

# THz Backward-Wave Oscillators for Plasma Diagnostic in Nuclear Fusion

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**Abstract**—Understanding of the anomalous transport attributed to short-scale length microturbulence through collective scattering diagnostics is key to the development of nuclear fusion energy. Signals in the subterahertz (THz) range (0.1–0.8 THz) with adequate power are required to map wider wavenumber regions. The progress of a joint international effort devoted to the design and realization of novel backward-wave oscillators at 0.346 THz and above with output power in the 1 W range is reported herein. The novel sources possess desirable characteristics to replace the bulky, high maintenance, optically pumped far-infrared lasers so far utilized in this plasma collective scattering diagnostic. The formidable fabrication challenges are described. The future availability of the THz source here reported will have a significant impact in the field of THz applications both for scientific and industrial applications, to provide the output power at THz so far not available.

**Index Terms**—Backward-wave oscillator (BWO), double-corrugated waveguide (DCW), double-staggered grating (DSG), plasma diagnostic, terahertz (THz).

## I. INTRODUCTION

TERAHERTZ (THz) vacuum electron devices are gaining significant consideration when the generation of relatively high power in the frequency range below 1 THz is

needed [1]–[9]. In particular, the backward-wave oscillator (BWO) is an effective solution to produce relatively high power and stable monochromatic THz signals. BWOs can be electronically tuned over a wide frequency range around the operating frequency and have high stability in frequency (up to  $10^{-7}$  to  $10^{-8}$  by phase locking). However, the only available BWOs are based on old technologies. No compact, affordable, and long-life THz BWOs is currently available.

Recently, the introduction of new high aspect-ratio fabrication processes derived from the MEMS technologies as the lithography, electroplating, and molding (LIGA) [5], [7] and advanced mechanical microfabrication as nano-CNC milling [10] provides an accuracy at the submicrometer level, which satisfies the demanding specifications of interaction structures or slow-wave structures (SWSs) to support the THz operation frequencies. Furthermore, the progress on simulation tools based on accurate 3-D electromagnetic and particle-in-cell (PIC) solvers permits now a reliable prediction of the THz vacuum source performance, to ease the fabrication phase. The development of innovative cathode materials [10] is leading to a novel generation of electron guns, with high current density and long lifetime, fundamental for the overall device mean time between failures and high-frequency stability. Nevertheless, the main challenges are to achieve a surface roughness below the skin depth (66 nm at 1 THz) for minimizing the ohmic losses and the assembly and alignment of the beam with tolerance in the order of tens of micrometers.

Nuclear fusion is one of the fields, where the availability of THz BWOs will have a relevant impact on improving the understanding of a critical phenomenon as the anomalous transport of the plasma. It still remains a fundamental area of investigation, which is essential for the development of fusion energy. The measurement technique is based on the collective Thomson scattering at THz frequency [12]–[15]. The plasma is illuminated by a THz beam that is scattered by the charged particles. The scattered beams due to the electron density fluctuations are detected by a receiver array and thereby map out the location, wavenumber spectrum, and strength of the turbulence.

In the recent upgrade of the high- $k$  scattering system at NSTX experiment at Princeton, the wavenumber  $k_{\theta}$  coverages, and it has been increased to target electron temperature gradient modes by increasing the probe frequency from 0.280 to 0.693 THz. The availability of solid-state sources is limited to about 0.3 THz and 30 mW, while about 100 mW are needed at 0.693 THz to excite a detectable scattering.

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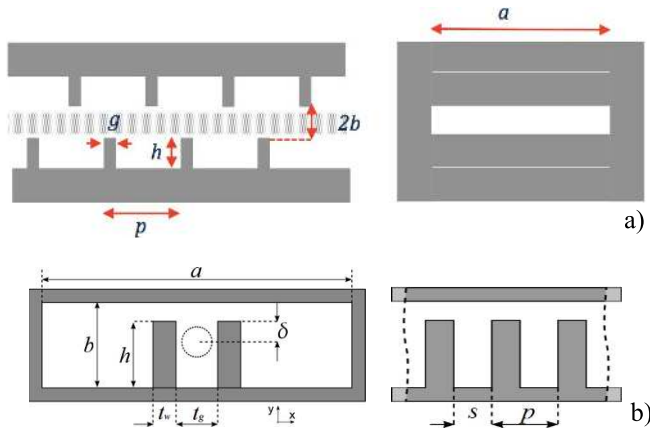


Fig. 1. SWS for 0.346-THz operation. (a) DSG. (b) DCW. The gray area in (a) and the dashed circle in (b) represent the sheet and cylindrical beam for the DSG and the DCW, respectively.

proposed for operation at THz frequencies and overcome the fabrication issues of the conventional structures.

