# THE CONTACT HOMOLOGY OF LEGENDRIAN SUBMANIFOLDS IN $\mathbb{R}^{2 n+1}$ 

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#### Abstract

We define the contact homology for Legendrian submanifolds in standard contact $(2 n+1)$-space using moduli spaces of holomorphic disks with Lagrangian boundary conditions in complex $n$-space. This homology provides new invariants of Legendrian isotopy which indicate that the theory of Legendrian isotopy is very rich. Indeed, in [4], the homology is used to detect infinite families of pairwise non-isotopic Legendrian submanifolds which are indistinguishable using previously known invariants.


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## 1. Introduction

The motivating problem for this paper is the classification of Legendrian submanifolds up to Legendrian isotopy. Here we restrict attention to the standard contact structure on $\mathbb{R}^{2 n+1}$. For $n=1$, the Legendrian isotopy problem has been extensively studied, $[\mathbf{2}, \mathbf{6}, \mathbf{9}, \mathbf{1 0}, \mathbf{1 1}]$, but there

[^0]have been few results for $n>1$. In this paper, we give a rigorous definition of contact homology, a potent new invariant originally described in [5]. This new invariant was applied in [4] to construct infinite families of non-Legendrian isotopic, Legendrian $n$-spheres, $n$-tori and surfaces of arbitrary genus. These are the first such high-dimensional examples. They also demonstrate that the analogues of rotation number and Thurston-Bennequin invariant (and diffeomorphism type) of a Legendrian submanifold are far from complete invariants of Legendrian isotopy. (See [4] for a definition of the high-dimensional analogues of the classical invariants.)

The goal of this paper is to define contact homology and prove that it is a Legendrian isotopy invariant.

Theorem. The contact homology of Legendrian submanifolds in $\mathbb{R}^{2 n+1}$ with the standard contact form is well defined. (It is invariant under Legendrian isotopy.)

We define the contact homology using punctured holomorphic disks in $\mathbb{C}^{n} \approx \mathbb{R}^{2 n}$ with boundary on the Lagrangian projection $\Pi_{\mathbb{C}}: \mathbb{C}^{n} \times \mathbb{R} \rightarrow$ $\mathbb{C}^{n}$ of the Legendrian submanifold, and which limit to double points of the projection at the punctures. This is analogous to the approach taken by Chekanov [2] in dimension 3 who was the first to prove that the classical invariants are not enough to distinguish isotopy classes. In dimension 3 , however, the entire theory can be reduced to combinatorics. As discussed in [4], our contact homology also fits into the over arching philosophy of Symplectic Field Theory outlined in [8]. There it goes by the name of the "relative contact homology" of the standard contact $(2 n+1)$-space.

In Section 2, we define contact homology more concretely and outline its invariance under Legendrian isotopy. If $L \subset \mathbb{R}^{2 n+1} \approx \mathbb{C}^{n} \times \mathbb{R}$ is a Legendrian submanifold we associate to $L$ a differential graded algebra (DGA), denoted $(\mathcal{A}, \partial)$, freely generated by the double points of $\Pi_{\mathbb{C}}(L) \subset \mathbb{C}^{n}$.

Since $L$ is embedded, one may distinguish upper and lower branches of $L$ at double points of $\Pi_{\mathbb{C}}(L)$ and using this structure, we associate a sign to every puncture of a holomorphic disk with boundary on $\Pi_{\mathbb{C}}(L)$. We define the differential of the DGA by counting punctured rigid holomorphic disks with boundary on $\Pi_{\mathbb{C}}(L)$ and with exactly one positive puncture. The contact homology of $L$ is defined to be $\operatorname{Ker} \partial / \operatorname{Im} \partial$. Thus, contact homology is similar to Floer homology of Lagrangian intersections. The proof of its invariance is similar in spirit to Floer's original approach $[\mathbf{1 3}, \mathbf{1 4}]$; we study bifurcations of moduli spaces of rigid holomorphic disks under variations of the Legendrian submanifold in a generic 1-parameter family of Legendrian submanifolds. Similar bifurcation analysis is also done in $[\mathbf{1 9}, \mathbf{2 1}, \mathbf{3 0}, 31]$.

In Section 6, the (formal) dimension of the moduli space of punctured holomorphic disks with boundary on an exact Lagrangian immersion which is an instant in a generic 1-parameter family is expressed in terms of its boundary data. We compute this dimension by relating the linearization of the $\bar{\partial}$-equation for punctured disks with boundary on the exact Lagrangian to the standard vector Riemann-Hilbert problem on the closed disk (i.e. the disk without punctures).

In Section 7, we show that for Legendrian submanifolds (and their 1-parameter families) in an open dense set in the space of such, the moduli-spaces of holomorphic disks are being transversely cut out. That is, we achieve transversality for the $\bar{\partial}$-equation without perturbing the complex structure on $\mathbb{C}^{n}$. The fact that we can keep the standard complex structure on $\mathbb{C}^{n}$ is important for computations of contact homology, see [4]. Similar transversality results were obtained by Oh [25] for closed holomorphic disks with Lagrangian boundary condition, under the additional assumption that the disks have an injective point on the boundary. In general, disks without such points cannot be excluded and we manage to prove transversality for disks involved in contact homology using the fact that they have only one positive puncture, and a technical result, established in Section 3, that all Legendrian submanifolds may be assumed real analytic close to the preimages of double points of $\Pi_{\mathbb{C}}$.

In Section 9, we show that moduli-spaces of holomorphic disks have certain compactness properties. We prove a version of Gromov compactness for punctured holomorphic disks with boundary on an immersed exact Lagrangian submanifold in $\mathbb{C}^{n}$. In particular, it follows that 0-dimensional moduli-spaces are compact and that 1-dimensional moduli-spaces have natural compactifications.

In Section 8, we establish gluing theorems. These are used to prove that the differential $\partial$ of the DGA $\mathcal{A}$ satisfies $\partial \circ \partial=0$, and that the homology of $(\mathcal{A}, \partial)$ is left unchanged by the two basic bifurcations which occur in generic 1-parameter families: appearance of disks of formal dimension -1 and self-tangency instances. The most technically difficult results are the so-called degenerate gluing theorems which are necessary to control the changes of the DGA under self-tangencies. Here, holomorphic disks with punctures at the self-tangency double point must be glued. To prove these gluing theorems, we use results from Section 4 which give the blow up rate of the constant in the elliptic estimate for the linearized $\bar{\partial}$-equation, as the transverse double point at one puncture approaches a self-tangency double point. To prove invariance under the appearance of disks of formal dimension -1 , we use an auxiliary Legendrian submanifold, see Section 10 and a method similar to the proof of Floer theory invariance which uses an elegant "homotopy of homotopies" argument (see, for example, [15, 29]).

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## 2. Contact Homology and Differential Graded Algebras

In this section, we describe how to associate to a Legendrian submanifold $L$ in standard contact $(2 n+1)$-space a differential graded algebra (DGA) $(\mathcal{A}, \partial)$. Up to a certain equivalence relation this DGA is an invariant of the Legendrian isotopy class of $L$. In Section 2.1, we recall the notion of Lagrangian projection and define the algebra $\mathcal{A}$. The grading on $\mathcal{A}$ is described in Sections 2.2. Sections 2.3 and 2.4 are devoted to the definition of $\partial$ and Section 2.5 proves the invariance of the homology of $(\mathcal{A}, \partial)$, which we call the contact homology. The main proofs of these three subsections rely on much analysis, which will be completed in the subsequent sections. In a sense, these last three subsections can be viewed as an overview of the remainder of the paper.
2.1. The algebra $\mathcal{A}$. Throughout this paper, we consider the standard contact structure $\xi$ on $\mathbb{R}^{2 n+1}=\mathbb{C}^{n} \times \mathbb{R}$ which is the hyperplane field given as the kernel of the contact 1 -form

$$
\begin{equation*}
\alpha=d z-\sum_{j=1}^{n} y_{j} d x_{j}, \tag{2.1}
\end{equation*}
$$

where $x_{1}, y_{1}, \ldots, x_{n}, y_{n}, z$ are Euclidean coordinates on $\mathbb{R}^{2 n+1}$. A Legendrian submanifold of $\mathbb{R}^{2 n+1}$ is an $n$-dimensional submanifold $L \subset \mathbb{R}^{2 n+1}$ everywhere tangent to $\xi$. We also recall that the standard symplectic structure on $\mathbb{C}^{n}$ is given by

$$
\omega=\sum_{j=1}^{n} d x_{j} \wedge d y_{j},
$$

and that an immersion $f: L \rightarrow \mathbb{C}^{n}$ of an $n$-dimensional manifold is Lagrangian if $f^{*} \omega=0$.

The Lagrangian projection projects out the $z$ coordinate:

$$
\begin{equation*}
\Pi_{\mathbb{C}}: \mathbb{R}^{2 n+1} \rightarrow \mathbb{C}^{n} ; \quad\left(x_{1}, y_{1}, \ldots, x_{n}, y_{n}, z\right) \mapsto\left(x_{1}, y_{1}, \ldots, x_{n}, y_{n}\right) . \tag{2.2}
\end{equation*}
$$

If $L \subset \mathbb{C}^{n} \times \mathbb{R}$ is a Legendrian submanifold, then $\Pi_{\mathbb{C}}: L \rightarrow \mathbb{C}^{n}$ is a Lagrangian immersion. Moreover, for $L$ in an open dense subset of all Legendrian submanifolds (with $C^{\infty}$ topology), the self intersection of $\Pi_{\mathbb{C}}(L)$ consists of a finite number of transverse double points. We call Legendrian submanifolds with this property chord generic.

The Reeb vector field $X$ of a contact form $\alpha$ is uniquely defined by the two equations $\alpha(X)=1$ and $d \alpha(X, \cdot)=0$. The Reeb chords of a Legendrian submanifold $L$ are segments of flow lines of $X$ starting and ending at points of $L$. We see from (2.1) that in $\mathbb{R}^{2 n+1}, X=\frac{\partial}{\partial z}$ and thus $\Pi_{\mathbb{C}}$ defines a bijection between Reeb chords of $L$ and double points of $\Pi_{\mathbb{C}}(L)$. If $c$ is a Reeb chord, we write $c^{*}=\Pi_{\mathbb{C}}(c)$.

Let $\mathcal{C}=\left\{c_{1}, \ldots, c_{m}\right\}$ be the set of Reeb chords of a chord generic Legendrian submanifold $L \subset \mathbb{R}^{2 n+1}$. To such an $L$, we associate an algebra $\mathcal{A}=\mathcal{A}(L)$ which is the free associative unital algebra over the group ring $\mathbb{Z}_{2}\left[H_{1}(L)\right]$ generated by $\mathcal{C}$. We write elements in $\mathcal{A}$ as

$$
\begin{equation*}
\sum_{i} t_{1}^{n_{1, i}} \ldots t_{k}^{n_{k, i}} \mathbf{c}_{i} \tag{2.3}
\end{equation*}
$$

where the $t_{j}$ 's are formal variables corresponding to a basis for $H_{1}(L)$ thought of multiplicatively and $\mathbf{c}_{i}=c_{i_{1}} \ldots c_{i_{r}}$ is a word in the generators. It is also useful to consider the corresponding algebra $\mathcal{A}_{\mathbb{Z}_{2}}$ over $\mathbb{Z}_{2}$. The natural map $\mathbb{Z}_{2}\left[H_{1}(L)\right] \rightarrow \mathbb{Z}_{2}$ induces a reduction of $\mathcal{A}$ to $\mathcal{A}_{\mathbb{Z}_{2}}$ (set $t_{j}=1$, for all $j$ ).
2.2. The grading on $\mathcal{A}$. Let $\Lambda_{n}$ be the Grassman manifold of Lagrangian subspaces in the symplectic vector space $\left(\mathbb{C}^{n}, \omega\right)$ and recall that $H_{1}\left(\Lambda_{n}\right)=\pi_{1}\left(\Lambda_{n}\right) \cong \mathbb{Z}$. There is a standard isomorphism

$$
\mu: H_{1}\left(\Lambda_{n}\right) \rightarrow \mathbb{Z}
$$

given by intersecting a loop in $\Lambda_{n}$ with the Maslov cycle $\Sigma$. To describe $\mu$ more fully, we follow [26] and refer the reader to this paper for proofs of the statements below.

Fix a Lagrangian subspace $\Lambda$ in $\mathbb{C}^{n}$ and let $\Sigma_{k}(\Lambda) \subset \Lambda_{n}$ be the subset of Lagrangian spaces that intersects $\Lambda$ in a subspace of $k$ dimensions. The Maslov cycle is

$$
\Sigma=\overline{\Sigma_{1}(\Lambda)}=\Sigma_{1}(\Lambda) \cup \Sigma_{2}(\Lambda) \cup \cdots \cup \Sigma_{n}(\Lambda) .
$$

This in an algebraic variety of codimension one in $\Lambda_{n}$. If $\Gamma:[0,1] \rightarrow \Lambda_{n}$ is a loop then $\mu(\Gamma)$ is the intersection number of $\Gamma$ and $\Sigma$. The contribution of an intersection point $t^{\prime}$ with $\Gamma\left(t^{\prime}\right) \in \Sigma$ to $\mu(\Gamma)$ is calculated as follows. Fix a Lagrangian complement $W$ of $\Lambda$. Then for each $v \in \Gamma\left(t^{\prime}\right) \cap \Lambda$ there exists a vector $w(t) \in W$ such that $v+w(t) \in \Gamma(t)$ for $t$ near $t^{\prime}$. Define the quadratic form $Q(v)=\left.\frac{d}{d t}\right|_{t=t^{\prime}} \omega(v, w(t))$ on $\Gamma\left(t^{\prime}\right) \cap \Lambda$ and observe that it is independent of the complement $W$ chosen. Without loss of generality, $Q$ can be assumed non-singular and the contribution of the
intersection point to $\mu(\Gamma)$ is the signature of $Q$. Given any loop $\Gamma$ in $\Lambda_{n}$, we say $\mu(\Gamma)$ is the Maslov index of the loop.

If $f: L \rightarrow \mathbb{C}^{n}$ is a Lagrangian immersion then the tangent planes of $f(L)$ along any loop $\gamma$ in $L$ gives a loop $\Gamma$ in $\Lambda_{n}$. We define the Maslov index $\mu(\gamma)$ of $\gamma$ as $\mu(\gamma)=\mu(\Gamma)$ and note that we may view the Maslov index as a map $\mu: H_{1}(L) \rightarrow \mathbb{Z}$. Let $m(f)$ be the smallest non-negative number that is the Maslov index of some non-trivial loop in $L$. We call $m(f)$ the Maslov number of $f$. When $L \subset \mathbb{C}^{n} \times \mathbb{R}$ is a Legendrian submanifold, we write $m(L)$ for the Maslov number of $\Pi_{\mathbb{C}}: L \rightarrow \mathbb{C}^{n}$.

Let $L \subset \mathbb{R}^{2 n+1}$ be a chord generic Legendrian submanifold and let $c$ be one of its Reeb chords with end points $a, b \in L, z(a)>z(b)$. Choose a path $\gamma:[0,1] \rightarrow L$ with $\gamma(0)=a$ and $\gamma(1)=b$. (We call such path a capping path of c.) Then $\Pi_{\mathbb{C}} \circ \gamma$ is a loop in $\mathbb{C}^{n}$ and $\Gamma(t)=d \Pi_{\mathbb{C}}\left(T_{\gamma(t)} L\right)$, $0 \leq t \leq 1$ is a path of Lagrangian subspaces of $\mathbb{C}^{n}$. Since $c^{*}=\Pi_{\mathbb{C}}(c)$ is a transverse double point of $\Pi_{\mathbb{C}}(L), \Gamma$ is not a closed loop.

We close $\Gamma$ in the following way. Let $V_{0}=\Gamma(0)$ and $V_{1}=\Gamma(1)$. Choose any complex structure $I$ on $\mathbb{C}^{n}$ which is compatible with $\omega$ $(\omega(v, I v)>0$ for all $v)$ and with $I\left(V_{1}\right)=V_{0}$. (Such an $I$ exists since the Lagrangian planes are transverse.) Define the path $\lambda\left(V_{1}, V_{0}\right)(t)=e^{t I} V_{1}$, $0 \leq t \leq \frac{\pi}{2}$. The concatenation, $\Gamma * \lambda\left(V_{1}, V_{0}\right)$, of $\Gamma$ and $\lambda\left(V_{1}, V_{0}\right)$ forms a loop in $\Lambda_{n}$ and we define the Conley-Zehnder index, $\nu_{\gamma}(c)$, of $c$ to be the Maslov index $\mu\left(\Gamma * \lambda\left(V_{1}, V_{0}\right)\right)$ of this loop. It is easy to check that $\nu_{\gamma}(c)$ is independent of the choice of $I$. However, $\nu_{\gamma}(c)$ might depend on the choice of homotopy class of the path $\gamma$. More precisely, if $\gamma_{1}$ and $\gamma_{2}$ are two paths with properties as $\gamma$ above then

$$
\nu_{\gamma_{1}}(c)-\nu_{\gamma_{2}}(c)=\mu\left(\gamma_{1} *\left(-\gamma_{2}\right)\right),
$$

where $\left(-\gamma_{2}\right)$ is the path $\gamma_{2}$ traversed in the opposite direction. Thus $\nu_{\gamma}(c)$ is well defined modulo the Maslov number $m(L)$.

Let $\mathcal{C}=\left\{c_{1}, \ldots, c_{m}\right\}$ be the set of Reeb chords of $L$. Choose a capping path $\gamma_{j}$ for each $c_{j}$ and define the grading of $c_{j}$ to be

$$
\left|c_{j}\right|=\nu_{\gamma_{j}}\left(c_{j}\right)-1,
$$

and for any $t \in H_{1}(L)$ define its grading to be $|t|=-\mu(t)$. This makes $\mathcal{A}(L)$ into a graded ring. Note that the grading depends on the choice of capping paths but, as we will see below, this choice will be irrelevant.

The above grading on Reeb chords $c_{j}$ taken modulo $m(L)$ makes $\mathcal{A}_{\mathbb{Z}_{2}}$ a graded algebra with grading in $\mathbb{Z}_{m(L)}$. (Note that this grading does not depend on the choice of capping paths.) In addition the map from $\mathcal{A}$ to $\mathcal{A}_{\mathbb{Z}_{2}}$ preserves gradings modulo $m(L)$.
2.3. The moduli spaces. As mentioned in the introduction, the differential of the algebra associated to a Legendrian submanifold is defined using spaces of holomorphic disks. To describe these spaces we need a few preliminary definitions.

Let $D_{m+1}$ be the unit disk in $\mathbb{C}$ with $m+1$ punctures at the points $p_{0}, \ldots p_{m}$ on the boundary. The orientation of the boundary of the unit disk induces a cyclic ordering of the punctures. Let $\partial \hat{D}_{m+1}=$ $\partial D_{m+1} \backslash\left\{p_{0}, \ldots, p_{m}\right\}$.

Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a Legendrian submanifold with isolated Reeb chords. If $c$ is a Reeb chord of $L$ with end points $a, b \in L, z(a)>z(b)$ then there are small neighborhoods $S_{a} \subset L$ of $a$ and $S_{b} \subset L$ of $b$ that are mapped injectively to $\mathbb{C}^{n}$ by $\Pi_{\mathbb{C}}$. We call $\Pi_{\mathbb{C}}\left(S_{a}\right)$ the upper sheet of $\Pi_{\mathbb{C}}(L)$ at $c^{*}$ and $\Pi_{\mathbb{C}}\left(S_{b}\right)$ the lower sheet. If $u:\left(D_{m+1}, \partial D_{m+1}\right) \rightarrow$ $\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$ is a continuous map with $u\left(p_{j}\right)=c^{*}$, then we say $p_{j}$ is positive (respectively negative) if $u$ maps points clockwise of $p_{j}$ on $\partial D_{m+1}$ to the lower (upper) sheet of $\Pi_{\mathbb{C}}(L)$ and points anti-clockwise of $p_{i}$ on $\partial D_{m+1}$ to the upper (lower) sheet of $\Pi_{\mathbb{C}}(L)$ (see Figure 1).


Figure 1. Positive puncture lifted to $\mathbb{R}^{2 n+1}$. The gray region is the holomorphic disk and the arrows indicate the orientation on the disk and the Reeb chord.

If $a$ is a Reeb chord of $L$ and if $\mathbf{b}=b_{1} \ldots b_{m}$ is an ordered collection (a word) of Reeb chords, then let $\mathcal{M}_{A}(a ; \mathbf{b})$ be the space, modulo conformal reparameterization, of maps $u:\left(D_{m+1}, \partial D_{m+1}\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$ which are continuous on $D_{m+1}$, holomorphic in the interior of $D_{m+1}$, and which have the following properties

- $p_{0}$ is a positive puncture, $u\left(p_{0}\right)=a^{*}$,
- $p_{j}$ are negative punctures for $j>0, u\left(p_{j}\right)=b_{j}^{*}$,
- the restriction $u \mid \partial \hat{D}_{m+1}$ has a continuous lift $\tilde{u}: \partial \hat{D}_{m+1} \rightarrow L \subset$ $\mathbb{C}^{n} \times \mathbb{R}$, and
- the homology class of $\tilde{u}\left(\partial D_{m+1}^{*}\right) \cup\left(\cup_{j} \gamma_{j}\right)$ equals $A \in H_{1}(L)$,
where $\gamma_{j}$ is the capping path chosen for $c_{j}, j=1, \ldots, m$. Elements in $\mathcal{M}_{A}(a ; \mathbf{b})$ will be called holomorphic disks with boundary on $L$ or sometimes simply holomorphic disks.

There is a useful fact relating heights of Reeb chords and the area of a holomorphic disk with punctures mapping to the corresponding double points. The action (or height) $\mathcal{Z}(c)$ of a Reeb chord $c$ is simply
its length and the action of a word of Reeb chords is the sum of the actions of the chords making up the word.

Lemma 2.1. If $u \in \mathcal{M}_{A}(a ; \mathbf{b})$, then

$$
\begin{equation*}
\mathcal{Z}(a)-\mathcal{Z}(\mathbf{b})=\int_{D_{m}} u^{*} \omega=\operatorname{Area}(u) \geq 0 . \tag{2.4}
\end{equation*}
$$

Proof. By Stokes theorem, $\int_{D_{m}} u^{*} \omega=\int_{\partial D_{m}} u^{*}\left(-\sum_{j} y_{j} d x_{j}\right)=$ $\int \tilde{u}^{*}(-d z)=\mathcal{Z}(a)-\mathcal{Z}(\mathbf{b})$. The second equality follows since $u$ is holomorphic and $\omega=\sum_{j=1}^{n} d x_{j} \wedge d y_{j}$. q.e.d.

Note that the proof of Lemma 2.1 implies that any holomorphic disk with boundary on $L$ must have at least one positive puncture. (In contact homology, only disks with exactly one positive puncture are considered.)

We now proceed to describe the properties of moduli spaces $\mathcal{M}_{A}(a ; \mathbf{b})$ that are needed to define the differential. We prove later that the moduli spaces of holomorphic disks with boundary on a Legendrian submanifold $L$ have these properties provided $L$ is generic among (belongs to a Baire subset of the space of) admissible Legendrian submanifolds ( $L$ is admissible if it is chord generic and it is real analytic in a neighborhood of all Reeb chord end points). For more precise definitions of these concepts, see Section 3, where it is shown that admissible Legendrian submanifolds are dense in the space of all Legendrian submanifolds. In Section 5 , we express moduli spaces $\mathcal{M}_{A}(a ; \mathbf{b})$ as 0 -sets of certain $C^{1}$-maps between infinite-dimensional Banach manifolds. We say a moduli space is transversely cut out if 0 is a regular value of the corresponding map.

Proposition 2.2. For a generic admissible Legendrian submanifold $L \subset \mathbb{C}^{n} \times \mathbb{R}$, the moduli space $\mathcal{M}_{A}(a ; \mathbf{b})$ is a transversely cut out manifold of dimension

$$
\begin{equation*}
d=\mu(A)+|a|-|\mathbf{b}|-1, \tag{2.5}
\end{equation*}
$$

provided $d \leq 1$. (In particular, if $d<0$ then the moduli space is empty.)
Proposition 2.2 is proved in Section 7.8. If $u \in \mathcal{M}_{A}(a ; \mathbf{b})$, we say that $d=\mu(A)+|a|-|\mathbf{b}|$ is the formal dimension of $u$, and if $v$ is a transversely cut out disk of formal dimension 0 we say that $v$ is a rigid disk.

The moduli spaces we consider might not be compact, but their lack of compactness can be understood. It is analogous to "convergence to broken trajectories" in Morse/Floer homology and gives rise to natural compactifications of the moduli spaces. This is also called Gromov compactness, which we cover in more detail in Section 9.

A broken holomorphic curve, $u=\left(u^{1}, \ldots, u^{N}\right)$, is a union of holomorphic disks, $u^{j}:\left(D_{m_{j}}, \partial D_{m_{j}}\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$, where each $u^{j}$ has exactly one positive puncture $p^{j}$, with the following property. To each $p^{j}$ with $j \geq 2$ is associated a negative puncture $q_{j}^{k} \in D_{m_{k}}$ for some $k \neq j$ such
that $u^{j}\left(p^{j}\right)=u^{k}\left(q_{j}^{k}\right)$ and $q_{j^{\prime}}^{k^{\prime}} \neq q_{j}^{k}$ if $j \neq j^{\prime}$, and such that the quotient space obtained from $D_{m_{1}} \cup \cdots \cup D_{m_{N}}$ by identifying $p^{j}$ and $q_{j}^{k}$ for each $j \geq 2$ is contractible. The broken curve can be parameterized by a single smooth $v:\left(D_{m}, \partial D\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$. A sequence $u_{\alpha}$ of holomorphic disks converges to a broken curve $u=\left(u^{1}, \ldots, u^{N}\right)$ if the following holds:

- For every $j \leq N$, there exists a sequence $\phi_{\alpha}^{j}: D_{m} \rightarrow D_{m}$ of linear fractional transformations and a finite set $X^{j} \subset D_{m}$ such that $u_{\alpha} \circ \phi_{\alpha}^{j}$ converges to $u^{j}$ uniformly with all derivatives on compact subsets of $D_{m} \backslash X^{j}$
- There exists a sequence of orientation-preserving diffeomorphisms $f_{\alpha}: D_{m} \rightarrow D_{m}$ such that $u_{\alpha} \circ f_{\alpha}$ converges in the $C^{0}$-topology to a parameterization of $u$.

Proposition 2.3. Any sequence $u_{\alpha}$ in $\mathcal{M}_{A}(a ; \mathbf{b})$ has a subsequence converging to a broken holomorphic curve $u=\left(u^{1}, \ldots, u^{N}\right)$. Moreover, $u^{j} \in \mathcal{M}_{A_{j}}\left(a^{j} ; \mathbf{b}^{j}\right)$ with $A=\sum_{j=1}^{N} A_{j}$ and

$$
\begin{equation*}
\mu(A)+|a|-|\mathbf{b}|=\sum_{j=1}^{N}\left(\mu\left(A_{j}\right)+\left|a^{j}\right|-\left|\mathbf{b}^{j}\right|\right) . \tag{2.6}
\end{equation*}
$$

Heuristically, this is the only type of non-compactness we expect to see in $\mathcal{M}_{A}(a ; \mathbf{b})$ : since $\pi_{2}\left(\mathbb{C}^{n}\right)=0$, no holomorphic spheres can "bubble off" at an interior point of the sequence $u_{\alpha}$, and since $\Pi_{\mathbb{C}}(L)$ is exact no disks without positive puncture can form either. Moreover, since $\Pi_{\mathbb{C}}(L)$ is compact, and since $\mathbb{C}^{n}$ has "finite geometry at infinity" (see Section 9 ), all holomorphic curves with a uniform bound on area must map to a compact set.

Proof. The main step is to prove convergence to some broken curve, which we defer to Section 9. The statement about the homology classes follows easily from the definition of convergence. Equation (2.6) follows from the definition of broken curves. q.e.d.

We next show that a broken curve can be glued to form a family of non-broken curves. For this, we need a little notation. Let $\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}$ be an ordered collection of words of Reeb chords. Let the length of (number of letters in) $\mathbf{c}^{j}$ be $l(j)$ and let $\mathbf{a}=a_{1} \ldots a_{k}$ be a word of Reeb chords of length $k>0$. Let $S=\left\{s_{1}, \ldots, s_{r}\right\}$ be $r$ distinct integers in $\{1, \ldots, k\}$. Define the word $\mathbf{a}_{S}\left(\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}\right)$ of Reeb chords of length $k-r+\sum_{j=1}^{r} l(j)$ as follows. For each index $s_{j} \in S$, remove $a_{s_{j}}$ from the word $\mathbf{a}$ and insert at its place the word $\mathbf{c}^{j}$.

Proposition 2.4. Let L be a generic admissible Legendrian submanifold. Let $\mathcal{M}_{A}(a ; \mathbf{b})$ and $\mathcal{M}_{B}(c ; \mathbf{d})$ be 0 -dimensional transversely cut out moduli spaces and assume that the $j$-th Reeb chord in $\mathbf{b}$ is $c$. Then there
exist a $\rho>0$ and an embedding

$$
G: \mathcal{M}_{A}(a ; \mathbf{b}) \times \mathcal{M}_{B}(c ; \mathbf{d}) \times(\rho, \infty) \rightarrow \mathcal{M}_{A+B}\left(a ; \mathbf{b}_{\{j\}}(\mathbf{d})\right) .
$$

Moreover, if $u \in \mathcal{M}_{A}(a ; \mathbf{b})$ and $u^{\prime} \in \mathcal{M}_{B}(c ; \mathbf{d})$ then $G\left(u, u^{\prime}, \rho\right)$ converges to the broken curve $\left(u, u^{\prime}\right)$ as $\rho \rightarrow \infty$, and any disk in $\mathcal{M}_{A}\left(a ; \mathbf{b}_{\{j\}}(\mathbf{d})\right)$ with image sufficiently close to the image of $\left(u, u^{\prime}\right)$ is in the image of $G$.

This follows from Proposition 8.1 and the definition of convergence to a broken curve.
2.4. The differential and contact homology. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a generic admissible Legendrian submanifold, let $\mathcal{C}$ be its set of Reeb chords, and let $\mathcal{A}$ denote its algebra. For any generator $a \in \mathcal{C}$ of $\mathcal{A}$ we set

$$
\begin{equation*}
\partial a=\sum_{\operatorname{dim} \mathcal{M}_{A}(a ; \mathbf{b})=0}\left(\# \mathcal{M}_{A}(a ; \mathbf{b})\right) A \mathbf{b}, \tag{2.7}
\end{equation*}
$$

where $\# \mathcal{M}$ is the number of points in $\mathcal{M}$ modulo 2 , and where the sum ranges over all words $\mathbf{b}$ in the alphabet $\mathcal{C}$ and $A \in H_{1}(L)$ for which the above moduli space has dimension 0 . We then extend $\partial$ to a map $\partial: \mathcal{A} \rightarrow \mathcal{A}$ by linearity and the Leibniz rule.

Since $L$ is generic admissible, it follows from Propositions 2.3 and 2.4 that the moduli spaces considered in the definition of $\partial$ are compact 0 -manifolds and hence consist of a finite number of points. Thus $\partial$ is well defined. Moreover,

Lemma 2.5. The map $\partial: \mathcal{A} \rightarrow \mathcal{A}$ is a differential of degree -1 . That is, $\partial \circ \partial=0$ and $|\partial(a)|=|a|-1$ for any generator $a$ of $\mathcal{A}$.

Proof. After Propositions 2.3 and 2.4, the standard proof in Morse (or Floer) homology [28] applies. It follows from (2.5) that $\partial$ lowers degree by 1 .
q.e.d.

The contact homology of $L$ is

$$
H C_{*}\left(\mathbb{R}^{2 n+1}, L\right)=\operatorname{Ker} \partial / \operatorname{Im} \partial
$$

It is essential to notice that since $\partial$ respects the grading on $\mathcal{A}$ the contact homology is a graded algebra.

We note that $\partial$ also defines a differential of degree -1 on $\mathcal{A}_{\mathbb{Z}_{2}}(L)$.

### 2.5. The invariance of contact homology under Legendrian iso-

 topy. In this section, we showProposition 2.6. If $L_{t} \subset \mathbb{R}^{2 n+1}, 0 \leq t \leq 1$ is a Legendrian isotopy between generic admissible Legendrian submanifolds, then the contact homologies $H C_{*}\left(\mathbb{R}^{2 n+1}, L_{0}\right)$, and $H C_{*}\left(\mathbb{R}^{2 n+1}, L_{1}\right)$ are isomorphic.

In fact we show something, that at least appears to be, stronger. Given a graded algebra $\mathcal{A}=\mathbb{Z}_{2}[G]\left\langle a_{1}, \ldots, a_{n}\right\rangle$, where $G$ is a finitely generated abelian group, a graded automorphism $\phi: \mathcal{A} \rightarrow \mathcal{A}$ is called elementary if there is some $1 \leq j \leq n$ such that

$$
\phi\left(a_{i}\right)= \begin{cases}A_{i} a_{i}, & i \neq j \\ \pm A_{j} a_{j}+u, & u \in \mathcal{A}\left(a_{1}, \ldots, a_{j-1}, a_{j+1}, \ldots, a_{n}\right), i=j\end{cases}
$$

where the $A_{i}$ are units in $\mathbb{Z}_{2}[G]$. The composition of elementary automorphisms is called a tame automorphism. An isomorphism from $\mathcal{A}$ to $\mathcal{A}^{\prime}$ is tame if it is the composition of a tame automorphism with an isomorphism sending the generators of $\mathcal{A}$ to the generators of $\mathcal{A}^{\prime}$. An isomorphism of DGA's is called tame if the isomorphism of the underlying algebras is tame.

Let $\left(\mathcal{E}_{i}, \partial_{i}\right)$ be a DGA with generators $\left\{e_{1}^{i}, e_{2}^{i}\right\}$, where $\left|e_{1}^{i}\right|=i,\left|e_{2}^{i}\right|=$ $i-1$ and $\partial_{i} e_{1}^{i}=e_{2}^{i}, \partial_{i} e_{2}^{i}=0$. Define the degree $i$ stabilization $S_{i}(\mathcal{A}, \partial)$ of $(\mathcal{A}, \partial)$ to be the graded algebra generated by $\left\{a_{1}, \ldots, a_{n}, e_{1}^{i}, e_{2}^{i}\right\}$ with grading and differential induced from $\mathcal{A}$ and $\mathcal{E}_{i}$. Two differential graded algebras are called stable tame isomorphic if they become tame isomorphic after each is stabilized a suitable number of times.

Proposition 2.7. If $L_{t} \subset \mathbb{R}^{2 n+1}, 0 \leq t \leq 1$ is a Legendrian isotopy between generic admissible Legendrian submanifolds, then the DGA's $\left(\mathcal{A}\left(L_{0}\right), \partial\right)$ and $\left(\mathcal{A}\left(L_{1}\right), \partial\right)$ are stable tame isomorphic.

Note that Proposition 2.7 allows us to associate the stable tame isomorphism class of a DGA to a Legendrian isotopy class of Legendrian submanifolds: any Legendrian isotopy class has a generic admissible representative and by Proposition 2.7, the DGA's of any two generic admissible representatives agree.

It is straightforward to show that two stable tame isomorphic DGA's have the same homology, see $[\mathbf{2}, \mathbf{1 1}]$. Thus Proposition 2.6 follows from Proposition 2.7. The proof of the later given below is, in outline, the same as the proof of invariance of the stable tame isomorphism class of the DGA of a Legendrian 1-knot in [2]. However, the details in our case require considerably more work. In particular, we must substitute analytic arguments for the purely combinatorial ones that suffice in dimension three.

In Section 3, we show that any two admissible Legendrian submanifolds of dimension $n>2$ which are Legendrian isotopic are isotopic through a special kind of Legendrian isotopy: a Legendrian isotopy $\phi_{t}: L \rightarrow \mathbb{C}^{n} \times \mathbb{R}, 0 \leq t \leq 1$, is admissible if $\phi_{0}(L)$ and $\phi_{1}(L)$ are admissible Legendrian submanifolds and if there exist a finite number of instants $0<t_{1}<t_{2}<\cdots<t_{m}<1$ and a $\delta>0$ such that the intervals $\left[t_{j}-\delta, t_{j}+\delta\right]$ are disjoint subsets of $(0,1)$ with the following properties.
(A) For $t \in\left[0, t_{1}-\delta\right] \cup\left(\bigcup_{j=1}^{m}\left[t_{j}+\delta, t_{j+1}-\delta\right]\right) \cup\left[t_{m}+\delta, 1\right], \phi_{t}(L)$ is an isotopy through admissible Legendrian submanifolds.
(B) For $t \in\left[t_{j}-\delta, t_{j}+\delta\right], j=1, \ldots, m, \phi_{t}(L)$ undergoes a standard self-tangency move. That is, there exists a point $q \in \mathbb{C}^{n}$ and neighborhoods $N \subset N^{\prime}$ of $q$ with the following properties. The intersection $N \cap \Pi_{\mathbb{C}}\left(\phi_{t}(L)\right)$ equals $P_{1} \cup P_{2}(t)$ which, up to biholomorphism looks like $P_{1}=\gamma_{1} \times P_{1}^{\prime}$ and $P_{2}=\gamma_{2}(t) \times P_{2}^{\prime}$. Here $\gamma_{1}$ and $\gamma_{2}(t)$ are subarcs around 0 of the curves $y_{1}=0$ and $x_{1}^{2}+\left(y_{1}-1 \pm t\right)^{2}=1$ in the $z_{1}$-plane, respectively, and $P_{1}^{\prime}$ and $P_{2}^{\prime}$ are real analytic Lagrangian ( $n-1$ )-disks in $\mathbb{C}^{n-1}=\left\{z_{1}=0\right\}$ intersecting transversely at 0 . Outside $N^{\prime} \times \mathbb{R}$ the isotopy is constant. See Figure 2. (The full definition of a standard self tangency move appears in Section 3. For simplicity, one technical condition there has been omitted at this point.)


Figure 2. Type B double point move.
Note that two Legendrian isotopic admissible Legendrian submanifolds of dimension 1 are in general not isotopic through an admissible Legendrian isotopy. In this case, one must allow also a "triple point move" see $[\mathbf{2}, \mathbf{1 1}]$.

To prove Proposition 2.7, we need to check that the differential graded algebra changes only by stable tame isomorphisms under Legendrian isotopies of type (A) and (B). We start with type (A) isotopies.

Lemma 2.8. Let $L_{t}, t \in[0,1]$ be a type (A) isotopy between generic admissible Legendrian submanifolds. Then the DGA's associated to $L_{0}$ and $L_{1}$ are tame isomorphic.

To prove this, we use a parameterized version of Proposition 2.2. If $L_{t}, t \in I=[0,1]$ is a type (A) isotopy, then the double points of $\Pi_{\mathbb{C}}\left(L_{t}\right)$ trace out continuous curves. Thus, when we refer to a Reeb chord $c$ of $L_{t^{\prime}}$ for some $t^{\prime} \in[0,1]$ this unambiguously specifies a Reeb chord for all $L_{t}$. For any $t$, we let $\mathcal{M}_{A}^{t}(a ; \mathbf{b})$ denote the moduli space $\mathcal{M}_{A}(a ; \mathbf{b})$ for $L_{t}$ and define

$$
\begin{equation*}
\mathcal{M}_{A}^{I}(a ; \mathbf{b})=\left\{(u, t) \mid u \in \mathcal{M}_{A}^{t}(a ; \mathbf{b})\right\} . \tag{2.8}
\end{equation*}
$$

As above "generic" refers to a member of a Baire subset, see Section 7.2 for a more precise formulation of this term for 1-parameter families.

Proposition 2.9. For a generic type (A) isotopy $L_{t}, t \in I=[0,1]$ the following holds. If $a, \mathbf{b}, A$ are such that $\mu(A)+|a|-|\mathbf{b}|=d \leq 1$, then the moduli space $\mathcal{M}_{A}^{I}(a ; \mathbf{b})$ is a transversely cut out d-manifold. If $X$ is the union of all these transversely cut out manifolds which are 0 -dimensional, then the components of $X$ are of the form $\mathcal{M}_{A_{j}}^{t_{j}}\left(a_{j}, \mathbf{b}_{j}\right)$, where $\mu\left(A_{j}\right)+\left|a_{j}\right|-\left|\mathbf{b}_{j}\right|=0$, for a finite number of distinct instances $t_{1}, \ldots, t_{r} \in[0,1]$. Furthermore, $t_{1}, \ldots, t_{r}$ are such that $\mathcal{M}_{B}^{t_{j}}(c ; \mathbf{d})$ is a transversely cut out 0 -manifold for every $c, \mathbf{d}, B$ with $\mu(B)+|c|-|\mathbf{d}|=1$.

Proposition 2.9 is proved in Section 7.9. At an instant $t=t_{j}$ in the above proposition, we say a handle slide occurs, and an element in $\mathcal{M}_{A_{j}}^{t_{j}}\left(a_{j}, \mathbf{b}_{j}\right)$ will be called a handle slide disk. (The term handle slide comes form the analogous situation in Morse theory.)

The proof of Lemma 2.8 is similar to that of Lemma 2.5. It uses the following compactness result.

Proposition 2.10. Any sequence $u_{\alpha}$ in $\mathcal{M}_{A}^{I}(a ; \mathbf{b})$ has a subsequence that converges to a broken holomorphic curve with the same properties as in Proposition 2.3.

The proof of this proposition is identical to that of Proposition 2.3, see Section 9.

We now prove Lemma 2.8 in two steps. First, consider type (A) isotopies without handle slides.

Lemma 2.11. Let $L_{t}, t \in[0,1]$ be a generic type (A) isotopy of Legendrian submanifolds for which no handle slides occur. Then, the boundary maps $\partial_{0}$ and $\partial_{1}$ on $\mathcal{A}=\mathcal{A}\left(L_{0}\right)=\mathcal{A}\left(L_{1}\right)$ satisfies $\partial_{0}=\partial_{1}$.

Proof. Proposition 2.10 implies that $\mathcal{M}_{A}^{I}(a ; \mathbf{B})$ is compact when its dimension is one. Since if a sequence in this space converged to a broken curve ( $u^{1}, \ldots, u^{N}$ ), then at least one $u^{j}$ would have negative formal dimension. This contradicts the assumptions that no handle slide occurs and that the type (A) isotopy is generic. Thus the corresponding 0 dimensional moduli spaces $\mathcal{M}_{A}^{0}$ and $\mathcal{M}_{A}^{1}$ used in the definitions of $\partial_{0}$ and $\partial_{1}$, respectively, form the boundary of a compact 1-manifold. Hence, their modulo 2 counts are equal. q.e.d.

In order to see what happens around a handle slide instant, we construct an auxiliary Legendrian submanifold of dimension one larger than $L$. The details of this can be found in Section 10 where the following is proved. Let $L_{t}, t \in[-\delta, \delta]$ and $\mathcal{M}_{A}^{0}(a ; \mathbf{b})$ be as $\mathcal{M}_{A}(a ; \mathbf{b})$ in Subsection 10.2. Let $\partial_{-}$denote the differential on $\mathcal{A}=\mathcal{A}\left(L_{-\delta}\right)$, and $\partial_{+}$the one on $\mathcal{A}=\mathcal{A}\left(L_{\delta}\right)$. For generators $c$ in $\mathcal{A}$, define

$$
\phi_{a}^{m}(c)= \begin{cases}c & \text { if } c \neq a \\ a+m A \mathbf{b} & \text { if } c=a\end{cases}
$$

where $m \in \mathbb{Z}_{2}$ and extend $\phi_{a}^{m}$ to a tame algebra automorphism of $\mathcal{A}$.
Lemma 2.12. There exists $m \in \mathbb{Z}_{2}$ so that the map $\phi_{a}^{m}: \mathcal{A} \rightarrow \mathcal{A}$ is a tame isomorphism from $\left(\mathcal{A}, \partial_{-}\right)$to $\left(\mathcal{A}, \partial_{+}\right)$.

This is Lemma 10.8.
Proof of Lemma 2.8. The lemma follows from Lemmas 2.11 and 2.12. q.e.d.

We consider elementary isotopies of type (B). Let $L_{t}, t \in I=[-\delta, \delta]$ be an isotopy of type (B) where two Reeb chords $\{a, b\}$ are born as $t$ passes through 0 . Let $o$ be the degenerate Reeb chord (double point) at $t=0$ and let $\mathcal{C}^{\prime}=\left\{a_{1}, \ldots, a_{l}, b_{1}, \ldots, b_{m}\right\}$ be the other Reeb chords. Again, we note that $c_{i} \in \mathcal{C}^{\prime}$ unambiguously defines a Reeb chord for all $L_{t}$ and $a$ and $b$ unambiguously define two Reeb chords for all $L_{t}$ when $t>0$. It is easy to see that (with the appropriate choice of capping paths) the grading on $a$ and $b$ differ by 1 , so let $|a|=j$ and $|b|=j-1$. Let $\left(\mathcal{A}_{-}, \partial_{-}\right)$and $\left(\mathcal{A}_{+}, \partial_{+}\right)$be the DGA's associated to $L_{-\delta}$ and $L_{\delta}$, respectively.

Lemma 2.13. The stabilized algebra $S_{j}\left(\mathcal{A}_{-}, \partial_{-}\right)$is tame isomorphic to $\left(\mathcal{A}_{+}, \partial_{+}\right)$.
Proof of Proposition 2.7 and 2.6. The first proposition follows from Lemmas 2.8 and 2.13 and implies in its turn the second. q.e.d.

We prove Lemma 2.13 in several steps below. Label the Reeb chords of $L_{t}$ so that

$$
\mathcal{Z}\left(b_{m}\right) \leq \ldots \leq \mathcal{Z}\left(b_{1}\right) \leq \mathcal{Z}(b)<\mathcal{Z}(a) \leq \mathcal{Z}\left(a_{1}\right) \leq \ldots \leq \mathcal{Z}\left(a_{l}\right),
$$

let $\mathcal{B}=\mathbb{Z}_{2}\left[H_{1}(L)\right]\left\langle b_{1}, \ldots, b_{m}\right\rangle$ and note that $\mathcal{B}$ is a subalgebra of both $\mathcal{A}_{-}$and $\mathcal{A}_{+}$. Then

Lemma 2.14. For $\delta>0$ small enough

$$
\partial_{+} a=b+v,
$$

where $v \in \mathcal{B}$.
Proof. Let $\mathbf{0} \in H_{1}(L)$ denote the zero element. In the model for the type (B) isotopy, there is an obvious disk in $\mathcal{M}_{\mathbf{0}}^{t}(a ; b)$ for $t>0$ small which is contained in the $z_{1}$-plane. We argue that this is the only point in the moduli space. We restrict attention to the neighborhood $N$ of $o^{*}$ that is biholomorphic to the origin in $\mathbb{C}^{n}$ as in the description of a type (B) move. Let $\pi_{i}: \mathbb{C}^{n} \rightarrow \mathbb{C}$ be the projection onto the $i^{t h}$ coordinate. If $u: D \rightarrow \mathbb{C}^{n}$ is a holomorphic map in $\mathcal{M}_{\mathbf{0}}^{t}(a ; b)$, then $\pi_{i} \circ u$ will either be constant or not. If $\pi_{i} \circ u$ is non-constant for $i>1$, then the image of $\pi_{1} \circ u$ intersected with $N$ has boundary on two transverse Lagrangian submanifolds. As such it will have a certain area $A_{i}$. Since $\mathcal{Z}(a)-\mathcal{Z}(b) \rightarrow 0$ as $t \rightarrow 0+$, we can choose $t$ small enough so that
$\mathcal{Z}(a)-\mathcal{Z}(b)<A_{i}$, for all $i>1$. Then, $\pi_{i} \circ u$ must be a point for all $i>1$ and for $i=1$, it can only be the obvious disk. Lemma 7.24 shows that $\mathcal{M}_{\mathbf{0}}^{t}(a ; b)$ is transversely cut out and thus contributes to $\partial_{+} a$. If $u \in \mathcal{M}_{A}^{t}(a ; b)$, where $A \neq \mathbf{0}$, then the image of $u$ must leave $N$. Thus, the above argument shows that $\mathcal{M}_{A}^{t}(a ; b)=\emptyset$ for $t$ small enough. Also, for $t>0$ sufficiently small $\mathcal{Z}(a)-\mathcal{Z}(b)<\mathcal{Z}\left(b_{m}\right)$. Hence, by Lemma 2.1, $v \in \mathcal{B}$.
q.e.d.

Define the elementary isomorphism $\Phi_{0}: \mathcal{A}_{+} \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$(on generators) by

$$
\Phi_{0}(c)= \begin{cases}e_{1}^{j} & \text { if } c=a \\ e_{2}^{j}+v & \text { if } c=b \\ c & \text { otherwise }\end{cases}
$$

The map $\Phi_{0}$ fails to be a tame isomorphism since it is not a chain map. However, we use it as the first step in an inductive construction of a tame isomorphism $\Phi_{l}: \mathcal{A}_{+} \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$. To this end, for $0 \leq i \leq l$, let $\mathcal{A}_{i}$ be the subalgebra of $\mathcal{A}_{+}$generated by $\left\{a_{1}, \ldots, a_{i}, a, b, b_{1}, \ldots, b_{m}\right\}$ (note that $\left.\mathcal{A}_{l}=\mathcal{A}_{+}\right)$. Then, with $\tau: S_{j}\left(\mathcal{A}_{-}\right) \rightarrow \mathcal{A}_{-}$denoting the natural projection and with $\partial_{-}^{s}$ denoting the differential induced on $S_{j}\left(\mathcal{A}_{-}\right)$, we have

Lemma 2.15.

$$
\begin{equation*}
\Phi_{0} \circ \partial_{+} w=\partial_{-}^{s} \circ \Phi_{0} w \tag{2.9}
\end{equation*}
$$

for $w \in \mathcal{A}_{0}$ and

$$
\begin{equation*}
\tau \circ \Phi_{0} \circ \partial_{+}=\tau \circ \partial_{-}^{s} \circ \Phi_{0} . \tag{2.10}
\end{equation*}
$$

Before proving this lemma, we show how to use it in the inductive construction which completes the proof of Lemma 2.13.
Proof of Lemma 2.13. The proof is similar to the proof of Lemmas 6.3 and 6.4 in [11] (cf. [2]). Define the map $H: S_{j}\left(\mathcal{A}_{-}\right) \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$on words $\mathbf{w}$ in the generators by

$$
H(\mathbf{w})= \begin{cases}0 & \text { if } \mathbf{w} \in \mathcal{A}_{-}, \\ 0 & \text { if } \mathbf{w}=\alpha e_{1}^{j} \beta \text { and } \alpha \in \mathcal{A}_{-} \\ \alpha e_{1}^{j} \beta & \text { if } \mathbf{w}=\alpha e_{2}^{j} \beta \text { and } \alpha \in \mathcal{A}_{-}\end{cases}
$$

and extend it linearly. Assume inductively that we have defined a graded isomorphism $\Phi_{i-1}: \mathcal{A}_{+} \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$so that it is a chain map when restricted to $\mathcal{A}_{i-1}$ and so that $\Phi_{i-1}\left(a_{k}\right)=a_{k}$, for $k>i-1$. (Note that $\Phi_{0}$ has these properties by Lemma 2.15.)

Define the elementary isomorphism $g_{i}: S_{j}\left(\mathcal{A}_{-}\right) \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$on generators by

$$
g_{i}(c)= \begin{cases}c & \text { if } c \neq a_{i} \\ a_{i}+H \circ \Phi_{i-1} \circ \partial_{+}\left(a_{i}\right) & \text { if } c=a_{i}\end{cases}
$$

and set $\Phi_{i}=g_{i} \circ \Phi_{i-1}$. Then $\Phi_{i}$ is a graded isomorphism. To see that $\Phi_{i}$ is a chain map when restricted to $\mathcal{A}_{i}$ observe the following facts: $\tau \circ H=0, \tau \circ g_{i}=\tau$, and $\tau \circ \Phi_{i}=\tau \circ \Phi_{0}$ for all $i$. Moreover, $\partial_{+} a_{i} \in \mathcal{A}_{i-1}$ and $\tau-\mathrm{id}_{S_{j}\left(\mathcal{A}_{-}\right)}=\partial_{-}^{s} \circ H+H \circ \partial_{-}^{s}$, where in the last equation, we think of $\tau: S_{j}\left(\mathcal{A}_{-}\right) \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$as $\tau: S_{j}\left(\mathcal{A}_{-}\right) \rightarrow \mathcal{A}_{-}$composed with the natural inclusion.

Using these facts, we compute

$$
\begin{aligned}
\partial_{-}^{s} g_{i}\left(a_{i}\right) & =\partial_{-}^{s}\left(a_{i}\right)+\left(\partial_{-}^{s} H\right) \Phi_{i-1} \partial_{+}\left(a_{i}\right) \\
& =\partial_{-}^{s}\left(a_{i}\right)+\left(H \partial_{-}^{s}+\tau+\mathrm{id}\right) \Phi_{i-1} \partial_{+}\left(a_{i}\right) \\
& =\partial_{-}^{s}\left(a_{i}\right)+\tau \Phi_{0} \partial_{+}\left(a_{i}\right)+\Phi_{i-1} \partial_{+}\left(a_{i}\right) \\
& =\Phi_{i-1} \partial_{+}\left(a_{i}\right) .
\end{aligned}
$$

Thus $\Phi_{i} \circ \partial_{+}\left(a_{i}\right)=\partial_{-}^{s} \circ g_{i}\left(a_{i}\right)=\partial_{-}^{s} \circ \Phi_{i}\left(a_{i}\right)$. Since $\Phi_{i}$ and $\Phi_{i-1}$ agree on $\mathcal{A}_{i-1}$ it follows that $\Phi_{i}$ is a chain map on $\mathcal{A}_{i}$. Continuing, we eventually get a tame chain isomorphism $\Phi_{l}: \mathcal{A}_{+} \rightarrow S_{j}\left(\mathcal{A}_{-}\right)$.
q.e.d.

The proof of Lemma 2.15 depends on the following two propositions.
Proposition 2.16. Let $L_{t}, t \in I=[-\delta, \delta]$ be a generic Legendrian isotopy of type (B) with notation as above (that is, o is the degenerate Reeb chord of $L_{0}$ and the Reeb chords $a$ and $b$ are born as $t$ increases past 0).

1) Let $\mathcal{M}_{A}^{0}(o, \mathbf{c})$ be a moduli space of rigid holomorphic disks. Then there exist $\rho>0$ and a local homeomorphism

$$
S: \mathcal{M}_{A}^{0}(o ; \mathbf{c}) \times[\rho, \infty) \rightarrow \mathcal{M}_{A}^{(0, \delta]}(a ; \mathbf{c})
$$

with the following property. If $u \in \mathcal{M}_{A}^{0}(o ; \mathbf{c})$, then any disk in $\mathcal{M}_{A}^{(0, \delta]}(a ; \mathbf{c})$ sufficiently close to the image of $u$ is in the image of $S$.
2) Let $\mathcal{M}_{A}^{0}(c, \mathbf{d})$ be a moduli space of rigid holomorphic disks. Let $S \subset\{1, \ldots, m\}$ be the subset of positions of $\mathbf{d}$ where the Reeb chord o appears (to avoid trivialities, assume $S \neq \emptyset$ ). Then there exists $\rho>0$ and a local homeomorphism

$$
S^{\prime}: \mathcal{M}_{A}^{0}(c, \mathbf{d}) \times[\rho, \infty) \rightarrow \mathcal{M}_{A}^{(0, \delta]}\left(c, \mathbf{d}_{S}(b)\right),
$$

with the following property. If $u \in \mathcal{M}_{0}(c ; \mathbf{d})$, then any disk in $\mathcal{M}_{A}^{(0, \delta]}\left(c ; \mathbf{d}_{S}(b)\right)$ sufficiently close to the image of $u$ is in the image of $S^{\prime}$.

This is a rephrasing of Theorem 8.2 and the following proposition is a restatement of Theorem 8.3.

Proposition 2.17. Let $L_{t}, t \in I=[-\delta, \delta]$ be a generic isotopy of type (B). Let $\mathcal{M}_{A_{1}}^{0}\left(o ; \mathbf{c}^{1}\right), \ldots, \mathcal{M}_{A_{r}}^{0}\left(o ; \mathbf{c}^{r}\right)$, and $\mathcal{M}_{B}^{0}(c ; \mathbf{d})$ be moduli spaces of rigid holomorphic disks. Let $S \subset\{1, \ldots, m\}$ be the subset of positions
in $\mathbf{d}$ where the Reeb chord o appears and assume that $S$ contains $r$ elements. Then there exists $\rho>0$ and an embedding
$G: \mathcal{M}_{B}^{0}(c ; \mathbf{d}) \times \Pi_{j=1}^{r} \mathcal{M}_{A_{j}}^{0}\left(o ; \mathbf{d}^{j}\right) \times[\rho, \infty) \rightarrow \mathcal{M}_{B+\sum A_{j}}^{[-\delta, 0)}\left(c ; \mathbf{d}_{S}\left(\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}\right)\right)$, with the following property. If $v \in \mathcal{M}_{0}(c ; \mathbf{d})$ and $u_{j} \in \mathcal{M}_{0}\left(o ; \mathbf{c}^{j}\right), j=$ $1, \ldots, r$ then any disk in $\mathcal{M}_{B+\sum A_{j}}^{[-\delta, 0)}\left(c ; \mathbf{d}_{S}\left(\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}\right)\right)$ sufficiently close to the image of $\left(v, u_{1}, \ldots, u_{r}\right)$ is in the image of $G$.

Proof of Lemma 2.15. Equation (2.9) follows from arguments similar to those in Lemma 2.8. Specifically, one can use these arguments to show that $\partial_{+} b_{i}=\partial_{-} b_{i}$. Then since $\partial_{+} b_{i} \in \mathcal{B}$ and since $\Phi_{0}$ is the identity on $\mathcal{B}$,

$$
\Phi_{0} \partial_{+} b_{i}=\partial_{+} b_{i}=\partial_{-} b_{i}=\partial_{-}^{s} \Phi_{0} b_{i} .
$$

We also compute

$$
\Phi_{0} \partial_{+} a=\Phi_{0}(b+v)=e_{2}^{j}+v+v=e_{2}^{j}=\partial_{-}^{s} \Phi_{0} a,
$$

and, since $\partial_{+} b$ and $\partial_{+} v$ both lie in $\mathcal{B}$,

$$
\Phi_{0} \partial_{+} b=\partial_{+} b, \quad \partial_{-}^{s} \Phi_{0} b=\partial_{-}^{s}\left(e_{1}^{j}+v\right)=\partial_{-} v=\partial_{+} v .
$$

Since $0=\partial_{+} \partial_{+} a=\partial_{+} b+\partial_{+} v$, we conclude that (2.9) holds.
To check (2.10), we write $\partial_{+} a_{i}=W_{1}+W_{2}+W_{3}$, where $W_{1}$ lies in the subalgebra generated by $\left\{a_{1}, \ldots, a_{l}, b_{1}, \ldots, b_{m}\right\}$, where $W_{2}$ lies in the ideal generated by $a$ and where $W_{3}$ lies in the ideal generated by $b$ in the subalgebra generated by $\left\{a_{1}, \ldots, a_{l}, b, b_{1}, \ldots, b_{m}\right\}$.

Let $u_{t}$, be a family of holomorphic disks with boundary on $L_{t}$. As $t \rightarrow 0, u_{t}$ converges to a broken disk $\left(u^{1}, \ldots, u^{N}\right)$ with boundary on $L_{0}$. This together with the genericity of the type (B) isotopy implies that for $t \neq 0$ small enough, there are no disks of negative formal dimension with boundary on $L_{t}$ since a broken curve which is a limit of a sequence of such disks would have at least one component $u^{j}$ with negative formal dimension.

Let $u_{s}: D \rightarrow \mathbb{C}^{n}, s \neq 0$ be rigid disks with boundary on $L_{s}$. If, the image $u_{-t}(\partial D)$ stays a positive distance away from $o^{*}$ as $t \rightarrow 0+$ then the argument above implies that $u_{-t}$ converges to a non-broken curve. Hence, $\partial_{-} a_{i}=W_{1}+W_{4}$ where for each rigid disk $u_{-t}: D \rightarrow \mathbb{C}^{n}$ contributing to a word in $W_{4}$, there exists points $q_{-t} \in \partial D$ such that $u_{-t}\left(q_{-t}\right) \rightarrow o^{*}$ as $t \rightarrow 0+$. The genericity assumption on the type (B) isotopy implies that no rigid disk with boundary on $L_{0}$ maps any boundary point to $o^{*}$, see Corollary 7.22. Hence, $u_{-t}$ must converge to a broken curve $\left(u^{1}, \ldots, u^{N}\right)$ which brakes at $o^{*}$. Moreover, by genericity and (2.6), every component $u^{j}$ of the broken curve must be a rigid disk with boundary on $L_{0}$. Proposition 2.17 shows that any such broken curve may be glued and Proposition 2.16 determines the pieces which
we may glue. It follows that $W_{4}=\hat{W}_{2}$ where $\hat{W}_{2}$ is obtained from $W_{2}$ by replacing each occurrence of $b$ with $v$. Therefore,

$$
\tau \Phi_{0} \partial_{+}\left(a_{i}\right)=\tau \Phi_{0}\left(W_{1}+W_{2}+W_{3}\right)=W_{1}+\hat{W}_{2}=\partial_{-}\left(a_{i}\right)=\tau \partial_{-}^{s} \Phi_{0}\left(a_{i}\right) .
$$

## 3. Admissible Legendrian submanifolds and isotopies

3.1. Chord genericity. Recall that a Legendrian submanifold $L \subset$ $\mathbb{R}^{2 n+1}$ is chord generic if all its Reeb chords correspond to transverse double points of the Lagrangian projection $\Pi_{\mathbb{C}}$. For a dense open set in the space of paths of Legendrian embeddings, the corresponding 1parameter families $L_{t}, 0 \leq t \leq 1$, are chord generic except for a finite number of parameter values $t_{1}, \ldots, t_{k}$ where $\Pi_{\mathbb{C}}\left(L_{t_{j}}\right)$ has one double point with self-tangency, and where for some $\delta>0 \Pi_{\mathbb{C}}\left(L_{t}\right),\left(t_{j}-\delta, t_{j}+\delta\right)$, is a versal deformation of $\Pi_{\mathbb{C}}\left(L_{t_{j}}\right)$, for $j=1, \ldots, k$. We call 1-parameter families $L_{t}$ with this property chord generic 1-parameter families.
3.2. Local real analyticity. For technical reasons, we require our Legendrian submanifolds to be real analytic in a neighborhood of the endpoints of their Reeb chords and that self-tangency instants in 1parameter families have a very special form.

Definition 3.1. A chord generic Legendrian submanifold $L \subset \mathbb{C}^{n} \times \mathbb{R}$ is admissible if for any Reeb chord $c$ of $L$ with endpoints $q_{1}$ and $q_{2}$, there are neighborhoods $U_{1} \subset L$ and $U_{2} \subset L$ of $q_{1}$ and $q_{2}$, respectively, such that $\Pi_{\mathbb{C}}\left(U_{1}\right)$ and $\Pi_{\mathbb{C}}\left(U_{2}\right)$ are real analytic submanifolds of $\mathbb{C}^{n}$.

