

The universal galaxy cluster pressure profile from a representative sample of nearby systems (REXCESS) and the $Y_{\text{SZ}} - M_{500}$ relation

M. Arnaud¹, G. W. Pratt^{1,2}, R. Piffaretti¹, H. Böhringer², J. H. Croston³, and E. Pointecouteau⁴

¹ Laboratoire AIM, IRFU/Service d'Astrophysique – CEA/DSM – CNRS – Université Paris Diderot, Bât. 709, CEA-Saclay, 91191 Gif-sur-Yvette Cedex, France
e-mail: Monique.Arnaud@cea.fr

² Max-Planck-Institut für extraterrestrische Physik, Giessenbachstraße, 85748 Garching, Germany

³ School of Physics and Astronomy, University of Southampton, Southampton, SO17 1BJ, UK

⁴ Université de Toulouse, CNRS, CESR, 9 Av. du colonel Roche, BP 44346, 31028 Toulouse Cedex 04, France

Received 7 October 2009 / Accepted 10 April 2010

ABSTRACT

We investigate the regularity of cluster pressure profiles with REXCESS, a representative sample of 33 local ($z < 0.2$) clusters drawn from the REFLEX catalogue and observed with *XMM-Newton*. The sample spans a mass range of $10^{14} M_{\odot} < M_{500} < 10^{15} M_{\odot}$, where M_{500} is the mass corresponding to a density contrast of 500. We derive an average profile from observations scaled by mass and redshift according to the standard self-similar model, and find that the dispersion about the mean is remarkably low, at less than 30 per cent beyond $0.2 R_{500}$, but increases towards the center. Deviations about the mean are related to both the mass and the thermo-dynamical state of the cluster. Morphologically disturbed systems have systematically shallower profiles while cooling core systems are more concentrated. The scaled profiles exhibit a residual mass dependence with a slope of ~ 0.12 , consistent with that expected from the empirically-derived slope of the $M_{500} - Y_X$ relation; however, the departure from standard scaling decreases with radius and is consistent with zero at R_{500} . The scatter in the core and departure from self-similar mass scaling is smaller compared to that of the entropy profiles, showing that the pressure is the quantity least affected by dynamical history and non-gravitational physics. Comparison with scaled data from several state of the art numerical simulations shows good agreement outside the core. Combining the observational data in the radial range $[0.03-1] R_{500}$ with simulation data in the radial range $[1-4] R_{500}$, we derive a robust measure of the universal pressure profile, that, in an analytical form, defines the physical pressure profile of clusters as a function of mass and redshift up to the cluster “boundary”. Using this profile and direct spherical integration of the observed pressure profiles, we estimate the integrated Compton parameter Y and investigate its scaling with M_{500} and L_X , the soft band X-ray luminosity. We consider both the spherically integrated quantity, $Y_{\text{sph}}(R)$, proportional to the gas thermal energy, and the cylindrically integrated quantity, $Y_{\text{cyl}}(R) = Y_{\text{SZ}} D_A^2$, which is directly related to the Sunyaev-Zel’dovich (SZ) effect signal. From the low scatter of the observed $Y_{\text{sph}}(R_{500}) - Y_X$ relation we show that variations in pressure profile shape do not introduce extra scatter into the $Y_{\text{sph}}(R_{500}) - M_{500}$ relation as compared to that from the $Y_X - M_{500}$ relation. The $Y_{\text{sph}}(R_{500}) - M_{500}$ and $Y_{\text{sph}}(R_{500}) - L_X$ relations derived from the data are in excellent agreement with those expected from the universal profile. This profile is used to derive the expected $Y_{\text{SZ}} - M_{500}$ and $Y_{\text{SZ}} - L_X$ relations for any aperture.

Key words. cosmology: observations – dark matter – galaxies: clusters: general – intergalactic medium – X-rays: galaxies: clusters

1. Introduction

Galaxy clusters provide valuable information on cosmology, from the nature of dark energy to the physics driving galaxy and structure formation. Clusters are filled with a hot ionised gas that can be studied both in X-ray and through the thermal Sunyaev-Zel’dovich (SZ) effect, a spectral distortion of the cosmic microwave background (CMB) generated via inverse Compton scattering of CMB photons by the free electrons. Its magnitude is proportional to the Compton parameter y , a measure of the gas pressure integrated along the line-of-sight, $y = (\sigma_T/m_e c^2) \int P dl$, where σ_T is the Thomson cross-section, c the speed of light, m_e the electron rest mass and $P = n_e T$ is the product of the electron number density and temperature. The total SZ signal, integrated over the cluster extent, is proportional to the integrated Compton parameter Y_{SZ} , $Y_{\text{SZ}} D_A^2 = (\sigma_T/m_e c^2) \int P dV$, where D_A is the angular distance to the system.

As the gas pressure is directly related to the depth of the gravitational potential, $Y_{\text{SZ}} D_A^2$ is expected to be closely related

to the mass. Numerical simulations (e.g., da Silva et al. 2004; Motl et al. 2005; Nagai 2006; Bonaldi et al. 2007) and analytical models (Reid & Spergel 2006) of cluster formation indicate that the intrinsic scatter of the $Y_{\text{SZ}} - M$ relation is low, regardless of the cluster dynamical state (see also Wik et al. 2008) or the exact details of the gas physics. However, the normalisation of the relation *does* depend on the gas physics (Nagai 2006; Bonaldi et al. 2007), as does the exact amount of scatter, the details of which are still under debate (Shaw et al. 2008). Given that this relation, and the underlying pressure profile, are key ingredients for the use of on-going or future SZ cluster surveys for cosmology, and provide invaluable information on the physics of the intra-cluster medium (ICM), it is important to calibrate these quantities from observations.

In recent years, SZ observational capability has made spectacular progress, from the first spatially resolved (single-dish) observations of individual objects (Pointecouteau et al. 1999, 2001; Komatsu et al. 1999, 2001) to the first discovery of new clusters with a blind SZ survey (Staniszewski et al. 2009). Spatially resolved SZE observations directly probe the mass

weighted temperature along the line of sight. By contrast, temperatures derived from X-ray spectra, by fitting an isothermal model to a multi-temperature plasma emission along the line of sight, are likely to be biased (Mathiesen & Evrard 2001). Although schemes to correct for this effect have been defined (Mazzotta et al. 2004; Vikhlinin 2006), it remains a potential source of systematics.

Stacking analysis of WMAP data around known X-ray clusters has allowed statistical detection of a scaled pressure profile (Afshordi et al. 2007) or a spatially resolved decrement (Lieu et al. 2006; Atrio-Barandela et al. 2008; Diego & Partridge 2010), showing clear discrepancies with the prediction of a simple isothermal β -model. Pressure or temperature profiles of individual clusters have started to be derived from combined analysis of X-ray and SZE imaging data, using non-parametric deprojection methods (Nord et al. 2009) or more realistic models than the β -model (Kitayama et al. 2004; Mroczkowski et al. 2009). Interestingly, the profiles are found to be consistent with profiles derived using X-ray spectroscopic data (see also Jia et al. 2008; Halverson et al. 2009). However, such studies are still restricted to a few test cases, particularly hot clusters.

The $Y_{SZ}-M$ relation has been recently derived by Bonamente et al. (2008), an important step forward as compared to previous work based on central decrement measurements using heterogeneous data sets (McCarthy et al. 2003; Morandi et al. 2007); however, quantities were estimated within $R_{2500} \sim 0.44 R_{500}$ ¹ and assuming an isothermal β -model, which may provide a biased estimate (Hallman et al. 2007). In addition, the first scaling relation using weak lensing masses, rather than X-ray hydrostatic masses, has now appeared (Marrone et al. 2009), although constraints from these data are currently weak.

In this context, statistically more precise, albeit indirect, information can be obtained from X-ray observations. A key physical parameter is Y_X , the X-ray analogue of the integrated Compton parameter, introduced by Kravtsov et al. (2006). Y_X is defined as the product of $M_{g,500}$, the gas mass within R_{500} and T_X , the spectroscopic temperature outside the core. The local $M_{500} - Y_X$ relation for relaxed clusters has recently been calibrated (Nagai et al. 2007; Maughan 2007; Arnaud et al. 2007; Vikhlinin et al. 2009), with excellent agreement achieved between various observations (e.g., see Arnaud et al. 2007). However, the link between Y_X and Y_{SZ} depends on cluster structure through

$$\frac{Y_{SZ} D_A^2}{Y_X} = \frac{\sigma_T}{m_e c^2} \frac{1}{\mu_e m_p} \frac{\langle n_e T \rangle}{\langle n_e \rangle_{R_{500}} T_X} \quad (1)$$

where the angle brackets denote volume averaged quantities. From Eq. (1), it is clear that an understanding of the radial pressure distribution and its scaling is important not only as a probe of the ICM physics, but also for exploitation of these data. High resolution measurements of the radial density and temperature distribution are now routinely available from X-ray observations but the pressure profile structure and scaling have been relatively little studied. The pressure profiles of groups have been studied by Finoguenov et al. (2006) and Johnson et al. (2009). In the cluster regime, Finoguenov et al. (2005) analysed the 2D pressure distribution in a flux-limited sample of 6 hot ($kT > 7$ keV) clusters at $z \sim 0.3$ showing fluctuations at the 30% level around

¹ Here and in the following, M_δ and R_δ are the total mass and radius corresponding to a density contrast, δ , as compared to $\rho_c(z)$, the critical density of the universe at the cluster redshift: $M_\delta = (4\pi/3)\delta\rho_c(z)R_\delta^3$. M_{500} corresponds roughly to the virialised portion of clusters, and is traditionally used to define the “total” mass.

the mean profile, scaled by temperature. To our knowledge, the only study of pressure profiles scaled by mass is that of Nagai et al. (2007), who used *Chandra* X-ray observations to derive a universal pressure profile, with the external slope derived from numerical simulations. However, their sample was restricted to hot ($kT > 5$ keV) relaxed clusters, which are all cool core systems, and contained five objects. For the reasons mentioned above, it is of considerable interest to extend this analysis to data from a larger and more representative sample of the cluster population.

In this paper we do this by investigating the regularity of cluster pressure profiles with REXCESS (Böhringer et al. 2007), a representative sample of 33 local ($z < 0.2$) clusters drawn from the REFLEX catalogue (Böhringer et al. 2004) and observed with *XMM-Newton*. We derive an average profile from observations scaled by mass and redshift according to the self-similar model and relate the deviations about the mean to both the mass and the thermo-dynamical state of the cluster (Sect. 3). Comparison with data from several state of the art numerical simulations (Sect. 4) shows good agreement outside the central regions, which is the most relevant aspect for the Y_{SZ} estimate. Combining the observational data in the radial range $[0.03-1] R_{500}$ with simulation data in the radial range $[1-4] R_{500}$ allows us to derive a robust measure of the universal pressure profile up to the cluster “boundary” (Sect. 5). Using this profile or direct spherical integration of the observed pressure profiles, we estimate the spherically and cylindrically integrated Compton parameter and investigate its scaling with Y_X , M_{500} and L_X , the soft band X-ray luminosity (Sect. 6).

We adopt a Λ CDM cosmology with $H_0 = 70$ km s⁻¹ Mpc⁻¹, $\Omega_M = 0.3$ and $\Omega_\Lambda = 0.7$. $h(z)$ is the ratio of the Hubble constant at redshift z to its present value, H_0 . T_X is the temperature measured in the $[0.15-0.75] R_{500}$ aperture. All scaling relations are derived using the BCES orthogonal regression method with bootstrap resampling (Akritas & Bershady 1996), and uncertainties are quoted throughout at the 68 per cent confidence level.

2. The REXCESS data set

A description of the REXCESS sample, including *XMM-Newton* observation details, can be found in Böhringer et al. (2007). The two clusters RXC J0956.4-1004 (the Abell 901/902 supercluster) and J2157.4-0747 (a bimodal cluster) are excluded from the present analysis. Cluster subsample classification follows the definitions described in Pratt et al. (2009): objects with center shift parameter $\langle w \rangle > 0.01 R_{500}$ are classified as morphologically disturbed, and those with central density $h(z)^{-2} n_{e,0} > 4 \times 10^{-2}$ cm⁻³ as cool core systems.

The gas density profiles, $n_e(r)$, were derived by Croston et al. (2008) from the surface brightness profiles using the non-parametric deprojection and PSF-deconvolution technique introduced by Croston et al. (2006). The density at any radius of interest is estimated by interpolation in the log-log plane. The procedure to extract the 2D temperature profiles is detailed in Pratt et al. (2009). The 3D profiles, $T(r)$, were derived by fitting convolved parametric models (Vikhlinin et al. 2006) to these data, taking into account projection and PSF effects (Pointecouteau et al. 2004) and weighting the contribution of temperature components to each ring as proposed by Vikhlinin (2006) to correct for the spectroscopic bias mentioned above. A Monte Carlo procedure is used to compute the errors, which are then corrected to take into account the fact that parametric models over-constrain the 3D profile. Full details will be given in a forthcoming paper. As the temperature profiles are

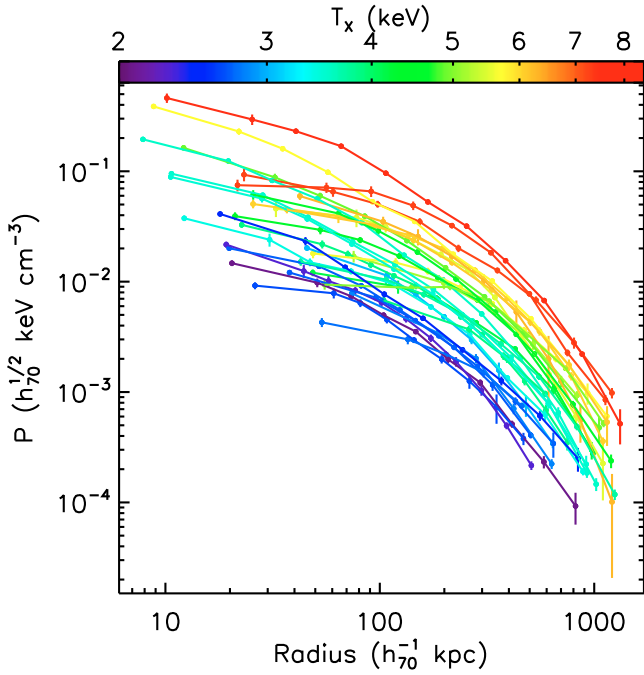


Fig. 1. The pressure profiles of the REXCESS sample. Pressures are estimated at the effective radii of the temperature profile (points with error bars). A line connects the data points for each cluster to guide the eye. The data are colour coded according to the spectroscopic temperature, T_X (see color bar).

measured on a lower resolution radial grid than the density profiles, the pressure profiles, $P(r) = n_e(r)T(r)$, are estimated at the weighted effective radii (Lewis et al. 2003) of each annular bin of the 2D temperature profiles. They are presented in Fig. 1.

Since the sample contains systems in a variety of dynamical states, we choose to use Y_X as a mass proxy rather than the hydrostatic mass. Extensive discussion of how this could affect our results is presented in Sect. 3.4. For each cluster, M_{500} is estimated iteratively from the $M_{500} - Y_X$ relation, as described in Kravtsov et al. (2006). We used the updated calibration of the $M_{500} - Y_X$ relation, obtained by combining the Arnaud et al. (2007) data on nearby relaxed clusters observed with *XMM-Newton* with new REXCESS data (Arnaud et al., in prep). The sample comprises 20 clusters: 8 clusters from Arnaud et al. (2007), excluding the two lowest mass clusters whose M_{500} estimate requires extrapolation, and the 12 relaxed REXCESS clusters with mass profiles measured at least down to $\delta = 550$. The derived $M_{500} - Y_X$ relation

$$h(z)^{2/5} M_{500} = 10^{14.567 \pm 0.010} \left[\frac{Y_X}{2 \times 10^{14} h_{70}^{-5/2} M_{\odot} \text{ keV}} \right]^{0.561 \pm 0.018} h_{70}^{-1} M_{\odot} \quad (2)$$

is consistent with the relation derived by Arnaud et al. (2007) but with improved accuracy on slope and normalization.

