

The Dirac Operator on Nilmanifolds and Collapsing Circle Bundles

by

Bernd Ammann and Christian Bär

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Abstract

We compute the spectrum of the Dirac operator on 3-dimensional Heisenberg manifolds. The behavior under collapse to the 2-torus is studied. Depending on the spin structure either all eigenvalues tend to $\pm\infty$ or there are eigenvalues converging to those of the torus. This is shown to be true in general for collapsing circle bundles with totally geodesic fibers. Using the Hopf fibration we use this fact to compute the Dirac eigenvalues on complex projective space including the multiplicities.

Finally, we show that there are 1-parameter families of Riemannian nilmanifolds such that the Laplacian on functions and the Dirac operator for certain spin structures have constant spectrum while the Laplacian on 1-forms and the Dirac operator for the other spin structures have nonconstant spectrum. The marked length spectrum is also constant for these families.

Keywords: Dirac operator, nilmanifolds, Heisenberg manifolds, circle bundles, collapse, isospectral deformation

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0 Introduction

The Dirac operator on a Riemannian spin manifold is a natural elliptic first order differential operator. If the underlying manifold is closed, then the spectrum is discrete. In order to understand the relation of the Dirac eigenvalues and geometric data it is desirable to have examples for which one can explicitly compute the spectrum. These examples can be used to test conjectures. For example,

they show that the spectrum does not determine the manifold, not even topologically. There are examples of Dirac isospectral pairs of spherical space forms with different fundamental groups [3].

In the following table we collect those closed Riemannian spin manifolds known to us for which the Dirac spectrum has been computed.

\mathbb{R}^n/Γ	flat tori	[15]
S^n	spheres of constant curvature	[33], [3], [34, 35], [11]
S^n/Γ	spherical space forms	[3]
S^{2m+1}	spheres with Berger metrics	[20] for $m = 1$ [4] for general m
S^3/\mathbb{Z}_k	3-dim. lens spaces with Berger metric	[2]
G	simply connected compact Lie groups	[14]
$\mathbb{C}P^{2m-1}$	complex projective spaces	[9, 10], [30]
$\mathbb{H}P^m$	quaternionic projective spaces	[8] for $m = 2$ [22] for general m
$Gr_2(\mathbb{R}^{2m})$	real Grassmannians	[31] for $m = 3$ [32] for general m
$Gr_{2p}(\mathbb{R}^{2m})$	real Grassmannians	[28, 29]
$Gr_2(\mathbb{C}^{m+2})$	complex Grassmannians	[23]
$G_2/SO(4)$		[28, 29]

We compute the Dirac spectrum for another class of Riemannian spin manifolds, for the 3-dimensional *Heisenberg manifolds*. They are of the form $M = \Gamma \backslash G$ where G is the 3-dimensional Heisenberg group and Γ is a cocompact lattice. The metric comes from a left invariant metric on G . We restrict to three dimensions mostly for simplicity. The methods also work in higher dimensions.

After collecting some facts from spin geometry and from the representation theory of nilpotent groups in the first two sections we perform the computation of the Dirac eigenvalues in Section 3 (Theorem 3.1). Our Heisenberg manifolds are circle bundles over the flat 2-torus. The explicit formulas for the Dirac eigenvalues allow us to study their behavior under collapse to the 2-torus, i.e. we let the length of the fibers tend to zero. It turns out that the behavior of the eigenvalues depends essentially on the spin structure. For some spin structures all eigenvalues tend to $\pm\infty$. For the other spin structures most of the eigenvalues also tend to $\pm\infty$ but some converge to exactly the eigenvalues of the torus (Corollary 3.3).

As it turns out in Section 4 this is a general fact for the Dirac operator on the total space of a collapsing circle bundle with totally geodesic fibers over any Riemannian manifold. There are two types of spin structures which we call *projectable* and *nonprojectable*. In the nonprojectable case all eigenvalues tend

to $\pm\infty$ (Theorem 4.5) whereas in the projectable case there are also eigenvalues converging to exactly the eigenvalues of the base (Theorem 4.1).

As an example we apply this to the Hopf fibration over the complex projective space $\mathbb{C}P^m$. Using explicit formulas for the Dirac eigenvalues on Berger spheres [4] we see that for m even (when $\mathbb{C}P^m$ is nonspin) all eigenvalues go to $\pm\infty$. If m is odd (when $\mathbb{C}P^m$ is spin) the limits of the eigenvalues give us a new computation of the Dirac spectrum of $\mathbb{C}P^m$ confirming (and simplifying) the results in [9, 10, 30]. We also obtain formulas for the multiplicities (Theorem 4.6) which have not been computed before.

In the last section we discuss the change of the spectrum of the Dirac operator on certain pairs or continuous families of Riemannian nilmanifolds. A *Riemannian nilmanifold* is a nilmanifold $\Gamma\backslash G$ endowed with a Riemannian metric whose lift to G is left invariant. Here we assume that G is a simply connected nilpotent Lie group and that Γ is a cocompact discrete subgroup. This generalizes Heisenberg manifolds. Besides the Dirac spectrum we look at several different spectral invariants which have already been studied by other authors: the spectrum of the Laplacian on functions, the spectrum of the Laplacian on forms, the length spectrum, i.e. the set of all lengths of closed geodesics, and the marked length spectrum, i.e. the set of all lengths of closed geodesics in each free homotopy class.

The questions behind the study of these examples are: “How far is the geometry on a Riemannian nilmanifold determined by these spectral invariants?” and “How are these invariants related to each other?” The class of nilmanifolds provides us with many examples for the discussion of these questions.

Milnor gave in [24] the first example of a pair of Riemannian manifolds (16-dimensional tori) which are nonisometric but isospectral for the Laplace operator on functions.

Gordon and Wilson [18] constructed a continuous family of nonisometric Riemannian 2-step nilmanifolds with the same spectrum of the Laplacian on functions and forms. This is the first example of a nontrivial Laplace-isospectral deformation. This family and other deformations of the Gordon-Wilson type, i.e. via almost inner automorphisms, do not only have the same spectrum for the Laplacian on functions and on forms, but also the same marked length spectrum. In the 2-step case with a left invariant metric Eberlein [13, Theorem 5.20] showed that two 2-step nilmanifolds having the same marked length spectrum arise via almost inner automorphisms and therefore have the same spectrum of the Laplacian on functions and on forms.

This statement is false if we drop the condition “2-step”. Gornet [17] constructed two families of deformations of 3-step Riemannian nilmanifolds which

are isospectral for the Laplacian on functions and for the marked length spectrum, but not for the spectrum of the Laplacian on 1-forms. We will show that these families are Dirac isospectral for the nonprojectable spin structures while they have nonconstant Dirac spectrum for the projectable spin structures.

1 Review of Spin Geometry

In this section we collect a few facts on spin manifolds and the Dirac operator which we will use later on. For details the reader may consult [6] or [21]. Let M be an oriented Riemannian manifold of dimension n . Let $P_{\text{SO}}(M)$ be the set of all oriented orthonormal bases of all tangent spaces of M . Obviously, $P_{\text{SO}}(M)$ is an $\text{SO}(n)$ -principal bundle. A *spin structure* on M consists of a $\text{Spin}(n)$ -principal bundle $P_{\text{Spin}}(M)$ over M together with a twofold covering map $\varphi : P_{\text{Spin}}(M) \rightarrow P_{\text{SO}}(M)$ such that the diagram

$$\begin{array}{ccc}
 P_{\text{Spin}}(M) \times \text{Spin}(n) & \rightarrow & P_{\text{Spin}}(M) \\
 \downarrow \varphi \times \Theta & & \downarrow \varphi \\
 P_{\text{SO}}(M) \times \text{SO}(n) & \rightarrow & P_{\text{SO}}(M)
 \end{array}
 \begin{array}{c}
 \searrow \\
 M \\
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 \end{array}$$

commutes. Here $\Theta : \text{Spin}(n) \rightarrow \text{SO}(n)$ is the standard twofold covering map and the horizontal arrows are given by the principal bundle structure.

A manifold which admits such a spin structure is called a *spin manifold*. Being spin is a global condition which is equivalent to the vanishing of the second Stiefel-Whitney class of the tangent bundle, $w_2(M) = 0$. The number of different spin structures, if any, is given by $\#H^1(M, \mathbb{Z}_2)$. In particular, a simply connected manifold has at most one spin structure.

Let us from now on assume that M is spin and that a spin structure on M is fixed. There is a unitary representation Σ_n of $\text{Spin}(n)$ of dimension $2^{\lfloor n/2 \rfloor}$, called the *spinor representation*. If n is odd, then Σ_n is irreducible. The associated hermitian vector bundle $\Sigma M := P_{\text{Spin}}(M) \times_{\text{Spin}(n)} \Sigma_n$ is called the *spinor bundle*.

The Levi-Civita connection on M induces a natural connection ∇^Σ on ΣM . It can be described as follows. Let e_1, \dots, e_n be a local orthonormal tangent frame defined over $U \subset M$. Let $\Gamma_{ij}^k : U \rightarrow \mathbb{R}$ be the corresponding Christoffel symbols, $\nabla_{e_i} e_j = \sum_{k=1}^n \Gamma_{ij}^k e_k$. Then (e_1, \dots, e_n) is a local section of $P_{\text{SO}}(M)$. Let q be a lift to $P_{\text{Spin}}(M)$, i.e. $\varphi \circ q = (e_1, \dots, e_n)$. Then q defines a trivialization of ΣM over

$U, \Sigma M|_U = U \times \Sigma_n$, with respect to which we have the following formula for ∇^Σ :

$$\nabla_{e_i}^\Sigma \sigma = \partial_{e_i} \sigma + \frac{1}{4} \sum_{j,k=1}^n \Gamma_{ij}^k \gamma(E_j) \gamma(E_k) \sigma \quad (1.1)$$

Since the spinor representation of $\text{Spin}(n)$ extends to a representation of the Clifford algebra $\text{Cl}(n)$ there is a well-defined *Clifford multiplication* $T_p M \otimes_{\Sigma_p} M \rightarrow \Sigma_p M$, $X \otimes \sigma \rightarrow \gamma(X)\sigma$. It satisfies the relations $\gamma(X)\gamma(Y) + \gamma(Y)\gamma(X) + 2\langle X, Y \rangle \text{id} = 0$.

For example in dimension $n = 3$, one can choose a basis for Σ_3 with respect to which

$$\gamma(E_1) = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}, \gamma(E_2) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \gamma(E_3) = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix},$$

where E_1, E_2, E_3 denotes the standard basis of $\mathbb{R}^3 \subset \text{Cl}(3)$.

The Dirac operator D acts on the sections of ΣM . It is defined by $D\sigma = \sum_{i=1}^n \gamma(e_i) \nabla_{e_i}^\Sigma \sigma$ where e_1, \dots, e_n is any orthonormal basis of the tangent space. If one has, in addition, a complex vector bundle E over M equipped with a connection one can form the *twisted Dirac operator* D^E acting on sections of $\Sigma M \otimes E$ by using the tensor product connection and Clifford multiplication on the first factor.

The Dirac operator is a formally self adjoint elliptic differential operator of first order. If the manifold M is closed, then D has discrete real spectrum.

Let G be a simply connected Lie group with a left invariant metric. We regard the elements of the Lie algebra \mathfrak{g} as left invariant vector fields. The choice of an oriented orthonormal basis of \mathfrak{g} yields a trivialization of the frame bundle $P_{\text{SO}}(G) = G \times \text{SO}(n)$. The unique spin structure can be written as $P_{\text{Spin}}(G) = G \times \text{Spin}(n)$, $\varphi = \text{id} \times \Theta$. Spinor fields are then simply maps $G \rightarrow \Sigma_n$.

Let $\Gamma \subset G$ be a lattice. Spin structures of $M = \Gamma \backslash G$ correspond to homomorphisms $\epsilon : \Gamma \rightarrow \mathbb{Z}/2\mathbb{Z} = \{-1, 1\}$. The corresponding spin structure is given by $P_{\text{Spin}, \epsilon}(\Gamma \backslash G) = G \times_{\Gamma} \text{Spin}(n)$ where $g_0 \in \Gamma$ acts on G by left multiplication and on $\text{Spin}(n)$ by multiplication with the central element $\epsilon(g_0)$. Spinor fields on M can then be identified with ϵ -equivariant maps $\sigma : G \rightarrow \Sigma_n$, i.e. $\sigma(g_0 g) = \epsilon(g_0) \sigma(g)$ for all $g \in G$, $g_0 \in \Gamma$. Denote the corresponding Hilbert space of square integrable spinors by $L^2(\Sigma_\epsilon M)$.

2 Review of Kirillov Theory

In this section we summarize some facts from the representation theory of nilpotent groups which will be of importance for our study of nilmanifolds. For details

see [12]. Let G be a simply connected nilpotent Lie group with Lie algebra \mathfrak{g} . The exponential map $\exp : \mathfrak{g} \rightarrow G$ is a global diffeomorphism whose inverse we denote by $\log : G \rightarrow \mathfrak{g}$.

