Temperature control of electromigration to form gold nanogap junctions

G. Esen and M. S. Fuhrer^{a)}

Department of Physics and Center for Superconductivity Research, University of Maryland, College Park, Maryland 20742-4111

(Received 3 August 2005; accepted 25 October 2005; published online 19 December 2005)

Controlled electromigration of gold nanowires of different cross-sectional areas to form nanogap junctions is studied using a feedback method. A linear correlation between the cross-sectional area of the gold nanowires and the power dissipated in the junction during electromigration is observed, indicating that the feedback mechanism operates primarily by controlling the temperature of the junction during electromigration. We also show that the role of the external feedback circuit is to prevent thermal runaway; minimization of series resistance allows controlled electromigration to a significant range of junction resistances with a simple voltage ramp. © 2005 American Institute of Physics. [DOI: 10.1063/1.2149174]

Electromigration (EM), the electrical current-induced diffusion of atoms in a thin metal film, is important as a failure mode in integrated circuit interconnects.¹ Failure of a narrow metal wire due to EM has also recently been utilized extensively to prepare stable electrical contact pairs with nanometer separation for single molecule electrical experiments.^{2,3} It has been suggested^{2–4} that the dominant failure mechanism in such electrically stressed gold nanowires is thermally assisted EM. Recently, feedback schemes for controlling the rate of EM were shown to allow control over the final junction resistance (and presumably gap size) in such nanogap junctions.^{5,6}

The empirical formula for the median time to failure MTTF= $Aj^{-n}e^{E_d/kT}$, known as the Black equation⁷ has been employed to determine interconnect reliability, where A is a sample dependent constant, *j* is the current density, *n* is an exponent empirically found between 1 and 7,⁸ E_a is an activation energy for atomic motion, *k* is the Boltzmann constant, and *T* is the temperature. However, in the formation of nanogap junctions by the EM of a short nanowire, the current and temperature (due to Joule heating) are both changing rapidly. Understanding the role of current and temperature is then critical to the design of circuits to produce nanogap junctions in a controlled manner.

In recent works,^{5,6} it was proposed that controllable EM occurs at constant applied power. Furthermore, it was stated⁵ that electromigration is triggered at a constant temperature. In this letter, by experimenting on different nanowire geometries, we confirm that EM occurs at constant temperature independent of geometry, and show that the feedback mechanism works primarily through controlling the temperature of the electromigrating junction. Furthermore, we estimate the temperature of the junction during EM to be only a few hundred Kelvins, low enough to allow the study of many molecular adsorbates without desorption or dissociation of the molecules. We also show that the role of the external feedback circuit is to prevent thermal runaway in the junction, allowing control in the region in which the temperature increases with increasing junction resistance. This region can be minimized by minimizing series resistance in the circuit, allowing significant control of the EM with a simple voltage ramp.

Our devices are fabricated in two steps using conventional electron-beam lithography (EBL) and lift off on SiO_2 (500 nm)/Si substrates. We first fabricate thin gold lines (with no adhesion layer) of width 40 nm to 100 nm and thickness 15 to 30 nm. In the second EBL step, we deposited 5 nm Cr and 70 nm Au to form contacts and bonding pads. A scanning electron microscope (SEM) micrograph of such a device is shown in the inset of Fig. 1.

To control the EM process, we used a computercontrolled feedback scheme similar to Ref. 5, consisting of the following steps: We first measure a reference conductance value at a voltage of 100 mV. We then increase the voltage until the conductance drops by a set fraction (typically 2-5%) of the reference conductance value. At this point, the voltage bias is ramped down by 50 to 100 mV (at



FIG. 1. (Color) Current vs bias voltage during the feedback-controlled electromigration of an Au wire at T=1.3 K. Part A is a smooth curve indicating than the EM has not begun, whereas in Part B the resistance of the line increases irreversibly due to EM. Both Parts A and B are recorded in a single voltage biasing process, producing a final resistance of ~120 Ω . At this point, the voltage was reduced to zero for some time. When the bias process was restarted in C, the wire resistance is the same, demonstrating that the EM process may be frozen by turning off the voltage. The inset shows the SEM micrograph of one of our devices. The scale bar in the inset is 2 micrometers long. Arrows indicate the progression of the curve.

