

THE DUST SUBLIMATION RADIUS AS AN OUTER ENVELOPE TO THE BULK OF THE NARROW FE K α LINE EMISSION IN TYPE 1 AGN

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ABSTRACT

The Fe K α emission line is the most ubiquitous feature in the X-ray spectra of active galactic nuclei (AGN), but the origin of its narrow core remains uncertain. Here, we investigate the connection between the sizes of the Fe K α core emission regions and the measured sizes of the dusty tori in 13 local Type 1 AGN. The observed Fe K α emission radii (R_{Fe}) are determined from spectrally resolved line widths in X-ray grating spectra, and the dust sublimation radii (R_{dust}) are measured either from optical/near-infrared reverberation time lags or from resolved near-infrared interferometric data. This direct comparison shows, on an object-by-object basis, that the dust sublimation radius forms an outer envelope to the bulk of the Fe K α emission. R_{Fe} matches R_{dust} well in the AGN with the best constrained line widths currently. In a significant fraction of objects without a clear narrow line core, R_{Fe} is similar to, or smaller than the radius of the optical broad line region. These facts place important constraints on the torus geometries for our sample. Extended tori in which the solid angle of fluorescing gas peaks at well beyond the dust sublimation radius can be ruled out. We also test for luminosity scalings of R_{Fe} , finding that Eddington ratio is not a prime driver in determining the line location in our sample. We discuss in detail potential caveats due to data analysis and instrumental limitations, simplistic line modeling, uncertain black hole masses, as well as sample selection, showing that none of these is likely to bias our core result. The calorimeter on board *Astro-H* will soon vastly increase the parameter space over which line measurements can be made, overcoming many of these limitations.

Subject headings: galaxies: active — X-rays: galaxies — Infrared: galaxies

1. INTRODUCTION

The Fe K α line is the most prominent emission feature observed in the spectra of accreting systems. It arises as a result of fluorescence in predominantly cold gas, and has been observed in X-ray binaries as well as active galactic nuclei (AGN). The broad component has been modeled as originating in the inner accretion disk as a result of relativistic broadening (Fabian et al. 1989). Despite many years of study, though, the origin of the *narrow* core of the line in AGN is still unclear. Previous works have found no strong correlation of the Fe line width with those seen in the broad line region (BLR), and suggest that there are likely multiple sites of origin, including the dusty torus, the BLR, as well as the outer accretion disk (Yaqoob & Padmanabhan 2004; Nandra 2006; Shu et al. 2010).

All modern AGN X-ray radiative transfer models include fluorescence emission computed self-consistently for a variety of geometries and obscuring column densities in the accretion disk as well as the torus (e.g. George & Fabian 1991; Nandra et al. 2007; Murphy & Yaqoob 2009; Brightman & Nandra 2011), and recent high quality X-ray spectra for many AGN – obscured as well as unobscured – bear out the correspondence between transmission, Compton scattering and fluorescence (e.g. Yaqoob 2012; Brenneman et al. 2014; Arévalo et al. 2014; Gandhi et al. 2014; Bauer et al. 2014; Gandhi et al. 2015). Luminosity and covering factor scalings between the infrared and X-ray regimes may also support a close connection between the line-emitting gas and distribution of warm dust (Gandhi et al. 2009; Ricci et al. 2014; Toba et al. 2014), but any spatial correspondence between the Fe K α emission zone and the torus still remains to be tested in detail.

The main issue hindering progress in pinning down the origin of the line core is lack of high quality data. In the X-rays, the broad and narrow components of the line must be disentangled, and the narrow core is only resolvable using grating spectra in the brightest AGN at present. In the infrared, the dust emission zones span sub-pc to pc size scales (e.g. Kishimoto et al. 2011a; Tristram et al. 2009; Burtscher et al. 2013), which require interferometric techniques in order to resolve directly. Infrared reverberation mapping is a growing field and provides an alternate means to infer dust emitter sizes (e.g. Suganuma et al. 2006; Koshida et al. 2014).

There have been recent advances on all these fronts, with a growing number of sources with commensurate data now available. Here we present a comparison of the directly resolved inner edge of the dusty tori (using either interferometry or reverberation mapping) with the sizes of the narrow Fe K α line core regions for a sample of 13 local AGN. We investigate the detailed connection between the two in individual objects, discuss the limitations in the current data, and make comparisons to other recent works. We also test whether the AGN Eddington accretion rates play a role in determining the Fe K α location. Our study, albeit on a small sample, sets the stage for much larger studies which will be possible once the calorimeter on board *Astro-H* (Takahashi et al. 2014) begins operation in 2016.

2. SAMPLE

2.1. X-ray grating spectra

Our starting sample is the High Energy Grating (HEG) sample of Shu et al. (2010), the largest sample to date of high spectral resolution data covering the ~ 6.4 keV energy regime. This includes 36 unique galaxies below $z=0.3$. In 27 of these, a measurement of the velocity full widths at half

maximum (v_{FWHM}) of the narrow line core was reported by the authors. This includes positive FWHM measurements, as well as unresolved line upper limits. The mean v_{FWHM} for this sample was found to be $2060 \pm 230 \text{ km s}^{-1}$. The HEG spectral resolution is 0.012 \AA , corresponding to a velocity resolution of $\approx 1860 \text{ km s}^{-1}$ at 6.4 keV , and is accounted for in the line modeling.

We have chosen measurements that are likely to best represent the narrow line core. In this respect, note that several new observations for NGC 4051 have become available since the work of Shu et al. (2010). We use the results discussed by Lobban et al. (2011) and Shu et al. (2012). The former work found that the presence of unresolved line component at 6.4 keV significantly improved the fit of the line complex in the time-averaged data from 12 observations totaling $\sim 300 \text{ ks}$ of exposure. Since this component is the one most likely associated with distant material, we treat the core as being unresolved. This leads to an upper limit on v_{FWHM} and a corresponding lower limit on R_{Fe} .

IC 4329A is also treated as upper limit on v_{FWHM} . Although a broad component to the line is detected by Shu et al. (2010, see also McKernan & Yaqoob 2004) with $v_{\text{FWHM}} = 18830_{-9620}^{+18590} \text{ km s}^{-1}$, the line peak energy in this case is found to be $E_{\text{peak}} = 6.305_{-0.096}^{+0.139} \text{ keV}$. An alternative fit with an unresolved line by the authors yielded a more plausible value of $E_{\text{peak}} = 6.399_{-0.005}^{+0.006} \text{ keV}$ closer to the expected neutral line energy, and the authors stress that these are more reliable measurements of the true narrow core.

For measurement of the Fe K α emission radii (R_{Fe}), we assume virial motion for the emitting clouds. Then,

$$R_{\text{Fe}} = \frac{GM_{\text{BH}}}{v^2} \quad (1)$$

where M_{BH} is the black hole mass, and $v = \sqrt{3}/2 v_{\text{FWHM}}$. The correction factor of $\sqrt{3}/2$ arises under the assumption of an isotropic velocity distribution and accounting for the fact that the line-of-sight velocity dispersion is half of the FWHM (cf. Netzer 1990; Peterson et al. 2004). Uncertainties related to unknown geometric projection factors are discussed below and in the Appendix. Where possible, M_{BH} values based upon reverberation mapping measurements are used, mostly based upon the H β time lag. This was the case for all sources except IRAS 13349+2438 which is derived from constraints on modeling its spectral energy distribution. In this case, we assumed a large uncertainty equal to the mass measurement itself in order to account for the lack of reverberation mapping. In addition, for NGC 4151, we use a recent measurement obtained from dynamical modeling. References are given in Table 1. Uncertainties on R_{Fe} were determined by joint Monte Carlo sampling of M_{BH} and v in log space.