We will require that self-tangency instants in 1-parameter families have the following special form. Consider $0 \in \mathbb{C}^{n}$ and coordinates $\left(z_{1}, \ldots, z_{n}\right)$ on $\mathbb{C}^{n}$. Let $P_{1}$ and $P_{2}$ be Lagrangian submanifolds of $\mathbb{C}^{n}$ passing through 0 . Let $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}$ and $y=\left(y_{1}, \ldots, y_{n}\right) \in \mathbb{R}^{n}$ be coordinates on $P_{1}$ and $P_{2}$, respectively. Let $R_{1} \subset P_{1}$ and $R_{2} \subset P_{2}$ be the boxes $\left|x_{j}\right| \leq 1$ and $\left|y_{j}\right| \leq 1, j=1, \ldots, n$. Let $B_{j}(2)$ and $B_{j}(2+\epsilon)$ for some small $\epsilon>0$ be the balls of radii 2 and $2+\epsilon$ around $0 \in P_{j}$, $j=1,2$. We require that in $R_{1}, P_{1}$ has the form

$$
\begin{equation*}
\gamma_{1} \times \hat{P}_{1} \tag{3.1}
\end{equation*}
$$

where $\gamma_{1}$ is an arc around 0 in the real line in the $z_{1}$-plane and where $\hat{P}_{1}$ is a Lagrangian submanifold of $\mathbb{C}^{n-1} \approx\left\{z_{1}=0\right\}$. We require that in $R_{2}, P_{2}$ has the form

$$
\begin{equation*}
\gamma_{2}(t) \times \hat{P}_{2}, \tag{3.2}
\end{equation*}
$$

where $\gamma_{2}$ is an arc around 0 in the unit-radius circle centered at $i$ in $z_{1-}$ plane and where $\hat{P}_{2}$ is a Lagrangian submanifold of $\mathbb{C}^{n-1} \approx\left\{z_{1}=0\right\}$, which meets $\hat{P}_{1}$ transversally at 0 .

If $q \in \mathbb{C}^{n}$, let $\lambda_{q}$ denote the complex line in $T_{q} \mathbb{C}^{n}$ parallel to the $z_{1}$-line. We also require that for every point $p \in B_{j}(2+\epsilon) \backslash B_{j}(2)$, the tangent plane $T_{p} P_{j}$ satisfies

$$
\begin{equation*}
T_{p} P_{j} \cap \lambda_{p}=0, \quad j=1,2 \tag{3.3}
\end{equation*}
$$

Definition 3.2. Let $L_{t}$ be a chord generic 1-parameter family of Legendrian submanifolds such that $L_{0}$ has a self-tangency. We say that the self-tangency instant $L_{0}$ is standard if there is some neighborhood $U$ of the self-tangency point and a biholomorphism $\phi: U \rightarrow V \subset \mathbb{C}^{n}$ such that

$$
\begin{equation*}
\phi\left(L_{t} \cap U\right)=P_{1} \cup P_{2}(t) \cap N, \tag{3.4}
\end{equation*}
$$

where $N$ is some neighborhood of $0 \in \mathbb{C}^{n}$, and where $P_{2}(t)$ is $P_{2}$ transalted $t$ units in the $y_{1}$-direction.

Definition 3.3. Let $L_{t}, 0 \leq t \leq 1$ be a chord generic 1-parameter family of Legendrian submanifolds. Let $t_{1}, \ldots, t_{k}$ be its self tangency instants. We say that $L_{t}$ is an admissible 1-parameter family if $L_{t}$ is admissible for all $t \neq t_{k}$, if there exists small disjoint intervals $\left(t_{j}-\delta, t_{j}+\delta\right)$ where the 1-parameter family is constant outside some small neighborhood $W$ of the self-tangency point, and if all self-tangency instants are standard.

Definition 3.4. A Legendrian submanifold $L \subset \mathbb{R} \times \mathbb{C}^{n}$ which is a self-tangency instant of an admissible 1-parameter family will be called semi-admissible.
3.3. Reducing the Legendrian isotopy problem. We prove a sequence of lemmas which reduce the classification of Legendrian submanifolds up to Legendrian isotopy to the classification of admissible Legendrian submanifolds up to admissible Legendrian isotopy.

We start with a general remark concerning lifts of Hamiltonian isotopies in $\mathbb{C}^{n}$. If $h$ is a smooth function with compact support in $\mathbb{C}^{n}$, then the Hamiltonian vector field

$$
X_{h}=-\frac{\partial h}{\partial y_{i}} \partial_{x_{i}}+\frac{\partial h}{\partial x_{i}} \partial_{y_{i}}
$$

associated to $h$ generates a 1-parameter family of diffeomorphisms $\Phi_{h}^{t}$ of $\mathbb{C}^{n}$. Moreover, the vector field $X_{h}$ lifts to a contact vector field

$$
\tilde{X}_{h}=-\frac{\partial h}{\partial y_{i}} \partial_{x_{i}}+\frac{\partial h}{\partial x_{i}} \partial_{y_{i}}+\left(h-y_{i} \frac{\partial h}{\partial y_{i}}\right) \partial_{z}
$$

on $\mathbb{C}^{n} \times \mathbb{R}$, which generates a 1-parameter family $\tilde{\Phi}_{h}^{t}$ of contact diffeomorphisms of $\mathbb{C}_{\tilde{n}}^{n} \times \mathbb{R}$ which is a lift of $\Phi_{h}^{t}$. We write $\Phi_{h}=\Phi_{h}^{1}$ and similarly $\tilde{\Phi}_{h}=\tilde{\Phi}_{h}^{1}$.

We note for future reference that in case the preimage of the support of $h$ in $L$ has more than one connected component, we may define a

Legendrian isotopy of $L$ by moving only one of these components (for a short time) using $\tilde{X}_{h}$ and keeping the rest of them fixed.

An $\epsilon$-isotopy is an isotopy during which no point moves a distance larger than $\epsilon>0$.

Lemma 3.5. Let $L$ be a Legendrian submanifold. Then, for any $\epsilon>$ 0 , there is an admissible Legendrian submanifold $L_{\epsilon}$ which is Legendrian $\epsilon$-isotopic to $L$.

Proof. As mentioned, we may after arbitrarily small Legendrian isotopy assume that $L$ is chord generic. Thus, it is enough to consider one transverse double point. We may assume that one of the sheets of $L$ close to this double point is given by $x \mapsto(x, d f(x), f(x))$ for some smooth function $f$. Let $g$ be a real analytic function approximating $f$ (e.g. its Taylor polynomial of some degree). Consider a Hamiltonian $h$ which is $h(x, y)=g(x)-f(x)$ in this neighborhood and 0 outside some slightly larger neighborhood. It is clear that the corresponding Hamiltonian vector field can be made arbitrarily small. Its flow map at time 1 is given by $\Phi_{h}^{1}(x, y)=(x, y+d g(x)-d f(x))$. Using this and suitable cut-off functions for the lifted Legendrian isotopies the lemma follows.
q.e.d.

Lemma 3.6. Let $L_{t}$ be any chord generic Legendrian isotopy from an admissible Legendrian submanifold $L_{0}$ to another one $L_{1}$. Then for any $\epsilon>0$, there is an admissible Legendrian isotopy $\epsilon$-close to $L_{t}$ connecting $L_{0}$ to $L_{1}$.

Proof. Let $t_{1}, \ldots, t_{k}$ be the self tangency instants of the isotopy. First, change the isotopy so that there exists small disjoint intervals $\left(t_{j}-\right.$ $\delta, t_{j}+\delta$ ) where the 1-parameter family is constant outside some small neighborhood $W$ of the self-tangency point. Consider the restriction of the isotopy to the self-tangency free regions. The 1-parametric version of the proof of Lemma 3.5 clearly applies to transform this part of the isotopy into one consisting of admissible Legendrian submanifolds. Then change the isotopy in the neighborhoods of the self tangency instants, using essentially the same argument as above, to a self-tangency of the from given in (3.1) and (3.2).

It remains to show how to fulfill the condition (3.3). To this end, consider a Lagrangian submanifold of the form (3.1). Locally it is given by $\left(x, d f(\hat{x})\right.$ ), where $\hat{x}=\left(x_{2}, \ldots, x_{n}\right)$. Let $\phi(x)$ be a function which equals 0 in $B\left(2-\frac{1}{2} \epsilon\right)$ and outside $B(2+2 \epsilon)$ and has $\frac{\partial^{2} \phi}{\partial x_{1} x_{j}} \neq 0$ for some $j$ at all points in $B(2+\epsilon) \backslash B(2)$. (For example, if $K$ is a small constant a suitable cut-off of the function $K x_{1}\left(x_{2}+\ldots x_{n}\right)$ has this property). We see as above that our original Legendrian is Legendrian isotopic to $(x, d f(x)+d \phi(x))$. The tangent space of the latter submanifold is
spanned by the vectors

$$
\begin{align*}
& \partial_{x_{1}}+\sum_{j} \frac{\partial^{2} \phi}{\partial x_{j} \partial x_{1}} \partial_{y_{j}},  \tag{3.5}\\
& \partial_{x_{r}}+\sum_{j} \frac{\partial^{2} \phi}{\partial x_{j} \partial x_{r}} \partial_{y_{j}}+\frac{\partial^{2} f}{\partial x_{j} \partial x_{r}} \partial_{y_{j}}, \quad 2 \leq r \leq n . \tag{3.6}
\end{align*}
$$

Any non-trivial linear combination of the last $n-1$ vectors projects non-trivially to the subspace $d x_{1}=d y_{1}=d y_{2}=\cdots=d y_{n}=0$. The first vector lies in the subspace $d x_{2}=\cdots=d x_{n}=0$; thus, since the first vector does not lie in the $z_{1}$-line because $\frac{\partial^{2} \phi}{\partial x_{1} x_{j}} \neq 0$ for some $j \neq 1$, no linear combination of the vectors does either. $P_{2}$ can be deformed in a similar manner.

After the self-tangency moment is passed, it is easy to Legendrian isotope back to the original family through admissible Legendrian submanifolds.
q.e.d.

## 4. Holomorphic disks

In this section, we establish notation and ideas that will be used throughout the rest of the paper.
4.1. Reeb chord notation. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a Legendrian submanifold and let $c$ be a Reeb chord of $L$. The $z$-coordinate of the upper and lower end points of $c$ will be denoted by $c^{+}$and $c^{-}$, respectively. See Figure 3. So as a point set $c=c^{*} \times\left[c^{-}, c^{+}\right]$and the action of $c$ is simply $\mathcal{Z}(c)=c^{+}-c^{-}$.


Figure 3. A Reeb chord in $\mathbb{R}^{3}$.
If $r>0$ is small enough so that $\Pi_{\mathbb{C}}^{-1}\left(B\left(c^{*}, r\right)\right)$ intersects $L$ is exactly two disk about the upper and lower end points of $c$, then we define $U\left(c^{ \pm}, r\right)$ to be the component of $\Pi_{\mathbb{C}}^{-1}\left(B\left(c^{*}, r\right)\right) \cap L$ containing $c^{*} \times c^{ \pm}$.
4.2. Definition of holomorphic disks. If $M$ is a smooth manifold, then let $\mathcal{H}_{k}^{\text {loc }}\left(M, \mathbb{C}^{n}\right)$ denote the Frechet space of all functions which agree locally with a function with $k$ derivatives in $L^{2}$. Let $\Delta_{m} \subset \mathbb{C}$ denote the unit disk with $m$ punctures on the boundary, let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a (semi-)admissible Legendrian submanifold.

Definition 4.1. A holomorphic disk with boundary on $L$ consists of two functions $u \in \mathcal{H}_{2}^{\text {loc }}\left(\Delta_{m}, \mathbb{C}^{n}\right)$ and $h \in \mathcal{H}_{\frac{3}{2}}\left(\partial \Delta_{m}, \mathbb{R}\right)$ such that

$$
\begin{align*}
\bar{\partial} u(\zeta) & =0, \text { for } \zeta \in \operatorname{int}\left(\Delta_{m}\right),  \tag{4.1}\\
(u(\zeta), h(\zeta)) & \in L, \text { for } \zeta \in \partial \Delta_{m}, \tag{4.2}
\end{align*}
$$

and such that for every puncture $p$ on $\partial \Delta_{m}$, there exists a Reeb chord $c$ of $L$ such that

$$
\begin{equation*}
\lim _{\zeta \rightarrow p} u(\zeta)=c^{*} \tag{4.3}
\end{equation*}
$$

When (4.3) holds, we say that ( $u, h$ ) maps the puncture $p$ to the Reeb chord $c$.

Remark 4.2. Since $u \in \mathcal{H}_{2}^{\text {loc }}\left(\Delta_{m}, \mathbb{C}^{n}\right)$, the restriction of $u$ to the boundary lies in $\mathcal{H}_{\frac{3}{2}}^{\text {loc }}\left(\partial \Delta_{m}, \mathbb{C}^{n}\right)$. Therefore, both $u$ and its restriction to the boundary are continuous. Hence, (4.2) and (4.3) make sense.

Remark 4.3. If $u \in \mathcal{H}_{2}^{\text {loc }}\left(\Delta_{m}, \mathbb{C}^{n}\right)$, then $\bar{\partial} u \in \mathcal{H}_{1}^{\text {loc }}\left(\Delta_{m}, T^{* 0,1} D_{m} \otimes\right.$ $\mathbb{C}^{n}$ ) and hence, the trace of $\bar{\partial} u$ (its restriction to the boundary $\partial \Delta_{m}$ ) lies in $\mathcal{H}_{\frac{1}{2}}^{\text {loc }}\left(\partial \Delta_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$. If $u$ is a holomorphic disk, then $\bar{\partial} u=0$ and hence, its trace $\bar{\partial} u \mid \partial \Delta_{m}$ is also 0 .

Remark 4.4. It turns out, see Section 9.5, that if $(u, f)$ is a holomorphic disk with boundary on a smooth $L$, then the function $u$ is in fact smooth up to and including the boundary and thus $f$ is also smooth. Hence, it is possible to rephrase Definition 4.1 in terms of smooth functions. (Also, it follows that the definition above agrees with that given in Section 2.3.) The advantage of the present definition is that it allows for Legendrian submanifolds of lower regularity. (The boundary condition makes sense for Legendrian submanifolds $L$ which are merely $C^{1}$.)
4.3. Conformal structures. We describe the space of conformal structures on $\Delta_{m}$ as follows. If $m \leq 3$, then the conformal structure is unique. Let $m>4$ and let the punctures of $\Delta_{m}$ be $p_{1}, \ldots, p_{m}$. Then fixing the positions of the punctures $p_{1}, p_{2}, p_{3}$, the conformal structure on $\Delta_{m}$ is determined by the position of the remaining $m-3$ punctures. In this way, we identify the space of conformal structures $\mathcal{C}_{m}$ on $\Delta_{m}$ with an open simplex of dimension $m-3$.
4.4. A family of metrics. Let $\Delta$ denote the unit disk in the complex plane. Consider $\Delta_{m}$ with $m$ punctures $p_{1}, \ldots, p_{m}$ on the boundary and conformal structure $\kappa$. Let $d$ be the smallest distance along $\partial \Delta$ between two punctures and take

$$
\delta=\min \left\{\frac{d}{100}, \frac{\pi}{100}\right\}
$$

Define $D(p, \delta)$ to be a disk such that $\partial \Delta(p, \delta)$ intersects $\partial \Delta$ orthogonally at two points $a_{+}$and $a_{-}$of distance $\delta$ (in $\partial \Delta$ ) from $p$.

Let $L_{p}$ be the oriented tangent-line of $\partial \Delta$ at $p$ and let $g_{p}$ be the unique Möbius transformation which fixes $p$, maps $a_{+}$to the point of distance $\delta$ from $p$ along $L_{p}$, maps $a_{-}$to the point of distance $-\delta$ from $p$ along $L_{p}$, and such that the image of $g_{p}(\Delta)$ intersects the component of $\mathbb{C}-L_{p}$ which intersects $\Delta$.

The function $h_{p}: D(p, \delta) \cap \Delta_{m} \rightarrow[0, \infty) \times[0,1]$ defined by

$$
h_{p}(\zeta)=-\frac{1}{\pi}\left(\log \left(-i \bar{p}\left(g_{p}(\zeta)-p\right)\right)-\log (\delta)\right)
$$

is a conformal equivalence. Let $g_{0}$ denote the Euclidean metric on $\mathbb{C}$. Then there exists a function $s:\left[0, \frac{1}{2}\right] \times[0,1] \rightarrow \mathbb{R}$ such that $h_{p}^{-1^{*}} g_{0}=$ $s(\zeta) g_{0}$ on $\left[0, \frac{1}{2}\right] \times[0,1]$. Let $\phi:[0, \infty) \rightarrow[0,1]$ be a smooth function which is 0 in a neighborhood of 0 and 1 in a neighborhood of $\frac{1}{2}$ for $\tau>\frac{1}{2}$. Let $g_{p}$ be the metric

$$
g_{p}(\tau+i t)=(\phi(\tau)+(1-\phi(\tau)) s(t+i t)) g_{0}
$$

on $[0, \infty) \times[0,1]$.
Now, consider $\Delta_{m}$ with the metric $g(\kappa)$ which agrees with $h_{p_{j}}^{*} g_{p_{j}}$ on $h_{p_{j}}^{-1}\left(\left[\frac{1}{2 \pi}, \infty\right) \times[0,1]\right)$ for each puncture $p_{j}$, and with $g_{0}$ on $\Delta_{m}-$ $\left(D\left(p_{1}, \delta\right) \cup \ldots, D\left(p_{m}, \delta\right)\right)$. Then $\left(\Delta_{m}, g(\kappa)\right)$ is conformally equivalent to $\left(\Delta_{m}, g_{0}\right)$.

We denote by $D_{m}(\kappa)$ the disk $\Delta_{m}$ with the metric $g(\kappa)$. If the specific $\kappa$ is unimportant or clear from context, we will simply write $D_{m}$. Also $E_{p_{j}} \subset D_{m}$ will denote the Euclidean neighborhood $[1, \infty) \times[0,1]$ of the $j^{\text {th }}$ puncture $p_{j}$ of $D_{m}$. We use coordinates $\zeta=\tau+i t$ on $E_{p_{j}}$ and let $E_{p_{j}}[M]$ denote the subset of $\tau+i t \in E_{p_{j}}$ with $|\tau| \geq M$.
4.5. Sobolev spaces. Consider $D_{m}$ with metric $g(\kappa)$ for some $\kappa \in \mathcal{C}_{m}$. Let $\hat{D}_{m}$ denote the open Riemannian manifold which is obtained by adding an open collar to $\partial D_{m}$ and extending the metric in a smooth and bounded way to $\hat{g}(c)$.

The Sobolev spaces $\mathcal{H}_{k}^{\text {loc }}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ are now defined in the standard way as the space of $\mathbb{C}^{n}$-valued functions (distributions) the restrictions of which to any open ball $B$ in any relatively compact coordinate chart $\approx \mathbb{R}^{2}$ lies in the usual Sobolev space $\mathcal{H}_{k}\left(B, \mathbb{C}^{n}\right)$.

Using the metric $\hat{g}(\kappa)$ and the finite cover

$$
\bigcup_{j} \operatorname{int}\left(\hat{E}_{p_{j}}[1]\right) \cup\left(\hat{D}_{m}-\cup_{j} \hat{E}_{p_{j}}[2]\right),
$$

where $\hat{E}_{p_{j}}$ is the union of $E_{p_{j}}$ and the corresponding part of the collar of $\hat{D}_{m}$, we define, for each integer $k$, the space $\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ as the subspace of all $f \in \mathcal{H}_{k}^{\text {loc }}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ with $\|f\|_{k}<\infty$.

We consider $\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ as a space of distributions acting on $C_{0}^{\infty}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$. We write

- $\mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$ for the space of restrictions to $\operatorname{int}\left(D_{m}\right) \subset \hat{D}_{m}$ of elements in $\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$, and
- $\dot{\mathcal{H}}_{k}\left(A, \mathbb{C}^{n}\right)$ for the set of distributions in $\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ supported in $A \subset \hat{D}_{m}$.
Then $\dot{\mathcal{H}}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$ is a closed subspace of $\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ and if $K_{m}=$ $\hat{D}_{m}-\operatorname{int}\left(D_{m}\right)$, then

$$
\mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right)=\mathcal{H}_{k}\left(\hat{D}_{m}, \mathbb{C}^{n}\right) / \dot{\mathcal{H}}_{k}\left(K_{m}, \mathbb{C}^{n}\right)
$$

We endow $\mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$ and $\dot{\mathcal{H}}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$ with the quotient- and induced-topology, respectively. Let $C_{0}^{\infty}\left(D_{m}, \mathbb{C}^{n}\right)$ denote the space of restrictions of elements in $C_{0}^{\infty}\left(\hat{D}_{m}, \mathbb{C}^{n}\right)$ to $D_{m}$.

Lemma 4.5. $C_{0}^{\infty}\left(D_{m}, \mathbb{C}^{n}\right)$ is dense in $\mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right), C_{0}^{\infty}\left(\operatorname{int}\left(D_{m}\right), \mathbb{C}^{n}\right)$ is dense in $\dot{\mathcal{H}}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$, and the spaces $\mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right)$ and $\dot{\mathcal{H}}_{-k}\left(D_{m}, \mathbb{C}^{n}\right)$ are dual with respect to the extension of the bilinear form

$$
\int_{D_{m}}\langle u, v\rangle d A
$$

where $u \in C_{0}^{\infty}\left(D_{m}, \mathbb{C}^{n}\right), v \in C_{0}^{\infty}\left(\operatorname{int}\left(D_{m}\right), \mathbb{C}^{n}\right)$ and $\langle$,$\rangle denotes the$ standard Riemannian inner product on $\mathbb{C}^{n} \approx \mathbb{R}^{2 n}$.

This is essentially Theorem B. 2.1 p. 479 in [20].
We will also use weighted Sobolev spaces: for $a \in \mathbb{R}$, let $e_{a}^{j}: D_{m} \rightarrow \mathbb{R}$ be a smooth function such that $e_{a}^{j}(\tau+i t)=e^{a \tau}$ for $\tau+i t \in E_{p_{j}}[3]$ and $e_{a}(\zeta)=1$ for $\zeta \in D_{m}-E_{p_{j}}[2]$. For $\mu=\left(\mu_{1}, \ldots, \mu_{m}\right) \in \mathbb{R}^{m}$, let $\mathbf{e}_{\mu}: D_{m} \rightarrow \mathbb{R} \otimes \mathrm{id} \subset \mathbf{G L}\left(\mathbb{C}^{n}\right)$ be

$$
\mathbf{e}_{\mu}(\zeta)=\Pi_{j=1}^{m} e_{\mu_{j}}^{j}(\zeta) \mathrm{id}
$$

Note that $\mathbf{e}_{\mu}(\zeta)$ preserves Lagrangian subspaces. We can now define $\mathcal{H}_{k, \mu}\left(D_{m}, \mathbb{C}^{n}\right)=\left\{u \in \mathcal{H}_{k}^{\text {loc }}\left(D_{m}, \mathbb{C}^{n}\right): \mathbf{e}_{\mu} u \in \mathcal{H}_{k}\left(D_{m}, \mathbb{C}^{n}\right)\right\}$, with norm $\|u\|_{k, \mu}=\left\|\mathbf{e}_{\mu} u\right\|_{k}$.
4.6. Asymptotics. Let $\Lambda_{0}$ and $\Lambda_{1}$ be (ordered) Lagrangian subspaces of $\mathbb{C}^{n}$. Define the complex angle $\theta\left(\Lambda_{0}, \Lambda_{1}\right) \in[0, \pi)^{n}$ inductively as follows:

If $\operatorname{dim}\left(\Lambda_{0} \cap \Lambda_{1}\right)=r \geq 0$, let $\theta_{1}=\cdots=\theta_{r}=0$ and let $\mathbb{C}^{n-r}$ denote the Hermitian complement of $\mathbb{C} \otimes \Lambda_{0} \cap \Lambda_{1}$ and let $\Lambda_{i}^{\prime}=\Lambda_{i} \cap \mathbb{C}^{n-r}$ for $i=0,1$. If $\operatorname{dim}\left(\Lambda_{0} \cap \Lambda_{1}\right)=0$, then let $\Lambda_{i}^{\prime}=\Lambda_{i}, i=0,1$ and let $r=0$. Then $\Lambda_{0}^{\prime}$ and $\Lambda_{1}^{\prime}$ are Lagrangian subspaces. Let $\alpha$ be smallest angle such that $\operatorname{dim}\left(e^{i \alpha} \Lambda_{0} \cap \Lambda_{1}\right)=r^{\prime}>0$. Let $\theta_{r+1}=\cdots=\theta_{r+r^{\prime}}=\alpha$. Now repeat the construction until $\theta_{n}$ has been defined. Note that $\theta\left(A \Lambda_{0}, A \Lambda_{1}\right)=$ $\theta\left(\Lambda_{0}, \Lambda_{1}\right)$ for every $A \in \mathbf{U}(n)$ since multiplication with $e^{i \alpha}$ commutes with everything in $\mathbf{U}(n)$.

Proposition 4.6. Let $(u, h)$ be a holomorphic disk with boundary on a (semi-) admissible Legendrian submanifold L. Let $p$ be a puncture on $D_{m}$ such that $p$ maps to the Reeb chord c. For $M>0$ sufficiently large, the following is true:

If $\Pi_{\mathbb{C}}(L)$ self-intersects transversely at $c^{*}$, then

$$
\begin{equation*}
|u(\tau+i t)|=\mathcal{O}\left(e^{-\theta \tau}\right), \quad \tau+i t \in E_{p}[M] \tag{4.4}
\end{equation*}
$$

where $\theta>0$ is the smallest complex angle of $c$.
If $\Pi_{\mathbb{C}}(L)$ has a self-tangency at $c^{*}$, then either there exists a real number $c_{0}$ such that

$$
\begin{equation*}
u(\tau+i t)=\left(\frac{ \pm 2}{c_{0}+\tau+i t}, 0, \ldots, 0\right)+\mathcal{O}\left(e^{-\theta \tau}\right) \quad \tau+i t \in E_{p}[M] \tag{4.5}
\end{equation*}
$$

or

$$
\begin{equation*}
|u(\tau+i t)|=\mathcal{O}\left(e^{-\theta \tau}\right), \quad \tau+i t \in E_{p}[M], \tag{4.6}
\end{equation*}
$$

where $\theta$ is the smallest non-zero complex angle of $L$ at $p$.
In particular, if the punctures $p_{1}, \ldots, p_{m}$ on $D_{m}$ map to Reeb chords $c_{1}, \ldots, c_{m}$ and if $f: D_{m} \rightarrow \mathbb{C}^{n}$ is any smooth function which is constantly equal to $c_{1}^{*}, \ldots, c_{m}^{*}$ in neighborhoods of $p_{1}, \ldots, p_{m}$, then $u-f \in$ $\mathcal{H}_{2}\left(D_{m}, \mathbb{C}^{n}\right)$.

Proof. Similar statements appear in [13] and [14]. (See also Theorem B in [27] for (4.4).) To see that (4.5) holds, we may assume that the selftangency point is $0 \in \mathbb{C}^{n}$ and that around $0, \Pi_{\mathbb{C}}(L)$ agrees with the local model in Definition 3.3. Elementary complex analysis (see Lemma 6.2 below) shows that for a standard self tangency the first component $u_{1}$ of a holomorphic disk is given by

$$
\begin{equation*}
u_{1}(\zeta)=\frac{ \pm 2}{\zeta-c_{0}+\sum_{n \in \mathbb{Z}} c_{n} \exp (n \pi \zeta)}, \tag{4.7}
\end{equation*}
$$

where $c_{j}$ are real constants, in $E_{p}[M]$. The remaining components $u^{\prime}$ of $u$ are controlled as above and the claim follows. The last statement follows immediately from the asymptotics at punctures.
q.e.d.

## 5. Functional analytic setup

As explained in Section 2, contact homology is built using modulispaces of holomorphic disks. In this section, we construct Banach manifolds of maps of punctured disks into $\mathbb{C}^{n}$ which satisfy certain boundary conditions. In this setting, moduli-spaces will appear as the zero-sets of bundle maps.

In Section 5.1, we define our Banach manifolds as submanifolds in a natural bundle of Banach spaces. To find atlases for our Banach manifolds, we proceed in the standard way: construct an "exponential map" from the proposed tangent space and show it is a diffeomorphism near the origin. To do this, in Section 5.2, we turn our attention to a special metric on the tangent bundle of the Legendrian submanifold. From this, we construct a family of metrics on $\mathbb{C}^{n}$ in Section 5.3 and use it to define a preliminary version of the "exponential map" for the Banach manifold. Section 5.4 contains some technical results needed to deal with families of Legendrian submanifolds. In Section 5.5, we show how to construct the atlas. Section 5.6 discusses how to invoke variations of the conformal structure of the source space into the present setup. In Section 5.7, we linearize the bundle map, the zero set of which is the moduli-space. Section 5.8 discusses some issues involving the semi-admissible case.
5.1. Bundles of affine Banach spaces. Let $L_{\lambda} \subset \mathbb{C}^{n} \times \mathbb{R}, \lambda \in \Lambda$, where $\Lambda$ is an open subset of a Banach space, denote a smooth family of chord generic admissible Legendrian submanifolds. That is, $\Lambda$ is smoothly mapped into the space of admissible Legendrian embeddings of $L$ endowed with the $C^{\infty}$-topology.

We also study the semi-admissible case. To this end, we also let $L_{\lambda}$, $\lambda \in \Lambda$, be a smooth family of semi-admissible Legendrian submanifolds. For simplicity, and since it will suffice for our applications, we assume that in this case, the self tangency point of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ remain fixed as $\lambda$ varies and that in a neighborhood of this point, the product structure of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ is preserved and the first components $\gamma_{1}$ and $\gamma_{2}$, shown in Figure 2 remain fixed as $\lambda$ varies.

Let $\mathbf{a}(\lambda)=\left(a_{1}(\lambda), \ldots, a_{m}(\lambda)\right), \lambda \in \Lambda$ be an ordered collection of Reeb chords of $L_{\lambda}$ depending continuously on $\lambda$. Consider $D_{m}$ with punctures $p_{1}, \ldots, p_{m}$, and a conformal structure $\kappa \in \mathcal{C}_{m}$.

Fix families, smoothly depending on $(\lambda, \kappa) \in \Lambda \times \mathcal{C}_{m}$, of smooth reference functions

$$
u_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa]: D_{m} \rightarrow \mathbb{C}^{n}
$$

such that $u_{\text {ref }}[\mathbf{a}(\lambda), \kappa]$ is constantly equal to $a_{k}^{*}$ in $E_{p_{k}}$, and

$$
h_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa]: \partial D_{m} \rightarrow \mathbb{R}
$$

such that $h_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa]$ is constantly equal to $a_{1}^{-}(\lambda)$ and $a_{1}^{+}(\lambda)$ on $[1, \infty) \subset$ $E_{p_{1}}$, and $[1, \infty)+i \subset E_{p_{1}}$, respectively, and, for $k \geq 2$, constantly equal to $a_{k}^{+}(\lambda)$ and $a_{k}^{-}(\lambda)$ on $[1, \infty) \subset E_{p_{k}}$, and $[1, \infty)+i \subset E_{p_{k}}$, respectively.

Let $\epsilon=\left(\epsilon_{1}, \ldots, \epsilon_{m}\right) \in[0, \infty)^{m}$. For $u: D_{m} \rightarrow \mathbb{C}^{n}$ and $h: \partial D_{m} \rightarrow \mathbb{R}$ consider the conditions

$$
\begin{align*}
& u-u_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa] \in \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right)  \tag{5.1}\\
& h-h_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa] \in \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m+1}, \mathbb{R}\right) \tag{5.2}
\end{align*}
$$

(Note that the $\kappa$-dependence of the right-hand sides in (5.1) and (5.2) has been dropped from the notation.) Define the affine Banach space

$$
\begin{aligned}
& \mathcal{F}_{2, \epsilon}(\mathbf{a}(\lambda), \kappa) \\
& \quad=\left\{(u, h): D_{m} \rightarrow \mathbb{C}^{n} \times \mathbb{R}: u \text { satisfies (5.1), } h \text { satisfies (5.2) }\right\}
\end{aligned}
$$

endowed with the norm which is the sum of the norms of the components. Let

$$
\mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)=\bigcup_{\lambda \in \Lambda} \mathcal{F}_{2, \epsilon}(\mathbf{a}(\lambda), \kappa)
$$

be the metric space with distance function

$$
\begin{align*}
d((v, f, \lambda),(w \cdot g, \mu))= & \left\|\left(v-u_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa]\right)-\left(w-u_{\mathrm{ref}}[\mathbf{a}(\mu), \kappa]\right)\right\|_{2, \epsilon}  \tag{5.3}\\
& +\left\|\left(f-h_{\mathrm{ref}}[\mathbf{a}(\lambda), \kappa]\right)-\left(g-h_{\mathrm{ref}}[\mathbf{a}(\mu), \kappa]\right)\right\|_{\frac{3}{2}}, \epsilon \\
& +|\lambda-\mu| .
\end{align*}
$$

We give $\mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$ the structure of a Banach manifold by producing an atlas as follows. Let $(w, f, \lambda) \in \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$. Let $\left(w_{\mu}, f_{\mu}, \mu\right)$ be any family such that $\left(w_{\lambda}, f_{\lambda}, \lambda\right)=(w, f, \lambda)$ and such that

$$
\mu \mapsto\left(w_{\mu}-u_{\mathrm{ref}}[\mathbf{a}(\mu), \kappa], f_{\mu}-h_{\mathrm{ref}}[\mathbf{a}(\mu), \kappa]\right)
$$

is a smooth map into $\mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right) \times \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m}, \mathbb{R}\right)$. Then a chart is given by

$$
\begin{align*}
& \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right) \times \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m}, \mathbb{R}\right) \times \Lambda \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa) ;  \tag{5.4}\\
& (g, r, \mu) \mapsto\left(w_{\mu}+g, f_{\mu}+r, \mu\right)
\end{align*}
$$

If $(u, h, \lambda) \in \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}(\lambda), \kappa)$ then $\bar{\partial} u \in \mathcal{H}_{1, \epsilon}\left(D_{m}, T^{* 0,1} D^{m} \otimes \mathbb{C}^{n}\right)$ and its trace $\bar{\partial} u \mid \partial D_{m}$ lies in $\mathcal{H}_{\frac{1}{2}}\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ ).

Definition 5.1. Let $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa) \subset \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$ denote the subset of elements $(u, h, \lambda)$ which satisfy

$$
\begin{align*}
& (u, h)(\zeta) \in L_{\lambda} \text { for all } \zeta \in \partial D_{m}  \tag{5.5}\\
& \int_{\partial D_{m}}\langle\bar{\partial} u, v\rangle d s=0, \text { for every } v \in C_{0}^{\infty}\left(\partial D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right) \tag{5.6}
\end{align*}
$$

where $\langle$,$\rangle denotes the inner product on T^{* 0,1} \otimes \mathbb{C}^{n}$ induced from the standard (Riemannian) inner product on $\mathbb{C}^{n}$.

Lemma 5.2. $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$ is a closed subset.
Proof. If ( $u_{k}, h_{k}, \lambda_{k}$ ) is a sequence in $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$ which converges in $\mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$, then $\lambda_{k} \rightarrow \lambda$ and the sequence $\left(u_{k} \mid \partial D_{m}, h_{k}\right)$ converges in sup-norm. Hence, (5.5) is a closed condition. Also, $\bar{\partial}$ is continuous as is the trace map. It follows that (5.6) is a closed condition as well. q.e.d.
5.2. The normal bundle of a Lagrangian immersion with a special metric. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an instant of a chord generic 1-parameter family of Legendrian submanifolds. Then $\Pi_{\mathbb{C}}: L \rightarrow \mathbb{C}^{n}$ is a Lagrangian immersion and the normal bundle of $\Pi_{\mathbb{C}}$ is isomorphic to the tangent bundle $T L$ of $L$. On the restriction $T_{L}(T L)$ of the tangent bundle $T(T L)$ of $T L$ to the zero-section $L$ there is a natural endomorphism $J: T_{L}(T L) \rightarrow T_{L}(T L)$ such that $J^{2}=-1$. It is defined as follows. If $p \in L$, then $T_{(p, 0)}(T L)$ is a direct sum of the space of horizontal vectors tangent to $L$ at $p$ and the space of vertical vectors tangent to the fiber of $\pi: T L \rightarrow L$ at $p$. If $v \in T_{L}(T L)$ is tangent to $L$ at $p \in L$, then $J v$ is the vector $v$ viewed as a tangent vector to the fiber $T_{p} L$ of $\pi: T L \rightarrow L$ at $(p, 0)$, and if $w$ is a vector tangent to the fiber of $\pi$ at $(p, 0)$ then $J w=-w$, where $-w$ is viewed as a tangent vector in $T_{p} L$. This defines $J$ on the two direct summands. Extend it linearly.

The immersion $\Pi_{\mathbb{C}}: L \rightarrow \mathbb{C}^{n}$ extends to an immersion $P$ of a neighborhood of the zero-section in $T L$ and $P$ can be chosen so that along $L, i \circ d P=d P \circ J$.

From a Riemannian metric $g$ on $L$, we construct a metric $\hat{g}$ on a neighborhood of the zero-section in $T L$ in the following way. Let $v \in T L$ with $\pi(v)=p$. Let $X$ be a tangent vector of $T L$ at $v$. The Levi-Civita connection of $g$ gives the decomposition $X=X^{H}+X^{V}$, where $X^{V}$ is a vertical vector, tangent to the fiber, and $X^{H}$ lies in the horizontal subspace at $v$ determined by the connection. Thus $X^{V}$ is a vector in $T_{p} L$ with its endpoint at $v$. It can be translated linearly to the origin $0 \in T_{p} L$. We use the same symbol $X^{V}$ to denote this vector translated to $0 \in T_{p} L$. Write $\pi X \in T_{p} L$ for the image of $X$ under the differential of the projection $\pi$ and let $R$ denote the curvature tensor of $g$.

Define the field of quadratic forms $\hat{g}$ on $T L$ as

$$
\begin{equation*}
\hat{g}(v)(X, Y)=g(p)(\pi X, \pi Y)+g(p)\left(X^{V}, Y^{V}\right)+g(p)(R(\pi X, v) \pi Y, v), \tag{5.7}
\end{equation*}
$$

where $v \in T L, \pi(v)=p$, and $X, Y \in T_{v}(T L)$.
Proposition 5.3. There exists $\rho>0$ such that $\hat{g}$ is a Riemannian metric on

$$
\{v \in T L: g(v, v)<\rho\} .
$$

In this metric, the zero section $L$ is totally geodesic and the geodesics in $L$ are exactly those in the metric $g$. Moreover, if $\gamma$ is a geodesic in $L$ and $X$ is a vector field in $T(T L)$ along $\gamma$, then $X$ satisfies the Jacobi equation if and only if $J X$ does.

Proof. Since $g(R(\pi X, v) \pi Y, v)=g(R(\pi Y, v) \pi X, v), \hat{g}$ is symmetric. When restricted to the 0 -section, $\hat{g}$ is non-degenerate. The first statement follows from the compactness of $L$.

In Lemmas 5.5 and 5.6 below, we show $L$ is totally geodesic and the statement about Jacobi-fields, respectively. q.e.d.

Let $x=\left(x_{1}, \ldots, x_{n}\right)$ be local coordinates around $p \in L$ and let $(x, \xi) \in \mathbb{R}^{2 n}$ be the corresponding coordinates on $T M$, where $\xi=\xi_{s} \partial_{s}$ (here, and in the rest of this section, we use the Einstein summation convention, repeated indices are summed over) where $\partial_{j}$ is the tangent vector of $T L$ in the $x_{j}$-direction. We write $\partial_{j^{*}}$ for the tangent vector of $T L$ in the $\xi_{j}$-direction. Let $\nabla, \hat{\nabla}$ denote the Levi-Civita connections of $g$ and $\hat{g}$, respectively. Let Roman and Greek indices run over the sets $\{1, \ldots, n\}$ and $\left\{1,1^{*}, 2,2^{*}, \ldots, n, n^{*}\right\}$, respectively and recall the following standard notation:

$$
\begin{aligned}
& g_{i j}=g\left(\partial_{i}, \partial_{j}\right), \quad \hat{g}_{\alpha \beta}=\hat{g}\left(\partial_{\alpha}, \partial_{\beta}\right) \\
& \nabla_{\partial_{i}} \partial_{j}=\Gamma_{i j}^{k} \partial_{k}, \quad \hat{\nabla}_{\partial_{\alpha}} \partial_{\beta}=\hat{\Gamma}_{\alpha \beta}^{\gamma} \partial_{\gamma} \\
& \left.R\left(\partial_{i}, \partial_{j}\right) \partial_{k}=R_{i j k}^{l} \partial_{l}, \quad g\left(R\left(\partial_{i}, \partial_{j}\right)\right) \partial_{k}, \partial_{r}\right)=R_{i j k r}
\end{aligned}
$$

Lemma 5.4. The components of the metric $\hat{g}$ satisfies

$$
\begin{align*}
\hat{g}_{i j}(x, \xi) & =g_{i j}(x)+\xi_{s} \xi_{t}\left(g_{k r}(x) \Gamma_{i s}^{k}(x) \Gamma_{j t}^{r}(x)+R_{i s j t}(x)\right)  \tag{5.8}\\
\hat{g}_{i^{*} j^{*}}(x, \xi) & =g_{i j}(x)  \tag{5.9}\\
\hat{g}_{i j^{*}}(x, \xi) & =\xi_{s} g_{j k}(x) \Gamma_{i s}^{k}(x) \tag{5.10}
\end{align*}
$$

Proof. Since $\partial_{j^{*}}$ is vertical, (5.9) holds. Note that the horizontal space at $(x, \xi)$ is spanned by the velocity vectors of the curves obtained by parallel translating $\xi$ along the coordinate directions through $x$. Let $V(t)$ be a parallel vector field through $x$ in the $\partial_{j}$-direction with $V(0)=$ $\xi$ and $\dot{V}(0)=a_{k} \partial_{k}$. Then

$$
V(t)=\left(\xi_{k}+t a_{k}+\mathcal{O}\left(t^{2}\right)\right) \partial_{k}
$$

and applying $\nabla_{\partial_{j}}$ to $V(t)$ we get

$$
0=\nabla_{\partial_{j}} V(t)=\xi_{s} \nabla_{\partial_{j}} \partial_{s}+a_{k} \partial_{k}+\mathcal{O}(t)
$$

Taking the limit as $t \rightarrow 0$, we find $a_{k} \partial_{k}=-\xi_{s} \Gamma_{j s}^{k}(x) \partial_{k}$. Hence, the horizontal space at $(x, \xi)$ is spanned by the vectors $\partial_{j}-\xi_{s} \Gamma_{j s}^{k} \partial_{k^{*}}, j=$ $1, \ldots, n$ and therefore,

$$
\partial_{j}^{V}=\xi_{s} \Gamma_{j s}^{k}(x) \partial_{k^{*}}
$$

Straightforward calculation gives (5.8) and (5.10). q.e.d.
Lemma 5.5. The Christofel symbols of the metric $\hat{g}$ at $(x, 0)$ satisfies

$$
\begin{align*}
& \hat{\Gamma}_{i j}^{k}(x, 0)=\hat{\Gamma}_{i j^{*}}^{k^{*}}(x, 0)=\Gamma_{i j}^{k}(x),  \tag{5.11}\\
& \hat{\Gamma}_{i j}^{k^{*}}(x, 0)=\hat{\Gamma}_{i j^{*}}^{k}(x, 0)=\Gamma_{i^{*} j^{*}}^{k^{*}}(x, 0)=0 . \tag{5.12}
\end{align*}
$$

Hence, if $\gamma$ is a geodesic in $(L, g)$, then it is also a geodesic in $(T L, \hat{g})$.
Proof. The equations

$$
\hat{\Gamma}_{\alpha \beta}^{\gamma}=\frac{1}{2} \hat{g}^{\gamma \delta}\left(\hat{g}_{\alpha \delta, \beta}+\hat{g}_{\beta \delta, \alpha}-\hat{g}_{\alpha \beta, \delta}\right),
$$

where $\hat{g}^{\alpha \beta}$ denotes the components of the inverse matrix of $\hat{g}$ and Lemma 5.4 together imply (5.11) and (5.12).

Let $x(t)$ be a geodesic in $(L, g)$. Then $\left(x, x^{*}\right)=(x(t), 0)$ satisfies

$$
\ddot{x}_{k}+\hat{\Gamma}_{i j}^{k} \dot{x}_{i} \dot{x}_{j}+\hat{\Gamma}_{i^{*} j}^{k} \dot{x}_{i^{*}} \dot{x}_{j}+\hat{\Gamma}_{i j^{*}}^{k} \dot{x}_{i} \dot{x}_{j^{*}}+\hat{\Gamma}_{i^{*} j^{*}}^{k} \dot{x}_{i^{*}} \dot{j}_{j^{*}}=\ddot{x}_{k}+\Gamma_{i j}^{k} \dot{x}_{i} \dot{x}_{j}=0,
$$

$$
\ddot{x}_{k^{*}}+\hat{\Gamma}_{i j}^{k^{*}} \dot{x}_{i} \dot{x}_{j}+\hat{\Gamma}_{i^{*} j}^{k^{*}} \dot{x}_{i^{*}} \dot{x}_{j}+\hat{\Gamma}_{i j^{*}}^{k^{*}} \dot{x}_{i} \dot{x}_{j^{*}}+\hat{\Gamma}_{i^{*} j^{*}}^{k^{*}} \dot{x}_{i^{*}} \dot{x}_{j^{*}}=0 .
$$

This proves the second statement. q.e.d.
Lemma 5.6. If $\gamma$ is a geodesic in ( $T L, \hat{g}$ ) which lies in $L$, then $X$ is a Jacobi-field along $\gamma$ if and only if JX is.

Proof. The proof appears in the Appendix. q.e.d.
5.3. A family of metrics on $\mathbb{C}^{n}$. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an instant of a chord generic 1-parameter family of Legendrian submanifolds and fix a Riemannian metric $g$ on $L$. Using the metric $\hat{g}$ on $T L$ (see Section 5.2), we construct a 1-parameter family of metrics $g(L, \sigma), 0 \leq \sigma \leq 1$, on $\mathbb{C}^{n}$ with good properties with respect to $\Pi_{\mathbb{C}}(L)$.

Let $c_{1}, \ldots, c_{m}$ be the Reeb chords of $L$. Fix $\delta>0$ such that all the $6 \delta$-balls $B\left(c_{j}^{*}, 6 \delta\right)$ are disjoint and such that the intersections $B\left(c_{j}^{*}, 6 \delta\right) \cap$ $\Pi_{\mathbb{C}}(L)$ are homeomorphic to two $n$-disks intersecting at a point.

Identify the normal bundle of the immersion $\Pi_{\mathbb{C}}$ with the tangent bundle $T L$. Consider the metric $\hat{g}$ on a $\rho$-neighborhood of the 0 -section in $T L\left(\rho>0\right.$ as in Proposition 5.3). Let $P: W \rightarrow \mathbb{C}^{n}$ be an immersion of a $\rho^{\prime}$-neighborhood $N\left(\rho^{\prime}\right)$ of the 0 -section $\rho^{\prime} \leq \rho$ such that $i \circ d P=d P \circ J$ along the 0 -section.

Consider the $P$-push-forward of the metric $\hat{g}$ to the image of $N\left(\rho^{\prime}\right)$ restricted to $L \backslash \bigcup_{j} U\left(c_{j}^{-}, \delta\right)$. Note that if $\rho^{\prime}>0$ is small enough this restriction of $P$ is an embedding and the push-forward metric is defined in a neighborhood of $\Pi_{\mathbb{C}}\left(L \backslash \bigcup_{j} U\left(c_{j}^{-}, 2 \delta\right)\right)$. Extend it to a metric $g^{1}$ on all of $\mathbb{C}^{n}$, which agrees with the standard metric outside a neighborhood of $\Pi_{\mathbb{C}}(L)$.

Consider the $P$-push-forward of the metric $\hat{g}$ to the image of the $\rho^{\prime}$-neighborhood of the 0 -section restricted to $L \backslash \bigcup_{j} U\left(c_{j}^{+}, \delta\right)$. This metric is defined in a neighborhood of $\Pi_{\mathbb{C}}\left(L \backslash \bigcup_{j} U\left(c_{j}^{+}, 2 \delta\right)\right)$ and can be
extended to a metric $g^{0}$ on all of $\mathbb{C}^{n}$, which agrees with the standard metric outside a neighborhood of $\Pi_{\mathbb{C}}(L)$.

Choose the metrics $g^{0}$ and $g^{1}$ so that they agree outside $\cup_{j} B\left(c_{j}^{*}, 3 \delta\right)$ and let $g^{\sigma}, 0 \leq \sigma \leq 1$ be a smooth 1-parameter family of metrics on $\mathbb{C}^{n}$ with the following properties:

- $g^{\sigma}=g^{0}$ in a neighborhood of $\sigma=0$,
- $g^{\sigma}=g^{1}$ in a neighborhood of $\sigma=1$,
- $g^{\sigma}$ is constant in $\sigma$ outside $\cup_{j} B\left(c_{j}^{*}, 4 \delta\right)$.

We take $g(L, \sigma)=g^{\sigma}$.
Remark 5.7. If $L_{\lambda}, \lambda \in \Lambda$ is a smooth family of chord generic Legendrian submanifolds then, as is easily seen, the above construction can be carried out in such a way that the family of 1-parameter families of metrics $g\left(L_{\lambda}, \sigma\right)$ becomes smooth in $\lambda$.

Given a vector field $v$ along a disk $u: D_{m} \rightarrow \mathbb{C}^{n}$ with boundary on $L$, we would like to be able to exponentiate $v$ to get a variation of $u$ among disks with boundaries on $L$. We will not be able to use a fixed metric $g^{\sigma}$ to do this. To solve this problem, let $\sigma: \mathbb{C}^{n} \times \mathbb{R} \rightarrow[0,1]$ be a smooth function which equals 0 on

$$
\mathbb{C}^{n} \times \mathbb{R}-\bigcup_{j} B\left(c_{j}^{*}, 5 \delta\right) \times\left[c_{j}^{+}-\frac{1}{2} \mathcal{Z}\left(c_{j}\right), c_{j}^{+}+1\right]
$$

and equals 1 on

$$
\bigcup_{j} B\left(c_{j}^{*}, 4 \delta\right) \times\left[c_{j}^{+}-\frac{1}{4} \mathcal{Z}\left(c_{j}\right), c_{j}^{+}+\frac{1}{2}\right]
$$

Let $\exp _{p}^{g}$ denote the exponential map of the metric $g$ at the point $p$. If $p \in L_{\lambda}$ and $v$ is tangent to $L_{\lambda}$ at $p$, then write $x(p)=\Pi_{\mathbb{C}}(p)$ and $\xi(v)=\Pi_{\mathbb{C}}(v)$. One may now easily prove the following lemma.

Lemma 5.8. Let $L_{\lambda}, \lambda \in \Lambda$ be a family of (semi-) admissible Legendrian submanifolds. Let $0 \in \Lambda$ and let $\sigma: \mathbb{C}^{n} \times \mathbb{R} \rightarrow[0,1]$ be the function constructed from $L_{0}$ as above. There exists $\rho>0$ and a neighborhood $W \subset \Lambda$ of 0 such that if $p$ is any point in $L_{\lambda}, \lambda \in W$ and $v$ any vector tangent to $L_{\lambda}$ at $p$ with $|\xi(v)|<\rho$ then

$$
\exp _{x}^{g\left(L_{\lambda}, \sigma(p)\right)} t \xi \in \Pi_{\mathbb{C}}\left(L_{\lambda}\right) \text { for } 0 \leq t \leq 1
$$

5.4. Extending families of Legendrian embeddings and their differentials. In the next subsection, we will need to exponentiate vector fields along a disk whose boundary is in $L_{0}(0 \in \Lambda)$ to get a disk with boundary in $L_{\lambda}$ for $\lambda$ near 0 . To accomplish this, we construct diffeomorphisms of $\mathbb{C}^{n}$.

Consider $L_{\lambda} \subset \mathbb{C}^{n} \times \mathbb{R}, \lambda \in \Lambda$ and let $0 \in \Lambda$. There exists a smooth family of Legendrian embeddings

$$
k_{\lambda}: L_{0} \rightarrow \mathbb{C}^{n} \times \mathbb{R}
$$

such that $k_{0}$ is the inclusion, $k_{\lambda}\left(L_{0}\right)=L_{\lambda}$, and $k_{\lambda}\left(c_{j}^{ \pm}(0)\right)=c_{j}^{ \pm}(\lambda)$ for each $j$.

As in Section 5.3, fix $\delta>0$ such that all the $6 \delta$-balls $B\left(c_{j}^{*}(0), 6 \delta\right)$ are disjoint and such that the intersections $B\left(c_{j}^{*}(0), 6 \delta\right) \cap \Pi_{\mathbb{C}}\left(L_{0}\right)$ are homeomorphic to two $n$-disks intersecting at a point.

Let $W \subset \Lambda$ be a neighborhood of 0 such that $c_{j}^{*}(\lambda) \in B\left(c_{j}^{*}(0), \delta\right)$ for $\lambda \in W$. We construct a smooth $\Lambda$-family $(\lambda \in W)$ of 1-parameter families of diffeomorphisms $\psi_{\lambda}^{\sigma}: \mathbb{C}^{n} \rightarrow \mathbb{C}^{n}, 0 \leq \sigma \leq 1, \lambda \in W$. Note that

$$
\begin{align*}
& K_{\lambda}^{1}=\Pi_{\mathbb{C}} \circ k_{\lambda}: L_{0}^{1}=L_{0} \backslash \bigcup_{j} U\left(c_{j}^{+}, 3 \delta\right) \rightarrow \mathbb{C}^{n},  \tag{5.13}\\
& K_{\lambda}^{0}=\Pi_{\mathbb{C}} \circ k_{\lambda}: L_{0}^{0}=L_{0} \backslash \bigcup_{j} U\left(c_{j}^{-}, 3 \delta\right) \rightarrow \mathbb{C}^{n}
\end{align*}
$$

are Lagrangian embeddings and that $K_{\lambda}^{1}\left(c_{j}^{*}(0)\right)=K_{\lambda}^{0}\left(c_{j}^{*}(0)\right)=c_{j}^{*}(\lambda)$, for each Reeb chord $c_{j}(0)$ of $L_{0}$.

Identify tubular neighborhoods of $L_{0}^{1}$ and $L_{0}^{0}$ with their respective tangent bundles so that $J$ along the 0 -section of the tangent bundles corresponds to $i$ in $\mathbb{C}^{n}$ (see Section 5.3). Define for $(p, v) \in T L_{0}^{\alpha} \subset \mathbb{C}^{n}$, $\alpha=0,1$,

$$
\begin{equation*}
\hat{K}_{\lambda}^{\alpha}(p, v)=K_{\lambda}^{\alpha}(p)+i d K_{\lambda}^{\alpha}(v) \tag{5.14}
\end{equation*}
$$

Then $\hat{K}_{\lambda}^{\alpha}$ is a diffeomorphism on some neighborhood of $L_{0}^{\alpha} \subset \mathbb{C}^{n}$, $\alpha=0,1$. Note that the diffeomorphisms $\hat{K}_{\lambda}^{0}$ and $\hat{K}_{\lambda}^{1}$ agree outside $\bigcup_{j} B\left(c_{j}^{*}, 4 \delta\right)$.

Extend $\hat{K}_{\lambda}^{0}$ and $\hat{K}_{\lambda}^{1}$ to diffeomorphisms on all of $\mathbb{C}^{n}$ in such a way that their extensions agree outside $\bigcup_{j} B\left(c_{j}^{*}, 4 \delta\right)$. Call these extensions $\psi_{\lambda}^{\alpha}, \alpha=0,1$.

Let $\psi_{\lambda}^{\sigma}, 0 \leq \sigma \leq 1$ be a $\Lambda$-family of 1-parameter families of diffeomorphisms which are constant in $\sigma$ near $\sigma=0$ and $\sigma=1$ and with the following properties. First, $\psi_{\lambda}^{\sigma}, 0 \leq \sigma \leq 1$ connects $\psi_{\lambda}^{0}$ to $\psi_{\lambda}^{1}$. Second, $\psi_{\lambda}^{\sigma}$ is constant in $\sigma$ outside $\cup_{j} B\left(c_{j}^{*}, 5 \delta\right)$ and in $\cup_{j}\left(B\left(c_{j}^{*}, 5 \delta\right) \backslash B\left(c_{j}^{*}, 4 \delta\right)\right) \cap L_{0}$. Third, $\psi_{\lambda}^{\sigma}\left(c_{j}^{*}(0)\right)=c_{j}^{*}(\lambda), 0 \leq \sigma \leq 1$.

For future reference, we let $Y_{\lambda}^{\sigma}$ denote the 1-parameter family of 1forms on $\Lambda$ with coefficients in smooth vector fields on $\mathbb{C}^{n}$ defined by

$$
\begin{equation*}
Y_{\lambda}^{\sigma}(x, \mu)=D_{\lambda} \psi_{\lambda}^{\sigma}(x) \cdot \mu, \quad \lambda \in \Lambda, \mu \in T_{\lambda} \Lambda, x \in \mathbb{C}^{n}, \sigma \in[0,1] \tag{5.15}
\end{equation*}
$$

By (5.14), $d \psi_{\lambda}^{\alpha}, \alpha=0,1$ are complex linear maps when restricted to the restriction of the tangent bundle of $\mathbb{C}^{n}$ to $L_{0}^{\alpha}$. Moreover, these maps fit together to a smooth $\Lambda$-family of maps $\hat{A}_{\lambda}: L_{0} \rightarrow \mathbf{G L}\left(\mathbb{C}^{n}\right)$ which is obtained as follows. Pick a smooth function $\beta$ on $L_{0}$ with values in $[0,1]$ which is 0 outside $U\left(c_{j}^{+}, 5 \delta\right)$ and 1 inside $U\left(c_{j}^{+}, 4 \delta\right)$ define

$$
\hat{A}_{\lambda}(p)=d \psi^{\beta(p)}\left(d \Pi_{\mathbb{C}}\left(T_{p} L\right)\right)
$$

Let $A_{\lambda}^{\sigma}: \mathbb{C}^{n} \rightarrow \mathbf{G L}\left(\mathbb{C}^{n}\right)$ be an $s$-parameter family of 1-parameter families of maps with the following properties.

- $A_{\lambda}^{\sigma}=\hat{A}_{\lambda}$ on $\Pi_{\mathbb{C}}\left(L_{0}\right) \backslash \Pi_{\mathbb{C}}\left(U\left(c_{j}^{+}, 5 \delta\right)\right)$
- $A_{\lambda}^{1}=\hat{A}_{\lambda}$ on $\Pi_{\mathbb{C}}\left(U\left(c_{j}^{+}, 4 \delta\right)\right)$
- $A_{\lambda}^{\sigma}$ is constant in $\sigma$ on $B\left(c_{j}^{*}, 5 \delta\right) \backslash B\left(c_{j}^{*}, 4 \delta\right) \cap L_{0}$
- $\bar{\partial} A_{\lambda}^{0}=0$ along $\Pi_{\mathbb{C}}\left(L_{0}\right) \backslash \Pi_{\mathbb{C}}\left(U\left(c_{j}^{+}, 4 \delta\right)\right)$ and $\bar{\partial} A_{\lambda}^{1}=0$ along $\Pi_{\mathbb{C}}\left(L_{0}\right) \backslash$ $\Pi_{\mathbb{C}}\left(U\left(c_{j}^{-}, 4 \delta\right)\right)$.
- $\left\|A_{\lambda}^{\sigma}-\mathrm{id}\right\|_{C^{\infty}} \leq 2\left\|\hat{A}_{\lambda}-\mathrm{id}\right\|_{C^{\infty}}$.
5.5. Local coordinates. We consider first the chord generic case. Let $L_{\lambda} \subset \mathbb{C}^{n} \times \mathbb{R}, \lambda \in \Lambda$ be a family of chord generic Legendrian submanifolds. We construct local coordinates on $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$.

Let $\sigma: \mathbb{C}^{n} \times \mathbb{R} \rightarrow[0,1]$ be the function constructed from $L_{0}, 0 \in \Lambda$. For $p \in L_{\lambda}$ and $v$ a tangent vector of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ at $q=\Pi_{\mathbb{C}}(p)$, write

$$
\exp _{q}^{g\left(L_{\lambda}, \sigma(p)\right)} v=\exp _{q}^{\lambda, \sigma} v
$$

Moreover, if $\rho>0$ is as in Lemma 5.8 and $|v| \leq \rho$, we write $z(p, v)$ for the $z$-coordinate of the endpoint of the unique continuous lift of the path $\exp _{q}^{\lambda, \sigma(p)} t v, 0 \leq t \leq 1$, to $L \subset \mathbb{C}^{n} \times \mathbb{R}$.

Let $(w, f) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$. Let $F: D_{m} \rightarrow \mathbb{R}$ be an extension of $f$ such that $F \in \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{R}\right)$ (in particular, $F$ is continuous) and such that $F$ is smooth with all derivatives uniformly bounded outside a small neighborhood of $\partial D_{m}$. Then $w \times F: D_{m} \rightarrow \mathbb{C}^{n} \times \mathbb{R}$. In the case that $w$ and $f$ are smooth, we take $F$ to be smooth. Furthermore, in the case that $w$ and $f$ are constant close to each puncture we take $F$ to be an affine parameterization of the corresponding Reeb chord in a neighborhood of each puncture where $w$ and $f$ are constant. The purpose of this choice of $F$ is that when we exponentiate a vector field at the disk $(w, f)$, we need $(w, F)$ to determine the metric.

For $r>0$, define

$$
\mathcal{B}_{2, \epsilon}((w, f), r) \subset \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right)
$$

as the intersection of the closed subspace of $v \in \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right)$ which satisfies

$$
\begin{align*}
& v(\zeta) \in \Pi_{\mathbb{C}}\left(T_{(f(\zeta), w(\zeta))} L\right), \text { for } \zeta \in \partial D_{m}  \tag{5.16}\\
& \int_{\partial D_{m}}\langle\bar{\partial} v, a\rangle d s=0, \text { for every } a \in C_{0}^{\infty}\left(\partial D_{m}, \mathbb{C}^{n}\right) \tag{5.17}
\end{align*}
$$

and the ball $\left\{u:\|u\|_{2, \epsilon}<r\right\}$.
When the parameter space $\Lambda$ is 0 -dimensional, we can define a coordinate chart around $(f, w, 0) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)(0 \in \Lambda)$ by

$$
\begin{aligned}
& \Phi[(w, f, 0)]: \mathcal{B}_{2, \epsilon}((w, f), r) \times \Lambda \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa) ; \\
& \Phi[(w, f, 0)](v, \lambda)=(u, l, \lambda)
\end{aligned}
$$

where

$$
\begin{aligned}
u(\zeta) & =\exp _{w(\zeta)}^{\lambda, \sigma(\zeta)}(v(\zeta)) \\
l(\zeta) & =z((w(\zeta), f(\zeta)), v(\zeta)), \quad \zeta \in \partial D_{m}
\end{aligned}
$$

When $\Lambda$ is not 0 -dimensional, we will need to use the maps $A_{\lambda}^{\sigma}$ to move the "vector field" $v$ from $L_{0}$ to $L_{\lambda}$. Moreover, to ensure our new maps are in the appropriate space of functions, we will also need to cut off the original map $w$. To this end, let $(w, f, \lambda) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$. Then, there exists $M>0$ and vector-valued functions $\xi_{j}, j=1, \ldots, m$ such that

$$
w(\tau+i t)=\exp _{a_{j}^{*}}^{\lambda, \omega(t)} \xi_{j}(\tau+i t), \quad \text { for } \tau+i t \in E_{p_{j}}[M],
$$

where $\omega:[0,1] \rightarrow[0,1]$ is a smooth approximation of the identity, which is constant in neighborhoods of the endpoints of the interval. Define $(w[M], f[M])$ as follows. Let

$$
w[M](\zeta)= \begin{cases}w(\zeta), & \text { for } \zeta \notin \cup_{j} E_{p_{j}}[M], \\ \exp _{a_{j}^{*}}^{\lambda, \omega(t)}\left(\alpha \xi_{j}\right), & \text { for } \zeta=\tau+i t \in E_{p_{j}}[M],\end{cases}
$$

where $\alpha: E_{p_{j}} \rightarrow \mathbb{C}$ is a smooth function which is 1 on $E_{p_{j}} \backslash E_{p_{j}}[M+1], 0$ on $E_{p_{j}}[2 M]$, and holomorphic on the boundary. Let $f[M]$ be the natural lift of the boundary values of $w[M]$. It is clear that $(w[M], f[M]) \rightarrow$ $(w, f)$ as $M \rightarrow \infty$. For convenience we use the notation $(w[\infty], f[\infty])$ to denote this limit. We write $F[M]$ for the extension of $f[M]$ to $D_{m}$.