87. 263101-1

^{a)}Author to whom correspondence should be addressed; electronic mail: mfuhrer@physics.umd.edu

^{© 2005} American Institute of Physics

Downloaded 27 Jan 2006 to 129.2.40.177. Redistribution subject to AIP license or copyright, see http://apl.aip.org/apl/copyright.jsp



FIG. 2. Power dissipated in the junction P_J vs the voltage drop at the junction V_J . The irreversible change in resistance due to EM starts at the point labeled (a). Inset shows the corresponding current vs bias voltage data. The starting nanowire has dimensions 830 nm long × 60 nm wide × 25 nm thick; the length and width of the nanowire is determined using SEM, and thickness by quartz crystal monitor during gold film deposition.

a rate of 50 mV/s) and a new reference conductance value is measured. We repeat this process until the desired conductance is reached. We performed our experiments in a gas flow 4 He cryostat at substrate temperatures from 1.3 K to room temperature.

A representative current versus voltage (I-V_{bias}) curve taken through a feedback-controlled EM process for one of the devices is shown in Fig. 1. The data labeled A show a smooth I versus V_{bias} curve indicating that EM has not begun in the gold wire. Although there is a resistance increase with increasing bias in A, we found that if we stop the voltage bias in Region A, this resistance increase is reversible. Such a reversible resistance increase shows that the gold wire heats up before EM begins. The data labeled B show that after this initial heating, the gold wire begins to change resistance irreversibly due to EM. The data labeled C show that one can stop and restart the voltage bias before the gold wire totally fails. I-Vbias curves of two bias processes perfectly match each other indicating that in the second biasing process the gold wire first heats up to the temperature where significant EM takes place, and then EM restarts.

We now discuss the temperature of the wire during EM. We assume the total resistance of the circuit *R* is the sum of two resistors, R_L , the lead resistance (equal to the total measured resistance at low bias) and R_J , the resistance of the "junction," the weak spot formed in the wire by EM; i.e., $R_J=0$ initially. The power dissipated in the junction is then $P_J=I^2R_J$. Note that R_J includes a contribution from the resistance change of the leads upon heating; and heating due to the resistance of the nanowire itself is ignored. Hence, P_J is a rough estimate of the power that is heating the junction, but should be valid when R_J is significantly non-zero.

Figure 2 shows P_J versus $V_J = IR_J$, where the inset graph is the corresponding I- V_{bias} curve of the data. After the junction begins to increase resistance due to EM [Point (a)], the power dissipated in the junction reaches a relatively constant value. Constant power dissipation in the junction is observed over an order of magnitude of junction voltage.



FIG. 3. Power dissipated in the junction during electromigration vs nanowire cross-sectional area. The power is the average power in the region of near-constant power seen in Fig. 2.

Figure 3 shows that the power dissipated in the junction during EM is proportional to the nanowire area. Since the thermal conductance of the wire is proportional to the crosssectional area, the power required to maintain a given wire temperature should also be proportional to area. Thus, we conclude from Fig. 3, and the relatively constant power observed in Fig. 2, that the EM rate is dominated by temperature, and the feedback scheme operates to control the wire temperature. This conclusion is reasonable from the Black equation, given that temperature enters exponentially, and voltage (through current density) only algebraically.

We now estimate the temperature of the junction T_J during EM. Here, we neglect the heat conduction to SiO₂ substrate by considering the relative magnitude of the thermal conductivity of gold and SiO₂ (Ref. 9) and consider the contacts as infinite heat sinks at T=1.5 K. The temperature at the midpoint of a wire with uniform power generation over its volume is $T_I = PL^2/8V\kappa$ where P, V, L, and κ are, respectively, the total power generated in the nanowire, the volume and length of nanowire, and the thermal conductivity of gold.¹⁰ For the wire in Fig. 2, the maximum power is estimated as ~0.67 mW (including P_I at point [Fig. 2(a)] and additional power generated due to the estimated resistance 11 Ω of the nanowire at T=1.5 K); using the thermal conductivity of gold as $\geq 320W/mK$ (the room-temperature value for bulk gold) we estimate $T_I \leq 145$ K. If we consider instead that all of the power is being generated at the center of the nanowire and carried out to the leads by the nanowire, then $T_J = PL^2/4V \kappa \leq 290$ K; which is still low enough to allow the study of many molecular adsorbates without desorption or dissociation of the molecules.

