the fourth Painlevé equation 1.11 are expressed multiplicatively as follows:

$$
\begin{equation*}
f_{0}=\frac{\tau_{0} s_{0}\left(\tau_{0}\right)}{\tau_{1} \tau_{2}}, \quad f_{1}=\frac{\tau_{1} s_{1}\left(\tau_{1}\right)}{\tau_{2} \tau_{0}}, \quad f_{2}=\frac{\tau_{2} s_{2}\left(\tau_{2}\right)}{\tau_{0} \tau_{1}} \tag{3.20}
\end{equation*}
$$

The relation between the $f$-variables and the six $\tau$-functions in Corollary 3.3 can be represented graphically as in Figure 1. Note also that 3.20 implies

$$
\begin{equation*}
\tau_{0}^{2} s_{0}\left(\tau_{0}\right)+\tau_{1}^{2} s_{1}\left(\tau_{1}\right)+\tau_{2}^{2} s_{2}\left(\tau_{2}\right)=3 x \tau_{0} \tau_{1} \tau_{2} \tag{3.21}
\end{equation*}
$$



Figure 1: Six $\tau$-functions

Example. As to the rational solution 2.9 the corresponding $\tau$-functions and their adjacent Bäcklund transformations are given by

$$
\begin{equation*}
\left(\tau_{0}, \tau_{1}, \tau_{2}\right)=(1,1,1), \quad\left(s_{0}\left(\tau_{0}\right), s_{1}\left(\tau_{1}\right), s_{2}\left(\tau_{2}\right)\right)=(x, x, x) \tag{3.22}
\end{equation*}
$$

As to the rational solution 2.10, we have

$$
\begin{align*}
& \left(\tau_{0}, \tau_{1}, \tau_{2}\right)=\left(x^{2}+1, x, 1\right) \\
& \left(s_{0}\left(\tau_{0}\right), s_{1}\left(\tau_{1}\right), s_{2}\left(\tau_{2}\right)\right)=\left(1, x^{2}-1, x^{4}+2 x^{2}-1\right) \tag{3.23}
\end{align*}
$$

These are examples of Okamoto polynomials which will be discussed in the next section.

Another corollary of Theorem 3.2 is the Toda equations for our $\tau$ functions.

Corollary 3.4. The fourth Painlevé equation 3.7 for the triple of $\tau$-functions $\tau_{0}, \tau_{1}, \tau_{2}$ implies the following equation of Toda type:

$$
\begin{equation*}
\left(\log \tau_{0}\right)^{\prime \prime}+x^{2}+\frac{\alpha_{1}-\alpha_{2}}{3}=\frac{s_{1}\left(\tau_{1}\right) s_{2}\left(\tau_{2}\right)}{\tau_{0}^{2}} \tag{3.24}
\end{equation*}
$$

namely,

$$
\begin{equation*}
\left(\frac{1}{2} D_{x}^{2}+x^{2}+\frac{\alpha_{1}-\alpha_{2}}{3}\right) \tau_{0} \cdot \tau_{0}=s_{1}\left(\tau_{1}\right) s_{2}\left(\tau_{2}\right) \tag{3.25}
\end{equation*}
$$

Proof. From Corollary 3.3, we have

$$
\begin{equation*}
f_{1} f_{2}=\frac{s_{1}\left(\tau_{1}\right) s_{2}\left(\tau_{2}\right)}{\tau_{0}^{2}} \tag{3.26}
\end{equation*}
$$

On the other hand, substitution of the formulas

$$
\begin{equation*}
F_{2}^{\prime}-F_{0}^{\prime}=f_{1}-x, \quad F_{0}^{\prime}-F_{1}^{\prime}=f_{2}-x, \quad F_{1}^{\prime}-F_{2}^{\prime}=2 x-f_{1}-f_{2} \tag{3.27}
\end{equation*}
$$

into 3.12 with $i=0$ gives

$$
\begin{equation*}
F_{0}^{\prime \prime}+x^{2}+\frac{\alpha_{1}-\alpha_{2}}{3}=f_{1} f_{2} \tag{3.28}
\end{equation*}
$$

Equating 3.26 and 3.28 we obtain the equation of Corollary as desired.

Our $\tau$-functions are slightly different from those introduced by K. Okamoto [7]. In our formulation, the $\tau$-function in the spirit of Okamoto, say $\tau^{\mathrm{ok}}$, can be defined through the integral of a "Hamiltonian" as follows:

$$
\begin{equation*}
H=\frac{1}{3}\left(f_{0} f_{1} f_{2}+\alpha_{1} f_{2}-\alpha_{2} f_{1}\right)=\left(\log \tau^{\mathrm{ok}}\right)^{\prime} \tag{3.29}
\end{equation*}
$$

Note that this implies $\left(\log \tau^{\mathrm{ok}}\right)^{\prime \prime}=H^{\prime}=f_{1} f_{2}$. Let us introduce the triple of $\tau$-functions of Okamoto type by

$$
\begin{equation*}
\left(\log \tau_{0}^{\mathrm{ok}}\right)^{\prime}=H_{0}, \quad\left(\log \tau_{1}^{\mathrm{ok}}\right)^{\prime}=H_{1},\left(\log \tau_{2}^{\mathrm{ok}}\right)^{\prime}=H_{2} \tag{3.30}
\end{equation*}
$$

where we define $H_{0}=H, H_{1}=\pi(H), H_{2}=\pi^{2}(H)$ by rotation. This implies

$$
\begin{align*}
& f_{0}=\left(\log \frac{\tau_{1}^{\mathrm{ok}}}{\tau_{2}^{\mathrm{ok}}}\right)^{\prime}+\alpha_{0} x, \quad f_{1}=\left(\log \frac{\tau_{2}^{\mathrm{ok}}}{\tau_{0}^{\mathrm{ok}}}\right)^{\prime}+\alpha_{1} x  \tag{3.31}\\
& f_{2}=\left(\log \frac{\tau_{0}^{\mathrm{ok}}}{\tau_{1}^{\mathrm{ok}}}\right)^{\prime}+\alpha_{2} x
\end{align*}
$$

(Compare these formulas with our definition 3.1.) From 3.28 we also see that

$$
\begin{equation*}
\tau_{0}^{\mathrm{ok}}=e^{x^{4} / 12+\left(\alpha_{1}-\alpha_{2}\right) x^{2} / 6} \tau_{0} \tag{3.32}
\end{equation*}
$$

up to the multiplication by the exponential of a linear function in $x$.

## §4. Rational solutions

In this section, we give an explicit description of the rational solutions of the fourth Painlevé equation 1.11 in terms of Schur functions.

Before discussing the rational solutions, we introduce a family of $\tau$ functions $\left(\tau_{m, n}\right)_{m, n \in \mathbb{Z}}$ for the fourth Painlevé equation 3.7. A similar treatment of the lattice of $\tau$-functions has been given by K. Okamoto [8]. We consider the elements

$$
\begin{equation*}
T_{1}=\pi s_{2} s_{1}, \quad T_{2}=s_{1} \pi s_{2} \tag{4.1}
\end{equation*}
$$

of the (extended) affine Weyl group $\widetilde{W}$. Note that these $T_{1}$ and $T_{2}$ represent the following parallel translations in the parameter space $V$ respectively:

$$
\begin{equation*}
T_{1} \cdot \mathbf{v}=\mathbf{v}+\left(\frac{2}{3},-\frac{1}{3},-\frac{1}{3}\right), \quad T_{2} \cdot \mathbf{v}=\mathbf{v}+\left(-\frac{1}{3}, \frac{2}{3},-\frac{1}{3}\right) \tag{4.2}
\end{equation*}
$$

for $\mathbf{v} \in V$. For the triple of $\tau$-functions $\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$ of the fourth Painlevé equation 3.7, we introduce an infinite family of dependent variables $\tau_{m, n}$ ( $m, n \in \mathbb{Z}$ ) as the Bäcklund transformations

$$
\begin{equation*}
\tau_{m, n}=T_{1}^{m} T_{2}^{n}\left(\tau_{0}\right) \quad(m, n \in \mathbb{Z}) \tag{4.3}
\end{equation*}
$$

Note that

$$
\begin{equation*}
T_{1}\left(\tau_{0}\right)=\tau_{1}, \quad T_{2}\left(\tau_{0}\right)=s_{1}\left(\tau_{1}\right) \quad \text { and } \quad T_{2} T_{1}\left(\tau_{0}\right)=T_{2}\left(\tau_{1}\right)=\tau_{2} \tag{4.4}
\end{equation*}
$$

By these formulas, we have

$$
\begin{equation*}
\tau_{0,0}=\tau_{0}, \quad \tau_{1,0}=\tau_{1}, \quad \tau_{1,1}=\tau_{2}, \quad \tau_{0,1}=s_{1}\left(\tau_{1}\right) \tag{4.5}
\end{equation*}
$$

The triple of $\tau$-functions $\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$ is transformed into $\left(\tau_{m, n}, \tau_{m+1, n}\right.$, $\left.\tau_{m+1, n+1}\right)$ by $T_{1}^{m} T_{2}^{n}$, and into $\left(\tau_{m, n}, \tau_{m, n+1}, \tau_{m+1, n+1}\right)$ by $T_{1}^{m} T_{2}^{n} s_{1}$, respectively. The following propositions are obtained immediately from the results of the previous section, by using the action of $\widetilde{W}$.