The DSG is conceived to support an elliptical sheet electron beam. The advantage of the sheet electron beam is the large cross section that permits to deliver a high beam current using a relatively beam current density. The sheet beam requires a careful design of the magnetic focusing system, but it is very promising to realize high-power vacuum electron devices by using low cathode loading guns.

The DCW is conceived to support a cylindrical electron beam. The DCW is of easy fabrication and assembly. The advantage of a cylindrical electron beam is that it is generated by the well-established Pierce gun technology and focused by the use of a conventional magnetic focusing system.

Both the SWSs are very promising to realize THz BWOs with a wide range of characteristics in terms of fabrication, output power, and cost.

### III. BWO DESIGN

The approach of using two different SWSs, the DSG and the DCW, to design a family of THz BWO is a breakthrough for tailored power generation at THz frequencies. Two BWOs based on the DSG and the DCW will be the first two devices for a new family of THz sources to cover a wide range of applications.

The main design targets are given in the following:

- 1) low cost;
- 2) easy assembly for high yield;
- 3) compact dimensions (200–300 cm<sup>3</sup>);
- 4) wide range of performance to potentially cover the sub-THz spectrum (0.1–1 THz);
- 5) tunable (at least 5%);
- 6) low beam voltage and compact power supply.

A 0.346-THz operating frequency is considered in the following for application in the plasma diagnostic in nuclear fusion, as described in Section I. The first design parameter defined was the beam voltage that determines the length of the period of the SWSs. A low beam voltage in the range of 12–18 kV favors to use of a compact and low-cost power supply. The resulting period was estimated to be in the range of the fabrication process.

Due to the different structures, different beam voltages were adopted. The DSG was designed to operate with 17-kV beam voltage, while the DCW was to support about 13 kV. The dimensions of the two SWSs are shown in Table I. It is notable that a period shorter than 200  $\mu\text{m}$  is required. In Fig. 2, the dispersion curve of the DCW with superimposed beam line at 12.8 kV is shown. The interaction impedance is typically low in the backward-wave mode.

1) *Couplers*: A detailed study based on 3-D electromagnetic simulation [18] on the coupler to transform the mode in the SWS in the TE<sub>10</sub> at the flange at the output port was carried out to maximize power transfer. The conductivity of copper considered in simulation is  $\sigma_{\text{cu}} = 3.9 \times 10^7$  S/m [7]. The coupler is a three-port network; one port is connected to the SWS, a second port is the beam tunnel to connect the gun, and the third port is the output port connected to the flange.

## II. TERAHERTZ SLOW-WAVE STRUCTURES

At microwave frequencies, helices are the most common SWSs, but as the frequency increases toward the millimeter-wave range, their dimensions are too small for fabrication, and new geometries must be adopted. The simple structure of the rectangular corrugated waveguide inspired different structures that can be realized by the available fabrication processes with the dimensions to support THz frequencies.

In particular, the double-staggered grating (DSG) [17] [Fig. 1(a)] and the double-corrugated waveguide (DCW) [18] [Fig. 1(b)] are the two SWSs that have been successfully

Presently, the only THz source to deliver  $\sim 100$  mW at 0.693 THz needed to provide the minimum scattered signal level to be detected by the receiver array is a bulky optically pumped far infrared. A second high-frequency source is needed to provide local oscillator (LO) power for the receiver array. The use of a second laser is not feasible. It has been chosen to use an array based on sensitive room temperature subharmonic mixers working roughly at half of the illumination source frequency, namely, 0.346 THz. It is required 3 to 5 mW of LO power per mixer.

This paper describes an international joint effort of three leading institutions in China, the U.K., and the U.S. to design and construct a novel family of BWOs, operating at the frequency of 0.346 THz, to satisfy the quest for LO power for the matrix array for the NSTX-U plasma diagnostic and other future plasma diagnostic systems [16]. The design target is to achieve an output power in the range of hundreds of milliwatt by lightweight, compact, affordable, low-operating cost BWOs to enable a wide matrix of receivers.

