

The “Top” Priority at the LHC

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(LHC Workshop, Princeton, March 21, 2007)

With Vernon Barger and Devin Walker [hep-ph/0612016](#)

Rakhi Mahbubani, Devin Walker and Lian-Tao Wang, to appear.

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– Top quark: A Window to New Physics

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(Tsinghua University, Jan. 8, 2008)

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An exciting time:

The triple “Coincidence”

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The triple “Coincidence”

- (1). A highly successful theory: the Standard Model.
- (2). Terascale new physics must exist!
- (3). Upcoming the LHC: Uncover the underlying physics.

Terascale Physics at the LHC

- Unitarity argument for $W_L W_L$ scattering
 - ⇒ New physics must show up at the Terascale: A Higgs boson $m_H < 1 \text{ TeV}$ or alike.
- Naturalness argument for a m_H or EW scale
 - ⇒ New physics needed beyond $H^0 \dots$
- Gauge coupling unification
 - ⇒ New threshold at the Terascale.
- Particle dark matter
 - ⇒ WIMP at the Terasacle natural.

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What exact form is it realized in nature?

Fundamental scalar in a weakly coupled theory? (SUSY or alike)

Composite Higgs and strongly interacting dynamics? (TC, Little Higgs)

Low scale string/gravity? (Large extrad dim., RS, ...)

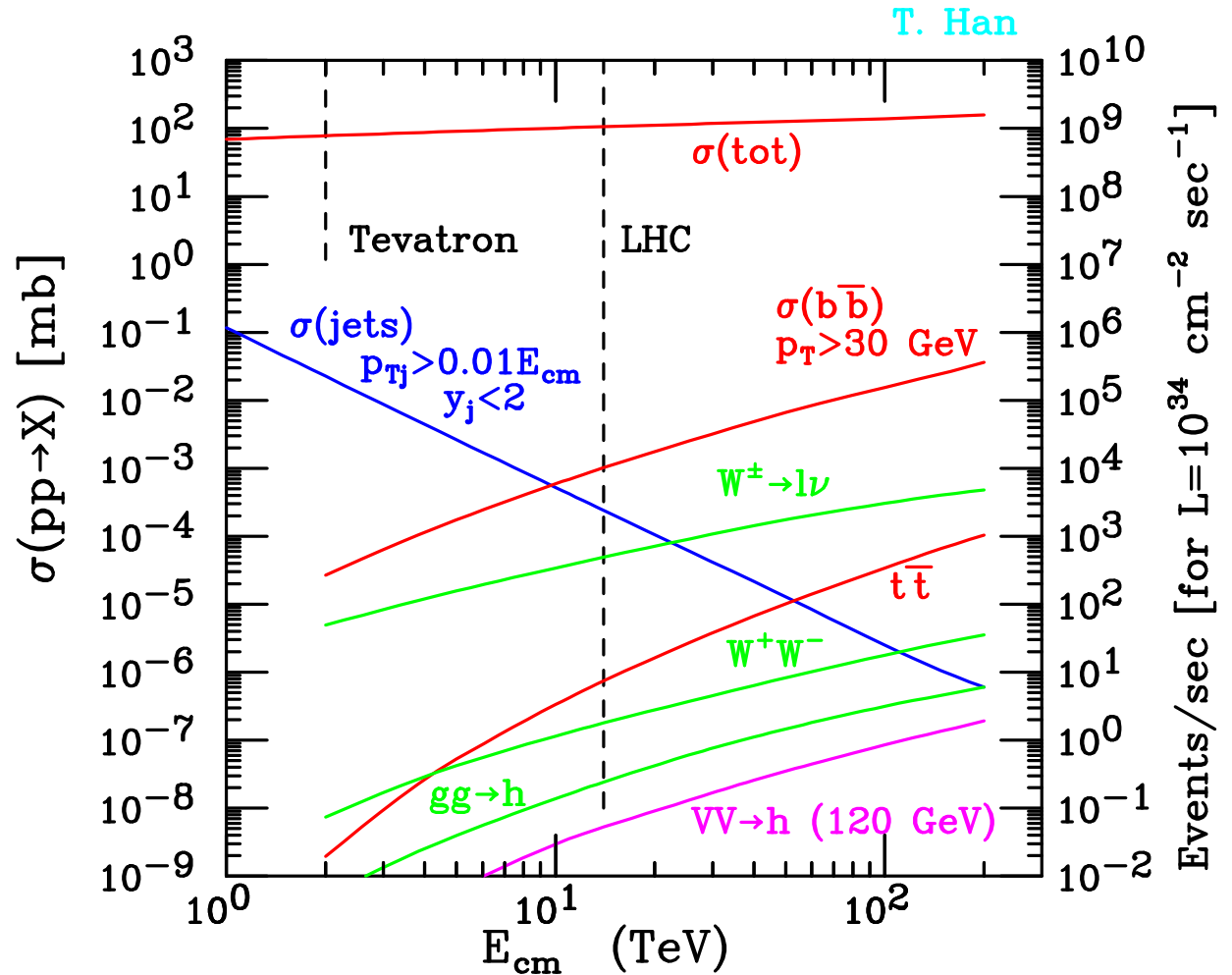
Dark matter connection? (WIMP: LSP ...)

The LHC will tell!

In anticipation of the LHC

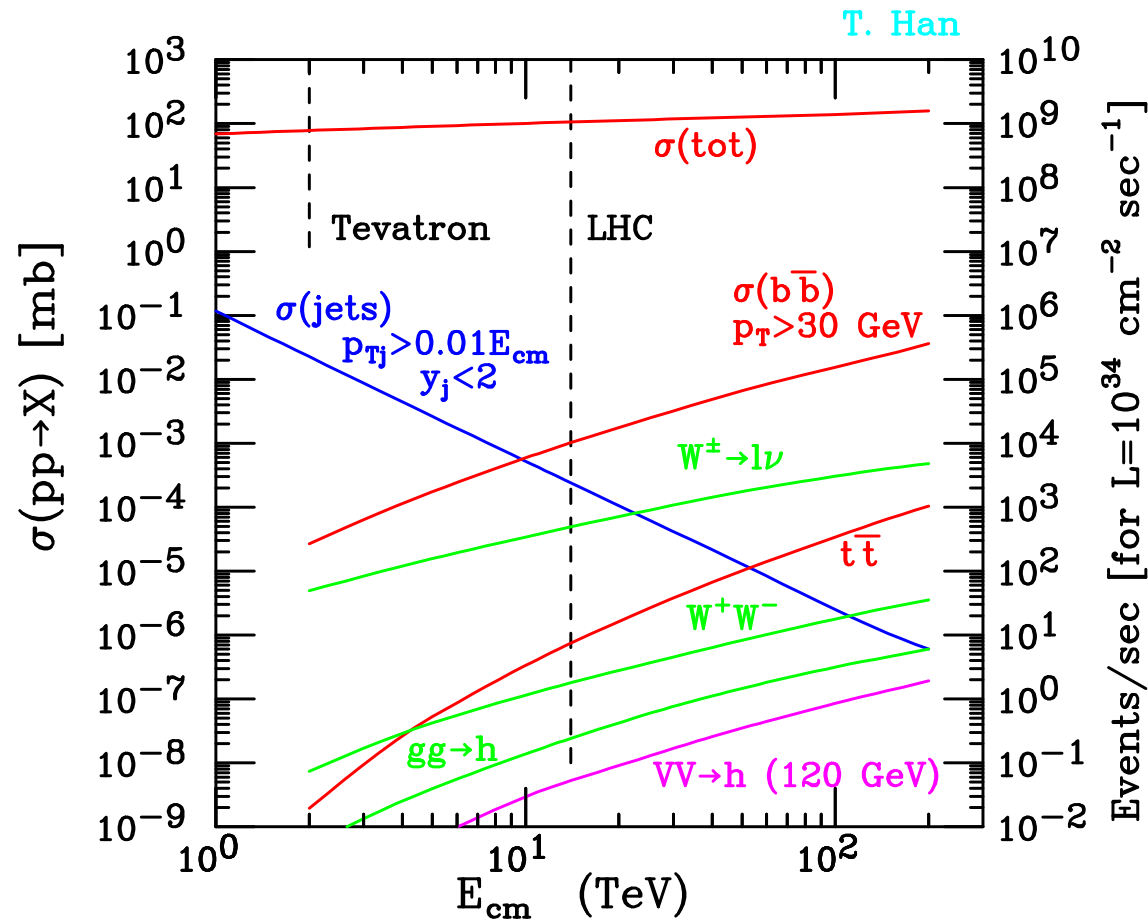
Re-discover the SM!

- DY: $Z \rightarrow \ell^+ \ell^-$, $W^\pm \rightarrow \ell^\pm \nu$
- Jet inclusive $E_T(j)$
- Heavy quarks Υ , J/ψ , $b\bar{b}$, $t\bar{t}$
- Gauge boson pairs WW , ZZ
- Large $\cancel{E}_T + X$



Top quarks at the LHC

LHC is a top factory:



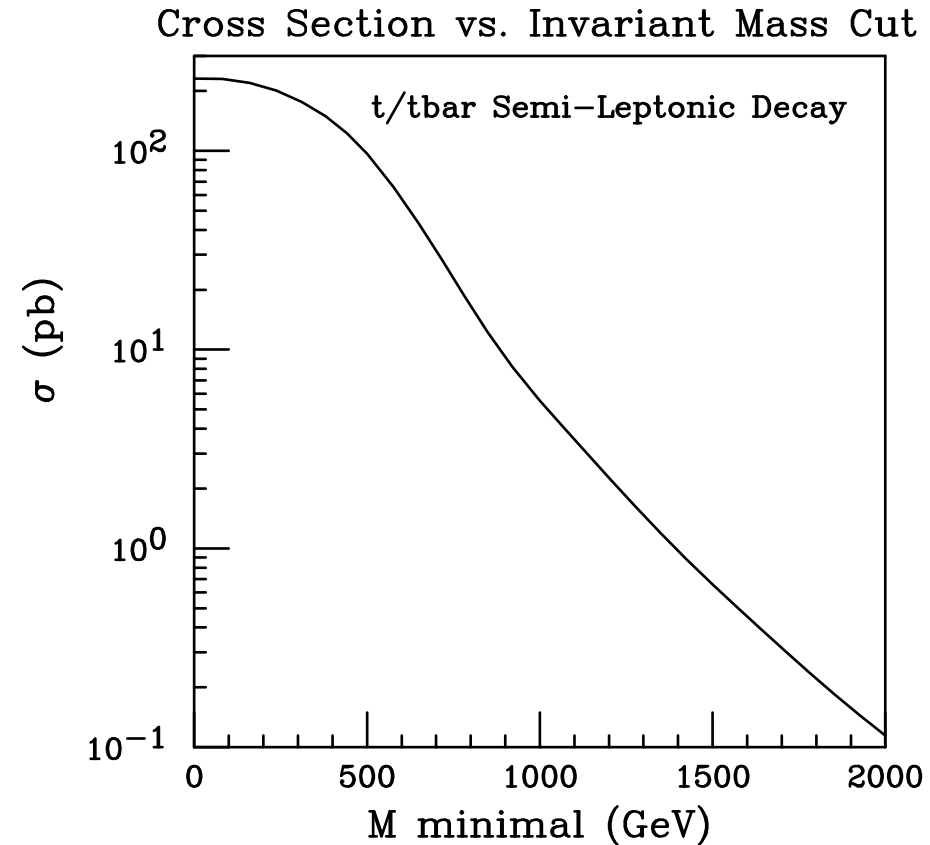
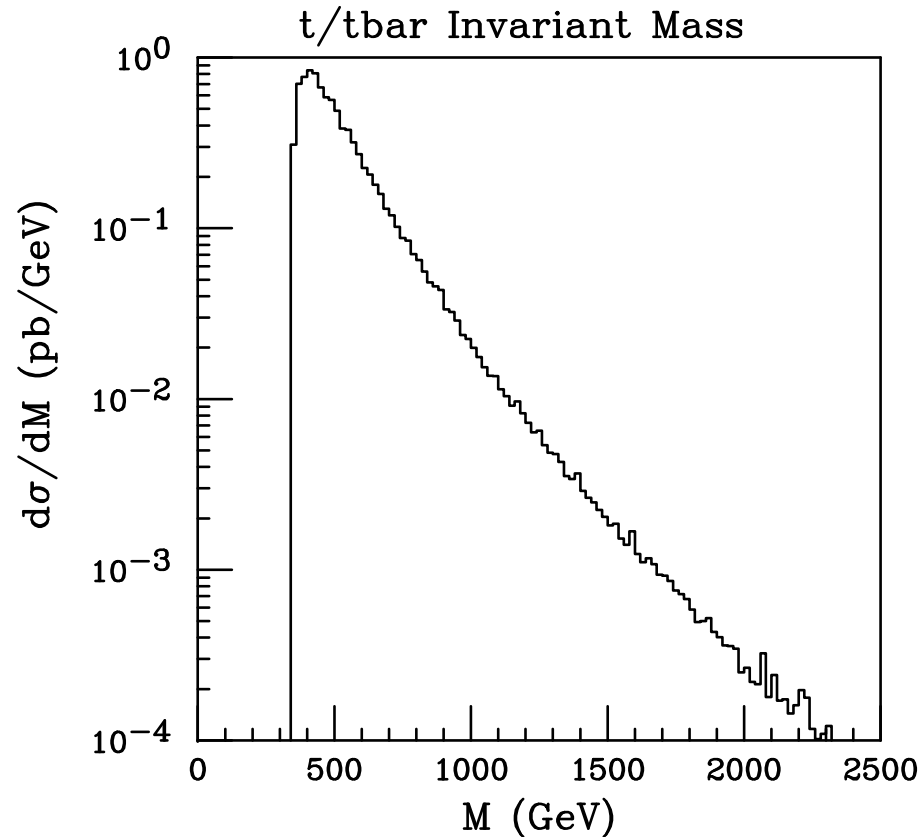
Event rate: **800K $t\bar{t}$ / fb⁻¹**, or **8 Hz @10³⁴/cm²/s !**

From Tevatron to LHC: $t\bar{t}$ increased by 100; EW increased by 10.

Top quarks at the LHC:

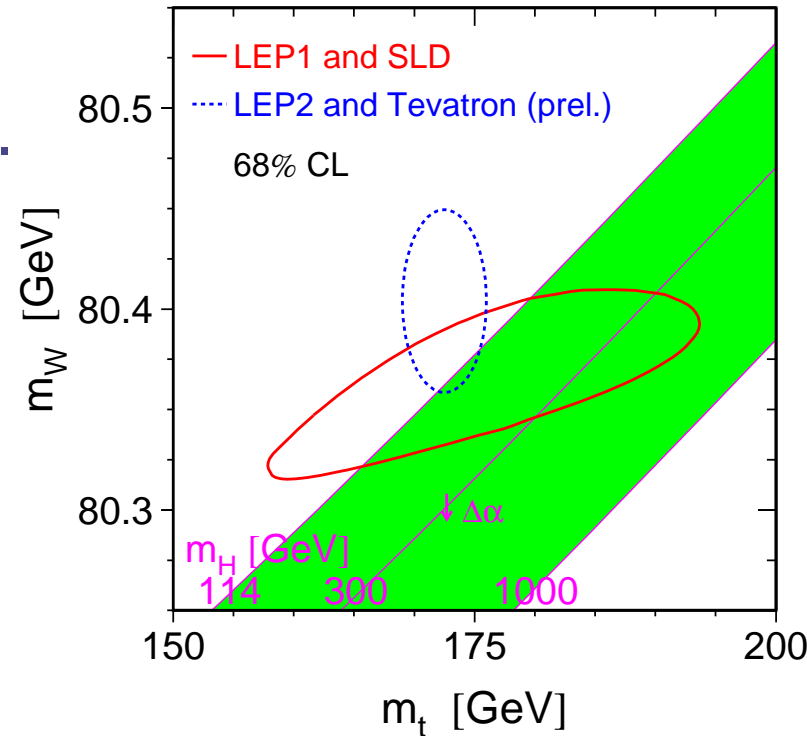
Production well predicted in the SM:

At the LHC: gg 90%, $q\bar{q}$ 10%;
(At the Tevatron: gg 10%, $q\bar{q}$ 90%.)



Why Top Quarks? bread & butter:

- Top quark exists, as the heaviest particle in the SM.
- m_t is the most precisely measured quark mass — Important for precision physics of the SM and beyond:
- Top decays before it hadronizes:
Good test ground for QCD, spin correlation, couplings, CP property...
- Possible deep connection to electroweak symmetry-breaking:
 $m_t \approx v/\sqrt{2}$.



Why Top Quarks? A window to new physics

- Largest Yukawa coupling (proportional to $m_t, \cot \beta$):
 $H, A \rightarrow t\bar{t}$.
- Strong Dynamics (TC, Topcolor/Top See-Saw, Little Higgs):
 $\rho_{TC}, \eta_{TC}, \pi_{TC}, Z_L \rightarrow t\bar{t}$.
- Extra-dimensions (warped and universal):
 $g_{KK}, G_{KK} \rightarrow t\bar{t}$.
- Flavor physics at high scale: $t \rightarrow Zc, \gamma c, gc (u)$.
- Supersymmetry (\tilde{t} often the lightest squark): $\tilde{t}_R \rightarrow t\tilde{\chi}^0$.
- LH with T-parity (theories with naturalness argument):
 $T \rightarrow tA^0$.
- To the least, precision couplings: $t\bar{t}Z, t\bar{t}H; t\bar{t}\gamma, g; t\bar{t}W \dots$

The Remainder of the Talk

- Search I: $t\bar{t}$ Resonant Production
 - Backgrounds
 - Reconstruction Methods
 - Search for New Physics
- Search II: $t\bar{t} + \cancel{E}_T$ Signal
 - Backgrounds
 - Signal Events Reconstruction
- Conclusions

Search I: $t\bar{t}$ Resonant Production

- “Bump searches” in the $M_{t\bar{t}}$ distribution.
- Representative features:

Model Class	Spin-0	Spin-1	Spin-2
Technicolor/Topcolor/RS	× (nrw/brd)	× (nrw/brd)	× (narrow)
MSSM	× (narrow)		
Little Higgs	× (narrow)	× (narrow)	

- A model-independent approach:
Parametrize each resonance with a few parameters:
 m , Γ , σ -normalization, chirality, CP violation

Strategy

To maximumly extract the resonant information:
(Spin, chirality couplings, CP properties ...)
 \implies Need full kinematics and top-ID.

- Using the clean channel: "Semi-leptonic"

$$t\bar{t} \rightarrow b\ell^\pm\nu, \quad bjj \rightarrow 2b \ 2j \ \ell^\pm \cancel{E}_T.$$

- Total Hadronic Channel: $\sigma_{t/t\bar{b}ar} \times (6/9)^2 \implies$ large background, no top-ID ...
 - Semi-Leptonic Channel: $\sigma_{t/t\bar{b}ar} \times 6/9 \times 2/9 \times 2 \implies$ current interest.
 - Pure leptonic Channel: $\sigma_{t/t\bar{b}ar} \times (2/9)^2 \implies$ small rate, incomplete kinematics ...
 - Semi-leptonic/hadronic ratio: 2/3
 - Leptonic/hadronic ratio: 1/9
- We propose new/refined top reconstruction methods.
 - Take advantage of the tops being highly boosted.

Background Considerations

- **W + jets, Z + jets, WW, WZ, ZZ backgrounds:**
 - Consider table of efficiencies reproduced from the ATLAS TDR (Volume II, p. 624). The expected events are in the last column.

Process	Efficiency with $p_T^l > 20$ GeV $E_{T\text{miss}} > 20$ GeV cuts	As before, with plus $N_{\text{jet}} \geq 4$ cut	As before, with plus $N_{\text{b-jet}} \geq 2$ cut	Events per 10 fb⁻¹
<i>t</i> \bar{t} signal	64.7	21.2	5.0	126,000
W + jets	47.9	0.1	0.002	1658
Z + jets	15.0	0.05	0.002	232
<i>WW</i>	53.6	0.5	0.006	10
<i>WZ</i>	53.8	0.5	0.02	8
<i>ZZ</i>	2.8	0.04	0.008	14
Total Background				1922
S/B				65

Background Cont'd

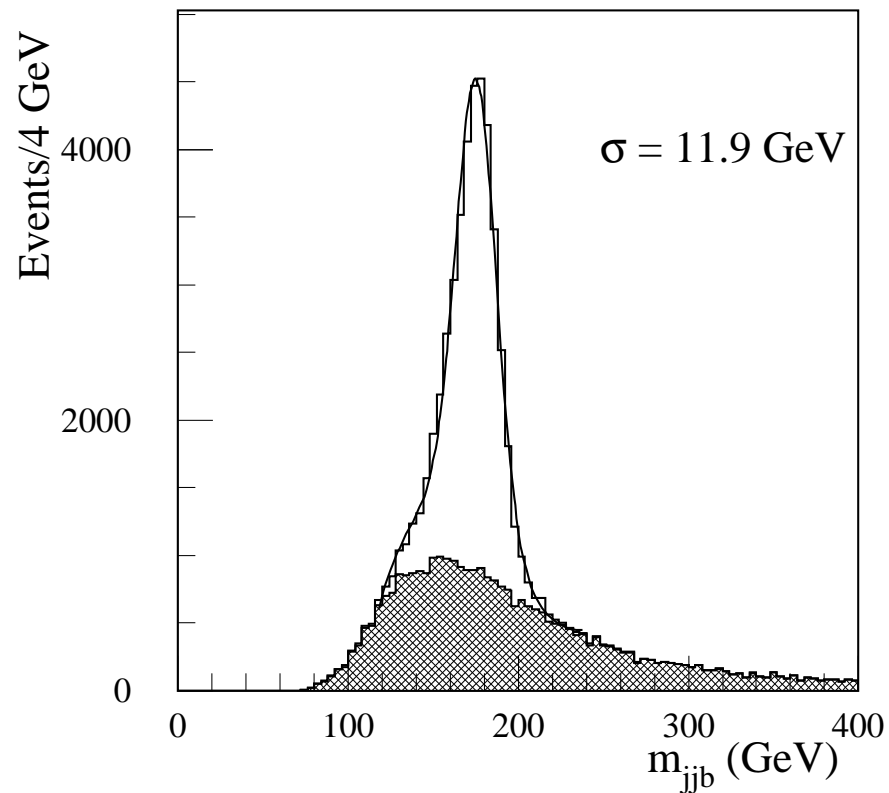
- **W + jets, W + 4 jets, Wbb + 2j, Wbb + 3j backgrounds:**

- Efficiencies from CMS TDR (Volume II, p. 238).

