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The Western Bulge of 162173 Ryugu Formed as a Result of a Rotationally Driven Deformation Process

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Abstract

162173 Ryugu, the target of *Hayabusa2*, has a round shape with an equatorial ridge, which is known as a spinning top shape. A strong centrifugal force is a likely contributor to Ryugu's top-shaped features. Observations by the Optical Navigation Camera on board *Hayabusa2* show a unique longitudinal variation in geomorphology; the western side of this asteroid, later called the western bulge, has a smooth surface and a sharp equatorial ridge, compared to the other side. Here, we propose a structural deformation process that generated the western bulge. Applying the mission-derived shape model, we employ a finite element model technique to analyze the locations that experience structural failure within the present shape. Assuming that materials are uniformly distributed, our model shows the longitudinal variation in structurally failed regions when the spin period is shorter than ~ 3.75 hr. Ryugu is structurally intact in the subsurface region of the western bulge while other regions are sensitive to structural failure. We infer that this variation is indicative of the deformation process that occurred in the past, and the western bulge is more relaxed structurally than the other region. Our analysis also shows that this deformation process might occur at a spin period between ~ 3.5 and ~ 3.0 hr, providing the cohesive strength ranging between ~ 4 and ~ 10 Pa.

Key words: minor planets, asteroids: individual (162173 Ryugu)

1. Introduction

When an asteroid has a round shape with equatorial ridges, it is called a spinning top-shaped asteroid. Ground radar observations have shown that spinning top shapes are common in the solar system (Benner et al. 2015). 162173 Ryugu (Watanabe et al. 2019) and 101955 Bennu (Scheeres et al. 2016, 2019) are examples of spinning top shapes. Such asteroids may have natural satellites (e.g., Ostro et al. 2006; Brozovic et al. 2011; Fang & Margot 2011; Becker et al. 2015; Michel et al. 2016; Naidu et al. 2016). Importantly, a spinning top shape is found to be independent of material composition (Benner et al. 2015).

The formation of a spinning top-shaped asteroid may result from a reaccumulation process after catastrophic disruption (Michel et al. 2018, 2019) and/or quasi-static spin-up of a spheroidal body driven by micrometeorite impacts or a solar radiation pressure-driven torque, which is the so-called YORP

effect (Rubincam 2000). Regarding the quasi-static spin-up process, there are two spinning top-shape formation scenarios depending on the internal strength distribution: surface mass movement if there is a strong interior covered by a weak layer (Walsh et al. 2008, 2012) and internal deformation if there is uniform structure (Hirabayashi & Scheeres 2014, 2019; Hirabayashi et al. 2015; Zhang et al. 2017). Note that even if the rotational period is relatively long, the uniform structure may also cause surface failure (Hirabayashi 2015). All these processes cause the shape to settle into the minimum potential configuration (Scheeres 2015; Walsh 2018).

Ryugu is the target of the *Hayabusa2* mission led by JAXA (Watanabe et al. 2019). Figure 1 describes images taken by the telescope camera, which is part of the Optical Navigation Camera (ONC) system, on board the spacecraft (Kameda et al. 2017; Suzuki et al. 2018; Tatsumi et al. 2018). This asteroid has the common features of a spinning top-shaped asteroid, and its spin axis is perpendicular to the equatorial ridge

