

Systematic analysis and optimization of early warning signals for critical transitions

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Abstract

Abrupt shifts between alternative regimes occur in complex systems, from cell regulation to brain functions to ecosystems. Several model-free Early Warning Signals (EWS) have been proposed to detect impending transitions, but failure or poor performance in some systems have called for better investigation of their generic applicability. In particular, there are still ongoing debates whether such signals can be successfully extracted from data. In this work, we systematically investigate properties and performance of dynamical EWS in different deteriorating conditions, and we propose an optimised combination to trigger warnings as early as possible, eventually verified on experimental data. Our results explain discrepancies observed in the literature between warning signs extracted from simulated models and from real data, provide guidance for EWS selection based on desired systems and suggest an optimised composite indicator to alert for impending critical transitions.

Highlights

- How to extract early warning signals (EWS) against critical transitions from data is still poorly understood
- A mathematical framework assesses and explains the performance of EWS in noisy deteriorating conditions
- Composite indicators are optimised to alert for impending shifts
- The results are applicable to wide classes of systems, as shown with models and on empirical data.

Keywords: Critical transition, Early warning signals, Bistable systems, Optimisation, Multivariate analysis, Bifurcation, Dynamics

1. Introduction

The dynamics of many complex systems is characterised by critical thresholds (tipping points) and abrupt shifts between alternative regimes (Scheffer et al., 2009; Ashwin and Zaikin, 2015). Various examples have been observed in diverse research fields and include collapses of ecosystems (Hirota et al., 2011; Wang et al., 2012a), sudden climate shifts (Lenton et al., 2012; Drixfhout et al., 2015) or financial crashes (Dmitriev et al., 2017; Diks et al., 2019). Abrupt regime shifts have particularly been theorised and observed in systems biology and medicine (Korolev et al., 2014; Trefois et al., 2015; Aihara et al., 2022), at the onset of certain disease states like atrial fibrillation (Quail et al., 2015) or epileptic seizures (Meisel and Kuehn, 2012), as well as in biological processes like regulation of gene networks (Angeli et al., 2004; Sharma et al., 2016) and cell fate decisions (Ghaffarizadeh et al., 2014; Mojtahedi et al., 2016), including epithelial-mesenchymal transitions (Lang et al., 2021). Correctly detecting and alerting for these critical changes allows to better understand complex developments and to anticipate dangerous outcomes. However, many such complex systems have not been fully characterised with mechanistic models, thus requiring simpler and more generic approaches to support data-driven estimates.

The critical transitions (CT) framework have been proposed to address tipping points using low-dimensional systems descriptions (Kuehn, 2011) and associated early warning signals (EWS), computed from statistical indicators extracted from data like increasing variance, autocorrelation or coefficient of variation (Drake and Griffen, 2010; Lade and Gross, 2012). These signs and derived indexes (Chen et al., 2012; Navid Moghadam et al., 2020; Matsu-mori et al., 2019), in principle generic for broad classes of systems, have been tested and applied on biological, epidemiological and medical data with alternate success (Carpenter et al., 2011; Dai et al., 2012; Wilkat et al., 2019; Proverbio et al., 2022a). Therefore, recent studies have recommended caution when attempting predictions based on EWS (Boettiger and Hastings, 2012; Clements and Ozgul, 2018; Dudney and Suding, 2020). Since there is an increasing interest for EWS in systems biology and biomedicine, it is thus compelling to provide a unified framework for the analysis and interpretation of such indicators, to determine in which cases they can be safely applied and to understand their limitations. In addition, going beyond univariate indicators will improve their performance in detecting and alerting for impending critical transitions.

In this work, we provide a systematic analysis of the CT framework and its associated EWS, to define their range of applicability and understand why discrepancies have been observed between theoretical predictions and experimental data (Kuehn et al., 2022; Cohen et al., 2022). Systems biology is characterised by two main paradigms (Mazzocchi, 2012): one investigating the single details of molecular combinations or regulatory networks, alike

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38 to “microstates” in statistical mechanics (Stumpf et al., 2017), and another looking for gen-
39 eral analytical models, built upon kinetic theories, to understand complicated biochemical
40 processes in simpler and general terms (Ferrell Jr et al., 2009). The latter allows to construct
41 classes of systems according to universal routes of dynamical development, regardless of the
42 microscopic details. We leverage this paradigm to make sense of critical transitions and
43 identify the most relevant classes pertaining to biological systems (Box 1). We also provide
44 guidance for EWS selection and optimisation, depending on realistic noise properties and
45 other notable features of classes of complex systems, developing new composite indicators.

46 Our work bridges mathematical insights and observations of real systems to classify
47 various tipping mechanisms. There are ongoing debates whether regime shifts in biological
48 systems, like cell-fate decision, are primarily driven by deterministic bifurcations (Andreucut
49 et al., 2011; Stanoev et al., 2021) or by random fluctuations (Wang et al., 2011; Stumpf
50 et al., 2017), which prompted several authors to question the old “Waddington landscape”
51 interpretation (Moris et al., 2016). By systematically analysing known regime shifts, we
52 classify the mathematical models to address various types of critical transitions, subject
53 to combinations of bifurcations and noise (Berglund and Gentz, 2006), and to develop a
54 method to extract systems’ robustness proxies from data (Box 2).

55 We first employ a framework based on dynamical manifolds, underpinning universal
56 routes to explosive transitions (Kuehn and Bick, 2021), to characterise the warning signals
57 associated to “noisy” bifurcations, and to study their dependency on noise properties and
58 other dynamical features like rapid approaches to threshold values. This way, we provide
59 general results about EWS robustness and sensitivity to dynamical features, to guide appli-
60 cations on various systems, understand their limitations and promote future developments.

61 Then, we focus on a critical transition sub-class of high biological relevance, the stochastic
62 saddle-node bifurcation (Ferrell Jr et al., 2009). For this tractable, yet realistic model of
63 complex biological processes, we develop a composite EWS indicator to optimise the leading
64 time of the alerts, i.e., how much in advance reliable signals are triggered, with respect to
65 an impending transition. The new indicator is optimised over realistic noise types using
66 the common genetic toggle switch model (Sharma et al., 2016), as representative of the
67 considered CT class. This way, we overcome the limitations of other EWS from literature,
68 which have mostly been developed over Gaussian noise while biological systems usually
69 feature correlated and state-dependent noise (Hasty et al., 2000; Dunlop et al., 2008; Zhang
70 et al., 2012). Thanks to this extension, the indicator also provides additional insights about
71 the systems under investigation, such as inference of noise type from data. The theoretical
72 results are finally tested and verified on publicly available experimental data, demonstrating
73 their potential for monitoring and interpreting diverse systems.

Box 1: Classification of critical transitions

Consider a dynamical system whose state (or regime) is usefully characterised by a set of dynamic variables $x \in \mathbb{R}^n$, whose relations to each other are modeled by a set of parameters $p \in \mathbb{R}^m$:

$$\frac{dx}{dt} = F(x(t), p), \quad (1)$$

where $F : \mathbb{R}^{n+m} \rightarrow \mathbb{R}^n$ is a system of sufficiently smooth functions. If p is not explicitly dependent on time, the system is termed *autonomous*; if $p = p(t)$, the system is called *non-autonomous*. The distinction between autonomous and non-autonomous can be supported when considering naturally fixed parameters (Maini et al., 1991), or when addressing timescale separation (“slow-fast system”) between biochemical processes, like mRNA transcription versus protein degradation times (Yasemi and Jolicœur, 2021). This results in sets of dynamical (for variables) and algebraic (for parameters, termed at quasi-steady state) equations (Del Vecchio et al., 2016). Together, variables and parameters define and shape a state space (or “landscape”) that, if $F(x, p)$ has elements of non-linearity, can be characterised by multiple attractors (MacArthur et al., 2009), i.e., region of stability for systems’s states. If parameters are allowed to change (either non-autonomously, or at quasi-steady state), the state space is dynamic and attractors can change, as opposed to static landscapes like Waddington’s.

The state space can be multidimensional. However, near bifurcation points, it can be aptly described using low-dimensional models associated to critical thresholds in the values of leading parameters (usually corresponding to the largest eigenvalues (Kuznetsov, 2013)). Such models are termed “normal forms” of a dynamical system, simplified minimal-order forms that determine the system’s behaviour and retain universal properties of generic bifurcations (see Kuehn and Bick (2021) and STAR Method C.1). Normal forms can be inferred from bistability properties (Angeli et al., 2004) or deduced from network models, if they are available for the considered systems (Gao et al., 2016; Tu et al., 2021).

In addition to bifurcation points, noise can characterise the system’s dynamics. Noise is ubiquitous in biology (Tsimring, 2014; Su et al., 2019) and can correspond to stochasticity in intrinsic biochemical processes or cell-cell variation (Zhang et al., 2012). Mathematically, noise variables can be modelled as fast degrees of freedom augmenting system (1), which is a dualistic representation to stochastic processes (Berglund and Gentz, 2006). Noise can push the system out of original attractors onto new ones, therefore causing random switches between phenotypic states even in the absence of dynamical bifurcations.

We propose to use the relative timescales between dynamical variables, parameters and noise to develop a systematic classification of transitions between system states. This way, we synthesise and improve the contributions of Thompson and Sieber (2011); Kuehn (2011); Ashwin et al. (2012); Shi et al. (2016) towards the establishment of a theory on critical transitions in real systems. To do so, extend and disentangle Eq. 1 to explicit the dependencies

on state variables $x \in \mathbb{R}^m$ and system parameters $p \in \mathbb{R}^n$, on the introduced stochastic variables $\xi \in \mathbb{R}^l$ and on the relative timescales modelled by time parameters τ_i , $i = \{x, p, \xi\}$. This results in a multiscale slow-fast system

$$\begin{cases} \tau_x \frac{dx}{dt} = f(x, p, \xi) \\ \tau_p \frac{dp}{dt} = g(x, p, \xi) \\ \tau_\xi \frac{d\xi}{dt} = h(x, p, \xi) . \end{cases} \quad (2)$$

Using this representation, tipping systems can be classified into three main classes of critical transitions on the basis of relative timescales: bifurcation-induced (“b-tipping”), noise-induced (“n-tipping”) and rate-induced (“r-tipping”), following the nomenclature introduced by Ashwin et al. (2012):

$$\begin{aligned} \text{b-tipping: } & \tau_p \gg \tau_x \gg \tau_\xi \\ \text{n-tipping: } & \tau_p \gg \tau_x \simeq \tau_\xi \\ \text{r-tipping: } & \tau_p \simeq \tau_x \gg \tau_\xi \end{aligned} \quad (3)$$

If $\tau_\xi > \tau_x$, the system becomes ergodic and visits the full state-space uniformly without displaying transitions (Shi et al., 2016).

The b-tipping class thus encompasses all those transitions primarily driven by bifurcations, i.e., slow changes in control mechanisms modelled as quasi-steady approaches of leading parameters to their threshold values. They modify the attractor landscape, in the presence of low noise-to-signal ratios, and can be further sub-classified according to dimension m and co-dimension n (Thompson and Sieber, 2011). In this work, we only consider low-dimensional ones, commonly found in cell dynamics studies. Examples include toggle-switch mechanisms for the lac-operon (Ozbudak et al., 2004), population collapses of microbiological colonies past threshold concentrations of stressors or nutrients (Dai et al., 2015), or epithelial-mesenchymal determination (Sarkar et al., 2019). Higher m and n yield more complex bifurcations associated to, e.g., neural network activity (Izhikevich, 2007).

The n-tipping class groups various transitions driven by stochastic fluctuations on fixed landscapes, including large, impactful and unexpected events (sometimes called “dragon kings” (Sornette, 2006)). Example range from enzymes crossing activation chemical barriers via “promoting vibrations” (Antoniou and Schwartz, 2011), “rebellious cells” undergoing contrasting development pathways during cell reprogramming (Mojtahedi et al., 2016), and other long-studied cases of noise-induced transitions (Horsthemke and Lefever, 1984).

