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## Thermal Shields for Heat Loss Reduction in Siemens-Type CVD Reactors

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The use of thermal shields to reduce radiation heat loss in Siemens-type CVD reactors is analyzed, both theoretically and experimentally. The potential savings from the use of the thermal shields is first explored using a radiation heat model that takes emissivity variations with wavelength into account, which is important for materials that do not behave as gray bodies. The theoretical calculations confirm that materials with lower surface emissivity lead to higher radiation savings. Assuming that radiation heat loss is responsible for around 50% of the total power consumption, a reduction of 32.9% and 15.5% is obtained if thermal shields with constant emissivities of 0.3 and 0.7 are considered, respectively. Experiments considering different thermal shields are conducted in a laboratory CVD reactor, confirming that the real materials do not behave as gray bodies, and proving that significant energy savings in the polysilicon deposition process are obtained. Using silicon as a thermal shield leads to energy savings of between 26.5-28.5%. For wavelength-dependent emissivities, the model shows that there are significant differences in radiation heat loss, of around 25%, when compared to that of constant emissivity. The results of the model highlight the importance of having reliable data on the emissivities within the relevant range of wavelengths, and at deposition temperatures, which remains a pending issue. © 2016 The Electrochemical Society. [DOI: 10.1149/2.0171603jss] All rights reserved.

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### Scope

90% of the polysilicon currently produced worldwide is demanded by the photovoltaic (PV) market, leaving the remaining amount for the microelectronics industry.<sup>1,2</sup> The chemical route -via chemical vapor deposition (CVD) of high purity trichlorosilane (TCS) on a hot filament, the so-called Siemens technology- currently dominates polysilicon production. High quality polysilicon is obtained, at the expense of high energy consumption.<sup>3,4</sup>

In the case of polysilicon for PV (also known as solar grade silicon) the process accounts for between a quarter and a third of the total energy consumption.<sup>5–7</sup> Thus, lowering the energy consumption of the Siemens process is essential to achieving the two wider objectives for silicon-based PV technology: low production cost and low energy payback time.

Furthermore, current price levels also press polysilicon producsers to reduce their production costs even more if they are seeking a sustainable business.<sup>8</sup>

As radiation heat loss is the major contributor to energy 37 consumption,<sup>9</sup> in this work the potential of thermal shields to re-38 duce radiation heat loss in an industrial Siemens reactor is studied. 39 Thermal (radiation) shields have been implemented in a number of 40 CVD reactors; e.g. for layer deposition in superconducting devices 41 or for the epitaxial growth of silicon layers.<sup>10,11</sup> Several proposals 42 have recently been made for polysilicon production,<sup>12-14</sup> but to our 43 knowledge quantitative analysis supported by experimental data has 44 not been provided in any of them. 45

Radiation heat loss as regards thermal shields in a Siemens-type CVD reactor is studied first here using a theoretical model. Then, the theoretical results are compared with the experimental results in a laboratory Siemens reactor. Discussion of the latter will offer insights into the accuracy of theoretical calculations depending on the thermal shield materials' optical properties, highlighting the relevance of the variation in optical properties with the wavelength for thermal shield materials.

### Potential to Reduce Radiation Heat Losses

First, the radiative heat transfer phenomenon is briefly described and the radiation heat loss model is presented. Then, theoretical radiation heat loss calculations for different thermal shields in an industrial Siemens reactor are put forward.

**Radiative heat transfer.**—Radiative heat transfer - also known as thermal radiation - describes the science of heat transfer caused by electromagnetic waves. These electromagnetic waves have the property of traveling through a vacuum or matter-containing media. The temperature of the radiant body governs the thermal radiation emission, and it occurs in the 0.1 to  $100 \,\mu$ m wavelength range.<sup>15,16</sup> It is not the aim of this section to explain the thermal radiation phenomenon in detail, but to describe a number of concepts and properties of the radiation heat transfer mechanism that will support the arguments we develop in this document.

As regards the radiation properties, four dimensionless magnitudes 69 are defined: absorptance ( $\alpha$ ), reflectance ( $\rho$ ), transmittance ( $\tau$ ) and 70 emissivity ( $\epsilon$ ). Absorptance, reflectance and transmittance are defined 71 as the ratio of the total amount of radiation absorbed, reflected or 72 transmitted by a surface to the total amount of radiation incident on 73 the surface, respectively. The emissivity<sup>a</sup> Emissivity is defined as the 74 ratio of the power per unit area radiated by a surface to the power 75 per unit area radiated by a black body at the same temperature. These 76 properties for real surfaces are dependent on temperature, direction 77 and wavelength. The relationship indicated in Equation 1 is obtained 78 by applying the energy balance to any real surface. 79

$$\alpha + \rho + \tau = 1 \tag{1}$$

In addition, according to Kirchhoff's law, all opaque surfaces ( $\tau = 80$ 0) reach  $\varepsilon_{\lambda}(\lambda, T) = \alpha_{\lambda}(\lambda, T)$ .<sup>15,16</sup>

