# THE OPTICAL/INFRARED COUNTERPART(S) OF IRAS 18333-2357 

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#### Abstract

Observations of the potential optical counterparts of the unusual source IRAS 18333-2357 show that this source is associated with an extraordinary planetary nebula system in the galactic globular cluster M22.

Three distinct optical objects were found within $2^{\prime \prime}$ of the IRAS 18333-2357 position as determined by precisely locating the $20 \mu \mathrm{~m}$ infrared source. One object is a red star with $m_{v} \approx 14.7 \mathrm{mag}$, which appears to be an unrelated background field star that is possibly significantly reddened beyond the line-of-sight reddening to M22. The second stellar object is a very blue star with $m_{v} \approx 14.3 \mathrm{mag}$ located about 1.3 south of the red star. Absorption lines of He II and possibly H are present in $4000-5000 \AA$ spectra of the stellar pair, similar to spectra of planetary nebula nuclei.

The third member of this optical triple is an extended emission line nebulosity approximately $10^{\prime \prime} \times 7^{\prime \prime}$ in size, centered about $1^{\prime \prime}$ east and south of the red star. The ionized gas in this nebulosity is extraordinarily oxygen-rich and neon-rich relative to both hydrogen and helium compared to the atmospheres of M22 red giants and is substantially oxygen-rich and neon-rich relative to hydrogen in comparison with typical planetary nebulae. This nebulosity is almost certainly in M22. The blue star is also very likely to be a member of M22, the source of ionizing photons for the nebulosity and probably the luminosity source for IRAS 18333-2357.

We suggest that the dust responsible for the strong infrared emission of IRAS $18333-2357$ is physically associated with the M22 nebulosity. In this case the total nebular mass, comprised of $3-10 \times 10^{-4} M_{\odot}$ of ionized gas plus $>6 \times 10^{-4} M_{\odot}$ of silicate or carbon-based grains, is possibly dominated by the dust component. The $\mathrm{O}, \mathrm{Mg}, \mathrm{Si}$ abundances in the case of silicate grains, or carbon in the case of carbon-based grains, may be enhanced relative to hydrogen by at least a factor of 1000 compared to solar abundances.

The relative abundances and mass of the M22 nebula are very unusual among known planetary nebulae. It is speculated that this system in M22 may be the result of the interaction within a close binary system containing at least one O-Ne white dwarf component, or perhaps related to planetary nebulae like A30 and A78. Subject headings: clusters: globular - infrared: sources - nebulae: planetary - stars: individual (IRAS 18333-2357)


## I. INTRODUCTION

The IRAS Point Source Catalog (1985, hereafter PSC) includes a bright source with unusual colors, IRAS 18333-2357, located about $1^{\prime}$ from the center of the metalpoor Galactic globular cluster M22 (NGC 6566). In the initial assessment of this source (Gillett et al. 1986, hereafter Paper I), its nature and evolutionary status were necessarily left uncertain. On the one hand, its infrared energy distribution, approximating that of a 100 K blackbody between 12 and $100 \mu \mathrm{~m}$, is quite rare in the PSC but similar to that of many planetary nebulae and bipolar nebulae (e.g., OH 0739-14) which may be precursors to planetary nebulae. On the other hand, the lack of an obvious optical or radio counterpart to IRAS 18333-2357 ruled out a typical planetary nebula. It was speculated in Paper I that this unusual object might be in a very early, short-lived stage of planetary nebulae evolution. Even the question of cluster membership was left unresolved; while the probability of a chance alignment of an infrared source similar to IRAS 18333-2357 within $1^{\prime}$ of the center of any Galactic globular

[^0]cluster was very low, $1.3 \times 10^{-3}$, this possibility could not be ruled out altogether.

There is only one known example of a planetary nebula in a globular cluster, K648 in M15; thus the confirmed presence in a globular cluster of such an unusual object as IRAS 18333-2357 is extremely rare, providing a well-understood distance, environment, initial composition and age for the further study and understanding of the phenomena involved.

This paper presents the initial ground-based assessment of IRAS 18333-2357. The observations, described in § II, were intended to address the issues of cluster membership and the possible evolutionary state and nature of the infrared source by precisely locating the infrared source, identifying potential optical counterpart(s), and investigating the optical/infrared properties of the potential counterparts.

A summary of the results of this study is given in § III. There are three distinct optical objects located within $\approx 2^{\prime \prime}$ of the infrared source: a red star, a very blue star, and an extended emission line nebulosity. Discussion of the observations of these three sources, the possible physical interrelationships between these objects and IRAS 18333-2357, and their potential association with M22 are presented in § IV.


FIG. 1.-(a) $g$ band CCD frame of region around IRAS 18333-2357. The ellipse indicates the nominal PSC position and position uncertainties of IRAS 18333-2357. The $20 \mu \mathrm{~m}$ source identified with IRAS 18333-2357 is located within $2^{\prime \prime}$ of the marked object just to the west and north of the nominal PSC position. Note the north-south elongation of this object in the $g$ band. The center of M22 is approximately $1^{\prime}$ north of the nominal IRAS $18333-2357$ position. The coordinates of M22-V8 and the $20 \mu \mathrm{~m}$ source identified with IRAS 18333-2357 are given in Table 1. The star just to the north of the stellar pair is 8.5 from the southern (blue) component. (b) $U$ band CCD frame of the region around IRAS 18333-2357. Note that the image of the potential optical counterpart is not elongated at this wavelength and is the brightest object in the field.

Speculations concerning the nature and origin of this combination of sources are presented in § V. It appears that IRAS 18333-2357 is indeed physically associated with M22, and while it probably should be considered a planetary nebula, its very small nebular mass and extreme abundance anomalies are very unusual among known planetary nebulae. IRAS 18333-2357 does not appear to be in an early stage of planetary nebula evolution, but instead may be in a late stage of planetary nebula evolution, with the lack of associated radio or $\mathrm{H} \alpha$ source the result of the abundance anomalies in the source.

## II. OBSERVATIONS

The ground-based observations of IRAS 18333-2357 were carried out in two stages: the initial step was a precise location of the infrared source and identification of potential optical
counterpart(s). Because IRAS $18333-2357$ is extremely red, with $f_{v}(25 \mu \mathrm{~m}) / f_{v}(12 \mu \mathrm{~m})>10$, the ground-based search was carried out at $20 \mu \mathrm{~m}$, where IRAS 18333-2357 is the dominant infrared source in the direction of M22.