Example. Let

$$G = \left\{ \left(\begin{array}{ccc} 1 & x & z \\ 0 & 1 & y \\ 0 & 0 & 1 \end{array} \right) \middle| x, y, z \in \mathbb{R} \right\}$$

be the 3-dimensional *Heisenberg group*. Its Lie algebra is

$$\mathfrak{g} = \left\{ \left(\begin{array}{ccc} 0 & x & z \\ 0 & 0 & y \\ 0 & 0 & 0 \end{array} \right) \middle| x, y, z \in \mathbb{R} \right\},$$

the 3-dimensional *Heisenberg algebra*. Write

$$g(x, y, z) := \begin{pmatrix} 1 & x & z \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix} \quad \text{and}$$

$$X(x, y, z) := \begin{pmatrix} 0 & x & z \\ 0 & 0 & y \\ 0 & 0 & 0 \end{pmatrix}.$$

Then the exponential map is given by

$$\begin{aligned} \exp(X(x, y, z)) &= g\left(x, y, z + \frac{1}{2}x \cdot y\right) \quad \text{and its inverse by} \\ \log(g(x, y, z)) &= X\left(x, y, z - \frac{1}{2}x \cdot y\right). \end{aligned}$$

Define generators of \mathfrak{g} by

$$\begin{aligned} X &:= X(1, 0, 0), \\ Y &:= X(0, 1, 0), \\ Z &:= X(0, 0, 1). \end{aligned}$$

Then $[X, Y] = Z$ is the only nontrivial commutator and Z spans the center of \mathfrak{g} .

Any Lie group acts on its Lie algebra via the *adjoint representation*

$$\text{Ad} : G \rightarrow \text{End}(\mathfrak{g}), \quad \text{Ad}_g(Y) = \left. \frac{d}{dt} (g \cdot \exp(tY) \cdot g^{-1}) \right|_{t=0}$$

and on the dual space \mathfrak{g}^* via the *coadjoint representation*

$$\text{Ad}^* : G \rightarrow \text{End}(\mathfrak{g}^*), \quad (\text{Ad}_g^* l)(Y) = l(\text{Ad}_{g^{-1}} Y), \quad g \in G, Y \in \mathfrak{g}, l \in \mathfrak{g}^*.$$

The orbits for the coadjoint representation are called *coadjoint orbits*. Given any $l \in \mathfrak{g}^*$ there exists, in the nilpotent case, a maximal isotropic subspace $\mathfrak{m} \subset \mathfrak{g}$ for

the antisymmetric bilinear form $(X, Y) \rightarrow l([X, Y])$ which is also a subalgebra. Such subalgebras are called *polarizing subalgebras*.

Example. Let G be the 3-dimensional Heisenberg group. Let $l \in \mathfrak{g}^*$. We have to distinguish two cases depending on whether l vanishes on the center or not.

Case 1: $l(Z) = 0$.

$$\begin{aligned} \left(\text{Ad}_{g(x,y,z)}^* l\right)(X(\xi, \eta, \zeta)) &= l\left(\text{Ad}_{g(-x,-y,-z+xy)} X(\xi, \eta, \zeta)\right) \\ &= l(X(\xi, \eta, \zeta + \xi y - x\eta)) \\ &= l(\xi \cdot X + \eta \cdot Y + (\zeta + \xi y - x\eta) \cdot Z) \\ &= l(\xi \cdot X + \eta \cdot Y) \\ &= l(X(\xi, \eta, \zeta)). \end{aligned}$$

Hence $\text{Ad}_g^* l = l$ for all $g \in G$, i.e. the orbit of l consists of l only. The bilinear form $(X_1, X_2) \rightarrow l([X_1, X_2])$ vanishes in this case. Hence $\mathfrak{m} = \mathfrak{g}$ is the unique polarizing subalgebra.

Case 2: $l(Z) = \tau \neq 0$.

Write $l = \alpha \cdot X^* + \beta \cdot Y^* + \tau \cdot Z^*$ where $X^*, Y^*, Z^* \in \mathfrak{g}^*$ is the basis dual to X, Y, Z . The computation above shows

$$\begin{aligned} \text{Ad}_{g(x,y,z)}^* l &= (\alpha + \tau y) \cdot X^* + (\beta - \tau x) \cdot Y^* + \tau \cdot Z^* \\ &= l + \tau y \cdot X^* - \tau x \cdot Y^*. \end{aligned}$$

Hence the coadjoint orbit through l is the affine hyperplane through l spanned by X^* and Y^* . As a polarizing subalgebra we can choose e.g. $\mathfrak{m} = \mathbb{R} \cdot X \oplus \mathbb{R} \cdot Z$ or $\mathfrak{m} = \mathbb{R} \cdot Y \oplus \mathbb{R} \cdot Z$.

For a general simply connected nilpotent Lie group G pick an $l \in \mathfrak{g}^*$. We will construct a continuous unitary representation of G associated to l . Choose a polarizing subalgebra $\mathfrak{m} \subset \mathfrak{g}$ for l . Let $M := \exp(\mathfrak{m})$ be the corresponding subgroup of G . By the definition

$$\tilde{\rho}_l(\exp(X)) := e^{2\pi i l(X)}$$

we obtain a well-defined 1-dimensional unitary representation $\tilde{\rho}_l$ of M . Induction yields a continuous unitary representation ρ_l of G , $\rho_l = \text{Ind}(M \uparrow G, \tilde{\rho}_l)$. One checks that ρ_l does not depend (up to equivalence) upon the choice of \mathfrak{m} , that ρ_{l_1} is equivalent to ρ_{l_2} if and only if l_1 and l_2 are in the same coadjoint orbit and that every irreducible unitary representation of G is equivalent to ρ_l for some l . In other words, Kirillov theory sets up a bijection between the set of coadjoint orbits and \hat{G} , the set of equivalence classes of irreducible unitary representations of G .

Example. Let G be the 3-dimensional Heisenberg group, let $l = \alpha X^* + \beta Y^* + \tau Z^* \in \mathfrak{g}^*$. Again, we have to distinguish the two cases $\tau = 0$ and $\tau \neq 0$.

Case 1. $\tau = 0$.

In this case $M = G$, hence $\rho_l = \tilde{\rho}_l$ is a 1-dimensional irreducible representation of G . Since coadjoint orbits of those l with $\tau = 0$ are just points the corresponding representations are parametrized by the two parameters $\alpha, \beta \in \mathbb{R}$, $\rho_l =: \rho_{\alpha, \beta}$, where

$$\begin{aligned} \rho_{\alpha, \beta}(g(x, y, z)) &= e^{2\pi i(\alpha X^* + \beta Y^* + \tau Z^*)(X(x, y, z - \frac{1}{2}xy))} \\ &= e^{2\pi i(\alpha x + \beta y)} \end{aligned} \quad (2.1)$$

Case 2. $\tau \neq 0$.

Since the representation ρ_l does not change if we replace l by another linear form in its coadjoint orbit we may assume $l = \tau \cdot Z^*$. Choose the polarizing subalgebra $\mathfrak{m} = \mathbb{R} \cdot Y \oplus \mathbb{R} \cdot Z$.

The representation space of the induced representation ρ_l is the space of L^2 -sections over the homogeneous space G/M in the line bundle associated with $\tilde{\rho}_l$, the Hilbert space $L^2(G/M, G \times_{\tilde{\rho}_l} \mathbb{C})$. Sections in $G \times_{\tilde{\rho}_l} \mathbb{C}$ can be identified with $\tilde{\rho}_l$ -equivariant maps $f : G \rightarrow \mathbb{C}$, i.e. $f(gm^{-1}) = \tilde{\rho}_l(m)f(g)$ for all $m \in M$, $g \in G$. There is a diffeomorphism $\mathbb{R} \rightarrow G/M$, $t \rightarrow [g(t, 0, 0)]$, where $[\cdot]$ denotes the equivalence class in G/M . Its inverse is simply given by $[g(x, y, z)] \rightarrow x$. Using this diffeomorphism we can construct out of a complex valued L^2 -function $u \in L^2(\mathbb{R}, \mathbb{C})$ a $\tilde{\rho}_l$ -equivariant map $f_u : G \rightarrow \mathbb{C}$ by the formula

$$\begin{aligned} f_u(g(x, y, z)) &= f_u(g(x, 0, 0)) \cdot g(0, y, z - xy) \\ &= \tilde{\rho}_l(g(0, y, z - xy)^{-1}) \cdot f_u(g(x, 0, 0)) \\ &= \tilde{\rho}_l(g(0, -y, -z + xy)) \cdot u(x) \\ &= e^{2\pi i l(X(0, -y, -z + xy))} \cdot u(x) \\ &= e^{-2\pi i \tau(z - xy)} \cdot u(x). \end{aligned}$$

Using the isomorphism $L^2(\mathbb{R}, \mathbb{C}) \rightarrow L^2(G/M, G \times_{\tilde{\rho}_l} \mathbb{C})$, $u \rightarrow f_u$, we can replace $L^2(G/M, G \times_{\tilde{\rho}_l} \mathbb{C})$ by $L^2(\mathbb{R}, \mathbb{C})$ as the representation space for ρ_l .

How does a given element $g(x, y, z) \in G$ act on $u \in L^2(\mathbb{R}, \mathbb{C})$ via ρ_l ? On the corresponding f_u it acts by

$$\begin{aligned} (\rho_l(g(x, y, z))f_u)(g(\tilde{x}, \tilde{y}, \tilde{z})) &= f_u(g(x, y, z)^{-1} \cdot g(\tilde{x}, \tilde{y}, \tilde{z})) \\ &= f_u(g(-x + \tilde{x}, -y + \tilde{y}, -z + \tilde{z} + xy - x\tilde{y})) \\ &= e^{-2\pi i \tau(-z + \tilde{z} - \tilde{x}(\tilde{y} - y))} \cdot u(\tilde{x} - x). \end{aligned}$$

Putting $\tilde{y} = \tilde{z} = 0$ yields the corresponding function in $L^2(\mathbb{R}, \mathbb{C})$

$$(\rho_l(g(x, y, z))u)(\tilde{x}) = e^{-2\pi i \tau(-z + \tilde{x}y)} \cdot u(\tilde{x} - x).$$

Denote the representation ρ_l considered as acting on $L^2(\mathbb{R}, \mathbb{C})$ by ρ_τ .

Summarizing we see that the 3-dimensional Heisenberg group has two families of irreducible unitary representations, the first one parametrized by $\alpha, \beta \in \mathbb{R}$ and acting on \mathbb{C} :

$$\rho_{\alpha, \beta}(g(x, y, z)) = e^{2\pi i(\alpha x + \beta y)}.$$

The second family is parametrized by $\tau \in \mathbb{R} - \{0\}$ and acts on $L^2(\mathbb{R}, \mathbb{C})$:

$$(\rho_\tau(g(x, y, z))u)(t) = e^{2\pi i\tau(z - ty)}u(t - x). \quad (2.2)$$

Note that ρ_τ is determined by the action of the center of G

$$\rho_\tau(g(0, 0, z)) = e^{2\pi i\tau z} \cdot \text{Id}.$$

We compute the action of the Heisenberg algebra \mathfrak{g} on $L^2(\mathbb{R}, \mathbb{C})$:

$$\begin{aligned} (\rho_\tau)_*(X)u(t) &= \frac{d}{ds}\rho_\tau(\exp sX)u(t)|_{s=0} \\ &= \frac{d}{ds}\rho_\tau(g(s, 0, 0))u(t)|_{s=0} \\ &= \frac{d}{ds}e^{2\pi i\tau \cdot 0}u(t - s)|_{s=0} \\ &= -u'(t). \end{aligned}$$

Similarly,

$$\begin{aligned} (\rho_\tau)_*(Y)u(t) &= -2\pi i\tau t u(t) \\ (\rho_\tau)_*(Z)u(t) &= 2\pi i\tau u(t). \end{aligned}$$

The Hilbert space $L^2(\mathbb{R}, \mathbb{C})$ has the well-known basis of *Hermite functions*

$$h_k(t) = e^{\frac{t^2}{2}} \left(\frac{d}{dt} \right)^k e^{-t^2}, \quad k \in \mathbb{N}_0.$$

They satisfy the relations

$$\begin{aligned} h'_k(t) &= t h_k(t) + h_{k+1}(t), \\ h_{k+2}(t) + 2t h_{k+1}(t) + 2(k+1)h_k(t) &= 0. \end{aligned}$$

We look at the L^2 -basis $u_k(t) = h_k(\sqrt{2\pi|\tau|}t)$. Then the relations translate into

$$2\pi|\tau|t u_k(t) + \sqrt{2\pi|\tau|} u_{k+1}(t) = u'_k(t), \quad (2.3)$$

$$u_{k+2}(t) + 2\sqrt{2\pi|\tau|}t u_{k+1}(t) + 2(k+1)u_k(t) = 0. \quad (2.4)$$

This basis will be used in the next section for the computation of the Dirac spectrum of 3-dimensional Heisenberg manifolds.

