

THE FOURSTAR GALAXY EVOLUTION SURVEY (ZFOURGE): ULTRAVIOLET TO FAR-INFRARED CATALOGS, MEDIUM-BANDWIDTH PHOTOMETRIC REDSHIFTS WITH IMPROVED ACCURACY, STELLAR MASSES, AND CONFIRMATION OF QUIESCENT GALAXIES TO $z\sim3.5^*$

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ABSTRACT

The FourStar galaxy evolution survey (ZFOURGE) is a 45 night legacy program with the FourStar nearinfrared camera on Magellan and one of the most sensitive surveys to date. ZFOURGE covers a total of 400 arcmin² in cosmic fields CDFS, COSMOS and UDS, overlapping CANDELS. We present photometric catalogs comprising >70,000 galaxies, selected from ultradeep K_s -band detection images (25.5–26.5 AB mag, 5σ , total), and >80% complete to $K_s < 25.3-25.9$ AB. We use 5 near-IR medium-bandwidth filters $(J_1, J_2, J_3, H_s, H_l)$ as well as broad-band K_s at 1.05–2.16 μ m to 25–26 AB at a seeing of ~ 0 ."5. Each field has ancillary imaging in 26–40 filters at 0.3-8 µm. We derive photometric redshifts and stellar population properties. Comparing with spectroscopic redshifts indicates a photometric redshift uncertainty $\sigma_z = 0.010, 0.009,$ and 0.011 in CDFS, COSMOS, and UDS. As spectroscopic samples are often biased toward bright and blue sources, we also inspect the photometric redshift differences between close pairs of galaxies, finding $\sigma_{z,pairs} = 0.01-0.02$ at 1 < z < 2.5. We quantify how $\sigma_{z,pairs}$ depends on redshift, magnitude, spectral energy distribution type, and the inclusion of FourStar medium bands. $\sigma_{z,pairs}$ is smallest for bright, blue star-forming samples, while red star-forming galaxies have the worst $\sigma_{z, pairs}$. Including FourStar medium bands reduces $\sigma_{z, pairs}$ by 50% at 1.5 < z < 2.5. We calculate star formation rates (SFRs) based on ultraviolet and ultradeep far-IR Spitzer/MIPS and Herschel/PACS data. We derive rest-frame U-V and V-J colors, and illustrate how these correlate with specific SFR and dust emission to z = 3.5. We confirm the existence of quiescent galaxies at $z \sim 3$, demonstrating their SFRs are suppressed by $> \times 15$.

Key words: cosmology: observations - galaxies: evolution - galaxies: high-redshift - infrared: galaxies

1. INTRODUCTION

Over the last few decades it has been possible to obtain new insights into the formation and evolution of galaxies in a statistically significant way by using large samples of sources from multiwavelength photometric surveys, for example with SDSS (York et al. 2000). Improved near-IR facilities on the ground, as well as advanced space-based instruments have enabled galaxy surveys probing the universe at higher resolution, fainter magnitudes and toward higher redshifts (z > 1.5) (e.g., Lawrence et al. 2007; Wuyts et al. 2008; Grogin et al. 2011; Koekemoer et al. 2011; Whitaker et al. 2011; Muzzin et al. 2013a; Skelton et al. 2014). These in turn have led to great progress in tracing the structural evolution of galaxies (e.g., Daddi et al. 2005; Franx et al. 2008; van Dokkum et al. 2008;

Bell et al. 2012; van der Wel et al. 2012, 2014; Wuyts et al. 2012), luminosity and stellar mass functions (e.g., Faber et al. 2007; Pérez-González et al. 2008; Marchesini et al. 2009; Muzzin et al. 2013b; Tomczak et al. 2014), the environmental effects on galaxy evolution (e.g., Postman et al. 2005; Papovich et al. 2010; Peng et al. 2010b; Cooper et al. 2012; Quadri et al. 2012; Kawinwanichakij et al. 2014; Allen et al. 2015) and the correlation between stellar mass and star formation rate (e.g., Noeske et al. 2007; Wuyts et al. 2011; Whitaker et al. 2012) over cosmic time.

The redshift range 1 < z < 3, when the universe was between 2.1 and 5.6 Gyr old, is an important epoch for studies of galaxy evolution. During this period 60% of all star formation took place (e.g., Madau et al. 1998; Sobral et al. 2013), an early population of quiescent galaxies started to appear (e.g., Daddi et al. 2005; Kriek et al. 2006; Marchesini et al. 2010) and galaxies evolved into the familiar elliptical and spiral morphologies that we see in the universe today (e.g., Bell

^{*} This paper contains data gathered with the 6.5 meter Magellan Telescopes located at Las Campanas observatory, Chile

Hubble Fellow.

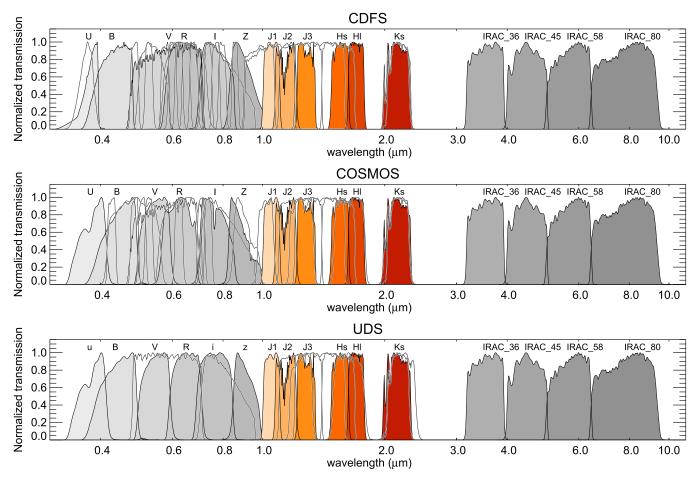


Figure 1. Normalized transmission corresponding to the FourStar medium-bandwidth and ancillary filters, each panel representing a different field. From top to bottom: CDFS, COSMOS and UDS. We show the FourStar J_1 , J_2 , J_3 , H_s , H_l and K_s medium-bandwidth filters with different shades of red. The UV to optical U, B, V, R, I and Z filters and the Spitzer/IRAC filters are shown with gray shaded curves. These correspond to different instruments in each field. The FourStar filters overlap with other broadband near-IR filters, e.g., HST/WFC3/F125W-F160W, while providing a higher resolution sampling. Atmospheric transmission was included in all FourStar filter curves. All filters are mentioned separately in Tables 2 (CDFS), 3 (COSMOS) and 4 (UDS).

Table 1
FourStar Observations

Cosmic field	Filter	Total Integration Time	5σ Depth
		(hr)	(AB mag)
CDFS	J_1	6.3	25.6
CDFS	J_2	6.5	25.5
CDFS	J_3	8.8	25.5
CDFS	H_{s}	12.2	24.9
CDFS	H_l	5.9	25.0
CDFS	K_s	5.0	24.8
COSMOS	J_1	13.9	26.0
COSMOS	J_2	16.0	26.0
COSMOS	J_3	13.8	25.7
COSMOS	H_{s}	12.1	25.1
COSMOS	H_l	12.1	24.9
COSMOS	K_s	13.4	25.3
UDS	J_1	7.9	25.6
UDS	J_2	8.7	25.9
UDS	J_3	9.3	25.6
UDS	H_{s}	11.0	25.1
UDS	H_l	10.4	25.2
UDS	K_s	3.9	24.7

et al. 2012). A fundamental observational limitation to understanding galaxy evolution is the availability of accurate distance estimates for mass-limited galaxy samples. These can be obtained with spectroscopy, but observations are limited to a biased population of galaxies: bright and most often star-forming, with strong emission lines.

Instead many galaxy surveys rely exclusively on the photometric sampling of the spectral energy distributions (SEDs) of galaxies to derive redshifts. Even when deep imaging spanning the optical and near-infrared is used to derive photometric redshifts, these surveys are generally hampered by systematic effects from the use of broadband filters. These can lead to large random errors, of the order of $\sigma_z/(1+z) \sim 0.1$. Moreover the photometric redshift accuracy is generally estimated by comparison to a small and unrepresentative spectroscopic sample, which does not allow for an analysis of the errors as a function of magnitude, redshift, or galaxy type. Redshift errors may introduce biases in derived luminosities and stellar masses (Chen et al. 2003; Kriek et al. 2008).

A better sampling of the SED improves the accuracy of the photometric redshifts greatly and can be obtained by the use of medium-bandwidth filters. These were first applied in the optical for the COMBO17 survey (Wolf et al. 2004). A notable

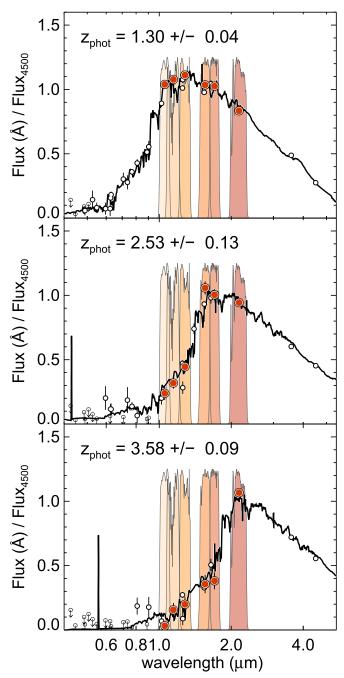


Figure 2. The FourStar filters provide detailed sampling of the Balmer/ 4000 Å break of galaxies at $z \gtrsim 1.5$. Here we show the SEDs of three observed galaxies in COSMOS with large Balmer/4000 Å breaks, at z=1.30, z=2.53 and z=3.58. With increasing redshift, the Balmer/4000 Å break moves through the range defined by the FourStar bands. Observed datapoints are shown as white or red dots with errorbars for ancillary and FourStar filters, respectively. Upper limits (mostly in the UV) are indicated with downward arrows. The solid curves are the EAZY best-fit SEDs (see Section 5). Observed and fitted SEDs are normalized at rest-frame 4500 Å.

feature in the SED of a galaxy is the Balmer/4000 Å break at rest-frame 4000 Å, which shifts into the near-IR at $z\gtrsim 1.5$. For high redshift surveys, it is therefore advantageous to split up the canonical broadband J and H filters into multiple near-IR medium-bandwith filters (van Dokkum et al. 2009), which stradle the Balmer/4000 Å break at $1.5\lesssim z\lesssim 3.5$. A set of near-IR medium-bandwidth filters was used for the NEWFIRM Medium-band Survey NMBS, a survey using NEWFIRM on

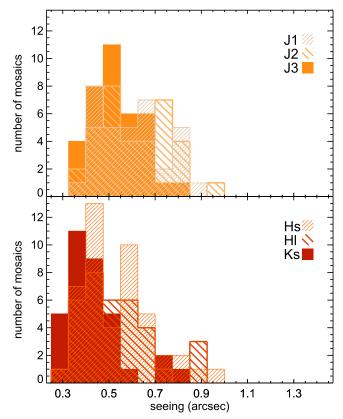


Figure 3. Seeing histograms of the FourStar single images, corresponding to $\sim 1-1.5$ hr observing blocks. Many of the images have a seeing of $\sim 0.14 - 0.15$.

the Kitt Peak Mayall 4 m Telescope, with a limiting 5σ depth in K of 23.5 AB mag for point sources and a photometric redshift accuracy of $\sigma_z/(1+z) \sim 1\%-2\%$ up to z=3 (Whitaker et al. 2011).

The FourStar Galaxy Evolution Survey (ZFOURGE) aims to advance further the study of intermediate to high redshift galaxies by pushing to much fainter limits (25-26 AB), well beyond the typical limits of groundbased spectroscopy. This provides a unique opportunity to study the higher redshift and lower mass galaxy population in unprecedented detail, at cutting edge mass completeness limits. The power of this deep survey is demonstrated by Tomczak et al. (2014), who showed the stellar mass functions of star-forming and quiescent galaxies can be accurately traced down to $10^9\,M_\odot$ at z=2, well below M^* . Papovich et al. (2015) showed that at this depth one can trace the evolution of progenitors of present-day M* galaxies (like M31 and the Milky Way Galaxy) out to $z \sim 3$. Furthermore Straatman et al. (2014) showed that a population of massive quiescent galaxies with $M > 10^{10.6}$ was already in place at $z \sim 4$, while Tilvi et al. (2013) used the FourStar medium-bandwidth filters to pinpoint Lyman Break galaxies at $z \sim 7$ and distinguish them from cool dwarf stars.

In this paper we present the ZFOURGE data products, 12 comprising 45 nights of observations with the FourStar near-infrared Camera on the 6.5 m Magellan Baade Telescope at Las Campanas in Chile (Persson et al. 2013). The survey was conducted over three extragalactic fields: CDFS (R.A. (J2000) = 03:32:30, decl. (J2000) = -27:48:30) (Giacconi et al. 2002), COSMOS (R.A. = 10:00:30, decl. = +02:17:30)

¹² Available for download at zfourge.tamu.edu.

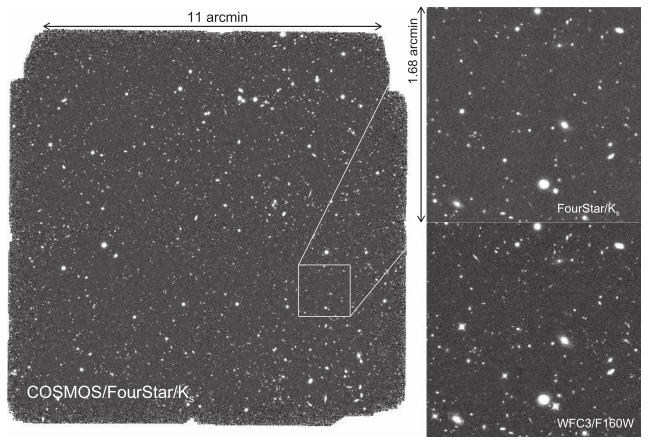


Figure 4. Left: the FourStar/ K_s -band reduced image in COSMOS. The FourStar footprint is $13' \times 13'$. Top right: zooming in on a $1'.68 \times 1'.68$ region in the COSMOS field. Bottom right: the same region with HST/WFC3/F160W.

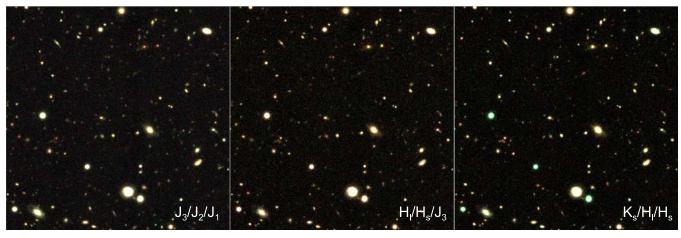


Figure 5. False color images of the same cutout region as shown in Figure 4, demonstrating the high quality obtained with the FourStar filters, as well as the usefulness of using medium-bandwidth filters to characterize the colors of galaxies within a classical *J* or *H* broad band. The filter combinations that were used in each panel are indicated at the bottom (red/green/blue).

(Scoville et al. 2007) and UDS (R.A. = 02:17:00, decl. = -05:13:00) (Lawrence et al. 2007), to reduce the effect of cosmic variance, and benefit from the large amount of public UV, optical and IR data already available. We present K_s -band selected near-IR catalogs, supplemented with public UV to IR data at $0.3-8~\mu m$, far-IR data from Spitzer/MIPS at $24~\mu m$ for all fields and Herschel/PACS at 100 and $160~\mu m$ for CDFS.

In Sections 2 and 3.1, we discuss the survey and image processing and optimization. In Section 3 we discuss source detection and photometry and include a description of the

ZFOURGE data products. In Section 4 we test the completeness limits of the survey. We derive photometric redshifts and rest-frame colors in Section 5 and stellar masses, stellar ages and star formation rates in Section 6. In Section 7 we show how to effectively distinguish quiescent from star-forming galaxies using a *UVJ* diagram, validating this classification with far-IR *Spitzer/* MIPS and *Herschel/*PACS data. A summary is provided in Section 8. Throughout, we assume a standard Λ CDM cosmology with $\Omega_{\rm M}=0.3$, $\Omega_{\Lambda}=0.7$ and $H_0=70~{\rm km~s^{-1}~Mpc^{-1}}$. The adopted photometric system is AB (Oke et al. 1995).

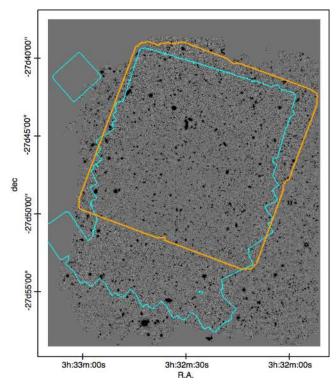


Figure 6. Deep K_s -band detection image in CDFS. The orange outline shows the ZFOURGE footprint. With cyan outlines we show the HST/WFC3/F160W footprint from CANDELS. North is up and east is to left.

2. DATA

2.1. ZFOURGE

The ZFOURGE (PI: I. Labbé) is a 45 night program with the FourStarinstrument (Persson et al. 2013) on the 6.5 m Magellan Baade Telescope at Las Campanas, Chile. FourStar has five near-IR medium bands: J_1 , J_2 , J_3 , H_s and H_l , covering the same range as the more classical J and H broadband filters, and a K_s -band. The central wavelengths of these filters range from 1.05 μ m (J_1) to 2.16 μ m (K_s).

The filter curves are shown in Figure 1; we have also added the filter curves of the ancillary dataset (see Section 2.4), showing that we cover the full UV to near-IR wavelength range. The FourStar filters overlap with broadband filters such as *Hubble Space Telescope (HST)*/WFC3/F125W, F140W and F160W in wavelength space, except they are narrower and sample the near-IR in more detail. The effective filter curves we use are modified to include the Lord et al. (1992) atmospheric transmission functions with a water column of 2.3 mm. The total integration time in each filter is shown in Table 1.

The sampling of the FourStar medium-bandwidth filters is illustrated in Figure 2, where we show the SEDs of observed galaxies in COSMOS with large Balmer/4000 Å breaks at $z \gtrsim 1.5$. The FourStar near-IR photometry is highlighted in red. The medium-band filters are shown in the background. They are particularly well suited to trace the Balmer/4000 Å break at higher redshifts, which is crucial to derive photometric redshifts.

2.2. FourStar Image Reduction

2.2.1. Pipeline

The FourStar data were reduced using a custom IDL pipeline written by one of the authors (I. Labbé) and also used

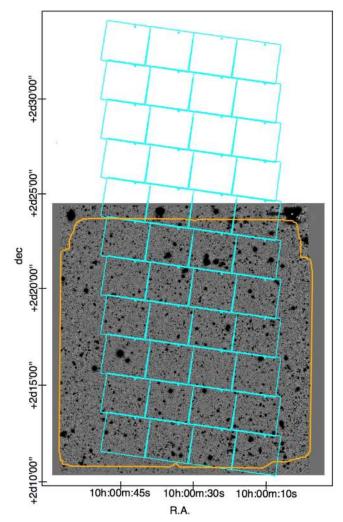


Figure 7. Deep K_s -band detection image in COSMOS and outlines as in Figure 6.

in the NMBS (Whitaker et al. 2011). It employs a two-pass sky subtraction scheme based on the IRAF package xdimsum.