This paper is organized as follows. The properties of two different SWSs suitable for THz BWO fabrication are reported in Section II. Section III describes the design aspects and the cold parameters. Section IV details the hot simulations and performance of the two BWOs. Challenges involved with microfabrication technologies of the proposed BWOs are discussed in Section V. Section VI reports on the gun and the window.

TABLE I

DSG		DCW	
Parameter	$\mu\text{m}$	Parameter	$\mu\text{m}$
$a$	483	$a$	1500
$b$	45	$b$	200
$h$	170	$h$	140
$p$	185	$p$	140
$g$	44	$s$	80
		$t_g$	120
		$t_w$	60

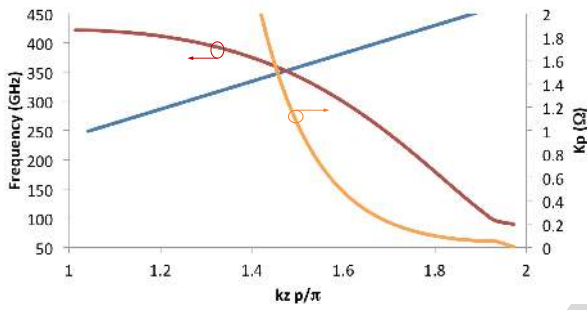


Fig. 2. DCW dispersion curve (brown curve), interaction impedance (orange curve), and beam line at 12.8 kV (blue line).

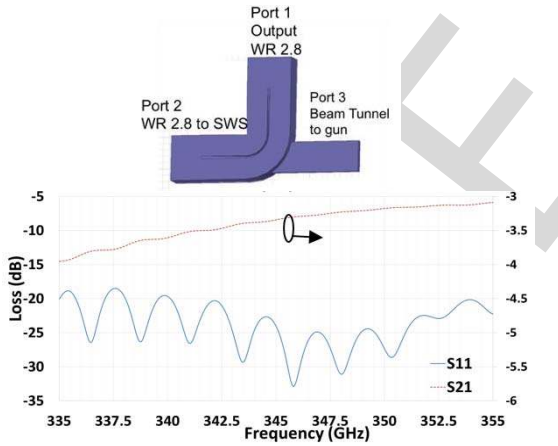


Fig. 3. DSG coupler S-parameters.

165 The coupler for the DSG is particularly challenging due to  
 166 the wide beam tunnel needed for the sheet beam. Having a  
 167 low cutoff frequency in the same range of the SWS, a ridge  
 168 is added to the bend (Fig. 3) to perturb the matching between  
 169 the SWS and the beam tunnel. The resulting  $S_{11}$  is better than  
 170  $-25$  dB over a wide frequency range around 0.346 THz.

171 A study of the coupler for the DCW was performed by  
 172 considering a back-to-back structure. First, a simple structure  
 173 including a tapered transition between a waveguide with  
 174 the same cross section of the DWG and the flanges is  
 175 designed, as shown in Fig. 4(a), to evaluate the effect of the

