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Reference

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TENTATIVE MODEL FOR THE INCOMMENSURATE AND FERROELECTRIC PHASES IN Pb₂CoWO₆

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<u>Abstract</u> The sequence of cubic, incommensurate and ferroelectric phases found in Pb_CoWO_ is described in the framework of the Landau theory. The model, which involves two different order-parameters, accounts for the first-order character of the transitions, the reentrant behaviour of the incommensurate phase and its coexistence with the ferroelectric phase.

INTRODUCTION

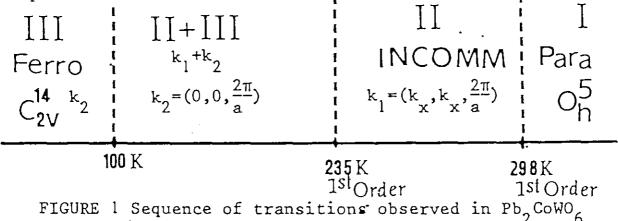
In order to clarify the controversial data previously reported in Pb_CoWO_(PCW) , large single crystals of this composition were synthesized and a comprehensive set of experiments were recently undertaken , which include polarized light microscopy , electric poling , birefringence measurements , and neutron scattering . These experiments confirmed the sequence of phase transitions which is schematized in Fig.1, namely:1) a strongly first-order transition at $T_{1}=298$ K from the high-temperature paraelectric phase, with the elpasolite-like structure(space-group O_{h}^{-1}), to an incommensurate (INC) phase(II) . The temperature dependence of the satellite peak intensities which appear at T_{1} , show a reentrant behaviour, i.e. a sharp increase at 285 K and a saturation around 240 K , followed by a continuous decrease with decreasing temperature. The satellites survive the transition to phase III and can be observed down to about 100 K.

2)At 235 K another first-order transition to a ferroelectric phase (III) takes place, characterized by the onset of a superstructure corresponding to a different wave-vect tor .In this phase one can observe a non-standard dependence on temperature, for the spontaneous polarization and hysteresis loops.

The aim of this paper is to present a tentative model which accounts for the preceding experimental results.A more detailed model will be published elsewhere. 590/[296]

SYMMETRY ANALYSIS

According to the X-ray³ and elastic and inelastic neutron diffraction data obtained for PCW³, at $T_T=298$ K the satellite peaks appear on the S-line of the fcc Brillouin-zone, which correspond to the waye-vector $k_1 = (k_1, k_2\pi/a)$ (see for example the tables by Zak et al'). By reducing the temperature the peaks show characteristic displacements revealing that the wave-vector remains on the Sline with a positioning parameter δ varying from 0.171Å at 285 K to 0.202Å⁴ at 220 K.As the symmetry group of k_1 is C_{2V}^{XY} , the star of k_1 has twelve arms, and consequently the order-parameter associated with the paraelectric I-INC II transition possesses twelve components.



At T_{12} =235 K the new satellite peaks, which reveal the onset of phase III, correspond to the wave-vector $k_2 = (0, 0, 2\pi/a)$ (the λ point of the fcc Brillouin-zone ,i.e.k in the Kovalev notation) and are associated with a doubling of the cell parameters in the x and y directions. Their position remains fixed as the temperature is lowered, by contrast to the satellites related to k with which they coexist down to about 100 K. The star of k_2 has three arms, as the symmetry group of k_2 is D_{4h} . A symmetry analysis of the irreducible representations (IR's) constructed from the ten small representations of k_2 shows that only the six-dimensional IR labelled τ_{10} in Kovalev's tables may lead to ferroelectric low-symmetry phases. Actually the polar character of phase III is attested by the whole set of dielectric measurements . More precisely, the domain pattern below 235 K is only compatible with an orthorhombic point-group C_{2v} , having an _ axis parallel to [110] These data designate unambiguously the space-group $C_{2v}(xy)$, with a four-fold multiplication of the cubic primitive cell, as the symmetry of phase III. The stabilization of this group, following our analysis , requires to take a into account sixth-degree invariants of the six-dimensional orderparameter in the thermodynamic potential associated with $\tau_{10}(k_2)$.

 $\eta_1 \eta_2 = 0 \eta_3 \eta_6 = 0 \eta_4 \eta_5 = 0$

(1)

From the preceding symmetry considerations, we are now able to propose a phenomenological model for the sequence of phases in PCW.