The slope differs from that expected in the standard self-similar model ($\alpha = 3/5$) by only $\sim 2\sigma$. We will thus also consider the $M_{500} - Y_X$ relation obtained by fixing the slope to its standard value:

$$h(z)^{2/5} M_{500} = 10^{14.561 \pm 0.009} \times \left[\frac{Y_X}{2 \times 10^{14} h_{70}^{-5/2} M_{\odot} \text{ keV}} \right]^{3/5} h_{70}^{-1} M_{\odot}. \quad (3)$$

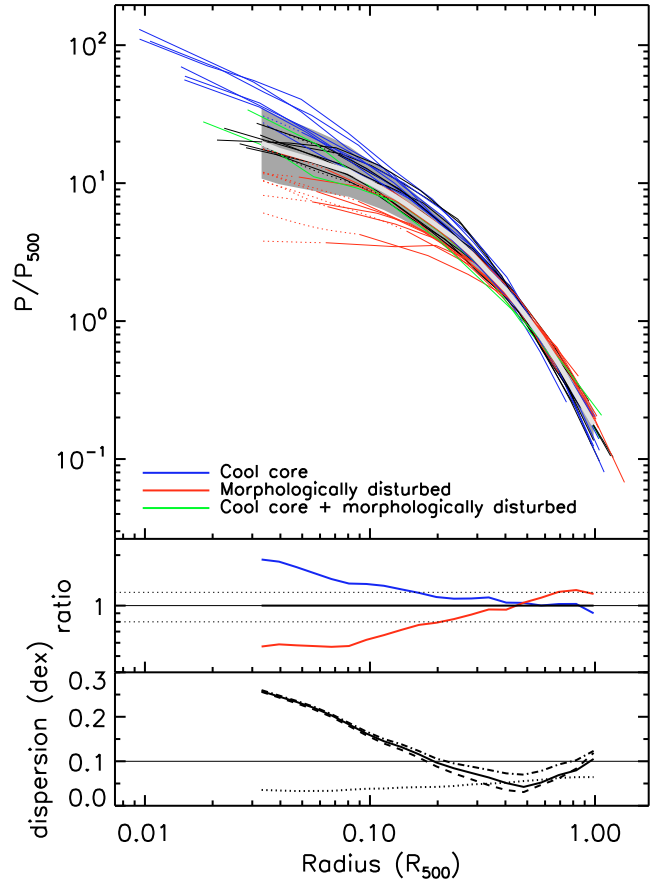


Fig. 2. The scaled pressure profiles of the REXCESS sample, colour coded according to the (thermo)dynamical state (see labels and Sect. 2). Black profiles denote clusters that are neither cool core nor morphologically disturbed. The radii are scaled to R_{500} and the pressure to P_{500} as defined in Eq. (5), with M_{500} estimated from the $M_{500} - Y_X$ relation (Eq. (2)). Full lines: measured pressure profile as in Fig. 1 with data points omitted for clarity. Dotted lines: extrapolated pressure (see text). The thick grey line is the average scaled profile and the grey area corresponds to the $\pm 1\sigma$ dispersion around it. *Middle panel:* ratio of the average profile of cool core (blue) and disturbed (red) systems to the overall average profile. *Bottom panel:* the solid line is the statistical dispersion as a function of scaled radius. Dotted line: additional dispersion expected from the intrinsic dispersion in the $M_{500} - Y_X$ relation. Dash-dotted line: quadratic sum of the two dispersions. Dashed line: dispersion obtained for M_{500} estimated from the standard slope $M_{500} - Y_X$ relation (Eq. (3)).

3. Scaled pressure profiles

3.1. Scaled profiles

The scaled pressure profiles

$$p(x) = \frac{P(r)}{P_{500}} \quad \text{where } x = \frac{r}{R_{500}} \quad (4)$$

are presented in Fig. 2. The pressure is normalised to the characteristic pressure P_{500} , reflecting the mass variation expected in the standard self-similar model, purely based on gravitation (Nagai et al. 2007, and Appendix A).

$$P_{500} = 1.65 \times 10^{-3} h(z)^{8/3} \times \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{2/3} h_{70}^2 \text{ keV cm}^{-3} \quad (5)$$

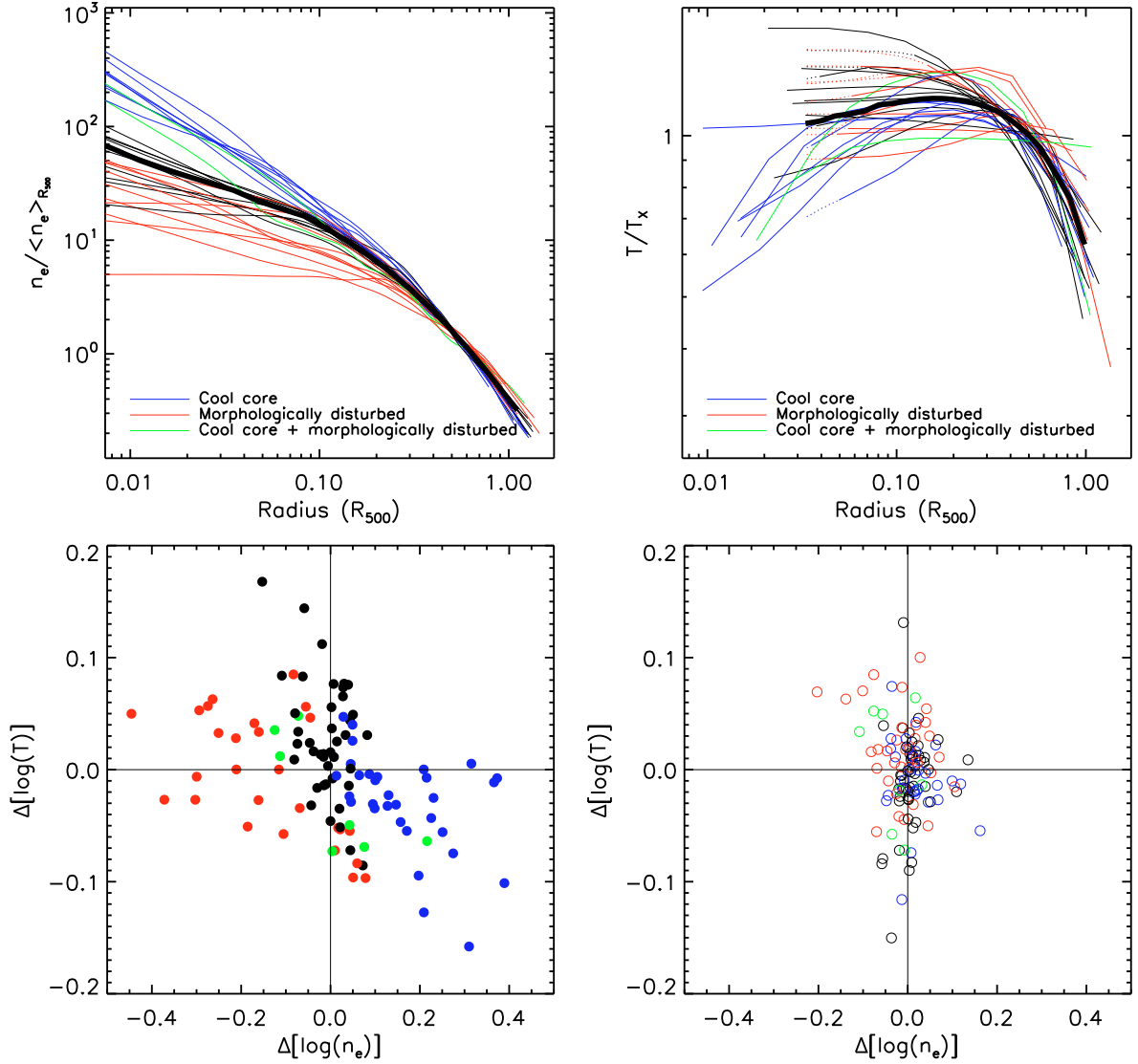


Fig. 3. The scaled density (*top left panel*) and temperature (*top right panel*) profiles of the REXCESS sample. Each profile is colour coded according to the cluster (thermo)dynamical state (see labels and Sect. 2). The radii are scaled to R_{500} , estimated from the $M_{500} - Y_X$ relation (Eq. (2)). The density is scaled to the mean density within R_{500} and the temperature to T_X , the spectroscopic temperature measured in the $[0.15-0.75] R_{500}$ aperture. In each panel, the thick black line is the average scaled profile. Dotted lines in the *top right panel*: extrapolated part of the profiles. *Bottom panels*: logarithmic deviation of the scaled temperatures from the average scaled profile versus the corresponding deviation for the density, at each effective radius of the temperature profile annular bins. Data corresponding to $r/R_{500} < 0.2$ and $r/R_{500} > 0.2$ are plotted in the *left and right panels*, respectively. The deviations are anti-correlated in the core.

For comparison we also plot in Fig. 3 the scaled temperature profiles, $t(x) = T(r)/T_X$ as well as the scaled density profiles, $\tilde{n}_e(x)$. Note that the density profiles have been normalised to the mean density within R_{500} , so that the dispersion is only due to variations in shape².

The resolution in the center and radial extent of the pressure profiles are determined by that of the temperature profiles, in practice the effective radius of the inner and outer annular temperature profile bins, which varies from cluster to cluster (see Fig. 2). In particular, the peaked emission of cool core clusters allows us to measure the profiles deeper into the core than for disturbed clusters, which have more diffuse emission (see also Sect. 3.3).

3.2. Average scaled pressure profile

We computed an average scaled pressure profile, $\mathbb{p}(x)$, from the median value of the scaled pressure in the radial range where data are available for at least 15 clusters without extrapolation (about $[0.03-1] R_{500}$). However, to avoid a biased estimate of the average profile in the core, where the dispersion is large and more peaked clusters are measured to lower radii (Fig. 2), it is important to include all clusters in the computation. For this purpose, we extrapolated the pressure profiles in the core using the best fitting temperature model used in the deprojection of the temperature profile. The extrapolated part of the profiles are plotted as dotted lines in Figs. 2 and 3. This extrapolation is only weakly model dependent since it essentially concerns disturbed clusters (Fig. 2), which are observed to have rather flat central temperature profiles (Fig. 3). The average profile is plotted as a thick line in Fig. 2. The dispersion around it is defined as the plus

² The normalisation of the density profiles, scaled according to the standard self-similar model, varies with mass as shown by Croston et al. (2008).

or minus standard deviation from the average profile, computed in the log – log plane.

3.3. Dispersion, radial structure and dynamical state

For a perfectly self-similar cluster population, the scaled profiles should coincide. The dispersion around the average scaled profile is less than 25% beyond the core ($r > 0.2 R_{500}$) and increases towards the center (Fig. 2, bottom panel). This dispersion reflects a variation of shape with cluster (thermo)dynamical state, as clearly seen in Fig. 2: shallower profiles, at all radii, are observed for morphologically disturbed clusters while the cooling core clusters have the most concentrated profiles. The typical difference between the average profiles of these two populations is $\sim 20\%$ in the outskirts and as high as a factor of four at $0.03 R_{500}$ (Fig. 2, middle panel).

When compared to the density profiles (Fig. 3, top-left panel) the pressure profiles are distinctly more regular and present less dispersion in the core. For instance, the dispersion at $0.04 R_{500}$ is 0.28 dex and 0.24 dex for the scaled density and pressure, respectively. The reason lies in the anti-correlation between the deviation of scaled temperatures and densities from their respective average scaled profiles, $\mathfrak{m}(x)$ and $\mathfrak{t}(x)$, as shown Fig. 3 (bottom-left panel). For data interior to $r < 0.2 R_{500}$, a Spearman rank test finds a probability of 10^{-7} that the anti-correlation between $\log(\bar{n}_e(x) - \log(\mathfrak{m}(x)))$ and $\log(t(x) - \log(\mathfrak{t}(x)))$ occurs by chance. The correlation disappears at large radii with a probability of 0.6 for $r > 0.2 R_{500}$ (Fig. 3, bottom-right panel). Qualitatively, this is the result of the well-known fact that cool core clusters have peaked density profiles (e.g., Jones & Forman 1984), with a temperature drop in the center, while unrelaxed objects have flatter density cores and constant or increasing temperature toward the center (Fig. 3, top panels).

3.4. Dependence on mass and mass-proxy relation

Since we derived M_{500} from the $M_{500} - Y_X$ relation, the scaling quantities R_{500} and P_{500} and the pressure profiles are not independent, as they are both related to the product of the gas density and temperature. We first examine how this may affect our results. From the definition of the pressure $P(r) = n_e(r)T(r)$, and noting that $P_{500} \propto M_{500}^{2/3}$ and that $Y_X = M_{g,500}T_X \propto \langle n_e(r) \rangle_{R_{500}} R_{500}^3 T_X \propto \langle n_e(r) \rangle_{R_{500}} T_X M_{500}$, where the angle brackets denote a volume average within R_{500} , the scaled pressure $p(x) = P(xR_{500})/P_{500}$ is proportional to

$$p(x) \propto \frac{P(xR_{500})}{\langle P(r) \rangle_{R_{500}}} \frac{\langle n_e(r)T(r) \rangle_{R_{500}}}{\langle n_e(r) \rangle_{R_{500}} T_X} \frac{Y_X}{M_{500}^{5/3}}. \quad (6)$$

This equation makes explicit the link between the scaled pressure profiles and the $M_{500} - Y_X$ relation. The first two dimensionless terms in the right hand part of the equation purely depend on the internal gas structure within R_{500} . They determine the average shape of the scaled profile. The third term depends on the global cluster scaling properties between Y_X and M_{500} and determine both the normalisation of the average scaled profile and the ‘‘typical’’ mass dependence of the profiles (discussed at the end of the section).

Using M_{500} values derived from the $M_{500} - Y_X$ relation, rather than the ‘‘true’’ M_{500} value, is equivalent to assuming a perfect correlation between M_{500} and Y_X , i.e. with no scatter. Provided that the correct $M_{500} - Y_X$ relation is used and that $\sigma_{\log,MY}$ does not depend on mass or dynamical state, use of the $M_{500} - Y_X$ relation will not introduce a systematic bias

into the scaled profiles, but their dispersion will be underestimated. Let us define the intrinsic scatter of the $M_{500} - Y_X$ relation, $\sigma_{\log,MY}$, as the standard deviation of $\log(M_{500})$ from the value from the best fitting relation at a given Y_X . We can estimate the additional dispersion due to $\sigma_{\log,MY}$ from the effect on the average scaled profile of a variation of $\log(M_{500})$ by $\pm\sigma_{\log,MY}$. Since $R_{500} \propto M_{500}^{1/3}$ and $P_{500} \propto M_{500}^{2/3}$, the profile is translated in the log – log plane by $\pm 1/3\sigma_{\log,MY}$ and $\pm 2/3\sigma_{\log,MY}$ along the x and y axis, respectively. Assuming $\sigma_{\log,MY} = 0.04$ (about 10%, Kravtsov et al. 2006; Arnaud et al. 2007), the additional dispersion (in dex units), computed from the difference between the translated profiles at a given scaled radius, is plotted in the bottom panel of Fig. 2. It is non-negligible beyond the core, but the total dispersion, estimated by summing quadratically this additional contribution, is expected to remain below 30%. It is negligible in the core, where the dispersion is dominated by structural variations.

Finally, the $M_{500} - Y_X$ relation being derived from mass estimated using the hydrostatic equilibrium, we expect an offset between that relation and the ‘‘true’’ $M_{500} - Y_X$ relation. The M_{500} used in this study are thus likely to be underestimated. The effect of such a bias is to translate all the scaled profiles together (provided that it is a simple factor independent of mass). This will not affect any shape or dispersion analysis but the normalisation of the mean scaled profile will be biased high. This is further discussed in Sects. 4.3 and 7.3.

We now turn to the question of the variation of the pressure profiles with mass. From the definition of P_{500} , any deviation from the standard self-similar scaling will appear as a variation of the scaled profiles $p(x)$ with mass, $p(x) \equiv p(x, M_{500})$. It will also translate into a non-standard slope α_{MY_X} for the $M_{500} - Y_X$ relation. From Eq. (6), and assuming that the shape does not vary with mass, we expect that the normalisation of $p(x)$ increases slightly with mass as $Y_X/M_{500}^{5/3}$, i.e. as $M_{500}^{\alpha_P}$ with:

$$\alpha_P = \frac{1}{\alpha_{MY_X}} - \frac{5}{3} = 0.12 \quad (7)$$

for the best fitting slope $\alpha_{MY_X} = 0.561$ (Eq. (2)).