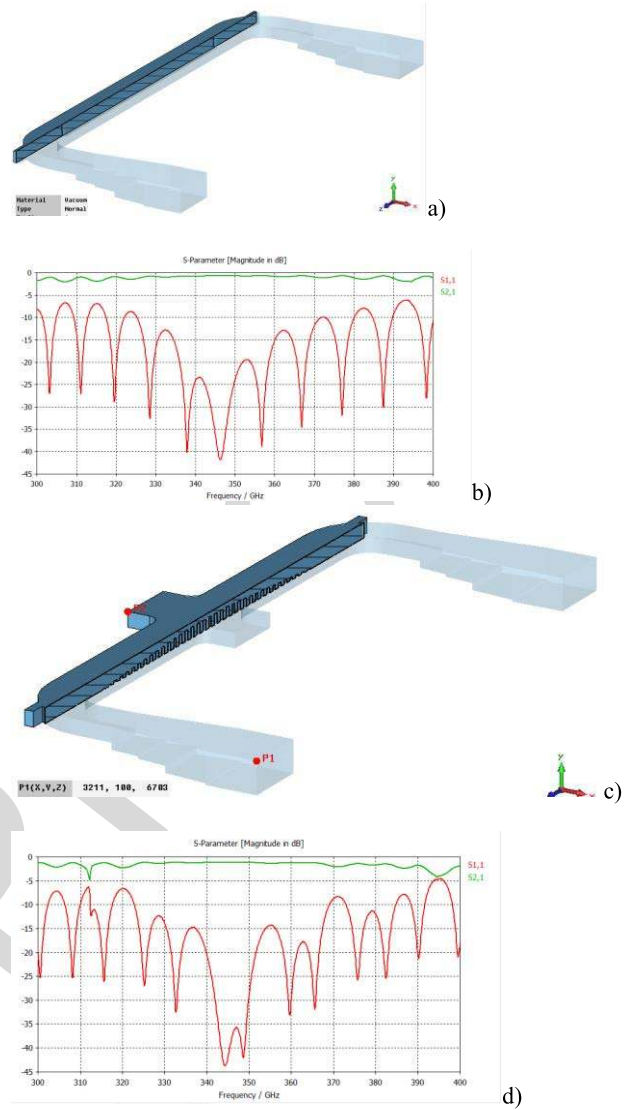


Fig. 4. DCW coupler. (a) Waveguide without DCW. (b) S-parameter structure. (c) Waveguide with DCW. (d) S-parameters.

waveguide tapering. Fig. 4(b) shows the obtained  $S_{11}$  better than  $-25$  dB in the operation range. Next, a second structure with similar topology, including three sections of pillars (two tapered sections and one short section with a nominal height), is designed for the best matching, as shown in Fig. 4(c). Results show that  $S_{11}$  in this case is better than  $-35$  dB in the region around the operating frequency [Fig. 4(d)].

Both the couplers' performance ensures the efficient propagation of the RF signal from the interaction structure to the flanges.

#### IV. LARGE SIGNAL SIMULATIONS

The design of the two BWOs is based on the definition of the critical length for oscillations to set a proper number of periods within the SWS. Results from this optimization process are shown in Table II.

Next, the 3-D PIC simulations performed to evaluate the electrical behavior of the BWOs. The DSG BWO supports an



TABLE II

BWO specifications	DSG	DCW
Beam Voltage	17.1 kV	12.8 kV
Beam current	14 mA	10 mA
Beam channel	483 x 90 $\mu\text{m}$	120 $\mu\text{m}$
Beam Aspect Ratio	5.4 : 1	Round
Beam Current Density	160 A/cm <sup>2</sup>	127 A/cm <sup>2</sup>
Magnetic field	0.35 T	0.5 T
No. Periods	65	116
Total length	~ 15 mm	~ 20 mm

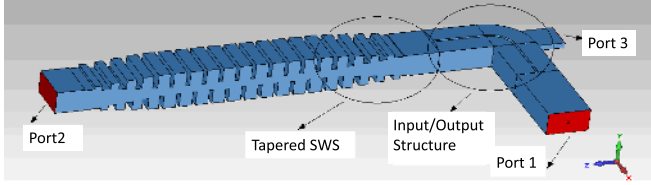


Fig. 5. PIC simulation setup for the DSG BWO.

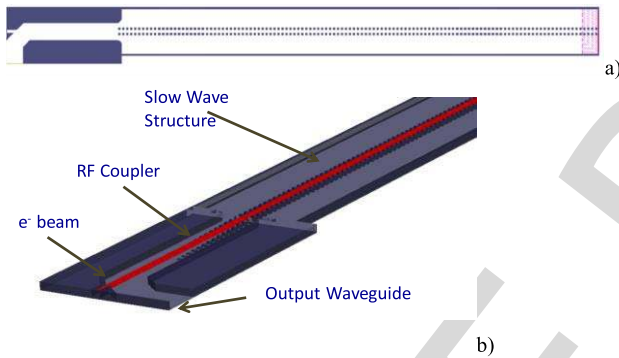


Fig. 6. PIC simulation setup for the DCW PIC simulations. (a) Top view. (b) Coupler detail.

193 elliptical electron beam with the cross section of  $400 \times 50 \mu\text{m}^2$   
 194 and the current of 14 mA (the aspect ratio of 5.4:1). The DCW  
 195 BWO supports a cylindrical electron beam of  $50\text{-}\mu\text{m}$  radius  
 196 and 10-mA current. Due to the different beam parameters  
 197 used, the overall device performance is different for the two  
 198 BWOs and should be evaluated in the context of the different  
 199 technological challenges required and the different application  
 200 needs. The DSG BWO is modeled by CST Particle Studio [19]  
 201 (simulation setup in Fig. 5) and the DCW BWO is modeled  
 202 by MAGIC3D [20] (simulation setup in Fig. 6).

