

**This is an electronic reprint of the original article.
This reprint *may differ* from the original in pagination and typographic detail.**

Author(s): Gabriel, Paternain; Salo, Mikko; Gunther, Uhlmann

Title: Tensor tomography on surfaces

Year: 2013

Version:

Please cite the original version:

Gabriel, P., Salo, M., & Gunther, U. (2013). Tensor tomography on surfaces. *Inventiones mathematicae*, 193(1), 229-247. <https://doi.org/10.1007/s00222-012-0432-1>

All material supplied via JYX is protected by copyright and other intellectual property rights, and duplication or sale of all or part of any of the repository collections is not permitted, except that material may be duplicated by you for your research use or educational purposes in electronic or print form. You must obtain permission for any other use. Electronic or print copies may not be offered, whether for sale or otherwise to anyone who is not an authorised user.

TENSOR TOMOGRAPHY ON SURFACES

GABRIEL P. PATERNAIN, MIKKO SALO, AND GUNTHER UHLMANN

ABSTRACT. We show that on simple surfaces the geodesic ray transform acting on solenoidal symmetric tensor fields of arbitrary order is injective. This solves a long standing inverse problem in the two-dimensional case.

1. INTRODUCTION

Let (M, g) be a compact oriented two-dimensional manifold with smooth boundary. We consider the geodesic ray transform acting on symmetric m -tensor fields on M . When the metric is Euclidean and $m = 0$ this transform reduces to the usual X-ray transform obtained by integrating functions along straight lines. More generally, given a symmetric (covariant) m -tensor field $f = f_{i_1 \dots i_m} dx^{i_1} \otimes \dots \otimes dx^{i_m}$ on M we define the corresponding function on SM by

$$f(x, v) = f_{i_1 \dots i_m} v^{i_1} \dots v^{i_m}.$$

Here $SM = \{(x, v) \in TM; |v| = 1\}$ is the unit circle bundle. Geodesics going from ∂M into M are parametrized by the set $\partial_+(SM) = \{(x, v) \in SM; x \in \partial M, \langle v, \nu \rangle \leq 0\}$ where ν is the outer unit normal vector to ∂M . For $(x, v) \in SM$ we let $t \mapsto \gamma(t, x, v)$ be the geodesic starting from x in direction v . We assume that (M, g) is nontrapping, which means that the time $\tau(x, v)$ when the geodesic $\gamma(t, x, v)$ exits M is finite for each (x, v) in SM .

The ray transform of f is defined by

$$If(x, v) = \int_0^{\tau(x, v)} f(\varphi_t(x, v)) dt, \quad (x, v) \in \partial_+(SM),$$

where φ_t denotes the geodesic flow of the Riemannian metric g acting on SM . If h is a symmetric $(m - 1)$ -tensor field, its inner derivative $d_\sigma h$ is a symmetric m -tensor field defined by $d_\sigma h = \sigma \nabla h$, where σ denotes symmetrization and ∇ is the Levi-Civita connection. A direct calculation in local coordinates shows that

$$d_\sigma h(x, v) = Xh(x, v),$$

where X is the geodesic vector field associated with φ_t . If additionally $h|_{\partial M} = 0$, then one clearly has $I(d_\sigma h) = 0$. The ray transform on symmetric m -tensors is said to be s -injective if these are the only elements in the kernel. The terminology arises from the fact that any tensor field f may be written uniquely as $f = f^s + d_\sigma h$, where f^s is a symmetric m -tensor with zero divergence and h is an $(m - 1)$ -tensor with $h|_{\partial M} = 0$ (cf. [17]). The tensor fields f^s and $d_\sigma h$ are called respectively the *solenoidal* and *potential* parts of the tensor f . Saying that I is s -injective is saying precisely that I is injective on the set of solenoidal tensors.

In this paper we will assume that (M, g) is *simple*, a notion that naturally arises in the context of the boundary rigidity problem [8]. We recall that a Riemannian manifold with boundary is said to be simple if the boundary is strictly convex and given any point p in M the exponential map \exp_p from $\exp_p^{-1}(M) \subset T_p M$ onto M is a diffeomorphism. In particular, a simple manifold is nontrapping.

The next result shows that the ray transform on simple surfaces is s -injective for tensors of any rank. This settles a long standing question in the two-dimensional case (cf. [13] and [17, Problem 1.1.2]).

Theorem 1.1. *Let (M, g) be a simple 2D manifold and let $m \geq 0$. If f is a smooth symmetric m -tensor field on M which satisfies $I f = 0$, then $f = d_\sigma h$ for some smooth symmetric $(m - 1)$ -tensor field h on M with $h|_{\partial M} = 0$. (If $m = 0$, then $f = 0$.)*

The geodesic ray transform is closely related to the boundary rigidity problem of determining a metric on a compact Riemannian manifold from its boundary distance function. See [24, 7] for recent reviews. The case $m = 0$, that is, the integration of a function along geodesics, is the linearization of the boundary rigidity problem in a fixed conformal class. The standard X-ray transform, where one integrates a function along straight lines, corresponds to the case of the Euclidean metric and is the basis of medical imaging techniques such as CT and PET. The case of integration along more general geodesics arises in geophysical imaging in determining the inner structure of the Earth since the speed of elastic waves generally increases with depth, thus curving the rays back to the Earth surface. It also arises in ultrasound imaging, where the Riemannian metric models the anisotropic index of refraction. Uniqueness and stability for the case $m = 0$ was shown by Mukhometov [9] on simple surfaces, and Fredholm type inversion formulas were given in [14].

The case $m = 1$ corresponds to the geodesic Doppler transform in which one integrates a vector field along geodesics. This transform appears in ultrasound tomography to detect tumors using blood flow measurements and also in non-invasive industrial measurements for reconstructing the velocity of a moving fluid. In the case $m = 1$, s -injectivity was shown by Anikonov and Romanov [1] and sharp stability estimates were proven in [23].

The integration of tensors of order two along geodesics arises as the linearization of the boundary rigidity problem and the linear problem is known as deformation boundary rigidity. Sharafutdinov [20] showed s -injectivity in this case for simple surfaces. Sharafutdinov's proof follows the outline of the proof by Pestov and Uhlmann [15] of the non-linear boundary rigidity problem for simple surfaces and it is certainly more involved than our proof of Theorem 1.1, which is independent of the solution to the non-linear problem.