One systematic uncertainty in converting velocities to sizes comes from the unknown geometric projection, which may result in underestimation of the true space velocities, leading to overestimated sizes. This is a well-known effect in BLR reverberation mapping and single-epoch estimates of M_{BH} , and is captured in the f -factor (e.g. Peterson et al. 2004). If we posit that conservation of angular momentum of the accreting matter leads to a flattened geometry, then we can expect that the projection effects are very similar for the X-ray emitting region and the BLR. In this case, size measurements of these regions will be affected in a similar way for any individual object, so comparisons between these quantities will not

be affected. On the sample level, the distribution of geometric inclinations will lead to a widening of any size-luminosity relation above the intrinsic dispersion by a factor of the order unity. Given that our sample covers four decades in luminosity and two decades in sizes, we can expect that the essence of the relation will be preserved even for unknown projection effects in individual sources. In the Appendix, we will discuss alternate size comparisons using reverberation-independent M_{BH} values – a test which avoids the above assumption.

2.2. Infrared and optical data

Near-IR (NIR) observations probe the emission region of the sublimation zone in the torus at temperatures of $\sim 1500 \text{ K}$. The size of this region can be inferred either from reverberation mapping or directly from interferometry. Although both types of measurements are qualitatively similar, the small quantitative differences between both are probably related to the detailed dust distribution (e.g. Kishimoto et al. 2009; Hönig & Kishimoto 2011; Koshida et al. 2014; Hönig et al. 2014). Here we collect and use both types of data from the literature (collectively referred to as R_{dust}). We note that R_{dust} measurements are available mostly for Type 1 AGN, because in Type 2 AGN the innermost hot dust is not easily visible.

The dust reverberation mapping radii ($R_{\text{dust,rev}}$) are the result of a long-term monitoring campaign with the MAGNUM telescope by Koshida et al. (2014) and have been inferred from V - and K -band light curve time lags (τ_{dust}) as $R_{\text{dust,rev}} = c\tau_{\text{dust}}/(1+z)$ where c denotes the speed of light and the $(1+z)$ factor corrects for cosmological time dilation. The interferometric sizes ($R_{\text{dust,intf}}$) are based on Keck Interferometer data and represent radii of a thin-ring model (Kishimoto et al. 2009, 2011a, 2013).

Of the 27 sources with v_{FWHM} from the sample of Shu et al. (2010), 7 have $R_{\text{dust,intf}}$ measurements, and 10 have published τ_{dust} values. Our final sample of objects with measured values of R_{Fe} as well as R_{dust} comprises 13 unique sources. These are listed in Table 1.

For these objects, we also compute the radii of their optical BLRs (R_{BLR}) for comparison to R_{Fe} . Measurements of the reverberation time lag of the H β emission line ($\tau_{\text{H}\beta}$) for 12 AGN of our sample – i.e. all except IRAS 13349+2438 – have been tabulated in Bentz et al. (2009), which we use here. Then, $R_{\text{BLR}} = c\tau_{\text{H}\beta}/(1+z)$.

We follow this by investigating any possible relation between the locations of emitting regions and the continuum emission. As a proxy of the latter, we consider the monochromatic continuum luminosities at 5500 \AA (L_{5500}), whose values are gathered from Kishimoto et al. (2011a), Suganuma et al. (2006) and Bentz et al. (2009). These are based upon fitting of their spectral energy distributions and corrected for starlight contamination. A mean uncertainty of 0.1 dex is assumed for L_{5500} . These are also listed in Table 1.

3. RESULTS

Fig. 1 shows the comparison between R_{Fe} and R_{dust} . For all objects, the measurements or limits on $R_{\text{dust,intf}}$ lie a little above $R_{\text{dust,rev}}$. For sources with measurements of both radii, $R_{\text{dust,intf}}$ is larger than $R_{\text{dust,rev}}$ by an average factor of $2.3 (\pm 0.3 \text{ mean standard deviation})$. Reverberation is sensitive to the fastest dust response with changing incident radiation, whereas the emission probed in interferometry is more sensitive to the average emitting surface area which is likely to peak at somewhat larger radii (e.g. Kishimoto et al. 2009;

Hönig & Kishimoto 2011; Koshida et al. 2014; Hönig et al. 2014). In any case, the trends discussed below are similar for both measures of R_{dust} .

For 9 of the 13 AGN, the measurements or limits on R_{Fe} are fully consistent with R_{dust} . These are NGC 3783, NGC 4151, NGC 4593, NGC 5548, Mrk 509, 3C 273, IRAS 13349+2438, NGC 4051 and IC 4329A. The remaining 4 sources (NGC 3516, NGC 7469, Mrk 590 and Fairall 9) have R_{Fe} values significantly smaller than R_{dust} . The ratio of $R_{\text{dust}}/R_{\text{Fe}}$ in these 4 ranges over 2.5 (NGC 3516) to 75 (Fairall 9).

The figure also plots R_{BLR} on the right-hand axis and reveals some interesting results for individual sources. Whereas the median value of R_{BLR} is 8 times smaller than R_{Fe} for the full sample, there are 6 objects (NGC 4051, NGC 4593, NGC 7469, Mrk 509, Fairall 9 and 3C 273) for which the R_{Fe} measurement or limit is entirely consistent with R_{BLR} . For NGC 3516, R_{Fe} lies below R_{dust} and the uncertainty estimates on R_{BLR} and R_{Fe} do not overlap, implying an Fe K α origin in an intermediate zone between the BLR and the torus. On the other hand, for Mrk 590, R_{Fe} is significantly smaller than even R_{BLR} . Finally, for Mrk 509 and Mrk 590, R_{BLR} lies within a factor of 1.5 of R_{dust} , implying a close proximity of the BLR clouds with the inner extent of the torus.

Examining the overall distribution of sources, the most important feature is the absence of sources significantly above the line of equality. Whereas there is significant scatter of sources to small values of R_{Fe} , for no source is R_{Fe} much greater than R_{dust} . The two lower limits are also consistent with this line. We discuss these results at length in the next section.

We next test whether the location of the Fe K α line is driven by fundamental luminosity scalings. We first plot R_{Fe} as a function of L_{5500} in Fig. 2. This shows a close correspondence with Fig. 1 in all aspects, with an absence of sources on size scales associated with those above the torus size–luminosity relation (Barvainis 1987; Kishimoto et al. 2011a), and a strong scatter of sources below. We tested for any relation between the R_{Fe} and L_{5500} with Monte Carlo resampling and computing the distribution of Spearman rank coefficients (ρ) for the randomized ensembles. We find $\bar{\rho} = 0.45^{+0.21}_{-0.25}$ (90%) and insignificant p-values over the entire ensemble, suggestive of only a weak positive correlation. Any stronger trend is hidden by the large scatter in the current (small) sample.

In Fig. 3, we remove black hole mass scaling from both axes by plotting R_{Fe} in units of the Gravitational radius ($R_g = GM_{\text{BH}}/c^2$) as a function of the Eddington ratio ($L_{\text{Bol}}/L_{\text{Edd}}$). The latter quantity is taken as a proxy of the (specific) accretion rate. Bolometric luminosities (L_{Bol}) are approximate estimates based upon a correction factor of 6 to L_{5500} (Scott et al. 2004; Richards et al. 2006). We cross-checked these L_{Bol} values with those reported from the broadband modeling carried out by Vasudevan & Fabian (2009), who include all of our sample objects except for IRAS 13349+2438, and found similar results. Objects with the smallest uncertainties on R_{Fe} occupy a narrow region of R_{Fe}/R_g around $\approx 20,000 \text{ km s}^{-1}$ (cf. Shu et al. 2011). There appears to be an intriguing hint that sources with high $L_{\text{Bol}}/L_{\text{Edd}}$ (above ≈ 0.06) show smaller detected values of R_{Fe}/R_g than sources at low Eddington ratios. But this is not a significant trend in the current sample at least, and for the comparison between R_{Fe}/R_g and $L_{\text{Bol}}/L_{\text{Edd}}$ we find $\bar{\rho} = -$

$0.10^{+0.34}_{-0.42}$ (90%). The second panel in the figure illustrates that dust is present on scales similar to R_{Fe}/R_g in more than half the sample.

4. DISCUSSION

4.1. Why does R_{dust} serve as an outer envelope to R_{Fe} ?

Our key result is that the dust sublimation radius forms an outer envelope to the Fe K α emission zone in Type 1 AGN, with uncertainties allowing R_{Fe} at most a factor of a few times R_{dust} for our sample. This is not an obvious prediction given that the gas distribution is likely to extend continuously on scales both larger and smaller than the innermost dust radius, and that tori may be clumpy with low optical depth ‘holes’ allowing radiation to penetrate well into its body. Our work is thus constraining for AGN torus models, as follows.