## Proposition 4.1. (1) For any $m, n \in \mathbb{Z}$, the triples

$$
\begin{equation*}
\left(\tau_{m, n}, \tau_{m+1, n}, \tau_{m+1, n+1}\right) \quad \text { and } \quad\left(\tau_{m, n}, \tau_{m, n+1}, \tau_{m+1, n+1}\right) \tag{4.6}
\end{equation*}
$$

represent the Bäcklund transformations of $\left(\tau_{0}, \tau_{1}, \tau_{2}\right)$ for the parameters

$$
\begin{align*}
& \left(\alpha_{0}+3 m, \alpha_{1}+3(n-m), \alpha_{2}-3 n\right) \quad \text { and }  \tag{4.7}\\
& \left(\alpha_{0}+\alpha_{1}+3 n,-\alpha_{1}+3(m-n), \alpha_{1}+\alpha_{2}-3 m\right),
\end{align*}
$$



Figure 2: $\tau$-Functions on the $A_{2}$-lattice
respectively.
(2) The corresponding $f$-variables are given respectively by

$$
\begin{align*}
& \left(\frac{\tau_{m, n} \tau_{m+2, n+1}}{\tau_{m+1, n} \tau_{m+1, n+1}}, \frac{\tau_{m+1, n} \tau_{m, n+1}}{\tau_{m+1, n+1} \tau_{m, n}}, \frac{\tau_{m+1, n+1} \tau_{m, n-1}}{\tau_{m, n} \tau_{m+1, n}}\right) \quad \text { and }  \tag{4.8}\\
& \left(\frac{\tau_{m, n} \tau_{m+1, n+2}}{\tau_{m, n+1} \tau_{m+1, n+1}}, \frac{\tau_{m, n+1} \tau_{m+1, n}}{\tau_{m+1, n+1} \tau_{m, n}}, \frac{\tau_{m+1, n+1} \tau_{m-1, n}}{\tau_{m, n} \tau_{m, n+1}}\right)
\end{align*}
$$

Proposition 4.2. (1) The family of $\tau$-functions $\tau_{m, n}(m, n \in \mathbb{Z})$ satisfies the following three types of bilinear equations:

$$
\begin{align*}
& \left(D_{x}+x\right) \tau_{m, n} \cdot \tau_{m+1, n}=\tau_{m, n-1} \tau_{m+1, n+1}, \\
& \left(D_{x}+x\right) \tau_{m, n} \cdot \tau_{m, n+1}=\tau_{m+1, n+1} \tau_{m-1, n}  \tag{4.9}\\
& \left(D_{x}+x\right) \tau_{m, n} \cdot \tau_{m-1, n-1}=\tau_{m-1, n} \tau_{m, n-1}
\end{align*}
$$

(2) The family of $\tau$-functions $\tau_{m, n}(m, n \in \mathbb{Z})$ satisfies the following three types of Toda equations:

$$
\begin{align*}
& \left(\frac{1}{2} D_{x}^{2}+x^{2}-\frac{2 \alpha_{1}+\alpha_{2}}{3}+2 m-n\right) \tau_{m, n} \cdot \tau_{m, n}=\tau_{m+1, n} \tau_{m-1, n} \\
& \left(\frac{1}{2} D_{x}^{2}+x^{2}+\frac{\alpha_{1}-\alpha_{2}}{3}-m+2 n\right) \tau_{m, n} \cdot \tau_{m, n}=\tau_{m, n+1} \tau_{m, n-1}  \tag{4.10}\\
& \left(\frac{1}{2} D_{x}^{2}+x^{2}+\frac{\alpha_{1}+2 \alpha_{2}}{3}-m-n\right) \tau_{m, n} \cdot \tau_{m, n}=\tau_{m-1, n-1} \tau_{m+1, n+1}
\end{align*}
$$

Remark. As we already remarked in the previous section, Bäcklund transformations for $\tau$-functions possibly become singular, when applied to particular solutions which are reducible to Riccati equations. In such cases, we need to restrict the indices $(m, n)$ for $\tau_{m, n}$ to a region of $\mathbb{Z}^{2}$ bounded by certain lines on which $\tau_{m, n}=0$.

All the rational solutions of 1.11 are obtained from

$$
\begin{align*}
& \text { (A) } \quad\left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)=(1,1,1 ; x, x, x), \quad \text { or }  \tag{4.11}\\
& \text { (B) } \quad\left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; f_{0}, f_{1}, f_{2}\right)=(3,0,0 ; 3 x, 0,0) .
\end{align*}
$$

by Bäcklund transformations. We will determine the $\tau$-functions $\tau_{m, n}(m, n$ $\in \mathbb{Z}$ ) for these rational solutions.

In the case of Bäcklund transformations of the rational solution (A) of 4.11, the $\tau$-functions $\tau_{m, n}(m, n \in \mathbb{Z})$ turn out to be polynomials, which we call the Okamoto polynomials. We recall that the $\tau$-functions for (A) are given by

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; \tau_{0}, \tau_{1}, \tau_{2}\right)=(1,1,1 ; 1,1,1) \tag{4.12}
\end{equation*}
$$

Theorem 4.3. The $\tau$-functions $\tau_{m, n}(x)$ for the solution 4.12 are polynomials in $x$. These polynomials $\tau_{m, n}(x)=Q_{m, n}(x)(m, n \in \mathbb{Z})$ are characterized by the Toda equations

$$
\begin{align*}
& \left(\frac{1}{2} D_{x}^{2}+x^{2}-1+2 m-n\right) Q_{m, n} \cdot Q_{m, n}=Q_{m+1, n} Q_{m-1, n} \\
& \left(\frac{1}{2} D_{x}^{2}+x^{2}-m+2 n\right) Q_{m, n} \cdot Q_{m, n}=Q_{m, n+1} Q_{m, n-1}  \tag{4.13}\\
& \left(\frac{1}{2} D_{x}^{2}+x^{2}+1-m-n\right) Q_{m, n} \cdot Q_{m, n}=Q_{m-1, n-1} Q_{m+1, n+1}
\end{align*}
$$

with initial condition

$$
\begin{equation*}
Q_{0,0}=Q_{1,0}=Q_{1,1}=1, \quad Q_{2,1}=x \tag{4.14}
\end{equation*}
$$

We remark that $Q_{m}(x)=Q_{m, 0}(x)$ and $R_{m}(x)=Q_{m+1,1}(x)(m \in \mathbb{Z})$ are the original Okamoto polynomials discussed in [1]. In fact, they are determined by the recurrence relations

$$
\begin{equation*}
\left(\frac{1}{2} D_{x}^{2}+x^{2}+2 m-1\right) Q_{m} \cdot Q_{m}=Q_{m+1} Q_{m-1} \quad(m \in \mathbb{Z}) \tag{4.15}
\end{equation*}
$$

with initial condition $Q_{0}=Q_{1}=1$, and by

$$
\begin{equation*}
\left(\frac{1}{2} D_{x}^{2}+x^{2}+2 m\right) R_{m} \cdot R_{m}=R_{m+1} R_{m-1} \quad(m \in \mathbb{Z}) \tag{4.16}
\end{equation*}
$$

with $R_{0}=1, R_{1}=x$, respectively. The fact that $\tau_{m, n}(x)$ are polynomials will be proved in Section 5 in the course of the proof of Theorem 4.5 below. The other statements in Theorem 4.3 are consequences of Proposition 4.2.