B-tipping and n-tipping directly link to the the aforementioned debates in systems biology about deterministic or stochastic drivers of critical changes. R-tipping refers to critical ramping of control parameters, not coped by the system, which has been so far observed in climate (Wieczorek et al., 2011) and engineering (Bonciolini et al., 2018) systems. The heat-shock response of plants to ramping temperature conditions (Moejes et al., 2017) may fall within this class, but further studies are required. The critical transition classes can be visualised on bifurcation diagrams or using quasi-potential landscapes (Zhou et al., 2012), which can be obtained as integrals of vector fields like Eq. 1 or inferred from data.

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Fig. 1 shows the classification between the transition classes, with illustrative examples of what can happen to systems within simplified attractors. Note that the hard-cut classification derives from the mathematical assumptions in Eq. 3: gradients between the transition classes may exist and call for deep investigation. In particular, our work focuses on “noisy bifurcations”, i.e., dynamics characterised by bifurcation points and the presence of low to moderate noise-to-signal ratio.

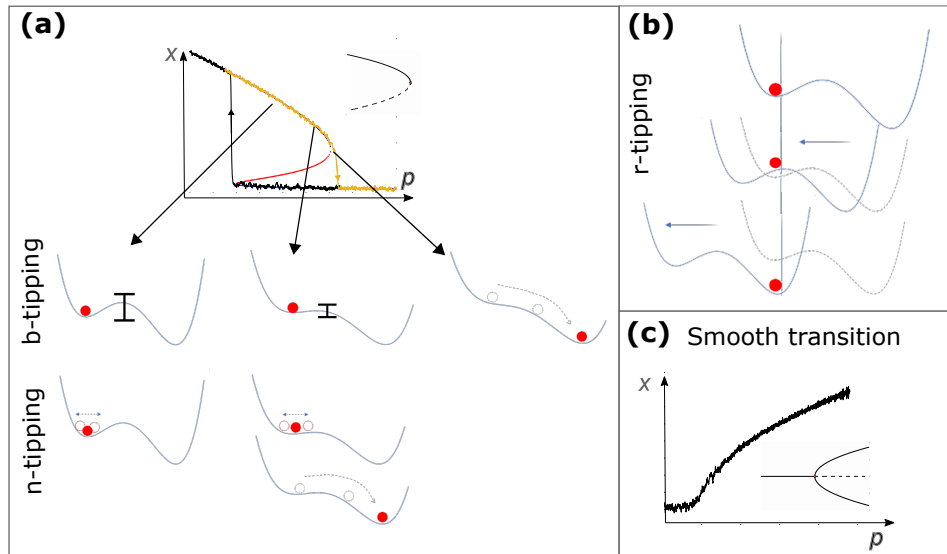


Figure 1: **Classification of transitions between states x of a dynamical system, controlled by a slow changing parameter p .** x and p may also correspond to network combinations of variables and parameters (Moris et al., 2016; Gao et al., 2016). (a): Illustration of b-tipping and n-tipping using a bistable system with saddle-node bifurcations (unstable branch in red; saddle-node template shown in inset). Hysteresis can occur, i.e., asymmetric routes to tipping from one stable state or from the other (orange, from up to down with increasing p ; black, from down to up with decreasing p). B-tipping: the system approaches the bifurcation point. The associated landscape is molded by p and the basin of attraction becomes shallower (as visualised by the bars) until disappearing; there, the system tips. N-tipping: if subject to strong fluctuations, depicted as wiggling of the red ball, the system can be pushed over the barrier onto an alternative attractor, even before the bifurcation point. (b) Illustration of r-tipping: rapid ramping of the control parameter makes it as if the landscape shifts and the systems does not manage to move along, therefore tipping onto another attractor “sliding” underneath. See Ashwin et al. (2012) for formal definitions. (c): Example of “smooth” transition without hysteresis, using a dynamical system close to a pitchfork bifurcation (inset) as template. To reproduce the plots, see STAR Method C.3.

Box 2: Bifurcations with noise and system robustness Among the critical transition classes described above, let us consider those primarily driven by bifurcations, with noise further influencing the dynamics. In this sense, we can speak of “noisy b-tipping”, with the first condition in Eq. 3 becoming

$$\tau_p \gg \tau_x > \tau_\xi , \quad (4)$$

that is, the noise-to-signal ratio is not negligible but the slow-fast condition between variables and parameters still applies.

For this class, normal forms can be used to analytically study systems’ robustness and derive early warning signals for impending tipping points (Kuehn, 2011). Normal forms are general and low-dimensional models $\dot{x} = f(x, p)$ that describe topologically equivalent systems within a bifurcation class, in the vicinity of critical points (Kuznetsov, 2013). They allow to extract analytical and generic results for wide classes of systems (Kuehn and Bick, 2021), at the price of neglecting homeostatic dynamics far from tipping points. As a result, they allow to focus on critical transition mechanisms across various systems, instead of studying the full evolution of a single system. Details about topological equivalence and construction of normal forms are in STAR Method C.1. Fig. 2 shows an example of reduction to normal forms for two simple models.

Here, we consider those normal forms of primary biological interest. The saddle-node bifurcation, often associated with population collapses (Scheffer et al., 2009; Dai et al., 2012) or biological state transitions (Alon, 2006), is defined by $f(x, p) = \pm p \pm x^2$. At $p = 0$, a stable ($\tilde{x}_s = \sqrt{p}$) and unstable ($\tilde{x}_u = -\sqrt{p}$) branch collide and vanish, resulting in a critical transition to an alternative branch (if it exists). Transcritical bifurcations $f(x, p) = px - x^2$ are characteristic, for instance, of epidemic outbreaks (Proverbio et al., 2022a). Here, the two equilibria $x_1 = 0$ and $x_2 = p$ meet at $p = 0$ and exchange stability. Finally, the family of pitchfork bifurcations $f(x, p) = px + lx^3$ describe branching processes from one to two states (or viceversa); $l > 0$ identifies subcritical bifurcations, associated to critical transitions, while $l < 0$ defines the supercritical case, with a continuous transition over mean values. This mechanism is identified in cell regulation processes (Moris et al., 2016).

Stochastically forced systems, associated to “noisy b-tipping”, can be written in the Itô form (Thompson and Sieber, 2011)

$$dx = f(x, p)dt + h(x, p)dW , \quad (5)$$

where dW is a Wiener process with variance σ and $f(x, p)$ is a suitable normal form from those described above. The term $h(x, p)$ allows to represent different noise types, to reflect modern knowledge of stochastic processes occurring in biological systems. Additive Gaussian noise with $h(x, p) = 1$ is usually associated to extrinsic cell-cell variability. State-dependent (multiplicative) noise $h(x, p) \neq \text{const}$ represents intrinsic stochasticity determined by, e.g., reaction rates, timescales or species concentrations of the underlying biochemical processes (O’Regan and Burton, 2018). Combinations of additive and multiplicative noise, with various ratios depending on different systems, are more realistic (Liu et al., 2009; Sidney et al., 2010) and fit experimental data better than Gaussian noise (Wang et al., 2012b).

If the microscopic kinetics is known, the noise terms can be exactly derived from the Master equation using Gillespie formalism (Gillespie, 2000). Alternatively, a diffusion approximation (Allen, 2010; Van Kampen, 1992) derives noise terms proportional to system state ($h(x, p) = x$), or to the drift term of Eq. 5, $h(x, p) \propto f(x, p)$. Here, for multiplicative noise, we consider $h(x, p) = \sqrt{f(x, p)}$ (O'Regan and Burton, 2018) and $h(x, p) = f(x, p)$, to reflect modelling of biological regulatory circuits (Hasty et al., 2000). This way, mechanistic and stochastic normal-form bifurcation models are examined to study the effects of intrinsic and extrinsic noise on statistical patterns of variability and related EWS.

Following the procedure detailed in STAR Method C.2, Eq. 5 is analysed by solving the slow dynamics, linearising around a trajectory inside the stable (attracting) manifold and changing the coordinates to highlight the residuals $y(t)$ around the linearization. This procedure gives

$$dy = \partial_x f(\tilde{x}_s(t), t)y dt + \sqrt{h^2(x)}dW \quad (6)$$

where \tilde{x}_s corresponds to the attracting part of the critical manifold (stable solutions). The linearised drift term corresponds to the leading eigenvalue of the deterministic normal form. Its magnitude $|\partial_x f(\tilde{x}_s(t), t)|$ is the asymptotic decay rate of a perturbation. It corresponds to the concept of engineering resilience (Holling, 1996), which is akin to that of robustness (Kitano, 2004). A change of notation $|\partial_x f(\tilde{x}_s(t), t)| = k$ makes explicit that Eq. 6 corresponds to a (possibly non-autonomous) Ornstein-Uhlenbeck process, with critical k given by $k_0 = 0$. It is a well-studied problem in stochastic processes theory, with analytical solutions for its statistics in different regimes (Allen, 2010; Gardiner, 1985). Eq. 6 can be regarded as a first order autoregressive model. However, its derivation from normal forms allows more nuanced interpretation: rather than being hypothesised as a statistical model to capture simple relationships, it is general for all models that can be reduced to normal forms.

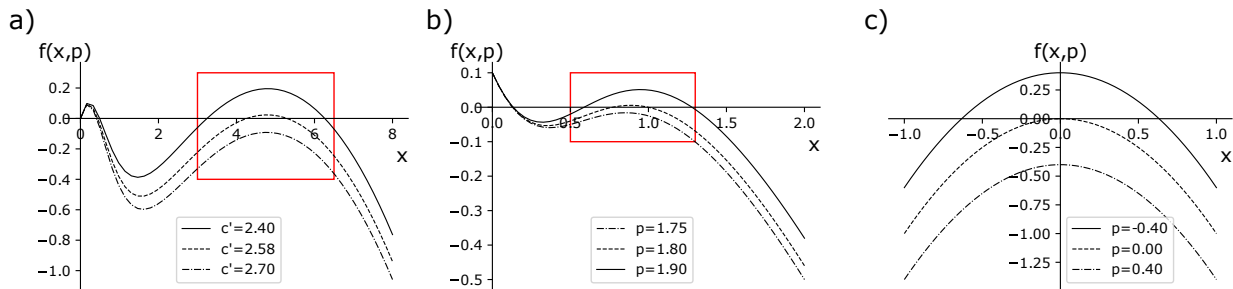


Figure 2: **Visual example of topological equivalence.** a) Plot for $dX/dt = f(x, p = c') = X(1 - X/K) - c'X^2/(X^2 + 1)$, a model of harvested ecological populations (Scheffer et al., 2009), also akin to Allee effects observed in microbiological colonies (Dai et al., 2012); X is the population density, K is the carrying capacity and c' is the maximum harvest rate. b) Plot for $f(x, c)$ of the autocatalytic loop model Eq. 15. c) Plot for $f(x, p)$ of the saddle node normal form $\dot{x} = -p - x^2$. The two realistic models are locally topologically equivalent to the normal form within the red rectangle (visual reference): they approach a bifurcation point, marked by $f(x, p)$ crossing the x -axis, as the parameter c' or c changes.

79 2. Results

80 2.1. Robustness of EWS for noisy bifurcations

81 Within the class of critical transitions induced by bifurcations characterised by small
 82 fluctuations, discussed in Box 1 and 2, we study the early warning signals associated to
 83 impending tipping points, considering different noise types that are better representative of
 84 biological dynamics than pure Gaussian noise (see Box 2).