For axisymmetric gas distributions, the emitted K α line intensity at any radius r depends upon the local solid angle $\Omega_{\text{K}\alpha}(r)$ of gas illuminated by the AGN (Krolik & Kallman 1987). The absence of sources with $R_{\text{Fe}} \gg R_{\text{dust}}$ then argues against geometries in which $\Omega_{\text{K}\alpha}(r \gg R_{\text{dust}}) > \Omega_{\text{K}\alpha}(R_{\text{dust}})$, i.e. where the line emitting area increases strongly with radius beyond R_{dust} . Line emission will not cease exactly at R_{dust} , which would be an unphysical scenario. Fig. 1 shows that the R_{Fe} values of up to a factor of a few times R_{dust} are allowed within the uncertainties. How do typical torus models match up to these constraints? In the geometry adopted by the widely-used MYTORUS model (Murphy & Yaqoob 2009), the torus is assumed to have a donut shape with a circular cross-section and a half opening angle of 60° . Radiation from the nucleus that directly illuminates the torus would impinge upon the donut surface out to distances of the surface tangent line. In the MYTORUS geometry, this distance is $1.7 R_{\text{dust}}$ where R_{dust} is the inner edge of the torus. This is consistent within the constraints from Fig. 1. On the other hand, tori with strong flaring in height as a function of radius would present large illuminated surface solid angles at large nuclear distances, and would not easily satisfy the above constraints on R_{Fe} relative to R_{dust} .

Radiation would also penetrate below the torus surface, but if the integrated column density through the torus¹ is Compton-thick as is generally thought to be the case, then absorption and Compton-scattering would quickly deplete line photons in the body of the torus, so $\Omega_{\text{K}\alpha}$ refers to the solid angle of the reflecting torus surface down to an effective optical depth $\tau_{\text{K}\alpha} \sim 1$. A natural prediction of this scenario is that sources in which the obscuring tori are likely to be *Compton-thin* (e.g. NGC 2110, Marinucci et al. 2015; NGC 7213, Ursini et al. 2015) need not follow the trend of Fig. 1, i.e. their Fe K α emission need not be restricted to lie around R_{dust} .

We also emphasize that our results do not exclude the presence of gas on extended scales altogether. If more distant gas is not excited by AGN radiation, it will not fluoresce and will remain invisible. Furthermore, studies based upon line width measurements will selectively target the strongest emission regions present on the innermost (fastest) scales. More distant (and fainter) emitting components could become apparent if our direct line-of-sight to the inner torus were obscured by material optically-thick to the line photons, as may be expected in Type 2 and Compton-thick AGN (cf. the detection of extended emission in NGC 4945 by Marinucci et al. 2012

¹ It should be noted that this is different from the line-of-sight column density, which would vary with inclination angle.

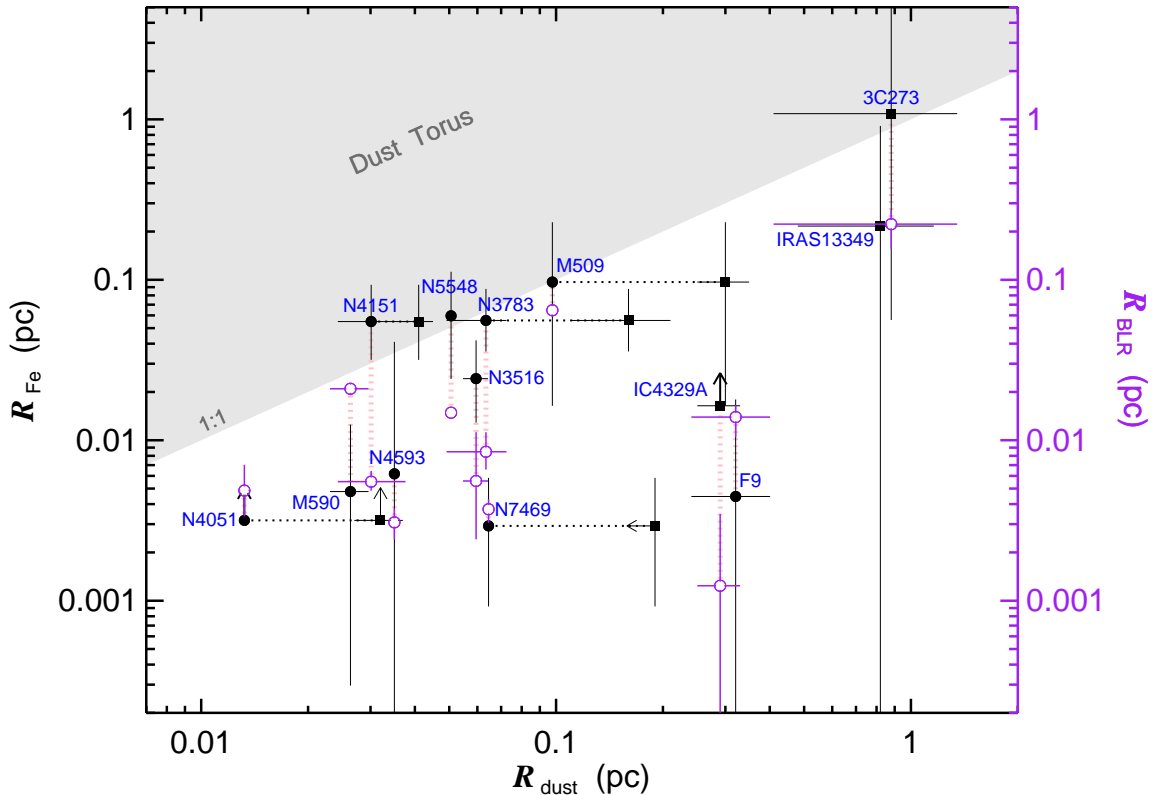


FIG. 1.— NIR hot dust sublimation radii (R_{dust}) vs. Fe $K\alpha$ line region radii (R_{Fe}) shown in black. Circles represent R_{dust} from reverberation lags and squares are interferometric measurements. The unfilled violet symbols are the BLR radii (R_{BLR}) based upon optical reverberation mapping. Dotted lines connect these various measurements for any one source. The gray shaded zone represents $R_{\text{Fe}} > R_{\text{dust}}$, i.e. scales commensurate with the body of the torus and beyond.

and in NGC 1068 by Bauer et al. 2014). Monte Carlo radiative transfer simulations can place detailed constraints on the radial gas geometry in the torus.

There are several potential complexities related to modeling of the Fe $K\alpha$ line and sample selection which must be investigated to test the robustness of the above constraints. These caveats are discussed at length in Section 4.4.

4.2. Fe $K\alpha$ emission radii much smaller than the torus

Conversely, the scatter in R_{Fe} to scales much smaller than R_{dust} argues for Fe $K\alpha$ emission from extended regions. This result is not new (cf. Yaqoob & Padmanabhan 2004; Nandra 2006; Shu et al. 2010), but is now inferred from the direct size comparison of Fig. 1. The absence of any obvious luminosity scaling in Fig. 2 supports this scenario, and Fig. 3 shows that any Eddington rate-driven physical components (e.g. accretion disk outflows) do not control the origin of Fe $K\alpha$, at least over the range of $L_{\text{Bol}}/L_{\text{Edd}}$ that we sample. The outer accretion disk and BLR clouds have been invoked as Fe $K\alpha$ emitters in several objects. For example, rapidly variable narrow Fe $K\alpha$ line from small scales have been found in Mrk 509 (Ponti et al. 2013), Mrk 841 (Petrucci et al. 2002) and NGC 7314 (Yaqoob et al. 2003). For NGC 7213, Bianchi et al. (2008) found that resolved Fe $K\alpha$ line width matches the optical H α line width and argued for a BLR origin in this case. Shu et al. (2010) found that a significant fraction of their sample showed Fe $K\alpha$ line widths consistent with, or broader than, the widths of typical optical and infrared

photoionized lines, concluding that the Fe $K\alpha$ emission zone location (relative to that of the BLR) genuinely varies from source to source. A BLR origin is also consistent for several sources in our sample.