A search on 1986 May 3, using the bolometer detector system with a $6^{\prime \prime}$ diameter beam on the NASA IRTF telescope, discovered an obvious $20 \mu \mathrm{~m}$ source coincident to within $\pm 2^{\prime \prime}$ with the optical object marked in Figure 1. The $20 \mu \mathrm{~m}$ source is very close to the nominal IRAS position, well within the PSC positional uncertainties, and the apparent optical counterpart is the closest significant optical source to the nominal IRAS position. We identify this source with IRAS 18333-2357, even though the measured 10 and $20 \mu \mathrm{~m}$ brightnesses are significantly less than those expected on the basis of the IRAS observations (see § IIc for further discussion). The coordinates of this


Fig. $1 b$
object, listed in Table 1, were determined by offsetting from a nearby SAO star and M22-V8. While the nearby cluster variable V8 and the field star V-1 were found to be comparable in brightness to IRAS $18333-2357$ at $10 \mu \mathrm{~m}$, neither was detectable at $20 \mu \mathrm{~m}$ at a $3 \sigma$ flux density one-fourth that of the source we associate with IRAS 18333-2357.

Visual inspection of this field under good seeing conditions clearly shows that the bulk of the optical emission emanates from two stars (hereafter designated $\mathrm{N}_{*}$ and $\mathrm{S}_{*}$ ) separated by about $1^{\prime \prime}$ in the north-south direction. In addition to the two
stars, the early spectroscopic observations of this region revealed a third source, an extended emission line region closely aligned with the stellar pair. The remainder of this section describes the preliminary observations of these three candidate optical counterparts to IRAS 18333-2357.

## a) Photometry

The results of conventional infrared photometry from 1 to $20 \mu \mathrm{~m}$, centered on the stellar pair and including both stellar components, carried out using the IRTF, the Palomar 5 m , and

TABLE 1
Position Measurements

| Source | Position |  | Uncertainties |  | References |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | R.A.(1950) | Decl.(1950) | R.A. | Decl. |  |
| M22 center | $18^{\mathrm{h}} 33^{\mathrm{m}} 20^{\text {s }} 4$ | $-23^{\circ} 56^{\prime} 56^{\prime \prime}$ | $20^{\prime \prime} \times$ |  | 1 |
| M22-V8 | 183317.55 | -23 5757.5 | $1^{\prime \prime} \times$ | $1^{\prime \prime}$ | 2 |
| IRAS 18333-2357 ........ | 183320.3 | -23 5752 | $27^{\prime \prime} \times$ | $6^{\prime \prime}(2 \sigma)$ | 3 |
| $20 \mu \mathrm{~m}$ source . . . $\ldots \ldots \ldots \ldots$. | 183320.03 | -23 5749.5 | $2^{\prime \prime} \times$ | $2^{\prime \prime}$ | 2 |

[^1]TABLE 2
Photometry Summary

| Parameter | Integrated Photometry $\left(\mathrm{N}_{*}+\mathrm{S}_{*}\right)$ (mag) |  |  |  | $\begin{aligned} & \text { Flux Ratio } \\ & \left(\mathrm{N}_{*} / \mathrm{S}_{*}\right) \\ & \mathrm{P} 1.5, \mathrm{CCD}^{d} \end{aligned}$ | Notes |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | M3.0, Bolo ${ }^{\text {a }}$ | P5.0, $\mathrm{InSb}^{\text {b }}$ | K1.3, $\mathrm{InSb}^{\text {a }}$ | P1.5, CCD ${ }^{\text {d }}$ |  |  |
| Filter: |  |  |  |  |  |  |
| $U$ | $\cdots$ | $\ldots$ | $\ldots$ | $13.5 \pm 0.2$ | ... | 1,2 |
| B | $\ldots$ | $\ldots$ | $\ldots$ | $14.6 \pm 0.1$ | $0.09 \pm 0.02$ | 1 |
| $g$.............. | $\ldots$ | $\ldots$ | $\ldots$ | $14.0 \pm 0.05$ | $0.43 \pm 0.09$ | 1,3 |
| $i$ | $\ldots$ | ... | ... | $13.8 \pm 0.1$ | $7.8 \pm 0.5$ | 1, 4 |
| $J$ | ... | 10.99 | 10.88 | ... | ... |  |
| H |  | 10.12 | 9.96 | $\ldots$ | $\ldots$ |  |
| K | $9.75 \pm 0.08$ | 9.82 | 9.72 | $\ldots$ | $\ldots$ |  |
| $\mathrm{H}_{2} \mathrm{O}$ index | ... | 0.09 | ... | $\cdots$ | $\ldots$ |  |
| CO index |  | 0.18 | ... | $\ldots$ | $\ldots$ |  |
| $L$ | $9.5 \pm 0.1$ | $9.5 \pm 0.3$ | $9.3 \pm 0.14$ | $\ldots$ | $\ldots$ |  |
| [10.2] | $6.3 \pm 0.1$ | ... | $\ldots$ | $\ldots$ | $\ldots$ |  |
| [12.5] | $4.9 \pm 0.12$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |  |
| [20] | 1.90 | ... | ... | $\ldots$ | $\ldots$ |  |

${ }^{\text {a }}$ IRTF 3.0 m on Mauna Kea with bolometer photometer; 6.0 beam diameter.
${ }^{\circ}$ Palomar 5.0 m with InSb photometer; $5 " 0$ beam diameter.
${ }^{\circ}$ Kitt Peak 1.3 m with InSb photometer; $11^{\prime \prime}$ beam diameter.
${ }^{\text {d }}$ P1.5, CCD: Palomar 1.5 m with RCA CCD $; 8.5$ beam diameter.
Notes.-(1) Relative to M22 stars IV-5, 6, 16, 18, 19 using B, $V$ magnitudes of Arp and Melbourne 1959 as quoted in Alcaino Atlas 1973. (2) B-V, $U-B$ relationship for M22 stars deduced from photometry by Eggen 1977. (3) Photometric system defined by Thuan and Gunn 1976. (4) Photometric system defined by Wade et al. 1979. With B-V,V-I relationship for M22 stars deduced from photometry by Eggen 1977.
the KPNO 1.3 m telescopes, are summarized in Table 2. Also included in Table 2 are the results of preliminary optical photometry of the pair in the $U, B, g$, and $i$ bands obtained from RCA CCD frames ( 0.47 per pixel) taken with the Palomar 1.5 m telescope. Reproductions of the central portion of the $g$ and $U$ band CCD frames are shown in Figures $1 a$ and $1 b$. Seeing at the time of observation was about 1 ".5, and therefore the stellar pair appears as an elongated north-south object in the $g$ frame. In order to estimate the relative contribution of the two components, the CCD images of this pair have been deconvolved by fitting point source profiles at the positions of the two stellar objects to the observed north-south profile. The derived flux ratios are also included in Table 2.