The $\tau$-functions for the rational solution (B) of 4.11 are given by

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2} ; \tau_{0}, \tau_{1}, \tau_{2}\right)=\left(3,0,0 ; e^{-x^{4} / 12}, e^{-x^{4} / 12+x^{2} / 2}, e^{-x^{4} / 12-x^{2} / 2}\right) \tag{4.17}
\end{equation*}
$$

Theorem 4.4. The $\tau$-functions $\tau_{m, n}(x)$ for the solution 4.17 are defined for $(m, n) \in \mathbb{Z}^{2}$ with $m \geq n \geq 0$. They can be written in the form

$$
\begin{equation*}
\tau_{m, n}(x)=\exp \left(-\frac{x^{4}}{12}+\frac{m-2 n}{2} x^{2}\right) H_{m-n, n} \quad(m \geq n \geq 0) \tag{4.18}
\end{equation*}
$$

for some polynomials $H_{m, n}(x)$. These polynomials $H_{m, n}(x)(m, n \geq 0)$ are characterized by the Toda equations

$$
\begin{align*}
& \left(\frac{1}{2} D_{x}^{2}+3 m\right) H_{m, n} \cdot H_{m, n}=H_{m+1, n} H_{m-1, n} \\
& \left(\frac{1}{2} D_{x}^{2}-3 n\right) H_{m, n} \cdot H_{m, n}=H_{m, n+1} H_{m, n-1} \tag{4.19}
\end{align*}
$$

with initial condition

$$
\begin{equation*}
H_{0,0}=H_{1,0}=H_{0,1}=1 \quad \text { and } \quad H_{1,1}=3 x \tag{4.20}
\end{equation*}
$$

We remark that $H_{m, 1}(x)$ and $H_{1, m}(m=0,1,2, \ldots)$ coincide with the Hermite polynomials up to rescaling. We will call $H_{m, n}(x)(m, n \geq 0)$ the generalized Hermite polynomials. The fact that $\tau_{m, n}(x)$ are expressed as in 4.18 will be proved in Section 5 in the course of the proof of Theorem 4.6 below.

The Okamoto polynomials $Q_{m, n}(x)(m, n \in \mathbb{Z})$ and the generalized Hermite polynomials $H_{m, n}(x)(m, n \geq 0)$ are in fact expressible in terms of Schur functions. We recall the definition of Schur functions in order to make this statement precise.

A partition $\lambda=\left(\lambda_{1}, \lambda_{2}, \ldots\right)$ (or a Young diagram) is a sequence of nonnegative integers such that $\lambda_{1} \geq \lambda_{2} \geq \cdots \geq 0$ and that $\lambda_{i}=0$ for $i \gg 0$.

The number of nonzero parts $\lambda_{i}$ is called the length of $\lambda$ and denoted by $l(\lambda)$. For each partition $\lambda$, we define the Schur function $S_{\lambda}(t)=S_{\lambda}\left(t_{1}, t_{2}, \ldots\right)$ by

$$
\begin{equation*}
S_{\lambda}(t)=\operatorname{det}\left(p_{\lambda_{i}-i+j}(t)\right)_{1 \leq i, j \leq l(\lambda)} \tag{4.21}
\end{equation*}
$$

where $p_{n}(t)$ are the polynomials in $t$ determined by the generating function

$$
\begin{equation*}
\exp \left(\sum_{k=1}^{\infty} t_{k} z^{k}\right)=\sum_{n=0}^{\infty} p_{n}(t) z^{n} \tag{4.22}
\end{equation*}
$$

(We set $p_{n}(t)=0$ for $n<0$.) Note that $p_{n}(t)$ can be defined equivalently by

$$
\begin{equation*}
p_{n}(t)=\sum_{k_{1}+2 k_{2}+\cdots+n k_{n}=n} \frac{t_{1}^{k_{1}} t_{2}^{k_{2}} \cdots t_{n}^{k_{n}}}{k_{1}!k_{2}!\cdots k_{n}!} \tag{4.23}
\end{equation*}
$$

We say that a subset $M \subset \mathbb{Z}$ is a Maya diagram if

$$
\begin{equation*}
m \in M \quad(m \ll 0) \quad \text { and } \quad m \notin M \quad(m \gg 0) \tag{4.24}
\end{equation*}
$$

To each Maya diagram $M=\left\{\ldots, m_{3}, m_{2}, m_{1}\right\} \quad\left(\cdots<m_{3}<m_{2}<m_{1}\right)$, one can associate a unique partition $\lambda=\left(\lambda_{1}, \lambda_{2}, \ldots\right)$ such that $m_{i}-m_{i+1}=$ $\lambda_{i}-\lambda_{i+1}+1$ for $i=1,2, \ldots$. Note that all the Maya diagrams $M+k=$ $\left\{\ldots, m_{2}+k, m_{1}+k\right\}(k \in \mathbb{Z})$ obtained from $M=\left\{\ldots, m_{3}, m_{2}, m_{1}\right\}$ by shifting define the same partition by this correspondence. For each pair $(m, n)$ of integers, we define the Maya diagram $M(m, n)$ as follows:

$$
\begin{equation*}
M(m, n)=3 D_{m} \cup\left(3 D_{n}+1\right) \cup\left(3 D_{0}+2\right) \tag{4.25}
\end{equation*}
$$

where

$$
\begin{equation*}
D_{l}=\{n \in \mathbb{Z} \mid n<l\} \quad(l \in \mathbb{Z}) \tag{4.26}
\end{equation*}
$$

We denote by $\lambda(m, n)$ the partition corresponding to $M(m, n)$. Partitions of the form $\lambda(m, n)(m, n \in \mathbb{Z})$ are called the 3-reduced partitions. We remark that a partition $\lambda$ is 3-reduced if and only if $\lambda$ has no hook with length of a multiple of 3. Also, Schur functions $S_{\lambda(m, n)}(t)$ for 3 -reduced partitions are called 3-reduced Schur functions. It is known that a Schur function $S_{\lambda}(t)$ is 3 -reduced if and only if

$$
\begin{equation*}
\partial_{t_{3 n}} S_{\lambda}(t)=0 \quad \text { for all } \quad n=1,2, \cdots \tag{4.27}
\end{equation*}
$$

ThEOREM 4.5. Each Okamoto polynomial $Q_{m, n}(x)(m, n \in \mathbb{Z})$ is a monic polynomial of degree $m^{2}+n^{2}-m n-m$ with integer coefficients. It is expressed by the 3-reduced Schur function as

$$
\begin{equation*}
Q_{m, n}(x)=N_{m, n} S_{\lambda(m, n)}\left(x, \frac{1}{2}, 0,0, \ldots\right) \tag{4.28}
\end{equation*}
$$

where $N_{m, n}$ is a positive integer determined by the hook-length formula.
Example. The Maya diagram $M(3,2)$ is obtained from $D_{3}, D_{2}, D_{0}$ as follows.


Hence we have $M(3,2)=\{\ldots,-2,-1,0,1,3,4,6\}$ and

$$
\begin{equation*}
\lambda(3,2)=(2,1,1)=\square \tag{4.29}
\end{equation*}
$$

In this case, the Schur function $S_{(2,1,1)}(t)$ and the Okamoto polynomials $Q_{3,2}(x)$ are

$$
\begin{align*}
S_{(2,1,1)}(t) & =\frac{1}{8} t_{1}^{4}-\frac{1}{2} t_{1}^{2} t_{2}-\frac{1}{2} t_{2}^{2}+t_{4}  \tag{4.30}\\
Q_{3,2}(x) & =x^{4}-2 x^{2}-1
\end{align*}
$$

respectively. A typical sequence of 3-reduced partitions is given by

$$
\begin{equation*}
\lambda(m, 0)=(2 m-2,2 m-4, \ldots, 2) \quad \text { for } \quad m>0 \tag{4.31}
\end{equation*}
$$

which corresponds to the Okamoto polynomials $Q_{m}(x)$ for $m>0$. For other examples, see Figure 3 in the next section.