Fig. 1 shows that in many of the sources with the best constraints on R_{Fe} (e.g. NGC 4151, NGC 5548, NGC 3783), R_{Fe} is well matched to R_{dust} . However, this does not imply that low grating data S/N is the cause of mismatching R_{Fe} and R_{dust} values in other objects. It is certainly true that the largest Fe $K\alpha$ error bars in Figs. 1 and 2 must be unphysical and are likely to be a result of the line fits modeling the underlying continuum (see, for example, the discussion on Fairall 9 by Shu et al. 2010). In addition, some of the sources with ill-constrained R_{Fe} also have relatively short exposure times (t_{exp}) – e.g. both Fairall 9 and NGC 4593 were observed for 80 ks with the HEG as compared to a median $t_{\text{exp}} = 150$ ks for the sample and a maximum of 890 ks in the case of NGC 3783. But as emphasized by Yaqoob & Padmanabhan (2004), the uncertainties are not simply a function of S/N, with the grating data showing that the peak Fe $K\alpha$ energy of 6.4 keV is not dominated by a narrow core in these sources.² Moreover, if we use stricter confidence intervals on v_{FWHM} (e.g. the 90% range from Shu et al. 2010), the sources with R_{Fe} significantly less than R_{dust} still remain discrepant in all

² Several of these including Fairall 9 and NGC 4593 belong to the ‘Group 2’ classification of objects in Yaqoob & Padmanabhan (2004) with weaker or absent narrow cores.

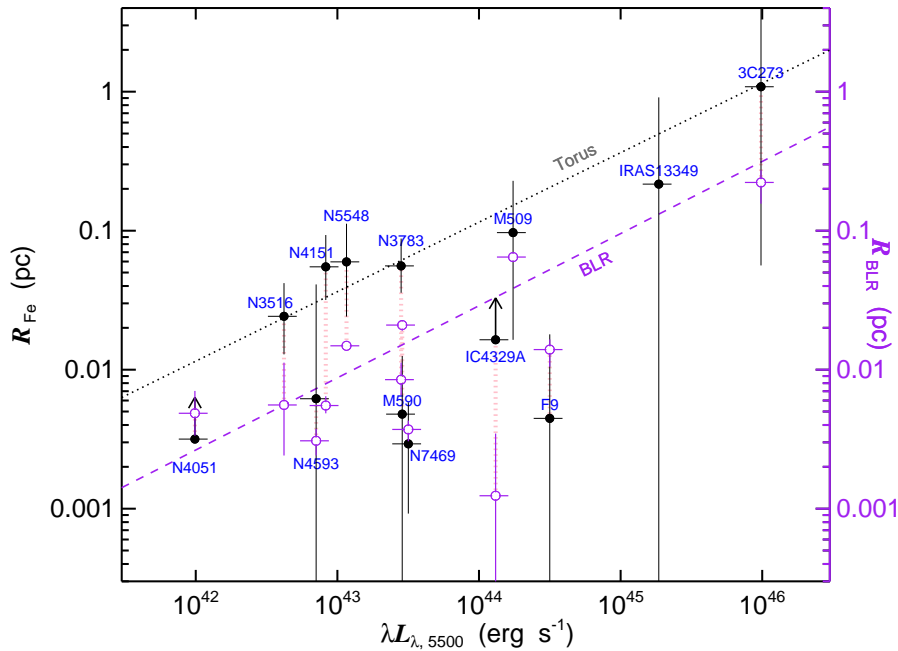


FIG. 2.— Optical AGN luminosity (L_{5500}) vs. R_{Fe} (filled) and R_{BLR} (unfilled). The dotted line is the torus size–luminosity relation Eq. 1 from Kishimoto et al. (2011b, originally from Suganuma et al. 2006) and the dashed line is the corresponding relation for the BLR from Bentz et al. (2009).

cases. The large uncertainties then appear to reflect an origin from multiple regions in the AGN environment including (fast) clouds present on many scales. Longer grating observations of the sources with large uncertainties on v_{FWHM} will help to confirm the absence of dominant narrow cores, or to disentangle fainter core components.

4.3. Comparison to other recent works

In two other recent works, Jiang et al. (2011) and Minezaki & Matsushita (2015) have investigated the use of the Fe K α line for measurement of black hole masses in AGN. Jiang et al. (2011) assumed a torus origin for the Fe K α line with radii derived from infrared reverberation and found a consistent trend between the masses predicted assuming isotropic motion of the Fe K α clouds with the masses estimated from optical reverberation mapping. However, a scaling factor was required for the two sets of masses to match each other – the scaling being equivalent to an average line core v_{FWHM} that is $2.6^{+0.9}_{-0.4}$ times broader than predicted from initial assumption of a torus origin.

Minezaki & Matsushita (2015) studied a restricted AGN sample with ‘best constraints’ on the FWHM of the Fe K α line, and found that generally the Fe K α emission zone is located between the reverberation radii of the broad H β emission line and the dust torus emission, the latter being estimated statistically (i.e. using a fixed ratio of R_{dust} to R_{BLR} based on K -band reverberation measurements). The criteria that constitute ‘best constraints’ are not quantified by Minezaki & Matsushita (2015) but their selection excludes most objects with the broadest Fe K α lines, as well as sources with unresolved Fe K α lines. The former category of sources includes objects in which the Fe K α emission zone is consistent with, or significantly smaller than the BLR (as we discussed in Section 4.2), while the unresolved sources could

potentially show Fe K α emission from outside the torus reverberation radius. Jiang et al. (2011), on the other hand, do include two sources with upper limits on v_{FWHM} ,³ and note the possibility that the narrow line core originates from radii smaller than the infrared radiation.

Our main result was derived entirely independent of the above works. While being consistent, a key distinction is that we go beyond average comparisons on a statistical basis. Importantly, from our object-by-object comparison, we make a prediction that R_{dust} forms an approximate outer envelope for normal Type 1 AGN. This prediction can be easily falsified if R_{Fe} were observed to be much larger than R_{dust} in future observations (see also discussion in next section). In this sense, it is much more powerful than inferences based upon statistical trends. This holds despite the fact that we have not restricted the base sample of Shu et al. (2010) according to FWHM constraints, and have included upper limits as well. We have then explored the consequences of such a postulation for the structure of AGN tori (Section 4.1) – aspects not mentioned in these other works whose focus is M_{BH} measurement. We note that since Minezaki & Matsushita (2015) do not include sources in which the peak Fe K α energy is not dominated by a narrow core,⁴ they do not probe the smallest scales that we discuss in Section 4.2 above.

Next, we will address various potential caveats which could affect our main result.

4.4. Potential caveats

Several potential caveats need to be kept in mind regarding current analyses of Fe K α line profiles, and these are discussed in detail below.

³ They treat Mrk 509 as an upper limit on v_{FWHM} . The resultant lower limit on R_{Fe} is entirely consistent with our trend between R_{dust} and R_{Fe} .

⁴ Except for their one ‘outlier’ NGC 7469.