85 Analytic expressions for key summary statistics indicators can be obtained from Eq. 6
 86 using standard approaches for stochastic processes (Allen, 2010; Gardiner, 1985). Their
 87 behaviour as the control parameter changes provides early warning signals for approaching
 88 noisy bifurcations (Scheffer et al., 2009). The lag- τ autocorrelation function does not depend
 89 on $h^2(x, p)$ but only on $|\partial_x f(\tilde{x}_s, p)| = k$:

$$AC(\tau) = e^{-k\tau}. \quad (7)$$

Hence, the common indicator lag-1 autocorrelation ($AC(1)$, with $\tau = 1$) only depends on the dampening rate. The power spectrum of the Fourier transforms and the variance, two common indicators, explicitly depend on $h^2(x, p)$:

$$S(\omega) = \frac{h^2(\tilde{x}, p)}{k^2 + \omega^2} \quad (8)$$

$$\text{Var} = \frac{h^2(\tilde{x}, p)}{2k}. \quad (9)$$

90 Coefficient of variation (CV) and Index of dispersion (ID), defined as

$$CV = \frac{\sqrt{\text{Var}}}{\tilde{x}_s}, \quad ID = \frac{\text{Var}}{\tilde{x}_s}, \quad (10)$$

91 also depend on $h^2(\tilde{x}, p)$. Other statistical moments, for stochastic processes with quasi-
 92 steady state parameter, can be expressed as

$$\langle y^\nu \rangle - \langle y \rangle^\nu = \int_{-\infty}^{\infty} (y' - \mu)^\nu P(y') dy' \quad (11)$$

93 where $P(y')$ is the probability density function from the associated Fokker-Plank equation
 94 (Gardiner, 1985) and μ is the expected average value. Skewness and kurtosis, sometimes
 95 suggested as indicators for EWS (Guttal and Jayaprakash, 2008), can be easily extracted
 96 from Eq. 11 as third and fourth moments ($\nu = 3$ and 4). Entropy-based indicators are more
 97 challenging to derive in case of multiplicative noise, as their defining integrals may not be
 98 solvable. Their derivation in case of Gaussian noise is described in STAR Method C.4; for
 99 the other cases, their behaviour is estimated below using computer simulations.

100 In all cases, the analytical results for each normal form can be obtained by substituting
 101 the corresponding dependency of the drift term to the control parameter: for the saddle-
 102 node, $k = 2\sqrt{p}$, for the transcritical $k = p$ and for the pitchforks $k = 2p$. In Fig. 3, the

103 effect of multiplicative noise on the trends of common indicators is shown using $h(\tilde{x}, p) = x$
 104 and $h(\tilde{x}, p) = x^2$.

105 Fig. 3 shows expected trends of common statistical indicators, for the three main normal
 106 forms and different noise types. Although the scaling induced by $k(p)$ differs, the qualitative
 107 trends are conserved across the bifurcations. This observation suggests genericity of EWS,
 108 but also difficulties to infer the existence of one or another bifurcation using statistical
 109 indicators alone. Other methods (e.g. Angeli et al. (2004)) are recommended to complement
 110 the inference.

111 For Gaussian noise, EWS are associated with increasing trends of statistical indicators
 112 (Dakos et al., 2015; Scheffer et al., 2009). However, multiplicative noise may alter or com-
 113 pletely disrupt them (as also noted by O'Regan and Burton (2018)), resulting in no early
 114 warnings prior to tipping points. Eq. 8 shows that even power spectrum trends can be
 115 subject to alterations from expected patterns, potentially resulting in spurious signals.

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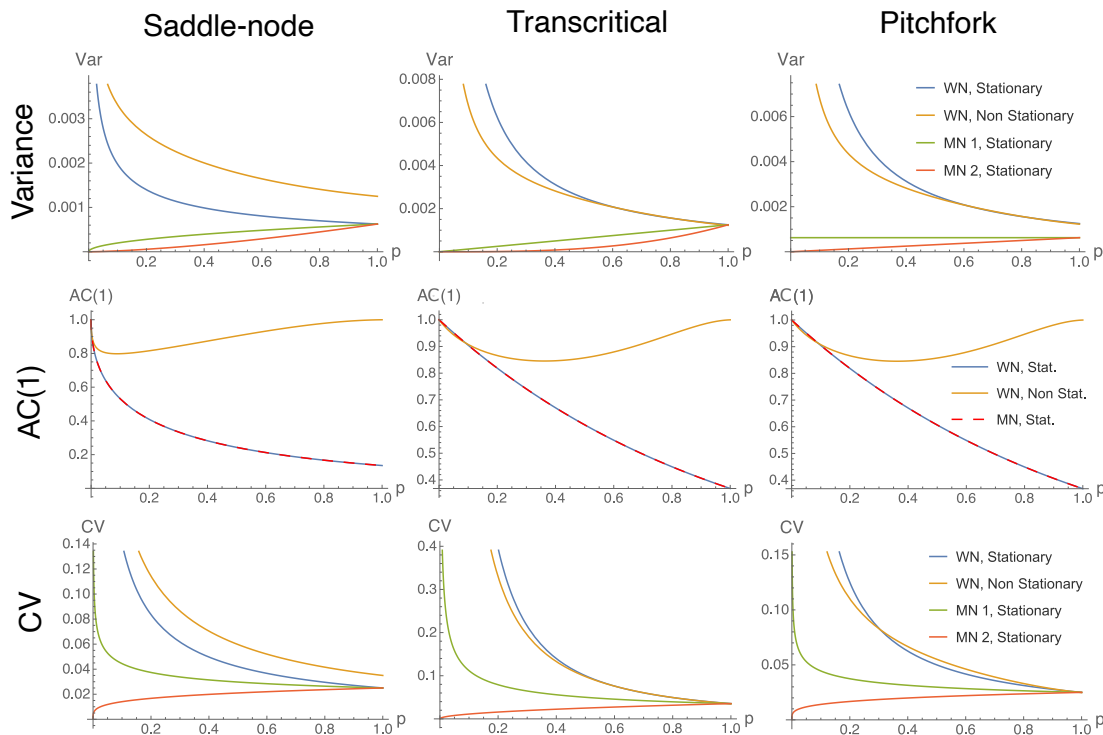


Figure 3: **Trends of common statistical indicators.** We consider Var, AC(1) and CV for saddle-node, transcritical and pitchfork bifurcations as $p \rightarrow 0$, in different dynamical contexts (combinations of noise characteristics and stationarity for the control parameter). WN: white (Gaussian) noise; MN 1: multiplicative (state-dependent) noise $h(\tilde{x}, p) = x$; MN 2: multiplicative noise $h(\tilde{x}, p) = x^2$. As the autocorrelation is independent on noise, only MN 1 is show and it overlaps with the white noise case.

A preliminary investigation on ramping parameters (Pavithran and Sujith, 2021) can

also be conducted. In this case, $\tau_x \simeq \tau_p$: the quasi-steady-state (stationary) assumption is relaxed, but r-tipping may not yet occur. Let us consider linear ramping as $k = k_0 - at$, where k_0 is any initial condition, a is a small rate coefficient and the ramping stops at the critical value $k = 0$. Both coefficients are set to 1 to represent commensurable time scales. Only Gaussian noise is considered. This is a particular case of inhomogeneous processes (Gardiner, 1985) for which statistical moment solutions exist in the form

$$\langle y(t) \rangle = e^{-\int_0^t k(t') dt'} \quad (12)$$

$$\langle y(t)y(t') \rangle = \frac{\sigma^2}{2k} e^{-2\int_0^t k(t'') dt''} + \sigma^2 \int_0^t e^{-2\int_{t'}^t k(s) ds} dt' . \quad (13)$$

117 Derived statistics are calculated analogously. Eq. 13 is solved using Mathematica software
 118 to tackle the rightmost integral yielding the non-elementary Error function Erf(t). Fig. 3
 119 shows that trends of common indicators may be modified by commensurable time scales of
 120 parameters evolution. Hence, raising reliable alerts becomes more challenging.

121
 122 Overall, this analysis demonstrates that theoretical early warning signals due to increas-
 123 ing trends of summary statistics are sensitive to the “dynamical context”, i.e. noise proper-
 124 ties and reciprocal time-scales. Hence, if the dynamical context is not carefully accounted
 125 for, spurious signals may be extracted from data, as observed in early findings from single
 126 systems (Brett et al., 2017; Proverbio et al., 2022a).

127 If the context is known, the current results suggest which indicators to use to obtain
 128 robust early warnings. The autocorrelation is robust against changing noise properties; the
 129 variance is more sensitive to multiplicative noise, but maintains its expected trends in case
 130 of ramping parameters. The coefficient of variation is also robust in case of commensurable
 131 time scales and copes well in case of certain types of multiplicative noise. Overall, what
 132 matters is the competition between changes in noise and changes in resilience: depending on
 133 which one is more rapid, the indicators and their associated EWS may perform as expected
 134 or fail to anticipate an impending critical transition.

135
 136 Measurement processes or details of realistic models may further influence EWS. Mea-
 137 surement uncertainties, assumed as Poisson processes associated with measuring instruments
 138 or procedures and thus independent of systems’ dynamics, can be introduced in the formu-
 139 las of statistical indicators by error propagation in quadrature (see STAR Method C.4 for
 140 details). In case of Gaussian noise and stationary processes, the expected trends of common
 141 indicators are not altered, hence, EWS can be in principle extracted even when using noisy
 142 measurements (*cf.* STAR Method C.4).

143 Single indicators may also be skewed in case realistic details are considered. For instance,
 144 on empirical data, normalising by the critical value and set a normal form around $p_0 = 0$
 145 and $\tilde{x}_s(p) = 0$ may be challenging, since such critical values are largely unknown. Hence,
 146 instead of computing $\tilde{x}_s(p) \rightarrow 0$ like on perfectly reconstructed normal forms, $\tilde{x}_s(p) \rightarrow x'_0$
 147 is often computed (Dai et al., 2015), where x'_0 corresponds to the critical value, unknown a

148 priori. Such case can be modelled as $\tilde{x}_s(p) = x'_0 + \sqrt{p}$. Hence, Eq. 10 becomes

$$CV_r = \frac{\sqrt{\text{Var}}}{x'_0 + \sqrt{p}}. \quad (14)$$

149 Here, other multiplicative noise forms may alter its behaviour and shadow possible early
150 warnings. Finally, skewness and kurtosis calculated from Eq. 11 display increasing trends
151 when $P(y')$ is symmetric (STAR Method C.4). However, this may not be true in case of
152 multiplicative noise (Sharma et al., 2016), resulting in distorted trends and early warnings.
153 In this sense, there is no ambiguity between the results of Guttal and Jayaprakash (2008),
154 proposing EWS from skewness, and Dai et al. (2012), observing flat and fluctuating trends
155 on experimental data: likely, the noise properties were different than what assumed.

156 2.2. Optimisation of EWS

157 Having assessed in which cases the proposed early warning signals are expected to work
158 for noisy b-tipping transitions, we now optimise their performance to provide significant and
159 as-early-as-possible alerts, in a range of dynamical contexts and for the most common tran-
160 sitions observed in systems biology. To this end, we focus on multistable systems (Sarkar
161 et al., 2019), develop and solve an optimisation problem using computer simulations to go
162 beyond the first-order approximation from Eq. 6 (see STAR Method C.2 for details), and
163 study a wide range of noise levels and types, to establish a composite indicator that is robust
164 and performing across multiple systems.

165
166 Multistable systems are systems whose deterministic landscape features at least two
167 attractors (Feng et al., 2016), and usually undergo either saddle-node bifurcations or n-
168 tipping. Bistability means local multistability across two attractors. Angeli et al. (2004)
169 provides necessary and sufficient conditions for bistability in a wide range of biological
170 systems. Among them, a feedback model with three-points I/O characteristic curves suffices.
171 A simple linear system with monotonic sigmoidal feedback can do the job, in a range of
172 parameters (Fig. 4). As a case study, the autocatalytic positive feedback loop derived from
173 Michelis-Menten kinetics (Sharma et al., 2016)

$$\dot{x} = f(x, c) + \eta(t) = K + c \frac{x^k}{1 + x^k} - x + \eta(t). \quad (15)$$

174 satisfies the bistability conditions, and can thus display transitions between attractors, if
175 $0 < K < 1/(3\sqrt{3})$ for $k = 2$ (Weber and Buceta, 2013). In Eq. 15, x is the concentration
176 of a transcriptional factor activator, activating its own transcriptions when bound to a
177 responsive element; K is the basal expression rate, c is the maximum production rate, k
178 is the Hill coefficient and $\eta(t)$ accounts for the stochastic terms. Eq. 15 comes from a
179 two-variable genetic toggle switch, assuming slow-fast timescale separation between the two
180 variables (Strogatz, 2015) and after a-dimensionalising the chemical details to retain the
181 dynamical scaffold. Notably, networks of Michelis-Menten regulators can be reduced to Eq.
182 15 after dimension reduction techniques (Gao et al., 2016).

