The dependence of star formation history and internal structure on stellar mass for 10⁵ low-redshift galaxies

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ABSTRACT

We study the relations between stellar mass, star formation history, size and internal structure for a complete sample of 122 808 galaxies drawn from the Sloan Digital Sky Survey. We show that low-redshift galaxies divide into two distinct families at a stellar mass of 3×10^{10} M_{\odot}. Lower-mass galaxies have young stellar populations, low surface mass densities and the low concentrations typical of discs. Their star formation histories are more strongly correlated with surface mass density than with stellar mass. A significant fraction of the lowest-mass galaxies in our sample have experienced recent starbursts. At given stellar mass, the sizes of low-mass galaxies are lognormally distributed with dispersion $\sigma(\ln R_{50}) \sim 0.5$, in excellent agreement with the idea that they form with little angular momentum loss through cooling and condensation in a gravitationally dominant dark matter halo. Their median stellar surface mass density scales with stellar mass as $\mu_* \propto M_*^{0.54}$, suggesting that the stellar mass of a disc galaxy is proportional to the three halves power of its halo mass. All of this suggests that the efficiency of the conversion of baryons into stars in low-mass galaxies increases in proportion to halo mass, perhaps as a result of supernova feedback processes. At stellar masses above $3 \times$ 10^{10} M_{\odot}, there is a rapidly increasing fraction of galaxies with old stellar populations, high surface mass densities and the high concentrations typical of bulges. In this regime, the size distribution remains lognormal, but its dispersion decreases rapidly with increasing mass and the median stellar mass surface density is approximately constant. This suggests that the star formation efficiency decreases in the highest-mass haloes, and that little star formation occurs in massive galaxies after they have assembled.

Key words: galaxies: evolution - galaxies: formation - galaxies: stellar content.

1 INTRODUCTION

Perhaps the most remarkable aspect of galaxies is their regularity – the fact that they can be classified into well-defined sequences. Hubble (1926) introduced the first scheme to categorize galaxies according to morphological type, and its basic concepts are still in use today. In its simplest form, three basic galaxy types are recognized: ellipticals, spirals and irregulars. These can be arranged in a linear sequence along which many properties vary coherently. The most striking trends are the variations in morphology (featureless ellipsoid to spiral disc to irregular appearance) and in colour and star formation activity. These trends were so striking that they formed the basis of Hubble's original scheme. Elliptical galaxies have red colours, little gas and no ongoing star formation; irregulars have blue colours, large gas fractions and are actively forming stars. An up-to-date review of star formation along the Hubble sequence is given in Kennicutt (1998).

It should be remembered that Hubble introduced his galaxy classification scheme at a time when distance estimates were available

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for only a handful of galaxies. Hubble could not study the properties of galaxies as a function of parameters that measure their absolute scale: sizes, absolute magnitudes and masses. Today it is known that the distribution of Hubble types depends strongly on galaxy luminosity. The typical elliptical galaxy is a factor of ~ 10 more luminous in the B band than the typical irregular galaxy (Roberts & Haynes 1994). As a result, the faint end of the galaxy luminosity function is dominated by late-type galaxies with strong emission lines, the bright end by early-type galaxies with little star formation (e.g. Loveday et al. 1992; Marzke et al. 1994, 1998; Lin et al. 1996; Zucca et al. 1997; Christlein 2000). These trends in star formation activity become even more pronounced when studied in the nearinfrared. This has led to suggestions that the ratio of the present to the past-averaged star formation rate may depend primarily on stellar mass and only secondarily on morphological type (Balogh et al. 2001; Boselli et al. 2001).

There have been many studies of stellar populations and star formation histories that focus on galaxies of a particular Hubble type. The star formation rates in most irregular and dwarf galaxies in the Local Group do not appear to have varied smoothly with time, but instead have undergone large fluctuations or 'bursts' (see Grebel 2000 for a recent review). There is evidence that the recent star formation histories of spirals depend on the surface brightness of their discs. Low surface brightness galaxies appear to have bluer colours and higher gas mass fractions than high surface brightness galaxies (e.g. de Blok, van der Hulst & Bothun 1995; de Blok, McGaugh & van der Hulst 1996; Bell & de Jong 2000). The situation for ellipticals is currently unclear. Although there is considerable evidence that the absorption-line strengths of early-type galaxies vary systematically with their velocity dispersions, there is much debate as to whether this reflects trends in stellar age, in metallicity or in the relative abundance of different heavy elements (e.g. Trager et al. 1998; Kuntschner et al. 2001).

The Sloan Digital Sky Survey will obtain u, g, r, i and z-band photometry, spectra and redshifts for at least 700 000 galaxies down to a limiting r-band magnitude of 17.77. This is the most ambitious survey of the local galaxy population ever undertaken and motivates a critical re-evaluation of trends in the star formation histories of galaxies as a function of mass, size, internal structure and environment.

We have recently developed a new method to constrain the star formation history, and to estimate the dust attenuation and stellar mass of galaxies (Kauffmann et al. 2003, hereafter Paper I). It is based on two stellar absorption-line indices, the 4000-Å break strength and the Balmer absorption-line index $H\delta_A$. These indices allow us to constrain the mean stellar age of a galaxy and the fraction of its stellar mass formed in bursts over the past few Gyr. A comparison with broad-band magnitudes then yields estimates of dust attenuation and of stellar mass. We have generated a large library of Monte Carlo realizations of different star formation histories, including starbursts of varying strength. We have used this library to generate likelihood distributions for parameters such as the burst mass fraction, dust attenuation strength, stellar mass and stellar mass-to-light ratio for every galaxy in a complete sample of 122 808 objects drawn from the Sloan Digital Sky Survey.