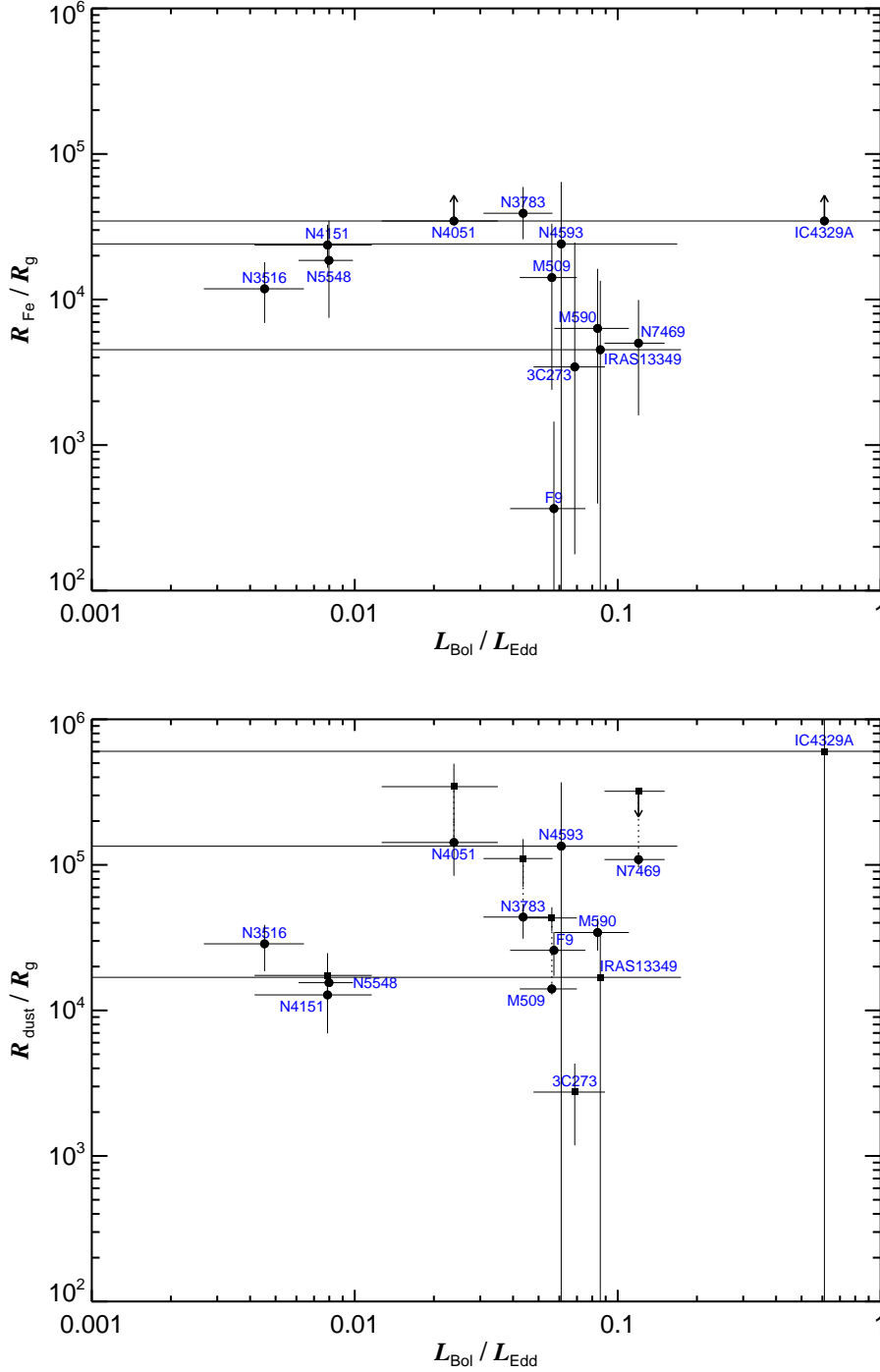


FIG. 3.— (Top) Eddington ratios ($L_{\text{Bol}}/L_{\text{Edd}}$) vs. R_{Fe} in units of R_g . (Bottom) $L_{\text{Bol}}/L_{\text{Edd}}$ vs. R_{dust} in units of R_g . Symbols are as in Fig. 1.

4.4.1. Combining multiple HEG datasets

For several sources (e.g. NGC 3783, NGC 4051) multiple HEG datasets have been combined in order to maximize the signal:noise and place the best constraints on the line width. This methodology ignores any potential change in the line emission zone between observations. From the individual measurements presented by Shu et al. (2010), there appears to be some variation in $\text{FWHM}(\text{Fe K}\alpha)$ for 2 of the 5 sources with multiple observations spaced by a few months to years

(these 2 sources are NGC 3516 and NGC 5548). But this is true only at the 68% confidence level, and the 90% confidence intervals show no such variation. In addition, the mean $\text{FWHM}(\text{Fe K}\alpha)$ measured *per source* (when combining multiple observations) for the sample of Shu et al. (2010) is entirely consistent with the mean $\text{FWHM}(\text{Fe K}\alpha)$ measured *per observation*. These facts suggest that combining multiple observations does not strongly affect our results. However, note that for two sources from our sample (3C 273 and

IRAS 13349+2438), the Fe K α line is detected in only one observation and it is unknown whether this is due to intrinsic and strong variability of the line.

The effective FWHM in a combined dataset can also be artificially broadened if the centroid energy changes between observations. Shu et al. (2010) find a sharp peak in the rest-frame centroid energies for their sample at ~ 6.400 keV and quote mean sample line energies centered on the neutral line to within -12 eV and $+3$ eV (in comparison to the HEG spectral resolution of ~ 40 eV), with the *per source* analysis fully consistent with the *per observation* one. Only in 2 cases out of 32 (3C 273 and 4C 74.26) did the centroid deviate strongly from the expected neutral line energy, although they note that these cases are marginal.

We also note that the above centroid energy range is equivalent to systemic velocity shifts of ≈ 140 km s $^{-1}$ (blueshift) up to ≈ 560 km s $^{-1}$ (redshifts) for the fluorescing material. In contrast, the median escape velocity at R_{dust} for our sources is ~ 2400 km s $^{-1}$ and the median line width for sources with an FWHM detection is 3400 km s $^{-1}$, arguing against any bias related to systemic outflows or inflows.

4.4.2. The limited HEG spectral resolution

What is the impact of the limiting spectral resolution of HEG for our results? The velocity resolution of $v_{\text{FWHM}} = 1860$ km s $^{-1}$ corresponds to $R_{\text{Fe}} = 0.08$ pc for $M_{\text{BH}} = 5 \times 10^7 M_{\odot}$ (the median M_{BH} for our sample). Whereas much narrower lines can be deconvolved given enough S/N in the data (e.g. see the case of NGC 3783 which has the best S/N and a resolved line even at 99% confidence), this may not be possible if the S/N is weaker. So are we biased to objects with broader lines, hence creating an artificial limit to R_{Fe} ? We argue that this is not the case, because for any given line flux, it is always easier to detect the line if it is unresolved. If the line is significantly broader than the spectral resolution, its peak decreases and the flux is spread over a range of energy bins, in which case detection naturally becomes more difficult. Despite this fact, most sources in our sample already show *resolved* lines with the HEG. And we emphasize that we are not excluding objects with unresolved lines which would show only upper limits to v_{FWHM} . Together, these argue against an artificial limit being created by the HEG spectral resolution in our analysis.

4.4.3. Biases due to multiple unmodeled line components

Another caveat is related to modeling the 6.4 keV feature as a single Gaussian and ignoring various intrinsic contributions to the feature such as the Compton shoulder which results from down-scattering of line photons off circumnuclear gas. However, the flux in this component is expected to be only ~ 10 – 20% of that in the narrow line core for reflection off Compton-thick material (Matt 2002), so a single Gaussian fit would only be slightly broader than in the absence of a Compton shoulder. We quantified this effect by simulating HEG Fe complex spectra for a range of column densities upto $N_{\text{H}} = 10^{25}$ cm $^{-2}$ using the MYTORUSL model (Murphy & Yaqoob 2009), which self-consistently models fluorescence and Compton scattering off a toroidal reflector. Another set of simulations using the GAUSS model alone for a single line were carried out for comparison. The intrinsic Fe K α line width was set to $\sigma = 1$ eV, well below the HEG resolution. We used the standard first order HEG canned re-

sponse matrices and effective area files⁵ for the simulations and covered a broad range in S/N similar to those in the real data. All spectra were then fit with single Gaussian lines. We found that the effective Gaussian line widths for the simulated MYTORUSL spectra were at most 18% broader than for the single GAUSS models. Correcting for this would increase R_{Fe} by at most 39% or 0.14 dex – at most a minor effect as can be seen in Fig. 1. An independent constraint comes from the fact that Shu et al. (2010) found a sharp centroid peak for their sample of close to the expect mean rest-frame line energy at 6.400 keV. If the Compton shoulder were strongly skewing the line fits, the centroid would be expected to lie at lower energies, but this not the case.