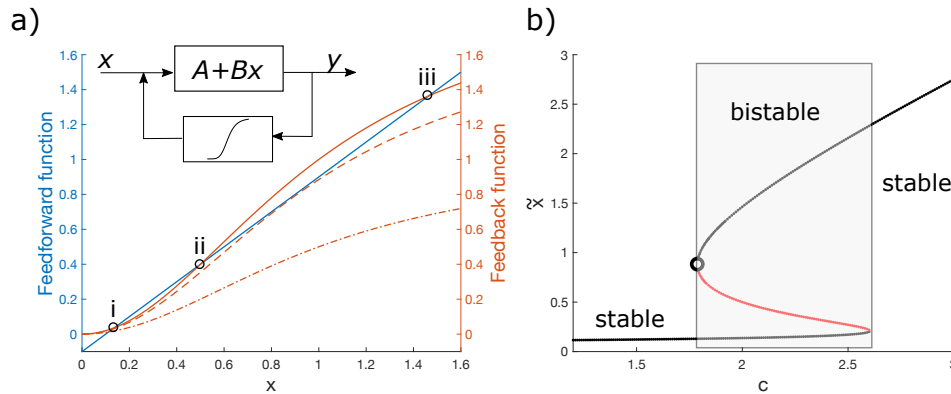


Figure 4: **Bistable systems, studied with (a) characteristic curves or (b) bifurcation diagram \tilde{x} (stable state) vs control parameter.** Among the systems undergoing saddle-node bifurcations, any linear system with nonlinear feedback and adequate feedback gain, such that the characteristic curve crosses the activation function in three points (two stable, one unstable), can display bistability. This example uses Eq. 15. a) The feedback function FB corresponds to the Hill function ($k = 2$), the feedforward FF to the linear part $-(K - x)$ with K in its appropriate range. The control parameter c tunes the FB function. Dashed-dotted line: c is not sufficient to promote bistability, corresponding to left stable region of (b). Dashed line: the critical value for which FB is tangent to FF , corresponding to saddle-node point, open circle in (b). Solid line: bistable system with three intersection points (stable, i and iii; unstable, ii). When studying the vector field $f(x, c)$ is easier than the characteristic curves, one can use the representation and interpretation in Fig. 2b. Note that the line styles have the same meaning in panel (a) and Fig. 2.

183 Eq. 15 displays bistability for a range of values c (the exact range depends on K and k
 184 (Proverbio et al., 2022b)) and, in particular, a saddle-node bifurcation between two alterna-
 185 tive steady states at a critical value c_0 of the parameter c , such that $\partial f / \partial x|_{(\tilde{x}, c_0)} = 0$:

$$c_0 = \frac{(x_0^k + 1)^2}{kx_0^{k-1}} \quad (16)$$

186 where x_0 is the tipping value for the system state. Therefore, system 15 can be used as
 187 a paradigmatic example of biological systems, within the saddle-node b-tipping class, to
 188 perform optimisation studies that go beyond the local and low-noise-to-signal-ratio approx-
 189 imation provided by normal forms.

190
 191 The quasi-steady state assumption is generally accepted for such systems (Del Vecchio
 192 et al., 2016), so we focus on dynamical contexts characterised by different types of noise,
 193 whether yielding n-tipping or possibly skewing statistical indicators due to multiplicative
 194 and/or additive nature. To model combinations of intrinsic and extrinsic noise, we set

$$\eta(t) = [\alpha + (1 - \alpha)h(x)]dW, \quad (17)$$

195 where α weights the white or multiplicative noise component ($\alpha = 1$ corresponds to purely
 196 additive Gaussian noise, $\alpha = 0$ to purely multiplicative); like above, $h(x) = x$ or $h(x) \propto$

197 $f(x) = x^k/(1+x^k)$ (Hasty et al., 2000) and dW is a Wiener process with variance σ . With-
 198 out loss of generality (Proverbio et al., 2022b), we set $k = 2$.

199
 200 As early warning signals are associated with increases of statistical indicators, we need
 201 to establish a measure of statistically significant increase, to rule out false positives and
 202 false negatives due to random fluctuations in the indicators. To do so, we employ the
 203 p-value analysis used in Proverbio et al. (2022b) (see STAR Method C.6 for details). It
 204 allows to measure at which value of the control parameter c , before c_0 , a significant signal
 205 is triggered, thus obtaining a “lead-parameter” $c_{sig}^{\mathcal{I}}(\sigma, \alpha)$ depending on noise properties and
 206 the considered indicator \mathcal{I} (see STAR Method C.6 for details). $c_{sig}^{\mathcal{I}}(\sigma, \alpha)$ is first computed
 207 for each indicator individually. Fig. 5a shows the results in case of white noise, while
 208 various functionals of multiplicative noise $h(x)$ (with $\alpha = 0$) are reported in Supplementary
 209 Figure S3. Each indicator yields various c_{sig} ; in Fig. 5a, Var, AC(1) and H_S maximise c_{sig}
 210 over various noise levels, while other indicators like skewness and kurtosis perform poorly,
 211 as anticipated by the analytical results. CV and ID are also rather poor, likely due to
 212 fluctuations of mean values and anticipating n-tipping (*cf.* also Supplementary Figure S2 and
 213 S3). For the case of multiplicative noise (Supplementary Figure S3), H_S keeps performing
 214 well while Var, as expected from the theoretical analysis, decreases its performance despite
 215 being still better than Skew and Kurt.

216 Complementing the analysis of the lead parameter requires understanding how many
 217 noise-induced tipping events occurred before it and assessing whether the increasing indi-
 218 cators alert for impending collapses or reflect transitions that have already happened. The
 219 analysis thus interprets warning indicators as “anticipating” or “just-on-time detecting” the
 220 tipping events. To do so, a counter \mathcal{C} quantifies, for each parameter value c and for each
 221 noise level σ , how many trajectories tip onto the alternative stable state. The results are
 222 in Fig. 5b: as σ increases, more n-tipping events occur before the bifurcation point. In
 223 particular for $\sigma > 0.42$, several noise-induced transitions occur at $c \simeq c_{sig}$. Hence, as noise
 224 increases, the indicators capture ongoing critical transitions but are not able anymore to pro-
 225 vide much earlier alerts. This likely explains the remarks from Dudney and Suding (2020),
 226 that EWS could not anticipate several transitions in real-world systems, in particular those
 227 characterised by high noise-to-signal ratios.

228
 229 The previous results are also employed to define an optimisation problem to maximise
 230 $c_{sig}^{\mathcal{I}}(\sigma, \alpha)$ for varying α . To do so, we define a composite indicator as linear combination of
 231 indicators

$$\mathcal{S} = \sum_k w_k \mathcal{I}_k \quad (18)$$

232 and look for a set of weights $\mathbf{w} = \{w_k\}$ that maximises all $c_{sig}^{\mathcal{I}}(\sigma, \alpha)$ as σ increases (to
 233 guarantee robustness against noise levels), for the various α :

$$\hat{\mathbf{w}} \text{ s.t. } \max_{\mathbf{w}} \mathcal{S} = \max_{\mathbf{w}} \left[\sum_l c_{sig}^{\mathcal{S}}(\mathbf{w}, \sigma_l) \right], \quad (19)$$

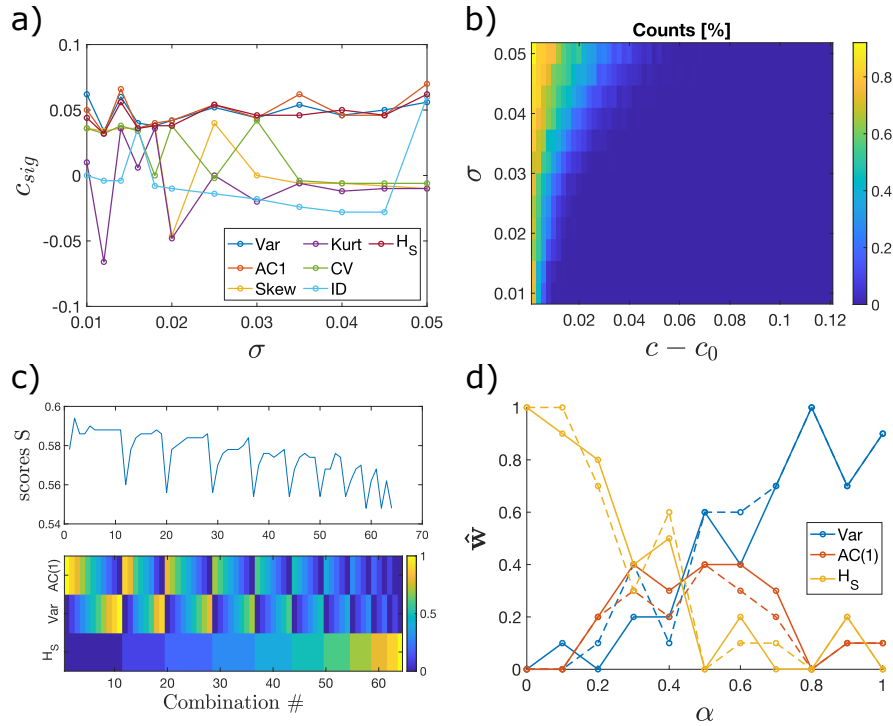


Figure 5: **Optimisation of leading indicators for EWS, according to lead parameter c_{sig} .** a) c_{sig} at various noise intensities σ , for all the most common indicators. b) The counter \mathcal{C} , normalised by all transitions to be interpreted as probability of n-tipping, at different noise intensities σ and distances $c - c_0$ from the bifurcation point. c) Scores S , corresponding to the argument of the cost function Eq. 19, for various combinations \mathcal{S} from Eq. 18. In the panel below, the color code shows the weights w_k for each indicator, in each combination. Results in panels a, b and c refer to $\alpha = 1$. For various types of $h(x)$ and $\alpha = 0$, see Supplementary Figure S3. d) Optimal weights $\hat{\mathbf{w}}$ for each indicator, as a function of noise mixing α . As a representative of multiplicative noise, we used $h(x) = x$. Other $h(x)$ conserve the trends, albeit changing the corresponding c_{sig} . It may happen that the optimisation is solved by multiple combinations (dashed lines).

234 where S are scores composed by sums of $c_{sig}^S(\mathbf{w}, \sigma_l)$ over all σ . In the set \mathcal{I} , we include those
 235 indicators that are expected to be robust and performing, first and foremost in the white
 236 noise case. Leveraging on the previous results, we therefore select Var, AC(1) and H_S . As
 237 the problem is non-convex (Fig. 5c), we perform a grid search for all combinations of w_k ,
 238 with a stride 0.1 and such that $\sum_k w_k = 1$. See Fig. Fig. 5c for the considered combinations
 239 to construct \mathcal{S} .

240 Figure 5d reports the results of the optimisation procedure. Combinations of Var and
 241 AC(1) make up for optimal indicators in case of white noise, $\hat{\mathbf{w}} = [0.9, 0.1, 0]$ for Var, AC(1)
 242 and H_S , respectively, in case of $\alpha = 0$. In this case, H_S is log-proportional to Var (see Eq.
 243 C.11) and does not add much information. In turn, combining the indicators maximises
 244 $c_{sig}^S(\sigma, \alpha)$ in case of mixed noise types. Finally, when multiplicative noise is prevalent in
 245 the system, using Shannon entropy is preferred ($\hat{\mathbf{w}} = [0, 0, 1]$ for $\alpha = 0$). Note that, as

246 the problem is non-convex, there may be more than one combination to create the optimal
247 \mathcal{S} . However, changes in weights w_k are always within $\Delta w_k \sim \pm 10\% w_k$ and the trends are
248 conserved (see dashed lines in Fig. 5d). Such small Δw_k yield changes of $\pm 4\%$ on the scores
249 S , on average over all α ($\Delta S \in [1.8; 6.5]\%$), while off-setting w_k by more than 50% (e.g.,
250 using full variance in case of multiplicative noise) worsens S (and consequently the optimal
251 lead parameter) up to more than 20%.