Paper I introduced our method and used it to measure how the total stellar mass of the Universe is distributed over galaxies as a function of their stellar mass, size, concentration, colour and surface mass density. In this paper, we focus on the *mass dependence* of the star formation histories, sizes and structural parameters of galaxies. We study to what extent the recent star formation histories of galaxies of different stellar mass can be characterized as continuous or

as burst-dominated. We also show how the distributions of structural parameters (size, surface mass density, concentration index) vary with stellar mass, and we demonstrate that variations in star formation history and in structural parameters are tightly coupled. Finally, we interpret our findings in the context of current theories of galaxy formation.

2 REVIEW OF THE OBSERVED AND DERIVED PARAMETERS OF GALAXIES

The sample of galaxies analysed in this paper is drawn from the Sloan Digital Sky Survey (York et al. 2000; Stoughton et al. 2002). We have included all galaxies with spectroscopic observations in the Data Release One (DR1) and with Petrosian *r*-band magnitudes in the range $14.5 < r^* < 17.77$ after correction for foreground galactic extinction. Details concerning the spectroscopic target selection may be found in Strauss et al. (2002). These data represent roughly 20 per cent of the projected survey total and include a total of 122 808 galaxies.

The reader is referred to Paper I for more information concerning the sample and a detailed description of the methods used to derive parameters such as stellar mass and burst mass fraction. We also discuss the uncertainties in our methods in Paper I. Here we provide a brief summary of the quantities that are studied in this paper.

(i) Absolute magnitudes. All magnitudes quoted in this paper are Petrosian magnitudes. In order to convert from apparent to absolute magnitude, we assume a Friedmann–Robertson–Walker cosmology with $\Omega = 0.3$, $\Lambda = 0.7$ and $H_0 = 70$ km s⁻¹ Mpc⁻¹. We have calculated the K-corrections K(z) for each galaxy using the routines in KCORRECT v1_11 (Blanton et al. 2002). In order to minimize the errors in this procedure, we K-correct the magnitudes of all galaxies in our sample to z = 0.1. More details concerning the SDSS photometric system may be found in Fukugita et al. (1996) and Smith et al. (2002). Information relating to the SDSS camera and photometric monitoring system can be found in Gunn et al. (1998) and Hogg et al. (2001). A technical paper describing the astrometry in the survey has recently been published by Pier et al. (2002).

(ii) Stellar spectral indices. We have adopted the narrow definition of the 4000-Å break introduced by Balogh et al. (1999), which we denote $D_n(4000)$. We also use the Balmer absorption index $H\delta_A$ as defined by Worthey & Ottaviani (1997) using a central bandpass bracketed by two continuum bandpasses. The evolution of the two indices have been calibrated using a new population synthesis code that incorporates high-resolution stellar libraries (Bruzual & Charlot 2003). Our measurements of H δ_A are corrected for contamination caused by nebular emission. As discussed in Paper I, the $D_n(4000)$ index is an excellent age indicator for young [<1 Gyr, $D_n(4000) <$ 1.5] stellar populations, but for older stellar populations, the index depends quite strongly on metallicity. Strong Balmer absorption occurs in galaxies that experienced a burst of star formation 0.1-1 Gyr ago. Analysed together, the two indices are a powerful probe of the recent star formation history of a galaxy. Because each is defined over a narrow wavelength interval, they are not sensitive to dust attenuation effects.

(iii) *Stellar masses.* The stellar masses are derived assuming a universal initial mass function (IMF) in the parametrization of Kroupa (2001). The typical 95 per cent confidence range in our estimated stellar masses is ± 40 per cent.

(iv) *Burst mass fractions.* We define the parameter F_{burst} as the fraction of the total stellar mass of a galaxy that formed in 'burst'

mode over the past 2 Gyr. In Paper I we showed that galaxies with continuous star formation histories occupy a very narrow locus in the $H\delta_A/D_n(4000)$ plane. Galaxies that experienced recent bursts have $H\delta_A$ values that are displaced from this locus. Because the typical observational error on the $H\delta_A$ index is large (1.4 Å) in our sample, the median value of the likelihood distribution of F_{burst} is sensitive to the mix of star formation histories in our model library, which functions as a Bayesian prior for our analysis. As a result, we focus below on galaxies for which our analysis implies $F_{\text{burst}} > 0$ with high confidence (97.5 per cent for our standard prior). The tests described in Paper I showed the resulting sample to be insensitive to the actual prior adopted.

(v) *Galaxy sizes*. We study the distribution functions of R50(z), the radius enclosing 50 per cent of the Petrosian z-band luminosity of a galaxy. We note that the sizes output by the current SDSS photometric pipeline have not been corrected for seeing effects (Stoughton et al. 2002).

(vi) *Concentration indices.* In order to be consistent with previous work (Blanton et al. 2001; Shimasaku et al. 2001; Strateva et al. 2001) we define the concentration index *C* as the ratio R90/R50, where R90 and R50 are the radii enclosing 90 and 50 per cent of the Petrosian *r*-band luminosity of the galaxy. It has been shown by Shimasaku et al. (2001) and Strateva et al. (2001) that for bright galaxies, there is a good correspondence between concentration parameter and 'by-eye' classification into Hubble type, with $C \sim 2.6$ marking the boundary between early- and late-type galaxies.

(vii) Surface brightnesses and surface mass densities. Following Blanton et al. (2001), we define the half-light surface brightness μ 50 to be the average *r*-band surface brightness within the *r*-band half-light radius *R*50 in mag arcsec⁻². The main spectroscopic survey is complete down to a limiting half-light surface brightness μ 50 = 24.5 mag arcsec⁻². We define the surface mass density μ_* as $0.5M_*/[\pi R50^2(z)]$, where R50(z) is the Petrosian half-light radius in the *z* band (we choose a radius defined in the *z* band rather than the *r* band, as this provides a better approximation to the radius enclosing half the total stellar mass).