Let $(w, f, 0) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)(0 \in \Lambda)$. For large $M>0$, consider $(w[M], F[M])$. To simplify notation, write $\sigma[M](\zeta)=\sigma(w[M](\zeta)$, $F[M](\zeta))$ and $w[M]_{\lambda}(\zeta)=\psi_{\lambda}^{\sigma[M](\zeta)}(w[M](\zeta))$. Define

$$
\begin{aligned}
& \Phi[(w, f, 0) ; M]: \mathcal{B}_{2, \epsilon}((w[M], f[M]), r) \times \Lambda \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa) ; \\
& \Phi[(w, f, 0) ; M](v, \lambda)=(u, l, \lambda)
\end{aligned}
$$

where

$$
\begin{aligned}
u(\zeta) & =\exp _{w[M]_{\lambda}(\zeta)}^{\lambda, \sigma[M(\zeta)}\left(A_{\lambda}^{\sigma[M](\zeta)} v(\zeta)\right) \\
l(\zeta) & =z\left(\left(w[M]_{\lambda}(\zeta), f[M]_{\lambda}(\zeta)\right), A_{\lambda}^{\sigma[M](\zeta)} v(\zeta)\right), \quad \zeta \in \partial D_{m}
\end{aligned}
$$

In the semi-admissible case, we use the above construction close to all Reeb chords except the chord $c_{0}$ at the self-tangency point. At $c_{0}^{*}$, we utilize the fact that we have a local product structure of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ which is assumed to be preserved in a rather strong sense under $\lambda \in \Lambda$, see Section 5.1. This allows us to construct the family of metrics $g_{\lambda}^{\sigma}$ as product metrics close to $c_{0}^{*}$. Once we have metrics with this property, we can apply the cut-off procedure above to the last $(n-1)$ coordinates of an element $(w, f, 0) \in \mathcal{W}_{2, \epsilon, \Lambda}$ and just keep the first coordinate of $w$ in a neighborhood of $c_{0}^{*}$ as it is. We use the same notation $(w[M], f[M])$ for
the map which results from this modified cut-off procedure from $(w, f)$ in the semi-admissible case.

Proposition 5.9. Let $\epsilon \in[0, \infty)^{m}$. Then there exists $r>0, M>0$, and a neighborhood $W \subset \Lambda$ of 0 such that the map

$$
\Phi[(w, f, 0)]: \mathcal{B}_{2, \epsilon}((w[M], f[M]), r) \times W \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)
$$

is $C^{1}$ and gives local coordinates on some open subset of $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$ containing ( $w, f, 0$ ). Moreover, if $\Lambda$ is 0 -dimensional, then we may take $M=\infty$.

Proof. Fix some small $r>0$. Consider the auxiliary map

$$
\begin{aligned}
& \Psi: \mathcal{B}_{2, \epsilon}((w[M], f[M]), r) \times \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m}, \mathbb{R}\right) \times \Lambda \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}), \\
& \Psi(v, r, \lambda)=\Phi[(w[M], f[M]), 0](v, \lambda)+(0,0, r)
\end{aligned}
$$

where $(u, h, \mu)+(0,0, r)=(u, h, \mu+r)$.
We show in Lemma 5.11 that $\Psi$ is $C^{1}$ with differential in a neighborhood of $(0,0,0)$ which maps injectively into the tangent space of the target and has closed images. These closed images have direct complements and hence, the implicit function theorem applies and shows that the image is a submanifold. Moreover, for $M$ large enough $(w, f, 0)$ is in the image.

We finally prove in Lemma 5.13, that $\mathcal{W}_{2, \epsilon}(\mathbf{a}, \kappa)$ lies inside the image and that it corresponds exactly to $r=0$ in the given coordinates. q.e.d.

Lemma 5.11 is a consequence of the following technical lemma.
Lemma 5.10. Let $\Lambda$ be an open neighborhood of 0 in a Banach space. Let $(w, f, \lambda) \in \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)$, and $v, u, q \in \mathcal{B}_{2, \epsilon}((w, f), r)$. Let $\zeta$ be a coordinate on $D_{m}$ and let $\epsilon \in[0, \infty)^{m}$.
(a) Let

$$
G: \mathbb{C}^{n} \times \mathbb{C}^{n} \times \mathbb{C}^{n} \times \mathbb{C}^{n} \times[0,1] \times \Lambda \rightarrow \mathbb{C}^{n}
$$

be a smooth function with all derivatives uniformly bounded and let $\sigma: \mathbb{C}^{n} \times \mathbb{R} \rightarrow[0,1]$ be a smooth function with the same property. If

$$
\begin{align*}
& G(x, 0,0, \theta, \sigma, \lambda)=0  \tag{5.18}\\
& G(x, \xi, 0, \theta, \sigma, 0)=0 \tag{5.19}
\end{align*}
$$

then there exists a constant $C$ (depending on $\|D w\|_{1, \epsilon},\|D F\|_{1, \epsilon}$ and $r)$ such that $G(\zeta, \lambda)=G(w(\zeta), v(\zeta), u(\zeta), q(\zeta), \sigma(F(\zeta)$, $w(\zeta)), \lambda)$ satisfies

$$
\begin{equation*}
\|G(\zeta, \lambda)\|_{2, \epsilon} \leq C\left(\|u\|_{2, \epsilon}+\|v\|_{2, \epsilon}+|\lambda|\right) . \tag{5.20}
\end{equation*}
$$

(b) Let

$$
G: \mathbb{C}^{n} \times \mathbb{C}^{n} \times \mathbb{C}^{n} \times[0,1] \times \Lambda \rightarrow \mathbb{C}^{n}
$$

be a smooth function with all derivatives uniformly bounded. If

$$
\begin{align*}
& G(x, 0,0, \sigma, \lambda)=0,  \tag{5.21}\\
& G(x, \xi, 0, \sigma, 0)=0,  \tag{5.22}\\
& D_{3} G(x, \xi, 0, \sigma, 0)=0, \text { and }  \tag{5.23}\\
& D_{5} G(x, \xi, 0, \sigma, 0)=0 \tag{5.24}
\end{align*}
$$

then there exists a constant $C$ (depending on $\|D w\|_{1, \epsilon},\|D F\|_{1, \epsilon}$ and $r)$ such that $G(\zeta, \lambda)=G(w(\zeta), v(\zeta), u(\zeta), \sigma(F(\zeta), w(\zeta)), \lambda)$ satisfies

$$
\|G(\zeta, \lambda)\|_{2, \epsilon} \leq C\left(\|u\|_{2, \epsilon}^{2}+|\lambda|^{2}\right) .
$$

Proof. The proof appears in the Appendix.
q.e.d.

In order to express the derivative of $\Psi$, we will use the function $K: \mathbb{C}^{n} \times \mathbb{C}^{n} \times[0,1] \times \Lambda \rightarrow \mathbb{C}^{n}$ defined by

$$
\begin{equation*}
K(x, \xi, \sigma, \lambda)=\exp _{\psi_{\lambda}^{\alpha}(x)}^{\lambda, \sigma} A_{\lambda}^{\sigma} \xi-\psi_{\lambda}^{\sigma}(x) . \tag{5.25}
\end{equation*}
$$

We will need to lift $K$ (at least on part of its domain) so that it maps to $\mathbb{C}^{n} \times \mathbb{R}$. We describe this lift.

Consider $L_{\lambda} \subset \mathbb{C}^{n} \times \mathbb{R}, \lambda \in \Lambda$. Let $K_{\lambda}: T L_{0} \rightarrow \mathbb{C}^{n} \times \mathbb{R}$ be an embedding extension of $k_{\lambda}$ (see Section 5.4). Consider the immersion $P_{\lambda}: V \subset T L_{0} \rightarrow \mathbb{C}^{n}$ which extends $\Pi_{\mathbb{C}} \circ k_{\lambda}$, where $V$ is a neighborhood of the 0 -section in $T L_{0}$. Choose $V$ and a neighborhood $W \subset \Lambda$ of 0 , so small that the self-intersection of $P_{\lambda}$ is contained inside $\bigcup_{j} B\left(c_{j}^{*}(0), 2 \delta\right)$. Consider the following subset $N$ of the product $\mathbb{C}^{n} \times[0,1]$.

$$
\begin{aligned}
N= & P(V) \backslash \bigcup_{j} B\left(c_{j}^{*}(0), 3 \delta\right) \times[0,1] \\
& \cup \bigcup_{j} P\left(V \mid L_{\lambda} \cap U\left(c_{j}^{+}, 4 \delta\right)\right) \times[1-\epsilon, 1] \\
& \cup \bigcup_{j} P\left(V \mid L_{\lambda} \cap U\left(c_{j}^{-}, 4 \delta\right)\right) \times[0, \epsilon] .
\end{aligned}
$$

We define a map $\psi_{\lambda}: N \rightarrow \mathbb{C}^{n} \times \mathbb{R}$ in the natural way, $\psi_{\lambda}(q, \sigma)=$ $K_{\lambda}\left(p_{\sigma}, v_{\sigma}\right)$ where ( $p_{\sigma}, v_{\sigma}$ ) is the preimage of $q$ under $P$ with $p \in U\left(c_{j}^{ \pm}, 4 \delta\right)$ where the sign is determined by $\sigma$.

Using this construction, we may do the following. If $W \subset \mathbb{C}^{n} \times \mathbb{C}^{n} \times$ $[0,1] \times \Lambda$ and $G: W \rightarrow \mathbb{C}^{n}$ is a function such that $(G, \sigma)(W) \in N$, then we may define a lift $\tilde{G}: W \rightarrow \mathbb{C}^{n} \times \mathbb{R}$.

We now use this construction to lift the function $K$ defined in (5.25). For $x$ sufficiently close to $L_{0}, \xi$ sufficiently small and $\sigma$ sufficiently close to 0 or 1 when $x$ is close to double points of $\Pi_{\mathbb{C}}\left(L_{0}\right)$ the lift $\tilde{K}$ of $K$ can be defined. Let $K_{\mathbb{R}}$ denote the $\mathbb{R}$-coordinate of $\tilde{K}$.

Lemma 5.11. If $\operatorname{dim} \Lambda>0$, let $M<\infty$. If $\operatorname{dim} \Lambda=0$, let $M=\infty$. The map

$$
\Psi: \mathcal{B}_{2, \epsilon}((w[M], f[M]), r) \times \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m}, \mathbb{R}\right) \times \Lambda \rightarrow \mathcal{F}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa)
$$

is $C^{1}$. Its derivative at $(v, h, \mu)$ is the map

$$
(u, l, \lambda) \mapsto\left(D_{2} K \cdot v+D_{4} K \cdot \lambda,\left\langle D_{2} K_{\mathbb{R}} \cdot v+D_{4} K_{\mathbb{R}} \cdot \lambda\right\rangle+l, \lambda\right)
$$

where all derivatives of $K$ and $\tilde{K}$ are evaluated at $\left(w[M]_{\lambda}, v, \sigma[M], \lambda\right)$ and where $\langle u\rangle$ denotes restriction of $u: D_{m} \rightarrow \mathbb{R}$ to the boundary.

Proof. Using local coordinates on $\mathcal{F}_{2, \epsilon, \Lambda}$ as described in Section 5.1, we write $\Psi=\left(\Psi_{1}, \Psi_{2}, \Psi_{3}\right)$ Statements concerning $\Psi_{3}$ are trivial. Note that $\Psi_{1}=K+\Psi_{\lambda}^{\sigma}(x)$. So, to see that $\Psi_{1}$ is continuous, we note that $K(x, 0, \sigma, \lambda)=0$ and apply Lemma 5.10 (a) to get that $K$ is Lipschitz in $v$ and $\lambda$ and hence continuous. To see that $\Psi_{1}$ is differentiable, we note that if

$$
\begin{aligned}
G(x, \xi, \eta, \sigma, \lambda)= & K(x, \xi+\eta, \sigma, \mu+\lambda)-K(x, \xi, \sigma, \mu) \\
& -\left(D_{2} K(x, \xi, \sigma, \mu) \cdot \eta+D_{4} K(x, \xi, \sigma, \mu) \cdot \lambda\right),
\end{aligned}
$$

then the conditions (5.21)-(5.24) are fulfilled and Lemma 5.10 (b) implies $\Psi_{1}$ is differentiable and has differential as claimed. Finally, applying Lemma 5.10 (a) to the map

$$
G(x, \xi, \eta, \sigma, \lambda)=D_{2} K(x, \xi, \sigma, \mu) \cdot \eta+D_{4} K(x, \xi, \sigma, \mu) \cdot \lambda
$$

shows $\Psi_{1}$ is $C^{1}$.
Using $\tilde{K}$, we can extend the $\mathbb{R}$-valued function $z\left(\left(w[M]_{\lambda}, f[M]_{\lambda}\right)\right.$, $\left.A_{\lambda}^{\sigma} v\right)$ to a small neighborhood of $\partial D_{m}$ in $D_{m}$. With this done the (nontrivial part) of the derivative of $\Psi_{2}$ can be handled exactly as above.
q.e.d.

Let $\zeta=x_{1}+i x_{2}$ be a complex local coordinate in $D_{m}$. Then, if $u: D_{m} \rightarrow \mathbb{C}^{n}$, we may view $\bar{\partial} u$ as $\partial_{1} u+i \partial_{2} u$. As in the proof of Lemma 5.11, we use local coordinates on $\mathcal{F}_{2, \epsilon, \Lambda}$ and write $\Psi=$ $\left(\Psi_{1}, \Psi_{2}, \Psi_{3}\right)$.

Lemma 5.12. Assume that $w: D_{m} \rightarrow \mathbb{C}^{n}$ and $v: D_{m} \rightarrow \mathbb{C}^{n}$ are smooth functions and let $g$ be any metric on $\mathbb{C}^{n}$ with Levi-Civita connection $\nabla$. Let $\zeta=\tau_{1}+i \tau_{2} \in D_{m}$, if $u(\zeta)=\exp _{w(\zeta)}(v(\zeta))$ then, $\bar{\partial} u=X_{1}(1)+i X_{2}(1)$, where $X_{j}, j=1,2$ are the Jacobi-fields along the geodesic $\exp _{w(\zeta)}(t v(\zeta)), 0 \leq t \leq 1$, with $X_{j}(0)=\partial_{j} w(\zeta)$ and $\nabla_{t} X_{j}(0)=\nabla_{\partial_{j}} v(\zeta)$.

In particular, there exists $r>0$ such that if $(v(\zeta), \lambda) \in \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right) \times$ $\Lambda,\|v\|_{2, \epsilon} \leq r$, then the restriction of $\bar{\partial} \Psi_{1}(v, \lambda)$ to $\partial D_{m}$ equals 0 if and only if the restriction of $\bar{\partial} v$ to the boundary equals 0 .

## Proof. Consider

$$
\alpha(s, t)=\exp _{w(\zeta+s)}(t v(\zeta+s)), \quad 0 \leq t \leq 1,-\epsilon \leq s \leq \epsilon
$$

Since for fixed $s, t \mapsto \alpha(s, t)$ is a geodesic, we find that $\partial_{s} \alpha(0, t)=$ $X_{1}(t)$ is a Jacobi field along the geodesic $t \mapsto \exp _{w(\zeta)}(t v(\zeta))$ with initial conditions

$$
\begin{aligned}
X_{1}(0) & =\partial_{s} \exp _{w(\zeta+s)}(0 \cdot v)=\partial_{1} w(\zeta), \\
\nabla_{t} X_{1}(0) & =\nabla_{t} \partial_{s} \alpha(0,0)=\nabla_{s} \partial_{t} \alpha(0,0)=\nabla_{s} v(\zeta) .
\end{aligned}
$$

Moreover,

$$
\exp _{w(\zeta+s)}(v(\zeta+s))=\alpha(s, 1)
$$

and hence,

$$
\partial_{1} \exp _{w(\zeta)}(v(\zeta))=\partial_{s} \alpha(0,1)=X_{1}(1)
$$

A similar analysis shows that

$$
\partial_{2} \exp _{w(\zeta)}(v(\zeta))=X_{2}(1)
$$

This proves the first statement.
Consider the second statement. Note that the metrics $g\left(L_{\lambda}\right.$, $\sigma(w[M](\zeta), F[M](\zeta)))$ are constant in $\zeta$ for $\zeta$ in a neighborhood of $\partial D_{m}$. Consider first the case that $w[M]$ and $v$ are smooth. Then the above result together with the Jacobi-field property of the metric $\hat{g}$ (see Lemma 5.6), from which $g\left(L_{\lambda}, \sigma\right)$ is constructed implies that for $\zeta=\tau_{1}+i \tau_{2} \in$ $\partial D_{m}, \bar{\partial}\left(\Psi_{1}(v, \lambda)\right)=X_{1}(1)+i X_{2}(1)$ equals the value of the Jacobi-field $X_{1}+J X_{2}=Y$ with initial condition $Y(0)=0, \nabla_{t} Y(0)=\left(\nabla_{\partial_{1}}+\right.$ $\left.i \nabla_{\partial_{2}}\right)\left(A_{\lambda}^{\sigma} v\right)$.

Note $Y(0)=0$ since, along the boundary, $\bar{\partial} w=0$. Let $u=A_{\lambda}^{\sigma} v$. We check that $\nabla_{t} Y(0)=\bar{\partial} u$. To this end, let $(x, \xi)=\left(x_{1}, \ldots, x_{n}, \xi_{1}, \ldots, \xi_{n}\right)$ be coordinates on $T^{*} L$, and use notation as in Lemma 5.5. Noting that $u$ is tangent to $L$ at $w(\zeta)$, we compute

$$
\begin{aligned}
& \nabla_{\partial_{\tau_{1}}} u=\frac{\partial u}{\partial \tau_{1}}+\left(\Gamma_{i j}^{k}\left(\frac{\partial w}{\partial \tau_{1}}\right)_{i}+\Gamma_{i^{*} j}^{k}\left(\frac{\partial w}{\partial \tau_{1}}\right)_{i^{*}}\right) u_{j} \partial_{k} \\
&+\left(\Gamma_{i j}^{k^{*}}\left(\frac{\partial w}{\partial \tau_{1}}\right)_{i}+\Gamma_{i^{*} j}^{k^{*}}\left(\frac{\partial w}{\partial \tau_{1}}\right)_{i^{*}}\right) u_{j} \partial_{k^{*}} \\
& J \nabla_{\partial_{\tau_{2}}} u=J \frac{\partial u}{\partial \tau_{2}}+\left(\Gamma_{i j}^{k}\left(\frac{\partial w}{\partial \tau_{2}}\right)_{i}+\Gamma_{i^{*} j}^{k}\left(\frac{\partial w}{\partial \tau_{1}}\right)_{i^{*}}\right) u_{j} \partial_{k^{*}} \\
&-\left(\Gamma_{i j}^{k^{*}}\left(\frac{\partial w}{\partial \tau_{2}}\right)_{i}+\Gamma_{i^{*} j}^{k^{*}}\left(\frac{\partial w}{\partial \tau_{2}}\right)_{i^{*}}\right) u_{j} \partial_{k} .
\end{aligned}
$$

Since $\bar{\partial} w(\zeta)=0$, we have

$$
\begin{aligned}
\left(\frac{\partial w}{\partial \tau_{1}}\right)_{j} & =\left(\frac{\partial w}{\partial \tau_{2}}\right)_{j^{*}} \\
\left(\frac{\partial w}{\partial \tau_{1}}\right)_{j^{*}} & =-\left(\frac{\partial w}{\partial \tau_{2}}\right)_{j}
\end{aligned}
$$

This together with Lemma 5.5 shows that $\nabla_{t} Y(0)=\bar{\partial} u$.
Hence, $\bar{\partial}\left(\Psi_{1}(v, \lambda)\right)=0$ for $\zeta \in \partial D_{m}$ if and only if the same is true for $v$ provided $v$ is shorter than the minimum of injectivity radii of $g\left(L_{\lambda}, \sigma\right)$. An approximation argument together with the continuity of $\Psi_{1}$ (also in $w[M]$, see the proof of Lemma 5.10 (a)), $\bar{\partial}$, and of restriction to the boundary gives the second statement in full generality. q.e.d.

Lemma 5.13. For $r>0$ small enough, the image of $\Psi$ is a submanifold of $\mathcal{F}_{2, \epsilon, \Lambda}$. Moreover, there exists $M>0, r>0$, and a neighborhood $U$ of $(w[M], f[M], 0)$ in $\mathcal{F}_{2, \epsilon, \Lambda}$ such that $(w, f) \in U$ and $U \cap \mathcal{W}_{2, \epsilon, \Lambda}$ is contained in the image of $\Psi$ and corresponds to the subset $h=0$ in the coordinates

$$
(v, h, \lambda) \in B_{2, \epsilon}((w[M], f[M]), r) \times \mathcal{H}_{\frac{3}{2}, \epsilon}\left(\partial D_{m}, \mathbb{R}\right) \times \Lambda
$$

Proof. Let $(w, f) \in \mathcal{W}_{2, \epsilon}(\mathbf{a}, \kappa)$. Let $K$ be as in Lemma 5.11. Then $D_{2} K(x, 0, \sigma, 0) \cdot \eta=\eta$ and $D_{4} K(x, 0, \sigma, 0)=0$. Hence, the differential of $\Psi$ at $(0,0,0)$ is

$$
d \Psi(0,0,0)=\left(\begin{array}{ccc}
\iota & 0 & 0 \\
\left\langle\iota_{\mathbb{R}}\right\rangle & \text { id } & 0 \\
0 & 0 & \text { id }
\end{array}\right)
$$

where $\iota$ denotes the inclusion of the tangent space of $B_{2, \epsilon}((f, w), r)$ into $\mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right)$ and $\left\langle\iota_{\mathbb{R}}\right\rangle v$ denotes the $\mathbb{R}$-component of the vector field $\tilde{v}$ which maps to $v$ under $\Pi_{\mathbb{C}}$ and is tangent to $L_{0}$. Note that the tangent space of $B_{2, \epsilon}((w, f), r)$ is a closed subspace of $\mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right)$.

Thus, $d \Psi(0,0,0)$ is an injective map with closed image. Since the first component of $\mathcal{F}_{2, \epsilon}$ is modeled on a Banach space which allow a Hilbert-space structure, we see that the image of the differential admits a direct complement. Moreover, applying Lemma 5.10 to the explicit differential in Lemma 5.11, we conclude that the norm of the differential of $\Psi$ is Lipschitz in $v$ and $\lambda$ with Lipschitz constant depending only on $\| D w[M]_{1, \epsilon}$ and $\|D F[M]\|_{1, \epsilon}$. Hence, the implicit function theorem shows that there exists $r>0$ and $W \subset \Lambda$ (independent of $M$ ) such that the image of $B((w[M], f[M]), r) \times W$ is a submanifold. From the norm-estimates on the differential, it follows that for $M$ large enough $(w, f)$ lies in this image.

The statement about surjectivity onto $U \cap \mathcal{W}_{2, \epsilon, \Lambda}$ follows from the fact that $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ is totally geodesic in the metric $g\left(L_{\lambda}, \sigma\right)$ and Lemma 5.12. The statement on coordinates is trivial.
q.e.d.
5.6. Bundle over conformal structures. The constructions above all depend on the conformal structure $\kappa$ on $D_{m}$. This conformal structure is unique if $m \leq 3$. Assume that $m>3$ and recall that we identified the space of conformal structures $\mathcal{C}_{m}$ on $D_{m}$ with an open simplex of dimension $m-3$.

The space

$$
\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a})=\bigcup_{\kappa \in \mathcal{C}_{m}} \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa),
$$

has a natural structure of a locally trivial Banach manifold bundle, over $\mathcal{C}_{m}$. To see this, we must present local trivializations.

Let $\Delta$ denote the unit disk in the complex plane and let $\Delta_{m}$ denote the same disk with $m$ punctures $p_{1}, \ldots, p_{m}$ on the boundary and conformal structure $\kappa$. Fixing the positions of $p_{1}, p_{2}, p_{3}$, this structure is determined by the positions of the remaining $m-3$ punctures. We coordinatize a neighborhood of the conformal structure $\kappa$ in $\mathcal{C}_{m}$ as follows. Pick $m-3$ vector fields $v_{1}, \ldots, v_{m-3}$, with $v_{k}$ supported in a neighborhood of $p_{k+3}, k=1, \ldots, m-3$ in such a way that $v_{k}$ generate a 1-parameter family of diffeomorphism $\phi_{p_{k+3}}^{\tau_{k}}: \Delta \rightarrow \Delta, \tau_{k} \in \mathbb{R}$ which is a rigid rotation around $p_{k+3}$ and which is holomorphic on the boundary. Let the supports of $v_{k}$ be sufficiently small so that the supports of $\phi_{p_{k+3}}^{\tau_{k}}, k=1, \ldots, m-3$ are disjoint. Then the diffeomorphisms $\phi_{p_{3}}^{\tau_{1}}, \ldots, \phi_{p_{m}}^{\tau_{m-3}}$ all commute. Define, for $\tau=\left(\tau_{1}, \ldots, \tau_{m-3}\right) \in \mathbb{R}^{m-3}$, $\phi^{\tau}=\phi_{p_{3}}^{\tau_{1}} \circ \cdots \circ \phi_{p_{m}}^{\tau_{m-3}}$ and a local coordinate system around $\kappa$ in $\mathcal{C}_{m}$ by

$$
\tau \mapsto\left(d \phi^{\tau}\right)^{-1} \circ j_{\kappa} \circ d \phi^{\tau}
$$

These local coordinate systems give an atlas on $\mathcal{C}_{m}$.
Using this family, we define the trivialization over $\mathbb{R}^{m-3} \approx U \subset \mathcal{C}_{m}$ by composition with $\phi^{-\tau}$. That is, a local trivialization over $U$ is given by

$$
\begin{aligned}
& \Phi: \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}, \kappa) \times U \rightarrow \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}) ; \\
& \Phi(w, f, \lambda, \tau)=\left(w \circ \phi^{-\tau}, f \circ \phi^{-\tau}, \lambda, \theta\right) .
\end{aligned}
$$

In a similar way, we endow the space

$$
\mathcal{H}_{1, \epsilon}\left(D_{m}, T^{*} D_{m}\right)=\bigcup_{\kappa \in \mathcal{C}_{m}} \mathcal{H}_{1, \epsilon}\left(D_{m}, T^{*} D_{m}, g(\kappa)\right),
$$

with its natural structure as a locally trivial Banach space bundle over $\mathcal{C}_{m}$.

Representing the space of conformal structures $\mathcal{C}_{m}$ in this way, we are led to consider its tangent space $T_{\kappa} \mathcal{C}_{m}$ as generated by $\gamma_{1}, \ldots, \gamma_{m-3}$, where $\gamma_{k}=i \cdot \bar{\partial} v_{k}$, in the following sense. If $\gamma$ is any section of $\operatorname{End}\left(T D_{m}\right)$ which anti-commutes with $j_{\kappa}$ and which vanishes on the boundary, then there exists unique numbers $a_{1}, \ldots, a_{m-3}$ and a unique vector field $v$ on
$\Delta_{m}$ which is holomorphic on the boundary and which vanish at $p_{k}$, $k=1, \ldots, m$ such that

$$
\begin{equation*}
\gamma=\sum_{k} a_{k} \gamma_{k}+i \bar{\partial} v \tag{5.26}
\end{equation*}
$$

The existence of such $v$ is a consequence of the fact that the classical Riemann-Hilbert problem for the $\overline{\overline{ }}$-operator on the unit disk with tangential boundary conditions has index 3 and is surjective (the kernel being spanned by the vector fields $\left.z \mapsto i z, z \mapsto i\left(z^{2}+1\right), z \mapsto z^{2}-1\right)$.

Going from the punctured disk $\Delta_{m}$ to $D_{m}$ with our standard metric, the behavior of the vector fields $v_{j}$ close to punctures where they are supported is easily described. In fact the vector fields can be taken as $\partial_{x}$ in coordinates $z=x+i y \in\left(\mathbb{C}_{+}, \mathbb{R}, 0\right)$ in a neighborhood of the puncture $p$ on $\partial \Delta_{m}$. The change of coordinates taking us to the standard end $[0, \infty) \times[0,1]$ is $\tau+i t=\zeta=-\frac{1}{\pi} \log z$ and we see the corresponding vector field on $[0, \infty) \times[0,1]$ is $\frac{1}{\pi} e^{\pi \zeta}$ (where we identify vector fields with complex valued functions). As in Proposition 6.13, we see that equation (5.26) holds on $D_{m}$ with $v$ in a Sobolev space with (small) negative exponential weights at the punctures.

Consider the space $\mathcal{H}_{1, \epsilon}\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ and the closed subspace $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ consisting of elements whose trace (restriction to the boundary) is 0 . The elements of this space are complex anti-linear maps $T D_{m} \rightarrow \mathbb{C}^{n}$ and so depend on the complex structure $j_{\kappa}$ on $D_{m}$. For simplicity, we keep the notation and consider $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ as a bundle over $\mathcal{C}_{m}$. To do this, we must find local trivializations of this bundle. To this end, we note that any complex structure $j_{\mu}$ in a neighborhood of a given complex structure $j_{\kappa}$ on $D_{m}$ can be written as $j_{\mu}=j_{\kappa}\left(\mathrm{id}+\gamma_{\mu}\right)\left(\mathrm{id}-\gamma_{\mu}\right)^{-1}$ where $\gamma$ is a section of $\operatorname{End}\left(T D_{m}\right)$ such that $\gamma \circ j_{\kappa}+j_{\kappa} \circ \gamma=0$. It is then easy to check that the map which takes $A \in \operatorname{Hom}\left(T D_{m}, \mathbb{C}^{n}\right)$ to $A \circ\left(\mathrm{id}+\gamma_{\mu}\right)$ identifies the ( $i, j_{\kappa}$ ) anti-linear maps with the $\left(i, j_{\mu}\right)$ anti-linear maps. We thus trivialize the bundle $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ over $U \subset \mathcal{C}_{m}$ around $\kappa \in U$ by taking the $\left(i, j_{\kappa}\right)$ anti-linear section $A$ to the $\left(i, j_{\mu}\right)$ anti-linear map $A\left(1+\gamma_{\mu}\right)$.

Note finally that in our local coordinates on $\mathcal{C}_{m}$ from above, we have $d \phi^{-\tau} \circ j_{\kappa} \circ d \phi^{\tau}=j_{\kappa}\left(1+\gamma_{\tau}\right)\left(1-\gamma_{\tau}\right)^{-1}$, where $\gamma_{\tau}=-\frac{1}{2} \bar{\partial} v$ up to first order in $\tau$.
5.7. The $\bar{\partial}$-map and its linearization. Consider the bundle $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ over $\mathcal{C}_{m}$ as in the previous section. We extend this bundle to a bundle over $\Lambda$ making it trivial in the $\Lambda$ directions and denote the result $\mathcal{H}_{1, \epsilon, \Lambda}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$.

The $\bar{\partial}$-map is the map

$$
\begin{aligned}
& \hat{\Gamma}: \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{a}) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right) ; \\
& \hat{\Gamma}(w, f, \kappa, \lambda)=\left(d w+i \circ d w \circ j_{\kappa}, \kappa, \lambda\right) .
\end{aligned}
$$

We will denote the first component of this map simply $\Gamma$. An element ( $w, f, \kappa, \lambda$ ) is thus holomorphic with respect to the complex structure $j_{\kappa}$ if and only if $\hat{\Gamma}(w, f, \kappa, \lambda)=(0, \kappa, \lambda)$. Hence, if $L_{\lambda}, \lambda \in \Lambda$ is a family of chord generic Legendrian submanifolds, then the (parameterized) moduli-space of holomorphic disks with boundary on $L_{\lambda}$, positive puncture at $a_{1}$, and negative punctures at $a_{2}, \ldots, a_{m}$ is naturally identified with the preimage under $\hat{\Gamma}$ of the 0 -section in $\mathcal{H}_{1, \epsilon, \Lambda}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes\right.$ $\mathbb{C}^{n}$ ) for sufficiently small $\epsilon \in[0, \infty)$.

We compute the linearization of the $\bar{\partial}$-map. As in Section 5.6, we think of tangent vectors $\gamma$ to $\mathcal{C}_{m}$ at $\kappa$ as sections of $\operatorname{End}\left(T D_{m}\right)$. For $\kappa \in \mathcal{C}_{m}$ and $u: D_{m} \rightarrow \mathbb{C}^{n}$, let $\bar{\partial}_{\kappa} u=d u+i \circ d u \circ j_{\kappa}$ and let $\partial_{\kappa} u=$ $d u-i \circ d u \circ j_{\kappa}$.

Let $(w, f, \kappa, 0) \in \mathcal{W}_{2,,, \Lambda}(\mathbf{a})$. Identify the tangent space of $\mathcal{W}_{2, \epsilon, \lambda}(\mathbf{a})$ at $(w, f, \kappa, 0)$ with $T \mathcal{B}_{2, \epsilon}((w, f), r) \times T_{\kappa} \mathcal{C}_{m} \times T_{0} \Lambda$.

Lemma 5.14. The differential of $\Gamma$ at $(w, f, \kappa, 0)$ is the map

$$
\begin{equation*}
d \Gamma(v, \gamma, \lambda)=\bar{\partial}_{\kappa} v+\bar{\partial}_{\kappa}\left(Y_{0}^{\sigma}(w, \lambda)\right)+i \circ \partial_{\kappa} w \circ \gamma \tag{5.27}
\end{equation*}
$$

Recall $Y_{0}^{\sigma}$ was defined in (5.15).
Proof. Assume first $w$ and $f$ are constant close to punctures. Let $\mathcal{B}_{2, \epsilon}((w, f), r) \times \mathcal{C}_{m} \times \Lambda$ be a local coordinates around $(w, f, \kappa, 0)$.

Let $K(x, \xi, \sigma, \lambda)=\psi_{\lambda}^{\sigma}(x)+\xi$. Then

$$
R(x, \xi, \sigma, \lambda)=\exp _{\psi_{\lambda}^{\sigma}(x)}^{\lambda, \sigma} A_{\lambda}^{\sigma} \xi-K(x, \xi, \sigma, \lambda)
$$

satisfies

$$
R(x, 0, \sigma, \lambda)=0, \quad D_{2} R(x, 0, \sigma, 0)=0, \quad D_{4} R(x, 0, \sigma, 0)=0
$$

thus, Lemma 5.10 (b) implies that

$$
\|R(w, v, \sigma, \lambda)\|_{2, \epsilon} \leq C\left(\|v\|_{2, \epsilon}^{2}+|\lambda|^{2}\right)
$$

Continuity of the linear operators

$$
\bar{\partial}_{\kappa+\gamma}: \mathcal{H}_{2, \epsilon}\left(D_{m}, \mathbb{C}^{n}\right) \rightarrow \mathcal{H}_{1, \epsilon}\left(D_{m}, T^{*} D_{m} \otimes \mathbb{C}^{n}\right)
$$

where we use local coordinates $\mathbb{R}^{m-3}$ on $\mathcal{C}_{m}, \kappa+\gamma \in \mathbb{R}^{m-3} \subset \mathcal{C}_{m}$, and the trivialization of $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ described in subsection 5.6, shows that

$$
\begin{equation*}
\left\|\bar{\partial}_{\kappa+\gamma} R(w, v, \sigma, \lambda)\right\|_{1, \epsilon} \leq C\left(\|v\|_{2, \epsilon}^{2}+|\lambda|^{2}\right) \tag{5.28}
\end{equation*}
$$

It is straightforward to check that

$$
\begin{align*}
& \left\|\bar{\partial}_{\kappa+\gamma} K(w, v, \kappa+\mu, \lambda)-\bar{\partial}_{\kappa} w-\left(\bar{\partial}_{\kappa} v+\bar{\partial}_{\kappa}\left(Y_{0}(w, \lambda)\right)+i \circ \partial_{\kappa} w \circ \gamma\right)\right\|_{1, \epsilon}  \tag{5.29}\\
& \quad \leq C\left(\|v\|_{2, \epsilon}^{2}+|\lambda|^{2}+|\gamma|^{2}\right) .
\end{align*}
$$

Equations (5.29) and (5.28) imply the lemma in the special case when $(w, f)$ is constant close to punctures (and in the general case if $\operatorname{dim}(\Lambda)=$ $0)$.

If $(w, f)$ is not constant close to punctures, consider the maps $(w[M]$, $f[M])$ which are constant close to punctures. We have $(w[M], f[M]) \rightarrow$ $(w, f)$ as $M \rightarrow \infty$. Since the local coordinates are $C^{1}$ a limiting argument proves (5.27) in the general case. q.e.d.
5.8. Auxiliary spaces in the semi-admissible case. In Section 7.9, we show that for a dense open set of semi-admissible Legendrian submanifolds $L$, no rigid holomorphic disks with boundary on $L$ have exponential decay at their degenerate corners. Once this has been shown, we know that if 0 is the degenerate corner and $L$ has the form (3.4) around 0 , then for any rigid holomorphic disk $u: D_{m} \rightarrow \mathbb{C}^{n}$ with puncture $p$ mapping to 0 , there exists $M>0$ and $c \in \mathbb{R}$ such that

$$
u(\zeta)=\left(-2(\zeta+c)^{-1}, 0, \ldots, 0\right)+\mathcal{O}\left(e^{-\theta|\zeta|}\right), \text { for } \zeta \in E_{p}[ \pm M]
$$

where $\theta>0$ is the smallest non-zero complex angle of the Reeb chord at 0 . (Here we implicitly assume that $P_{2}$ in our standard self tangency model lies above $P_{1}$ in the $z$-direction, and that neighborhoods of positive (negative) punctures are parameterized by $[1, \infty) \times[0,1]$ $((-\infty,-1] \times[0,1])$.) To study disks of this type, we use the following construction.

Let $a_{0}$ denote the Reeb chord at 0 . Assume that a has the Reeb chord $a_{0}$ in $k$ positions. For $C=\left(c_{1}, \ldots, c_{k}\right) \in \mathbb{R}^{k}$, fix a smooth reference function which equals

$$
u_{\mathrm{ref}}^{C}(\zeta)=\left(-2\left(\zeta+c_{j}\right)^{-1}, 0, \ldots, 0\right)
$$

in a neighborhood of the $j^{\text {th }}$ puncture mapping to $a_{0}$ and also a smooth function $F_{\text {ref }}^{C}: \partial D_{m} \rightarrow \mathbb{R}$ so that $\left(u_{\text {ref }}, F_{\text {ref }}\right) \mid \partial D_{m}$ maps to $L$.

Let $L_{\lambda}, \lambda \in \Lambda$ be a family of semi-admissible Legendrian submanifolds. We construct for $\epsilon \in[0, \infty)^{m}$, with those components $\epsilon_{j}$ of $\epsilon$ which correspond to punctures mapping to the degenerate corner satisfying $0<\epsilon_{j}<\theta$ and for fixed $C \in \mathbb{R}^{k}$, the spaces

$$
\mathcal{F}_{2, \epsilon}^{C}(\mathbf{a})
$$

by using reference functions looking like $u_{\text {ref }}^{C}$ for $C \in \mathbb{R}^{k}$ in neighborhoods of punctures mapping to $a_{0}^{*}=0 \in \mathbb{C}^{n}$. We construct local coordinates as in Section 5.1 taking advantage of the fact that $\lambda \in \Lambda$ fixes
$a_{0}^{*}$. Also we consider the space

$$
\mathcal{W}_{2, \epsilon, \Lambda}^{C}(\mathbf{a}),
$$

which is defined in the same way as before. We note that the construction giving local coordinates on this space in Section 5.5 can still be used since in the semi-admissible case, we need not cut-off the first component of $w$ in $(w, f)$ close to punctures mapping to $c_{0}$ since $\lambda \in \Lambda$ are assumed to preserve the product structure and $\gamma_{1}$ and $\gamma_{2}$.

With this done, we consider the bundle

$$
\begin{equation*}
\widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}=\bigcup_{C \in \mathbb{R}^{k}} \mathcal{W}_{2, \epsilon, \Lambda}^{C}(\mathbf{a}) \tag{5.30}
\end{equation*}
$$

which is a locally trivial bundle over $\mathbb{R}^{k}$.
In the case that a has $\geq 3$ elements, we fix for $C \in \mathbb{R}^{k}$ the diffeomorphism $\phi^{C}: D_{m} \rightarrow D_{m}$ which equals to $\zeta \rightarrow \zeta+c_{j}$ in $E_{p_{j}}[M]$ for any puncture $p_{j}$ mapping to $c_{0}$, equals the identity on $D_{m} \backslash \bigcup_{j} E_{p_{j}}[M-2]$, and is holomorphic on the boundary. (Since we often reduce the few punctured cases to the many punctured case, see Section 8.6, the following two constructions will not be used in the sequel, we add them here for completeness.) In case a has length 1 , we think of $D_{1}$ as of the upper half-plane $\mathbb{C}_{+}$with the puncture at $\infty$. The map $z \mapsto-\frac{1}{\pi} \log z$ identifies the region $\left\{z \in \mathbb{C}_{+}:|z|>R\right\}$ with the strip $\left[\frac{1}{\pi} \log R, \infty\right) \times[0,1]$ where we think of the latter space as a part of $E_{p}$, where $p$ is the puncture of $D_{1}$. Also, this map takes the conformal reparameterization $z \mapsto e^{\pi C} z$ to $\phi^{C}: \zeta \mapsto \zeta+C$ in $E_{p}$ and we identify $\mathbb{R}$ with this set of conformal reparameterizations $\left\{\phi^{C}\right\}_{C \in \mathbb{R}}$. In case a has length 2 , we think of $D_{2}$ as the strip $\mathbb{R} \times[0,1]$ and identify $\mathbb{R}$ with the conformal reparameterizations $\phi^{C}(\zeta)=\zeta \mapsto \zeta+C$.

We construct local coordinates

$$
\Phi: B_{2, \epsilon}(0, r) \times \mathcal{C}_{m} \times \mathbb{R}^{k} \times \Lambda \rightarrow \widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}
$$

in a neighborhood of $u_{\text {ref }}^{C}$. For fixed $(v, c) \in B_{2, \epsilon}(0, r) \times \mathbb{R}^{k}, \Phi$ is depends on $\kappa \times \Lambda$ exactly as above. We therefore fix $(\kappa, 0) \in \mathcal{C}_{m} \times \Lambda$ and describe the dependence of the remaining factors. For $c \in \mathbb{R}^{k}$ in a neighborhood of $C$, define the map $A_{c, C}: D_{m} \rightarrow \operatorname{End}\left(C^{n}\right)$,

$$
\left[A_{c, C}(\zeta)\right]\left(v_{1}, \ldots, v_{n}\right)= \begin{cases}\left(v_{1}, \ldots, v_{n}\right) & \text { if } \zeta \notin E_{p_{j}} \\ \left(d\left(u_{\mathrm{ref}}^{c}\right)_{1} \circ d\left(u_{\mathrm{ref}}^{C}\right)_{1}^{-1} v_{1}, v_{2}, \ldots, v_{n}\right) & \text { if } \zeta \in E_{p_{j}}\end{cases}
$$

where $\left(u_{\text {ref }}^{c}\right)_{1}$ denotes the first component of $u_{\text {ref }}^{c}$, and let

$$
\Phi(v, c)=\exp _{u_{\mathrm{ref}}(\zeta)}^{\sigma(F(\zeta))}\left(A_{c, C}(\zeta) v(\zeta)\right)
$$

Note that if $w: D_{m} \rightarrow \mathbb{C}^{n}$ is any holomorphic disk with boundary on $L$ which is asymptotic to some $u_{\text {ref }}^{C}$ at the degenerate corner, then we may
define coordinates on a neighborhood of $w$ in $\tilde{\mathcal{W}}_{2, \epsilon}$ by replacing $u_{\text {ref }}^{c}$ in the above formulas by $w \circ \phi^{c}$.

Using these coordinates, we find that the linearization of the $\bar{\partial}$-map $\Gamma$ at a holomorphic ( $w, f, \kappa, 0,0$ ) equals

$$
\begin{align*}
d \Gamma(v, \gamma, c, \lambda)= & \bar{\partial}_{\kappa} v+\bar{\partial}_{\kappa}\left(\left(Y_{0}(w, \lambda)\right)\right)  \tag{5.31}\\
& +i \circ \partial_{\kappa} w \circ \gamma+\bar{\partial}_{\kappa}\left(d w \cdot\left(\left.\frac{\partial \phi^{C}}{\partial C}\right|_{C=0}\right) \cdot c\right) .
\end{align*}
$$

Here $c=\left(c_{1}, \ldots, c_{k}\right)$ is a tangent vector to $\mathbb{R}^{k}$ written in the basis $\left\{\hat{C}_{1}, \ldots, \hat{C}_{k}\right\}$ where $\hat{C}_{j}$ is a unit vector in the tangent space to $C_{j} \in \mathbb{R}$. We notice that the second term in (5.31) lies in $\mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m}\right)$ because of the special assumptions on $L_{\lambda}$ in a neighborhood of $c_{0}^{*}$ and that the last term does as well since the holomorphicity of $w$ allows us to control its higher derivatives and since it vanishes in the region where $\phi^{C}$ is just a translation.
5.9. Homology decomposition. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a (semi-)admissible Legendrian submanifold. Let $\mathbf{c}=c_{0} c_{1} \ldots c_{m}$ be a word of Reeb chords of $L$. If $(u, f) \in \mathcal{W}_{2, \epsilon}(\mathbf{c})$, then the homotopy classes of the paths induced by $\left(u \mid \partial D_{m}, f\right)$ in $L$ connecting the Reeb chord endpoints determines the path component of $(u, f) \in \mathcal{W}_{2, \epsilon}(\mathbf{c})$.

Let $A \in H_{1}(L)$ and let $\mathcal{W}_{2, \epsilon}(\mathbf{c} ; A) \subset \mathcal{W}_{2, \epsilon}(\mathbf{c})$ be the union of those path components of $\mathcal{W}_{2, \epsilon}(\mathbf{c})$ such that the homology class of the loop $f\left(\partial D_{m}\right) \cup\left(\bigcup_{j} \gamma_{j}\right)$ equals $A$, where $\gamma_{j}$ is the capping path chosen for the Reeb chord $c_{j}$ endowed with the appropriate orientation, see Section 2.3. For fixed conformal structure $\kappa$, we write $\mathcal{W}_{2, \epsilon}(\mathbf{c}, \kappa ; A)$ and in the chord semi generic case $\widetilde{\mathcal{W}}_{2, \epsilon}(\mathbf{c} ; A)$ and interpret these notions in a similar way.

## 6. Fredholm properties of the linearized equation

In this section, we study properties of the linearized $\bar{\partial}$-equation. In particular, we determine the index of the $\bar{\partial}$-operator with Legendrian boundary conditions. It will be essential for our geometric applications to use weighted Sobolev spaces and to understand how constants in certain elliptic estimates depend on the weights.

Our presentation has two parts: the "model" case where the domain is a strip or half-plane; and the harder case where the domain is $D_{m}$.

In Section 6.1, we discuss the existence of smooth representatives of cokernel elements. In Section 6.2, we derive expansions for the kernel and cokernel elements. We use these two subsections in Sections 6.3 through 6.5, to prove the elliptic estimate for the model problem, as well as derive a formula for the index. In Sections 6.6, we set up the boundary conditions for the linearized problem with domain $D_{m}$. In Sections 6.7 through 6.10, we prove the Fredholm properties for the $D_{m}$
case. In Sections 6.10 and 6.11 , we connect the index formula to the Conley-Zehnder index of Section 2.
6.1. Cokernel regularity. To control the cokernels of the operators studied below, we use the following regularity lemma.

For this subsection only, we use coordinates $(x, y)$ for the half-plane $\mathbb{R}_{+}^{2}=\{(x, y): y \geq 0\}$. Let $A: \mathbb{R} \rightarrow \mathbf{G L}\left(\mathbb{C}^{n}\right)$ be a smooth map with $\operatorname{det}(A)$ uniformly bounded away from 0 and all derivatives uniformly bounded. We also simplify notation for this subsection only and define the following Sobolev spaces: let $\mathcal{H}_{k}=\mathcal{H}_{k}\left(\mathbb{R}^{2}, \mathbb{C}^{n}\right)$; let $\overline{\mathcal{H}}_{k}$ denote the space of restrictions of elements in $\mathcal{H}_{k}$ to $\operatorname{int}\left(\mathbb{R}_{+}^{2}\right)$; let $\dot{\mathcal{H}}_{k}$ denote the subspace of elements in $\mathcal{H}_{k}$ with support in $\mathbb{R}_{+}^{2}$; let $\overline{\mathcal{H}}_{1}[0]$ denote the subspace of all elements in $\overline{\mathcal{H}}_{1}$ which vanish on the boundary; and let $\overline{\mathcal{H}}_{2}[A]$ denote the subspace of elements $u$ in $\overline{\mathcal{H}}_{2}$ such that $u(x, 0) \in$ $A(x) \mathbb{R}^{n}$ and such that the trace of $\bar{\partial} u$ (its restriction to the boundary) equals 0 in $\mathcal{H}_{\frac{3}{2}}\left(\mathbb{R}, \mathbb{C}^{n}\right)$.

An element $\xi$ in the cokernel of $\bar{\partial}$ will be in the dual space of $\mathcal{H}_{1}[0]$. The dual of $\mathcal{H}_{1}$ is $\mathcal{H}_{-1}$ and thus the dual of $\mathcal{H}_{1}[0]$ is the quotient space

$$
\begin{equation*}
\mathcal{H}_{-1} / \mathcal{H}_{1}[0]^{\perp} \tag{6.1}
\end{equation*}
$$

where $\mathcal{H}_{1}[0]^{\perp}$ denotes the annihilator of $\mathcal{H}_{1}[0]$ in $\mathcal{H}_{-1}$. As usual, let $\langle$, denote the standard Riemannian inner product on $\mathbb{C}^{n} \approx \mathbb{R}^{n}$.

Lemma 6.1. Fix $h>0$ and assume that $v \in \dot{\mathcal{H}}_{-1}$ satisfies

$$
\begin{equation*}
\int_{\mathbb{R} \times[0, h)}\langle\bar{\partial} u, v\rangle d x \wedge d y=0 \tag{6.2}
\end{equation*}
$$

for all $u \in \overline{\mathcal{H}}_{2}[A]$ with compact support in $\mathbb{R} \times[0, h)$. Then, for every $\epsilon$ with $0<\epsilon<h$ and every $k>0$, the class $[v] \in \dot{\mathcal{H}}_{-1} / \overline{\mathcal{H}}_{1}[0]^{\perp}$ of $v$ contains an element $v_{0}$ which is $C^{k}$ in $\mathbb{R} \times[0, \epsilon)$, up to and including the boundary.

Proof. The proof is standard and therefore omitted. q.e.d.
6.2. Kernel and cokernel elements. Consider the strip $\mathbb{R} \times[0,1] \subset$ $\mathbb{C}$ endowed with the standard flat metric, the corresponding complex structure and coordinates $\zeta=\tau+i t$. For $k \geq 0$, let

$$
\mathcal{H}_{k}=\mathcal{H}_{k}\left(\mathbb{R} \times[0,1], \mathbb{C}^{n}\right)
$$

and for $k \leq 0$, let $\mathcal{H}_{k}$ denote the $L^{2}$-dual of $\mathcal{H}_{-k}$. We also use the notions $\mathcal{H}_{k}^{\text {loc }}$ which are to be understood in the corresponding way.

If $u \in \mathcal{H}_{k}^{\text {loc }}$, then the restriction of $u$ to $\partial(\mathbb{R} \times[0,1])=\mathbb{R} \cup \mathbb{R}+i$ lies in $\mathcal{H}_{k-\frac{1}{2}}^{\text {loc }}\left(\mathbb{R} \cup \mathbb{R}+i, \mathbb{C}^{n}\right)$. For $u \in \mathcal{H}_{1}^{\text {loc }}$, consider the boundary conditions

$$
\begin{array}{cl}
\int_{\mathbb{R}}\langle u, v\rangle d \tau=0 & \text { for all } v \in C_{0}^{\infty}\left(\mathbb{R}, i \mathbb{R}^{n}\right) \\
\int_{\mathbb{R}+i}\langle u, v\rangle d \tau=0 & \text { for all } v \in C_{0}^{\infty}\left(\mathbb{R}+i, \mathbb{R}^{n}\right) \tag{6.4}
\end{array}
$$

Let $f: \mathbb{R} \times[0,1] \rightarrow \mathbb{C}^{n}$ be a smooth function satisfying (6.3) and (6.4). Define the function $f^{d}: \mathbb{R} \times[0,2] \rightarrow \mathbb{C}^{n}$ as

$$
f^{d}(\tau+i t)= \begin{cases}f(\tau+i t) & \text { for } 0 \leq t \leq 1 \\ -\bar{f}(\tau+i(2-t)) & \text { for } 1<t \leq 2\end{cases}
$$

where $\bar{w}$ denotes the complex conjugate of $w \in \mathbb{C}^{n}$. Then, $f^{d}$ and $\partial_{\tau} f^{d}$ are continuous, $\partial_{t} f^{d}$ may have a jump discontinuity over the line $\mathbb{R}+i$, $f^{d}(\tau+0 i)=-f^{d}(\tau+2 i)$, and $\left\|f^{d}\right\|_{1}=2\|f\|_{1}$. Hence, we can define the double $u^{d} \in \mathcal{H}_{1}^{\text {loc }}(\mathbb{R} \times[0,2])$ of any $u \in \mathcal{H}_{1}^{\text {loc }}$ which satisfies (6.3) and (6.4). For $u \in \mathcal{H}_{k}^{\text {loc }}$, let $\bar{\partial} u=\left(\partial_{\tau}+i \partial_{t}\right) u$ and $\partial u=\left(\partial_{\tau}-i \partial_{t}\right) u$.

Lemma 6.2. If $u \in \mathcal{H}_{1}^{\text {loc }}$ satisfies (6.3) and (6.4) and
(a) $\bar{\partial} u=0$ in the interior of $\mathbb{R} \times[0,1]$ then

$$
u(\zeta)=\sum_{n \in \mathbb{Z}} C_{n} \exp \left(\left(\frac{\pi}{2}+n \pi\right) \zeta\right),
$$

where $C_{n} \in \mathbb{R}$.
(b) $\partial u=0$ in the interior of $\mathbb{R} \times[0,1]$ then

$$
u(\zeta)=\sum_{n \in \mathbb{Z}} C_{n} \exp \left(\left(\frac{\pi}{2}+n \pi\right) \bar{\zeta}\right),
$$

where $C_{n} \in \mathbb{R}$.
Moreover, if $u$ satisfies (a) or (b) and $u \in \mathcal{H}_{k}$ for some $k \in \mathbb{Z}$, then $u=0$.

Proof. We prove (a), (b) is proved in the same way. Clearly, it is enough to consider one coordinate at a time. So assume the target is $\mathbb{C}$ and let $u$ be as in the statement.

Consider $u^{d}$, then $\bar{\partial} u^{d}$ is an element of $\mathcal{H}_{0}^{\text {loc }}(\mathbb{R} \times[0,2], \mathbb{C})$ with support on $\mathbb{R}+i \cup \partial(\mathbb{R} \times[0,2])$. Such a distribution is a three-term linear combination of tensor products of a Dirac-delta in the $t$-variable and a distribution on $\mathbb{R}$ and hence lies in $\mathcal{H}_{0}(\mathbb{R} \times[0,2], \mathbb{C})$ only if it is zero. Thus $\bar{\partial} u=0$ and we may use elliptic regularity to conclude that $u$ is smooth in the interior of $\mathbb{R} \times[0,2]$. (In fact, doubling again and using the same argument, we find that $u$ is smooth also on the boundary.)

We may now Fourier expand $u^{d}(\tau, \cdot)$ in the eigenfunctions $\phi$ of the operator $i \partial_{t}$ which satisfy the boundary condition $\phi(0)=-\phi(2)$. These eigenfunctions are

$$
t \mapsto \exp \left(i\left(\frac{\pi}{2}+n \pi\right) t\right), \text { for } n \in \mathbb{Z}
$$

We find

$$
u^{d}=\sum_{n} c_{n}(\tau) \exp \left(i\left(\frac{\pi}{2}+n \pi\right) t\right)
$$

where, by the definition of $u^{d}, c_{n}(\tau)$ are real valued functions and

$$
\bar{\partial} u^{d}=\sum_{n}\left(c_{n}^{\prime}(\tau)-\left(\frac{\pi}{2}+n \pi\right) c_{n}(\tau)\right) \exp \left(i\left(\frac{\pi}{2}+n \pi\right) t\right)
$$

Hence,

$$
u(\zeta)=\sum_{n} C_{n} \exp \left(\left(\frac{\pi}{2}+n \pi\right) \zeta\right) .
$$

Assume that $u \in \mathcal{H}_{k}$ for some $k \in \mathbb{Z}$. Then, since for $j \geq 0$ the restriction of any $v \in \mathcal{H}_{j}(\mathbb{R} \times[0,2], \mathbb{C})$ to $\mathbb{R} \times[0,1]$ lies in $\mathcal{H}_{j}$,

$$
\lambda_{u}(v)=\int_{\mathbb{R} \times[0,2]}\left\langle v, u^{d}\right\rangle d \tau \wedge d t
$$

is a continuous linear functional on $\mathcal{H}_{j}(\mathbb{R} \times[0,2], \mathbb{C})$ for $j=k$ if $k \geq 0$ or $j=-k$ if $k<0$.

Let $\psi: \mathbb{R} \rightarrow[0,1]$ be a smooth function equal to 1 on $[0,1]$ and 0 outside $[-1,2]$. For $n, r \in \mathbb{Z}$ let

$$
\alpha_{n, r}(\tau+i t)=\psi(\tau+r) \exp \left(i\left(\frac{\pi}{2}+n \pi\right) t\right)
$$

Then, $\alpha_{n, r} \in \mathcal{H}_{j}(\mathbb{R} \times[0,2], \mathbb{C})$ and $\left\|\alpha_{n, r}\right\|_{j}=K(n)$ for some constant $K(n)$ and all $r$. It is straightforward to see that

$$
\lambda_{u}\left(\alpha_{n, r}\right)=2 C_{n} \int_{r-1}^{r+2} \psi(\tau+r) \exp \left(\left(\frac{\pi}{2}+n \pi\right) \tau\right) d \tau=l_{n, r}
$$

The set $\left\{l_{n, r}\right\}_{r \in \mathbb{Z}}$ is unbounded unless $C_{n}=0$. Hence, $\lambda_{u}$ is continuous only if each $C_{n}=0$.
q.e.d.
6.3. The right angle model problem. As mentioned, we will use weighted Sobolev spaces. The weight functions are functions on $\mathbb{R} \times[0,1]$ which are independent of $t$ and have the following properties.

For $a=\left(a^{+}, a^{-}\right) \in \mathbb{R}^{2}$ and $\theta \in[0, \pi)$, let

$$
\begin{equation*}
m(\theta, a)=\min \left\{\left|n \pi+\theta+a^{+}\right|,\left|n \pi+\theta+a^{-}\right|\right\}_{n \in \mathbb{Z}} \tag{6.5}
\end{equation*}
$$

For $a \in \mathbb{R}^{2}$ with $m\left(\frac{\pi}{2}, a\right)>0$, let $e_{a}: \mathbb{R} \rightarrow \mathbb{R}$ be a smooth positive function with the following properties:
P1. There exists $M>0$ such that $e_{a}(\tau)=e^{a^{+} \tau}$ for $\tau \geq M$ and $e_{a}(\tau)=e^{a^{-} \tau}$ for $\tau \leq-M$.
$\mathbf{P} 2$. The logarithmic derivative of $e_{a}, \alpha(\tau)=\frac{e_{a}^{\prime}(\tau)}{e_{a}(\tau)}$, is (weakly) monotone and $\alpha^{\prime}(\tau)=0$ if and only if $\alpha(\tau)$ equals the global maximum or minimum of $\alpha$.
P3. The derivative of $\alpha$ satisfies $\left|\alpha^{\prime}(\tau)\right|<\frac{1}{5} m\left(\frac{\pi}{2}, a\right)^{2}$ for all $\tau \in \mathbb{R}$.
Let

$$
\mu=\left(\mu_{1}, \ldots, \mu_{n}\right)=\left(\mu_{1}^{+}, \mu_{1}^{-}, \ldots, \mu_{n}^{+}, \mu_{n}^{-}\right) \in \mathbb{R}^{2 n}
$$

be such that $m\left(\frac{\pi}{2}, \mu_{j}\right)>0$, for $j=1, \ldots, n$. Define the $(n \times n)$-matrix valued function $\mathbf{e}_{\mu}$ on $\mathbb{R}$ as

$$
\mathbf{e}_{\mu}(\tau)=\operatorname{Diag}\left(e_{\mu_{1}}(\tau), \ldots, e_{\mu_{n}}(\tau)\right)
$$

Define the weighted Sobolev spaces

$$
\mathcal{H}_{k, \mu}=\left\{u \in \mathcal{H}_{k}^{\text {loc }}: \mathbf{e}_{\mu} u \in \mathcal{H}_{k}\right\}, \text { with norm }\|u\|_{k, \mu}=\left\|\mathbf{e}_{\mu} u\right\|_{k}
$$

To make the doubling operation used in Section 6.2 work on $\mathcal{H}_{2}$, we impose further boundary conditions. If $u \in \mathcal{H}_{1}^{\text {loc }}$, then its trace lies in $\mathcal{H}_{\frac{1}{2}}^{\text {loc }}\left(\mathbb{R} \cup \mathbb{R}+i, \mathbb{C}^{n}\right)$. We say that $u$ vanishes on the boundary if

$$
\begin{equation*}
\int_{\mathbb{R} \cup \mathbb{R}+i}\langle u, v\rangle d \tau=0 \text { for every } v \in \mathbb{C}_{0}^{\infty}\left(\mathbb{R} \cup \mathbb{R}+i, \mathbb{C}^{n}\right) \tag{6.6}
\end{equation*}
$$

Define

$$
\begin{gathered}
\mathcal{H}_{2, \mu}(\underbrace{\frac{\pi}{2}, \ldots, \frac{\pi}{2}}_{n})=\left\{u \in \mathcal{H}_{2, \mu}: u\right. \text { satisfies (6.3), (6.4), } \\
\quad \text { and } \bar{\partial} u \text { satisfies }(6.6)\}, \\
\mathcal{H}_{1, \mu}[0]=\left\{u \in \mathcal{H}_{1, \mu}: u \text { satisfies }(6.6)\right\} .
\end{gathered}
$$

Proposition 6.3. If $m\left(\frac{\pi}{2}, \mu_{j}\right)>0$ for $j=1, \ldots, n$ then the operator

$$
\bar{\partial}: \mathcal{H}_{2, \mu}\left(\frac{\pi}{2}, \ldots, \frac{\pi}{2}\right) \rightarrow \mathcal{H}_{1, \mu}[0]
$$

is Fredholm with index

$$
\sum_{j=1}^{n} \sharp\left(-\frac{\mu_{j}^{-}}{\pi}-\frac{1}{2},-\frac{\mu_{j}^{+}}{\pi}-\frac{1}{2}\right)-\sharp\left(\frac{\mu_{j}^{-}}{\pi}-\frac{1}{2}, \frac{\mu_{j}^{+}}{\pi}-\frac{1}{2}\right)
$$

where $\sharp(a, b)$ denotes the number of integers in the interval $(a, b)$.
Moreover, if $\mu_{j}^{+}=\mu_{j}^{-}$for all $j$ and $M(\mu)=\min \left\{m\left(\frac{\pi}{2}, \mu_{1}\right), \ldots\right.$, $\left.m\left(\frac{\pi}{2}, \mu_{n}\right)\right\}$, then $u \in \mathcal{H}_{2, \mu}\left(\frac{\pi}{2}, \ldots, \frac{\pi}{2}\right)$ satisfies

$$
\begin{equation*}
\|u\|_{2, \mu} \leq C(\mu)\|\bar{\partial} u\|_{1, \mu} \tag{6.7}
\end{equation*}
$$

where $C(\mu) \leq \frac{K}{M(\mu)}$, for some constant $K$.
Proof. The problem studied is split and it is clearly sufficient to consider the case $n=1$. We first determine the dimensions of the kernel and cokernel. It is immediate from Lemma 6.2 that the kernel of $\bar{\partial}$ is finite dimensional on $\mathcal{H}_{2, \mu}\left(\frac{\pi}{2}\right)$ and that the number of linearly independent solutions is exactly $\sharp\left(-\frac{\mu^{-}}{\pi}-\frac{1}{2},-\frac{\mu^{+}}{\pi}-\frac{1}{2}\right)$.