Similarly, blended lines could produce artificial broadening in single Gaussian fits. Such lines may arise from multiple ionization stages of Fe, and we here quantify their influence on the measurement of the effective FWHM, and hence on the derivation of R_{Fe} . Firstly, it is worth emphasizing that the observed line centroid energy distribution peaking at ≈ 6.40 keV for the full HEG sample argues strongly for lowly ionized states (typically below Fe XVII), as discussed by Shu et al. (2010). Upon examining the centroid energies (E) reported by Shu et al. (2010) for our sample of 13 AGN, we find a median $E = 6.403$ keV (± 0.010 keV mean standard deviation). The upper end of this range ($E = 6.413$ keV) corresponds closely to an ionization stage of Fe XIV (cf. Liu et al. 2010; Palmeri et al. 2003; Mendoza et al. 2004). A further constraint on the strongest ionization stage comes from the distribution of E values for the resolved lines in the longest HEG exposures. There are 5 objects with exposure time of more than ≈ 150 ks: NGC 3516, NGC 3783, NGC 4151, NGC 5548 and IRAS 13349+2438. All except IRAS 13349+2438⁶ show $E < 6.415$ keV at the conservative 90% confidence level, again consistent with ionization stages lower than Fe XIV–XV. Based upon this upper energy range, we assess a worst case scenario where all ionization stages from Fe I to Fe XV contribute maximally to artificially broadening the line. All stages are taken to possess the same component line strength and the same intrinsic FWHM. The centroid energies for these stages range over $E \approx 6.394$ – 6.415 keV (cf. Fig. 2 of Liu et al. 2010), or a full range of $\Delta E = 21$ eV. This corresponds to a velocity broadening of 980 km s $^{-1}$. This velocity should be subtracted in quadrature from the measured FWHM values when modeling the 6.4 keV feature as a single Gaussian. The largest effect will be on the source with the smallest FWHM for which the line is resolved, which is NGC 3783 (FWHM = 1750 km s $^{-1}$; see Table 1), and in this case, the intrinsic FWHM of the individual ionization stages is only 20% smaller, at 1450 km s $^{-1}$. The corresponding R_{Fe} increases by 45% and is still fully consistent with R_{dust} . The effect of line blending for the other sources in our sample is much less.

4.4.4. Sample representativeness

The parent sample of Shu et al. (2010) is certainly not a complete sample in a physical sense that is ideally suited for statistical studies. Despite being the largest sample with high quality Fe K α measurements to date, there is some unavoidable bias of HEG observations towards sources with a previously known presence of Fe K α (either broad or narrow) and

⁵ <http://cxc.harvard.edu/caldb/propplan/grating>

⁶ This source is already noted by Shu et al. (2010) as one of the few anomalous cases with potential line variability between observations. The data for this source also yield only very weak constraints on the overall line properties including flux and FWHM (see Fig. 1).

to objects that were observed with other aims in mind (see also discussion on this point in Yaqoob & Padmanabhan 2004 and Shu et al. 2010).

On the other hand, we note that there is no bias in terms of physical source properties or classification. The sample includes optical Type 1, 1.5 and 2 AGN and also some Narrow Line Seyfert 1s, all below $z=0.3$. However, the sample excludes obscured sources with line-of-sight gas column densities $N_{\text{H}} > 10^{23} \text{ cm}^{-2}$ in order to avoid spectral modeling complexities. While a signal-to-noise requirement is imposed so that the total counts in the full HEG bandpass must be more than 1500, Shu et al. (2010) state that relaxing this criterion would have admitted only 2 more objects, so this does not appear to be a major restriction unless Fe $K\alpha$ properties evolve very rapidly with redshift (i.e. to fainter source fluxes) – this is an unlikely prospect in our opinion.

Of the 36 sources studied by Shu et al. (2010), the statistical quality of the data was good enough for FWHM measurements in 27 objects. In the remaining 9 objects, the signal:noise of the data around the Fe line region is so poor that the line itself is not detected at more than 95% confidence and the Gaussian model fit can become unphysically broad as it begins to model the continuum itself (see Section 3 of Shu et al. 2010). This is the case for NGC 526a, Mrk 705, NGC 3227, Mrk 766, PDS 456, IRAS 18325–5926, NGC 7314, Ark 564 and MR 2251–178. Examining the reason for the poor data quality of these 9 sources, we note that their mean HEG exposure time is only 85 ks, as opposed to 192 ks for the 27 objects with FWHM measurements. In addition, the two sources with the longest exposure times (PDS 456 and MR 2251–178 with exposures of 145.2 ks and 148.7 ks, respectively) are both known to be relatively distant, obscured quasars showing highly complex continua and evidence of significant absorption lines resulting from strong outflows (e.g. Gibson et al. 2005; Hagino et al. 2015). In other words, the absence of a significant Fe $K\alpha$ emission line in these 9 sources is most likely a combination of relatively short exposure times resulting in low signal:noise, and complex spectra in which a uniform analysis does not immediately resolve a narrow line core.

The main restriction imposed by the infrared selection is one of target visibility – most (though not all) of the dust reverberation lags have been carried out from the MAGNUM telescope in Hawaii. There may be an implicit preference for observing bright sources and those sources with known (significant) levels of optical variability, but it is not clear how this would bias our sample. Interferometric size measurements have been carried out mostly from Keck also in Hawaii, but the sample of AGN where $R_{\text{dust,inf}}$ has been measured is currently limited. So there is no obvious physical bias in the source selection, albeit our sample of only 13 AGN from the parent sample remains small.

4.4.5. Alternate black hole mass estimates

Finally, one may question the robustness and suitability of our adopted M_{BH} measurements. As discussed in Section 2.1, certain assumptions about the geometric correction (f -factor) are implicit in virial M_{BH} estimates based upon reverberation mapping, and we have assumed that similar correction factors apply for the Fe $K\alpha$ emission zone.

The influence of making such an assumption can be judged by using alternate black hole mass measurements, which are entirely independent of reverberation mapping. In the Appendix, we have compiled such measurements in Table 2. We

present a figure showing the resultant comparison of R_{Fe} with R_{dust} in this case, and demonstrate that our main inference of R_{dust} forming an envelope to R_{Fe} remains unchanged. We refer the reader to the Appendix for this test.

To summarize, we have discussed and quantified some caveats to our work, including those from combining multiple datasets, limited instrumental resolution and biases related to simplistic (single Gaussian) modeling of a likely complex line feature. None of these issues appears to cause a bias strong enough to overwhelm our main result of R_{dust} acting as an outer envelope to R_{Fe} . We point out current shortfalls in terms of sample completeness, although the apparent heterogeneous nature of sample selection argues against any strong selection bias. We also explore the use of alternate black hole masses and their influence on our main result, which do not change our main result.

Although we cannot rule out the presence of a *combination* of these various biases, it would be surprising if they were to conspire to produce a limit of R_{Fe} around R_{dust} – which is a completely independent quantity. However, we acknowledge the current limited sample size and the paucity of line variability measurements. More complete sampling of AGN in both X-rays and in the infrared are clearly important for drawing robust conclusions, and we discuss future possibilities in the next section.

4.5. Future perspectives

This field is expected to leap forward with the imminent launch of the *Astro-H* mission (Takahashi et al. 2014). With a spectral resolution of $\approx 4\text{--}7 \text{ eV}$, the Soft X-ray Spectrometer (SXS) employing calorimetric photon detection will provide an improvement in spectral resolution of $\approx 6\text{--}10$ as compared to present grating instruments. Since R_{Fe} has an inverse quadratic scaling with line velocity, SXS will extend the radius out to which the $K\alpha$ emission can be localized by a factor of $\gtrsim 40$ in any individual object. This assumes a single Gaussian feature. Individual ionization stages will not all be necessarily separable by the SXS, though some deconvolution will be possible. In terms of the overall sensitivity to line parameter measurement, this can be quantified in terms of a ‘figure of merit’, which combines the effective area and the spectral resolution. Relative to HEG, the overall SXS sensitivity to line detection is expected to be improved by a factor of ≈ 7 , extending our reach to correspondingly faint (and more distant) systems. However, the improvement in the corresponding figure of merit for detection of line *broadening* (which is key to our work) is about two orders of magnitude relative to the HEG.⁷

In the next decade, the *Athena* mission (Nandra et al. 2013) will have a much larger collecting area than *Astro-H*, pushing such studies out to high redshift. Reverberation of the narrow Fe $K\alpha$ core for more objects would also be an independent constraint on R_{Fe} . Variable narrow lines in some cases have already been mentioned. In addition, Liu et al. (2010) found Fe line reverberation on a timescale of $\sim 20\text{--}40$ day in NGC 5548. This would push R_{Fe} a factor of $\sim 1.5\text{--}3$ lower than our estimate in Table 1, but still consistent within the uncertainty on R_{Fe} .