The case of tensor fields of rank four describes the perturbation of travel times of compressional waves propagating in slightly anisotropic elastic media; see Chapter 7 of [17]. For tensor fields of order higher than two, Sharafutdinov gave explicit reconstruction formulas for the solenoidal part [17] in the case when the underlying metric is Euclidean. For results obtained under curvature restrictions see [17, 12] and [19, 2] where non-convex boundaries are also considered.

In the absence of curvature assumptions, the problem for higher order tensors has been considered in [18] on certain surfaces of revolution. For general simple manifolds, it is shown in [21] that I^*I is a pseudodifferential operator of order -1 on a slightly larger simple manifold, and it is elliptic on solenoidal tensor fields. Here I^* denotes the adjoint of I with respect to natural inner products. Thus one can recover the wave front set of a distribution solution of $If = 0$. Using this result, the analysis of [25, 23] and Theorem 1.1 it is straightforward to derive stability estimates.

We give two proofs of Theorem 1.1. These proofs are partially inspired by the proof of the Kodaira vanishing theorem in Complex Geometry [4]. The theorem states that if M is a Kähler manifold, K_M is its canonical line bundle and E is a *positive* holomorphic line bundle, then the cohomology groups $H^q(M, K_M \otimes E)$ vanish for any $q > 0$. The positivity of E means that there exists a Hermitian metric on E such that iF_∇ is a positive differential form, where F_∇ is the curvature of the canonical connection ∇ induced by the Hermitian metric. Via an L^2 energy identity, this positivity of the curvature implies the vanishing of the relevant harmonic forms. Tensoring with E will be translated in our setting as introducing an attenuation given by a suitable connection into the relevant transport equation. This attenuation will play the role of ∇ above and a version of the Pestov identity will play the role of the L^2 energy identity. Of course, this is just an analogy and the technical details are very different in the two settings, but the analogy is powerful enough to provide the key idea for dealing with the transport equation. One actually has a choice of different possible connections as attenuations, and this is what produces the two different proofs. This approach was already employed in [11], but it was surprising to us to discover that it could also be successfully used to solve the tensor tomography problem for simple surfaces. We remark that it is still an open problem to establish Theorem 1.1 when $\dim M \geq 3$ and $m \geq 2$.

After some preliminaries we provide in Section 3 a new point of view on the Pestov identity which makes its derivation quite natural. We also explain here why there is an essential difference between the cases $m = 0, 1$ and $m \geq 2$. In Section 4 we prove Theorem 1.1 choosing as attenuation a primitive of the area form. We give an alternative proof in Section 5 choosing as attenuation the Levi-Civita connection. This alternative proof will also yield a more general result. In order to state it, let I_m denote the ray transform acting on symmetric m -tensors and let $C_\alpha^\infty(\partial_+(SM))$ denote the set of functions $h \in C^\infty(\partial_+(SM))$ such that the function $h_\psi(x, v) = h(\varphi_{-\tau(x, -v)}(x, v))$ is smooth on SM . In natural L^2 inner products, the adjoint of I_0 is the operator

$$I_0^* : C_\alpha^\infty(\partial_+(SM)) \rightarrow C^\infty(M), \quad I_0^*h(x) = \int_{S_x} h_\psi(x, v) dS_x(v).$$

Here $S_x = \{(x, v) \in TM; |v| = 1\}$ and dS_x is the volume form on S_x . For more details see [15], where it is also proved that I_0^* is surjective on any simple manifold.

Theorem 1.2. *Let (M, g) be a compact nontrapping surface with strictly convex smooth boundary. Suppose in addition that I_0 and I_1 are s -injective and that I_0^* is surjective. If f is a smooth symmetric m -tensor field on M , $m \geq 1$, which satisfies $I_m f = 0$, then $f = d_\sigma h$ for some smooth symmetric $(m - 1)$ -tensor field h on M with $h|_{\partial M} = 0$.*

Given this result, it seems natural to conjecture that s -injectivity on tensors should hold on nontrapping surfaces.

Conjecture 1.3. *Let (M, g) be a compact nontrapping surface with strictly convex boundary. If f is a smooth symmetric m -tensor field on M which satisfies $If = 0$, then $f = d_\sigma h$ for some smooth symmetric $(m - 1)$ -tensor field h on M with $h|_{\partial M} = 0$. (If $m = 0$, then $f = 0$.)*

Acknowledgements. M.S. was supported in part by the Academy of Finland, and G.U. was partly supported by NSF and a Rothschild Distinguished Visiting Fellowship at the Isaac Newton Institute. The authors would like to express their gratitude to the Newton Institute and the organizers of the program on Inverse Problems in 2011 where this work was carried out.

2. PRELIMINARIES

Let (M, g) be a compact oriented two dimensional Riemannian manifold with smooth boundary ∂M . As usual SM will denote the unit circle bundle which is a compact 3-manifold with boundary given by $\partial(SM) = \{(x, v) \in SM : x \in \partial M\}$.

Let X denote the vector field associated with the geodesic flow φ_t . Since M is assumed oriented there is a circle action on the fibers of SM with infinitesimal generator V called the *vertical vector field*. It is possible to complete the pair X, V to a global frame of $T(SM)$ by considering the vector field $X_\perp := [X, V]$. There are two additional structure equations given by $X = [V, X_\perp]$ and $[X, X_\perp] = -KV$ where K is the Gaussian curvature of the surface. Using this frame we can define a Riemannian metric on SM by declaring $\{X, X_\perp, V\}$ to be an orthonormal basis and the volume form of this metric will be denoted by $d\Sigma^3$. The fact that $\{X, X_\perp, V\}$ are orthonormal together with the commutator formulas implies that the Lie derivative of $d\Sigma^3$ along the three vector fields vanishes, in other words, the three vector fields preserve the volume form $d\Sigma^3$. We refer the reader to [22] for more details on the assertions in this paragraph.

