

THE ADAMS-NOVIKOV SPECTRAL SEQUENCE FOR THE SPHERES

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The Adams spectral sequence has been an important tool in research on the stable homotopy of the spheres. In this note we outline new information about a variant of the Adams sequence which was introduced by Novikov [7]. We develop simplified techniques of computation which allow us to discover vanishing lines and periodicity near the edge of the E_2 -term, interesting elements in $E_2^{2,*}$, and a counterexample to one of Novikov's conjectures. In this way we obtain independently the values of many low-dimensional stems up to group extension. The new methods stem from a deeper understanding of the Brown-Peterson cohomology theory, due largely to Quillen [8]; see also [4]. Details will appear elsewhere; or see [11].

When p is odd, the p -primary part of the Novikov sequence behaves nicely in comparison with the ordinary Adams sequence. Computing the E_2 -term seems to be as easy, and the Novikov sequence has many fewer nonzero differentials (in stems ≤ 45 , at least, if $p = 3$), and periodicity near the edge. The case $p = 2$ is sharply different. Computing E_2 is more difficult. There are also hordes of nonzero differentials d_3 , but they form a regular pattern, and no nonzero differentials outside the pattern have been found. Thus the diagram of E_4 ($= E_\infty$ in dimensions ≤ 17) suggests a vanishing line for E_∞ much lower than that of E_2 of the classical Adams spectral sequence [3].

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1. The spectral sequence. The construction of the classical Adams spectral sequence for the spheres [1] works equally well if the spec-

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trum $K(Z_p)$ representing ordinary cohomology is replaced by an arbitrary ring spectrum X . If X satisfies certain conditions, the E_2 -term of the resulting sequence will be isomorphic to

$$\text{Ext}_{A^X}(\Lambda^X, \Lambda^X),$$

where $A^X = X^*(X)$ is the algebra of operations in X -cohomology theory and $\Lambda^X = \pi_*(X)$ is the coefficient ring. Novikov showed [7] that if $X = MU$ (the spectrum representing complex cobordism) this multiplicative spectral sequence converges to the stable homotopy ring π_*^S :

$$E_\infty^{s,t} \cong F^s \pi_{t-s}^S / F^{s+1} \pi_{t-s}^S,$$

where F^* is a filtration of π_*^S . Furthermore, if $X' = BP_p$, the Brown-Peterson spectrum [4] for the prime p , the resulting spectral sequence $\{ {}_p E_r, {}_p d_r \}$ is exactly the p -primary part $\{ E_r \otimes Q_p, d_r \otimes Q_p \}$ of the MU spectral sequence (Q_p is the ring of rational numbers with denominators prime to p .)

Not much is known about the MU spectral sequence, because even limited computations of E_2 have been difficult. This is regrettable, since what is known indicates that the Novikov sequence has certain a priori advantages over the usual one. The nonzero terms are sparse, for example: ${}_p E_2^{s,t} = 0$ if $t \not\equiv 0 \pmod{2(p-1)}$. Furthermore, almost all of the image of the J -homomorphism [2], [9] lies on the line $s=1$, in the following sense. According to Novikov, $E_2^{1,2t} = Z_m(t) \langle \alpha_t \rangle$, a cyclic group with generator α_t , isomorphic to the image of J in dimension $2t-1$ (isomorphic to Z_2 if $2t-1 \equiv 5 \pmod{8}$). There is a map $q_1: \pi_n^S \rightarrow E_2^{1,n+1}$ such that an element of $E_2^{1,n+1}$ survives to E_∞ iff it belongs to $\text{im } q_1$. Furthermore, if \tilde{q}_1 denotes the restriction of q_1 to $\text{im } J$, then [7, Chapters 10 and 11]