The optical photometry is significantly uncertain for two reasons: (1) uncertainty in the relative contributions of $\mathrm{N}_{*}$ and $\mathrm{S}_{*}$ and (2) the photometry is calibrated relative to M22 cluster stars with $B$ and $B-V$ photometry (Arp and Melbourne 1959) as quoted by Alcaino (1973). Extrapolation of this B, V calibration to the $U, g$, and $i$ bands was accomplished using UBVRI color relations for M22 stars (Eggen 1977) and the $g, i$ band definitions from Thuan and Gunn (1976) and Wade et al. (1979). The quoted uncertainties reflect the star-to-star spread in the conversion factors.
The region around IRAS 18333-2357 was also imaged at the KPNO 4 m prime focus using the RCA3 CCD on 1987 March 31. One 900 s exposure (Fig. 2a) was obtained through an interference filter centered at $5006 \AA$ with a bandpass of 31 $\AA$. A second exposure of 240 s was taken through a line-free filter centered at $5279 \AA$ with a $273 \AA$ bandpass. While the nebulosity can be identified in the $5006 \AA$ frame, it can be seen much more easily when the line-free image is subtracted from the on-line image (Fig. 2b). Residual flux from some of the very red, very blue, and saturated stars is evident in the difference frame.

## b) Spectroscopy

Spectroscopic observations in the vicinity of the stellar pair were obtained with the double spectrograph (Oke and Gunn 1982) attached to the Palomar 5 m telescope on 1986 May 28, 29, and 31 and 1987 May 6. The 1986 May spectra of this region were obtained using both $1^{\prime \prime}$ and $2^{\prime \prime}$ slit widths with corresponding spectral resolutions of $4 \AA$ and $8 \AA$ in the $3775-$ $5485 \AA$ range and $12 \AA$ and $24 \AA$ in the $5500-10060 \AA$ range. Exposure times ranged from 100 to 1800 s . The higher resolution 1987 May spectrum was centered on the stellar pair and covered $3858 \AA$ to $4306 \AA$ with $1.4 \AA$ resolution, and 6331 $\AA$ to $6978 \AA$ with $2 \AA$ resolution.

Figure 3 is a section of a blue spectrum showing the $4600-$ $5100 \AA$ region with the slit aligned north-south and located so as to minimize the stellar contribution to the nebular spectrum. Figure $4 a$ shows a $4000-5000 \AA$ spectrum of the stellar pair, while Figure $4 b$ shows the $3860-4300 \AA$ spectrum obtained in 1987 May.

## III. Results

## a) Nebulosity

Figure 3 illustrates the fundamental property of the nebulosity discovered during these observations. The spectrum is dominated by the $5007 \AA$ and $4959 \AA$ lines of [O III]. Conspicuous by their absence are the $4861 \AA \mathrm{H} \beta$ hydrogen recombination line and the $4686 \AA$ recombination line of ionized helium. Indeed, the only other nebular emission line apparent in the $3800-9000 \AA$ wavelength range is the $3869,3967 \AA$ doublet of [ Ne III ] (Fig. 4b). A weak emission feature near the expected position of the $4363 \AA$ [O III] line is indistinguishable in these spectra from the mercury night sky line at $4361 \AA$. We do not believe that this feature is the nebular line. Table 3 lists the intensity of [ Ne III] $3869 \AA$ relative to [O III] $5007 \AA$ as well


Fig. 2.-5006 $\AA$ CCD frame of the region around IRAS $18333-2357$. North is up, and east is to the left. The stellar pair ( $\mathrm{N}_{*}$ and $\mathrm{S}_{*}$ ) appear as the elongated image near the center of the field. (b) Difference image. The same field as in (a) with the $5279 \AA$ image subtracted from the $5006 \AA$ image. The difference has been balanced for redder stars in the field.

TABLE 3
Nebular Spectroscopy Summary

| Line | Relative Line Intensity ( $\mathrm{I}_{5007}=100$ ) |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | M22 Nebula |  | Planetary Nebulae |  |  |  |
|  |  |  | Excitation Temperature ${ }^{\text {b }}$ |  |  | K648 ${ }^{\text {c }}$ |
|  | Measured | Corrected for Extinction ${ }^{\text {a }}$ | High | Medium | Low |  |
| [Ne III], $3869 \AA \ldots \ldots \ldots .$. | $28 \pm 14$ | $37 \pm 18$ | 48 | 12 | 6 | 3.8 |
| $[\mathrm{O} \text { III }], 4363 \AA$ | $<3.0$ | $<3.5$ | 1.2 | 1.0 | 0.5 | 1.3 |
| Не II, $4686 \AA \ldots \ldots \ldots . .$. | $<1.6$ | $<1.75$ | 5.5 | 0.5 | $\ldots$ | $<1$ |
| $\mathrm{H} \beta, 4861 \AA \ldots \ldots \ldots \ldots .$. | <1.3 | $<1.3$ | 9 | 8 | 14 | 48 |
|  | 31 | 31 | 32 | ... | $\ldots$ | 33 |
| [ O III], $5007 \AA \ldots \ldots \ldots . .$. | 100 | 100 | 100 | 100 | 100 | 100 |
| He ı, 5876 $\AA \ldots \ldots \ldots . . . .$. | $<0.8$ | $<0.65$ | 0.5 | 1.1 | 2.0 | 6.2 |
| H $\alpha, 6563 \AA \ldots \ldots \ldots \ldots \ldots$. | <1.4 | $<1.0$ | 26 | 23.2 | 42 | 143 |

[^2]

Fig. 3.-Sample $4600-5100 \AA$ spectrum of the M22 nebula. The spectrum is not flux-calibrated. The wavelengths of the $4861 \AA \mathrm{H} \beta$ and the $4686 \AA$ He il recombination lines are indicated.
as the upper limits for a few relevant lines. Line ratios corrected for line-of-sight reddening to M 22 of $E(B-V)=0.32 \mathrm{mag}$ (Alcaino and Liller 1983) are also included in Table 3, as are typical line ratios for low-, moderate-, and high-excitation planetary nebulae (Aller and Czyzak 1983). The hydrogen recombination lines in the M22 nebula are extremely weak, by a factor of 20 or more for $\mathrm{H} \alpha$, relative to [O III] in typical planetary nebulae. The helium recombination lines in the M22 nebula are at least somewhat, perhaps a factor of 2 , weaker relative to [ O III ]. Lower limits to the relative abundance of oxygen and neon in the M22 nebula are derived in § IVc.

The radial velocity determined from the [ O II ] lines is $-162 \pm 25 \mathrm{~km} \mathrm{~s}^{-1}$. The [ O III] and [ Ne III] lines widths are indistinguishable from the instrumental line width, and the deduced upper limit to the expansion velocity of the nebular region is $75 \mathrm{~km} \mathrm{~s}^{-1}$.

The morphology and extent of the nebulosity is evident in Figure $2 b$. The [ O mII ] emission is asymmetrically distributed relative to the stellar pair. The line emission is concentrated to the east of the stellar pair peaking about $1^{\prime \prime}$ south and east of $\mathrm{N}_{*}$. The nebular extent is approximately $10^{\prime \prime} \times 7^{\prime \prime}$ with the long dimension in the north-south direction.