In this paper, we will often use bivariate density distributions $\phi(x, y)$ in the parameters x and y, defined such that $\phi(x, y) dx dy$ is the number of galaxies per unit volume with x in the interval (x, x + dx) and y in the interval (y, y + dy). Such bivariate distributions are calculated by weighting each galaxy in the sample by $1/V_{\text{max}}$, where V_{max} is the volume corresponding to the total redshift range over which the galaxy would pass our sample selection criteria. Because our sample is extremely large, errors arising from large-scale structure are expected to be small. In many cases we will be interested in trends in the distribution of the parameter y as a function of the parameter x (for example, in the distribution of concentration or size as a function of stellar mass). In such cases, we make the trends more visible by plotting the conditional distribution of y given x, i.e. $\phi(x, y)/\int \phi(x, y') dy'$.

3 THE OBSERVED CORRELATIONS

3.1 The dependence of star formation history on stellar mass

In Fig. 1, we present the conditional density distributions of our two stellar age indicators $D_n(4000)$ and $H\delta_A$ as a function of stellar mass and as a function of *g*-band absolute magnitude. The grey-scale indicates the fraction of galaxies in a given logarithmic mass (or magnitude) bin that fall into each age-indicator bin. The contours are separated by factors of 2 in population density.

It is clear that $D_n(4000)$ and $H\delta_A$ depend strongly on both stellar mass and absolute magnitude. The main effect of transforming from luminosity to mass is to produce a more regular variation of the indices. Fig. 1 shows that low-mass galaxies are 'young' – they typically have low values of $D_n(4000)$ and high values of $H\delta_A$. At a mass of $M_* \sim 3 \times 10^{10} \,\mathrm{M_{\odot}}$, a transition towards older stellar populations begins to take effect. Almost all of the most massive galaxies in our sample have high values of $D_n(4000)$ and low values of $H\delta_A$.

Fig. 2 shows 'slices' through the $D_n(4000)-M_*$ distribution. Recall that the measurement error on the $D_n(4000)$ index is small, typically around 0.04, or a few per cent of the total range of values spanned by galaxies in our sample. Again, one sees a striking trend towards older stellar populations for galaxies with larger stellar masses. Both in Fig. 1 and in these plots one has the impression of two separate and relatively well-defined populations. The relative weight of the 'old' population increases strongly at stellar masses above $\sim 10^{10} \text{ M}_{\odot}$.

The bivariate density distribution of $H\delta_A$ and $D_n(4000)$ is shown in Fig. 3 for galaxies in eight ranges of stellar mass spanning four orders of magnitude from $M_* = 10^8 \text{ M}_{\odot}$ to $M_* = 10^{12} \text{ M}_{\odot}$. In these plots, the grey-scale indicates the fraction of galaxies that fall into each bin of $H\delta_A$ and $D_n(4000)$. It is striking how galaxies move diagonally across the $H\delta_A/D_n(4000)$ plane as their masses increase.

From Fig. 3, one can see clearly that low-mass galaxies with $D_n(4000) < 1.5$ (characteristic of stellar populations with mean ages of less than a few Gyr) have a larger scatter in H δ_A equivalent width than high-mass galaxies with similar break strengths. This means that the fraction of low-mass galaxies that have experienced recent bursts is higher than that for high-mass galaxies, even when the two populations are compared at similar mean stellar age. We illustrate this in detail in Fig. 4, where we plot the median value of H δ_A as a function of D_n(4000) for low-mass and for high-mass galaxies. The solid and dotted errorbars indicate the 25th to 75th and fifth to 95th percentile ranges of the distribution of H δ_A values. For comparison, we have overplotted 200 different model galaxies with continuous star formation histories. The models span a wide range in exponential time-scale ($\tau = 0.1$ to 10 Gyr) and formation redshift $(z_{\text{form}} = 20 \text{ to } 0.2)$. The reader is referred to Paper I for more details concerning the models. Solar metallicity models are plotted in the right-hand panel. In the left-hand panel, we have plotted 20 per cent solar models, which are more appropriate for galaxies with masses between 10^8 and 10^9 M_{\odot} (Tremonti et al. in preparation) For massive galaxies, the data agree extremely well with the model predictions. The scatter in H δ_A at given $D_n(4000)$ is consistent with that expected from observational errors. Note that although we plot the distribution of H δ_A down to very low D_n(4000) values, Fig. 2 shows that less than 5 per cent of galaxies with stellar masses in the range 10^{10} – 10^{11} M_{\odot} have values D_n(4000) too low to be consistent with any of our continuous models. For low-mass galaxies, this is no longer true. Fig. 2 shows that \sim 30 per cent of all galaxies have $D_n(4000) < 1.2$. As discussed in Paper I, many of these galaxies are likely to be experiencing a starburst at the present day. Fig. 4 shows that for low-mass galaxies, the data are also displaced upwards relative to the predictions of the continuous models and the scatter is too large to be consistent with observational errors. As discussed in Paper I, galaxies with moderate values of $D_n(4000)$ and very high values of H δ_A are probably in a post-starburst phase.

In Paper I, we introduced models in which galaxies formed stars in two different modes: (i) a 'continuous' mode, which we parametrized by an exponential law with time-scale τ and starting time t_{form} and (ii) a superposed 'burst' mode. A burst was defined



Figure 1. Conditional density distributions showing trends in the stellar age indicators $D_n(4000)$ and $H\delta_A$ as functions of the logarithm of stellar mass and of *g*-band absolute magnitude. Galaxies have been weighted by $1/V_{\text{max}}$ and the bivariate distribution function has been normalized to a fixed number of galaxies in each bin of log M_* or M(g). Here and in all subsequent contour plots, each contour represents a factor of 2 change in density.

to be an episode of star formation lasting between 3×10^7 and 3×10^8 yr in which a fraction F_{burst} of the total stellar mass of the galaxy was formed. In Paper I, we also introduced a method that allowed us to calculate the a posteriori likelihood distribution of F_{burst} for each galaxy in our sample, given its observed absorption-line indices and the measurement errors on these indices. In Fig. 5, we plot the distribution of galaxies as a function of F_{burst} for eight different mass ranges. The distribution of the median value of F_{burst} [which we denote as $F_{\text{burst}}(50 \text{ per cent})$] for all galaxies in the given mass range is plotted as a solid histogram. We define a subsample of high-confidence bursty galaxies as those objects with $F_{\text{burst}}(2.5 \text{ per cent}) > 0$, where $F_{\text{burst}}(2.5 \text{ per cent})$ is the lower 2.5 percentile point of the likelihood distribution of F_{burst} . The distribution of $F_{\text{burst}}(50 \text{ per cent})$ for this subsample is plotted as a dotted histogram in Fig. 5.

