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Theory of Magnetic Domain Phases in Ferromagnetic Superconductors

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Recently discovered superconducting P-doped EuFe_2As_2 compounds reveal the situation when the superconducting critical temperature substantially exceeds the ferromagnetic transition temperature. The main mechanism of the interplay between magnetism and superconductivity occurs to be an electromagnetic one, and a short-period magnetic domain structure was observed just below Curie temperature [V. S. Stolyarov *et al.*, *Sci. Adv.* **4**, eaat1061 (2018)]. We elaborate a theory of such a transition and demonstrate how the initial sinusoidal magnetic structure gradually transforms into a solitonlike domain one. Further cooling may trigger a first-order transition from the short-period domain Meissner phase to the self-induced ferromagnetic vortex state, and we calculate the parameters of this transition. The size of the domains in the vortex state is basically the same as in the normal ferromagnet, but with the domain walls which should generate the set of vortices perpendicular to the vortices in the domains.

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The coexistence of magnetism and singlet superconductivity has always been of great interest because of their competing nature. Already in 1956, Ginzburg [1] showed that uniform magnetism in bulk systems may destroy superconductivity due to the electromagnetic (EM) mechanism (so-called, orbital effect), i.e., generation of the screening Meissner currents. In addition, the exchange field tends to align electron spins parallel to each other, which prevents the formation of Cooper pairs with the opposite spin directions [exchange (EX) mechanism] [2]. As a result, the coexistence of *uniform* ferromagnetism and superconductivity becomes possible primarily in thin-film structures with the damped orbital effect [1], spin-triplet uranium-based superconductors [3], or artificial superconductor-ferromagnet hybrids [4–6].

In contrast, *nonuniform* magnetic states may peacefully coexist with the superconducting ordering. The typical example is the antiferromagnetic superconductors RRh_4B_4 and RMO_6S_8 with the rare-earth element R [7], where the net magnetic moment at the scale of the superconducting coherence length ξ is zero and, thus, does not influence Cooper pairs. Somewhat similar cryptoferrimagnetic phases were predicted for the ferromagnetic superconductors (FSs) [8] and was later observed in ErRh_4B_4 [9] and HoMo_6S_8 [10] together with the reentrant superconductivity (see, e.g., Ref. [11] for a review).

The early theories of nonuniform magnetism in FSs accounted for only the EM interaction [12]. For the isotropic compounds, the EM effect favors the spiral magnetic texture in the superconducting phase instead of

the ferromagnetism [13,14] while magnetic anisotropy should trigger the formation of domain structures (DSs) [12,15] which can coexist with Abrikosov vortices [16–18]. However, further investigation of these intriguing phenomena in FSs with purely EM interaction was interrupted because it turned out that even a small exchange field producing negligible contribution to the magnetic energy should dramatically affect the magnetic texture of FSs [11]. In the late 1980s there were no FSs where EM interaction could dominate, and the research became mainly focused on the EX mechanism [8,19–23]. In principle, the effects of EM interaction could be observed in triplet FSs. However, in all three known triplet FSs [3] the Curie temperature θ is well above the superconducting critical temperature T_c so that below T_c the magnetic structure is already frozen and insensitive to the superconductivity.

Recently, the interest in FSs with purely EM interaction has been unexpectedly renewed with the discovery of the P-doped EuFe_2As_2 compound where $\theta < T_c$ and superconductivity coexists with ferromagnetism in a broad temperature interval [24–29]. The rather large critical temperature $T_c \sim 20\text{--}30$ K in the ferroarcentide family [30] and the robustness of superconductivity towards a disorder strongly support the *s*-wave character of the superconducting pairing. In such a case, the exchange field h_{EX} generated by Eu atoms in the low temperature ferromagnetic phase should be rather weak. The upper critical field $H_{c2}(T)$ in EuFe_2As_2 compounds is characterized by a large slope at T_c : $dH_{c2}/dT \sim 3$ T/K [31], which corresponds to the small superconducting coherence