We show in the left-top panel of Fig. 4 the scaled profiles colour coded as a function of T_X . There is some indication that hotter (thus more massive) clusters lie above cooler systems. To better quantify this, the right-top panel of the figure shows the variation with M_{500} of the scaled pressure, $p(x)$, for different scaled radii, $x = r/R_{500}$. At each radius, we fitted the data with a power law $p(x) \propto M_{500}^{\alpha(x)}$. The pivot of the power law, where the pressure equals the average scaled value, $p(x) = \mathbb{P}(x)$, is about $M_{500} \sim 3 \times 10^{14} M_{\odot}$. Although the slopes at various radii are consistent within the errors (Fig. 4, right top panel), we note a systematic decrease with radius from $\alpha(x) = 0.22 \pm 0.16$ at $r = 0.1 R_{500}$ to $\alpha(x) = -0.01 \pm 0.16$ at R_{500} . This variation can be adequately represented by the analytical expression, $\alpha(x) = \alpha_P + \alpha'_P(x)$, with:

$$\alpha'_P(x) = 0.10 - (\alpha_P + 0.10) \frac{(x/0.5)^3}{1. + (x/0.5)^3} \quad (8)$$

yielding to scaled profiles varying with mass as:

$$p(x, M_{500}) = \mathbb{P}(x) \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{\alpha_P + \alpha'_P(x)}. \quad (9)$$

This corresponds to a break of self-similarity in shape: the departure from standard mass scaling, likely to be due to the effects

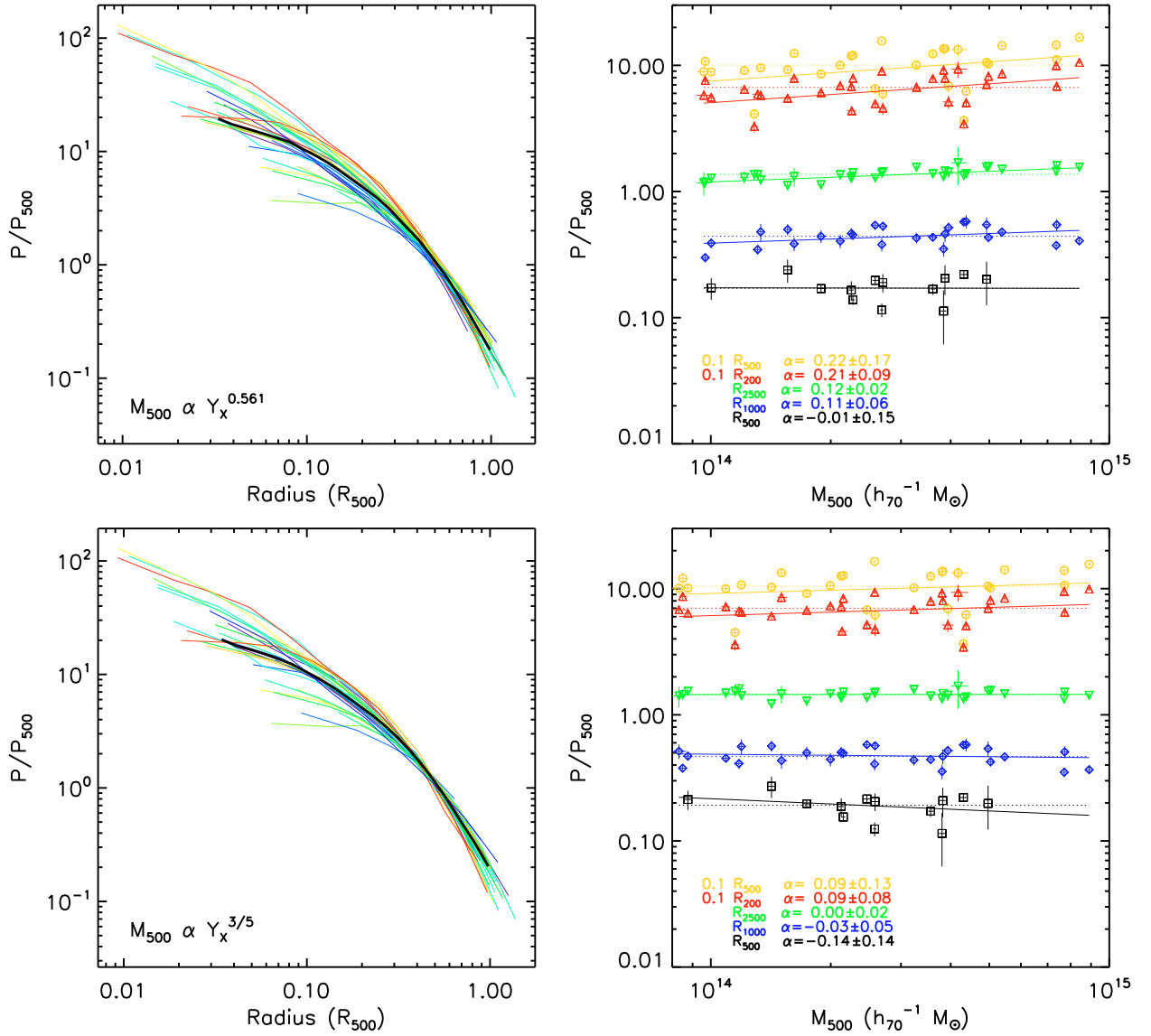


Fig. 4. The scaled pressure profiles of the REXCESS sample, colour coded according to spectroscopic temperature measured in the [0.15–0.75] R_{500} aperture (*left panels*, same coding as in Fig. 1). *Right panels*: corresponding scaled pressure estimated at different values of scaled radii plotted as a function of cluster mass. Full lines: power law fit at each scaled radius, with the best fitting slope given in the labels. Dotted line: value for the average scaled profile at that radius. The cluster masses, M_{500} , are estimated from the $M_{500} - Y_X$ relation, either the best fitting empirical relation (Eq. (2), *top panels*) or the relation obtained from fixing the slope to its standard value (Eq. (3), *bottom panels*).

of non-gravitational processes, becomes less pronounced as we move towards the cluster outskirts and is consistent with zero at R_{500} . Such a behaviour was also noticed in the entropy profiles (Nagai et al. 2007; Pratt et al. 2010).

Note, however, that the mass dependence is weaker for the pressure than for the entropy: the pressure slopes are about two times smaller than those of the entropy (Fig. 4 and Pratt et al. 2010, their Fig. 3) and the break of self-similarity has a low significance. The comparison with a constant slope model gives a F-test probability of 0.2. The mean slope (0.10 ± 0.02) and the slope $\alpha(x)$ at all radii are consistent with the expected 0.12 value. In a good approximation, the mass dependence of the scaled profiles can then be modelled by a simple variation in normalisation:

$$p(x, M_{500}) = \mathbb{P}(x) \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{\alpha_p=0.12} \quad (10)$$

where $\mathbb{P}(x)$ is the average scaled profile derived in Sect. 3.2.

We then compared to the results obtained using M_{500} derived from the self-similar $M_{500} - Y_X$ relation with slope 3/5 (Eq. (3)). The scaled profiles are plotted in the bottom panel of Fig. 4. In this case, we do not expect any dependence of $p(x)$ with M_{500} , and this is indeed the case: the slopes $\alpha_p(x)$ are consistent with zero at all radii (right bottom panel). The dispersion in scaled profiles is also smaller (see Fig. 2 bottom panel). In that case, the dispersion is only due to structural variations, while in the non-standard case, the mass dependence of $p(x)$ also contributes to the dispersion.

3.5. Comparison to Chandra results for relaxed clusters

In Fig. 5, we plot the pressure profiles presented in Nagai et al. (2007), derived from *Chandra* data analyzed by Vikhlinin et al. (2006). We only consider clusters with measured M_{500} values, excluding MKW4 ($T = 1.4$ keV) and A2390 ($z = 0.23$) which fall outside the T_X and z range of REXCESS, respectively. We

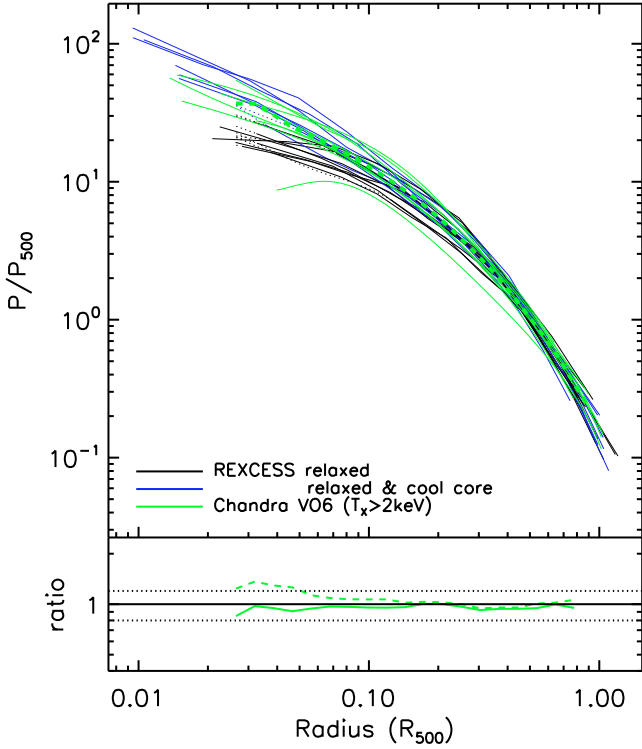


Fig. 5. The scaled pressure profiles (green lines) derived from Vikhlinin et al. (2006) *Chandra* data on relaxed clusters compared to the scaled profiles of the REXCESS sample excluding morphologically disturbed clusters (same colour code as in Fig. 2). The thick green dotted line is the average *Chandra* profile. *Bottom panel:* ratio of that average *Chandra* profile to that of REXCESS for all morphologically undisturbed objects (dotted line) or only cool core clusters (full line).

used the published M_{500} values, derived from the hydrostatic equilibrium (HSE) equation, and computed the pressure from the best fitting parametric models of the density and temperature profile given in Vikhlinin et al. (2006), in the radial range of the observed temperature profile. Since the *Chandra* data set only contains relaxed clusters, they are compared to the REXCESS profiles excluding morphologically disturbed objects.

All *Chandra* profiles, except one³, lie within the range of the REXCESS profiles. The bottom panel of Fig. 5 shows the ratio of the average *Chandra* profile to the average REXCESS profile. The agreement between the average profiles, both in shape and normalisation, is nearly perfect beyond the core, where the dispersion of the scaled profiles is lower. However, on average, the *Chandra* profiles are slightly more peaked towards the center (dotted line in bottom panel of Fig. 5) and have a smaller dispersion than the “relaxed” REXCESS clusters. Better agreement is found with the average REXCESS profile for cool core clusters (full line in bottom panel of Fig. 5). This is not surprising, since all clusters in the *Chandra* data set present the central temperature drop characteristic of cool core clusters.

This good agreement is an indication of the robustness of scaled pressure profile measurements with current X-ray satellites. The comparison also illustrates the importance of

considering a representative cluster sample to measure the average profile and dispersion in the core.

4. Comparison with numerical simulations

4.1. The data set

We consider three large samples of simulated clusters at redshift zero extracted from Λ CDM cosmological N -body/hydrodynamical simulations ($\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$). The data set includes the samples from Borgani et al. (2004, hereafter BO), Piffaretti & Valdarnini (2008, PV) and Nagai et al. (2007, NA). All simulations include treatment of radiative cooling, star formation, and energy feedback from supernova explosions. The three simulated data sets are fully independent and derived using different numerical schemes and implementations of the gas physics (see references above for full description). This allows us to check the robustness of the theoretical predictions of the pressure profiles by comparing the three simulated data sets. The fact that the NA simulation was undertaken on a mesh-based Eulerian code, while the PV and BO simulations were derived from particle-based Lagrangian codes is particularly relevant, considering some well known cluster-scale discrepancies between the numerical approaches, such as is seen in the entropy profiles (see, e.g., Voit et al. 2005; Mitchell et al. 2009, and references therein). The star formation algorithm and the SN feedback model are also quite different both in implementation and in feedback efficiency.

In order to avoid comparison with inappropriately low mass objects we impose the REXCESS lower mass limit $M_{500} \geq 10^{14} M_\odot$, leading to a final number of simulated clusters of 93, 88, and 14 for the BO, PV, and NA samples, respectively. We computed the pressure profile for each cluster using the mass-weighted gas temperature, since the deprojection of the observed profile takes into account the spectroscopic bias (Sect. 2). The assumed baryon densities are $\Omega_b = 0.039, 0.049, 0.043$ for the BO, PV, and NA samples, respectively. The assumed baryon fraction, $f_b = \Omega_b / \Omega_m$ has a direct impact on the gas density and thus pressure profile at a given total mass. We thus corrected the gas profiles by the ratio between the assumed f_b value and the WMAP5 value (Dunkley et al. 2009) for each sample. To scale each individual pressure profile we consider both the “true” R_{500}^{Sim} and M_{500}^{Sim} values and the hydrostatic values R_{500}^{HSE} and $M_{500}^{\text{HSE}} = M_{500}^{\text{HSE}}(< R_{500}^{\text{HSE}})$. The former are derived from the total mass distribution in the simulation. The latter was derived from the gas density and temperature profiles and the hydrostatic equilibrium equation, using the same procedure for all clusters. As in previous work (e.g., Piffaretti & Valdarnini 2008, and references therein), we find that M_{500}^{HSE} underestimates the true mass. We find a mean bias for the whole sample of -13 per cent with a dispersion of ± 16 per cent; the average bias estimated for the different simulations agrees within a few percent at all radii larger than $0.1 R_{500}$.

4.2. Comparison of numerical simulations

We derive the average scaled profile for each simulation, and the dispersion around it, from the median value and 16 and 84 percentiles of the scaled pressure distribution at a given scaled radius. We also compute an average simulation profile. Since the average profile computed from the total sample would be biased by the number of objects in the largest data set, we average the three mean profiles from each simulation data set, and calculate

³ The outlier is A133, a relaxed cooling core cluster (Vikhlinin et al. 2005). This cluster appears to present a general deficit of gas. Its pressure at all radii is low as compared to other clusters, as well as its gas mass fraction ($f_{\text{gas},500} = 0.083 \pm 0.006$ to be compared to a weighted mean for the rest of *Chandra* sample of $0.115 \pm 0.010[\text{stdev}]$).

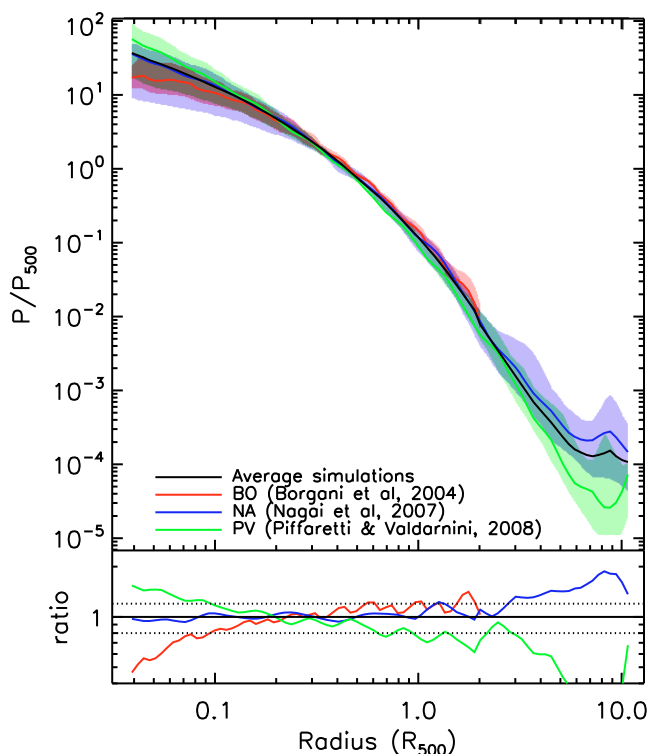


Fig. 6. The scaled pressure profiles derived from numerical simulations of Borgani et al. (2004) (pink), Nagai et al. (2007) (blue) and Piffaretti & Valdarnini (2008) (green). Black line: overall average profile (see text). coloured lines: average profile for each simulation with the coloured area corresponding to the dispersion around it. *Bottom panel:* ratio of each simulation average profile to the overall average profile.

the dispersion from all available profiles. The results derived using the true mass are shown in Fig. 6.

Taking into account that the profiles vary by more than 5 orders of magnitude from the cluster center to the outskirts, the agreement between the three simulations is exceptionally good. The profiles agree within 20% between ~ 0.1 and $\sim 3 R_{500}^{\text{Sim}}$ (Fig. 6 lower panel). As expected, larger differences are found in the core, where non-gravitational processes are more important and where the differences in their implementation in the codes will become more evident. The BO profiles are available only up to the “virial” radius, $\approx 2.03 R_{500}^{\text{Sim}}$ but the PV and NA profiles are traced up to $10 R_{500}^{\text{Sim}}$, where they deviate significantly, but still agree within the dispersion. However, the differences are systematic with the PV profiles lying below the NA profiles. This may hint at a difference in the way in which Lagrangian and Eulerian codes behave in the IGM-WHIM regime. Note also the flattening of the pressure profile in the outskirts, around $5 R_{500}^{\text{Sim}}$, which is likely to define the actual boundary of the cluster, where it meets the intergalactic medium. In the following we will use this boundary to compute the total integrated SZ signal, Y_{SZ} . In spite of the difference in the pressure in the outskirts, there is good agreement on Y_{SZ} between the simulations: the SZ signal within $5 R_{500}$ computed from the average PV and NA profiles differ by -15% , and $+9\%$, respectively, from the value computed using the average simulation profile.