In our lowest-mass range $(10^8-3 \times 10^8 \text{ M}_{\odot})$, we find that half of all galaxies have $F_{\text{burst}}(50 \text{ per cent}) > 0$. However, it is only for 10 per cent of these objects that we can state with high (>97.5 per cent) confidence that a burst *did occur* in the recent past. The fraction of galaxies that have experienced recent bursts decreases very strongly with increasing stellar mass. This trend is apparent both in the fraction of galaxies with $F_{\text{burst}}(50 \text{ per cent}) > 0$ and in the fraction with $F_{\text{burst}}(2.5 \text{ per cent}) > 0$. There are very few high-confidence bursty

galaxies in our high-mass bins. Out of our entire sample of 120 000 galaxies, we only pick out around 500 galaxies with masses comparable to that of the Milky Way and with $F_{\text{burst}}(2.5 \text{ per cent}) > 0$. For our full sample, the median value of the burst mass fraction correlates with stellar mass. The same is not true for the high-confidence bursty galaxies - these systems have burst mass fractions that are independent of stellar mass. Note that our requirement that a burst be detected at high confidence will, by definition, bias this subsample towards high-burst fractions. The trend in burst mass fraction for our full sample is in some sense expected, because it is well known that more massive galaxies have lower gas mass fractions and thus contain less fuel for star formation (e.g. Roberts & Haynes 1994; Bell & de Jong 2001; Boselli et al. 2001). What about the influence of our adopted prior? The analysis in Paper I showed that decreasing the number of bursty galaxies in the model library resulted in lower estimated values of F burst (50 per cent). Our standard prior assumes that the fraction of bursty galaxies is constant, in apparent contradiction with the results shown in Fig. 5. Our choice of prior would therefore tend to weaken, rather than strengthen, any true decrease in F_{burst} towards high masses. Paper I also demonstrated that the definition and inferred properties of our highconfidence subset of bursty galaxies in insensitive to the adopted prior.



Figure 2. Histograms showing the fraction of galaxies as a function of $D_n(4000)$ in eight different ranges of stellar mass. The numbers in the upper right-hand corner of each panel list, from top to bottom, the fifth, 25th, 50th, 75th and 95th percentiles of the distribution. The number in the lower right-hand corner is the number of galaxies contributing to the histogram.

In Figs 6 and 7 we compare the distributions of H α emissionline equivalent widths for galaxies with F_{burst} (2.5 per cent) > 0 and for galaxies with $F_{\text{burst}}(50 \text{ per cent}) = 0$. In Fig. 6, we have subdivided the two samples into different bins in $D_n(4000)$ in order to take into account any systematic mean age difference between galaxies with and without recent bursts. For galaxies with low values of $D_n(4000)$, there is a tendency for bursty galaxies to have stronger H α emission than non-bursty galaxies with similar break strengths. These are galaxies where the present-day star formation rates are enhanced over other systems of similar mean stellar age (recall that the H α emission is coming from stars less than 10^7 yr in age). For galaxies with high values of $D_n(4000)$, the trend is reversed. These high $D_n(4000)$ galaxies are presumably 'post-starburst' systems that experienced a strong burst some time in the past, but have subsequently stopped forming any stars. In Fig. 7, we compare galaxies with $F_{\text{burst}}(2.5 \text{ per cent}) > 0$ and galaxies with $F_{\text{burst}}(50 \text{ per cent}) = 0$ in fixed bins of stellar mass. Lowmass bursty galaxies are skewed to higher H α equivalent widths, while high-mass bursty galaxies exhibit a wider range in equivalent widths than their counterparts with $F_{burst}(50 \text{ per cent}) = 0$. Because information concerning H α was not used when selecting our sample of bursty galaxies, we regard these differences in emissionline properties as evidence that our method does pick out a sample of galaxies with very strong variations in recent star formation history.

3.2 The dependence of structural parameters on stellar mass

In Fig. 8, we present conditional density distributions for concentration index *C*, *r*-band half-light surface brightness μ 50, and surface mass density μ_* as functions of stellar mass and of *r*-band absolute magnitude. Fig. 8 demonstrates that the structural parameters of galaxies correlate both with absolute magnitude and with stellar mass, but once again we find that the trends are smoother when plotted as a function of stellar mass.

The surface mass density exhibits a strikingly tight correlation with stellar mass. μ_* increases by nearly two orders of magnitude from ~10⁷ M_☉ kpc⁻² for galaxies with $M_* ~ 10^8$ M_☉ to ~10⁹ M_☉ kpc⁻² for the most massive galaxies with $M_* = 10^{12}$ M_☉. In contrast, the *r*-band surface brightness only increases by a factor of 4 as M(r) increases by 8 mag. There is a sharp change in slope in the μ_*-M_* relation at a stellar mass of ~3 × 10¹⁰ M_☉. A transition



Figure 3. The bivariate distribution function of $H\delta_A$ and $D_n(4000)$ in eight ranges of stellar mass.

at the same stellar mass is also seen in the $C-M_*$ relation. Interestingly, this is also the stellar mass at which galaxies switch from low $D_n(4000)$ to high $D_n(4000)$ in Fig. 1. It is important to note that virtually all galaxies in our sample with $M_* > 10^8 \text{ M}_{\odot}$ have $\mu_{1/2} < 23.5 \text{ mag arcsec}^{-2}$. The main spectroscopic survey is complete to surface brightnesses a magnitude fainter than this, so our results should not be biased by surface brightness selection effects.