length $\xi \sim 1.5$ nm. A very strong polarization of the Eu subsystem is achieved at fields $\gtrsim 1$ T [29], but it does not result in any observable decrease of the transition temperature [31] and we may conclude that $h_{\text{EX}} \ll T_c$ (hereinafter we put the Boltzmann constant $k_B = 1$). In addition, the time resolved magneto-optical measurements in EuFe_2As_2 [32] reveal a very slow relaxation time for Eu^{+2} spin $\tau \sim 100$ ps, which implies the exchange interaction $h_{\text{EX}} \sim \hbar/\tau \sim 0.1$ K. Moreover, the Mössbauer studies [26,33] and density-functional band-structure calculations [34] indicate that the exchange interaction between Eu atoms and superconducting electrons in EuFe_2As_2 and similar compounds is very weak, $h_{\text{EX}} \lesssim 1$ K, and then the exchange Ruderman-Kittel-Kasuya-Yosida contribution into the magnetic energy is $\theta_{\text{EX}} \sim N(E_F)h_{\text{EX}}^2 \sim 10^{-3}$ K [for the electron density of states $N(E_F) \sim 2-3$ states/eV per one Eu atom [30]].

The strong spin-orbit scattering in EuFe_2As_2 is likely to suppress the paramagnetic mechanism of superconductivity destruction. When the spin-orbit electron scattering mean free path l_{so} is of the order of the ordinary mean free path l ($l_{\text{so}} \gtrsim l$), the EM interaction dominates over the exchange one in the nonuniform magnetic structure formation, if $\theta_{\text{EX}} < \theta(a\xi^2/l^3)$, where a is of the order of interatomic distance [35]. The small value of the superconducting coherence length in EuFe_2As_2 ensures the domination of the EM mechanism.

The unusual relation $\theta < T_c$ in EuFe_2As_2 provides access to the almost unexplored situation when the ferromagnetism nucleates in a fully developed superconducting state. Recent pioneering experiments on the high-resolution visualization of the magnetic texture in EuFe_2As_2 [36,37] provide the first direct evidence of the transitions from the short-period domain Meissner state to the phase where magnetic domains coexist with Abrikosov vortices. Interestingly, these transitions reveal hysteresis behavior when varying the temperature [37]. However, despite the rapid experimental progress, the theory of the magnetic state's evolution in anisotropic FSs with purely EM interaction is still lacking.

In this Letter we present the theory of magnetic domain phases in FSs with low Curie temperature and purely EM interaction. We demonstrate how magnetic domains evolve

from a sinusoidal profile to the steplike structures when cooling the sample. Additionally, we calculate the key parameters of the first-order transition to the phase with coexisting domains and vortices and suggest an explanation for the hysteresis behavior of the domain structure in EuFe_2As_2 [36,37]. Finally, we show that the domain walls favor the generation of unusual vortices with the cores perpendicular to the vortices in the domains.

Before going into detail we briefly overview possible regimes in the evolution of magnetic texture with the variation of temperature T (see Fig. 1). In the cooling process at $T = T_c$ the superconducting Meissner phase appears. The well-developed superconductivity prevents the nucleation of uniform ferromagnetism at $T = \theta$. As a result, the magnetic order emerges at the temperature $\theta_m < \theta$ and the magnetization has the sine profile with only one spatial harmonic. While cooling below θ_m the nonlinear effects give rise to other spatial harmonics and the magnetization evolves towards the well-developed steplike DS with increasing domain size. At temperature T_{M-V} the growing amplitude of the magnetization makes the uniform superconducting phase less favorable than the phase with coexisting DS and vortex lattice. However, in the cooling regime the immediate vortex entry at $T = T_{M-V}$ is prevented by the Bean-Livingston-like barrier. This results in overcooling of the Meissner state and the first-order phase transition to the vortex state (VS) occurs only at $T = T_{cm} < T_{M-V}$ when the barrier vanishes. At the same time, in the heating regime the system stays in the VS until $T = T_{M-V}$, thus demonstrating the hysteresis behavior.

As is well known, in a ferromagnetic thin-film sample with perpendicular anisotropy the domain structure appears in order to minimize the stray field. The period of such a structure depends on the thickness of the sample and usually exceeds micron size, and it is much larger than that in the domain Meissner phase. Note that the formation of the short-period domain Meissner state is related to the volume effect of the interaction between magnetism and superconductivity, while the existence of the domains in normal ferromagnets is related to its demagnetization factor (shape effect). In the considered case, these ferromagnetic domains will be in the vortex state.