	Semi-leptonic $t\bar{t}$	Other $t\bar{t}$	W+4j	Wbb+2j	Wbb+3j	S/B
Before cuts	365k	1962k	82.5k	109.5k	22.5k	5.9
L1+HLT Trigger	62.2%	5.3%	24.1%	8.35%	8.29%	7.8
Four jets $E_T > 30$ GeV	25.4%	1.01%	4.1%	1.48%	3.37%	9.9
p_T^{lepton} cut	24.8%	0.97%	3.9%	1.41%	3.14%	10.3
b-tag criteria	6.5%	0.24%	0.064%	0.52%	0.79%	25.4
Kinematic fit	6.3%	0.23%	0.059%	0.48%	0.72%	26.7
Cross section (pb)	5.21	1.10	0.10	0.08	0.05	26.7
Scaled to $\mathcal{L} = 1$ fb⁻¹	5211	1084	104	82	50	26.7

Background Cont'd

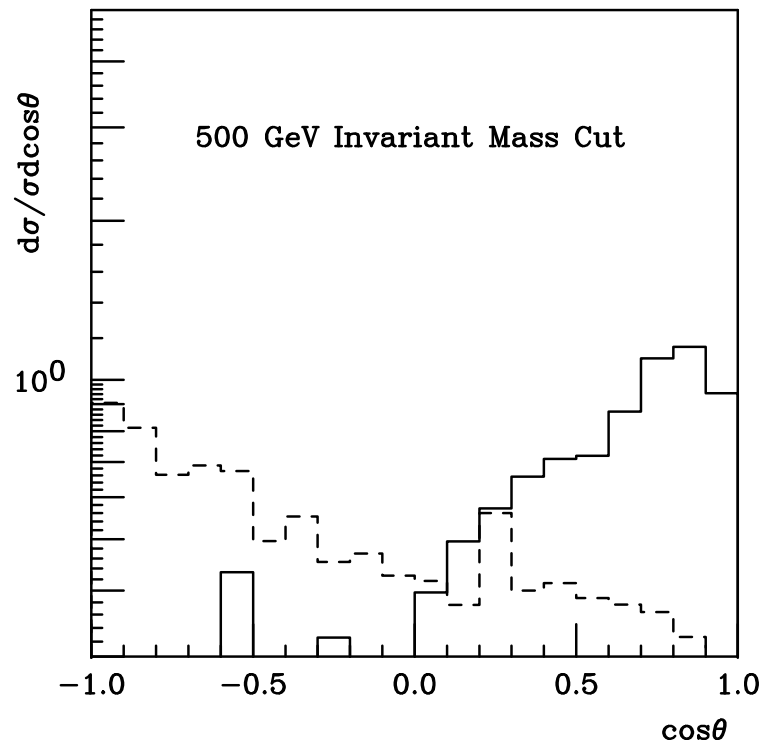
- Reconstructing the hadronic decay (\bar{t} , say):
 - From ATLAS collaboration (ATLAS TDR, Volume II, p. 625): The \bar{t} is reconstructed via the hadronic decay $\bar{b}jj$. The wrong b may have some contamination (shaded area).



Background Cont'd

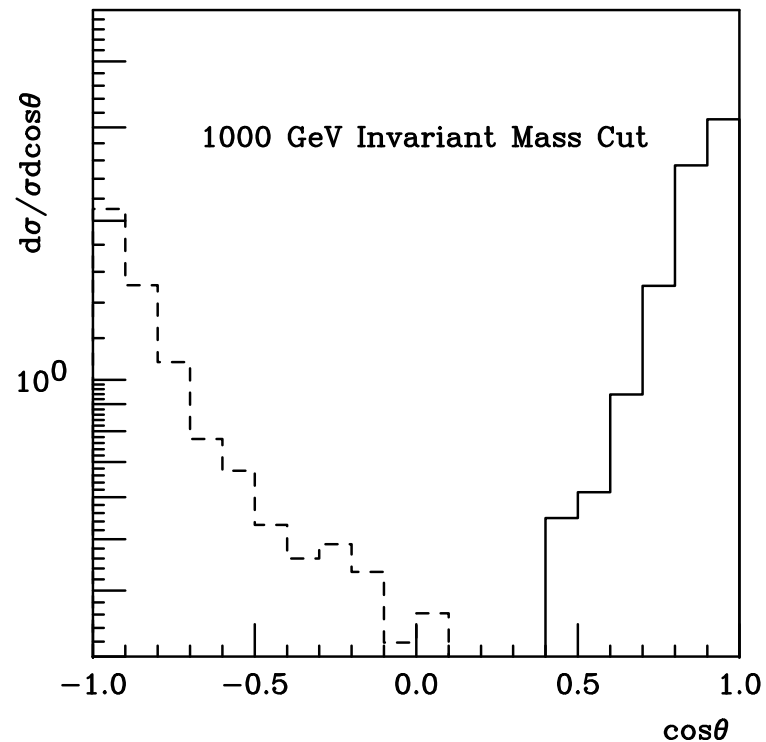
- At high $M_{t\bar{t}}$, the tops are boosted.
That helps select $(\ell^+ b)$, rather than $(\ell^+ \bar{b})$.

Cos θ between lepton and b quarks



Solid – Lepton/leptonic b quark $\cos\theta$
Dashed – Lepton/hadronic b quark $\cos\theta$

Cos θ between lepton and b quarks

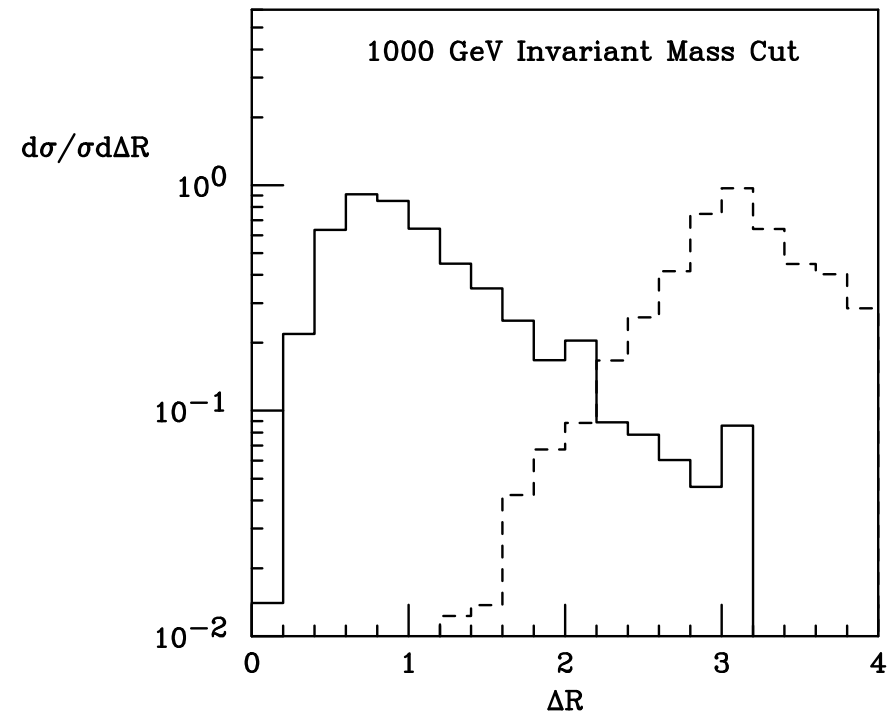
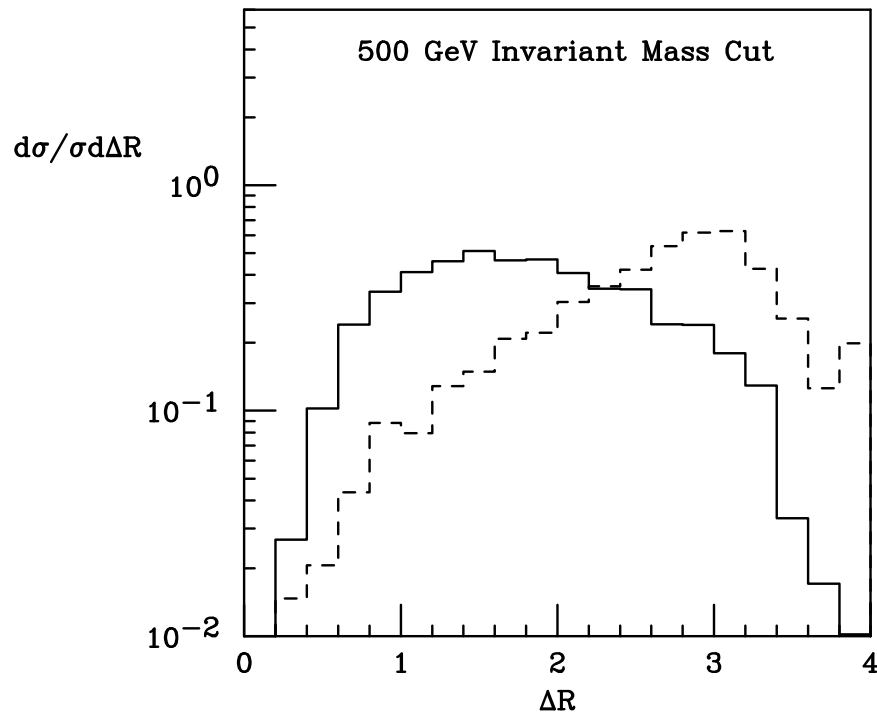


Solid – Lepton/leptonic b quark $\cos\theta$
Dashed – Lepton/hadronic b quark $\cos\theta$

Background Cont'd

- Lepton isolation does not effect the signal appreciably:

$$\Delta R > 0.4$$



Event Selection

- When indicated we apply the following acceptance cuts:
 - Jets: $p_T > 20 \text{ GeV}$, $|\eta_j| \leq 2.5$
 - Leptons: $p_T \geq 20 \text{ GeV}$, $|\eta_l| \leq 2.5$, $\Delta R > 0.4$
 - $\cancel{E}_T \geq 20 \text{ GeV}$ (tightened the ATLAS cuts)
 - Minimum Transverse Mass: $M_T(t\bar{t}) > 600 \text{ GeV}$
- Detector effects (Gaussian smearing) taken into account:

$$\frac{\Delta E}{E} = \frac{a}{\sqrt{E/\text{GeV}}} \oplus b.$$

- CMS ECAL resolution: $a = 0.03$, $b = 0.005$
- CMS HCAL resolution: $a = 1.1$, $b = 0.05$
- pT CMS resolution: $a = 1.5 \times 10^{-4}$, $b = 0.05$

ATLAS/CMS simulations:

SM processes not a threat to the top-quark sample.