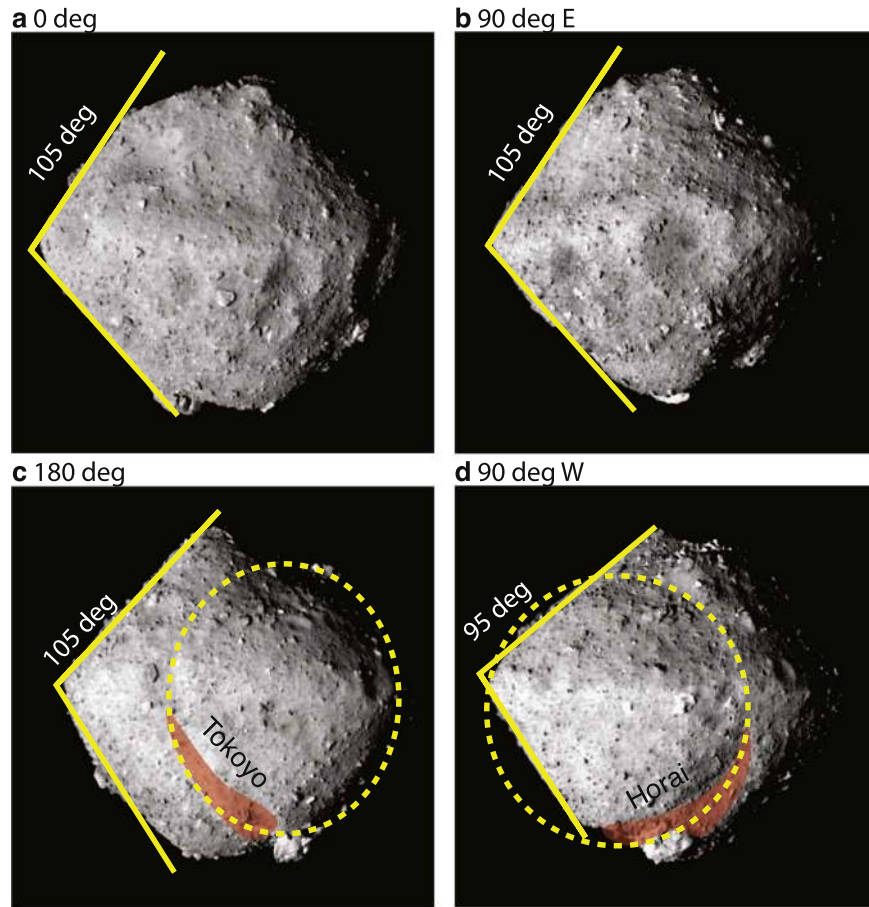


Figure 1. Ryugu seen from the *Hayabusa2* spacecraft at different views. (a) 0° (image: hyb2_onc_20180630_074839_tvf_12a). (b) 90° East (image: hyb2_onc_20180630_140543_tvf_12a). (c) 180° (image: hyb2_onc_20180630_120959_tvf_12a). (d) 90° West (image: hyb2_onc_20180630_101759_tvf_12a). The yellow circles indicate the western bulge. The red regions are Tokoyo and Horai Fossae. The ridge angles are determined by only considering the sunlit sides and by averaging the slopes at the middle latitudes. Image credit: JAXA/UTokyo/Kochi U./Rikkyo U./Nagoya U./Chiba Ins. Tech/Meiji U./U. Aizu/AIST.

(Watanabe et al. 2019). The currently reported spin period is 7.63262 hr (Watanabe et al. 2019), which is much longer than the ~ 2.3 hr spin barrier, the typical spin limit of small rubble pile asteroids (Pravec et al. 2008). The bulk density was estimated as ~ 1.19 g cm $^{-3}$ from the orbital motion of the *Hayabusa2* spacecraft (Watanabe et al. 2019).

The surface condition of Ryugu was observed to be longitudinally divided into the eastern and western regions (Sugita et al. 2019). These areas are distinguished by Tokoyo and Horai Fossae, a trough system that is widely placed in the southern hemisphere and possibly spreads toward the northern hemisphere (Figure 1). The western region (160° E–70° W), later known as the western bulge, is apparently less affected by crater bombardment and thus is smoother than the other side (Sugita et al. 2019). Furthermore, while Ryugu’s spinning top shape is axisymmetric in general; however, it is not strictly so (Figure 1). The ridge angle, which defines an angle between the surfaces in the northern and southern hemispheres at the equatorial ridge, or Ryujin Dorsum, varies when seen from different views. The ridge angle of the western bulge is 95° (panel (d) in Figure 1) and is sharper than that of other sides, which is 105° (panels (a) through (c) in Figure 1). This feature was also discussed in Watanabe et al. (2019) in detail.

This study tests a hypothesis that this longitudinal dichotomy results from shape deformation when Ryugu was rotating with a short rotational period in the past. The outline of this paper is as follows. First, we employ the finite element model (FEM)

analysis (Hirabayashi & Scheeres 2019) by assuming that the material distribution is uniform, which is compatible with observations (Sugita et al. 2019; Watanabe et al. 2019). Second, we use a dynamical analysis model to determine at what spin period Ryugu would have experienced the shape deformation process that caused the western bulge. Our work extends the structural analysis by Watanabe et al. (2019), who only focused on the general internal deformation process at spin periods of 3.75 and 3.5 hr, to analyze the deformation process that implies the formation of the western bulge at fast rotation.