Recall that an element in the cokernel of $\bar{\partial}$ is an element $\xi$ in the dual space of $\mathcal{H}_{1, \mu}[0]$. The dual of $\mathcal{H}_{1, \mu}$ is $\mathcal{H}_{-1,-\mu}$ and thus, as in (6.1), the dual of $\mathcal{H}_{1, \mu}[0]$ is the quotient space

$$
\mathcal{H}_{-1,-\mu} / \mathcal{H}_{1, \mu}[0]^{\perp} .
$$

Lemma 6.1 implies that any element in the cokernel has a smooth representative. Let $v$ be a smooth representative. Then

$$
\int_{\mathbb{R} \times[0,1]}\langle\bar{\partial} u, v\rangle d \tau \wedge d t=0
$$

for any smooth compactly supported function $u$ which meets the boundary conditions (6.3), (6.4), and (6.6). Using partial integration, we conclude

$$
\begin{equation*}
\int_{\mathbb{R} \times[0,1]}\langle u, \partial v\rangle d \tau \wedge d t=0 \tag{6.8}
\end{equation*}
$$

Thus $\partial v=0$ in the interior. Noting that for any two functions $\phi_{0} \in$ $C_{0}^{\infty}(\mathbb{R}, \mathbb{R})$ and $\phi_{1} \in C_{0}^{\infty}(\mathbb{R}, i \mathbb{R})$, there exists a function $u \in \mathbb{C}_{0}^{\infty}(\mathbb{R} \times$ $[0,1], \mathbb{C})$ such that $\bar{\partial} u|\partial(\mathbb{R} \times[0,1])=0, u| \mathbb{R}=\phi_{0}$, and $u \mid \mathbb{R}+i=\phi_{1}$, we find that $i v$ satisfies (6.3) and (6.4). Lemma 6.2 then implies that the cokernel has dimension $\#\left(\frac{\mu_{j}^{-}}{\pi}-\frac{1}{2}, \frac{\mu_{j}^{+}}{\pi}-\frac{1}{2}\right)$.

We now prove that the image of $\bar{\partial}$ is closed, and in doing so also establish (6.7). Let

$$
A(\tau)=\exp \left(\int_{0}^{\tau} \alpha(\sigma) d \sigma\right)
$$

Then multiplication with $A$ defines a Banach space isomorphism $A: \mathcal{H}_{k, \mu}$ $\rightarrow \mathcal{H}_{k}$. The inverse $A^{-1}$ of $A$ is multiplication with $A(\tau)^{-1}$. These isomorphisms give the following commutative diagram

where $\mathcal{H}_{2}\left(\frac{\pi^{*}}{2}\right)$ is defined as $\mathcal{H}_{2}\left(\frac{\pi}{2}\right)$ except that instead of requiring that $\bar{\partial} u$ vanishes on the boundary, we require that $(\bar{\partial}-\alpha) u$ does. We prove that the operator $\bar{\partial}-\alpha$ on the right in the above diagram has closed range and conclude the corresponding statement for the operator on the left. Note that if $u \in \mathcal{H}_{2}\left(\frac{\pi}{2}{ }^{*}\right)$, then both $\partial_{\tau} u$ and $\partial_{t} u$ satisfy (6.3) and (6.4). Hence, the doubling operation described in Section 6.2 induces a map $\mathcal{H}_{2}\left(\frac{\pi}{2}^{*}\right) \rightarrow \mathcal{H}_{2}(\mathbb{R} \times[0,2])$ with $\left\|u^{d}\right\|_{2}=2\|u\|_{2}$.

Let

$$
\begin{equation*}
S(\mu)=\left\{n \in \mathbb{Z}:-\frac{\mu^{-}}{\pi}-\frac{1}{2}<n<-\frac{\mu^{+}}{\pi}-\frac{1}{2}\right\} \tag{6.9}
\end{equation*}
$$

(Note that $S(\mu)=\emptyset$ if $\mu^{+} \geq \mu^{-}$.) The map $\gamma_{n}: \mathcal{H}_{2}\left(\frac{\pi}{2}^{*}\right) \rightarrow \mathcal{H}_{2}(\mathbb{R}, \mathbb{R})$,

$$
\begin{equation*}
u \mapsto c_{n}(\tau)=\int_{0}^{2} u^{d}(\tau, t) \exp \left(-i\left(\frac{\pi}{2}+n \pi\right) t\right) d t \tag{6.10}
\end{equation*}
$$

is continuous. Let $W_{2} \subset \mathcal{H}_{2}\left(\frac{\pi}{2}{ }^{*}\right)$ be the closed subspace

$$
\begin{equation*}
W_{2}=\bigcap_{n \in S(\mu)} \operatorname{ker}\left(\gamma_{n}\right) . \tag{6.11}
\end{equation*}
$$

Using the Fourier expansion of $u^{d}$, we see that $W_{2}$ has a direct complement

$$
\begin{equation*}
V_{2}=\bigcap_{n \notin S(\mu)} \operatorname{ker}\left(\gamma_{n}\right) . \tag{6.12}
\end{equation*}
$$

(Note that if $\mu^{+}>\mu^{-}$, then $W_{2}=\mathcal{H}_{2}\left(\frac{\pi}{2}^{*}\right)$ and $V_{2}=\emptyset$.)
Similarly, we view the maps $\gamma_{n}$ defined by (6.10) as maps $\mathcal{H}_{1}[0] \rightarrow$ $\mathcal{H}_{1}(\mathbb{R}, \mathbb{R})$ and get the corresponding direct sum decomposition $\mathcal{H}_{1}[0]=$ $W_{1} \oplus V_{1}$. If $u \in \mathcal{H}_{2}\left(\frac{\pi}{2}^{*}\right)$, then the Fourier expansion of $u^{d}$ is

$$
u^{d}(\tau+i t)=\sum_{n} c_{n}(\tau) e^{i\left(\frac{\pi}{2}+n \pi\right) t}
$$

Hence,

$$
\begin{equation*}
(\bar{\partial}-\alpha) u^{d}(\tau+i t)=\sum_{n}\left(c_{n}^{\prime}(\tau)-\left(\alpha(\tau)+\frac{\pi}{2}+n\right) c_{n}(\tau)\right) e^{i\left(\frac{\pi}{2}+n \pi\right) t} . \tag{6.13}
\end{equation*}
$$

It follows that $\bar{\partial}\left(W_{2}\right) \subset W_{1}$ and $\bar{\partial}\left(V_{2}\right) \subset V_{1}$.
Let $w \in W_{2}$. Fourier expansion of $w^{d}$ gives

$$
\begin{align*}
& 2\|(\bar{\partial}-\alpha) w\|_{0}^{2}  \tag{6.14}\\
& \quad=\sum_{n \notin S(\mu)} \int_{\mathbb{R}}\left(\left|c_{n}^{\prime}\right|^{2}+\left(\left(\frac{\pi}{2}+n \pi+\alpha(\tau)\right)^{2}+\alpha^{\prime}\right)\left|c_{n}\right|^{2}\right) d \tau \\
& \quad \geq 2 C\|w\|_{1}^{2},
\end{align*}
$$

where the constant $C$ is obtained as follows. If $\mu^{+}>\mu^{-}$, then $\mathbf{P 2}$ implies that the coefficients of $\left|c_{n}\right|^{2}$ are strictly positive, and if $\mu^{+}<\mu^{-}$, then P3 implies that the coefficients in front of $\left|c_{n}\right|^{2}$ are larger than $\frac{4}{5} m\left(\frac{\pi}{2}, \mu\right)^{2}$ since $n \notin S(\mu)$. Finally, if $\mu^{-}=\mu^{+}$, then $\alpha^{\prime}=0$ and the coefficients in front of $\left|c_{n}\right|^{2}$ are larger than $m\left(\frac{\pi}{2}, \mu\right)^{2}$ for all $n$.

If $w \in \mathcal{H}_{2}\left(\frac{\pi}{2}{ }^{*}\right)$, then $\partial_{\tau} w$ and $i \partial_{t} w$ satisfies (6.3) and (6.4) and the Fourier coefficients $c_{n}(\tau)$ of their doubles vanish for $n \in S(\mu)$. Thus, the same argument applies to these functions and the following estimates are obtained

$$
\begin{aligned}
\left\|(\bar{\partial}-\alpha) \partial_{\tau} w\right\|_{0} & \geq C\left\|\partial_{\tau} w\right\|_{1} \\
\left\|(\bar{\partial}-\alpha) \partial_{t} w\right\|_{0} & \geq C\left\|\partial_{t} w\right\|_{1} .
\end{aligned}
$$

If $\mu_{+}=\mu_{-}$, then $\alpha^{\prime}=0$ and $\bar{\partial}-\alpha$ commutes with both $\partial_{t}$ and $\partial_{\tau}$. Hence,

$$
\begin{aligned}
\|(\bar{\partial}-\alpha) w\|_{1} & \geq \frac{1}{2}\left(\|(\bar{\partial}-\alpha) w\|_{0}+\left\|\partial_{\tau}(\bar{\partial}-\alpha) w\right\|_{0}+\left\|\partial_{t}(\bar{\partial}-\alpha) w\right\|_{0}\right) \\
& \geq C\left(\|w\|_{1}+\left\|\partial_{\tau} w\right\|_{1}+\left\|\partial_{t} w\right\|_{1}\right) \geq C\|w\|_{2}
\end{aligned}
$$

where $C=K m\left(\mu, \frac{\pi}{2}\right)$. This proves (6.7).
If $\mu_{+} \neq \mu_{-}$, then $\partial_{\tau}(\bar{\partial}-\alpha) w=(\bar{\partial}-\alpha) \partial_{\tau} w-\alpha^{\prime} w$, and with $K>0$, we conclude from the triangle inequality

$$
\begin{aligned}
& K\|(\bar{\partial}-\alpha) w\|_{0}+\left\|\partial_{\tau}(\bar{\partial}-\alpha) w\right\|_{0}+\left\|\partial_{t}(\bar{\partial}-\alpha) w\right\|_{0} \\
& \quad \geq K C\|w\|_{1}+C\left\|\partial_{\tau} w\right\|_{1}-\left\|\alpha^{\prime} w\right\|_{0}+C\left\|\partial_{t} w\right\|_{1} \\
& \quad \geq\left(K C-\frac{m\left(\frac{\pi}{2}, \mu\right)^{2}}{5}\right)\|w\|_{1}+C\left\|\partial_{\tau} w\right\|_{1}+C\left\|\partial_{t} w\right\|_{1}
\end{aligned}
$$

since $\left|\alpha^{\prime}\right|<\frac{m\left(\frac{\pi}{2}, \mu\right)^{2}}{5}$. Thus, choosing $K$ sufficiently large, we find that there exists a constant $K_{1}$ such that for $w \in W$

$$
\begin{equation*}
\|w\|_{2} \leq K_{1}\|(\bar{\partial}-\alpha) w\|_{1} \tag{6.15}
\end{equation*}
$$

Thus, if $\mu^{+}>\mu^{-}$, we conclude that the range of $\bar{\partial}-\alpha$ is closed. If $\mu^{+}<\mu^{-}$, we need to consider also $V_{2}$.

For $v \in V_{2}$, we have

$$
v^{d}(\tau, t)=\sum_{n \in S(\mu)} c_{n}(\tau) \exp \left(i\left(\frac{\pi}{2}+n \pi\right) t\right)
$$

Let $V_{2}^{\perp}$ be the space of functions in $V_{2}$ which, under doubling, map to the orthogonal complement of the doubles $\phi_{n}^{d}$ of the functions $\phi_{n}(\zeta)=$ $\exp \left(\left(\frac{\pi}{2}+n \pi\right) \zeta+\int \alpha d \tau\right), n \in S(\mu)$ with respect to the $L^{2}$-pairing on $\mathcal{H}_{2}(\mathbb{R} \times[0,2], \mathbb{C})$. Then $V_{2}^{\perp}$ is a closed subspace of finite codimension in $V_{2}$.

We claim there exists a constant $K_{2}$ such that for all $v^{\perp} \in V_{2}^{\perp}$

$$
\begin{equation*}
\left\|v^{\perp}\right\|_{2} \leq K_{2}\left\|(\bar{\partial}-\alpha) v^{\perp}\right\|_{1} \tag{6.16}
\end{equation*}
$$

Assume that this is not the case. Then there exists a sequence $v_{j}^{\perp}$ of elements in $V_{2}^{\perp}$ such that

$$
\begin{align*}
& \left\|v_{j}^{\perp}\right\|_{2}=1  \tag{6.17}\\
& \left\|(\bar{\partial}-\alpha) v_{j}^{\perp}\right\|_{1} \rightarrow 0 . \tag{6.18}
\end{align*}
$$

Let $P>M$ be an integer (see condition P1) and let $v^{\perp} \in V^{\perp}$. Consider the restriction of $v^{\perp}$ and $\bar{\partial} v^{\perp}$ to $\Theta_{P}=\{\tau+i t:|\tau| \geq P\}$. Using Fourier expansion as in (6.14), partial integration, and the fact
that $\alpha^{\prime}(\tau)=0$ for $|\tau|>M$, we find

$$
\begin{align*}
& 2\left\|(\bar{\partial}-\alpha) v^{\perp} \mid \Theta_{P}\right\|_{1}  \tag{6.19}\\
& \geq C\left(\left\|v^{\perp} \mid \Theta_{P}\right\|_{2}+\sum_{n \in S(\mu)} \mu^{+}\left(\left|c_{n}(P)\right|^{2}+\left|c_{n}^{\prime}(P)\right|^{2}\right)\right. \\
& \left.\quad-\mu^{-}\left(\left|c_{n}(-P)\right|^{2}+\left|c_{n}^{\prime}(-P)\right|^{2}\right)\right)
\end{align*}
$$

By a compact Sobolev embedding, we find for each positive integer $P$ a subsequence $\left\{v_{j(P)}^{\perp}\right\}$ which converges in $\mathcal{H}_{1}([-P, P] \times[0,1], \mathbb{C})$. Moreover, we may assume that these subsequences satisfies $\left\{v_{j(P)}^{\perp}\right\} \supset\left\{v_{j(Q)}^{\perp}\right\}$ if $P<Q$.

Let $\left(c_{n}\right)_{j}$ be the sequence of Fourier coefficient functions associated to the sequence $v_{j}^{\perp}$. The estimates

$$
\begin{equation*}
\|c\|_{k} \leq C\left(\|c\|_{k-1}+\left\|\left(\frac{d}{d \tau}-\left(\frac{\pi}{2}+n \pi+\alpha\right)\right) c\right\|_{k-1}\right) \tag{6.20}
\end{equation*}
$$

and (6.18) implies that $\left(c_{n}\right)_{j(P)}$ converges to a smooth solution of the equation $\left(\frac{d}{d \tau}-\left(\frac{\pi}{2}+n \pi+\alpha\right)\right) c=0$ on $[-P, P]$. Hence, $v_{j(P)}^{\perp}$ converges to a smooth solution of $(\bar{\partial}-\alpha) u=0$ on $\Theta_{P}$ satisfying the boundary conditions (6.3) and (6.4). Such a solution has the form

$$
\sum_{n \in S(\mu)} k_{n} \phi_{n}(\zeta)
$$

where $k_{n}$ are real constants.
We next show that in fact all $k_{n}$ must be zero. Note that by Morrey's theorem and (6.17), we get a uniform $C^{0}$-bound $\left|v_{j}^{\perp}\right| \leq K$. Therefore, $\left|\left(c_{n}\right)_{j}\right| \leq 2 K$ and hence,

$$
\begin{aligned}
& \int\left\langle\left(v_{j}^{\perp}\right)^{d},\left(\phi_{n}\right)^{d}\right\rangle d \tau \wedge d t \\
&= \int_{\mathbb{R}}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi\right) \tau+\int \alpha d \tau\right) d \tau \\
&= \int_{-P}^{P}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi\right) \tau+\int \alpha d \tau\right) d \tau \\
&+\int_{P}^{\infty}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi+\mu^{+}\right) \tau\right) d \tau \\
&+\int_{-\infty}^{-P}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi+\mu^{+}\right) \tau\right) d \tau
\end{aligned}
$$

But

$$
\begin{aligned}
& \left|\int_{P}^{\infty}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi+\mu^{+}\right) \tau\right) d \tau\right| \\
& \quad+\left|\int_{-\infty}^{-P}\left(c_{n}\right)_{j} \exp \left(\left(\frac{\pi}{2}+n \pi+\mu^{-}\right) \tau\right) d \tau\right| \\
& \quad \leq \frac{2 K}{m\left(\frac{\pi}{2}, \mu\right)}\left(\exp \left(\left(\frac{\pi}{2}+n \pi+\mu^{+}\right) P\right)\right. \\
& \left.\quad+\exp \left(-\left(\frac{\pi}{2}+n \pi+\mu^{-}\right) P\right)\right) \rightarrow 0 \text { as } P \rightarrow \infty
\end{aligned}
$$

We conclude from this that unless $k_{n}=0, v_{j(P)}^{\perp}$ violates the orthogonality conditions for $P$ and $j(P)$ sufficiently large.

Consider (6.19) applied to elements in the sequence $\left\{v_{j}^{\perp}\right\}$. As $j \rightarrow \infty$ the term on the left-hand side and the sum in the right-hand side tends to 0 . Hence, $\left\|v_{j}^{\perp} \mid \Theta_{P}\right\|_{2} \rightarrow 0$. Applying (6.20) to $\left(c_{n}\right)_{j}$ and noting that both terms on the right-hand side goes to 0 , we conclude that also $\left\|v_{j}^{\perp} \mid[-P \times P] \times[0,1]\right\|_{2} \rightarrow 0$. This contradicts (6.17) and hence (6.16) holds.

The estimates (6.15) and (6.16) together with the direct sum decompositions $\mathcal{H}_{2}\left(\frac{\pi}{2}^{*}\right)=W_{2} \oplus V_{2}$ and $\mathcal{H}_{1}[0]=W_{1} \oplus V_{1}$, and the fact that $\bar{\partial}-\alpha$ respects this decomposition shows that the image of $\bar{\partial}-\alpha$ is closed also in the case $\mu^{+}<\mu^{-}$.
q.e.d.

Remark 6.4. In many cases, the first statement in Proposition 6.3 still holds with weaker assumptions on the weight function than P1P3. For example, if $\mu_{+}<\mu_{-}$, then we need only know that $\max \left\{\alpha^{\prime}, 0\right\}$ is sufficiently small compared to $\left(\frac{\pi}{2}+n \pi+\alpha\right)^{2}$ for $n \notin S(\mu)$ to derive (6.15) and the derivation of (6.16) is quite independent of $\alpha^{\prime}$ as long as $\alpha$ eventually becomes constant.
6.4. The model problem with angles. We study more general boundary conditions than those in Section 6.3. Recall $\left(x_{1}+i y_{1}, \ldots, x_{n}+\right.$ $\left.i y_{n}\right)$ are coordinates on $\mathbb{C}^{n}$. Let $\partial_{j}$ denote the unit tangent vector in the $x_{j}$-direction, for $j=1, \ldots, n$. For $\theta=\left(\theta_{1}, \ldots, \theta_{n}\right) \in[0, \pi)^{n}$, let $\Lambda(\theta)$ be the Lagrangian subspace of $\mathbb{C}^{n}$ spanned by the vectors $e^{i \theta_{1}} \partial_{1}, \ldots, e^{i \theta_{n}} \partial_{n}$. Consider the following boundary conditions for $u \in \mathcal{H}_{1}^{\text {loc }}$.

$$
\begin{gather*}
\int_{\mathbb{R}}\langle u, v\rangle d \tau=0 \text { for all } v \in C_{0}^{\infty}\left(\mathbb{R}, i \mathbb{R}^{n}\right)  \tag{6.21}\\
\int_{\mathbb{R}+i}\langle u, v\rangle d \tau=0 \text { for all } v \in C_{0}^{\infty}(\mathbb{R}+i, i \Lambda(\theta)) . \tag{6.22}
\end{gather*}
$$

If $m\left(\theta_{j}, \mu_{j}\right)>0$ (see (6.5)) for all $j$, then define
$\mathcal{H}_{2, \mu}(\theta)=\left\{u \in \mathcal{H}_{2, \mu}: u\right.$ satisfies (6.21), (6.22), and $\bar{\partial} u$ satisfies (6.6) $\}$,
$\mathcal{H}_{1, \mu}[0]=\left\{u \in \mathcal{H}_{1, \mu}: u\right.$ satisfies (6.6) $\}$.

Proposition 6.5. If $m\left(\theta_{j}, \mu_{j}\right)>0$ for $j=1, \ldots, n$ then the operator

$$
\bar{\partial}: \mathcal{H}_{2, \mu}(\theta) \rightarrow \mathcal{H}_{1, \mu}[0]
$$

is Fredholm of index

$$
\begin{equation*}
\sum_{j=1}^{n} \sharp\left(-\frac{\mu_{j}^{-}+\theta_{j}}{\pi},-\frac{\mu_{j}^{+}+\theta_{j}}{\pi}\right)-\sharp\left(\frac{\mu_{j}^{-}+\theta_{j}}{\pi}-1, \frac{\mu_{j}^{+}+\theta_{j}}{\pi}-1\right) . \tag{6.23}
\end{equation*}
$$

Moreover, if $\mu_{j}^{+}=\mu_{j}^{-}$for all $j$ and $M(\mu)=\min \left\{m\left(\mu_{1}, \theta_{1}\right), \ldots\right.$, $\left.m\left(\mu_{n}, \theta_{n}\right)\right\}$, then $u \in \mathcal{H}_{2, \mu}\left(\theta_{1}, \ldots, \theta_{n}\right)$ satisfies

$$
\begin{equation*}
\|u\|_{2, \mu} \leq C(\mu)\|\bar{\partial} u\|_{1, \mu} \tag{6.24}
\end{equation*}
$$

where $C(\mu) \leq \frac{K}{M(\mu)}$, for some constant $K$.
Proof. Consider the holomorphic $(n \times n)$-matrix

$$
\mathbf{g}_{\theta}(\zeta)=\operatorname{Diag}\left(\left(\exp \left(\frac{\pi}{2}-\theta_{1}\right) \zeta\right), \ldots,\left(\exp \left(\frac{\pi}{2}-\theta_{n}\right) \zeta\right)\right)
$$

Multiplication with $\mathbf{g}_{\theta}$ defines isomorphisms

$$
\begin{aligned}
\mathcal{H}_{2, \mu}(\theta) & \rightarrow \mathcal{H}_{2, \lambda}\left(\frac{\pi}{2}, \ldots, \frac{\pi}{2}\right) \quad \text { and } \\
\mathcal{H}_{1, \mu}[0] & \rightarrow \mathcal{H}_{1, \lambda}[0]
\end{aligned}
$$

where $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$ and $\lambda_{j}^{ \pm}=\mu_{j}^{ \pm}-\frac{\pi}{2}+\theta_{j}$. Since $\mathbf{g}_{\theta}$ is holomorphic it commutes with $\bar{\partial}$. The proposition now follows from Proposition 6.3.
q.e.d.

### 6.5. Smooth perturbations of the model problem with angles.

 Let $B: \mathbb{R} \times[0,1] \rightarrow \mathbf{U}(n)$ be a smooth map such that$$
\begin{equation*}
\bar{\partial} B \mid \partial \mathbb{R} \times[0,1]=0 \tag{6.25}
\end{equation*}
$$

Let $\theta \in[0, \pi)$ and consider the following boundary conditions for $u \in$ $\mathcal{H}_{1}^{\text {loc }}$ :

$$
\begin{equation*}
\int_{\mathbb{R}}\langle u, v\rangle d \tau \wedge d t=0 \tag{6.26}
\end{equation*}
$$

for all $v \in C_{0}^{\infty}\left(\mathbb{R}, \mathbb{C}^{n}\right)$ such that $v(\tau) \in i B(\tau) \mathbb{R}^{n}$,

$$
\int_{\mathbb{R}+i}\langle u, v\rangle d \tau \wedge d t=0
$$

$$
\begin{equation*}
\text { for all } v \in C_{0}^{\infty}\left(\mathbb{R}+i, \mathbb{C}^{n}\right) \text { such that } v(\tau+i) \in i B(\tau) \Lambda(\theta) \tag{6.27}
\end{equation*}
$$

For $\mu=\left(\mu^{+}, \mu^{-}\right) \in \mathbb{R}^{2}$ let $\lambda(\mu)=\left(\mu^{+}, \mu^{-}, \mu^{+}, \mu^{-}, \ldots, \mu^{+}, \mu^{-}\right) \in \mathbb{R}^{2 n}$ define

$$
\begin{aligned}
& \mathcal{H}_{2, \mu}(\theta, B)=\left\{u \in \mathcal{H}_{2}^{\text {loc }}: u \text { satisfies }(6.26) \text { and }(6.27)\right. \\
& \left.\bar{\partial} u \text { satisfies }(6.6) \text { and } \mathbf{e}_{\lambda(\mu)} u \in \mathcal{H}_{2}\right\}
\end{aligned}
$$

Proposition 6.6. If $m\left(\theta_{j}, \mu\right)>0$ for $j=1, \ldots, n$, then there exists $\delta>0$ such that for all $B$ satisfying (6.25) with $\|B-\mathrm{id}\|_{C^{2}}<\delta$, the operator

$$
\bar{\partial}: \mathcal{H}_{2, \mu}(\theta, B) \rightarrow \mathcal{H}_{1, \mu}[0]
$$

is Fredholm of index

$$
\begin{equation*}
\sum_{j=1}^{n} \sharp\left(-\frac{\mu^{-}+\theta_{j}}{\pi},-\frac{\mu^{+}+\theta_{j}}{\pi}\right)-\sharp\left(\frac{\mu^{-}+\theta_{j}}{\pi}-1, \frac{\mu^{+}+\theta_{j}}{\pi}-1\right) . \tag{6.28}
\end{equation*}
$$

Proof. Multiplication with $B$ and $B^{-1}$ defines Banach space isomorphisms

$$
\mathcal{H}_{2, \mu}(\theta) \xrightarrow{\times B} \mathcal{H}_{2, \mu}(\theta, B),
$$

and

$$
\mathcal{H}_{1, \mu}[0] \xrightarrow{\times B^{-1}} \mathcal{H}_{1, \mu}[0] .
$$

Thus up to conjugation, the operator considered is the same as

$$
\bar{\partial}+B^{-1} \bar{\partial} B: \mathcal{H}_{2, \mu}(\theta) \rightarrow \mathcal{H}_{1, \mu}[0] .
$$

The theorem now follows from Proposition 6.5, and the fact that the subspace of Fredholm operators is open and that the index is constant on path components of this subspace.
q.e.d.
6.6. Boundary conditions. In the upcoming subsections, we study the linearized $\bar{\partial}$-problem on a disk $D_{m}$ with $m$ punctures. Refer back to Section 4.4 for notation concerning $D_{m}$.

Definition 6.7. A smooth map $A: \partial D_{m} \rightarrow \mathbf{U}(n)$ will be called small at infinity if there exists $M>1$ such that for each $j=1, \ldots, m$, the restriction of $A$ to $\partial E_{p_{j}}[M]$ approaches a constant map in the $C^{2}$-norm on each component of $\partial E_{p_{j}}\left[M^{\prime}\right]$ as $M^{\prime} \rightarrow \infty$. It will be called constant at infinity if there exists $M>1$ such that for each $j=1, \ldots, m$, the restriction of $A$ to each component of $\partial E_{p_{j}}[M]$ is constant.

Let $A: \partial D_{m} \rightarrow \mathbf{U}(n)$ be small at infinity. For $u \in \mathcal{H}_{1}^{\text {loc }}\left(D_{m}, \mathbb{C}^{n}\right)$, consider the boundary condition:

$$
\begin{align*}
\int_{\partial D_{m}}\langle u, v\rangle d s=0, & \text { for all } v \in C_{0}^{0}\left(\partial D_{m}, \mathbb{C}^{n}\right)  \tag{6.29}\\
& \text { such that } v(\zeta) \in i A(\zeta) \mathbb{R}^{n} \text { for all } \zeta \in \partial D_{m} .
\end{align*}
$$

In previous subsections, coordinates $\zeta=\tau+$ it on $\mathbb{R} \times[0,1]$ were used and we implicitly considered the bundle $T^{* 0,1} \mathbb{R} \times[0,1]$ as trivialized by the form $d \bar{\zeta}$, and sections in this bundle as $\mathbb{C}^{n}$-valued functions. We do not want to specify any trivialization of $T^{* 0,1} D_{m}$ and so we view the $\bar{\partial}$ operator as a map from $\mathcal{H}_{2}$-functions into $\mathcal{H}_{1}$-sections of $T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}$. Consider, for $u \in \mathcal{H}_{1}^{\text {loc }}\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$, the boundary condition

$$
\begin{equation*}
\int_{\partial D_{m}}\langle u, v\rangle d s=0, \text { for all } v \in C_{0}^{0}\left(\partial D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right) \tag{6.30}
\end{equation*}
$$

Henceforth, to simplify notation, if the source space $X$ in a Sobolev space $\mathcal{H}_{k}(X, Y)$ is $D_{m}$, we will drop it from the notation. If $u \in$ $\mathcal{H}_{2}^{\text {loc }}\left(\mathbb{C}^{n}\right)$, then $\bar{\partial} u \in \mathcal{H}_{1}^{\text {loc }}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$. Define
$\mathcal{H}_{2}\left(\mathbb{C}^{n} ; A\right)=\left\{u \in \mathcal{H}_{2}\left(\mathbb{C}^{n}\right): u\right.$ satisfies (6.29) and $\bar{\partial} u$ satisfies (6.30) $\}$,
and
$\mathcal{H}_{1}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right)=\left\{u \in \mathcal{H}_{1}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right): u\right.$ satisfies $\left.(6.30)\right\}$.
Define

$$
\begin{array}{r}
\mathcal{H}_{2, \mu}\left(\mathbb{C}^{n} ; A\right) \\
=\left\{u \in \mathcal{H}_{2}^{\text {loc }}\left(\mathbb{C}^{n}\right): u \text { satisfies (6.29), } \bar{\partial} u \text { satisfies }(6.30),\right. \\
\text { and } \left.\mathbf{e}_{\mu} u \in \mathcal{H}_{2}\left(\mathbb{C}^{n}\right)\right\}
\end{array}
$$

and

$$
\begin{aligned}
& \mathcal{H}_{1, \mu}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right) \\
& =\left\{u \in \mathcal{H}_{1}^{\text {loc }}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right): u \text { satisfies }(6.30)\right. \\
& \left.\quad \text { and } \mathbf{e}_{\mu} u \in \mathcal{H}_{1}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)\right\}
\end{aligned}
$$

Let $p_{j}$ be a puncture of $D_{m}$. The orientation of $D_{m}$ induces an orientation of $\partial D_{m}$. Let $A_{j}^{0}$ and $A_{j}^{1}$ denote the constant maps to which $A$ converges on the component of $\partial E_{p_{j}}$ close to $p_{j}$ corresponding to $\mathbb{R}$ and $\mathbb{R}+i$, respectively. Define

$$
\theta(j)=\theta\left(A_{j}^{0} \mathbb{R}^{n}, A_{j}^{1} \mathbb{R}^{n}\right)
$$

Then, there are unique unitary complex coordinates

$$
z(j)=\left(x(j)_{1}+i y(j)_{1}, \ldots, x(j)_{n}+i y(j)_{n}\right)
$$

in $\mathbb{C}^{n}$ such that

$$
\begin{aligned}
& A_{j}^{0} \mathbb{R}^{n}=\operatorname{Span}\left\langle\partial(j)_{1}, \ldots, \partial(j)_{n}\right\rangle \\
& A_{j}^{1} \mathbb{R}^{n}=\operatorname{Span}\left\langle e^{i \theta(j)_{1}} \partial(j)_{1}, \ldots, e^{i \theta(j)_{n}} \partial(j)_{n}\right\rangle
\end{aligned}
$$

Proposition 6.8. Let $A: \partial D_{m} \rightarrow \mathbf{U}(n)$ be small at infinity. If $\mu$ satisfies $\mu_{j} \neq-\theta(j)_{r}+k \pi$ for $j=1, \ldots, m, r=1, \ldots, n$, and every $k \in \mathbb{Z}$, then the operator

$$
\begin{equation*}
\bar{\partial}: \mathcal{H}_{2, \mu}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1, \mu}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right) \tag{6.31}
\end{equation*}
$$

is Fredholm.
Proof. Assume that for $M>0, A \mid \partial E_{p_{j}}[M-1]$ is sufficiently close to a constant map (see Proposition 6.6). Choose smooth complex-valued functions $\alpha_{0}, \alpha_{1}, \ldots, \alpha_{m}$ with the following properties: $\alpha_{j}$ is constantly 1 on $E_{p_{j}}[M+2]$; the sum $\sum_{j} \alpha_{j}$ is close to the constant function 1 ,
$\bar{\partial} \alpha_{j}=0$ on $\partial D_{m}$; and $\alpha_{j}$ is constantly equal to 0 on $D_{m}-E_{p_{j}}[M+1]$, for $j=1, \ldots, m$.

Glue to each $E_{p_{j}}[M]$ a half-infinite strip $(-\infty, M] \times[0,1]$ and denote the result $\bar{E}_{p_{j}}$. Extend the boundary conditions from $E_{p_{j}}[M]$ to $\bar{E}_{p_{j}}$ keeping them close to constant. Let the weight in the weight function remain constant. Glue to $D_{m}-\bigcup_{j} E_{p_{j}}[M+2], m$ half disks and extend the boundary conditions smoothly. Denote the result $\bar{D}_{m}$. Note that the boundary value problem on $\bar{D}_{m}$ is the vector-Riemann-Hilbert problem, which is known to be Fredholm, and that the weighted norm on this compact disk is equivalent to the standard norm.

Now let $u \in \mathcal{H}_{2, \mu}\left(\mathbb{C}^{n} ; A\right)$. Then $\alpha_{j} u$ is in the appropriate Sobolev space for the extended boundary value problem on $\bar{E}_{p_{j}}\left(\bar{D}_{m}\right.$ if $\left.j=0\right)$ and because the elliptic estimate holds for all of these problems and since all of them except possibly the one on $\bar{D}_{m}$ has no kernel, there exists a constant $C$ such that

$$
\begin{align*}
\|u\|_{2, \mu} & \leq\left\|\alpha_{0} u\right\|_{2, \mu}+\sum_{j=1}^{n}\left\|\alpha_{j} u\right\|_{2, \mu}  \tag{6.32}\\
& \leq C\left(\left\|\alpha_{0} u\right\|_{1, \mu}+\sum_{j=0}^{n}\left\|\bar{\partial}\left(\alpha_{j} u\right)\right\|_{1, \mu}\right) \\
& \leq C\left(\sum_{j=0}^{n}\left\|\bar{\partial} \alpha_{j} u\right\|_{1, \mu}+\left\|\alpha_{0} u\right\|_{1, \mu}+\sum_{j=0}^{n}\left\|\alpha_{j} \bar{\partial} u\right\|_{1, \mu}\right)
\end{align*}
$$

We shall show that (6.32) implies that every bounded sequence $u_{r}$ such that $\bar{\partial} u_{r}$ converges has a convergent subsequence. This implies that $\bar{\partial}$ has a closed image and a finite dimensional kernel ([20] Proposition 19.1.3). Clearly, it is sufficient to consider the case $\bar{\partial} u_{r} \rightarrow 0$. Consider the restrictions of $u_{r}$ to a compact subset $K$ of $D_{m}$ such that

$$
\operatorname{supp}\left(\alpha_{0}\right) \cup \operatorname{supp}\left(\bar{\partial} \alpha_{0}\right) \cup \cdots \cup \operatorname{supp}\left(\bar{\partial} \alpha_{m}\right) \subset K
$$

A compact Sobolev embedding argument gives a subsequence $\left\{u_{r^{\prime}}\right\}$ which converges in $\mathcal{H}_{1}\left(K, \mathbb{C}^{n}\right)$. Thus, (6.32) implies that $\left\{u_{r^{\prime}}\right\}$ is a Cauchy sequence in $\mathcal{H}_{2, \mu}\left(A ; \mathbb{C}^{n}\right)$ and hence it converges.

It remains to prove that the cokernel is finite dimensional. Lemma 6.1 shows that any element in the cokernel of $\bar{\partial}$ can be represented by a smooth function $v$ on $D_{m}$. Partial integration implies this function satisfies $\partial v=0$ with boundary conditions given by the matrix function $i A$. Assume first that $A$ is constant at infinity. Then, Lemma 6.2 and conjugation with the holomorphic $(n \times n)$-matrix $\mathbf{g}_{\theta}$ as in the proof of Proposition 6.5 gives explicit formulas for the restrictions of these smooth functions to $E_{p_{j}}[M]$, for each $j$. It is straightforward to check from these local formulas that $v$ lies in $\mathcal{H}_{2,-\mu}\left(\mathbb{C}^{n}, i A\right)$. Thus, repeating
the argument above with $\partial$ replacing $\bar{\partial}$ shows that the cokernel is finite dimensional. The lemma follows in the case when $A$ is constant at infinity. The general case then follows by an approximation argument as in the proof of Proposition 6.6.
q.e.d.
6.7. Index-preserving deformations. We compute the index of the operator in (6.31). Using approximations, it is easy to see that it is sufficient to consider the case when $A: \partial D_{m} \rightarrow \mathbf{U}(n)$ is constant at infinity. Thus, let $A$ be such a map which is constant on $\partial E_{p_{j}}[M]$ for every $j$ and consider the Fredholm operator

$$
\begin{equation*}
\bar{\partial}: \mathcal{H}_{2, \mu}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1, \mu}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right) \tag{6.33}
\end{equation*}
$$

where $\mu=\left(\mu_{1}, \ldots, \mu_{m}\right) \in \mathbb{R}^{m}$ satisfies

$$
\begin{equation*}
\mu_{j} \neq-\theta(j)_{r}+n \pi \text { for every } j, r, n \tag{6.34}
\end{equation*}
$$

Lemma 6.9. Let $B_{s}: D_{m} \rightarrow \mathbf{U}(n), s \in[0,1]$, be a continuous family of smooth maps such that

$$
\begin{align*}
& B_{s} \text { is bounded in the } C^{2} \text {-norm, }  \tag{6.35}\\
& B_{s} \mid \partial E_{p_{j}}[M] \text { is constant in } \tau+\text { it, } \\
& \bar{\partial} B_{s} \mid \partial D_{m}=0, \text { and } \\
& B_{0} \equiv \mathrm{id} .
\end{align*}
$$

Let $\lambda:[0,1] \rightarrow \mathbb{R}^{m}$ be a continuous map such that $\lambda(0)=\mu$ and $\lambda(s)$ satisfies (6.34) for every $s \in[0,1]$. Then the operator

$$
\bar{\partial}: \mathcal{H}_{2, \lambda(1)}\left(\mathbb{C}^{n} ; B_{1} A\right) \rightarrow \mathcal{H}_{1, \lambda(1)}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right)
$$

has the same Fredholm index as the operator in (6.33).
Proof. The Fredholm operator

$$
\bar{\partial}: \mathcal{H}_{2, \lambda(s)}\left(\mathbb{C}^{n} ; B_{s} A\right) \rightarrow \mathcal{H}_{1, \lambda(s)}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right)
$$

is conjugate to

$$
\bar{\partial}-B_{s} \bar{\partial} B_{s}^{-1}: \mathcal{H}_{2, \mu}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1, \mu}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right) .
$$

The family $\bar{\partial}-B_{s} \bar{\partial} B_{s}^{-1}$ is then a continuous family of Fredholm operators.
q.e.d.

In order to apply Lemma 6.9 , we shall show how to deform given weights and boundary conditions into other boundary conditions and weights keeping the Fredholm index constant using the conditions in Lemma 6.9. We accomplish this in two steps: first deform the problem so that the boundary value matrix is diagonal; then change the weights and angles at the ends into a special form where compactification is possible.

Lemma 6.10. Let $A: \partial D_{m} \rightarrow \mathbf{U}(n)$ be constant at infinity. Then, there exists a continuous family $B_{s}: D_{m} \rightarrow \mathbf{U}(n), 0 \leq s \leq 1$, of maps satisfying (6.35) such that

$$
B_{1}(\zeta) A(\zeta)=\operatorname{Diag}\left(b_{1}(\zeta), \ldots, b_{n}(\zeta)\right), \zeta \in \partial D_{m}
$$

Proof. We first make $A$ diagonal on the ends where it is constant. Note that in canonical coordinates $z(j)$ on the end $E_{p_{j}}[M]$ the matrix $A$ is diagonal. Let $B_{j} \in \mathbf{U}(n)$ be the matrix which transforms the complex basis $\partial(j)_{1}, \ldots, \partial(j)_{n}$ to the standard basis. Let $B_{j}(s)$ be a smooth path in $\mathbf{U}(n)$, starting at id and ending at $B_{j}$. Define $B_{s}=B_{j}(s)$ on $E_{p_{j}}[M]$ for each $j$.

We need to extend this map to all of $D_{m}$. To this end, consider the loop on the boundary of $S=D_{m}-E_{p_{j}}[M]$. There exists a 1parameter family of functions $B_{s}: S \rightarrow \mathbf{U}(n)$ such that $B_{0}=\mathrm{id}$ and $B_{1} A$ is diagonal, since any loop is homotopic to a loop of diagonal matrices. The loops $B_{s}$ can be smoothly extended to all of $D_{m}$.

Finally, we need that $\bar{\partial} B_{s}=0$ on the boundary. We get this as follows: let $C$ be a collar on the boundary with coordinates $\tau$ along the boundary and $t$ orthogonal to the boundary, $0 \leq t \leq \epsilon$ and let $\phi:[0, \epsilon] \rightarrow \mathbb{R}$ be a smooth function which equals the identity on $\left[0, \frac{\epsilon}{4}\right]$ and 0 for $t \geq \frac{\epsilon}{2}$. Redefine $B_{s}$ on the collar as

$$
\tilde{B}_{s}=B_{s}(\zeta) \exp \left(i \phi(t) B_{s}^{-1}(\zeta) \bar{\partial} B_{s}(\zeta)\right)
$$

Then $\tilde{B}_{s}$ satisfies the boundary conditions and equals $B_{s}$ on the boundary and in the complement of the collar.

Consider the loop on the boundary of $S=D_{m}-\partial E_{p_{j}}[M]$. There exists a 1-parameter family of functions $B_{s}: S \rightarrow \mathbf{U}(n)$ such that $B_{0}=$ id and $B_{1} A$ is diagonal, since any loop is homotopic to a loop of diagonal matrices. The loops $B_{s}$ can be smoothly extended to all of $D_{m}$ and the above trick makes $B_{s}$ satisfy the boundary conditions. q.e.d.

Now let $A: D_{m} \rightarrow \mathbf{U}(n)$ take values in diagonal matrices. Assume that $A$ is constant near the punctures and that $\mu=\left(\mu_{1}, \ldots, \mu_{m}\right) \in \mathbb{R}^{m}$ satisfies (6.34).

Lemma 6.11. There are continuous families of smooth maps $B_{s}: D_{m}$ $\rightarrow$ Diag $\subset \mathbf{U}(n)$ and $\lambda:[0,1] \rightarrow \mathbb{R}$ which satisfy (6.35) and (6.34) (where the $\theta(j)$ are computed w.r.t. $B_{s}$ ) respectively such that

$$
B_{1} A=\mathrm{id}
$$

in a neighborhood of each puncture.
Proof. Let $M>0$ be such that $A$ is constant in $E_{p_{j}}[M]$ for each $j$. Let $\phi:[0,1] \rightarrow[0,1]$ be an approximation of the identity which is constant near the endpoints of the interval. Let $\psi:[M, \infty) \rightarrow[0,1]$ be
a smooth increasing function which is identically 0 on $[M, M+1]$ and identically 1 on $[M+2, \infty)$. For $\alpha=\left(\alpha_{1}, \ldots, \alpha_{m}\right) \in(-\pi, \pi)^{m}$, let

$$
\tilde{g}_{\alpha}(\zeta)= \begin{cases}1 & \text { for } \zeta \in D_{m}-\bigcup_{j} E_{p_{j}}[M] \\ e^{i \psi(\tau) \alpha_{j} \phi(t)} & \text { for } \zeta=\tau+i t \in E_{p_{j}}[M]\end{cases}
$$

and let $g_{\alpha}$ be a function which agrees with $\tilde{g}_{\alpha}$ except on $E_{p_{j}}[M]-$ $E_{p_{j}}[M+2]$ and which satisfies $\bar{\partial} g_{\alpha} \mid \partial D_{m}=0$.

Consider the complex angle $\theta(j) \in[0, \pi)^{n}$ and the weight $\mu_{j}$. Assume first that $\mu_{j} \neq k \pi$ for all $k \in \mathbb{Z}$ and $j=1, \ldots, m$. Let $m_{j}$ be the unique number $0 \leq m_{j} \leq \pi$ such that $m_{j}=k \pi-\mu_{j}$ for some $k \in \mathbb{Z}$. By (6.34) $\theta(j)_{r} \neq m_{j}$ for all $r$. If $\theta(j)_{r}>m_{j}$, define $\alpha_{r}=\pi-\theta(j)_{r}$, and if $\theta(j)_{r}<m_{j}^{*}$, define $\alpha(j)_{r}=-\theta(j)_{r}$. Define

$$
B_{s}=\operatorname{Diag}\left(g_{s \alpha_{1}}, \ldots, g_{s \alpha_{n}}\right)
$$

and let $\lambda(s) \equiv \mu$.
Assume now that $\mu_{j}=k \pi$ for some $j$. For $0 \leq s \leq \frac{1}{2}$, let $B_{s}=\mathrm{id}$ and take $\lambda_{j}=\mu_{j}-\epsilon s$ for some sufficiently small $\epsilon>0$. Repeat the above construction to construct $B_{s}$ for $s \leq \frac{1}{2} \leq 1$.
q.e.d.
6.8. The Fredholm index of the standardized problem. Consider $D_{m}$ with $m$ punctures on the boundary, conformal structure $\kappa$ and metric $g(\kappa)$ as above and neighborhoods $E_{p_{j}}$ of the punctures $p_{1}, \ldots, p_{m}$.

Let $\Delta_{m}$ denote the representative of the conformal structure $\kappa$ on $D_{m}$ which is the unit disk in $\mathbb{C}$ with $m$ punctures at $1, i,-1, q_{3}, \ldots, q_{m}$ with the flat metric. Then there exists a conformal and therefore holomorphic map $\Gamma: D_{m} \rightarrow \Delta_{m}$. We study the behavior of $\Gamma$ on $E_{p_{j}}$. Let $p=p_{j}$ and let $q$ be the puncture on $\Delta_{m}$ to which $p$ maps. After translation and rotation in $\mathbb{C}$, we may assume that the point $q=0$ and that $\Delta_{m}$ is the disk of radius 1 centered at $i$. We may then find a holomorphic function on a neighborhood $U \subset \Delta_{m}$ of $q=0$ which fixes 0 and maps $\partial \Delta_{m} \cap U$ to the real line. Composing with this map, we find that $\Gamma$ maps $\infty$ to $0, \tau+0 i$ to the negative real axis and $\tau+i$ to the positive real axis for $\tau>M$ for some $M$. Thus this composition equals $C \exp (-\pi \zeta)$ where $C<0$ is some negative real constant. Thus, up to a bounded holomorphic change of coordinates on a neighborhood of $q$, the map $\Gamma$ on $E_{p_{j}}$ looks like $\Gamma(\zeta)=\exp (-\pi \zeta)$ and its inverse $\Gamma^{-1}$ in these coordinates satisfies $\Gamma^{-1}(z)=-\frac{1}{\pi} \log (z)$.

Let $A: \partial D_{m} \rightarrow$ Diag $\subset \mathbf{U}(n)$ be a smooth function which is constantly equal to id close to each puncture. We may now think of $A$ as being defined on $\partial \Delta_{m}$. We extend $A$ smoothly to $\partial \Delta$ by defining its extension $\hat{A}$ at the punctures as $\hat{A}\left(p_{j}\right)=\mathrm{id}$ for each $j$.

Consider the following boundary condition for $u \in \mathcal{H}_{2}\left(\Delta, \mathbb{C}^{n}\right)$ :

$$
\begin{align*}
& \int_{\partial \Delta}\langle u, v\rangle d s=0 \text { for all } v \in C_{0}^{\infty}\left(\partial \Delta, \mathbb{C}^{n}\right)  \tag{6.36}\\
& \qquad \text { with } v(z) \in i \hat{A}(z) \mathbb{R}^{n} \text { for all } z \in \partial \Delta .
\end{align*}
$$

For $u \in \mathcal{H}_{1}\left(\Delta, T^{* 0,1} \Delta \otimes \mathbb{C}^{n}\right)$, consider the boundary conditions

$$
\begin{equation*}
\int_{\partial \Delta}\langle u, v\rangle d s=0 \text { for all } v \in C_{0}^{\infty}\left(\partial \Delta, T^{0,1} \Delta \otimes \mathbb{C}^{n}\right) \tag{6.37}
\end{equation*}
$$

Define

$$
\begin{gathered}
\mathcal{H}_{2}\left(\Delta, \mathbb{C}^{n} ; \hat{A}\right)=\left\{u \in \mathcal{H}_{2}(\Delta, \mathbb{C}): u \text { satisfies }(6.36)\right. \\
\text { and } \bar{\partial} u \text { satisfies }(6.37)\} \\
\mathcal{H}_{1}\left(\Delta, T^{* 0,1} \Delta ;[0]\right)=\left\{u \in \mathcal{H}_{2}(\Delta, \mathbb{C}): u \text { satisfies }(6.37)\right\} .
\end{gathered}
$$

Lemma 6.12. The operator

$$
\bar{\partial}: \mathcal{H}_{2}\left(\Delta, \mathbb{C}^{n} ; \hat{A}\right) \rightarrow \mathcal{H}_{1}\left(\Delta, T^{* 0,1} \Delta ;[0]\right)
$$

is Fredholm of index $n+\mu(\hat{A})$, where $\mu(\hat{A})$ denotes the Maslov index of the loop $z \mapsto A(z) \mathbb{R}^{n}, z \in \partial \Delta$, of Lagrangian subspaces in $\mathbb{C}^{n}$.

Proof. This is (a direct sum of) classical Riemann-Hilbert problems. q.e.d.

Let $\lambda(a)=(a, \ldots, a) \in \mathbb{R}^{m}$.
Proposition 6.13. For $-\pi<a<0$, the Fredholm index of the operator

$$
\bar{\partial}: \mathcal{H}_{2, \lambda(a)}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1, \lambda(a)}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right)
$$

equals $n+\mu(A)$.
Proof. The holomorphic map $\Gamma: D_{m} \rightarrow \Delta_{m}$ and its holomorphic inverse commute with the $\bar{\partial}$ operator. Any solution on $D_{m}$ must look like $\sum_{n \leq 0} c_{n} e^{\pi n \zeta}$ (the negative weights allows for $c_{0} \neq 0$ ) in canonical coordinates close to each puncture. Thus, $\Gamma^{-1}$ pulls back solutions on $D_{m}$ to solutions on $\Delta_{m}$. Using also $\Gamma$, we see that the kernels are isomorphic.

Elements in the cokernel on $D_{m}$ are of the form $\left(\sum_{n<0} c_{n} e^{n \pi \bar{\zeta}}\right) d \bar{\zeta}$ (the positive weight implies $c_{0}=0$ ). Pulling back with $\Gamma^{-1}$ gives elements of the form $\left(\sum_{n>0} \bar{z}^{n}\right) \frac{d \bar{z}}{\bar{z}}$ which are in the cokernel of the $\bar{\partial}$ on $\Delta$. So, the cokernels are also isomorphic.
q.e.d.
6.9. The index of the linearized problem. In this subsection, we determine the Fredholm indices of the problems which are important in our applications to contact geometry.

Let $\left.A: \partial D_{m} \rightarrow \mathbf{U}(n)\right)$ be a map which is small at infinity. Assume that $A_{j}^{0} \mathbb{R}^{n}$ and $A_{j}^{1} \mathbb{R}^{n}$ are transverse for all $j$. For $0 \leq s \leq 1$, let $\mathbf{f}_{j}(s) \in \mathbf{U}(n)$ be the matrix which in the canonical coordinates $z(j)$ is represented by the matrix

$$
\operatorname{Diag}\left(e^{-i\left(\pi-\theta(j)_{1}\right) s}, \ldots, e^{-i\left(\pi-\theta(j)_{n}\right) s}\right)
$$

If $p$ and $q$ are consecutive punctures on $\partial D_{m}$, then let $I(a, b)$ denote the (oriented) path in $\partial D_{m}$ which connects them. Define the loop $\Gamma_{A}$ of Lagrangian subspaces in $\mathbb{C}^{n}$ by letting the loop

$$
\left(A \mid I\left(p_{1}, p_{2}\right)\right) * \mathbf{f}_{2} *\left(A \mid I\left(p_{2}, p_{3}\right)\right) * \mathbf{f}_{3} * \cdots *\left(A \mid I\left(p_{m}, p_{1}\right)\right) * \mathbf{f}_{1}
$$

of elements of $\mathbf{U}(n)$ act on $\mathbb{R}^{n} \subset \mathbb{C}^{n}$.
Proposition 6.14. For $A$ as above, the index of the operator

$$
\bar{\partial}: \mathcal{H}_{2}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1}\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n} ;[0]\right)
$$

equals $n+\mu\left(\Gamma_{A}\right)$ where $\mu$ is the Maslov index.
Proof. Using Lemmas 6.10 and 6.11, we deform $A$ to put the problem into standardized form with weight $-\epsilon$ at each corner, without changing the index. Call the new matrix $B$. We need to consider how $B$ is constructed from $A$. The key step to understand is the point where we make $B$ equal the identity on the ends. This is achieved by first introducing a small negative weight and then rotating the space

$$
A_{j}^{1} \mathbb{R}^{n}=\operatorname{Span}\left\langle e^{i \theta(j)_{1}} \partial_{1}, \ldots, e^{i \theta(j)_{n}} \partial_{n}\right\rangle
$$

to $A_{j}^{0}$ according to

$$
\begin{equation*}
\operatorname{Span}\left\langle e^{i\left(\theta(j)_{1}+s \phi(\tau)\left(\pi-\theta(j)_{1}\right)\right.} \partial_{1}, \ldots, e^{i\left(\theta(j)_{n}+s \phi(\tau)\left(\pi-\theta(j)_{n}\right)\right.} \partial_{n}\right\rangle, \tag{6.38}
\end{equation*}
$$

where $0 \leq s \leq 1$ and $\phi:[M, \infty) \rightarrow[0,1]$ equals 1 on $[M+2, \infty)$ and 0 on $[M, M+1]$.

We now calculate the Maslov-index $\mu(B)$. Since as we follow $\mathbb{R}+i$ along the negative $\tau$-direction from $M+1$ to $M, B$ experiences the inverse of the rotation (6.38), the proposition follows. q.e.d.

We now consider the simplest degeneration at a corner. Compare this with Theorem 4.A of [13] or the appendix of [30]. Let $\epsilon>0$ be a small number. Let $A_{s}: D_{m} \rightarrow \mathbf{U}(n), 0 \leq s \leq 1$ be a family matrices which are small at infinity and constant in $s$ near each puncture in $S \subset\{1, \ldots, m\}$, where each component of the complex angle is assumed to be positive. At $p_{r}, r \notin S$, assume that $\theta(r)_{s}=\left(\pi-s, \theta_{2}(r), \ldots, \theta_{n}(r)\right)$, where $\theta_{j}(r) \neq 0, j=2, \ldots, n$.. Let $\beta(\epsilon) \in \mathbb{R}^{m}$ satisfy $\beta(\epsilon)_{r}=0$ if $r \in S$ and $\beta(\epsilon)_{r}=-\epsilon$ if $r \notin S$.

Proposition 6.15. The index of the operators

$$
\bar{\partial}: \mathcal{H}_{2}\left(\mathbb{C}^{n} ; A_{s}\right) \rightarrow \mathcal{H}_{1}\left(T^{0,1} D_{m} ;[0]\right)
$$

for $s>0$ and of the operator

$$
\bar{\partial}: \mathcal{H}_{2, \beta(\epsilon)}\left(\mathbb{C}^{n} ; A_{0}\right) \rightarrow \mathcal{H}_{1,(-\epsilon, 0, \ldots, 0)}\left(T^{0,1} D_{m} ;[0]\right)
$$

are the same.
Proof. This is a consequence of Lemma 6.9. q.e.d.
Finally, we show how the index is affected if the weight is changed.
Proposition 6.16. Let $A: D_{m} \rightarrow \mathbf{U}(n)$ be constant at infinity and suppose that the complex angle at each puncture except possibly $p_{1}$ has positive components. Assume that $0 \leq \pi-\theta(1)_{1}<\pi-\theta(1)_{2}<\cdots<$ $\pi-\theta(1)_{n}$. Let $\epsilon>0$ be smaller than $\min _{r}\left(\pi-\theta(r)_{r}\right)$, and let $\pi-\theta(1)_{j}<$ $\delta<\pi-\theta(1)_{j-1}$ Then, the index of the problem

$$
\bar{\partial}: \mathcal{H}_{2,(-\epsilon, 0, \ldots, 0)}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1,(-\epsilon, 0, \ldots, 0)}\left(T^{* 0,1} D_{m} ;[0]\right)
$$

is $j$ larger than that of

$$
\bar{\partial}: \mathcal{H}_{2,(\delta, 0, \ldots, 0)}\left(\mathbb{C}^{n} ; A\right) \rightarrow \mathcal{H}_{1,(\delta, 0, \ldots, 0)}\left(T^{* 0,1} D_{m} ;[0]\right)
$$

Proof. First, deform the matrix into diagonal form without changing the weights. If $n>1$, this can be done in such a way that the index corresponding to the first component is positive. Then put the first component in standardized form. We must consider the index difference arising from the first component as the weight changes from negative to positive. The condition that a solution lies in $\mathcal{H}_{2, \delta}$ means that the corresponding solution on $\Delta$ vanishes at $p_{1}$. Thus, the dimension of the kernel increases by 1. The cokernel remains zero-dimensional. This argument can then be repeated for other components. To handle the 1-dimensional case, one may either use similar arguments for cokernels or reduce to the higher dimensional case by adding extra dimensions.

> q.e.d.
6.10. The index and the Conley-Zehnder index. We translate Proposition 6.14 into a more invariant language. Recall from Section 2.2 that we denote by $\nu_{\gamma}(c)$ the Conley-Zehnder index of Reeb chord $c$ with capping path $\gamma$. In the following proposition, we suppress $\gamma$ from the notation.

Proposition 6.17. Let $(u, f) \in \mathcal{W}_{2}(\mathbf{c}, \kappa ; B)$ be a holomorphic disk with boundary on an admissible L, and with $j$ positive punctures at Reeb chords $a_{1}, \ldots, a_{j}$ and $k$ negative punctures at Reeb chords $b_{1}, \ldots, b_{k}$. Then the index of $d \Gamma_{(u, f)}$ equals

$$
\begin{equation*}
\mu(B)+(1-j) n+\sum_{r=1}^{j} \nu\left(a_{r}\right)-\sum_{r=1}^{k} \nu\left(b_{r}\right) \tag{6.39}
\end{equation*}
$$

Remark 6.18. Note that (6.39) is independent of the choices of capping paths.

Proof. We simply translate the result of Proposition 6.14. At a positive puncture $p$, the tangent space corresponding to $\mathbb{R}+0 i(\mathbb{R}+i)$ in $\partial E_{p}$ is the lower (upper) one and at a negative puncture the situation is reversed. We must compare the rotation path $\lambda\left(V_{1}, V_{0}\right)$ used in the definition of the Conley-Zehnder index with the rotation used in the construction of the $\operatorname{arcs} \mathbf{f}_{i}$ in Proposition 6.14. At a negative puncture, the path $\mathbf{f}_{i}$ is the inverse path of $\lambda\left(V_{1}, V_{0}\right)$. Hence, the contribution to the Maslov index of $\mathbf{f}_{i}$ at a negative corner equals minus the contribution from $\lambda$. Consider the situation at positive puncture mapping to $a^{*}$. Let $\lambda\left(V_{1}, V_{0}\right)$ be the path used in the definition of the Conley-Zehnder index. Then, $\lambda\left(V_{1}, V_{0}\right)$ rotates the lower tangent space $V_{1}$ of $\Pi_{\mathbb{C}}(L)$ at $a^{*}$ to the upper $V_{0}$ according to $e^{s I}, 0 \leq s \leq \frac{\pi}{2}$, where $I$ is a complex structure compatible with $\omega$. Let $\lambda\left(V_{0}, V_{1}\right)$ be the path which rotates $V_{0}$ to $V_{1}$ in the same fashion. Then, the path $\mathbf{f}_{j}$ is the inverse path of $\lambda\left(V_{0}, V_{1}\right)$ and hence the contribution to the Maslov index of $\mathbf{f}_{j}$ equals the contribution from $\lambda\left(V_{1}, V_{0}\right)$ minus $n$.

To get the loop $B$ from $\Gamma_{A}$ (see Proposition 6.14), the arcs $\mathbf{f}_{i}$ must be removed and replaced by the arcs $\Gamma_{i}$, induced from the capping paths of the Reeb chords. A straightforward calculation gives

$$
\mu\left(\Gamma_{A}\right)=\mu(B)+\sum_{r=1}^{j} \nu\left(a_{r}\right)-n j-\sum_{s=1}^{k} \nu\left(b_{s}\right) .
$$

Hence,

$$
n+\mu\left(\Gamma_{A}\right)=\mu(B)+(1-j) n+\sum_{r=1}^{j} \nu\left(a_{r}\right)-\sum_{s=1}^{k} \nu\left(b_{s}\right) .
$$

q.e.d.
6.11. The index and the Conley-Zehnder index at a self tangency. In this section, we prove the analog of Proposition 6.17 for semiadmissible submanifolds. First, we need a definition of the ConleyZehnder index of a degenerate Reeb chord. Let $L \subset \mathbb{R} \times \mathbb{C}^{n}$ be a chord semi generic Legendrian submanifold. Let $c$ be the Reeb chord of $L$ such that $\Pi_{\mathbb{C}}(L)$ has a double point with self tangency along one direction at $c^{*}$. Let $c^{+}$and $c^{-}$be the end points of $c, z\left(c^{+}\right)>z\left(c^{-}\right)$. Let $V_{0}=d \Pi_{\mathbb{C}}\left(T_{c^{+}} L\right)$ and $V_{1}=d \Pi_{\mathbb{C}}\left(T_{c^{-}} L\right)$. Then $V_{0}$ and $V_{1}$ are Lagrangian subspaces of $\mathbb{C}^{n}$ such that $\operatorname{dim}_{\mathbb{R}}\left(V_{0} \cap V_{1}\right)=1$. Let $W \subset \mathbb{C}^{n}$ be the 1dimensional complex linear subspace containing $V_{0} \cap V_{1}$ and let $\mathbb{C}^{n-1}$ be the Hermitian orthogonal complement of $W$. Then $V_{0}^{\prime}=V_{0} \cap \mathbb{C}^{n-1}$ and $V_{1}^{\prime}=V_{1} \cap \mathbb{C}^{n-1}$ are transverse Lagrangian subspaces in $\mathbb{C}^{n-1}$. Pick a complex structure $I^{\prime}$ on $\mathbb{C}^{n-1}$ compatible with $\omega \mid \mathbb{C}^{n-1}$ such that $I^{\prime} V_{1}^{\prime}=V_{0}^{\prime}$. Define $\lambda\left(V_{1}, V_{0}\right)$ to be the path of Lagrangian planes
$s \mapsto V_{0} \cap V_{1} \times e^{s I^{\prime}} V_{1}^{\prime}$. Also pick a capping path $\gamma:[0,1] \rightarrow L$ with $\gamma(0)=c^{+}$and $\gamma(1)=c^{-}$. Then $\gamma$ induces a path $\Gamma$ of Lagrangian subspaces of $\mathbb{C}^{n}$. Define the Conley-Zehnder index of $c$ as

$$
\nu_{\gamma}(c)=\mu\left(\Gamma * \lambda\left(V_{1}, V_{0}\right)\right)
$$

Let $0<\epsilon<\theta$, where $\theta$ is the smallest non-zero complex angle of $L$ at $c$. Let $\mathcal{W}_{2, \epsilon}(\mathbf{c} ; \kappa)$ denote the space of maps with boundary conditions constructed from the Sobolev space with weight $\epsilon$ at each puncture mapping to $c$ and define $\widetilde{\mathcal{W}}_{2, \epsilon}(\mathbf{c} ; \kappa)$ as in Section 5.8. If $a$ is a Reeb chord of $L$, then let $\delta(a, c)=0$ if $a \neq c$ and $\delta(c, c)=1$. Again we suppress capping paths from the notation.