The generalized Hermite polynomials $H_{m, n}(x)$ are expressed by the Schur functions for rectangular Young diagrams $\lambda=\left(n^{m}\right)=(n, n, \ldots, n, 0$, $0, \ldots$.

THEOREM 4.6. Each generalized Hermite polynomial $H_{m, n}(x)(m, n \geq$ $0)$ is a polynomial of degree $m n$ with rational coefficients. It can be written as

$$
\begin{equation*}
H_{m, n}(x)=C_{m, n} S_{\left(n^{m}\right)}\left(x, \frac{1}{6}, 0,0, \ldots\right) \tag{4.32}
\end{equation*}
$$

where the normalization constant is given by

$$
\begin{equation*}
C_{m, n}=(-1)^{n(n-1) / 2} 3^{(m+n)(m+n-1) / 2}(m+n-1)^{!} \tag{4.33}
\end{equation*}
$$

with $n!=n!(n-1)!(n-2)!\cdots 2!1$ !.
The relationship between the sequences $H_{1, n}(x), H_{m, 1}(x)$ and the ordinary Hermite polynomials $H_{n}(x)$ is obvious since

$$
\begin{align*}
& \sum_{n=0}^{\infty} S_{(n)}\left(x, \frac{1}{6}, 0,0, \ldots\right) z^{n}=\exp \left(x z+\frac{1}{6} z^{2}\right) \\
& \sum_{m=0}^{\infty} S_{\left(1^{m}\right)}\left(x, \frac{1}{6}, 0,0, \ldots\right) z^{m}=\exp \left(x z-\frac{1}{6} z^{2}\right) \tag{4.34}
\end{align*}
$$

while the Hermite polynomials have the generating function

$$
\begin{equation*}
\sum_{n=0}^{\infty} \frac{1}{n!} H_{n}(x) z^{n}=\exp \left(2 x z-z^{2}\right) \tag{4.35}
\end{equation*}
$$

The proof of Theorems 4.5 and 4.6 will be given in the next section.

## $\S 5$. Proof of Theorems 4.5 and 4.6

The Hirota bilinear equations for our $\tau$-functions (Theorem 3.1) arise naturally from the so-called modified KP hierarchy [2] by certain similarity reduction (see also [5], [10]). This fact is the key to the proof of Theorems 4.5 and 4.6.

Consider two functions $G_{0}\left(t_{1}, t_{2}\right)$ and $G_{1}\left(t_{1}, t_{2}\right)$ in the two variables $\left(t_{1}, t_{2}\right)$, and suppose that they satisfy the following Hirota bilinear equation

$$
\begin{equation*}
\left(D_{t_{1}}^{2}+D_{t_{2}}\right) G_{0}\left(t_{1}, t_{2}\right) \cdot G_{1}\left(t_{1}, t_{2}\right)=0 \tag{5.1}
\end{equation*}
$$

Introducing the degrees of $t_{1}, t_{2}$ by $\operatorname{deg} t_{1}=1$ and $\operatorname{deg} t_{2}=2$, we assume that each $G_{i}$ is homogeneous of degree $d_{i} \in \mathbb{C}$ for $i=0,1$ :

$$
\begin{equation*}
\left(t_{1} \partial_{t_{1}}+2 t_{2} \partial_{t_{2}}\right) G_{i}\left(t_{1}, t_{2}\right)=d_{i} G_{i}\left(t_{1}, t_{2}\right) \tag{5.2}
\end{equation*}
$$

Fixing a constant $k \in \mathbb{C}^{\times}$, we define the functions $\tau_{0}(x), \tau_{1}(x)$ in one variable by

$$
\begin{equation*}
\tau_{i}(x)=G_{i}(x, k) \quad\left(k \in \mathbb{C}^{\times} ; i=0,1\right) \tag{5.3}
\end{equation*}
$$

Formally, each $G_{i}$ is recovered by the formula

$$
\begin{equation*}
G_{i}\left(t_{1}, t_{2}\right)=\left(\frac{t_{2}}{k}\right)^{\frac{d_{i}}{2}} \tau_{i}\left(\frac{t_{1}}{\sqrt{t_{2} / k}}\right) \tag{5.4}
\end{equation*}
$$

Then it is easy to check
LEMMA 5.1. Under the similarity condition 5.2, the equation 5.1 for the pair $G_{0}\left(t_{1}, t_{2}\right), G_{1}\left(t_{1}, t_{2}\right)$ is equivalent to the Hirota bilinear equation

$$
\begin{equation*}
\left(2 k D_{x}^{2}-x D_{x}+d_{0}-d_{1}\right) \tau_{0}(x) \cdot \tau_{1}(x)=0 \tag{5.5}
\end{equation*}
$$

for $\tau_{0}(x), \tau_{1}(x)$.
From this lemma with $k=1 / 2$, we immediately have
Proposition 5.2. The fourth Painlevé equation 3.7 for the triple of $\tau$-functions $\tau_{0}(x), \tau_{1}(x), \tau_{2}(x)$ is equivalent to the similarity reduction of the Hirota equations

$$
\begin{equation*}
\left(D_{t_{1}}^{2}+D_{t_{2}}\right) G_{i}\left(t_{1}, t_{2}\right) \cdot G_{i+1}\left(t_{1}, t_{2}\right)=0 \quad(i=0,1,2) \tag{5.6}
\end{equation*}
$$

for three functions $G_{i}\left(t_{1}, t_{2}\right)(i=0,1,2)$ in two variables. The similarity condition is given by

$$
\begin{equation*}
G_{i}\left(t_{1}, t_{2}\right)=\left(2 t_{2}\right)^{d_{i} / 2} \tau_{i}\left(\frac{t_{1}}{\sqrt{2 t_{2}}}\right) \quad(i=0,1,2) \tag{5.7}
\end{equation*}
$$

and the parameters are related by

$$
\begin{align*}
& \alpha_{0}=1-2 d_{0}+d_{1}+d_{2}, \\
& \alpha_{1}=1+d_{0}-2 d_{1}+d_{2}  \tag{5.8}\\
& \alpha_{2}=1+d_{0}+d_{1}-2 d_{2}
\end{align*}
$$

Recall that the (first) modified KP hierarchy [2] is the following system of Hirota bilinear equations for a pair of $\tau$-functions $\tau_{0}(t)$ and $\tau_{1}(t)$ in infinite time variables $t=\left(t_{1}, t_{2}, \ldots\right)$ :

$$
\begin{equation*}
\sum_{n=0}^{\infty} p_{n}(-2 s) p_{n+2}\left(\widetilde{D}_{t}\right) \exp \left(\sum_{m=1}^{\infty} s_{m} D_{t_{m}}\right) \tau_{0}(t) \cdot \tau_{1}(t)=0 \tag{5.9}
\end{equation*}
$$

where $s=\left(s_{1}, s_{2}, \ldots\right)$ are parameters and $\widetilde{D}_{t}=\left(D_{t_{1}} / 1, D_{t_{2}} / 2, \ldots\right)$. The constant term of 5.9 with respect to $s$ implies the bilinear equation

$$
\begin{equation*}
\left(D_{t_{1}}^{2}+D_{t_{2}}\right) \tau_{0}(t) \cdot \tau_{1}(t)=0 \tag{5.10}
\end{equation*}
$$

which is nothing but the equation 5.1 discussed above. For the proof of Theorems 4.5 and 4.6 , we will recall the following fact about Schur functions from the theory of KP hierarchy. Let $X_{m}=X_{m}\left(t ; \partial_{t}\right)(m \in \mathbb{Z})$ be the vertex operators of the KP hierarchy defined by the generating function

$$
\begin{equation*}
X(z)=\sum_{m \in \mathbb{Z}} X_{m} z^{m}=\exp \left(\sum_{k=1}^{\infty} t_{k} z^{k}\right) \exp \left(-\sum_{k=1}^{\infty} \frac{z^{-k}}{k} \partial_{t_{k}}\right) \tag{5.11}
\end{equation*}
$$