203 The results of output power for the two devices are shown  
 204 in Fig. 7. The DSG BWO [Fig. 7(a)] provides about 1 W  
 205 and the DCW BWO about 0.4 W [Fig. 7(b)]. The electron  
 206 energy distribution along the longitudinal direction for both  
 207 the BWOs is shown in Fig. 8. The spectral response at the  
 208 output port of the DCW BWO is shown in Fig. 9, showing the  
 209 highly monochromatic generation of signal at the frequency of  
 210 interest.

211 In order to demonstrate the tunability of the BWO designs,  
 212 the tuning range for the DSG and the DCW is shown  
 213 in Fig. 10(a) and (b), respectively. It can be noted that a

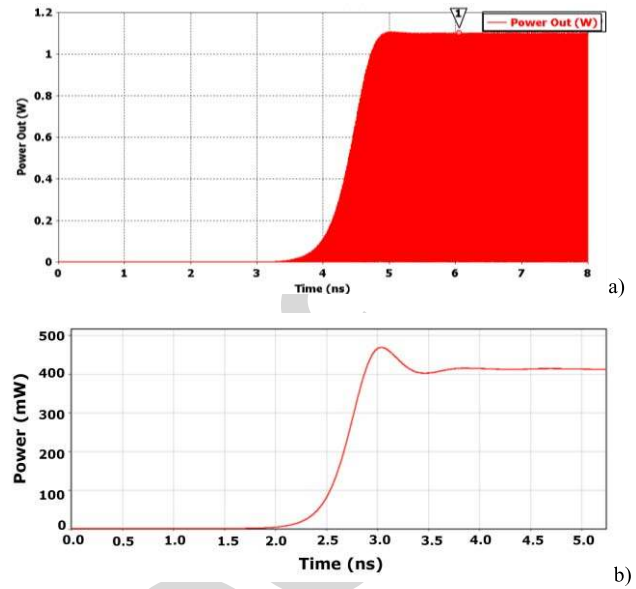


Fig. 7. Average output power for (a) DSG and (b) DCW BWOs.

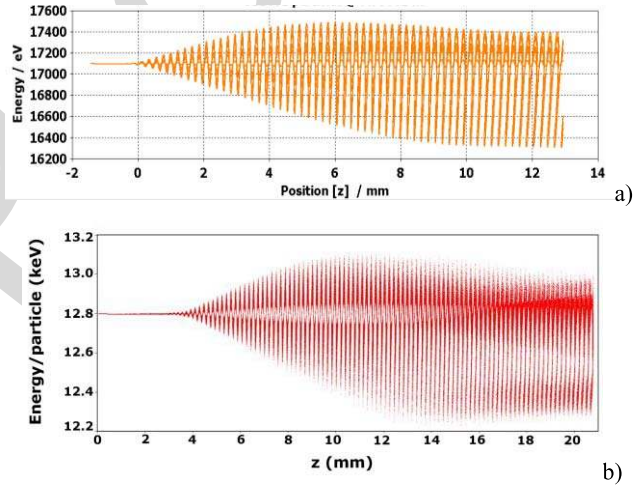


Fig. 8. Electron energy. (a) DSG BWO. (b) DCW BWO.

214 variation of beam voltage in the range of 15.8–17.8 kV permits  
 215 a variation in frequency of 12 GHz for the DSG and that the  
 216 same relative change in the nominal beam voltage allows a  
 217 tuning of 14 GHz for the DCW.

218 The BWO performances so far presented are at the state of  
 219 the art. The high power level and tuning features, not achiev-  
 220 able by any other technology today, represent a breakthrough  
 221 in the field.