The flux in the [O III] $\lambda 5007$ line has been estimated from the narrow band difference image by comparison to the standard stars BD $+33^{\circ} 2642$ and Kopff 27 which were observed imme-
diately prior to M22 through the same filters. The derived flux has been adjusted for residual contribution from the blue star since the subtraction is balanced for red stars on average. The blue star is estimated to contribute $20 \%$ of the total measured flux, yielding a net reddening-corrected $5007 \AA \AA$ flux from these data of $4.5 \times 10^{-16} \mathrm{~W} \mathrm{~m}^{-2}$.

The $5007 \AA$ line flux has also been estimated from the spectroscopic data. The line flux measured with a $2^{\prime \prime}$ wide slit centered on the emission peak is about $3.3 \times 10^{-17} \mathrm{~W} \mathrm{~m}^{-2}$. The integrated line flux from the entire nebula, based on stepped $1^{\prime \prime}$ slit spectra, is estimated to be perhaps 2 times larger, leading to a reddening-corrected integrated $5007 \AA$ line flux $2 \times 10^{-16} \mathrm{~W}$ $\mathrm{m}^{-2}$.
In the following discussion, we adopt the mean value of the above estimates, $3 \times 10^{-16} \mathrm{~W} \mathrm{~m}^{-2}$ with an estimated uncertainty of about a factor of 2 for the integrated [O III] $5007 \AA$ line flux from the M22 nebula. The corresponding upper limit to the integrated $\mathrm{H} \alpha$ line flux, corrected for cluster reddening, is $3 \times 10^{-18} \mathrm{~W} \mathrm{~m}^{-2}$. Nebular mass estimates are discussed in $\S$ IV $c$.

## b) Stellar Objects

The CCD observations clearly show that the two stellar objects have very different colors, with the northern component $\left(\mathrm{N}_{*}\right)$ dominating at the longest wavelength and the


Fig. 4.-(a) Sample $4000-5000 \AA$ spectrum of the stellar objects $N_{*}$ and $S_{*}$ from 1986 May observations. The spectrum is not flux-calibrated. The wavelengths of several hydrogen and helium lines are indicated. (b) $3860-4300 \AA$ spectra of stellar pair from 1987 May.
southern component ( $\mathrm{S}_{*}$ ) increasingly dominant towards the blue until $\mathrm{N}_{*}$ is not detectable in the $U$ band frame. The separation between the two components, as measured by the positional difference between the $i$ and $U$ observations, is $1^{\prime \prime} 3$, essentially in the north-south direction. All of the photometry includes both components, as do the stellar spectra.

The adopted magnitudes for the two stellar components are

TABLE 4
Adopted Component Magnitudes

| Filter | Observed |  | Corrected for $E(B-V)=0.32^{\mathrm{a}}$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{N}_{*}$ | $S_{*}$ | $\mathbf{N}_{*}$ | $S_{*}$ |
| $U$ |  | 13.5 (0.2) |  | 12.0 (0.2) |
| B | 16.0 (0.1) | 14.8 (0.1) | 15.3 (0.2) | 13.5 (0.2) |
| $(V)^{\text {b }}$ | 14.7 (0.2) | 14.3 (0.2) | 13.7 (0.2) | 13.3 (0.2) |
| $J$... | 11.04 | ... | 10.76 | ... |
| H | 10.09 | $\ldots$ | 9.92 | $\ldots$ |
| K | 9.81 | $\ldots$ | 9.70 | $\ldots$ |

[^3]shown in Table 4, including component magnitudes corrected for the line-of-sight reddening to M22.
$$
\text { i) Southern Component }\left(\mathbf{S}_{*}\right)
$$

The stellar pair is relatively faint compared to nearby cluster stars at the longer wavelengths, but $S_{*}$ is the brightest object in the $2^{\prime}$ field of the $U$ band frame. $S_{*}^{*}$ has nearly the same $B$ magnitude and $B-g$, and thus presumably $B-V$, color as the nearby blue horizontal-branch cluster stars IV-16 and IV-18 (Arp and Melbourne 1959); however, it is about 1.5 mag bluer in $U-B$ than these stars.

Presumably this extremely blue star has escaped previous notice for two reasons. First, it is located within $1^{\prime}$ of the center of M22, well within the core radius of the cluster, while most searches for unusually blue or UV bright stars have been carried out in the less crowded outer regions of the cluster (e.g., Zinn, Newell, and Gibson 1972). Second, the proximity of the red star, $\mathrm{N}_{*}$, tends to mask the extreme nature of $\mathrm{S}_{*}$.

The $4000-5000 \AA$ spectrum of the stellar pair, Figures $4 a$ and $4 b$, also indicates that $S_{*}$ is an extremely hot star. Clearly present are absorption lines of He II at 4686, 4541, 4200, 4026 and $3923 \AA$, while the $4471 \AA$ line of He I is not visible. Additional He II lines of similar strength to the above isolated He II absorption lines would be expected about $2 \AA$ to the blue of the hydrogen Balmer lines. The observations presented here are

inadequate to discriminate unambiguously between He II or H as the principal absorber or to separate the contributions of the two stars to the principal features near these wavelengths.

The $4000-5000 \AA$ spectrum of the stellar pair is qualitatively similar to those of very hot planetary nebulae nuclei (Mendez, Kudritzki, and Simon 1985). The absorption features appear to be somewhat weak, particularly since $\mathrm{N}_{*}$ might be expected to contribute significantly to the hydrogen absorptions.

The radial velocity of $S_{*}$, using only the 1987 May observations of the isolated $\mathrm{He}_{\text {II }}$ lines, is $-157 \pm 15 \mathrm{~km} \mathrm{~s}^{-1}$. The 1987 May results are much more reliable than the earlier data because of the higher $S / N$ and spectral resolution of the observations. The nature of $S_{*}$ and its possible relationship to the other sources in this complex is discussed briefly in $\S$ IV $a$ and $V$.

## ii) Northern Component $\left(\mathbf{N}_{*}\right)$

The northern component, $\mathrm{N}_{*}$, dominates the emission at near-infrared wavelengths. The near-infrared magnitudes of $\mathbf{N}_{*}$ in Table 4 have been obtained by applying small corrections to the observed integrated magnitude of the stellar pair, where the $\mathrm{S}_{*}$ contribution has been estimated by extrapolation beyond $V$ assuming a constant magnitude after correction for cluster reddening. The properties of $\mathrm{N}_{*}$ in comparison with
field giant stars and M22 cluster giants are illustrated in Figures $5 a$ and $5 b$. The broad-band colors of $\mathrm{N}_{*}$, corrected for cluster reddening, are consistent with an early $\mathbf{M}$ giant spectral type. The observed CO and $\mathrm{H}_{2} \mathrm{O}$ indices are also consistent with such a classification. All the broad-band colors of $\mathrm{N}_{*}$ are redder than those of known M22 cluster stars and its CO index is larger than other cluster stars. Based on the broad-band and $\mathrm{CO} / \mathrm{H}_{2} \mathrm{O}$ photometry, a reasonable interpretation of $\mathrm{N}_{*}$ is that of a background early M giant. However, since the colors of cluster giants are not grossly different from those of field giants, the colors of $\mathrm{N}_{*}$ could also be consistent with a very cool cluster star.