And in a similar vein, enlarged NIR coverage will greatly help to better fill the R_{Fe} vs. R_{dust} parameter space. There are currently a total of 31 AGN with measurements of R_{dust} .

⁷ <http://astro-h.isas.jaxa.jp/ahqr.pdf>

This includes sources from the references cited in Table 1 as well as a few other published and unpublished sources (e.g. GQ Com; Sitko et al. 1993). We have included 13 of these (42%) in our X-ray cross-matched sample. Although Keck Interferometer is no longer available, more size measurements will be available soon (Kishimoto et al. 2015 submitted) making use of the three-telescope beam combiner AMBER at the VLTI. This will boost the present sample and address projection effects by invoking closure phase data. And although the MAGNUM telescope is now decommissioned, there are many other world-wide efforts ongoing to obtain AGN NIR time lags (e.g. Pozo Nuñez et al. 2015). LSST could also serve as a dust reverberation machine providing thousands of time lag measurements (Hönig 2014). Finally, in the distant future, we can expect to directly resolve the torus in many more sources with larger interferometers such as the Planet Formation Imager (Monnier et al. 2014).

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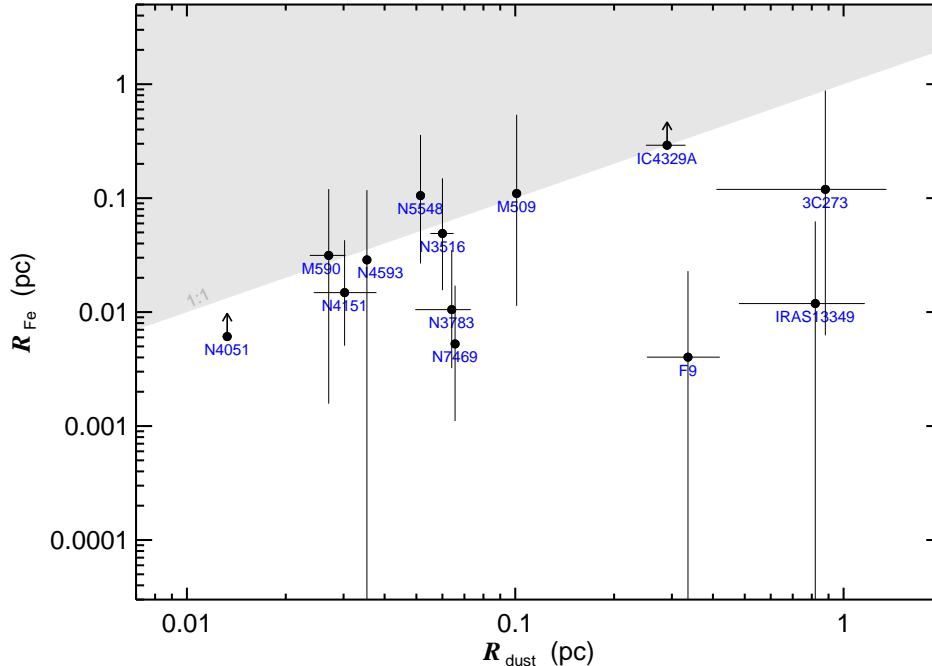


FIG. 4.— The comparison of R_{Fe} with $R_{\text{dust,rev}}$ as in Fig. 1, but this time with the computation of R_{Fe} based upon black hole masses obtained from reverberation-independent methods.

APPENDIX

BLACK HOLE MASSES INDEPENDENT OF REVERBERATION MAPPING

Our adopted values of M_{BH} in Table 1 are mostly virial estimates based upon reverberation mapping of the BLR, which implicitly assume a certain correction for geometric projection effects. The size comparisons in Fig. 1 effectively presuppose that the same projection effects apply for the Fe $K\alpha$ emission regions and the BLR, which may, or may not, be the case. Although the correction factors are expected to be small (see Section 2.1), it is worth investigating how strongly the above assumption affects our results. In order to assess this, we compile black hole masses which are independent of reverberation mapping. Any distinct projection effect for the Fe $K\alpha$ and $H\beta$ emission regions will then be retained in the emitter size comparisons, and the resultant change in the R_{Fe} distribution (relative to Fig. 1) will give an estimate of the maximal impact that the unknown projection effects can have.

Table 2 presents our compilation of alternate M_{BH} values.⁸ These are mostly based upon host galaxy stellar velocity dispersion (σ_*) measurements and the relation between M_{BH} and σ_* from Gültekin et al. (2009). Where this was not possible, we used scaling relations between X-ray variability amplitude (Zhou et al. 2010; McHardy 2013) and M_{BH} , or the empirical relation between optical continuum luminosity and the BLR radius (Vestergaard 2002) combined with the FWHM of the $H\beta$ emission line. There is broad agreement between the alternate masses tabulated here and those listed in Table 1, with the alternate masses (mostly based upon $M_{\text{BH}}-\sigma_*$) being larger than the virial estimates by a median factor of 1.8 but with a significant spread of 0.7 dex (standard deviation) between the two sets of masses.

The resultant values of R_{Fe} are also listed in the table, and are plotted in Fig. 4 which shows that although the location of the Fe $K\alpha$ emission radii relative to R_{dust} (and thus also relative to R_{BLR}) do change in some individual objects as compared to Fig. 1 (e.g. $R_{\text{Fe}}(\text{NGC } 3783)$ is now smaller than R_{dust} , and $R_{\text{Fe}}(\text{Mrk } 590)$ is now consistent with R_{dust}), our main inference of R_{dust} serving as an approximate envelope to R_{Fe} still holds. IC 4329A is the only source in which the limit to R_{Fe} apparently sits very close to (but still consistent with) R_{dust} . As we discussed in Section 4.1, Fe $K\alpha$ emission will not cease exactly at R_{dust} , and R_{Fe} values of up to a few times R_{dust} are allowed by Fig. 1 and also by typical compact torus models. *Astro-H* observations of this source will help to pinpoint its location on the $R_{\text{Fe}}-R_{\text{dust}}$ plane.

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⁸ Although our primary adopted value of $M_{\text{BH}}(\text{NGC } 4151)$ from Hönig et al. (2014) is not based upon reverberation mapping, we use an alternate M_{BH} value based upon the $M_{\text{BH}}-\sigma_*$ method here as a cross-check.