3 Heisenberg Manifolds

In this section let G denote the 3-dimensional Heisenberg group. Let r be a positive integer. Define

$$\Gamma_r := \{g(r \cdot x, y, z) \in G \mid x, y, z \in \mathbb{Z}\}.$$

Then Γ_r is a uniformly discrete cocompact lattice in G . For every lattice in G there is an automorphism of G mapping the lattice to some Γ_r , [19, §2]. We want to compute the Dirac spectrum of the Heisenberg manifold $M = \Gamma_r \backslash G$.

For positive reals $d, T > 0$ we equip G and M with the left invariant metric for which

$$\begin{aligned} e_1 &:= -d \cdot X, \\ e_2 &:= -d \cdot Y, \\ e_3 &:= T^{-1} \cdot Z \end{aligned}$$

form an orthonormal frame. Here X, Y, Z denote the standard generators of \mathfrak{g} (compare last section). Using the Koszul formula we compute the Christoffel symbols

$$\Gamma_{12}^3 = \Gamma_{23}^1 = \Gamma_{32}^1 = -\Gamma_{21}^3 = -\Gamma_{31}^2 = -\Gamma_{13}^2 = \frac{d^2 T}{2}, \Gamma_{ij}^k = 0 \quad \text{otherwise.}$$

To emphasize the dependence of M on the parameters r, d , and T we also write $M(r, d, T)$ instead of M . Note that d and T are differential geometric parameters whereas r is topological. The fundamental groups $\pi_1(M(r, d, T)) = \Gamma_r$ are nonisomorphic for different values of r . We regard spinor fields on the universal covering as maps $G \rightarrow \Sigma_3$. Choose a basis of $\Sigma_3 \cong \mathbb{C}^2$ such that Clifford multiplication of the standard basis $E_1, E_2, E_3 \in \mathbb{R}^3$ is given by

$$\gamma(E_1) = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}, \gamma(E_2) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \gamma(E_3) = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix},$$

compare first section. The spinor connection on G with respect to the left invariant metric described above is given by (1.1)

$$\begin{aligned} \nabla_{e_1}^\Sigma \sigma &= \partial_{e_1} \sigma + \frac{1}{4} \cdot \left\{ \frac{d^2 T}{2} \gamma(E_2) \gamma(E_3) - \frac{d^2 T}{2} \gamma(E_3) \gamma(E_2) \right\} \cdot \sigma \\ &= \partial_{e_1} \sigma + \frac{d^2 T}{4} \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \cdot \sigma, \\ \nabla_{e_2}^\Sigma \sigma &= \partial_{e_2} \sigma + \frac{d^2 T}{4} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \cdot \sigma, \\ \nabla_{e_3}^\Sigma \sigma &= \partial_{e_3} \sigma + \frac{d^2 T}{4} \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} \cdot \sigma. \end{aligned}$$

For the Dirac operator we obtain

$$D = \sum_{i=1}^3 \gamma(e_i) \nabla_{e_i}^\Sigma = \sum_{i=1}^3 \gamma(e_i) \partial_{e_i} - \frac{d^2 T}{4}.$$

Spin structures of $M(r, d, T)$ are given by homomorphisms $\epsilon : \Gamma_r \rightarrow \mathbb{Z}/2\mathbb{Z} = \{-1, 1\}$. Since $\mathbb{Z}/2\mathbb{Z}$ is abelian ϵ must factor through $\Gamma_r/[\Gamma_r, \Gamma_r]$. The commutator subgroup is $[\Gamma_r, \Gamma_r] = \{g(0, 0, rz) \mid z \in \mathbb{Z}\}$. Hence $\Gamma_r/[\Gamma_r, \Gamma_r] \cong r\mathbb{Z} \oplus \mathbb{Z} \oplus (\mathbb{Z}/r\mathbb{Z})$. The homomorphism ϵ is determined by its image on generators of $\Gamma_r/[\Gamma_r, \Gamma_r]$, i.e. by

$$\begin{aligned} \delta_1 &:= \epsilon(g(r, 0, 0)), \\ \delta_2 &:= \epsilon(g(0, 1, 0)), \\ \delta_3 &:= \epsilon(g(0, 0, 1)). \end{aligned}$$

Put $\delta := (\delta_1, \delta_2, \delta_3)$. The first two components δ_1 and δ_2 can take the values 1 and -1 . The third component δ_3 can also take both values if r is even, but if r is odd δ_3 is necessarily $\delta_3 = 1$. Summarizing, we see that spin structures on $M(r, d, T)$ are given by triples $\delta = (\delta_1, \delta_2, \delta_3) \in (\mathbb{Z}/2\mathbb{Z})^3$ where δ_3 must be $+1$ in case r is odd. The corresponding homomorphism is then

$$\epsilon(g(rx, y, z)) = \delta_1^x \cdot \delta_2^y \cdot \delta_3^z \in \mathbb{Z}/2\mathbb{Z}$$

and spinor fields are maps $\sigma : G \rightarrow \Sigma_3 = \mathbb{C}^2$ such that

$$\sigma(g_0 g) = \epsilon(g_0) \cdot \sigma(g), \quad g_0 \in \Gamma_r, \quad g \in G. \quad (3.1)$$

The Heisenberg group G acts on the Hilbert space $L^2(\Sigma_\epsilon M)$ of L^2 -spinor fields on M by the right regular representation R :

$$(R(g_0)\sigma)(g) = \sigma(g g_0).$$

At the same time the Clifford algebra $\text{Cl}(3)$ acts on this Hilbert space by pointwise Clifford multiplication γ . Note that this action is compatible with condition (3.1). Moreover, the actions of G and of $\text{Cl}(3)$ commute. Modules with commuting G - and $\text{Cl}(3)$ -operations will be called G - $\text{Cl}(3)$ -bimodules.

Differentiation of a spinor field is given by the derived right regular action:

$$\partial_X \sigma = \left. \frac{d}{dt} \right|_{t=0} R(\exp tX) \sigma = R_*(X) \sigma, \quad X \in \mathfrak{g}.$$

We will decompose the Hilbert space $L^2(\Sigma_\epsilon M)$ into closed subspaces invariant under the actions R and γ . These subspaces are then also left invariant by the Dirac operator. Expressed in terms of R and γ the Dirac operator is given by

$$D = \sum_{i=1}^3 R_*(e_i) \otimes \gamma(E_i) - \frac{d^2 T}{4}. \quad (3.2)$$

Let the spin structure corresponding to $\delta = (\delta_1, \delta_2, \delta_3) \in (\mathbb{Z}/2\mathbb{Z})^3$ be fixed. Let us start with the case $\delta_3 = 1$. Then for fixed $g \in G$ and spinor field $\sigma : G \rightarrow \Sigma_3$ the map

$$\varphi_g : \mathbb{R} \rightarrow \Sigma_3, \varphi_g(z) = \left(R(g(0, 0, z))\sigma \right)(g),$$

is 1-periodic, $\varphi_g(z + 1) = \varphi_g(z)$. We expand φ_g into a Fourier series

$$\varphi_g(z) = \sum_{\tau \in \mathbb{Z}} \varphi_\tau(g) \cdot e^{2\pi i \tau z},$$

where the Fourier coefficients are given by $\varphi_\tau(g) = \int_0^1 R(g(0, 0, t))\sigma(g)e^{-2\pi i \tau t} dt$.

Putting $z = 0$ we obtain

$$\sigma(g) = \varphi_g(0) = \sum_{\tau \in \mathbb{Z}} \varphi_\tau(g).$$

We found a first decomposition of $L^2(\Sigma_\varepsilon M) = \bigoplus_{\tau \in \mathbb{Z}} H_\tau = H_0 \oplus \bigoplus_{\tau \in \mathbb{Z} - \{0\}} H_\tau$. On H_τ central elements $g(0, 0, z)$ act via R by multiplication by $e^{2\pi i \tau z}$. Hence H_τ decomposes into copies of the irreducible representations ρ_τ . The subspace H_0 is the part on which the center of G acts trivially. Thus it decomposes under R into 1-dimensional representations of the form $\rho_{\alpha, \beta}$. We now determine this decomposition.

For $\sigma \in H_0$ the map $\mathbb{R}^2 \mapsto \Sigma_3, (x, y) \rightarrow \sigma(g(x, y, 0))$, is periodic for the lattice $2(r\mathbb{Z} \oplus \mathbb{Z})$. Hence we can expand it into the Fourier series

$$\sigma(g(x, y, z)) = \sum_{\substack{\alpha \in \frac{1}{2r}\mathbb{Z} \\ \beta \in \frac{1}{2}\mathbb{Z}}} a_{\alpha, \beta} \cdot e^{2\pi i(\alpha x + \beta y)}.$$

The equivariance property (3.1) imposes further restrictions on α and β , namely

$$\sigma(g(x, y, z)) = \sum_{\substack{\alpha \in \frac{1}{2r}\mathbb{Z}, e^{2\pi i r \alpha} = \delta_1 \\ \beta \in \frac{1}{2}\mathbb{Z}, e^{2\pi i \beta} = \delta_2}} a_{\alpha, \beta} \cdot e^{2\pi i(\alpha x + \beta y)}.$$

This yields the decomposition

$$H_0 = \bigoplus_{\substack{\alpha \in \frac{1}{2r}\mathbb{Z}, e^{2\pi i r \alpha} = \delta_1 \\ \beta \in \frac{1}{2}\mathbb{Z}, e^{2\pi i \beta} = \delta_2}} H^{\alpha, \beta}$$

where each summand $H^{\alpha, \beta}$ is isomorphic to $\mathbb{C} \otimes \Sigma_3 = \Sigma_3$ with G acting via R by $\rho_{\alpha, \beta}$ on the first factor (see 2.1 for the definition) and $\text{Cl}(3)$ acting via γ on the second.

How does the Dirac operator D (3.2) act on $H^{\alpha,\beta}$? We compute

$$\begin{aligned}
D &= -dR_*(X) \cdot \gamma(E_1) - dR_*(Y) \cdot \gamma(E_2) \\
&\quad + T^{-1}R_*(Z) \cdot \gamma(E_3) - \frac{d^2T}{4} \\
&= -d \cdot 2\pi i\alpha \cdot \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} - d \cdot 2\pi i\beta \cdot \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \\
&\quad + T^{-1} \cdot 0 \cdot \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} - \frac{d^2T}{4} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
&= \begin{pmatrix} -\frac{d^2T}{4} & 2\pi id(\beta - i\alpha) \\ -2\pi id(\beta + i\alpha) & -\frac{d^2T}{4} \end{pmatrix}.
\end{aligned}$$

The eigenvalues are $-\frac{d^2T}{4} \pm 2\pi d\sqrt{\alpha^2 + \beta^2}$.

Next we decompose H_τ , $\tau \neq 0$, under the actions of G via R and of $\text{Cl}(3)$ via γ . Let $\sigma \in H_\tau$. Then σ is determined by $f_\sigma : \mathbb{R}^2 \rightarrow \mathbb{C}$, for $f_\sigma(x, y) = \sigma(g(x, y, 0))$, because

$$\sigma(g(x, y, z)) = e^{2\pi i\tau z} \cdot f_\sigma(x, y).$$

Fixing $x \in \mathbb{R}$ we expand $y \rightarrow f_\sigma(x, y)$ into a Fourier series and using (3.1) we obtain

$$f_\sigma(x, y) = \sum_{\beta \in \frac{1}{2}\mathbb{Z}, e^{2\pi i\beta} = \delta_2} b_\beta^\sigma(x) \cdot e^{2\pi i\beta y}.$$

Using (3.1) once more we get the periodicity property of b_β^σ ,

$$b_\beta^\sigma(x + \xi) = \delta_1^{\xi/r} \cdot b_{\beta + \tau\xi}^\sigma(x), \quad x \in \mathbb{R}, \xi \in r \cdot \mathbb{Z}.$$

Hence we have $|\tau| \cdot r$ independent functions $b_{1+\varepsilon}^\sigma, \dots, b_{|\tau|r+\varepsilon}^\sigma$, with $\varepsilon = 1/2$ for $\delta_2 = -1$ and $\varepsilon = 0$ for $\delta_2 = 1$. The map

$$\begin{aligned}
H_\tau &\rightarrow \bigoplus_{j=1}^{|\tau|r} L^2(\mathbb{R}, \Sigma_3) = \bigoplus_{j=1}^{|\tau|r} L^2(\mathbb{R}, \mathbb{C}) \otimes \Sigma_3 \\
\sigma &\mapsto (b_{1+\varepsilon}^\sigma, \dots, b_{|\tau|r+\varepsilon}^\sigma)
\end{aligned}$$

is an isomorphism. This decomposition is left invariant by the actions R and γ . Each summand carries a $G\text{-Cl}(3)$ -bimodule structure given by the action of R on the first factor and γ on the second. This $G\text{-Cl}(3)$ -bimodule structure is equivalent to the $G\text{-Cl}(3)$ -bimodule structure given by ρ_τ and γ (for the definition of ρ_τ see 2.2).