252 2.3. Verification on experimental data

253 The theoretical predictions are verified and used to interpret experimental data from a
254 previous publication (Dai et al., 2012). The data are sampled from controlled experiments
255 of budding yeast population collapse. Budding yeast cooperatively breaks down the sucrose
256 necessary for its survival, thus inducing a density-dependent dynamics that realises the Allee
257 effect of bistable population dynamics (*cf.* Fig. 2b). Repeated experiments empirically
258 reproduced a saddle-node bifurcation by measuring population density (state variable) as
259 a function of dilution factors (DF, control parameters) affecting the sucrose environment.
260 Various EWS for population collapse can be estimated using distributional data. More
261 details about data collection and analysis are in STAR Method C.7. Testing our theoretical
262 results on a different system than Eq. 15, yet still belonging to the saddle-node driven
263 b-tipping class, would thus assess their generic applicability within this class.

264 Fig. 6 shows trends of each indicator individually, as function of the dilution factor
265 (with critical value at 1600). The error bars are estimated from bootstrapping (STAR
266 Method C.7). Fig. 6 reproduces the results from Dai et al. (2012) and includes the additional
267 indicators considered in this paper. The mean is used to reconstruct the upper stable
268 branch of a saddle-node bifurcation diagram (see Fig. 1), reconstructed from data (the
269 full diagram can be found in the original publication). However, it can not be used as
270 proper EWS as decreasing mean values could signify smooth changes rather than critical
271 transitions, if the transition type and critical parameter are not known. Skewness and
272 kurtosis fluctuate around 0 and 3, respectively, without providing EWS, as one expects in
273 case of symmetric potentials (see Eq. C.17 and C.18). $AC(1)$ and the autocorrelation time
274 (defined as $-1/\log[AC(1)]$ (Dai et al., 2012)) first drop before increasing sharply just before
275 the critical value. Comparing it with Fig. 3, we speculate that there are commensurable
276 time scales between the intake of sugar by yeast cells and their evolution in density. Further
277 experiments are suggested to check for this intriguing hypothesis.

278 Even in this case, as expected, Var, Entropy (H_S), CV and ID display monotonous in-
279 creasing trends close to the bifurcation point. The increases are thus assessed using the
280 p-value test (*cf.* STAR Method C.6) to check whether they are significant or associated
281 with fluctuations. To trigger a significant early warning signal, we require a conservatory
282 significant p-value < 0.01 . This way, we estimate the significant dilution factor DF_{sig} for
283 each indicator. For variance, $DF_{sig} = 1133$, for the others $DF_{sig} = 1000$. Comparing with
284 the optimisation results (from the previous section and Fig. S3), we infer the presence of
285 multiplicative noise in the system's dynamics. Note that entropy showcases the smallest
286 p-value at $DF_{sig} = 1000$; it is also the most robust when changing the repetitions in the

287 bootstrapping procedure (STAR Method C.7).

288

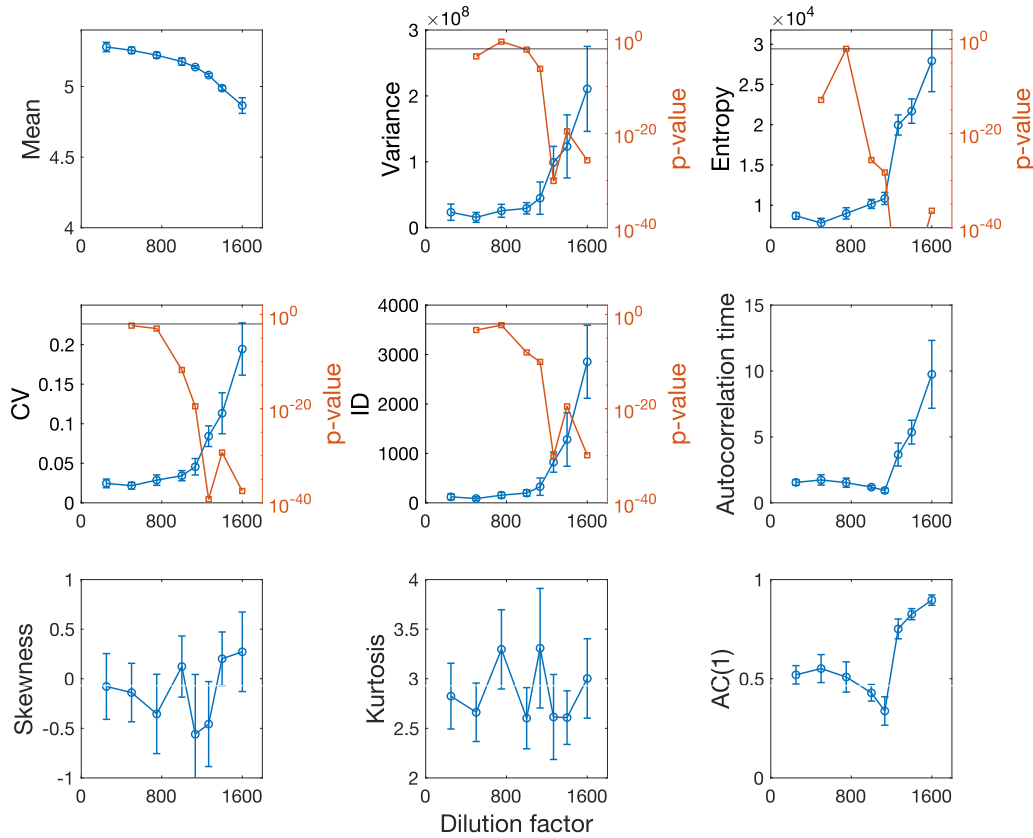


Figure 6: **Statistical indicators calculated on empirical data.** Data are from Dai et al. (2012), as functions of dilution factor (DF). Their corresponding p-values are estimated when the trend is increasing while approaching the bifurcation point (rightwards point). All statistical moments of degree γ have units of measure $(\text{cells}/\mu\text{l})^\gamma$. The autocorrelation time is in days. The mean reproduces the upper stable branch of a saddle-node bifurcation diagram (*cf.* Fig. 1) until the empirically estimated bifurcation point at $\text{DF}=1600$. Horizontal solid lines mark $p\text{-value} = 0.01$.

289 To test the hypothesis of association between EWS performance and noise type, we test
 290 combined indicators with H_S and Var . According to the optimisation above, the higher the
 291 variance content in the mixture, the lower the significance of the increasing trend. This is
 292 what is observed in Fig. 7: having a balance between Var and H_S yields $\text{DF}_{sig} = 1000$, but
 293 with a higher p-value than when reducing the ratio Var/H_S or when comparing with the
 294 case of entropy alone (from Fig. 6).

295 Finally, we test combining CV and H_S , since Fig. 6 suggests that CV could perform
 296 well. Indeed the new combined indicator yields $\text{DF}_{sig} = 750$ (Fig. 7, right), one dilution
 297 step before the others. This is not in contrast with the optimisation analysis: CV is, in fact,
 298 expected to be as performing as H_S if the noise levels are relatively high (see Fig. S3). We

299 recall that CV was not included in the optimisation analysis to be generic and robust across
 300 noise types and levels. However, if high σ in state-dependent noise is known, constructing a
 301 composite indicator using both CV and H_S may improve the alerting performance.

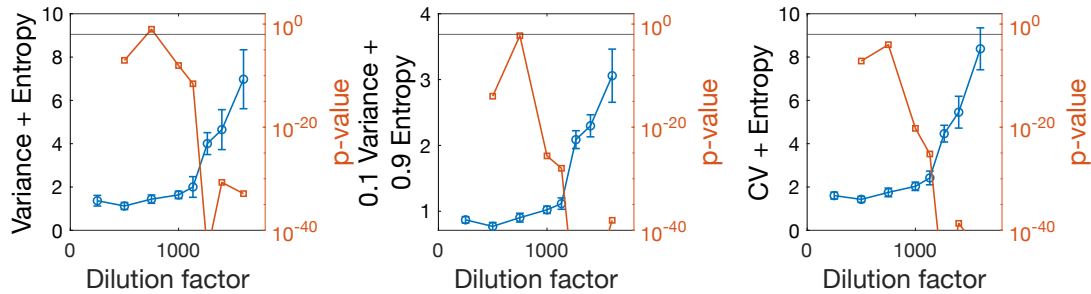


Figure 7: **Combined indicators calculated on empirical data.** The analysis is analogous to that in Fig. 6, using combinations of indicators. The horizontal solid lines mark $p\text{-value} = 0.01$.

302 3. Discussion

303 The paper provides a systematic classification of tipping mechanisms, highlights their
 304 underlying modelling assumptions, and bridges mathematical insights and observations of
 305 real systems to classify various tipping mechanisms, towards quantitative understanding and
 306 prediction of such relevant phenomena. The work shifts the focus from studying specific sys-
 307 tems, that may undergo some transitions, to studying transitions, along with their classes
 308 and properties, which can accommodate the behaviour of different systems. An interesting
 309 question for future studies will be to develop data-driven methods to classify each system
 310 within its corresponding class, much like those developed to distinguish stochastic or chaotic
 311 signals (Rosso et al., 2007). This will dramatically help the understanding of biological pro-
 312 cesses and guide the selection of EWS or other methods to anticipate critical transitions, as
 313 well as informing methods to reconstruct cell developmental trajectories like those proposed
 314 by Eugenio et al. (2014).

315 Moreover, we systematically investigate early warning signals associated to noisy bifurcation-
 316 induced transitions, key dynamical routes for the regulation and control of many natural
 317 processes. So far, EWS have been mostly studied in highly controlled computational settings,
 318 or checked on empirical data with alternate success. Our results make sense of previous ob-
 319 servations, help to define their range of applicability to reliably predict systems' behaviours,
 320 and allow to understand why spurious signals may be triggered in certain cases. We also
 321 assess whether and when EWS can be interpreted as anticipating or just-on-time detecting
 322 critical transitions in the presence of noise. By carefully analysing noise types and param-
 323 eter dynamics, we also extend previous results to more realistic settings, to guide real-world
 324 applications.

325 Using both analytical and computational methods, we observe that the variance – a
 326 highly employed indicator for EWS – may be sensitive to state-dependent noise, while $AC(1)$

327 can be skewed by ramping control parameters. Both are good indicators in case of quasi-
328 steady-state dynamics and Gaussian noise, with the ability to provide information about
329 augmented risk of tipping events. In the other cases, Shannon entropy is the most robust
330 and performing indicator and is suggested for applications in case of uncertain settings.
331 If precise information about noise type and intensity are available, constructing composite
332 indicators can improve the early-alerting performance, e.g., by combining CV and H_S .

333 The optimisation of composite indicators points to the use of machine learning methods
334 when abundant data are available (Bury et al., 2021), but also opens important caveats
335 for their application in real life: feature combinations may be optimised for certain settings
336 (e.g., noise intensity or type) but may be hardly generalisable for others. Our results remark
337 that training should be performed considering all possible combinations, or by first assessing
338 which critical transition class is being considering. Otherwise, misleading signals may be
339 triggered and wrong conclusions reached. On the other hand, our results can be used
340 for feature selection of more interpretable machine learning algorithms that leverage the
341 proposed composite indicators, insofar defined for a-priori assessment of systems that lack
342 big data.

343 This work provides results and guidelines for the application of early warning signals from
344 the critical transitions framework, but some points should still be covered by future stud-
345 ies. They include more refined analytical derivations of indicators in case of inhomogeneous
346 processes as well as closed formulae for entropy in exotic settings. Further investigations on
347 realistic systems, including non-autonomous transitions currently understudied in systems
348 biology, are thus suggested as extensions of our work. Another limitation of the present
349 study is the restriction to low dimensional systems. In principle, they are representative of
350 any system after dimension reduction techniques are applied, but it is necessary to assess if
351 and how the latter induce performance drops. Extending the analysis to high dimensional
352 systems, e.g., by testing multivariate indicators (Weinans et al., 2021) or further refining
353 EWS performance when multiple independent variables can be observed, is thus suggested
354 to future studies. Finally, our theoretical results have been verified on empirical data from
355 literature, but we acknowledge the need of performing additional experiments to continu-
356 ously validate our predictions. In particular, we suggest to design new experiments to test
357 the quantitative predictions about lead parameters and to assess what happens in case of
358 rapidly ramping parameters.