4.3. Comparison of REXCESS profiles with simulations

Figure 7 compares the observed scaled profiles with the prediction of the simulations. We first consider the simulated profiles

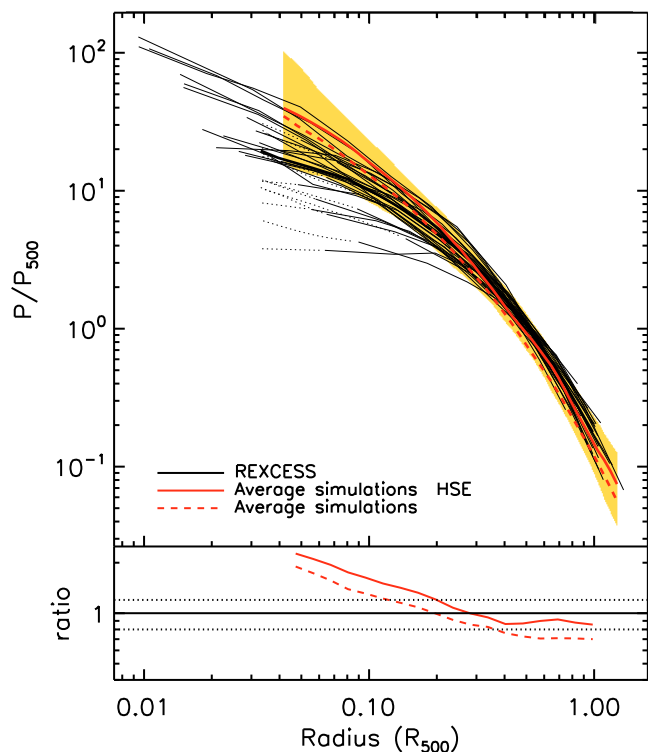


Fig. 7. Comparison of the REXCESS scaled profiles with the prediction of numerical simulations. Black lines: REXCESS data (as in Fig. 2). Thick black line: average REXCESS scaled profile. Red line: average simulation profile and dispersion around it (orange area) using the hydrostatic mass. Dotted red line: same using the true mass. *Bottom panel:* ratio of these average simulation profiles to the REXCESS average profile.

scaled using the hydrostatic quantities R_{500}^{HSE} and M_{500}^{HSE} , since the observations rely on hydrostatic mass estimates. Note that we used the $M_{500} - Y_X$ relation calibrated from a sample of relaxed clusters, while for the simulations we used R_{500}^{HSE} and M_{500}^{HSE} for the whole sample. However we checked that, when considering only relaxed clusters, the median bias on M_{500} changes by only 2%, the main effect being a factor of 2 decrease in its dispersion.

The simulation prediction and the REXCESS data agree well in the external part ($r \gtrsim 0.2 R_{500}$), with the observed profiles lying within the dispersion around the average simulation profile (Fig. 7). Remarkably, the observed and simulated average profiles are parallel above $0.4 R_{500}$ (i.e. they have the same shape), with a normalisation offset of only $\sim 10\%$ (Fig. 7, bottom panel). The slight underestimate of the pressure in the simulations is similar to the offset observed for the $M_{500} - Y_X$ relation and may be due, at least in part, to over-condensation of hot gas in the cold dense phase (see discussion in Arnaud et al. 2007). As we move towards the center, the agreement progressively degrades, the simulations predicting more peaked profiles than those observed (Fig. 7 bottom panel). This behaviour was also noticed by Nagai et al. (2007) when comparing their simulations with *Chandra* relaxed clusters, and it is also observed for the temperature profiles (see Pratt et al. 2007). As mentioned above, the core properties are most sensitive to non-gravitational processes and these discrepancies are again likely to reflect the fact that modelling of the processes is still inadequate.

The average simulation profile derived using the true mass for each simulated cluster is also shown in the figure (dotted lines). As compared to the scaling based on R_{500}^{HSE} and M_{500}^{HSE} ,

the scaled profile of each cluster is translated to the left and to the bottom in the log – log plane. The average profile lies below the profile based on the hydrostatic values, as expected from the mean bias between M_{500}^{Sim} and M_{500}^{HSE} . The offset with the observed profile in the outer region becomes more significant, about 30%.

In conclusion, there is an excellent agreement in shape between the simulated and observed profiles for the cluster outer regions, which is the most relevant aspect for the Y_{SZ} estimate. The better agreement in normalisation with the simulations when using the hydrostatic mass suggests that the hydrostatic X-ray masses used to scale the observed profiles are indeed underestimated.

5. The universal pressure profile

As pointed out by Nagai et al. (2007), an analytic cluster pressure profile model is useful both for analysis of SZ observations and for theoretical studies. Of prime interest is a model for the average scaled profile of the entire cluster population. For nearby clusters it can be derived from the present data, the REXCESS sample being a representative sample.

We considered the generalized NFW (GNFW) model proposed by Nagai et al. (2007):

$$\mathbb{P}(x) = \frac{P_0}{(c_{500}x)^\gamma [1 + (c_{500}x)^\alpha]^{(\beta-\gamma)/\alpha}}. \quad (11)$$

The parameters (γ, α, β) are respectively the central slope ($r \ll r_s$), intermediate slope ($r \sim r_s$) and outer slope ($r \gg r_s$), where $r_s = R_{500}/c_{500}$, and they are highly correlated with r_s . In order to constrain the parameters, it is essential to consider a wide radial range, including both the core ($r < 0.1 R_{500}$) and the cluster periphery ($r > R_{500}$). In particular, β remains essentially unconstrained when considering only data within $r < R_{500}$, resulting in large uncertainties in the profile model beyond R_{500} and thus on the corresponding integrated SZ signal.

Taking advantage of the good agreement between observations and simulations in the outer cluster regions, we thus defined an hybrid average profile, combining the profiles from observations and simulations. It is defined by the observed average scaled profile in the radial range $[0.03-1] R_{500}$ derived in Sect. 3.2 and the average simulation profile in the $[1-4] R_{500}$ region. For the simulations, we used the profile based on the hydrostatic quantities and renormalised it by +10% to correct for the observed offset with the observations at $r > 0.4 R_{500}$. We fitted this hybrid profile with the GNFW model in the log – log plane, weighting the “data” points according to the dispersion. The best fitting model is plotted in Fig. 8, with parameters:

$$[P_0, c_{500}, \gamma, \alpha, \beta] = [8.403 h_{70}^{-3/2}, 1.177, 0.3081, 1.0510, 5.4905]. \quad (12)$$

Using the dimensionless “universal” profile, $\mathbb{P}(x)$ (Eqs. (11) and (12)), and taking into account the mass dependence established in Sect. 3.4, we can describe the physical pressure profile of clusters as a function of mass and redshift (assuming standard evolution):

$$\begin{aligned} P(r) &= P_{500} \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_\odot} \right]^{\alpha_p + \alpha'_p(x)} \mathbb{P}(x) \\ &= 1.65 \times 10^{-3} h(z)^{8/3} \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_\odot} \right]^{2/3 + \alpha_p + \alpha'_p(x)} \\ &\quad \times \mathbb{P}(x) h_{70}^2 \text{ keV cm}^{-3} \end{aligned} \quad (13)$$

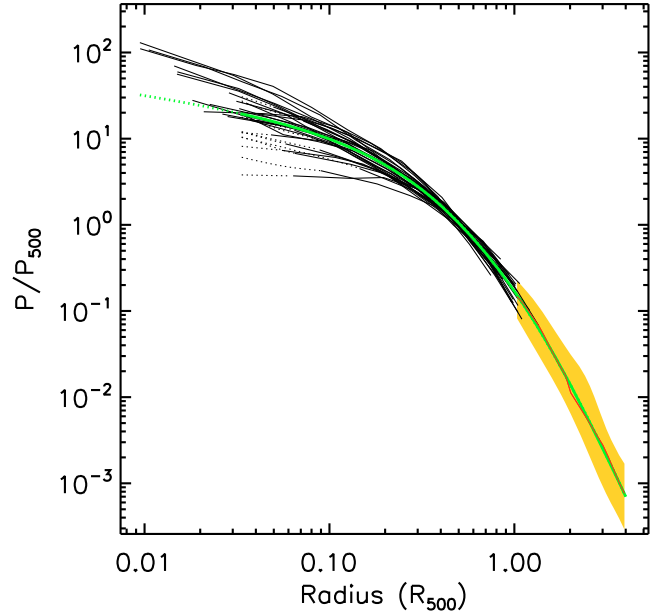


Fig. 8. GNFW model of the universal pressure profile (green line). It is derived by fitting the observed average scaled profile in the radial range $[0.03-1] R_{500}$, combined with the average simulation profile beyond R_{500} (red line). Black lines: REXCESS profiles. Orange area: dispersion around the average simulation profile.

with $x = r/R_{500}$, α_p and $\alpha'_p(x)$ from Eqs. (7) and (8), and $\mathbb{P}(x)$ from Eq. (11) with parameters from Eq. (12). The second term in the mass exponent, α_p , corresponds to a modification of the standard self-similarity (i.e., the steeper mass dependence of the profile), while the third term, $\alpha'_p(x)$ (Eq. (8)), introduces a break in self-similarity (i.e., a mass dependence of the shape). The latter is a second order effect, which can be neglected in first approximation.

We also fitted each individual observed cluster profile with the GNFW model, fixing the β value to that derived above (Eq. (12)), as well as the average scaled profile of the cool core and morphologically disturbed clusters. The best fitting parameters are listed in Appendix C, where we also provide plots of each individual cluster profile with its best fitting model.

6. Integrated Compton parameter scaling relations

6.1. Definitions and method

In this section we discuss scaling relations directly relevant for SZE studies. We will consider the volume integrated Compton parameter Y , for both cylindrical and spherical volumes of integration. The spherically integrated quantity, $Y_{\text{sph}}(R)$, proportional to the gas thermal energy, is defined as:

$$Y_{\text{sph}}(R) = \frac{\sigma_T}{m_e c^2} \int_0^R 4\pi P(r) r^2 dr \quad (14)$$

and the cylindrically integrated quantity, $Y_{\text{cyl}}(R) = Y_{\text{SZ}} D_A^2$, directly related to the SZ signal within an aperture $\theta = R/D_A$, is:

$$\begin{aligned} Y_{\text{cyl}}(R) &= \frac{\sigma_T}{m_e c^2} \int_0^R 2\pi r dr \int_r^{R_b} \frac{2 P(r') r' dr'}{\sqrt{r'^2 - r^2}} \\ &= Y_{\text{sph}}(R_b) - \frac{\sigma_T}{m_e c^2} \int_R^{R_b} 4\pi P(r) \sqrt{r^2 - R^2} r dr \end{aligned} \quad (15)$$

where R_b is the cluster radial extent. In the following, we adopt $R_b = 5 R_{500}$, as suggested by numerical simulations (Sect. 4.2). Note that the total SZ signal is then equivalently $Y_{\text{sph}}(5 R_{500})$ or $Y_{\text{cyl}}(5 R_{500})$.

For each cluster, the spherically integrated Compton parameter can be readily computed from the observed pressure profile. The Y_{sph} scaling relations can then be directly derived from the data for integration radii up to R_{500} , the observed radial range. They are presented below in Sect. 6.2. Such a derivation is not possible for Y_{cyl} (or the total Y_{SZ} signal): it involves integration along the line of sight up to $R_b = 5 R_{500}$, i.e., beyond the observed radial range. However, using the universal pressure profile, we can compute the volume integrated Compton parameter, Y , for any region of interest, and derive the corresponding scaling relations (presented below in Sect. 6.3). The two approaches give fully consistent results, as shown below.

Finally, for convenience, we also define a characteristic Compton parameter, Y_{500} , corresponding to the characteristic pressure P_{500} (see Appendix A):

$$Y_{500} = \frac{\sigma_{\text{T}}}{m_e c^2} \frac{4\pi}{3} R_{500}^3 P_{500} \\ = 2.925 \times 10^{-5} h(z)^{2/3} \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{5/3} h_{70}^{-1} \text{Mpc}^2. \quad (16)$$

6.2. Observed $Y_{\text{sph}} - Y_{\text{X}}$ and $Y_{\text{sph}} - M_{500}$ relations

The values for $Y_{\text{sph}}(R_{2500})$ and $Y_{\text{sph}}(R_{500})$, derived from the observed pressure profiles, are given in Table C.1. R_{2500} is defined as $R_{2500} = 0.44 R_{500}$ from the scaling relations presented in Arnaud et al. (2005). The integration was performed using the MC deconvolved density and model temperature profiles, allowing us to propagate the statistical errors, including that on R_{500} . We checked that using instead the best fitting GFW model for each profile gives consistent results within the statistical errors. Note that the errors on M_{500} take into account the statistical errors on the relevant X-ray data, but not the uncertainties on the $M_{500} - Y_{\text{X}}$ relation itself. The latter are therefore not included in the statistical errors on the slope and normalisation of the relations.

Figure 9 shows the $Y_{\text{sph}} - Y_{\text{X}}$ relations with $Y_{\text{X}} = M_{\text{g},500} T_{\text{X}}$, together with the best fitting power law. We normalised Y_{X} by:

$$C_{\text{XSZ}} = \frac{\sigma_{\text{T}}}{m_e c^2} \frac{1}{\mu_e m_p} = 1.416 \times 10^{-19} \frac{\text{Mpc}^2}{M_{\odot} \text{keV}} \quad (17)$$

for $\mu_e = 1.148$, the mean molecular weight of electrons for a 0.3 solar abundance plasma. Note that the $Y_{\text{sph}} - Y_{\text{X}}$ relation depends only weakly on the assumed $M_{500} - Y_{\text{X}}$ relation, via the estimate of R_{500} only. For some clusters, the computation of $Y_{\text{sph}}(R_{500})$ requires extrapolation: by more than 20% for 8 clusters and, in the worst case, RXC J2157.4-0747, the profile of which is measured only up to $R_{\text{det}} \sim 0.6 R_{500}$, $Y_{\text{sph}}(R_{500})$ is larger by a factor 1.8 than the value within R_{det} . However, the best fitting $Y_{\text{sph}}(R_{500}) - Y_{\text{X}}$ relation is stable to the inclusion or exclusion of clusters requiring extrapolation, the best fitting parameters being consistent within the errors.

As mentioned in the introduction, the $Y_{\text{sph}} - Y_{\text{X}}$ relation depends on the internal cluster structure (Eq. (1)). For $Y_{\text{sph}}(R_{2500})$, we obtained:

$$Y_{\text{sph}}(R_{2500}) = 10^{-0.272 \pm 0.097} \left[\frac{C_{\text{XSZ}} Y_{\text{X}}}{h_{70}^{-5/2} \text{Mpc}^2} \right]^{1.036 \pm 0.020} h_{70}^{-5/2} \text{Mpc}^2. \quad (18)$$

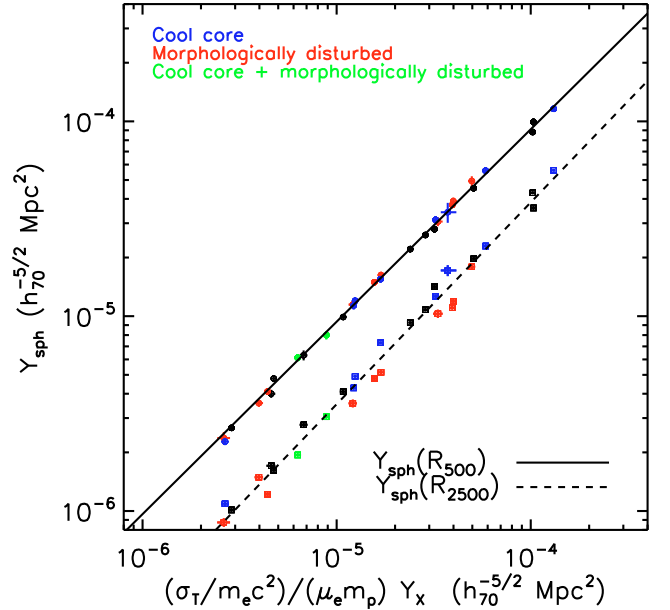


Fig. 9. The $Y_{\text{sph}} - Y_{\text{X}}$ relations from REXCESS data. $Y_{\text{sph}}(R)$ is the spherically integrated Compton parameter, within R_{2500} (squares) and R_{500} (circles). $Y_{\text{X}} = M_{\text{g},500} T_{\text{X}}$ is the product of the gas mass within R_{500} and the spectroscopic temperature T_{X} . Data points are colour-coded according to cluster dynamical state. Lines: best fitting power law.

The best fitting slope is slightly greater than one (a 2σ effect), reflecting the stronger mass dependence of the pressure profile in the center ($r < R_{2500}$) as compared the expectation from the $M_{500} - Y_{\text{X}}$ relation (Fig. 4 and Sect. 3.4). The intrinsic dispersion is $\sigma_{\log_{10}, Y} = 0.054 \pm 0.006$, with the morphologically disturbed clusters lying below the mean relation and the relaxed clusters lying above it, a consequence of the shallower profile of the former as compared to the latter (Fig. 2). When we move to $Y_{\text{sph}}(R_{500})$, the best fitting slope (1.003 ± 0.008) becomes consistent with unity, i.e. the shape variation with mass, when averaged within R_{500} , has essentially no effect (see also below). The intrinsic dispersion is no longer measurable, the dispersion is consistent with that expected from the statistical errors. This is a direct consequence of the high similarity of the pressure profiles beyond the core ($r \gtrsim 0.2 R_{500}$), while the core typically contributes by less than 10% to $Y_{\text{sph}}(500)$ (see below and Fig. 11). Fixing the slope to one, the best fitting normalisation gives:

$$\frac{Y_{\text{sph}}(R_{500})}{C_{\text{XSZ}} Y_{\text{X}}} = 0.924 \pm 0.004. \quad (19)$$

Note that this ratio is nothing more than the ratio, $T_{\text{mg}}/T_{\text{X}}$, of the gas mass weighted temperature to T_{X} . It is less than unity, as found in other studies (Vikhlinin et al. 2006), and as expected for decreasing temperature profiles.