Then we have
LEmma 5.3. For any partition $\lambda$ and $k \in \mathbb{Z}$, the pair $\tau_{0}(t)=S_{\lambda}(t)$ and $\tau_{1}(t)=X_{k} S_{\lambda}(t)$ solves the first modified KP hierarchy 5.9. In particular we have

$$
\begin{equation*}
\left(D_{t_{1}}^{2}+D_{t_{2}}\right) \tau_{0}(t) \cdot \tau_{1}(t)=0 \tag{5.12}
\end{equation*}
$$

We will give a proof of this lemma in Appendix for completeness.
All the Schur functions $S_{\lambda}(t)$ are obtained from $S_{\emptyset}(t)=1$ by applying vertex operators repeatedly:

$$
\begin{equation*}
S_{\lambda}(t)=X_{\lambda_{1}} \cdots X_{\lambda_{n}} .1 \tag{5.13}
\end{equation*}
$$

for any partition $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}, 0, \ldots\right)$. The action of vertex operators on Schur functions can be computed by 5.13 together with the commutation relations

$$
\begin{equation*}
X_{k} X_{l}=-X_{l-1} X_{k+1}, \quad X_{k} \cdot 1=0 \quad(k<0), \quad X_{0} \cdot 1=1 \tag{5.14}
\end{equation*}
$$

where $k, l \in \mathbb{Z}$. (See Appendix.) A more systematic way is to use Maya diagrams. For a given Maya diagram $M$, let $\lambda$ be the corresponding partition and suppose that $l(\lambda) \leq n$. Then we have

$$
X_{k} \cdot S_{\lambda}(t)=\left\{\begin{array}{cl} 
\pm S_{\mu}(t) & \text { if } \quad k+n \notin M  \tag{5.15}\\
0 & \text { if } \quad k+n \in M
\end{array}\right.
$$

for each $k \in \mathbb{Z}$. Here $\mu$ stands for the partition corresponding to the Maya diagram $M \cup\{k+n\}$. The sign in this formula is determined by the parity of the number of integers $m \in M$ such that $m>k+n$.

Proof of Theorem 4.5 for Okamoto polynomials. By using the formula 5.15 , one can compute how the 3 -reduced Schur functions are transformed by vertex operators.

Lemma 5.4. As to the action of the vertex operators, we have the following two types of cyclic relations among 3-reduced Schur functions

$$
\begin{align*}
X_{2 m-n} \cdot S_{\lambda(m, n)}(t) & = \pm S_{\lambda(m+1, n)}(t) \\
X_{2 n-m} \cdot S_{\lambda(m+1, n)}(t) & = \pm S_{\lambda(m+1, n+1)}(t),  \tag{5.16}\\
X_{-m-n} \cdot S_{\lambda(m+1, n+1)}(t) & = \pm S_{\lambda(m, n)}(t)
\end{align*}
$$

and

$$
\begin{align*}
X_{2 n-m+1} \cdot S_{\lambda(m, n)}(t) & = \pm S_{\lambda(m, n+1)}(t) \\
X_{2 m-n-1} \cdot S_{\lambda(m, n+1)}(t) & = \pm S_{\lambda(m+1, n+1)}(t)  \tag{5.17}\\
X_{-m-n} \cdot S_{\lambda(m+1, n+1)}(t) & = \pm S_{\lambda(m, n)}(t)
\end{align*}
$$

for any $m, n \in \mathbb{Z}$.
Example. Consider the 3-reduced partitions $\lambda(3,1)=(3,1), \lambda(4,1)=$ $(5,3,1)$ and $\lambda(4,2)=(4,2,1,1)$. For this triple of partitions, we have a 'cycle' of 3-reduced Schur functions

which is an example of 5.16 for $(m, n)=(3,1)$. Notice also that the index of each vertex operator $X_{k}$ represents the difference of degrees of Schur functions.

From Lemma 5.4 together with Lemma 5.3, we obtain two types of triples of 3-reduced Schur functions satisfying the bilinear equations of Proposition 5.2. Note that the 3-reduced Schur function $S_{\lambda(m, n)}$ is homogeneous of degree

$$
\begin{equation*}
d_{m, n}=|\lambda(m, n)|=m^{2}+n^{2}-m n-m \tag{5.18}
\end{equation*}
$$

with respect to the degree defined by $\operatorname{deg} t_{i}=i(i=1,2, \ldots)$. Set

$$
\begin{equation*}
s_{m, n}(x)=S_{\lambda(m, n)}\left(x, \frac{1}{2}, 0,0, \ldots\right) \tag{5.19}
\end{equation*}
$$

for any $m, n \in \mathbb{Z}$. Then, by combining Lemma 5.4 and Proposition 5.2, we have

Proposition 5.5. (1) For any $m, n \in \mathbb{Z}$, the triple

$$
\begin{equation*}
\left(s_{m, n}(x), s_{m+1, n}(x), s_{m+1, n+1}(x)\right) \tag{5.20}
\end{equation*}
$$

solves the fourth Painlevé equation 3.7 with the parameters

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2}\right)=(3 m+1,3(n-m)+1,-3 n+1) \tag{5.21}
\end{equation*}
$$

(2) For any $m, n \in \mathbb{Z}$, the triple

$$
\begin{equation*}
\left(s_{m, n}(x), s_{m, n+1}(x), s_{m+1, n+1}(x)\right) \tag{5.22}
\end{equation*}
$$

solves the fourth Painlevé equation 3.7 with parameters

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2}\right)=(3 n+2,3(m-n)-1,-3 m+2) \tag{5.23}
\end{equation*}
$$

We remark that, in the coordinates $\left(v_{1}, v_{2}, v_{3}\right)$ of the parameter space $V$ as in 1.10, the triples of $\tau$-functions in this proposition give rise to solutions with parameters

$$
\begin{equation*}
\left(v_{1}, v_{2}, v_{3}\right)=\left(\frac{1}{3}, 0,-\frac{1}{3}\right)-m\left(\frac{2}{3},-\frac{1}{3},-\frac{1}{3}\right)-n\left(-\frac{1}{3}, \frac{2}{3},-\frac{1}{3}\right) \tag{5.24}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(v_{1}, v_{2}, v_{3}\right)=\left(0, \frac{1}{3},-\frac{1}{3}\right)-m\left(-\frac{1}{3}, \frac{2}{3},-\frac{1}{3}\right)-n\left(\frac{2}{3},-\frac{1}{3},-\frac{1}{3}\right) \tag{5.25}
\end{equation*}
$$

respectively.
It is clear that each triple of $\tau$-functions of Proposition 5.5 defines a rational solution of 1.11 in $f$-variables. For each $\left(\alpha_{0}, \alpha_{1}, \alpha_{2}\right)$ of this proposition, the fourth Painlevé equation 1.11 has a unique rational solution by [4]. Hence we conclude that each $s_{m, n}(x)$ is a constant multiple of the $\tau$-function $\tau_{m, n}(x)$ for the solution (A) of 4.11. This shows that $\tau_{m, n}(x)$ are in fact polynomials in $x$. The assertion that the Okamoto polynomials $\tau_{m, n}(x)=Q_{m, n}(x)$ are monic polynomials with integer coefficients follows either from the Bäcklund transformations or from the Toda equations of Proposition 4.2.