The optical spectra do not exhibit any of the strong spectral features, e.g., TiO band heads at 6160,7050 and $7590 \AA$, which are typical of solar abundance $\mathbf{M}$ stars. The absorption feature near $4225 \AA$ (Fig. $4 b$ ) is attributed to Ca I absorption in $\mathrm{N}_{*}$ as is the $\mathrm{H} \alpha$ absorption. This spectral appearance would be consistent with that of a cool metal-poor cluster star. However, the $2.2 \mu \mathrm{~m}$ magnitude of $\mathrm{N}_{*}$ is about 3 mag fainter than that of the reddest cluster stars, and is located in an area of the colormagnitude diagram not occupied by cluster stars. If, on the other hand, $\mathbf{N}_{*}$ were a field star with near-solar abundances, its spectral type, as indicated from the optical spectra, would be significantly earlier than $M$, say $\sim K 2-K 3$, and an additional


Fig. $5 a$


Fig. $5 b$
Fig. 5.-(a) $(J-K)_{0}$ vs. $(V-K)_{o}$ plot. The solid dots are $\mathbf{M} 22$ giant stars (Frogel et al. 1983). The solid line indicates the field giant color-color relation from Frogel et al. (1978). The position of $\mathbf{N}_{*}$, corrected for line-of-sight cluster reddening, is indicated by the open circle labeled $\mathbf{N}_{*}$. The arrow shows the effect of a hypothetical local absorption of $A_{v}=1.0 \mathrm{mag}$ (see § IVb for discussion). (b) CO index ${ }_{0}$ vs. ( $\left.V-K\right)_{0}$. The definition of symbols is as for (a). The horizontal bar through the $\mathrm{N}_{*}$ position indicates the uncertainty in its $(V-K)_{0}$ color,
reddening component of interstellar or circumstellar origin would be required in order to account for its very red broad band colors. This case is illustrated in Figures $5 a$ and $5 b$, where the reddening arrow indicates the effect of an additional extinction of $A_{v}=1.0 \mathrm{mag}$. This topic is revisited in § IV $a$.
The radial velocity of $\mathrm{N}_{*}$, based on the Ca I and $\mathrm{H} \alpha$ absorption features, is $-86 \pm 10 \mathrm{~km} \mathrm{~s}^{-1}$.
c) 10 and 20 Micron Observations

While the $20 \mu \mathrm{~m}$ source is centered on the stellar pair to within $2^{\prime \prime}$, and this emission has been identified with the IRAS 18333-2357 source (§ II), the present observations are inade-
quate to discriminate positionally between the three optical sources as to the origin of the infrared excess.

The ground-based measurements of 10 and $20 \mu \mathrm{~m}$ flux density in a $6^{\prime \prime}$ diameter beam are a factor of 2-3 less than the interpolation/extrapolation of the 12 and $25 \mu \mathrm{~m}$ IRAS flux densities (cf. Fig. 6). The flux difference at $10 \mu \mathrm{~m}$ is likely to be the result of several objects of similar brightness, including M22-V8 and V-1, in the IRAS $12 \mu \mathrm{~m}$ beam. However, at 20 and $25 \mu \mathrm{~m}$, IRAS $18333-2357$ is the dominant source in the vicinity of M22. There is no indication from either the groundbased IRTF or IRAS observations (Paper I) that the emitting source around $20 \mu \mathrm{~m}$ is significantly extended. The $20 \mu \mathrm{~m}$


Fig. 6. -v $f_{v}$ vs. wavelength for $S_{*}, \mathrm{~N}_{*}$, and IRAS 18333-2357, (filled circles): Color-corrected PSC fluxes from Paper I. (circled crosses): Ground-based observations with $6^{\prime \prime}$ beam size. (open squares): Near-infrared observations of $\mathrm{N}_{*}$. (crosses): CCD observations of $\mathrm{N}_{*}$. (open circles): CCD observations of $\mathrm{S}_{*}$. The $\mathrm{N}_{*}$ and $\mathrm{S}_{*}$ fluxes have been corrected for cluster reddening. Solid lines indicate blackbody flux distribution fitted to the three components. Note that a wide range of temperature greater than about $35,000 \mathrm{~K}$ is consistent with the photometry of $\mathrm{S}_{*}$.
source core size is no more than about $4^{\prime \prime}$ in diameter, based on $20 \mu \mathrm{~m}$ IRTF measurements offset in the cardinal directions from the peak position by $2^{\prime \prime}$ and $3^{\prime \prime}$. Explanations for the 20 $\mu \mathrm{m}$ flux discrepancy could be (1) a low-level extended halo outside the $6^{\prime \prime}$ beam used for the ground-based observations, (2) a continuum energy distribution very strongly peaked beyond about $25 \mu \mathrm{~m}$, (3) line emission in the $25 \mu \mathrm{~m}$ IRAS bandpass but not in the $20 \mu \mathrm{~m}$ IRTF filter/atmospheric bandpass, or (4) source variability.

## Iv. DISCUSSION

## a) Cluster Membership of Components

The radial velocity measurements are summarized in Table 5. The radial velocity of M22 itself is $-153 \pm 2.6 \mathrm{~km} \mathrm{~s}^{-1}$ (Webbink 1981), quite distinct from that of field stars in this direction. Since the estimated central escape velocity from M22 is about $35 \mathrm{~km} \mathrm{~s}^{-1}$ (Webbink 1981), radial velocities between -118 and $-188 \mathrm{~km} \mathrm{~s}^{-1}$ are very likely to indicate cluster membership, while radial velocities outside this range are not consistent with cluster membership. On this basis, the nebulosity is a virtually certain cluster member, $S_{*}$ is likely to be a cluster member, and $\mathrm{N}_{*}$ cannot be a cluster member.

We argue that $S_{*}$ is physically associated with the nebulosity. The probability of a chance alignment of $S_{*}$, a very hot star as evidenced by the presence of He II absorption lines and its $U-B$ color, within about $2^{\prime \prime}$ of the center of an ionized nebulosity is completely negligible. Therefore we conclude that the two objects are physically associated and in M22.