The pipeline processes the data, which consist of dithered frames for each of the four FourStar detectors, separately for each $\sim 1-1.5$ hr observing block. Observed frames taken with each of the detectors were reduced and subsequently combined into a single mosaic.

Linearity corrections from the FourStar website¹³ were applied to the raw data. Dark current was determined to be variable so we did not remove any dark pattern. We also found constant bias levels along columns and rows in the raw data. We therefore subtracted the median of a column/row from itself.

Master flat field data were produced using twilight observations. For the K_s -band, where thermal contributions play a role, we attempted to mitigate the impact of illumination from the warm telescope. By combining multiple dithered observations of a blank field at the end of a night when the telescope had cooled, we were able to characterize the telescope illumination pattern. Shortly afterwards we took twilight flats and subtracted the telescope illumination pattern from each exposure. The flats with the telescope contribution removed were normalized and combined into the master K_s -band flats.

http://instrumentation.obs.carnegiescience.edu/FourStar/calibration.html

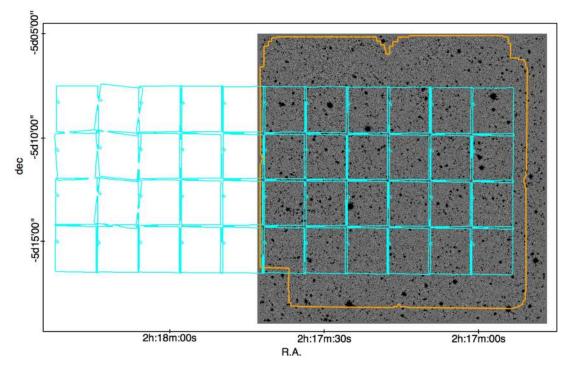


Figure 8. Deep K_s -band detection image in UDS and outlines as in Figure 6.

Sky background models were then subtracted from individual science exposures. The sky background was computed by averaging up to eight images taken before and after that exposure. Masking routines were run to remove: (1) bad pixels via a static mask from the FourStar website (2) satellite trails (3) guider cameras entering the field of view and (4) persistence from saturated objects in previous exposures. Bad pixels make up between 0.3% and 1.7% of the detectors (Persson et al. 2013). In addition, the individual exposures were visually screened for any remaining tracking issues, asteroids, airplanes and satellites.

Corrections for geometric distortion and absolute astrometric solutions were computed by crossmatching sources using astrometric reference images. In COSMOS we used the Canada–France–Hawaii Telescope (CFHT)/i-band as reference (Erben et al. 2009; Hildebrandt et al. 2009), in CDFS we used ESO/MPG/WFI/I from the ESI survey (Hildebrandt et al. 2006; Erben et al. 2009) and in UDS the UKIDDS data release eight K_s -band image (O. Almaini 2016, in preparation). The observations were interpolated onto a pixel grid with a resolution of 0."15 pix $^{-1}$, which is close to the native scale of FourStar of 0."159 pix $^{-1}$. The new grid shares the WCS tangent point (CRVAL) with the CANDELS HST images (Grogin et al. 2011; Koekemoer et al. 2011) and places CRVAL at a half-integer pixel position (CRPIX).

To optimize the signal-to-noise ratio (S/N) of the images for each observing block (and for the final mosaics), they were weighted by their seeing, sky background levels and ellipticity of the PSF before they are combined. The seeing conditions at Las Campanas were extraordinarily good, with a median seeing FWHM for the entire set of observations of 0.75 as shown in the histogram in Figure 3. Since the K_s -band cannot be observed with the HST, we paid special attention to this filter and only observed the K_s -band when the seeing was excellent. This resulted in a a very low median seeing for the FourStar/ K_s -filter of 0.74.

Finally, we subtracted a background in the final mosaics using Source Extractor (SE; Bertin & Arnouts 1996) to ensure any remaining structure in the background did not impact the aperture photometry. In short, SE iteratively estimated the median of the distribution of pixel values in areas of 48×48 pixels in CDFS and COSMOS and 96×96 pixels in UDS. The dimensions of these areas were chosen to avoid overestimating the background near bright sources. These estimates were smoothed on a scale of $3\times$ the background area, after which the background for the full images was calculated using a bicubic spline interpolation.

2.2.2. Photometric Calibration

Here we describe how we derived the near-IR photometric zeropoints of the final mosaics. Since these vary significantly with changes in local precipitable water vapor and airmass, we employed a differential photometric calibration scheme, using secondary standard stars. First, we selected a nearby standard star. We selected relatively faint ($K_s = 14.5-17 \text{ mag}$) spectrophotometric standard stars from the CALSPEC Calibration Database. 14 We then observed this primary standard star under photometric conditions immediately before or after a science observation in a particular filter. The science dataset was reduced and photometrically calibrated using the primary standard star observations and using an atmospheric watercolumn of 2.3 mm. Second, we then selected bright, unsaturated stars in each of the chips of the science field for use as secondary standard stars. All other science observations of an observing block were then calibrated to the primary standard star via the secondary standard stars within each of the science fields.

¹⁴ http://www.stsci.edu/hst/observatory/crds/calspec.html

Table 2
CDFS Passband Parameter

CDFS Passband Parameters Filter **FWHM** Galactic Zeropoint Offset (μm) (") (AB mag) Extinction \overline{B} 0.4318 0.73 -0.03222.097 -0.02922.151 0.019 -0.0140.7693 0.73 I R 0.6443 0.65 27.321 -0.148-0.0200.3749 25.932 -0.181-0.037U 0.81 V0.5919 0.73 22.968 -0.010-0.022 \mathbf{Z} 0.9036 0.73 21.378 0.041 -0.011-0.004 H_{s} 1.5544 0.60 26.618 -0.031 H_l 1.7020 0.50 26.588 -0.051-0.0041.0540 0.59 26.270 -0.041-0.009 J_1 J_2 1.1448 0.62 26.558 -0.043-0.0061.2802 0.56 -0.067 J_3 26.521 -0.0060.46 -0.083-0.003 K_s 2.1538 26.851 NB118 1.1909 0.47 24.668 0.000 -0.006NB209 2.0990 0.45 24.786 0.000 -0.00325.670 F098M 0.26 0.9867 0.011 -0.008F105W 1.0545 0.24 26.259 -0.002-0.007F125W 1.2471 0.26 26.229 0.004 -0.005F140W 1.3924 0.27 26.421 -0.027-0.004F160W 1.5396 0.27 25.942 -0.000-0.004F814W 0.8057 0.22 25.931 -0.004-0.011IA484 0.4847 0.81 25.463 -0.013-0.024IA527 0.5259 0.87 25.639 -0.059-0.0220.5763 25.543 IA574 1.01 -0.148-0.019IA598 0.6007 25.962 -0.040-0.0180.69 IA624 0.6231 0.67 25.887 0.014 -0.017IA651 0.6498 0.67 26.072 -0.062-0.016IA679 0.6782 0.86 26.105 -0.080-0.015IA738 0.7359 0.83 26.003 -0.003-0.013IA767 0.7680 0.77 26.000 -0.028-0.012IA797 0.7966 0.74 25.986 -0.022-0.012IA856 0.8565 0.74 25.713 -0.007-0.010 WFI_{V} 0.5376 0.96 23.999 -0.076-0.021WFI Rc -0.0160.6494 0.84 24.597 -0.038 WFI_U38 0.3686 0.98 21.587 -0.291-0.032tenisK 2.1574 0.86 24.130 0.233 -0.002KsHI 2.1748 0.45 31.419 0.022 -0.003IRAC_36 1.50 20.054 3.5569 -0.0160.000 IRAC_45 4.5020 1.50 20.075 0.005 0.000

Note. Zeropoints are the effective zeropoints. These have galactic extinction and zeropoint corrections derived in Section 5 incorporated, i.e., $zp = zp_I + offset + GE$, with zp_I representing the photometrically derived zeropoint of image I, offset the zeropoint correction and GE the galactic extinction value. The corrections (in units of AB magnitude) are indicated in separate columns.

20.626

21.803

0.023

0.022

0.000

0.000

1.90

2.00

IRAC 58

IRAC_80

5.7450

7.9158

In Section 5 we derive additional corrections to the zeropoints, that are typically of the order of 0.05 mag. We added these to the photometric zeropoints calculated here.

2.2.3. Image Depths

We measured the depths of the FourStar images by determining the root mean square (rms) of the background pixels. Since pixels may be correlated on small scales, e.g., due to confusion or systematics introduced during the reduction process, we used a method in which we randomly placed 5000 apertures of 0.6 diameter in each background subtracted

Table 3
COSMOS Passband Parameters

Filter	λ_c	FWHM	Zeropoint	Offset	Galactic
	(μm)	(")	(AB mag)		Extinction
\overline{B}	0.4448	0.61	31.129	-0.195	-0.076
G	0.4870	0.90	26.290	-0.015	-0.069
I	0.7676	0.77	25.759	0.091	-0.034
IA427	0.4260	0.79	31.119	-0.202	-0.079
IA484	0.4847	0.75	31.214	-0.116	-0.069
IA505	0.5061	0.82	31.252	-0.083	-0.065
IA527	0.5259	0.74	31.281	-0.058	-0.061
IA624	0.6231	0.72	31.348	-0.002	-0.050
IA709	0.7074	0.81	31.343	-0.015	-0.042
IA738	0.7359	0.80	31.347	-0.014	-0.039
R	0.6245	0.79	25.903	0.023	-0.047
U	0.3828	0.82	24.913	-0.235	-0.086
V	0.5470	0.80	31.418	0.077	-0.059
Rp	0.6276	0.83	31.453	0.100	-0.047
Z	0.8872	0.74	24.859	0.121	-0.030
Zp	0.9028	0.90	31.557	0.187	-0.030
H_l	1.7020	0.60	26.624	0.033	-0.010
H_{s}	1.5544	0.54	26.673	0.062	-0.012
J_1	1.0540	0.57	26.358	0.026	-0.020
J_2	1.1448	0.55	26.590	0.038	-0.018
J_3	1.2802	0.53	26.573	0.011	-0.016
K_s	2.1538	0.47	26.918	-0.011	-0.006
NB118	1.1909	0.58	24.637	0.000	-0.018
NB209	2.0990	0.52	24.849	0.000	-0.006
F125W	1.2471	0.26	26.236	-0.000	-0.011
F140W	1.3924	0.26	26.455	-0.000	-0.010
F160W	1.5396	0.26	25.948	-0.000	-0.008
F606W	0.5921	0.20	26.437	-0.016	-0.038
F814W	0.8057	0.21	25.951	0.032	-0.024
$UVISTA_J$	1.2527	0.82	30.052	0.062	-0.011
$UVISTA_H$	1.6433	0.81	29.995	0.003	-0.008
$UVISTA_Ks$	2.1503	0.79	30.028	0.035	-0.006
$UVISTA_Y$	1.0217	0.85	30.045	0.061	-0.016
IRAC_36	3.5569	1.70	21.530	-0.051	0.000
IRAC_45	4.5020	1.70	21.537	-0.044	0.000
IRAC_58	5.7450	1.90	21.577	-0.004	0.000
IRAC_80	7.9158	2.00	21.520	-0.061	0.000

image. Due to the dither pattern the images have less coverage from individual frames at the edges. We therefore considered only regions with coverage within 80% of the maximum exposure. Sources were also masked, based on the SE segmentation maps after object detection (see Section 3.2).

The resulting aperture flux distributions, representing the variation in the noise, were fit with a Gaussian, from which we derived the standard deviation (σ) . We then applied the point-spread-functions (PSFs) derived from bright stars (further explained in Section 3.1), to determine a flux correction for missing light outside of the aperture. σ was then multiplied by 5 and converted to magnitude using the effective zeropoint (the photometrically derived zeropoints as desribed above, with a correction applied) of each FourStarmosaic, to obtain an estimate of the 5σ limiting depth. The resulting depth in AB magnitude can thus be summarized as

$$depth(5\sigma) = zp - 2.5 \log_{10}[5\sigma apcorr], \tag{1}$$

with zp the zeropoint of the image and apcorr the aperture flux correction (typically factors of 1.7–2.6, depending on the seeing). The 5σ depths are summarized in Table 1 and have

Table 4
UDS Passband Parameters

Filter	$\lambda_c \ (\mu { m m})$	FWHM (")	Zeropoint (AB mag)	Offset	Galactic Extinction
и	0.3828	1.06	24.905	-0.268	-0.089
В	0.4408	0.91	24.803	-0.123	-0.074
V	0.5470	0.93	24.870	-0.072	-0.058
R	0.6508	0.96	24.914	-0.038	-0.049
i	0.7656	0.98	24.986	0.021	-0.035
Z	0.9060	0.99	24.974	0.001	-0.027
J_1	1.0540	0.55	26.121	-0.036	-0.022
J_2	1.1448	0.53	26.408	-0.029	-0.019
J_3	1.2802	0.51	26.481	-0.019	-0.015
H_s	1.5544	0.49	26.591	-0.000	-0.011
H_l	1.7020	0.51	26.448	-0.036	-0.010
K_s	2.1538	0.44	26.804	-0.067	-0.006
$J^{"}$	1.2502	0.91	30.863	-0.052	-0.015
Н	1.6360	0.89	31.262	-0.108	-0.010
K	2.2060	0.86	31.825	-0.059	-0.006
F125W	1.2471	0.26	26.214	-0.000	-0.016
F140W	1.3924	0.26	26.439	-0.000	-0.013
F160W	1.5396	0.26	25.935	-0.000	-0.011
F606W	0.5893	0.20	26.383	-0.054	-0.054
F814W	0.8057	0.23	25.926	0.015	-0.033
Y	1.0207	0.58	27.004	0.026	-0.022
KsHI	2.1748	0.46	27.520	0.026	-0.006
IRAC_36	3.5569	1.70	21.539	-0.042	0.000
IRAC_45	4.5020	1.70	21.556	-0.025	0.000
IRAC_58	5.7450	1.90	21.458	-0.123	0.000
IRAC_80	7.9158	2.00	21.522	-0.059	0.000

typical values of 25.5–26.0 AB mag in J_1 , J_2 , J_3 and 24.9–25.2 AB mag in H_s , H_l and 24.7–25.3 AB mag K_s .

In Figure 4 we show as an example the FourStar/ K_s -band image in COSMOS. We also compare with the near-IR CANDELS/HST/WFC3/F160W observations, with FWHM = 0.19 and a limiting 5σ depth of 26.4 AB mag. The deeper space-based F160W image has a higher resolution, but as a result of the very deep magnitude limits combined with excellent seeing conditions we can achieve almost a similar quality for our near-IR ground-based observations. The K_s -band images in CDFS and UDS have similar depth. To highlight the wealth of information provided by the fine spectral sampling of the FourStar medium-bandwidth filters we show again in Figure 5 the same cut-out region of Figure 4, using different filter combinations.

2.3. K_s-band Detection Images

We combine our FourStar/ K_s -band observations with deep pre-existing K-band imaging to create super-deep detection images. In CDFS we use Very Large Telescope (VLT)/HAWK-I/K from HUGS (with natural seeing between 0."3 and 0."5) (Fontana et al. 2014), VLT/ISAAC/K (v2.0) from GOODS, including ultra deep data in the HUDF region (seeing = 0."5 (Retzlaff et al. 2010), CFHST/WIRCAM/K from TENIS (seeing = 0."9) (Hsieh et al. 2012), and Magellan/PANIC/K in HUDF (seeing = 0."4) (PI: I. Labbé). In COSMOS we add VISTA/K from UltraVISTA (DR2) (seeing = 0."7) (McCracken et al. 2012) and in UDS we add imaging with UKIRT/WFCAM/K from UKIDSS (DR10) (seeing = 0."7) (O. Almaini et al 2016, in preparation) and also natural seeing VLT/HAWK-I/K imaging from HUGS.

Using sources common to the images a distortion map was determined. Subsequent bicubic spline interpolation was used to register the images to better than 0.03 across the field. We then determined the background rms flux variation ($\sigma_{\rm rms}$) and the seeing in each image, and we used these to assign a weight using weight = $1/(\sigma_{\rm rms} \times {\rm seeing}^2)$. Note that the images were not PSF-matched prior to combining. The final combined image stacks were obtained by a weighted average of the individual Kand K_s -band science images. Weight maps were obtained by averaging the individual exposure maps in the same way as the science images. The final K_s -band stacks have maximum limiting depths at 5σ significance of 25.5 and 25.7 AB mag in COSMOS and UDS, respectively, which are 0.2 and 1.0 mag deeper than the individual FourStar/ K_s -band observations. The depth in CDFS varies between 26.2 and 26.5, 1.4 to 1.7 mag deeper than the FourStar/ K_s -band image only. The average seeing in the three fields is FWHM = 0.45, 0.58 and 0.60. We use these images for source detection (Section 3), after calculating and subtracting the background. They are shown in Figures 6 to 8, with the ZFOURGE footprint indicated as well as the HST/WFC3/F160W footprint from CANDELS.

2.4. Ancillary Data: UV, Optical, NIR, and IR Imaging

In addition to the six FourStar filters, we incorporate imaging in 20–34 filters into each catalog, from publicly available surveys at $0.3-8~\mu m$. In CDFS we have a total of 40 bands, in COSMOS a total of 37 and in UDS a total of 26. These are summarized in Tables 2–4, where we additionally show, for every image, the central wavelength, PSF FWHM (see Section 3.1), effective zeropoint, galactic extinction value and zeropoint offset derived in Section 5. The galactic extinction values were calculated using the E(B-V) values from Schlafly & Finkbeiner (2011), interpolated between the given bandpasses and the central wavelengths of our filterset.

The CDFS UV-to-optical filters include VLT/VIMOS/*U*, *R*-imaging (Nonino et al. 2009), *HST*/Advanced Camera for Surveys (ACS)/*B*, *V*, *I*, *Z*-imaging (Giavalisco et al. 2004; Wuyts et al. 2008), ESO/MPG/WFI/*U*₃₈, *V*, *R*_c-imaging (Erben et al. 2005; Hildebrandt et al. 2006), *HST*/WFC3/*F098M*, *F105W*, *F125W*, *F140W*, *F160W* and *HST*/ACS*F606W*, *F814W*-imaging (Grogin et al. 2011; Koekemoer et al. 2011; Windhorst et al. 2011; Brammer et al. 2012), 11 Subaru/Suprime-Cam optical medium bands (Cardamone et al. 2010) with seeing <1."1 (from a set of 18, including seeing >1."1 images) and CFHT/WIRCAM/*K*-band imaging (Hsieh et al. 2012).

In COSMOS we added CFHT/*u*, *g*, *r*, *i*, *z*-imaging (Erben et al. 2009; Hildebrandt et al. 2009), Subaru/Suprime-Cam/*B*, *V*, *r*+, *z*+-imaging and 7 Subaru/Suprime-Cam optical medium-bandwidth filters (Taniguchi et al. 2007) with seeing <1."1 (from a set of 12, including seeing >1."1 images), *HST*/WFC3/F125W, F140W, F160W and *HST*/ACSF606W, *F814W*-imaging (Grogin et al. 2011; Koekemoer et al. 2011; Brammer et al. 2012) and UltraVISTA/*Y*, *J*, *H*, *K*_s-imaging (McCracken et al. 2012).