Proof of Theorem 4.6 for generalized Hermite polynomials. By the formula 5.15 , we have

LEMMA 5.6. Under the action of vertex operators, we have the following relations among Schur functions for rectangular Young diagrams:

$$
\begin{align*}
X_{-m} \cdot S_{\left((n+1)^{m}\right)}(t) & =(-1)^{m} S_{\left(n^{m}\right)}(t) \\
X_{n-m} \cdot S_{\left((n+1)^{m}\right)}(t) & =(-1)^{m} S_{\left(n^{(m+1)}\right)}(t),  \tag{5.26}\\
X_{n} \cdot S_{\left(n^{m}\right)}(t) & =S_{\left(n^{(m+1)}\right)}(t),
\end{align*}
$$

for $m, n \geq 0$.
For each $m, n \geq 0$, let

$$
\begin{equation*}
h_{m, n}(x)=S_{\left(n^{m}\right)}\left(x, \frac{1}{6}, 0,0, \ldots\right) \tag{5.27}
\end{equation*}
$$

be the specialization of the Schur function $S_{\lambda}(t)=S_{\left(n^{m}\right)}(t)$ associated with rectangular Young diagram $\lambda=\left(n^{m}\right)$. Then by Lemma 5.1 (with $k=1 / 6$ ), we have

Lemma 5.7.

$$
\begin{align*}
& \left(D_{x}^{2}-3 x D_{x}+3 m\right) h_{m, n+1} \cdot h_{m, n}=0 \\
& \left(D_{x}^{2}-3 x D_{x}+3(m-n)\right) h_{m, n+1} \cdot h_{m+1, n}=0  \tag{5.28}\\
& \left(D_{x}^{2}-3 x D_{x}-3 n\right) h_{m, n} \cdot h_{m+1, n}=0
\end{align*}
$$

These relations do not fit directly for the triple of $\tau$-functions as in 3.7 since they do not make a 'cycle'. This problem can be repaired however by changing the normalization of $h_{m, n}$ as follows:

$$
\begin{equation*}
u_{m, n}(x)=\exp \left(-\frac{x^{4}}{12}+\frac{m-n}{2} x^{2}\right) h_{m, n}(x) \tag{5.29}
\end{equation*}
$$

Then we have

## Proposition 5.8.

$$
\begin{align*}
& \left(D_{x}^{2}-x D_{x}+m+2 n+1\right) u_{m, n+1} \cdot u_{m, n}=0 \\
& \left(D_{x}^{2}-x D_{x}+m-n\right) u_{m+1, n} \cdot u_{m, n+1}=0  \tag{5.30}\\
& \left(D_{x}^{2}-x D_{x}-2 m-n-1\right) u_{m, n} \cdot u_{m+1, n}=0
\end{align*}
$$

Namely, the triple

$$
\begin{equation*}
\left(\tau_{0}, \tau_{1}, \tau_{2}\right)=\left(u_{m, n}, u_{m+1, n}, u_{m, n+1}\right) \tag{5.31}
\end{equation*}
$$

solves the fourth Painlevé equation 3.7 with parameters

$$
\begin{equation*}
\left(\alpha_{0}, \alpha_{1}, \alpha_{2}\right)=(3(m+n+1),-3 m,-3 n) \tag{5.32}
\end{equation*}
$$

Proof. Note that the Hirota bilinear equations have the following formulas of Leibniz type:

$$
\begin{aligned}
& D_{x}\left(g_{1} u_{1} \cdot g_{2} u_{2}\right)=D_{x}\left(g_{1} \cdot g_{2}\right) u_{1} u_{2}+g_{1} g_{2} D_{x}\left(u_{1} \cdot u_{2}\right), \\
& D_{x}^{2}\left(g_{1} u_{1} \cdot g_{2} u_{2}\right)=D_{x}^{2}\left(g_{1} \cdot g_{2}\right) u_{1} u_{2}+2 D_{x}\left(g_{1} \cdot g_{2}\right) D_{x}\left(u_{1} \cdot u_{2}\right) \\
& \\
& +g_{1} g_{2} D_{x}^{2}\left(u_{1} \cdot u_{2}\right) .
\end{aligned}
$$

Applying these to $g_{i}=\exp \left(x^{4} / 12+a_{i} x^{2} / 2\right)(i=1,2)$, we have

$$
\begin{aligned}
& \left(D_{x}^{2}-3 x D_{x}+\beta\right)\left(g_{1} u_{1} \cdot g_{2} u_{2}\right) \\
= & g_{1} g_{2}\left\{D_{x}^{2}+\left(2 a_{12}-3\right) x D_{x}+\left(2-3 a_{12}+a_{12}^{2}\right) x^{2}+\left(a_{1}+a_{2}\right)+\beta\right\} u_{1} \cdot u_{2}
\end{aligned}
$$

where $a_{12}=a_{1}-a_{2}$. The equations 5.30 can be checked easily by using this formula.

We remark that, in the coordinates $\left(v_{1}, v_{2}, v_{3}\right)$ of $V$, the solutions of Proposition 5.8 have the parameters

$$
\begin{equation*}
\left(v_{1}, v_{2}, v_{3}\right)=-m\left(\frac{2}{3},-\frac{1}{3},-\frac{1}{3}\right)+n\left(-\frac{1}{3},-\frac{1}{3}, \frac{2}{3}\right) \quad(m, n=0,1,2, \ldots) \tag{5.33}
\end{equation*}
$$

As in the case of Okamoto polynomials, we see that each $\tau_{m, n}(m \geq$ $n \geq 0$ ) for the solution (B) of 4.11 is a constant multiple of $u_{m-n, n}$ by comparing the parameters. Hence we see that $\tau_{m, n}(x)$ has the expression of 4.18. The only problem remaining is to fix the constant factors. The leading coefficient of the polynomial $H_{m, n}(x)$ is given by

$$
\begin{equation*}
(-1)^{n(n-1) / 2}(m-1)^{!}(n-1)^{!} 3^{(m+n)(m+n-1) / 2} \tag{5.34}
\end{equation*}
$$

which can be determined inductively by the Toda equations of Theorem 4.4. On the other hand, the leading coefficient of $h_{m, n}(x)$ is determined as

$$
\begin{equation*}
(m-1)^{!}(n-1)^{!} /(m+n-1)^{!} \tag{5.35}
\end{equation*}
$$



Figure 3: Okamoto polynomials on the $A_{2}$-lattice
by the hook-length formula. Hence we have Theorem 4.6.

We show in Figures 3 and 4 below how the $\tau$-functions for rational solutions are arranged on the $A_{2}$-lattice. Also, we include some examples of Okamoto polynomials and generalized Hermite polynomials of small degrees.
Okamoto polynomials. In the following, we use the notation $Q_{\lambda}(x)=$ $Q_{m, n}(x)$ for the Okamoto polynomial associated with the 3-reduced partition $\lambda=\lambda(m, n)$. We give below some examples of Okamoto polynomials $Q_{\lambda}(x)$.

$$
\begin{aligned}
& Q_{(0)}=Q_{1}=1, \quad Q_{(1)}=R_{1}=x, \quad Q_{(2)}=Q_{2}=1+x^{2} \\
& Q_{(1,1)}=-1+x^{2}, \quad Q_{(3,1)}=R_{2}=-1+2 x^{2}+x^{4} \\
& Q_{(2,1,1)}=-1-2 x^{2}+x^{4}, \quad Q_{(3,1,1)}=-5 x+x^{5} \\
& Q_{(4,2)}=Q_{3}=5+5 x^{2}+5 x^{4}+x^{6}, \quad Q_{(2,2,1,1)}=-5+5 x^{2}-5 x^{4}+x^{6}
\end{aligned}
$$



Figure 4: Generalized Hermite polynomials on the $A_{2}$-lattice

$$
\begin{aligned}
& Q_{(4,2,1,1)}=-7-14 x^{4}+x^{8}, \quad Q_{(5,3,1)}=R_{3}=-35 x+14 x^{5}+8 x^{7}+x^{9}, \\
& Q_{(5,3,1,1)}=25-75 x^{2}-50 x^{4}-10 x^{6}+5 x^{8}+x^{10} \\
& Q_{(6,4,2)}=Q_{4}=175+350 x^{2}+175 x^{4}+140 x^{6}+65 x^{8}+14 x^{10}+x^{12} \\
& Q_{(7,5,3,1)}=R_{4}=1225-4900 x^{2}-4900 x^{4}-980 x^{6} \\
& \quad+350 x^{8}+420 x^{10}+140 x^{12}+20 x^{14}+x^{16} \\
& Q_{(8,6,4,2)}=Q_{5}=67375+134750 x^{2}+202125 x^{4}+107800 x^{6} \\
& \quad+42350 x^{8}+20020 x^{10}+8050 x^{12}+2200 x^{14}+355 x^{16}+30 x^{18}+x^{20}
\end{aligned}
$$

Note that the original Okamoto polynomials are given by

$$
\begin{array}{ll}
Q_{n}=Q_{(2 n-2,2 n-4, \ldots, 4,2)} & (n>0), \\
R_{n}=Q_{(2 n-1,2 n-3, \ldots, 3,1)} & (n>0),
\end{array} \quad R_{-n}=Q_{(n, n, \ldots, 2,2,1,1)} \quad(n \geq 0), ~ Q_{(n, n+1, n+1, \ldots, 1,1)} \quad(n \geq 0) .
$$