A black body is defined as any body that emits and absorbs the maximum possible radiation in all wavelengths, that is:  $\alpha = 1$ ,  $\rho =$  $\tau = 0$ . Plack's law defines the spectral radiated power of a black body. In addition, according to Stefan-Boltzmann's law the expression for the total radiation emitted per unit area of a black body is indicated in Equation 2; where *T* is the temperature and  $\sigma$  the Boltzmann constant.

$$E_b(T) = \sigma T^4$$
 [2]

However, the majority of the surfaces do not behave as black bodies; thus, the gray body concept arises. A gray body is any opaque body ( $\tau = 0$ ,  $\alpha + \rho = 1$ ) whose reflectance, absorptance and emissivity properties are non dependent on the wavelength. The behavior of many real surfaces can be approximated to that of a gray body; in Equation 3 the expression of the total radiation emitted per unit area of a gray body is presented.

$$E_{\varrho}(T) = \varepsilon_{\varrho} \sigma T^4$$
[3]

The parameter  $\varepsilon_g$  corresponds to the emissivity of a gray body. But, being more rigorous, real surfaces do not necessary behave as gray bodies, and their properties vary with the wavelength for a

<sup>&</sup>lt;sup>a</sup>Some authors refer to this parameter as 'emittance'. In this work emissivity and emittance are the same concept; however, there is a subtle difference between the two.<sup>16</sup>

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<sup>98</sup> given temperature. These surfaces radiate a different fraction  $\varepsilon_{\lambda}$  at <sup>99</sup> each wavelength; thus, the expression of the total radiation emitted <sup>100</sup> per unit area of a real surface is indicated in Equation 4. Note that the <sup>101</sup> parameter  $\varepsilon_r$  in Equation 4 is calculated by means of Equation 5; that <sup>102</sup> is, integrating  $\varepsilon_{\lambda}$  along all the radiation spectrum.

$$E_{real}(T) \cong \varepsilon_r \sigma T^4 \tag{4}$$

$$\varepsilon_r = \frac{\int_0^\infty \varepsilon_\lambda E_{b\lambda} \, d\lambda}{\int_0^\infty E_{b\lambda} \, d\lambda}$$
[5]

Real material properties.—As said previously, the radiative prop-104 erties of real materials are not necessarily those of gray bodies. The 105 106 difficulty is how to characterize the radiative properties of a selected material under working conditions. The reflectance  $(\rho_{\lambda})$  and transmit-107 tance  $(\tau_{\lambda})$  of real surfaces can be determined by means of the Fourier 108 transform infrared spectroscopy (FTIR);<sup>17</sup> thus, from Equation 1 ab-109 sorptivity  $(\alpha_{\lambda})$  can be obtained. But these measurements are typically 110 performed at room temperature; there are no overall techniques for 111 112 the measurement of radiative properties at high temperatures. It is true that for certain materials, in particular for some metals, it is accept-113 able to consider that their radiative properties remain constant with 114 temperature, although this cannot be easily generalized.<sup>16,18,19</sup> 115

Radiation heat loss model.—A radiation heat loss model for heat 116 loss calculations in a Siemens-type reactor was presented and de-117 scribed in detail in Ref. 20, and validated in Ref. 21. It is further 118 developed within the framework of this research to broaden its appli-119 cability and account for materials that do not behave as gray bodies. 120 One parameter needs to be defined for radiation heat loss calcu-121 lations: radiosity (J), the rate of outgoing radiant heat per unit area 122 from a surface. It is the sum of the directly emitted heat flux (E) and 123 the reflected incoming radiant heat flux from the surface (G). The 124 fraction of heat flux from one surface to another is determined by the 125 so-called configuration factor, or geometrical factor. The calculation 126 of the configuration factors  $(F_{i-i})$  is made using a geometric Hottel's 127 crossed-string method.<sup>22</sup> In the present case note that the rods and the 128 reactor wall have a cylindrical geometry. 129

If the material properties, the geometrical arrangement, the surface temperatures and the incoming and directly emitted radiant heat flux are known, the net heat flux exchanged (Q) in Watts from any surface ( $S_i$ ), is obtained from the difference between the radiosity and the incoming radiant heat flux. Then, the net radiation heat flux exchanged for a certain surface *i* can be expressed as shown in Equation 6.

$$Q_{i} = S_{i} \cdot (J_{i} - G_{i}) = S_{i} \cdot J_{i} - \sum_{j=1}^{n} S_{j} \cdot F_{j-i} \cdot J_{j}$$
[6]