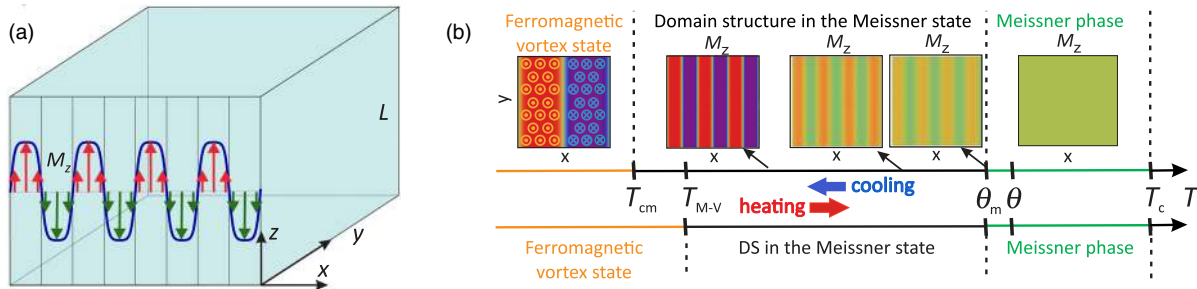


FIG. 1. (a) Ferromagnetic superconductor (FS) with the easy-axis magnetic anisotropy along the z axis. (b) The phase diagram demonstrating the temperature evolution of the coexisting phases in FS with dominant EM mechanism.

To support the above qualitative picture we calculate the temperature evolution of the magnetic texture in the FS with $\theta < T_c$ using the Ginzburg-Landau approach. The free-energy functional describing the FS with the strong easy-axis anisotropy along the z axis reads [11]

$$F(B_z, M_z) = \frac{(B_z - 4\pi M_z)^2}{8\pi} + \frac{A^2}{8\pi\lambda^2} + \frac{n\tilde{\theta}}{M_0^2} \left[\tau M_z^2 + \frac{b M_z^4}{2M_0^2} + a^2 \left(\frac{\partial M_z}{\partial x} \right)^2 \right]. \quad (1)$$

Here, $\tau = (T - \theta)/\theta$ is the reduced temperature, $\mathbf{M} = M_z(x)\mathbf{z}$ is the magnetization, $\mathbf{B} = B_z(x)\mathbf{z}$ is the magnetic field with the corresponding vector potential $\mathbf{A} = A(x)\mathbf{y}$, λ is the London penetration depth, n is the concentration of magnetic atoms, M_0 is the saturation magnetization at $T = 0$, and $\tilde{\theta} = M_0^2/n$. The estimates for the coefficients a and b give $b \sim 1$ and $a \ll \lambda$.

In the cooling regime the first sinusoidal harmonic of magnetization characterized by the wave vector q_m emerges at $T = \theta_m \lesssim \theta$. To calculate the shift $\tau_m = (\theta - \theta_m)/\theta$ one may neglect the term $\propto M_z^4$ in Eq. (1) and make the Fourier expansion $M_z(x) = \int M_q \exp(iqx) dq/2\pi$. Then using Maxwell equations we rewrite the averaged free energy $\bar{F}V = \int F dV$:

$$\bar{F} = \sum_q |M_q|^2 [2\pi/(1 + \lambda^2 q^2) + \Gamma(\tau + a^2 q^2)], \quad (2)$$

where $\Gamma = n\tilde{\theta}/M_0^2$. The condition $\bar{F} = 0$ defines the dependence $\tau(q)$ whose minimum corresponds to the actual temperature shift τ_m . The result depends on the value λ . If $\lambda < a\sqrt{\Gamma/2\pi}$, only the uniform state with $q_m = 0$ should appear, while for $\lambda > a\sqrt{\Gamma/2\pi}$, the free-energy minimum corresponds to the sinusoidal profile $M_z(x)$ with $q_m = (2\pi/\Gamma\lambda^2 a^2)^{1/4}$, and we find $\tau_m = (2a/\lambda)\sqrt{2\pi/\Gamma}$. The period of the emerging magnetic structure is smaller than λ , which makes this structure compatible with superconductivity due to the weak Meissner screening.