Kinematics Reconstruction

- Present two schemes:

Reconstruction of the missing neutrino!

- $(M_W, m_t)^*$ scheme: the masses as known inputs.
- Small angle scheme: t decay products collimated.

*Similar to the three-constraint kinematic fit used at the Tevatron to determine the top mass.

CDF Collaboration, Phys. Rev. Lett. 80 (1998) 2767;

D0 Collaboration, Phys. Rev. D27 (1983) 052001.

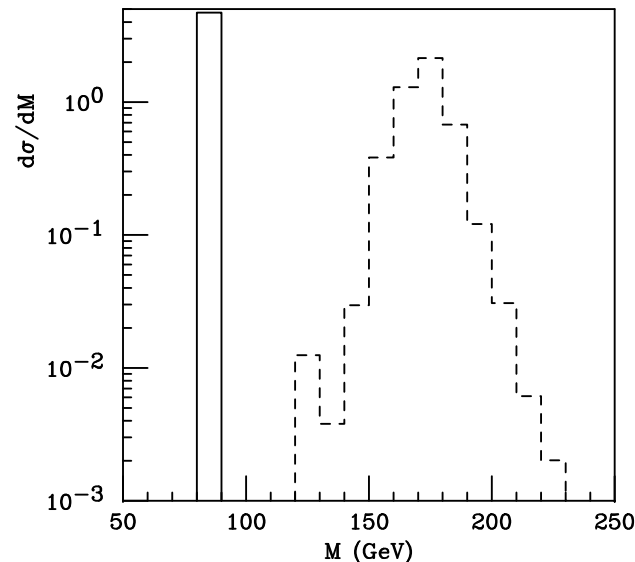
(M_W, m_t) Scheme

- **Step 1:** Reconstruct W boson with the constraint: $M_W^2 = m_{l\nu}^2$.
Yields two-fold ambiguity for the neutrino longitudinal momentum:

$$p_{\nu L} = \frac{A p_{eL} \pm E_e \sqrt{A^2 - 4 \vec{p}_{eT}^2 \cancel{E}_T^2}}{2 p_{eT}^2}, \quad \text{where } A = M_W^2 + 2 \vec{p}_{eT} \cdot \cancel{E}_T.$$

- If $A^2 - 4 \vec{p}_{eT}^2 \cancel{E}_T^2 \geq 0$, then choose the value that reconstructs $m_t^2 = m_{l\nu b}^2$.
- If not, go to Step 2.

Expected distributions from this reconstruction alone:



(M_W, m_t) Scheme Cont'd

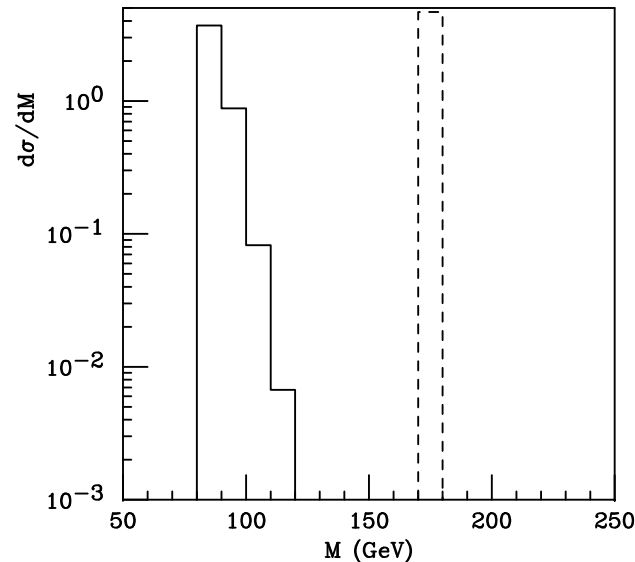
- **Step 2:** If Step 1 gives complex solutions, first reconstruct top quark with: $m_t^2 = m_{l\nu b}^2$.
Yields two-fold ambiguity for the neutrino longitudinal momentum:

$$p_{\nu L} = \frac{A' p_{blL} \pm \sqrt{p_{blL}^2 A'^2 + (E_{bl}^2 - p_{blL}^2) (A'^2 - 4E_{bl}^2 \cancel{E}_T^2)^2}}{2(E_{bl}^2 - p_{blL}^2)},$$

$$A' = m_t^2 - M_{bl}^2 + 2\vec{p}_{blT} \cdot \vec{\cancel{E}}_T.$$

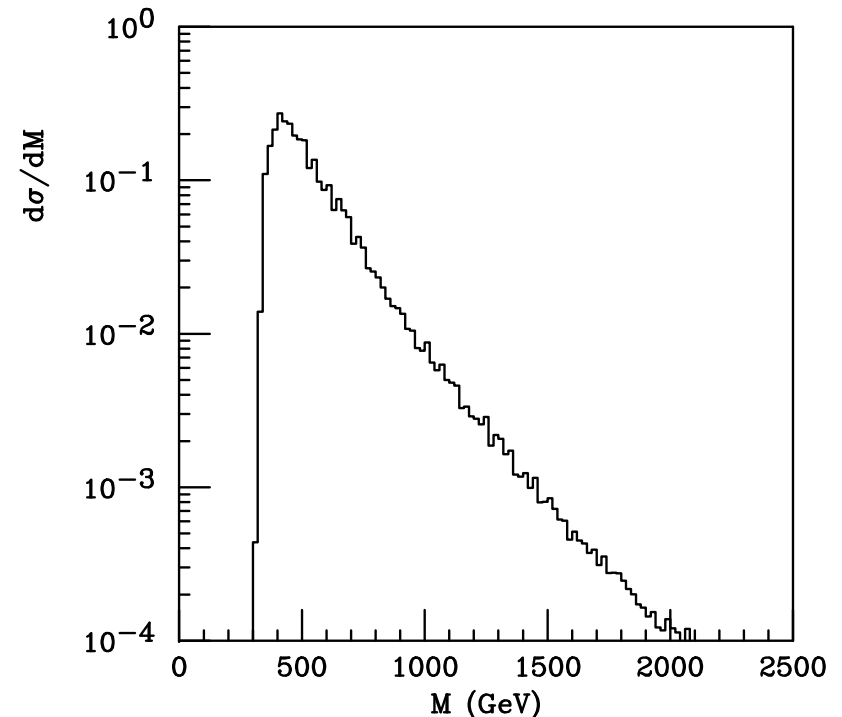
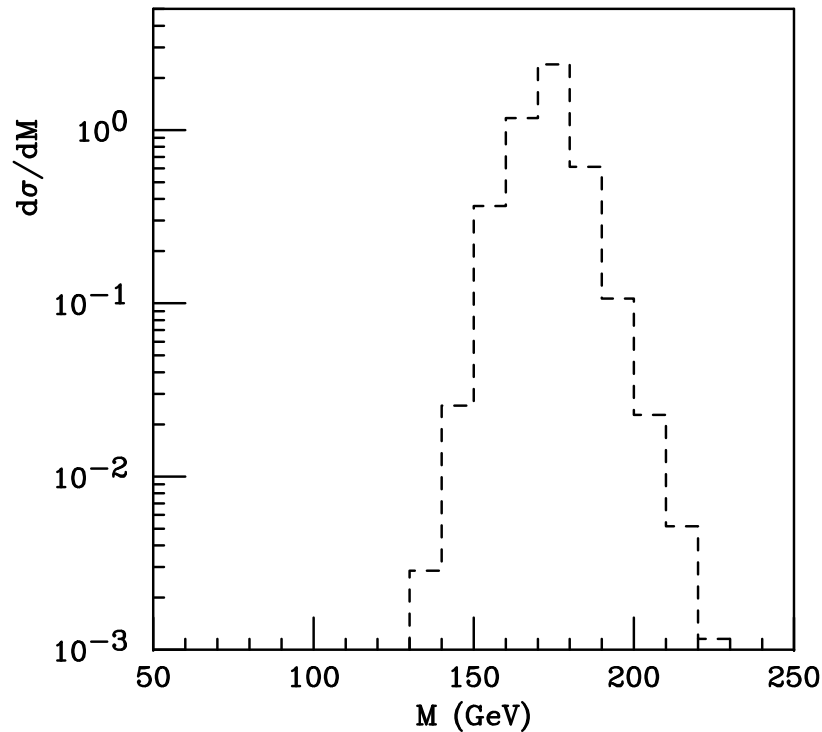
- If $p_{blL}^2 A'^2 + (E_{bl}^2 - p_{blL}^2) (A'^2 - 4E_{bl}^2 \cancel{E}_T^2)^2 \geq 0$, then choose $M_W^2 = m_{l\nu}^2$.
- If not, go to Step 3.

Expected distributions from this reconstruction alone:



(M_W, m_t) Scheme Cont'd

- **Step 3:** If the solution fails to reconstruct for both times, the event is discarded.
 - The discarded event rate $\sim 16\%$.
 - Choosing to keep the discarded solutions (by taking the real components) results in a distortion of high/low invariant mass tails.
- Resulting reconstructions of m_t and M_{tt} :