2. How Did Ryugu Deform at a Short Spin Period?

We employ an FEM analysis (Hirabayashi et al. 2016; Hirabayashi & Scheeres 2019) to investigate the structural failure mode of Ryugu at a given spin period. This technique is the same as that done by Watanabe et al. (2019). In this work, we use the shape model developed from the stereophotoclinometry (SPC) technique (Watanabe et al. 2019). The version used is SHAPE_SPC_3M_v20180731.obj in which there are 1,579,015 vertices and 3,145,728 facets (Watanabe et al. 2019). To avoid computational intensity, we reduce the size of the shape model by using the Quadric Edge Collapse Decimation technique available in Meshlab (<http://www.meshlab.net/>). The newly derived shape model consists of 3072 facets and 1538 vertices.

We use an FEM solver available in ANSYS Mechanical APDL (18.1) under the license owned by Auburn University’s Samuel Ginn College of Engineering. We develop a three-dimensional, 10-node FEM mesh for the current shape of Ryugu by using TetGen, an open tetrahedron-mesh generator (Si 2015), and a mesh-refining tool in ANSYS (Hirabayashi & Scheeres 2019). The generated FEM mesh consists of 8308 elements and 15,038 nodes. This model assumes that spin-up acceleration is so small that dynamic change in stress is negligible. Thus, we solve the equilibrium equation of structure:

$$\frac{\partial \sigma_{ij}}{\partial x_j} + \rho b_i = 0, \quad (1)$$

where ρ is the bulk density, σ_{ij} is the stress component ($i, j = 1, 2, 3$), and b_j is the body force. Also, we use the Einstein notation in this equation. For elastic deformation, we use a linear-elastic rule, i.e., Hookean deformation. To describe inelastic deformation, we assume a plastic flow to be associated with a yield criterion where the behavior of materials becomes inelastic (the associate flow rule; Hirabayashi et al. 2016). A plastic flow is given as a function of a partial derivative of a yield criterion with respect to stress (Chen & Han 1988). To model the yield condition of regolith in Ryugu, we use the Drucker–Prager yield criterion, which depends on the cohesive strength and the angle of friction. In this work, the angle of friction is fixed at 35° , which is a typical value for regolith and soil materials (Lambe & Whitman 1969). On the other hand, we consider the cohesive strength to be a free parameter to investigate the failure mode at a given spin period (Hirabayashi & Scheeres 2019). The bulk density is fixed at 1.2 g cm^{-3} .

We introduce two critical parameters in our analysis. The first parameter is the smallest cohesive strength that can hold the present shape, which is later known as the minimum cohesive strength (Hirabayashi & Scheeres 2019). This parameter is an ideal value that could induce structural failure somewhere in Ryugu’s body. Thus, if the actual cohesive strength is lower than the minimum cohesive strength, structural failure should happen, causing deformation. We describe the minimum cohesive strength as a function of the spin period, which is called a failure mode diagram (FMD; Hirabayashi & Scheeres 2019). The second parameter is the stress ratio, a ratio of the stress state to the yield stress. If an element has a stress ratio reaching unity, it fails structurally.

We compute the failure mode of the present shape at nine spin periods: 2.5, 3.0, 3.5, 3.75, 4.0, 5.0, 6.0, 7.6, and 8.0 hr. Note that Watanabe et al. (2019) only introduced the cases at spin periods of 3.5 and 3.75 hr. Figure 2 shows the FMD of Ryugu. At a spin period longer than 3.75 hr, inelastic deformation may appear due to shear stress under compression if materials are nearly cohesionless. Such sensitive regions depend on local topography and are located in small regions on the surface (Figure 2). Thus, we rule out the scenario that inelastic deformation at slow rotation produced the western bulge. At fast rotation, on the other hand, because the centrifugal force plays a dominant role in the stress distribution, Ryugu needs high cohesive strength to hold the present shape. Failed regions start to appear in the interior of the body when the spin period is shorter than 3.75 hr (Watanabe et al. 2019). Also, the minimum cohesive strength monotonically increases as the spin period becomes short (Figure 2).