Fig. 9 presents the size distributions for galaxies in eight disjoint ranges of stellar mass. We plot the fraction of galaxies as a function of the natural logarithm of the half-light radius in the z band. These size distributions are extremely well described by a lognormal function

$$P(R) dR = \frac{1}{\sqrt{2\pi\sigma}} \exp\left\{-\frac{\left[\ln(R/R_{\rm med})\right]^2}{2\sigma^2}\right\} \frac{dR}{R},\tag{1}$$

where $R_{\rm med}$ is the median value of the size distribution and σ is the dispersion in ln *R*. We have indicated the value of σ and $R_{\rm med}$ for the lognormal fit in each panel in Fig. 8. The *shape* of the size distribution is largely independent of mass for galaxies with $M_* < 3 \times 10^{10} \,{\rm M_{\odot}}$ and is well fitted by $\sigma \sim 0.4$ –0.5. The median size $R_{\rm med}$ scales with stellar mass as $R_{\rm med} \propto M_*^{0.18}$. For stellar masses larger than $3 \times 10^{10} \,{\rm M_{\odot}}$, the shape of the size distribution is still lognormal, but the dispersion decreases. $R_{\rm med}$ also increases more rapidly with mass $(\propto M_*^{0.33})$ in this regime.

Figs 10 and 11 show the distributions of concentration index *C* and surface mass density μ_* for these same eight stellar mass ranges. The distribution of *C* is nearly independent of stellar mass for galaxies with $M_* < 3 \times 10^{10} \text{ M}_{\odot}$. The median value of *C* for these systems is around 2.3. Only approximately 10 per cent of low-mass galaxies have C > 2.6, the value that marks the transition from late-type to



Figure 4. The median value of $H\delta_A$ is plotted as a function of $D_n(4000)$ in two different stellar mass ranges. The solid and dotted errorbars indicate the 25th to 75th and fifth to 95th percentile ranges of the distribution of $H\delta_A$. Symbols indicate the locus occupied by galaxies with continuous star formation histories. In the left-hand panel, we have plotted 25 per cent solar models and in the right-hand panel we plot solar metallicity models.



Figure 5. The solid histograms show the distributions of galaxies as a function of median value of the likelihood distribution of F_{burst} in eight stellar mass ranges. The dotted histograms show the distributions of the median value of F_{burst} for the high-confidence bursty galaxies in each mass range. The numbers on each panel indicate the fraction of galaxies in each mass bin with $F_{\text{burst}}(50 \text{ per cent}) > 0$ (upper) and $F_{\text{burst}}(2.5 \text{ per cent}) > 0$ (lower).



Figure 6. The distribution of H α equivalent widths for galaxies with $F_{\text{burst}}[2.5 \text{ per cent}) > 0$ (dashed)] is compared with that for galaxies with $F_{\text{burst}}(50 \text{ per cent}) = 0$ (solid) for four different ranges in the value of the D_n(4000) index. The numbers listed in the panels indicate the number of objects in the bursty sample (top) and in the continuous sample (bottom).

early-type morphologies (Strateva et al. 2001). At larger masses, the distribution shifts to progressively higher concentrations. In our highest-mass bin 90 per cent of galaxies have *C* index values larger than 2.6. Fig. 11 shows that for galaxies with $M_* < 3 \times 10^{10} \text{ M}_{\odot}$, the median surface mass density scales with stellar mass as $\mu_* \propto M_*^{0.63}$. The μ_* distribution is fairly broad and its shape does not appear to depend very strongly on stellar mass. At values of M_* larger than $3 \times 10^{10} \text{ M}_{\odot}$, the scaling of the median value of μ_* with stellar mass becomes much weaker. The shape of the distribution function skews and is eventually strongly peaked at values of μ_* around $10^9 \text{ M}_{\odot} \text{ kpc}^{-2}$.

We note that all structural parameters discussed in this section have been calculated within circular apertures. We have estimated the additional scatter in size resulting from variations in axis ratio and find that this has a very small (<10 per cent) effect on the values of the dispersion σ that we estimate. Because of seeing effects, R50(z) may be overestimated for galaxies with small angular sizes. We have split the galaxies in each stellar mass range into two equal samples according to redshift, and we have compared the size distribution of the nearby sample with that of the full sample. This is illustrated in two of the panels of Fig. 9. As can be seen, the effects appear to be quite small. In addition, we have checked the scalings using R90(z) and find that our conclusions remain unchanged.

3.3 The connection between star formation history and structural parameters

In the previous two subsections we showed that both the star formation histories and the structural parameters of galaxies depend strongly on stellar mass. It is therefore not surprising that the star formation histories of galaxies correlate with μ_* and *C*. This is illustrated in Fig. 12 where we present the conditional density distributions of the two stellar absorption indices $D_n(4000)$ and $H\delta_A$ as functions of surface mass density and of concentration. Galaxies with low surface density have low values of $D_n(4000)$ and galaxies with high surface density have high values of $D_n(4000)$, with a strong transition at $\mu_* \sim 3 \times 10^8 \text{ M}_{\odot} \text{ kpc}^{-2}$. There is also a striking transition in the values of $D_n(4000)$ and $H\delta_A$ at *C* values around 2.6. We note that this corresponds very well to the value of *C* recommended by Strateva et al. (2001) for optimum separation between early and late Hubble types, based on an analysis of a small sample of galaxies classified by eye.