It will be convenient for later purposes and for the sake of completeness to explicitly write down the three vector fields locally. We can always choose isothermal coordinates (x_1, x_2) so that the metric can be written as $ds^2 = e^{2\lambda}(dx_1^2 + dx_2^2)$ where λ is a smooth real-valued function of $x = (x_1, x_2)$. This gives coordinates (x_1, x_2, θ) on SM where θ is the angle between a unit vector v and $\partial/\partial x_1$. In these coordinates we may write $V = \partial/\partial\theta$ and

$$X = e^{-\lambda} \left(\cos \theta \frac{\partial}{\partial x_1} + \sin \theta \frac{\partial}{\partial x_2} + \left(-\frac{\partial \lambda}{\partial x_1} \sin \theta + \frac{\partial \lambda}{\partial x_2} \cos \theta \right) \frac{\partial}{\partial \theta} \right),$$

$$X_\perp = -e^{-\lambda} \left(-\sin \theta \frac{\partial}{\partial x_1} + \cos \theta \frac{\partial}{\partial x_2} - \left(\frac{\partial \lambda}{\partial x_1} \cos \theta + \frac{\partial \lambda}{\partial x_2} \sin \theta \right) \frac{\partial}{\partial \theta} \right).$$

Given functions $u, v : SM \rightarrow \mathbb{C}$ we consider the inner product

$$(u, v) = \int_{SM} u \bar{v} d\Sigma^3.$$

Since X, X_\perp, V are volume preserving we have $(Vu, v) = -(u, Vv)$ for $u, v \in C^\infty(SM)$, and if additionally $u|_{\partial(SM)} = 0$ or $v|_{\partial(SM)} = 0$ then also $(Xu, v) = -(u, Xv)$ and $(X_\perp u, v) = -(u, X_\perp v)$.

The space $L^2(SM)$ decomposes orthogonally as a direct sum

$$L^2(SM) = \bigoplus_{k \in \mathbb{Z}} H_k$$

where H_k is the eigenspace of $-iV$ corresponding to the eigenvalue k . A function $u \in L^2(SM)$ has a Fourier series expansion

$$u = \sum_{k=-\infty}^{\infty} u_k,$$

where $u_k \in H_k$. Also $\|u\|^2 = \sum \|u_k\|^2$, where $\|u\|^2 = (u, u)^{1/2}$. The even and odd parts of u with respect to velocity are given by

$$u_+ := \sum_{k \text{ even}} u_k, \quad u_- := \sum_{k \text{ odd}} u_k.$$

In the (x, θ) -coordinates previously introduced we may write

$$u_k(x, \theta) = \left(\frac{1}{2\pi} \int_0^{2\pi} u(x, t) e^{-ikt} dt \right) e^{ik\theta} = \tilde{u}_k(x) e^{ik\theta}.$$

Observe that for $k \geq 0$, u_k may be identified with a section of the k -th tensor power of the canonical line bundle; the identification takes u_k into $\tilde{u}_k e^{k\lambda} (dz)^k$ where $z = x_1 + ix_2$.

Definition 2.1. A function $u : SM \rightarrow \mathbb{C}$ is said to be holomorphic if $u_k = 0$ for all $k < 0$. Similarly, u is said to be antiholomorphic if $u_k = 0$ for all $k > 0$.

We will also employ the fiberwise Hilbert transform $H : C^\infty(SM) \rightarrow C^\infty(SM)$, defined in terms of Fourier coefficients as

$$Hu_k := -i \operatorname{sgn}(k) u_k.$$

Here $\operatorname{sgn}(k)$ is the sign of k , with the convention $\operatorname{sgn}(0) = 0$. Thus, u is holomorphic iff $(\operatorname{Id} - iH)u = u_0$ and antiholomorphic iff $(\operatorname{Id} + iH)u = u_0$. The commutator of the Hilbert transform and the geodesic vector field was computed in [15]:

Proposition 2.2. If $u \in C^\infty(SM)$, then $[H, X]u = X_\perp u_0 + (X_\perp u)_0$.

The next proposition concerns holomorphic solutions of transport equations. In the case of simple surfaces this was proved in [16, Proposition 5.1], however here we note that it is true whenever the ray transform on 0-tensors and 1-tensors is s -injective.

Proposition 2.3. Let (M, g) be a compact nontrapping surface with strictly convex boundary, and assume that I_0 and I_1 are s -injective. If f is a smooth holomorphic (antiholomorphic) function on SM and if $u \in C^\infty(SM)$ satisfies

$$Xu = -f \text{ in } SM, \quad u|_{\partial(SM)} = 0,$$

then u is holomorphic (antiholomorphic) and $u_0 = 0$.

Proof. Suppose that f is holomorphic (the antiholomorphic case is analogous). Using Proposition 2.2, the function $w = (\text{Id} - iH)u \in C^\infty(SM)$ satisfies

$$Xw = (\text{Id} - iH)Xu + i[H, X]u = -(\text{Id} - iH)f + iX_\perp u_0 + i(X_\perp u)_0.$$

But $(\text{Id} - iH)f = f_0$, so splitting in even and odd parts gives

$$Xw_- = -f_0 + i(X_\perp u)_0, \quad Xw_+ = iX_\perp u_0.$$

Since $u|_{\partial(SM)} = 0$, we have $w_\pm|_{\partial(SM)} = 0$. Using this boundary condition we have $IXw_\pm = 0$, so the ray transforms of the 0-tensor $f_0 - i(X_\perp u)_0$ and the 1-tensor $-iX_\perp u_0$ vanish. By our injectivity assumption, $f_0 - i(X_\perp u)_0 = 0$ and $-iX_\perp u_0 = 0$ (note that $-iX_\perp u_0$ corresponds to a solenoidal 1-form). In particular

$$Xw_\pm = 0.$$

By the boundary condition for w_\pm , we obtain $w_\pm = 0$. This implies $(\text{Id} - iH)u = 0$, so u is holomorphic and $u_0 = 0$. \square

Without a doubt, the result that makes everything possible is the existence of holomorphic integrating factors for the transport equation. This is true on nontrapping surfaces whenever I_0^* is surjective. The special case of simple surfaces was covered in [11, Theorem 4.1], following a similar result for 0-forms in [16].

Proposition 2.4. *Let (M, g) be a compact nontrapping surface with strictly convex boundary, and assume that $I_0^* : C_\alpha^\infty(\partial_+(SM)) \rightarrow C^\infty(M)$ is surjective. If $A_j dx^j$ is a smooth 1-form on M and $A(x, v) = A_j v^j$, then there exist a holomorphic $w \in C^\infty(SM)$ and an antiholomorphic $\tilde{w} \in C^\infty(SM)$ such that $Xw = X\tilde{w} = -A$.*

Proof. We only do the holomorphic case (the other case is analogous), and prove that there is a holomorphic $w \in C^\infty(SM)$ with $Xw = -A$. First note that the nontrapping condition implies that (M, g) is simply connected. This follows from a well known fact in Riemannian geometry: a compact connected and non-contractible Riemannian manifold with strictly convex boundary must have a closed geodesic in its interior [26]. In our case this violates the assumption of being nontrapping unless M is contractible and hence simply connected.