Proposition 6.19. Let $(u, f) \in \mathcal{W}_{2, \epsilon}(\mathbf{c}, \kappa ; B)$ and $(v, g) \in \widetilde{\mathcal{W}}_{2, \epsilon}(\mathbf{c}, \kappa$; $B)$ be holomorphic disks. If $\mathbf{c}=\left(a ; b_{1}, \ldots, b_{m}\right)$ where $a \neq c$, then the index of $d \Gamma_{(u, f)}$ equals

$$
\mu(B)+\nu(a)-\sum_{r=1}^{k}\left(\nu\left(b_{j}\right)+\delta\left(b_{j}, c\right)\right)
$$

and the index of $d \Gamma_{(v, g)}$ equals

$$
\mu(B)+\nu(a)-\sum_{r=1}^{k} \nu\left(b_{j}\right) .
$$

If $\mathbf{c}=\left(c ; b_{1}, \ldots, b_{m}\right)$, then the index of $d \Gamma_{(u, f)}$ equals

$$
\mu(B)+\nu(c)-\sum_{r=1}^{k}\left(\nu\left(b_{j}\right)+\delta\left(b_{j}, c\right)\right)
$$

and the index of $d \Gamma_{(v, g)}$ equals

$$
\mu(B)+(\nu(c)+1)-\sum_{r=1}^{k} \nu\left(b_{j}\right) .
$$

Remark 6.20. Note again that the index computations are independent of the choices of capping paths.

Proof. The proof is similar to the proof of Proposition 6.17. Consider first the $\bar{\partial}$-operator with boundary conditions determined by $(u, f)$ and acting on a Sobolev space with small negative weight. Again we need to compute the Maslov index contributions from the paths $\mathbf{f}_{i}$ in the loop $\Gamma_{A}$, where $\mathbf{f}_{i}$ fixes the common direction in the tangent spaces at a self tangency double point. Note that at a positive puncture $c$ the contribution is now the contribution of $\lambda\left(V_{1}, V_{0}\right)$ minus $(n-1)$. At a negative puncture it is again minus the contribution of $\lambda\left(V_{1}, V_{0}\right)$. Applying Proposition 6.16, the first and third index calculations above follow. Noting that the tangent space of $\widetilde{\mathcal{W}}_{2, \epsilon}(\mathbf{c}, \kappa, B)$ is obtained from that of
$\mathcal{W}_{2, \epsilon}(\mathbf{c}, \kappa ; B)$ by adding one $\mathbb{R}$-direction for each puncture mapping to $c$ the other index formulas follow as well. q.e.d.

## 7. Transversality

In this section, we show how to achieve transversality (or "surjectivity") for the linearized $\bar{\partial}$ equation by perturbing the Lagrangian boundary condition. When proving transversality for some Floer-type theory, it is customary to show that solution-maps are "somewhere injective" (see $[\mathbf{2 2}, \mathbf{1 6}]$, for example). One then constructs a small perturbation, usually of the almost complex structure or the Hamiltonian term, which is supported near points where the map is injective. With a partial integration argument, these perturbations eliminate non-zero elements of the cokernel of $\bar{\partial}$.

For our set-up, we perturb the Lagrangian boundary condition. In Sections 7.1 through 7.4, we describe the space of perturbations for the chord generic, one-parameter chord generic, and chord semi-generic cases. Although we do not have an injective (boundary) point, we exploit the fact that there is only one positive puncture, and hence, by Lemma 2.1, the corresponding double point can represent a corner only once. Of course other parts of the boundary can map to this corner elsewhere, but not at other boundary punctures. With this observation, we prove transversality in Sections 7.7 and 7.9 first for the open set of non-exceptional maps, defined in Section 7.6 and from this for all maps provided the expected kernel has sufficiently low dimension. We also prove some results in Sections 7.10 and 7.12 which will be useful later for the degenerate gluing of Section 8 .

### 7.1. Perturbations of admissible Legendrian submanifolds. Let

 $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an admissible Legendrian submanifold. Let $a(L)$ denote the minimal distance between the images under $\Pi_{\mathbb{C}}$ of two distinct Reeb chords of $L$ and let $A(L)$ be such that $\Pi_{\mathbb{C}}(L)$ is contained in the ball $B(0, A(L)) \subset \mathbb{C}^{n}$. Fix $\delta>0$ and $R>0$ such that $\delta \ll a(L)$ and such that $R \gg A(L)$.Definition 7.1. Let $\operatorname{Ham}(L, \delta, R)$ be the linear space of smooth functions $h: \mathbb{C}^{n} \rightarrow \mathbb{R}$ with support in $B(0, R)$ and satisfying the following two conditions for any Reeb chord $c$.
(i) The restriction of $h$ to $B\left(c^{*}, \delta\right)$ is real analytic,
(ii) The differential of $h$ satisfies $D h\left(c^{*}\right)=0$ and also $h\left(c^{*}\right)=0$.

We are going to use Hamiltonian vector fields of elements in $\operatorname{Ham}(L, \delta$, $R$ ) to perturb $L$. Condition (i) ensures that $L$ stays admissible, and (ii) that the set of Reeb chords $\left\{c_{0}, \ldots, c_{m}\right\}$ of $L$ remains fixed.

Lemma 7.2. The space $\operatorname{Ham}(L, \delta, R)$ with the $C^{\infty}$-norm is a Banach space.

Proof. Using the characterization of real analytic functions as smooth functions, the derivatives of which satisfy certain uniform growth restrictions, one sees that the limit of a $C^{\infty}$-convergent sequence of real analytic functions on an open set is real analytic. q.e.d.

Lemma 7.3. If $L$ is admissible and $h \in \operatorname{Ham}(L, \delta, R)$, then $\tilde{\Phi}_{h}(L)$ (see Section 3) is admissible.

Proof. For each Reeb chord $c$, the Hamiltonian vector field is real analytic in $B\left(c^{*}, \delta\right)$. Also, $\Phi_{h}\left(c^{*}\right)=c^{*}$ and hence, there exists a neighborhood $W$ of $c^{*}$ such that $\Phi_{h}^{t}(W) \subset B\left(c^{*}, \delta\right)$ for $0 \leq t \leq 1$. A well-known ODE-result implies that the flow of a real analytic vector field depends in a real analytic way on its initial data. This shows that $\tilde{\Phi}_{h}(L)$ is admissible.
q.e.d.
7.2. Perturbations of 1-parameter families of admissible submanifolds. Let $L_{t}, t \in[0,1]$ be an admissible 1-parameter family of Legendrian submanifolds without self-tangencies. Let $a=\min _{0 \leq t \leq 1} a\left(L_{t}\right)$ and $A=\max _{0 \leq t \leq 1} A\left(L_{t}\right)$. Fix $\delta>0$ and $R>0$ such that $\delta \ll a$ and $R \gg A$.

We define a continuous family of isomorphisms $\operatorname{Ham}\left(\delta, R, L_{0}\right) \rightarrow$ $\operatorname{Ham}\left(\delta, R, L_{t}\right), 0 \leq t \leq 1$. Let $\left(c_{1}(t), \ldots, c_{m}(t)\right)$ be the Reeb chords of $L_{t}$. Then $\left(c_{1}^{*}(t), \ldots, c_{m}^{*}(t)\right), 0 \leq t \leq 1$ is a continuous curve in $\left(\mathbb{C}^{n}\right)^{m}$. Let $\psi^{t}: B(0, R) \rightarrow B(0, R)$ be a continuous family of compactly supported diffeomorphisms which when restricted to $B\left(c_{j}^{*}(0), \delta\right)$, $j=1, \ldots, m$ agree with the map

$$
z \mapsto z+\left(c_{j}^{*}(t)-c_{j}^{*}(0)\right)
$$

Composition with $\psi^{t}$ can be used to give the space

$$
\operatorname{pHam}\left(L_{t}, \delta, R\right)=\bigcup_{0 \leq t \leq 1} \operatorname{Ham}\left(L_{t}, \delta, R\right)
$$

the structure of a Banach manifold which is a trivial bundle over $[0,1]$. We note that if $(h, t)$ in $\mathrm{pHam}\left(L_{t}, \delta, R\right)$, then Lemma 7.3 implies that $\tilde{\Phi}_{a}\left(L_{t}\right)$ is admissible.
7.3. Bundles over perturbations. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an admissible chord generic Legendrian submanifold. Above we constructed a smooth map of the Banach space $\operatorname{Ham}(L, \delta, R)$ into the space of admissible chord generic Legendrian embeddings of $L$ into $\mathbb{C}^{n} \times \mathbb{R}$.

Let $\mathbf{c}=\left(c_{0}, c_{1}, \ldots, c_{m}\right)$ be Reeb chords of $L$ and let $\epsilon \in[0, \infty)^{m}$ and consider, as in Section 5.1, the space

$$
\begin{equation*}
\mathcal{W}_{2, \epsilon, \operatorname{Ham}(L, \delta, R)}(\mathbf{c}) \tag{7.1}
\end{equation*}
$$

and its tangent space

$$
\begin{equation*}
T_{(w, f, \kappa, a)} \mathcal{W}_{2, \epsilon, \operatorname{Ham}(L, \delta, R)} \approx T_{(w, f)} \mathcal{W}_{2, \epsilon} \oplus T_{\kappa} \mathcal{C}_{m} \oplus \operatorname{Ham}(L, \delta, R) \tag{7.2}
\end{equation*}
$$

In a similar way, we consider for a 1-parameter family $L_{t}$ the space

$$
\begin{equation*}
\mathcal{W}_{2, \epsilon, \mathrm{p} H a m}\left(L_{t}, \delta, R\right)(\mathbf{c}) \tag{7.3}
\end{equation*}
$$

and its tangent space.
For $\Lambda=\operatorname{Ham}(L, \delta, R)$ or $\Lambda=\mathrm{pHam}\left(L_{t}, \delta, R\right)$, consider also the bundle map $\Gamma: \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c}) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}[0]\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$. Here, we are thinking of the spaces as bundles over $\operatorname{Ham}(L, \delta, R)$ and we denote projection onto this space by pr. To emphasize this, we will write ( $\Gamma$, pr) instead of just $\Gamma$ in the sequel. The differential $d \Gamma$ was calculated in Lemma 5.14.
7.4. Perturbations in the semi-admissible case. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a semi-admissible Legendrian submanifold. Let $\left(c_{0}, \ldots, c_{m}\right)$ be the Reeb chords of $L$. Assume that the self tangency Reeb chord is $c_{0}$, that $c_{0}^{*}=0$, and that $L$ has standard form in a neighborhood of 0 , see Definition 3.3.

Let $a(L)$ denote the minimal distance between the images under $\Pi_{\mathbb{C}}$ of two distinct Reeb chords of $L$. Fix $\delta>0$ such that $\delta \ll a(L)$. For $r>0$, let $C(r)=\mathbb{C} \times B^{\prime}(0, r) \subset \mathbb{C}^{n}$, where $B^{\prime}(0, r)$ is the $r$-ball in $\mathbb{C}^{n-1} \approx\left\{z_{1}=0\right\}$, where as, always, $\left(z_{1}, \ldots, z_{n}\right)=\left(x_{1}+i y_{1}, \ldots, x_{n}+i y_{n}\right)$ are coordinates on $\mathbb{C}^{n}$.

Definition 7.4. Let $\operatorname{Ham}_{0}(L, \delta)$ be the linear space of smooth functions $h: \mathbb{C}^{n} \rightarrow \mathbb{R}$ with support in $C(10 \delta) \cup \bigcup_{j \geq 1} B\left(c_{j}^{*}, 10 \delta\right)$ and satisfying the following conditions.
(i) The restriction of $h$ to $B\left(c_{j}^{*}, \delta\right) 1 \leq j \leq m$ is real analytic,
(ii) In $C(10 \delta), \frac{\partial h}{\partial x_{1}}=0=\frac{\partial h}{\partial y_{1}}$ and the restriction of $h$ to $C(\delta)$ is real analytic.
(iii) The differential of $h$ satisfies $D h\left(c_{j}^{*}\right)=0$ and also $h\left(c_{j}^{*}\right)=0$, for all $j$.

Lemma 7.5. The space $\operatorname{Ham}_{0}(L, \delta)$ with the $C^{\infty}$-norm is a Banach space.

Proof. See Lemma 7.2 and note that the restriction of $h$ to $C(10 \delta)$ can be identified with a function of ( $n-1$ )-complex variables supported in $B^{\prime}(0,10 \delta)$.
q.e.d.

Let $\tilde{\Phi}_{h}$ be the Legendrian isotopy which is defined by using the flow of $h$ locally around the Reeb chords of $L$. This is well-defined for $h$ sufficiently small. Let $\operatorname{Ham}_{0}(L, \delta, s)$ denote the $s$-ball around 0 in $\operatorname{Ham}_{0}(L, \delta)$.

Lemma 7.6. There exists $s>0$ such that for $h \in \operatorname{Ham}_{0}(L, \delta, s)$, $\tilde{\Phi}_{h}(L)$ is an admissible chord semi-generic Legendrian submanifold.

Proof. Note that the product structure in $C(10 \delta)$ is preserved since $h$ does not depend on $\left(x_{1}, y_{1}\right)$. Moreover, the isotopy is fixed in the region $B(0,2+\epsilon) \backslash B(0,2)$ for $s$ and $\delta$ sufficiently small. q.e.d.

We have defined a smooth map of $\operatorname{Ham}_{0}(L, \delta, s)$ into the space of admissible chord semi-generic Legendrian submanifolds and this maps fulfills the conditions on $\Lambda$ in Section 5.8. We can therefore construct the spaces

$$
\begin{equation*}
\mathcal{W}_{2, \epsilon, \operatorname{Ham}_{0}(L, \delta, s)}, \quad \text { and } \quad \widetilde{\mathcal{W}}_{2, \epsilon, \operatorname{Ham}_{0}(L, \delta)}, \tag{7.4}
\end{equation*}
$$

see Sections 5.1 and 5.8 , respectively. Moreover, as there, we will consider the $\bar{\partial}$-map and its linearization.
7.5. Consequences of real analytic boundary conditions. For $r>0$, let $E_{+}=\{z \in \mathbb{C}:|z|<r, \operatorname{Im}(z) \geq 0\}$. If $w:\left(E_{+}, \partial E_{+}\right) \rightarrow$ $\left(\mathbb{C}^{n}, M\right)$ where $M$ is a real analytic Lagrangian submanifold and $w$ is holomorphic in the interior and continuous on the boundary, then by Schwartz-reflection principle, $w$ extends in a unique way to a holomorphic map $w^{d}: E \rightarrow \mathbb{C}^{n}$ mapping $\operatorname{Im}(z)=0$ to $M$, where $E=\{z \in$ $\mathbb{C}:|\zeta|<r\}$ for $r$ sufficiently small. We call $w^{d}$ the double of $w$.

Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a chord (semi-) generic Legendrian submanifold.
Lemma 7.7. Let $p$ be a point in $U \subset L$ such that $\Pi_{\mathbb{C}}(U)$ is real analytic, where $U \subset L$ is a neighborhood of $p$ on which $\Pi_{\mathbb{C}}$ is injective. Assume that

$$
w:\left(E_{+}, \partial E_{+}, 0\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(U), p\right)
$$

is holomorphic. Then, there is a holomorphic function $u$ with Taylor expansion at 0,

$$
u(z)=a_{0}+a_{1} z+\ldots, a_{0} \neq 0
$$

such that $w(z)=p+z^{k} u(z)$ for some integer $k>0$.
Proof. The double $w^{d}$ has a Taylor expansion. q.e.d.
Lemma 7.8. Let $p, U$ and $L$ satisfy the conditions of Lemma 7.7. Assume that $w: D_{m} \rightarrow \mathbb{C}^{n}$ is holomorphic with boundary on $\Pi_{\mathbb{C}}(L)$. Then $w^{-1}(p) \cap \partial D_{m}$ is a finite set.

Proof. Using Lemma 4.6, we may find $M>0$ such that there are no preimages of $p$ in $\cup_{j} E_{p_{j}}[M]$. Since the complement of $\cup_{j} E_{p_{j}}[M]$ is compact, the lemma now follows from Lemma 7.7. q.e.d.

Lemma 7.9. Let $p \in L$ satisfy the conditions of Lemma 7.7, and let

$$
\begin{equation*}
w_{1}, w_{2}:\left(E_{+}, \partial E_{+}, 0\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L), p\right) \tag{7.5}
\end{equation*}
$$

be holomorphic maps such that $w_{2}$ maps one of the components $I$ of $\partial E_{+} \backslash\{0\}$ to $w_{1}(I)$. Then there exists a map $\hat{w}: E \rightarrow \mathbb{C}^{n}$ and integers $k_{j} \geq 1$ such that $w_{j}^{d}(z)=\hat{w}\left(z^{k_{j}}\right), j=1,2$.

Proof. As above, we may reduce to the case when $\Pi_{\mathbb{C}}(L)=\mathbb{R}^{n} \subset \mathbb{C}^{n}$. The images $C_{j}=w_{j}^{d}(E), j=1,2$ are analytic subvarieties of complex dimension 1 which intersects in a set of real dimension 1. Hence, they agree. Projection of $C=C_{1}=C_{2}$ onto a generic complex line through
$p$ identifies $C$ (locally) with the standard cover of the disk possibly branched at 0 . This gives the map $\hat{w}$. q.e.d.
7.6. Exceptional holomorphic maps. Let $\Lambda$ be one of the spaces $\operatorname{Ham}(L, \delta, R), \operatorname{pHam}(L, \delta, R)$, or $\operatorname{Ham}_{0}(L, \delta, s)$. Let $(w, f, \lambda) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})$ (or $\widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}(\mathbf{c})$ ) be a holomorphic disk and let $q$ be a point on $\partial D_{m}$ such that $w(q)$ lies in a region where $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ is real analytic. Assume that $d w(q)=0$. Since $w$ has a Taylor expansion around $q$ in this case, we know there exists a half-disk neighborhood $E$ of $q$ in $D_{m}$ such that $q$ is the only critical point of $w$ in $E$. The boundary $\partial E$ is subdivided by $q$ into two arcs $\partial E \backslash\{q\}=I_{+} \cup I_{-}$. We say that $q$ is an exceptional point of $(w, f)$ if there exists a neighborhood $E$ as above such that $w\left(I_{+}\right)=w\left(I_{-}\right)$.

Definition 7.10. Let $(w, f, \lambda) \in \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})$, where $\mathbf{c}=\left(c_{0}(\lambda), c_{1}(\lambda)\right.$, $\left.\ldots, c_{m}(\lambda)\right)$ and $c_{0}(\lambda)$ is the Reeb chord on $L_{\lambda}$ of the positive puncture of $D_{m+1}$. Let $B_{1}(\lambda)$ and $B_{2}(\lambda)$ be the two local branches of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$ at $c_{0}^{*}(\lambda)$. Then $(w, f)$ is exceptional holomorphic if it has two exceptional points $q_{1}$ and $q_{2}$ with $w\left(q_{1}\right)=w\left(q_{2}\right)=c_{0}^{*}(\lambda)$ and if a neighborhood in $\partial D_{m}$ of $q_{j}$ maps to $B_{j}(\lambda), j=1,2$.

Definition 7.11. Let $\mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c})\left(\widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c})\right)$ denote the complement of the closure of the set of all exceptional holomorphic maps in $\mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})$ $\left(\widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}(\mathbf{c})\right)$.

We note that $\mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c})$ is an open subspace of a Banach manifold and hence a Banach manifold itself.

### 7.7. Transversality on the complement of exceptional holomorphic maps in the admissible case.

Lemma 7.12. For $L$ admissible (respectively $L_{t}$ a 1-parameter family of admissible submanifolds) the bundle map, see Section 5.7

$$
(\Gamma, \operatorname{pr}): \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c}) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}[0]\left(D_{m}, T^{*} D_{m} \otimes \mathbb{C}^{n}\right)
$$

where $\Lambda=\operatorname{Ham}(L, \delta, R)\left(\right.$ respectively $\left.\Lambda=\mathrm{pHam}\left(L_{t}, \delta, R\right)\right)$ is transverse to the 0-section.

Proof. The proof for 1-parameter families $L_{t}$ is only notationally more difficult. We give the proof in the stationary case. We must show that if $w: D_{m} \rightarrow \mathbb{C}^{n}$ is a (non-exceptional) holomorphic map (in the conformal structure $\kappa$ on $D_{m}$ ) which represents a holomorphic disk $(w, f)$ with boundary on $L=L_{\lambda}$ (without loss of generality, we take $\lambda=0$ below) then

$$
d \Gamma\left(T_{((w, f), \kappa, 0)} \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})\right)=\mathcal{H}_{1, \epsilon}\left(D_{m}, T^{*} D_{m} \otimes \mathbb{C}^{n}\right)
$$

i.e., $d \Gamma$ is surjective. To show this, it is enough to show that

$$
\left\{d \Gamma\left(T_{((w, f), \kappa, 0)} \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})\right)\right\}^{\perp}=\{0\}
$$

where $V^{\perp}$ denotes the annihilator with respect to the $L^{2}$-pairing of $V \subset \mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{*} D_{m} \otimes \mathbb{C}^{n}\right)$ in its dual space.

An element $u$ in this annihilator satisfies

$$
\begin{equation*}
\int_{D_{m}}\langle\bar{\partial} v, u\rangle d A=0 \tag{7.6}
\end{equation*}
$$

for all $v \in T_{w} \mathcal{B}_{2, \epsilon}(0, r)$. Here $d A$ is the area form on $D_{m}$. Lemma 6.1 implies that $u$ can be represented by a $C^{2}$-function which is anti-holomorphic.

We note that integrals of the form

$$
\begin{equation*}
\int_{D_{m}}\langle\phi, \psi\rangle d A \tag{7.7}
\end{equation*}
$$

where $\langle$,$\rangle is the inner product on T^{*} D_{m}$ and where $\phi$ and $\psi$ are sections, are conformally invariant. We may therefore compute integrals of this form in any conformal coordinate system on the disk $D_{m}$.

Restrict attention to a small neighborhood of the image of the positive corner at $c_{0}^{*}$. Recall that $w$ is assumed non-exceptional and consider a branch $B$ of $\Pi_{\mathbb{C}}(L)$ at $c_{0}^{*}$ such that $w$ does not have an exceptional point mapping to $c_{0}^{*} \in B$. Since $B$ is real analytic we may biholomorphically identify ( $\mathbb{C}^{n}, B, c_{0}^{*}$ ) with ( $\left.\mathbb{C}^{n}, \mathbb{R}^{n}, 0\right)$.

Let $p$ be the positive puncture on $D_{m}$. For $M$ large enough, by Proposition 4.6, the image of the component of $\partial E_{p}[M]$ which lies in $B$ is a regular oriented curve. Denote it by $\gamma$. For simplicity we assume that the component mapping to $\gamma$ is $[M, \infty) \times\{0\} \subset E_{p_{0}}[M]$ and we let $E_{0}=[M, \infty) \times\left[0, \frac{1}{2}\right)$.

Let $p_{1}, \ldots, p_{r}$ be the preimages under $w$ of $c_{0}^{*}$ with the property that one of the components of a punctured neighborhood of $p_{j}$ in $\partial D_{m}$ maps to $\gamma$. Note that $r<\infty$ by Proposition 4.6 and Lemma 7.8 and that by shrinking $\gamma$, we may assume that all these images are exactly $\gamma$.

We say that a point $p_{j}$ is positive if close to $p_{j}, w$ and the natural orientation on the boundary of $\partial D_{m}$ induce the positive orientation on $\gamma$. Otherwise, we say it is negative.

The image of the other half of the punctured neighborhood of $p_{1}$ in $\partial D_{m}$ maps to a curve $\gamma^{\prime}$ under $w$. Our assumption that $w$ is nonexceptional guarantees that $\gamma$ and $\gamma^{\prime}$ are distinct.

Let $w_{j}$ denote the restriction of $w$ to a small neighborhood of $p_{j}$. Let $E=\{z \in \mathbb{C}:|z|<r\}$, let $E_{+}=\{z \in E: \operatorname{Im}(z) \geq 0\}$, and let $E_{-}=\{z \in E: \operatorname{Im}(z) \leq 0\}$. Lemma 7.9 implies that we can find a map $\hat{w}: E \rightarrow \mathbb{C}^{n}$ and coordinate neighborhoods $\left(E_{ \pm}(j), \partial E_{ \pm}(j)\right)$ of $p_{j}$ (where the sign $\pm$ is that of $p_{j}$ ) such that $w_{j}^{d}(z)=\hat{w}\left(z^{k_{j}}\right)$ for each $j$. Note that $w$ non-exceptional implies all $k_{j}$ are odd.

Let $k=k_{1} k_{2} \ldots k_{r}$ and let $\hat{k}_{j}=\frac{k}{k_{j}}$. Let $\phi_{j}: E \rightarrow E(j)$ be the map $z \mapsto z^{\hat{k}_{j}}$. Consider the restrictions $u_{j}$ of the anti-holomorphic map $u$
to the neighborhoods $\left(E_{ \pm}(j), \partial E_{ \pm}(j)\right)$. Because of the real analytic boundary conditions (recall that $\left(B, \mathbb{C}^{n}\right)$ is biholomorphically identified with $\left(\mathbb{R}^{n}, \mathbb{C}^{n}\right)$ ), these maps can be doubled using Schwartz reflection principle. Use $\phi_{j}$ to pull-back the maps $u_{j}$ and $w_{j}$ to $E$.

Let $a: \mathbb{C}^{n} \rightarrow \mathbb{R}$ be any smooth function with support in a small ball around a point $q^{\prime} \in \gamma^{\prime}$, where $q^{\prime}$ is chosen so that no point outside $\bigcup_{j} E_{ \pm}(j)$ in $\partial D_{m}$ maps to $q^{\prime}$. (There exists such a point because of the asymptotics of $w$ at punctures and Lemma 7.7.) Let $Y_{a}$ be the Hamiltonian vector field associated with $a$, see Section 3.3.

If $v$ is a smooth function with support in $\bigcup_{j} E_{ \pm}(j)$ which is real and holomorphic on $\bigcup_{j} \partial E_{ \pm}(j)$, if $\xi+i \eta$ are coordinates on $E_{ \pm}(j)$, and if the support of $a$ is sufficiently small, then

$$
\begin{equation*}
0=\int_{D_{m}}\left\langle\bar{\partial}\left(Y_{a}+v\right), u\right\rangle d A \tag{7.8}
\end{equation*}
$$

$$
\begin{equation*}
=-\sum_{j} \int_{E_{ \pm}(j)}\left\langle Y_{a}+v, \partial u\right\rangle d \xi \wedge d \eta+\sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \xi \tag{7.9}
\end{equation*}
$$

$$
\begin{equation*}
=\sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \xi . \tag{7.10}
\end{equation*}
$$

The equality in (7.8) follows since $u$ is an element of the annihilator and since $a$ can be arbitrarily well $C^{2}$-approximated by elements in $\operatorname{Ham}(L, \delta, R)$. The equality in (7.9) follows by partial integration and the restrictions on the supports of $a$ and $v$. The equality in (7.10) follows from $\partial u=0$. Taking $a=0$, we see, since we are free to choose $v$, that $u$ must be real valued on $\partial E_{ \pm}(j)$ for every $j$.

We then take $v=0$ and express the integral in (7.10) as an integral over $I_{+}=\{x+0 i: x>0\} \subset E$. Note that if $\xi+i \eta$ are coordinates on $E(j)$ then under the identification by $\phi_{j}, d \xi=d x^{\hat{k}_{j}}=\hat{k}_{j} x^{\hat{k}_{j}-1} d x$ and

$$
\begin{align*}
& \sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i Y_{a}, u\right\rangle d \xi  \tag{7.11}\\
& \quad=\int_{I_{+}}\left\langle-i Y_{a}\left(\hat{w}\left(z^{k}\right)\right), \sum_{j} \sigma(j) \hat{k}_{j} \bar{z}^{\hat{k}_{j}-1} u_{j}\left(z^{\hat{k}_{j}}\right)\right\rangle d x
\end{align*}
$$

where $\sigma(j)= \pm 1$ equals the sign of $p_{j}$. Thus, if

$$
\alpha(z)=\sum_{j} \sigma(j) \hat{k}_{j} z^{\hat{k}_{j}-1} u_{j}\left(z^{\hat{k}_{j}}\right),
$$

then $\alpha$ is antiholomorphic and by varying $a$, we see that $\alpha$ vanishes in the $\mathbb{R}^{n}$-directions along an arc in $I_{+}$. Therefore, $\alpha$ vanishes identically on $E$.

Pick now instead $a$ supported in a small ball around $q$ in $\gamma$. With the same arguments as above, we find

$$
\begin{align*}
0= & \int_{D_{m}}\left\langle\bar{\partial}\left(Y_{a}+v\right), u\right\rangle d A  \tag{7.12}\\
= & -\int_{E_{0}}\left\langle Y_{a}+v, \partial u\right\rangle d \tau \wedge d t-\sum_{j} \int_{E_{ \pm}(j)}\left\langle Y_{a}+v, \partial u\right\rangle d \xi \wedge d \eta \\
& +\int_{[M, \infty)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \tau+\sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \xi \\
= & \int_{[M, \infty)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \tau+\sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i\left(Y_{a}+v\right), u\right\rangle d \xi
\end{align*}
$$

and conclude that $u(\tau, 0) \in \mathbb{R}^{n}$ for $\tau \in[M, \infty)$ as well.
Again taking $v=0$, we get for the last integral in (7.12)

$$
\begin{equation*}
\left.\sum_{j} \int_{\partial E_{ \pm}(j)}\left\langle-i Y_{a}, u\right\rangle d \xi=\int_{I_{-}}\left\langle-i Y_{a}\left(\hat{w}\left(z^{k}\right)\right), \alpha(z)\right)\right\rangle d x=0 \tag{7.13}
\end{equation*}
$$

where $I_{-}=\{x+0 i: x<0\} \subset E$, and where the last equality follows since $\alpha=0$. Equations (7.13) and (7.12) together implies (by varying $a$ ) that $u$ must vanish along an arc in $[M, \infty)$. Since $u$ is antiholomorphic, it must then vanish everywhere. This proves the annihilator is 0 and the lemma follows.

> q.e.d.

Remark 7.13. In the case that $w$ has an injective point on the boundary, the above argument can be shortened. Namely, under this condition, there is an arc $A$ on the boundary of $D_{m}$ where $w$ is injective and varying $v$ and $a$, there we see that $u$ must vanish along $A$ and therefore everywhere. Oh achieves transversality using boundary perturbations assuming an injective point [25].

Corollary 7.14. Let $\mathbf{c}=a b_{1} \ldots b_{m}$. For a Baire set of $h \in \operatorname{Ham}(L, \delta$, $R)=\Lambda, \Gamma^{-1}(0) \cap \operatorname{pr}^{-1}(h) \cap \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c} ; A)$ is a finite dimensional smooth manifold of dimension

$$
\mu(A)+\nu_{\gamma}(a)-\sum_{j} \nu_{\gamma}\left(b_{j}\right)+\max (0, m-2)
$$

For a Baire set of $h \in \operatorname{pHam}(L, \delta, R)=\Lambda \Gamma^{-1}(0) \cap \operatorname{pr}^{-1}(h) \cap \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c} ; A)$ is a finite dimensional smooth manifold of dimension

$$
\mu(A)+\nu_{\gamma}(a)-\sum_{j} \nu_{\gamma}\left(b_{j}\right)+\max (0, m-2)+1
$$

Proof. Let $Z \subset \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c} ; A)$ denote the inverse image of the 0 -section in $\mathcal{H}_{1, \epsilon, \Lambda}[0]\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)$ under ( $\left.\Gamma, \operatorname{pr}\right)$. By the implicit function theorem and Lemma 7.12, $Z$ is a submanifold. Consider the restriction of the projection $\pi: Z \rightarrow \Lambda$. Then $\pi$ is a Fredholm map of index equal to the index of the Fredholm section $\Gamma$. An application of the Sard-Smale theorem shows that for generic $\lambda \in \Lambda, \pi^{-1}(\lambda)$ is a submanifold of dimension given by the Fredholm index of $\Gamma$. Note that in the first case, the restriction of $d \Gamma$ to the complement of the $\max (0, m-2)$-dimensional subspace $T \mathcal{C}_{m} \subset T \mathcal{W}_{2, \epsilon}(\mathbf{c} ; A)$ is an operator of the type considered in Proposition 6.17. Thus, the proposition follows in this first case. In the second case, we restrict to a $(\max (0, m-2)+1)$-codimensional subspace instead.
q.e.d.
7.8. General transversality in the admissible case. If $\mathbf{c}$ is a collection of Reeb chords, we define $l(\mathbf{c})$ as the number of elements in c. We note that if $(f, w)$ is a holomorphic disk with boundary on $L$ with $r$ punctures, then, if $r \leq 2$, the kernel of $d \Gamma$ at $(f, w)$ is at least ( $3-r$ )-dimensional. This is a consequence of the existence of conformal reparameterizations in this case.

Theorem 7.15. For a dense open set of $h \in \operatorname{Ham}(L, \delta, R)(h \in$ $\mathrm{pHam}(L, \delta, R)), \Gamma^{-1}(0) \cap \mathrm{pr}^{-1}(h) \subset \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})$ is a finite dimensional $C^{1}$-smooth manifold of dimension as in Corollary 7.14, provided this dimension is $\leq 1$ if $l(\mathbf{c}) \geq 3$ and $\leq 1+(3-l(\mathbf{c}))$ otherwise.

Proof. After Corollary 7.14, we need only exclude holomorphic disks in the closure of exceptional holomorphic disks. Let $a \in \operatorname{Ham}(a \in$ pHam ) be such that $\Gamma^{-1}(0)$ is regular. Then the same is true for $\tilde{a}$ in a neighborhood of $a$. Now assume there exists a holomorphic disk in the closure of exceptional holomorphic disks at $a$. Then there must exists an exceptional holomorphic disk for some $\tilde{a}$ in the neighborhood. However, such a disk $w$ has $k \geq 2$ points mapping to the image of the positive puncture and with $w\left(I_{+}\right)=w\left(I_{-}\right)$. It is then easy to construct (by "moving the branch point") a $k$-parameter $(k+(3-l(\mathbf{c}))$ parameter if $l(\mathbf{c}) \leq 2$ ) family of distinct (since the location of the branch point changes) non-exceptional holomorphic disks with boundary on $L(\tilde{a})$. This contradicts the fact that the dimension of $\Gamma^{-1}(0)$ is $<k$ $(<k+(3-l(\mathbf{c}))$ for every $\tilde{a}$ in the neighborhood. q.e.d.
Proof of Proposition 2.2. If the number of punctures is $\geq 3$, the proposition is just Theorem 7.15. The case of fewer punctures can be reduced to that of many punctures as in Section 8.6. q.e.d.

Corollary 7.16. For chord generic admissible Legendrian submanifolds in a Baire set of such manifolds, no rigid holomorphic disk with boundary on $L$ decays faster than $e^{-(\theta+\delta)|\tau|}$ close to any of its punctures mapping to a Reeb chord c. Here $\theta$ is the smallest complex angle of the

Reeb chord $c, \delta>0$ is arbitrary, and $\tau+i t$ are coordinates near the puncture.

Proof. Such a holomorphic disk would lie in $\mathcal{W}_{2, \epsilon}(\mathbf{c})$, where the component of $\epsilon$ corresponding to the puncture mapping to the Reeb chord $c$ is larger than $\theta$. By Proposition 6.16, this change of weight lowers the Fredholm index of $d \Gamma$ by at least 1 . Since the Fredholm index of $d \Gamma$ with smaller weight (e.g., 0 -weight) is the minimal which allows for existence of disks, the lemma follows from Theorem 7.15. q.e.d.

Proof of Proposition 2.9. The first statement in the proposition follows exactly as above. To see that handle slides appear at distinct times, let ( $a_{1} \mathbf{b}_{1} ; A_{1}$ ) and ( $a_{2} \mathbf{b}_{2} ; A_{2}$ ) be such that

$$
\mu\left(A_{1}\right)+\left|a_{1}\right|-\left|\mathbf{b}_{1}\right|=\mu\left(A_{2}\right)+\left|a_{2}\right|-\left|\mathbf{b}_{2}\right|=0
$$

and consider the bundle $\mathcal{W}_{2, \Lambda}\left(a_{1} \mathbf{b}_{1} ; A_{1}\right) \tilde{\times} \mathcal{W}_{2, \Lambda}\left(a_{2} \mathbf{b}_{2} ; A_{2}\right)$. Here $\tilde{\times}$ denotes the fiberwise product where, in the fibers, the deformation coordinates $\left(t_{1}, t_{2}\right)$ are restricted to lie in the diagonal: $t_{1}=t_{2}=t$. This is a bundle over $\Lambda$, and $\Gamma$ induces a bundle map to the bundle $\mathcal{H}_{1, \Lambda}\left(D_{m_{1}}, \mathbb{C}^{n}\right) \times \mathcal{H}_{1, \Lambda}\left(D_{m_{2}}, \mathbb{C}^{n}\right)$, where $\times$ denotes fiberwise product. It is then easy to check that $\Gamma$ is a Fredholm section of index -1 . As in Theorem 7.15, we see that $d \Gamma$ is surjective and that the inverse image of the 0 -section intersected with $\mathrm{pr}^{-1}(h)$ is empty for generic $h$. This shows that the handle slides appear at distinct times.

The statement about all rigid disks being transversely cut out at a handle slide instant can be proved in a similar way: let ( $a_{1} \mathbf{b}_{1} ; A_{1}$ ) be as above and let $\left(a_{3} \mathbf{b}_{3} ; A_{3}\right)$ be such that

$$
\mu\left(A_{3}\right)+\left|a_{3}\right|-\left|\mathbf{b}_{3}\right|=1 .
$$

Consider the bundle

$$
\mathcal{W}_{2, \Lambda}\left(a_{3} \mathbf{b}_{3} ; A_{3}\right) \tilde{\times} \mathcal{W}_{2, \Lambda}\left(a_{3} \mathbf{b}_{3} ; A_{3}\right) \tilde{\times} \mathcal{W}_{2, \Lambda}\left(a_{1} \mathbf{b}_{1} ; A_{1}\right),
$$

and the bundle map $\Gamma$ defined in the natural way with target

$$
\mathcal{H}_{1, \Lambda}\left(D_{m_{3}}, \mathbb{C}^{n}\right) \times \mathcal{H}_{1, \Lambda}\left(D_{m_{3}}, \mathbb{C}^{n}\right) \times \mathcal{H}_{1, \Lambda}\left(D_{m_{1}}, \mathbb{C}^{n}\right)
$$

Then the map $\Gamma$ has Fredholm index 0 and as above, we see $d \Gamma$ is surjective. Hence, $\Gamma^{-1}(0) \cap \mathrm{pr}^{-1}(h)$ is a transversely cut out 0 -manifold for generic $h$. We show that this implies that if $t$ is such that $\mathcal{M}_{A_{1}}^{t}\left(a_{1} ; \mathbf{b}_{1}\right)=$ $\{(v, g)\} \neq \emptyset$, then $\mathcal{M}_{A_{3}}^{t}\left(a_{3} ; \mathbf{b}_{3}\right)$ is transversally cut out. Let $(u, f) \in$ $\mathcal{M}_{A_{3}}^{t}\left(a_{3} ; \mathbf{b}_{3}\right)$ and assume the differential $d \Gamma_{(u, f)}^{t}$, which is a Fredholm operator of index 0 is not surjective. Then it has a cokernel of dimension $d>0$. Furthermore, the image of the tangent space to the fiber under the differential $d \Gamma$ at the point $(((u, f),(u, f),(v, g)), h)$ is contained in a subspace of codimension $\geq 2 d-1$ in the tangent space to the fiber of the target space. This contradicts $\Gamma^{-1}(0) \cap \mathrm{pr}^{-1}(h)$ being transversely cut out.
q.e.d.

### 7.9. Transversality in the semi-admissible case.

Lemma 7.17. Suppose $L$ is admissible chord-semi-generic and $\Lambda=$ $\operatorname{Ham}_{0}(L, \delta, s)$, then the bundle maps

$$
\begin{aligned}
& (\Gamma, \operatorname{pr}): \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c}) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}\left(D_{m}, T^{0,1} D_{m} \otimes \mathbb{C}^{n}\right), \\
& (\Gamma, \operatorname{pr}): \widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c}) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}\left(D_{m}, T^{0,1} D_{m} \otimes \mathbb{C}^{n}\right)
\end{aligned}
$$

are transverse to the 0 -section.
Proof. We proceed as in the proof of Lemma 7.12. Let $u$ be an element in the annihilator. The argument of Lemma 7.12 still applies up to the point where we conclude $\alpha \mid I_{+}$equals 0 . In the present setup not all Hamiltonian vector fields are allowed (see Definition 7.4). However, the ones that are allowed can be used exactly as in the proof of Lemma 7.12 to conclude the last $(n-1)$ components of $u$ must vanish identically.

Since $D_{m}$ is conformally equivalent to the unit disk $\Delta_{m}$ with $m$ punctures on the boundary and since integrals as in (7.7) are conformally invariant, we have for any smooth compactly supported $v$ with appropriate boundary conditions

$$
\begin{equation*}
0=\int_{\Delta_{m}}\langle\bar{\partial} v, u\rangle d A=\int_{\Delta_{m}}\langle v, \partial u\rangle d A+\int_{\partial \Delta_{m}}\left\langle u, e^{-i \theta} v\right\rangle d \theta . \tag{7.14}
\end{equation*}
$$

As usual, the first term in (7.14) vanishes and we find that $u$ is orthogonal to $e^{i \theta} T_{w\left(e^{i \theta}\right)} \Pi_{\mathbb{C}}(L)$.

Now the boundary of the holomorphic disk must cross the region $X=B(0,2+\epsilon) \backslash B(0,2)$, and the inverse image of this region contains an $\operatorname{arc} A$ in the boundary. The intersection between the tangent plane of $T_{p} \Pi_{\mathbb{C}}(L), p \in X$ and the $z_{1}$-line equals 0 and the $z_{1}$-line is invariant under multiplication by $e^{i \theta}$. Hence, the orthogonal complement of $e^{i \theta} T_{w\left(e^{i \theta)}\right.} \Pi_{\mathbb{C}}(L)$ intersects the $z_{1}$-line trivially as well (for $\theta \in A$ ). We conclude that the first component of $u$ must vanish identically along $A$ and by anti-analytic continuation vanish identically. It follows that $u$ is identically zero.
q.e.d.

In analogy with Corollary 7.14, we get (with $c$ denoting the degenerate Reeb chord of $L$ )

Corollary 7.18. For a dense open set of $h \in \operatorname{Ham}_{0}(L, \delta, s), \Gamma^{-1}(0) \cap$ $\operatorname{pr}^{-1}(h) \subset \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c} ; A)$ and $\Gamma^{-1}(0) \cap \operatorname{pr}^{-1}(h) \subset \widetilde{\mathcal{W}_{2, \epsilon, \Lambda}^{\prime}}(\mathbf{c} ; A)$ are finite dimensional manifolds. If $\mathbf{c}=a b_{1} \ldots b_{m}$ with $a \neq c$, then the dimensions are

$$
\begin{aligned}
& \mu(A)+\nu(a)-\sum_{r=1}^{m}\left(\nu\left(b_{j}\right)+\delta\left(b_{j}, c\right)\right)+\max (0, m-2) \text { and } \\
& \mu(A)+\nu(a)-\sum_{r=1}^{m}\left(\nu\left(b_{j}\right)\right)+\max (0, m-2), \text { respectively. }
\end{aligned}
$$

If $\mathbf{c}=c b_{1} \ldots b_{k}$, then the dimensions are

$$
\begin{aligned}
& \mu(B)+\nu(c)-\sum_{r=1}^{m}\left(\nu\left(b_{j}\right)\right)+\max (0, m-2) \text { and } \\
& \mu(B)+\nu(c)+1-\sum_{r=1}^{m}\left(\nu\left(b_{j}\right)\right)+\max (0, m-2), \text { respectively. }
\end{aligned}
$$

The same argument as in the proof of Theorem 7.15 gives
Theorem 7.19. For a dense open set of $h \in \operatorname{Ham}_{0}(L, \delta, s), \Gamma^{-1}(0) \cap$ $\operatorname{pr}^{-1}(h) \cap \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c})$ and $\Gamma^{-1}(0) \cap \operatorname{pr}^{-1}(h) \cap \widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}(\mathbf{c})$ are finite dimensional manifolds of dimensions given by the dimension formula in Corollary 7.18 , provided this dimension is $\leq 1$ if $l(\mathbf{c}) \geq 3$ and $\leq 1+(3-l(\mathbf{c}))$ otherwise.

Remark 7.20. Note that the expected dimension of the set of disks with dimension count in $\tilde{\mathcal{W}}_{2, \epsilon}$ equal to 1 in $\mathcal{W}_{2, \epsilon}$ is equal to $-k$ (or $-k+(3-l(\mathbf{c}))$ if $l(\mathbf{c} \leq 2)$ ), where $k$ is the number of punctures mapping to the self-tangency Reeb chord. Therefore for a dense open set in the space of chord semi-generic Legendrian submanifolds, this space is empty. Since any disk with exponential decay at the self-tangency point has a neighborhood in $\mathcal{W}_{2, \epsilon}$, we see that generically such disks do not exist, provided their dimension count in $\tilde{\mathcal{W}}_{2, \epsilon}$ is as above.
7.10. Enhanced transversality. Let $L$ be a (semi-)admissible submanifold. If $q \in L$ and $\zeta_{0} \in \partial D_{m}$, then define

$$
\mathcal{W}_{2, \epsilon}\left(\mathbf{c}, \zeta_{0}, p\right)=\left\{(w, f) \in \mathcal{W}_{2, \epsilon}(\mathbf{c}):(w, f)\left(\zeta_{0}\right)=p\right\}
$$

and in the semi-admissible case also $\widetilde{\mathcal{W}}_{2, \epsilon}\left(\mathbf{c}, \zeta_{0}, p\right)$ in a similar way.
If $\mathrm{ev}_{\zeta_{0}}: \mathcal{W}_{2, \epsilon}(\mathbf{c}) \rightarrow L$ denotes the map $\mathrm{ev}_{\zeta_{0}}(w, f)=(w, f)\left(\zeta_{0}\right)$. Then $\mathrm{ev}_{\zeta_{0}}$ is smooth and transverse to $p$ (as is seen by using local coordinates on $\left.\mathcal{W}_{2, \epsilon}(\mathbf{c})\right)$. Moreover, $\operatorname{ev}_{\zeta_{0}}^{-1}(p)=\mathcal{W}_{2, \epsilon}(\mathbf{c}, p)$ and hence, $\mathcal{W}_{2, \epsilon}(\mathbf{c}, p)$ is a closed submanifold of $\mathcal{W}_{2, \epsilon}(\mathbf{c})$ of codimension $\operatorname{dim}(L)$. Note that the tangent space $T_{(w, f)} \mathcal{W}_{2, \epsilon}\left(\mathbf{c}, p, \zeta_{0}\right)$ is the closed subspace of elements $(v, \gamma)$ in the tangent space $T_{(w, f)} \mathcal{W}_{2, \epsilon}(\mathbf{c})$ which are such that $v: D_{m} \rightarrow$ $\mathbb{C}^{n}$ satisfies $v\left(\zeta_{0}\right)=0$.

We consider

$$
\mathcal{W}_{2, \epsilon}(\mathbf{c}, p)=\bigcup_{\zeta \in \partial D_{m}} \mathcal{W}_{2, \epsilon}(\mathbf{c}, \zeta, p)
$$

as a locally trivial bundle over $\partial D_{m}$. Local trivializations are given compositions with suitable diffeomorphisms which move the boundary point $\zeta$ a little.

We define perturbation spaces as the closed subspaces $\operatorname{Ham}^{p}(L, \delta, R)$ $\subset \operatorname{Ham}(L, \delta, R)$ and $\operatorname{Ham}_{0}^{p}(L, \delta) \subset \operatorname{Ham}_{0}(L, \delta)$ of functions $h$ such that $h(p)=0$ and $D h(p)=0$. Thus, $\tilde{\Phi}_{h}$ fixes $p$. (Note that if $p$ is the
projection of a Reeb chord, this is no additional restriction.) If $\Lambda$ denotes one of these perturbation spaces, we form the bundles

$$
\begin{aligned}
& \mathcal{W}_{2, \epsilon, \Lambda}(\mathbf{c}, p)=\bigcup_{L_{\lambda}, \lambda \in \Lambda} \mathcal{W}_{2, \epsilon}(\mathbf{c}, p), \\
& \widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}(\mathbf{c}, p)=\bigcup_{L_{\lambda}, \lambda \in \Lambda} \mathcal{W}_{2, \epsilon}(\mathbf{c}, p)
\end{aligned}
$$

with local coordinates as before.
As before, let ' denote exclusion of exceptional holomorphic maps.
Lemma 7.21. Assume that $p \in L$ has a neighborhood $U$ such that $\Pi_{\mathbb{C}}(U)$ is real analytic. Then the bundle maps

$$
\begin{align*}
& (\Gamma, \operatorname{pr}): \mathcal{W}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c}, p) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}\left(D_{m}, T^{0,1^{*}} D_{m} \otimes \mathbb{C}^{n}\right)  \tag{7.15}\\
& (\Gamma, \operatorname{pr}): \widetilde{\mathcal{W}}_{2, \epsilon, \Lambda}^{\prime}(\mathbf{c}, p) \rightarrow \mathcal{H}_{1, \epsilon, \Lambda}\left(D_{m}, T^{0,1^{*}} D_{m} \otimes \mathbb{C}^{n}\right)
\end{align*}
$$

are transverse to the 0-section.
Proof. The proof is the same as the proof of Lemma 7.12 in the admissible case and the same as that of Lemma 7.17 in the semi-admissible case, provided the arcs $\gamma$ and $\gamma^{\prime}$ used there do not contain the special point $p$. On the other hand, if one of these arcs does contain $p$, we may shorten it until it does not. (The key point is that the condition that the Hamiltonian vanishes at a point does not destroy the approximation properties of the elements in the perturbation space for smooth functions supported away from this point).
q.e.d.

Corollary 7.22. Let $n>1$. For $L$ in a Baire subset of the space of (semi-)admissible Legendrian n-submanifolds, no rigid holomorphic disk passes through the end point of any Reeb chord of L.

Note, when $n=1$, this corollary is not true.
Proof. The proof of Theorem 7.15 shows that for a Baire set, there are no exceptional holomorphic disks. The Sard-Smale theorem in combination with Lemma 7.21 implies that for a Baire subset of this Baire set, the dimension of the space of rigid holomorphic disks with some point mapping to the end point of a specific Reeb chord is given by the Fredholm index of the operator $d \Gamma$ corresponding to $\Gamma$ in (7.15). Since the source space of this operator is the sum of a copy of $\mathbb{R}$ (from the movement of $\zeta$ on the boundary) and a closed codimension $\operatorname{dim} L$ subspace of the source space of $d \Gamma$ in Lemma 7.12 which has minimal index for disks to appear generically, we see the index in the present case is too small. This implies that the subset is generically empty. Taking the intersection of these Baire subsets for the finite collection of Reeb chord endpoints of $L$, we get a Baire subset with the required properties. q.e.d.

Corollary 7.23. If $L$ is as in Corollary 7.22, then there are no rigid holomorphic disks with boundary on $L$ which are nowhere injective on the boundary.

Proof. Let $w: D_{m+1} \rightarrow \mathbb{C}^{n}$ represent a holomorphic disk with boundary on $L$. By Corollary 7.22 , we may assume that no point in the boundary of $\partial D_{m}$ maps to an intersection point of $\Pi_{\mathbb{C}}(L)$.

Assume that $w$ has no injective point on the boundary and let the punctures of $D_{m+1}$ map to the Reeb chords $\left(c_{0}, \ldots, c_{m}\right)$ where the positive puncture maps to $c_{0}$. Let $C$ be the holomorphic chain which is the closure of image $w\left(D_{m}\right)$ of $w$ with local multiplicity 1 everywhere. Then

$$
\begin{equation*}
\operatorname{Area}(C)<\operatorname{Area}(w) \tag{7.16}
\end{equation*}
$$

since close to the point in $C$ most distant from the origin in $\mathbb{C}^{n}, w$ has multiplicity at least two.

The corners of $C$ is a subset $S$ of $c_{0}^{*}, \ldots, c_{m}^{*}$ and by integrating $\sum_{j} y_{j} d x_{j}$ along the boundary $\partial C$ of $C$ which lies in the exact Lagrangian $\Pi_{\mathbb{C}}(L)$ we find

$$
\begin{equation*}
\operatorname{Area}(C)=\mathcal{Z}\left(c_{0}\right)-\sum_{c_{j}^{*} \in S, j>0} \mathcal{Z}\left(c_{j}\right) \tag{7.17}
\end{equation*}
$$

where the first term must be present (i.e. $C$ must have a corner at $c_{0}^{*}$ ) since, otherwise, the area of $C$ would be negative contradicting the fact that $C$ is holomorphic. On the other hand

$$
\begin{equation*}
\operatorname{Area}(w)=\mathcal{Z}\left(c_{0}\right)-\sum_{j>0} \mathcal{Z}\left(c_{j}\right) \tag{7.18}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\operatorname{Area}(C) \geq \operatorname{Area}(w) \tag{7.19}
\end{equation*}
$$

which contradicts (7.16). This contradiction finishes the proof. q.e.d.
7.11. Transversality in a split problem. In this section, we discuss transversality for disks, with one or two punctures, lying entirely in one complex coordinate plane. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an admissible Legendrian submanifold. Let $\Delta \subset \mathbb{R}^{2}$ denote the standard simplex. Let $\Delta_{1}\left(\Delta_{2}\right)$ be the subsets of $\mathbb{R}^{2}$ which is bounded by $\partial \Delta$, smoothened at two (one) of its corners. Let $\left(z_{1}, \ldots, z_{n}\right)$ be coordinates on $\mathbb{C}^{n}$. Let $\pi_{i}: \mathbb{C}^{n} \rightarrow \mathbb{C}$ denote projection to the $i$-th coordinate and let $\hat{\pi}_{i}: \mathbb{C}^{n} \rightarrow \mathbb{C}^{n-1}$ denote projection to the Hermitian complement of the $z_{i}$-line. Finally, if $\gamma(t)$, $t \in I \subset \mathbb{R}$ is a one parameter family of lines, then we let $\int_{\gamma} d \theta$ denote the (signed) angle $\gamma(t)$ turns as $t$ ranges over $I$.

Lemma 7.24. Let $(u, h) \in \mathcal{W}_{2}(a b ; A), \mu(A)+|a|-|b|=1$, be a holomorphic disk with boundary on $L$ such that $\hat{\pi}_{1} \circ u$ is constant and such that $\pi_{1} \circ u=f \circ g$, where $g: \Delta_{2} \rightarrow D_{2}$ is a diffeomorphism and
$f: \Delta_{2} \rightarrow \mathbb{C}$ is an immersion. Furthermore, if $t_{1}, t_{2}$ are coordinates along components of $\partial D_{2}$, assume that the paths $\Gamma(t)=d \Pi_{\mathbb{C}}\left(T_{(u, h)\left(t_{j}\right)} L\right)$ of Lagrangian subspaces are split: $\Gamma\left(t_{j}\right)=\gamma\left(t_{j}\right) \times \hat{V}_{j}$, where $\gamma(t) \subset \mathbb{C}$ is a (real line) and $\hat{V}_{j} \subset \mathbb{C}^{n-1}, j=1,2$, are transverse Lagrangian subspaces. Then, $d \Gamma_{(u, h)}$ is surjective. (In other words, $(u, h)$ is transversely cut out.)

Proof. The Fredholm index of $d \Gamma$ at $(u, h)$ equals 1 . If $v$ is the vector field on $D_{2}$ which generates the 1-parameter family of conformal automorphisms of $D_{2}$ (the vector field $\partial_{\tau}$ in coordinates $\tau+i t \in \mathbb{R} \times[0,1]$ on $D_{2}$ ), then $\xi=d u \cdot v$ lies in the kernel of $d \Gamma$ and $d \hat{\pi} \cdot \xi=0$.

Since the boundary conditions are split, we may consider them separately. It follows from Section 6.4 that the $\hat{\pi}_{1} d \Gamma$ with boundary conditions given by the two transverse Lagrangian subspaces $\hat{V}_{1}$ and $\hat{V}_{2}$ has index 0 , no kernel and no cokernel.

Let $\theta_{1}$ and $\theta_{2}$ be the interior angles at the corners of the immersion $f$. Since $f\left(\partial \Delta_{2}\right)$ bounds an immersed disk, we have

$$
\int_{\gamma_{1}} d \theta+\int_{\gamma_{2}} d \theta+\left(\pi-\theta_{1}\right)+\left(\pi-\theta_{2}\right)=2 \pi
$$

If $\eta_{1}=\pi_{1} \circ \eta$, where $\eta$ is in the kernel of $d \Gamma$ then, thinking of $D_{2}$ as $\mathbb{R} \times[0,1]$, we find that, asymptotically, for some integers $n_{1} \geq 0$ and $n_{2} \geq 0$

$$
\eta_{1}(\tau+i t)= \begin{cases}c_{1} e^{-\left(\theta_{1}+n_{1} \pi\right)(\tau+i t)}, & \text { for } \tau \rightarrow+\infty \\ c_{2} e^{\left(\theta_{2}+n_{2} \pi\right)(\tau+i t)}, & \text { for } \tau \rightarrow-\infty\end{cases}
$$

where $c_{1}$ and $c_{2}$ are real constants. Cutting $D_{2}$ off at $|\tau|=M$ for some sufficiently large $M$, we thus find a solution of the classical RiemannHilbert problem with Maslov-class

$$
\frac{1}{\pi}\left(\theta_{1}+\theta_{2}-\theta_{1}-\theta_{2}-\left(n_{1}+n_{2}\right) \pi\right)
$$

Since the classical Riemann-Hilbert problem has no solution if the Maslov class is negative and exactly one if it is 0 , we see that the solution $\xi=\xi_{1}$ produced above is unique up to multiplication with real constants.
q.e.d.

Lemma 7.25. Let $(u, h) \in \mathcal{W}_{2}(a ; A), \mu(A)+|a|=1$, be a holomorphic disk with boundary on $L$ such that $\hat{\pi}_{1} \circ u$ is constant and such that $\pi_{1} \circ u=f \circ g$, where $g: \Delta_{1} \rightarrow D_{2}$ is a diffeomorphism and $f: \Delta_{1} \rightarrow \mathbb{C}$ is an immersion. Furthermore, if $t$ is a coordinate along $\partial D_{1}$, assume that the path $\Gamma(t)=d \Pi_{\mathbb{C}}\left(T_{(u, h)\left(t_{j}\right)} L\right)$ of Lagrangian subspaces is split: $\Gamma(t)=$ $\gamma_{1}\left(t_{j}\right) \times \gamma_{2}(t) \times \cdots \times \gamma_{n}(t)$, where $\gamma_{j}(t) \subset \mathbb{C}$ is a (real line) such that

$$
\int_{\gamma_{j}} d \theta<0, \text { for } 2 \leq j \leq n
$$

Then $d \Gamma_{(u, h)}$ is surjective.

Proof. The proof is similar to the one just given. Using asymptotics and the classical Riemann-Hilbert problem, it follows that the kernel of $d \Gamma$ is spanned by two linearly independent solutions $\xi^{j}, j=1,2$, with $\hat{\pi}_{1} \xi^{j}=0$. q.e.d.

### 7.12. Auxiliary tangent-like spaces in the semi-admissible case.

Let $L$ be a chord semi-admissible Legendrian submanifold and assume that $L$ lies in the open subset of such manifolds where the modulispace of rigid holomorphic disks with corners at $\mathbf{c}$ is 0 -dimensional (and compact by Theorem 9.2). Now if $(w, f)$ is a holomorphic disk with boundary on $L$, then by Lemma 7.17, we know that the operator

$$
\begin{equation*}
d \Gamma: T_{(w, f)} \widetilde{\mathcal{W}}_{2, \epsilon}(\mathbf{c}) \rightarrow \mathcal{H}_{1, \epsilon}\left(D_{m}, T^{*} D_{m} \otimes \mathbb{C}^{n}\right) \tag{7.20}
\end{equation*}
$$

is surjective.
For any ( $w, f$ ) with $m+1$ punctures which maps the punctures $p_{1}, \ldots p_{k}$ to the self tangency Reeb chord of $L_{h}$ let $\hat{\epsilon} \in[0, \infty)^{m+1-k} \times$ $(-\delta, 0)^{k}$, where $\delta>0$ is small compared to the complex angle of the self tangency Reeb chord and the components of $\hat{\epsilon}$ which are negative correspond to the punctures $p_{1}, \ldots, p_{k}$. Define the tangent-like space

$$
T_{(w, f, h)} \mathcal{W}_{2, \hat{\epsilon}}(\mathbf{c})
$$

as the linear space of elements $(v, \gamma)$ where $\gamma \in T_{\kappa} \mathcal{C}_{m+1}$ and where $v \in \mathcal{H}_{2, \hat{\epsilon}}\left(D_{m+1}, \mathbb{C}^{n}\right)$ satisfies

$$
\begin{aligned}
& v(\zeta) \in \Pi_{\mathbb{C}}\left(T_{(w, f)(\zeta)} L\right) \text { for all } \zeta \in \partial D_{m} \\
& \quad \int_{\partial D_{m}}\langle\bar{\partial} v, u\rangle d s=0 \text { for all } u \in \mathbb{C}_{0}^{\infty}\left(\partial D_{m}, \mathbb{C}^{n}\right)
\end{aligned}
$$

and consider the linear operator

$$
\begin{equation*}
d \hat{\Gamma}(v, \gamma)=\bar{\partial}_{\kappa} v+i \circ d w \circ \gamma . \tag{7.21}
\end{equation*}
$$

The index of this Fredholm operator equals that of the operator in (7.20) and moreover, by asymptotics of solutions to these equations (close to the self-tangency Reeb chord, we can use the same change of coordinates in the first coordinate as in the non-linear case, see Section 4.6 to determine the behavior of solutions), we find that the kernels are canonically isomorphic. Thus, since the operator in (7.20) is surjective so is the operator in (7.21).

## 8. Gluing theorems

In this section, we prove the gluing theorems used in Sections 2.3 and 2.5. In Section 8.1, we state the theorems. Our general method of gluing curves is the standard one in symplectic geometry. However, some of our specific gluings require a significant amount of analysis. We first "preglue" the pieces of the broken curves together. For the stationary case, this is done in Section 8.5 and for the self-tangency case in Sections
8.10 and 8.15. We then apply Picard's Theorem, stated in Section 8.2. Picard's Theorem requires a sequence of uniformly bounded right inverses of the linearized $\bar{\partial}$ map. We prove the bound for the stationary case in Section 8.7 and for the self-tangency case in Sections 8.13 and 8.19. Picard's Theorem also requires a bound on the non-linear part of the expansion of $\bar{\partial}$, which we discuss in Section 8.20. To handle disks with too few boundary punctures, we show in Sections 8.6 through 8.6.2, how by marking boundary points the disks can be thought of as sitting inside a moduli space of disks with many punctures.