Generalized Hermite polynomials. The polynomials $H_{n, 1}(x)$ and $H_{1, n}(x)$ coincide with the ordinary Hermite polynomials up to rescaling.
$H_{0,0}=1, \quad H_{1,0}=1, \quad H_{2,0}=3, \quad H_{3,0}=2^{1} 3^{3}, \quad H_{4,0}=2^{2} 3^{7}$,

$$
\begin{aligned}
& H_{0,1}=1, \quad H_{1,1}=3 x, \quad H_{2,1}=3^{3}\left(-\frac{1}{3}+x^{2}\right) \\
& H_{3,1}=2^{1} 3^{6}\left(-x+x^{3}\right), \quad H_{4,1}=2^{2} 3^{11}\left(\frac{1}{3}-2 x^{2}+x^{4}\right), \\
& H_{0,2}=-3, \quad H_{1,2}=-3^{3}\left(\frac{1}{3}+x^{2}\right), \quad H_{2,2}=-3^{6}\left(\frac{1}{3}+x^{4}\right), \\
& H_{3,2}=-2^{1} 3^{10}\left(\frac{1}{3}+x^{2}-x^{4}+x^{6}\right), \quad H_{4,2}=-2^{2} 3^{16}\left(\frac{5}{9}+\frac{10}{3} x^{4}-\frac{8}{3} x^{6}+x^{8}\right), \\
& H_{0,3}=-2^{1} 3^{3}, \quad H_{1,3}=-2^{1} 3^{6}\left(x+x^{3}\right), \\
& H_{2,3}=-2^{1} 3^{10}\left(-\frac{1}{3}+x^{2}+x^{4}+x^{6}\right), \quad H_{3,3}=-2^{2} 3^{15}\left(\frac{-5}{3} x+2 x^{5}+x^{9}\right), \\
& H_{4,3}=-2^{3} 3^{22}\left(\frac{25}{27}-\frac{50}{9} x^{2}-\frac{25}{9} x^{4}-\frac{20}{9} x^{6}+5 x^{8}-2 x^{10}+x^{12}\right) \\
& H_{0,4}=2^{2} 3^{7}, \quad H H_{1,4}=2^{2} 3^{11}\left(\frac{1}{3}+2 x^{2}+x^{4}\right), \\
& H_{2,4}=2^{2} 3^{16}\left(\frac{5}{9}+\frac{10}{3} x^{4}+\frac{8}{3} x^{6}+x^{8}\right), \\
& H_{3,4}=2^{3} 3^{22}\left(\frac{25}{27}+\frac{50}{9} x^{2}-\frac{25}{9} x^{4}+\frac{20}{9} x^{6}+5 x^{8}+2 x^{10}+x^{12}\right), \\
& H_{4,4}=2^{4} 3^{30}\left(\frac{875}{243}+\frac{3500}{81} x^{4}-\frac{50}{9} x^{8}+\frac{20}{3} x^{12}+x^{16}\right) .
\end{aligned}
$$

## §A. Appendix

In this Appendix, we give a brief summary of relevant facts on Schur functions and their relation to KP-hierarchy for the sake of reference.

## A.1. Schur functions.

A partition $\lambda=\left(\lambda_{1}, \lambda_{2}, \ldots\right)$ is a sequence of non-negative integers such that $\lambda_{1} \geq \lambda_{2} \geq \cdots \geq 0$ and that $\lambda_{i}=0$ for $i \gg 0$. The number of nonzero $\lambda_{i}$ is called the length of $\lambda$ and denoted by $l(\lambda)$. For each partition $\lambda$, the Schur function $S_{\lambda}(t)=S_{\lambda}\left(t_{1}, t_{2}, \ldots\right)$ is defined as follows:

$$
\begin{equation*}
S_{\lambda}(t)=\operatorname{det}\left(p_{\lambda_{i}-i+j}(t)\right)_{1 \leq i, j \leq l(\lambda)} \tag{A.1}
\end{equation*}
$$

where $p_{n}(t)$ are the polynomials defined by the generating function

$$
\begin{equation*}
\exp \left(\sum_{k=1}^{\infty} t_{k} z^{k}\right)=\sum_{n=0}^{\infty} p_{n}(t) z^{n} \tag{A.2}
\end{equation*}
$$

Usually, the Schur functions are defined as the following character polynomials of the general linear group $G L(N, \mathbb{C})(N \geq l(\lambda))$ :

$$
\begin{equation*}
s_{\lambda}\left(x_{1}, \ldots, x_{N}\right)=\frac{\operatorname{det}\left(x_{j}^{\lambda_{i}+\delta_{i}}\right)}{\operatorname{det}\left(x_{j}^{\delta_{i}}\right)} \tag{A.3}
\end{equation*}
$$

where $\delta_{i}=N-i(i=1, \ldots, N)$. The polynomials $S_{\lambda}(t)$ and $s_{\lambda}(x)$ are related by $S_{\lambda}(t)=s_{\lambda}(x)$, where $t_{k}=\sum_{i=1}^{n}\left(x_{i}^{k}\right) / k$. In this context, the formula A. 1 above is the Jacobi-Trudi formula representing $s_{\lambda}(x)$ in terms of complete homogeneous symmetric functions.

The coefficients of $S_{\lambda}(t)$ with respect to the $t$-variables are related with irreducible character $\pi_{\lambda}$ of the symmetric group $\mathfrak{S}_{n}$ of degree $n=|\lambda|=$ $\sum_{i} \lambda_{i}$ as follows:

$$
\begin{equation*}
S_{\lambda}(t)=\sum_{m_{1}, m_{2}, \ldots \geq 0} \pi_{\lambda}\left(1^{m_{1}} 2^{m_{2}} \cdots\right) \frac{t_{1}^{m_{1}}}{m_{1}!} \frac{t_{2}^{m_{2}}}{m_{2}!} \cdots \tag{A.4}
\end{equation*}
$$

where $\pi_{\lambda}\left(1^{m_{1}} 2^{m_{2}} \cdots\right)$ is the character value on the conjugate class of cycle type $\left(1^{m_{1}} 2^{m_{2}} \cdots\right)$. In particular, the coefficient of $t_{1}^{n}$ is given by the hooklength formula

$$
\begin{equation*}
\frac{\pi_{\lambda}\left(1^{n}\right)}{n!}=\prod_{s \in \lambda} \frac{1}{h(s)} \tag{A.5}
\end{equation*}
$$

where $h(s)=\lambda_{i}+\lambda_{j}^{\prime}-i-j+1, \lambda^{\prime}$ being the conjugate partition, denotes the hook-length of $\lambda$ at $s=(i, j)$.

## A.2. (Modified) KP hierarchy

In the following, we use the notation

$$
\begin{equation*}
\xi(z, t)=\sum_{n=1}^{\infty} t_{n} z^{n}, \quad \xi\left(z^{-1}, \widetilde{\partial}_{t}\right)=\sum_{n=1}^{\infty} \frac{z^{-n}}{n} \partial_{t_{n}} \tag{A.6}
\end{equation*}
$$

Consider the operators $V_{k}=V_{k}(z, t)(k \in \mathbb{Z})$ defined by

$$
\begin{equation*}
V_{k}=e^{k \xi(z, t)} e^{-k \xi\left(z^{-1}, \widetilde{\partial}_{t}\right)} \tag{A.7}
\end{equation*}
$$

For each $m \in \mathbb{Z}$, we define the operators $X_{m}$ and $X_{m}^{*}$ as the coefficient of $z^{m}$ in $V_{1}$ and $V_{-1}$, respectively:

$$
\begin{align*}
& V_{1}(z, t)=X(z, t)=\sum_{m \in \mathbb{Z}} X_{m} z^{m} \\
& V_{-1}(z, t)=X^{*}(z, t)=\sum_{m \in \mathbb{Z}} X_{m}^{*} z^{m} \tag{A.8}
\end{align*}
$$