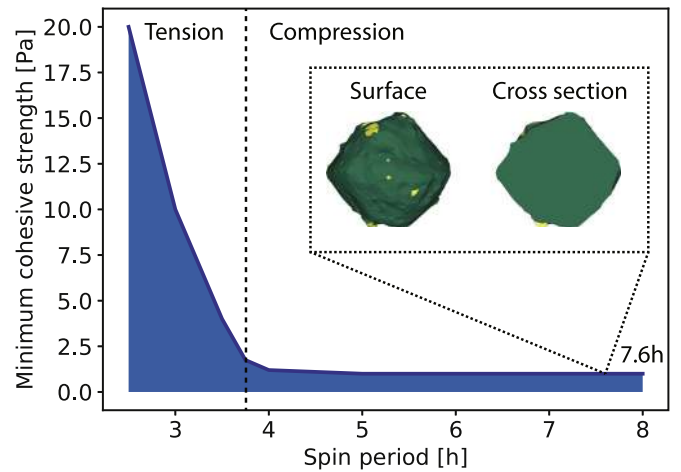


Figure 2. Failure mode diagram of Ryugu. The x axis is the spin period, while the y axis is the minimum cohesive strength. The dashed line indicates the boundary between compression-dominant deformation and tension-dominant deformation at a spin period of 3.75 hr. The shaded region indicates that Ryugu cannot structurally exist because the cohesive strength is below the minimum cohesive strength. The top right panel describes the failure mode of Ryugu at a spin period of 7.6 hr. The green and yellow regions are structurally intact and failed areas, respectively.

In this paper, to discuss the deformation process that made the western bulge, we employ the following approach. Our technique only accounts for the present shape and provides the failure mode that this shape experiences, meaning that we do not directly show the failure mode that the precursor shape had. To infer the deformation process in the past, we search for structurally intact points in the present shape. When a deformation process occurs at a given spin period, the deformed area may settle into a configuration in which the potential of the shape becomes minimum (Scheeres 2015; Walsh 2018). Therefore, when Ryugu encounters a fast-rotation condition once again, only the elements that did not experience structural failure previously become sensitive to failure. The originally failed regions, on the other hand, may not experience structural failure because they are structurally relaxed.

If the body is axisymmetric, the deformation process should be symmetric because the centrifugal and gravity forces are the main loadings. If an asteroid is usually irregular in shape, the failure mode becomes complex and asymmetric (Hirabayashi & Scheeres 2019). Our FEM results show that the structurally intact region corresponds to the western bulge. Figure 3 gives the failed regions viewed from the spin axis direction. Lines AA' , BB' , and CC' are the cross-sections at three different longitudes. Lines AA' and BB' cross the western bulge while Line CC' does not. While the equatorial plane and surface widely experience structural failure at spin periods of both 3.75 and 3.5 hr, the structurally intact regions appear at longitudes between 120° E and 90° W . We note that the 3.75 hr spin period gives wider failed regions than the 3.5 hr spin period because a transition between compression and tension happens at a spin period of ~ 3.75 hr, causing structurally failed regions to be larger (Hirabayashi 2015).

We also show the failed regions on the vertical cross-sections (Figure 4). Along AA' and BB' , structurally intact regions spread in the subsurface of the western bulge. On the other hand, the distribution of the failed region along CC' is almost symmetric along the spin axis direction. The