We compute the eigenvalues of the Dirac operator on $L^2(\mathbb{R}, \mathbb{C}) \otimes \Sigma_3$.

$$\begin{aligned}
D &= -d(\rho_\tau)_*(X) \cdot \gamma(E_1) - d(\rho_\tau)_*(Y) \cdot \gamma(E_2) \\
&\quad + T^{-1}(\rho_\tau)_*(Z) \cdot \gamma(E_3) - \frac{d^2T}{4} \\
&= -d \cdot \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \cdot \left(-\frac{d}{dt}\right) - d \cdot \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} (-2\pi i \tau t) \\
&\quad + T^{-1} \cdot \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \cdot 2\pi i \tau - \frac{d^2T}{4} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
&= \begin{pmatrix} -2\pi T^{-1}\tau - \frac{d^2T}{4} & -2\pi i d\tau t + id\frac{d}{dt} \\ 2\pi i d\tau t + id\frac{d}{dt} & 2\pi T^{-1}\tau - \frac{d^2T}{4} \end{pmatrix}.
\end{aligned}$$

We use the L^2 -basis u_0, u_1, \dots for $L^2(\mathbb{R}, \mathbb{C})$ where $u_k(t) = h_k(\sqrt{2\pi|\tau|}t)$, see Section 2. In case $\tau > 0$ we apply D to $\begin{pmatrix} u_k \\ 0 \end{pmatrix}$ and obtain

$$\begin{aligned}
D \begin{pmatrix} u_k \\ 0 \end{pmatrix} &= \begin{pmatrix} \left(-2\pi T^{-1}\tau - \frac{d^2T}{4}\right) u_k \\ id(2\pi\tau t u_k + u'_k) \end{pmatrix} \\
(2.3) \quad &= \begin{pmatrix} \left(-2\pi T^{-1}\tau - \frac{d^2T}{4}\right) u_k \\ id(4\pi\tau t u_k + \sqrt{2\pi\tau} u_{k+1}) \end{pmatrix}
\end{aligned}$$

$$\begin{aligned}
(2.4) \quad &= \begin{pmatrix} \left(-2\pi T^{-1}\tau - \frac{d^2T}{4}\right) u_k \\ -id\sqrt{2\pi\tau} \cdot 2k u_{k-1} \end{pmatrix}.
\end{aligned}$$

Similarly, $D \begin{pmatrix} 0 \\ u_{k-1} \end{pmatrix} = \begin{pmatrix} id\sqrt{2\pi\tau} u_k \\ \left(2\pi T^{-1}\tau - \frac{d^2T}{4}\right) u_{k-1} \end{pmatrix}$. Hence $\begin{pmatrix} u_0 \\ 0 \end{pmatrix}$ is an eigenvector for D with eigenvalue $-2\pi T^{-1}\tau - \frac{d^2T}{4}$ and for $k \in \mathbb{N}$ the vectors $\begin{pmatrix} u_k \\ 0 \end{pmatrix}$ and $\begin{pmatrix} 0 \\ u_{k-1} \end{pmatrix}$ span a D -invariant subspace. With respect to this basis D is given by the 2×2 -matrix

$$D \sim \begin{pmatrix} -2\pi T^{-1}\tau - \frac{d^2T}{4} & id\sqrt{2\pi\tau} \\ -2k id\sqrt{2\pi\tau} & 2\pi T^{-1}\tau - \frac{d^2T}{4} \end{pmatrix}.$$

The eigenvalues are $-\frac{d^2T}{4} \pm 2\sqrt{\pi^2\tau^2 T^{-2} + k\pi d^2\tau}$. Analogous reasoning for the case $\tau < 0$ yields an eigenvalue $-2\pi T^{-1}|\tau| - \frac{d^2T}{4}$ and for $k \in \mathbb{N}$ the eigenvalues $-\frac{d^2T}{4} \pm 2\sqrt{\pi^2\tau^2 T^{-2} + k\pi d^2|\tau|}$.

It remains to consider the case $\delta_3 = -1$. In this case we get the decomposition

$$L^2(\Sigma_\varepsilon M) = \bigoplus_{\tau \in \mathbb{Z} + \frac{1}{2}} H_\tau,$$

the representations $\rho_{\alpha,\beta}$ do not occur. The computation of the eigenvalues of D on H_τ is the same as above. We summarize

THEOREM 3.1 *Spin structures on the 3-dimensional Heisenberg manifold $M(r, d, T)$, $d, T > 0$, $r \in \mathbb{N}$, correspond to triples $\delta = (\delta_1, \delta_2, \delta_3) \in (\mathbb{Z}/2\mathbb{Z})^3$ where the case $\delta_3 = -1$ occurs if and only if r is even. The eigenvalues of the Dirac operator are as follows*

A. *If $\delta_3 = +1$:*

- a) $\lambda_{\alpha,\beta}^\pm := -\frac{d^2T}{4} \pm 2\pi d\sqrt{\alpha^2 + \beta^2}$
with multiplicity 1 for each $\alpha \in \frac{1}{2r} \cdot \mathbb{Z}$ with $e^{2\pi i r \alpha} = \delta_1$ and each $\beta \in \frac{1}{2} \cdot \mathbb{Z}$ with $e^{2\pi i \beta} = \delta_2$.
- b) $\lambda_\tau := -\frac{d^2T}{4} - 2\pi \tau T^{-1}$
with multiplicity $2\tau r$ for each $\tau \in \mathbb{N}$.
- c) $\lambda_{\tau,k}^\pm := -\frac{d^2T}{4} \pm 2\sqrt{\pi^2 \tau^2 T^{-2} + k \pi d^2 \tau}$
with multiplicity $2\tau r$ for all $\tau, k \in \mathbb{N}$.

B. *If $\delta_3 = -1$:*

- a) $\lambda_\tau := -\frac{d^2T}{4} - 2\pi \tau T^{-1}$
with multiplicity $2\tau r$ for each $\tau \in \mathbb{N}_0 + \frac{1}{2}$.
- b) $\lambda_{\tau,k}^\pm := -\frac{d^2T}{4} \pm 2\sqrt{\pi^2 \tau^2 T^{-2} + k \pi d^2 \tau}$
with multiplicity $2\tau r$ for each $\tau \in \mathbb{N}_0 + \frac{1}{2}$ and each $k \in \mathbb{N}$.

Note that the Dirac spectrum for the spin structure $\delta = (\delta_1, \delta_2, -1)$ is independent of δ_1 and of δ_2 .

The eigenvalue $\lambda_{\tau,k}^+$ is zero for suitable $T > 0$. Its multiplicity is $2\tau r$ hence unbounded if we let $\tau \rightarrow \infty$. Thus we have

COROLLARY 3.2 *For each $N \in \mathbb{N}$ the 3-dimensional Heisenberg manifold $M = \Gamma_r \backslash G$ with spin structure corresponding to δ possesses a left invariant metric such that the dimension of the space of harmonic spinors is at least N .*

A similar observation for Berger metrics on S^3 is due to Hitchin [20]. It is known that every closed 3-manifold with a fixed spin structure admits a Riemannian metric such that the corresponding Dirac operator has nontrivial kernel [4]. It is conjectured that it can be made arbitrarily large.

The manifold $M(r, d, T)$ is a circle bundle over the torus $T^2 = \mathbb{R}^2/r\mathbb{Z} \oplus \mathbb{Z}$. If one shrinks the fiber length ($T \rightarrow 0$) then $M(r, d, T)$ is said to *collapse* to T^2 .

It is interesting to examine the behavior of the eigenvalues under this collapse. We see that eigenvalues of the type λ_τ or $\lambda_{\tau,k}^\pm$ tend to $\pm\infty$. But the eigenvalues $\lambda_{\alpha,\beta}^\pm$ converge to $\pm 2\pi d\sqrt{\alpha^2 + \beta^2}$ which are precisely the Dirac eigenvalues of the torus T^2 .

COROLLARY 3.3 *If the Heisenberg manifold $M(r, d, T)$ collapses to the 2-torus, i.e. if $T \rightarrow 0$, then the Dirac eigenvalues behave as follows:*

- A.** *If $\delta_3 = +1$:
There are eigenvalues converging to those of the 2-torus. All other eigenvalues tend to $\pm\infty$.*
- B.** *If $\delta_3 = -1$:
All eigenvalues tend to $\pm\infty$.*

This phenomenon will be studied in larger generality in the next section.

To conclude this section we mention that a computation of the Dirac eigenvalues on Heisenberg manifolds can also be carried out in higher dimensions. It just becomes notationally more complicated, see [1] for more details. In all dimensions ≥ 5 it is possible to find Heisenberg manifolds with nonisomorphic fundamental groups having the same Dirac spectrum. See Section 5 for a general discussion.

4 S^1 -Actions and Collapse

In this section we will study the asymptotic behavior of the spectrum of the Dirac operator on a larger class of collapsing manifolds. This will generalize the observations we made in the last section for 3-dimensional Heisenberg manifolds. Our discussion in this section is similar in spirit to the one in [5] for the Laplace operator on functions.

We suppose that S^1 acts freely and isometrically on the compact connected Riemannian spin manifold (M, \tilde{g}) of dimension $n+1$. We view M^{n+1} as the total space of an S^1 -principal bundle over the base space $N^n = M^{n+1}/S^1$.

This base space shall carry the unique metric g such that the projection

$$\pi : (M, \tilde{g}) \longrightarrow (N, g)$$

is a Riemannian submersion.

This S^1 -principal bundle has a unique connection 1-form $i\omega : TM \rightarrow i\mathbb{R}$ such that $\ker \omega|_m$ is orthogonal to the fibers for all $m \in M$ with respect to \tilde{g} .

The S^1 -action induces a Killing vector field K . To keep the discussion simple we will assume that the length $|K|$ is constant on M . This is equivalent to saying that the fibers of π are totally geodesic. One can relax this assumption, compare Remark 4.2.

By rescaling the metric \tilde{g} on M along the fibers while keeping it the same on $\ker\omega$ we obtain a 1-parameter family of metrics \tilde{g}_ℓ on M for which $\pi : M \rightarrow N$ is a Riemannian submersion, $\omega(K) \equiv 1$, and $2\pi\ell$ is the length of the fibers. The length of K with respect to \tilde{g}_ℓ is ℓ .

As the metric on M is completely characterized by the connection 1-form $i\omega$, the fiber length $2\pi\ell$ and the metric g on N , we can express the Dirac operator in terms of ω , ℓ , and g . This will allow us to analyze the behavior of the spectrum for collapsing S^1 -fibers ($\ell \rightarrow 0$). It will turn out that the spin structure on M determines whether there are converging eigenvalues or not.

The S^1 -action on M induces an S^1 -action on $P_{\text{SO}}(M)$. A spin structure $\tilde{\varphi} : P_{\text{Spin}}(M) \rightarrow P_{\text{SO}}(M)$ will be called *projectable* if this S^1 -action on $P_{\text{SO}}(M)$ lifts to $P_{\text{Spin}}(M)$. Otherwise it will be called *nonprojectable*.

Any projectable spin structure induces a spin structure on N as follows: Let $\tilde{\varphi} : P_{\text{Spin}}(M) \rightarrow P_{\text{SO}}(M)$ denote the projectable spin structure on M and let $P_{\text{SO}(n)}(M)$ consist of those frames over M having K/ℓ as the first vector. We can identify $P_{\text{SO}}(N)$ with $P_{\text{SO}(n)}(M)/S^1$. Now $\tilde{\varphi}^{-1}(P_{\text{SO}(n)}(M))/S^1$ is a $\text{Spin}(n)$ -bundle over N and $\tilde{\varphi}$ induces a corresponding spin structure on N .

On the other hand, any spin structure on N canonically induces a projectable spin structure on M via pull-back: Let $\varphi : P_{\text{Spin}}(N) \rightarrow P_{\text{SO}}(N)$ be a spin structure of N . Then $\pi^*\varphi : \pi^*P_{\text{Spin}}(N) \rightarrow \pi^*P_{\text{SO}}(N) =: P_{\text{SO}(n)}(M)$ is a $\Theta_n : \text{Spin}(n) \rightarrow \text{SO}(n)$ equivariant map. Enlarging the structure group to $\text{Spin}(n+1)$ by

$$\tilde{\varphi} := \pi^*\varphi \times_{\Theta_n} \Theta_{n+1} : \pi^*P_{\text{Spin}}(N) \times_{\text{Spin}(n)} \text{Spin}(n+1) \rightarrow P_{\text{SO}(n)}(M) \times_{\text{SO}(n)} \text{SO}(n+1)$$

yields a spin structure on M .

Projectable spin structures and projectable spinors have already been studied by Moroianu [25].

Throughout this section we use the following convention: When we say that an operator has the eigenvalues μ_1, μ_2, \dots we suppose that each eigenvalue be repeated according to its multiplicity.

THEOREM 4.1 *Let M be a closed Riemannian spin manifold. Let S^1 act isometrically on M . We assume that the orbits have constant length. Let $N = M/S^1$ carry the induced Riemannian metric and let \tilde{g}_ℓ be the metric on M described above. Let $E \rightarrow N$ be a Hermitian vector bundle with a metric connection ∇^E .*

We suppose that the spin structure on M is projectable and that N carries the induced spin structure. Let μ_1, μ_2, \dots be the eigenvalues of the twisted Dirac operator D^E on $\Sigma N \otimes E \rightarrow N$.