359 Our results can be readily tested and applied on real-world monitoring systems and can
360 inform the development of new indicators to address specific problems like cancer onset,
361 much like previous works (Chen et al., 2012) did using less performing measurements. In
362 addition, leveraging the sensitivity of indicators' trends to noise type and parameter dy-
363 namics can provide new methods to infer the latter from empirical data. For instance,
364 comparing Fig. 6 with Fig. 3 supports hypothesis of commensurable time scales between
365 intake of sucrose (affected by the dilution factor) and cells' growth in yeast experiments (Dai
366 et al., 2012); such hypothesis, to be confirmed using controlled experiments, could advance
367 our knowledge beyond the current slow-fast approximations (Del Vecchio et al., 2016). Sim-
368 ilarly, the prevalence of certain noise types can be inferred by comparing data and theory.
369 Overall, we connect theory and data, such that knowledge about the dynamical settings al-

370 lows optimising early warning signals, and analysis of statistical indicators enables inference
371 of dynamical properties.

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376 **Author contributions**

377 Conceptualization, D.P and A.S. and J.G.; methodology and investigation, D.P.; manuscript
378 writing, D.P and A.S. and J.G.; supervision, A.S. and J.G.; funding acquisition, A.S. and
379 J.G.

380 **Declaration of interests**

381 The authors declare no competing interests.

382 **Figure legends**

383 **Figure 1: Classification of transitions between states x of a dynamical sys-**
384 **tem, controlled by a slow changing parameter p .** x and p may also correspond to
385 network combinations of variables and parameters (Moris et al., 2016; Gao et al., 2016). (a):
386 Illustration of b-tipping and n-tipping using a bistable system with saddle-node bifurcations
387 (unstable branch in red; saddle-node template shown in inset). Hysteresis can occur, i.e.,
388 asymmetric routes to tipping from one stable state or from the other (orange, from up to
389 down with increasing p ; black, from down to up with decreasing p). B-tipping: the system
390 approaches the bifurcation point. The associated landscape is molded by p and the basin
391 of attraction becomes shallower (as visualised by the bars) until disappearing; there, the
392 system tips. N-tipping: if subject to strong fluctuations, depicted as wiggling of the red
393 ball, the system can be pushed over the barrier onto an alternative attractor, even before
394 the bifurcation point. (b) Illustration of r-tipping: rapid ramping of the control parameter
395 makes it as if the landscape shifts and the systems does not manage to move along, therefore
396 tipping onto another attractor “sliding” underneath. See Ashwin et al. (2012) for formal
397 definitions. (c): Example of “smooth” transition without hysteresis, using a dynamical sys-
398 tem close to a pitchfork bifurcation (inset) as template. To reproduce the plots, see STAR
399 Method C.3.

400
401 **Figure 2: Visual example of topological equivalence.** a) Plot for $dX/dt = f(x, p =$
402 $c') = X(1 - X/K) - c'X^2/(X^2 + 1)$, a model of harvested ecological populations (Scheffer
403 et al., 2009), also akin to Allee effects observed in microbiological colonies (Dai et al., 2012);
404 X is the population density, K is the carrying capacity and c' is the maximum harvest rate.

405 b) Plot for $f(x, c)$ of the autocatalytic loop model Eq. 15. c) Plot for $f(x, p)$ of the saddle
 406 node normal form $\dot{x} = -p - x^2$. The two realistic models are locally topologically equivalent
 407 to the normal form within the red rectangle (visual reference): they approach a bifurcation
 408 point, marked by $f(x, p)$ crossing the x-axis, as the parameter c' or c changes.

409
 410 **Figure 3: Trends of common statistical indicators.** We consider Var, AC(1) and
 411 CV for saddle-node, transcritical and pitchfork bifurcations as $p \rightarrow 0$, in different dynamical
 412 contexts (combinations of noise characteristics and stationarity for the control parameter).
 413 WN: white (Gaussian) noise; MN 1: mult noise with $h(\tilde{x}, p) = x$; MN 2: mult. noise with
 414 $h(\tilde{x}, p) = x^2$. As the autocorrelation is independent on noise, only MN 1 is show and it
 415 overlaps with the white noise case.

416
 417 **Figure 4: Bistable systems, studied with (a) characteristic curves or (b) bifur-**
 418 **cation diagram \tilde{x} (stable state) vs control parameter.** Among the systems undergoing
 419 saddle-node bifurcations, any linear system with nonlinear feedback and adequate feedback
 420 gain, such that the characteristic curve crosses the activation function in three points (two
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 422 function FB corresponds to the Hill function ($k = 2$), the feedforward FF to the linear part
 423 $-(K - x)$ with K in its appropriate range. The control parameter c tunes the FB function.
 424 Dashed-dotted line: c is not sufficient to promote bistability, corresponding to left stable
 425 region of (b). Dashed line: the critical value for which FB is tangent to FF , corresponding
 426 to saddle-node point, open circle in (b). Solid line: bistable system with three intersection
 427 points (stable, i and iii; unstable, ii). When studying the vector field $f(x, c)$ is easier than
 428 the characteristic curves, one can use the representation and interpretation in Fig. 2b. Note
 429 that the line styles have the same meaning in panel (a) and Fig. 2.

430
 431 **Figure 5: Optimisation of leading indicators for EWS, according to lead pa-**
 432 **rameter c_{sig} .** a) c_{sig} at various noise intensities σ , for all the most common indicators. b)
 433 The counter \mathcal{C} , normalised by all transitions to be interpreted as probability of n-tipping,
 434 at different noise intensities σ and distances $c - c_0$ from the bifurcation point. c) Scores
 435 S , corresponding to the argument of the cost function Eq. 19, for various combinations \mathcal{S}
 436 from Eq. 18. In the panel below, the color code shows the weights w_k for each indicator,
 437 in each combination. Results in panels a, b and c refer to $\alpha = 1$. For various types of $h(x)$
 438 and $\alpha = 0$, see Supplementary Figure S3. d) Optimal weights \hat{w} for each indicator, as a
 439 function of noise mixing α . As a representative of multiplicative noise, we used $h(x) = x$.
 440 Other $h(x)$ conserve the trends, albeit changing the corresponding c_{sig} . It may happen that
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442
 443 **Figure 6: Statistical indicators calculated on empirical data.** Data are from Dai
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 446 statistical moments of degree γ have units of measure $(\text{cells}/\mu\text{l})^\gamma$. The autocorrelation time
 447 is in days. The mean reproduces the upper stable branch of a saddle-node bifurcation dia-

448 gram (*cf.* Fig. 1) until the empirically estimated bifurcation point at $DF=1600$. Horizontal
449 solid lines mark $p\text{-value} = 0.01$.

450

451 **Figure 7: Combined indicators calculated on empirical data.** The analysis is
452 analogous to that in Fig. 6, using combinations of indicators. The horizontal solid lines
453 mark $p\text{-value} = 0.01$.

454 **STAR Method A. Key Resources**

455 Software: MATLAB R2021b (Matworks), Mathematica v12 (Wolfram)

456 Analysis and figures script: GitHub ([https://github.com/daniele-proverbio/EWS_o](https://github.com/daniele-proverbio/EWS_optimise)
457 `ptimise`)

458 **STAR Method B. Resource availability**

459 *STAR Method B.1. Material availability*

460 This study did not generate new materials

461 *STAR Method B.2. Data and code availability*

- 462 • All original code has been deposited at GitHub, [https://github.com/daniele-pro](https://github.com/daniele-proverbio/EWS_optimise)
463 `verbio/EWS_optimise`, and is publicly available.
- 464 • All data used are publicly available on Zenodo: Dai, L., Vorselen, D., Korolev, K.
465 S., Gore, J. (2012). Generic Indicators for Loss of Resilience Before a Tipping Point
466 Leading to Population Collapse. *Science*. [https://doi.org/10.1126/science.1219](https://doi.org/10.1126/science.1219805)
467 805

468 **STAR Method C. Method details**

469 *STAR Method C.1. Topological equivalence and normal forms*

470 Bifurcations model drastic changes in the qualitative behaviour of dynamical systems,
471 such as shifts in equilibria and regimes (Kuznetsov, 2013; Kuehn and Bick, 2021). Before
472 delving into bifurcations and their representation as normal forms, recall the concept of
473 topological equivalence.

474 Local topological equivalence between two dynamical systems $\{\mathcal{T}, \mathbb{R}^n, \phi^t\}$ and $\{\mathcal{T}, \mathbb{R}, \psi^t\}$
475 is established if there exist a homeomorphism $h : \mathbb{R}^n \rightarrow \mathbb{R}^n$ that maps orbits of the first
476 system to orbits of the second one, and the direction of time is preserved. Local topo-
477 logically equivalence near an equilibrium \hat{u} is, in turn, established between a dynamical
478 system $\{\mathcal{T}, \mathbb{R}^n, \phi^t\}$ and a dynamical system $\{\mathcal{T}, \mathbb{R}, \psi^t\}$ near an equilibrium \hat{y} if there exist a
479 homeomorphism $h : \mathbb{R}^n \rightarrow \mathbb{R}^n$ that is defined in a small neighborhood $U \in \mathbb{R}^n$ of \hat{u} , satisfies
480 $\hat{y} = h(\hat{u})$, and maps orbits of the $\{\mathcal{T}, \mathbb{R}^n, \phi^t\} \in U$ onto orbits of $\{\mathcal{T}, \mathbb{R}, \psi^t\} \in V = h(U) \subset \mathbb{R}^n$
481 while preserving the direction of time.

482

483 A bifurcation consists in the appearance of a topologically non-equivalent phase portrait
484 under variation of parameters. The difference between the dimension of the parameter space
485 and the dimension of the corresponding bifurcation boundary is called “codimension”.

486 To determine a system’s behaviour near bifurcations, minimal-order forms, called “nor-
487 mal forms”, can be employed. In fact, the normal form of the bifurcation is locally topologi-
488 cally equivalent near an equilibrium to all systems exhibiting that certain type of bifurcation
489 (Haragus and Iooss, 2010).

490 Consider a dynamical system

$$\dot{x} = f(x, p'), \quad x \in \mathbb{R}^n, \quad p' \in \mathbb{R}^n \quad (\text{C.1})$$

491 and a polynomial model

$$\dot{\zeta} = g(\zeta, p; \beta), \quad \zeta \in \mathbb{R}^n, \quad p \in \mathbb{R}^k, \quad \beta \in \mathbb{R}^l \quad (\text{C.2})$$

492 having dimension n , codimension k and polynomial order l . Without loss of generality,
493 a change of coordinates can set the bifurcation point occurs at $(x, p) = (0, p_0)$ (Strogatz,
494 2015). System C.2 is thus called a *topological normal form* for a given bifurcation if any
495 generic system C.1 with the equilibrium $x = 0$ satisfying the same bifurcation conditions at
496 $p' = 0$ is locally topologically equivalent near the origin to model (C.2) for some values of
497 the coefficients β_i . Using normal forms, it is thus possible to study classes of bifurcations
498 using simple polynomials. If the system satisfies certain conditions on $\partial^j f / \partial \varphi^j|_{(0, p_0)}$ around
499 the critical point, where j is the derivative order and $\varphi = \{x, p\}$, it is called “generic”. The
500 nondegeneracy conditions $\partial^j f / \partial x^j$ are related to the “criticality” of a bifurcation (Kuehn,
501 2011), while the transversality conditions $\partial^j f / \partial p^j$ govern the bifurcation unfolding and thus
502 its genericity (the bifurcation exists even after small perturbations). The saddle-node in-
503 vestigated in the main text (*cf.* Fig. 4) is the most common generic normal form with
504 dimension 1 and codimension 1 (Haragus and Iooss, 2010).