For a Siemens reactor of n-1 rods, a n-equations system needs to 136 be solved, as the reactor wall is considered as an additional surface. 137 The radiosities of each surface (Ji) are the unknowns of the system. 138 The temperature of the rod surfaces and the reactor wall is known, as 139 is the corresponding surface emissivities. Once the  $J_i$  is obtained for 140 the *n* surfaces, the incoming radiant heat flux per unit area  $(G_i)$  is also 141 known. Thus, the net radiation heat exchanged by each surface  $(Q_i)$ 142 is obtained by substituting  $J_i$  and  $G_i$  in Equation 6. 143

To account for emissivity variations with the wavelength, radiation heat loss is obtained by means of Equations 7, 8, 9 and 10, which are solved independently for each wavelength

$$S_{i} \cdot \frac{1}{1 - \varepsilon_{i}(\lambda)} \cdot J_{i}(\lambda) - \sum_{j=1}^{n} S_{i} \cdot F_{i-j} \cdot J_{j}(\lambda) = S_{i} \cdot \frac{\varepsilon_{i}(\lambda)}{1 - \varepsilon_{i}(\lambda)} \cdot \sigma \cdot T_{i}^{4}$$
[7]

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$$E_i(\lambda) = \varepsilon_i(\lambda) \cdot \sigma \cdot T_i^4$$
[8]

$$G_i(\lambda) = \frac{1}{1 - \varepsilon_i(\lambda)} \cdot (J_i(\lambda) - E_i(\lambda))$$
[9]

$$Q_i(\lambda) = S_i \cdot (J_i(\lambda) - G_i(\lambda))$$
[10]

where i = 1, ..., n.

The net heat flux exchanged  $(Q_i)$  in Watts by any surface  $(S_i)$ , 151 is obtained by integrating  $Q_i(\lambda)$  along all the radiation spectrum. In Equation 11 the net heat flux exchanged by a surface is presented; 153  $E_b(\lambda)$  is the total radiation emitted per unit area of a black body indicated in Equation 2.

$$Q_i = \frac{\int_0^\infty Q_i(\lambda) E_b(\lambda) \, d\lambda}{\int_0^\infty E_b(\lambda) \, d\lambda}$$
[11]

This radiative model allows extra surfaces in the Siemens reactor to be considered and their positive or negative effect on heat savings studied. This can be the case of a thermal shield. A thermal shield is a cylinder surrounding the polysilicon rods and placed between them and the reactor wall. The presence of this shield may block a significant part of the radiated heat that otherwise would be lost through the reactor wall.

Now, the net heat flux exchanged  $(Q_i)$  in Watts by any surface ( $S_i$ ), is again obtained by integrating  $Q_i(\lambda)$  along all the radiation spectrum; but by replacing Equation 7 with Equations 12–15 (where i = 1, ..., m - 1, and m is the number of thermal shields considered).

$$S_{i} \cdot \frac{1}{1 - \varepsilon_{i}(\lambda)} \cdot J_{i}(\lambda) - \sum_{j=1}^{m} S_{i} \cdot F_{i-j} \cdot J_{j}(\lambda) = S_{i} \cdot \frac{\varepsilon_{i}(\lambda)}{1 - \varepsilon_{i}(\lambda)} \cdot \sigma \cdot T_{i}^{4}$$
[12]

$$S_m \cdot \frac{1}{1 - \varepsilon_m(\lambda)} \cdot J_m(\lambda) - \sum_{j=1}^m S_m \cdot F_{m-j} \cdot J_j(\lambda) = S_m \cdot \frac{\varepsilon_m(\lambda)}{1 - \varepsilon_m(\lambda)} \cdot \sigma \cdot T_m^4$$
[13]

$$\left(S_m \cdot \frac{\varepsilon_s(\lambda)}{1 - \varepsilon_s(\lambda)} + \frac{1}{\gamma(\lambda)}\right) \cdot \sigma T_m^4 - S_m \cdot \frac{\varepsilon_s(\lambda)}{1 - \varepsilon_s(\lambda)} \cdot J_m(\lambda) = \frac{\sigma T_n^4}{\gamma(\lambda)}$$
[14]

$$\gamma(\lambda) = \frac{1}{S_m \cdot \varepsilon_s(\lambda)} + \frac{1}{S_n} \cdot \left(\frac{1}{\varepsilon_n(\lambda)} - 1\right) + \left(\frac{2}{\varepsilon_s(\lambda)} - 1\right) \cdot \sum_{i=m+1}^{n-1} \frac{1}{S_i}$$
[15]

Note that even if the emissivity values now considered may be 171 wavelength dependent, materials still are considered opaque ( $\tau = 0$ ). 172

*Theoretical calculations.*—The potential of different thermal shields for radiation heat savings in an industrial Siemens reactor is studied here. The equations presented above are applied to a 36-rod, state-of-the-art Siemens reactor, and as a first approach, the emissivity of the materials is considered constant and wavelength independent. The initial and final diameter of the polysilicon rods is 0.7 and 13 cm, respectively, and their length is 2 m.