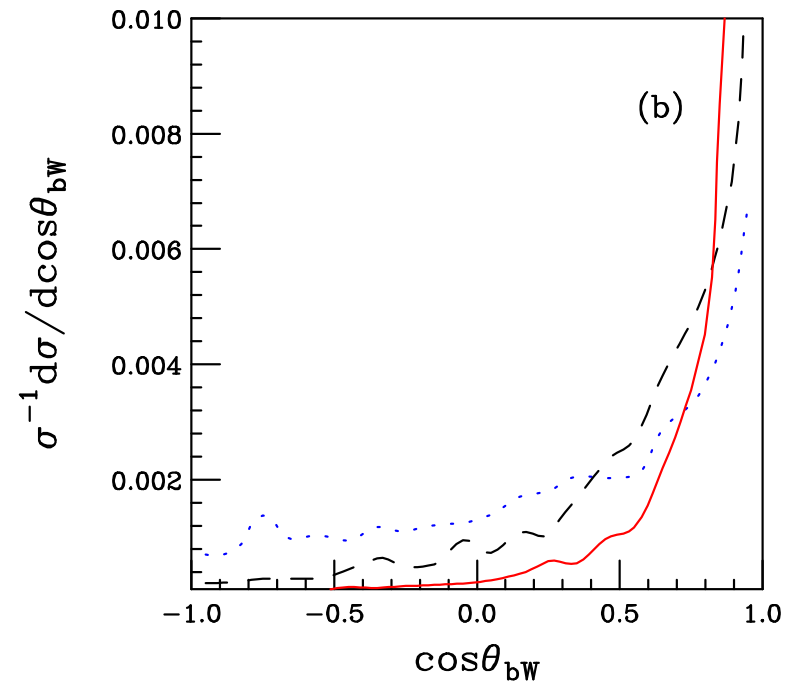
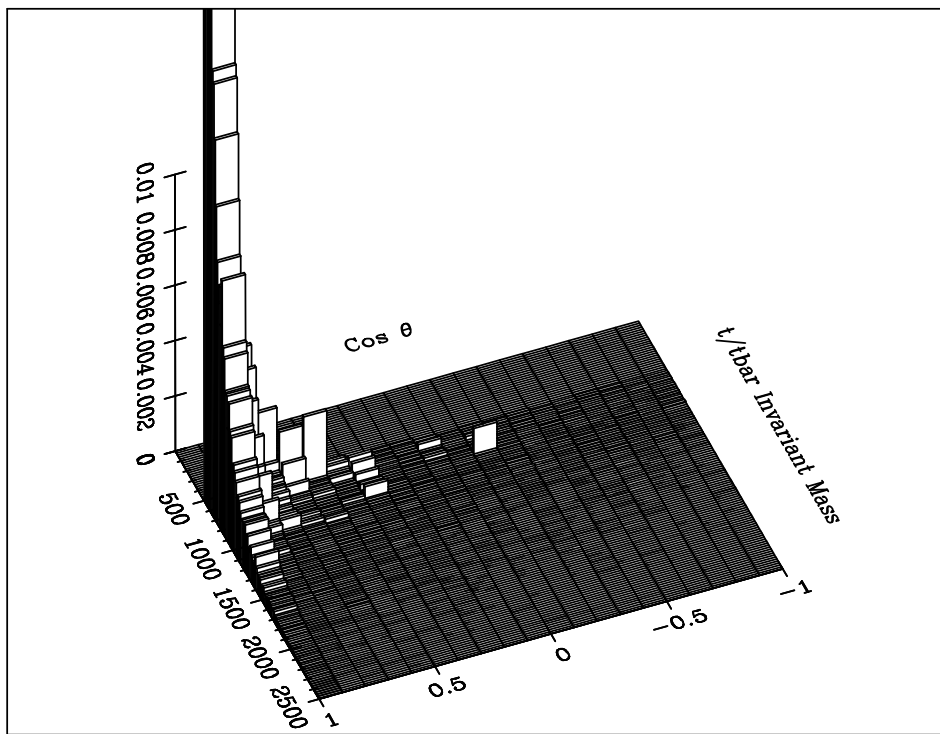
We have the full $t\bar{t}$ kinematics!

Small Angle Selection Scheme

Wish not to rely on m_t input, because ...

- High $t\bar{t}$ invariant mass limit: Tops quarks highly energetic/boosted:
- Expect a small angle θ_{lv} , or in turn, θ_{bW} to give the correct solution.

Consider $\cos\theta_{bW}$:

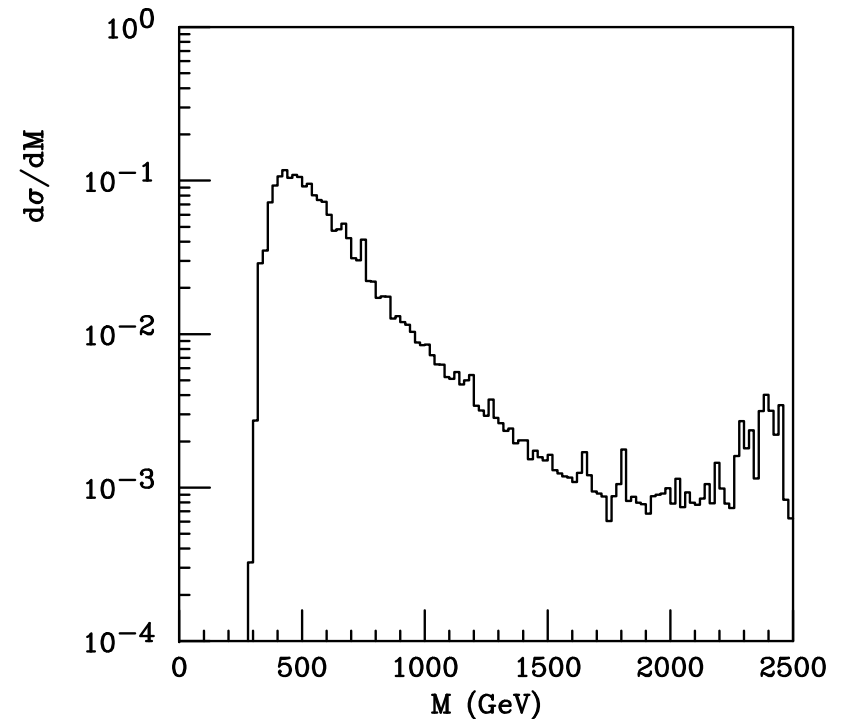
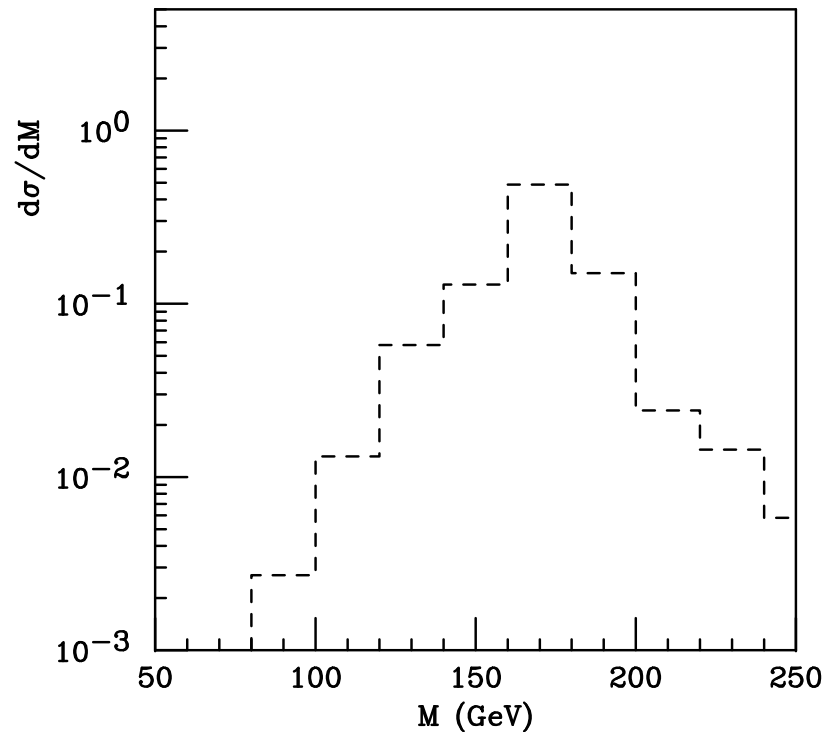


$M_T > 0, 600, \text{ and } 1000 \text{ GeV.}$

Small Angle Selection Scheme Cont'd

Thus,

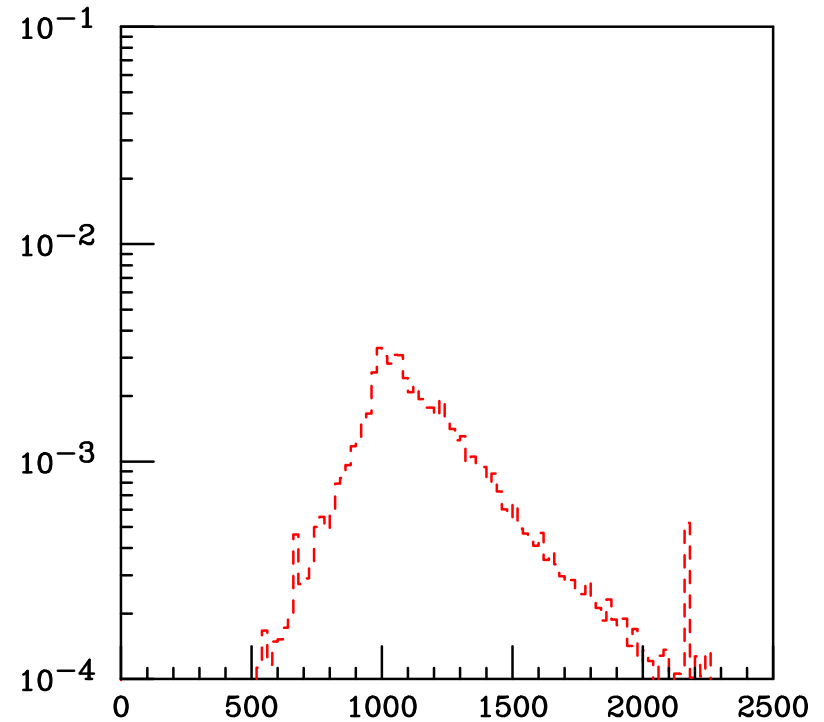
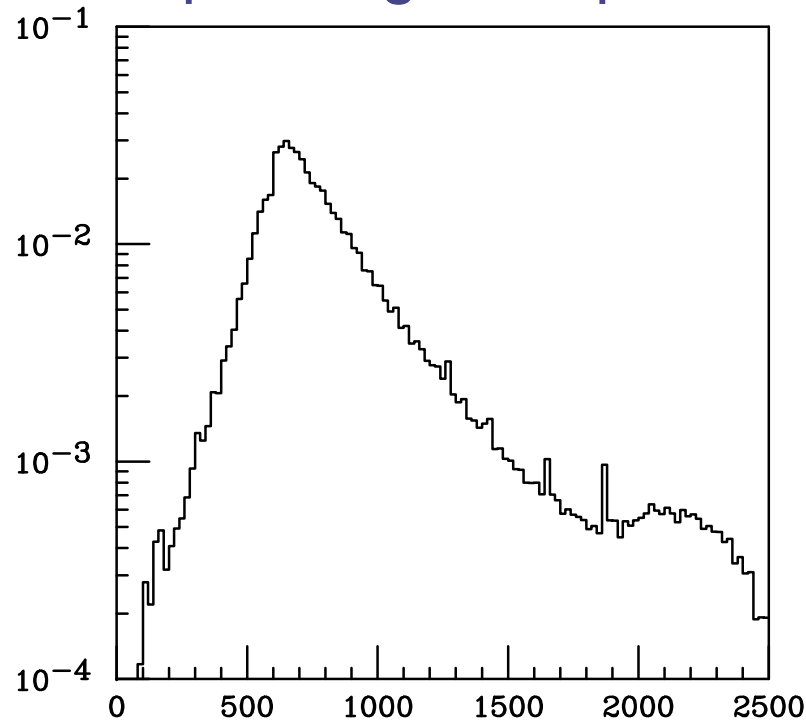
- Again, $M_w^2 = m_{l\nu}^2 \implies$ two-fold ambiguity in $p_{\nu L}$ and thus p_W .
- The scheme: selecting the solution with the smaller angle between the b and the W .
- Resulting reconstruction of m_t and m_{tt} :



Note the leakage at the high m_{tt} region...

Small Angle Selection Scheme Cont'd

- High Invariant mass tail is due to smearing and the wrong solution in small angle selection.
- A solution: Provide an incrementally higher transverse mass cut depending on expected resonance.



$M_T > 650$ and 950 GeV.