In UDS the additional filters are CFHT/MegaCam/*U* (O. Almaini 2016, in preparation; S. Foucaud 2016, in preparation), Subaru/Surpime-Cam/*B*, *V*, *R*, *i*, *z* (Furusawa et al. 2008), UKIRT/WFCAM/*J*, *H*, *K_s* (O. Almaini 2016, in preparation), *HST*/WFC3/*F125W*, *F140W*, *F160W* and *HST*/ACS*F606W*, *F814W* (Grogin et al. 2011; Koekemoer et al. 2011; Brammer et al. 2012) and VLT/HAWK-I/*Y* (Fontana et al. 2014).

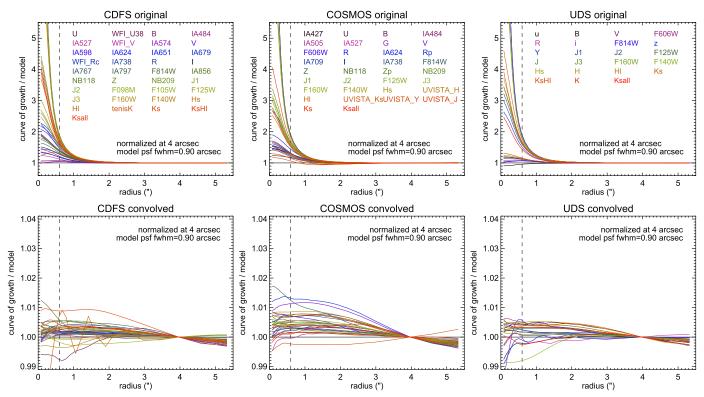


Figure 9. Top: curves of growth of the median stacked PSF of stars in the unconvolved images, normalized at 4'' radius and divided by the curve of growth of the target moffat PSF. The vertical dashed lines represent the radius at which we measure flux (Section 3.4). The spread in integrated flux is very large between different images, which would lead to biased color measurements. Bottom: here we show the curves of growth of the convolved images, where each PSF is convolved to match a Moffat profile. The correspondence with the target PSF is almost one-on-one, with at most a 1.5% deviation at $r = 0\rlap.{''}6$.

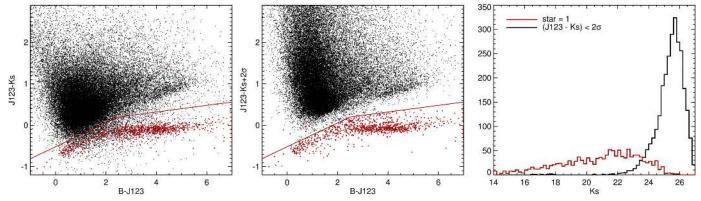


Figure 10. Left and middle panel: B - J123 vs. $J123 - K_s$ for stars (red) and galaxies (black). The red line indicates the star/galaxy separation. In the first panel some sources are scattered below the red line. These are effectively removed from the star classification by selecting in $J123 - K_s$ at $>2\sigma$ confidence (middle panel). Right panel: K_s -band magnitudes of stars (red) and sources that scatter below the red line (black). By only selecting at $>2\sigma$ confidence in $J123 - K_s$, we exclude the faintest sources from the star sample.

In CDFS and UDS we have additionally available FourStar narrow-bandwidth data at $1.18 \, \mu m$ (FourStar/NB18) and $2.09 \, \mu m$ (FourStar/NB209) (Lee et al. 2012). The narrowbands are sensitive to emission line flux. Small bandwidths in combination with high S/N for some galaxies may lead to biased photometric redshift and stellar mass estimates, because the models we use for determining redshifts and stellar population parameters do not contain well-calibrated strong emission lines. As such, they are incorporated into the catalogs, but are not used to derive photometric redshifts or stellar masses. The images have 5σ image depths of 25.2 and

24.8 AB mag in NB118 and CDFS and COSMOS, respectively and 24.4 and 24.0 AB mag in NB209.

The Spitzer/IRAC/3.6 and $4.5~\mu m$ images used in CDFS are the ultradeep mosaics from the IUDF (PI: Labbé), using data from the cycle seven IUDF program, IGOODS (PI: Oesch), GOODS (PI: Dickinson), ERS (PI: Fazio), S-CANDELS (PI: Fazio), SEDS (PI: Fazio) and UDF2 (PI: Bouwens). In CDFS we further use Spitzer/IRAC/5.8 and $8.0~\mu m$ images from GOODS (Dickinson et al. 2003). In COSMOS and UDS we use the 3.6 and $4.5~\mu m$ images from SEDS (Ashby et al. 2013). The 5.8 and $8.0~\mu m$ data in COSMOS are from S-COSMOS

 Table 5

 Explanation of The Photometric Catalog Header

id	ID number
x,y	pixel coordinates (scale: 0".15/pixel)
ra, dec	right ascension, declination (J2000)
SEflags	Source Extractor flags
iso_area	isophotal area above Source Extractor analysis treshold (pix2)
fap_Ksall ^a	(convolved) K_s -band aperture flux within a 1"2 diameter circular
	aperture
eap_Ksall	uncertainty on fap_Ksall
apcorr	aperture correction applied to fauto_Ksall to obtain f_Ksall (f_Ksall = fauto_Ksall * apcorr)
Ks_ratio	ratio between fap_Ksall and f_Ksall (Ks_ratio = f_Ksall/ fap_Ksall)
fapcirc0D_Ksall ^b	aperture flux measured within a D'' diameter (seeing dependent) ^b circular aperture
eapcirc0D_Ksall	uncertainty on fapcirc0D_Ksall
apcorr0D	aperture correction applied to fapcirc0D_Ksall to obtain fcirc0D_Ksall (fcirc0D_Ksall = fapcirc0D_Ksall * apcorr0D)
fcirc0D_Ksall ^{a,b}	total (aperture corrected) K_s -band flux within a D'' diameter (seeing dependent) ^b circular aperture
ecirc0D_Ksall	uncertainty on fcirc0D_Ks
fauto_Ksall	K_s -band flux within a Kron-like elliptical aperture
flux50_radius	radius (pixels) enclosing 50% of the K_s -band flux
a_vector	major axis of a Kron-like elliptical aperture
b_vector	minor axis of a Kron-like elliptical aperture
kron_radius	radius of a circularized Kron-like elliptical aperture
f_Ksall ^a	total (aperture corrected) K_s -band flux within a Kron-like elliptical aperture
e_Ksall	uncertainty on f_Ksall
w_Ksall	weight corresponding to f_Ksall, median normalized
f_[]	(convolved) aperture flux in filter [] within a 1."2diameter circular aperture, corrected to total (fap_[] = f_[]/Ks_ratio)
e_[]	uncertainty on f_[] (also scaled with Ks_ratio)
w_[]	weight corresponding to f_[], median normalized
wmin_optical	minimum w_[] of groundbased optical filters
wmin_hst_optical	minimum w_[] of HST optical filters
wmin_fs	minimum w_[] of FourStar filters
wmin_jhk	minimum w_[] of broadband J, H, & K filters
wmin_hst wmin_irac	minimum w_[] of <i>HST</i> near-IR filters minimum w_[] of <i>Spitzer/IRAC</i> filters
wmin_all	minimum w_[] of all filters
star	this flag is set to 1 if the source is likely to be a star, to 0 otherwise,
Sta	following the criteria described in Section 3.7
nearstar	this flag is set to 1 if the source is located within
	r('') < 10 - (m - 16) of a bright star with m the
use ^c	apparent magnitude of the star and $m < 17.5$ in J_1, J_2, J_3, J or K_s sources that pass the following criteria are set to 1:
	- star = 0 $- nearstar = 0$
	- snr ≥ 5
	- wmin_fs > 0.1 (A minimum exposure time of at least $0.1 \times$ the median exposure in the FourStar bands) ^d
	- wmin_optical > 0 (coverage in all optical bands)
	- not a catastrophic EAZY fit: χ^2 (reduced) $\leq 1000^{\text{e}}$
	- not a catastrophic FAST fit, i.e., a finite and positive stellar mass estimate above $10^6 \ M_{sun}$
	- consistent flux ratios between similar bands of different instru- ments, namely the <i>J-</i> , <i>H-</i> , and <i>K-</i> bands of FourStar and VISTA,
	and F814W- and groundbased <i>I</i> -bands - no 5σ detection at wavelengths bluer than the restframe 912 Å Lyman limit
	Lyman limit - not at $z < 0.1$
snr	signal-to-noise (=fapcirc0D_Ksall/eapcirc0D_Ksall)

Notes.

(Sanders et al. 2007) and in UDS from spUDS (J. Dunlop et al. 2016, in preparation).

The ancillary images are registered and interpolated to the same grid as the FourStar mosaics, using the program wregister in IRAF. Backgrounds for the UV, optical and near-IR images were estimated with SE and manually subtracted.

We further supplement the optical/near-IR catalogs with deep far-IR imaging from <code>Spitzer/MIPS</code> at 24 μm (GOODS-S: PI Dickinson, COSMOS: PI Scoville, UDS: PI Dunlop). Median 1σ flux uncertainties in 24 μm for the COSMOS and UDS pointings are roughly 10 μJy . The CDFS pointing is deeper with a median 1σ flux uncertainty of 3.9 μJy . In CDFS we additionally make use of public <code>Herschel/PACS</code> observations from PEP (Magnelli et al. 2013) at 100 μm and 160 μm , with 1σ flux uncertainties of 205 and 354 μJy respectively. In COSMOS and UDS deep <code>Herschel/PACS</code> data are not yet publicly released.

3. PHOTOMETRY

3.1. PSF Matching

The full UV/optical to near-IR dataset contains images of varying seeing quality. The FWHMs of the PSF corresponding to each image varies between 0.12 for the HST bands to 1.105 for some of the UV/optical images. To measure aperture fluxes consistently over the full wavelength range, i.e., measuring the same fraction of light per object in each filter, the images have to be convolved so that the PSFs match. To achieve a consistent PSF we first characterize the PSF in all individual images, we then define a theoretical model PSF as a reference, and finally convolve all bands to match the reference PSF.

The average PSF for each image was produced by selecting unsaturated stars with high S/N (>150; see Section 3.7, in which we describe how stars were identified in the images), in postage stamps of 10.65×10.65 . For each star we measured a curve of growth, i.e the total integrated light as a function of radius, with nearby objects masked using the SE segmentation map. Outliers, such as saturated stars, were then determined based on the shape of their light profile compared with the median curve of growth, and rejected from the sample. We median averaged the remaining stars, and, after normalizing the flux, used this to fill in masked regions. After renormalizing each tile by the total integrated flux at sufficiently large radius (25 pixels or 3."75) we again stacked the postage stamps to obtain a median star. Finally, to obtain a clean sample, we again compared the light profiles of individual stars against the median light profile, and iteratively rejected stars if the average deviation from the median curve of growth squared exceeded 5%. The result was a tightly homogeneous sample of stars, from which we obtained the final median two-dimensional PSF.

We generated as a reference PSF a model Moffat profile (Moffat 1969) with FWHM = 0.99 and β = 2.5. The advantage of using a model PSF rather than the average PSF from an image, is that a theoretical model is noiseless. To convolve the images to match the target model PSF, we first derive a kernel for each image individually. For this we use a deconvolution code developed by I. Labbé which fits a series of Gaussianweighted Hermite polynomials to the Fourier transform of the PSF. The original images were then convolved with this kernel to match the target PSF. This method results in very low residuals and is optimal for images with either a smaller PSF, or a PSF that is at most slightly larger (\lesssim 15%); further details

^a Note that these K_s -band fluxes are derived from the superdeep combined K_s -band images. Within the catalogs only f_Ks corresponds to FourStar/ K_s .

^b In CDFS D = 0.77, in COSMOS and UDS D = 0.99 (i.e., 1.5× the seeing FWHM).

 $^{^{\}rm c}$ A standard selection of galaxies can be obtained by selecting sources with use = 1.

d Effectively this means that every source has at least 20 minutes exposure in each FourStar band. Because of the dither pattern, sources with lower weight that are removed by this criterion lie at the edges of the images. For wmin_fs we used wmin_ksall instead of the weight of the FourStar Ks-band. Based on an empirical estimation from inspecting many fits.

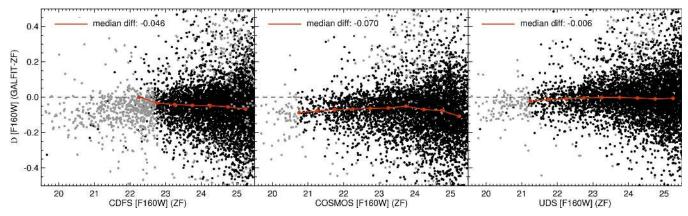


Figure 11. The difference between ZFOURGE and GALFIT F160W magnitudes plotted as a function of ZFOURGE magnitude, for galaxies with use = 1 and SEflag = 0 (excluding blended or contaminated sources). We show sources with GALFIT flag = 1 (a suspicious fit) in gray and sources with GALFIT flag = 0 (a good fit) with black datapoints. Bad fits (GALFIT flag >1) were ignored. The median magnitude difference for galaxies with GALFIT flag = 0 is shown by the red solid line and filled bulletpoints in bins of 0.5 mag. We also indicate the median offset in the legend. We find slightly brighter magnitudes with GALFIT, of 0.006-0.070 mag on average, with the difference presumably attributable to different techniques to derive total magnitudes and potential color gradients.

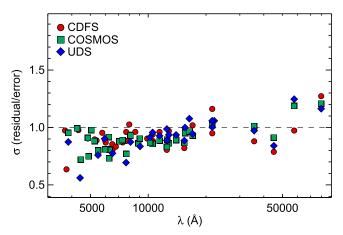


Figure 12. NMAD scatter in the error-normalized flux residuals as a function of wavelength, for the three ZFOURGE fields. σ is close to unity for most filters, indicating the photometric errors are accurate.

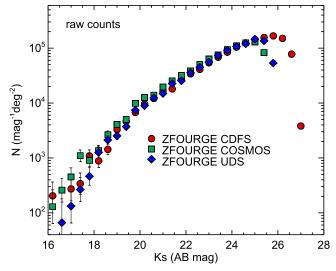


Figure 13. K_s -band total magnitude number counts of sources with use = 1. We observe a turn-over in the histograms around 25–26 mag, with less detections for fainter sources.

are shown in Appendix A. We find that 12% of the images have a PSF that is broader than our target PSF. Our method improves the accuracy of the final convolved PSFs, compared with e.g., maximum likelihood algorithms. For example, Skelton et al. (2014) find <1% accuracy when convolving HST/WFC3 images, using the same technique as employed here, compared to e.g., Williams et al. (2009) and Whitaker et al. (2011), who use maximum likelihood methods to match point sources to within 2%-5% accuracy.

PSF curves of growth before and after convolution are shown in Figure 9, normalized by the model PSF. For each convolved image we obtain excellent agreement, within 1.5% at r < 0."6. IRAC photometry, with FWHM > 1."5 is treated separately in Section 3.6.

3.2. Source Detection

We created detection images from the superdeep background subtracted K_s -band images, as described in Section 2.3, by noise equalizing the images, i.e., multiplying the images with the square root of the corresponding weight images. We then ran SE to create a list of sources and their locations. We optimized source detection by setting the deblending parameters of SE to DEBLEND_THRESH = 64 and DEBLEND_MINCONT = 0.0000001 and the clean parameter (CLEAN) to N. We also generated a segmentation map with SE representing the location and area of each source. The total number of sources in the catalogs is 30,911 in CDFS 20,786 in COSMOS and 22,093 in UDS. Our SE parameter files are included in the ZFOURGE data release.

3.3. K_s-band Total Flux Determination

To measure the K_s -band total flux, SE was run in dual image mode on the superdeep K_s -band images, using the noise equalized images (Section 3.2) for source detection. We used a flexible elliptical aperture (Kron 1980), to obtain SE's FLUX_AUTO.

This estimate is not yet the total K_s -band flux and we have to account for missing flux outside the aperture. We derived a correction factor from the stacked K_s -band PSF separately for each field. This aperture correction varies between sources and

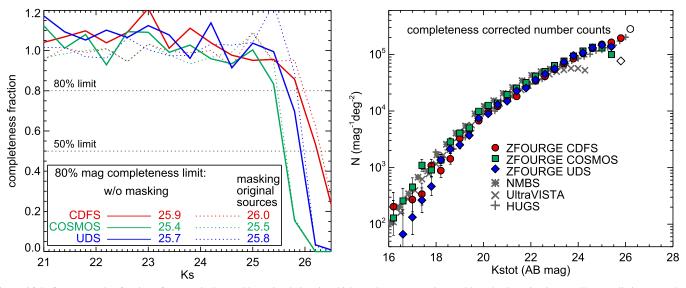


Figure 14. Left: we test what fraction of sources is detected by a simulation, in which mock sources are inserted into the detection images. For a realistic approach, in which sources are allowed to overlap, we recover 80% down to very deep magnitudes of $K_s = 25.4$ –25.9. Right: completeness corrected number counts. We correct the observed counts in Figure 13 using the completeness estimates in each field. Here we omit sources at magnitudes fainter than the 50% completeness limits. Sources with <80% completeness are shown with open symbols. For comparison, the K_s -band completeness corrected number counts of similar galaxy surveys are indicated in grayscale.

Table 6Completeness Test Results

	With Masking		w/o Masking	
	80%	50%	80%	50%
CDFS	26.0	26.3	25.9	26.2
COSMOS	25.5	25.6	25.4	25.6
UDS	25.8	26.0	25.7	25.9

 Table 7

 Explanation of EAZY Photometric Redshift Catalog Header

id	ID number
z_spec	spectroscopic redshift (if no redshift available, z_spec is set to -1)
z_a	photometric redshift derived without a K luminosity prior
z_m1	weighted redshift derived without a K luminosity prior
chi_a	minimum χ^2 derived without a K luminosity prior
z_p	best-fit redshift after applying the prior
chi_p	minimum χ^2 after applying the prior
z_m2	weighted redshift after applying the prior
odds	parameter indicating presence of second χ^2 minimum (1 if no minimum)
168,u68	1 sigma confidence interval
195,u95	2 sigma confidence interval
199,u99	3 sigma confidence interval
nfilt	number of filters used in the fit
q_z	quality parameter
z_peak	default derived photometric redshift
peak_prob	peak probability
z_mc	randomly drawn redshift value from redshift probability distribution

is a function of the size of the auto-aperture that was estimated by SE.

We determined the aperture correction by using the curve of growth of the PSF. Total K_s -band fluxes were then calculated

using

$$F_{Ks,\text{tot}} = F_{Ks,\text{auto}} \frac{F_{PSF}(<4'')}{F_{PSF}(< r_{Kron})}$$
 (2)

(Labbé et al. 2003b; Quadri et al. 2007), with $F_{Ks,\text{tot}}$ the total K_s -band flux, $F_{Ks,\text{auto}}$ the flux within the auto-aperture, i.e., FLUX_AUTO from SE, F_{PSF} (<4'') the flux of the PSF within a 4'' radius and F_{PSF} ($< r_{Kron}$) the flux within the circularized Kron radius.

We additionally measured the total flux using a fixed circular aperture, of $\sim 1.5 \times$ the PSF FWHM of the deep K_s -band images. In CDFS we therefore used a 0."7 diameter aperture and in COSMOS and UDS a 0."9 diameter aperture. These aperture fluxes were also corrected for flux outside of the aperture.