- (1) if $n = 8k + 1$, $E_\infty^{1,n+1} = E_\infty^{1,n+1} = Z_2$;
 - (2) if $n = 8k + 3$ ($k > 0$), then $\text{im } q_1 = \text{im } \tilde{q}_1$ has index 2 in $E_2^{1,n+1} = Z_m(4k+2)$, and \tilde{q}_1 has kernel Z_2 ; in fact, $d_3 \alpha_{4k+2} = h^3 \alpha_{4k} \neq 0$;
 - (3) if $n = 8k + 5$, $E_2^{1,n+1} = Z_2$ does not survive to E_∞ ; in fact, $d_3 \alpha_{4k+3} = h^3 \alpha_{4k+1} \neq 0$;
 - (4) if $n = 8k + 7$, $\text{im } \tilde{q}_1 = Z_m(4k+4) = E_2^{1,n+1} = E_\infty^{1,n+1}$.
- Here $h = \alpha_1$.

2. Quillen's algebra. Novikov knew that, given a prime p , the algebra $A^{BP} = BP^*(BP)$ was much simpler than $A^{MU} \otimes Q_p$, but he did not have complete information about A^{BP} . Later, Quillen [8] discovered an idempotent ϵ , which split the spectrum MUQ_p into a sum of suspensions of the spectrum BP_p [4]. Now

$$\pi_*(BP) = Q_p[k_1, k_2, \dots], \quad H_*(BP) = Q_2[m_1, m_2, \dots],$$

with $|k_i| = -|m_i| = -2(p^i - 1)$. We can take $m_i = (1/p^i)h\epsilon[CP^{p^i-1}]$; the Hurewicz homomorphism h is monic, and may be computed using Quillen's formal-group techniques [11] or standard methods. Thanks to the idempotent ϵ , Quillen and Adams were able to write down explicit formulas for the Hopf-algebra structure of the algebra of operations A^{BP} ($= A$, for short).

First, there is a coalgebra R of operations, free as a Q_p -module on generators r_E , where E runs over all finitely nonzero sequences (e_1, e_2, \dots) of nonnegative integers and $|r_E| = 2(\sum (p^i - 1)e_i)$. The diagonal map is given by $\phi^*r_E = \sum_{E'+E''=E} r_{E'} \otimes r_{E''}$. Then $\Lambda' = \pi_*(BP)$ is an algebra over the coalgebra R , with action given (via the Hurewicz map) by $r_E m_n = m_{n-i}$ if $e_i = p^{n-i}$ and all other e_j are zero, and $r_E m_n = 0$ otherwise. Moreover, multiplication by an element λ of Λ' is also a BP -cohomology operation, and in fact every operation can be written as a (possibly infinite) sum $\sum \lambda_i r_{E_i}$ in which the degree of each $\lambda_i r_{E_i}$ is a constant independent of i . Unfortunately, the composition $r_E r_F$ of two operations in R does not usually lie in R ; however, it can be written uniquely as a finite sum $r_E r_F = \sum_K c_K r_K$ with $c_K \in \Lambda'$, using the methods of [11] or those of [4]. This enables us to express compositions $(\lambda r_E)(\lambda' r_F)$ in the form $\sum \lambda_i r_{E_i}$. Thus the algebra A of all operations is the completed tensor product $\Lambda' \hat{\otimes} R$.

PROPOSITION 1. *Let $\bar{\Lambda}$ be the two-sided ideal in A generated by all elements of Λ of negative degree. Let $\mathcal{Q}_p/(Q_0)$ be the algebra of reduced Steenrod p th powers [6]. Then there is an isomorphism $f: A/\bar{\Lambda} \cong \mathcal{Q}_p/(Q_0)$.*

PROOF. Let $\text{Th}: BP_p \rightarrow K(Z_p)$ be the Z_p Thom class. Then

$$\begin{array}{ccc} \tilde{f} = \text{Th}_*: [BP, BP] & \rightarrow & [BP, K(Z_p)] \\ \parallel & & \parallel \\ A & & H^*(BP; Z_p) \\ & & \parallel \\ & & \mathcal{Q}_p/(Q_0) \end{array}$$

satisfies

$$\begin{aligned} \tilde{f}(k^E r_F) &= c(\mathcal{O}^F), & E = 0 \text{ [6];} \\ &= 0, & \text{otherwise;} \end{aligned}$$

where c is the canonical antiautomorphism. The map \tilde{f} induces the required f on $A/\bar{\Lambda}$.