Figure 10 shows the $Y_{\text{sph}}(R_{500}) - M_{500}$ data together with the best fitting relation:

$$h(z)^{-2/3} Y_{\text{sph}}(R_{500}) = 10^{-4.739 \pm 0.003} \\ \times \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{1.790 \pm 0.015} h_{70}^{-5/2} \text{Mpc}^2. \quad (20)$$

Since M_{500} is derived from the $M_{500} - Y_{\text{X}}$ relation, this expression does not contain more information than the $Y_{\text{sph}}(R_{500}) - Y_{\text{X}}$ relation, combined with the calibration of the $M_{500} - Y_{\text{X}}$ relation. As expected, the normalisation and slope are consistent with that

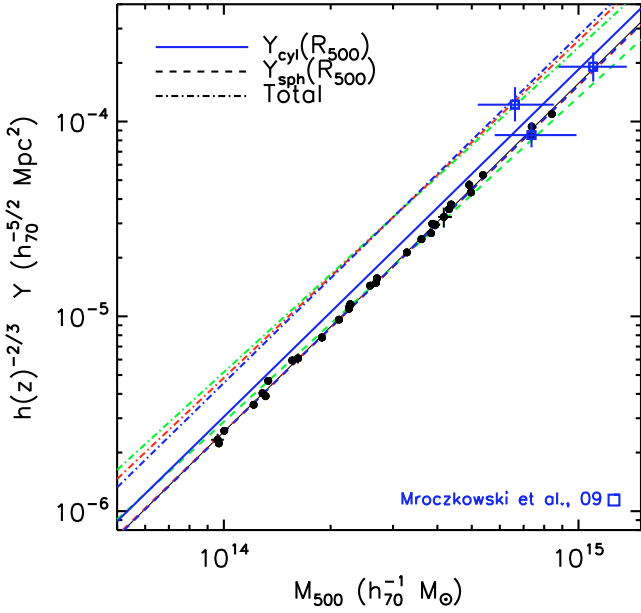


Fig. 10. The $Y_{SZ} - M_{500}$ relations. Thin black line: power law $Y_{sph}(R_{500}) - M_{500}$ relation best fitting the REXCESS data (black points). Lines: $Y_{sph}(R_{500}) - M_{500}$ relation (dashed line), $Y_{cyl}(R_{500}) - M_{500}$ relation (relation between $Y_{SZ} D_A^2$ within a R_{500} aperture and M_{500} ; full line) and $Y_{cyl}(5R_{500}) - M_{500}$ relation (relation between the total $Y_{SZ} D_A^2$ signal and M_{500} ; dash-dotted line) derived from the universal GFW scaled pressure profile and for different M_{500} scaling: standard self-similar scaling ($\alpha = 5/3$; green), modified scaling taking into account the non-standard slope of the $M_{500} - Y_X$ relation ($\alpha = 1.78$; blue), and further taking into account the break of self-similarity of the pressure profile shape (see text, red). Blue squares: $(Y_{sph}(R_{500}), M_{500})$ measurements for 3 clusters: from top to bottom, A1835 ($z = 0.25$), A1914 ($z = 0.17$), and CL J1226.9+3332 ($z = 0.89$). They were derived by Mroczkowski et al. (2009) from a joint analysis of SZA and X-ray observations using a GFW model.

obtained by combining Eqs. (2) and (19), and, similar to the $Y_{sph}(R_{500}) - Y_X$ relation, the scatter is consistent with the statistical scatter.

6.3. Scaling relations from the universal pressure profile

6.3.1. $Y_{sph} - M_{500}$ and $Y_{SZ} D_A^2 - M_{500}$ relations

Let us first consider Y_{sph} derived from the universal pressure profile. Combining Eqs. (14), (13) and (16):

$$Y_{sph}(R) = Y_{500} \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{\alpha_p} \times \int_0^x 3f(u, M_{500}) \mathbb{P}(u) u^2 du \quad (21)$$

with $f(u, M_{500}) = (M_{500}/3 \times 10^{14} h_{70}^{-1} M_{\odot})^{\alpha'_p(u)}$. This term in the integral reflects the break of self-similarity in the pressure profile (Sect. 5). Neglecting this effect, the corresponding $Y_{sph} - M_{500}$ relation, for any integration radius, is a power law of slope $\alpha = 5/3 + \alpha_p = 1/\alpha_{MY_X}$ (Eqs. (7) and (16)). Taking into account this effect, the relation is no longer a simple power law. Following the behavior of the pressure profiles - $\alpha'_p(u)$ decreases with radius or equivalently the departure from standard mass scaling becomes less and less pronounced as we move towards the cluster outskirts - the relation is expected to become shallower with increasing integration radius, closer to the standard self-similar

relation ($\alpha = 5/3$). The relations for various mass scalings can be compared in Fig. 10, for $Y_{sph}(R_{500})$ and $Y_{sph}(5 R_{500})$ (i.e. the total Y_{SZ} signal). The effect of the self-similarity break is small. In the mass range $M_{500} = [10^{14} M_{\odot}, 10^{15} M_{\odot}]$, Y_{sph} varies, as compared to the value computed neglecting this effect, by $[-7\%, +8\%]$, $[-1\%, +0.5\%]$ and $[+6\%, -6\%]$, for an integration radius of R_{2500} , R_{500} and $5 R_{500}$, respectively. When taking into account the self-similarity break, the corresponding effective slopes of the $Y_{sph} - M_{500}$ for that mass range are 1.84, 1.78 and 1.73, as compared to $1/\alpha_{MY_X} = 1.78$ ignoring the effect. The effect is fully negligible for the $Y_{sph}(R_{500}) - M_{500}$ relation, as found above directly from the data; it is at most equal to the statistical uncertainty on $1/\alpha_{MY_X} = 1.78 \pm 0.06$ (Eq. (2)) and we will neglect it in the following.

In that case, and combining Eqs. (21), (7) and (16), the $Y_{sph} - M_{500}$ relation for an integration radius of $x R_{500}$ can be written as:

$$h(z)^{-2/3} Y_{sph}(x R_{500}) = A_x \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{\alpha} \quad (22)$$

where

$$\alpha = 1.78; \quad A_x = 2.925 \times 10^{-5} I(x) h_{70}^{-1} \text{Mpc}^2 \quad (23)$$

$$I(x) = \int_0^x 3\mathbb{P}(u) u^2 du \quad (24)$$

with $\mathbb{P}(u)$ from Eqs. (11) and (12). Numerical values for $I(x)$ of particular interest are $I(1) = 0.6145$ and $I(5) = 1.1037$. The former gives the normalisation of the $Y_{sph}(R_{500}) - M_{500}$ relation, $\log(A_x) = -4.745$. It is in excellent agreement (1% difference) with the normalisation derived from a direct fit to the data (Eq. (20)). The latter gives the normalisation, $\log(A_x) = -4.491$, of the relation for the total $Y_{SZ} D_A^2$ signal, assuming a cluster radial extent of $5 R_{500}$.

Similarly, the relation for the SZ signal within an aperture of $x R_{500}$ is obtained from Eqs. (15), (13) and (16):

$$h(z)^{-2/3} Y_{SZ}(x R_{500}) D_A^2 = B_x \left[\frac{M_{500}}{3 \times 10^{14} h_{70}^{-1} M_{\odot}} \right]^{\alpha} \quad (25)$$

with

$$\alpha = 1.78; \quad B_x = 2.925 \times 10^{-5} J(x) h_{70}^{-1} \text{Mpc}^2 \quad (26)$$

$$J(x) = I(5) - \int_x^5 3\mathbb{P}(u) \sqrt{u^2 - x^2} u du \quad (27)$$

for a cluster extent of $5 R_{500}$. For an aperture of R_{500} , $J(1) = 0.7398$ or $\log(B_x) = -4.665$. The corresponding $Y_{SZ} D_A^2 - M_{500}$ relation is plotted in Fig. 10. We also show measurements for A1835 ($z = 0.25$), A1914 ($z = 0.17$), and CL J1226.9+3332 ($z = 0.89$), derived by Mroczkowski et al. (2009) from a joint analysis of SZA and X-ray observations using a GFW pressure profile model. Although the measurement errors are still large, the consistency with the present scaling relation is an encouraging sign of the validity of our determination of the scaling relations. Since the clusters cover a wide redshift range, it further suggests a standard self-similar evolution, as assumed in Eq. (25).

Uncertainties on the above relations, that are established combining observational and theoretical data, cannot be assessed rigorously. Rough estimates of the statistical errors can be derived by combining the errors on the $Y_{sph}(R_{500}) - Y_X$ and $M_{500} - Y_X$ relations, with the latter largely dominant. This gives $\alpha = 1.78 \pm 0.06$ or $\alpha = 1.78 \pm 0.08$, further adding quadratically the systematic effect of the pressure self-similarity break discussed above. The logarithmic error on the normalisation at the pivot is ± 0.024 ($\pm 6\%$).

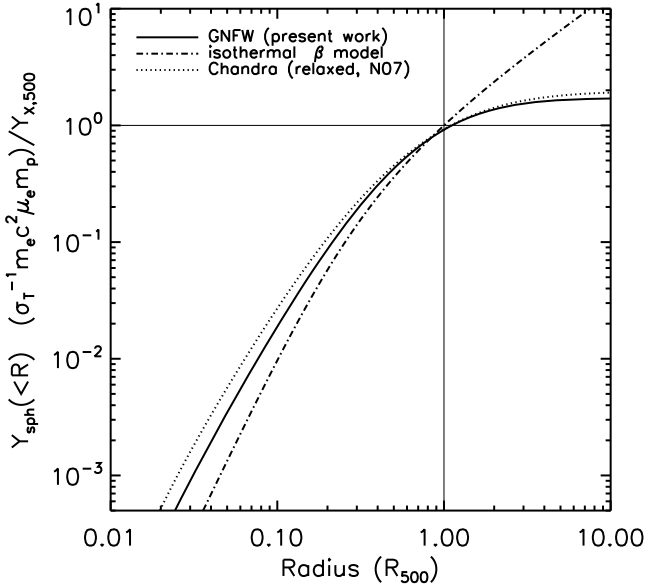


Fig. 11. Spherically integrated Compton parameter, $Y_{\text{sph}}(R)$, as a function of scaled integration radius, normalised to $C_{\text{XSZ}}Y_{\text{X}}$. Full line: universal GNFW scaled pressure profile (Eqs. (11) and (12)). Dotted line: GNFW model obtained by Nagai et al. (2007) for *Chandra* relaxed ($T > 5$ keV) clusters. Dashed-dotted line: isothermal β -model with $\beta = 2/3$ and a core radius of $0.2 R_{500}$.

6.3.2. Behavior of $Y_{\text{sph}}(R)$ and comparison with the isothermal β -model

It is instructive to study in more detail the radial dependence of Y_{sph} . $Y_{\text{sph}}(R)$ varies with radius as $I(x)$ (Eq. (24) with $\mathbb{p}(u)$ from Eqs. (11) and (12)). By construction its normalisation scales with mass as Y_{X} . Figure 11 shows the variation of $Y_{\text{sph}}(R)$ with scaled integration radius, normalised to $C_{\text{XSZ}}Y_{\text{X}}$, so that we are effectively probing $Y_{\text{sph}}(R)$ at fixed mass.

At large radii, the integrand in $I(x)$ varies as $\mathbb{p}(u)u^2 du \propto u^{-2}$ for an outer pressure profile slope of $\beta \sim 5$. As a result, $Y_{\text{sph}}(R)$ converges rapidly beyond R_{500} and the total SZ signal is not very sensitive to the assumption on cluster extent. Assuming a cluster extent of $4 R_{500}$, $6 R_{500}$ or even $100 R_{500}$, rather than $5 R_{500}$, changes the total signal by only -2% , $+1.3\%$ and $+4\%$, respectively. On the other hand, the figure shows the dominant contribution of the external regions to Y_{sph} : 50% of the contribution to $Y_{\text{sph}}(R_{500})$ comes from $R \geq 0.53 R_{500}$ while the region within $0.1 R_{500}$ and $0.2 R_{500}$ contributes by only 2% and 9.5% respectively. This will be even more pronounced for the Y_{SZ} signal (integration within a cylindrical volume).

We also plot $Y_{\text{sph}}(R)$ for the GNFW model obtained by Nagai et al. (2007) from *Chandra* data (for the corrected parameters, [12.2, 1.3, 0.4, 0, 9, 5.0], published by Mroczkowski et al. 2009). It is slightly larger in the center, as expected from the more peaked nature of the scaled *Chandra* profiles (Sect. 3.5). The agreement⁴ is very good in the outskirts, as it is for the profiles (Fig. 5), with a slightly higher asymptotic value due the slightly smaller value of β .

We also compare with the result obtained with an isothermal β -model, with $\beta = 2/3$ and a core radius of $r_c = 0.2 R_{500}$ (Arnaud et al. 2002). The difference is only 10% at R_{500} but the model diverges at high radii. This clearly shows that the total

⁴ Note, however, that (Nagai et al. 2007) assumed a standard self-similar mass scaling of the pressure profile. The $Y - M_{500}$ relations derived from their profiles would differ from ours in terms of slope.

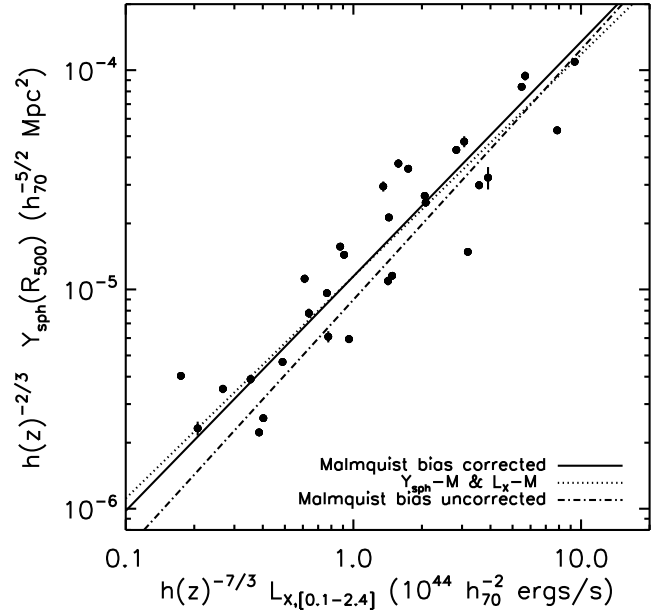


Fig. 12. The $Y_{\text{sph}}(R_{500}) - L_{\text{X}}$ relations. L_{X} is the [0.1–2.4] keV luminosity within R_{500} . Full line: power law relation best fitting the REXCESS data for L_{X} corrected for Malmquist bias (black points). Dotted line: relation computed by combining the $Y_{\text{sph}}(R_{500}) - M_{500}$ relation derived from the universal pressure profile (Eq. (22)) and the $L_{\text{X}} - M_{500}$ relation. Dash-dotted Line: best fitting $Y_{\text{sph}}(R_{500}) - L_{\text{X}}$ for uncorrected L_{X} .

Y_{SZ} signal derived assuming an isothermal β -model is very sensitive to the assumed extent of the cluster. It will also be always overestimated by such a model, as emphasized by Hallman et al. (2007). As an illustration, assuming a cluster extent of $2.03 R_{500}$, the top-hat virial radius often used in the literature, the β -model gives a total Y_{SZ} signal 1.7 higher than the universal pressure profile. This over-estimate depends on the choice of the β -model shape parameters. It decreases with decreasing core radius and increasing β value. It is still a factor of 1.4 for $r_c = 0.1 R_{500}$ and $\beta = 0.75$ and reaches a factor of 2.1 for $r_c = 0.3 R_{500}$ and $\beta = 0.6$.

6.4. The $Y - L_{\text{X}}$ relations

The scaling between the SZ signal and the X-ray luminosity, L_{X} is an important relation for comparing X-ray surveys such as the *ROSAT* All Sky Survey and future or on going SZE surveys, such as the *Planck* survey. The luminosity within R_{500} and in the soft-band [0.1–2.4] keV, most relevant for X-ray Surveys, has been estimated for REXCESS clusters by Pratt et al. (2009); here we used the values both corrected and uncorrected for Malmquist bias. Figure 12 shows the corresponding $Y_{\text{sph}}(R_{500}) - L_{\text{X}}$ relations.

We fitted the REXCESS data with a power law:

$$h(z)^{-2/3} Y_{\text{sph}}(R_{500}) = C \left[\frac{h(z)^{-7/3} L_{\text{X}}}{10^{44} h_70^{-2} \text{ erg s}^{-1}} \right]^{\alpha} h_70^{-5/2} \text{ Mpc}^2. \quad (28)$$

The best fitting parameters are given in Table 1. The intrinsic scatter around the relation is important, more than 50%, reflecting the important scatter, at given Y_{X} , of the soft band luminosity computed without excising the core (see Pratt et al. 2009). The best fitting relation is consistent with the relation expected from combining the $Y_{\text{sph}}(R_{500}) - M_{500}$ relation derived from the universal pressure profile (Eq. (22)) and the $L_{\text{X}} - M_{500}$ relation. For

Table 1. $h(z)^{-2/3} Y_{\text{sph}}(R_{500}) - h(z)^{-7/3} L_X$ and updated $h(z)^{-7/3} L_X - M_{500}$ relations (see text).