Therefore we have two estimates for the total flux, one using the auto aperture flux, and one using a fixed circular aperture. For small, low S/N sources, the autoscaling aperture size may be very small, leading to extreme aperture corrections. Therefore, we only considered the circular aperture measurements for sources if their circularized Kron radius was very small, i.e., smaller than the circular aperture radius.

3.4. Aperture Fluxes

In addition to the total K_s -band flux, we derived flux estimates in all filters in the three ZFOURGE fields. We ran SE in dual image mode, using the combined K_s -band images for source detection and the PSF matched images to measure photometry. We use the PSF matched images to make sure the captured light within the apertures is consistent over all the images. We also included the convolved versions of the deep K_s -band stacks. We use circular apertures of 1.12 diameter, which are sufficiently large to capture most of the light (the PSFs of the convolved images have a FWHM = 0.19), but small enough to optimize S/N.

We correct all aperture fluxes to total, using the ratio between the total flux in the original deep K_s -band stacked images to the aperture flux in the PSF matched K_s -band stack,

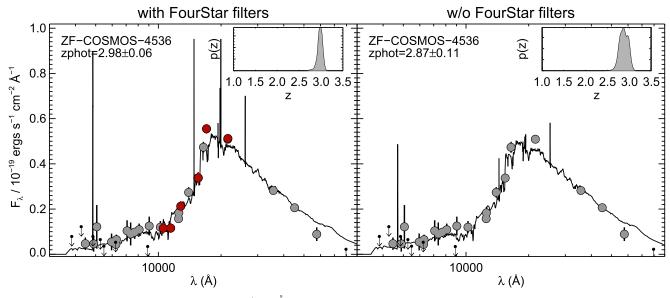


Figure 15. Example galaxy at z=2.1 with a large Balmer/4000 Å break, traced by the FourStar filters (indicated in red). Gray datapoints represent flux in ancillary filters, with downward pointing arrows representing upper limits. In the left panel we show the best-fit SED template derived using the FourStar near-IR medium-bandwidth filters. In the right panel we show the best fit, without the FourStar bands. The insets show the redshift probability functions corresponding to the fits. Including the FourStar filters leads to a factor two better constraint on the photometric redshift.

i.e.,:

$$F_{F,\text{tot}} = F_{F,\text{aper}} * \frac{F_{Ks,\text{tot}}}{F_{Ks,\text{aper}}}.$$
 (3)

Here, $F_{F,\text{tot}}$ is the aperture flux in filter F scaled to total, $F_{F,\text{aper}}$ the unscaled aperture flux, $F_{Ks,\text{tot}}$ the total K_s -band flux described in Section 3.3 and $F_{Ks,\text{aper}}$ the aperture flux from the PSF-matched K_s -band image stacks.

3.5. Flux Uncertainties

The uncertainty on the flux measured in an aperture has contributions from the background, the Poisson noise of the source, and the instrument read noise. The relative contribution from the latter two effects will be very small for the faint galaxies and medium band filters used in this study (Persson et al. 2013). If the adjacent pixels in an image are uncorrelated, the background noise $\sigma_{\rm rms}$ measured in an aperture containing N pixels will scale in proportion to \sqrt{N} . In a more realistic scenario, pixels are expected to be correlated on small scales due to interpolation or PSF smoothing and on large scale due to imperfect background subtraction, flux from extended objects, undetected sources, or systematics introduced in the reduction process, such as flat field errors. For perfectly correlated pixels, the background noise is expected to scale as $\sigma_{\rm rms} \propto N$. The actual scaling of the noise in an image lies somewhere in between and can be parameterized by

$$\sigma_{\text{NMAD}} = \sigma_1 \alpha N^{\beta/2} \tag{4}$$

with $\sigma_{\rm NMAD}$ the normalized median absolute deviation and β taking on a value between $1 < \beta < 2$. α is a normalization parameter and σ_1 is the standard deviation of the background pixels. (Labbé et al. 2003a; Quadri et al. 2007; Whitaker et al. 2011) We estimated the noise as a function of aperture size empirically by placing circular apertures of varying diameter at 2000 random locations in each image that was used for photometry. These are the convolved images for the aperture

fluxes and the unconvolved K_s -band stacks that were used to measure total flux. We used the SE segmentation map to mask sources. We also excluded regions with low weight, such as the edges of the FourStar detectors.

For each aperture diameter, we fit a Gaussian to the measured flux distribution and obtained the standard deviation ($\sigma_{\rm rms}$). We then fit Equation (4) to the various estimates of $\sigma_{\rm rms}$ as a function of N pixels in each aperture, to obtain σ_1 , α and β .

For circular apertures with radius r pixels, the uncertainty (e_F) on the flux measurement in filter F is

$$e_F = \sigma_{\text{NMAD}}(r) / \sqrt{w_F} = \sigma_1 \alpha (\pi r^2)^{\beta/2} / \sqrt{w}$$
 (5)

with w_F the median normalized weight. We did not include a Poisson error in our flux uncertainties, as faint sources are background-limited, while uncertainties on bright sources are dominated by systematics.

Weights were obtained from the median normalized exposure images and were measured as the median in apertures with sizes corresponding to those used to measure flux.

The radius r used in Equation (5) was chosen to match the aperture size used for the different flux determinations. 1."2 diameter apertures are used for the aperture fluxes and, for the total flux, we use SE's KRON_RADIUS, which is based on autoscaling kron-like apertures.

The aperture flux uncertainties obtained from Equation (5) were scaled to total for a consistent relative error.

3.6. IRAC and MIPS Photometry

The *Spitzer*/IRAC and MIPS images (in all fields) and *Herschel*/PACS images (available for CDFS only) have much broader PSFs than the UV, optical and near-IR images and source blending is a significant effect. The FWHM in the IRAC images is typically >1."5 and in MIPS >4". To obtain photometry, we use a source fitting routine that models and subtracts profiles of neighboring objects prior to measuring photometry for a target (Labbé et al. 2006; Wuyts et al. 2008; Whitaker et al. 2011; Skelton et al. 2014; Tomczak et al. 2016).

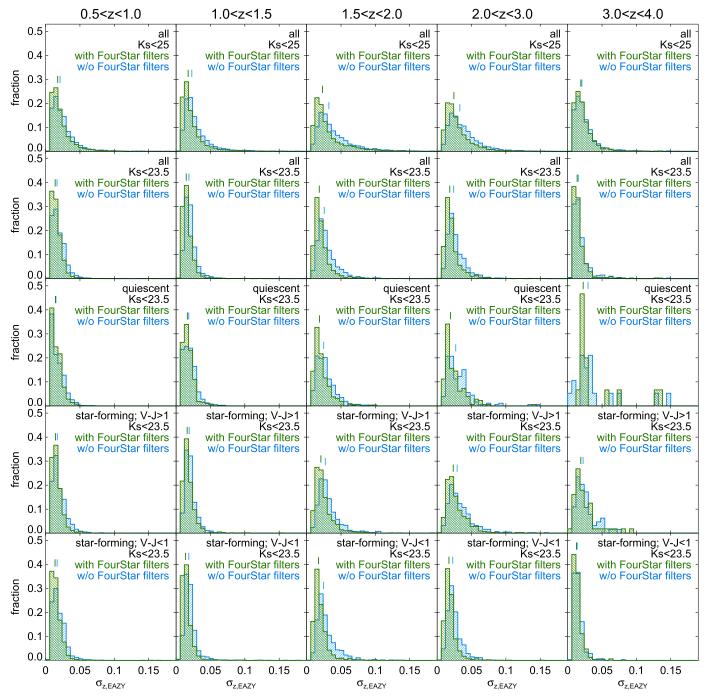


Figure 16. Redshift error, $\sigma_{z, \text{EAZY}}$ histograms in redshift bins (from left to right), normalized to the total number of sources in each bin. In the first and second row we inspect general magnitude-limited samples, with $K_s < 25$ AB and $K_s < 23.5$. In the third to last rows of panels we show, respectively, the error histograms of quiescent galaxies, red star-forming galaxies with $V - J \geqslant 1$, and blue star-forming galaxies with V - J < 1. The median $\sigma_{z, \text{EAZY}}$ is indicated just above the histograms in each panel, using the respective colors (green or blue) of the EAZY fits with and without near-IR medium bands. The photometric uncertaintes are systematically smaller if we include the FourStar near-IR medium-bandwidth filters when fitting SED templates (green histograms).

The position and extent of each source was based on the SE segmentation maps derived from the super deep K_s -band detection images. The K_s -band images are assumed to provide a good prior for the location and extent of the unresolved far-IR flux, as sources that are bright in K are also typically bright at redder infrared wavelengths. Each source in the K_s -band image was extracted using the segmentation map and convolved to match the PSF of the lower resolution far-IR image, assuming negligible morphological corrections. All sources were then fit simultaneously to create a model for the lower resolution image. Next, for each

source in the lower resolution image, the modeled light of neighboring sources was subtracted, after which we measured the flux on the cleaned maps within circular apertures with diameter D, using D=1.9 for IRAC and D=79 for MIPS.

To correct the far-IR aperture flux to total, the measurements were multiplied by the ratio of the total K_s -band flux to the D=1.8 aperture flux on the PSF convolved K_s -band template image. Because the MIPS PSF has significant power in the wings at large radii, which are not represented in the convolution kernel, we apply an additional fixed correction

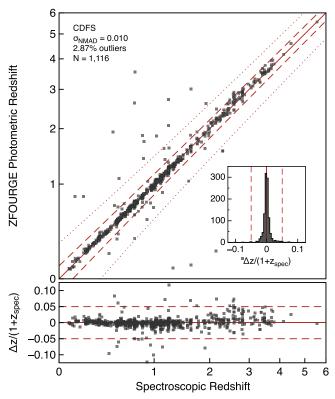


Figure 17. Top: photometric redshifts from ZFOURGE vs. spectroscopic redshifts in CDFS. The NMAD scatter, the fraction of objects with $\Delta z/(1+z_{\rm spec})>0.15$, and the number of galaxies with matches in both catalogs are shown in the upper left of the plot, while the histograms of $\Delta z/(1+z_{\rm spec})$ are shown as an inset in the bottom right of the plot. Bottom: the residual between the photometric and spectroscopic redshifts, divided by $1+z_{\rm spec}$. The red solid, dashed and dotted lines indicate, respectively, $\Delta z/(1+z_{\rm spec})=0\pm0,\pm0.05,$ and $\pm0.15.$

of $\times 1.2$ to account for missing flux at r > 15'' (using values for point-sources from the MIPS instrument handbook).

Flux uncertainties were estimated from background maps. These were individually generated for each source on scales of three times the 30" tile size used for the modeling, using the cleaned tiles. From these we measured rms variations using apertures at random locations. Variations on larger scales were corrected by spatially adjusting the zeropoint using the iterative procedure described in Section 5.

3.7. Stars

The majority of stars were identified by their observed B-J123 and $J123-K_s$ colors. J123 here is derived as the median of the flux in the J1, J2 and J3 filters. Stars form a tight sequence in $J123-K_s$ compared to galaxies. In the first two panels of Figure 10 our selection criterion is indicated as a red line, with stars having:

$$(J123 - K_s) < 0.288(B - J123) - 0.52[(B - J123) < 2.5]$$

 $(J123 - K_s) < 0.08(B - J123)$ [(B - J123) > 2.5].

Here we only classified sources as stars if they are below the red line at $>2\sigma$ confidence in $J123-K_s$. By selecting at $>2\sigma$ confidence, we automatically reject faint sources that scatter below the red line from the star sample. This is illustrated by the histograms in the third panel of Figure 10, where we show the

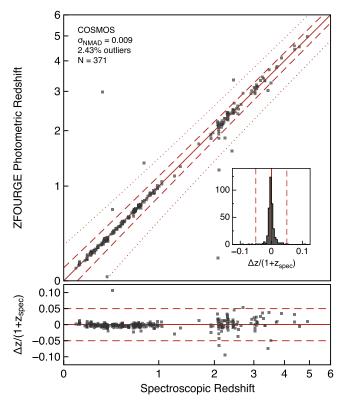


Figure 18. Photometric vs. spectroscopic redshifts for COSMOS (see caption of Figure 17).

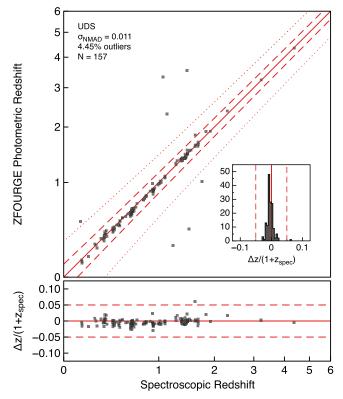


Figure 19. Photometric vs. spectroscopic redshifts for UDS (see caption of Figure 17).

magnitude counts of stars against sources that are not now classified as stars, but would have been otherwise. These have a distribution of magnitudes that peaks around $K_s = 25.5-26$ mag.

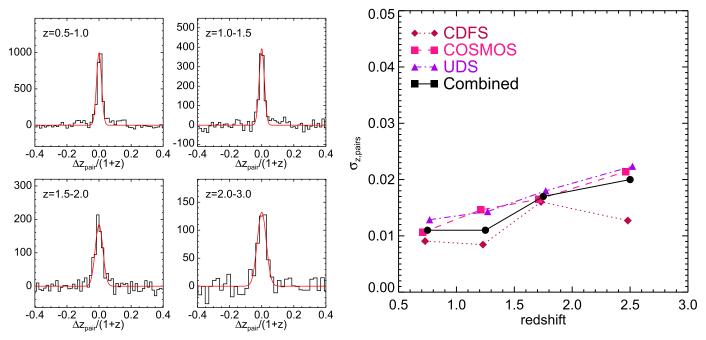


Figure 20. Analysis of photometric redshift accuracy using close pairs of galaxies. Left: distribution of $\Delta z_{\text{pairs}}/(1+z)$ for galaxies with $K_s < 23.5$. We fit Gaussians to each histogram (red lines), from which we derive $\sigma_{z,\text{pairs}}$, the average uncertainty for individual galaxies, which is the standard deviation of the pair distribution divided by $\sqrt{2}$ (1 + z_{mean}). Right: $\sigma_{z,\text{pairs}}$ as determined in the left panel as a function of redshift (black solid line). We also show the results for the individual fields, as indicated in the legend. $\sigma_{z,\text{pairs}}$ increases with redshift, from $\sigma_{z,\text{pairs}} = 0.01$ to 0.02.

For sources that were not covered by the J1, J2 and J3 bands, we used broadband J or F125W where available, but only considered sources brighter than 25 mag in K_s . For sources without B-band coverage, we used $J123 - K_s < 0$ to classify stars, considering only sources brighter than 22 mag in K_s . An finally, if sources did have B-band coverage, but were saturated in B, we also used $J123 - K_s < 0$.

Red cool stars may not be selected in this way, as they have red J-K colors. To ensure that we cover all types of stars, we fit the observed SEDs of all sources with EAZY using the speX stellar library. ¹⁵ For a few sources that were not flagged already by their B-J123 and J123-K colors, the reduced χ^2 indicated a stellar template was a better fit to the data than the best-fit galaxy template (Section 5) and we flagged these sources as stars as well.

A source that is not selected by any of the methods above, is considered a saturated star if it is brighter than 16 mag in J_1 or K_s and at the same time could not be fit to a galaxy template, having a large reduced χ^2 , which we empircally estimated by inspecting many SEDs to be $\chi^2 > 3000$.

In total, 1.8% of the sources in the catalogs are classified as stars.

3.8. Catalog Format

We provide separate photometric catalogs for each cosmic field. These contain the co-ordinates, total fluxes, flux uncertainties, weight estimates, flags and S/N estimates of each source. Individual sources are indicated by their ID, starting at ID = 1. A description of the columns is given in Table 5.

The CDFS catalog contains 30,911 sources, the COSMOS catalog 20,786 and the UDS catalog contains 22,093 sources. Magnitudes for each source can be obtained by applying a

zeropoint of 25 in the AB system (corresponding to a flux density of $3.631 \times 10^{-30} \, \mathrm{erg \, s^{-1} \, Hz^{-1} \, cm^{-2}}$ or $0.3631 \, \mu \mathrm{Jy}$). e.g., the stacked K_s -band total magnitude is $25 - 2.5 \times \log_{10} \mathrm{f_Ksall}$.

All fluxes in the catalogs are scaled to total. They can be converted back to aperture flux (1."2 diameter) by dividing by Ks_ratio for each source. The exceptions are fap_Ksall. fauto_Ksall and fapcirc0D_Ksall. The first is the actual (convolved) K_s -band aperture flux, and can only be converted in the other direction, toward total. The second is the auto aperture flux from SE, and we need only to apply the aperture correction, apcorr, described in Section 3.3, to obtain f_Ksall. The last is an alternative to fauto_Ksall, and is measured in a fixed circular aperture with diameter D, instead of the flexible elliptical Kron-like aperture from SE (using apertures of D = 0."7 in CDFS and D = 0."9 in COSMOS and UDS). From fapcirc0D_Ksall we can obtain the total fcirc0D_Ksall by multiplying with apcorr0D.

Each flux measurement of each source in each filter has been assigned a weight, reflecting the depth in the images at the source locations. The weights are normalized to the median of the corresponding weight images. In the catalogs we also indicated the minimum weight for sets of filters. For example, the lowest weight of the FourStar filters is indicated by wmin_fs. If this value is greater than 0, it means a positive weight in all FourStar images.

In addition to photometric catalogs, we provide the EAZY (Section 5) and FAST (Section 6) output files, containing the photometric redshifts and stellar population properties.

3.9. A Standard Selection of Galaxies

For convenient use of the catalogs, we have designed a use flag. This takes into account S/N, the star/galaxy classifications described above and the depth of the images at the respective source locations. This flag also includes sources that

¹⁵ http://pono.ucsd.edu/~adam/browndwarfs/spexprism/

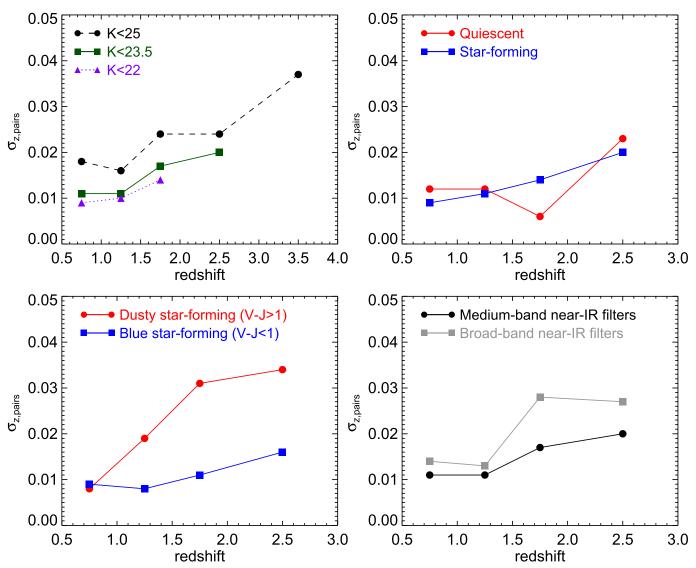


Figure 21. $\sigma_{z,pairs}$ vs. redshift, investigating trends with magnitude limits (top-left) or SED types (second and third panels), and investigating the effect of using the FourStar filters (top-right). $\sigma_{z,pairs}$ tends to be smaller for brighter galaxies are considered, and for blue star-forming galaxies. $\sigma_{z,pairs}$ is clearly smaller if the near-IR medium band filters are used (compared to near-IR broadband), especially at 1.5 < z < 2.5.

are well within the FourStar footprint and are observed with each of the near-IR medium-bandwidth filters. The K_s -band stacks cover a somewhat larger area, especially in CDFS, which means that not all sources in the catalogs have FourStar imaging (although the majority do). A standard selection of galaxies can be obtained by selecting on use = 1 (see full definition in Table 5) from the catalogs.