We now ask whether it is possible to pinpoint the primary cause of variations in the star formation history of galaxies - do galaxies have low values of $D_n(4000)$ and high values of $H\delta_A$ because their masses are small or because their surface mass densities are low? We attempt to address this question in Fig. 13 where we plot the fractions of galaxies with $D_n(4000) > 1.55$ and with $H\delta_A < 2.0$ as a function of surface mass density in narrow ranges of stellar mass (left-hand panels), and as a function of stellar mass in narrow ranges of surface mass density (right-hand panels). We only plot these fractions for values of M_* and μ_* where there are more than 100 galaxies per bin. $D_n(4000) = 1.55$ and $H\delta_A = 2$ are chosen as natural places to divide the sample, because this is where clear transitions occur in the $D_n(4000)/C$ and $H\delta_A/C$ plots in Fig. 12. (Note that the C index is determined directly from the photometric data, independently of both stellar mass and surface mass density.)



Figure 7. The distribution of H α equivalent widths for galaxies with $F_{burst}(2.5 \text{ per cent}) > 0$ (dashed) is compared with that for galaxies with $F_{burst}(50 \text{ per cent}) = 0$ (solid) for four different ranges in stellar mass. The numbers listed in the panels indicate the number of objects in the bursty sample (top) and in the continuous sample (bottom).

Fig. 13 shows that the two indices are more fundamentally related to μ_* than to M_* . The fraction of galaxies with $D_n(4000) > 1.55$ or with $H\delta_A < 2$ is largely independent of stellar mass at a given value of μ_* , although the most massive galaxies with $M_* > 10^{11}$ M_{\odot} clearly deviate in that they have larger values of $D_n(4000)$ and smaller values of $H\delta_A$ at given μ_* . At given stellar mass, however, the fraction of galaxies with $D_n(4000) > 1.55$ or with $H\delta_A < 2$ is a strongly increasing function of μ_* . These results are in qualitative agreement with those of Bell & de Jong (2000), who studied trends in the optical and infrared colours of a sample of nearby *spiral* galaxies and found that the colours most sensitive to star formation history correlate best with surface density. It also fits in well with numerous studies that have demonstrated a clear correlation between star formation *rate* and gas surface density for samples of spirals (e.g. Kennicutt 1983; Wong & Blitz 2002).

3.4 Summary

In the previous subsections, we demonstrated that there are strong correlations between the star formation histories, stellar masses and structural parameters of galaxies. Here we provide a summary of our results.

(i) Low-mass galaxies have low values of $D_n(4000)$ and high values of $H\delta_A$, indicative of young stellar populations. High-mass galaxies have high values of $D_n(4000)$ and low values of $H\delta_A$, indicative of old stellar populations. A sharp transition from young to old occurs at a stellar mass of around $3 \times 10^{10} M_{\odot}$.

(ii) Low-mass galaxies have low surface mass densities and low concentrations. High-mass galaxies have high surface mass den-

sities and high concentrations. An abrupt change in the slopes of the μ_*-M_* and $C-M_*$ relations occurs at a stellar mass of 3×10^{10} M_☉.

(iii) The size distribution function of galaxies less massive than $3 \times 10^{10} \text{ M}_{\odot}$ is well characterized by a lognormal distribution with dispersion $\sigma = 0.4$ –0.5. Galaxies more massive than $3 \times 10^{10} \text{ M}_{\odot}$ also have a lognormal size distribution, but the dispersion is smaller.

(iv) Below the transition mass the median size of galaxies increases as $M_*^{0.18}$. Above the transition mass the increase is more rapid, $R \propto M_*^{0.3}$.

(v) Galaxies with low surface mass densities and low concentrations have $D_n(4000)$ and $H\delta_A$ values indicative of young stellar populations. Galaxies with high surface mass densities and high concentrations have $D_n(4000)$ and $H\delta_A$ values indicative of old stellar populations. The transition from young to old stellar populations occurs at $\mu_* \sim 3 \times 10^8 \text{ M}_{\odot} \text{ kpc}^{-2}$ and $C \sim 2.6$.

(vi) The star formation histories of low-mass galaxies, as traced by $D_n(4000)$ and $H\delta_A$, appear to be more fundamentally related to surface mass density than to stellar mass.

(vii) The fraction of galaxies that have experienced recent starbursts decreases strongly with increasing stellar mass (and increasing surface mass density).

4 INTERPRETATION AND DISCUSSION

Our analysis has demonstrated that there is a sharp transition in the physical properties of galaxies at a stellar mass of $\sim 3 \times 10^{10} M_{\odot}$.



Figure 8. Conditional density distributions showing trends in the structural parameters μ_* , $\mu_{1/2}$ and C = R90/R50 as a function the logarithm of stellar mass and as a function of *r*-band absolute magnitude. Galaxies have been weighted by $1/V_{\text{max}}$ and the bivariate distribution function has been normalized to a fixed number of galaxies in each bin of log M_* and of *r*-band absolute magnitude. The line in the bottom left-hand panel indicates the surface brightness completeness limit of the SDSS survey.

Galaxies less massive than this have low surface mass densities, low concentration indices typical of discs and young stellar populations. More massive galaxies have high surface mass densities, high concentration indices typical of bulges, and predominantly old stellar populations.

The size distribution of disc galaxies is of particular interest, because it can be compared directly to simple theoretical models. In a classic paper in 1980, Fall & Efstathiou considered the formation of a disc by the condensation of gas in a gravitationally dominant dark matter halo. If the gas initially has the same specific angular momentum as the dark matter and conserves its angular momentum during its contraction, then the characteristic radius of the disc will scale as

$$R_{\rm d} \propto \lambda R_{\rm halo},$$
 (2)

where R_{halo} is the virial radius of the halo. The spin parameter λ is defined as

$$\lambda = J|E|^{1/2}G^{-1}M_{\text{halo}}^{-5/2},\tag{3}$$

where E is the total energy of the halo, J is its total angular momentum and M is its total mass (see Mo, Mao & White 1998 for a detailed analytic model).