Search for New Resonances in m_{tt}

- Search for integer spin resonances via

$$gg \rightarrow \phi_0 \rightarrow \bar{t} + t$$

$$q\bar{q} \rightarrow V_1 \rightarrow \bar{t} + t$$

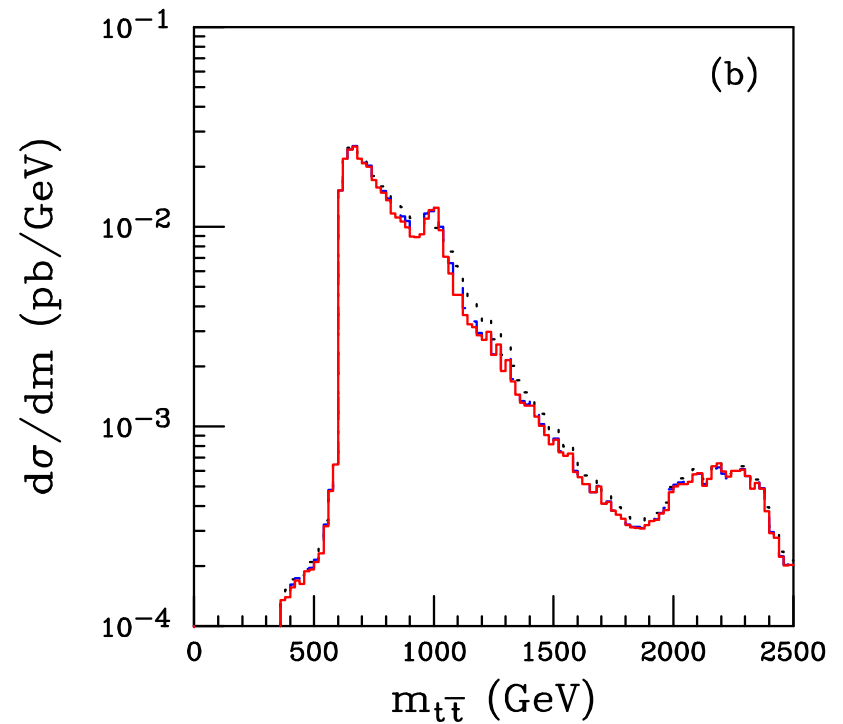
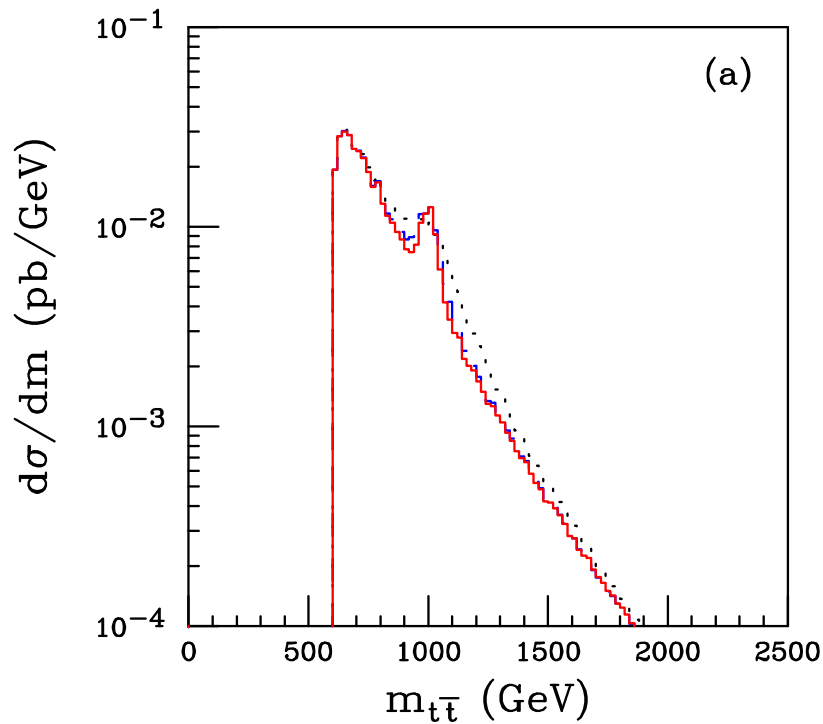
$$q\bar{q}, gg \rightarrow \tilde{h}_2 \rightarrow \bar{t} + t$$

where ϕ , V and \tilde{h} are $J = 0, 1, 2$ resonances.

- Parametrize each interaction with **five** parameters:
 - m – mass of the resonance (1 TeV for benchmark study)
 - Γ – total width
 - $\Gamma_\phi = 0.5(m_\phi/TeV)^3$, $\Gamma_V = 5\%m_V$, $\Gamma_{\tilde{h}} = 1.2\%m_{\tilde{h}}$
 - ω^2 – cross section normalization factor
 - $\omega_\phi = 1$ recovers SM higgs
 - $\omega_V = 1$ recovers Z' with electroweak couplings
 - $\omega_{\tilde{h}} = 1$ recovers RS graviton
 - Chirality CP violation ...

Resonance Distributions in $m_{t\bar{t}}$

- (M_W, m_t) and small angle selection, respectively:

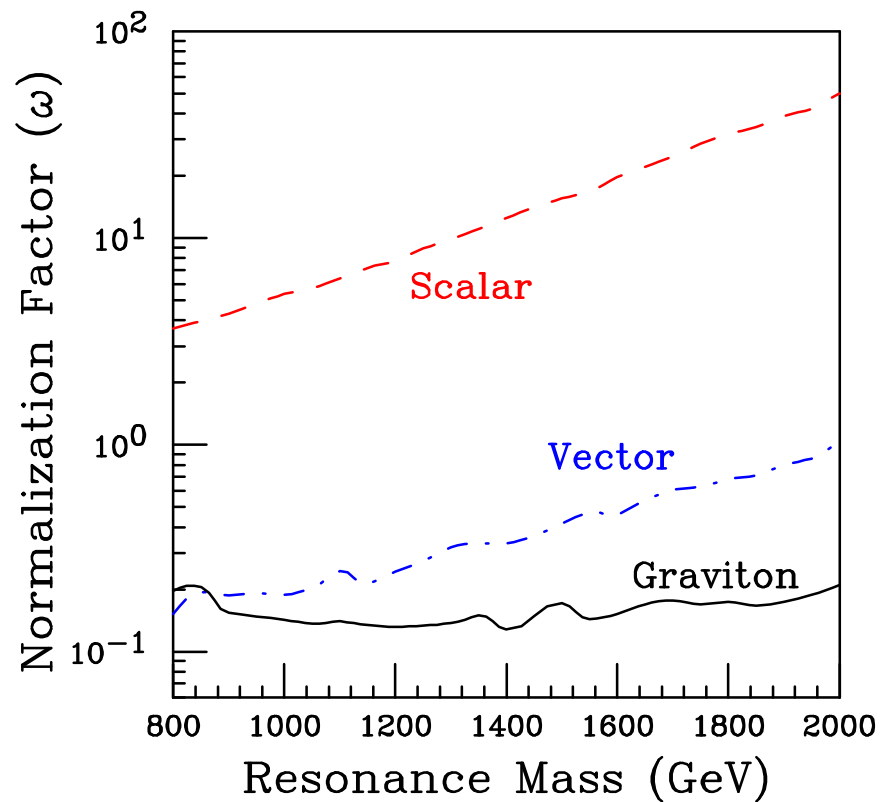


$$\Gamma_{tot} = 2\% m, 5\% m, 20\% m.$$

Discovery Potential

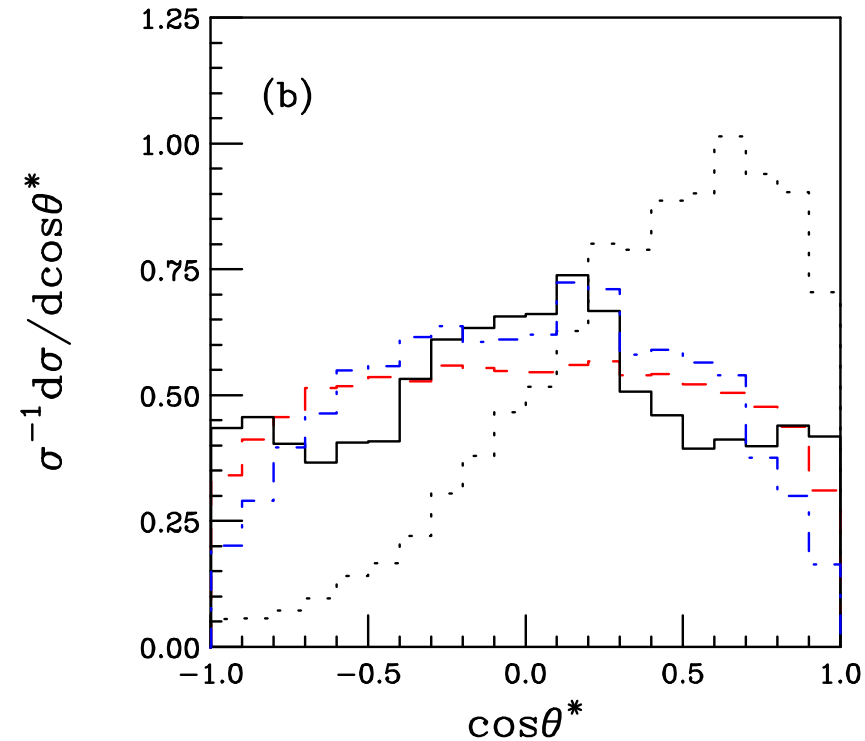
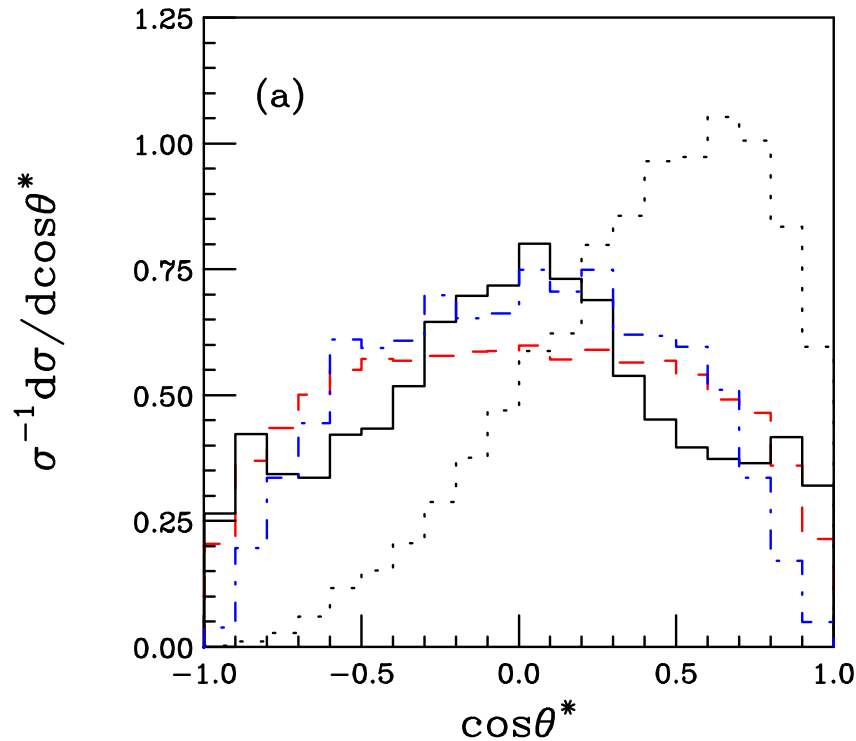
- Determine minimal ω for a 5σ discovery

$$S/\sqrt{B+S} = 5$$



Angular Distributions in $t\bar{t}$ c.m. frame

- (M_W, m_t) and small angle selection, respectively



Red dashed: scalar \rightarrow flat;