Since M is simply connected, there is a Hodge decomposition $A_j dx^j = da + \star db$ for some $a, b \in C^\infty(M)$ (\star is the Hodge star operator) and we have $A = Xa + X_\perp b$. Replacing w by $w - a$, it is enough to consider the case where $A = X_\perp b$.

Let us try a solution of the form $w = (\text{Id} + iH)\hat{w}$ where $\hat{w} \in C^\infty(SM)$ is even. By Proposition 2.2,

$$Xw = (\text{Id} + iH)X\hat{w} - i[H, X]\hat{w} = (\text{Id} + iH)X\hat{w} - iX_\perp \hat{w}_0.$$

Now it is sufficient to find \hat{w} even with $X\hat{w} = 0$ and $\hat{w}_0 = -ib$. Using the surjectivity of I_0^* , there is some $h \in C_\alpha^\infty(\partial_+(SM))$ with $I_0^* h = -2\pi ib$. But if $w' \in C^\infty(SM)$ is the function with $Xw' = 0$ in SM and $w'|_{\partial_+(SM)} = h$, we have $(w')_0 = \frac{1}{2\pi} I_0^* h = -ib$. It is enough to take $\hat{w} = w'_+$. \square

We will also make use of a regularity result from [11, Proposition 5.2].

Proposition 2.5. *Let (M, g) be a compact nontrapping manifold with strictly convex boundary, let $f \in C^\infty(SM)$, and define*

$$u(x, v) := \int_0^{\tau(x, v)} f(\varphi_t(x, v)) dt, \quad (x, v) \in SM.$$

If $I f = 0$, then $u \in C^\infty(SM)$.

We conclude this section by explaining the identification between real-valued symmetric m -tensor fields and certain smooth functions on SM . Given such a tensor $f = f_{i_1 \dots i_m} dx^{i_1} \otimes \dots \otimes dx^{i_m}$ we consider the corresponding function on SM (henceforth referred to as the restriction) defined by

$$f(x, v) = f_{i_1 \dots i_m} v^{i_1} \dots v^{i_m}.$$

Then clearly

$$f = \sum_{k=-m}^m f_k$$

where $\bar{f}_k = f_{-k}$. Moreover if m is even (resp. odd) all the odd (resp. even) Fourier coefficients vanish.

Conversely suppose that we are given a smooth real-valued function $f \in C^\infty(SM)$ such that $f_k = 0$ for $|k| \geq m+1$. Suppose in addition that if m is even (resp. odd) then $f = f_+$ (resp. $f = f_-$). Since f is real-valued $\bar{f}_k = f_{-k}$. For each $k \geq 1$, the function $f_{-k} + f_k$ gives rise to a unique real-valued symmetric k -tensor F_k whose restriction to SM is precisely $f_{-k} + f_k$. This can be seen as follows: recall that a smooth element f_k can be identified with a section of $T^*M^{\otimes k}$ hence, its real part defines a symmetric k -tensor. (For $k = 0$, $\bar{f}_0 = f_0$ is obviously a real-valued 0-tensor.) More explicitly, in the coordinates (x, θ) , given $f_k = \tilde{f}_k e^{ik\theta}$ we define

$$F_k := 2\text{Re}(\tilde{f}_k e^{k\lambda}(dz)^k).$$

It is straightforward to check that these local expressions glue together to give a real-valued symmetric k -tensor whose restriction to SM is $F_k(x, v) = f_{-k}(x, v) + f_k(x, v)$.

By tensoring with the metric tensor g and symmetrizing it is possible to raise the degree of a symmetric tensor by two. Hence if σ denotes symmetrization, $\alpha F_k := \sigma(F_k \otimes g)$ will be a symmetric tensor of degree $k+2$ whose restriction to SM is again $f_k + f_{-k}$ since g restricts as the constant function 1 to SM . Now consider the symmetric m -tensor

$$F := \sum_{i=0}^{\lfloor m/2 \rfloor} \alpha^i F_{m-2i}.$$

It is easy to check that the restriction of F to SM is precisely f .

3. PESTOV IDENTITY

In this section we consider the Pestov identity, which is the basic energy identity that has been used since the work of Mukhometov [9] in most injectivity proofs of ray transforms in the absence of real-analyticity or special symmetries. The Pestov

identity often appears in a somewhat ad hoc way, but here we give a new point of view which makes its derivation more transparent.

The easiest way to motivate the Pestov identity is to consider the injectivity of the ray transform on functions. The first step is to recast the injectivity problem as a uniqueness question for the partial differential operator P on SM where

$$P := VX.$$

This involves a standard reduction to the transport equation.

Proposition 3.1. *Let (M, g) be a compact oriented nontrapping surface with strictly convex smooth boundary. The following statements are equivalent.*

- (a) *The ray transform $I : C^\infty(M) \rightarrow C(\partial_+(SM))$ is injective.*
- (b) *Any smooth solution of $Pu = 0$ in SM with $u|_{\partial(SM)} = 0$ is identically zero.*

Proof. Assume that the ray transform is injective, and let $u \in C^\infty(SM)$ solve $Pu = 0$ in SM with $u|_{\partial(SM)} = 0$. This implies that $Xu = -f$ in SM for some smooth f only depending on x , and we have $0 = u|_{\partial_+(SM)} = If$. Since I is injective one has $f = 0$ and thus $Xu = 0$, which implies $u = 0$ by the boundary condition.