Then we can number the eigenvalues $(\lambda_{j,k}(\ell))_{j \in \mathbb{N}, k \in \mathbb{Z}}$ of the twisted Dirac operator \tilde{D}^ℓ for \tilde{g}_ℓ on $\Sigma M \otimes \pi^* E \rightarrow M$ such that they depend continuously on ℓ and such that for $\ell \rightarrow 0$:

(1) For any $j \in \mathbb{N}$ and $k \in \mathbb{Z}$

$$\ell \cdot \lambda_{j,k}(\ell) \rightarrow k.$$

In particular, $\lambda_{j,k}(\ell) \rightarrow \pm\infty$ if $k \neq 0$.

(2) If $n = \dim N$ is even, then

$$\lambda_{j,0}(\ell) \rightarrow \mu_j.$$

On the other hand if $n = \dim N$ is odd, then

$$\begin{aligned} \lambda_{2j-1,0}(\ell) &\rightarrow \mu_j \\ \lambda_{2j,0}(\ell) &\rightarrow -\mu_j \end{aligned}$$

In both cases, the convergence of the eigenvalues $\lambda_{j,0}(\ell)$ is uniform in j .

REMARK 4.2 The theorem also holds for a family of nonconstant length functions $\ell : N \rightarrow \mathbb{R}^+$ provided ℓ and $(1/\ell) \text{grad } \ell$ converge uniformly to 0. The proof is analogous to the one below except that some formulas contain an additional term depending on $\text{grad } \ell$.

In the proof of the theorem we will assume for simplicity that E is a trivial line bundle. For nontrivial E every spinor bundle on N or M simply has to be twisted by E or $\pi^* E$ resp.

In order to prove Theorem 4.1 we write the Dirac operator \tilde{D}^ℓ as a sum of a vertical Dirac operator, a horizontal Dirac operator, and a zero order term.

For the definition of the vertical Dirac operator we need the Lie derivative of a spinor along the S^1 -fibers. The action of S^1 on $P_{\text{Spin}} M$ induces an action of S^1 on $\Sigma M = P_{\text{Spin}}(M^{n+1}) \times_{\text{Spin}(n+1)} \Sigma_{n+1}$ which we denote by κ . A spinor with base point m will be mapped by $\kappa(e^{it})$ to a spinor with base point $m \cdot e^{it}$. We define the *Lie derivative* of a smooth spinor Ψ in the direction of the Killing field K by

$$\mathcal{L}_K(\Psi)(m) = \left. \frac{d}{dt} \right|_{t=0} \kappa(e^{-it})(\Psi(m \cdot e^{it})).$$

See [7] for a general discussion of Lie derivatives for spinors. Since \mathcal{L}_K is the differential of a representation of the Lie group S^1 on $L^2(\Sigma M)$, we get the decomposition

$$L^2(\Sigma M) = \bigoplus_{k \in \mathbb{Z}} V_k$$

into the eigenspaces V_k of \mathcal{L}_K for the eigenvalue ik , $k \in \mathbb{Z}$. The S^1 -action commutes with the Dirac operator on M , hence this decomposition is preserved by the Dirac operator.

We now want to calculate the difference between the covariant derivative and the Lie derivative in the direction of K . We always use the convention that any r -form α acts on a spinor Ψ by

$$\gamma(\alpha)\Psi := \sum_{i_1 < \dots < i_r} \alpha(e_{i_1}, \dots, e_{i_r}) \gamma(e_{i_1}) \cdots \gamma(e_{i_r}) \Psi$$

where the e_i form an orthonormal basis of the tangent space.

LEMMA 4.3 *If $\tilde{\Psi}$ is a smooth section of ΣM , then*

$$\nabla_K \tilde{\Psi} - \mathcal{L}_K \tilde{\Psi} = \frac{\ell^2}{4} \gamma(d\omega) \tilde{\Psi}$$

Proof First we fix some notations. In the rest of this section, we denote the horizontal lift of $X \in TN$ by $\tilde{X} \in TM$. Let (f_1, \dots, f_n) be a local orthonormal frame of N on the open subset U and let A be a lift to the spin structure. Then $(e_0 := K/\ell, e_1 := \tilde{f}_1 \circ \pi, \dots, e_n := \tilde{f}_n \circ \pi)$ is a local orthonormal frame on $\pi^{-1}(U)$ with lift π^*A . Using the Koszul formula we can express the Christoffel symbols $\tilde{\Gamma}_{ij}^k$ for (e_0, \dots, e_n) in terms of the Christoffel symbols Γ_{ij}^k of (f_1, \dots, f_n) , the S^1 -bundle curvature form $d\omega$ and the length ℓ . For $i, j, k \in \{1, \dots, n\}$ we get

$$\tilde{\Gamma}_{ij}^k = \Gamma_{ij}^k \quad - \tilde{\Gamma}_{ij}^0 = \tilde{\Gamma}_{i0}^j = \tilde{\Gamma}_{0i}^j = \frac{\ell}{2} d\omega(e_i, e_j) \quad (4.1)$$

$$\tilde{\Gamma}_{i0}^0 = \tilde{\Gamma}_{0i}^0 = \tilde{\Gamma}_{00}^i = \tilde{\Gamma}_{00}^0 = 0$$

With respect to these local frames on M and N we apply (1.1) for M and we obtain

$$\begin{aligned} \frac{1}{\ell} \nabla_K \tilde{\Psi} &= \frac{1}{\ell} \partial_K \tilde{\Psi} + \frac{1}{4} \sum_{j,k=0}^n \tilde{\Gamma}_{0j}^k \gamma(e_j) \gamma(e_k) \tilde{\Psi} \\ &= \frac{1}{\ell} \partial_K \tilde{\Psi} + \frac{\ell}{8} \sum_{j,k=1}^n d\omega(e_j, e_k) \gamma(e_j) \gamma(e_k) \tilde{\Psi} \\ &= \frac{1}{\ell} \partial_K \tilde{\Psi} + \frac{\ell}{4} \gamma(d\omega) \tilde{\Psi}. \end{aligned}$$

On the other hand the choice of the frames implies that

$$\mathcal{L}_K \tilde{\Psi} = \partial_K \tilde{\Psi},$$

thus the lemma follows. \square

Now we associate to the S^1 -bundle $M \rightarrow N$ the complex line bundle $L := M \times_{S^1} \mathbb{C}$ with the natural connection given by $i\omega$.

LEMMA 4.4 *Let n be even. There is a homothety of Hilbert spaces*

$$Q_k : L^2(\Sigma N \otimes L^{-k}) \rightarrow V_k$$

such that the horizontal covariant derivative is given by

$$\nabla_{\tilde{X}} Q_k(\Psi) = Q_k(\nabla_X \Psi) + \frac{\ell}{4} \gamma(K/\ell) \gamma(\tilde{V}_X) Q_k(\Psi)$$

where V_X is the vector field on N satisfying $d\omega(\tilde{X}, \cdot) = \langle \tilde{V}_X, \cdot \rangle$. Clifford multiplication is preserved, i.e.

$$Q_k(\gamma(X)\Psi) = \gamma(\tilde{X})Q_k(\Psi).$$

Proof Since n is even, we have $\Sigma_n = \Sigma_{n+1}$. We define a vector bundle map

$$\begin{aligned} \Pi_k : \Sigma M &\rightarrow \Sigma N \otimes L^{-k} \\ (m, [\pi^* A, \sigma]) &\mapsto (\pi(m), [A, \sigma] \otimes [m, 1]^{-k}), \end{aligned}$$

where $[\pi^* A, \sigma]$ (resp. $[A, \sigma], [m, 1]$) denotes the equivalence class of $(\pi^* A, \sigma)$ (resp. $(A, \sigma), (m, 1)$) in $\pi^* P_{\text{Spin}}(N) \times_{\text{Spin}(n)} \Sigma_n = \Sigma M$ (resp. $P_{\text{Spin}}(N) \times_{\text{Spin}(n)} \Sigma_n = \Sigma N$, resp. $M \times_{S^1} \mathbb{C} = L$). This bundle map is fiberwise a vector space isomorphism preserving Clifford multiplication. Therefore for any spinor $\Psi : N \rightarrow \Sigma N \otimes L^{-k}$ there is a unique $Q_k(\Psi)$ such that

$$\begin{array}{ccc} M & \xrightarrow{Q_k(\Psi)} & \Sigma M \\ \pi \downarrow & & \downarrow \Pi_k \\ N & \xrightarrow{\Psi} & \Sigma N \otimes L^{-k} \end{array}$$

commutes. Hence Q_k is a well defined homomorphism of Hilbert spaces and injectivity of Q_k follows from the surjectivity of π .

For any section $\tilde{\Psi}$ of ΣM we get the equivalences

$$\begin{aligned}
& \tilde{\Psi} \in \text{image}(Q_k) \\
\iff & \Pi_k \circ \tilde{\Psi}(m) = \Pi_k \circ \tilde{\Psi}(me^{it}) \quad \forall m \in M, e^{it} \in S^1 \\
\iff & \tilde{\Psi}(m) = e^{-ikt} \kappa(e^{-it}) \tilde{\Psi}(me^{it}) \quad \forall m \in M, e^{it} \in S^1 \\
\iff & ik \tilde{\Psi} = \mathcal{L}_K \tilde{\Psi}.
\end{aligned}$$

Thus the image of Q_k is precisely V_k . The map Q_k is an isometry up to the factor ℓ .

Finally it remains to show the formula for the covariant derivative in the horizontal directions. For this it suffices to prove it locally and we can restrict to the case $X = f_i$ with the notations as in the previous lemma. As before we use (1.1) and (4.1).

$$\begin{aligned}
\nabla_{\tilde{f}_i} Q_k(\Psi) &= Q_k(\nabla_{f_i} \Psi) + \frac{1}{4} \sum_{j=1}^n \tilde{\Gamma}_{i0}^j \gamma(K/\ell) \gamma(e_j) Q_k(\Psi) \\
&\quad + \frac{1}{4} \sum_{j=1}^n \tilde{\Gamma}_{ij}^0 \gamma(e_j) \gamma(K/\ell) Q_k(\Psi) \\
&= Q_k(\nabla_{f_i} \Psi) + \frac{\ell}{4} \gamma(K/\ell) \gamma(\tilde{V}_{f_i}) Q_k(\Psi)
\end{aligned}$$

□

Proof of Theorem 4.1 for n even. We define the *horizontal Dirac operator* as the unique closed linear operator $D_h : L^2(\Sigma M) \rightarrow L^2(\Sigma M)$ on each V_k given by

$$D_h := Q_k \circ D \circ Q_k^{-1}$$

where D is the (twisted) Dirac operator on $\Sigma N \otimes L^{-k}$.

We define the *vertical Dirac operator*

$$D_v := \gamma(K/\ell) \mathcal{L}_K$$

and the zero order term

$$Z := -(1/4) \gamma(K/\ell) \gamma(d\omega).$$

Using Lemma 4.3 and Lemma 4.4 we can express the Dirac operator \tilde{D}^ℓ as a sum

$$\begin{aligned}
\tilde{D}^\ell &= \sum_{i=0}^n \gamma(e_i) \nabla_{e_i} \\
&= \frac{1}{\ell} D_v + D_h + \ell Z.
\end{aligned}$$

Since D_h , $\gamma(K/\ell)$ and Z commute with the S^1 -action they also commute with \mathcal{L}_K . On the other hand, we have $\gamma(K/\ell)Q_k(\Psi) = Q_k(c\gamma(\text{dvol}_n)\Psi)$ with $c \in \{1, i, -1, -i\}$ depending on n and the representation of Cl_{n+1} . Since $\gamma(\text{dvol}_n)$ anticommutes with any twisted Dirac operator on N , we know that $\gamma(K/\ell)$ anticommutes with D_h . Let Ψ be a common eigenspinor for \mathcal{L}_K and D_h for the eigenvalues ik and μ resp.

On $U := \text{span}\{\Psi, \gamma(K/\ell)\Psi\}$ the operator $A_\ell := (1/\ell)D_v + D_h$ is represented by the matrix

$$\frac{1}{\ell} \begin{pmatrix} 0 & -ik \\ ik & 0 \end{pmatrix} + \begin{pmatrix} \mu & 0 \\ 0 & -\mu \end{pmatrix} = \begin{pmatrix} \mu & -ik/\ell \\ ik/\ell & -\mu \end{pmatrix}.$$

Thus for $k = 0$ the restriction $A_\ell|_U$ has eigenvalues $\pm\mu$. For $k \neq 0$ the eigenvalues of $A_\ell|_U$ are the square roots of $(k/\ell)^2 + \mu^2$. Therefore the eigenvalues $(\lambda_{j,k}^0(\ell))_{j \in \mathbb{N}, k \in \mathbb{Z}}$ of A_ℓ can be numbered such that they are continuous in ℓ and satisfy properties (1) and (2) of Theorem 4.1. The additional term ℓZ does not change this behavior because $Z : L^2(\Sigma M) \rightarrow L^2(\Sigma M)$ is bounded. \square

Proof of Theorem 4.1 for n odd. In the case “ n odd” the proof is similar but we have to do some modifications as the vector space equality $\Sigma_{n+1} = \Sigma_n$ is no longer valid.