Recall the following notation. Bold-face letters will denote ordered collections of Reeb chords. If $\mathbf{c}$ denotes a non-empty ordered collection $\left(c_{1}, \ldots, c_{m}\right)$ of Reeb chords, then we say that the length of $\mathbf{c}$ is $m$. We say that the length of the empty ordered collection is 0 . Let $\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}$ be an ordered collection of ordered collections of Reeb chords. Let the length of $\mathbf{c}^{j}$ be $l(j)$ and let $\mathbf{a}=\left(a_{1}, \ldots, a_{k}\right)$ be an ordered collection of Reeb chords of length $k>0$. Let $S=\left\{s_{1}, \ldots, s_{r}\right\}$ be $r$ distinct integers in $\{1, \ldots, k\}$. Define the ordered collection $\mathbf{a}_{S}\left(\mathbf{c}^{1}, \ldots, \mathbf{c}^{r}\right)$ of Reeb chords of length $k-r+\sum_{j=1}^{r} l(j)$ as follows. For each index $s_{j} \in S$, remove $a_{s_{j}}$ from the ordered collection a and insert at its place the ordered collection $\mathbf{c}^{j}$.

Recall that if $a$ is a Reeb chord and $\mathbf{b}$ is a collection of Reeb chords of a Legendrian submanifold, then $\mathcal{M}_{A}(a ; \mathbf{b})$ denotes the moduli space of holomorphic disks with boundary on $L$, punctures mapping to $(a, \mathbf{b})$, and boundary in $L$ which after adding the chosen capping paths represents the homology class $A \in H_{1}(L)$. After Theorem 7.15, we know that if the length of $\mathbf{b}$ is at least 2 , then $\mathcal{M}_{A}(a ; \mathbf{b})$ is identified with the inverse image of the regular value 0 of the $\bar{\partial}$-map $\Gamma$ in Section 5.7. If the length of $\mathbf{b}$ is 0 or 1 , then $\mathcal{M}_{A}(a ; \mathbf{b})$ is identified with the quotient of $\Gamma^{-1}(0)$ under the group of conformal reparameterizations of the source of the holomorphic disk.

Similarly, if $L_{\lambda}, \lambda \in \Lambda$ is a 1-parameter family of chord generic Legendrian submanifolds, we write $\mathcal{M}_{A}^{\Lambda}(a ; \mathbf{b})$ for the parameterized moduli space of rigid holomorphic disks with boundary in $L_{\lambda}$, and punctures at $(a(\lambda), \mathbf{b}(\lambda)), \lambda \in \Lambda$. We also write $\mathcal{M}_{A}^{\lambda}(a, \mathbf{b})$ to denote the moduli space for a fixed $L_{\lambda}, \lambda \in \Lambda$.

Finally if $K \subset \mathbb{C}^{n}$ and $\delta>0$, then $B(K, \delta)$ denotes the subset of all points in $\mathbb{C}^{n}$ of distance less than $\delta$ from $K$.
8.1. The Gluing Theorems. In this section, we state the various gluing theorems.
8.1.1. Stationary gluing. Let $L$ be an admissible Legendrian submanifold. Recall that a holomorphic disk with boundary on $L$ is defined as a pair of functions $(u, f)$, where $u: D_{m} \rightarrow \mathbb{C}^{n}$ and $f: \partial D_{m} \rightarrow \mathbb{R}$. Below we will often drop the function $f$ from the notation and speak of the holomorphic disk $u$. Let $\mathcal{M}_{A}(a ; \mathbf{b})$ and $\mathcal{M}_{C}(c ; \mathbf{d})$ be moduli spaces
of rigid holomorphic disks, where $\mathbf{b}$ has length $m, 1 \leq j \leq m$, and $\mathbf{d}$ has length $l$.

Theorem 8.1. Assume that the $j$-th Reeb chord in $\mathbf{b}$ equals $c$. Then there exists $\delta>0, \rho_{0}>0$ and an embedding

$$
\begin{aligned}
\mathcal{M}_{A}(a ; \mathbf{b}) \times \mathcal{M}_{C}(c ; \mathbf{d}) \times\left[\rho_{0}, \infty\right) & \rightarrow \mathcal{M}_{A+C}\left(a ; \mathbf{b}_{\{j\}}(\mathbf{d})\right) ; \\
(u, v, \rho) & \mapsto u \sharp \rho v,
\end{aligned}
$$

such that if $u \in \mathcal{M}_{A}(a ; \mathbf{b})$ and $v \in \mathcal{M}_{C}(c ; \mathbf{d})$ and the image of $w \in$ $\mathcal{M}_{A+C}\left(a ; \mathbf{b}_{j}(\mathbf{d})\right)$ lies inside $B\left(u\left(D_{m+1}\right) \cup v\left(D_{l+1}\right) ; \delta\right)$, then $w=u \sharp \rho v$ for some $\rho \in\left[\rho_{0}, \infty\right)$.

Proof. The theorem follows from Lemmas 8.5, 8.9, and 8.16 and Proposition 8.4. q.e.d.
8.1.2. Self tangency shortening and self tangency gluing. Let $L_{\lambda}, \lambda \in(-1,1)=\Lambda$ be an admissible 1-parameter family of Legendrian submanifolds such that $L_{0}$ is semi-admissible with self-tangency Reeb chord $a$. For simplicity (see Section 3), we assume that all Reeb chords outside a neighborhood of $a$ remain fixed under $\Lambda$. We take $\Lambda$ so that if $\lambda>0$, then $L_{-\lambda}$ has two new-born Reeb chords $a^{+}$and $a^{-}$, where $\mathcal{Z}\left(a^{+}\right)>\mathcal{Z}\left(a^{-}\right)$. Assume that all moduli spaces of rigid holomorphic disks with boundary on $L_{\lambda}$ are transversely cut out for all fixed $\lambda \in \Lambda$, that for all $\lambda \in \Lambda$, there are no disks with negative formal dimension, and that all rigid disks with a puncture at $a$ satisfy the non-decay condition of Lemma 4.6 (see Remark 7.20).

Theorem 8.2. Let $\Lambda^{-}=(-1,0)$. Let $\mathcal{M}_{A}^{0}(a, \mathbf{b})$ be a moduli space of rigid holomorphic disks where the length of $\mathbf{b}$ is $l$. Then there exist $\rho_{0}>0, \delta>0$ and a local homeomorphism

$$
\begin{aligned}
\mathcal{M}_{A}^{0}(a ; \mathbf{b}) \times\left[\rho_{0}, \infty\right) & \rightarrow \mathcal{M}_{A}^{\Lambda^{-}}\left(a^{+} ; \mathbf{b}\right) ; \\
(u, \rho) & \mapsto \sharp \rho,
\end{aligned}
$$

such that if $u \in \mathcal{M}_{A}^{0}(a ; \mathbf{b})$ and the image of $w \in \mathcal{M}_{A}^{\Lambda^{-}}\left(a^{+} ; \mathbf{b}\right)$ lies inside $B\left(u\left(D_{l+1}\right) ; \delta\right)$, then $w=\not \sharp_{\rho} u$ for some $\rho \in\left[\rho_{0}, \infty\right)$.

Let $\mathcal{M}_{C}^{0}(c, \mathbf{d})$ be a moduli space of rigid holomorphic disks where the length of $\mathbf{d}$ is $m$. Let $S \subset\{1, \ldots, m\}$ be the subset of positions of $\mathbf{d}$ where the Reeb chord a appears (to avoid trivialities, assume $S \neq \emptyset$ ). Then there exists $\rho_{0}>0$ and $\delta>0$ and a local homeomorphism

$$
\begin{aligned}
\mathcal{M}_{C}^{0}(c, \mathbf{d}) \times\left[\rho_{0}, \infty\right) & \rightarrow \mathcal{M}_{C}^{\Lambda^{-}}\left(c, \mathbf{d}_{S}\left(a^{-}\right)\right) ; \\
(u, \rho) & \mapsto \sharp_{\rho} u,
\end{aligned}
$$

such that if $u \in \mathcal{M}_{C}^{0}(c ; \mathbf{d})$ and the image of $w \in \mathcal{M}_{C}^{\Lambda^{-}}\left(c ; \mathbf{d}_{S}\left(a^{-}\right)\right)$lies inside $B\left(u\left(D_{m+1}\right) ; \delta\right)$, then $w=\not \sharp_{\rho} u$ for some $\rho \in\left[\rho_{0}, \infty\right)$.

Proof. Consider the first case, the second follows by a similar argument. Applying Proposition 8.4 and Lemmas 8.10, 8.11 and 8.17, we find a homeomorphism $\mathcal{M}_{A}^{0}(a ; \mathbf{b}) \rightarrow \mathcal{M}_{A}^{\lambda_{-}}\left(a^{+}, \mathbf{b}\right)$ for $\lambda^{-}<0$ small enough. The proof of Corollary 7.14 implies that $\mathcal{M}_{\Lambda^{-}}\left(a^{+}, \mathbf{b}\right)$ is a $1-$ dimensional manifold homeomorphic to $\mathcal{M}_{\lambda_{-}}(a ; \mathbf{b}) \times \Lambda_{-}$, the theorem follows.
q.e.d.

Theorem 8.3. Let $\Lambda^{+}=(0,1)$ and let $\mathcal{M}_{A_{1}}^{0}\left(a ; \mathbf{b}^{1}\right), \ldots, \mathcal{M}_{A_{r}}^{0}\left(a ; \mathbf{b}^{r}\right)$ and $\mathcal{M}_{C}^{0}(c ; \mathbf{d})$ be a moduli spaces of rigid holomorphic disks where the length of $\mathbf{b}^{j}$ is $l(j)$, and the length of $\mathbf{d}$ is $m$. Let $S \subset\{1, \ldots, m\}$ be the subset of positions of $\mathbf{d}$ where the Reeb chord a appears and assume that $S$ contains $r$ elements. Then there exists $\delta>0, \rho_{0}>0$ and an embedding

$$
\begin{aligned}
\mathcal{M}_{C}^{0}(c ; \mathbf{d}) \times \Pi_{j=1}^{r} \mathcal{M}_{A_{j}}^{0}\left(a ; \mathbf{b}^{j}\right) \times\left[\rho_{0}, \infty\right) & \rightarrow \mathcal{M}_{C+\sum_{j} A_{j}}^{\Lambda+}\left(c ; \mathbf{d}_{S}\left(\mathbf{b}^{1}, \ldots, \mathbf{b}^{r}\right)\right) ; \\
\left(v, u_{1}, \ldots, u_{r}, \rho\right) & \mapsto v \not \sharp_{\rho}\left(u_{1}, \ldots, u_{r}\right),
\end{aligned}
$$

such that if $v \in \mathcal{M}_{C}^{0}(c ; \mathbf{d})$ and $u_{j} \in \mathcal{M}_{A_{j}}^{0}\left(a ; \mathbf{b}^{j}\right), j=1, \ldots, r$ and the image of $w \in \mathcal{M}_{C+\sum_{j} A_{j}}^{\Lambda^{+}}\left(c ; \mathbf{d}_{S}\left(\mathbf{b}^{1}, \ldots, \mathbf{b}^{r}\right)\right)$ lies inside $B\left(v\left(D_{m+1}\right) \cup\right.$ $\left.\left.u_{1}\left(D_{l(1)+1}\right) \cup \cdots \cup u_{r}\left(D_{l(r)+1}\right)\right) ; \delta\right)$ then $w=v \not \sharp_{\rho}\left(u_{1}, \ldots, u_{r}\right)$ for some $\rho \in\left[\rho_{0}, \infty\right)$.

Proof. Apply Proposition 8.4 and Lemmas 8.13, 8.15, and 8.18 and reason as above.
q.e.d.
8.2. Floer's Picard lemma. The proofs of the theorems stated in the preceding subsections are all based on the following.

Proposition 8.4. Let $f: B_{1} \rightarrow B_{2}$ be a smooth map between Banach spaces which satisfies

$$
f(v)=f(0)+d f(0) v+N(v),
$$

where $d f(0)$ is Fredholm and has a right inverse $G$ satisfying

$$
\|G N(u)-G N(v)\| \leq C(\|u\|+\|v\|)\|u-v\|
$$

for some constant $C$. Let $B(0, \epsilon)$ denote the $\epsilon$-ball centered at $0 \in B_{1}$ and assume that

$$
\|G f(0)\| \leq \frac{1}{8 C}
$$

Then for $\epsilon<\frac{1}{4 C}$, the zero-set of $f^{-1}(0) \cap B(0, \epsilon)$ is a smooth submanifold of dimension $\operatorname{dim}(\operatorname{Ker}(d f(0)))$ diffeomorphic to the $\epsilon$-ball in $\operatorname{Ker}(d f(0))$.

Proof. See [12].
q.e.d.

In our applications of Proposition 8.4, the map $f$ will be the $\bar{\partial}$-map, see Section 5.7.
8.3. Notation and cut-off functions. To simplify notation, we deviate slightly from our standard notation for holomorphic disks. We use the convention that the neighborhood $E_{p_{0}}$ of the positive puncture $p_{0}$ in the source $D_{m}$ of a holomorphic disk $(u, f)$ will be parameterized by $[1, \infty) \times[0,1]$ and that neighborhoods of negative punctures $E_{p_{j}}, j \geq 1$ are parameterized by $(-\infty,-1] \times[0,1]$.

In the constructions and proofs below, we will use certain cut-off functions repeatedly. Here we explain how to construct them. Let $K>0, a<b$, and let $\phi:[a, b+K+1] \rightarrow[0,1]$ be a smooth function which equals 1 on $[a, b]$ and equals 0 in $[b+K, b+K+1]$. It is easy to see there exists such functions with $\left|D^{k} \phi\right|=\mathcal{O}\left(K^{-k}\right)$ for $k=1,2$. Let $\epsilon>0$ be small. Let $\psi:[0,1] \rightarrow \mathbb{R}$ be a smooth function such that $\psi(0)=\psi(1)=0, \psi^{\prime}(0)=\psi^{\prime}(1)=1$, with $|\psi| \leq \epsilon$. We will use cut-off functions $\alpha:[a, b+K+1] \times[0,1] \rightarrow \mathbb{C}$ of the form

$$
\alpha(\tau+i t)=\phi(\tau)+i \psi(t) \phi^{\prime}(\tau) .
$$

Note that $\alpha \mid \partial([a, b+K+1] \times[0,1])$ is real-valued and $\partial \alpha=0$ on $\partial([a, b+K+1] \times[0,1])$. Also, $\left|D^{k} \alpha\right|=\mathcal{O}\left(K^{-1}\right)$ for $k=1,2$.
8.4. A gluing operation. Let $L$ be a chord generic Legendrian submanifold. Let $(u, f) \in \mathcal{W}_{2, \epsilon}(a, \mathbf{b})$ where $\mathbf{b}=\left(b_{1}, \ldots, b_{m}\right)$ and let $\left(v_{j}, h_{j}\right)$ $\in \mathcal{W}_{2, \epsilon}\left(b_{j}, \mathbf{c}^{j}\right), j \in S \subset\{1, \ldots, m\}$. Denote the punctures on $D_{m+1}$ by $p_{j}, j=0, \ldots, m$ and the positive puncture on $D_{l(j)+1}$ by $q_{j}$.

Let $g^{\sigma}, \sigma \in[0,1]$ be a 1-parameter family of metrics as in Section 5.3. Then for $M>0$ large enough, there exists unique functions

$$
\begin{aligned}
& \xi: E_{p_{j}}[-M] \rightarrow T_{b_{j}^{*}} \mathbb{C}^{n} \\
& \eta_{j}: E_{q_{j}}[M] \rightarrow T_{b_{j}^{*}} \mathbb{C}^{n},
\end{aligned}
$$

such that

$$
\begin{aligned}
\exp _{b_{j}^{*}}^{t}(\xi(\tau+i t)) & =u(\tau+i t) \\
\exp _{b_{j}^{*}}^{t}\left(\eta_{j}(\tau+i t)\right) & =v_{j}(\tau+i t)
\end{aligned}
$$

where $\exp ^{\sigma}$ denotes the exponential map of the metric $g^{\sigma}$. Note that by our special choice of metrics the functions, $\xi$ and $\eta$ are tangent to $\Pi_{\mathbb{C}}(L)$ and holomorphic on the boundary.

For large $\rho>0$, let $D_{r}^{S}(\rho), r=1+m+\sum_{j \in S}(l(j)-1)$ be the disk obtained by gluing to the end of

$$
D_{m+1} \backslash \bigcup_{j \in S} E_{p_{j}}[-\rho]
$$

corresponding to $p_{j}$ a copy of

$$
D_{l(j)+1}-E_{q_{j}}[\rho]
$$

by identifying $\rho \times[0,1] \subset E_{p_{j}}$ with $-\rho \times[0,1] \subset E_{q_{j}}$, for each $j \in S$. Note that the metrics (and the complex structures $\kappa_{1}$ and $\kappa_{2}(j)$ ) on $D_{m+1}$
and $D_{l(j)+1}$ glue together to a unique metric (and complex structure $\kappa_{\rho}$ ) on $D_{r}^{S}(\rho)$. We consider $D_{m+1} \backslash \bigcup_{j \in S} E_{p_{j}}[-\rho]$ and $D_{l(j)+1} \backslash E_{q_{j}}[\rho]$ as subsets of $D_{r}^{S}(\rho)$.

For $j \in S$, let $\Omega_{j} \subset D_{r}^{S}(\rho)$ denote the subset

$$
E_{q_{j}}[\rho-2, \rho] \cup E_{p_{j}}[-\rho,-\rho+2] \approx[-2,2] \times[0,1]
$$

of $D_{r}^{S}(\rho)$. Let $z=\tau+i t$ be a complex coordinate on $\Omega_{j}$ and let $\alpha^{ \pm}: \Omega_{j} \rightarrow$ $\mathbb{C}$ be cut-off functions which are real valued and holomorphic on the boundary and with $\alpha^{+}=1$ on $[-2,-1] \times[0,1], \alpha^{+}=0$ on $[0,2] \times[0,1]$, $\alpha^{-}=1$ on $[1,2] \times[0,1]$, and $\alpha^{-}=0$ on $[-2,0] \times[0,1]$. Define the function $\Sigma_{\rho}^{S}\left(u, v_{1}, \ldots, v_{r}\right): D_{r} \rightarrow \mathbb{C}^{n}$ as

$$
\begin{aligned}
& \Sigma_{\rho}^{S}\left(u, v_{1}, \ldots, v_{r}\right)(\zeta) \\
& \quad= \begin{cases}v_{j}(\zeta), \quad \zeta \in D_{l(j)+1} \backslash E_{q_{j}}[\rho-2], \\
u(\zeta), \quad \zeta \in D_{m+1} \backslash \bigcup_{j \in S} E_{p_{j}}[-\rho+2], \\
\exp _{b_{j}^{*}}^{t}\left(\alpha^{-}(z) \xi_{j}(z)+\alpha^{+}(z) \eta_{j}(z)\right), \quad z=\tau+i t \in \Omega_{j} .\end{cases}
\end{aligned}
$$

8.5. Stationary pregluing. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be an admissible Legendrian submanifold. Let $u: D_{m+1} \rightarrow \mathbb{C}^{n}$ be a holomorphic disk with its $j$-th negative puncture $p$ mapping to $c,(u, f) \in \mathcal{W}_{2}(a, \mathbf{b})$, and let $v: D_{l+1} \rightarrow \mathbb{C}^{n}$ be a holomorphic disk with the positive puncture $q$ mapping to $c,(v, h) \in \mathcal{W}_{2}(c, \mathbf{d})$. Define

$$
\begin{equation*}
w_{\rho}=\Sigma_{\rho}^{\{j\}}(u, v) . \tag{8.1}
\end{equation*}
$$

Lemma 8.5. The function $w_{\rho}$ satisfies $w_{\rho} \in \mathcal{W}_{2}\left(a, \mathbf{b}_{\{j\}}(\mathbf{d})\right)$ and

$$
\begin{equation*}
\left\|\bar{\partial} w_{\rho}\right\|_{1}=\mathcal{O}\left(e^{-\theta \rho}\right), \tag{8.2}
\end{equation*}
$$

where $\theta$ is the smallest complex angle at the Reeb chord c. In particular, $\left\|\bar{\partial} w_{\rho}\right\|_{1} \rightarrow 0$ as $\rho \rightarrow \infty$.

Proof. The first statement is trivial. Outside $\Omega_{j}, w_{\rho}$ agrees with $u$ or $v$ which are holomorphic. Thus it is sufficient to consider the restriction of $w_{\rho}$ to $\Omega_{j}$. To derive the necessary estimates, we Taylor expand the exponential map at $c^{*}$. To simplify notation, we let $c^{*}=0 \in \mathbb{C}^{n}$ and let $\xi \in \mathbb{R}^{2 n}$ be coordinates in $T_{0} \mathbb{C}^{n}$ and $x \in \mathbb{R}^{2 n}$ coordinates around $0 \in \mathbb{C}^{n}$. Then

$$
\begin{equation*}
\exp _{0}^{t}(\xi)=\xi-\Gamma_{i j}^{k}(t) \xi^{i} \xi^{k} \partial_{k}+\mathcal{O}\left(|\xi|^{3}\right) \tag{8.3}
\end{equation*}
$$

This implies the inverse of the exponential map has Taylor expansion

$$
\begin{equation*}
\xi=x+\Gamma_{i j}^{k}(t) x^{i} x^{k} \partial_{k}+\mathcal{O}\left(x^{3}\right) \tag{8.4}
\end{equation*}
$$

From (8.3) and (8.4), we get

$$
\begin{equation*}
\exp _{0}^{t}\left(\alpha^{+} \xi_{j}\right)=\alpha^{+} u+\left(\left(\alpha^{+}\right)^{2}-\alpha^{+}\right) \Gamma_{i j}^{k}(t) u^{i} u^{j} \partial_{k}+\mathcal{O}\left(|u|^{3}\right) \tag{8.5}
\end{equation*}
$$

and a similar expression for $\alpha^{-} \eta_{j}$ in terms of $v_{j}$. Lemma 4.6 implies that $u$ and $D u$ are $\mathcal{O}\left(e^{-\theta \rho}\right)$ in $E_{p_{j}}[\rho]$, which together with (8.5) implies (8.2).

> q.e.d.
8.6. Marked points. In order to treat disks with less than three punctures (i.e., disks with conformal reparameterizations) in the same way as disks with more than three punctures, we introduce special points which we call marked points on the boundary. When disks with few punctures and marked points are glued to a disk with many punctures there arises a disk with many punctures and marked points and we must study also that situation.

Remark 8.6. Below we will often write simply $\mathcal{W}_{2, \epsilon}$ to denote spaces like $\mathcal{W}_{2, \epsilon}(\mathbf{c})$, dropping the Reeb chords from the notation.

Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a (semi-)admissible Legendrian submanifold and let $u: D_{m} \rightarrow \mathbb{C}^{n}$ represent $(u, f) \in \mathcal{W}_{2, \epsilon}(\kappa)$ where $\kappa$ is a fixed conformal structure on $D_{m}$. Let $U_{r} \subset \Pi_{\mathbb{C}}(L), r=1, \ldots, k$ be disjoint open subsets where $\Pi_{\mathbb{C}}(L)$ is real analytic and let $q_{r} \in \partial D_{m}$ be points such that $u\left(q_{r}\right) \in U_{r}$ and $d u\left(q_{r}\right) \neq 0$. After possibly shrinking $U_{r}$, we may biholomorphically identify ( $\left.\mathbb{C}^{n}, U_{r}, u\left(q_{r}\right)\right)$ with ( $\left.\mathbb{C}^{n}, V \subset \mathbb{R}^{n}, 0\right)$. Let $\left(x_{1}+i y_{1}, \ldots, x_{n}+i y_{n}\right)$ be coordinates on $\mathbb{C}^{n}$ and assume these coordinates are chosen so $d u\left(q_{r}\right) \cdot v_{0}=\partial_{1}$, where $v_{0} \in T_{q_{r}} D_{m}$ is a unit vector tangent to the boundary. Let $H_{r} \subset \mathbb{R}^{n}$ denote an open neighborhood of 0 in the submanifold $\left\{x_{1}=0\right\}$.

Let $S$ denote the cyclically ordered set of points $S=\left\{p_{1}, \ldots, p_{m}, q_{1}\right.$, $\left.\ldots, q_{k}\right\}$ where $p_{i} \in \partial D_{m}$ are the punctures of $D_{m}$. Fix three points $s_{1}, s_{2}, s_{3} \in\left\{p_{1}, p_{2}, p_{3}, q_{1}, \ldots, q_{k}\right\}$, then the positions of the other points in $S$ parameterizes the conformal structures on $\Delta_{m+k}$. As in Section 5.6, we pick vector fields $\tilde{v}_{j}, j=1, \ldots, m+k-3$ supported around the nonfixed points in $S$. Given a conformal structure on $\Delta_{m+k}$, we endow it with the metric which makes a neighborhood of each puncture $p_{j}$ look like the strip and denote disks with such metrics $\tilde{D}_{m, k}$.

If $(u, f), u: \tilde{D}_{m, k} \rightarrow \mathbb{C}^{n}$ and $f: \partial \tilde{D}_{m, k} \rightarrow \mathbb{R}$ are maps and $\tilde{\kappa}$ is a conformal structure on $\tilde{D}_{m+k}$ then forgetting the marked points $q_{1}, \ldots, q_{k}$, we may view the maps as defined on $D_{m}$ and the conformal structure $\tilde{\kappa}$ gives a conformal structure $\kappa$ on $D_{m}$. Note though that the standard metrics on $D_{m}$ corresponding to $\kappa$ may be different from the metric corresponding to $\tilde{\kappa}$ (this happens when one of the punctures $q_{j}$ is very close to one of the punctures $p_{r}$ ). However, the metrics differ only on a compact set and thus using this forgetful map, we define for a fixed conformal structure $\tilde{\kappa}$ on $\tilde{D}_{m, k}$ the space

$$
\mathcal{W}_{2, \epsilon}^{S}(\tilde{\kappa}) \subset \mathcal{W}_{2, \epsilon}(\kappa)
$$

as the subset of elements represented by maps $w: D_{m} \rightarrow \mathbb{C}^{n}$ such that $w\left(q_{r}\right) \in H_{r}$ for $r=1, \ldots, k$. Using local coordinates on $\mathcal{W}_{2, \epsilon}(\kappa)$ around
$(u, f)$ we see that for some ball $B$ around $(u, f), \mathcal{W}_{2, \epsilon}^{S}(\tilde{\kappa}) \cap B$ is a codimension $k$ submanifold with tangent space at $(w, g)$ the closed subset of $T_{(w, g)} \mathcal{W}_{2, \epsilon}$ consisting of $v: D_{m} \rightarrow \mathbb{C}^{n}$ with $\left\langle v\left(q_{r}\right), \partial_{1}\right\rangle=0$. We call $\tilde{D}_{m, k}$ a disk with $m$ punctures and $k$ marked points.

The diffeomorphisms $\tilde{\phi}_{j}^{\sigma_{j}}, \sigma_{j} \in \mathbb{R}$ generated by $\tilde{v}_{j}$ gives local coordinates $\sigma=\left(\sigma_{1}, \ldots, \sigma_{m+k-3}\right) \in \mathbb{R}^{m+k-3}$ on the space of conformal structures on $\tilde{D}_{m, k}$ and the structure of a locally trivial bundle to the space

$$
\mathcal{W}_{2, \epsilon}^{S}=\bigcup_{\tilde{\kappa} \in \mathcal{C}_{m+k}} \mathcal{W}_{2, \epsilon}^{S}(\tilde{\kappa})
$$

The $\bar{\partial}$-map is defined in the natural way on this space and we denote it $\tilde{\Gamma}$.
8.6.1. Marked points on disks with few punctures. Let $L \subset$ $\mathbb{C}^{n} \times \mathbb{R}$ be a (semi-)admissible submanifold, let $m \leq 2$ and consider a holomorphic disk $(u, f)$ with boundary on $L$, represented by a map $u: D_{m} \rightarrow \mathbb{C}^{n}$. We shall put $3-m$ marked points on $D_{m}$.

Pick $U_{r} \subset \Pi_{\mathbb{C}}(L), 1 \leq r \leq 3-m$ as disjoint open subsets in which $\Pi_{\mathbb{C}}(L)$ is real analytic and let $q_{r} \in \partial D_{m}$ be points such that $u\left(q_{r}\right) \in U_{r}$ and $d u\left(q_{r}\right) \neq 0$. Such points exists by Lemma 7.7. As in Section 8.6, we then consider the $q_{r}$ as marked points and as there we use the notation $H_{r}$ for the submanifold into which $q_{r}$ is mapped.

Then the class in the moduli space of holomorphic disks of every holomorphic disk $(w, g)$ which is sufficiently close to $(u, f)$ in $\mathcal{W}_{2, \epsilon}$ has a unique representative $(\hat{w}, \hat{g}) \in \mathcal{W}_{2, \epsilon}^{S}$. Namely, any $\operatorname{such}(w, g)$ must intersect $H_{r}$ in a point $q_{r}^{\prime}$ close to $q_{r}, 1 \leq r \leq 3-m$. If $\psi$ denotes the unique conformal reparameterization of $D_{m}$ which takes $q_{r}$ to $q_{r}^{\prime}$, $1 \leq r \leq 3-m$ then $\hat{w}(\zeta)=w(\psi(\zeta))$. Moreover, if

$$
\begin{equation*}
d \Gamma_{(u, f)}: T_{(u, f)} \mathcal{W}_{2, \epsilon} \rightarrow \mathcal{H}_{1, \epsilon}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right) \tag{8.6}
\end{equation*}
$$

has index $k$ (note $k \geq 3-m$ since the space of conformal reparameterizations of $D_{m}$ is $(3-m)$-dimensional), then the restriction of $d \Gamma_{(u, f)}$ to $T_{(u, f)} \mathcal{W}_{2, \epsilon}^{S}$ has index $k-(3-m)$. In particular, if $d \Gamma_{(u, f)}$ is surjective, so is its restriction.

We conclude from the above that to study the moduli space of holomorphic disks in a neighborhood of a given holomorphic disk, we may (and will) use a neighborhood of that disk in $\mathcal{W}_{2, \epsilon}^{S}$ and $\tilde{\Gamma}$ rather than a neighborhood in the bigger space $\mathcal{W}_{2, \epsilon}$ and $\Gamma$.
8.6.2. Marked points on disks with many punctures. Let $L \subset$ $\mathbb{C}^{n} \times \mathbb{R}$ be as above, let $m \geq 3$ and consider a holomorphic disk $(u, f)$ with boundary on $L$, represented by a map $u: D_{m} \rightarrow \mathbb{C}^{n}$. We shall put $k$ marked points on $D_{m}$.

Pick $U_{r} \subset \Pi_{\mathbb{C}}(L), 1 \leq r \leq k$ as disjoint open subsets in which $\Pi_{\mathbb{C}}(L)$ is real analytic and let $q_{r} \in \partial D_{m}$ be points such that $u\left(q_{r}\right) \in U_{r}$ and
$d u\left(q_{r}\right) \neq 0$. As in Section 8.6, we then consider the $q_{r}$ as marked points and as there we use the notation $H_{r}$ for the submanifold into which $q_{r}$ is mapped.

Note that $(u, f)$ lies in $\mathcal{W}_{2, \epsilon}$ as well as in $\mathcal{W}_{2, \epsilon}^{S}$. We define a map

$$
\Omega: U \subset \mathcal{W}_{2, \epsilon}^{S} \rightarrow \mathcal{W}_{2, \epsilon} ; \quad \Omega\left(\left(w, g, \tilde{\phi}^{s}\right)\right)=\left(\hat{w}, \hat{g}, \tilde{\phi}^{t}\right)
$$

where $U$ is a neighborhood of $((u, f), \tilde{\kappa})$ as follows.
Consider the local coordinates $\omega \in \mathbb{R}^{m+k-3}$ on $\mathcal{C}_{m+k}$ around $\tilde{\kappa}$ and the product structure

$$
\mathbb{R}^{m+k-3}=\mathbb{R}^{m-3} \times \mathbb{R}^{j} \times \mathbb{R}^{k-j}
$$

where $\mathbb{R}^{m-3}$ is identified with the diffeomorphisms

$$
\phi^{\tau}=\tilde{\phi}_{p_{4}}^{\tau_{1}} \circ \cdots \circ \tilde{\phi}_{p_{m}}^{\tau_{m-3}}, \quad \tau=\left(\tau_{1}, \ldots, \tau_{m-3}\right) \in \mathbb{R}^{m-3},
$$

where $j$ is the number of elements in $\left\{s_{1}, s_{2}, s_{3}\right\} \backslash\left\{p_{1}, p_{2}, p_{3}\right\}$, and where $\mathbb{R}^{k-j}$ is identified with the diffeomorphisms

$$
\phi^{\sigma}=\tilde{\phi}_{\hat{s}_{1}}^{\sigma_{1}} \circ \cdots \circ \tilde{\phi}_{\hat{s}_{k-j}}^{\sigma_{k-j}}, \quad \sigma=\left(\sigma_{1}, \ldots, \sigma_{k-j}\right) \in \mathbb{R}^{k-j}
$$

where $\left\{\hat{s}_{1}, \ldots, \hat{s}_{k-j}\right\}=S \backslash\left(\left\{p_{4}, \ldots, p_{m}\right\} \cup\left\{s_{1}, s_{2}, s_{3}\right\}\right)$.
For $\tilde{\theta}$ near $\tilde{\kappa}$, let $\left\{s_{1}^{\prime}, \ldots, s_{m+k-3}^{\prime}\right\}$ denote the corresponding positions of punctures and marked points in $\partial \Delta$ and let $\psi: \Delta \rightarrow \Delta$ be the unique conformal reparameterization such that $\psi\left(p_{j}\right)=p_{j}^{\prime}$ for $j=1,2,3$ and note that we may view $\psi$ as a map from $D_{m}$ to $\tilde{D}_{m, k}$. Let $s_{l}^{\prime \prime}=\psi^{-1}\left(s_{l}^{\prime}\right)$ for $3 \leq l \leq k+m-3$ and $s_{l} \neq p_{i}, i=1,2,3$ let $(\tau, \sigma) \in \mathbb{R}^{m+k-3-j}$ be the unique element such that $\phi^{\tau} \circ \phi^{\sigma}\left(s_{l}\right)=s_{l}^{\prime \prime}$. Define

$$
\Omega(w, \tilde{\theta})=\left(w \circ \psi \circ \phi^{\sigma},\left(\phi^{\tau}\right)^{-1}\right),
$$

where $\left(\phi^{\tau}\right)^{-1}$ is interpreted as a conformal structure on $D_{m}$ in a neighborhood of $\kappa$ in local coordinates given by $\phi^{\tau}, \tau \in \mathbb{R}^{m-3}$ and where we drop the boundary function from the notation since it is uniquely determined by the $\mathbb{C}^{n}$-function component of $\Omega(w, \tilde{\theta})$ and $g$.

Lemma 8.7. The map $\Omega$ maps $U \cap \tilde{\Gamma}^{-1}(0)$ into $\Gamma^{-1}(0)$. Moreover, $\Omega$ is a $C^{1}$-diffeomorphism on a neighborhood of $(u, f)$.

Proof. Assume that $(w, g) \in \tilde{\Gamma}^{-1}(0)$. Then $w$ is holomorphic in the conformal structure $\tilde{\theta}_{\dot{\tilde{\theta}}}$. Since $\psi$ is a conformal equivalence and since the conformal structure $\tilde{\theta}$ is obtained from $\tilde{\kappa}$ by action of the inverses of $\phi^{\tau} \circ \phi^{\sigma}$ this implies

$$
0=d w \circ d \psi+i \circ(d w \circ d \psi) \circ\left(d \phi^{\sigma}\right) \circ\left(d \phi^{\tau}\right) \circ j_{\kappa} \circ\left(d \phi^{\tau}\right)^{-1} \circ\left(d \phi^{\sigma}\right)^{-1} .
$$

Thus
$0=\left(d w \circ d \psi \circ d \phi^{\sigma}+i \circ\left(d w \circ d \psi \circ d \phi^{\sigma}\right) \circ\left(d \phi^{\tau}\right) \circ j_{\kappa} \circ\left(d \phi^{\tau}\right)^{-1}\right) \circ\left(d \phi^{\sigma}\right)^{-1}$, and $w \circ \psi \circ \phi^{\sigma}$ is holomorphic in the conformal structure $d \phi^{\tau} j_{\kappa}\left(d \phi^{\tau}\right)^{-1}$ as required.

For the last statement we use the inverse function theorem. It is clear that the $\operatorname{map} \Omega$ is $C^{1}$ and that the differential of $\Omega$ at $(u, f)$ is a Fredholm operator. In fact, on the complement of all conformal variations on $\tilde{D}_{m, k}$ not supported around any of $p_{3}, \ldots, p_{m}$, it is just an inclusion into a subspace of codimension $k$, which consists of elements $v$ which vanish at $q_{1}, \ldots, q_{k}$. Since $d u\left(q_{r}\right) \neq 0$ for all $r$, it follows easily that the image of the remaining $k$ directions in $T_{(u, f)} \mathcal{W}_{2, \epsilon}^{S}$ spans the complement of this subspace. q.e.d.

It is a consequence of Lemma 8.7 that if $(u, f)$ is a holomorphic disk with boundary on $L$ and more than 3 punctures, then we may view a neighborhood of $(u, f)$ in the moduli space of such disks either as a submanifold in $\mathcal{W}_{2, \epsilon}^{S}$ or in $\mathcal{W}_{2, \epsilon}$ in a neighborhood of $(u, f)$.

Remark 8.8. Below we extend the use of the notion $\mathcal{W}_{2, \epsilon}$ to include also spaces $\mathcal{W}_{2, \epsilon}^{S}$, when this is convenient. The point being that after Sections 8.6.1 and 8.6.2, we may always assume the number of marked points and punctures is $\geq 3$, so that the moduli space of holomorphic disks (locally) may be viewed as a submanifold of $\mathcal{W}_{2, \epsilon}$.

### 8.7. Uniform invertibility of the differential in the stationary

case. Let

$$
\Gamma: \mathcal{W}_{2} \rightarrow \mathcal{H}_{1}[0]\left(D_{m}, T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)
$$

be the $\bar{\partial}$-map defined in Section 5.7 (see Remark 8.6 for notation). Let $u: D_{m+1} \rightarrow \mathbb{C}^{n}$ and $v: D_{l+1} \rightarrow \mathbb{C}^{n}$ be as in Section 8.1.1 and consider the differential $d \Gamma_{\rho}$ at $\left(w_{\rho}, \kappa_{\rho}\right)$, where $w_{\rho}$ is as in Lemma 8.5 and $\kappa_{\rho}$ is the natural metric (complex structure) on $D_{r}(\rho), r=m+l$. After Sections 8.6.1 and 8.6.2, we know that after adding $3-m$ or $3-l$ marked points on holomorphic disks with $\leq 2$ punctures, we may assume that $m \geq 2$ and $l \geq 2$ below.

Lemma 8.9. There exist constants $C$ and $\rho_{0}$ such that if $\rho>\rho_{0}$, then there are continuous right inverses

$$
G_{\rho}: \mathcal{H}_{1}[0]\left(T^{* 0,1} D_{r}(\rho) \otimes \mathbb{C}^{n}\right) \rightarrow T_{\left(w_{\rho}, \kappa_{\rho}\right)} \mathcal{W}_{2}
$$

of $d \Gamma_{\rho}$ with

$$
\left\|G_{\rho}(\xi)\right\| \leq C\|\xi\|_{1}
$$

Proof. The kernels

$$
\begin{aligned}
& \operatorname{ker}\left(d \Gamma_{\left(u, \kappa_{1}\right)}\right) \subset T_{u} \mathcal{W}_{2} \oplus T_{\kappa_{1}} \mathcal{C}_{m+1} \\
& \operatorname{ker}\left(d \Gamma_{\left(v, \kappa_{2}\right)}\right) \subset T_{v} \mathcal{W}_{2} \oplus T_{\kappa_{2}} \mathcal{C}_{l+1}
\end{aligned}
$$

are both 0-dimensional. As in Section 8.6, we view elements $\gamma_{1} \in$ $T_{\kappa_{1}} \mathcal{C}_{m+1}\left(\gamma_{2} \in T_{\kappa_{2}} \mathcal{C}_{l+1}\right)$ as linear combinations of sections of $\operatorname{End}\left(T D_{m+1}\right)$ (End $\left.\left(T D_{l+1}\right)\right)$ supported in compact annular regions close to all punctures and marked points, except at three. Since these annular regions
are disjoint from the regions affected by the gluing of $D_{m+1}$ and $D_{l+1}$, we get an embedding

$$
T_{\kappa_{1}} \mathcal{C}_{m+1} \oplus T_{\kappa_{2}} \mathcal{C}_{l+1} \rightarrow T_{\kappa_{\rho}} \mathcal{C}_{r} .
$$

In fact, using this embedding,

$$
T_{\kappa_{\rho}} \mathcal{C}_{r}=T_{\kappa_{1}} \mathcal{C}_{m+1} \oplus T_{\kappa_{2}} \mathcal{C}_{l+1} \oplus \mathbb{R},
$$

where the last summand can be taken to be generated by a section $\gamma_{0}$ of $\operatorname{End}\left(T D_{r}(\rho)\right)$ supported in an annular region around a puncture (marked point) in $D_{r}(\rho)$ where there was no conformal variation before the gluing. Then $\gamma_{0}$ spans a subspace of dimension 1 in $T_{\left(w_{\rho}, \kappa_{\rho}\right)} \mathcal{W}_{2}$. Let $a$ be a coordinate along this 1-dimensional subspace. We prove that for $(\xi, \gamma) \in W_{\rho}=\{a=0\}$, we have the estimate

$$
\begin{equation*}
\|(\xi, \gamma)\| \leq C\left\|d \Gamma_{\rho}(\xi, \gamma)\right\|_{1, \epsilon}, \tag{8.7}
\end{equation*}
$$

for all sufficiently large $\rho$. Since the Fredholm-index of $d \Gamma_{\rho}$ equals 1 , this shows $d \Gamma_{\rho}$ are surjective and with uniformly bounded inverses $G_{\rho}$ as claimed and thus finishes the proof.

Assume (8.7) is not true. Then there exists a sequence of elements $\left(\xi_{N}, \gamma_{N}\right) \in W_{\rho(N)}, \rho(N) \rightarrow \infty$ as $N \rightarrow \infty$ with

$$
\begin{align*}
& \left\|\left(\xi_{N}, \gamma_{N}\right)\right\|=1  \tag{8.8}\\
& \left\|\bar{\partial}_{\kappa_{\rho(N)}} \xi_{N}+i \circ \partial_{\kappa} w_{\rho(N)} \circ \gamma_{N}\right\|_{1} \rightarrow 0 . \tag{8.9}
\end{align*}
$$

As in Section 8.4, we glue a negative puncture at $p$ to a positive one at $q$. Note that on the strip

$$
\begin{equation*}
\Theta_{\rho}=\left(E_{p}[-1] \backslash E_{p}[-\rho]\right) \cup\left(E_{q}[1] \backslash E_{q}[\rho]\right) \approx[-\rho, \rho] \times[0,1] \subset D_{r}(\rho) \tag{8.10}
\end{equation*}
$$

the conformal structure $\kappa_{\rho}$ is the standard one and therefore $\bar{\partial}_{\kappa_{\rho}}$ is just the standard $\bar{\partial}$ operator. Also, since $\gamma_{N}$ does not have support in $\Theta_{\rho}$ the second term in (8.9) equals 0 when restricted to $\Theta_{\rho}$.

Let $\alpha_{\rho}: \Theta_{\rho} \rightarrow \mathbb{C}$ be cut-off functions which are real and holomorphic on the boundary, equal 1 on $[-2,2] \times[0,1]$, equal 0 outside $\left[-\frac{1}{2} \rho, \frac{1}{2} \rho\right] \times$ $[0,1]$, and satisfy $\left|D^{k} \alpha_{\rho}\right|=\mathcal{O}\left(\rho^{-1}\right), k=1,2$.

Then $\alpha_{\rho(N)} \xi_{N}$ is a sequence of functions on $\mathbb{R} \times[0,1]$ which satisfy boundary conditions converging to two transverse Lagrangian subspaces. Just as we prove in Lemma 6.9 that the (continuous) index is preserved under small perturbations, we conclude that the (upper semicontinuous) dimension of the kernel stay zero for large enough $\rho(N)$; thus, there exists a constant $C$ such that
$\left\|\xi_{N} \mid[-2,2] \times[0,1]\right\|_{2} \leq\left\|\alpha \xi_{N}\right\|_{2} \leq C\left(\left\|\alpha_{\rho(N)}\left(\bar{\partial} \xi_{N}\right)\right\|_{1}+\left\|\left(\bar{\partial} \alpha_{\rho(N)}\right) \xi_{N}\right\|_{1}\right)$.
As $N \rightarrow \infty$, both terms on the right-hand side in (8.11) approaches 0 . Hence,

$$
\begin{equation*}
\left\|\xi_{N} \mid[-2,2] \times[0,1]\right\|_{2} \rightarrow 0, \text { as } N \rightarrow \infty \tag{8.12}
\end{equation*}
$$

Pick cut-off functions $\beta_{N}^{+}$and $\beta_{N}^{-}$on $D_{r}(\rho)$ which are real valued and holomorphic on the boundary and have the following properties. On $D_{m+1} \backslash E_{p}[-\rho+1], \beta_{N}^{+}=1$ and on $D_{l+1} \backslash E_{q}[\rho], \beta_{N}^{+}=0$. On $D_{l+1} \backslash E_{q}[\rho-1], \beta_{N}^{-}=1$ and on $D_{m+1} \backslash E_{p}[-\rho], \beta_{N}^{-}=0$. Let $\left(\xi_{N}, \gamma_{N}\right)^{ \pm}=$ $\left(\beta_{N}^{ \pm} \xi_{N}, \beta_{N}^{ \pm} \gamma_{N}\right)$. Using the invertibility of $d \Gamma_{+}=d \Gamma_{\left(u, \kappa_{1}\right)}$ and $d \Gamma_{-}=$ $d \Gamma_{\left(v, \kappa_{2}\right)}$, we find a constant $C$ such that

$$
\begin{align*}
\left\|\left(\xi_{N}, \Gamma_{N}\right)^{ \pm}\right\| & \leq C\left\|d \Gamma_{ \pm}\left(\xi_{N}, \gamma\right)^{ \pm}\right\|_{1}  \tag{8.13}\\
& \leq C\left(\left\|\beta_{N}^{ \pm} d \Gamma_{\rho}\left(\xi_{N}, \gamma_{N}\right)\right\|_{1}+\left\|\left(\bar{\partial} \beta_{N}^{ \pm}\right) \xi_{N}\right\|_{1}\right)
\end{align*}
$$

The first term in the last line of (8.13) tends to 0 as $N \rightarrow \infty$ by (8.9), the second term tends to 0 by (8.12). Hence, the left-hand side of (8.13) tends to 0 as $N \rightarrow \infty$. Thus, (8.12) and (8.13) contradict (8.8) and we conclude (8.7) holds.

> q.e.d.

### 8.8. Self-tangencies, coordinates and genericity assumptions.

Let $z=x+i y=\left(z_{1}, \ldots, z_{n}\right)=\left(x_{1}+i y_{1}, \ldots, x_{n}+i y_{n}\right)$ be coordinates on $\mathbb{C}^{n}$. Let $L \subset \mathbb{C}^{n} \times \mathbb{R}$ be a semi-admissible Legendrian submanifold with self-tangency double point at 0 . We assume that the self-tangency point is standard, see Section 3.

Theorems 7.19 and 9.2 imply that the moduli-space of rigid holomorphic disks with boundary on $L$ is a 0 -dimensional compact manifold. Moreover, because of the enhanced transversality discussion in Section 7.10, we may assume that there exists $r_{0}>0$ such that for all $0<r<r_{0}$, if $u: D_{m} \rightarrow \mathbb{C}^{n}$ is a rigid holomorphic disk with boundary on $L$, then $\partial D_{m} \cap u^{-1}(B(0, r))$ is a disjoint union of subintervals of $\partial E_{p_{j}}[ \pm M]$, for some $M>0$ and some punctures $p_{j}$ on $\partial D_{m}$ mapping to 0 .

By Lemma 4.6, if $u: D_{m} \rightarrow \mathbb{C}^{n}$ is a rigid holomorphic disk with $q^{+}$a positive ( $q^{-}$a negative) puncture mapping to 0 , then there exists $c \in \mathbb{R}$ such that for $\zeta=\tau+i t \in E_{q^{ \pm}}[ \pm M]$

$$
u(\zeta)=\left(-2(\zeta+c)^{-1}, 0, \ldots, 0\right)+\mathcal{O}\left(e^{-\theta|\tau|}\right)
$$

for some $\theta>0$. For simplicity, we assume below that coordinates on $E_{q^{ \pm}}[M]$ are chosen in such a way that $c=0$ above.
8.9. Perturbations for self-tangency shortening. For $0<a<1$, with $a$ very close to 1 and $R>0$ with $R^{-1} \ll r_{0}$, let $b_{R}:[0, \infty) \rightarrow \mathbb{R}$ be a smooth non-increasing function with support in $\left[0, R^{-1}\right)$ and with the following properties

$$
\begin{align*}
b_{R}(r) & =\left(R+R^{a}\right)^{-2} \text { for } r \in\left[0,\left(R+\frac{1}{2} R^{a}\right)^{-1}\right]  \tag{8.14}\\
\left|D b_{R}(r)\right| & =\mathcal{O}\left(R^{-a}\right) \\
\left|D^{2} b_{R}(r)\right| & =\mathcal{O}\left(R^{2-2 a}\right) \\
\left|D^{3} b_{R}(r)\right| & =\mathcal{O}\left(R^{4-3 a}\right)
\end{align*}
$$

$$
\left|D^{4} b_{R}(r)\right|=\mathcal{O}\left(R^{6-4 a}\right)
$$

The existence of such a function is easily established using the fact that the length of the interval where $D b_{R}$ is supported equals

$$
R^{-1}-\left(R+\frac{1}{2} R^{a}\right)^{-1}=\frac{1}{2} R^{a-2}+\mathcal{O}\left(R^{2(a-2)}\right)
$$

Let

$$
\begin{equation*}
h_{R}(z)=-x_{1}(z) b_{R}(|z|) . \tag{8.15}
\end{equation*}
$$

Let $L^{1}$ and $L^{2}$ be the two branches of the local Lagrangian projection near the self-tangency, see Section 3 or Figure 2. For $s>0$, let $\Psi_{R}^{s}$ denote the time $s$ Hamiltonian flow of $h_{R}$ and let $L_{R}(s)$ denote the Legendrian submanifold which results when $\Psi_{R}^{s}$ is lifted to a contact flow on $\mathbb{C}^{n} \times \mathbb{R}$ (see Section 3) which is used to move $L^{2}$. Let $L_{R}^{2}(s)=\Psi_{R}^{s}\left(L^{2}\right)$. Let $g(R, s, \sigma)$ be a 3 -parameter family of metrics on $\mathbb{C}^{n}$ such that $L^{1}$ is totally geodesic for $g(R, s, 0), L_{R}^{2}(s)$ is totally geodesic for $g(R, s, 1)$ and such that $g(R, s, 0)$ and $g(R, s, 1)$ have properties as the metrics constructed in Section 5.3.

Note that $L_{R}^{2}(1) \cap L^{1}$ consists of exactly two points with coordinates $\left( \pm\left(R+R^{a}\right)^{-1}+\mathcal{O}\left(R^{-3}\right), 0, \ldots, 0\right)$.

We will use $\Psi_{R}^{s}$ to deform holomorphic disks below. It will be important for us to know they remain almost holomorphic in a rather strong sense, for which we need to derive some estimates on the flow $\Psi_{R}^{s}$ and its derivatives. Let $X_{R}$ denote the Hamiltonian vector field of $h_{R}$. Then if $D$ denotes derivative with respect to the variables in $\mathbb{C}^{n}$ and $\cdot$ denotes contraction of tensors

$$
\begin{equation*}
\frac{d}{d s} D \Psi_{R}^{s}=D X_{R} \cdot D \Psi_{R}^{s} ; \quad D \Psi_{\mu}^{0}=\mathrm{id} \tag{8.17}
\end{equation*}
$$

$$
\begin{align*}
\frac{d}{d s} D^{2} \Psi_{R}^{s}= & D^{2} X_{R} \cdot D \Psi_{R}^{s} \cdot D \Psi_{R}^{s}+D X_{R} \cdot D^{2} \Psi_{R}^{s} ; \quad D^{2} \Psi_{R}^{0}=0  \tag{8.18}\\
\frac{d}{d s} D^{3} \Psi_{R}^{s}= & D^{3} X_{R} \cdot D \Psi_{R}^{s} \cdot D \Psi_{R}^{s} \cdot D \Psi_{R}^{s}+2 D^{2} X_{R} \cdot D^{2} \Psi_{R}^{s} \cdot D \Psi_{R}^{s} \\
& +D^{2} X_{R} \cdot D \Psi_{R}^{s} \cdot D^{2} \Psi_{R}^{s}+D X_{R} \cdot D^{3} \Psi_{R}^{s} ; \quad D^{3} \Psi_{R}^{0}=0 . \tag{8.19}
\end{align*}
$$

Since $X_{R}=i \cdot D h_{R}$ and $x_{1}(z)=\mathcal{O}\left(R^{-1}\right)$ for $|z|$ in the support of $b_{R}$, (8.14) implies

$$
\begin{align*}
\left|X_{R}\right| & =\mathcal{O}\left(R^{-(1+a)}\right)  \tag{8.20}\\
\left|D X_{R}\right| & =\mathcal{O}\left(R^{(1-2 a)}\right)  \tag{8.21}\\
\left|D^{2} X_{R}\right| & =\mathcal{O}\left(R^{(3-3 a)}\right)  \tag{8.22}\\
\left|D^{2} X_{R}\right| & =\mathcal{O}\left(R^{(5-4 a)}\right) \tag{8.23}
\end{align*}
$$

If $0 \leq s \leq 1$, then
F0. (8.16) and (8.20) imply $\left|\Psi_{R}^{s}-\mathrm{id}\right|=\mathcal{O}\left(R^{-(1+a)}\right)$.
F1. (8.17) and (8.21) first give $\left|D \Psi_{R}^{s}\right|=\mathcal{O}(1)$. This together with (8.21) imply $\left|D \Psi_{R}^{s}-\mathrm{id}\right|=\mathcal{O}\left(R^{1-2 a}\right)$.

F2. (8.18), (8.21), (8.22), F1, and Duhamel's principle imply $\left|D^{2} \Psi_{R}^{s}\right|$ $=\mathcal{O}\left(R^{3-3 a}\right)$.
F3. In a similar way, as in F2, we derive $\left|D^{3} \Psi_{R}^{s}\right|=\mathcal{O}\left(R^{5-4 a}\right)$.
Let $u: \mathbb{R} \times[0,1] \rightarrow \mathbb{C}^{n}$ be a holomorphic function and and let $\omega:[0,1]$ $\rightarrow[0,1]$ be a smooth non-decreasing surjective approximation of the identity which is constant in a $\delta$-neighborhood of the ends of the interval. Consider the function $u_{R}(\tau+i t)=\Psi_{R}^{\omega(t)}(u(\tau+i t))$. We want estimates for $u_{R}, \bar{\partial} u_{R}$ and $D \bar{\partial} u_{R}$ and $\partial_{\tau} D \bar{\partial} u_{R}$.

F0 implies

$$
\begin{equation*}
u_{R}=u+\mathcal{O}\left(R^{-(1+a)}\right) \tag{8.24}
\end{equation*}
$$

For the estimates on $\bar{\partial} u_{R}$ and its derivatives, we note

$$
\begin{equation*}
\bar{\partial} u_{R}=D \Psi_{R}^{\omega(t)} \frac{\partial u}{\partial \tau}+i\left(D \Psi_{R}^{\omega(t)} \frac{\partial u}{\partial t}+\frac{d \omega}{d t} X_{R}(u)\right) \tag{8.25}
\end{equation*}
$$

By (8.20), (8.21), F1, and the holomorphicity of $u$,

$$
\begin{equation*}
\left|\bar{\partial} u_{R}\right|=\mathcal{O}\left(R^{1-2 a}\right)|D u|+\mathcal{O}\left(R^{-(1+a)}\right) . \tag{8.26}
\end{equation*}
$$

Taking derivatives of (8.25) with respect to $\tau$ and $t$, we find (using F0-3 and (8.20)-(8.23))

$$
\begin{align*}
\left|D \bar{\partial} u_{R}\right|= & \mathcal{O}\left(R^{1-2 a}\right)\left|D^{2} u\right|+\mathcal{O}\left(R^{3-3 a}\right)|D u|^{2} \\
& +\mathcal{O}\left(R^{1-2 a}\right)|D u|+\mathcal{O}\left(R^{-(1+a)}\right), \tag{8.27}
\end{align*}
$$

$$
\left|\partial_{\tau} D \bar{\partial} u_{R}\right|=\mathcal{O}\left(R^{1-2 a}\right)\left|D^{3} u\right|+\mathcal{O}\left(R^{3-3 a}\right)\left|D u \| D^{2} u\right|+\mathcal{O}\left(R^{5-4 a}\right)|D u|^{3}
$$

$$
\begin{equation*}
+\mathcal{O}\left(R^{1-2 a}\right)\left|D^{2} u\right|+\mathcal{O}\left(R^{3-3 a}\right)|D u|^{2}+\mathcal{O}\left(R^{-(1+a)}\right)|D u| \tag{8.28}
\end{equation*}
$$

Finally, let $\theta:[0,1] \rightarrow \mathbb{R}$ be a smooth function supported in a $\frac{1}{2} \delta$ neighborhood of the endpoints of the interval with $\theta^{\prime}(0)=\theta^{\prime}(1)=1$. Define

$$
\begin{equation*}
\hat{u}_{R}(\tau+i t)=u_{R}(\tau+i t)+i \theta(t) \bar{\partial} u_{R}(\tau+i t) . \tag{8.29}
\end{equation*}
$$

Then $u_{R}=\hat{u}_{R}$ on the boundary and $\hat{u}_{R}$ is holomorphic on the boundary. Also for some constant $C$

$$
\begin{align*}
\left|\hat{u}_{R}\right| & \leq C\left(\left|u_{R}\right|+\left|\bar{\partial} u_{R}\right|\right),  \tag{8.30}\\
\left|\bar{\partial} \hat{u}_{R}\right| & \leq C\left(\left|\bar{\partial} u_{R}\right|+\left|D \bar{\partial} u_{R}\right|\right),  \tag{8.31}\\
\left|D \bar{\partial} \hat{u}_{R}\right| & \leq C\left(\left|\bar{\partial} u_{R}\right|+\left|D \bar{\partial} u_{R}\right|+\left|\partial_{\tau} D \bar{\partial} u_{R}\right|\right) . \tag{8.32}
\end{align*}
$$

8.10. Self-tangency preshortening. Let $u: D_{m+1} \rightarrow \mathbb{C}^{n}$ be a rigid holomorphic disk with boundary on $L$ and with negative punctures $p_{1}, \ldots, p_{k}$ mapping to 0 . (The case of one positive puncture mapping to 0 is completely analogous to the case of one negative puncture so for simplicity, we consider only the case of negative punctures.)

For large $\rho>0$, let $R=R(\rho)$ be such that the intersection points of $L^{1}$ and $L_{R}^{2}(1)$ are $a^{ \pm}=\left( \pm\left(\rho+\rho^{a}\right)^{-1}, 0, \ldots, 0\right)$. Then $R(\rho)=\rho+\mathcal{O}\left(\rho^{-1}\right)$. Define

$$
u_{\rho}(\zeta)= \begin{cases}u(\zeta) & \text { for } \zeta \in D_{m+1} \backslash\left(\bigcup_{j=1}^{k} E_{p_{j}}\left[-\frac{1}{2} \rho\right]\right), \\ \hat{u}_{R(\rho)}(\tau+i t) & \text { for } \zeta=\tau+i t \in E_{p_{j}}\left[-\frac{1}{2} \rho\right]\end{cases}
$$

Then there exist unique functions

$$
\xi_{R}(j): E_{p_{j}}[-\rho] \rightarrow T_{a^{-}} \mathbb{C}^{n}
$$

such that

$$
\exp ^{R, t}\left(\xi_{R}(j)(\zeta)\right)=u_{\rho}(\zeta), \quad \zeta \in E_{p_{j}}[-\rho]
$$

where $\exp ^{R, t}$ denotes the exponential map in the metric $g(R, \omega(t), t)$ at $a^{-}$.

Let $\alpha_{\rho}:(-\infty,-\rho] \times[0,1] \rightarrow \mathbb{C}$ be a smooth cut-off function, real valued and holomorphic on the boundary and such that $\alpha_{\rho}(\tau+i t)=1$ for $\tau$ in a small neighborhood of $-\rho, \alpha_{\rho}(\tau+i t)=0$ for $\tau \leq-\rho-\frac{1}{2} \rho^{a}$, and $\left|D^{k} \alpha_{\rho}\right|=\mathcal{O}\left(\rho^{-a}\right), k=1,2$. Define $w_{\rho}: D_{m+1} \rightarrow \mathbb{C}^{n}$ as
$w_{\rho}(\zeta)= \begin{cases}u_{\rho}(\zeta) & \text { for } \zeta \in D_{m+1} \backslash\left(\bigcup_{j} E_{p_{j}}[-\rho]\right), \\ \exp ^{R, t}\left(\alpha_{\rho}(\zeta) \xi_{R}(j)(\zeta)\right) & \text { for } \zeta=\tau+i t \in E_{p_{j}}[-\rho], j=1, \ldots, k, \\ a_{-} & \text {for } \zeta \in \bigcup_{j} E_{p_{j}}\left[-\rho-\frac{1}{2} \rho^{a}\right] .\end{cases}$
8.11. Weight functions for shortened disks. Let $u: D_{m+1} \rightarrow \mathbb{C}^{n}$ be a rigid holomorphic disk with boundary on $L$. Let $\epsilon>0$ be small and let $e_{\rho}: D_{m+1} \rightarrow \mathbb{R}$ be a function which equals $e^{-\epsilon|\tau|}$ for $\tau+i t \in$ $E_{p_{j}} \backslash E_{p_{j}}[-\rho]$ and is constantly equal to $e^{-\epsilon \rho}$ for $\tau+i t \in E_{p_{j}}[-\rho]$. Define $\mathcal{W}_{2,-\epsilon, \rho}$ just as in Section 5.8, but replacing the weight function $e_{\epsilon}$ with the new weight function $e_{\rho}$. The corresponding weighted norms will be denoted $\|\cdot\|_{2,-\epsilon, \rho}$. We also write $\mathcal{H}_{1,-\epsilon, \rho}[0]\left(D_{m+1}, T^{* 0,1} \otimes \mathbb{C}^{n}\right)$ to denote the subspace of elements in the Sobolev space with the weight function $e_{\rho}$ which vanishes on the boundary.

### 8.12. Estimates for self-tangency preshortened disks.

Lemma 8.10. The function $w_{\rho}$ in (8.33) lies in $\mathcal{W}_{2,-\epsilon, \rho}$ (see Remark 8.6 for notation) and there exists a constant $C$ such that

$$
\left\|\bar{\partial} w_{\rho}\right\|_{1,-\epsilon, \rho} \leq C e^{-\epsilon \rho} \rho^{-1-\frac{1}{2} a} .
$$

Proof. The first statement is obvious. Consider the second. In $D_{m+1} \backslash$ $\left(E_{p_{j}}[-\rho]\right), w_{\rho}$ equals $u$ which is holomorphic. It thus remains to check $E_{p_{j}}[-\rho] \approx(-\infty,-\rho] \times[0,1]$.