Black dots: Chiral vector $\rightarrow d_{11}^1 \Rightarrow (1 + \cos\theta^*)^2$;

Blue dash-dots: graviton from $gg \rightarrow d_{2\pm 1}^2 \Rightarrow \sin^2\theta^*$;

Black solid: graviton from $q\bar{q} \rightarrow d_{1\pm 1}^2 \Rightarrow \sin^4\theta^* + \dots$

- Forward/Backward Asymmetry for Parity-Violation:

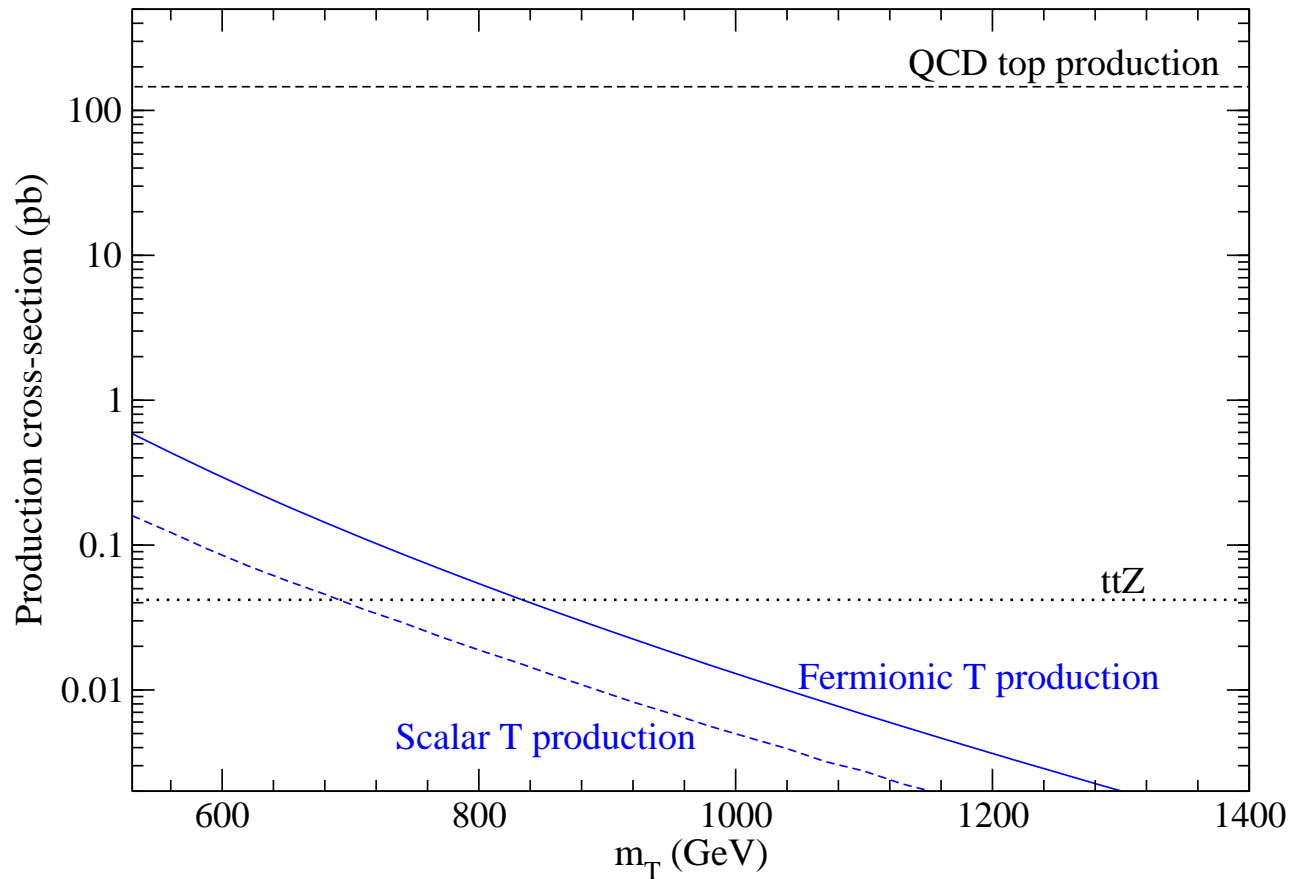
$$A_f^{\text{had}} = \frac{N_F - N_B}{N_F + N_B}$$

- N_F (N_B) is the number of events with the top quark momentum \vec{p}_{top} in the forward (backward) direction defined relative to the quark moving direction \vec{p}_q ,
 - “Forward” for the final state top is thus defined relative to the boost direction from the resonance c.m. frame (because the valence quarks tend to carry a higher-momentum fractions than the sea (anti-) quarks.)
 - Gluon contributions are homogeneous and subtracted out.
- Similarly, one may consider **CP-asymmetry**.
(work in progress.)

Search II: $t\bar{t} + \cancel{E}_T$ Signal

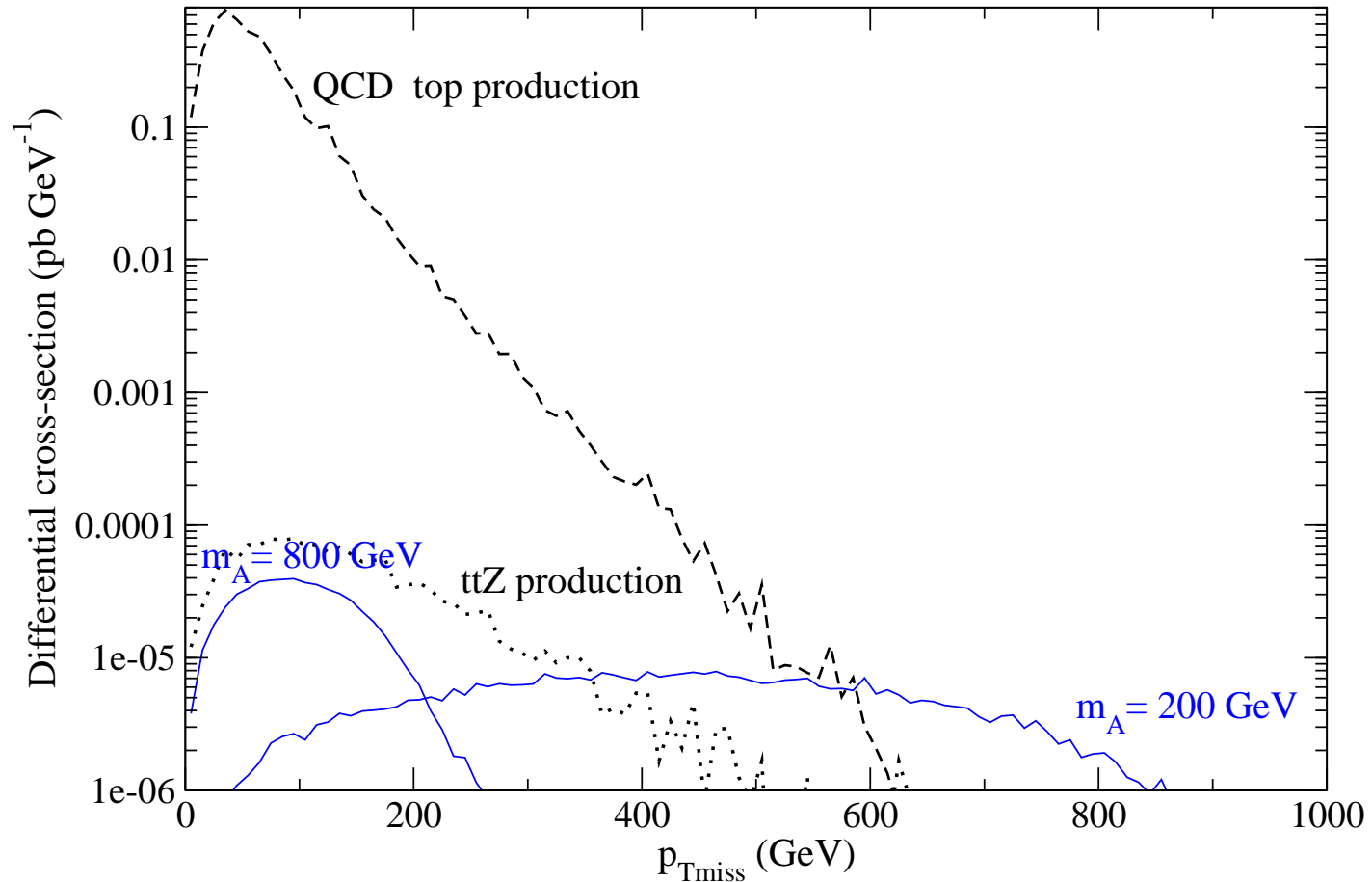
Quite generically,

$$\begin{aligned} pp &\rightarrow T\bar{T}X \rightarrow t\bar{t} A^0 A^0 X \\ &\rightarrow t\bar{t} A^0 A^0 X \rightarrow bj_1 j_2 \bar{b} \ell^- \bar{\nu} A^0 A^0 X + c.c. \end{aligned}$$