The use flag allows for a straightforward sample selection, representing galaxies with good photometry, i.e., high S/N sources from well exposed regions of the images. For specific science goals a different selection may be optimal. We also warn that the use = 1 sample may still contain problematic sources, with for example uncertain photo-z's and poorly constrained EAZY or FAST fits, and we recommend to always inspect the individual SEDs. However, the use = 1 sample should be a reliable representation of the galaxy population in large statistical studies.

The total area of the Ks-band detection images is 280.9^2 for CDFS, 176.5^2 for COSMOS and 189.3^2 for UDS. Selecting only galaxies with wmin_fs > 0.1 that are not near bright stars, reduces the area to 132.2^2 , 139.2^2 and 135.6^2 .

3.10. Quality Verification

3.10.1. Flux Comparisons

Here we test whether the total fluxes derived above are reliable, by (1) comparing our magnitudes in various bands to independent estimates by a different survey directly, and (2) by comparing our total magnitudes in the detection band to a completely different method to measure total flux. For the former we used the 3D-HST data set, as both surveys use many of the same images and cover similar fields. Many of the same basic image reduction methods were used to derive photometry for 3DHST, but we performed our own alignment and registration to the 0."15 pix⁻¹ scale of FourStar and an independent estimate of the background. In detail, photometric methods between the two surveys differ significantly, and in addition we derived total flux from the ultra-deep K_s band images, whereas source detection and total flux derivation for 3D-HST was based on HST/WFC3/F160W images. In general we find excellent correspondence between the two surveys. We show diagnostic plots in Appendix B.

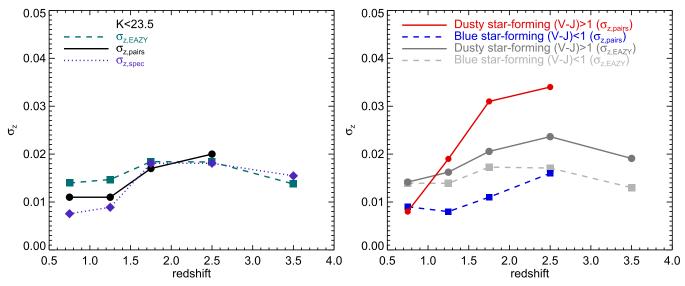


Figure 22. Left: comparing the three different redshift quality tests, σ_z from testing against spectroscopic samples, $\sigma_{z,pairs}$ from the redshift pair test, and the photometric redshift uncertainty from the p(z) from EAZY, $\sigma_{z,EAZY}$. These correspond well at all redshifts. Right: comparing $\sigma_{z,pairs}$ with $\sigma_{z,EAZY}$ tends to be underestimated for red star-forming galaxies.

We tested our method of extracting total flux through SE by comparing to total flux derived with GALFIT (Peng et al. 2010a), a program which fits two-dimensional model light profiles to galaxy imaging. The fitting process benefits from high resolution imaging, so we make use of the HST/ WFC3/F160W size catalogs from van der Wel et al. (2014), based on the source catalog of 3D-HST, which contains parameters derived with GALFIT. Galaxies may have different morphologies in different bands. However, as the F160W and K_s -band filters lie very closely together in wavelength space, we assume that the correction to total in our catalogs, which is based on the ratio between K_s -band aperture and total flux, also produces an accurate approximation of total F160W-band flux. As the comparison with 3DHST shows (Appendix B), these magnitudes are accurate to within \sim 1%, with magnitude offsets of 0.017, 0.005 and -0.011 in CDFS, COSMOS and UDS, respectively.

The comparison with GALFIT magnitudes is shown in Figure 11. We use the goodness of fit flag included in the size catalogs to select sources with a good (GALFIT flag = 0) or suspicous fit (GALFIT flag = 1), but not sources with bad fits (GALFIT flag >1). We find a median offset between ZFOURGE total F160W magnitude and GALFIT magnitude of -0.046,-0.070 and -0.006 mag, for CDFS, COSMOS and UDS, respectively. Skelton et al. (2014) show the same comparison, with similar trends with magnitude, and find magnitude offsets for the three fields of -0.03, -0.04 and 0.00. The small offsets that we find between GALFIT magnitude and magnitude derived with SE, are likely attributable to details of our photometric procedure to determine total fluxes, which are somewhat dependent on galaxy profile and S/N (see Labbé et al. 2003a; Skelton et al. 2014) and possible color gradients.

3.10.2. Flux Uncertainty Verification

Here we test the accuracy of the flux uncertainties derived in Section 3.5. We used the outcome of the SED fitting described below in Section 5. The residual between the best-fit template

flux and the observed flux in a filter should reflect the photometric errors in the catalogs. If these are accurate, then normalizing the distribution of the residuals by the photometric error, should result in a Gaussian with a width of unity. We derived the NMAD of the distribution of the error-normalized residuals, and show the scatter (σ) , for each filter in the catalog, as a function of wavelength in Figure 12. Overall these look very good, with the average $\sigma_{\rm NMAD}$ very close to unity, and the vast majority (>90%) of bands within 20% of unity.

3.10.3. Close Pair Contamination

It is naturally expected that some sources lie in close angular proximity of each other, and may contaminate the aperture flux of their close neighbor. For UV to near-IR photometry, this may lead to systematic errors on the aperture photometry of a source, especially if the neighbor is much brighter (for IRAC and MIPS photometry we used a source fitting routine that takes into account flux from neighboring sources; see Section 3.6). The aperture diameter that we used above is 1"2. We inspected the catalogs for pairs of galaxies that lie closer than 1"2 distance away from each other. We only looked at sources that are not already classified as stars or as being located in the neighborhood of a bright star, as we already accounted for these sources that their flux estimate may be affected. The percentage of sources with a neighbor at <1.2distance is 3.8% in CDFS, 4.1% in COSMOS, and 4.4% in UDS. If only the fainter part of a projected galaxy pair is affected, we estimate that $\sim 2\%$ of the sources in each field may suffer flux effects from nearby sources.

4. COMPLETENESS

We counted the number of sources with use = 1 per Ks-band total magnitude bin in each catalog. This result, taking into account the effective area corresponding to the use flag, is shown in Figure 13. For the different fields, the histograms turn over at 25.5-26 mag, indicating it becomes more difficult to detect fainter sources.

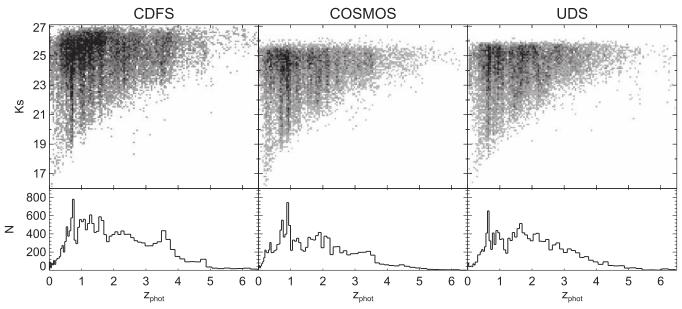


Figure 23. Top panels: distribution of K_s -band magnitudes as function of redshift. The grayscale indicates the density in each point, with darker colors for higher densities. Bottom panels: photometric redshift (z_peak) distribution. Higher density peaks are clearly visible, for example the ZFOURGE identified cluster at z = 2.095 in COSMOS (Spitler et al. 2012; Yuan et al. 2014).

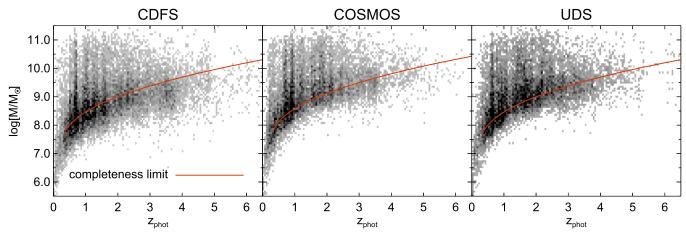


Figure 24. Similar to the top panels of Figure 23, but with stellar mass instead of K_s -band magnitude. The solid red line indicates the 80% mass completeness limit in each field.

 Table 8

 Explanation of The FAST Stellar Population Catalog Header

id	ID number	
z	=z_peak (or z_spec if available)	
ltau	$\log[\tan yr^{-1}]$	
metal	metallicity (fixed to 0.020)	
lage	log[age yr ⁻¹]	
Av	dust reddening	
lmass	$\log[M/M_{\odot}]$	
lsfr	$\log[SFR/(M_{\odot} \text{ yr}^{-1})]$	
lssfr	$log[sSFR(yr^{-1})]$	
la2t	$\log[{ m age}/ au]$	
chi2	minimum χ^2	

To test how well sources are recovered from the images, we perform completeness tests, using the super-deep K_s -band detection images. We drop 10,000 mock sources, obtained from median stacking low S/N (9 < S/N_{K_s} < 11) sources with use = 1, in the detection images. The stacks were scaled to a

Table 9
Explanation of The SFR Catalog Header

id	ID number
Z	phometric redshift (or spectroscopic redshift if available)
f_24	Spitzer/MIPS 24 μ m flux (mJy)
e_24	Spitzer/MIPS 24 μ m flux error (mJy)
f_100 ^a	Herschel/PACS 100 μ m flux (mJy)
e_100ª	Herschel/PACS 100 μ m flux error (mJy)
f_160 ^a	Herschel/PACS 160 μm flux (mJy)
e_160 ^a	Herschel/PACS 160 μ m flux error (mJy)
L_IR	total integrated IR luminosity L_{\odot}
L_UV	total UV luminosity L_{\odot}
SFR	star formation rate (Equation (6))

Note.

magnitude range of 18 < mag(AB) < 27.5. We used a powerlaw distribution of magnitudes, matching the slope of the number counts in Figure 13 between $K_s = 21 \text{ AB}$ and

^a Herschel/PACS data only available in CDFS.

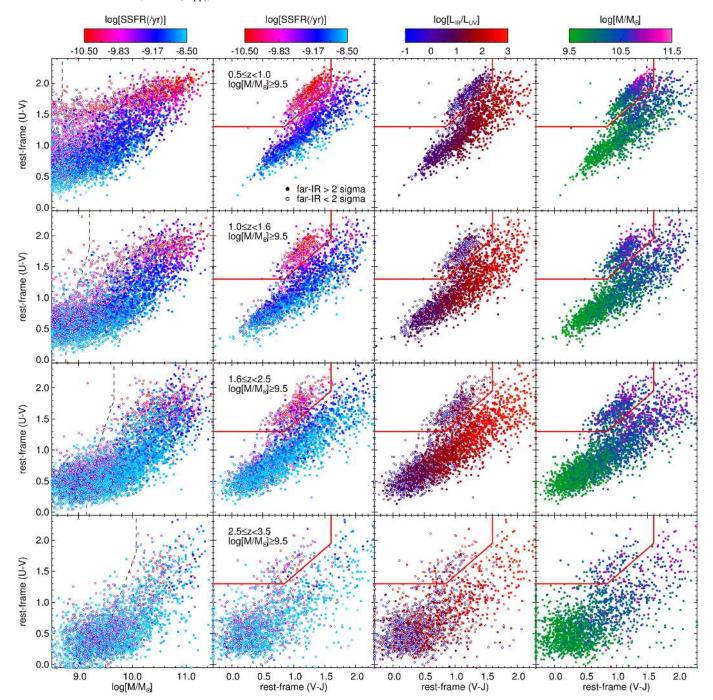


Figure 25. Rest-frame U-V vs. stellar mass (left columns) or vs. V-J (second, third and fourth columns) for a mass complete sample of galaxies with use = 1, $S/N_{K_S} > 10$ and stellar mass $M>10^{9.5}M_{\odot}$, in four redshift bins (top to bottom). The vertical dashed lines indicate our stellar mass completeness limit. The red solid line in the UVJ diagrams separates quiescent (top left) from star-forming (bottom left to top right) galaxies. Galaxies that are undetected in the far-IR at $<2\sigma$ are shown with open symbols. The sSFRs decrease toward blue V-J and red U-V colors, reflecting a gradient in age (see also Figure 26). Galaxies with the lowest sSFRs are located in the quiescent region of the diagram, with red U-V and blue V-J colors. Galaxies span a large range in $\log[L_{\rm IR}/L_{\rm UV}]$, ranging from -1 for the bluest UVJ star-forming galaxies to 3 for the dustiest sources, and quiescent galaxies having low $\log[L_{\rm IR}/L_{\rm UV}]$. A mass sequence is also visible, where massive galaxies tend to be redder.

 $K_s = 25$ AB. The distribution follows $d \log N/K_s = 0.24$, i.e., a factor 1.7 more sources per unit magnitude, with N the number of sources and K_s the total K_s -band magnitude, in agreement with previously determined values (e.g., Fontana et al. 2014). We ran SE using the same input parameters used to generate the catalogs. We measured the observed magnitude of the input sources that were retrieved with SE. We then compared these with the input source distribution to calculate

the correction as a function of observed magnitude, accounting for both completeness and scatter.

We performed the simulation in two ways. First by simply dropping mock sources randomly in the images, only excluding a few small areas around a few very bright stars. To prevent artificial crowding of simulated sources, we only dropped in 500 sources per run, and repeated the simulations a large number of times.

Next we investigated what fraction of incompleteness is due to crowding, where bright sources prevent the detection of fainter sources nearby. We masked all detected sources, using the segmentation map from SE and constrained the location of the simulated sources, such that they do not overlap. In this way, we purely tested if sources can be detected above the noise level in the images.

We show the results of the two tests in the left panel of Figure 14. Even if only stars are masked and sources are allowed to overlap (solid lines) we recover at least 80% down to very deep K_s -band magnitudes of 25.4–25.9 and 50% down to 25.6–26.2. These values correspond well with the turnover in K_s -band number counts in Figure 13 and the stacked K_s -band image depths (Section 2.3). The 50% and 80% completeness limits of both tests are tabulated in Table 6. The slight elevation with a higher than 100% completeness fraction at magnitudes <24.5 for the non-masking case is due to confusion with bright sources.

We correct the number counts from Figure 13 using the completeness estimates as function of observed magnitude from the more conservative test (obtained w/o masking, i.e., the solid curves) in each field and show these in the right panel of Figure 14. We also include similar results from the NMBS (Whitaker et al. 2011), UltraVISTA (Muzzin et al. 2013a) and HUGS (Fontana et al. 2014) surveys. NMBS and UltraVISTA have shallower depths, but much larger areas than ZFOURGE. Our number counts agree with these earlier results from the literature. They also show that ZFOURGE is one of the most sensitive surveys to date, comparable to HUGS and 1–2 mag deeper in K_s than earlier groundbased surveys. Similar to NMBS, we find an excess of sources at brighter K_s -band magnitudes in COSMOS.

5. PHOTOMETRIC REDSHIFTS

5.1. Template Fitting

Photometric redshifts were derived with EAZY (Brammer et al. 2008), by fitting linear combinations of nine spectral templates to the observed SEDs. Of these, seven are the default templates described by Brammer et al. (2008), five of which are from a library of PÉGASE stellar population synthesis models (Fioc & Rocca-Volmerange 1999), one represents a young and dusty galaxy and another is that of an old, red galaxy (see also Whitaker et al. 2011). The final two templates represent an old and dusty galaxy and a strong emission line galaxy (Erb et al. 2010). The code has the option to include a template error function, which we use, to account for systematic wavelength-dependent uncertainties in the templates. We also make use of a luminosity prior, based on the apparent magnitude calculated from the total K_s -band flux.

Offsets in the zeropoints may systematically affect the measured flux and therefore also the derived photometric redshifts. We correct for zeropoint offsets, by iteratively fitting EAZY templates to the full optical-near-IR observed SEDs. This procedure is described in detail by Whitaker et al. (2011) and Skelton et al. (2014). Similar to Skelton et al. (2014), we use all sources in the fits, including those without a spectroscopic redshift available. We also use a two step process in which we first only vary the zeropoints of the *HST*-bands and then, keeping these fixed, we vary the zeropoints of the groundbased and *Spitzer*/IRAC data.

During this iterative fitting procedure, both the zeropoints and the templates were modified. These are separable

corrections, as the templates are modified after shifting both the data and the best-fit SEDs to the rest-frame. Due to the wide range of galaxy redshifts and large number of filters in the catalogs, each part of the spectrum is sampled by a number of photometric bands. In small bins of rest-frame wavelength, we determined systematic offsets between the data and the templates and updated the templates. This allows the templates to reflect subtle features not initially included, such as the dust-absorption feature at 2175 Å. After adjusting the templates, zeropoint corrections are calculated in the observed frame. The process is repeated until zeropoint corrections in all bands except U or the IRAC bands become less than 1% and this typically happens after three or four iterations.

The zeropoint offsets are listed in Tables 2–4. The zeropoints in these tables are the effective zeropoints, with galactic extinction and the zeropoint offsets incorporated. The offsets are typically of the order of 0.05 mag. The largest offsets occur for the COSMOS and UDS U-bands, which are known to have uncertain zeropoints (Erben et al. 2009; Whitaker et al. 2011; Skelton et al. 2014). Template and zeropoint errors are hardest to separate from each other for the U- and IRAC 8 μ m bands, as these lie at the blue and red ends of the spectra, without bracketing filters.

The residuals between the best-fit templates and observed SEDs are excellent tracers of spatial variations in the zeropoint. We found small variations for all images. In particular, we were able to pinpoint small offsets between the different quadrants of the FourStar images in UDS. To alleviate the spatial effect, our final derivation for every filter includes two runs of the fitting process. After the first run we remove a two-dimensional polynomial fit to the spatial residuals. This is directly incorporated into the catalogs, i.e., we apply a correction to all sources as a function of their *x*- and *y*-coordinates in the images and using the corresponding two-dimensional offsets in each filter. Finally the fitting process is repeated in the way described above to obtain the final zeropoint offsets.

The spatial variations in zeropoint of the VLT/VIMOS/*R*-band image are larger than in the other images and could not be described by a polynomial function. This image is very deep, so we do not wish to discard it. We therefore impose a minimum error on the flux of 5%. In Figure 34 in Appendix C we show the residual maps after subtracting the polynomial fits.

We use the output parameter z_peak from EAZY as indicator of the photometric redshift. z_peak is estimated by marginalizing over the redshift probability distribution function, p(z). If p(z) has more than one peak, z_peak only marginalizes over the peak with the largest integrated probability.

We provide the full EAZY photometric redshift catalogs. See Table 7 for an explanation of the catalog header.