By using the formula
(A.9) $V_{k}(z, t) V_{l}(w, t)=\left(1-\frac{w}{z}\right)^{k l} e^{k \xi(z, t)+l \xi(w, t)} e^{-k \xi\left(z^{-1}, \widetilde{\partial}_{t}\right)-l \xi\left(w^{-1}, \widetilde{\partial}_{t}\right)}$,
we obtain
Lemma A.1. The vertex operators $X_{m}$ and $X_{m}^{*}(m \in \mathbb{Z})$ satisfy the following anti-commutation relations:

$$
\begin{align*}
& X_{m} X_{n}+X_{n-1} X_{m+1}=0 \\
& X_{m}^{*} X_{n}^{*}+X_{n-1}^{*} X_{m+1}^{*}=0  \tag{A.10}\\
& X_{m} X_{n}^{*}+X_{n+1}^{*} X_{m-1}=\delta_{m+n, 0}
\end{align*}
$$

Proposition A.2. For any partition $\lambda=\left(\lambda_{1}, \lambda_{2}, \ldots\right)$ of length $l(\lambda) \leq$ l, we have

$$
\begin{equation*}
S_{\lambda}(t)=X_{\lambda_{1}} \cdots X_{\lambda_{l}} .1 \tag{A.11}
\end{equation*}
$$

Proof. By using A.9, we have
(A.12) $X\left(z_{1}, t\right) \cdots X\left(z_{l}, t\right) .1=\prod_{1 \leq i<j \leq l}\left(1-\frac{z_{j}}{z_{i}}\right) \prod_{i=1}^{l} \exp \left(\sum_{n=1}^{\infty} t_{n} z_{i}^{n}\right)$.

By taking the coefficient of $z^{\lambda}=z_{1}^{\lambda_{1}} \cdots z_{l}^{\lambda_{l}}$ of this expression, we obtain the formula A.1.

The KP hierarchy is a system of nonlinear partial differential equations for an unknown function $\tau(t)=\tau\left(t_{1}, t_{2}, \ldots\right)$ including the Hirota bilinear equation

$$
\begin{equation*}
\left(D_{t_{1}}^{4}-4 D_{t_{1}} D_{t_{3}}+3 D_{t_{2}}^{2}\right) \tau(t) \cdot \tau(t)=0 \tag{A.13}
\end{equation*}
$$

The whole system of the KP hierarchy is represented by the following bilinear relation:

$$
\begin{equation*}
\oint \frac{d z}{2 \pi i} X^{*}\left(z, t^{\prime}\right) \tau\left(t^{\prime}\right) X(z, t) \tau(t)=0 \tag{A.14}
\end{equation*}
$$

Proposition A.3. For any partition $\lambda$, the Schur function $S_{\lambda}(t)$ is a solution of the KP hierarchy.

Proof. Note first that the bilinear equation A. 14 can be rewritten as follows:

$$
\begin{equation*}
\left(\sum_{m+n=-1} X_{m}^{*} \otimes X_{n}\right) \tau \otimes \tau=0 \tag{A.15}
\end{equation*}
$$

Here $\tau \otimes \tau=\tau\left(t^{\prime}\right) \tau(t)$ is regarded as an element of $\mathbb{C}\left[\left[t^{\prime}\right]\right] \otimes \mathbb{C}[[t]]$. By the anti-commutation relation A.1, one has

$$
\begin{align*}
& \left(\sum_{m+n=-1} X_{m}^{*} \otimes X_{n}\right) X_{k} \otimes X_{k}  \tag{A.16}\\
& \quad=X_{k+1} \otimes X_{k-1}\left(\sum_{m+n=-1} X_{m}^{*} \otimes X_{n}\right)-1 \otimes X_{k-1} X_{k}
\end{align*}
$$

and the last term $X_{k-1} X_{k}$ vanishes. Hence, by applying the operator $X_{k+1} \otimes X_{k-1}$ to A.15, it follows that $X_{k} \tau(t)$ is also a solution of the KP hierarchy. Starting from the solution $\tau(t)=1$, we see that all the Schur functions are solutions of KP hierarchy by Proposition A. 2

Proposition A.4. Let $\tau_{0}(t)=\tau(t)$ be any solution of the KP hierarchy, and put

$$
\begin{equation*}
\tau_{1}(t)=X(w, t) \tau(t) \tag{A.17}
\end{equation*}
$$

Then, we have

$$
\begin{equation*}
\oint \frac{d z}{2 \pi i} z X^{*}\left(z, t^{\prime}\right) \tau_{0}\left(t^{\prime}\right) X(z, t) \tau_{1}(t)=0 \tag{A.18}
\end{equation*}
$$

Proof. Apply $X(w, t)$ to the second factor $X(z, t) \tau(t)$ of the bilinear relation A.14. Then one obtains A. 18 by using the relation

$$
\begin{equation*}
X(w, t) X(z, t) \tau(t)=-\frac{z}{w} X(z, t) X(w, t) \tau(t)=-\frac{z}{w} X(z, t) \tau_{1}(t) \tag{A.19}
\end{equation*}
$$

as desired.
The formula A. 18 is the bilinear relation of the first modified KP hierarchy.

By the change of variables $t \rightarrow t-s$ and $t^{\prime} \rightarrow t+s$, the relations A. 14 and A. 18 can be rewritten into the following systems of Hirota bilinear
equations

$$
\begin{align*}
& \sum_{n=0}^{\infty} p_{n}(-2 s) p_{n+1}\left(\widetilde{D}_{t}\right) \exp \left(\sum_{m=1}^{\infty} s_{m} D_{t_{m}}\right) \tau(t) \cdot \tau(t)=0  \tag{A.20}\\
& \sum_{n=0}^{\infty} p_{n}(-2 s) p_{n+2}\left(\widetilde{D}_{t}\right) \exp \left(\sum_{m=1}^{\infty} s_{m} D_{t_{m}}\right) \tau_{0}(t) \cdot \tau_{1}(t)=0 \tag{A.21}
\end{align*}
$$

where $\widetilde{D}_{t_{n}}=D_{t_{n}} / n$. These are the Hirota bilinear equations for the $\tau$ functions of the KP hierarchy and the first modified KP hierarchy, respectively.

## References

[1] S. Fukutani, K. Okamoto and H. Umemura, Special polynomials associated with the rational solutions and the Hirota bilinear relations of the 2nd and the 4th Painlevé equations, preprint.
[2] M. Jimbo and T. Miwa, Solitons and infinite dimensional Lie algebras, Publ. RIMS, Kyoto Univ., 19 (1983), 943-1001.
[3] M. Kac and P. van Moerbeke, On an explicitly soluble system of nonlinear differential equations related to certain Toda lattices, Advances Math., 16 (1975), 160-169.
[4] Y. Murata, Rational solutions of the second and the fourth equations of Painlevé, Funkcialaj Ekvacioj, 28 (1985), 1-32.
[5] T. Nishitani and M. Tajiri, On similarity solutions of Boussinesq equation, Phys. Lett., 89A (1982), 379-380.
[6] M. Noumi and K. Okamoto, Irreducibility of the second and the fourth Painlevé equations, Funkcialaj Ekvacioj, 40 (1997), 139-163.
[7] K. Okamoto, Studies on Painlevé equations III, Second and fourth Painlevé equations $P_{I I}$ and $P_{I V}$, Math. Ann., 275 (1986), 221-255.
[8] K. Okamoto, Algebraic relations among six adjacent $\tau$-functions related to the fourth Painlevé system, Kyushu J. Math., 50 (1996), 513-532.
[9] H. Umemura, On the irreducibility of the first differential equations of Painlevé.
[10] G.R.W. Quispel, F. W. Nijhoff and H. W. Capel, Linearization of the Boussinesq equation and the modified Boussinesq equation, Phys. Lett., 91A (1982), 143-145.

Masatoshi Noumi<br>Department of Mathematics<br>Kobe University Rokko, Kobe 657-8501<br>Japan<br>noumi@math.kobe-u.ac.jp

Yasuhiko Yamada<br>Department of Mathematics<br>Kobe University Rokko, Kobe 657-8501<br>Japan<br>yamaday@math.kobe-u.ac.jp


[^0]:    Received November 4, 1996.
    ${ }^{1}$ After completing this paper, the authors were informed by K. Kajiwara and Y. Ohta that they obtained independently the expression of Okamato polynomials in terms of Schur functions.