For the standard basis $E_1, \dots, E_n \in \mathbb{R}^n$ and $E_0, \dots, E_n \in \mathbb{R}^{n+1}$ we define the *volume element* by

$$\nu_n := i^{\frac{n+1}{2}} \gamma(E_1) \cdots \gamma(E_n) \in \text{Cl}(n+1).$$

Its square is the identity and hence has the eigenvalues ± 1 . Let

$$\Sigma_{n+1} = \Sigma_n^{(+)} \oplus \Sigma_n^{(-)}$$

be the splitting into the eigenspaces. These $\Sigma_n^{(\pm)}$ are the two irreducible representations of $\text{Cl}(n)$. The Clifford multiplication $\gamma(E_0)$ anticommutes with ν_n and therefore restricts to an isomorphism of $\Sigma_n^{(+)} \rightarrow \Sigma_n^{(-)}$. This isomorphism changes the sign of $\gamma(E_i)$ for $i > 0$ as $\gamma(E_0)\gamma(E_i) = -\gamma(E_i)\gamma(E_0)$.

Denote the corresponding spinor bundles by $\Sigma^{(\pm)}N$. Lemma 4.4 is now valid with

$$Q_k : L^2\left(\left(\Sigma^{(+)}N \oplus \Sigma^{(-)}N\right) \otimes L^{-k}\right) \rightarrow V_k.$$

The same arguments as in the even case can be used to prove that the eigenvalues of $\tilde{D}^\ell|_{V_k}$ converge to $\pm\infty$ for $k \neq 0$ and to the eigenvalues of the Dirac operator on $\Sigma^{(+)}N \oplus \Sigma^{(-)}N$ for $k = 0$. The Dirac operator on $\Sigma^{(+)}N$ has the same eigenvalues as the Dirac operator on $\Sigma^{(-)}N$ but with opposite signs. This yields the statement of Theorem 4.1. \square

In the case of a nonprojectable spin structure we get a similar result. Here the variable k takes no longer integer values but

$$k \in \mathbb{Z} + \frac{1}{2}.$$

THEOREM 4.5 *Let M be a closed Riemannian spin manifold. Let S^1 act isometrically on M . We assume that the orbits have constant length. Let $N = M/S^1$ carry the induced Riemannian metric and let \tilde{g}_ℓ be the metric on M described above. Let $E \rightarrow N$ be a Hermitian vector bundle with a metric connection ∇^E .*

We suppose that the spin structure on M is nonprojectable.

Then we can number the eigenvalues $(\lambda_{j,k}(\ell))_{j \in \mathbb{N}, k \in \mathbb{Z} + (1/2)}$ of the twisted Dirac operator \tilde{D}^ℓ for \tilde{g}_ℓ on $\Sigma M \otimes \pi^ E \rightarrow M$ such that they depend continuously on ℓ and such that for $\ell \rightarrow 0$*

$$\ell \cdot \lambda_{j,k}(\ell) \rightarrow k \quad \forall j = 1, 2, \dots$$

In particular, $\lambda_{j,k}(\ell) \rightarrow \pm\infty$.

Proof The proof is a variation of the proof of Theorem 4.1. We will restrict to the case n even and E trivial.

Let $\tilde{\varphi} : P_{\text{Spin}}(M) \rightarrow P_{\text{SO}}(M)$ be a nonprojectable spin structure on M . In this case N may or may not be spin. But as before the principal $\text{SO}(n+1)$ -bundle $P_{\text{SO}}(M)$ can be restricted to a principal $\text{SO}(n)$ -bundle $P_{\text{SO}(n)}(M)$. Moreover $P := \tilde{\varphi}^{-1}(P_{\text{SO}(n)}(M))$ is a principal $\text{Spin}(n)$ -bundle over M .

The action of $S^1 \cong (\mathbb{R}/2\pi\mathbb{Z})$ does not lift to P , but the double covering of S^1 , i.e. $S^1 \cong (\mathbb{R}/4\pi\mathbb{Z})$, does act on P . We define $\text{Spin}^{\mathbb{C}}(n)$ to be $\text{Spin}(n) \times_{\mathbb{Z}_2} S^1$ where $-1 \in \mathbb{Z}_2$ identifies $(-A, c)$ with $(A, -c)$. The Lie group $\text{Spin}^{\mathbb{C}}(n)$ naturally acts on Σ_n . The actions of $\text{Spin}(n)$ and S^1 on P induce a free action of $\text{Spin}^{\mathbb{C}}(n)$ on P and we can view P as a principal $\text{Spin}^{\mathbb{C}}(n)$ -bundle over N . Then we can form the bundle $P \times_{\text{Spin}^{\mathbb{C}}(n)} \Sigma_n$.

If N is spin, this bundle is just $\Sigma N \otimes L^{\frac{1}{2}}$. If N is not spin, then neither ΣN nor $L^{\frac{1}{2}}$ exist but $\Sigma N \otimes L^{\frac{1}{2}}$ does exist.

Again we get a splitting

$$L^2(\Sigma M) = \bigoplus_{k \in \mathbb{Z} + \frac{1}{2}} V_k$$

into eigenspaces V_k for \mathcal{L}_K to the eigenvalue ik . The rest of the proof of Theorem 4.5 is the same as the one for the case $k \neq 0$ in Theorem 4.1. \square

Example We apply Theorems 4.1 and 4.5 to the Hopf fibration $S^{2m+1} \rightarrow \mathbb{C}P^m$. The Dirac eigenvalues for the metric \tilde{g}_ℓ on S^{2m+1} have been computed by the second author in [4, Theorem 3.1]:

The Dirac operator of S^{2m+1} with Berger metric \tilde{g}_ℓ has the following eigenvalues:

- (i) $\lambda(\ell) = \frac{1}{\ell} \left(a + \frac{m+1}{2} \right) + \frac{\ell m}{2}$, $a \in \mathbb{N}_0$, with multiplicity $\binom{m+a}{a}$.
- (ii) $\lambda(\ell) = (-1)^{m+1} \cdot \left\{ \frac{1}{\ell} \left(a + \frac{m+1}{2} \right) + \frac{\ell m}{2} \right\}$, $a \in \mathbb{N}_0$, with multiplicity $\binom{m+a}{a}$.
- (iii) $\lambda_{1,2}(\ell) = (-1)^j \frac{\ell}{2} \pm \left\{ \left[\frac{\ell}{2} (m - 2j - 1) + \frac{1}{\ell} \left(a_1 - a_2 + \frac{m-1}{2} - j \right) \right]^2 + 4(m - j + a_1)(j + 1 + a_2) \right\}^{1/2}$

$a_1, a_2 \in \mathbb{N}_0$, $j = 0, \dots, m-1$, with multiplicity

$$\frac{(a_1 + m)! (a_2 + m)! (a_1 + a_2 + m + 1)}{a_1! a_2! (a_1 + m - j)! (a_2 + j + 1)! m! j! (m - j - 1)!}$$

We observe for the eigenvalues of type (i) and (ii)

$$\lim_{\ell \rightarrow 0} \ell \cdot \lambda(\ell) = \pm \left(a + \frac{m+1}{2} \right).$$

We see that the spin structure on S^{2m+1} is projectable if and only if m is odd. In fact, $\mathbb{C}P^m$ is spin if and only if m is odd.

To obtain a finite limit of $\lambda(\ell)$ itself we must necessarily have an eigenvalue of type (iii) and $a_1 - a_2 + \frac{m-1}{2} = j$. In particular, m must be odd again. Plugging this expression for j into the formula for the eigenvalue of type (iii) and passing to the limit $\ell \rightarrow 0$ yields the eigenvalues of the complex projective space:

THEOREM 4.6 *The Dirac operator on $\mathbb{C}P^m$, m odd, has the eigenvalues*

$$\pm 2 \sqrt{\left(a_1 + \frac{m+1}{2} \right) \left(a_2 + \frac{m+1}{2} \right)}$$

where $a_1, a_2 \in \mathbb{N}_0$, $|a_1 - a_2| \leq \frac{m-1}{2}$ with multiplicity

$$\frac{(a_1 + m)! (a_2 + m)! (a_1 + a_2 + m + 1)}{a_1! a_2! \left(a_1 + \frac{m+1}{2} \right)! \left(a_2 + \frac{m+1}{2} \right)! m! \left(a_1 - a_2 + \frac{m-1}{2} \right)! \left(a_2 - a_1 + \frac{m-1}{2} \right)!}$$

The formula for the eigenvalues coincides with the results in [9, 10, 30] after the right parameter substitutions. Our expression however is simpler than the previous ones. The formula for the multiplicities is new.

5 Isospectral Deformations

We now study the isospectrality question for the Dirac operator on Riemannian nilmanifolds $\Gamma \backslash G$. The spectrum of the Dirac operator does not only depend on the Riemannian structure but also on the spin structure, i.e. on the homomorphism $\varepsilon : \Gamma \rightarrow \{-1, 1\}$. For $\varepsilon \equiv 1$ it will turn out in the examples under consideration that the spectrum of the Dirac operator behaves much like the spectrum of the Laplacian on 1-forms. For other ε the behavior can be different.

At first, let $L_\varepsilon^2(\Gamma \backslash G)$ be the space of all complex-valued locally square-integrable functions on G , satisfying the ε -equivariance condition

$$f(g_0g) = \varepsilon(g_0)f(g),$$

for all $g_0 \in \Gamma$ and $g \in G$. The nilpotent group G acts on this space via the right regular representation

$$\begin{aligned} R : G &\rightarrow \text{End}(L_\varepsilon^2(\Gamma \backslash G)) \\ g_0 &\mapsto (R(g_0) : f \mapsto f(\cdot g_0)). \end{aligned}$$

Let Γ' be another cocompact discrete subgroup of G , $\varepsilon' : \Gamma' \rightarrow \{-1, 1\}$ a homomorphism. We say that (Γ, ε) and (Γ', ε') are *representation equivalent* if the right regular representations on $L_\varepsilon^2(\Gamma \backslash G)$ and on $L_{\varepsilon'}^2(\Gamma' \backslash G)$ are isomorphic as G -representations.

As already mentioned at the end of Section 1, sections of the spinor bundle $\Sigma_\varepsilon(\Gamma \backslash G)$ may be identified with Γ -equivariant maps $G \rightarrow \Sigma_n$. Elements of Σ_n will also be interpreted as constant maps $G \rightarrow \Sigma_n$ and therefore as sections of ΣG .

Hence $f \otimes s \mapsto f \cdot s$ defines an isomorphism of Hilbert spaces $L_\varepsilon^2(\Gamma \backslash G) \otimes_{\mathbb{C}} \Sigma_n \rightarrow L^2(\Sigma_\varepsilon(\Gamma \backslash G))$. In the rest of this section we will identify via this isomorphism. Clifford multiplication is performed on the second factor

$$\gamma(X)(f \otimes s) = f \otimes (\gamma(X)s)$$

for all $X \in \mathfrak{g}$, $f \in L_\varepsilon^2(\Gamma \backslash G)$, $s \in \Sigma_n$.

THEOREM 5.1 *If (Γ, ε) and (Γ', ε') are representation equivalent, then the corresponding Dirac operators have the same spectrum for any left-invariant Riemannian metric on G .*

Examples:

- Milnor's 16-dimensional tori [24] satisfy the condition of the theorem for $\varepsilon \equiv \varepsilon' \equiv 1$.
- The Gordon-Wilson deformations [18] preserve the spectrum of the Dirac operator. In fact, if Ψ_t is a family of almost inner automorphisms of G , Γ a cocompact discrete subgroup and $\varepsilon : \Gamma \rightarrow \{-1, 1\}$ a homomorphism, then $(\Psi_t(\Gamma), \varepsilon \circ \Psi_t^{-1})$ are pairwise representation equivalent. This deformation preserves the length spectrum, the spectrum of the Laplacian on functions and on forms, and the Dirac spectrum for all spin structures.

Since the spectrum of the Dirac operator determines the volume, we get

COROLLARY 5.2 *If (Γ, ε) and (Γ', ε') are representation equivalent, then $\Gamma \backslash G$ and $\Gamma' \backslash G$ have the same volume for any left invariant volume form on G .*

Proof of Theorem 5.1 Since $L^2_\varepsilon(\Gamma \backslash G)$ and $L^2_{\varepsilon'}(\Gamma' \backslash G)$ are isomorphic as representations of the Lie group G they are isomorphic as representations of its Lie algebra \mathfrak{g} . Let

$$\vartheta : L^2_\varepsilon(\Gamma \backslash G) \rightarrow L^2_{\varepsilon'}(\Gamma' \backslash G)$$

be such an isomorphism, i.e. $\vartheta \circ R = R' \circ \vartheta$ and $\vartheta \circ R_* = R'_* \circ \vartheta$. In order to prove the theorem we show that $\vartheta \otimes \text{id}$ maps eigenspinors of the Dirac operator in $L^2(\Sigma_\varepsilon(\Gamma \backslash G))$ to eigenspinors of the Dirac operator in $L^2(\Sigma_{\varepsilon'}(\Gamma' \backslash G))$. It is important to keep in mind that these two Dirac operators are just restrictions of the Dirac operator on ΣG to spinors equivariant under ε (resp. ε').