5.2. Photometric Redshift Uncertainties Determined by EAZY

As a result of the use of near-IR medium-bandwidth filters, spectral features such as the Balmer/4000 Å break are better sampled for galaxies at 1.5 < z < 3.5. In Figure 15 we illustrate the ability of the FourStar medium-bandwidth filters to constrain galaxy SEDs and redshift probability distibutions. We can determine a photometric redshift error due to the fitting process, using the 16th–84th percentiles from p(z). For better constrained redshifts, p(z) will be narrower and the error on $z_{\rm phot}$ will be smaller. We show the SED of a galaxy at a redshift of $z=2.98\pm0.06$, with the uncertainty derived

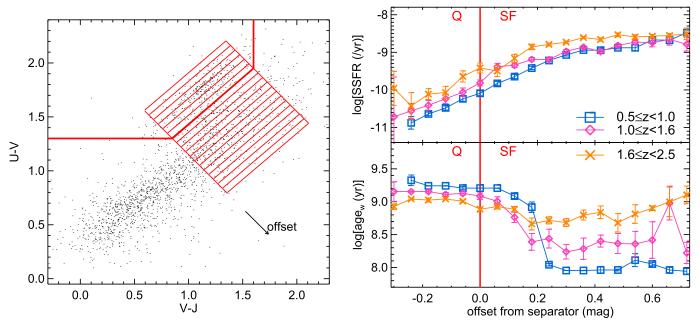


Figure 26. Left: bins in UVJ space, set orthogonally to the red line that separates quiescent from star-forming galaxies. In these bins we measure the bootstrapped median SSFR and weighted age: age_w . The underlying datapoints are a randomly chosen subsample from Figure 25 to illustrate the general locus of galaxies in the diagram. Right: gradient in sSFR (top) and age_w (bottom) with increasing distance from the separator. Datapoints on the left of the red line represent quiescent galaxies ("Q") and datapoints on the right star-forming galaxies ("SF"). sSFR increases toward bluer U - V and V - J, while age_w decreases.

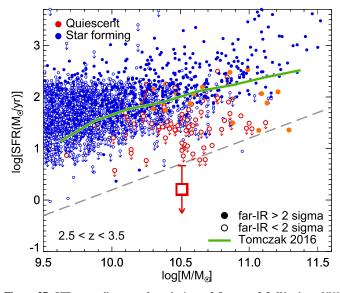


Figure 27. SFR vs. stellar mass for galaxies at 2.5 < z < 3.5. We show *UVJ* selected star-forming galaxies in blue and *UVJ* selected quiescent galaxies in orange/red. Galaxies with $<2\sigma$ measurements in the far-IR are shown with open symbols. Downward pointing arrows are $<1\sigma$ limits on SFR. The large square symbol represents the median SFR_{IR} of 21 undetected quiescent galaxies in CDFS. The thick green line is the median SFR-stellar mass relation for star-forming galaxies from Tomczak et al. (2016), and the gray dashed line represents the criterion for quiescence at z=3, derived using sSFR $<(3t_h)^{-1}$.

from the 68th percentile of the p(z). This galaxy has a strong 4000 Å/Balmer feature, well sampled by the FourStar medium-bandwidth filters. The photometric redshift derived without the use of medium-bandwidth filters in the near-IR, i.e., using only the available broadband groundbased Y, Y, Y, Y and Y, Y or spacebased F125W, F140W and F160W filters, is $Z = 2.87 \pm 0.11$. The galaxy has a broader redshift

probability distribution, p(z), without the FourStar filters, i.e., the redshift is less tightly constrained in the fit.

In Figure 16 we show histograms of the errors, $\sigma_{z, \text{EAZY}} = p68(z)/(1+z)$, with p68(z) the error from the 68th percentile of p(z), in bins between z=0.5 and z=4. This is the redshift region where we expect the impact of the medium-bandwidth filters to be greatest. We show the histograms for a magnitude-limited sample, with $K_s < 25$ in the top row, and with $K_s < 23.5$ in the second row. We also show the histograms of different galaxy types, by splitting up the sample into quiescent and star-forming galaxies, using the UVJ technique (e.g., Whitaker et al. 2011). The star-forming galaxies were additionally split into blue and red by their restframe U-V and V-J colors, which we explain further in Section 7.

The histograms indicate that over a large range in redshift, the errors on the photometric redshifts are smaller if we include the FourStar filters. This holds for all galaxy types. The effect is especially clear around z = 2, and is noticable for higher redshifts as well. For example, at 1.5 < z < 2, the median uncertainty is 40% higher without FourStar filters, with $\sigma_{z,EAZY} = 0.036$ compared $\sigma_{z,EAZY} = 0.025$. Whitaker et al. (2011) find a similar trend with redshift, for the medium bands of NMBS. The peak of the histograms shifts toward higher $\sigma_{z, \text{EAZY}}$ with increasing redshift, up to z = 3, except for blue star-forming galaxies (with blue U - V and V - J colors, see Section 7), for which σ_z , EAZY actually improves. A notable spectral feature for these glaxies is the Lyman Break at rest-frame 912 Å, which is moving through the optical medium-bandwidth filters at this redshift.

5.3. Comparison with Spectroscopic Redshifts

A common comparison in the literature is to compare the photometric redshifts with spectroscopic redshifts. In Figures 17

to 19 we do this, using the compilation of publicly available spectroscopic redshifts in these fields provided by Skelton et al. (2014), with a matching radius of 1". We also included the first release from the MOSDEF survey (Kriek et al. 2015) and the VIMOS Ultra-Deep Survey (Tasca et al. 2016). The overall correspondence is excellent, as indicated by the scatter in the difference between photometric and spectroscopic redshifts. We quantify the errors in the photometric redshifts, σ_z , using the NMAD of $\Delta z/(1+z)$, i.e., 1.48× the median absolute deviation of $|z_{\rm phot}-z_{\rm spec}|/(1+z_{\rm spec})$. In CDFS $\sigma_z=0.010$, in COSMOS $\sigma_z=0.009$ and in UDS $\sigma_z=0.011$. Only a small percentage are outliers, with $\Delta z/(1+z)>0.15$. In CDFS 2.9% are outliers, in COSMOS 2.4% and in UDS 4.5%. At z>1.5 we find $\sigma_z=0.020$, $\sigma_z=0.022$ and $\sigma_z=0.013$ in CDFS, COSMOS and UDS, respectively.

We have also compared with the unpublished redshifts of the ZFIRE survey (Nanayakkara et al. 2016), and find an NMAD of $\sigma_z = \sim 0.02$. Full results are shown in Nanayakkara et al. (2016).

5.4. Redshift Pair Analysis

The drawback of comparing to spectroscopic samples is that these are usually biased toward bright ($K_s < 22$) star-forming galaxies, or unusual sources, such as active galactic nucleus. Therefore these comparisons are not representative of the full photometric catalog and do not allow a careful study of how photometric redshift errors depend on galaxy properties. Here we present an alternative statistical analysis by looking at galaxy pairs. This method was first described and validated by Quadri & Williams (2010). It does not rely on spectroscopic information and can be applied to the full catalogs, including faint sources. Therefore this technique provides us with a more representative photometric redshift uncertainty than possible by comparing to spectroscopic redshifts.

Due to clustering, close pairs of galaxies on the sky have a significant probability of being physically associated, and of lying at the same redshift. Other galaxy pairs will actually be chance projections along the line of sight, but this contamination by random pairs can be accounted for statistically, by randomizing the galaxy positions and repeating the analysis. Each true galaxy pair will give an independent estimate of the true redshift, and we can take the mean of the two values as our best estimate of the true redshift. The distribution of $\Delta z_{\text{pairs}}/(1+z_{\text{mean}})$ of the pairs of galaxies can then be used to estimate the average photometric redshift uncertainties. It is a narrow distribution for robustly derived redshifts, or broader if the redshifts are very uncertain.

For illustration, we show the distributions of $\Delta z_{\rm pairs}/(1+z_{\rm mean})$ in the left panel of Figure 20, for pairs of galaxies with use = 1 and total K_s -band magnitude <23.5, in four redshift bins. The pairs have angular separations between 2".5 and 15". To each distribution we fit a Gaussian and determined the standard deviation. As this is the standard deviation for the redshift differences, we divide by $\sqrt{2}$ to obtain the average redshift uncertainty for individual galaxies, $\sigma_{z,pairs}$, for a particular redshift bin, i.e., $\sigma_{z,pairs}$ is obtained from $\Delta z_{pairs}/(\sqrt{2}(1+z_{\rm mean}))$. In the right panel we show $\sigma_{z,pairs}$ as a function of redshift. $\sigma_{z,pairs}$ increases with redshift, but in general is excellent: varying from 1% to 2% going from z=0.5 to z=2.5. Calculating $\sigma_{z,pairs}$ requires fairly large samples. This partly explains the scatter between results on individual ZFOURGE fields. Other reasons for differences

between the fields are different image filter sets and image depths.

An analysis of $\sigma_{z,pairs}$ can be affected by systematic errors in the photometric redshifts, leading to underestimates of the true redshift uncertainty. For example, because of systematic photometric errors, many sources could be fit with similiar, but wrong, redshifts. This is discussed in more detail by Quadri & Williams (2010). Furthermore, it is important to keep in mind that if all redshifts are systematically overestimated or underestimated, this will not be detected by this method. We tested this scenario by inspecting pairs with at least one spectroscopic redshift available. We derived similar results, indicating that $\sigma_{z,pairs}$ for photometric pairs is not systematically affected. The three fields in the survey also provide constraints, as we make use of different filtersets and systematics introduced between different filters will not be the same in each field. However, if systematics are introduced due to a particular choice of template, this is likely to go unnoticed.

We expect $\sigma_{z,pairs}$ to be sensitive to various parameters, including the type of galaxy, the magnitude and redshift. As a first exploration we will here characterize how our photometric redshift uncertainty depends on these parameters. In the first panel of Figure 21 we show $\sigma_{z,pairs}$ versus z for three different magnitude-limited samples, with $K_s < 22$, $K_s < 23.5$ (as above) and $K_s < 25$. For the brightest galaxies, with $K_s < 22$, the uncertainty is very small, around 1% up to z = 2. However, the uncertainty increases by roughly a factor toward fainter magnitudes up to $K \sim 25$, which is near our completeness limit.

We have additionally investigated the dependence of $\sigma_{z,pairs}$ on galaxy type, using the same UVJ selected samples of quiescent, red star-forming and blue-starforming galaxies as in Section 5.2. The results are shown in the second and third panels of Figure 21. Interestingly, the photometric redshifts of star-forming galaxies and quiescent galaxies are equally well constrained at most redshifts. The exception occurs at intermediate redshift (1.5 < z < 2), where instead we find much smaller redshift uncertainties for quiescent galaxies. This is the redshift range where the Balmer/4000 Å break is moving through the J_1 , J_2 and J_3 medium-bandwidth filters. In contrast, Quadri & Williams (2010) used shallower broadband photometry—with fewer optical filters—and found that quiescent galaxies have significantly better photometric redshifts at all redshifts. This emphasizes that the characteristics of photometric redshifts are dataset-dependent.

Comparing blue and red (dusty) star-forming galaxies, we find that red galaxies have a factor 2–3 worse $\sigma_{z,pairs}$, than do blue galaxies, but these uncertainties are still small: 3%–3.5% at z>2. The redshifts of these galaxies are more difficult to constrain, even with medium band photometry, as they have relatively featureless SEDs, and a degeneracy between redshift and the color of the reddest template allowed in the EAZY set (e.g., Marchesini et al. 2010; Spitler et al. 2014). Here we have split the sample at rest-frame V-J=1, but the effect will be stronger for dustier galaxies at redder V-J. This is a significant issue for star-forming galaxies with high mass or high star formation rates (SFRs), which often tend to be quite dusty.

In the last panel of Figure 21 we compare $\sigma_{z,pairs}$ for the case where we have not included the near-IR medium band FourStar filters in the EAZY fits (but note that two of our three fields still include medium band filters in the optical). For

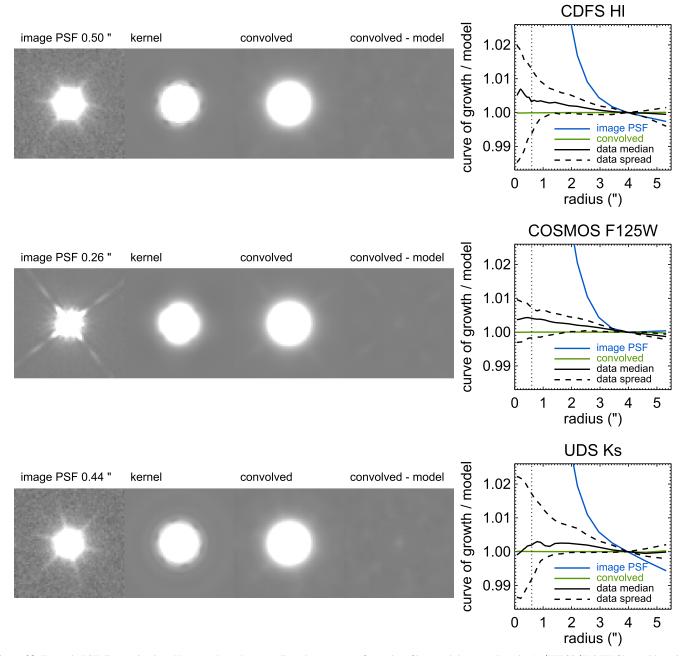


Figure 28. Example PSF diagnostic plots. Here we show the groundbased FourStar/ H_l and K_s filters and the space based HST/WFC3/F125W filter, with various PSF widths. Postage stamps $(10\rlap...65 \times 10\rlap...65)$ from left to right are: the median stacked PSF of the original science images, with their FWHM indicated at the top; a kernel derived using the deconvolution code developed by I. Labbé; the PSF from the leftmost panel, convolved using the corresponding kernel; the convolved PSF minus the Moffat model. The rightmost panels show curves of growth for a number of cases, divided by the model curve of growth and normalized at $r=4\rlap...$ A perfect comparison with the model means the ratio of curves of growth will be one at all radii. In blue lines we show the ratio of the curve of growth of the original PSF (leftmost postage stamp) and the model PSF and in green the same, but after convolving. The black lines represent the median and 1σ scatter of curve of growth ratios for individual point sources in the convolved images that were used for photometry.

the entire range considered here, the photometric redshifts are better derived if we do use the FourStar medium bands. The effect is strongest at z=1.5–2.5 where $\sigma_{z, \mathrm{pairs}}$ using the FourStar medium bands is ~50% smaller compared to $\sigma_{z, \mathrm{pairs}}$ with the FourStar filters removed. This confirms the efficacy of the near-IR medium bands at intermediate to high redshift.

The pairs analysis also provides an interesting way to verify whether the redshift uncertainties that come from the EAZY template fits are reasonable. In the left panel of Figure 22 we compare $\sigma_{z, \text{pairs}}$ to $\sigma_{z, \text{EAZY}}$, and find that they provide

heartening agreement. This figure also shows that σ_z , the uncertainty estimated from comparing the photometric to the spectroscopic redshifts, provides a good estimate of the true uncertainties for the K < 23.5 sample.

Although the redshift uncertainties estimated by EAZY appear quite reliable for the general population of galaxies, we find that they are underestimated by a factor of 2 for dusty star-forming galaxies. In the right panel of Figure 22 we compare $\sigma_{z,pairs}$ to $\sigma_{z,EAZY}$ from the EAZY fits. The pair redshifts of blue star-forming galaxies are better than we expect from the EAZY p(z). However, for red and dusty star-forming galaxies the pair

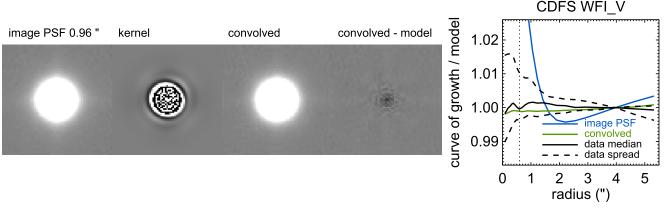


Figure 29. Example of a case where the original PSF was worse than the target PSF. See Figure 28 for a description of the panels. Here the image was deconvolved. A small residual is left after subtracting the model PSF from the deconvolved PSF, but the integrated flux at r = 0.% 6 corresponds well, to within 1%.

redshifts are 50% worse than the EAZY p(z). This effect increases with redshift and toward fainter magnitudes (not shown here). It indicates that with current methods and state-of-the-art surveys, the degeneracy between rest-frame color and redshift for dusty galaxies cannot yet be accurately resolved, and we caution that photometric redshift uncertainties for faint dusty galaxies at z > 1.5 are generally underestimated.

5.5. Redshift Distributions

By improving the accuracy of the photometric redshifts, we can derive improved stellar masses and start identifying large scale structure. In Figure 23 we plot the K_s -band magnitudes as function of z_{peak} (or Z_{spec} where available). The ZFOURGE redshift distributions for each field reveal density peaks corresponding to known overdensities, e.g., at z < 1 in COSMOS (e.g., Kovač et al. 2010; Knobel et al. 2012). These include an overdensity at $z \sim 2.1$, identified by Spitler et al. (2012) using ZFOURGE photometric redshifts. This overdensity was spectroscopically confirmed at a redshift of z = 2.095, with $\sigma_z/(1+z) = \sim 2\%$ (Yuan et al. 2014).

6. STELLAR MASSES AND STAR FORMATION RATES

We estimated 80% mass completeness limits using a method similar to Quadri et al. (2012) (see also Marchesini et al. 2009). We selected galaxies within the range $K_s = 24.0$ –24.7 mag and scaling their fluxes to $K_s = 25.0$ mag. Then we determined the 80th percentile mass rank in narrow redshift bins. Galaxies above this value are the most massive objects that could plausibly fall below the K_s -band selection limit. A smooth function to these values is shown in Figure 24, as function of redshift. At z = 2 we reach a completeness limit of $\sim 10^9 \, M_{\odot}$, and at z = 4 we are complete above $\sim 10^{9.5} \, M_{\odot}$. Beyond z = 4 the completeness limit is extrapolated.

Stellar population properties (stellar mass, SFR, dust extinction, and age) were derived by fitting Bruzual & Charlot (2003) models with FAST (Kriek et al. 2009), assuming a Chabrier (2003) initial mass function, exponentially declining star formation histories with timescale τ , solar metallicity and a dust law as described in Calzetti et al. (2000). For each source the redshift is fixed to the photometric redshift (z_peak) derived with EAZY, or the spectroscopic redshift if known. We limit dust extinction to $0 \le A_V \le 4$, age to $7.5 \le \log_{10}(\text{age}) \le 10.1 \, \text{Gyr}$ and τ to $7 \le \tau \le 11 \, \text{Gyr}$. We

provide the full FAST stellar population catalogs. See Table 8 for an explanation of the catalog header.

SFRs, dust attenuations, ages and star formation histories of galaxies derived from SED fitting to UV, optical and near-IR photometry may be uncertain, especially if galaxies are highly dust-obscured. A different estimate of the SFRs can be obtained by inferring the total infrared luminosity ($L_{\rm IR} \equiv L_{8-1000~\mu m}$) of galaxies and combining this with the luminosity emitted in the UV ($L_{\rm UV}$ at rest-frame 2800 Å). $L_{\rm UV} + L_{\rm IR}$ provides an estimate of the total bolometric luminosity, which can be converted to SFR under the assumption that the galaxy is continuously forming stars (Kennicutt 1998; Bell et al. 2005). In addition to the FAST catalogs, we provide catalogs with the net observed $L_{\rm UV} + L_{\rm IR}$ SFRs (see Table 9 for a description).