Relation	$\log_{10} C$	α	$\sigma_{\log,i}$
$Y_{\text{sph}}(R_{500}) - L_X - \text{MB}$	-4.940 ± 0.036	1.07 ± 0.08	0.190 ± 0.025
$L_X - M_{500} - \text{MB}$	0.193 ± 0.034	1.76 ± 0.13	0.199 ± 0.035
$Y_{\text{sph}}(R_{500}) - L_X$	-5.047 ± 0.037	1.14 ± 0.08	0.184 ± 0.024
$L_X - M_{500}$	0.274 ± 0.032	1.64 ± 0.12	0.183 ± 0.032

Notes. L_X is the [0.1–2.4] keV luminosity within R_{500} . MB: relations corrected for Malmquist bias. For each observable set, (B, A) , we fitted a power law relation of the form $B = C(A/A_0)^\alpha$, with $A_0 = 10^{44} h_{70}^{-2}$ ergs/s and $3 \times 10^{14} h_{70}^{-1} M_\odot$ for L_X and M_{500} , respectively. $\sigma_{\log,i}$: intrinsic scatter about the best fitting relation in the log – log plane.

consistency, the latter was updated (parameters given in Table 1), using present M_{500} values derived from the updated $M_{500} - Y_X$ relation (Eq. (2)). The slope and normalisation (taking into account the different pivot used) are consistent with those published in Pratt et al. (2009).

For practical purposes, the scaling of $Y_{\text{SZ}}(x R_{500}) D_A^2$ or that of the total SZ signal with L_X is of more direct interest than the $Y_{\text{sph}}(R_{500}) - L_X$ relation. In view of the good agreement of the latter with the universal profile model, the $Y - L_X$ relation, for any integration region of interest, can be safely derived by correcting the normalisation in Eq. (28) by the model ratio of Y to $Y_{\text{sph}}(R_{500})$. This ratio is simply $I(x)/I(1)$ for the spherically integrated Compton parameter, e.g., $I(5)/I(1) = 1.796$ for the total SZ signal, and $J(x)/I(1)$ for the $Y_{\text{SZ}}(x R_{500}) D_A^2$ signal.

6.5. Comparison with standard self-similar relations

The $Y - M_{500}$ relations derived above do not seem to deviate much from standard self-similarity (Fig. 10). A fully consistent standard (ST) model, with standard slope $Y - M_{500}$ relations, is obtained when using the standard slope $M_{500} - Y_X$ relation (Eq. (3)), as shown in Appendix B. The universal profile and scaling relations obtained in that case are given in the Appendix, together with a detailed comparison of the presently derived scaling relations with the ST relations. In summary, the difference for the $Y - M_{500}$ relations mirrors that for the $M_{500} - Y_X$ relation. As compared to values derived from the ST relation, Y is lower at low mass and higher at high mass. Typically, the difference for the total Y_{SZ} signal ranges from -19% to $+6\%$ in the $[10^{14} - 10^{15}] M_\odot$ mass range. On the other hand, the $Y - L_X$ relations, which only depend on cluster internal structure, are essentially the same in the two models: the difference is less than 5% in the $[0.1 - 10] 10^{44}$ ergs/s luminosity range.

7. Discussion

7.1. Departure from standard self-similarity

The present work is based on a representative sample of nearby clusters. The sample, REXCESS, was chosen by X-ray luminosity alone, without regard to morphology or dynamical state. As for the entropy (Pratt et al. 2010), the depth of the observations allowed us to probe the scaling behavior of the pressure profiles out to R_{500} . Both points are essential for a complete picture of the modification of the standard self-similarity due to non-gravitational processes, including its radial behavior.

The behaviour of the pressure profiles, with respect to standard self-similarity with zero dispersion, resembles that

generally found for other quantities such as the entropy or density: 1) regularity in shape outside the core 2) increased dispersion inside the core linked to cooling effects and dynamical state and 3) departure from standard mass scaling that becomes less pronounced towards the cluster outskirts. However, the latter two deviations are less pronounced than for the entropy and/or density, showing that the pressure is the quantity least affected by dynamical history and non-gravitational physics. This further supports the view that Y_{SZ} is indeed a good mass-proxy.

7.2. Robustness of the universal profile

We combined observational and simulation data to derive the universal pressure profile. The convergence of various approaches to determine scaled cluster profiles supports the robustness of our determination of the universal profile, particularly of its shape. This includes the agreement between independent simulations, between these simulations and the present observed data based on a representative cluster sample, and also the agreement between the present *XMM-Newton* data and published *Chandra* data for clusters of similar thermo-dynamical state. As a result, we believe that quantities which purely depend on the universal profile shape are particularly robust and well converged. This includes the typical SZ decrement profile or relations between the Compton parameter estimated in various apertures.

However, the universal profile beyond R_{500} is purely based on simulations and thus less secure than within R_{500} . While the standard non-gravitational processes currently implemented in numerical simulations are known to introduce small modifications of the profiles at large radii with respect to the adiabatic case (Nagai et al. 2007), other less explored processes may affect the profiles. In particular, the electron-proton equilibration time is larger than the Hubble time in the outskirts (Fox & Loeb 1997) and if the electron temperature is indeed smaller than the ion temperature, this will affect the pressure profile and lead to a decrease in the total Y_{SZ} signal (Rudd & Nagai 2009). The pressure profile interior to R_{500} is directly based on observations but derived from temperatures estimated using azimuthally averaged spectra. These have been corrected for the spectroscopic bias due to projection but not for azimuthal variations. High resolution SZ data with improved sensitivity are needed to probe any remaining systematic effects due to the spectroscopic bias, and to directly observe the shape of the pressure profile beyond R_{500} , which is out of reach of current X-ray observatories.

7.3. $Y - M_{500}$ relations

The cluster masses have been estimated using the mass-proxy Y_X . The absolute normalisation and slope of the $Y - M_{500}$ relations, derived using the universal profile, thus rely on the underlying observationally defined $M_{500} - Y_X$ relation. Initial comparison with $Y_{\text{SZ}}(R_{500})$ data for 3 high mass systems, measured with SZA by Mroczkowski et al. (2009) and analysed with a realistic analytic pressure profile, indicates good agreement. A key point is to extend this type of analysis to larger samples and to include lower mass systems.

We further emphasize that the $M_{500} - Y_X$ relation was calibrated from hydrostatic mass estimates using relaxed objects. The $Y - M_{500}$ relation we derive is technically a Y -X-ray mass relation and is expected to differ from the “true” $Y - M_{500}$ relation by the offset between the “true” mass and the hydrostatic mass for *relaxed* objects.

With the present study based on a mass proxy, we cannot assess the intrinsic scatter of the “true” $Y_{\text{sph}}(R_{500}) - M_{500}$ relation. However, an upper limit is the quadratic sum of the scatter of the $Y_X - M_{500}$ relation and that of the $Y_{\text{sph}}(R_{500}) - Y_X$ relation. Note that the latter is purely due to variations in pressure profile shapes. Our measure of the $Y_{\text{sph}}(R_{500}) - Y_X$ relation, using $Y_{\text{sph}}(R_{500})$ directly derived from spherical integration of the pressure profiles, exhibits dispersion consistent with the $<5\%$ statistical scatter. Our study thus does show that variations in pressure profile shapes do not introduce an extra scatter into the $Y_{\text{sph}}(R_{500}) - M_{500}$ relation as compared to that of the $Y_X - M_{500}$ relation. Actually, the scatter of the $Y_{\text{sph}} - M_{500}$ relation might even be smaller than that of the $Y_X - M_{500}$ relation: the tightness of these relations seems to arise from the empirical evidence that density and temperature are anti-correlated and Y_{sph} depends on their local products as opposed to a global product for Y_X .

8. Conclusions

The present work is the first examination of the properties of the ICM pressure for a *representative* sample of nearby clusters covering the mass range $10^{14} < M_{500} < 10^{15} M_{\odot}$. Scaling the individual pressure profiles by mass and redshift according to the standard self-similar model, we derived an average scaled pressure profile for the cluster population and relate the deviations about the mean to both the mass and the thermo-dynamical state of the cluster:

- Cool core systems exhibit more peaked pressure profiles, while morphologically disturbed systems have shallower profiles.
- As a result, the dispersion is large in the core region, reaching approximately 80 per cent at $0.03 R_{500}$. However, as compared to the density, the pressure exhibits less scatter, a result of the anticorrelation of the density and temperature profiles interior to $0.2 R_{500}$. Outside the core regions, the dispersion about the average profile is remarkably low, at less than 30 per cent beyond $0.2 R_{500}$.
- We find a residual mass dependence of the scaled profiles, with a slope of ~ 0.12 , consistent with that expected from the empirical non-standard slope of the $M_{500} - Y_X$ relation. However, there is some evidence that the departure from standard scaling decreases with radius and is consistent with zero at R_{500} . We provide an analytical correction to the mean slope that accounts for this second order effect.

The observational data are compared to and combined with simulated data to derive the universal ICM pressure profile:

- Simulated scaled profiles from three independent sets of state of the art numerical simulations show excellent agreement, within 20%, between 0.1 and $3 R_{500}$, for pressures varying by 4 orders of magnitude in that radial range.
- Comparison with observed scaled data shows good agreement outside the core regions, which is the most relevant aspect for the Y_{SZ} estimate. The average simulation profile lies parallel to the observed data, with only a slight offset (~ 10 per cent) when the simulated profiles are scaled using the hydrostatic mass.
- This motivates us to combine the average observed scaled profile in the $[0.03-1] R_{500}$ radial range with the average simulated profile in the $[1-4] R_{500}$ range. This hybrid profile is fitted by a generalised NFW model, which allows us to define a dimensionless universal ICM pressure profile.

Combined with the empirical mass scaling of the profiles, this universal profile defines the physical pressure profile of clusters, up to the cluster boundary, as a function of mass and redshift, assuming self-similar evolution.

This universal profile is then used to predict the scaling relations involving the integrated Compton parameter Y :

- The expected $Y_{\text{sph}}(xR_{500}) - M_{500}$ or $Y_{\text{SZ}}(xR_{500}) - M_{500}$ relations are derived for any aperture. The slope is the inverse of the empirical slope of the $M_{500} - Y_X$ relation. The normalisation is given by the dimensionless integral of the universal profile within the region of interest expressed in scaled radius. The corresponding $Y_{\text{SZ}} - L_X$ relations can be derived by combining the relevant $Y - M_{500}$ relation with the empirical $L_X - M_{500}$ relation.
- The $Y_{\text{sph}}(R_{500}) - M_{500}$ and $Y_{\text{sph}}(R_{500}) - L_X$ relations derived directly from the individual profiles are in excellent agreement with those expected from the universal profile.
- We confirm that the isothermal β -model over-estimates the Y signal at given mass. This overestimate depends strongly on the assumption on cluster extent and reaches a factor of nearly two at $2 R_{500}$.

As a matter of practical application, the universal pressure profile is given in Eq. (11) with parameters in Eq. (12). For clusters of given mass M_{500} and z , the physical pressure profile can then be derived from Eq. (13) and the spherical $Y_{\text{sph}}(R)$ or cylindrical $Y_{\text{cyl}}(R)$ quantities can be estimated for any radius of interest using Eqs. (22)–(24) and Eqs. (25)–(27), respectively. These equations can be used as is when M_{500} is estimated for relaxed systems using the HSE equation, and for all clusters using M_{500} derived from mass-proxy relations. The preferred relations would be the $M_{500} - Y_X$ and the $M_{500} - L_X$, where L_X is the core-excised bolometric luminosity (Pratt et al. 2009), as both these relations display low scatter, compared to the relation between M_{500} and the full aperture soft band L_X . A typical application would be to predict the SZ signal of a known X-ray cluster with measured L_X or M_{500} , or to estimate the mass and thus X-ray properties of newly discovered SZ clusters. Other applications include the analysis of low S/N and/or poor resolution SZ observation of X-ray clusters, e.g., allowing to optimise the integration aperture and use a realistic decrement shape. On the other hand, care is needed when knowledge of the “true” mass is important, e.g., in predicting cluster number counts for future SZ surveys or in SZ selection function modelling. The above total $Y_{\text{SZ}} - M_{500}$ relation should be corrected by the bias between the true mass and the HSE mass at R_{500} , which is typically $\sim 13\%$ as determined from comparison with current numerical simulations.

A major open issue is the pressure evolution. With the present study based on a local cluster sample, we could only assume standard self-similar evolution. Because the SZ signal is not subject to redshift dimming, on going SZ surveys are expected to detect many new clusters at high z . Of particular interest is the *Planck* survey, which, thanks to its All-Sky coverage, will detect massive, thus rare, clusters, the best objects for precise cosmology with clusters. SZ follow-up, at the best possible resolution, and sensitive X-ray follow-up (particularly with *XMM-Newton*) will be crucial to assess possible evolution of pressure profile shape and measure the evolution of the $M_{500} - Y_X$ and $Y_{\text{SZ}} - M_{500}$ relations. Further progress, in particular on the mass bias and on the intrinsic scatter of the $Y - M$ relation, is expected from the wealth of high quality multi-wavelength data that will be available in the coming years.

Acknowledgements. We would like to thank Stefano Borgani, Daisuke Nagai, and Riccardo Valdarnini for providing us with the simulated data and for helpful discussions and useful comments on the manuscript. We thank J. B. Melin for enlightening discussions related to SZ observations. The present work is based on observations obtained with *XMM-Newton*, an ESA science mission with instruments and contributions directly funded by ESA Member States and the USA (NASA). E.P. acknowledges the support of grant ANR-06-JCJC-0141.