The a priori probability that $\mathrm{N}_{*}$ is a field star is quite low. The surface density of field stars in the vicinity of M22, brighter than $V=+15$ mag and redder than $B-V=1.0 \mathrm{mag}$, is about $0.4 \mathrm{arcmin}^{-2}$. This estimate has been obtained from the M22 color-magnitude diagram (Alcaino and Liller 1983) assuming

TABLE 5
Radial Velocity Summary

| Object | Radial Velocity ( $\mathrm{km} \mathrm{s}^{-1}$ ) | Source |
| :---: | :---: | :---: |
| M22 cluster | $-152.5 \pm 2.6$ | Webbink 1981 |
| Field stars | $1.5 \pm 9$ | Lloyd-Evans 1978 |
| Nebulosity | $-162 \pm 25$ | [ $\mathrm{OHII}^{\text {II }} 5007 \AA, 4959 \AA$ |
| $\mathrm{S}_{*}$ | $-157 \pm 15$ | He if absorption (May 1987) |
| $\mathrm{N}_{*}$ | $-86 \pm 10$ | Ca 1 4226, $\mathrm{H} \alpha$ |

that stars not closely associated with the M22 giant branch are field stars. Thus one would expect only about one such field star within $1^{\prime}$ of the center of M22 and the probability of a chance alignment of a field star within $2^{\prime \prime}$ of the M22 nebula-S ${ }_{*}$ system is only about $10^{-3}$. In spite of this, the radial velocity of $\mathrm{N}_{*}$ would appear to rule out cluster membership. If $\mathrm{N}_{*}$ is a field star, then its broad-band colors are substantially redder, by $\sim 1-1.5 \mathrm{mag}$ in $V-K$, than expected for field stars with strong Ca I 4226 and no TiO features. This additional reddening could be due to interstellar extinction in or beyond M22 or to possible circumstellar extinction local to $\mathbf{N}_{*}$ itself. However, the measured CO index (Fig. 5b) is not consistent with an early spectral type for $\mathrm{N}_{*}$. Thus this star appears to be quite peculiar in its properties, and no consistent picture of this object has emerged.

## b) Relationship to IRAS 18333-2357

In Paper I it was suggested, among other possibilities, that IRAS 18333-2357 may be related to planetary nebula-like objects because of its unusual infrared energy distribution. Its continuum energy distribution is quite similar to that of many planetary nebulae in the PSC and unlike that of most circumstellar dust shells around cool stars. This argument also strongly suggests a physical association of IRAS 18333-2357 and the M22 nebulosity- $S_{*}$ pair.

A luminosity comparison also supports this point of view. Figure 6 shows the observed energy distributions for IRAS 18333-2357, $\mathrm{N}_{*}$, and $\mathrm{S}_{*}$. The optical and near-infrared observations have been corrected for a line-of-sight reddening of $E(B-V)=0.32$. In this plot, blackbodies with the same peak $v f_{v}$ have the same integrated flux, and therefore the same luminosity, if at the same distance. The integrated flux of $\mathrm{N}_{*}$ is only about 0.15 that of IRAS 18333-2357; thus it appears that this star cannot be the luminosity source for IRAS 18333-2357 unless it is very heavily attenuated by local absorbing dust. The integrated flux of $\mathrm{N}_{*}$, even with the assumption of an additional $A_{v}=1 \mathrm{mag}$, as indicated in Figures $5 a$ and $5 b$, is still a factor of 3 less than that of IRAS 18333-2357.
On the other hand, there is some evidence that the integrated flux of $S_{*}$, while very uncertain, may be substantially larger than that of IRAS 18333-2357. Blackbody energy distributions for temperatures of 35,000 and $80,000 \mathrm{~K}$, adjusted to agree with the flux density of $S_{*}$ at $B$, are shown in Figure 6, illustrating the cases of equal and $10 \times$ the integrated flux density of IRAS 18333-2357. Either temperature is consistent with the broad-band photometry of $\mathrm{S}_{*}$, while the spectrum of $\mathrm{S}_{*}$ is roughly similar to those of planetary nebula nuclei with effective temperatures in the range $60,000-100,000 \mathrm{~K}$ (Mendez, Kudritzki, and Simon 1985). The luminosity of $\mathrm{S}_{*}$ is approximately given by $700\left[T\left(\mathrm{~S}_{*}\right) / 35,000\right]^{3} L_{\odot}$; thus it would be the most luminous star in M22 if its effective temperature exceeds about $55,000 \mathrm{~K}$.

An additional argument supporting the association of IRAS 18333-2357 and the $\mathrm{S}_{*}-\mathrm{M} 22$ nebula may be derived from estimating the grain equilibrium temperature in a dust cloud coincident with the [ O III ] $5007 \AA$ nebula heated by $\mathrm{S}_{*}$. Assuming $0.1 \mu \mathrm{~m}$ amorphous carbon grains with $Q \sim \lambda^{-1}$ for $\lambda>0.1 \mu \mathrm{~m}$, a dust cloud radius of $4^{\prime \prime}$, and the $L, T$ relation given above, one derives equilibrium dust temperatures of 63 $\mathrm{K}, 77 \mathrm{~K}$, and 103 K for values of $T\left(\mathrm{~S}_{*}\right)=35,000 \mathrm{~K}, 50,000 \mathrm{~K}$, and $80,000 \mathrm{~K}$, respectively. This picture is self-consistent in terms of the expected temperature range for $S_{*}$ and the $\sim 100$

K far-infrared energy distribution of IRAS 18333-2357 (Paper I).

We conclude that the most likely luminosity source for IRAS 18333-2357 is $S_{*}$. This conclusion is consistent with association of IRAS $18333-2357$ with the $\mathrm{S}_{*}-\mathrm{M} 22$ nebula system.
c) Nebular Properties
i) Abundances

Lower limits to the relative abundance of oxygen and neon can be obtained from the observed line ratio limits, given assumptions about the electron density and temperature. An upper limit of $20,000 \mathrm{~K}$ for the electron temperature can be determined from the lower limit to the [ O III] $(5007+4959)$ / [ O III] 4363 line ratio. In the following discussion, we assume that the electron temperature is $10,000 \mathrm{~K}$, a value typical of planetary nebulae.

With the relations from Aller (1985) and the upper limits to the $\mathrm{H} \alpha, \mathrm{He}$ II 4686 and He I 5876 lines, then $N\left(\mathrm{O}^{++}\right)$/ $N\left(\mathrm{H}^{+}\right)>8 \times 10^{-3}, \quad N\left(\mathrm{O}^{++}\right) / N\left(\mathrm{He}^{++}\right)>2 \times 10^{-2}, \quad$ and $N\left(\mathrm{O}^{++}\right) / N\left(\mathrm{He}^{+}\right)>5.5 \times 10^{-3}$. If all the oxygen is in the form of O III, then its relative abundance compared to solar (Aller and Czyzak 1983) is $[\mathrm{O} / \mathrm{H}]>+1.0$ and $[\mathrm{O} / \mathrm{He}]>-0.1$; where $[X / \mathrm{H}]=\log (X / \mathrm{H})-\log (X / \mathrm{H})$ Sun. Similarly, the relative intensity of [ Ne mI ] $3869 \AA$ corresponds to $[\mathrm{Ne} /$ $\mathrm{H}]>+1.9$ and $[\mathrm{Ne} / \mathrm{He}]>+0.8$.