N-body simulations show that in hierarchical clustering cosmologies, the distribution of λ for dark matter haloes has a lognormal form with parameters that depend very weakly on cosmology or on halo mass (Barnes & Efstathiou 1987; Warren et al. 1992; Cole & Lacey 1996; Lemson & Kauffmann 1999) with typical values $\lambda_{med} \simeq 0.04$ and $\sigma_{\lambda} \simeq 0.5$. The fact that for low-mass galaxies the shape and width of the size distributions in Fig. 9 agree so well with this simple theory is truly remarkable. We note that our results are in qualitative agreement with those of Syer, Mao & Mo (1999), who find that the distribution of spin parameter λ inferred from a sample of nearby disc galaxies is in excellent agreement with the predictions of cosmological simulations. On the other hand, our results appear to disagree with those of de Jong & Lacey (2000), who find that the distribution of disc sizes is narrower than predicted by simple models. We note that neither of these two studies addressed possible variations in the size distributions of disc galaxies as a function of luminosity or of mass.

It is not obvious how the stellar mass of a galaxy should scale with the mass of its dark matter halo. For simplicity, one might assume that a fixed fraction of the baryonic mass ends up in disc stars, in which case $M_* \propto R_{halo}^3$ and $R_d \propto M_*^{1/3}$ (this assumes that baryons



Figure 9. Solid histograms show the fraction of galaxies as a function of the natural logarithm of the half-light radius in the *z* band [R50(z)] in eight different ranges of stellar mass. The curves show lognormal fits to the distribution and the numbers in each panel are the best-fitting values of σ (top) and R_{med} (bottom). The dotted histograms in the panels with $9 < \log M_* < 9.5$ and $11 < \log M_* < 11.5$ are for the subsample of galaxies with $V/V_{max} < 0.5$.

follow the dark matter, at least initially). From Fig. 9, we find that median size increases *more slowly* than predicted. This means that the stellar mass fraction must *increase* with halo mass.

We now derive the scaling between the stellar mass of a disc galaxy and the mass of its dark matter halo from the *slope* of the μ_*-M_* relation. If we adopt the assumptions of the simple model, we can write

$$\mu_* \propto M_* / R_d^2 \propto \epsilon M_{\text{halo}} / R_{\text{halo}}^2 \propto \epsilon M_{\text{halo}}^{1/3} \propto \epsilon^{2/3} M_*^{1/3}, \qquad (4)$$

where $\epsilon \propto M_*/M_{halo}$ is the efficiency with which the disc has converted the available gas into stars. The Fall & Efstathiou arguments are only directly relevant to the formation of disc galaxies, while the conditional density distributions shown in Fig. 9 are derived using all galaxies in the sample. At stellar masses greater than $10^{10} \,\mathrm{M_{\odot}}$, a large fraction of the galaxies are bulge-dominated systems. In order to clarify the scaling properties of our two principal classes of galaxy, we define two subsamples that eliminate most of the mixed systems: (i) a sample of late-type galaxies with C < 2.6 and $D_n(4000) < 1.55$ (41 312 galaxies); and (ii) a sample of

early-type galaxies with C > 2.6 and $D_n(4000) > 1.55$ (53119 galaxies). Fig. 12 demonstrates that these cuts divide young, disc-dominated systems from old, bulge-dominated systems surprisingly cleanly.

The μ_*-M_* relation of our late-type subsample is shown in the left-hand panel of Fig. 14, while the right-hand panel shows the corresponding relation for early types. It is striking that the μ_*-M_* relation for the late-type subsample can be described quite well by a single power law. The μ_*-M_* relation for earlytype galaxies still appears to undergo a transition in slope. (Note that there are few late-type galaxies with stellar masses above 2×10^{11} M_{\odot} and few early types with M_* below 10^{10} M_{\odot}.) Fig. 14 also shows that early-type galaxies of given stellar mass have higher stellar surface densities than late-type galaxies of the same mass. This difference is largest for the lower-mass objects and disappears for the most massive systems in our samples.

Fitting to the dominant part of the population of late-type galaxies (i.e. over the stellar mass range 10^8-10^{11} M_{\odot}), we find $\mu_* \propto$



Figure 10. Histograms showing the distribution of galaxies as a function of *C* index for eight different ranges of stellar mass. The numbers in each panel list, from top to bottom, the fifth, 25th, 50th, 75th and 95th percentiles of the distribution.

 $M_*^{0.535\pm0.03}$, so that $\epsilon \propto M_*^{0.30\pm0.05} \propto M_{\rm halo}^{0.43\pm0.09}$ and $R_{\rm d} \propto M_*^{0.23\pm0.02}$. The star formation efficiency for disc galaxies apparently increases with halo mass. Such an increase is expected in a model where supernova feedback controls the dynamics of the interstellar medium (e.g. McKee & Ostriker 1977; Efstathiou 2000). In most feedback models, supernovae inhibit the formation of stars in low-mass haloes, leading to small values of M_*/M_{halo} (e.g. White & Frenk 1991; Kauffmann, White & Guiderdoni 1993; Cole et al. 1994). However, because we do not have any information concerning the gas fractions of the galaxies in our sample, we cannot say whether the suppression of star formation occurs because gas is expelled from the galaxy and from its halo, as is often assumed in these models, or because the conversion of gas into stars is slowed in low-mass systems. It is interesting that Tremonti et al. (in preparation) find the gas phase metallicity of low-mass galaxies to increase with stellar mass as $Z \propto M_{*}^{0.4-0.5}$ (depending on which metallicity indicator they use), suggesting a similar dependence of overall star formation efficiency on stellar mass.