In Figure 1 the heat loss due to radiation in Watts (W) throughout 180 a polysilicon deposition process for a constant surface temperature of 181 1150°C is shown; the curves correspond to the case with no thermal 182 shield and four cases with thermal shields. The emissivities of the 183 thermal shields are 0.3, 0.45, 0.55 and 0.7. In Table I the theoretical 184 radiation heat loss savings for the aforementioned thermal shields are 185 presented. The radiation heat loss savings, compared to the heat loss 186 if no thermal shield is considered, are 65.8, 52.6, 44.3 and 30.5% for 187 thermal shield emissivities ( $\epsilon$ ) of 0.3, 0.45, 0.55 and 0.7, respectively. 188 This means, assuming that the radiation heat loss is responsible for 189 around 50% of the total power consumption, that with a thermal shield 190 with  $\varepsilon = 0.3$  a reduction in power consumption of 32.9% is obtained, 191 while for  $\varepsilon = 0.7$  the reduction would be of 15.5%. 192

The temperature reached by the different thermal shields depending on their emissivity is presented in Figure 2. In all cases, and from the beginning of the process, these temperatures are above 850°C, which will result in polysilicon deposition on these surfaces. Thus, 196

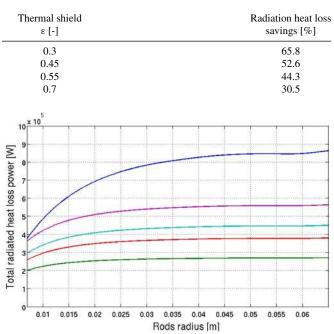


Table I. Theoretical radiation heat loss savings for different

thermal shields.

**Figure 1.** Radiation heat loss for a 36-rod Siemens reactor considering different thermal shields. No thermal shield (blue),  $\varepsilon = 0.7$  (purple),  $\varepsilon = 0.55$  (cyan),  $\varepsilon = 0.45$  (red),  $\varepsilon = 0.3$  (green).

<sup>197</sup> after a few minutes into the deposition process the thermal shield's
<sup>198</sup> surface emissivity will be 0.7, that of silicon at high temperatures.<sup>23</sup>
<sup>199</sup> Furthermore, contamination issues can arise unless the shields are of
<sup>200</sup> a highly pure material. One way to overcome these drawbacks would
<sup>201</sup> be to use a thermal shield made of purified silicon.<sup>12</sup> Not only will it
<sup>202</sup> avoid contamination, but one can also collect the silicon deposited on
<sup>203</sup> the shields, adding it to the silicon produced in a batch.

The potential of thermal shields can be compared to the use of a 204 polished or a reflective-coated inner wall of a reactor, which will lower 205 the wall emissivity. For a given growth rate, and knowing the power 206 consumption throughout a deposition process, and the initial and the 207 final diameters of the polysilicon rods, the energy consumption in 208 kWh/kg can be calculated. In Figure 3 the kWh/kg ratio for the case 209 of a reflective-coated wall is compared to those considering a silicon 210 211 thermal shield, no thermal shield and a thermal shield of  $\varepsilon = 0.3$ . For the calculations in Figure 3 the emissivity of the wall and the thermal 212

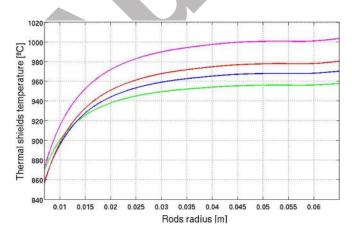
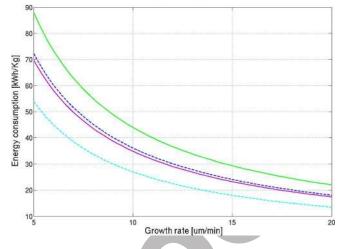


Figure 2. The temperature of thermal shields depending on their emissivity ( $\epsilon$ ) throughout a deposition process. Thermal shield emissivities:  $\epsilon = 0.7$  (purple),  $\epsilon = 0.55$  (red),  $\epsilon = 0.45$  (blue),  $\epsilon = 0.3$  (green).