Conversely, assume that the only smooth solution of $Pu = 0$ in SM which vanishes on $\partial(SM)$ is zero. Let $f \in C^\infty(M)$ be a function with $If = 0$, and define the function

$$u(x, v) := \int_0^{\tau(x, v)} f(\gamma(t, x, v)) dt, \quad (x, v) \in SM.$$

This function satisfies the transport equation $Xu = -f$ in SM and $u|_{\partial(SM)} = 0$ since $If = 0$, and also $u \in C^\infty(SM)$ by Proposition 2.5. Since f only depends on x we have $Vf = 0$, and consequently $Pu = 0$ in SM and $u|_{\partial(SM)} = 0$. It follows that $u = 0$ and also $f = -Xu = 0$. \square

We now focus on proving a uniqueness statement for solutions of $Pu = 0$ in SM . For this it is convenient to express P in terms of its self-adjoint and skew-adjoint parts in the $L^2(SM)$ inner product as

$$P = A + iB, \quad A := \frac{P + P^*}{2}, \quad B := \frac{P - P^*}{2i}.$$

Here the formal adjoint P^* of P is given by

$$P^* := XV.$$

In fact, if $u \in C^\infty(SM)$ with $u|_{\partial(SM)} = 0$, then

$$\begin{aligned} (1) \quad \|Pu\|^2 &= ((A + iB)u, (A + iB)u) = \|Au\|^2 + \|Bu\|^2 + i(Bu, Au) - i(Au, Bu) \\ &= \|Au\|^2 + \|Bu\|^2 + (i[A, B]u, u). \end{aligned}$$

This computation suggests to study the commutator $i[A, B]$. We note that the argument just presented is typical in the proof of L^2 Carleman estimates [6].

By the definition of A and B it easily follows that $i[A, B] = \frac{1}{2}[P^*, P]$. By the commutation formulas for X, X_\perp and V , this commutator may be expressed as

$$\begin{aligned} [P^*, P] &= XVVX - VXXV = VXVX + X_\perp VX - VXVX - VXX_\perp \\ &= V[X_\perp, X] - X^2 = -X^2 + VKV. \end{aligned}$$

Consequently

$$([P^*, P]u, u) = \|Xu\|^2 - (KVu, Vu).$$

If the curvature K is nonpositive, then $[P^*, P]$ is positive semidefinite. More generally, one can try to use the other positive terms in (1). Note that

$$\|Au\|^2 + \|Bu\|^2 = \frac{1}{2}(\|Pu\|^2 + \|P^*u\|^2).$$

The identity (1) may then be expressed as

$$\|Pu\|^2 = \|P^*u\|^2 + ([P^*, P]u, u).$$

Moving the term $\|Pu\|^2$ to the other side, we have proved the version of the Pestov identity which is most suited for our purposes. The Pestov identity has a long history (see for instance [17]) and also this version is known, but the main point in our proof was that the identity boils down to a standard L^2 estimate based on separating the self-adjoint and skew-adjoint parts of P and on computing one commutator, $[P^*, P]$.

Proposition 3.2. *If (M, g) is a compact oriented surface with smooth boundary, then*

$$\|XVu\|^2 - (KVu, Vu) + \|Xu\|^2 - \|VXu\|^2 = 0$$

for any $u \in C^\infty(SM)$ with $u|_{\partial(SM)} = 0$.

It is known (cf. proof of [3, Proposition 7.2]) that on a simple surface, one has

$$\|XVu\|^2 - (KVu, Vu) \geq 0, \quad u \in C^\infty(SM), \quad u|_{\partial(SM)} = 0.$$

Also, if $Xu = -f$ where $f = f_0 + f_1 + f_{-1}$ is the sum of a 0-form and 1-form (identified with the corresponding functions on SM), we have

$$\|Xu\|^2 - \|VXu\|^2 = \|f_0\|^2 \geq 0.$$

These two facts together with the Pestov identity give the standard proof of s -injectivity of the ray transform for 0-forms and 1-forms on simple surfaces. It is easy to see where this proof breaks down if $m \geq 2$: the Fourier expansion $f = \sum_{k=-m}^m f_k$ implies

$$\|Xu\|^2 - \|VXu\|^2 = \|f_0\|^2 - \sum_{2 \leq |k| \leq m} (k^2 - 1) \|f_k\|^2.$$

This term may be negative, and the Pestov identity may not give useful information unless there is some extra positivity like a curvature bound.

Finally, we consider the Pestov identity in the presence of attenuation given by $A(x, v) = A_j(x)v^j$ where $A_j dx^j$ is a purely imaginary 1-form on M . We write A both for the 1-form and the function on SM . The geometric interpretation is that $d + A$ is a unitary connection on the trivial bundle $M \times \mathbb{C}$ (it is customary to also call A a connection), and its curvature is the 2-form

$$F_A := dA + A \wedge A = dA.$$

Then $\star F_A$ is a function on M where \star is the Hodge star. We consider the operator

$$P := V(X + A).$$

Since $\bar{A} = -A$, the formal adjoint of P in the $L^2(SM)$ inner product is

$$P^* = (X + A)V.$$

The same argument leading to Proposition 3.2, based on computing the commutator $[P^*, P]$, gives the following Pestov identity proved also in [11, Lemma 6.1].

Proposition 3.3. *If (M, g) is a compact oriented surface with smooth boundary and if A is a purely imaginary 1-form on M , then*

$$\|(X + A)Vu\|^2 - (KVu, Vu) + \|(X + A)u\|^2 - \|V(X + A)u\|^2 + (\star F_A Vu, u) = 0$$

for any $u \in C^\infty(SM)$ with $u|_{\partial(SM)} = 0$.

Using the Fourier expansion of u , the last term in the identity is given by

$$\sum_{k=-\infty}^{\infty} ik(\star F_A u_k, u_k)$$

This shows that if u is holomorphic and $i \star F_A > 0$, or if u is antiholomorphic and $i \star F_A < 0$, one gains an additional positive term in the Pestov identity. This will be crucial in absorbing negative contributions from the term $\|(X + A)u\|^2 - \|V(X + A)u\|^2$ when proving s -injectivity on tensor fields.

4. FIRST PROOF

The proof of s -injectivity of the ray transform reduces to proving the next result. We say that $f \in C^\infty(SM)$ has degree m if $f_k = 0$ for $|k| \geq m + 1$.

Proposition 4.1. *Let (M, g) be a simple surface, and assume that $u \in C^\infty(SM)$ satisfies $Xu = -f$ in SM with $u|_{\partial(SM)} = 0$. If $f \in C^\infty(SM)$ has degree $m \geq 1$, then u has degree $m - 1$. If f has degree 0, then $u = 0$.*

Proof of Theorem 1.1. Let f be a symmetric m -tensor field on M , identified with its restriction to SM , and suppose that $If = 0$. We write

$$u(x, v) := \int_0^{\tau(x, v)} f(\varphi_t(x, v)) dt, \quad (x, v) \in SM.$$

Then $u|_{\partial(SM)} = 0$, and also $u \in C^\infty(SM)$ by Proposition 2.5.