## V. BWO MICROFABRICATION

222 The main challenge in the THz frequency range is the fab-  
 223 rication of SWSs with the expected electromagnetic behavior  
 224 while establishing a reliable and repeatable process. For the  
 225 DSG circuit, vane height is the most sensitive dimension which  
 226 determines the bandwidth of the device. The period of the  
 227 structure affects the central operating frequency, whereas the  
 228 width of the DSG controls the dispersion curve. The DCW  
 229 structure is more sensitive to the  $h$  and  $p$  values driving the  
 230

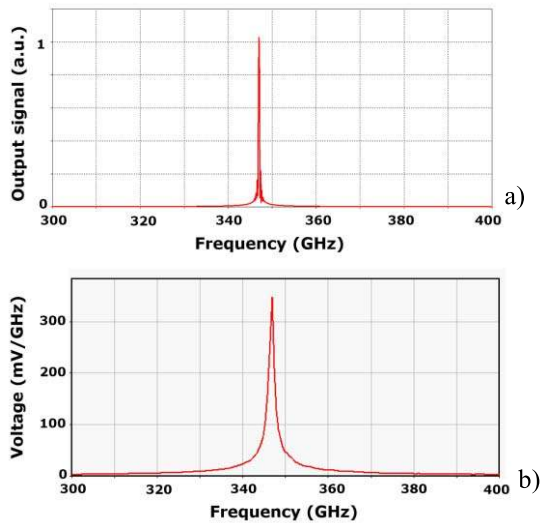


Fig. 9. Spectrum of (a) DSG and (b) DCW BWOs.

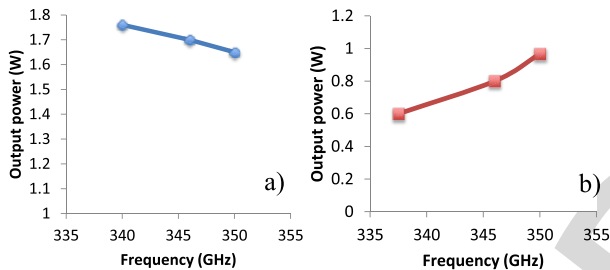


Fig. 10. Instantaneous power and tuning range of (a) DSG and (b) DCW BWOs for the beam voltage 15.8–17.8 and 11.75–13.3 kV, respectively.

231 dispersion curve. Machining tolerances are expected to be  
 232  $\pm 1 \mu\text{m}$ , which is sufficient to achieve the desired performance.

233 Photolithographic techniques, such as UV-LIGA, are  
 234 demonstrated suitable for the dimension accuracy required  
 235 for the two SWSs considered. However, especially for a  
 236 small number of pieces, the fabrication of the mold and  
 237 the electroforming process is not convenient. Furthermore,  
 238 a relevant effort to achieve a level of surface roughness better  
 239 than the skin depth (about 110 nm at 0.346 THz) to reduce  
 240 ohmic losses is required. CNC milling offers high flexibility  
 241 and possibility of patterning the third dimension. The state-  
 242 of-the-art prototype nano-CNC milling machine, developed  
 243 by DTL, a subsidiary of DMG-Mori-Seki, permits one to  
 244 achieve performance at the state of the art, with reduced  
 245 cost and high repeatability for dimensions suitable for THz  
 246 regime structures [10]. The high accuracy of the nanomilling  
 247 machine was proved to obtain levels of surface roughness  
 248 down to 40 nm, well below the skin depth at 0.346 THz. The  
 249 NN1000 nano/micromilling machine has a maximum spindle  
 250 speed of 50 000 r/min; the chip load is kept below 0.001-mm  
 251 feed per tool flute rpm. The proper setting of the machining  
 252 parameters is fundamental in achieving excellent surface finish  
 253 and tool lifetime.

254 In the case of the DSG and the DGW at 0.346 THz,  
 255 the dimensions shown in Table I represent a formidable  
 256 fabrication challenge. A test of feasibility for the fabrication  
 257 was performed realizing the DSG and the DCW in aluminum