References

- Afshordi, N., Lin, Y.-T., Nagai, D., & Sanderson, A. J. 2007, *MNRAS*, 378, 293
Akritas, M. G., & Bershadsky, M. A. 1996, *ApJ*, 470, 706
Arnaud, M., Aghanim, N., & Neumann, D. 2002, *A&A*, 389, 1
Arnaud, M., Pointecouteau, E., & Pratt, G. W. 2005, *A&A*, 441, 893
Arnaud, M., Pointecouteau, E., & Pratt, G. W. 2007, *A&A*, 474, L37
Atrio-Barandela, F., Kashlinsky, A., Kocevski, D., & Ebeling, H. 2008, *ApJ*, 675, L57
Bonaldi, A., Tormen, G., Dolag, K., & Moscardini, L. 2007, *MNRAS*, 278, 1248
Bonamente, M., Joy, M., LaRoque, S., et al. 2008, *ApJ*, 675, 106
Böhringer, H., Schuecker, P., Guzzo, L., et al. 2004, *A&A*, 425, 367
Böhringer, H., Schuecker, P., Pratt, G. W., et al. 2007, *A&A*, 469, 363
Borgani, S., Murante, G., Springel, V., et al. 2004, *MNRAS*, 348, 1078
Croston, J. H., Arnaud, M., Pointecouteau, E., & Pratt, G. W. 2006, *A&A*, 459, 1007
Croston, J. H., Pratt, G. W., Böhringer, H., et al. 2008, *A&A*, 487, 431
da Silva, A. C., Kay, S. T., Liddle, A. R., & Thomas, P. 2004, *MNRAS*, 348, 1401
Diego, J. M., & Partridge, B. 2010, *MNRAS*, 402, 1179
Dunkley, J., Komatsu, E., Nolta, M. R., et al. 2009, *ApJS*, 180, 306
Finoguenov, A., Böhringer, H., & Zhang, Y.-Y. 2005, *A&A*, 442, 827
Finoguenov, A., Davis, D. S., Zimer, M., & Mulchaey, J. 2006, *ApJ*, 646, 143
Fox, D., & Loeb, A. 1997 *ApJ*, 491, 459
Hallman, E. J., Burns, J. O., Motl, P. M., & Norman, M. L. 2007, *ApJ*, 665, 911
Halverson, N., Lanting, T., Ade, P., et al. 2009, *ApJ*, 701, 42
Jia, S., Böhringer, H., Pointecouteau, E., Chen, Y., & Zhang, Y. 2008, *A&A*, 489, 1
Johnson, S., Ponman, T., & Finoguenov, A. 2009, *MNRAS*, 395, 1287
Jones, C., & Forman, W. 1984, *ApJ*, 276, 38
Kitayama, T., Komatsu, E., Ota, N., et al. 2004, *PASJ*, 56, 17
Komatsu, E., Kitayama, T., Suto, Y., et al. 1999, *ApJ*, 516, L1
Komatsu, E., Matsuo, H., Kitayama, T., et al. 2001, *PASJ*, 53, 57
Kravtsov, A. V., Vikhlinin, A., & Nagai, D. 2006, *ApJ*, 650, 128
Lewis, A. D., Buote, D. A., & Stocke, J. T. 2003, *ApJ*, 586, 135
Lieu, R., Mittaz, J., & Zhang, S.-N. 2006, *ApJ*, 648, 176
Mathiesen, B. F., & Evrard, A. E. 2001, *ApJ*, 546, 100
Maughan, B. J. 2007, *ApJ*, 668, 772
Marrone, D., Smith, G. P., Richard, J., et al. 2009, *ApJ*, 701, L114
Mazzotta, P., Rasia, E., Moscardini, L., & Rormen, G. 2004, *MNRAS*, 354, 10
McCarthy, I., Holder, G., Babul, A., & Balogh, M. 2003, *ApJ*, 591, 526
Mitchell, N. L., McCarthy, I. G., Bower, R. G., Theuns, T., & Crain, R. A. 2009, *MNRAS*, 395, 180
Motl, P. M., Hallman, E. J., Burns, J. O., & Norman, M. L. 2005, *ApJ*, 623, L63
Morandi, A., Ettori, S., & Moscardini, L. 2007, *MNRAS*, 379, 518
Mroczkowski, T., Bonamente, M., Carlstrom, J., et al. 2009, *ApJ*, 694, 1034
Nagai, D. 2006, *ApJ*, 650, 538
Nagai, D., Kravtsov, A. V., & Vikhlinin, A. 2007, *ApJ*, 668, 1
Nord, M., Basu, K., Pacaud, F., et al. 2009, *A&A*, 506, 623
Piffaretti, R., & Valdarnini, R. 2008, *A&A*, 491, 71
Pointecouteau, E., Giard, M., Benoit, A., et al. 1999, *ApJ*, 519, L115
Pointecouteau, E., Giard, M., Benoit, A., et al. 2001, *ApJ*, 552, 42
Pointecouteau, E., Arnaud, M., Kaastra, J., & de Plaa, J. 2004, *A&A*, 423, 33
Pointecouteau, E., Arnaud, M., & Pratt, G. W. 2005, *A&A*, 435, 1
Pratt, G. W., Böhringer, H., Croston, J. H., et al. 2007, *A&A*, 461, 71
Pratt, G. W., Croston, J. H., Arnaud, M., & Böhringer, H. 2009, *A&A*, 498, 361
Pratt, G. W., Arnaud, M., Piffaretti, R., et al. 2010, *A&A*, 511, A85
Rudd, D., & Nagai, D. 2009, *ApJ*, 701, L16
Reid, B., & Spergel, D. 2006, *ApJ*, 651, 643
Shaw, L., Holder, G., & Bode, P. 2008, *ApJ*, 686, 206
Staniszewski, Z., Ade, P. A., Aird, K. A., et al. 2009, *ApJ*, 701, 32
Vikhlinin, A. 2006, *ApJ*, 640, 710
Vikhlinin, A., Markevitch, M., Murray, S. S., et al. 2005, *ApJ*, 628, 655
Vikhlinin, A., Kravtsov, A., Forman, W., et al. 2006, *ApJ*, 640, 691
Vikhlinin, A., Burenin, R., Ebeling, H., et al. 2009, *ApJ*, 692, 1033
Voit, G. M. 2005, *Rev. Mod. Phys.*, 77, 207
Voit, M., Kay, S., & Bryan, G. 2005, *MNRAS*, 364, 909
Wik, D. R., Sarazin, C. L., Ricker, P. M., & Randall, S. W. 2008, *ApJ*, 680, 17

Appendix A: Characteristic self-similar quantities

Following Nagai et al. (2007) and Voit (2005) the characteristic quantities, P_{500} and Y_{500} , used in the present work, are defined from a simple self-similar model. The characteristic temperature is $kT_{500} = \mu m_p G M_{500}/2 R_{500}$, the temperature of a singular isothermal sphere with mass M_{500} . Here, μ is the mean molecular weight and m_p , the proton mass. We recall that M_{500} is defined as the mass within the radius R_{500} at which the mean mass density is 500 times the critical density, $\rho_c(z)$, of the universe at the cluster redshift: $M_{500} = (4\pi/3) R_{500}^3 500 \rho_c(z)$ with $\rho_c(z) = 3H(z)^2/(8\pi G)$. $H(z)$ is the Hubble constant, $H(z) = H(0) \sqrt{\Omega_M (1+z)^3 + \Omega_\Lambda}$ and G is the Newtonian constant of gravitation. The characteristic gas density is $\rho_{g,500} = 500 f_B \rho_c(z)$, i.e., the ratio of the gas density to the dark matter density is that of the Universe baryon fraction f_B . The electron density is $n_{e,500} = \rho_{g,500}/(\mu_e m_p)$ where μ_e is the mean molecular weight per free electron.

The characteristic pressure, P_{500} , is then defined as:

$$P_{500} = n_{e,500} kT_{500} \quad (\text{A.1})$$

$$= \frac{3}{8\pi} \left[\frac{500 G^{-1/4} H(z)^2}{2} \right]^{4/3} \frac{\mu}{\mu_e} f_B M_{500}^{2/3} \quad (\text{A.2})$$

and the corresponding characteristic integrated Compton parameter is:

$$Y_{500} = \frac{\sigma_T}{m_e c^2} \frac{4\pi}{3} R_{500}^3 P_{500} = \frac{\sigma_T}{m_e c^2} \frac{f_B M_{500} kT_{500}}{\mu_e m_p} \quad (\text{A.3})$$

$$= \frac{\sigma_T}{m_e c^2} \left[\frac{\sqrt{500} G H(z)}{4} \right]^{2/3} \frac{\mu}{\mu_e} f_B M_{500}^{5/3}. \quad (\text{A.4})$$

Numerical coefficients given in the corresponding Eqs. (5) and (16) are obtained for $f_B = 0.175$, $\mu = 0.59$ and $\mu_e = 1.14$, the values adopted by Nagai et al. (2007), allowing a direct comparison with their best fitting GNFW model. Note that the exact choice for these parameters does not matter, and does not need to reflect ‘‘true’’ values, as long as the same convention is used throughout the study (e.g., when comparing observed and theoretical scaled profiles or observed scaled profiles from different samples or instruments).

Appendix B: The standard self-similar case

In this Appendix, we summarise results (hereafter ST results) obtained when M_{500} is estimated for each REXCESS clusters using the $M_{500} - Y_X$ relation with a standard slope (Eq. (3)). The other physical parameters are consistently estimated, R_{500} , Y_X and T_X simultaneously in the iteration process used to derive M_{500} (Sect. 2), and $Y_{\text{sph}}(R_{500})$ from integration of the pressure profiles up to R_{500} . For practical purposes, the baseline parameters obtained using the best fitting empirical $M_{500} - Y_X$ relation (Eq. (2)) can be converted to the ST values using the power law relations given in Table B.1. The luminosity L_X is kept unchanged, the difference in R_{500} values (at most 4.5%), having a negligible impact due to the steep drop of emission with radius.

In the ST case, the scaled pressure profiles do not show any significant dependence on mass, as shown in Sect. 3.4. In other words, the pressure profiles follow a standard self-similar mass scaling:

$$P(r) = P_{500} \mathbb{D}(r/R_{500}) \quad (\text{B.1})$$

Table B.1. Power law relations to convert physical parameters of REXCESS clusters from those derived using the empirical $M_{500} - Y_X$ relation (Eq. (2)) to those derived using the standard slope relation (Eq. (3))*.

Relation	C	α
$M_{500}^{\text{ST}} - M_{500}$	0.968	1.089
$T_X^{\text{ST}} - T_X$	1.002	0.992
$Y_X^{\text{ST}} - Y_X$	0.995	1.017
$Y_{\text{sph}}^{\text{ST}}(R_{500}) - Y_{\text{sph}}(R_{500})$	0.991	1.031

Notes. (*) For each observable, Q , the conversion follows the form $(Q^{\text{ST}}/Q_0) = C(Q/Q_0)^\alpha$ where the pivot, Q_0 is $3 \times 10^{14} M_\odot$, 5 keV, $2 \times 10^{14} M_\odot$ keV and $2 \times 10^{-5} \text{ Mpc}^2$ for M_{500} , T_X , Y_X and $Y_{\text{sph}}(R_{500})$, respectively.

Table B.2. $h(z)^{-2/3} Y_{\text{sph}}(R_{500}) - h(z)^{-7/3} L_X$ and $h(z)^{-7/3} L_X - M_{500}$ relations for M_{500} estimated using the standard slope $M_{500} - Y_X$ relation (Eq. (3))*.

Relation	$\log_{10} C$	α	$\sigma_{\log,i}$
$Y_{\text{sph}}(R_{500}) - L_X - \text{MB}$	-4.947 ± 0.037	1.08 ± 0.08	0.192 ± 0.025
$L_X - M_{500} - \text{MB}$	0.215 ± 0.035	1.61 ± 0.12	0.199 ± 0.035
$Y_{\text{sph}}(R_{500}) - L_X$	-5.056 ± 0.038	1.16 ± 0.08	0.184 ± 0.024
$L_X - M_{500}$	0.295 ± 0.032	1.50 ± 0.11	0.183 ± 0.032

Notes. (*) Same notations as in Table 1.

with P_{500} defined by Eq. (5). The GFW parameters of the universal profile $\mathbb{p}(x)$, derived as described in Sect. 5, are:

$$[P_0, c_{500}, \gamma, \alpha, \beta] = [8.130 h_{70}^{-3/2}, 1.156, 0.3292, 1.0620, 5.4807]. \quad (\text{B.2})$$

As a result, the integrated Compton parameters also follow standard self-similarity, $Y \propto M_{500}^{5/3}$. The $Y - M_{500}$ relations derived from the universal pressure profile can be written as:

$$\begin{aligned} Y_{\text{sph}}(x R_{500}) &= Y_{500} I(x) \\ Y_{\text{SZ}}(x R_{500}) D_A^2 &= Y_{500} J(x) \end{aligned} \quad (\text{B.3})$$

with Y_{500} given by Eq. (16) and $I(x)$ or $J(x)$ defined by Eqs. (24) and (27), respectively. For the GFW parameters given by Eq. (B.2), the numerical values of $I(1)$, $I(5)$ and $J(1)$ are 0.6552, 1.1885 and 0.7913, respectively. The $Y_{\text{sph}}(x R_{500}) - M_{500}$ relation derived from a direct fit to the data has a slope of 1.663 ± 0.013 , fully consistent with $5/3$. Over the $[10^{14} - 10^{15}] M_\odot$ mass range, it differs by less than 0.8% from that derived from the universal profile (Eq. (B.3)).

We also derived the observed $Y_{\text{sph}}(R_{500}) - L_X$ relation, as well as the $L_X - M_{500}$ corresponding to the modified M_{500} values. The best fitting power law parameters are given in Table B.2. The former is consistent with the relation expected from combining the $L_X - M_{500}$ relation with the $Y_{\text{sph}}(R_{500}) - M_{500}$ relation derived from the universal pressure profile (Eq. (B.3)). The $Y - L_X$ relation, for any integration region of interest, can be derived by correcting the normalisation of the $Y_{\text{sph}}(x R_{500}) - L_X$ given in Table B.2 by the model ratio of Y to $Y_{\text{sph}}(R_{500})$, as described in Sect. 6.4.

Figure B.1 compares the scaling relations derived in the paper with the ST relations derived in this section. The empirical slope of the $M_{500} - Y_X$ relation being smaller than the standard value, M_{500} at a given Y_X is higher at low Y_X and smaller at high Y_X (top panel). Equivalently, Y_X at given mass is smaller at low mass, by $\sim -16\%$ at $M_{500} = 10^{14} M_\odot$, and higher at high

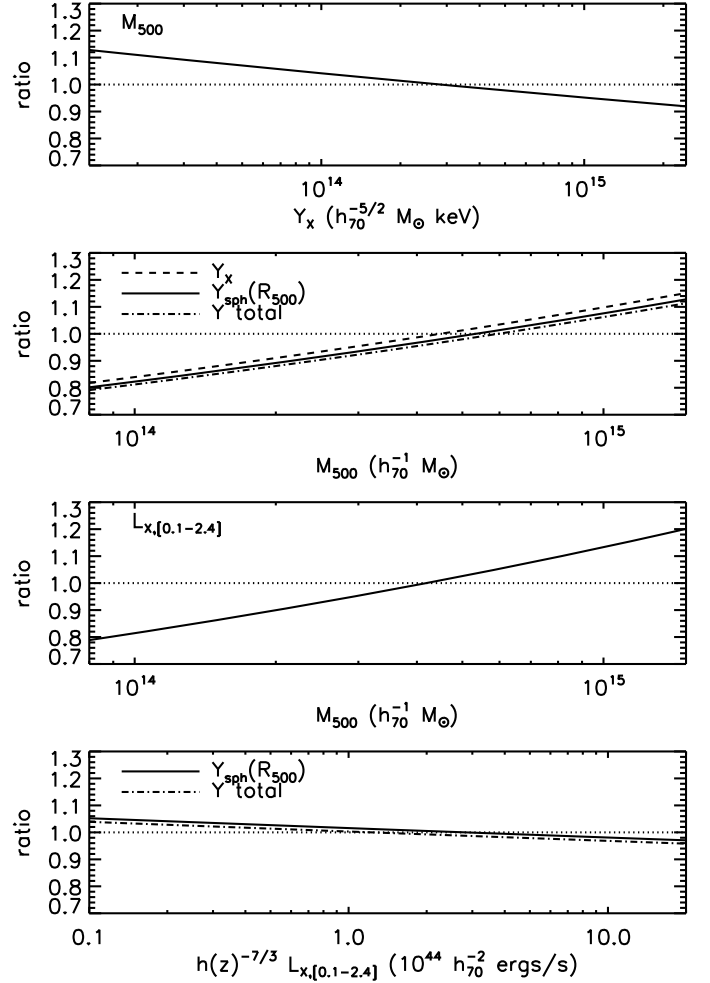


Fig. B.1. Ratio of the scaling relations derived using the empirical $M_{500} - Y_X$ relation (Eq. (2)) to those derived using the standard slope relation (Eq. (3)). From top to bottom: M_{500} as a function of Y_X ; Y_X , $Y_{\text{sph}}(R_{500})$ and total Y_{SZ} as a function of M_{500} ; L_X as a function of M_{500} ; $Y_{\text{sph}}(R_{500})$ and total Y_{SZ} as a function of L_X .

mass, by $\sim +10\%$ at $M_{500} = 10^{15} M_\odot$ (second panel). The behavior of Y_{SZ} closely follows that of Y_X (same panel) simply because the ratio of the two purely depends on the shape of the universal profile. This shape is barely affected by the small difference in R_{500} values used to scale the physical pressure profiles. Similarly, the $Y_{\text{SZ}} - L_X$ relation only depends on cluster internal structure and is essentially the same in the two models (bottom panel). $Y_{\text{SZ}}(R_{500})$ is slightly higher/lower at low/high L_X following the change of R_{500} at given L_X . As the $M_{500} - Y_X$ is shallower than the ST relation, the $M_{500} - L_X$ is also shallower (thus higher R_{500} at low L_X) or equivalently the $L_X - M_{500}$ is steeper (third panel).

Appendix C: Pressure profiles and best fitting model

Here we list the physical cluster properties and the parameters of the GFW model best fitting each profile (Table C.1). Individual profiles and their best fitting model are plotted in Figs. C.1–C.3. We also provide the GFW parameters for the average scaled profile of the cool core and morphologically disturbed clusters in Table C.2.

Table C.1. Cluster physical parameters.