**Figure 3.** Total power consumption of a 36-rod Siemens reactor for different growth rates considering: no thermal shield  $\epsilon_{wall} = 0.5$ - (green), silicon thermal shield  $\epsilon = 0.7$ - (purple), thermal shield with  $\epsilon = 0.3$  (cyan) and no thermal shield and polished reactor wall  $\epsilon_{wall} = 0.3$ - (blue).

shields is considered constant throughout a deposition process; and the 213 radiation heat loss is 50% of the total power consumption. The lowest 214 kWh/kg ratio is obtained for a low emissivity thermal shield, and the 215 kWh/kg ratio for that with a silicon thermal shield and a polished 216 inner wall are quite close. However, note that the low emissivity 217 thermal shield and the polished walls will not maintain their initial 218 emissivities for more than a short period of time, as silicon or a silane-219 based compound will deposit. After a few minutes into the deposition 220 process the blue curve will start to move slowly upwards until it 221 reaches the green curve; and the cyan curve will quickly move to 222 behave like the purple curve. Thus, the effect of a thermal shield is 223 more efficient in terms of energy savings than considering a polished 224 reactor wall; this statement is true even when considering a high initial 225 emissivity value for the thermal shield (e.g.,  $\varepsilon = 0.7$ ). 226

### Laboratory Scale Experiments

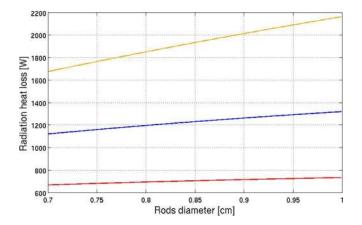
A number of experiments considering thermal shields are conducted in a laboratory Siemens reactor,<sup>24</sup> and the effect on radiation heat savings obtained is put forward.

Since the temperature of the thermal shield in the laboratory reac-231 tor will be lower than in the industrial case, the laboratory prototype 232 allows us to test the effect of thermal shields with different emissivi-233 ties. The key parameter for the selection of the thermal shield material 234 is the emissivity ( $\varepsilon$ ); but also, the material selected must be easily 235 machinable, and available with the geometries and thickness required 236 for its assembly inside the reactor chamber, so its mechanical strength 237 must be assured. The following materials are evaluated: molybdenum, 238 boron nitride, stainless steel, aluminum oxide (alumina), zirconium, 239 graphite foil and silicon. Some of the relevant properties of these 240 materials are presented in Table II; the values shown are considered 241 wavelength independent since this dependence is unknown. 242

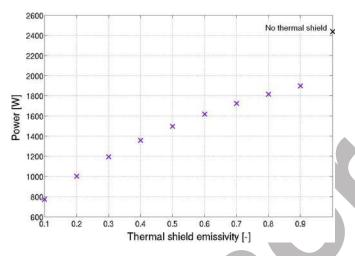
# Table II. Properties of different materials considered for the thermal shields. $^{25,26}$

Material	$\epsilon [-] (T = 25^{\circ}C)$	$\epsilon \ [-] \ (T \sim 600^\circ C)$	Ease of machining
Molybdenum	-	0.8-0.9	Medium
Stainless steel	0.6-0.8	0.7-0.9	Low
Alumina	-	0.3-0.4	Medium
Boron nitride	0.9-0.95	-	Medium
Zirconium	-	0.1-0.3	High
Graphite foil	0.7-0.9	0.4-0.6	Low
Silicon	-	0.7	Medium

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**Figure 4.** Radiation heat loss in the laboratory Siemens reactor for a 7-rod configuration considering a silicon thermal shield (blue), a low emissivity thermal shield  $-\epsilon = 0.3$ - (red) and without thermal shield (orange).



**Figure 5.** Laboratory Siemens reactor power consumption (P) predicted in theory for different thermal shield emissivities and for the case of no thermal shield considering a 7-rod configuration.

First, the radiation heat loss equations with thermal shields are applied to the laboratory Siemens reactor. The radiation heat loss for a 7-rod configuration with a low emissivity shield, with a silicon thermal shield and without thermal shield is presented in Figure 4; it can be seen how the lowest radiation heat losses are obtained for a low emissivity thermal shield. The temperatures reached by the thermal shields are in the range of 600-750°C.

The power consumption predicted by the model for different thermal shield emissivities and for that of no thermal shield, are presented in Figure 5. For these calculations a constant deposition temperature of 1100°C, the same growth rate and the same duration of the deposition processes, is considered, thus averaging the measured data. It can be seen that the lower emissivity of the thermal shield, the higher radiation heat savings.

**Experiments with thermal shields.**—A 7-rod configuration is chosen as a compromise solution between a dense compactness - a large number of rods - and the size of the reactor chamber. The length of the rods is 10 cm and their initial diameter is around 0.7 cm.

From the thermal shield materials listed in Table II, the following have been selected for testing: silicon, alumina and stainless steel. Different thickness of the selected materials are considered, and in some cases the outer surface of the thermal shields is silver coated<sup>b</sup>.

 Table III. Experiments conducted with 7-rod configuration in the laboratory Siemens reactor.