With further cooling below θ_m the emerging higher harmonics M_q result in the crossover from the sine magnetization profile to the steplike domains $M_z(x) = \pm M$ with the increasing size. The wave vector Q of the domain structure is determined by the balance between the energy of Meissner currents tending to increase Q and the domain walls' energy, which favors small Q values. In the limit $Q\lambda \gg 1$, the first contribution is proportional to $M^2/(\lambda^2 Q^2)$ while the estimate for the energy of the linear domain walls appearing in the systems with strong magnetic anisotropy gives $\Gamma a\sqrt{|\tau|}M^2 Q$. The minimization of the resulting free energy shows that the wave vector $Q \sim q_m(\tau_m/|\tau|)^{1/6}$ decreases when cooling the sample.

The above conclusion is perfectly supported by the accurate calculations. In the easy-axis ferromagnets one

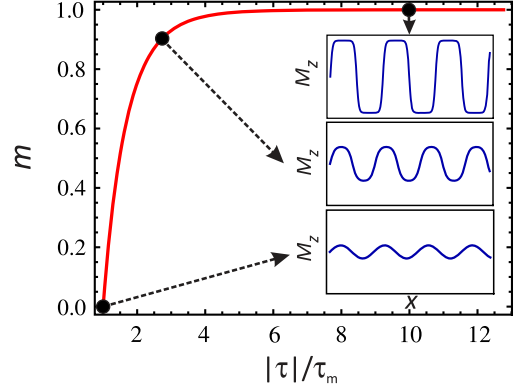


FIG. 2. The dependence of the parameter m describing the form of magnetic domains on temperature $\tau = (T - \theta)/\theta$.

can choose the ansatz for the magnetization in the form of the elliptic sine function: $M_z(x) = M \text{sn}[2K(m)Qx/\pi]$. Here $K(m)$ is the elliptic integral and the parameter m controls the shape of the $M_z(x)$ profile. Such an ansatz perfectly describes the gradual transition between the sine magnetization ($m = 0$) and the steplike one ($m \rightarrow 1$). Substituting the Fourier components of $M_z(x)$ into Eq. (2), restoring the term $\propto M_z^4$, and minimizing the resulting functional, we obtain the analytical expressions reflecting the temperature evolution of the values m , Q , and M [38]. The results confirm the very fast emergence of the well-developed DS below θ_m (see Fig. 2) with the increasing domain size (see Fig. 3). In the most interesting case of the well-developed DS, i.e., $m \rightarrow 1$, the wave vector Q , magnetization M , and free energy \bar{F} take the form:

$$M = \frac{M_0 \sqrt{\tau_m}}{\sqrt{b}} \left[\frac{|\tau|}{\tau_m} - \frac{2^{2/3}}{3} \left(\frac{|\tau|}{\tau_m} \right)^{1/3} \right]^{1/2}, \quad (3)$$

$$Q = q_m \frac{\pi}{2^{4/3}} \left(\frac{\tau_m}{|\tau|} \right)^{1/6}, \quad \bar{F} = -\Gamma \frac{b M^4}{2 M_0^2}. \quad (4)$$

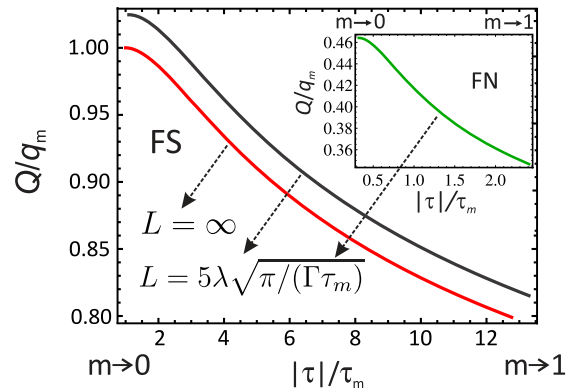


FIG. 3. The dependence of DS wave vector Q on temperature $\tau = (T - \theta)/\theta$ for infinite FS (red curve) and FS of the thickness L (black curve). Inset: The same for FN slab.