A generator r_E is *indecomposable* if it cannot be expressed as a finite sum $r_E = \sum \lambda_i R_i R_i'$, where $\lambda_i \in \Lambda'$; $R_i, R_i' \in R$; and $|R_i|, |R_i'| > 0$.

THEOREM 2. *The generator r_E of R is indecomposable if and only if $E = (p^i, 0, 0, \dots), i \geq 0$. Moreover, $pr_{(p^i, 0, 0, \dots)}$ is decomposable.*

The proof is obtained by noticing certain pleasant properties of the multiplication table for R and applying them in the proper sequence.

3. Resolutions over A . To compute Ext we must construct resolutions over A , which seems difficult at first glance since R is not an algebra, A is not connected, and the ground ring Q_p is not a field. The next proposition shows how to circumvent some of these difficulties. Define the filtrations $F^s \Lambda' = \sum_{i \leq 2s} (\Lambda')^i$, $F^s A = F^s \Lambda' \hat{\otimes} A$, and $F^s M = (F^s A)M$ if M is an A -module. We have

$$0 \rightarrow F^1 M \xrightarrow{i} M \xrightarrow{j} \text{cok } i \rightarrow 0.$$

Write JM for $\text{cok } i$; then J is easily made into a functor on the category of A -modules.

PROPOSITION 3. *There exist complexes*

$$C: \dots \rightarrow C_i \xrightarrow{d_i} C_{i-1} \rightarrow \dots \rightarrow C_1 \xrightarrow{d_1} C_0 = A \rightarrow \Lambda' \rightarrow 0$$

satisfying

- (1) $C_1 = \sum Au_j$ with $d_{1u_j} = r_{(p^i, 0, 0, \dots)}$;
- (2) $C_i = \prod_j Aw_j^{(i)}$ is locally finitely generated as an A -module, $i > 1$;
- (3) $\ker(Jd_i) \subset j(\text{im } d_{i+1})$ in JC_i for all $i, n \geq 0$.

Any such C is an A -projective resolution of Λ' .

The proof is straightforward. Notice that the infinite direct product $\prod Aw_j^{(i)}$ is not necessarily free over A ; it is projective, however. As a further aid to computation there is

LEMMA 4. *If $\{C_i, d_i\}$ is any A -projective resolution of Λ' , write $C_i^* = \text{Hom}_A^*(C_i, \Lambda')$, $d_i^* = \text{Hom}_A^*(d_i, \Lambda')$. Then*

$$\text{Ext}_A^{s,t}(\Lambda', \Lambda') = \text{Tors}(\text{cok}(d_s^*)^t), \quad (s, t) \neq (0, 0).$$

PROOF. This follows from the fact that $\text{Ext}_A^{s,t}$ is finite for $(s, t) \neq (0, 0)$ [7, Corollary 2.1].

Thus in determining Ext we need know just the boundaries, and not the cycles too. In fact we can even work over Z_{p^r} for suitable f .

Now we can prove

PROPOSITION 5. $\text{Ext}^{0,t} = 0$ unless $t = 0$; $\text{Ext}^{0,0} = Z$.

THEOREM 6. $\text{Ext}^{2,t}$ contains a direct summand isomorphic to Z_p for $t = 2p^i(p-1)$ ($i \geq 1$) and $t = 2(p^i+1)(p-1)$ ($i > 1$).

THEOREM 7. For $p = 2$, the element of $\text{Ext}^{2,2^i}$ found in Theorem 6 maps to the Arf-invariant element h_i^2 of the classical Adams spectral sequence [5].

PROOF. Apply the Thom map (Proposition 1) to a suitable A -resolution.