We now discuss why the feedback process is feasible. The thermal time constant¹¹ $\tau_{th} = L^2 \rho C_p / \pi^2 \kappa$, where C_p is the specific heat, and ρ is the density, is less than 1 ns in our wires; much faster than our external feedback circuit. This suggests that the EM process itself must occur very slowly. If this is the case, is feedback needed at all? i.e., could V_{bias} simply be turned to zero at the desired R? The answer is no—at constant junction temperature (i.e., following the $I-V_{\text{bias}}$ curve in, e.g., Fig. 1) I, and hence R, is multiple valued at a given V_{bias} . Stated another way, the change in temperature with junction resistance dT_J/dR_J during EM must be negative to prevent thermal runaway once EM begins. Assuming that T_J is proportional to I^2R_J , for our simple se-

Downloaded 27 Jan 2006 to 129.2.40.177. Redistribution subject to AIP license or copyright, see http://apl.aip.org/apl/copyright.jsp



FIG. 4. (Color) Current vs bias voltage during the EM of three similar gold wires $(600-700 \text{ nm } \log \times 40 \text{ nm } \text{wide} \times 15 \text{ nm } \text{thick})$ at T=1.3 K. For the blue and red curves, the external feedback is turned off at the points marked by the blue and red arrows respectively. For the black curve, no external feedback was used.

ries circuit model $dT_J/dR_J < 0$ implies $R_J > R_L$, corresponding to the stable branch of the $I-V_{\text{bias}}$ curve, where $dV_{\text{bias}}/dR > 0$ during EM (the voltage increases as EM progresses).

Figure 4 illustrates this instability. We turned off the feedback at various points during the EM process, solely ramping the voltage upwards at a fixed rate. The red curve shows the feedback turned off while on the stable $I-V_{\text{bias}}$ branch (positive dV_{bias}/dR); the current decreases smoothly with increasing voltage from this point. However, when the feedback is turned off on the unstable $I-V_{\text{bias}}$ branch (negative dV_{bias}/dR ; blue and black curves), the current drops rapidly to the stable branch at the same V_{bias} . Thus, the feedback scheme is only necessary to produce final resistances $R_J < R_L$; with suitable circuit design (minimization of R_L ; i.e., short nanowires with highly conducting leads) small final R_J 's may be produced using a simple voltage ramp. Note that in some circuits (e.g., Fig. 1) the stable $I-V_{\text{bias}}$ branch does

not extend beyond the unstable branch; in such cases, a simple voltage ramp causes abrupt failure of the wire by melting (as observed via post-mortem SEM), resulting in large (>10 nm) gaps.

In conclusion, we performed controllable EM on nanowires with different cross sections. We found that the average power dissipated in the junction during EM increases linearly with the area of the junction indicating temperature control of the process and confirming that the mechanism is thermally assisted EM. Using the maximum power dissipated in a typical device, we estimate the junction temperature during EM performed at T=1.5 K to be only a few hundred Kelvins. We also note that the role of the feedback process in controlling EM is to prevent thermal runaway in the region of positive dT_J/dR_J . This region can be reduced by reducing the series resistance in the circuit, allowing controlled EM with a simple voltage ramp.

This work has been supported by the U.S. Department of Energy under Grant No. DE-FG02-01ER45939 and by the NSF-Materials Research Science and Engineering Center under Grant No. DMR-00-80008.

- ¹J. R. Lloyd, Semicond. Sci. Technol. **12**, 1177 (1997), and the references therein.
- ²H. Park, A. K. L. Kim, J. Park, A. P. Alivisatos, and P. L. McEuen., Appl. Phys. Lett. **75**, 301 (1999).
- ³H. Park, J. Park, A. K. L. Lim, E. H. Anderson, A. P. Alivisatos, and P. L. McEuen, Nature (London) **407**, 57 (2000).
- ⁴M. F. Lambert, M. F. Goffman, J. P. Bourgoin, and P. Hesto, Nanotechnology **14**, 772 (2003).
- ⁵D. R. Strachan, D. E. Smith, D. E. Johnston, T. H. Park, M. J. Therien, D. A. Bonnell, and A. T. Johnson, Appl. Phys. Lett. **86**, 043109 (2005).
- ⁶A. A. Houck, J. Labaizewicz, E. K. Chan, J. A. Folk, and I. L. Chuang, Nano Lett. 5, 1685 (2005).
- ⁷J. R. Black, Proc. IEEE **57**, 1587 (1969).
- ⁸D. Young and A. Christou, IEEE Trans. Reliab. 43, 186 (1994).
- ⁹M. G. Burzo, P. L. Komarov, and P. E. Raad, *IEEE 2002 Inter Society Conference on Thermal Phenomena* (2000), pp. 142–149.
- ¹⁰F. P. Incorpera and D. P. DeWitt, *Introduction to Heat Transfer*, 4th ed. (Wiley, New York, 2002).
- ¹¹L. Lu, W. Yi, and D. L. Zhang, Rev. Sci. Instrum. 72, 2996 (2001).