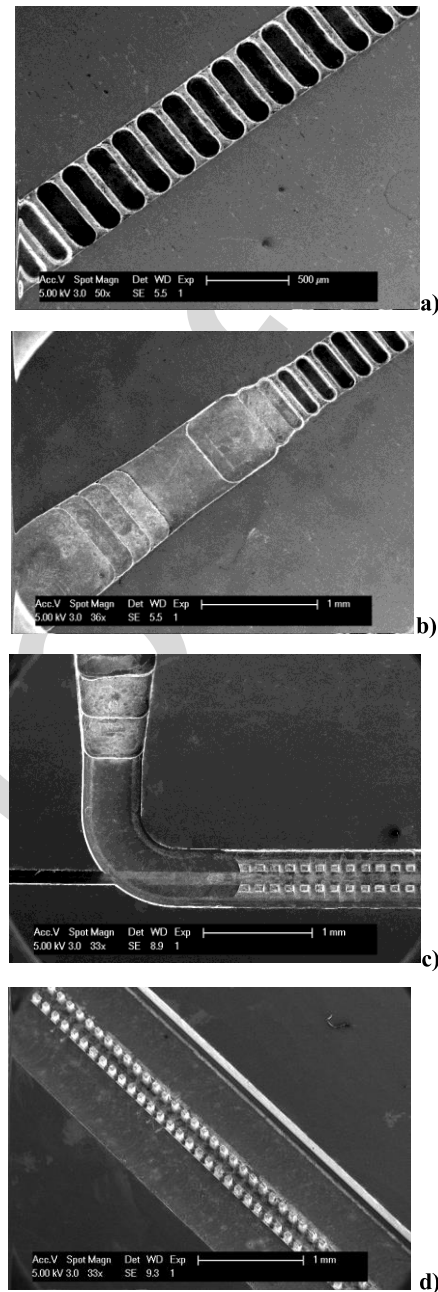


Fig. 11. SEM images of the SWSs fabricated by nano-CNC milling. (a) and (b) DSG. (c) and (d) DCW.

258 with the dimensions shown in Table I. Four different SEM  
 259 views of the DSG and the DCW realized by nano-CNC  
 260 milling are shown in Fig. 11. The high level of accuracy  
 261 for the very small dimensions is readily observed. The high  
 262 definition of the pillars is notable. Due to the characteristics  
 263 of aluminum, the surface roughness achieved was higher than  
 264 expected. The fabricated samples were primarily built to test  
 265 the microfabrication process in terms of dimensions achieved.  
 266 However, the measurements of the S-parameters were carried  
 267 out. The setup for the DSG measurement is shown in Fig. 12.  
 268 It consists of two halves assembled together by a system of  
 269 alignment pins. The setup for the DCW is similar to a lid  
 270 to close the waveguide that does not require a very accurate  
 271 alignment procedure.

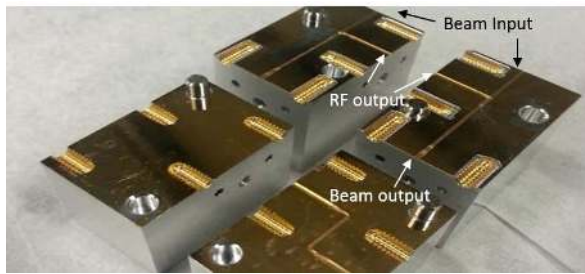


Fig. 12. Full DSG assembly with alignment pins and flanges.

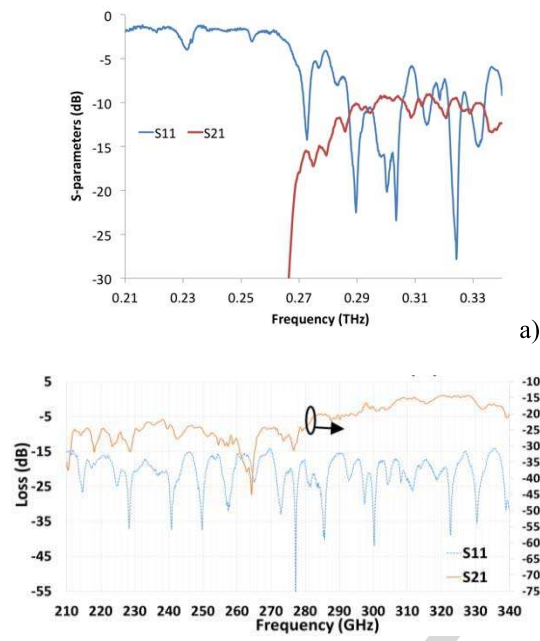


Fig. 13. Measurements of the S-parameters of (a) DCW and (b) DSG in the low range of the band.

272 The S-parameters of the fabricated structures were measured. The  $S_{11}$  and  $S_{21}$  for the fabricated DSG and DCW are shown in Fig. 13(a) and (b), respectively. The measurements are limited to the range of frequency below 0.34 THz due to the available frequency range of the vector network analyzer. 277 The relatively high value of the transmission losses ( $S_{21}$ ) is due to the difficulty to machining aluminum in this initial fabrication test. An improved surface will be obtained by the use of a different tooling and replacing aluminum with copper. The transmission parameter of the DSG circuit is lower than what was predicted in simulation models, and this is due to poor surface roughness. The first DSG grating circuit has surface roughness with  $R_a$  (arithmetic mean surface roughness) of about 500 nm, and it is expected that this can be improved to well below 100 nm by implementing diamond tooling. However, the fabricated samples demonstrated the CNC milling as a suitable process for SWS in the sub-THz range.