We use the conversion from Bell et al. (2005) to calculate SFRs from our data, scaled to a Chabrier (2003) IMF

SFR
$$[M_{\odot} \text{ yr}^{-1}] = 1.09 \times 10^{-10} (L_{IR} + 2.2L_{UV}).$$
 (6)

To derive $L_{\rm IR}$ we use our extracted 24–160 μ m photometry (Section 3.6), to which we fit a model spectral template to calculate the total luminosity. The model template is the averaged template from Wuyts et al. (2008) (hereafter W08), generated by averaging the logarithm of the templates from the library of Dale & Helou (2002). The motivation of this approach was to introduce a simple conversion of flux to luminosity, first proposed by W08 and later validated by Wuyts et al. (2011) and Tomczak et al. (2016). So far public Herschel/PACS data only exists or the CDFS field. For the other fields we fitted the W08 template only to the $Spitzer/MIPS/24~\mu m$ photometry.

Total IR luminosities are then obtained by integrating these fits between 8 and 1000 μm in the rest-frame. $L_{\rm UV}=1.5\nu L_{\nu,2800}$ is the estimated rest-frame 1216–3000 Å UV luminosity, that we derived with EAZY. Both $L_{\rm IR}$ and $L_{\rm UV}$ are in units of L_{\odot} . This conversion assumes that the total IR luminosity reflects the amount of obscured UV light from young stellar populations. Thus by adding its contribution to that of the unobscured UV luminosity ($L_{\rm UV}$) the net star formation rate for galaxies can be measured.

7. FIRST VALIDATION OF THE *UVJ* DIAGRAM AT z = 3

Using EAZY, we derived various rest-frame colors, for example in the Johnson / U and V-bands (Maíz Apellániz 2006),

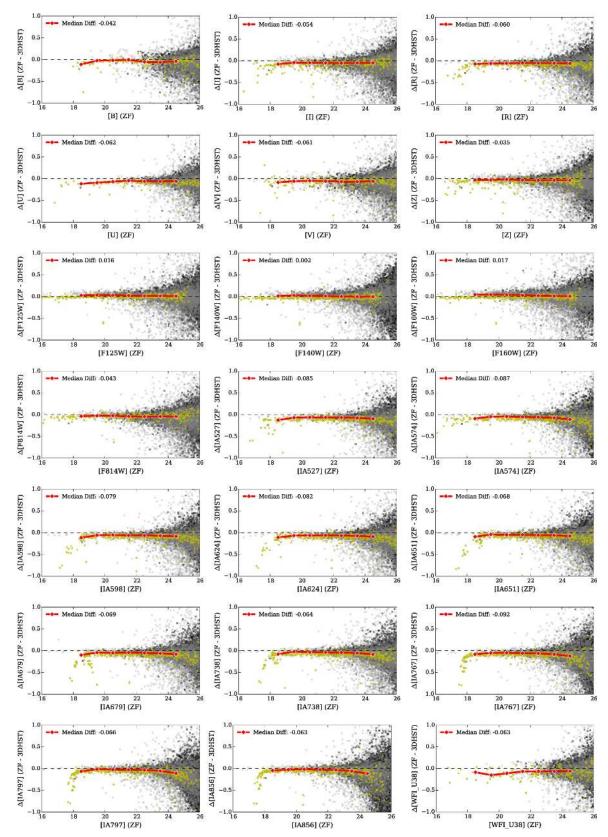


Figure 30. The difference between ZFOURGE and 3D-HST total magnitudes plotted as a function of ZFOURGE magnitudes in CDFS for each band in common. Galaxies with use = 1 are shown as black points, point sources with star = 1 are shown as yellow points, and blended sources with SEflag = 2 (and use = 1) are shown as grey points. The median magnitude difference for all galaxies is shown by the red solid line and large red diamond symbols in bins of 1 mag.

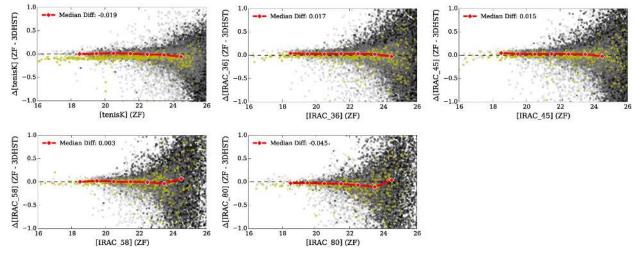


Figure 30. (Continued.)

in the J-filter from the Two Micron All Sky Survey, and at 2800 Å (using a tophat shaped transmission curve). Rest-frame colors were calculated by integrating the redshifted rest-frame filter bandpasses of the best-fit template for each individual source. The process is described in more detail by Brammer et al. (2011), see also Whitaker et al. (2011). Rest-frame 2800 Å luminosity can be used as a proxy of the UV luminosity of a galaxy, which in turn can be used to derive the unobscured SFR (see Section 6). Rest-frame U - V is historically used to distinguish between blue, late-type galaxies with active star formation and red early type galaxies with low star formation. At low redshift these quiescent "red and dead" galaxies dominate at high mass and form a tight "red sequence" versus luminosity or stellar mass (e.g., Tully et al. 1982; Baldry et al. 2004). In Figure 25 (left column) we present the restframe U - V versus stellar mass relation of our mass complete sample, color coded by specific star formation rate (sSFR = SFR/ M^* , with M^* stellar mass).

Clearly, at high redshift there are also large numbers of red massive galaxies with active star formation. These galaxies are red due to dust attenuation (see e.g., Brammer et al. 2009), so a single U - V color cannot be used anymore to separate quiescent and star-forming galaxies. We therefore use a twocolor diagram, rest-frame U-V versus rest-frame V-J(hereafter UVJ), to efficiently separate quiescent from starforming galaxies (see e.g., Labbé et al. 2005; Wuyts et al. 2007; Williams et al. 2009). In Figure 25 (second column from the left), we show the UVJ diagram for the same galaxies. Red star-forming galaxies are now well separated from red quiescent galaxies, and a clear red sequence in UVJ is present up to z = 3.5. At this high redshift a bimodality in color space was also found by (e.g., Whitaker et al. 2011). Here we aim to validate the UVJ selection criterion for such high redshifts by inspecting the sSFRs based on Herschel data.

Several trends with UVJ color can be seen in the data. First, sSFRs show a gradient in color space, such that redder U-V and bluer V-J colors correspond to lower sSFR. This can be interpreted as a stellar "age" gradient. We inspected the stellar ages derived from SED fitting for these galaxies in bins toward redder V-J and bluer U-V, shown in Figure 26. Here the ages are weighted by the timescale τ from the exponentially declining star formation histories in the models to better represent the age of

the bulk of the stellar mass in the galaxy (see also Förster Schreiber et al. 2004). Interestingly, this gradient can even been seen among star-forming galaxies with the highest sSFRS.

Second, star-forming galaxies span a large range in colors due to dust attenuation. This can be seen in the third column of Figure 25, which shows the infrared excess IRX $= \log_{10} L_{\rm IR}/L_{\rm UV}$, the ratio between dust absorbed emission over unattenuated UV emission from star formation. The IRX ranges from IRX ~ -1 for the bluest UVJ colors (colors typically found in dropout selected samples, (e.g., Bouwens et al. 2014) to IRX \sim 3 for the dustiest sources with the reddest U-V and V-J colors. Quiescent galaxies are also characterized by very low $L_{\rm IR}/L_{\rm UV}$ ratios, despite their red U-V colors and faint UV fluxes. SED fitting shows a similar trend with best-fit dust attenuation A_V . Galaxies at the tip of the red star-forming sequence also appear to have lower sSFRs, so the colors of the very reddest galaxies appear to be a combination of dust and old age (see also Forrest et al. 2016; Fumagalli et al. 2016; A. Bedregal 2016, in preparation). However, we caution that these galaxies are also expected to have the least accurate photometric redshifts (Section 5.4).

Finally, there is a clear trend with stellar mass, such that massive galaxies tend to be redder. Up to the highest redshifts probed here $(z \sim 3.5)$ massive $M > 10^{10.5} M_{\odot}$ galaxies are predominantly quiescent and/or dusty (see e.g., Spitler et al. 2014; Straatman et al. 2014). This is not a selection effect, as only galaxies are shown above our mass completeness limit. The U-V versus stellar mass (left column) shows an absence of low mass quiescent galaxies at z=2-3, with massive galaxies quenching first, and lower mass quiescent galaxies rapidly building up from z=2 to z=0.5 (see Tomczak et al. 2014).

The usage of the UVJ technique to identify quiescent galaxies has been validated to $z \lesssim 2.0$, by inspecting stacked WFC3 grism spectra (Whitaker et al. 2013) and stacked MIPS 24 μ m data (Fumagalli et al. 2014). However, verification at $z \gtrsim 3$ has proved difficult. Straatman et al. (2014) and Spitler et al. (2014) identified quiescent galaxies in the UVJ diagram at even at higher redshifts, to $z \sim 4$, but at these extreme redshifts the IR observations are not deep enough, and the samples are too small, to rule out significant obscured star formation. Here, we place the strongest constraints yet on the SFRs of quiescent galaxies at $z \sim 3$ by studying the infrared fluxes from ultradeep

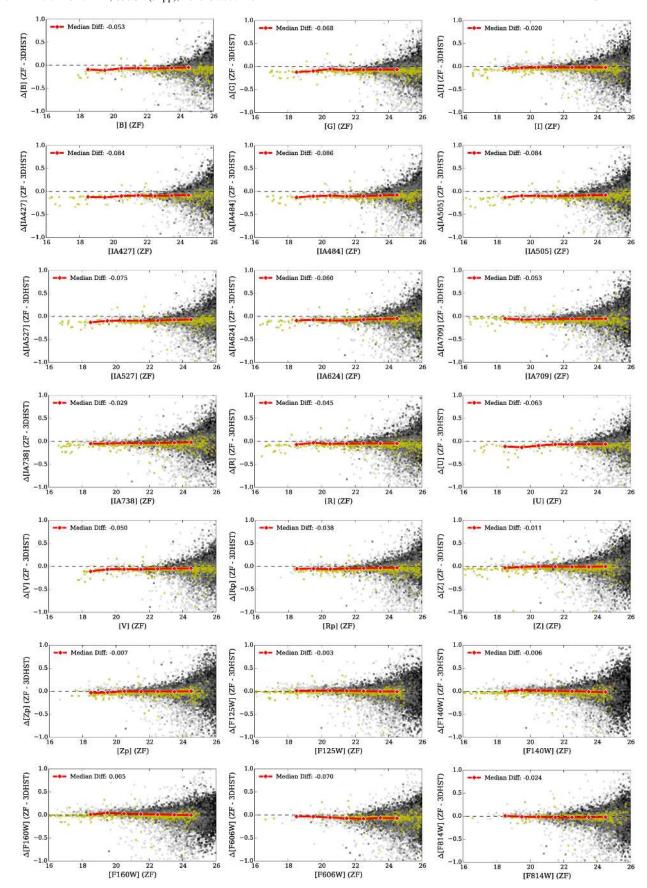
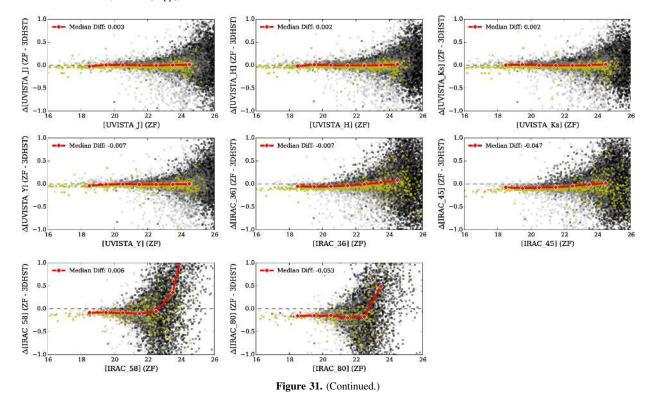


Figure 31. The difference between ZFOURGE and 3D-HST total magnitudes plotted as a function of ZFOURGE magnitudes in COSMOS for each band in common. Symbols are the same as in Figure 30.



PACS Herschel data from GOODS-Herschel and PEP. Using the net UV+IR SFR based on the ultradeep Herschel data, we show the SFR versus stellar mass diagram in Figure 27. The individual quiescent galaxies at 2.5 < z < 3.5 are shown in red

and orange. 80% of the quiescent galaxies are not detected in

the far-IR at $<2\sigma$ (open red symbols).

To confirm a galaxy as quiescent based on sSFR, we adopt the definition of Damen et al. (2009) which states that a galaxy is quiescent if their sSFR is $<(3t_h)^{-1}$, with t_h the Hubble time. At z = 3 this limit corresponds to a sSFR $< 1.6 \times 10^{-10} \text{ yr}^{-1}$. The observed constraints on the sSFR of individual galaxies are not very strong, due to the limited sensitivity in the IR at such high redshift. At z = 3 our best estimate of the SFR comes from the combination of ultradeep Spitzer/MIPS 24 µm and Herschel/Pacs/100 and 160 μ m data in the CDFS (see Section 6). The 1σ limiting SFR for individual sources is SFR_{IR} < 15 M_{\odot} yr⁻¹. This means that for galaxies with stellar mass $5 \times 10^{10} M_{\odot}$ at 2.5 < z < 3.5, we can constrain their individual sSFRs to be $<3.0 \times 10^{-10}$ yr⁻¹. To place firmer constraints, we follow the procedure of Straatman et al. (2014) and stack the 24–160 µm fluxes of UVJ-selected quiescent galaxies, limiting our sample to $M > 10^{10} M_{\odot}$, to gain more sensitivity in the far-IR. We calculate the median of the stack of 21 IR undetected ($<2\sigma$) galaxies in CDFS, finding a SFR_{IR} = $1.6 \pm 3.1 \, M_{\odot} \, \text{yr}^{-1}$, where the errors are derived by bootstrapping. At a mean stellar mass of $3.2 \times 10^{10} M_{\odot}$, this translates into a sSFR = $0.5 \pm 1.0 \times 10^{-10} \, \mathrm{yr}^{-1}$, easily meeting the Damen et al. (2009) criterion for quiescence.

This shows that UVJ selected quiescent galaxies at $z \sim 3$ indeed harbour very low levels of star formation, a factor of $>15\times$ lower at 95% confidence than the median of star-forming galaxies at the same redshift (Tomczak et al. 2016). Quenching of star formation was thus very efficient, even at such early times.

8. SUMMARY

In this paper we have presented the data products and public release of ZFOURGE, a near-IR galaxy survey with the FourStar instrument aimed to increase our understanding of the evolution of galaxies at intermediate to high redshift. We made use of five medium-bandwidth near-IR filters $(J_1, J_2, J_3, H_s, H_l)$, and ultra-deep K_s -band detection images, obtained by co-adding the FourStar/ K_s -band images with publicly available K_s -band data. The medium-bandwidth filters are particularly well suited to constrain the photometric redshifts of sources in the redshift range 1.5 < z < 3.5.

The K_s -band detection images have 5σ depths between 25.5 and 26.5 AB mag, and we detected >70,000 galaxies in the three extragalactic fields CDFS, COSMOS and UDS. The ZFOURGE FourStar observations were augmented with images over a large range in wavelength, $0.3-160\,\mu\text{m}$, and fluxes were consistently derived using accurate PSF modeling and deblending methods. The ZFOURGE data release comprises source catalogs with photometry, photometric redshifts and stellar population properties, such as stellar mass derived with SED modeling, and SFRs derived from UV and IR luminosities.

Photometric redshifts were tested against spectroscopic redshifts from the literature, resulting in $\sigma_z=0.010$ in CDFS, $\sigma_z=0.009$ in COSMOS and $\sigma_z=0.011$ in UDS. As spectroscopic samples of galaxies are often biased toward bright and blue galaxies, we performed another, independent test of the robustness of the photometric redshifts, by inspecting galaxy pairs. We found excellent results, with $\sigma_z=0.01$ –0.02 for a K<23.5 mag limited sample, between z=0.5 and z=2.5, and showed that photometric redshifts are better constrained by 50% if the near-IR FourStar medium-bandwidth filters are included, compared to SED fitting with the FourStar filters removed.

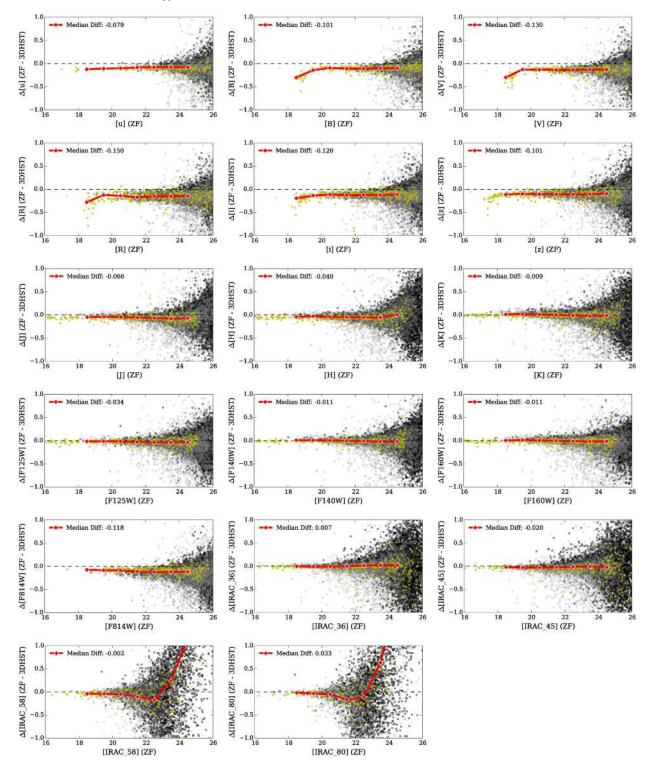


Figure 32. The difference between ZFOURGE and 3D-HST total magnitudes plotted as a function of ZFOURGE magnitudes in UDS for each band in common. Symbols are the same as in Figure 30.

We investigated the efficacy of the *UVJ* diagram to classify galaxies beyond z=2 and illustrated how *UVJ* colors correlate with sSFR and infrared luminosity excess (dust attenuation). Using the *UVJ* diagram, we selected a sample of quiescent galaxies at 2.5 < z < 3.5 and investigated their sSFR properties. We confirmed that these were indeed quiescent, with an average sSFR = $0.5 \pm 1.0 \times 10^{-10}$ and $>15 \times$ supressed SFRs, relative to the average stellar mass versus SFR relation

of star-forming galaxies, thereby for the first time validating the UVJ classification to z=3.5.

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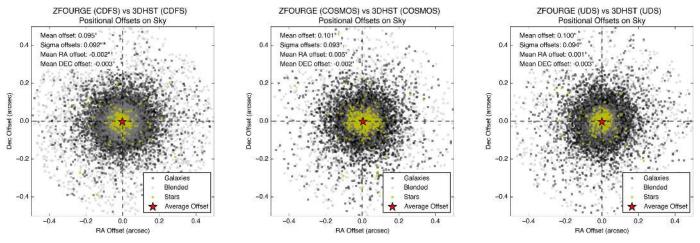


Figure 33. Positional offsets between source locations in ZFOURGE and 3D-HST, using the same symbols as in Figures 31 to 32 (Cont.). The median offsets are indicated by red stars.