Let $s = \sum_i f_i s_i$ be an arbitrary eigenspinor in $L^2(\Sigma_\varepsilon(\Gamma \backslash G))$ for the eigenvalue λ with $f_i \in L^2_\varepsilon(\Gamma \backslash G)$ and $s_i \in \Sigma_n$. Let E_1, \dots, E_n an orthonormal basis of \mathfrak{g} . Using the Leibniz rule

$$\nabla_{E_j}(f_i s_i) = (R_*(E_j)f_i)s_i + f_i \nabla_{E_j}s_i \quad (5.1)$$

we get

$$\begin{aligned} \lambda(\vartheta \otimes \text{id})(s) &= (\vartheta \otimes \text{id})(Ds) \\ &= \sum_{ij} (\vartheta \otimes \text{id}) \left((\text{id} \otimes \gamma(E_j)) \nabla_{E_j}(f_i s_i) \right) \\ &= \sum_{ij} \left(\vartheta(R_*(E_j)f_i) \gamma(E_j)s_i + \vartheta(f_i) \gamma(E_j) \nabla_{E_j}s_i \right) \\ &= \sum_{ij} \left(R'_*(E_j)(\vartheta(f_i)) \gamma(E_j)s_i + \vartheta(f_i) \gamma(E_j) \nabla_{E_j}s_i \right) \\ &= \sum_{ij} (\text{id} \otimes \gamma(E_j)) \nabla_{E_j}(\vartheta(f_i)s_i) \\ &= D((\vartheta \otimes \text{id})(s)). \end{aligned}$$

Thus $\vartheta \otimes \text{id}$ is an isomorphism of Hilbert spaces that preserves eigenspinors and eigenvalues and therefore the Dirac spectra coincide. \square

To get further applications of Theorem 5.1 we need criteria telling us under which conditions (Γ, ε) and (Γ', ε') are representation equivalent. In the examples of Theorem 5.1 we already mentioned that almost inner automorphisms give us representation equivalent families. This criterion is sufficient but not necessary. However, in the special case when G is strictly nonsingular, we will get a necessary and sufficient criterion. A nilpotent Lie group G with center $Z(G)$ will be called *strictly nonsingular* if for every $z \in Z(G)$ and $x \in G - Z(G)$ there exists an $a \in G$ such that

$$xax^{-1}a^{-1} = z.$$

Note that the Heisenberg group is strictly nonsingular.

Let Γ be a cocompact discrete subgroup of G . The quotient $G/Z(G)$ will be denoted by \overline{G} and the image of Γ under the canonical projection $\pi : G \rightarrow \overline{G}$ by $\overline{\Gamma}$. According to [27, Proposition 2.17] and [12, Lemma 5.1.4] $\overline{\Gamma}$ is a cocompact discrete subgroup of \overline{G} . If $\varepsilon|_{\Gamma \cap Z(G)} \equiv 1$ then there is an $\overline{\varepsilon}$ such that

$$\begin{array}{ccc} \Gamma & \xrightarrow{\varepsilon} & \{-1, 1\} \\ \downarrow & \nearrow \overline{\varepsilon} & \\ \overline{\Gamma} & & \end{array}$$

commutes. Such an ε will be called *projectable*. For $\dim Z(G) = 1$ this definition agrees with our definition of projectable spin structures in Section 4.

THEOREM 5.3 *Let G be a simply connected, strictly nonsingular nilpotent Lie group with left invariant metric. Let Γ_1, Γ_2 be cocompact, discrete subgroups of G such that*

$$\Gamma_1 \cap Z(G) = \Gamma_2 \cap Z(G) \quad (=:\Gamma_Z)$$

and let $\varepsilon_i : \Gamma_i \rightarrow \{-1, 1\}$ be homomorphisms. We define the volume quotient

$$v := \frac{\text{vol}(\overline{\Gamma}_1 \backslash \overline{G})}{\text{vol}(\overline{\Gamma}_2 \backslash \overline{G})} = \frac{\text{vol}(\Gamma_1 \backslash G)}{\text{vol}(\Gamma_2 \backslash G)} \in \mathbb{Q}^+.$$

Then $(\Gamma_1, \varepsilon_1)$ and $(\Gamma_2, \varepsilon_2)$ are representation equivalent if and only if one of the following conditions is satisfied:

(1) $\varepsilon_1|_{\Gamma_Z} \equiv \varepsilon_2|_{\Gamma_Z} \equiv 1$ and $(\bar{\Gamma}_1, \bar{\varepsilon}_1)$ and $(\bar{\Gamma}_2, \bar{\varepsilon}_2)$ are representation equivalent for \bar{G} .

(2) $\varepsilon_1|_{\Gamma_Z} \equiv \varepsilon_2|_{\Gamma_Z} \not\equiv 1$ and $v = 1$.

REMARK 5.4 If under the assumptions of the above theorem $\varepsilon_1|_{\Gamma_Z} \equiv \varepsilon_2|_{\Gamma_Z} \not\equiv 1$ but $v \neq 1$, then $(\Gamma_1, \varepsilon_1)$ and $(\Gamma_2, \varepsilon_2)$ are “representation equivalent up to multiplicity v ”, i.e. $L_{\varepsilon_1}^2(\Gamma_1 \backslash G)$ and $L_{\varepsilon_2}^2(\Gamma_2 \backslash G)$ have the same irreducible components and for every irreducible component H with multiplicities m_1, m_2 we get $v = m_1/m_2$.

Let us compare this result to Theorem 3.1. Using the notation of Section 3 for the 3-dimensional Heisenberg groups write $M(r_i, d, T) = \Gamma_i \backslash G$. Then $v = r_1/r_2$. The remark explains why the eigenvalues of the Dirac operator in Theorem 3.1 B do not depend on r and why the multiplicity is proportional to r .

Proof of Theorem 5.3 and Remark 5.4 If $(\Gamma_1 \backslash G, \varepsilon_1)$ and $(\Gamma_2 \backslash G, \varepsilon_2)$ are representation equivalent, then $L_{\varepsilon_1}^2(\Gamma_1 \backslash G)$ and $L_{\varepsilon_2}^2(\Gamma_2 \backslash G)$ are also equivalent as Γ_Z -modules. Since the action of $g_0 \in \Gamma_Z$ on $L_{\varepsilon_i}^2(\Gamma_i \backslash G)$ is just multiplication by $\varepsilon_i(g_0)$ we get

$$\varepsilon_1|_{\Gamma_Z} \equiv \varepsilon_2|_{\Gamma_Z}.$$

Let \mathcal{T} be the set of coadjoint orbits of \mathfrak{g}^* . Via Kirillov theory \mathcal{T} parametrizes the set of irreducible unitary representations of G . We write elements of \mathcal{T} in the form $[\tau]$ with $\tau \in \mathfrak{g}^*$. The irreducible G -module corresponding to $[\tau] \in \mathcal{T}$ will be denoted by $H_{[\tau]}$. We write $L_{\varepsilon_i}^2(\Gamma_i \backslash G)$ as the discrete sum of irreducible representations

$$L_{\varepsilon_i}^2(\Gamma_i \backslash G) = \bigoplus_{[\tau] \in \mathcal{T}} m_i([\tau]) H_{[\tau]}$$

where $m_i([\tau]) \in \mathbb{N}_0$ denotes the multiplicity of $H_{[\tau]}$. Let \mathfrak{z} be the Lie algebra of the center $Z(G)$. The action of $Z(G)$ on $H_{[\tau]}$ is trivial if and only if $\tau|_{\mathfrak{z}} \equiv 0$. Therefore

$$H_i^{\text{proj}} := \bigoplus_{\substack{[\tau] \in \mathcal{T} \\ \tau|_{\mathfrak{z}} \equiv 0}} m_i([\tau]) H_{[\tau]} \quad i = 1, 2$$

is the space of L^2 -functions acted trivially upon by $Z(G)$. The orthogonal complement is

$$H_i^\perp := \bigoplus_{\substack{[\tau] \in \mathcal{T} \\ \tau|_{\mathfrak{z}} \not\equiv 0}} m_i([\tau]) H_{[\tau]} \quad i = 1, 2.$$

The G -modules $L_{\varepsilon_i}^2(\Gamma_i \backslash G)$ for $i = 1, 2$ are isomorphic if and only if H_1^{proj} is isomorphic to H_2^{proj} and H_1^\perp is isomorphic to H_2^\perp . In a first step we investigate

under which conditions H_1^{proj} and H_2^{proj} are isomorphic and then we treat H_1^\perp and H_2^\perp .

In the case $\varepsilon_i|_{\Gamma_Z} \not\equiv 1$ we have $H_i^{\text{proj}} = \{0\}$. Therefore H_1^{proj} and H_2^{proj} are trivially isomorphic G -modules.

In the case $\varepsilon_i|_{\Gamma_Z} \equiv 1$ elements in H_i^{proj} are exactly those functions in $L^2(\Gamma_i \backslash G)$ which are pullbacks of functions in $L^2_{\bar{\varepsilon}_i}(\bar{\Gamma}_i \backslash \bar{G})$. Thus we get a module-isomorphism (for both groups G and \bar{G})

$$H_i^{\text{proj}} \xrightarrow{\sim} L^2_{\bar{\varepsilon}_i}(\bar{\Gamma}_i \backslash \bar{G}).$$

Therefore in this case H_1^{proj} and H_2^{proj} are isomorphic G -modules if and only if $(\bar{\Gamma}_1, \bar{\varepsilon}_1)$ and $(\bar{\Gamma}_2, \bar{\varepsilon}_2)$ are representation equivalent for \bar{G} .

It remains to study the H_i^\perp . In Lemma 5.5 below we will show that $m_i([\tau]) = c_{G,\tau} \cdot \text{vol}(\bar{\Gamma}_i \backslash \bar{G})$ for any $[\tau] \in \mathcal{T}$ with $\tau|_{\mathfrak{z}} \not\equiv 0$ where $c_{G,\tau}$ only depends on G , on τ , and on the volume form on \bar{G} but not on the lattice Γ_i . This will complete the proof of the theorem: Suppose that condition (1) or (2) in the theorem is fulfilled. Because of Corollary 5.2 we then know that the volume quotient v is 1, and therefore H_1^\perp and H_2^\perp are isomorphic G -modules. Using the statement for $v \neq 1$ we get Remark 5.4. \square

To complete the proof of Theorem 5.3 and Remark 5.4 it remains to show Lemma 5.5. The lemma is a consequence of [26]. In order to simplify notation we drop the index i in the formulation from now on.

For any $\rho \in \mathfrak{z}^*$ we define a bilinear form on $\bar{\mathfrak{g}}$ by

$$b_\rho(\bar{X}, \bar{Y}) := \rho([X, Y]),$$

where X and Y are vectors in \mathfrak{g} and \bar{X} and \bar{Y} their images in $\bar{\mathfrak{g}}$. As \mathfrak{g} is strictly nonsingular, b_ρ is nondegenerate for all $\rho \neq 0$ [26, 16]. Therefore $\dim \bar{\mathfrak{g}}$ is even and we suppose $\dim \bar{\mathfrak{g}} = 2d$. We define the affine lattice $L_{\Gamma,\varepsilon}$ by

$$L_{\Gamma,\varepsilon} := \left\{ \rho \in \mathfrak{z}^* \mid e^{2\pi i \rho(z)} = \varepsilon(z) \quad \forall z \in \log(\Gamma_Z) \right\}$$

LEMMA 5.5 *If $\rho := \tau|_{\mathfrak{z}} \in L_{\Gamma,\varepsilon} - \{0\}$ then*

$$m([\tau]) = \left| \int_{\bar{\Gamma} \backslash \bar{G}} \underbrace{b_\rho \wedge \dots \wedge b_\rho}_{d\text{-times}} \right| = \text{vol}(\bar{\Gamma} \backslash \bar{G}) \cdot c_{G,\tau}$$

where $c_{G,\tau}$ only depends on G , on τ and on the volume form on \bar{G} .

If $\rho := \tau|_{\mathfrak{z}} \in \mathfrak{z}^ - (L_{\Gamma,\varepsilon} \cup \{0\})$ then $m([\tau]) = 0$.*

The lemma does not make any statement for $\rho = 0$. The reader may notice that for strictly nonsingular Lie groups the orbit of $\tau|_{\mathfrak{z}} \neq 0$ is the set of all $\tau' \in \mathfrak{g}^*$ with $\tau|_{\mathfrak{z}} \equiv \tau'|_{\mathfrak{z}}$.