For early-type galaxies, we fit over the range $3 \times 10^{10} - 3 \times 10^{12}$ M_{\odot} and obtain $\mu_* \propto M_*^{0.0\pm 0.02}$. If we assume that $R_d \propto R_{halo}$ for these galaxies, this would imply that $\epsilon \propto M_{\star}^{-1/3}$, i.e. the star formation efficiency for ellipticals decreases with halo mass. Clues as to why this is the case may come from recent Chandra and XMM observations, which show that cooling of gas to temperatures below 1 keV is suppressed at the centres of rich clusters, where the most massive ellipticals are located (e.g. Bohringer et al. 2001; Fabian et al. 2001). It seems very probable that the assembly of bulge-dominated galaxies involved chaotic processes such as mergers or inhomogeneous, rapidly star-forming collapse, which, for most systems, must have completed well before the present day. Under these circumstances, it is not clear that our assumed scaling between galaxy radius and dark matter halo radius is valid. Even if the simple scaling laws were to apply, the halo sizes and surface densities relevant for early-type galaxies are probably to be defined at higher effective redshift than those relevant for late-type galaxies. This may explain why early-type galaxies in Fig. 14 have higher surface densities than late-type galaxies of similar mass. The scatter in surface density (and size) is also smaller for early-type than for late-type galaxies, as is very evident in Fig. 14 and in Figs 8 and 9.



Figure 11. Histograms showing the distribution of galaxies as a function of log μ_* for eight different ranges in stellar mass. The numbers in each panel list, from top to bottom, the fifth, 25th, 50th, 75th and 95th percentiles of each distribution.

We note that the Lyman break galaxies (LBGs) at high redshift can plausibly be identified with the progenitors of early-type galaxies. Pettini et al. (2001) have derived a mean effective dynamical mass of 1.3×10^{10} M_{\odot} within a mean half-light radius of 2.3 kpc for a sample of LBGs at $z \sim 3$. Shapley et al. (2001) and Papovich, Dickinson & Ferguson (2001) have estimated the stellar masses of LBGs using their observed rest-frame ultraviolet and optical spectral energy distributions. These agree well with the dynamical mass estimates. Lyman break galaxies thus have characteristic surface mass densities of $\sim 10^9$ M_{\odot} kpc⁻², quite similar to those of presentday early-type galaxies. Moreover, if the star formation associated with the Lyman break galaxies is integrated from $z \sim 2$ to 6, one finds that they could have produced most of the stars in present-day early-type galaxies (Steidel et al. 1999; Thompson, Weymann & Storrie-Lombardi 2001).

Within the framework of hierarchical models of galaxy formation, the shape of the size distributions of low-mass galaxies is easy to understand, but the trend of younger mean stellar ages with decreasing stellar mass is puzzling, since lower-mass dark matter haloes typically form at earlier epochs. If all the gas cooled off and formed stars when these haloes were assembled, lower-mass galaxies would have older stellar populations (e.g. Kauffmann et al. 1993). We have seen, however, that non-gravitational processes such as supernova feedback may play an important role in regulating the rate at which stars form in these systems. To explain the observations, these processes must slow the conversion of gas into stars so that low-mass systems remain gas rich until the present day and have low overall efficiencies for turning gas into stars. We have seen that the star formation histories of low-mass galaxies are correlated more strongly with surface density than with stellar mass. If feedback processes are able to prevent the gas in low surface density galaxies from dissipating to form dense molecular clouds, then their star formation time-scales will be long.

Conversely, it may appear puzzling that star formation has terminated in the majority of high-mass galaxies, because the hierarchical model claims that these systems were assembled relatively late. In order to explain the observations, star formation must have been both efficient and rapid in the progenitors of these systems, so that



Figure 12. Conditional density distributions showing trends in the stellar age indicators $D_n(4000)$ and $H\delta_A$ as functions of the logarithm of the surface mass density μ_* and of the concentration index *C*.

the final stages of their assembly involved rather little dissipation or star formation. In current models, galaxies that fall into groups and clusters are stripped of their reservoir of gas. As a result, their star formation rates decline and their colours redden (see, for example, Diaferio et al. 2001). This is unlikely to explain why the *majority* of massive galaxies do not form stars at the present day, because many do not reside in rich groups and clusters. In such systems, gas must somehow be prevented from cooling and forming stars after the bulge has formed. Again, the suppressed cooling at the centres of observed galaxy clusters may offer a clue to the mechanism responsible for this.

Finally, it will be important to understand the physical origin of the characteristic mass scale that is so strikingly imprinted on the galaxy population. The structural properties and formation paths of dark matter haloes vary smoothly with mass. The abrupt transition in the stellar ages, star formation histories and structure of galaxies that we find at $3 \times 10^{10} \text{ M}_{\odot}$ can only be explained by star formation and/or feedback processes. We believe that this tension between simplified phenomenological models of galaxy evolution and the trends seen in the data will lead to new insights into the physical processes that regulate how galaxies form and evolve. We intend to address these issues in more detail in future work.

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Figure 13. Top left: the fraction of galaxies with $D_n(4000) > 1.55$ as a function of log μ_* in fixed ranges in stellar mass. Solid, short-dashed, long-dashed and dotted lines are for galaxies with log M_* in the ranges 8–9, 9–10, 10–11 and 11–12, respectively. Top right: the fraction of galaxies with $D_n(4000) > 1.55$ as a function of log M_* in fixed ranges in stellar surface mass density. Solid, short-dashed, long-dashed and dotted lines are for galaxies with log μ_* in the ranges 7.0–7.8, 7.8–8.3, 8.3–8.8 and 8.8–9.3, respectively. Bottom left: the fraction of galaxies with $H\delta_A < 2$ as a function of log μ_* in the same ranges in M_* . Bottom right: the fraction of galaxies with $H\delta_A < 2$ as a function of log M_* in the same ranges in M_* .



Figure 14. Conditional density distributions showing trends in μ_* as a function the logarithm of stellar mass for galaxies with $D_n(4000) < 1.55$ and C < 2.6 (left) and with $D_n(4000) > 1.55$ and C > 2.6 (right).

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