Now f has degree m , and u satisfies $Xu = -f$ in SM with $u|_{\partial(SM)} = 0$. Proposition 4.1 implies that u has degree $m - 1$ (and $u = 0$ if $m = 0$). We let $h := -u$. As we explained in Section 2, h gives rise to a symmetric $(m - 1)$ -tensor still denoted by h . Since $X(h) = f$, this implies that $d_\sigma h$ and f agree when restricted to SM and since they are both symmetric tensors of the same degree it follows that $d_\sigma h = f$. This proves the theorem. \square

Proposition 4.1 is in turn an immediate consequence of the next two results. Note that these are generalizations of Proposition 2.3.

Proposition 4.2. *Let (M, g) be a simple surface, and assume that $u \in C^\infty(SM)$ satisfies $Xu = -f$ in SM with $u|_{\partial(SM)} = 0$. If $m \geq 0$ and if $f \in C^\infty(SM)$ is such that $f_k = 0$ for $k \leq -m - 1$, then $u_k = 0$ for $k \leq -m$.*

Proposition 4.3. *Let (M, g) be a simple surface, and assume that $u \in C^\infty(SM)$ satisfies $Xu = -f$ in SM with $u|_{\partial(SM)} = 0$. If $m \geq 0$ and if $f \in C^\infty(SM)$ is such that $f_k = 0$ for $k \geq m + 1$, then $u_k = 0$ for $k \geq m$.*

We will only prove Proposition 4.2, the proof of the other result being completely analogous. As a preparation, following Guillemin and Kazhdan in [5] we introduce the first order operators $\eta_{\pm} : C^{\infty}(SM) \rightarrow C^{\infty}(SM)$ given by

$$\eta_+ := (X + iX_{\perp})/2, \quad \eta_- := (X - iX_{\perp})/2.$$

These operators are elliptic [5], and clearly $X = \eta_+ + \eta_-$. Let $\Omega_k := C^{\infty}(SM) \cap H_k$. The commutation relations $[-iV, \eta_+] = \eta_+$ and $[-iV, \eta_-] = -\eta_-$ imply that $\eta_{\pm} : \Omega_k \rightarrow \Omega_{k\pm 1}$. If $A(x, v) = A_j(x)v^j$ where A is a purely imaginary 1-form on M , we also split $A = A_+ + A_-$ where $A_{\pm} \in \Omega_{\pm 1}$ and write

$$\mu_+ := \eta_+ + A_+, \quad \mu_- := \eta_- + A_-.$$

Note that $X + A = \mu_+ + \mu_-$. One has $\mu_{\pm} : \Omega_k \rightarrow \Omega_{k\pm 1}$, and

$$(2) \quad (\mu_{\pm}u, v) = -(u, \mu_{\mp}v), \quad u, v \in C^{\infty}(SM) \text{ with } u|_{\partial(SM)} = 0 \text{ or } v|_{\partial(SM)} = 0.$$

The following commutator formula was proved in [10, Lemma 4.3]:

$$(3) \quad [\mu_+, \mu_-]u = \frac{i}{2}(KVu + (\star F_A)u).$$

Proof of Proposition 4.2. Assume that $u \in C^{\infty}(SM)$ satisfies $Xu = -f$ in SM and $u|_{\partial(SM)} = 0$, where $f_k = 0$ for $k \leq -m - 1$. Splitting in even and odd parts with respect to velocity gives that $Xu_{\mp} = -f_{\pm}$ in SM , $u_{\mp}|_{\partial(SM)} = 0$. Therefore it is enough to prove the result in two cases: when f is even (so m is even) and u is odd, and when f is odd (so m is odd) and u is even. The proofs for these cases are very similar, and we will only consider the first case.

Assume that f is even, m is even, and u is odd. Let ω_g be the area form of (M, g) and choose a real valued 1-form φ with $d\varphi = \omega_g$. Consider the unitary connection

$$A(x, v) := is\varphi_j(x)v^j$$

where $s > 0$ is a fixed number to be chosen later. Then $i\star F_A = -s$. By Proposition 2.4, there exists a holomorphic $w \in C^{\infty}(SM)$ satisfying $Xw = -i\varphi$. We may assume that w is even. The functions $\tilde{u} := e^{sw}u$ and $\tilde{f} := e^{sw}f$ then satisfy

$$(X + A)\tilde{u} = -\tilde{f} \text{ in } SM, \quad \tilde{u}|_{\partial(SM)} = 0.$$

Using that e^{sw} is holomorphic, we have $\tilde{f}_k = 0$ for $k \leq -m - 1$. Also, since e^{sw} is even, \tilde{f} is even and \tilde{u} is odd. We now define

$$v := \sum_{k=-\infty}^{-m-1} \tilde{u}_k.$$

Then $v \in C^{\infty}(SM)$, $v|_{\partial(SM)} = 0$, and v is odd. Also, $((X + A)v)_k = \mu_+v_{k-1} + \mu_-v_{k+1}$. If $k \leq -m - 2$ one has $((X + A)v)_k = ((X + A)\tilde{u})_k = 0$, and if $k \geq -m + 1$ then $((X + A)v)_k = 0$ since $v_j = 0$ for $j \geq -m$. Also $((X + A)v)_{-m-1} = 0$ because v is odd. Therefore the only nonzero Fourier coefficient is $((X + A)v)_{-m}$, and

$$(X + A)v = \mu_+v_{-m-1} \text{ in } SM, \quad v|_{\partial(SM)} = 0.$$

We apply the Pestov identity in Proposition 3.3 with attenuation A to v , so that

$$\|(X + A)Vv\|^2 - (KVv, Vv) + (\star F_A Vv, v) + \|(X + A)v\|^2 - \|V(X + A)v\|^2 = 0.$$

It was proved in [11, Lemma 6.5] that if (M, g) is simple and $v|_{\partial(SM)} = 0$, then

$$(4) \quad \|(X + A)Vv\|^2 - (KVv, Vv) \geq 0.$$

We also have

$$(5) \quad (\star F_A Vv, v) = - \sum_{k=-\infty}^{-m-1} i|k|(\star F_A v_k, v_k) = s \sum_{k=-\infty}^{-m-1} |k| \|v_k\|^2.$$