Features: No Bump, but much \cancel{E}_T

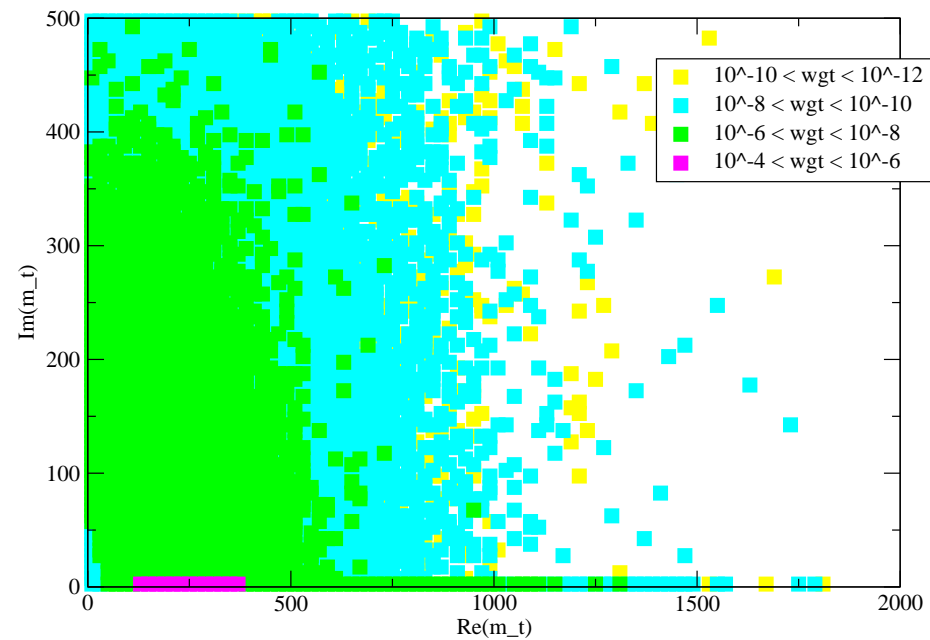
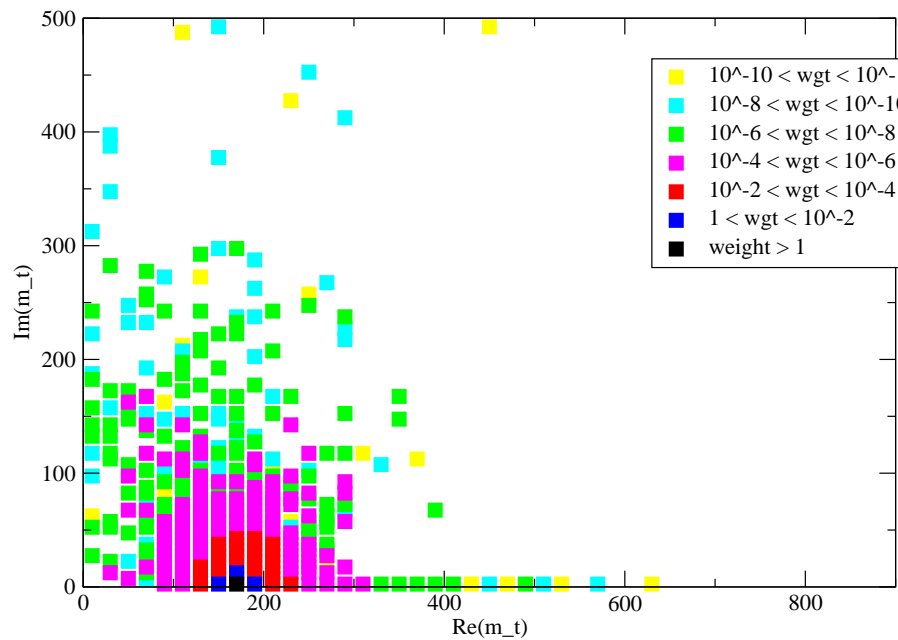
Due to more missing particles from both T and \bar{T} , no p_ν can be reconstructed. Instead, may lead to larger \cancel{E}_T :



$$m_T = 1 \text{ TeV.}$$

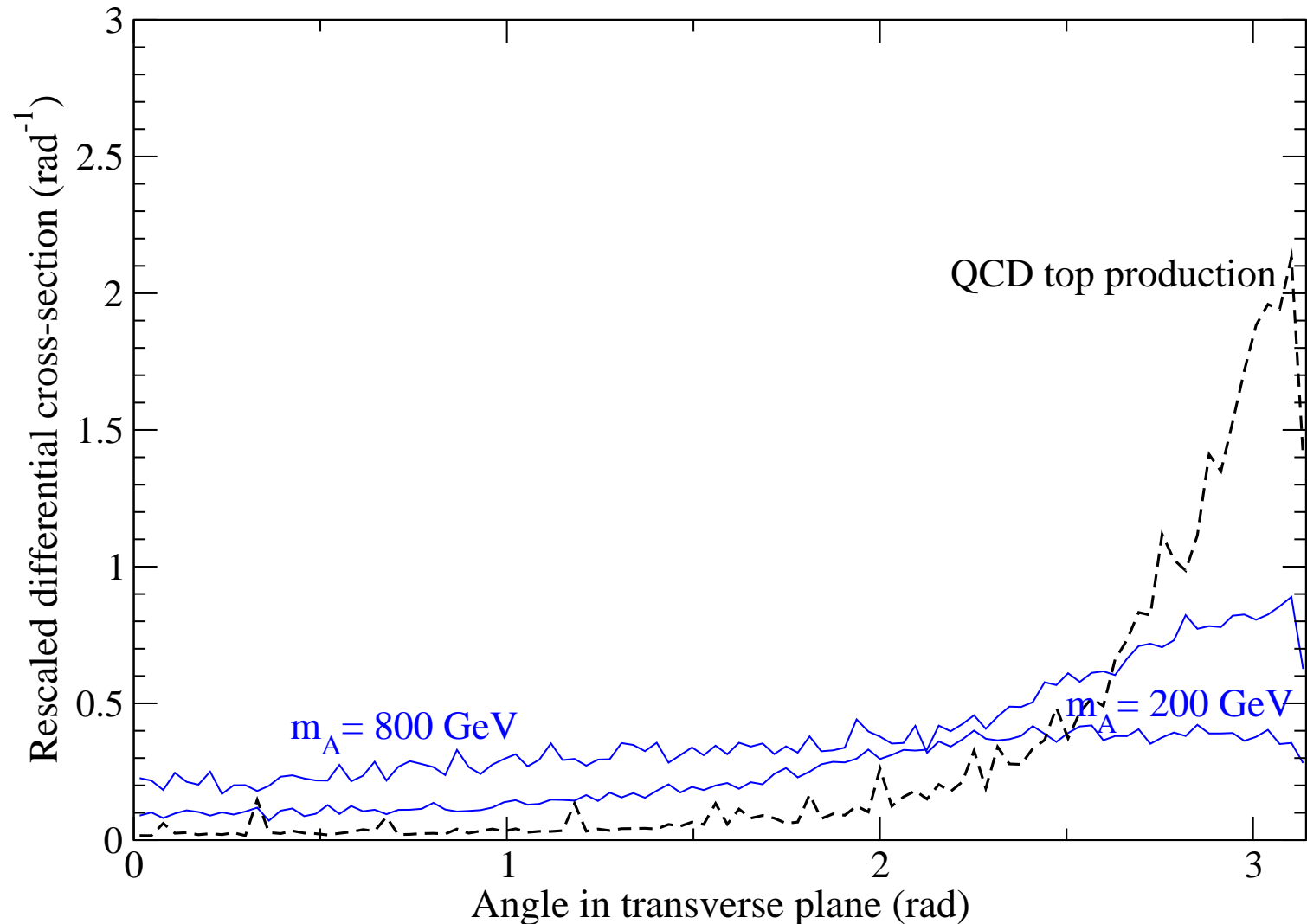
It's Good or Bad: A complex m_t^r

p_ν reconstruction by M_W may yield complex solutions:



Back-to-back $t - \bar{t}$

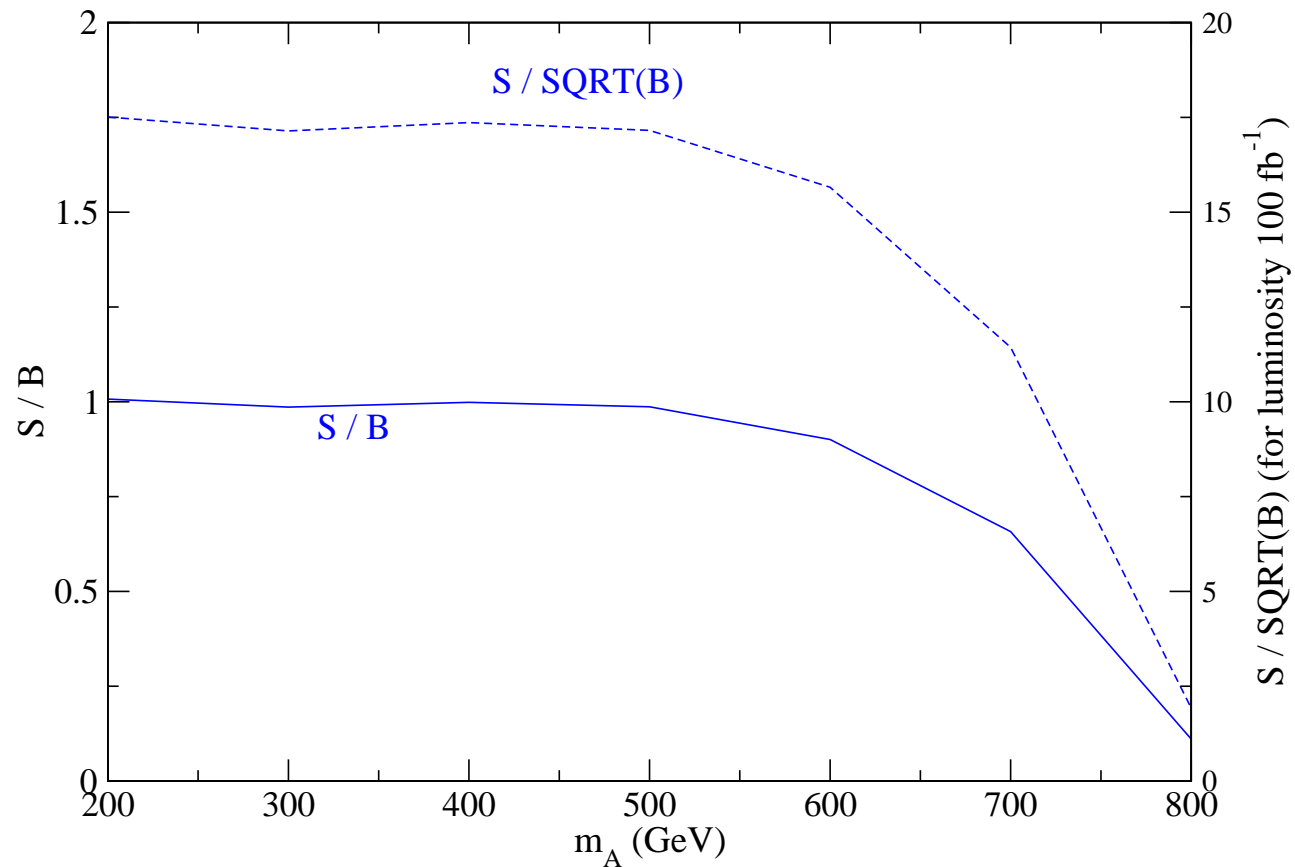
SM $t - \bar{t}$ are back-to-back in the transverse plane, while those from $T\bar{T}$ are kicked randomly:



LHC Reach for $T\bar{T}$ Signal

After judicious cuts, plus

$$|m_t - m_t^r|^2 > 110 \text{ GeV},$$



$$m_T = 1 \text{ TeV}.$$

Conclusions

- LHC is a top factory – providing **8 million $t\bar{t}$'s per 10 fb^{-1}** .
Good channel to probe physics beyond SM.
May serve as an **early indicator for new physics**.
- For the resonant signal $t\bar{t}$: Two methods to reconstruct **semi-leptonic $t\bar{t}$ events** at high-invariant mass,
to study **resonant spin, Chiral couplings, CP properties ...**
- For non-resonant signal $t\bar{t} + \cancel{E}_T$: Observation of the semi-leptonic channel promising, but kinematics difficult:
No information for the missing particle mass m_A, m_{χ^0} .

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Top quark studies are of high priority!