Experiment name	Description		
No shield (No)	Without any thermal shield		
Silicon shield (Si1)	Multi-crystalline silicon thermal		
	shield (290 µm/layer; 3 layers)		
Silicon shield (Si2)	Mono + Multi-crystalline silicon thermal		
	shield (400 + 290 µm; 1 + 1 layers)		
Silicon shield (Si3)	Mono + Multi-crystalline silicon therma		
	shield $(400 + 290 \times 4 \mu\text{m}; 1 + 4 \text{layers})$		
Alumina shield (Alu1)	Alumina shield		
	(1 mm thick)		
Alumina shield (Alu2)	Alumina shield silver coated		
	(1 mm thick)		
Steel shield (Ste)	Stainless steel shield		
	(1 mm thick)		

 Table IV. Experimental data obtained for the 7-rod configuration experiments: 'silicon shields'.

Experiment	(No)	(Si1)	(Si2)	(Si3)
T <sub>deposition</sub> [°C]	1106	1101	1108	1108
Si deposited [gr]	50.7	61.9	59.3	59.8
Poweraverage	2343	1979	2042	2122
Time [min]	392	406	385	375
$\mathbf{T}_{wall} \ [^{\circ}\mathbf{C}]$	280	233	184	181
T <sub>shield</sub> [°C]	-	678	641	616
Growth rate [µm/min]	2.9	3.5	3.6	3.6
Consumption [kWh/kg]	311	216	221	222
Energy savings [%]	-	28.4	26.8	26.5

The emissivity of silver is very low ( $\varepsilon \sim 0.02$ -0.05), so if this coating withstands the process temperatures, it will act as a mirror making a non-opaque body behave almost as if it were. 267

The relevant data related to these experiments is presented in the following tables. First, the different thermal shields are described and related to their corresponding label in Table III. Then, the experimental results are grouped together in 'silicon shields' and 'alumina and stainless steel shields'; Tables IV and V, respectively. 270

From the data presented in Table IV, the energy savings obtained 273 with the different silicon thermal shields are similar. The reduction 274 in the kWh/kg ratio obtained considering thermal shields related to 275 experiment (No) are between 26.5 and 28.4%. All these experiments 276 were conducted under similar conditions and their duration is similar. 277 Despite the fact that the deposition surface temperature is in all cases 278 around 1100°C, there is a difference in the growth rate obtained in 279 experiment (No). This is so because the presence of a thermal shield 280 changes the distribution of the gas temperature, and higher tempera-281 tures are achieved in the gas surrounding the silicon rods. 282

From the data presented in Table V, the energy savings in kWh/kg, 283 compared with experiment (No), vary between 15.1 and 30.7%. The 284

# Table V. Experimental data obtained for the 7-rod configuration experiments: 'alumina and stainless steel shields'.

Experiment	(No)	(Alu1)	(Alu2)	(Ste)
T <sub>deposition</sub> [°C]	1106	1107	1108	1098
Si deposited [gr]	50.7	65.3	53.7	49.3
Poweraverage	2343	2333	1669	1915
Time [min]	392	430	404	388
$T_{wall} [^{\circ}C]$	280	142	175	152
$T_{shield}$ [°C]	-	736	705	570
Growth rate [µm/min]	2.9	3.6	3.2	3.0
Consumption [kWh/kg]	311	256	205	251
Energy savings [%]	-	15.1	30.7	16.8

kWh/kg values in the laboratory scale reactor are several times higher than those found in industrial processes, mainly since the process pressure is 6-7 times lower. Comparing experiments (Alu1) and (Alu2), the silver coating seems to be effective; however its behavior differs from that expected from its theoretical ε (further explanations will be presented in Discussion on energy savings section).

Lastly, in experiments conducted with silicon thermal shields etch-291 ing is detected on the surface of the shields. This is attributed to the 292 presence of SiCl<sub>4</sub> as a by-product of the reduction reaction. The oc-293 currence of this phenomenon versus polysilicon deposition depends 294 on the mol fraction of SiCl<sub>4</sub>, which will depend on the deposition sur-295 face temperature.<sup>27,28</sup> High SiCl<sub>4</sub> concentrations and low temperatures 296 favor the etching. However, as already explained, under industrial de-297 position conditions the temperature of the thermal shields will be 298 such that polysilicon will be deposited on the thermal shields, and no 299 etching is expected. 300

**Discussion on energy savings.**—From the above, energy savings have been confirmed for the 7-rod configuration experiments considering different thermal shields.

If the experimental data from Tables IV and V (average power consumption and energy savings) is compared with the theoretical calculations for different thermal shields (Figure 5), a good agreement for the case of no thermal shields is obtained; differences between both values are under 3.8%. Note that our calculations consider constant deposition conditions, while the experimental conditions of the deposition process vary slightly from one experiment to another.