505
506 For low-dimensional systems, their associated normal forms can be explicitly obtained
507 using e.g. Taylor expansion methods over both nondegeneracy and transversality conditions
508 (Strogatz, 2015). For high-dimensional systems, numerical methods like XPP-AUT (<http://www.math.pitt.edu/~bard/xpp/whatis.html>) or network reduction techniques (Gao
509 et al., 2016; Tu et al., 2021) can be employed to infer or derive the normal forms. Obtaining
510 analytical results for any system is still an open research field.

512 *STAR Method C.2. Analysis of slow dynamics*

513 The fluctuations around the stable manifold of Eq. 5 can be analysed by studying the
514 fast-slow dynamics around it and determining stochastic equations for the residuals (Kuehn,
515 2011; Berglund and Gentz, 2006; O’Regan and Burton, 2018). Here, we briefly recall the
516 procedure to derive Eq. 6. Recall the normal form of a generic fold bifurcation:

$$\dot{x} = p + x^2, \quad (\text{C.3})$$

517 It has two steady states:

$$\hat{x}_1 = -\sqrt{-p-0} \quad (\text{stable}) \quad (\text{C.4})$$

518

$$\hat{x}_2 = -\sqrt{-p-0} \quad (\text{unstable}) \quad (\text{C.5})$$

519 where the term “ -0 ” explicits the distance from the bifurcation point $x_0 = 0$ (by definition).
 520 Consider a neighborhood of the attractor (stable fixed point) \hat{x}_1 and see what happens after
 521 small perturbations. To do so, perform a local linearization by considering $\delta x = (x - \hat{x}_1)$.

522 Thus:

$$\frac{d\delta x}{dt} \simeq f(\hat{x}_1) + \left. \frac{\partial f}{\partial x} \right|_{\hat{x}_1} \delta x + O(\delta x^2). \quad (\text{C.6})$$

523 So, using Eq. C.3 and Eq. C.4, we obtain:

$$\frac{d\delta x}{dt} \simeq 2\sqrt{-p}\delta x. \quad (\text{C.7})$$

524 This deterministic form can be augmented by a Wiener process with variance σ arbitrary
 525 multiplied by $h(x)$, representing non-Gaussian noise properties. This modelling choice con-
 526 verts the family of ODEs into SDEs (Berglund and Gentz, 2006; Namachchivaya and Leng,
 527 1990; Khas’minskii, 1966). A change $\delta x \rightarrow y$ makes the notation lighter into:

$$dy = 2\sqrt{-p}y dt + \sqrt{h^2(x)}dW. \quad (\text{C.8})$$

528 The equation describes a system evolving under small noise in a neighbourhood of the stable
 529 equilibrium, when this is not far away from the bifurcation point.

530 The term $(\hat{x}_1 - 0) = \sqrt{-p}$ is the distance of the stable equilibrium from the bifurcation
 531 point and depends on the leading parameter p . We can thus rescale it to a new variable $-k$:

$$dy = -k y dt + \sqrt{h^2(x)}dW \quad (\text{C.9})$$

532 The sign “ $-$ ” in “ $-k$ ” is included so that Eq. C.9 is interpreted as the associated Langevin
 533 equation to a Ornstein-Uhlenbeck process (Gardiner, 1985). The term multiplying the de-
 534 terministic drift can thus be interpreted as $-\partial V/\partial x$ where $V(x)$ is the potential governing
 535 the drift of the particle subjected to random noise. In our case, thanks to the choices made,

$$V = \frac{1}{2}k y^2, \quad (\text{C.10})$$

536 that is, a quadratically shaped adjoining potential typical of an overdamped oscillator under
 537 noise, of which k represents the depth. The working hypothesis is that boundary of the
 538 ideal potential V can grasp the boundary of the attracting basin of the original model after
 539 sufficiently long time. Eq. C.9 is analytically tractable to understand the main qualitative
 540 features of more complicated critical transitions. However, it requires ad hoc extensions
 541 when studying system-specific quantitative details like observability boundaries and lead
 542 times. Gardiner (1985) also extends Eq. C.9 to inhomogeneous processes with ramping
 543 parameters, used in Eq. 13.

544 *STAR Method C.3. Reproduce Fig 1*

545 Fig. 1 displays examples of a bistable system with critical transitions and hysteresis as
 546 well as smooth transitions. Panel (a) corresponds to the bifurcation diagram of Eq. 15,
 547 flipped along the vertical axis to highlight the hysteresis.

548 Panel (c) shows the bifurcation diagram, over an unfolded supercritical pitchfork bifur-
 549 cation, of $\dot{x} = q + p(x - 1) - (x - 1)^3$, which corresponds to the bifurcation normal form,
 550 shifted (to better visualize the diagram) and modified by a small perturbing term $q = 0.01$
 551 unfolding the bifurcation (Thompson and Sieber, 2011) into a smooth branch. In brief, an
 552 *unfolding* of a dynamical system under static equivalence is one that exhibits all possible
 553 bifurcations of the equilibrium (rest) points, up to topological equivalence of the set of equi-
 554 libria (Kuznetsov, 2013). In other terms, it investigates what happens when small terms are
 555 added to the original bifurcation, mimicking extra parameters, small offsets or “impurities”.

556 The illustrative attractors in panel (a) and (b) are two-well potentials associated, e.g., to
 557 the cusp bifurcation (aka “organising centre” Thompson and Sieber (2011); Eugenio et al.
 558 (2014)), a generic bifurcation described by $\dot{x} = a + bx - x^3$, where the combination of a and
 559 b determine bistability and the route to a saddle-node bifurcation.

560 *STAR Method C.4. Supporting analytical results*

561 *STAR Method C.4.1. Entropy in case of Gaussian noise*

Within a symmetric potential, elicited by a (locally) quadratic normal form, consider a
 Gaussian distributed variable $y \sim \mathcal{N}(\mu, \text{Var})$. Its entropy is:

$$\begin{aligned}
 H_S(y) &= - \int p(y') \log p(y') dy' = \\
 &= -\mathbb{E} [\log \mathcal{N}(\mu, \text{Var})] = \mathbb{E} \left[\log \left[\frac{1}{\sqrt{2\pi\text{Var}}} \exp \left(-\frac{1}{2\text{Var}}(x - \mu)^2 \right) \right] \right] = \\
 &= \frac{1}{2} \log (2\pi\text{Var}) + \frac{1}{2\text{Var}} \mathbb{E} [(x - \mu)^2] = \\
 &= \frac{1}{2} [\log(2\pi\text{Var}) + 1] ,
 \end{aligned} \tag{C.11}$$

562 That is, for the case of Gaussian noise, H_S is directly proportional to the variance and
 563 displays similar trends, that can be used to derive EWS.

564 *STAR Method C.4.2. Measurement noise*

565 Consider a measurement process with uncertainties σ_m^2 , independent from system vari-
 566 ance (Eq. 9). The resulting expected error, obtained from summing the two standard
 567 deviation in quadrature (Taylor, 1997), is:

$$\sigma_{\text{tot}}^2 = \text{Var} + \sigma_m^2 . \tag{C.12}$$

568 To derive the autocorrelation, combine its definition

$$AC(\tau) = \frac{\text{Cov}(x(t)x(t + \tau))}{\sqrt{\text{Var}(x(t))\text{Var}(x(t + \tau))}} = e^{-k \cdot |\tau|} \text{ for } t \rightarrow \infty \tag{C.13}$$

569 (where Cov indicates the covariance and Var the variance) with Eq. C.12 (substituting
 570 $Var = \sigma_{tot}^2$). In principle, we can explicitly consider multiplicative noise like in the main
 571 text. However, the goal in this case is to compute if notable discrepancies exist between
 572 ideal measurements (no uncertainty) and realistic measurements (with some uncertainty,
 573 that can be filtered to correspond to white noise). Hence, only the case of white process
 574 noise is currently considered. This results in:

$$AC(1)_m = \frac{\frac{\sigma^2}{2k} e^{-k}}{\sqrt{\left(\frac{\sigma^2}{2k} + \sigma_m^2\right)^2}}. \quad (C.14)$$

575 Obviously, $\lim_{\sigma_m^2 \rightarrow 0} AC(1)_m = AC(1)$. From Eq. C.14, we can immediately see that
 576 measurement uncertainties σ_m induce small scaling but do not alter the functional. Only
 577 relatively high measurement uncertainty levels change the absolute values of expected lag-1
 578 autocorrelation, but maintain the increasing patterns close to critical points.

579 *STAR Method C.4.3. Skewness and kurtosis*

580 For certain simulated systems, the third statistical moment (skewness) has been sug-
 581 gested to provide useful early warnings (Guttal and Jayaprakash, 2008). However, experi-
 582 mental results (Dai et al., 2012) were not able to confirm the expectations, estimating flat
 583 and fluctuating trends before a tipping point.

584 For a stochastic process with quasi-steady state parameter, its statistical moments are

$$\langle y^n \rangle - \langle y \rangle^n = \int_{-\infty}^{\infty} (y' - \mu)^n P(y') dy' \quad (C.15)$$

585 where $P(y')$ is the associated probability density function and μ is the expected average
 586 value.

587 For odd n , if $\mu = 0$ and $P(y')$ is symmetric, the integral equals 0 by definition. Symmetric
 588 probability density functions are associated, for instance, with quadratic potentials (Eq.
 589 C.10) that are typical of bifurcation normal forms under white noise, for which (Gardiner,
 590 1985)

$$P(y) = \sqrt{\frac{k}{\pi\sigma^2}} \text{Exp} \left[-\frac{2}{\sigma^2} U(y) \right] = \sqrt{\frac{k}{\pi\sigma^2}} \text{Exp} \left[-\frac{ky^2}{\sigma^2} \right] \quad (C.16)$$

591 Consequently, the normal forms – in particular, the saddle-node – considered above are
 592 expected to display a flat skewness.

593 On the other hand, the integral C.16 may be non-zero, and even depend on the drift
 594 parameter k , if $\mu \neq 0$ or if $P(y)$ is asymmetrical. In the first case, solving Eq. C.16 yields
 595 (provided that $\text{Re}[k] > 0$):

$$\text{Skew} = -\frac{\mu(3 + 2k\mu^2)}{2k}. \quad (C.17)$$

596 In this case, as $k \rightarrow 0$, the skewness is expected to increase, potentially providing an early
 597 warning

598 On the other hand, an asymmetric potential can be obtained in case of multiplicative
 599 noise (Gardiner, 1985; Sharma et al., 2016). Depending on the specific form, it may be
 600 possible to observe increasing trends associated to EWS, but they may be system-specific
 601 and not generalisable. In this sense, there is no ambiguity between the results of Guttal
 602 and Jayaprakash (2008) and Dai et al. (2012): they were studying systems with different
 603 properties, using an indicator that is not particularly performing and generalisable.

604
 605 As for the kurtosis, in case of $\mu = 0$ (typical white noise), kurtosis = 3Var^2 . This can
 606 be obtained by solving Eq. C.16. If $\mu \neq 0$, or for other exotic noise forms, and if $\text{Re}[k] > 0$:

$$\text{Kurt} = \frac{3 + 4k\mu^2(3 + k\mu^2)}{4k^2}, \quad (\text{C.18})$$

607 whose leading term for $0 < k < 1$ still equals Var^2 . Hence, the variance is already represen-
 608 tative of higher moments, which are not expected to improve EWS unless system-specific
 609 noise and drift forms are considered. Note that, for both Eq. C.17 and Eq. C.18, the
 610 constant noise level σ is normalised to 1 for ease of notation.

611 *STAR Method C.5. Computational simulations*

612 In all computer simulations of Eq. 15, $K = 0.1$ to set bistability. The analysis con-
 613 centrates on the upper stable branch of the bifurcation diagram (Fig. 4, right) to compare
 614 with white noise results. In this case, multiplicative noise corresponds to intrinsic regulatory
 615 mechanisms (Hasty et al., 2000; Norman et al., 2015) rather than stochasticity due to small
 616 numbers (Gillespie, 2000). Simulations are performed in MATLAB (R2021b) using the Mil-
 617 stein method with a time step of 0.01 (arbitrary units). For quasi-steady state simulations,
 618 distributional data for each c from far to close the bifurcation point are computed upon
 619 stable values of system's state, after a transient.