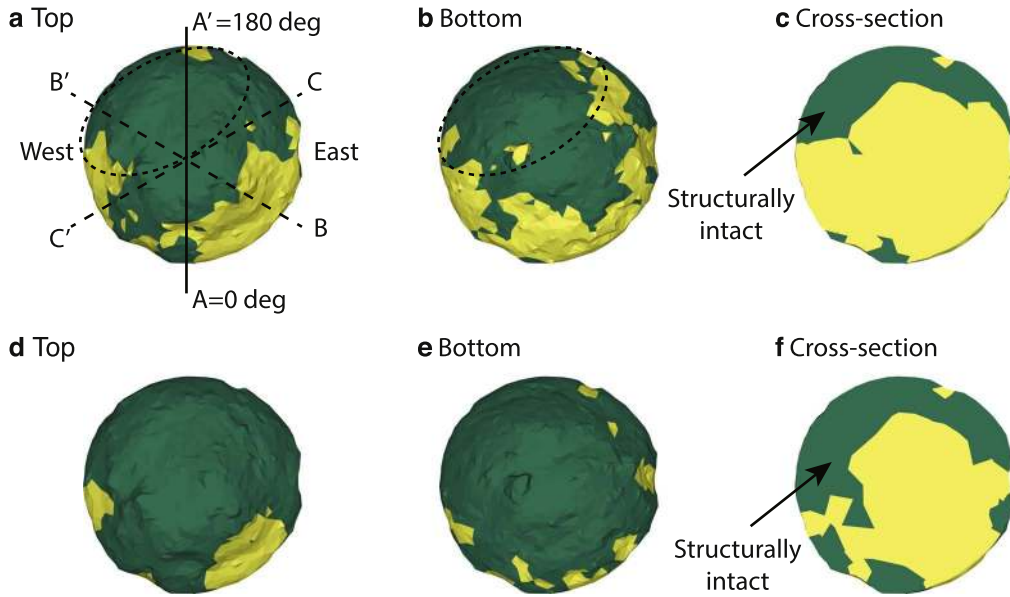


Figure 3. Distributions of the failed regions (yellow) at spin periods of 3.75 hr (top) and 3.5 hr (bottom). The left, middle, and right columns show the top surface, the bottom surface, and the cross-section, respectively. The circle with a dashed line corresponds to the western bulge (Sugita et al. 2019). Lines AA' , BB' , and CC' define the cross-sections, which are cut through longitudes of 0° and 180° , 60° E and 120° W, and 120° E and 60° W.

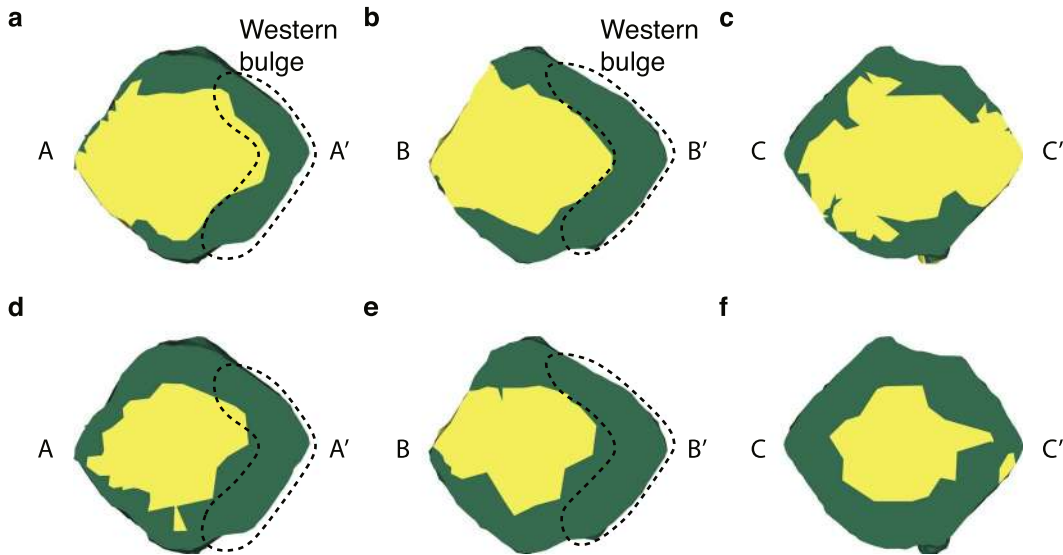


Figure 4. Distributions of the failed regions at spin periods of 3.75 hr (panels (a) through (c)) and 3.5 hr (panels (d) through (f)). Panels (a) and (d) show the failure mode on AA' , which is defined in Figure 3. Similarly, panels (b) and (e) indicate the failure mode on BB' , and panels (c) and (f) describe that on CC' . The dotted lines show the structurally intact regions that appear asymmetrically in the subsurface of the western bulge. Note that panels (c) and (f) are given in Watanabe et al. (2019).

FEM-derived minimum cohesive strength at spin periods of 3.75 hr and 3.5 hr are 1.75 Pa and 4 Pa, respectively (Watanabe et al. 2019). Our analysis indicates that the existence of the intact region beneath the western bulge cannot be explained by symmetric deformation. We propose that this results from a deformation process that made the western bulge.

3. How Fast did Ryugu Spin?

In the previous section, we argued that the western bulge might result from a deformation process. In this section, we analyze the rotational state that triggered this deformation process. During the deformation, which occurs at fast spin, fluidized materials with weak mechanical strength

(Melosh 2011) may reach the equatorial region. If the centrifugal force is strong, such materials should be lofted from there. Because the centrifugal force is proportional to the radial distance from an asteroid’s center of mass, this material shedding process can create a circular ridge (Scheeres et al. 2006), which supports the observed circularity of Ryugu (Watanabe et al. 2019).