For the experiments conducted with thermal shields, the averaged 311 power consumption and energy savings obtained vary between 1667-312 2333 W and 15.1-30.7%, respectively. According to data presented in 313 Figure 5, the previous values correspond to thermal shield emissivities 314 above 0.6. In the case of the silicon thermal shields, the energy savings 315 obtained correspond to  $\varepsilon = 0.7$ –0.8, for the alumina shields to  $\varepsilon >$ 316 0.9, for the silver coated alumina shield to  $\varepsilon = 0.6-0.7$ ; and for the 317 stainless steel shield to  $\varepsilon > 0.9$ . These  $\varepsilon$  values do not correspond 318 to those found in the bibliography assuming the gray body approach, 319 which is no surprise since the gray body approach simplifies much of 320 the radiative behavior of real bodies. 321

Reflectance, transmittance and emissivity measurements.—With 322 the aim of clarifying the real emissivity of the thermal shield materials 323 tested in the laboratory Siemens reactor, reflectance ( $\rho$ ) and transmit-324 tance ( $\tau$ ) measurements for different  $\lambda$  are taken. Both,  $\rho(\lambda)$  and 325  $\tau(\lambda)$ , can be measured directionally or integrated; in the present case 326 integrated measurements are suitable since the materials considered 327 do not have specular surfaces. These measurements are conducted at 328 room temperature. 329

In Figure 6 the integrated transmittance measurements, within the wavelength range  $\lambda \in (2.5-20) \ \mu$ m, for different thermal shields are presented. In all cases, noticeably for the silicon shield, the integrated transmittance is  $\tau \neq 0$ . Measurements for a silicon, alumina and stainless steel thermal shields are presented in Figure 6. The integrated transmittance measured is on average 41.3, 8.1 and 0.5% for the

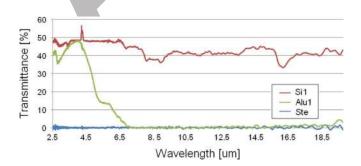
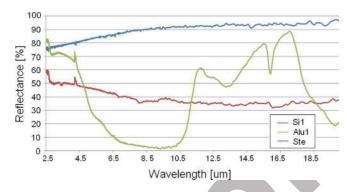


Figure 6. Integrated transmittance ( $\tau$ ) of: 290  $\mu$ m multi-crystalline silicon (red), 1 mm alumina (green) and 1 mm stainless steel (cyan).



**Figure 7.** Integrated reflectance ( $\rho$ ) of: 290  $\mu$ m multi-crystalline silicon (red), 1 mm alumina (green) and 1 mm stainless steel (cyan).

290 μm multi-crystalline, 1 mm alumina and 1 mm stainless steel 336 samples, respectively.

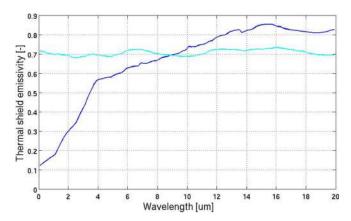
Integrated reflectance measurements are also conducted;  $\rho(\lambda)$  for  $\lambda \in (2.5-20) \,\mu\text{m}$  for silicon, alumina and stainless steel are presented in Figure 7. From Figure 7, the averaged reflectance of the silicon sample is 40%, while the respective values for that of alumina and the stainless steel samples are 48.1 and 92.8%, respectively. 339

From the average values of the aforementioned transmittance and 343 reflectance integrated measurements, only the stainless steel sample 344 presents a very low transmittance. Materials experimentally tested in 345 the laboratory Siemens reactor at room temperature definitely do not 346 behave as gray bodies, and similar behavior can be expected at higher 347 temperatures.<sup>18,19</sup> The latter explains the differences between the pre-348 dicted energy savings and the empirically obtained ones. The next 349 section discusses the effect that the wavelength-dependent emissivi-350 ties can have on the radiation heat losses. 351

### Discussion on the Contribution to the Radiation Heat Loss Model 352

The model for radiation heat loss is applied here for the radiation heat loss calculations of a 36-rod industrial Siemens reactor, considering thermal shields that do not behave as gray bodies. Two hypothetical thermal shields with an averaged  $\varepsilon(\lambda) = 0.7$  are considered, with an emissivity variation presented in Figure 8. It can be seen that  $\varepsilon(\lambda)$  of material 1 is approximately constant, while  $\varepsilon(\lambda)$  of material 2 is heavily dependent on the wavelength.

The radiation heat loss for  $\lambda \in (0.1, 20) \,\mu$ m, calculated for a 36-rod industrial Siemens reactor, is presented in Figure 9. The two scenarios presented; hereinafter scenarios 1 and 2, correspond to material 1 and material 2 thermal shields, respectively. In both cases, the radiation 363



**Figure 8.** Emissivity  $\varepsilon(\lambda)$  for two different thermal shield materials: material 1 (cyan) and material 2 (blue). In both cases, the averaged  $\varepsilon(\lambda) = 0.7$ .