620 The Milstein method runs Monte Carlo chains over Itô-Taylor expanded stochastic dif-
 621 ferential equations for any variable z , up to second order:

$$z(t_i + 1) = z(t_i) + f(z(t_i))\Delta t + g(z(t_i))\Delta W_i + \frac{1}{2}g(z(t_i))g'(z(t_i)) [(\Delta W_i)^2 - \Delta t]; \quad (\text{C.19})$$

622 It better converges to the true Itô integral and was proven to have improved accuracy
 623 (Bayram et al., 2018). When $g(z(t)) = \text{const}$ (only additive noise without state-dependency),
 624 it is equivalent to the common Euler-Maruyama scheme.

625
 626 Setting simulation parameters of noise intensity and distance to critical points require
 627 understanding their reciprocal scales. To do so, we employ a methodology introduced in
 628 (Kuehn, 2011; Proverbio et al., 2022b), that is, to look for significant changes in the Kramers'
 629 escape rates out of bistable potentials. The Kramers escape rate is (Gardiner, 1985)

$$\tau = 2\pi(\sqrt{|U''(\tilde{x}_1)U''(\tilde{x}_2)|})^{-1} \exp[(U(\tilde{x}_2) - U(\tilde{x}_1))/\sigma] \quad (\text{C.20})$$

630 and measures the average expected rate of escape of multiple noisy particles from attracting
 631 wells. For any saddle-node bifurcation $\dot{x} = p - x^2$ equipped with additive noise, $U(\tilde{x}_2) -$
 632 $U(\tilde{x}_1) = 32/3p^{3/2}$ and $|U''(\tilde{x}_{1,2})| = 2\sqrt{p}$. Hence,

$$\tau \simeq \mathcal{O} \left(\exp \left[\frac{p^{3/2}}{\sigma} \right] \right) \quad (\text{C.21})$$

633 Comparable ranges of control parameters and noise levels are studied in (Proverbio et al.,
 634 2022b) and reproduced in Supplementary Figure S4. We use those results to distinguish two
 635 regimes, one where few noise-induced transitions might occur and another regime primar-
 636 ily determined by the approach to the bifurcation. We set values of $c - c_0$ (distance from
 637 bifurcation point) and σ (noise intensity) accordingly, to span both regimes and see what
 638 changes when n-tipping becomes more frequent.

639 Finally, the statistical indicators are computed using their standard definitions, using
 640 their corresponding MATLAB functions. For example, variance and Shannon entropy H_S
 641 are:
 642

$$\text{Var}_j = \frac{1}{N-1} \sum_{r=1}^N (B_{j,r} - \hat{B}_j)^2 \quad (\text{C.22})$$

$$H_S = - \sum p_j \log p_j \quad (\text{C.23})$$

643 for any point j corresponding to a single parameter value, with N data B distributed
 644 around a mean value \hat{B} and probability density function p_j . Other statistical moments and
 645 indicators can be computed similarly.
 646

647 *STAR Method C.6. p-value assessment of significant increase and optimisation*

648 By theory, an early warning signal is triggered when an increasing trend of suitable sta-
 649 tistical indicators is observed (Scheffer et al., 2009). However, during real-time monitoring,
 650 it is often challenging to say whether a measured increase of mean values is significant or
 651 not, due to random fluctuations and uncertainties that may occur. If increasing trends are
 652 not quantified properly, spurious signals may be triggered (Boettiger and Hastings, 2012).
 653 For analysis performed using moving windows over time-series data, the Kendall's τ score of
 654 monotonous increases have been proposed (Boettiger and Hastings, 2012; Proverbio et al.,
 655 2022a), as well as threshold of confidence intervals, with respect to baseline values (Drake
 656 and Griffen, 2010).

657 Since we work with distributional data, we propose to employ significance levels on
 658 Welch's p-value scores (non-equal variances allowed between the populations), which relate
 659 to threshold in confidence intervals and are readily interpretable (Proverbio et al., 2022b).
 660 They also allow to estimate the sensitivity to noise intensities and the expected lead pa-
 661 rameter for detection or anticipation of critical transitions. The idea is to compare the full
 662 distributions at each parameter value c with a reference one, usually taken far from the
 663 bifurcation point and without n-tipping, and check whether they are significantly separated
 664 towards increasing values. The p-value scores are used to assess the significance. This

665 method can still be sensitive to fluctuating scores (hence, a smoothing is employed), but it
666 has the advantage of relying on a-prioristic values, e.g. significant p-value $p_{sig} = 0.05$. Of
667 course, a p-value does not distinguish between increasing or decreasing trends: it is thus
668 coupled with simple visualization of the direction of the trends.

669 Examples of the three methods are provided in Supplementary Figure S1.

670
671 Quantifying the significance of increasing trends is leveraged as follows: we extract
672 at which value of the control parameter c the p-value crosses the significance threshold
673 $p_{sig} = 0.05$ as a reference. Other common thresholds $p = 0.1$ or $p = 0.01$ can be used,
674 yielding consistent results. When p-value $< p_{sig}$, it means that an indicator has significantly
675 increased more than the baseline, triggering a warning signal. Consider all c_i tested during
676 the simulations, $i = 1..N$ with $N = (c_{max} - c_{min})/0.002$; c_{max} and c_{min} are two arbitrary
677 values greater and lower than the bifurcation value c_0 , within the bistable region, and 0.002
678 is the simulation step $|c_i - c_{i-1}|$. Out of all c_i , estimate $c_{sig} = c_j$, where j is the first index
679 at which p-value $_j < p_{sig}$ stably, *i.e.*, without considering small fluctuating values (for that,
680 a smoothing is employed). This is performed for each indicator \mathcal{I} and each noise level σ .
681 Hence, the analysis estimates

$$c_{sig}^{\mathcal{I}}(\sigma) = c_j \text{ s.t. } \text{p-value}_j(\mathcal{I}) < p_{sig} \wedge \min(j) . \quad (\text{C.24})$$

682 The optimisation problem described in the main text aims at maximising the combination
683 of all $c_{sig}^{\mathcal{I}}(\sigma_l)$ obtained at different noise levels σ_l , so that the results are robust against a
684 range of signal-to-noise ratios. As described in the main text, the analysis is complemented
685 with a counter \mathcal{C} to quantify how many tipping events occurred before the bifurcation point,
686 for each σ .

687
688 A final comment regards the set of considered indicators \mathcal{I} . In principle, CV could
689 be included among the as its performance improves in case of multiplicative noise (see
690 Supplementary Figure S3. However, the optimisation procedure does not strongly select it,
691 preferring the combinations in Fig. 5d. Hence, it has been removed altogether, to improve
692 the computational speed when using more fine-grained steps for the grid search.

693 *STAR Method C.7. Data collection and analysis*

694 Experimental data were collected and curated by the original study (Dai et al., 2012).
695 We refer to it for details about the experimental protocols. The publicly available data
696 correspond to ensemble of replicate populations, at each observation time corresponding to
697 input dilution factors altering the environmental sucrose concentration. The eight dilution
698 are 250, 500, 750, 1000, 1133, 1266, 1400 and 1600. Population densities were recorded by
699 measuring optical density at 620 nm using a Thermo Scientific Multiskan FC microplate
700 photometer. The values used in the analysis represent cell numbers, estimated from optical
701 densities converted through calibration curves described in the original publication. For
702 each observation time, several statistical indicators were calculated over the ensembles as
703 explained in the previous section.

704 The standard errors and confidence intervals of the indicators were given by bootstrap.
705 In bootstrap, the replicates are resampled by combining the data over 5 days (observation
706 lag for one dilution factor) into a single distribution. Resampling was performed by 50 to
707 1000 repetitions, to check the robustness of final p-values against bootstrapping hyper-
708 parameters and to confirm consistency with the original results. Since there are, on average,
709 60 data entries for each dilution factor value, we eventually employ bootstrapping with 50
710 repetitions, to avoid biases in the p-values due to random over-repetitions of some data.

711 The p-values to quantify significant increases in the distributions of indicators are cal-
712 culated as described in STAR Method C.6, using the distribution at dilution factor 250
713 (the smallest and furthest from the bifurcation point) as baseline, and comparing all other
714 distributions against it, making sure that the mean value was increasing before drawing
715 conclusions.

716 STAR Method D. Supplementary information titles and legends

717 **Figure S1: Quantitative definition of EWS.** Left: Example of looking for trends
718 past thresholds of confidence intervals. In this case, past the 2σ interval (dashed line) over
719 the uncertainty of the rightmost point, used as baseline far from the bifurcation value c_0 .
720 Centre: Example of Kendall's τ estimation. Compare the trends within two sliding win-
721 dows. If the new one is monotonously increasing with respect to the old one, $\tau > 0$, while
722 no increase corresponds to $\tau = 0$; the steeper the trend, the higher τ . Right: Example of
723 p-value between two distributions corresponding to different parameter values: the baseline,
724 corresponding to the rightmost c , and another generic c' . Each distribution corresponds to
725 an average value of the statistical indicator (superimposed and shifted for visualization pur-
726 poses). p-values's significance can be checked with standard statistical methods, to assess
727 whether the registered increase is significant or not. All figures use variance computed from
728 simulations of Eq. 15 in main text, with $n = 2$, $K = 0.1$ and $\sigma = 0.02$. c_0 is the critical
729 value for bifurcation point.

730
731 **Figure S2: Trends of notable indicators before and after the bifurcation point**
732 c_0 . It is displayed as a function of the control parameter c from Eq. 15 of main text. The in-
733 creasing trends yield early warning signals. The violet ribbon represents confidence intervals
734 of 2 standard deviations, estimated from repeated simulations. Indicators are: Variance,
735 lag-1 Autocorrelation, Skewness, Kurtosis, Coefficient of Variation, Index of Dispersion,
736 Shannon Entropy (H_S). Note that some of them peak at the transition point, while others
737 don't due to noise-induced transitions altering their expected trends. All simulations are
738 performed with white noise, $\sigma = 0.012$.

739
740 **Figure S3: Dependency of c_{sig} (Eq. C.24 of main text) for each considered**
741 **indicator \mathcal{I} and noise intensity σ .** It is displayed with the corresponding counting \mathcal{C} of
742 noise-induced transitions happening before the bifurcation point, at each noise intensity σ .
743 Different multiplicative noise types are considered (*cf.* Eq. 17 of main text): a) $h(x) = x$.
744 b) $h(x) = x^2$. c) $h(x) = x^2/(1 + x^2)$. Due to differing fluctuation types, the indicators have

745 different performances in identifying the lead parameter. Conserved patterns are: entropy
746 H_S is normally the best, particularly for high σ ; Skew and Kurt perform poorly. $AC(1)$
747 follows H_S closely, but with slightly lower c_{sig} . Var and ID are normally worse than CV,
748 as they are less sensitive to mean values. Notably, CV works better than in the case of
749 white noise (compare with Fig. 5 of main text) but it still lags behind H_S , particularly in
750 case of low σ . Note that several n-tipping occur before the bifurcation point as σ increases,
751 except for $h(x) = x^2/(1 + x^2)$ that better buffers the system variability, as also noted in
752 (Proverbio et al., 2022b). Particularly for this case, the main indicators provide anticipating
753 signals (around $c_{sig} \geq 0.05$) while n-tipping starts around $c \simeq 0.02$. In the other cases, the
754 indicators are normally providing early warnings, except in the case $\sigma > 0.046$ for which
755 they may just-on-time detect the few n-tipping events already happening.

756
757 **Figure S4: Kramers' escape rate τ as a function of noise level σ and p (distance**
758 **from bifurcation point).** Its analytical form is in Eq. C.21 of main text of the main text.
759 We use the boundary colored in yellow as a proxy to set commensurable magnitudes between
760 control parameter and noise intensity in computer simulations.

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