PROPOSITION 8. The two-primary part ${}_2\text{Ext}^{*,t}$ has the following "edge" values:

$$\begin{aligned} {}_2\text{Ext}^{n,2(n+k)} &= 0, & k < 0; \\ &= Z_2, & k = 0, n \geq 1 \text{ (generated by } h^n); \\ &= 0, & k = 1, n \geq 2; \\ &= Z_2, & 2 \leq k \leq 5, n \geq 4 \text{ (generated by } h^{n-1}\alpha_{k+1}). \end{aligned}$$

Further computations of the additive structure of ${}_2\text{Ext}^{*,*}$ in low dimensions are given in Figure 1. Thanks to Proposition 8, the first three nonzero Novikov differentials $d_3\alpha_i = h^3\alpha_{i-1}$, $i = 3, 6, 7$, give rise to infinite towers of nonzero d_3 's. Moreover, every other differential in the range $t-s \leq 17$ must be zero for dimensional reasons. Finally, ${}_2E_\infty$ has a vanishing line considerably lower than that of the E_∞ -term of the classical Adams spectral sequence in this range of dimensions. We conjecture that the preceding four sentences are also true without restriction on the dimensions.

Similar computations for $p = 3$ disclose striking edge properties like Proposition 8, but many fewer differentials. Contrary to Novikov's conjecture [7], there is a nonzero differential $d_5: E_2^{2,36} \rightarrow E_2^{7,40}$ for $p = 3$. This differential, whose existence is inferred from Toda's result [10], also gives rise to an infinite family of nonzero differentials. It is encouraging that there is only one nonzero differential in the range $t-s \leq 40$, as compared to 17 in the classical 3-primary Adams spectral sequence.

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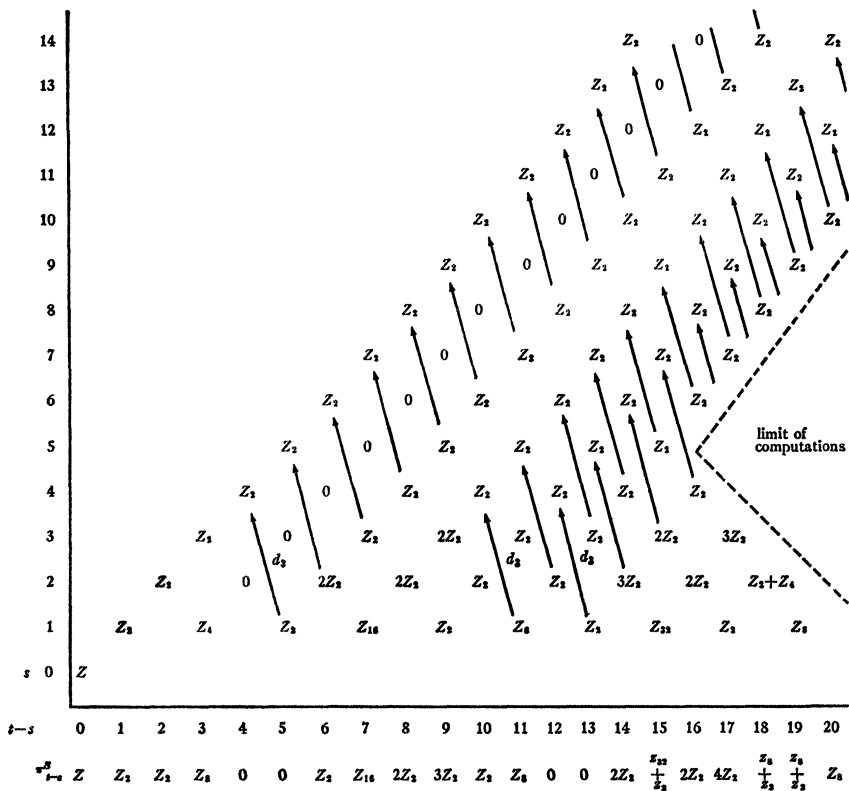


FIGURE 1. ${}_2\text{Ext}^{s,t}$ for the Novikov sequence.