To determine this spin period condition, we compute the locations of the dynamical equilibrium points at which the total force is balanced in the rotating frame. If these points touch the surface, fluidized materials should be lofted from there. The force balance condition is described as:

$$f_g(\mathbf{x}) + f_c(\mathbf{x}) = 0, \quad (2)$$

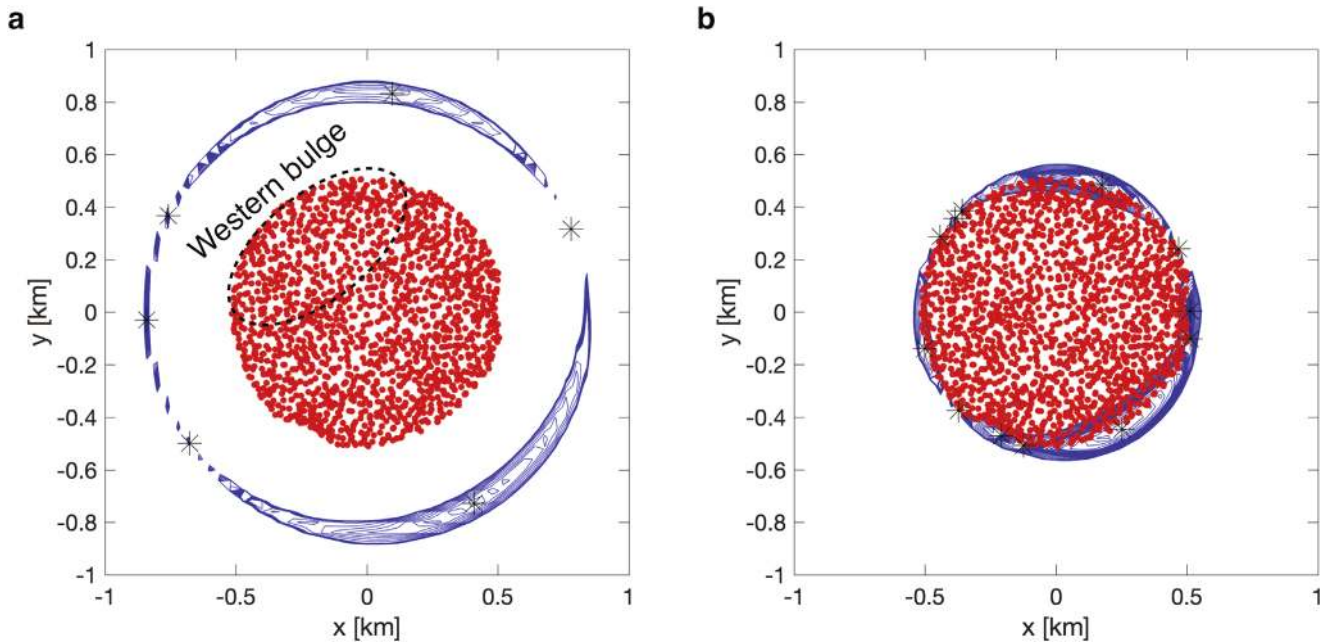


Figure 5. Locations of the dynamical equilibrium points. Panel (a) shows the case at a spin period of 7.6 hr. Panel (b) describes the case of a spin period of 3.5 hr. The red points indicate the shape of Ryugu. The star markers are the locations of the dynamical equilibrium points. The contour outside the shape describes the zero-velocity curves. The red dots describe the shape of Ryugu viewed from the spin axis. The dashed line indicates the location of the western bulge.

where f_g is the gravity force, and f_c is the centrifugal force. We compute the gravity force vector by using the algorithm developed by Werner & Scheeres (1997) and a refined shape model that consists of 3072 facets and 1538 vertices. These force terms are functions of the dynamical equilibrium points, \mathbf{x} . We numerically compute the dynamical equilibrium points by using Newton’s implicit equation solver (e.g., Press et al. 2007).

We also compute the zero-velocity curves, which show the pseudo-energy levels that correlate with the locations of the dynamical equilibrium points (Murray & Dermott 2008). Energetically, these points become either saddle points or local extrema. Seen from Ryugu’s center of mass, the region outside the dynamical equilibrium points may experience strong centrifugal forces, while the gravity is dominant in the area within these points. For spinning top-shaped asteroids, there are usually more than four dynamical equilibrium points because of their axisymmetric shapes (Scheeres et al. 2006).