The above abundance limits for the M22 nebula are grossly unlike the envelope composition of M22 stars. M22 is a very metal-poor globular cluster, with $[\mathrm{Fe} / \mathrm{H}]=-1.8$ to -1.9 (Cohen 1981; Alcaino and Liller 1983 and references therein). Stars in metal-poor clusters tend to be somewhat overabundant in oxygen relative to Fe (see, e.g., Pilachowski, Sneden, and Wallerstein 1983); thus [O/H] for M22 stars may be about -1.5 . Therefore, the oxygen abundance (and that of Ne ) in the M22 nebula appears to be enhanced by at least a factor of 500 with respect to hydrogen and 30 with respect to helium when compared to typical cluster stars.
$\mathrm{H}_{\text {iI }}$ regions with an overabundance of heavy elements are generally characterized by electron temperatures that are lower than $10,000 \mathrm{~K}$ since the heavy ions, particularly those of oxygen, are the dominant cooling mechanism (e.g., Price 1981). Assuming a lower electron temperature for the M22 nebula would lead to even larger overabundances of oxygen and neon relative to helium and hydrogen. In fact, nebulae such as A30, having a large overabundance of heavy ions, should have very low electron temperatures if the only heating mechanism is photoionization. If this were true, the [ O mI ] $\lambda 5007$ line would be very weak, which it is not. Harrington and Feibelman (1984) find a similar situation for A30 and attribute the necessary heating to the interaction between the nebula and a strong stellar wind from the central star.
ii) Nebular Mass

In this discussion, the M22 nebula is treated as a simple, uniform density sphere of radius $3^{\prime \prime}$, an electron temperature of $10,000 \mathrm{~K}$, and an integrated [ O III ] $5007 \AA$ line flux of $3 \times 10^{-16} \mathrm{~W} \mathrm{~m}{ }^{-2}$. With these simplifying assumptions, the mass of $O \mathrm{mI}$ and the total ionized gas mass are still substantially uncertain because of (1) unknown ionization equilibrium and (2) uncertain electron density in the nebula as a result of nondetection of hydrogen and helium recombination lines.

In order to illustrate the range of potential nebular mass, two examples are evaluated: (1) H and He abundances corre-

TABLE 6
Inferred Nebular Properties

| Case | $\underset{\left(\mathrm{cm}^{-3}\right)}{N_{\mathrm{e}}}$ | Ionized Mass ( $\left.1 \times 10^{-4} M_{\odot}\right)$ |  |  |  | $\begin{aligned} & f_{v}(20 \mathrm{~cm}) \\ & \left(10^{-6} \mathrm{Jy}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $M(\mathrm{H})$ | $M(\mathrm{He})$ | M( $\mathrm{O}_{\text {iII }}$ ) | $M(\mathrm{Ne} \mathrm{miI)}$ |  |
| $\mathrm{H}, \mathrm{He}_{\mathrm{I}}, \mathrm{He}_{\text {II }}$ lines at upper limits | 44 | 1.1 | 7.8 | 0.1 | 0.1 | 9 |
| Pure O ili, Ne ili | 4.6 | $\ldots$ | $\ldots$ | 1.3 | 2.2 | 0.17 |

sponding to the observed recombination line upper limits, and (2) abundances of all other elements and ionization stages negligible relative to O iII and Ne III. The first case leads to an upper limit on the ionized mass and electron density, while the second case results in a minimum ionized gas mass and electron density. The results are listed in Table 6. The ionized gas mass of the M22 nebula and corresponding electron density is very small, about $3-10 \times 10^{-4} M_{\odot}$ and 5-50 $\mathrm{cm}^{-3}$, respectively. Indeed, if, as suggested above, the dust responsible for the infrared emission of IRAS 18333-2357 is associated with the M22 nebulosity, then the total mass in the nebular region (ionized gas plus dust) is dominated by the solid component. The minimum dust mas associated with IRAS 18333-2357 is about $6 \times 10^{-4} M_{\odot}$ for either high absorption efficiency silicate material or amorphous carbon grains (Paper I). Thus the minimum dust mass estimate is approximately the same as the estimated mass of ionized gas; in this case the dust-to-gas ratio in the M22 nebula would be at least two orders of magnitude larger than the canonical interstellar value of 0.01 .

Neither the IRAS observations nor the observations presented here provide direct information on the composition of the dust responsible for the infrared emission. Amorphous carbon or silicate grains would appear to be most likely from the point of view of element abundances and frequency of occurrence. A strong $30 \mu \mathrm{~m}$ emission feature has been found in some planetary nebulae and carbon stars (Forrest, Houck, and McCarthy 1981). Such an emission feature near $30 \mu \mathrm{~m}$ could explain the discrepancy between the $20 \mu \mathrm{~m}$ ground-based photometry and the $25 \mu \mathrm{~m}$ IRAS measurements; however, this feature is thought to be due to compounds like $\mathrm{Fe}_{3} \mathrm{C}$ (Forrest, Houck, and McCarthy 1981) or MgS (Goebel et al. 1980) and would require even more extreme overabundances of the heavy elements involved.

## V. NATURE AND ORIGIN OF THE M22 NEBULA

There is no evidence in the M22 nebula for high-velocity flows which are commonly associated with novae or supernovae events. In addition, there is an obvious candidate
progenitor/ionization source, $\mathrm{S}_{*}$; thus the nebulosity, the infrared source, and $S_{*}$ could constitute a planetary nebula system in M22. Indeed, in Paper I, it was suggested that IRAS 18333-2357 may be related to planetary nebulae, but the lack of associated radio or $\mathrm{H} \alpha$ source required that it not be a typical planetary nebula. Only one other planetary nebula, K648 in M15, is known to be associated with a Galactic globular cluster.

The $L, T$ relation for $S_{*}$ intercepts the region occupied by planetary nebulae nuclei in the temperature range 30,000 to about $70,000 \mathrm{~K}$ (cf. Pottasch 1984). Thus $\mathbf{S}_{*}$ could be similar to other planetary nebulae nuclei. If $\mathrm{S}_{*}$ is the origin of the M22 nebula, one would expect it to exhibit envelope abundances similar to or more extreme than that of the nebula. It is surprising, given this scenario, that $S_{*}$ shows photospheric absorption lines of He II and possible H . However, the He II and $\mathrm{H} / \mathrm{He}$ II absorption features are somewhat weak and may be consistent with substantial envelope deficiencies for both elements.

The time scale for expansion of the M22 nebula to its current size is roughly $3000\left(20 \mathrm{~km} \mathrm{~s}^{-1} / v_{\text {exp }}\right) \mathrm{yr}$, where $v_{\text {exp }}$ is the expansion velocity. This time scale is quite reasonable compared to the expected duration of the planetary nebula phase, provided the expansion velocity is of order $20 \mathrm{~km} \mathrm{~s}^{-1}$.

The M22 nebula is compared to K648 and other planetary nebulae in Table 7. Its properties are sufficiently distinct from those of typical planetary nebulae to suggest that its origin or progenitor star or both are substantially different from those responsible for most, if not all, other planetary nebulae.