TWO ORDER-PARAMETER MODEL FOR THE PHASE DIAGRAM OF Pb CoWO

In order to work out the qualitative features of the phase diagram of PCW, one can restrict to consider the effective thermodynamic potential associated with the sequence of phases in this compound. It can be written here as:

 $\mathbf{F} = \frac{\alpha}{2} \rho^{2} + \frac{\beta}{4} \rho^{4} + \frac{a}{2} (\eta_{1}^{2} + \eta_{3}^{2}) + \frac{b}{4} \mathbf{1} (\eta_{1}^{4} + \eta_{3}^{4}) + \frac{b}{2} 2 \eta_{1}^{2} \eta_{3}^{2} + \frac{c}{6} (\eta_{1}^{6} + \eta_{3}^{6}) + \delta \rho^{2} (\eta_{1}^{2} + \eta_{3}^{2})$ (2)

where ρ is the order-parameter modulus corresponding to the INC phase, and (n_1, n_2) are the non-zero components of the six-dimensional order-parameter associated with the ferroelectric phase. The δ term represents the lower degree coupling between the two order-parameters. The minimization of F with respect to ρ and (η_1, η_2) leads for suitable values of the coefficients in F (in particular for $\delta > -2/\beta$ b,) to the phase diagram represented in Fig.2. Thus, one can see that for the thermodynamic path indicated in dashed line in Fig.2, one goes from the paraelectric phase $(\rho=0,\eta=0)$ to the INC phase $(\rho=0,\eta=0)$ across a line of first-order transitions, then to a region of coexistence of the two phases $(0 \neq 0, \eta \neq 0)$ across another line of firstorder transitions. The reentrant behaviour observed for the INC phase can be obtained within this picture by assuming that the thermodynamic path corresponds to decreasing values of ρ when approaching the INC-ferroelectric transition line and within the coexistence region. The continuous vanishing of the INC sätellite peaks observed in the vicinity of 100 K may result from the fact that the limit stability region between phases II and III (dotted line in Fig.2) has been reached.

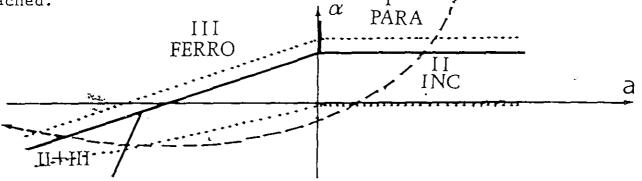


FIGURE 2 Phase diagram corresponding to potential(2) in the (α, a) plane.Full and dotted lines represent first-order and stability lines respectively.The thermodynamic path suggested for PCW is shown by a dashed line.

DIELECTRIC PROPERTIES IN THE FERROELECTRIC PHASE

The dielectric behaviour of PCW below 235 K is obtained, following the standard procedure by considering the coupling between the sixcomponent order-parameter and the spontaneous polarization. In this respect one finds that the lower degree coupling term can be written:

$$\delta_1 P_x (\eta_2 \eta_3 \eta_6 - \eta_1 \eta_4 \eta_5) + \delta_2 P_y (\eta_2 \eta_4 \eta_5 - \eta_1 \eta_3 \eta_6)$$
(3)

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which denotes an unusual faintness index n=3. The standard minimization procedure provides the following expressions for the spontaneous polarization components below T_{++} :

$$P_x \approx \chi_0 \delta_1 \eta_1 \eta_3^2$$
, $P_y \approx \chi_0 \delta_2 \eta_1 \eta_3^2$, $P_z = 0$ (4)

and a dielectric susceptibility of the form:

$$\chi = \chi_{0} \left(1 + \frac{\chi_{0}^{0} \eta^{2}}{\beta - \alpha/\eta^{2} - \chi_{0}^{0} \delta \eta^{2}} \right)$$
(5)

However while the experimental curves obtained for the polarization and dielectric permittivity below 135 K reflect the theoretical temperature dependences given by (4) and (5), between 135 K and 235 K the experimental results are strongly influenced by the INC regime. In particular one has to take into account the domain effects introduced by the coexistence of the INC and ferroelectric phases.

REFERENCES

- 1.V.S.Filip'ev, and E.G.Fesenko, Sov. Phys. Crystallography 9,231(1964)
- 2.V.A.Bokov, S.A.Kizhaev, I.E.Myl'nikova, and A.G.Tutov, Sov.Phys. Solid State 6,2419(1965)
- 3.H.Tamura, Ferroelectrics 21, 449(1978)
- 4.W.Brixel, R.Bouteiller, and H.Schmid, to be published in J.Crystal Growth (1987)
- 5.W.Brixel, M.LWerk, P.Fisher, W.Bührer, J.P.Rivera, P.Tissot, and Schmid, Japonese Jour. Appl. Phys. 24, Suppl. 24-2, 242 (1985)
- 6.W.Brixel, Thesis, University of Geneva(1987)
- 7.W.Bührer, W.Brixel, and H.Schmid, in "Phonons 85", World Scientific Pub.Singapor(1985)
- 8.W.Bührer, H.Schmid, W.Brixel, F. Maaroufi, P. Toledano, to be published
- 9.J.Zak, A. Casher, H. Glück, and Y. Gur, "The Irreducible Representations of Space Groups", Benjamin, New York (1969)
- 10.0.V.Kovalev, "Irreducible Representations of the Space Groups", Gordon and Breach, New York (1965)
- 11.K.Aizu, J.Phys.Soc.Japan 34, 121(1973)