Note that well below θ the growing magnetization $M(|\tau|)$ may become large enough to induce Abrikosov vortices. The resulting coexistence phase should emerge through the first-order transition. In the presence of the vortex lattice the screening Meissner currents are small and, thus, at temperatures below the transition point instead of DS one should have the uniform magnetization. The thermodynamic critical temperature T_{M-V} of such a transition can be obtained from the comparison between the DS and VS free energies. For the well-developed steplike profile $M_z(x)$, the former energy takes the form Eq. (4) while the latter one reads [39]

$$F_v = \frac{(B_z - 4\pi M_z)^2}{8\pi} + \Gamma \left(\tau M_z^2 + \frac{b M_z^4}{2 M_0^2} \right) + \frac{B_z \tilde{H}_{c1}}{4\pi}. \quad (5)$$

Here $\tilde{H}_{c1} = H_{c1} \ln[\beta d/\xi]/\ln[\lambda/\xi]$, where H_{c1} is the lower critical field, ξ is the superconducting coherence length, $d^2 = 2\Phi_0/(\sqrt{3}B_z)$, Φ_0 is the superconducting flux quantum, and $\beta = 0.381$ is the geometrical factor relevant for the triangular vortex lattice.

In the case $\tilde{H}_{c1} \ll 4\pi M_z$, which is typical for FSs minimizing Eq. (5) with respect to the magnitude of the uniform magnetization M_z , we obtain the free energy of the VS:

$$F_v = -\Gamma M_0^2 \tau^2 / (2b) + M_0 \tilde{H}_{c1} \sqrt{|\tau|/b}. \quad (6)$$

For the reasonable choice of parameters we expect $|\tau_{M-V}| = (\theta - T_{M-V})/\theta \gg \tau_m$; thus, the temperature of the phase transition between the DS and the VS is

$$|\tau_{M-V}| = \frac{\tau_m}{4} \left(\frac{24\pi\sqrt{b}}{\Gamma} \frac{1}{\tau_m^{3/2}} \frac{\tilde{H}_{c1}}{4\pi M_0} \right)^{6/5}. \quad (7)$$

However, the evolution of the magnetic order near the temperature T_{M-V} should reveal hysteresis behavior. Indeed, in the cooling process the vortices cannot enter the sample at T_{M-V} because of the Bean-Livingston barrier which vanishes only at the temperature $T_{cm} < T_{M-V}$. The profile of this barrier $U(x)$ is determined by the interplay between the energies of the vortex interaction with Meissner currents and with the antivortex located in the neighboring domain. For the steplike profile of the magnetization, the profile $U(x)$ has the form

$$U(x) = -\frac{\Phi_0}{4\pi} \int_0^x \frac{dB_v(2\tilde{x})}{d\tilde{x}} d\tilde{x} - \frac{4\Phi_0 M}{\pi\lambda^2 Q^2} \sum_{n=0}^{\infty} \frac{\sin[Qx(2n+1)]}{(2n+1)^3}, \quad (8)$$

where B_v is the magnetic field produced by the antivortex. The condition of the barrier vanishing $(dU/dx)|_{x \rightarrow \xi} = 0$ allows us to calculate the value $\tau_{cm} = (T_{cm} - \theta)/\theta$ using Eqs. (4) and (3):

$$|\tau_{cm}| = \frac{[\sqrt{1+9\gamma}+1]^{3/2}}{\sqrt{54}} \tau_m, \quad \gamma = \frac{4\pi b}{\Gamma \tau_m^2} \left(\frac{H_{cm}}{4\pi M_0} \right)^2, \quad (9)$$

where H_{cm} is the thermodynamic critical field.

Taking the parameters of EuFe_2As_2 [29], we may estimate the ratio between T_{M-V} and T_{cm} . Since τ_m is rather small $\tau_m \sim 10^{-3} \ll 1$, we have $\gamma \gg 1$ and $|\tau_{cm}| \approx (\gamma^{3/4}/\sqrt{2})\tau_m$, and, thus, T_{M-V} significantly exceeds T_{cm} . However, the calculated value of T_{cm} may be substantially smaller than the temperature of the vortex entry in experiment. Indeed, the presence of accidental vortices trapped, e.g., at the defects, increases T_{cm} since the Meissner currents near the vortex cores are of the order of the depairing current, which favors the formation the additional vortex-antivortex pairs. Thus, in real type-II superconductors we expect $T_{cm} \lesssim T_{M-V}$.