The two most obvious differences are the mass and abundances of the M22 nebula. The M22 nebula (ionized gas plus dust) is about 100 times less massive than a typical planetary nebula and about 10 times less massive than K648 in M15. Thus the lack of a significant associated radio source ( $<1 \mathrm{mJy}$ at 20 cm ; Hamilton, Helfand, and Becker 1985) and $\mathrm{H} \alpha$ source (less than that of K648; Peterson 1984) is a direct result of the small ionized mass of the nebula. The calculated upper limit to the 20 cm free-free emission from the M22 nebula is only $9 \mu \mathrm{Jy}$

TABLE 7
M22 Nebula/Planetary Nebulae Comparison

| Obiect | $N_{e}$ | $T_{e}$ | $\begin{gathered} \text { MASS } \\ \left(10^{-4} M_{\odot}\right) \end{gathered}$ | Abundances Relative to Solar ${ }^{\text {a }}$ |  |  |  |  | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | [C/H] | [O/H] | [ $\mathrm{Ne} / \mathrm{H}$ ] | [ $\mathrm{O} / \mathrm{He}$ ] | [ $\mathrm{Ne} / \mathrm{He}$ ] |  |
| M22 Nebula | 5-44 | <20,000 | 3-10 | ? | $>1.0$ | $>1.9$ | $>-0.1$ | $>0.8$ | 1 |
| K648 | 8000 | 12,500 | 110 | +0.06 | -1.25 | $-1.5$ | -1.25 | -1.5 | 2 |
| Typical PN | 4000 | 11,000 | 2000 | +0.20 | -0.25 | -0.03 | -0.21 | $+0.01$ | 3 |
| A78 | 200 | 14,000 | 10 | 2-3 | 0.5-1.0 | 1-2 | -1.2 | -1 to 0 | 4 |
| A78 (center) | 1500 | 17,500 | 10 | ... | 2.3 | 2.5 | 0.8 | 1.0 | 5 |
| A30 ........ | ... | 18,300 | $\ldots$ | $\ldots$ | ... | $\ldots$ | -0.5 | 0.8 | 6 |

[^4](Table 6), and the measured upper limit to the $\mathrm{H} \alpha$ line intensity is more than a factor of 1000 less than that of K648.

The very small mass of the M22 nebula may be related to the abundance anomaly, in that hydrogen and helium could be viewed as being very underabundant. Indeed, the mass of oxygen and neon in the M22 nebula is roughly comparable to that in typical planetary nebulae and may be ten times more than that in K648. In addition, a comparison of IRAS 18333-2357 with the planetary nebula NGC 2392 (Paper I) shows that these two sources have very similar energy distribution in the $12-100 \mu \mathrm{~m}$ range and that the deduced radiating dust mass, if the dust composition is the same for both sources, is about a factor of 6 greater for the M22 source.

Both the M22 nebula and K648 are in very metal-poor globular clusters; $[\mathrm{Fe} / \mathrm{H}] \sim-1.8$ in $\mathbf{M} 22$ (Cohen 1981; Alcaino and Liller 1983) and $[\mathrm{Fe} / \mathrm{H}] \sim-2.0$ for M15 (Zinn 1980). K648 shows clear evidence of substantial overabundance of carbon relative to hydrogen compared to M15 cluster stars and probably a small enhancement of oxygen (Adams et al. 1984). The M22 nebula exhibits a much more extreme ( $>100$ times) enhancement of oxygen abundance relative to both hydrogen and helium. Neon is also grossly overabundant, and, depending on the composition of the dust associated with the nebula, carbon or Mg and Si would appear to be vastly overabundant as well.

The carbon enhancement in K648 has been interpreted as a result of mixing of material involved in helium burning into the surface layers of the progenitor star prior to ejection as a planetary nebula (Adams et al. 1984). The more extreme abundances in the M22 nebula appear to indicate that the material comprising the M22 nebula has experienced hydrogen, helium, and possibly carbon burning with both the hydrogen-rich and helium-rich outer layers of the progenitor star expelled or otherwise removed before the formation of the current nebula.

Models of the origin of the M22 nebula must also account
for both the age of the progenitor star, presumably the age of M22, about 15 Gyr (Alcaino and Liller 1983), as well as the extreme abundance anomalies noted above.

One possibility is that the M22 nebula originated from a close binary system containing a carbon-oxygen white dwarf. Iben and Tutukov (1984) have outlined how such binary systems can be formed by one or more mass exchanges between intermediate mass stars. If the binary spacing is appropriately small, the orbit may decay as a result of gravitational radiation leading to a final interaction on a time scale of $10^{10}$ yr. While the above authors were concerned with mechanisms for formation of Type I supernovae, interaction in a somewhat less massive binary system may potentially lead to a much less explosive result.

An alternative explanation is that the M22 nebula represents an older, more diffuse, form of the enriched material seen in the hydrogen deficient knots in the planetary nebulae Abell 30 and Abell 78 (Jacoby and Ford 1983; Harrington and Feibelman 1984; Manchado, Pottasch, and Manpuso 1988). The physical characteristics of these nebulae compare reasonably well with the M22 nebula (see Table 7), except that the densities of the Abell 30 and Abell 78 material are about 10 times greater. The M22 nebula would then be the third member of a rare class of planetary nebulae in which the central star has experienced a final helium shell flash leading to the ejection of processed material (Iben et al. 1983).

We wish to thank Jordan Caurasco for assistance with the 200 inch observations. The IRTF is supported by NASA. Partial support by NASA through JPL is acknowledged by F. C. Gillett, G. Neugebauer, and B T. Soifer.

Infrared astronomy at Caltech is supported by a grant from the National Science Foundation. J. G. Cohen is grateful for support from the Caltech Recycling Center.

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[^0]:    ${ }^{1}$ Operated by the Association of Universities for Research in Astronomy, Inc., under contract with the National Science Foundation.

[^1]:    References.-(1) Hertz and Grindlay 1983; (2) This paper; (3) Paper I.

[^2]:    ${ }^{\text {a }}$ Assuming $E(B-V)=0.32$ mag.
    ${ }^{\text {b }}$ Average line ratios for classes as defined by Aller and Czyzak 1983.
    ${ }^{\text {c Adams et al. } 1984 .}$

[^3]:    ${ }^{2} A_{\lambda} / E(B-V)$ from Rieke and Lebofsky 1985.
    ${ }^{\mathrm{b}}$ Inferred using $g, V, B-V$ relation from Thuan and Gunn 1976.

[^4]:    ${ }^{2}$ Solar abundances from Aller and Czyzak 1983.
    References.-(1) Adams et al. 1984; (2) Aller and Czyzak 1983; (3) Jacoby and Ford 1983; (4) Manchado et al. 1988; (5) Harrington and Feibelman 1984.