For the remaining two terms, we compute

$$\|(X + A)v\|^2 - \|V(X + A)v\|^2 = \|\mu_+ v_{-m-1}\|^2 - m^2 \|\mu_+ v_{-m-1}\|^2.$$

If $m = 0$, then this expression is nonnegative and we obtain from the energy identity that $v = 0$. Assume from now on that $m \geq 2$. Using (2), (3), and the fact that $v_k|_{\partial(SM)} = 0$ for all k , we have

$$\begin{aligned} \|\mu_+ v_k\|^2 &= \|\mu_- v_k\|^2 + \frac{i}{2}(KVv_k + (\star F_A)v_k, v_k) \\ &= \|\mu_- v_k\|^2 - \frac{s}{2}\|v_k\|^2 - \frac{k}{2}(Kv_k, v_k). \end{aligned}$$

If $k \leq -m - 1$ we also have

$$\mu_+ v_{k-1} + \mu_- v_{k+1} = ((X + A)v)_k = 0.$$

We thus obtain

$$\begin{aligned} \|(X + A)v\|^2 - \|V(X + A)v\|^2 &= -(m^2 - 1)\|\mu_+ v_{-m-1}\|^2 \\ &= -(m^2 - 1) \left[\|\mu_- v_{-m-1}\|^2 - \frac{s}{2}\|v_{-m-1}\|^2 + \frac{m+1}{2}(Kv_{-m-1}, v_{-m-1}) \right] \\ &= -(m^2 - 1) \left[\|\mu_+ v_{-m-3}\|^2 - \frac{s}{2}\|v_{-m-1}\|^2 + \frac{m+1}{2}(Kv_{-m-1}, v_{-m-1}) \right] \\ &= -(m^2 - 1) \left[\|\mu_- v_{-m-3}\|^2 - \frac{s}{2}(\|v_{-m-1}\|^2 + \|v_{-m-3}\|^2) \right. \\ &\quad \left. + \frac{m+1}{2}(Kv_{-m-1}, v_{-m-1}) + \frac{m+3}{2}(Kv_{-m-3}, v_{-m-3}) \right]. \end{aligned}$$

Continuing this process, and noting that $\mu_- v_k \rightarrow 0$ in $L^2(SM)$ as $k \rightarrow -\infty$ (this is valid since $\|\mu_- v\|^2 = \sum_{k=-\infty}^{-m-2} \|(\mu_- v)_k\|^2 < \infty$, where $(\mu_- v)_k = \mu_- v_{k+1}$ and $\mu_- v \in L^2(SM)$), we obtain

$$(6) \quad \|(X + A)v\|^2 - \|V(X + A)v\|^2 = \frac{m^2 - 1}{2}s \sum \|v_k\|^2 - \frac{m^2 - 1}{2} \sum |k|(Kv_k, v_k).$$

Collecting (4)–(6) and using them in the Pestov identity implies that

$$0 \geq \frac{m^2 - 1}{2}s \sum \|v_k\|^2 + \left(s - \frac{m^2 - 1}{2} \sup_M K \right) \sum |k| \|v_k\|^2.$$

If we choose $s > \frac{m^2 - 1}{2} \sup_M K$, then both terms above are nonnegative and therefore have to be zero. It follows that $v = 0$, so $\tilde{u}_k = 0$ for $k \leq -m - 1$ and also $u_k = 0$ for $k \leq -m - 1$ since $u = e^{-sw}\tilde{u}$ where e^{-sw} is holomorphic. \square

5. SECOND PROOF

In this section we provide a short alternative proof to the key proposition in the injectivity result.

Second proof of Proposition 4.2. Suppose that u is a smooth solution of $Xu = -f$ in SM where $f_k = 0$ for $k \leq -m-1$ and $u|_{\partial(SM)} = 0$. We choose a nonvanishing function $h \in \Omega_m$ and define

$$A := -h^{-1}Xh.$$

Observe that since $h \in \Omega_m$, then $h^{-1} \in \Omega_{-m}$. Also $Xh = \eta_+h + \eta_-h \in \Omega_{m+1} \oplus \Omega_{m-1}$ which implies that $A \in \Omega_1 \oplus \Omega_{-1}$. It follows that A is the restriction to SM of a (complex-valued) 1-form on M .

The function hu solves the problem

$$(X + A)(hu) = -hf \text{ in } SM, \quad hu|_{\partial(SM)} = 0.$$

Note that hf is a holomorphic function. Next we employ a holomorphic integrating factor: by Proposition 2.4 there exists a holomorphic $w \in C^\infty(SM)$ with $Xw = A$. The function $e^w hu$ then satisfies

$$X(e^w hu) = -e^w hf \text{ in } SM, \quad e^w hu|_{\partial(SM)} = 0.$$

The right hand side $e^w hf$ is holomorphic. Now Proposition 2.3 implies that the solution $e^w hu$ is also holomorphic and $(e^w hu)_0 = 0$. Looking at Fourier coefficients shows that $(hu)_k = 0$ for $k \leq 0$, and therefore $u_k = 0$ for $k \leq -m$ as required. \square

Finally, let us explain the choice of h and A in the proof in more detail. Since M is a disk we can consider global isothermal coordinates (x, y) on M such that the metric can be written as $ds^2 = e^{2\lambda}(dx^2 + dy^2)$ where λ is a smooth real-valued function of (x, y) . As explained in Section 2 this gives coordinates (x, y, θ) on SM where θ is the angle between a unit vector v and $\partial/\partial x$. Then Ω_m consists of all functions $a(x, y)e^{im\theta}$ where $a \in C^\infty(M)$. We choose the specific nonvanishing function

$$h(x, y, \theta) := e^{im\theta}.$$

Recall that in the (x, y, θ) coordinates the geodesic vector field X is given by:

$$X = e^{-\lambda} \left(\cos \theta \frac{\partial}{\partial x} + \sin \theta \frac{\partial}{\partial y} + \left(-\frac{\partial \lambda}{\partial x} \sin \theta + \frac{\partial \lambda}{\partial y} \cos \theta \right) \frac{\partial}{\partial \theta} \right).$$

The connection $A = -Xh/h$ has the form

$$A = ime^{-\lambda} \left(-\frac{\partial \lambda}{\partial y} \cos \theta + \frac{\partial \lambda}{\partial x} \sin \theta \right) = im \left(-\frac{\partial \lambda}{\partial y} dx + \frac{\partial \lambda}{\partial x} dy \right).$$

Here as usual we identify A with $A(x, v)$ where $(x, v) \in SM$. This shows that the connection A is essentially the Levi-Civita connection of the metric g on the tensor power bundle $TM^{\otimes m}$, and since $(X + A)h = 0$ we have that h corresponds to a section of the pull-back bundle $\pi^*(TM^{\otimes m})$ whose covariant derivative along the geodesic vector field vanishes (here $\pi : SM \rightarrow M$ is the standard projection).