Similar to the FEM analysis, we fix the bulk density at 1.2 g cm^{-3} . Figure 5 shows the locations of the dynamical equilibrium points and the zero-velocity curves at spin periods of 7.6 and 3.5 hr. The zero-velocity curves are drawn at pseudo-energy levels near the dynamical equilibrium points to enhance the energy conditions at these points. We find that at a spin period of 3.5 hr, 12 dynamical equilibrium points appear and are about to touch the surface. At this spin condition, the zero-velocity curves are circular, the size of which is consistent with Ryujin Dorsum.

We also give constraints on the cohesive strength of Ryugu. Structural failure should occur in Ryugu at a spin period at 3.5 hr if the cohesive strength is $\sim 4 \text{ Pa}$ (see Figure 2). This spin period condition may be an upper bound. When Ryugu deforms, its moment of inertia changes due to deformation. If this asteroid rotates at a given spin period before deformation, it should change its spin period after deformation in a manner that the angular momentum is conserved. Because of the

rotationally induced deformation mode of a spinning top-shaped asteroid (e.g., Hirabayashi & Scheeres 2014), the moment of inertia increases as deformation continues. Thus, before Ryugu’s precursor deforms to become the current shape at the 3.5 hr spin period, it may have a shorter spin period than that period.

To consider a lower bound of the spin period of deformation, we conduct the following thought experiment. Once Ryugu’s precursor structurally failed at a spin period shorter than 3.5 hr, the shape eventually settled into the current configuration at the 3.5 hr spin period. Assuming that it was spinning along the maximum moment of inertia axis, we conservatively choose a sphere that has the smallest moment of inertia to be the shape of the precursor body. Here, the spin period of this object before deformation is considered to be the lower bound. We assume that the total mass and volume are constant, ignoring the amount of shedding materials. Because the mean radius of Ryugu is 448.4 m (Watanabe et al. 2019), we obtain this spin period as 3.0 hr and the minimum cohesive strength as $\sim 10 \text{ Pa}$ (Figure 2). We conclude that the cohesive strength of Ryugu ranges between ~ 4 and $\sim 10 \text{ Pa}$.

4. Discussion

We showed that the subsurface structure of the western bulge is not sensitive to structural failure at present, implying that this region is likely to be relaxed structurally. We proposed that Ryugu experienced a deformation process that generated the western bulge. This process might occur only on the side of the bulge. Such an asymmetric deformation process was numerically predicted for a spheroidal object by Sánchez & Scheeres (2016), who pointed out that even small structural heterogeneity can cause an asymmetric deformation process, although theoretical studies missed this point (e.g., Holsapple 2010; Hirabayashi 2015). Our hypothesis explains the east–west geomorphological dichotomy. The smooth, less-cratered terrain on the western bulge can be generated due to mass movement.

Also, the deformation process can produce void space in the interior, followed by a material infilling process, which explains the formation of Tokoyo and Horai Fossae.

While our discussions were focused on rotationally induced failure, we note alternative possibilities. First, impact craters may contribute to deformation processes in Ryugu. However, we rule out this possibility as a direct contributor because this does not produce the circularity of Ryujin Dorsum. Second, a head-on collision of two similar-sized objects may produce the east–west dichotomy (Leleu et al. 2018). However, although the contact point should be highly compressed and have a ridge-like feature (Leleu et al. 2018), Ryugu does not have such a feature. Instead, we observe depression features, Horai and Tokoyo Fossae.


We note that our work does not determine whether internal deformation (e.g., Hirabayashi & Scheeres 2014, 2019; Scheeres et al. 2016; Zhang et al. 2017) or surface mass movement without internal deformation (e.g., Walsh et al. 2008, 2012; Zhang et al. 2017) is the primary cause of the considered asymmetric deformation process. If the internal structure is strong, Ryugu may have surface mass movement. On the other hand, if the interior is uniform, this asteroid may have internal deformation. We, however, rule out large density inhomogeneity between the western bulge and other regions because the center of mass of Ryugu derived by *Hayabusa2* corresponds to numerical prediction based on the mission-derived shape model with constant density within error (Watanabe et al. 2019). It is also possible that Ryugu may have a complex deformation processes if its rubble pile structure possesses internal discontinuities that formed during the reaccumulation process after catastrophic disruption of its parent body.

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