Cluster	R_{500} (Mpc)	Y_X ($10^{14} M_\odot \text{ keV}$)	$Y_{\text{sph}}(R_{2500})$ (10^{-5} Mpc^2)	$Y_{\text{sph}}(R_{500})$ (10^{-5} Mpc^2)	P_{500} ($10^{-3} \text{ keV cm}^{-3}$)	P_0	c_{500}	α	γ	$\chi^2/\text{d.o.f.}$
RXC J0003.8+0203	0.879	0.763 ± 0.030	0.410 ± 0.009	0.990 ± 0.036	1.466	3.93	1.33	1.41	0.567	0.3/9
RXC J0006.6-3443	1.075	2.35 ± 0.13	1.030 ± 0.050	3.06 ± 0.16	2.292	3.27	1.10	1.41	0.408	0.0/1
RXC J0020.7-2542	1.056	2.253 ± 0.072	1.419 ± 0.034	2.80 ± 0.11	2.331	20.26	2.16	1.37	0.035	3.7/7
RXC J0049.4-2931	0.800	0.477 ± 0.022	0.277 ± 0.010	0.630 ± 0.037	1.254	8.58	1.31	1.07	0.422	0.2/4
RXC J0145.0-5300	1.112	2.819 ± 0.097	1.193 ± 0.029	3.89 ± 0.18	2.461	9.73	1.06	1.06	0.000	1.1/4
RXC J0211.4-4017	0.684	0.203 ± 0.006	0.101 ± 0.003	0.267 ± 0.010	0.902	8.97	1.04	0.93	0.267	3.0/6
RXC J0225.1-2928	0.683	0.185 ± 0.014	0.087 ± 0.004	0.237 ± 0.017	0.832	19.28	1.19	0.88	0.000	5.4/5
RXC J0345.7-4112	0.685	0.188 ± 0.009	0.109 ± 0.003	0.227 ± 0.009	0.836	3.68	1.65	1.67	0.690	1.1/7
RXC J0547.6-3152	1.148	3.59 ± 0.11	1.976 ± 0.037	4.54 ± 0.14	2.799	8.52	1.74	1.51	0.260	3.8/6
RXC J0605.8-3518	1.059	2.285 ± 0.070	1.264 ± 0.025	3.13 ± 0.14	2.338	4.23	0.88	0.96	0.659	1.1/6
RXC J0616.8-4748	0.947	1.194 ± 0.044	0.515 ± 0.014	1.627 ± 0.060	1.784	4.06	1.16	1.43	0.234	1.4/3
RXC J0645.4-5413	1.302	7.291 ± 0.248	3.60 ± 0.11	9.93 ± 0.47	3.722	11.10	0.94	0.89	0.265	2.5/5
RXC J0821.8+0112	0.753	0.325 ± 0.017	0.171 ± 0.007	0.400 ± 0.019	1.053	1.72	1.37	2.01	0.860	1.5/1
RXC J0958.3-1103	1.076	2.64 ± 0.25	1.72 ± 0.11	3.42 ± 0.40	2.553	4.13	1.77	2.07	0.719	0.0/3
RXC J1044.5-0704	0.939	1.189 ± 0.024	0.732 ± 0.010	1.550 ± 0.051	1.820	7.08	1.27	1.05	0.644	13.7/7
RXC J1141.4-1216	0.893	0.879 ± 0.018	0.491 ± 0.007	1.199 ± 0.046	1.597	4.42	1.08	1.08	0.652	15.3/6
RXC J1236.7-3354	0.758	0.335 ± 0.011	0.162 ± 0.003	0.479 ± 0.020	1.062	47.76	0.72	0.61	0.000	3.2/4
RXC J1302.8-0230	0.850	0.625 ± 0.020	0.305 ± 0.007	0.800 ± 0.039	1.349	3.63	1.09	1.21	0.519	14.8/6
RXC J1311.4-0120	1.351	9.27 ± 0.17	5.610 ± 0.084	11.60 ± 0.30	4.169	23.13	1.16	0.78	0.399	17.1/7
RXC J1516+0005	1.010	1.689 ± 0.050	0.927 ± 0.013	2.211 ± 0.083	2.035	4.48	1.52	1.65	0.474	4.1/5
RXC J1516.5-0056	0.932	1.105 ± 0.038	0.479 ± 0.015	1.494 ± 0.054	1.740	2.57	1.09	1.51	0.465	1.2/4
RXC J2014.8-2430	1.176	4.133 ± 0.097	2.293 ± 0.056	5.59 ± 0.23	2.971	4.94	0.75	0.82	0.684	8.8/7
RXC J2023.0-2056	0.740	0.281 ± 0.014	0.149 ± 0.005	0.358 ± 0.016	0.968	4.00	1.36	1.41	0.515	0.2/2
RXC J2048.1-1750	1.095	2.782 ± 0.084	1.104 ± 0.024	3.73 ± 0.12	2.542	4.34	1.33	1.76	0.000	10.7/3
RXC J2129.8-5048	0.903	0.856 ± 0.043	0.357 ± 0.016	1.147 ± 0.051	1.508	9.21	0.94	1.00	0.000	0.2/0
RXC J2149.1-3041	0.891	0.864 ± 0.024	0.429 ± 0.009	1.135 ± 0.051	1.585	9.96	0.71	0.71	0.446	3.3/6
RXC J2157.4-0747	0.753	0.311 ± 0.012	0.122 ± 0.005	0.411 ± 0.015	1.007	1.46	1.24	2.54	0.491	0.1/1
RXC J2217.7-3543	1.031	2.023 ± 0.050	1.079 ± 0.021	2.611 ± 0.077	2.260	27.70	1.18	0.81	0.133	0.2/5
RXC J2218.6-3853	1.147	3.51 ± 0.14	1.796 ± 0.049	4.94 ± 0.29	2.751	27.29	1.06	0.82	0.000	1.0/4
RXC J2234.5-3744	1.307	7.22 ± 0.17	4.300 ± 0.075	8.82 ± 0.25	3.647	25.04	2.01	1.23	0.000	10.6/5
RXC J2319.6-7313	0.793	0.445 ± 0.018	0.194 ± 0.004	0.612 ± 0.026	1.207	338.9	0.17	0.33	0.065	1.9/5

Notes. Columns (2), (3): R_{500} is the radius corresponding to a density contrast of 500, estimated iteratively from the $M_{500} - Y_X$ relation (Eq. (2)), where $Y_X = M_{g,500} T_X$ is the product of the gas mass within R_{500} and the spectroscopic temperature T_X . Columns (4) and (5): spherically integrated Compton parameter within R_{2500} and R_{500} , respectively. Column (6): P_{500} as defined by Eq. (5). Columns (7) to (10) give the best fitting GNF parameters for the pressure profiles (Eq. (11)). The external slope parameter β has been fixed to 5.49 (see text). Redshift z and M_{500} values can be found in Table 1 of Pratt et al. (2010).

Table C.2. Best fitting GNF parameters for the average pressure profile of the REXCESS sub-samples of cool core and morphologically disturbed clusters (Eq. (11)).

Sub-sample	P_0	c_{500}	α	γ
Cool-core	3.249	1.128	1.2223	0.7736
Morphologically disturbed	3.202	1.083	1.4063	0.3798

Notes. The external slope parameter β has been fixed to 5.49.

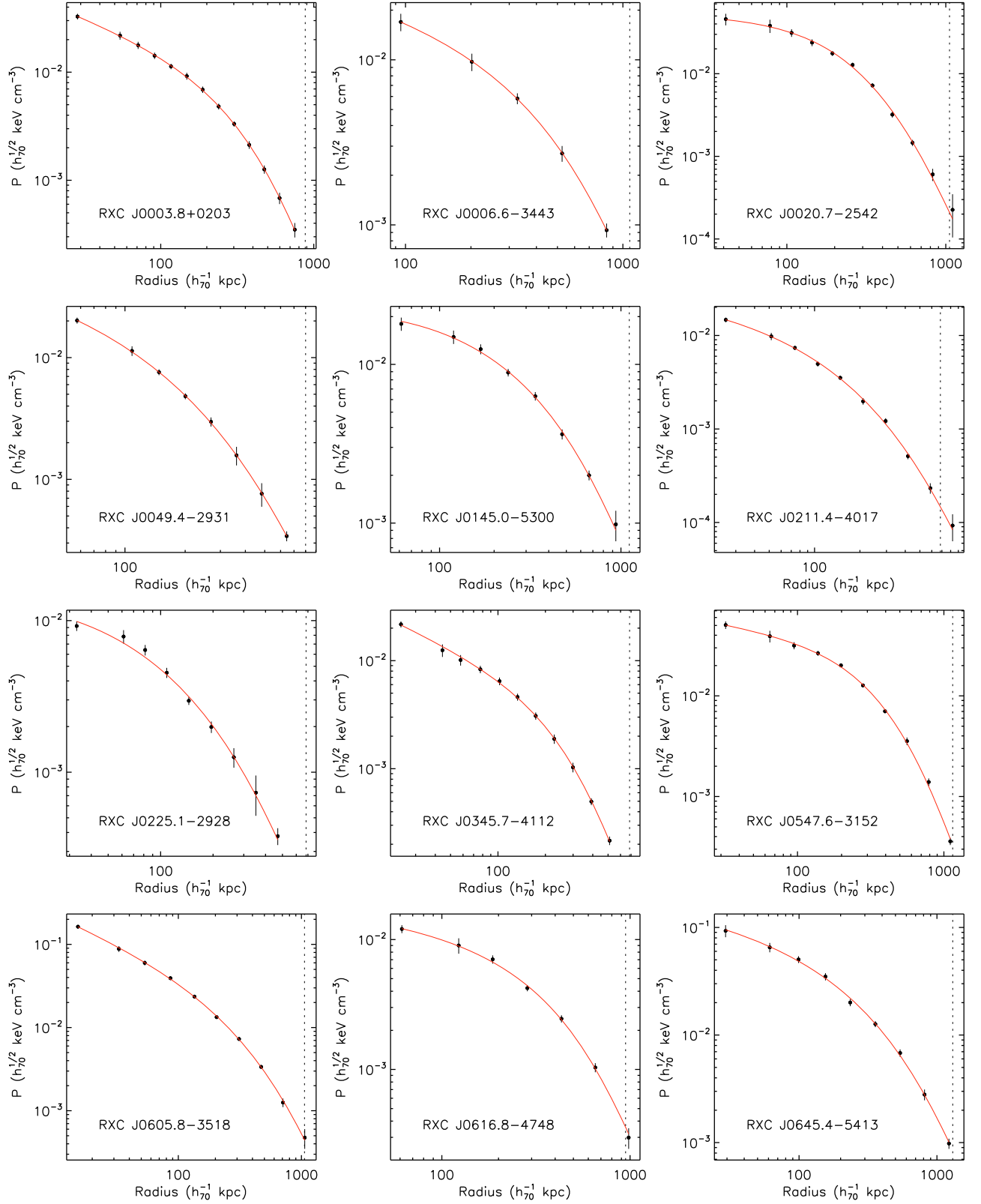


Fig. C.1. Pressure profiles for the entire REXCESS sample with the best fitting GFW model (red line). The dotted vertical line indicates R_{500} for each cluster.

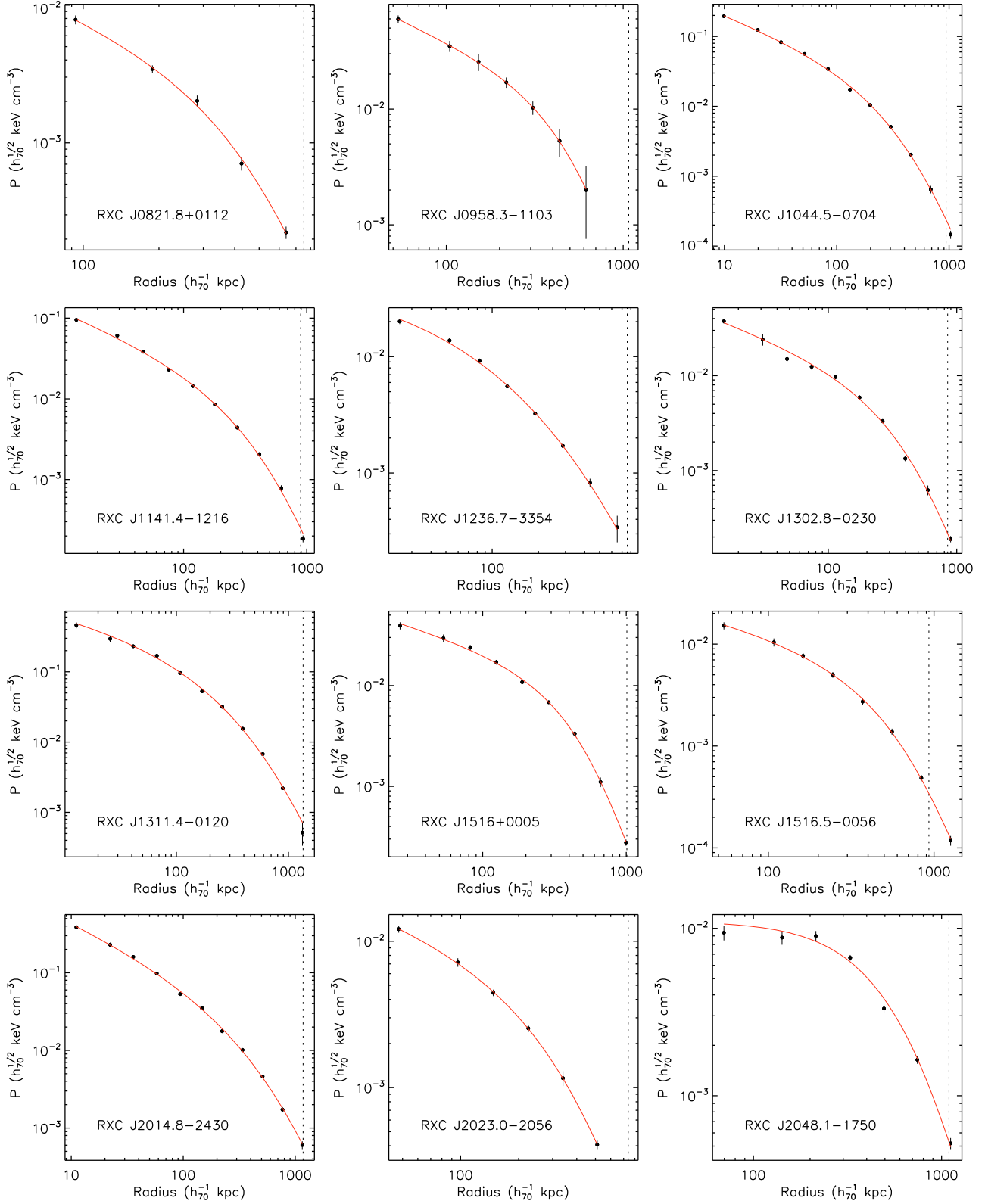


Fig. C.2. continued.

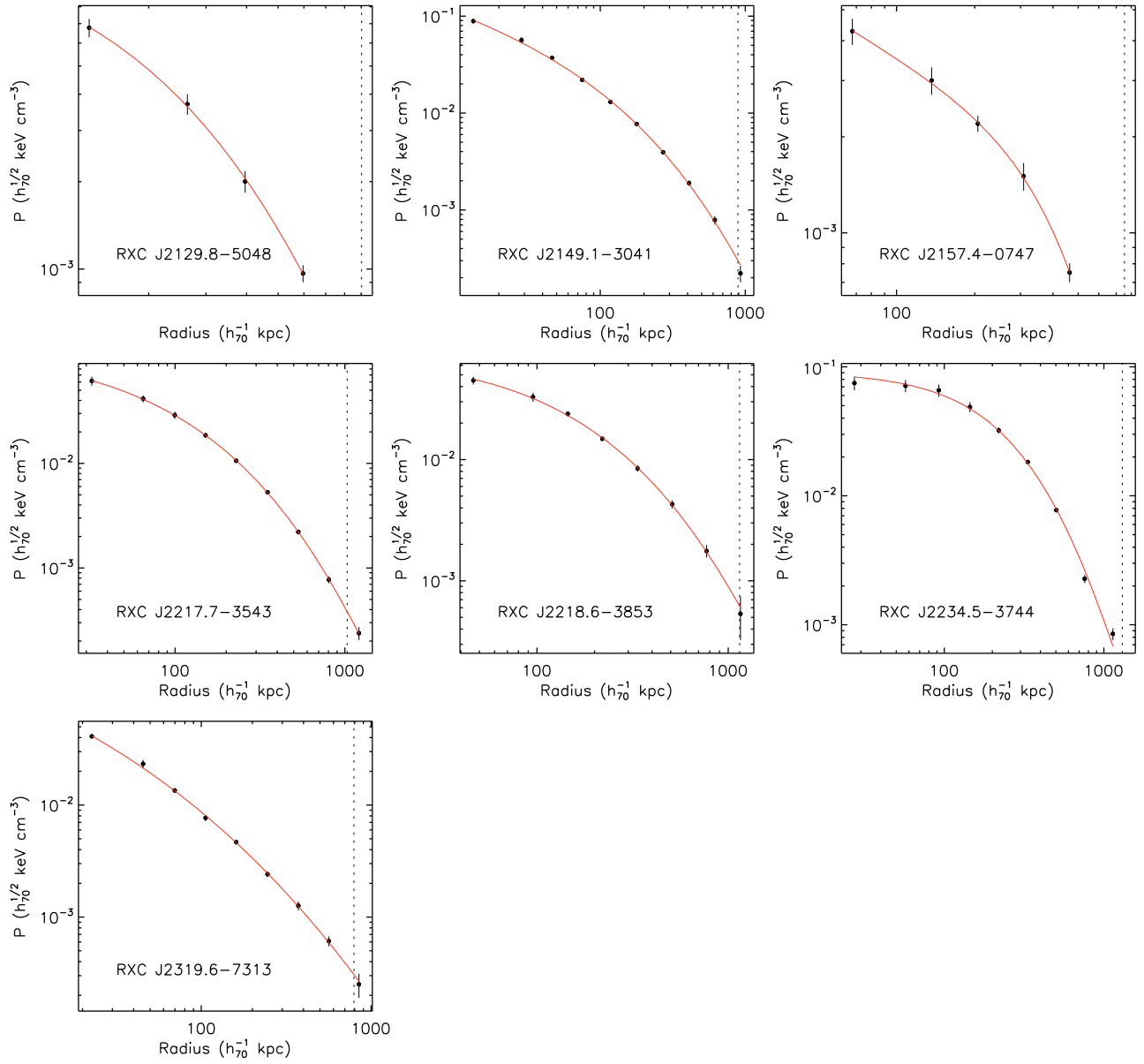


Fig. C.3. continued.