Taylor expansion of $\exp ^{R, t}$ gives

$$
\begin{equation*}
\exp ^{R, t} \xi=\xi-\Gamma_{i j}^{k}(R, t) \xi^{i} \xi^{j} \partial_{k}+\mathcal{O}\left(|\xi|^{3}\right) \tag{8.34}
\end{equation*}
$$

The Taylor expansion of the inverse then gives

$$
\begin{equation*}
\xi_{R}=\hat{u}_{R}+\Gamma_{i j}^{k}(R, t) \hat{u}_{R}^{i} \hat{u}_{R}^{j} \partial_{k}+\mathcal{O}\left(\left|\hat{u}_{R}\right|^{3}\right) \tag{8.35}
\end{equation*}
$$

Thus in $(-\infty,-\rho] \times[0,1]$, we have

$$
\begin{equation*}
w_{\rho}=\alpha_{\rho} \hat{u}_{R}+\left(\alpha_{\rho}-\alpha_{\rho}^{2}\right) \Gamma_{i j}^{k}(R, t) \hat{u}_{R}^{i} \hat{u}_{R}^{j} \partial_{k}+\mathcal{O}\left(\left|\hat{u}_{R}\right|^{3}\right) . \tag{8.36}
\end{equation*}
$$

Now $R=\rho+\mathcal{O}\left(\rho^{-1}\right)$ from Section 8.10, $\left|D^{k} \alpha_{\rho}\right|=\mathcal{O}\left(\rho^{-a}\right)$ for all cut-off functions, and by Lemma $4.6\left|D^{k} u\right|=\mathcal{O}\left(\rho^{-(1+k)}\right)$, in $(-\infty,-\rho] \times[0,1]$; thus, applying (8.30) through (8.32) to (8.36), we get

$$
\left|\bar{\partial} w_{\rho}\right|+\left|D \bar{\partial} w_{\rho}\right|=\mathcal{O}\left(\rho^{-(1+a)}\right) .
$$

Noting that $\bar{\partial} w_{\rho}$ is supported on an interval of length $\frac{1}{2} \rho^{a}$, so multiplying with the weight function, we find

$$
\left\|\bar{\partial} w_{\rho}\right\|_{1,-\epsilon, \rho} \leq C e^{-\epsilon \rho} \rho^{-1-\frac{1}{2} a}
$$

q.e.d.
8.13. Controlled invertibility for self-tangency shortening. Let $d \Gamma_{\rho}$ denote the differential of the map

$$
\Gamma_{\rho}: \mathcal{W}_{2,-\epsilon, \rho} \rightarrow \mathcal{H}_{1,-\epsilon, \rho}[0]\left(D_{m+1}, T^{* 0,1} \otimes \mathbb{C}^{n}\right)
$$

Referring to Sections 8.6.1 and 8.6.2, we assume that $m \geq 2$ and $l(j) \geq 2$ for each $j$.

Lemma 8.11. There exist constants $C$ and $\rho_{0}$ such that if $\rho>\rho_{0}$, then there is a continuous right inverse $G_{\rho}$ of $d \Gamma_{\rho}$

$$
G_{\rho}: \mathcal{H}_{1,-\epsilon, \rho}[0]\left(T^{* 0,1} D_{r}(\rho) \otimes \mathbb{C}^{n}\right) \rightarrow T_{\left(w_{\rho}, \kappa_{\rho}, 0\right)} \mathcal{W}_{2,-\epsilon, \rho}
$$

such that for any $\delta>0$

$$
\begin{equation*}
\left\|G_{\rho}(\xi)\right\| \leq C \rho^{1+\delta}\|\xi\|_{1,-\epsilon, \rho} \tag{8.37}
\end{equation*}
$$

Proof. The kernel

$$
\operatorname{ker}\left(d \Gamma_{u}\right) \subset T_{u} \mathcal{W}_{2,-\epsilon} \oplus T_{\kappa} \mathcal{C}_{m+1}
$$

has dimension 0 . By the invertibility of $d \Gamma_{u}$, we conclude there is a constant $C$ such that for $\xi \in T_{u} \mathcal{W}_{\epsilon, 2, \rho}$ we have

$$
\begin{equation*}
\|\xi\| \leq C\left\|d \Gamma_{u, \rho} \xi\right\|_{1,-\epsilon} \tag{8.38}
\end{equation*}
$$

Assume that (8.37) is not true. Then there exists a sequence $\xi_{N} \in$ $T_{w_{\rho}} \mathcal{W}_{2,-\epsilon, \rho(N)}$ with $\rho(N) \rightarrow \infty$ as $N \rightarrow \infty$ such that

$$
\begin{align*}
& \left\|\xi_{N}\right\|=1  \tag{8.39}\\
& \left\|d \Gamma_{\rho} \xi_{N}\right\|_{1,-\epsilon, \rho(N)} \leq C \rho^{-1-\frac{\delta}{2}} \tag{8.40}
\end{align*}
$$

Let $\alpha: D_{m+1} \rightarrow \mathbb{C}$ be a smooth function which equals 0 on $E_{p_{j}}[-\rho-$ $\left.\frac{1}{4} \rho^{a}\right]$ and equals 1 on $D_{m+1} \backslash\left(\bigcup E_{p_{j}}[-\rho-10]\right)$, which is real valued and holomorphic on the boundary and with $\left|D^{k} \alpha\right|=\mathcal{O}\left(\rho^{-a}\right), k=1,2$. Then (8.38) implies

$$
\begin{equation*}
\left\|\alpha \xi_{N}\right\| \leq C\left(\left\|(\bar{\partial} \alpha) \xi_{N}\right\|_{1,-\epsilon}+\left\|\alpha d \Gamma_{u, \rho} \xi_{N}\right\|_{1,-\epsilon}\right)=\mathcal{O}\left(\rho^{-a}\right) \tag{8.41}
\end{equation*}
$$

Finally, we let $\hat{\phi}:\left(-\infty,-\rho+\rho^{a}\right] \rightarrow \mathbb{C}$ be the function which equals $\theta(\rho)-\theta(\tau)$, where $\theta(\tau)$ denotes the angle between the tangent line of $L_{\rho}^{2}(1)$ intersected with the $z_{1}$-plane (the plane of the first coordinate in $\left.\mathbb{C}^{n}\right)$ at $u(\tau+i)$ and the real line in that plane. From Lemma 4.6, we calculate that $\left|D^{k} \hat{\phi}\right|=\mathcal{O}\left(\rho^{a-2}\right), 0 \leq k \leq 2$. Using the same procedure as for cut-off functions, we extend it to a function $\phi:\left(-\infty,-\rho+\rho^{a}\right) \times[0,1]$ which is holomorphic on the boundary, which equals $\hat{\phi}$ on $(-\infty,-\rho+$ $\left.\rho^{a}\right) \times\{1\}$ and which equals 0 on $\left(-\infty,-\rho+\rho^{a}\right) \times\{0\}$ and with the same derivative estimates. Let $\mathbf{M}=\operatorname{Diag}(\phi, 1, \ldots, 1)$.

Let $\alpha$ be a cut-off function which is 0 in $D_{m+1} \backslash E_{p_{j}}\left[-\rho+\rho^{a}\right]$ and 1 on $E_{p_{j}}[-\rho]$. Having frozen the angle away from 0 , we can use Lemmas 6.8 and 6.9 (assuming that $\epsilon$ is smaller than the smallest component of the complex angle) to get

$$
\begin{equation*}
\left\|e^{-\epsilon \rho} \alpha \mathbf{M} \xi_{N}\right\| \leq C \rho\left(\left\|e^{-\epsilon \rho}(\bar{\partial} \alpha \mathbf{M}) \xi_{N}\right\|+\left\|e^{-\epsilon \rho} \alpha \mathbf{M} d \Gamma_{\rho} \xi_{N}\right\|\right) \tag{8.42}
\end{equation*}
$$

The first term on the right-hand side inside the brackets is $\mathcal{O}\left(\rho^{a-2}\right)+$ $\mathcal{O}\left(\rho^{-2 a}\right)$ the second term is $\mathcal{O}\left(\rho^{-1-\delta}\right)$. Hence as $\rho \rightarrow \infty$ the right-hand side goes to 0 . This together with (8.41) contradicts (8.39) and we conclude the lemma holds.
8.14. Perturbations for self-tangency gluing. For $R>0$ with $R^{-1} \ll r_{0}$, let $a_{R}:[0, \infty) \rightarrow \mathbb{R}$ be a smooth non-increasing function with support in $\left[0, \frac{1}{2} R^{-1}\right)$ and with the following properties

$$
\begin{align*}
a_{R}(r) & =R^{-1} \text { for } r \in\left[0, R^{-2}\right],  \tag{8.43}\\
\left|D a_{R}(r)\right| & =\mathcal{O}(1) \\
\left|D^{2} a_{R}(r)\right| & =\mathcal{O}(R)
\end{align*}
$$

The existence of such functions is easily established. Let $h_{R}: \mathbb{C}^{n} \rightarrow \mathbb{C}^{n}$ be given by

$$
\begin{equation*}
h_{R}(z)=x_{1}(z) a_{R}\left(\left|z_{1}\right|\right) . \tag{8.44}
\end{equation*}
$$

For $s>0$, let $\Phi_{R}^{s}$ denote the time $s$ Hamiltonian flow of $h_{R}$ and let $L_{R}(s)$ denote the Legendrian submanifold which results when $\Phi_{R}^{s}$ is lifted to a local contact flow on $\mathbb{C}^{n} \times \mathbb{R}$ which is used to move $L^{2}$. (Note that $\Phi_{R}^{s}$ fixes the last $n-1$ coordinates and has small support in the $z_{1}$-direction and so its lift can be extended to the identity outside $L^{2}(s)$.) Let $L_{R}^{2}(s)=\Phi_{R}^{s}\left(L^{2}\right)$. We pick $a_{R}$ so that $L_{R}^{2}(s) \cap L^{1}=\emptyset$, for $0<s \leq(K R)^{-1}$ for some fixed $K>4$.

As in Section 8.9, we derive the estimates

$$
\begin{align*}
& \left|\Phi_{R}^{s}-\mathrm{id}\right| \leq \mathcal{O}\left(R^{-2}\right)  \tag{8.45}\\
& \left|D \Phi_{R}^{s}-\mathrm{id}\right|=\mathcal{O}\left(R^{-1}\right)  \tag{8.46}\\
& \left|D^{2} \Phi_{R}^{s}\right|=\mathcal{O}(1) \tag{8.47}
\end{align*}
$$

For convenient notation, we write $\gamma_{R}^{2}(s)$ for the curve in which $L_{R}^{2}(s)$ intersects the $z_{1}$-line in a neighborhood of 0 .
8.15. Self-tangency pregluing. Let $u: D_{m} \rightarrow \mathbb{C}^{n}$ be a rigid holomorphic disk with boundary on $L$ and with negative punctures $p_{1}, \ldots, p_{k}$ (as above, we write $\left.S=\left\{p_{1}, \ldots, p_{k}\right\}\right)$ mapping to 0 . Let $v_{j}: D_{l(j)+1} \rightarrow \mathbb{C}_{n}$ be rigid holomorphic disks with positive punctures $q_{j}$ mapping to 0 .

For $0<\rho<\infty$, let $R=\rho, s=(K \rho)^{-1}$ and let $L_{\rho}$ be the Legendrian submanifold which results when $\Phi_{R}^{s}$ is applied. Consider the region $\Xi_{\rho}$ in the $z_{1}$-line bounded by the curves $\gamma_{R}^{2}(s), \gamma_{R}^{1}(s), u^{1}(\rho+i t), 0 \leq t \leq 1$, and $v(j)^{1}(\rho+i t), 0 \leq t \leq 1$. By the Riemann mapping theorem there exists a holomorphic map from a rectangle $\phi_{\rho}:[-A(\rho), A(\rho)] \times[0,1] \rightarrow \mathbb{C}$ which parameterizes this region in such a way that $[-A(\rho), A(\rho)] \times\{j-1\}$ maps to $\gamma_{j}(s), j=1,2$. Moreover, since $\Xi_{\rho}$ is symmetric with respect to reflections in the $\operatorname{Im} z_{1}=y_{1}$-axis we have $\phi_{\rho}(0+i[0,1]) \subset\left\{\operatorname{Re} z_{1}=\right.$ $\left.x_{1}=0\right\}$.

Lemma 8.12. The shape of the rectangle depends on $\rho$. More precisely, there exists constants $0<K_{1}<K_{2}<\infty$ such that $K_{1} \rho \leq A(\rho) \leq$ $K_{2} \rho$ for all $\rho$.

Proof. Identify the $z_{1}$-line with $\mathbb{C}$. Consider the region $\Theta_{\rho}$ bounded by the circles of radius 1 and $1+4 \rho^{-2}$ both centered at $i \in \mathbb{C}$, and the lines through $i$ which intersects the $x_{1}$-axis in the points $\pm 2(\rho)^{-1}$. Mark the straight line segments of its boundary. The conformal modulus of this region is easily seen to be $\rho+\mathcal{O}\left(\rho^{-1}\right)$.

On the other hand, using (8.45) and (8.46) one constructs a ( $K+$ $\mathcal{O}\left(\rho^{-1}\right)$ )-quasi conformal map from $\Theta_{\rho}$ to $\Xi_{\rho}$, for some $K>0$ independent of $\rho$. This implies the conformal modulus $m_{\rho}$ of $\Theta_{\rho}$ satisfies

$$
\begin{equation*}
\left(K+\mathcal{O}\left(\rho^{-1}\right)\right)^{-1}\left(\rho+\mathcal{O}\left(\rho^{-1}\right)\right) \leq m_{\rho} \leq\left(K+\mathcal{O}\left(\rho^{-1}\right)\right)\left(\rho+\mathcal{O}\left(\rho^{-1}\right)\right) \tag{8.48}
\end{equation*}
$$

and the lemma follows.
q.e.d.

Let $u^{1}$ and $v_{j}^{1}$ denote the $z_{1}$-components of the maps $u$ and $v_{j}$. Since $\Phi_{R}^{s}$ fixes $\gamma_{2}$ outside $\left|x_{1}\right| \leq(2 \rho)^{-1}$, we note that
(8.49) $u^{1}$ maps the region $E_{p_{j}}[-\rho] \backslash E_{p_{j}}[-2 \rho]$ into $\Theta_{\rho} \backslash \phi_{\rho}(0 \times[0,1])$.
and that
$v_{j}^{1}$ maps the region $E_{p_{j}}[\rho] \backslash E_{p_{j}}[2 \rho]$ into $\Theta_{\rho} \backslash \phi_{\rho}(0 \times[0,1])$.
Fix $0<a<\frac{1}{4}$. Using $u^{1}, v_{j}^{1}$, the conformal map $\phi_{\rho}$ and their inverses, we construct a complex 1-dimensional manifold $D_{r}(\rho)$ by gluing $\Omega_{j}(\rho)=[-A(\rho), A(\rho)] \times[0,1]$ to $D_{m+1} \backslash\left(\bigcup_{j} E_{p_{j}}[-(1+a) \rho]\right)$ and $D_{l(j)+1} E_{q}[(1+a) \rho]$. Note that, by construction, $D_{r}(\rho)$ comes equipped with a holomorphic function

$$
\begin{equation*}
w_{\rho}^{1}: D_{r}(\rho) \rightarrow \mathbb{C}, \tag{8.51}
\end{equation*}
$$

which equals $u^{1}$ on $D_{m+1} \backslash \bigcup_{j} E_{p_{j}}[-(1+a) \rho]$, which equals $v_{j}^{1}$ on $D_{l(j)+1} \backslash$ $E_{q_{j}}[(1+a) \rho]$, and which equals $\phi_{\rho}$ on $\Omega_{j}$, for all $j$.

We next exploit the product structure of $\Pi_{\mathbb{C}} L$ in a neighborhood of 0 . If $u^{\prime}$ and $v_{j}^{\prime}$ denotes the remaining components of $u$ and $v_{j}$ so that $u=$ $\left(u^{1}, u^{\prime}\right)$ and $v_{j}=\left(v_{j}^{1}, v_{j}^{\prime}\right)$, then in some neighborhood of the punctures $q_{j}$ and $p_{j}, v_{j}^{\prime}$ and $u_{j}^{\prime}$ are holomorphic functions with boundary on the two transverse Lagrangian manifolds $P_{1}$ and $P_{2}$, see Section 3. As in Section 5.3 , we find a 1 -parameter family $g(\sigma)$ of metrics on $\mathbb{C}^{n-1} \approx\left\{z_{1}=0\right\}$. Then, for $M$ sufficiently large, there exist unique vector valued functions $\xi_{j}^{\prime}$ and $\eta_{j}^{\prime}$ such that

$$
\begin{gather*}
\exp _{0}^{t} \xi_{j}^{\prime}(\tau+i t)=u^{\prime}(\tau+i t), \quad \tau+i t \in E_{p_{j}}[-M],  \tag{8.52}\\
\exp _{0}^{t} \eta_{j}^{\prime}(\tau+i t)=v_{j}^{\prime}(\tau+i t), \quad \tau+i t \in E_{q_{j}}[M] . \tag{8.53}
\end{gather*}
$$

Now pick a cut-off function $\alpha^{+}$which equals 1 on $D_{m+1} \backslash \bigcup_{j} E_{p_{j}}[-\rho+5]$ and 0 on $E_{p_{j}}[-\rho+3]$. Pick similar cut-off functions $\alpha^{-}$on $D_{l(j)+1}$. Define $w_{\rho}^{\prime}: D_{r}(\rho) \rightarrow \mathbb{C}^{n-1}$ by

$$
w_{\rho}^{\prime}(\zeta)= \begin{cases}u^{\prime}(\zeta), & \zeta \in D_{m+1} \backslash \bigcup E_{p_{j}}[-\rho+5]  \tag{8.54}\\ v_{j}^{\prime}(\zeta), & \zeta \in D_{l(j)+1} \backslash E_{q_{j}}[\rho-5] \\ \exp _{0}^{t}\left(\alpha^{+}(\zeta) \xi_{j}(\zeta)\right), & \zeta \in E_{p_{j}}[-\rho+5] \backslash E_{p_{j}}[-\rho], \\ \exp _{0}^{t}\left(\alpha^{-}(\zeta) \eta_{j}(\zeta)\right), & \zeta \in E_{q_{j}}[\rho-5] \backslash E_{p_{j}}[\rho], \\ 0, & \zeta \in \Omega_{j} .\end{cases}
$$

Finally, combining (8.51) and (8.54), we define

$$
\begin{equation*}
w_{\rho}=\left(w_{\rho}^{1}, w_{\rho}^{\prime}\right) . \tag{8.55}
\end{equation*}
$$

8.16. Weight functions and Sobolev norms for self tangency gluing. Consider $D_{r}(\rho)$ from the previous section, $\epsilon>0$, and a smooth function $f: D_{r}(\rho) \rightarrow \mathbb{C}^{n}$. Let

- $f^{+}$denote the restriction of $f$ to

$$
\operatorname{int}\left(D_{m+1} \backslash \bigcup_{j} E_{p_{j}}[-(1+a) \rho]\right)
$$

which we consider as a subset of $D_{m+1}$.

- $f^{-}$denote the restriction of $f$

$$
\bigcup_{j} \operatorname{int}\left(D_{l(j)+1} \backslash E_{q_{j}}[(1+a) \rho]\right),
$$

which we consider as subset of the disjoint union $\bigcup_{j} D_{l(j)+1}$

- $f^{0}$ denote the restriction of $f$ to the disjoint union $\bigcup_{j} \operatorname{int}\left(\Omega_{j}(\rho)\right)$.

For $\epsilon>0$, let $e_{\epsilon}^{-}$denote the weight function on $\bigcup_{j} D_{l(j)+1}$ which equals 1 on $D_{l(j)+1} \backslash E_{q_{j}}$ and equals $e^{\delta|\tau|}$ in $E_{q_{j}}$, each $j$. Let $\|\cdot\|_{k, \epsilon,-}$ denote the Sobolev norm with weight $e_{\delta}^{-}$. Let $e_{\epsilon}^{0}$ denote the weight function on $\Omega_{j}$ which equals $e^{\epsilon(A(\rho)+\rho+\tau)}$ and $\|\cdot\|_{k, \epsilon, 0}$ denote the Sobolev norm with this weight. Finally, let $e_{\epsilon}^{+}$be the function on $D_{m+1}$ which equals $e^{2 \epsilon(A(\rho)+\rho)}$ on $\left.D_{m+1} \backslash \bigcup_{j} E_{p_{j}}\right]$ and equals $e^{2 \epsilon(A(\rho)+\rho)-\epsilon|\tau|}$ in $E_{p_{j}}$. Let $\|\cdot\|_{k, \epsilon,+}$ denote the corresponding norm.

Define

$$
\begin{equation*}
\|f\|_{k, \epsilon, \rho}=\left\|f_{+}\right\|_{k, \epsilon,+}+\left\|f_{0}\right\|_{k, \epsilon, 0}+\left\|f_{-}\right\|_{k, \epsilon,-} . \tag{8.56}
\end{equation*}
$$

Using this norm, we define as in the shortening case the spaces $\mathcal{W}_{2, \epsilon, \rho}$ and $\mathcal{H}_{1, \epsilon, \rho}[0]\left(D_{m+1}, T^{* 0,1} \otimes \mathbb{C}^{n}\right)$. The $\bar{\partial}$-map

$$
\Gamma: \mathcal{W}_{2, \epsilon, \rho} \rightarrow \mathcal{H}_{1, \epsilon, \rho}[0]\left(T^{* 0,1} D_{m} \otimes \mathbb{C}^{n}\right)
$$

is defined in the natural way.

### 8.17. Estimates for self tangency glued disks.

Lemma 8.13. The function $w_{\rho}$ in (8.55) lies in $\mathcal{W}_{2, \epsilon, \rho}$ and there exists a constant $C$ such that

$$
\left\|\bar{\partial} w_{\rho}\right\|_{1, \epsilon, \rho} \leq C e^{\left(-\theta+2 K_{2} \epsilon\right) \rho}
$$

where $\theta \gg \epsilon$ is the smallest non-zero complex angle at the self tangency point 0 and where $K_{2}$ is as in Lemma 8.12.

Proof. Note that the first coordinate of $w_{\rho}$ is holomorphic and that the support of $\bar{\partial} w_{\rho}$ is disjoint from $\Omega_{j}$. Using the asymptotics of $u^{\prime}$ and $v_{j}^{\prime}$, the proof of Lemma 8.5 applies to give the desired estimate once we note that the weight function is $\mathcal{O}\left(e^{K_{2} \epsilon \rho}\right)$ by Lemma 8.12. q.e.d.
8.18. Estimates for real boundary conditions. In order to prove the counterpart of Lemma 8.11 in the self tangency gluing case, we study an auxiliary non-compact counterpart of the gluing region.

Let $\Omega(\rho)=[-A(\rho), A(\rho)] \times[0,1]$ and let $M_{\rho}$ be the complex manifold which results when $(-\infty,-(1-a) \rho] \times[0,1]$ and $[(1-a) \rho, \infty) \times[0,1]$ are glued to $\Omega(\rho)$ with the holomorphic gluing maps $u^{1} \circ\left(\phi_{\rho}\right)^{-1}$ and $v_{j}^{1} \circ\left(\phi_{\rho}\right)^{-1}$, respectively. (That is, the maps which were used to construct $D_{r}(\rho)$.) We consider Sobolev norms on $M_{\rho}$ similar to those used above. For $\epsilon>0$, let

- $e_{\epsilon}^{0}::[-A(\rho), A(\rho)] \times[0,1] \rightarrow \mathbb{R}$ be the function $e_{\epsilon}^{0}(\tau+i t)=e^{\epsilon \tau}$,
- $e_{\epsilon}^{-}:(-\infty,-(1-a) \rho] \times[0,1] \rightarrow \mathbb{R}$ be the function $e_{\epsilon}^{-}(\tau+i t)=$ $e^{\epsilon \epsilon \rho-A(\rho)+\tau)}$
- $e_{\epsilon}^{+}:[(1-a) \rho, \infty) \times[0,1] \rightarrow \mathbb{R}$ be the function $e_{\epsilon}^{+}(\tau+i t)=$ $e^{\epsilon(-\rho+A(\rho)+\tau)}$
If $f: M_{\rho} \rightarrow \mathbb{C}$ is function we let as above $f^{-}, f^{0}, f^{+}$denote the restrictions of $f$ to the interiors of the pieces from which $M_{\rho}$ was constructed and define the Sobolev norm

$$
\begin{equation*}
\|f\|_{k, \rho, \epsilon}=\left\|f^{-}\right\|_{k, \epsilon}+\left\|f^{0}\right\|_{k, \epsilon}+\left\|f^{+}\right\|_{k, \epsilon} . \tag{8.57}
\end{equation*}
$$

Lemma 8.14. There are constants $C$ and $\rho_{0}$ if $\rho>0$ and if $f: M_{\rho} \rightarrow$ $\mathbb{C}$ is function which is real valued and holomorphic on the boundary and has $\|f\|_{k, \rho, \epsilon} \leq \infty$, then

$$
\begin{equation*}
\|f\|_{k, \rho, \epsilon} \leq C\|\bar{\partial} f\|_{k-1, \rho, \epsilon}, \tag{8.58}
\end{equation*}
$$

for $k=1,2$.
Proof. To prove the lemma, we first study the gluing functions. Let $\psi:[-\rho,-(1-a) \rho) \times[0,1] \rightarrow[-\mathcal{A}(\rho), 0] \times[0,1]$ be the function $u^{1} \circ\left(\phi_{\rho}\right)^{-1}$. Note that $\psi$ is holomorphic and that by (8.49) has a holomorphic extension (still denoted $\psi$ ) to $[-\rho, 0) \times[0,1]$.

To simplify notation, we change coordinates and think of the source $[-\rho, 0) \times[0,1]$ as $[0, \rho) \times[0,1]$ and of the target $[-A(\rho), 0] \times[0,1]$ as $[0, A(\rho)] \times[0,1]$. Thus

$$
\begin{equation*}
\psi:[0, \rho] \times[0,1] \rightarrow[0, A(\rho)] \times[0,1] \tag{8.59}
\end{equation*}
$$

is a holomorphic map. Consider the complex derivative $\frac{\partial \psi}{\partial z}$. This is again a holomorphic function which is real on the boundary of $[0, \rho) \times$ $[0,1]$. In analogy with Lemma 6.2, we conclude that

$$
\begin{equation*}
\frac{\partial \psi}{\partial z}=\sum_{n \in Z} c_{n}^{\prime} e^{n \pi z} \tag{8.60}
\end{equation*}
$$

for some real constants $c_{n}^{\prime}$. Integrating this and using $\psi(0)=0$, we find

$$
\begin{equation*}
\psi(z)=c_{0} z+\sum c_{n} e^{n \pi z} \tag{8.61}
\end{equation*}
$$

for some real constants $c_{n}$. Then

$$
\begin{equation*}
i=\psi(i)=c_{0} i+\sum c_{n} e^{n \pi i} \tag{8.62}
\end{equation*}
$$

and we conclude $c_{0}=1$. Moreover, if $\psi^{d}$ denotes the double of $\psi$ (which has the same Fourier expansion), then since $\psi^{d}(i t)$ is purely imaginary for $0 \leq t \leq 2$, we find that $c_{n}=-c_{n}$ for all $n \neq 0$. Thus

$$
\begin{equation*}
\psi(z)=z+\sum_{n} c_{n}\left(e^{n \pi z}-e^{-n \pi z}\right) \tag{8.63}
\end{equation*}
$$

The area of the image of $\psi^{d}$ is $\mathcal{O}(\rho)$ by Lemma 8.12. Since this area equals the $L^{2}$-norm of the derivative of $\psi^{d}$, we conclude that

$$
\begin{equation*}
2 \int_{0}^{\rho} 1^{2} d \tau+\sum_{n \in Z} \int_{0}^{\rho} n^{2} \pi^{2}\left|c_{n}\right|^{2} e^{2 n \pi \tau} d \tau=\mathcal{O}(\rho) \tag{8.64}
\end{equation*}
$$

Integrating, we find there exists a constant $K$ and $0<\delta \ll 1$ such that

$$
\begin{equation*}
\left|c_{n}\right| \leq K \rho(n)^{-\frac{1}{2}} e^{-n \pi \rho} \leq K e^{-n(\pi-\delta) \rho} \tag{8.65}
\end{equation*}
$$

for each $n \neq 0$. Thus, in the gluing region $[0, a \rho)$, we find

$$
\begin{equation*}
|\psi(z)-z| \leq K \sum_{n>0} e^{-n(\pi-\delta-a) \rho} \leq K^{\prime} e^{-(\pi-2(\delta+a)) \rho}=K^{\prime} e^{-\eta \rho} \tag{8.66}
\end{equation*}
$$

where $\eta>0$. Similarly, one shows $|D \psi-\mathrm{id}| \leq K e^{-\frac{1}{2} \eta \rho}$ and $\left|D^{2} \psi\right| \leq$ $K e^{-\frac{1}{2} \eta \rho}$.

Assume (8.58) is not true, then there exists a sequence $f_{j}$ of functions on $M_{\rho(j)}, \rho(j) \rightarrow \infty$ as $j \rightarrow \infty$, with

$$
\begin{align*}
& \left\|f_{j}\right\|_{2, \rho, \epsilon}=1,  \tag{8.67}\\
& \left\|\bar{\partial} f_{j}\right\|_{1, \rho, \epsilon} \rightarrow 0, \quad \text { as } j \rightarrow \infty \tag{8.68}
\end{align*}
$$

Let $\gamma:(-\infty,-(1-a) \rho] \times[0,1]$ be a cut-off function which equals 1 on $\left(-\infty,-\left(1-\frac{1}{4} a\right) \rho\right] \times[0,1]$ which equals 0 on $\left[-\left(1-\frac{1}{2} a\right) \rho,-(1-a) \rho\right)$, has $\left|D^{k} \gamma\right|=\mathcal{O}\left(\rho^{-1}\right), k=1,2$, and is real valued and holomorphic on the boundary. Then by uniform invertibility of the $\bar{\partial}$-operator on the strip with constant weight $\epsilon$, we find

$$
\begin{equation*}
\|\gamma f\|_{2, \epsilon} \leq C\left(\|(\bar{\partial} \gamma) f\|_{1, \epsilon}+\|\gamma \bar{\partial} f\|_{1, \epsilon}\right) \tag{8.69}
\end{equation*}
$$

Here both terms on the right-hand side goes to 0 as $\rho \rightarrow \infty$. In a similar way, we conclude

$$
\begin{equation*}
\|\beta f\|_{2, \epsilon} \rightarrow 0 \tag{8.70}
\end{equation*}
$$

for $\beta$ a cut-off function on $[(1-a) \rho, \infty)$.
Now let $\alpha$ be a cut-off function on $[-A(\rho), A(\rho)] \times[0,1]$ which equals 1 on $[-A(\rho)+2, A(\rho)-2] \times[0,1]$ and equals 0 outside $[-A(\rho)+1, A(\rho)-$ $1] \times[0,1]$. We find

$$
\begin{equation*}
\|\alpha f\|_{2, \epsilon} \leq C\left(\|(\bar{\partial} \alpha) f\|_{1, \epsilon}+\|\alpha \bar{\partial} f\|_{1, \epsilon}\right) \tag{8.71}
\end{equation*}
$$

Here the second term on the right-hand side goes to 0 as $\rho \rightarrow \infty$ by (8.68). The first goes to 0 as well since $\|\gamma f\| \rightarrow 0$ and $\|\beta f\| \rightarrow 0$ and since the transition functions are very close to the identity for $\rho$ large.

In conclusion, we find $\|f\|_{2, \rho, \epsilon} \rightarrow 0$, contradicting (8.67), and (8.58) holds.
q.e.d.
8.19. Uniform invertibility for self tangency gluing. Let $d \Gamma_{\rho}$ denote the differential of the map

$$
\Gamma: \mathcal{W}_{2, \epsilon, \rho} \rightarrow \mathcal{H}_{1, \epsilon, \rho}[0]\left(D_{m+1}, T^{* 0,1} \otimes \mathbb{C}^{n}\right)
$$

at $w_{\rho}$. Referring to Sections 8.6.1 and 8.6.2, we assume that $m \geq 2$ and $l(j) \geq 2$ for each $j$.

Lemma 8.15. There exist constants $C$ and $\rho_{0}$ such that if $\rho>\rho_{0}$ and then there is a continuous right inverse $G_{\rho}$ of $d \Gamma_{\rho}$

$$
G_{\rho}: \mathcal{H}_{1, \epsilon, \rho}[0]\left(T^{* 0,1} D_{r}(\rho) \otimes \mathbb{C}^{n}\right) \rightarrow T_{\left(w_{\rho}, \kappa_{\rho}\right)} \mathcal{W}_{2, \epsilon, \rho}
$$

such that

$$
\left\|G_{\rho}(\xi)\right\| \leq C\|\xi\|_{1, \epsilon, \rho}
$$

Proof. Recall $0<\epsilon \ll \theta$, where $\theta>0$ is the smallest non-zero complex angle at the self-tangency point. Assume we glue $k$ disks $v_{1}, \ldots, v_{k}$ to $u$. The kernels

$$
\begin{array}{r}
d \Gamma_{\left(u, \kappa_{1}\right)} \subset T_{\left(u, \kappa_{1}\right)} \mathcal{W}_{2,-\epsilon}, \\
d \Gamma_{\left(v_{j}, \kappa_{2}(j)\right)} \subset T_{\left(v_{j}, \kappa_{2}(j)\right)} \mathcal{W}_{2, \epsilon}, \tag{8.72}
\end{array}
$$

are both of dimension 0 and $d \Gamma_{\left(u, \kappa_{1}\right)}$ and $d \Gamma_{\left(v_{j}, \kappa_{2}(j)\right)}$ are invertible.
As usual, we consider the embedding

$$
\begin{equation*}
T_{\kappa_{1}} \mathcal{C}_{m+1} \oplus \bigoplus_{j=1}^{k} T_{\kappa_{2}(j)} \mathcal{C}_{l(j)+1} \rightarrow T_{\kappa_{\rho}} \mathcal{C}_{r}, \tag{8.73}
\end{equation*}
$$

which identifies the left-hand side with a subspace of codimension $k$ in $T_{\kappa_{\rho}} \mathcal{C}_{r}$. Let $W_{\rho}$ denote the complement of this subspace in $T_{\left(w_{\rho}, \kappa_{\rho}\right)} \mathcal{W}_{2, \epsilon, \rho}$. We show that there exists a constant $C$ such that for $\rho$ large enough and $(\xi, \gamma) \in W_{\rho}$

$$
\begin{equation*}
\|(\xi, \gamma)\| \leq C\left\|d \Gamma_{\rho}(\xi, \gamma)\right\| \tag{8.74}
\end{equation*}
$$

Assume (8.74) is not true, then there exists a sequence $\left(\xi_{N}, \gamma_{N}\right) \in$ $W_{\rho(N)}$, where $\rho(N) \rightarrow \infty$ as $N \rightarrow \infty$ with

$$
\begin{align*}
& \left\|\left(\xi_{N}, \gamma_{N}\right)\right\|=1,  \tag{8.75}\\
& \left\|d \Gamma_{\rho(N)}\left(\xi_{N}, \gamma_{N}\right)\right\| \rightarrow 0, \quad \text { as } N \rightarrow \infty . \tag{8.76}
\end{align*}
$$

Let $\beta_{\rho}^{0}: D_{r}(\rho) \rightarrow \mathbb{C}$ be a cut-off function which equals 1 on $D_{m+1} \backslash$ $\left(\bigcup_{j} E_{p_{j}}\left[-\frac{1}{2} \rho\right]\right)$, equals 0 outside $D_{m+1} \backslash\left(\bigcup_{j} E_{p_{j}}\left[-\frac{3}{4} \rho\right]\right)$, with $\left|D^{k} \beta_{\rho}^{0}\right|=$
$\mathcal{O}\left(\rho^{-1}\right), k=1,2$. By the uniform invertibility of $d \Gamma_{\left(u, \kappa_{1}\right)}$ we find

$$
\begin{align*}
& \left\|\beta_{\rho(j)}^{0}\left(\xi_{N}, \gamma_{N}\right)\right\|_{2, \epsilon, \rho}  \tag{8.77}\\
& \leq C\left\|d \Gamma_{\left(u, \kappa_{1}\right)} \beta_{\rho(N)}^{0}\left(\xi_{N}, \Gamma_{N}\right)\right\|_{1, \epsilon, \rho} \\
& \leq C\left(\left\|\left(\bar{\partial} \beta_{\rho(N)}^{0}\right) \xi_{N}\right\|_{1, \epsilon, \rho}+\left\|\beta_{\rho(N)}^{0} d \Gamma_{\rho}\left(\xi_{N}, \gamma_{N}\right)\right\|_{1, \epsilon, \rho}\right)
\end{align*}
$$

Both terms on the right-hand side go to 0 as $N \rightarrow \infty$. Hence

$$
\begin{equation*}
\left\|\beta_{\rho(N)}^{0}\left(\xi_{N}, \gamma_{N}\right)\right\|_{2, \epsilon, \rho} \rightarrow 0, \quad \text { as } N \rightarrow \infty \tag{8.78}
\end{equation*}
$$

Similarly, with $\beta_{\rho}^{j}: D_{r}(\rho) \rightarrow \mathbb{C}$ a cut-off function which equals 1 on $D_{l(j)+1} \backslash E_{q_{j}}\left[\frac{1}{2} \rho\right]$, equals 0 outside $D_{l(j)+1} \backslash E_{q_{j}}\left[\frac{3}{4} \rho\right]$, with $\left|D^{k} \beta_{\rho}^{0}\right|=$ $\mathcal{O}\left(\rho^{-1}\right), k=1,2$, we find, by the uniform invertibility of $d \Gamma_{\left(v_{j}, \kappa_{2}(j)\right)}$ that

$$
\begin{equation*}
\left\|\beta_{\rho(N)}^{j}\left(\xi_{N}, \gamma_{N}\right)\right\|_{2, \epsilon, \rho} \rightarrow 0, \quad \text { as } N \rightarrow \infty \text { for all } j \tag{8.79}
\end{equation*}
$$

For $1 \leq j \leq k$, we consider the region

$$
\begin{align*}
& \Theta_{j}(\rho)  \tag{8.80}\\
&=\left(E_{p_{j}} \backslash E_{p_{j}}[-(1+a) \rho]\right) \cup_{\left(\left(\phi_{\rho}\right)^{-1} \circ u^{1}\right)} \Omega_{j}  \tag{8.81}\\
& \cup_{\left(\left(\phi_{\rho}\right)^{-1} \circ v_{j}^{1}\right)}\left(E_{q_{j}} \backslash E_{q_{j}}[(1+a) \rho]\right)
\end{align*}
$$

Note that there is a natural inclusion $\Theta_{j}(\rho) \subset M_{\rho}$, where $M_{\rho}$ is as in Lemma 8.14. Also note that the boundary conditions of the linearized equation over $\Omega_{j}(\rho)$ splits into a 1-dimensional problem corresponding to the first coordinate and an $(n-1)$-dimensional problem with boundary conditions converging to two transverse Lagrangian subspaces in the remaining $(n-1)$ coordinate directions.

Let $\alpha_{\rho}^{+}$be a cut-off function on $\Theta_{j}(\rho)$ which equals 1 on

$$
\begin{equation*}
E_{p_{j}}\left[-\frac{1}{4} \rho\right] \backslash E_{p_{j}}\left[-\left(1+\frac{1}{2} a\right) \rho\right] \tag{8.82}
\end{equation*}
$$

equals 0 outside

$$
\begin{equation*}
E_{p_{j}}\left[-\frac{1}{8} \rho\right] \backslash E_{p_{j}}\left[-\left(1+\frac{2}{3} a\right) \rho\right] \tag{8.83}
\end{equation*}
$$

with $\left|D^{k} \alpha^{+}\right|=\mathcal{O}\left(\rho^{-1}\right), k=1,2$, and which is real valued and holomorphic on the boundary. Note that over the region where $\alpha^{+}$is supported, the boundary conditions of $w_{\rho}$ agrees with those of $u$. Thus the angle between the line giving the boundary conditions of $w_{\rho}$ and the real line is $\mathcal{O}\left(\rho^{-1}\right)$ and it is easy to construct a unitary diagonal matrix function $\mathbf{M}$ on the support $\alpha^{+}$with $\left|D^{k} \mathbf{M}\right|=\mathcal{O}\left(\rho^{-1}\right), k=1,2$ with the property
that $\mathbf{M} \xi_{N}$ has the boundary conditions of $w_{\rho}$ in the last $(n-1)$ coordinates and has real boundary conditions in the first coordinate. Thus Lemma 8.14 implies that

$$
\begin{equation*}
\left\|\alpha^{+} \xi_{N}\right\|_{2, \rho, \epsilon} \leq C\left\|\mathbf{M} \alpha^{+} \xi_{N}\right\| \leq C\left(\left\|\left(\bar{\partial} \alpha^{+} \mathbf{M}\right) \xi\right\|_{1, \epsilon, \rho}+\left\|\mathbf{M} \bar{\partial} \xi_{N}\right\|_{1, \epsilon, \rho}\right) . \tag{8.84}
\end{equation*}
$$

Here both terms in the right-hand side goes to 0 as $N \rightarrow \infty$.
In exactly the same way, we show that

$$
\begin{equation*}
\left\|\alpha^{-} \xi_{N}\right\| \rightarrow 0 \quad \text { as } \rho \rightarrow \infty \tag{8.85}
\end{equation*}
$$

for a cut-off function $\alpha^{-}$with support on the other end of $\Theta_{\rho}$.
Let $\alpha^{0}$ be a cut-off function which equals 1 on $[-A(\rho)+2, A(\rho)-2] \times$ $[0,1]$ and equals 0 outside $[-A(\rho)+1, A(\rho)-1] \times[0,1]$ and with the usual properties. Then the function

$$
\begin{equation*}
(\tau+i t) \mapsto\left(d \Phi_{R(N)}^{t s(N)}\left(\phi_{\rho}(\tau+i t)\right)\right)^{-1} \cdot \alpha^{0}(\tau+i t) \xi_{N}(\tau+i t) \tag{8.86}
\end{equation*}
$$

has the boundary conditions of $w_{\rho}$ in the last $(n-1)$ coordinates (two transverse Lagrangian subspaces in this region) and has real boundary conditions in the first coordinate.

Lemma 8.14 implies

$$
\begin{align*}
\left\|\alpha^{0} \xi_{N}\right\|_{1, \rho, \epsilon} \leq & C\left\|\left(d \Phi_{R}^{t s(N)}(N)\right)^{-1} \cdot \alpha^{0} \xi_{N}\right\|_{1, \rho, \epsilon}  \tag{8.87}\\
\leq & C\left(\| \bar{\partial}\left(\alpha^{0} d \Phi_{R}^{t s(N)}(N)\right)^{-1}\right) \cdot \xi_{N} \|_{0, \rho, \epsilon} \\
& \left.\left.+\|\left(\alpha^{0} d \Phi_{R}^{t s(N)}(N)\right)^{-1}\right) \cdot \bar{\partial} \xi_{N} \|_{0, \rho, \epsilon}\right) .
\end{align*}
$$

Using (8.46) and (8.47) in combination with (8.84) and (8.85), we find that the first term in (8.87) goes to 0 as $N \rightarrow 0$. By (8.76), so does the second. Hence

$$
\begin{equation*}
\left\|\alpha^{0} \xi_{N}\right\|_{1, \rho, \epsilon} \rightarrow 0 \tag{8.88}
\end{equation*}
$$

Applying the same argument to $\partial_{\tau} \xi_{N}$ and $i \partial_{t} \xi_{N}$, we conclude that

$$
\begin{equation*}
\left\|\alpha^{0} \xi_{N}\right\|_{2, \rho, \epsilon} \rightarrow 0 \tag{8.89}
\end{equation*}
$$

Now (8.79), (8.78), (8.84), (8.85), and (8.89) contradict (8.75) and we find that (8.74) holds.

To finish the proof, we let $\mu_{j}=\bar{\partial} \frac{\partial \phi^{C_{j}}}{\partial C_{j}}$, see Section 7.9. Then $\mu_{j}$ anti-commutes with $j_{\kappa_{\rho}}$ and we consider the $\mu_{j}$ as newborn conformal variations spanning the complement of $W_{\rho}$ in $T_{\left(w_{\rho}, \kappa_{\rho}\right)} \mathcal{W}_{2, \epsilon, \rho}$.

The images of $\mu_{j}, j=1, \ldots, k$ under $d \Gamma_{\rho}$ are clearly linearly independent since they have mutually disjoint supports. We show that their images stays a uniformly bounded distance away from the subspace $d \Gamma_{\rho}\left(W_{\rho}\right)$. Assume not, then there exists a sequence of elements $\left(\xi_{\rho}, \gamma_{\rho}\right)$ in $W_{\rho}$ with

$$
\begin{equation*}
\left\|d \Gamma_{\rho}\left(\left(\xi_{\rho}, \gamma_{\rho}\right)-\mu_{j}\right)\right\|_{1, \epsilon, \rho} \rightarrow 0 \quad \text { as } \rho \rightarrow \infty \tag{8.90}
\end{equation*}
$$

Since $\left\|d \Gamma_{\rho} \mu_{j}\right\|_{1, \epsilon, \rho}=\mathcal{O}(1)$, we conclude from (8.74) that $\left\|\left(\xi_{\rho}, \gamma_{\rho}\right)\right\|_{2, \epsilon, \rho}=$ $\mathcal{O}(1)$. Then, with the cut-off function $\beta_{\rho}^{j}$ from above and notation as in Section 7.9 we find

$$
\begin{align*}
& \left\|d \Gamma_{\left(v_{j}, \kappa_{2}(j)\right)}\left(\beta_{\rho}^{j}\left(\xi_{\rho}, \gamma_{\rho}\right)-\hat{C}_{j}\right)\right\|_{1, \epsilon}  \tag{8.91}\\
& \quad=\left\|d \Gamma_{\rho}\left(\beta_{\rho}^{j}\left(\xi_{\rho}, \gamma_{\rho}\right)-\mu_{j}\right)\right\|_{1, \epsilon, \rho} \\
& \quad \leq\left\|\beta_{\rho}^{j}\left(d \Gamma_{\rho}\left(\xi_{\rho}, \gamma_{\rho}\right)-\mu_{j}\right)\right\|_{1, \epsilon, \rho}+\left\|\left(\bar{\partial} \beta_{\rho}^{j}\right)\left(\left(\xi_{\rho}, \gamma_{\rho}\right)-\mu_{j}\right)\right\|_{1, \epsilon, \rho} \tag{8.92}
\end{align*}
$$

The right-hand side of the above equation goes to 0 as $\rho \rightarrow \infty$. Hence, so does the left-hand side. This, however, contradicts the invertibility of $d \Gamma_{\left(v_{j}, \kappa_{2}(j)\right)}$ and we conclude $d \Gamma_{\rho}\left(W_{\rho}\right)$ stays a bounded distance away from $d \Gamma_{\rho}\left(\mu_{j}\right)$. Thus, defining $G_{\rho}\left(d \Gamma_{\rho} \mu_{j}\right)=\mu_{j}$ finishes the proof. q.e.d.
8.20. Estimates on the non-linear term. In Section 5.7, we linearized the map $\Gamma$ using local coordinates $B$ around $(w, f) \in \mathcal{W}_{2, \epsilon}$. To apply Floer's Picard lemma, we must study also higher order variations of $\Gamma$.

For $(w, f) \in \mathcal{W}_{2, \epsilon}, w: D_{m} \rightarrow \mathbb{C}^{n}$ and conformal structure $\kappa$ on $D_{m}$, we take as in Section 5.5 local coordinates $(v, \kappa) \in B_{2, \epsilon} \times \mathbb{R}^{m-3}$ on $\mathcal{W}_{2, \epsilon}$ around ( $w, f$ ) and write (in these coordinates)

$$
\Gamma(v \gamma)=\bar{\partial}_{\kappa} v+i \circ \partial_{\kappa} w \circ \gamma+N(v, \gamma) .
$$

We refer to $N(v, \gamma)$ as the non-linear term.
We first consider stationary gluing
Lemma 8.16. There exists a constant $C$ such that the non-linear term $N(v, \gamma)$ of $\Gamma$ in a neighborhood $w_{\rho}$, where $w_{\rho}$ is as in Section 8.5 satisfies

$$
\begin{align*}
& \|N(u, \beta)-N(v, \gamma)\|_{1}  \tag{8.93}\\
& \leq C\left(\|u\|_{2}+|\beta|+\|v\|_{2}+|\gamma|\right)\left(\|u-v\|_{2}+|\beta-\gamma|\right) .
\end{align*}
$$

Proof. With notation as in Section 5.5, we have

$$
\Gamma(v, \gamma)=\bar{\partial}_{\kappa+\gamma}\left(\exp _{w_{\rho}(\zeta)}^{\sigma(\zeta)} v(\zeta)\right)
$$

We prove (8.93) first in the special case $\gamma=\beta=0$. We perform our calculation in coordinates $x+i y$ on $D_{r}(\rho)$, which agree with the standard coordinates on the ends and in the gluing region on $D_{r}(\rho)$. On the remaining parts of the disk, the metric of these coordinates differs from the usual metric by a conformal factor, but since the remaining part is compact, the estimates are unaffected by this change of metric. In these coordinates, we write $\bar{\partial}_{\kappa}=\partial_{x}+i \partial_{y}$. Now, as in Lemma 5.12, we find

$$
\partial_{x} \exp _{w_{\rho}}^{\sigma} v=J\left[w_{\rho}, v, \partial_{x} w_{\rho}, \partial_{x} v, \sigma\right](1)+\partial_{\sigma}\left(\exp _{w_{\rho}}^{\sigma} v\right) \cdot \partial_{x} \sigma
$$

where $J\left[x, \xi, x^{\prime}, \xi^{\prime}, \sigma\right]$ denotes the Jacobi field in the metric $g(\sigma)$ along the geodesic $\exp _{x}^{\sigma} t \xi$ with initial conditions $J(0)=x^{\prime}, J^{\prime}(0)=\xi^{\prime}$. Of course a similar equation holds for $\partial_{y} \exp _{w_{\rho}}^{\sigma} v$.

Let $G:\left(\mathbb{C}^{n}\right)^{4} \times[0,1] \times \mathbb{R} \rightarrow \mathbb{C}^{n}$ be the function

$$
G\left(x, \xi, x^{\prime}, \xi^{\prime}, \sigma, \sigma^{\prime}\right)=J\left[x, \xi, x^{\prime}, \xi^{\prime}, \sigma\right](1)-x^{\prime}-\xi^{\prime}+\partial_{\sigma} \exp _{x}^{\sigma} \xi \cdot \sigma^{\prime}
$$

(unrelated to the earlier right inverses $G_{\rho}$ ) then with $w_{\rho}=w$,

$$
N(v)=G\left(w, v, \partial_{x} w, \partial_{x} v, \sigma, \partial_{x} \sigma\right)+i G\left(w, v, \partial_{y} w, \partial_{y} v, \sigma, \partial_{y} \sigma\right) .
$$

Moreover, the function $G$ is smooth with uniformly bounded derivatives, it is linear in $x^{\prime}, \xi^{\prime}, \sigma^{\prime}$, and satisfies

$$
\begin{align*}
& G\left(x, 0, x^{\prime}, \xi^{\prime}, \sigma, \sigma^{\prime}\right)=0  \tag{8.94}\\
& D_{2} G\left(x, 0, x^{\prime}, \xi^{\prime}, \sigma, \sigma^{\prime}\right)=0
\end{align*}
$$

where the last equation follows from Taylor expansion of the exponential map and the Jacobi field.

With this established, the arguments needed to prove estimates on integral norms in the lemma are similar to those given in the proof of Lemma 5.10 and will be omitted. Finally, we remark that the input of the space of conformal structures is easily controlled since this space is finite dimensional.
q.e.d.

In the self-tangency shortening case, the estimate is somewhat changed since we work in Sobolev spaces with negative exponential weights in the gluing region. Here we have

Lemma 8.17. There exists a constant $C$ such that the non-linear term $N(v, \gamma, \lambda)$ of $\Gamma$ in a neighborhood of $w_{\rho}$, where $w_{\rho}$ is as in Section 8.10 satisfies

$$
\begin{aligned}
& \|N(u, \beta)-N(v, \kappa)\|_{1,-\epsilon, \rho} \\
& \quad \leq C e^{\epsilon \rho}\left(\|u\|_{2,-\epsilon, \rho}+|\beta|+\|v\|_{2,-\epsilon, \rho}+|\gamma|\right)\left(\|u-v\|_{2,-\epsilon, \rho}+|\beta-\gamma|\right) .
\end{aligned}
$$

Proof. The proof is exactly the same as the proof of Lemma 8.16. We must only take into account in what way the weights affect the estimates. Starting with the $L^{2}$-norm, we see that the norm $\|\cdot\|_{2, \rho,-\epsilon}$ does not control the sup-norm uniformly in $\rho$. But it does control $e^{-\epsilon \rho}$ times the sup-norm. Thus we conclude instead of the usual $L^{2}$-estimate that

$$
\begin{equation*}
\|N(u)-N(v)\| \leq C e^{\epsilon \rho}\left(\|u\|_{2,-\epsilon, \rho}+\|v\|_{2,-\epsilon, \rho}\right)\|u-v\|_{2,-\epsilon, \rho} . \tag{8.95}
\end{equation*}
$$

Similarly, we loose this factor in the other estimates where we use the sup-norm. Let $e_{\rho}$ denote the weight function from Section 8.11. When
we use the $L^{4}$-estimate, we have instead the following

$$
\begin{aligned}
& \int_{D_{m}}(|D u|+|D v|)^{2}|D u-D v|^{2} e_{\rho}^{2} d A \\
& \quad \leq e^{2 \epsilon \rho} \int_{D_{m}}(|D u|+|D v|)^{2}|D u-D v|^{2} e_{\rho}^{4} d A \\
& \quad \leq e^{2 \epsilon \rho}\left(\int_{D_{m}}(|D u|+|D v|)^{4} e_{\rho}^{4}\right)^{\frac{1}{2}}\left(\int_{D_{m}}(|D u-D v|)^{4} e_{\rho}^{4}\right)^{\frac{1}{2}} \\
& \quad \leq C e^{2 \epsilon \rho}\left(\|u\|_{2,-\epsilon, \rho}+\|v\|_{2,-\epsilon, \rho}\right)^{2}\left(\|u-v\|_{2,-\epsilon, \rho}\right)^{2} .
\end{aligned}
$$

We conclude finally

$$
\|N(u)-N(v)\|_{1,-\epsilon, \rho} \leq C e^{\epsilon \rho}\left(\|u\|_{2,-\epsilon, \rho}+\|v\|_{2,-\epsilon, \rho}\right)\|u-v\|_{2,-\epsilon, \rho} .
$$

The conformal structures can be handled as in Lemma 8.16. q.e.d.
Finally, we consider self-tangency gluing, where we have a large weight function which does not interfere (destructively) with the sup-norm and the $L^{4}$ estimates.

Lemma 8.18. There exists a constant $C$ such that the non-linear term $N(v, \gamma)$ of $\Gamma$ in a neighborhood of $w_{\rho}$, where $w_{\rho}$ is as Section 8.15 satisfies

$$
\begin{aligned}
& \|N(u, \beta)-N(v, \kappa)\|_{1} \\
& \quad \leq C\left(\|u\|_{2, \epsilon, \rho}+|\beta|+\|v\|_{2, \epsilon, \rho}+|\gamma|\right)\left(\|u-v\|_{2, \epsilon, \rho}+|\beta-\gamma|\right) .
\end{aligned}
$$

Proof. See the proof of Lemma 8.16
q.e.d.

## 9. Gromov compactness

In this section, we prove a version of the Gromov compactness theorem. In Section 9.2, we discuss the compactification of the space of conformal structures which is done in [17]. In Section 9.3, we translate the notions of convergence and (limiting) broken curves from [22] to our setting. There are two notions of convergence we must prove: a strong local convergence and a weak global convergence. In Sections 9.5 and 9.6 , we discuss how to adopt Floer's original approach [13] to prove the strong local convergence. Local convergence implies that our holomorphic disks, away from the punctures, are smooth up to and including the boundaries, see Remark 9.5. To prove the weak global convergence in Section 9.7, we analyze where the area (or energy) of a sequence of disks accumulates, and construct an appropriate sequence of reparameterizations of the domain to recover this area.

We note that although our holomorphic curves map to a non-compact space, $\mathbb{C}^{n}$, the set of curves we consider lives in a compact subset. This follows because $\mathbb{C}^{n}$ is a symplectic manifold with "finite geometry at
infinity": a holomorphic curve with a non-compact image must contain infinite area. And the area of any disk we consider is bounded above by the action of the chords mapped to at its corners. Thus, we can prove the Gromov compactness theorem in this non-compact set-up. For a review of finite geometry at infinity (also known as "tame"), see [1] Chapter 5, as well as $[\mathbf{7}, \mathbf{1 8}, \mathbf{3 1}]$.
9.1. Notation and conventions for this section. Unlike in the other sections, we need to consider Sobolev spaces with derivatives in $L^{p}$ for $p \neq 2$. We define in the obvious way the spaces $W_{k}^{p, \text { loc }}\left(\Delta_{m}, \mathbb{C}^{n}\right)$ to indicate $\mathbb{C}^{n}$-valued functions on $\Delta_{m}$ whose first $k$ derivatives are locally $L^{p}$-integrable. For this section only, we denote the corresponding norm by $\|\cdot\|_{k, p}$.

In order to define broken curves in the next subsection, we will need to extend the disk continuously to the boundary punctures. Of course the extra Legendrian boundary condition, $h$, does not extend continuously. For this reason, we will only extend $u$ to $\bar{\Delta}_{m}$, the closure of $\Delta_{m}$; thus, $u: \bar{\Delta}_{m} \rightarrow \mathbb{C}^{n}$. Note that $\|u\|_{p, k}$ might still blow up at these punctures. We sometimes only consider $u$ and $u \mid \partial \bar{\Delta}_{m}$ in which case we write $u$ : $\left(\bar{\Delta}_{m}, \partial \bar{\Delta}_{m}\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$. For $X \subset \bar{\Delta}_{m}$, let $\|u\|_{k, p: X}=\|u \mid X\|_{k, p}$, and $\|u\|_{k, p: \epsilon}$ denote the norm restricted to some disk (or half-disk) of radius $\epsilon$.

Because we sometimes change the number of boundary punctures, we will denote by $D$ the unit disk in $\mathbb{C}^{n}$.
9.2. Compactification of space of conformal structures. Recall $\mathcal{C}_{m}$ is the space of conformal structures (modulo conformal reparameterizations) on the unit disk in $\mathbb{C}$ with $m$ boundary punctures.

When $m \geq 3$, we define a stable cusp disk representative with $m$ marked boundary points, $\left(\Sigma ; p_{1}, \ldots, p_{m}\right)$, to be a connected, simplyconnected union of unit disks in $\mathbb{C}$ where pairs of disks may overlap at isolated boundary points (which we call double points of $\Sigma$ ) and each disk in $\Sigma$ has at least 3 points, called marked points, which correspond either to double points or the original boundary marked points. When $m=1$ or 2 , the stable cusp disk representative shall be a single disk. Two stable cusp disk representatives are equivalent if there exist a conformal reparameterization of the disks taking one set of marked points to the other. We define a stable (cusp) disk with $m$ marked points to be an equivalence class of stable disk representatives with $m$ marked points.

In Section 10 of $[\mathbf{1 7}]$, Fukaya and Oh prove that $\overline{\mathcal{C}}_{m}$, the compactification of $\mathcal{C}_{m}$, is the space of stable disks with $m$ marked points.
9.3. The statement. A broken curve $(u, h)=\left(\left(u^{1}, h^{1}\right), \ldots\left(u^{N}, h^{N}\right)\right)$ is a connected union of holomorphic disks, $\left(u^{j}, h^{j}\right)$, (recall $u^{j}$ is extended to $\bar{\Delta}_{m_{j}}$ ) where each $u^{j}$ has exactly one positive puncture and except for
one disk, say $u^{1}$, the positive puncture of $u^{j}$ agrees with the negative puncture of some other $u^{j^{\prime}}$. One may easily check that a broken curve can be parameterized by a single smooth $v:\left(D_{m}, \partial D\right) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$, such that $v^{-1}$ is finite except at points where two punctures were identified, here $v^{-1}$ is an arc in $\Delta_{m}$.

Definition 9.1. A sequence of holomorphic disks ( $u_{\alpha}, h_{\alpha}$ ) converges to a broken curve $(u, h)=\left(\left(u^{1}, h^{1}\right), \ldots,\left(u^{N}, h^{N}\right)\right)$ if the following holds

1) (Strong local convergence) For every $j \leq N$, there exists a sequence $\phi_{\alpha}^{j}: D \rightarrow D$ of linear fractional transformations and a finite set $X^{j} \subset D$ such that $u_{\alpha} \circ \phi_{\alpha}^{j}$ converges to $u^{j}$ uniformly with all derivatives on compact subsets of $D \backslash X^{j}$
2) (Weak global convergence) There exists a sequence of orientationpreserving diffeomorphisms $f_{\alpha}: D \rightarrow D$ such that $u_{\alpha} \circ f_{\alpha}$ converges in the $C^{0}$-topology to a parameterization of $u$.

Henceforth, to simplify notation when passing to a subsequence, we will not change the indexing.

Theorem 9.2. Assume $\left(u_{\alpha}, h_{\alpha}\right) \in \mathcal{M}\left(a ; b_{1}, \ldots, b_{m}\right)$ is a sequence of holomorphic disks with $L_{\alpha}$ Legendrian boundary condition. Let $\kappa_{\alpha} \in$ $\mathcal{C}_{m+1}$ denote the conformal structure on the domain of $u_{\alpha}$. Assume $L_{\alpha}$ converges to an embedded Legendrian $L$ in the $C^{\infty}$-topology. Then, there exists a subsequence $\left(u_{\alpha}, h_{\alpha}, \kappa_{\alpha}\right)$ such that $\kappa_{\alpha}$ converges to $\kappa \in \overline{\mathcal{C}}_{m+1}$ and $\left(u_{\alpha}, h_{\alpha}\right)$ converges to a broken curve $(u, h)$ whose domain is a stable disk representative of $\kappa$.

Note that using the strong local convergence property a posteriori, this compactness result proves that all derivatives of a holomorphic disk $(u, h)$ are locally integrable away from the finite set of points. In particular, such disks are smooth at the boundary away from these points. See Remark 9.5.

We also remark that, the appropriately modified, Theorem 9.2 holds if the disks have more than one positive puncture.
9.4. Area of a disk. For holomorphic $u:(D, \partial D) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)$, recall that $\operatorname{Area}(u)=\int u^{*} \omega$, where $\omega=\sum_{i} d x_{i} \wedge d y_{i}$, denotes its (signed) area.

Lemma 9.3. Consider an admissible Legendrian isotopy parameterized by $\lambda \in \Lambda$. We assume $\Lambda \subset \mathbb{R}$ is compact. Denote by $L_{\lambda}$ the moving Legendrian submanifold. There exists a positive upper semi-continuous function $\hbar: \Lambda \rightarrow \mathbb{R}^{+}$such that for any non-constant holomorphic map $u:(D, \partial D) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}\left(L_{\lambda}\right)\right)$, Area $(u) \geq \hbar(\lambda)$.

Proof. We need the following statement from Proposition 4.3.1 (ii) of Sikorav in [1]: There are constants $r_{1}, k$ (depending only on $\mathbb{C}^{n}$ ) such that if $r \in\left(0, r_{1}\right]$ and $u: \Sigma \rightarrow B(x, r)$ is a holomorphic map of a

Riemann surface containing $x$ in its image and with $u(\partial \Sigma) \subset \partial B(x, r)$, then $\operatorname{Area}(\Sigma) \geq k r^{2}$.

Since $u$ is non-constant, Stokes Theorem implies $u$ must have boundary punctures. Choose $r>0$, an upper semi-continuous function of $\lambda$, such that:

- for all Reeb chords $c, \Pi_{\mathbb{C}}\left(L_{\lambda}\right) \cap B\left(c^{*}, r\right)$ is real analytic and diffeomorphic either to $\mathbb{R}^{n} \times\{0\} \cup\{0\} \times \mathbb{R}^{n}$ or the local picture of the singular moment in a standard self-tangency move (see Definition 3.3).
- for all distinct Reeb chords $c_{1}, c_{2}, B\left(c_{1}^{*}, r\right) \cap B\left(c_{2}^{*}, r\right)=\emptyset$ and
- $r<r_{1}$.