TABLE 1
SAMPLE PROPERTIES

Source	M_{BH} $\times 10^7 M_{\odot}$	Ref.(M_{BH})	τ_{dust} days	$R_{\text{dust,rev}}$ $\times 0.1 \text{ pc}$	$R_{\text{dust,intf}}$ $\times 0.1 \text{ pc}$	Ref.(R_{dust})	R_{BLR} $\times 0.01 \text{ pc}$	L_{Bol} $10^{46} \text{ erg s}^{-1}$	$v_{\text{FWHM}}^{\text{Fe}}$ km s^{-1}	R_{Fe} $\times 0.1 \text{ pc}$
Fairall9	255.0 \pm 56.0	i	400.0 \pm 100.0	3.21 \pm 0.80	—	1	1.40 $^{+0.26}_{-0.35}$	1.90	18100 $^{+76840}_{-12390}$	0.045 $^{+0.135}_{-0.045}$
Mrk590	15.8 \pm 3.4	ii	32.2 \pm 4.0	0.26 \pm 0.03	—	2	2.09 $^{+0.16}_{-0.19}$	0.17	4350 $^{+6060}_{-2830}$	0.048 $^{+0.077}_{-0.042}$
NGC3516	42.7 \pm 14.6	i	71.5 \pm 5.8	0.59 \pm 0.05	—	2	0.56 $^{+0.17}_{-0.32}$	0.025	3180 $^{+670}_{-1113}$	0.242 $^{+0.142}_{-0.113}$
NGC3783	29.8 \pm 5.4	i	76.3 $^{+10.9}_{-17.2}$	0.64 $^{+0.09}_{-0.14}$	1.60 \pm 0.50	3, 4	0.84 $^{+0.27}_{-0.19}$	0.17	1750 $^{+360}_{-360}$	0.558 $^{+0.314}_{-0.203}$
NGC4051	1.9 \pm 0.8	i	15.8 \pm 0.5	0.13 \pm 0.004	0.32 \pm 0.05	2, 5	0.49 $^{+0.22}_{-0.08}$	0.0059	< 1860	> 0.032
NGC4151	48.5 \pm 20.0	iii	36.0 $^{+9.0}_{-7.0}$	0.30 $^{+0.08}_{-0.06}$	0.41 \pm 0.04	6, 5	0.55 $^{+0.08}_{-0.07}$	0.050	2250 $^{+400}_{-360}$	0.549 $^{+0.377}_{-0.228}$
3C273	6590.0 \pm 1300.0	iv	—	—	8.80 \pm 4.70	7	22.23 $^{+4.96}_{-6.59}$	58.80	5900 $^{+8640}_{-5830}$	10.857 $^{+67.554}_{-10.240}$
NGC4593	5.4 \pm 9.4	i	42.1 \pm 0.9	0.35 \pm 0.01	—	2	0.31 $^{+0.07}_{-0.07}$	0.042	2230 $^{+5180}_{-6020}$	0.062 $^{+0.346}_{-0.217}$
IRAS13349+2438	1000.0 \pm 1000.0	v	—	—	8.20 \pm 3.40	8	—	11.16	5150 $^{+2810}_{-2810}$	2.162 $^{+2.162}_{-0.162}$
IC4329A	9.9 \pm 17.9	i	—	—	2.90 \pm 0.40	8	0.12 $^{+0.22}_{-0.15}$	0.79	< 1860	> 0.164
NGC5548	67.1 \pm 2.6	i	61.3 \pm 0.3	0.51 \pm 0.002	—	2	1.49 $^{+0.05}_{-0.05}$	0.07	2540 $^{+1140}_{-820}$	0.596 $^{+0.545}_{-0.358}$
Mrk509	143.0 \pm 12.0	i	120.3 \pm 1.1	0.98 \pm 0.01	3.00 \pm 0.50	2, 9	6.46 $^{+0.50}_{-0.44}$	1.04	2910 $^{+2590}_{-1250}$	0.968 $^{+1.337}_{-0.801}$
NGC7469	12.2 \pm 1.4	i	78.1 \pm 0.1	0.65 \pm 0.0004	< 1.90	2, 8	0.37 $^{+0.06}_{-0.07}$	0.19	4890 $^{+2770}_{-1700}$	0.029 $^{+0.029}_{-0.029}$

Notes: Uncertainties are quoted for 68% confidence. M_{BH} ref.: i: Peterson et al. (2004); ii: Kaspi et al. (2000); iii: Hönig et al. (2014); iv: Paltani & Türler (2005); v: Lee et al. (2013). R_{dust} ref.: 1: Clavel et al. (1989); 2: Koshida et al. (2014); 3: Lira et al. (2011); 4: Weigelt et al. (2009); 5: Kishimoto et al. (2012); 6: Hönig et al. (2009); 7: Kishimoto et al. (2011a); 8: Kishimoto et al. (2015, submitted); 9: Kishimoto et al. (2013).

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TABLE 2
ALTERNATE BLACK HOLE MASS ESTIMATES

Source	Method(M_{BH})	Measurement	Ref.	$\log M_{\text{BH}}$ M_{\odot}	$\Delta \log M_{\text{BH}}$ M_{\odot}	R_{Fe} $\times 0.1 \text{ pc}$
(1)	(2)	(3)	(4)	(5)	(6)	(7)
Fairall 9	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 228 \pm 18 \text{ km s}^{-1}$	Oliva et al. (1999)	8.36 ± 0.46	$+0.05 \pm 0.47$	$0.040^{+0.186}_{-0.040}$
Mrk 590	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 189 \pm 6 \text{ km s}^{-1}$	Nelson et al. (2004)	8.02 ± 0.44	-0.82 ± 0.45	$0.314^{+0.890}_{-0.298}$
NGC 3516	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 181 \pm 5 \text{ km s}^{-1}$	Nelson et al. (2004)	7.94 ± 0.44	-0.31 ± 0.47	$0.490^{+1.016}_{-0.335}$
NGC 3783	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 95 \pm 10 \text{ km s}^{-1}$	Onken et al. (2004)	6.75 ± 0.48	$+0.73 \pm 0.49$	$0.105^{+0.241}_{-0.073}$
NGC 4051	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 86 \pm 3 \text{ km s}^{-1}$	Nelson et al. (2004)	6.57 ± 0.44	-0.28 ± 0.47	> 0.061
NGC 4151	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 116 \pm 3 \text{ km s}^{-1}$	Onken et al. (2014)	7.12 ± 0.44	$+0.57 \pm 0.48$	$0.148^{+0.288}_{-0.098}$
3C 273	X-ray variance	$\log(\text{NVA}) = 0.20 \pm 0.08$	McHardy (2013)	8.86 ± 0.08	$+0.96 \pm 0.12$	$1.193^{+7.775}_{-1.129}$
NGC 4593	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 135 \pm 6 \text{ km s}^{-1}$	Nelson et al. (2004)	7.40 ± 0.45	-0.67 ± 0.88	$0.287^{+0.903}_{-0.287}$
IRAS 13349+2438	$L_{5100}-R_{\text{BLR}}$	$L_{5100} = 10^{44.64} \text{ erg s}^{-1}$	Wang et al. (2004); Grupe et al. (2004)	7.74 ± 0.50	$+1.26 \pm 0.66$	$0.119^{+0.474}_{-0.119}$
IC 4329A	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 225 \pm 9 \text{ km s}^{-1}$	Oliva et al. (1999)	8.24 ± 0.44	-1.25 ± 0.90	> 2.911
NGC 5548	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 195 \pm 13 \text{ km s}^{-1}$	Woo et al. (2010)	8.07 ± 0.46	-0.25 ± 0.46	$1.053^{+2.659}_{-0.788}$
Mrk 509	X-ray variance	$\log(\sigma_{\text{rms}}^2) = -3.24 \pm 0.59$	Zhou et al. (2010)	8.21 ± 0.62	-0.05 ± 0.62	$1.098^{+4.413}_{-0.993}$
NGC 7469	$M_{\text{BH}}-\sigma_*$	$\sigma_* = 131 \pm 5 \text{ km s}^{-1}$	Nelson et al. (2004)	7.34 ± 0.45	-0.25 ± 0.45	$0.053^{+0.122}_{-0.042}$

Column (2) denotes the method used to determine M_{BH} . Three methods are used: the $M_{\text{BH}}-\sigma_*$ relation from Gültekin et al. (2009); the X-ray variance scaling with M_{BH} either from McHardy (2013, who characterizes this in terms of the Normalized Variability Amplitude or NVA) or from Zhou et al. (2010, who provide measurements of the rms variability σ_{rms}^2); the empirical relation of L_{5100} with R_{BLR} from Vestergaard (2002). Column (3) lists the relevant observable used in the method. Column (4) gives the reference for the measurement of the observable. Column (5) lists the resultant M_{BH} value, and Column (6) the difference in $\log M_{\text{BH}}$ with respect to the value quoted and used in the main paper (Table 1). Finally, Column (7) lists the corresponding value of R_{Fe} using the alternate M_{BH} estimate. For IRAS 13349+2438, we assume a large error of 0.5 dex in M_{BH} as in De Marco et al. (2013). Note that for this source $v_{\text{FWHM}}^{\text{H}\beta} = 2800 \pm 180 \text{ km s}^{-1}$ is reported in Grupe et al. (2004).

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