Proof of Theorem 1.2. An inspection of the proof above reveals that it also proves Theorem 1.2 since we have assumed precisely the hypotheses that allow the use of

Propositions 2.3 and 2.4. Also observe that, as pointed out in the proof of Proposition 2.4, a compact nontrapping surface with strictly convex boundary must be simply connected and hence a disk. □

REFERENCES

- [1] Yu. Anikonov, V. Romanov, *On uniqueness of determination of a form of first degree by its integrals along geodesics*, J. Inverse Ill-Posed Probl. **5** (1997), 467–480.
- [2] N.S. Dairbekov, *Integral geometry problem for nontrapping manifolds*, Inverse Problems **22** (2006), 431–445.
- [3] D. Dos Santos Ferreira, C.E. Kenig, M. Salo, G. Uhlmann, *Limiting Carleman weights and anisotropic inverse problems*, Invent. Math. **178** (2009), 119–171.
- [4] P. Griffiths, J. Harris, *Principles of Algebraic Geometry*, Reprint of the 1978 original Wiley Classics Library. John Wiley & Sons, Inc., New York, 1994.
- [5] V. Guillemin, D. Kazhdan, *Some inverse spectral results for negatively curved 2-manifolds*, Topology **19** (1980), 301–312.
- [6] L. Hörmander, *The analysis of linear partial differential operators. IV. Fourier integral operators*. Reprint of the 1994 edition. Classics in Mathematics. Springer-Verlag, Berlin, 2009.
- [7] S. Ivanov, *Volume comparison via boundary distances*. Proceedings of the International Congress of Mathematicians, vol. II, 769–784, New Delhi, 2010.
- [8] R. Michel, *Sur la rigidité imposée par la longueur des géodésiques*, Invent. Math. **65** (1981), 71–83.
- [9] R.G. Mukhometov, *The reconstruction problem of a two-dimensional Riemannian metric, and integral geometry (Russian)*, Dokl. Akad. Nauk SSSR **232** (1977), no.1, 32–35.
- [10] G.P. Paternain, *Transparent connections over negatively curved surfaces*, J. Mod. Dyn. **3** (2009), 311–333.
- [11] G.P. Paternain, M. Salo, G. Uhlmann, *The attenuated ray transform for connections and Higgs fields*, arXiv:1108.1118, to appear in Geom. Funct. Anal.
- [12] L. Pestov, *Well-Posedness Questions of the Ray Tomography Problems*, (Russian), Siberian Science Press, Novosibirsk, 2003.
- [13] L. Pestov, V.A. Sharafutdinov, *Integral geometry of tensor fields on a manifold of negative curvature*, Siberian Math. J. **29** (1988), 427–441.
- [14] L. Pestov, G. Uhlmann, *On characterization of the range and inversion formulas for the geodesic X-ray transform*, Int. Math. Res. Not. (2004), 4331–4347.
- [15] L. Pestov, G. Uhlmann, *Two dimensional compact simple Riemannian manifolds are boundary distance rigid*, Ann. of Math. **161** (2005), 1089–1106.
- [16] M. Salo, G. Uhlmann, *The attenuated ray transform on simple surfaces*, J. Diff. Geom. **88** (2011), 161–187.
- [17] V.A. Sharafutdinov, *Integral geometry of tensor fields*, Inverse and Ill-Posed Problems Series. VSP, Utrecht, 1994.
- [18] V.A. Sharafutdinov, *Integral geometry of a tensor field on a surface of revolution*, Siberian Math. J. **38** (1997), 603–620.
- [19] V.A. Sharafutdinov, *A problem in integral geometry in a nonconvex domain*, Siberian Math. J. **43** (2002), 1159–1168.
- [20] V.A. Sharafutdinov, *Variations of Dirichlet-to-Neumann map and deformation boundary rigidity of simple 2-manifolds*, J. Geom. Anal. **17** (2007), 147–187.
- [21] V.A. Sharafutdinov, M. Skokan, G. Uhlmann, *Regularity of ghosts in tensor tomography*, J. Geom. Anal. **15** (2005), 517–560.
- [22] I.M. Singer, J.A. Thorpe, *Lecture notes on elementary topology and geometry*. Reprint of the 1967 edition. Undergraduate Texts in Mathematics. Springer-Verlag, New York-Heidelberg, 1976.
- [23] P. Stefanov, G. Uhlmann, *Stability estimates for the X-ray transform of tensor fields and boundary rigidity*, Duke Math. J. **123** (2004), 445–467.

- [24] P. Stefanov, G. Uhlmann, *Boundary and lens rigidity, tensor tomography and analytic microlocal analysis*, in Algebraic Analysis of Differential Equations, Festschrift in Honor of Takahiro Kawai, edited by T. Aoki, H. Majima, Y. Katei and N. Tose, pp. 275–293 (2008).
- [25] P. Stefanov, G. Uhlmann, *Linearizing non-linear inverse problems and its applications to inverse backscattering*, J. Funct. Anal. **256** (2009), 2842–2866.
- [26] G. Thorbergsson, *Closed geodesics on non-compact Riemannian manifolds*, Math. Z. **159** (1978), 249–258.

DEPARTMENT OF PURE MATHEMATICS AND MATHEMATICAL STATISTICS, UNIVERSITY OF CAMBRIDGE, CAMBRIDGE CB3 0WB, UK.

E-mail address: `g.p.paternain@dpms.cam.ac.uk`

DEPARTMENT OF MATHEMATICS AND STATISTICS, UNIVERSITY OF JYVÄSKYLÄ.

E-mail address: `mikko.j.salo@jyu.fi`

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WASHINGTON AND UC IRVINE.

E-mail address: `gunther@math.washington.edu`