Proof of Lemma 5.5 In the case $\varepsilon \equiv 1$ the lemma is just a consequence of [26, Theorem 7]. In the case $\varepsilon \neq 1$ we set $\Gamma' := \ker(\varepsilon)$ and we decompose

$$L_1^2(\Gamma' \backslash G) = L_1^2(\Gamma \backslash G) \oplus L_\varepsilon^2(\Gamma \backslash G)$$

as an orthogonal sum of G -modules. We already now the lemma for $(\Gamma, 1)$ and $(\Gamma', 1)$. Using the fact that every factor occurs only finitely many times, we can calculate the multiplicity in $L_\varepsilon^2(\Gamma \backslash G)$ as the difference of the multiplicities of $L_1^2(\Gamma' \backslash G)$ and $L_1^2(\Gamma \backslash G)$. If ε is projectable, then $L_{\Gamma, \varepsilon} = L_{\Gamma, 1} = L_{\Gamma', 1}$, but $\text{vol}(\overline{\Gamma'} \backslash \overline{G}) = 2 \text{vol}(\overline{\Gamma} \backslash \overline{G})$ and if ε is nonprojectable then $L_{\Gamma', 1} = L_{\Gamma, 1} \dot{\cup} L_{\Gamma, \varepsilon}$ but $\text{vol}(\overline{\Gamma'} \backslash \overline{G}) = \text{vol}(\overline{\Gamma} \backslash \overline{G})$. This proves the lemma for general ε . \square

THEOREM 5.6 *There is a family of nilmanifolds having*

1. constant spectrum of the Laplacian on functions,
2. nonconstant spectrum of the Laplacian on 1-forms
3. constant marked length spectrum
4. nonconstant spectrum of the Dirac operator for projectable spin structures,
5. constant spectrum of the Dirac operator for nonprojectable spin structures.

Proof Example I and II of [17] will have all these properties for suitably chosen lattices. Properties 1–3 have been shown in [17]. We will prove 4 and 5 for Example I. The arguments for Example II are analogous. To describe Example I let \mathfrak{g} be the 7-dimensional Lie algebra with orthonormal basis $\{X_1, X_2, X_3, X_4, Z_1, Z_2, \mathcal{Z}\}$ and Lie bracket

$$\begin{aligned} [X_1, X_2] &= [X_3, X_4] = Z_1 + \mathcal{Z} \\ [X_1, X_3] &= [X_4, X_2] = Z_2 \\ [X_2, X_3] &= [X_1, Z_1] = [Z_2, X_4] = \mathcal{Z}. \end{aligned}$$

All other Lie brackets of basis vectors shall be zero. This Lie algebra is strongly nonsingular and \mathcal{Z} generates the center of \mathfrak{g} . The matrix

$$\begin{pmatrix} \cos s & 0 & 0 & \sin s & 0 & 0 & 0 \\ 0 & \cos 2s & -\sin 2s & 0 & 0 & 0 & 0 \\ 0 & \sin 2s & \cos 2s & 0 & 0 & 0 & 0 \\ -\sin s & 0 & 0 & \cos s & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \cos s & -\sin s & 0 \\ 0 & 0 & 0 & 0 & \sin s & \cos s & 0 \\ 0 & 0 & 0 & 0 & -1 + \cos s & -\sin s & 1 \end{pmatrix}$$

defines a family of Lie algebra automorphisms φ_s of \mathfrak{g} and induces a family of Lie group automorphisms Φ_s of G , i.e. $\Phi_{s*} = \varphi_s$.

The center $Z(G)$ is fixed under the automorphism Φ_s . Therefore Φ_s induces an automorphism $\overline{\Phi}_s$ of $\overline{G} = G/Z(G)$ with commuting diagram

$$\begin{array}{ccc} G & \xrightarrow{\Phi_s} & G \\ \downarrow \pi & & \downarrow \pi \\ \overline{G} & \xrightarrow{\overline{\Phi}_s} & \overline{G}. \end{array}$$

These $\overline{\Phi}_s$ also act isometrically.

Let Γ be an arbitrary discrete cocompact subgroup of G . We write $\Gamma_s := \Phi_s(\Gamma)$, $\overline{\Gamma} = \pi(\Gamma)$ and $\overline{\Gamma}_s := \overline{\Phi}_s(\overline{\Gamma})$.

We now prove that the family $M_s := \Gamma_s \backslash G$, $s \in \mathbb{R}$, satisfies the theorem. Gornet showed in [17] that this family is isospectral for the Laplacian on functions, but not for the Laplacian on 1-forms. Furthermore she showed that Φ_s marks an isomorphism of the length spectrum from M_0 to M_s .

We have to show 4 and 5 and we start with 5. Let $\varepsilon : \Gamma \rightarrow \{-1, 1\}$ describe a nonprojectable spin structure on M_0 . As we want the spin structure to be chosen continuously in s , the spin structure on M_s is given by $\varepsilon \circ \Phi_s^{-1}$, and therefore also nonprojectable. Since $\overline{\Phi}_s$ is a family of isometric automorphisms, the volume of $\overline{\Gamma}_s \backslash \overline{G}$ is independent of s . Hence the second condition in Theorem 5.3 is satisfied and therefore $(\Phi_s \circ \Gamma, \varepsilon \circ \Phi_s^{-1})$ is a pairwise representation equivalent family. Theorem 5.1 implies that the spectrum of the Dirac operator does not depend on s .

To show 4 let M_0 carry a projectable spin structure described by $\varepsilon : \Gamma \rightarrow \{-1, 1\}$. Let $\overline{\varepsilon} : \overline{\Gamma} \rightarrow \{-1, 1\}$ describe the induced spin structure on $\overline{\Gamma} \backslash \overline{G}$. By continuity as above $\varepsilon_s := \varepsilon \circ \Phi_s^{-1}$ and $\overline{\varepsilon}_s := \overline{\varepsilon} \circ \overline{\Phi}_s^{-1}$ describe the spin structures on M_s and $\overline{M}_s := \overline{\Gamma}_s \backslash \overline{G}$ resp.

Just as in the proof of Theorem 5.1 we write the space of L^2 -spinors on M_s as

$$L^2(\Sigma_{\varepsilon_s}(\Gamma_s \backslash G)) = L^2_{\varepsilon_s}(\Gamma_s \backslash G) \otimes_{\mathbb{C}} \Sigma_7.$$

We can assume for our calculations that

$$\omega := \gamma(E_1) \cdot \gamma(E_2) \cdot \dots \cdot \gamma(E_7)$$

acts as $-\text{Id}$ on Σ_7 . As in the proof of Theorem 5.3 we write

$$L^2_{\varepsilon_s}(\Gamma_s \backslash G) = \bigoplus_{[\tau] \in \mathcal{T}} m_s([\tau]) H_{[\tau]}.$$

The Leibniz rule 5.1 shows that $H_{[\tau]} \otimes \Sigma_7$ is a subspace invariant under $\nabla_{E_j} : L^2(\Sigma_{\varepsilon_s}(\Gamma_s \backslash G)) \rightarrow L^2(\Sigma_{\varepsilon_s}(\Gamma_s \backslash G))$ and therefore it is also invariant under the Dirac operator $D = \sum_{j=1}^7 (\text{id} \otimes \gamma(E_j)) \cdot \nabla_{E_j}$. Hence

$$\text{spec } D^{M_s} = \bigcup_{[\tau] \in \mathcal{T}} \text{spec } D|_{H_{[\tau]} \otimes \Sigma_7}$$

and the multiplicity of each eigenvalue λ of D^{M_s} is the sum of $m_s([\tau])$ times the multiplicity of the eigenvalue λ in $D|_{H_{[\tau]} \otimes \Sigma_7}$.

To prove that $\text{spec } D^{M_s}$ is not constant in s it is sufficient to show that there is one eigenvalue λ_s in $\text{spec } D^{M_s}$ which is continuous and nonconstant in s .

Therefore we will show that there is a $\tau \in \mathfrak{g}^*$ vanishing on $[\mathfrak{g}, \mathfrak{g}]$ such that $\tau_s := \tau \circ \Phi_s^{-1}$ satisfies $m_s([\tau_s]) > 0$ and $D|_{H_{[\tau_s]} \otimes \Sigma_7}$ has at least one nonconstant eigenvalue depending continuously on s .

Because of $\tau_s|_{[\mathfrak{g}, \mathfrak{g}]} \equiv 0$ we know that $H_{[\tau_s]}$ is one-dimensional and $\exp X \in G$ acts on $H_{[\tau_s]}$ by multiplication with $e^{2\pi i \tau_s(X)}$. Elements f of $H_{[\tau_s]}$ satisfy $f(\exp X) = c \cdot e^{2\pi i \tau_s(X)}$ for all $X \in \mathfrak{g}$ with $c \in \mathbb{C}$. The multiplicity $m_s([\tau_s])$ is 1 if $e^{2\pi i \tau_s(X)} = \varepsilon_s(\exp X)$ for any $\exp X \in \Gamma_s$ and otherwise the multiplicity is 0. This condition is independent of s .

We will compute the determinant of $D|_{H_{[\tau_s]} \otimes \Sigma_7}$. Writing $E_1 := X_1, \dots, E_4 := X_4, E_5 := Z_1, E_6 := Z_2$ and $E_7 := \mathcal{Z}$, we apply (1.1) and get for $S \in \Sigma_7$ and $f \in H_{[\tau_s]}$:

$$\begin{aligned} D^{M_s}(fS) &= \sum_i (R_*(E_i)f) \gamma(E_i)S + \frac{1}{4} f \sum_{i,j,k} \Gamma_{ij}^k \gamma(E_i) \gamma(E_j) \gamma(E_k) S \\ &= \sum_{i=1}^4 2\pi i \tau_s(X_i) f \gamma(E_i) S + f A(S) \end{aligned}$$

where $A := (1/4) \sum_{i,j,k} \Gamma_{ij}^k \gamma(E_i) \gamma(E_j) \gamma(E_k) \in \text{End}(\Sigma_7)$ is constant in s . Using the Koszul formula one can express A in a suitable basis by the matrix

$$A = \frac{1}{4} \begin{pmatrix} -2 & 0 & 0 & 1 & 2 & i & i & -2 \\ 0 & 0 & 1 & 0 & -i & 0 & 0 & -i \\ 0 & 1 & 0 & 0 & i & 0 & 0 & i \\ 1 & 0 & 0 & 2 & 2 & -i & -i & -2 \\ 2 & i & -i & 2 & 2 & 0 & 0 & -1 \\ -i & 0 & 0 & i & 0 & 0 & -1 & 0 \\ -i & 0 & 0 & i & 0 & -1 & 0 & 0 \\ -2 & i & -i & -2 & -1 & 0 & 0 & -2 \end{pmatrix}$$

Then

$$\det D|_{H_{[\tau_s]} \otimes \Sigma_7} = \det \left[A + 2\pi i \sum_{i=1}^4 \tau_s(X_i) \gamma(E_i) \right]$$

and

$$\begin{aligned} \delta_\tau &:= \left. \frac{d}{ds} \right|_{s=0} \det D|_{H_{[\tau_s]} \otimes \Sigma_7} \\ &= -\frac{1}{64} \pi^2 (-\tau(X_1) \tau(X_2) + \tau(X_3) \tau(X_4)) \\ &\quad \left(4096 \pi^4 \tau(X_1)^4 + 4096 \pi^4 \tau(X_2)^4 + 4096 \pi^4 \tau(X_3)^4 + 4096 \pi^4 \tau(X_4)^4 \right. \\ &\quad + 8192 \pi^4 \tau(X_1)^2 \tau(X_2)^2 + 8192 \pi^4 \tau(X_1)^2 \tau(X_3)^2 + 8192 \pi^4 \tau(X_2)^2 \tau(X_3)^2 \\ &\quad + 8192 \pi^4 \tau(X_1)^2 \tau(X_4)^2 + 8192 \pi^4 \tau(X_2)^2 \tau(X_4)^2 + 8192 \pi^4 \tau(X_3)^2 \tau(X_4)^2 \\ &\quad + 512 \pi^2 \tau(X_3) \tau(X_1) + 128 \pi^2 \tau(X_4)^2 + 512 \pi^2 \tau(X_4) \tau(X_2) + 128 \pi^2 \tau(X_2)^2 \\ &\quad \left. + 128 \pi^2 \tau(X_1)^2 + 128 \pi^2 \tau(X_3)^2 + 1 \right). \end{aligned}$$

The set

$$\left\{ (\tau(X_1), \tau(X_2), \tau(X_3), \tau(X_4)) \mid \tau \in \mathfrak{g}^*, \tau|_{[\mathfrak{g}, \mathfrak{g}]} \equiv 0, m([\tau]) > 0 \right\}$$

is an affine lattice in \mathbb{R}^4 (i.e. up to translation a discrete cocompact subgroup). But the only polynomial vanishing on an affine lattice of \mathbb{R}^4 is the zero polynomial. Therefore there exist $\tau \in \mathfrak{g}^*$ with $\tau|_{[\mathfrak{g}, \mathfrak{g}]} \equiv 0$, $m([\tau]) > 0$ and $\delta_\tau \neq 0$. \square

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Authors' address:

Mathematisches Institut
Universität Freiburg
Eckerstr. 1
79104 Freiburg
Germany

e-mail:

`ammann@mathematik.uni-freiburg.de`
`baer@mathematik.uni-freiburg.de`