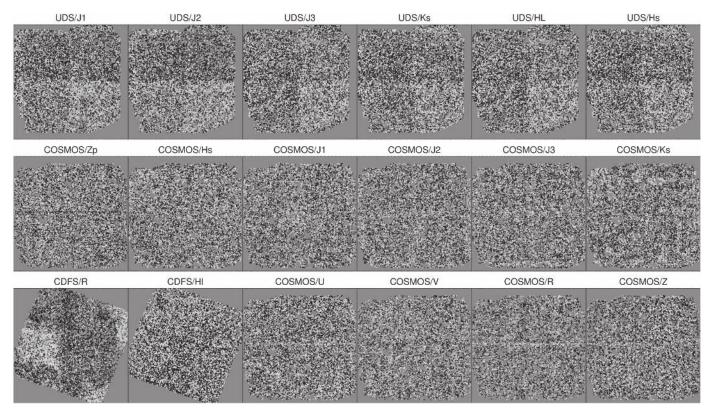


Figure 34. Example spatial zeropoint residual maps of after subtracting the 2-D polynomial fit. The grayscale ranges from 0.95 to 1.05, i.e., a 5% flux deviation. CDFS/R had a complicated structure with large zeropoint variations, even after corrections. We resolved this by adding a 5% systematic uncertainty as an error floor when fitting photometry.

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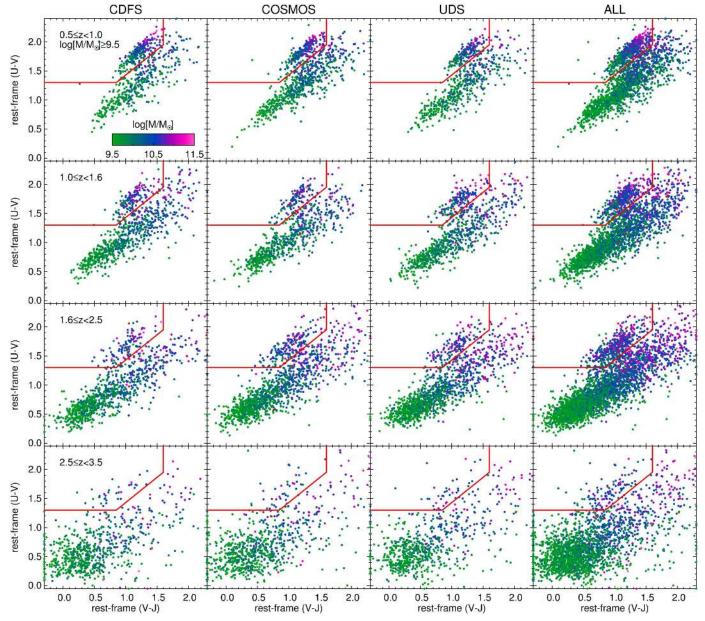


Figure 35. Rest-frame U - V vs. V - J diagrams of galaxies with use = 1, $S/R_{K_s} > 10$ and stellar mass $M > 10^{9.5} M_{\odot}$. In the first three columns we show the three ZFOURGE fields. In the last column these are combined. From top to bottom we show bins of increasing redshift. The color scaling indicates stellar mass. The rest-frame U - V and V - J colors in each field show the same pattern, and the same location for the red sequence, indicating consistent photometry.

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APPENDIX A PSF CONVOLUTION

In Section 3.1 we have explained how images are convolved, such that their average point source profile matches a FWHM = 0.9 Moffat PSF, with the goal of obtaining consistent aperture photometry over all filters. Here we show example diagnostic plots (Figure 28) of the kernel derivation. The four columns represent the original image PSF, the kernel used for convolution, the PSF after convolving with this kernel and the residual after subtracting the model Moffat profile. curves of growth representing the PSFs before and after convolution are shown in the rightmost panels. The green curve in particular shows the

light profile of the convolved PSF divided by that of the model, which is close to unity. Whereas in Figure 9 we show the curves of growth measured on the median PSF derived from the convolved images, we inspect here the median and 1σ scatter of the curves of growth of individual stars, finding similar residuals of <2% compared with the model PSFs.

Mathematically, we are able to produce PSF matched images in case of *deconvolution*, when the image has a worse quality, i.e., a broader PSF, than the target FWHM = $0.^{\prime\prime}9$ Moffat profile. We have applied deconvolution in some cases with care, as we cannot make any assumptions on the true underlying light profile of galaxies. We find that the strongest residuals occur once the image PSF becomes much larger than the target PSF. After inspecting the residuals by eye and taking into account the aperture radius with a diameter of $1.^{\prime\prime}2$, we

include images that have up to 15% broader PSFs than the target PSF. We show an example in Figure 29. While the deconvolved PSF shows some residual compared to the model, the curves of growth indicate that we capture the same amount of light within 1% at r=0.%, the aperture radius that we use to derive the catalog fluxes with. In total 11 of 92 UV to near-IR images were deconvolved.

APPENDIX B COMPARISON TO THE 3D-HST PHOTOMETRIC CATALOGS

In this section we compare of the total magnitudes measured by ZFOURGE and those measured by 3D-HST, who make use of many of the same ancillary images. They also largely use the same methods to derive photometry. For each filter in common we calculate the difference in magnitude between crossmatched sources and show this versus total magnitude as per the ZFOURGE catalogs in Figures 30 and 32. For crossmatching we used a maximum angular separation of 1". We separately indicated sources that were flagged as possibly blended or contaminated by neighbors by SE (SEflags \geqslant 2). \triangle mag has somewhat more scatter for these sources at faint magnitudes, but overall the correspondence is quite good between the surveys, with \triangle mag close to zero. The most notable exceptions are the V, R, i and z-bands in UDS, which tend to be \sim 0.1 magnitudes brighter in ZFOURGE.

In Figure 33 we show the positional offsets between source locations in ZFOURGE and the 3D-HST survey. The median offsets are ≤ 0.005 in R.A. or decl. in all fields, indicating the images are uniformly calibrated and can be reliably used for inter-survey comparisons.

APPENDIX C SPATIAL VARIATION IN THE ZEROPOINTS

In Figure 34 we show example spatial zeropoint residual maps derived with EAZY, by comparing the best-fit templates to the observed galaxy SEDs. The residuals are of the order of <5%. We fitted a two-dimensional polynomial to each offset map and used these to derive a correction to the flux of a specific filter, for all sources as a function of their x- and y-position. This was done for the full dataset. With this method we were able to trace systematic offsets of $\pm \sim 4\%$ between the four FourStar detectors in UDS and correct these to $\pm \sim 2\%$. In one image strong, nonlinear spatial effects stand out: this is in VIMOS/R in CDFS. In this image the strong spatial varation could not be removed by this first order correction. Due to its large depth, we have kept the image in our sample, but applied a minimum error floor to the V-band flux of 5% to take into account uncertainties on the zeropoint of the image.

APPENDIX D UVJ DIAGRAM FIELD COMPARISON

In this section we show the *UVJ* diagram color-coded by stellar mass (Figure 35), as in the third column of Figure 25. We show the same redshift bins, but split the diagrams into the three ZFOURGE fields. By comparing the rest-frame colors in the different fields we can look for inconsistencies, for example if the median locii of the datapoints are offset relative to each other. Here this is not the case, indicating consistent photometry.

REFERENCES

```
Allen, R. J., Kacprzak, G. G., Spitler, L. R., et al. 2015, ApJ, 806, 3
Ashby, M. L. N., Willner, S. P., Fazio, G. G., et al. 2013, ApJ, 769, 80
Baldry, I. K., Glazebrook, K., Brinkmann, J., et al. 2004, ApJ, 600, 681
Bell, E. F., Papovich, C., Wolf, C., et al. 2005, ApJ, 625, 23
Bell, E. F., van der Wel, A., Papovich, C., et al. 2012, ApJ, 753, 167
Bertin, E., & Arnouts, S. 1996, A&AS, 117, 393
Bouwens, R. J., Illingworth, G. D., Oesch, P. A., et al. 2014, ApJ, 793, 115
Brammer, G. B., van Dokkum, P. G., & Coppi, P. 2008, ApJ, 686, 1503
Brammer, G. B., van Dokkum, P. G., Franx, M., et al. 2012, ApJS, 200, 13
Brammer, G. B., Whitaker, K. E., van Dokkum, P. G., et al. 2009, ApJL,
Brammer, G. B., Whitaker, K. E., van Dokkum, P. G., et al. 2011, ApJ, 739, 24
Bruzual, G., & Charlot, S. 2003, MNRAS, 344, 1000
Calzetti, D., Armus, L., Bohlin, R. C., et al. 2000, ApJ, 533, 682
Cardamone, C. N., van Dokkum, P. G., Urry, C. M., et al. 2010, ApJS,
   189, 270
Chabrier, G. 2003, PASP, 115, 763
Chen, H.-W., Marzke, R. O., McCarthy, P. J., et al. 2003, ApJ, 586, 745
Cooper, M. C., Griffith, R. L., Newman, J. A., et al. 2012, MNRAS, 419, 3018
Daddi, E., Renzini, A., Pirzkal, N., et al. 2005, ApJ, 626, 680
Dale, D. A., & Helou, G. 2002, ApJ, 576, 159
Damen, M., Labbé, I., Franx, M., et al. 2009, ApJ, 690, 937
Dickinson, M., Giavalisco, M., & GOODS Team 2003, in The Mass of
   Galaxies at Low and High Redshift, ed. R. Bender & A. Renzini (Berlin:
   Springer), 324
Elbaz, D., Dickinson, M., Hwang, H. S., et al. 2011, A&A, 533, A119
Erb, D. K., Pettini, M., Shapley, A. E., et al. 2010, ApJ, 719, 1168
Erben, T., Hildebrandt, H., Lerchster, M., et al. 2009, A&A, 493, 1197
Erben, T., Schirmer, M., Dietrich, J. P., et al. 2005, AN, 326, 432
Faber, S. M., Willmer, C. N. A., Wolf, C., et al. 2007, ApJ, 665, 265
Fioc, M., & Rocca-Volmerange, B. 1999, arXiv:astro-ph/9912179
Fontana, A., Dunlop, J. S., Paris, D., et al. 2014, A&A, 570, A11
Forrest, B., Tran, K.-V. H., Tomczak, A. R., et al. 2016, ApJL, 818, L26
Förster Schreiber, N. M., van Dokkum, P. G., Franx, M., et al. 2004, ApJ,
   616, 40
Franx, M., van Dokkum, P. G., Schreiber, N. M. F., et al. 2008, ApJ, 688, 770
Fumagalli, M., Franx, M., van Dokkum, P., et al. 2016, ApJ, 822, 1
Fumagalli, M., Labbé, I., Patel, S. G., et al. 2014, ApJ, 796, 35
Furusawa, H., Kosugi, G., Akiyama, M., et al. 2008, in ASP Conf. Ser. 399,
   Panoramic Views of Galaxy Formation and Evolution, ed. T. Kodama,
   T. Yamada, & K. Aoki (San Francisco, CA: ASP), 131
Giacconi, R., Zirm, A., Wang, J., et al. 2002, ApJS, 139, 369
Giavalisco, M., Ferguson, H. C., Koekemoer, A. M., et al. 2004, ApJL,
   600, L93
Grogin, N. A., Kocevski, D. D., Faber, S. M., et al. 2011, ApJS, 197, 35
Hildebrandt, H., Erben, T., Dietrich, J. P., et al. 2006, A&A, 452, 1121
Hildebrandt, H., Pielorz, J., Erben, T., et al. 2009, A&A, 498, 725
Hsieh, B.-C., Wang, W.-H., Hsieh, C.-C., et al. 2012, ApJS, 203, 23
Kawinwanichakij, L., Papovich, C., Quadri, R. F., et al. 2014, ApJ, 792, 103
Kennicutt, R. C., Jr. 1998, ARA&A, 36, 189
Knobel, C., Lilly, S. J., Iovino, A., et al. 2012, ApJ, 753, 121
Koekemoer, A. M., Faber, S. M., Ferguson, H. C., et al. 2011, ApJS, 197, 36
Kovač, K., Lilly, S. J., Knobel, C., et al. 2010, ApJ, 718, 86
Kriek, M., van der Wel, A., van Dokkum, P. G., Franx, M., &
   Illingworth, G. D. 2008, ApJ, 682, 896
Kriek, M., van Dokkum, P. G., Franx, M., et al. 2006, ApJL, 649, L71
Kriek, M., van Dokkum, P. G., Labbé, I., et al. 2009, ApJ, 700, 221
Kriek, M., van Dokkum, P. G., Whitaker, K. E., et al. 2011, ApJ, 743, 168
Kriek, M., Shapley, A. E., Reddy, N. A., et al. 2015, ApJS, 218, 15
Kron, R. G. 1980, ApJS, 43, 305
Kroupa, P. 2001, MNRAS, 322, 231
Labbé, I., Bouwens, R., Illingworth, G. D., & Franx, M. 2006, ApJL, 649, L67
Labbé, I., Franx, M., Rudnick, G., et al. 2003a, AJ, 125, 1107
Labbé, I., Huang, J., Franx, M., et al. 2005, ApJL, 624, L81
Labbé, I., Rudnick, G., Franx, M., et al. 2003b, ApJL, 591, L95
Lawrence, A., Warren, S. J., Almaini, O., et al. 2007, MNRAS, 379, 1599
Lee, J. C., Ly, C., Spitler, L., et al. 2012, PASP, 124, 782
Lord, S. D., Hollenbach, D. J., Colgan, S. W. J., et al. 1992, BAAS, 24, 1182
Madau, P., Pozzetti, L., & Dickinson, M. 1998, ApJ, 498, 106
Magnelli, B., Popesso, P., Berta, S., et al. 2013, A&A, 553, A132
Maíz Apellániz, J. 2006, AJ, 131, 1184
Maraston, C. 2005, MNRAS, 362, 799
Marchesini, D., van Dokkum, P. G., Förster Schreiber, N. M., et al. 2009, ApJ,
   701, 1765
```

```
Marchesini, D., Whitaker, K. E., Brammer, G., et al. 2010, ApJ, 725, 1277
McCracken, H. J., Milvang-Jensen, B., Dunlop, J., et al. 2012, A&A,
   544, A156
Moffat, A. F. J. 1969, A&A, 3, 455
Muzzin, A., Marchesini, D., Stefanon, M., et al. 2013a, ApJS, 206, 8
Muzzin, A., Marchesini, D., Stefanon, M., et al. 2013b, ApJ, 777, 18
Nanayakkara, T., Glazebrook, K., Kacprzak, G. G., et al. 2016, arXiv:1607.
Noeske, K. G., Weiner, B. J., Faber, S. M., et al. 2007, ApJL, 660, L43
Nonino, M., Dickinson, M., Rosati, P., et al. 2009, ApJS, 183, 244
Oke, J. B., Cohen, J. G., Carr, M., et al. 1995, PASP, 107, 375
Papovich, C., Labbé, I., Quadri, R., et al. 2015, ApJ, 803, 26
Papovich, C., Momcheva, I., Willmer, C. N. A., et al. 2010, ApJ, 716, 1503
Peng, C. Y., Ho, L. C., Impey, C. D., & Rix, H.-W. 2010a, AJ, 139, 2097
Peng, Y.-j., Lilly, S. J., Kovač, K., et al. 2010b, ApJ, 721, 193
Pérez-González, P. G., Rieke, G. H., Villar, V., et al. 2008, ApJ, 675, 234
Persson, S. E., Murphy, D. C., Smee, S., et al. 2013, PASP, 125, 654
Postman, M., Franx, M., Cross, N. J. G., et al. 2005, ApJ, 623, 721
Quadri, R., Marchesini, D., van Dokkum, P., et al. 2007, AJ, 134, 1103
Quadri, R. F., & Williams, R. J. 2010, ApJ, 725, 794
Quadri, R. F., Williams, R. J., Franx, M., & Hildebrandt, H. 2012, ApJ, 744, 88
Retzlaff, J., Rosati, P., Dickinson, M., et al. 2010, A&A, 511, A50
Sanders, D. B., Salvato, M., Aussel, H., et al. 2007, ApJS, 172, 86
Schlafly, E. F., & Finkbeiner, D. P. 2011, ApJ, 737, 103
Scoville, N., Aussel, H., Brusa, M., et al. 2007, ApJS, 172, 1
Skelton, R. E., Whitaker, K. E., Momcheva, I. G., et al. 2014, ApJS, 214, 24
Sobral, D., Smail, I., Best, P. N., et al. 2013, MNRAS, 428, 1128
```

```
Spitler, L. R., Labbé, I., Glazebrook, K., et al. 2012, ApJL, 748, L21
Spitler, L. R., Straatman, C. M. S., Labbé, I., et al. 2014, ApJL, 787, L36
Straatman, C. M. S., Labbé, I., Spitler, L. R., et al. 2014, ApJL, 783, L14
Taniguchi, Y., Scoville, N., Murayama, T., et al. 2007, ApJS, 172, 9
Tasca, L. A. M., Le Fevre, O., Ribeiro, B., et al. 2016, arXiv:1602.01842
Tilvi, V., Papovich, C., Tran, K.-V. H., et al. 2013, ApJ, 768, 56
Tomczak, A. R., Quadri, R. F., Tran, K.-V. H., et al. 2014, ApJ, 783, 85
Tomczak, A. R., Quadri, R. F., Tran, K.-V. H., et al. 2016, ApJ, 817, 118
Tully, R. B., Mould, J. R., & Aaronson, M. 1982, ApJ, 257, 527
van der Wel, A., Bell, E. F., Häussler, B., et al. 2012, ApJS, 203, 24
van der Wel, A., Franx, M., van Dokkum, P. G., et al. 2014, ApJ, 788, 28
van Dokkum, P. G., Franx, M., Kriek, M., et al. 2008, ApJL, 677, L5
van Dokkum, P. G., Labbé, I., Marchesini, D., et al. 2009, PASP, 121, 2
Whitaker, K. E., Labbé, I., van Dokkum, P. G., et al. 2011, ApJ, 735, 86
Whitaker, K. E., van Dokkum, P. G., Brammer, G., & Franx, M. 2012, ApJL,
  754, L29
Whitaker, K. E., van Dokkum, P. G., Brammer, G., et al. 2013, ApJL, 770, L39
Williams, R. J., Quadri, R. F., Franx, M., van Dokkum, P., & Labbé, I. 2009,
    ApJ, 691, 1879
Windhorst, R. A., Cohen, S. H., Hathi, N. P., et al. 2011, ApJS, 193, 27
Wolf, C., Meisenheimer, K., Kleinheinrich, M., et al. 2004, A&A, 421, 913
Wuyts, S., Labbé, I., Franx, M., et al. 2007, ApJ, 655, 51
Wuyts, S., Labbé, I., Schreiber, N. M. F., et al. 2008, ApJ, 682, 985
Wuyts, S., Förster Schreiber, N. M., Genzel, R., et al. 2012, ApJ, 753, 114
Wuyts, S., Förster Schreiber, N. M., Lutz, D., et al. 2011, ApJ, 738, 106
York, D. G., Adelman, J., Anderson, J. E., Jr., et al. 2000, AJ, 120, 1579
Yuan, T., Nanayakkara, T., Kacprzak, G. G., et al. 2014, ApJL, 795, L20
```