Up to now we have not accounted for the finite size of the sample. For the slab of the finite thickness L , the domain structure produces the stray magnetic field which decays at distances $\propto Q^{-1}$ from the slab surface. The corresponding contribution to the free energy tends to shrink the domains, thus making the vector Q higher for the thinner samples at a fixed T (see Fig. 3) [38]. However, even for the thin-film FSs the stray field effect does not qualitatively change the dependence $Q(\tau)$ and the domain size remains increasing upon cooling the sample.

Interestingly, in the finite samples the emergence of the dense vortex lattice below T_{cm} on top of magnetic DS should result in substantial growth in the domain size. In this coexistence phase the Meissner screening is almost destroyed and the situation is fully analogous to the nonsuperconducting ferromagnetic film (FN). In such a film the ferromagnetic domain structure should appear to minimize the stray field. Calculating the dependence $Q(\tau)$ for this case we find that at a given T the domain is significantly larger if the superconductivity is destroyed (see the inset in Fig. 3). Such an increase of the domain size associated with the transition to the VS was clearly observed in Refs. [36,37]. Note that a similar effect should exist in FSs with $T_c < \theta$ [40–42].

On experiment (see Supplemental Material to Ref. [37]) the observed in EuFe_2As_2 low temperature domain structure is very similar to the branched domain patterns in normal ferromagnets (see, for example, Sec. 5.2 in Ref. [43]). The dense vortex structure in ferromagnetic domains in EuFe_2As_2 makes them equivalent to that in the normal ferromagnets. The recent observation of the low temperature domain structure in the overdoped normal EuFe_2As_2 shows that it is basically the same as in the optimally doped superconducting EuFe_2As_2 [44].

The reason for the branching of the domains in ferromagnets is related to the fact that the stray field energy increases with the increase of the domain size faster than the domain wall energy and starting some critical thickness

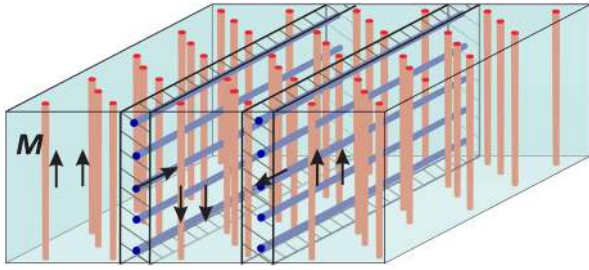


FIG. 4. “Perpendicular” vortices appearing at the domain walls of the Bloch type.

the branching of the domain becomes energetically favorable [45]. At low temperature the internal magnetic field in EuFe_2As_2 is rather large $B(0) \approx 9$ kOe (which strongly exceeds the low critical field) [37] and the period of the Abrikosov lattice is of the order of 40 nm. So it is much smaller than the characteristic size of the domain pattern and then the vortices simply decorate without changing the usual mechanism of the formation of rria-cost magnetic domain structure observed in EuFe_2As_2 [37] at low temperature. Note that the calculation of the branched domain patterns in the ferromagnets is beyond the scope of this Letter.

Interestingly, in the VS near the domain walls the vortices can become oriented perpendicular to the vortices in domains (see Fig. 4). This phenomenon originates from the transformation of the linear domain walls with only one magnetization component M_z to the Bloch-type domain walls having an additional component M_y . This component favors the vortex core directed along the y axis near the domain walls. The comparison between the free energies [38] shows that for EuFe_2As_2 the linear domain walls exist in the temperature interval $\Delta T \sim 0.3$ K near θ , while for lower T the Bloch domain walls appear. The subsequent emergence of the perpendicular vortices is favorable if $4\pi M \geq 4\pi M_{\min} = (2\lambda/\pi\delta)H_{c1}$, where δ is the domain wall width [38]. For EuFe_2As_2 the estimates give $4\pi M_{\min} \sim 0.4$ kOe, while $4\pi M \sim 10^4$ Oe [29,37]. Thus, the condition of the perpendicular vortices’ appearance is fulfilled.

In summary, we have described the temperature evolution of the magnetic domain structures and vortex-domain coexistence in ferromagnetic superconductors with purely electromagnetic interaction. The developed theory not only describes the transition from the short-period Meissner domain phase to the vortex phase recently observed in EuFe_2As_2 , but also predicts the formation of “perpendicular vortices” at the domain walls.

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