Let $\theta_{\lambda}$ be the smallest angle among all the complex angles associated to all the tranverse double points of $\Pi_{\mathbb{C}}\left(L_{\lambda}\right)$. Now set

$$
\hbar(\lambda)=\min \left\{\min _{c \in \mathcal{C}\left(L_{\lambda}\right)} \mathcal{Z}(c), \frac{k r^{2} \cos ^{2} \theta_{\lambda}}{8}\right\}>0 .
$$

Suppose $u$ maps all $n$ of its punctures to the same double point $c^{*}$, then by (2.4)

$$
\operatorname{Area}(u) \geq \mathcal{Z}(c) \geq \hbar
$$

(Note the number of positive punctures of $u$ must be larger than the number of negative ones since $u$ is not constant.)

Otherwise, assume $u$ maps boundary punctures to at least two distinct double points $c_{1}^{*}, c_{2}^{*}$ where $c_{1}^{*}$ is a non-degenerate double point. Then, $c_{2}^{*} \notin \bar{B}\left(c_{1}^{*}, r\right)$ implies that there exists a point $x \in u(D) \cap$ $\Pi_{\mathbb{C}}(L) \cap \partial B\left(c_{1}^{*}, \frac{r}{2}\right)$. Moreover, $B\left(x, \frac{r \cos \theta_{\lambda}}{2}\right) \subset B(x, r)$ intersects $\Pi_{\mathbb{C}}(L)$ in only one sheet. Using the real-analyticity of the boundary, we double $u(D) \cap B\left(x, \frac{r \cos \theta_{\lambda}}{2}\right)$ and apply the proposition of Sikorav to conclude

$$
\operatorname{Area}(u) \geq \operatorname{Area}\left(u(D) \cap B\left(x, \frac{r \cos \theta_{\lambda}}{2}\right)\right) \geq \frac{k r^{2} \cos ^{2} \theta_{\lambda}}{8} \geq \hbar
$$

q.e.d.

We introduce one more area-related notion, again borrowed from [22]. Given a sequence of holomorphic maps $u_{\alpha}$, we say $z \in D$ is a point mass of $\left\{u_{\alpha}\right\}$ with mass $m$ if there exists a sequence $z_{\alpha} \in D$ converging to $z \in D$ such that

$$
\lim _{\epsilon \rightarrow 0} \lim _{\alpha \rightarrow \infty} \operatorname{Area}\left(u \mid B_{\epsilon}\left(z_{\alpha}\right) \cap D\right)=m .
$$

9.5. Strong local convergence I: bootstrapping. In this subsection, we formulate the following "bootstrap" elliptic estimate: if we know a holomorphic curve lies locally in $W_{k}^{p}$ with $p>2, k \geq 1$, then the $\|\cdot\|_{p^{\prime}, k^{-}}$(local)-norm controls the $\|\cdot\|_{p^{\prime}, k+1^{-}}$(local)-norm for $p^{\prime} \in[2, p)$.

Let $A \subset \mathbb{C}$ denote the open disk or half-disk with boundary on the real line. Let $W_{k}^{p}\left(A, \mathbb{C}^{n}\right)$ denote the closure, under the $\|\cdot\|_{k, p}$-norm, of the set of all smooth compactly supported functions from $A$ to $\mathbb{C}^{n}$.

Theorem 9.4. Fix $k \geq 1$ and (not necessarily small) $\delta_{k-1}>\delta_{k} \geq 0$. For any compact $K \subset A$, there exists a "constant" $C_{1}=C_{1}\left(\|u\|_{k, 2+\delta_{k-1}}\right)$ depending continuously on $\|u\|_{k, 2+\delta_{k-1}}$ such that for all holomorphic maps $u \in W_{k}^{2+\delta_{k-1}}\left(A, \mathbb{C}^{n}\right)$ with $u(\partial A) \subset \Pi_{\mathbb{C}}(L)$, we have

$$
\begin{equation*}
\|u\|_{k+1,2+\delta_{k}: K} \leq C_{1}\|u\|_{k, 2+\delta_{k}: A} \tag{9.1}
\end{equation*}
$$

Moreover, if $u_{\alpha}$ is a sequence of holomorphic maps in $W_{k}^{2+\delta_{k-1}}\left(A, \mathbb{C}^{n}\right)$ such that $u_{\alpha}(\partial A) \subset \Pi_{\mathbb{C}}(L)$ and $\left\|u_{\alpha}\right\|_{k, 2+\delta_{k-1}}$ is uniformly bounded, then there exists a subsequence $u_{\alpha}$ converging in $W_{k}^{2+\delta_{k}}\left(K, \mathbb{C}^{n}\right)$ to some holomorphic map $u: K \rightarrow \mathbb{C}^{n}$.

Remark 9.5. Note how we can use the Sobolev embedding theorem to conclude that all derivatives of the curve lie in $L^{2}$ locally, assuming we have a finite local $\|\cdot\|_{1,2+\delta_{0}}$ norm to begin with. In particular, a holomorphic disk $(h, u)$ with boundary punctures becomes smooth at the boundary away from the punctures. We did not have to assume this smoothness a priori.

This proof first appeared as Lemma 2.3 in [13] and later corrected as Proposition 3.1 [24]. Floer and Oh both prove the $k=1$ case and state the general case. Since there are no new techniques here, we omit the proof for the general case. Instead, we simply formulate a key lemma for the proof.

Lemma 9.6. For every $l>k, l-2 / q>k-2 / p$, there exists $a$ constant $C$ such that if $\xi \in W_{k}^{p}\left(A, \mathbb{C}^{n}\right)$ is compactly supported, $\xi \mid \partial A \subset$ $\mathbb{R}^{n}$, and $\bar{\partial} \xi \in W_{l-1}^{q}\left(A, \mathbb{C}^{n}\right)$ then

$$
\begin{equation*}
\|\xi\|_{l, q} \leq C\|\bar{\partial} \xi\|_{l-1, q} . \tag{9.2}
\end{equation*}
$$

This is stated as Lemma 2.2 of [13] and Lemma 3.2 of [24]. Floer attributes this result to Theorem 20.1.2 of [20]. However, we were unable to deduce Lemma 9.6 for $k>1$ from Hörmander's theorem. Alternatively, one can use the Seeley extension theorem (see [23], section 1.4 for example) to extend the map to the full disk (in the case of the half disk) and then use the well-known full disk version of Lemma 9.6.
9.6. Strong local convergence II: uniformly bounding higher

Sobolev norms. In order to apply Theorem 9.4, we need a uniform bound on the $\|\cdot\|_{k, 2+\delta}$-norm where $\delta>0$ might be large. Our holomorphic disk only come with a bound on the $\|\cdot\|_{1,2}$-norm in terms of the action. In this subsection, we indicate how the latter norm controls the former.

Theorem 9.7. Consider the sequence of holomorphic disks $\left(u_{\alpha}, h_{\alpha}\right) \in$ $\mathcal{M}\left(a ; b_{1}, \ldots, b_{m}\right)$. There exists a finite number of points $z_{1}, \ldots, z_{l} \in$ $\partial \Delta_{m}$ and a "constant" $C_{11}=C_{11}(K, p, k)$ such that for any positive
integer $k$, for any $p \in \mathbb{R}$ with $k>\frac{2}{p}$, and for any compact set $K \subset$ $\Delta_{m} \backslash\left\{z_{1}, \ldots, z_{l}\right\}$,

$$
\left\|D^{k} u_{\alpha}\right\|_{0, p: K} \leq C_{11} .
$$

Proof. Theorem 2 in [13], and later Proposition 3.3 in [24], prove this result when $m=2, k=1$ and the boundary conditions are two embedded Lagrangians instead of one immersed Lagrangian. So for the sake of brevity, we only sketch the ideas.

Consider a small ball $B_{\alpha}$ centered at $z_{\alpha}$ where $\left\|D^{k} u_{\alpha}\right\|_{0, p: B_{\alpha}}$ is unbounded. Zoom in by changing coordinates $z \mapsto \frac{z-z_{\alpha}}{\epsilon_{\alpha}}$ where $\epsilon_{\alpha}$ converges to 0 at some appropriate rate. We remark that when rescaling variables $(s, t) \mapsto(\beta s, \beta t)$, the $(1,2)$-norm is conformally invariant, whereas the ( $k, p$ )-norm changes like

$$
\begin{equation*}
\left\|D^{k} f\right\|_{p} \rightarrow\left(\left(\beta^{-k}\right)^{p}\left(\beta^{2}\right)\right)^{\frac{1}{p}}\left\|D^{k} f\right\|_{p}=\beta^{-k+\frac{2}{p}}\left\|D^{k} f\right\|_{p} \tag{9.3}
\end{equation*}
$$

thus, the cases $(k, p), k>\frac{2}{p}$ and $(1, p), p>2$ from [13] are identical.
If $z_{\alpha}$ converges to an interior point, one can construct, by zooming in, a holomorphic sphere which contradicts $\int_{S^{2}} \omega=0$. If $z_{\alpha}$ converges to a boundary point $z_{1}$, one can construct, by zooming in, a holomorphic disk

$$
w:(D, \partial D) \rightarrow\left(\mathbb{C}^{n}, \Pi_{\mathbb{C}}(L)\right)
$$

for which by Lemma $9.3,\|w\|_{0,2} \geq \hbar$.
The limit point $z_{1}$ is an example of a point mass for the sequence $u_{\alpha}$. We repeat this for another $B_{\alpha}^{\prime}$, where $B_{\alpha}^{\prime} \cap B_{\alpha}=\emptyset$ for large $\alpha$ and $\left\|D^{k} u_{\alpha}\right\|_{0, p: B_{\alpha}^{\prime}}$ is unbounded. This produces a separate point mass $z_{2}$. Since the area of $u_{\alpha}$ divided by $\hbar$ is bounded (uniformly) from above, this process can be repeated only a finite number of times. q.e.d.
9.7. Recovering the bubbles. The goal of this subsection is to construct a (not necessarily conformal) reparameterization of $\bar{\Delta}_{m}$ which recovers all disks which bubble off. This reparameterization implies the second convergence in Definition 9.1.

Consider a sequence ( $u_{\alpha}, h_{\alpha}$ ) which converges strongly on any compact $K \subset \Delta_{m} \backslash\left\{z_{1}, \ldots, z_{l}\right\}$. By the proof of Theorem 9.7, we can assume that $z_{1}$ is a point mass with mass $m_{1}>0$.

Let $\mathbb{C}_{+} \subset \mathbb{C}$ denote the upper-half plane. Let $B_{r}=\left\{z \in \mathbb{C}_{+}:\|z\|<\right.$ $r\}$ and $C_{r}=\partial B_{r}$. Define the conformal map

$$
\psi_{\alpha}: \mathbb{C}_{+} \rightarrow \bar{\Delta}_{m}, \quad \psi_{\alpha}(z)=\frac{-z+i R_{\alpha}^{2}}{z+i R_{\alpha}^{2}} \cdot z_{1}
$$

where $R_{\alpha} \in \mathbb{R}$ is such that

$$
\operatorname{Area}\left(u_{\alpha} \mid \psi_{\alpha}\left(B_{R_{\alpha}}\right)\right)=m_{1} \text {. }
$$

Pass to a subsequence and assume $\alpha<\alpha^{\prime}$ implies $R_{\alpha}<R_{\alpha^{\prime}}$, which can be done since by the definition of point mass, $\lim _{\alpha \rightarrow \infty} R_{\alpha}=\infty$. Note that

$$
\begin{equation*}
\lim _{\alpha \rightarrow \infty} \psi_{\alpha}\left(B_{R_{\alpha}}\right)=\lim _{\alpha \rightarrow \infty} \psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}}\right)=z_{1} \tag{9.4}
\end{equation*}
$$

Assume $\alpha$ is large enough so that $\psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}}\right)$ contains no other point masses of the sequence $u_{\alpha}$. However, $\psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}}\right)$ might contain boundary punctures.

After passing to a subsequence, we can use Theorems 9.4 and 9.7 to assume that $u_{\alpha}$ converges to some $u$ on any compact set in $\Delta_{m} \backslash$ $\left(\left\{z_{2}, \ldots, z_{l}\right\} \cup \psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}}\right)\right)$.

The definition of $R_{\alpha}$ and (9.4) imply

$$
\begin{equation*}
\lim _{\alpha \rightarrow \infty} \operatorname{Area}\left(u_{\alpha} \mid \psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}} \backslash B_{R_{\alpha}}\right)\right)=0 \tag{9.5}
\end{equation*}
$$

Use (9.5) and argue as in the previous subsection to find some half circle $C_{R_{\alpha}^{\prime}} \subset \mathbb{C}_{+}$, with $R_{\alpha}^{\prime} \in\left(R_{\alpha}^{3 / 2}-1, R_{\alpha}^{3 / 2}\right]$ such that

$$
\left\|u_{\alpha} \circ \psi_{\alpha}\right\|_{C^{0}: C_{R_{\alpha}^{\prime}}} \rightarrow 0
$$

Define the center of mass of $u_{\alpha} \circ \psi_{\alpha}$ to be

$$
z_{\alpha}=x_{\alpha}+i y_{\alpha}=\frac{1}{m_{1}} \int_{B_{R_{\alpha}}}\left|D\left(u_{\alpha} \circ \psi_{\alpha}\right)\right|^{2}(x+i y) d x \wedge d y \in B_{R_{\alpha}}
$$

where $x+i y$ are coordinates on $\mathbb{C}_{+}$. Define the conformal map $\phi_{\alpha}$ which sends $i$ to $z_{\alpha}$ :

$$
\phi_{\alpha}: \mathbb{C}_{+} \rightarrow \mathbb{C}_{+}, \quad \phi_{\alpha}(z)=y_{\alpha} z+x_{\alpha}
$$

Note that although $\phi_{\alpha}^{-1}\left(C_{R_{\alpha}}\right)$ might remain bounded, $\phi_{\alpha}^{-1}\left(C_{R_{\alpha}^{\prime}}\right)$ converges to $\infty$ because

$$
\left|\phi_{\alpha}^{-1}\left(R_{\alpha}^{\prime} e^{i \theta}\right)\right| \geq \frac{\left(R_{\alpha}^{3 / 2}-1\right)}{\left|y_{\alpha}\right|} \max \{|\cos \theta|,|\sin \theta|\}
$$

and $y_{\alpha}<R_{\alpha}$.
Define the conformal map

$$
\Psi: D \rightarrow \mathbb{C}_{+}, \quad \Psi(z)=\frac{z-1}{i z+i}
$$

where $D \subset \mathbb{C}$ is the unit disk. Note that $\Psi^{-1} \phi_{\alpha}^{-1}\left(C_{R_{\alpha}^{\prime}}\right) \rightarrow-1$ and that $u_{\alpha} \circ \psi_{\alpha} \circ \phi_{\alpha} \circ \Psi$ all have center of mass at $0 \in D$. (Recall that the center of mass uses the Euclidean metric on $\mathbb{C}_{+}$, not on $D$.)

Since

$$
\left\|u_{\alpha} \circ \psi_{\alpha} \circ \phi_{\alpha} \circ \Psi\right\|_{C^{0}: \Psi^{-1} \circ \phi_{\alpha}^{-1}\left(C_{R_{\alpha}^{\prime}}\right)} \rightarrow 0
$$

pass to a subsequence as before and conclude that $u_{\alpha} \circ \psi_{\alpha} \circ \phi_{\alpha} \circ \Psi$ converges to some holomorphic $w$ on compact sets outside of some boundary
point masses and punctures, as well as -1 (since $u_{\alpha} \circ \psi_{\alpha} \circ \phi_{\alpha} \circ \Psi$ is not defined at -1 ).

As before, $w$ can be continuously extended to -1 . We claim that under this reparameterization, -1 is not a point mass of $u_{\alpha} \circ \psi_{\alpha} \circ \phi_{\alpha} \circ \Psi$. Otherwise, in the $\mathbb{C}_{+}$set-up, as some mass escaped to $\infty$, the center of mass would have to go to $\infty$ as well, contradicting the fact that it is fixed at $i \in \mathbb{C}_{+}$.

Because $u_{\alpha}$ converges to $u$ outside of $\psi_{\alpha}\left(B_{R_{\alpha}^{3 / 2}}\right)$, and because no area is "unaccounted" for by (9.5), we can continuously extend $u$ to $z_{1}$ so that $u\left(z_{1}\right)=w(-1)$. Considering how $u$ and $w$ were obtained from $u_{\alpha}$, it is easy to see that the sign of the punctures ( $z_{1}$ for $u$ and -1 for $w$ ) will be opposite. Thus since each of $u$ and $w$ must have a positive puncture, each will have exactly one. Repeat the above argument at all the other point masses $z_{j}$. Then repeat for any new point masses in the sequences defining the holomorphic disks $w_{j}$ associated to $z_{j}$. Continuing until all point masses have been dealt with, we see no holomorphic curves were overlooked in the reparameterization.
9.8. Proof of Theorem 9.2. Let $\Pi_{\mathbb{C}}(L)$ denote the limiting Lagrangian boundary condition. Let $\hbar=\hbar\left(\Pi_{\mathbb{C}}(L)\right)$ be the minimal area of non-constant maps defined in Section 9.4. Use the discussion in Section 9.2 to pass to a subsequence whose conformal structures converge to a stable disk.

We wish to apply Theorem 9.4 to derive strong local convergence. To achieve the required uniform bound on $\left\|u_{\alpha}\right\|_{k, 2+\delta_{k-1}: K}$ for some compact set $K \subset \Delta_{m}$ which lies away from point masses, we apply Theorem $9.7 k$ times to bound $\left\|u_{\alpha}\right\|_{i, 2+\delta_{k-1}: K}$ for $i=1, \ldots, k$. The reparameterizations $\phi_{\alpha}^{j}$ in Definition 9.1 come from the discussion in Section 9.7.

The weak global convergence follows readily from Section 9.7. q.e.d.

## 10. Handle slides

In this section, we will prove Lemma 2.12 which was used to show that the stable tame isomorphism type of the contact homology algebra associated to a Legendrian submanifold does not change under handle slide instances in generic 1-parameter families of Legendrian subamnifolds.
10.1. An auxiliary Legendrian submanifold. We associate to a 1parameter family of Legendrian embeddings $\phi_{t}: L \rightarrow \mathbb{C}^{n} \times \mathbb{R}, 0 \leq t \leq 1$, Legendrian embeddings $\Phi_{f}^{\delta}: L \times \mathbb{R} \rightarrow \mathbb{C}^{n+1} \times \mathbb{R}$, depending on $\delta>0$ and a positive Morse function $f: \mathbb{R} \rightarrow \mathbb{R}$.

Let $\phi_{t}: L \subset \mathbb{C}^{n} \times \mathbb{R}, t \in[-1,1]$ be a Legendrian isotopy. For small $\delta>0$, fix smooth non-decreasing functions

$$
\begin{equation*}
\alpha^{\delta}:[-1,1] \rightarrow[-\delta, \delta] \tag{10.1}
\end{equation*}
$$

such that $\alpha^{\delta}( \pm t)= \pm \delta$ for $\frac{3}{4} \leq t \leq 1$, and such that $\alpha^{\delta}(t)=\delta t$ for $-\frac{1}{4} \leq t \leq \frac{1}{4}$. Note that $\alpha^{\delta} \rightarrow 0$ as $\delta \rightarrow 0$.

Fix standard coordinates

$$
\left(\left(x_{1}, y_{1}, \ldots, x_{n}, y_{n}\right), z\right)=(x, y, z)
$$

on $\mathbb{C}^{n} \times \mathbb{R}$. Define $\phi_{t}^{\delta}, t \in \mathbb{R}$ as

$$
\phi_{t}^{\delta}= \begin{cases}\phi_{-\delta} & \text { for } t \in(-\infty,-1] \\ \phi_{\alpha^{\delta}(t)} & \text { for } t \in[-1,1] \\ \phi_{\delta} & \text { for } t \in[1, \infty)\end{cases}
$$

Write

$$
\phi_{t}^{\delta}(q)=\left(x_{t}(q), y_{t}(q), z_{t}(q)\right), \quad q \in L .
$$

Fix a positive Morse function $f: \mathbb{R} \rightarrow \mathbb{R}$ and $\delta>0$. Let $f^{\prime}(t)=\frac{d f}{d t}$ denote the derivative of $f$. Define $\Phi_{f}^{\delta}: \mathbb{R} \times L \rightarrow \mathbb{C} \times \mathbb{C}^{n} \times \mathbb{R}$,

$$
\Phi_{f}^{\delta}(t, q)=\left(x_{0}(t, q), y_{0}(t, q), x(t, q), y(t, q), z(t, q)\right), \quad(t, q) \in \mathbb{R} \times L,
$$

where

$$
\begin{aligned}
x_{0}(q, t) & =t \\
y_{0}(q, t) & =f(t)\left(\frac{\partial z_{t}}{\partial t}+y_{j_{t}}(q) \frac{\partial x_{j_{t}}}{\partial q}\right)+f^{\prime}(t) z_{t}(q) \\
x(q, t) & =x_{t}(q) \\
y(q, t) & =f(t) y_{t}(q) \\
z(q, t) & =f(t) z_{t}(q)
\end{aligned}
$$

It is straightforward to check that $\Phi_{f}^{\delta}$ is a Legendrian embedding.
Assume that the Morse function $f: \mathbb{R} \rightarrow \mathbb{R}$ above has local minima at $\pm 1$ and no critical points in the region $(-\infty,-1) \cup(1, \infty)$. Then the $x_{0}$-coordinate $c_{0}$ of each Reeb chord $c$ of $\Phi$ satisfies $\left|c_{0}\right| \leq 1$.

Let $u$ be a holomorphic disk with boundary on $\Phi_{f}^{\delta}$ and one positive puncture.

Lemma 10.1. If the positive puncture of $u$ maps to a Reeb chord $c$ of $\Phi$ with $c_{0}= \pm 1$. Then the image of $u$ lies in $\left\{x_{0}= \pm 1\right\}$.

Proof. For definiteness, assume the positive puncture of $u$ maps to $c$ with $c_{0}=1$. Project $\Phi_{f}^{\delta}$ to the $\left(x_{0}, y_{0}\right)$-plane. The image of this projection is contained in the region

$$
\left\{-\alpha\left|x_{0}-1\right| \leq y_{0} \leq \alpha\left|x_{0}-1\right|\right\}
$$

for some $\alpha$. If the projection $u_{0}$ of $u$ to the ( $x_{0}, y_{0}$ )-plane is non-constant, then it covers at least one of the regions

$$
\left\{-\alpha\left|x_{0}-1\right|>y_{0}\right\} \cap B_{r}((1,0)) \quad \text { or } \quad\left\{\alpha\left|x_{0}-1\right|<y_{0}\right\} \cap B_{r}((1,0)),
$$

for some ball $B_{r}((1,0))$. Since $u$ has boundary on $\Phi_{k}^{\delta}, u_{0}$ takes no boundary point to the line $\left\{x_{0}=1\right\}$. This and the above covering property contradicts $u_{0}$ being bounded in the $y_{0}$-direction. The lemma follows.
q.e.d.

Lemma 10.2. The image of every holomorphic disk with boundary on $\Phi_{f}^{\delta}$ is contained in the region $\left\{\left|x_{0}\right| \leq 1\right\}$.

Proof. Arguing as in the proof of Lemma 10.1, we find that the projection to the $\left(x_{0}, y_{0}\right)$-plane of a holomorphic disk with boundary on $\Phi_{f}^{\delta}$ cannot intersect the lines $\left\{x_{0}= \pm 1\right\}$ in interior points. It follows that the image of any disk lies entirely in one of the regions $\left\{x_{0} \leq-1\right\}$, $\left\{\left|x_{0}\right| \leq 1\right\}$, or $\left\{x_{0} \geq 1\right\}$. However, a disk with image in $\left\{x_{0} \geq 1\right\}$ $\left(\left\{x_{0} \leq-1\right\}\right)$ must have its positive puncture at a Reeb chord $c$ with $c_{0}=1\left(c_{0}=-1\right)$. The lemma follows from Lemma 10.1. q.e.d.

Assume now that $\phi_{\delta}(L)$ is generic for each $\delta \neq 0$. Then, (see the proof of Lemma 7.24) any rigid disk with boundary on $\Phi_{f}^{\delta}$ and positive corner at some Reeb chord $c$ with $c_{0}= \pm 1$ is transversely cut out. Moreover, by Lemma 7.12 , transversality of the $\bar{\partial}$-equation can be achieved by perturbation near the positive puncture of a disk and it follows that there exists (arbitrarily small) perturbations of $\Phi_{f}^{\delta}$ which are supported in the region $\left\{\left|x_{0}\right|<1\right\}$ and which makes every moduli space (of formal dimension $\leq 1$ ) transversely cut out. We fix such a perturbation of $\Phi_{f}^{\delta}$, but keep the notation $\Phi_{f}^{\delta}$ for the perturbed Legendrian embedding.

Let $\mathcal{A}\left(\Phi_{f}^{\delta}\right)$ denote the algebra over $\mathbb{Z}_{2}\left[H_{1}(\mathbb{R} \times L)\right]=\mathbb{Z}_{2}\left[H_{1}(L)\right]$ generated by the Reeb chords of $\Phi_{f}^{\delta}$ as in Subsection 2 and define the map (differential) $\partial$ of $\mathcal{A}\left(\Phi_{f}^{\delta}\right)$ as there.

Lemma 10.3. The map $\partial: \mathcal{A}\left(\Phi_{f}^{\delta}\right) \rightarrow \mathcal{A}\left(\Phi_{f}^{\delta}\right)$ satisfies $\partial \circ \partial=0$.
Proof. In the light of Lemma 10.2, a word by word repetition of the proof of Lemma 2.5 establishes the lemma. q.e.d.
10.2. Invariance under handle slides. Let $\phi_{t}: L \rightarrow \mathbb{C}^{n} \times \mathbb{R},-1 \leq$ $t \leq 1$ be a Legendrian isotopy such that $L_{0}$ is a generic handle slide moment. That is, there exists one handle slide disk in some $\mathcal{M}_{A}(a ; \mathbf{b})$, which is the only non-empty moduli space of formal negative dimension, that all moduli spaces of holomorphic disk with boundary on $\phi_{t}(L)=L_{t}$, $t \neq 0$ of negative formal dimension are empty, and that all moduli spaces of rigid disks are transversally cut out. We choose notation so that $\left\{b_{1}, \ldots, b_{r}, a, c_{1}, \ldots, c_{s}\right\}$ are the Reeb chords of $L_{0}$ and so that

$$
\mathcal{Z}\left(b_{1}\right) \leq \cdots \leq \mathcal{Z}\left(b_{r}\right) \leq \mathcal{Z}(a) \leq \mathcal{Z}\left(c_{1}\right) \leq \cdots \leq \mathcal{Z}\left(c_{s}\right)
$$

Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be a positive Morse function with local minima at $\pm 1$, no critical points in the region $(-\infty,-1) \cup(1, \infty)$, and one local maximum at 0 .

Lemma 10.4. For all sufficiently small $\delta>0$, the Reeb chords of $\Phi_{f}^{\delta}$ are
$\left\{b_{j}[-1], b_{j}[1], b_{j}[0]\right\}_{j=1}^{r} \cup\{a[-1], a[1], a[0]\} \cup\left\{c_{j}[-1], c_{j}[1], c_{j}[0]\right\}_{j=1}^{s}$,
where for any Reeb chord $c$ of $L_{0},|c[-1]|=|c[1]|=|c[0]|-1=|c|$.
Proof. It is easy to see that for $\delta=0$, the Reeb chords are as described above and that the corresponding double points in $\mathbb{C} \times \mathbb{C}^{n}$ are transverse. This shows that the Reeb chords are as claimed for all sufficiently small $\delta$.

The second statement in the lemma is a straightforward consequence of the grading formula (for example, the front projection formula, see Lemma 3.4 in [4]). q.e.d.

We call $b_{j}[0], a[0]$, and $c_{j}[0],[0]$-Reeb chords, and $b_{j}[ \pm 1], a[ \pm 1]$, and $c_{j}[ \pm 1],[ \pm 1]$-Reeb chords. As above, we perturb $\Phi_{f}^{\delta}$ slightly in the region $\left|x_{0}\right|<1$ to make it generic with respect to holomorphic disks. Note that the $x_{0}$-coordinate of $[ \pm 1]$-Reeb chord equals $\pm 1$ and that the $x_{0}$ coordinate of a [0]-Reeb chord is very close to 0 for small $\delta>0$.

Consider a sequence of functions $f_{k}$ as above with $f_{k} \rightarrow 1$ as $k \rightarrow \infty$ (i.e. each $f_{k}$ has a non-degenerate maximum at 0 and non-degenerate local minima at $\pm 1$ ). Fix $k$ and pick $\delta>0$ sufficiently small so that $\Phi_{f_{k}}^{\delta}$ satisfies Lemma 10.4. Let $\Phi_{k}^{\delta}=\Phi_{f_{k}}^{\delta}$.

We next note that as $\delta \rightarrow 0, \Phi_{k}^{\delta} \rightarrow \Phi_{k}^{0}$ where

$$
\Phi_{k}^{0}(t, q)=\left(t, f_{k}^{\prime}(t) z(q), x(q), f_{k}(t) y(q), f_{k}(t) z(q)\right)
$$

with $(x(q), y(q), z(q))=\phi_{0}(q)$.
Lemma 10.5. There exists $k_{0}$ such that for all $k>k_{0}$, there exists a $\delta_{k}>0$ such that for all $\delta<\delta_{k}$ and any Reeb chord $c$, the following holds. The moduli spaces $\mathcal{M}(c[0], c[1])$ and $\mathcal{M}(c[0], c[-1])$ of holomorphic disks with boundary on $\Phi_{k}^{\delta}$ consist of exactly one point which is a transversely cut out rigid disk.

Proof. First consider the case $\delta=0$. It is easy to find rigid disks in the $\left(x_{0}, y_{0}\right)$-plane with positive puncture at $c[0]$ and negative puncture at $c[ \pm 1]$. Moreover, by Lemma 7.24, these disks are transversely cut out.

To see that these are the only disks, let $U$ and $V$ be neighborhoods of the endpoints of the Reeb chord $c$ in $L_{0}$ and consider the projections of $\Phi_{k}^{0}([-1,1] \times U)$ and $\Phi_{k}^{0}([-1,1] \times V)$ to $\mathbb{C}^{n}$. For sufficiently large $k$, these projections intersect only at 0 and it follows that there exists a positive $h>0$ such that the area of the projection of any disk with boundary on $\Phi_{k}^{0}$, positive puncture at $c[0]$, and negative at $c[ \pm 1]$ is either equal to zero or larger than $h$. Since $\mathcal{Z}(c[0]) \rightarrow \mathcal{Z}(c[ \pm 1])$ as $k \rightarrow \infty$, it follows that for $k$ large enough the disks in the ( $x_{0}, y_{0}$ )-plane are the only ones.

Finally, we note that the fact that the moduli space $\mathcal{M}(c[0], c[ \pm 1])$ corresponding to $\Phi_{k}^{0}$ is transversely cut out implies that the statement of the lemma holds also for $\Phi_{k}^{\delta}$ for all sufficiently small $\delta$ (where the smallness depends on $k$ ).

We next note that as $k \rightarrow \infty, \Phi_{k}^{0}$ approaches the Legendrian submanifold

$$
\Phi(t, q)=(t, 0, x(q), y(q), z(q))
$$

The projection of this Legendrian submanifold to $\mathbb{C}$ is simply the $x_{0}$-axis and its projection to $\mathbb{C}^{n}$ agrees with that of $L_{0}$.

Lemma 10.6. There exists $k_{0}$ such that for all $k>k_{0}$, there exists a $\delta_{k}>0$ such that for all $\delta<\delta_{k}$ and any Reeb chord $c \neq a$, the following holds. If the moduli space $\mathcal{M}_{A}(c[0] ; \mathbf{e})$, where $\mathbf{e}$ is a word constant in the $[0]$-generators and $\mathbf{e} \neq c[ \pm 1]$, has formal dimension 0 then it is empty.

Proof. Again we start with the case $\delta=0$. Consider a disk $u$ as above with boundary on $\Phi_{k}^{0}$. As $k \rightarrow \infty, \Phi_{k}^{0} \rightarrow \Phi$ and the projection of $u$ converges to a broken disk $\left\{v^{j}\right\}_{j=1}^{m}$ with boundary on $L_{0}$. The components $v^{j}$ of such a broken disk either have formal dimension at least 0 , or equals the handle slide disk. Also, any Reeb chord $b$ appearing as a puncture of some $v^{j}$ has $\mathcal{Z}(b) \leq \mathcal{Z}(c)$ and exactly one component of the broken disk must have its positive puncture at $c \neq a$. This component has formal dimension at least 0 (since it is not the handle slide disk). For a disk $v^{j}$, let $\left|v_{+}^{j}\right|$ be the grading of its positive puncture and $\left|v_{-}^{j}\right|$ the sum of gradings of its negative punctures and the negative of the grading of the homology data. Then the formal dimension of (the moduli space of) $v^{j}$ is $\left|v_{+}^{j}\right|-\left|v_{-}^{j}\right|-1$. The above implies that

$$
N=\sum_{j=1}^{m}\left(\left|v_{+}^{j}\right|-\left|v_{-}^{j}\right|\right) \geq 1
$$

Since the positive puncture of $u$ is its only [0]-puncture, it follows that the formal dimension of $u$ equals $N$. The statement of the lemma follows for $\delta=0$. Since emptiness of a moduli space is an open condition, the lemma follows in general. q.e.d.

Let $\Omega$ be the map from $\mathcal{A}\left(\Phi_{k}^{\delta}\right)$ to $\mathcal{A}\left(L_{0}\right)$ which maps $c[ \pm 1]$ to $c$ and $c[0]$ to 0 for any Reeb chord $c$ of $L$.

Lemma 10.7. There exists $k_{0}$ such that for all $k>k_{0}$, there exists a $\delta_{k}>0$ such that for all $\delta<\delta_{k}$, the following holds. If $u$ is a holomorphic disk with boundary on $\Phi_{k}^{\delta}$ in $\mathcal{M}_{C}(a[0]$, $\mathbf{e})$, where $\mathbf{e}$ is a word constant in the [0]-generators and $\mathbf{e} \neq a[ \pm 1]$, and if this moduli space has formal dimension 0 , then $C=A$ and $\Omega \mathbf{e}=\mathbf{b}$.

Proof. Consider first the case $\delta=0$. Taking the limit as $k \rightarrow \infty$ and arguing as in the proof of Lemma 10.6, we see that the projection of $u$ converges to a broken disk $\left\{v^{j}\right\}_{j=1}^{m}$, that all Reeb chords $b$ appearing as a puncture of some $v^{j}$ satisfies $\mathcal{Z}(a) \geq \mathcal{Z}(b)$, and that there is a unique component with its positive puncture at $a$. If this component is not the handle slide disk, then the argument in the proof of Lemma 10.6 shows that the formal dimension of $u$ is at least 1 . If, on the other hand, this component is the handle slide disk, then the formal dimension of $u$ equals 0 only if the broken disk has no other components. This shows the lemma for $\delta=0$. Again since the condition that a moduli space is empty is open, the lemma follows in general. q.e.d.

Fix $k$ sufficiently large and $\delta>0$ sufficiently small so that Lemmas $10.5,10.6$, and 10.7 holds for $\Phi_{k}^{\delta}$. We also assume that $\Phi_{k}^{\delta}$ is generic with respect to holomorphic disks. Let $\Phi=\tilde{\Phi}_{k}^{\delta}$

Let $\hat{\mathcal{A}}=\mathcal{A}(\Phi)$. We denote the differential of $\hat{\mathcal{A}}$ by $\Delta$, see Lemma 10.3. There are natural inclusions $\mathcal{A}_{ \pm}=\mathcal{A}\left(L_{ \pm \delta}\right) \subset \hat{\mathcal{A}}$. Lemma 10.1 implies that this is an inclusion of DGA's in other words,

$$
\Delta c[ \pm 1]=\Gamma_{ \pm}\left(\partial_{ \pm} c\right)
$$

where $\Gamma_{ \pm}: \mathcal{A}_{ \pm} \rightarrow \hat{\mathcal{A}}$ is the map defined on generators by $\Gamma_{ \pm}(c)=c[ \pm 1]$, and where $\partial_{ \pm}$is the differential on $\mathcal{A}_{ \pm}$. For generators $b_{j}[0]$, we have by Lemmas 10.5 and 10.6

$$
\begin{equation*}
\Delta b_{j}[0]=b_{j}[1]+b_{j}[-1]+\beta_{1}^{j}+\mathcal{O}(2), \tag{10.2}
\end{equation*}
$$

where $\beta_{1}^{j}$ is linear in the $c[0]$-generators and $\mathcal{O}(2)$ denotes a linear combination of monomials which are at least quadratic in the [0]-generators. For the generator $a[0]$, we have by Lemmas 10.5 and 10.7

$$
\begin{equation*}
\Delta a[0]=a[1]-a[-1]+\epsilon+\alpha_{1}+\mathcal{O}(2) \tag{10.3}
\end{equation*}
$$

where $\Omega(\epsilon)=m A \mathbf{b}$, where $m \in \mathbb{Z}_{2}$ and where $\alpha_{1}$ is linear in the [0]generators. For generators $c_{j}[0]$ we have by Lemmas 10.5 and 10.6

$$
\begin{equation*}
\Delta c_{j}[0]=c_{j}[1]+c_{j}[-1]+\gamma_{1}^{j}+\delta_{1}^{j}(a[0])+\mathcal{O}(2), \tag{10.4}
\end{equation*}
$$

where $\gamma_{1}^{j}+\delta_{1}^{j}(a[0])$ is linear in the $[0]$-generators, where $\delta_{1}^{j}(a[0])$ lies in the ideal generated by $a[0]$, and where $\gamma_{1}^{j}$ is constant in the $a[0]$ generator.

Below we will consider $\partial_{+}$and $\partial_{-}$as different differentials on the algebra $\mathcal{A}$. Let $\epsilon$ be as in (10.3) and write $\theta=\Omega(\epsilon)$. Consider the stable tame isomorphism $\psi$ of $\left(\mathcal{A}, \partial_{+}\right)$defined on generators as

$$
\psi(c)= \begin{cases}c & \text { if } c \neq a \\ a+\theta & \text { if } c=a\end{cases}
$$

If $v \in A$ and $c$ is a generator of $\mathcal{A}$, then let

$$
\left(\frac{v}{c}\right) \bullet: \mathcal{A} \rightarrow \mathcal{A}
$$

be the map defined on monomials by replacing each occurrence of $c$ by $v$. Then with $\partial_{+}^{\psi}=\psi^{-1} \circ \partial_{+} \circ \psi$ denoting the induced differential, a straightforward calculation gives

$$
\partial_{+}^{\psi}(c)= \begin{cases}\left(\frac{a-\theta}{a}\right) \bullet\left(\partial_{+} c\right) & \text { if } c \neq a  \tag{10.5}\\ \partial_{+} a+\partial_{+} \theta & \text { if } c=a\end{cases}
$$

Lemma 10.8. The algebra $\left(\mathcal{A}, \partial_{-}\right)$is isomorphic to the algebra $\left(\mathcal{A}, \partial_{+}^{\psi}\right)$.

Proof. We prove that the two differentials agree on generators. By Lemma 10.3, $\Delta^{2}=0$. Thus, summing the terms constant in the [0]generators after acting by $\Delta$ in (10.2), we find

$$
\begin{equation*}
0=\partial_{+} b_{j}[1]+\partial_{-} b_{j}[-1]+\left(\Delta \beta_{1}\right)_{0}, \tag{10.6}
\end{equation*}
$$

where $\left(\Delta \gamma_{1}\right)_{0}$ denotes the part of $\Delta \gamma_{1}$ which is constant in the [0]generators. Since the constant part of $\Delta b_{k}[0]$ equals $b_{k}[1]+b_{k}[-1]$, it follows that

$$
\Omega\left(\Delta \beta_{1}\right)_{0}=0 .
$$

Therefore, applying $\Omega$ in (10.6), we conclude

$$
\begin{equation*}
\partial_{-} b_{j}=\partial_{+} b_{j}=\left(\frac{a-\theta}{a}\right) \bullet\left(\partial_{+} b_{j}\right), \tag{10.7}
\end{equation*}
$$

since no monomial in $\partial_{+} b_{j}$ contains $a$.
Applying $\Delta$ to (10.3), we find similarly

$$
\begin{equation*}
\partial_{-} a[-1]=\partial_{+} a[1]+\Delta \epsilon+\left(\Delta \alpha_{1}\right)_{0} . \tag{10.8}
\end{equation*}
$$

The first equality in (10.7) implies that

$$
\Omega(\Delta \epsilon)=\partial_{+} \theta
$$

Since every [0]-generator in $\alpha_{1}$ is for the form $b_{j}[0]$, we find, as with $\beta_{1}$ above, that $\Omega\left(\Delta \alpha_{1}\right)_{0}=0$. We conclude

$$
\begin{equation*}
\partial_{-} a=\partial_{+} a+\partial_{+} \theta . \tag{10.9}
\end{equation*}
$$

Applying $\Delta$ to (10.4) gives

$$
\begin{equation*}
\partial_{-} c_{j}[-1]=\partial_{+} c_{j}[1]+\left(\Delta \gamma_{1}^{j}\right)_{0}+\left(\Delta \delta_{1}^{j}(a[0])\right)_{0} . \tag{10.10}
\end{equation*}
$$

Applying $\left(\frac{a[-1]+\epsilon}{a[1]}\right) \bullet$ to both sides in (10.10) and noting that no monomial in $\partial_{-} c_{j}[-1]$ contains an $a[1]$ generator, we get

$$
\begin{align*}
\partial_{-} c_{j}[-1]= & \left(\frac{a[-1]+\epsilon}{a[1]}\right) \bullet\left(\partial_{+} c_{j}[1]\right)+\left(\frac{a[-1]+\epsilon}{a[1]}\right) \bullet\left(\Delta \gamma_{1}^{j}\right)_{0}  \tag{10.11}\\
& +\left(\frac{a[-1]+\epsilon}{a[1]}\right) \bullet\left(\Delta \delta_{1}^{j}(a[0])\right)_{0} .
\end{align*}
$$

Each term in $\left(\Delta \delta_{1}^{j}(a[0])\right)_{0}$ arises by replacing $a[0]$ in every monomial $\xi a[0] \eta$ of $\delta_{1}^{j}(a[0])$ with $(a[1]-a[-1]+\epsilon)$ yielding $\xi(a[1]+a[-1]+\epsilon) \eta$. When $\left(\frac{a[1]+\epsilon}{a[-1]}\right) \bullet$ is a applied to $\xi(a[1]+a[-1]+\epsilon) \eta$, the result is

$$
\xi(a[-1]+\epsilon+a[-1]+\epsilon) \eta=0 .
$$

Thus, the last term in (10.11) vanishes. Since the [0]-generator of any monomial in $\gamma_{1}^{j}$ equals either $c_{k}[0]$ for some $k$, or $b_{r}[0]$ for some $r$ and since the constant part of $\Delta c_{k}[0]$ equals $c_{k}[1]+c_{k}[-1]$, and the constant part of $\Delta b_{k}[0]$ equals $b_{k}[1]+b_{k}[-1]$, we conclude that

$$
\Omega\left(\frac{a[-1]+\epsilon}{a[1]}\right) \cdot\left(\Delta \gamma_{1}^{j}\right)_{0}=0
$$

Thus, applying $\Omega$ to (10.11), we arrive at

$$
\begin{equation*}
\partial_{-} c_{j}=\left(\frac{a+\theta}{a}\right) \bullet \partial_{+} c_{j} . \tag{10.12}
\end{equation*}
$$

The lemma follows from (10.7), (10.9), (10.12).
q.e.d.

Corollary 10.9. If $L_{t} \subset \mathbb{C}^{n} \times \mathbb{R},-1 \leq t \leq 1$, is a Legendrian isotopy with a generic handle slide at $t=0$ as above, then the stable tame isomorphism classes of $\left(\mathcal{A}\left(L_{-1}, \partial_{-1}\right)\right)$ and $\left(\mathcal{A}\left(L_{1}\right), \partial_{1}\right)$ are the same.

## Appendix

In this appendix, we present the proofs of two technical lemmas.
Proof of Lemma 5.6. We establish the following two properties of the metric $\hat{g}$ and the endomorphism $J$. If $\gamma$ is a curve in $L$ with tangent vector $T$ and $X$ is any vector field in $T(T L)$ along $\gamma$, then

$$
\begin{equation*}
\hat{\nabla}_{T} J X=J \hat{\nabla}_{T} X \tag{10.13}
\end{equation*}
$$

If $X, Y$, and $Z$ are tangent vectors to $T L$ at $(p, 0) \in L$ such that $Y$ and $Z$ are horizontal (i.e. tangent to $L$ ) and if $\hat{R}$ denotes the curvature tensor of $\hat{g}$ at $(p, 0)$ then

$$
\begin{equation*}
\hat{R}(J X, Y) Z=J \hat{R}(X, Y) Z \tag{10.14}
\end{equation*}
$$

For (10.13), use local coordinates and write, for $\gamma(t)=x(t), T(x)=$ $a_{k}(x) \partial_{k}, X(x)=b_{j}(x) \partial_{j}+b_{j^{*}}(x) \partial_{j^{*}}$. By Lemma 5.5,
$\hat{\nabla}_{T} J X=a_{k} \hat{\nabla}_{\partial_{k}}\left(-b_{j^{*}} \partial_{j}+b_{j} \partial_{j^{*}}\right)$
$=a_{k}\left[-\left(\partial_{k} b_{j^{*}}\right) \partial_{j}+\left(\partial_{k} b_{j}\right) \partial_{j^{*}}\right.$

$$
\left.-b_{j^{*}}\left(\hat{\Gamma}_{k j}^{r} \partial_{r}+\hat{\Gamma}_{k j}^{r^{*}} \partial_{r^{*}}\right)+b_{j}\left(\hat{\Gamma}_{k j^{*}}^{r} \partial_{r}+\hat{\Gamma}_{k j^{*}}^{r^{*}} \partial_{r^{*}}\right)\right]
$$

$$
=a_{k}\left[-\left(\partial_{k} b_{j^{*}}\right) \partial_{j}+\left(\partial_{k} b_{j}\right) \partial_{j^{*}}-b_{j^{*}} \hat{\Gamma}_{k j}^{r} \partial_{r}+b_{j} \hat{\Gamma}_{k j^{*}}^{r^{*}} \partial_{r^{*}}\right]
$$

$$
=J a_{k}\left[\left(\partial_{k} b_{j^{*}}\right) \partial_{j^{*}}+\left(\partial_{k} b_{j}\right) \partial_{j}+b_{j^{*}} \hat{\Gamma}_{k j^{*}}^{r^{*}} \partial_{r^{*}}+b_{j} \hat{\Gamma}_{k j}^{r} \partial_{r}\right]=J \hat{\nabla}_{T} X .
$$

For (10.14), introduce normal coordinates $x$ around $p$. Then,

$$
\begin{equation*}
g_{i j}(0)=\delta_{i j}, \quad \Gamma_{i j}^{k}(0)=0 \tag{10.15}
\end{equation*}
$$

for all $i, j, k$, and hence Lemma 5.4 implies,

$$
\hat{g}_{i j}(0, \xi)=\delta_{i j}+\mathcal{O}\left(\xi^{2}\right), \quad \hat{g}_{i^{*} j}(0, \xi)=0, \quad \hat{g}_{i^{*} j^{*}}(0, \xi)=\delta_{i j} .
$$

Therefore,

$$
\begin{equation*}
\hat{g}^{i j}(0, \xi)=\delta^{i j}+\mathcal{O}\left(\xi^{2}\right), \quad \hat{g}^{i^{*} j}(0, \xi)=0, \quad \hat{g}^{i^{*} j^{*}}(0, \xi)=\delta^{i j} \tag{10.16}
\end{equation*}
$$

We show that, in these normal coordinates,

$$
\begin{equation*}
\hat{R}\left(\partial_{i^{*}}, \partial_{j}\right) \partial_{k}=J \hat{R}\left(\partial_{i}, \partial_{j}\right) \partial_{k} \tag{10.17}
\end{equation*}
$$

at ( 0,0 ). Since $\hat{R}$ is a tensor field, (10.17) implies (10.14).
Lemma 5.5 implies that all Christofel symbols of $\hat{g}$ vanishes at $(x, \xi)=$ $(0,0)$ and also that $\partial_{i} \Gamma_{j k}^{r_{k}^{*}}(x, 0)=0$ all $i, j, k, r^{*}$. Hence,

$$
\begin{aligned}
\hat{R}\left(\partial_{i}, \partial_{j}\right) \partial_{k} & =\hat{\nabla}_{\partial_{i}} \hat{\nabla}_{\partial_{j}} \partial_{k}-\hat{\nabla}_{\partial_{j}} \hat{\nabla}_{\partial_{i}} \partial_{k} \\
& =\left(\partial_{i} \hat{\Gamma}_{j k}^{r}\right) \partial_{r}+\left(\partial_{i} \hat{\Gamma}_{j k}^{*}\right) \partial_{r^{*}}-\left(\partial_{j} \hat{\Gamma}_{i k}^{r}\right) \partial_{r}-\left(\partial_{j} \hat{\Gamma}_{i k}^{r^{*}}\right) \partial_{r^{*}} \\
& =\left(\partial_{i} \Gamma_{j k}^{r}\right) \partial_{r}-\left(\partial_{j} \Gamma_{i k}^{r}\right) \partial_{r},
\end{aligned}
$$

and thus,

$$
\begin{equation*}
J \hat{R}\left(\partial_{i}, \partial_{j}\right) \partial_{k}=\left(\partial_{i} \Gamma_{j k}^{r}\right) \partial_{r^{*}}-\left(\partial_{j} \Gamma_{i k}^{r}\right) \partial_{r^{*}} . \tag{10.18}
\end{equation*}
$$

We compute the left-hand side of (10.17):

$$
\begin{align*}
\hat{R}\left(\partial_{i^{*}}, \partial_{j}\right) \partial_{k} & =\hat{\nabla}_{\partial_{i^{*}}} \hat{\nabla}_{\partial_{j}} \partial_{k}-\hat{\nabla}_{\partial_{j}} \hat{\nabla}_{\partial_{i^{*}}} \partial_{k} \\
& =\hat{\nabla}_{\partial_{i^{*}}}\left(\hat{\Gamma}_{j k}^{r} \partial_{r}+\hat{\Gamma}_{j k}^{r_{k}^{*}} \partial_{r^{*}}\right)-\hat{\nabla}_{\partial_{j}}\left(\hat{\Gamma}_{i^{*} k}^{r} \partial_{r}+\hat{\Gamma}_{i^{*} k}^{r^{*}} \partial_{r^{*}}\right) . \tag{10.19}
\end{align*}
$$

Lemma 5.5 gives $\partial_{j} \hat{\Gamma}_{i^{*} k}^{r}=0$, and Lemma 5.4 in combination with (10.16) give $\partial_{i^{*}} \hat{\Gamma}_{j k}^{r}=0$. Hence,

$$
\begin{equation*}
\hat{R}\left(\partial_{i^{*}}, \partial_{j}\right) \partial_{k}=\left(\partial_{i^{*}} \hat{\Gamma}_{j k}^{r^{*}}\right) \partial_{r^{*}}-\left(\partial_{j} \hat{\Gamma}_{i k^{*}}^{r^{*}}\right) \partial_{r^{*}}=\left(\partial_{i^{*}} \hat{\Gamma}_{j k}^{r^{*}}\right) \partial_{r^{*}}-\left(\partial_{j} \Gamma_{i k}^{r}\right) \partial_{r^{*}} . \tag{10.20}
\end{equation*}
$$

It thus remains to compute $\partial_{i^{*}} \hat{\Gamma}_{j k}^{r^{*}}$.

$$
\begin{align*}
\partial_{i^{*}} \hat{\Gamma}_{j k}^{r^{*}} & =\frac{1}{2} \partial_{i^{*}}\left(\hat{g}^{r^{*} l^{*}}\left(\hat{g}_{j l^{*}, k}+\hat{g}_{k l^{*}, j}-\hat{g}_{j k, l^{*}}\right)+\hat{g}^{r^{*} l}\left(\hat{g}_{j l, k}+\hat{g}_{k l, j}-\hat{g}_{j k, l}\right)\right)  \tag{10.21}\\
& =\frac{1}{2} \hat{g}^{r l} \partial_{i^{*}}\left(\hat{g}_{j l^{*}, k}+\hat{g}_{k l^{*}, j}-\hat{g}_{j k, l^{*}}\right) \quad[\text { by }(10.16)] \\
& =\frac{1}{2} g^{r l}\left(\left(\partial_{k} \Gamma_{j i}^{m}\right) g_{m l}+\left(\partial_{j} \Gamma_{k i}^{m}\right) g_{m l}-\left(R_{j i k l}+R_{j l k i}\right)\right) \tag{10.15}
\end{align*}
$$

$$
=\frac{1}{2}\left(\partial_{k} \Gamma_{j i}^{r}+\partial_{j} \Gamma_{k i}^{r}-\left(R_{j i k r}+R_{j r k i}\right)\right) \quad[(10.15)] .
$$

But

$$
R_{j i k r}=g\left(\nabla_{\partial_{j}} \nabla_{\partial_{i}} \partial_{k}-\nabla_{\partial_{i}} \nabla_{\partial_{j}} \partial_{k}, \partial_{r}\right)=\partial_{j} \Gamma_{i k}^{r}-\partial_{i} \Gamma_{j k}^{r}=\partial_{j} \Gamma_{k i}^{r}-\partial_{i} \Gamma_{j k}^{r},
$$

and

$$
R_{j r k i}=R_{k i j r}=\partial_{k} \Gamma_{i j}^{r}-\partial_{i} \Gamma_{k j}^{r}=\partial_{k} \Gamma_{j i}^{r}-\partial_{i} \Gamma_{j k}^{r}
$$

Hence

$$
\partial_{i^{*}} \hat{\Gamma}_{j k}^{r^{*}}=\partial_{i} \Gamma_{j k}^{r}
$$

which together with (10.19) and (10.20) imply (10.17).
Consider a geodesic of ( $T L, \hat{g}$ ) in $L$ with tangent vector $T$. By (10.13) and (10.14),

$$
\begin{equation*}
\hat{\nabla}_{T} \hat{\nabla}_{T} J X+\hat{R}(J X, T) T=J\left(\hat{\nabla}_{T} \hat{\nabla}_{T} X+\hat{R}(X, T) T\right) \tag{10.22}
\end{equation*}
$$

Thus $X$ is a Jacobi field if and only if $J X$ is. q.e.d.

Proof of Lemma 5.10. For simplicity, we suppress intermediate functions in the notation, e.g., we write $\sigma(\zeta)$ for $\sigma(w(\zeta), F(\zeta))$. Consider (a). Assume that $w, v, u, q, F$ are smooth functions. By (5.18)

$$
\begin{equation*}
|G(\zeta, \lambda)| \leq C(|v|+|u|) \tag{10.23}
\end{equation*}
$$

since the derivatives of $G$ are uniformly bounded.
(For simplicity, we will use the letter $C$ to denote many different constants in this proof. This (constant!) change of notation will not be pointed out each time.)

Let $\hat{G}(\zeta)=G(\zeta, \lambda)$. We write $(w, v, u, q, \sigma)=\left(x_{1}, x_{2}, x_{3}, x_{4}, x_{5}\right)$ and use the Einstein summation convention. The derivative of $\hat{G}(\zeta)$ is

$$
D \hat{G}(\zeta)=D_{j} \hat{G} \cdot D x_{j}
$$

where $D$ without subscript refers to derivatives with respect to $\zeta$, and $D_{j} \hat{G}$ refers to the derivative of $\hat{G}$ with respect to its $j$-th argument. We use the following notation for functions $\left(y_{1}, \ldots, y_{l}\right)$,

$$
\left|D^{j_{1}} y\right|^{k_{1}} \ldots\left|D^{j_{m}} y\right|^{k_{m}}=\sum_{\alpha \in A} \Pi_{k=1}^{l}\left|D^{j_{1}} y_{k}\right|^{\alpha_{k}^{1}} \ldots\left|D^{j_{m}} y_{k}\right|^{\alpha_{k}^{m}}
$$

where $A=\left\{\alpha \in\left(\mathbb{Z}_{\geq 0}\right)^{l m}: \alpha_{1}^{r}+\cdots+\alpha_{l}^{r}=k_{r}\right\}$.
Let $(w, F, q)=\left(y_{1}, y_{2}, y_{3}\right)$ and $(v, u)=\left(z_{1}, z_{2}\right)$, then by (5.18)

$$
\begin{aligned}
& \left|D_{j} \hat{G}\right| \leq C|z|, \quad j \in\{1,4,5\} \\
& \left|D_{j} \hat{G}\right| \leq C, \quad j \in\{2,3\} \\
& D \sigma=D_{1} \sigma \cdot D F+D_{2} \sigma \cdot D w, \text { hence, }|D \sigma| \leq C|D y| .
\end{aligned}
$$

Then

$$
\begin{equation*}
|D \hat{G}(\zeta)|^{2} \leq C\left(|z|^{2}|D y|^{2}+|z||D z \| D y|+|D z|^{2}\right) . \tag{10.24}
\end{equation*}
$$

The second derivative of $\hat{G}(\zeta)$ is

$$
D^{2} \hat{G}(\zeta)=D_{i} D_{j} \hat{G} \cdot D x_{i} \cdot D x_{j}+D_{j} \hat{G} \cdot D^{2} x_{j} .
$$

By (5.18),

$$
\begin{aligned}
&\left|D_{i} D_{j} \hat{G}\right| \leq C|z|, \quad i, j \in\{1,4,5\} \\
&\left|D_{i} D_{j} \hat{G}\right| \leq C, \quad j \in\{2,3\} \\
& D^{2} \sigma=D_{1}^{2} \sigma \cdot D F \cdot D F+2 D_{2} D_{1} \sigma \cdot D w \cdot D F \\
&+D_{2}^{2} \sigma D w \cdot D w+D_{1} \sigma \cdot D^{2} F+D_{2} \sigma \cdot D^{2} w, \\
& \text { hence }\left|D^{2} \sigma\right| \leq C\left(|D y|^{2}+\left|D^{2} y\right|\right) .
\end{aligned}
$$

Thus,

$$
\left|D^{2} \hat{G}(\zeta)\right|^{2} \leq C\left(\left|z^{2}\right|\left(|D y|^{4}+|D y|^{2}\left|D^{2} y\right|\right)+|z||D z||D y|\left|D^{2} y\right|\right.
$$

$$
\begin{equation*}
\left.+|D z|^{4}+|D z|\left|D^{2} z\right||D y|+|D z|^{2}\left|D^{2} z\right|+\left|D^{2} z\right|^{2}\right) \tag{10.25}
\end{equation*}
$$

Note that by (5.21) and (5.23), $r$, which the constant $C$ absorbs, controls the $q$ (or $y_{3}$ ) norms. Moreover, the remaining $y_{1}$ and $y_{2}$ norms are also absorbed by $C$. Thus, using (10.23), (10.24), and (10.25) we derive the estimate

$$
\begin{equation*}
\|\hat{G}(\zeta)\|_{2, \epsilon} \leq C\left(\|u\|_{2, \epsilon}+\|v\|_{2, \epsilon}\right) \tag{10.26}
\end{equation*}
$$

as follows. The Sobolev-Gagliardo-Nirenberg theorem implies $\|D y\|_{L^{4}} \leq$ $C\|D y\|_{1,2}$ (and the corresponding statement for $u$ and $v$ ). Morrey's theorem implies that $\|u\|_{2, \epsilon}$ controls the sup-norm of $u$ (and the corresponding statement for $v$ ). These facts together with Hölder's inequality give (10.26).

It is now straightforward to prove (a). Let $\Omega=(x, \xi, \eta, \theta, \sigma)$, then

$$
\begin{equation*}
G(\Omega, \lambda)=G(\Omega, 0)+D_{6} G(\Omega, 0) \cdot \lambda+R(\Omega, \lambda) \cdot \lambda \cdot \lambda . \tag{10.27}
\end{equation*}
$$

Differentiating (10.27) twice with respect to $\lambda$ of and applying (5.18), we find $R(x, 0,0, \theta, \sigma, \lambda)=0$. Applying the argument above to $D_{6} G$ and $R$, and to $G(\Omega, 0)$ but using (5.19) and $u$ instead of (5.18) and ( $u, v$ ), (5.20) follows.

The proof of (b) is similar. We first use (5.22) and (5.23) to conclude

$$
\begin{equation*}
\hat{G}=G(w, v, u, \sigma, 0) \leq C|u|^{2} \tag{10.28}
\end{equation*}
$$

The derivative of $\hat{G}(\zeta)$ is

$$
D \hat{G}(\zeta)=D_{j} \hat{G} \cdot D x_{j}
$$

and with $(w, F, v)=\left(y_{1}, y_{2}, y_{3}\right)$

$$
\begin{aligned}
& \left|D_{j} \hat{G}\right| \leq C|u|^{2}, \quad j \in\{1,2,4\} \\
& \left|D_{3} \hat{G}\right| \leq C|u| \\
& D \sigma=D_{1} \sigma \cdot D F+D_{2} \sigma \cdot D w, \text { hence }|D \sigma| \leq C|D y|
\end{aligned}
$$

Thus

$$
\begin{equation*}
|D \hat{G}(\zeta)|^{2} \leq C\left(|u|^{4}|D y|^{2}+|u|^{3}|D u||D y|+|u|^{2}|D u|^{2}\right) . \tag{10.29}
\end{equation*}
$$

The second derivative of $\hat{G}(\zeta)$ is

$$
D^{2} \hat{G}(\zeta)=D_{i} D_{j} \hat{G} \cdot D x_{i} \cdot D x_{j}+D_{j} \hat{G} \cdot D^{2} x_{j}
$$

We have

$$
\begin{aligned}
& \left|D_{i} D_{j} \hat{G}\right| \leq C|u|^{2}, \quad i, j \in\{1,2,4\} \\
& \left|D_{i} D_{3} \hat{G}\right| \leq C|u|, \quad i \in\{1,2,4\}, \\
& \left|D_{3}^{2} \hat{G}\right| \leq C .
\end{aligned}
$$

This implies

$$
\begin{align*}
\left|D^{2} \hat{G}(\zeta)\right|^{2} \leq & C\left(|u|^{4}\left(|D y|^{4}+|D y|^{2}\left|D^{2} y\right|+\left|D^{2} y\right|^{2}\right)\right.  \tag{10.30}\\
& +|u|^{3}\left(|D u \| D y|^{3}+|D u||D y|\left|D^{2} y\right|\right)+|u|^{2}|D u|^{2}|D y|^{2} \\
& \left.+|u|^{2}|D u|^{4}+|D u||D y|\left|D^{2} y\right|+\left|D^{2} u\right|\left|D^{2} y\right|+|u|^{2}\left|D^{2} u\right|^{2}\right) .
\end{align*}
$$

In the same way as above, we derive from (10.28), (10.29), and (10.30) the estimate

$$
\begin{equation*}
\|\hat{G}(\zeta)\|_{2, \epsilon} \leq C\|u\|_{2, \epsilon}^{2} \tag{10.31}
\end{equation*}
$$

The proof of (b) can now be completed as follows. Write $\Omega=(x, \xi, \eta, \sigma)$ then
$G(\Omega, \lambda)=G(\Omega, 0)+D_{5} G(\Omega, 0) \cdot \lambda+D_{5}^{2} G(\Omega, 0) \cdot \lambda \cdot \lambda+R(\Omega, \lambda) \cdot \lambda \cdot \lambda \cdot \lambda$, and differentiation gives $R(x, 0,0, \sigma, \lambda)=0$. For $G(\Omega, 0)$, we use (10.31). The term $D_{5} \hat{G}(\zeta)$ can be estimated as in (a) by $C\|u\|_{2, \epsilon}$. The two remaining terms are also estimated as in (a) by $C\left(\|u\|_{2, \epsilon}+\|v\|_{2, \epsilon}\right)$.

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