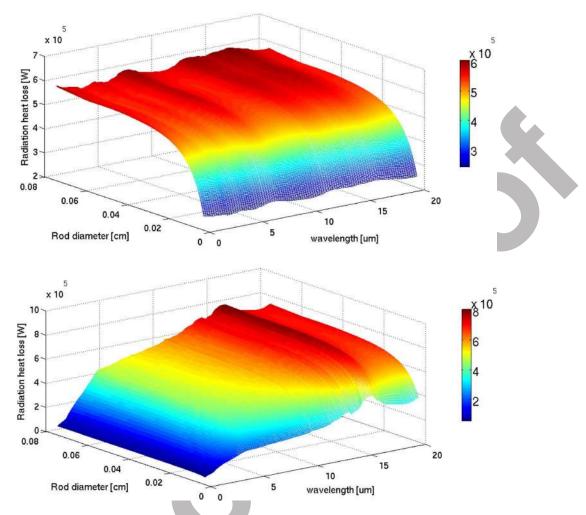


Figure 9. Radiation heat loss for different wavelengths for two different thermal shield materials: material 1 (top) and material 2 (bottom).

heat loss variation with  $\lambda$  is similar to the corresponding variation in 364  $\varepsilon(\lambda)$  of the shield material considered. 365

When the surfaces presented in Figure 9 are integrated along the 366 entire radiation spectrum, the radiation heat loss values presented in 367 Figure 10 for scenarios 1 and 2 are obtained. This curves are pre-368

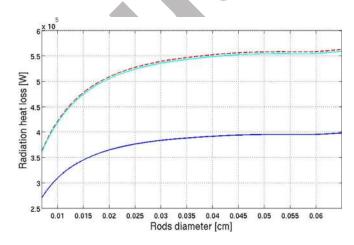


Figure 10. Radiation heat loss for different thermal shield materials: material 1 (cyan), material 2 (blue), obtained by integrating Figure 9 along all the radiation spectrum. The case of a material with a constant  $\varepsilon(\lambda) = 0.7$  is presented for comparison (red).

sented together with the corresponding curve if a constant emissivity 369 for the thermal shield  $\varepsilon(\lambda) = 0.7$  is considered - scenario 3. It can 370 be appreciated how the scenarios 1 and 3 are quite close, but great 371 differences in radiation heat loss are obtained between scenarios 1 and 372 3, and scenario 2; the averaged differences are above 25%. As regards 373 the thermal shield temperature, results obtained for scenarios 1 and 3 374 are also quite close; the temperature of the shields is around 870°C 375 at the beginning of the deposition process, increasing rapidly until it 376 reaches around 1000°C at the end of the process. In scenario 2, the 377 thermal shield temperature has a similar behavior with temperature 378 values that go from 860 to 975°C. 379

The aforementioned differences above are explained since not all 380 wavelengths contribute to the same extent to the radiation heat loss; 381 in particular for these three scenarios, the greatest contribution of  $\varepsilon(\lambda)$ 382 occurs in the range  $\lambda \in (1-6) \ \mu m$ . 383

These results highlight the importance of having reliable data on 384 the emissivities in the relevant range of wavelengths, and for the 385 application of silicon CVD, at deposition process temperatures, which 386 remains pending. 387

#### Conclusion

A radiation model for heat loss calculations in a Siemens-type reactor has been presented, in which the fraction of energy leaving a 390 certain surface that arrives at another surface is evaluated using the 391 geometric Hottel crossed-string method, and the effect of the emissiv-392 ity variation with the wavelength is taken into account. A significant 393

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potential for reducing radiation heat loss in Siemens reactors has been 394 identified, considering different thermal shields. The model shows 395 that materials with lower surface emissivities lead to higher radiation 396 heat loss savings. The effect of a thermal shield is also more efficient 397 in terms of energy savings than considering a polished reactor wall, 398 even for a thermal shield with a high initial emissivity value. 399

Experiments considering different thermal shields are conducted 400 in a laboratory Siemens reactor. It has been experimentally shown 401 that significant energy savings in the polysilicon deposition process 402 are obtained. 403

Silicon thermal shields have some advantages in terms of pre-404 venting contamination and collecting the silicon deposited on them, 405 and energy savings of between 26.5-28.5% have been experimentally 406 407 proven.

Reflectance and transmittance measurements as a function of 408 wavelength are taken for the materials tested, proving that they do 409 not behave as gray bodies at room temperature, and similar behavior 410 can be expected at higher temperatures. Results highlight the impor-411 tance of having reliable emissivity data on the materials involved at 412 deposition temperatures, which remains pending. 413

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