## 290 VI. GUN AND WINDOW

291 The design of the electron gun for the cylindrical beam is based on a conventional Pierce gun and does not present 292

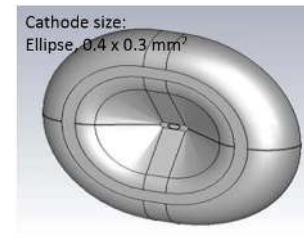


Fig. 14. Electron gun schematic.

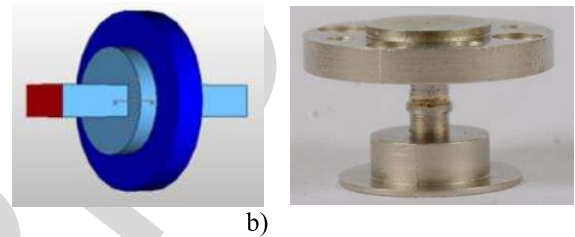
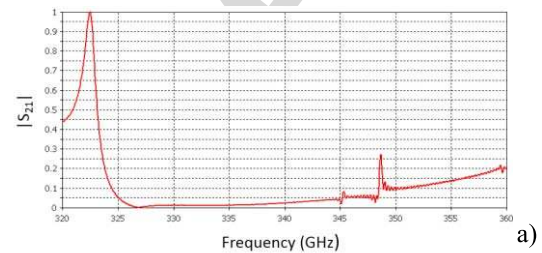


Fig. 15. 0.346-THz window. (a) Simulated S-parameter ( $S_{21}$ ). (b) Schematic. (c) Prototype.

293 specific novelty. On the contrary, the sheet beam requires an accurate design of the gun and the magnetic focusing system. 294

295 A planar cathode is considered to generate the cylindrical electron beam. A preliminary simulation and test was performed, where a beam voltage of 17.4 kV and a current of 14 mA have been obtained. The elliptical electron beam has a ratio 5.4:1 with a current density of 94 A/cm<sup>2</sup> and a 50% fill factor [6]. The schematic for the gun is shown in Fig. 14. 300

301 Different solutions of magnetic focusing systems based on solenoidal structures are under investigation to obtain up to 1.3 T for the full length of the DSG BWO. Based on PIC analysis performed in CST, 98.5% beam transmission efficiency is expected. The solenoid magnet structure has a radial component of magnetic field of 1.3 T and a longitudinal component of 0.35 T. External dimensions of the magnetic structure are  $62 \times 32 \times 35.4$  mm<sup>3</sup>. Engineering estimates predict that the weight of the full system, including magnets, will be under 10 pounds. 310

311 A window, suitable for both the DSG and DWG BWOs, was designed and simulated using CST MWS [Fig. 15(a)] and tested in the frequency range 327–347 GHz. The window is a pillbox-type with MPCVD diamond as the disk. The MCVCP diamond dielectric constant is 5.6 with the loss tangent of 0.003 in the simulation. The thickness of the disk is 0.3 mm, the diameter is 2 mm, and the depth and the diameter of the circular waveguides are 0.7 and 1.2 mm, respectively. 318 The flange connecting the internal SWS and the outer load is WR2.8 rectangular waveguide. Further refinements are 320



in progress. Fig. 15(b) and (c) shows the schematic and a first prototype of the window, respectively.

## VII. CONCLUSION

An international collaboration of leading institutions in vacuum electronics in China, the U.K., and the U.S. is working on building a new family of THz vacuum electron devices for medium power generation. The availability of these sources will permit to enable a new high- $k$  plasma diagnostic to improve the understanding of plasma turbulence in nuclear fusion reactor and many other applications in the THz range. Two different topologies of SWSs have been adopted to design the 0.346-THz BWOs. The DSG and the DCW have been demonstrated to be suitable interaction structures to provide a wide range of performance, with a tailored design. The fabrication of the SWS is a formidable challenge. The samples realized by CNC milling have proved the high accuracy of the process.

The fabrication of all the parts for the final assembly of the BWOs is in progress.

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# THz Backward-Wave Oscillators for Plasma Diagnostic in Nuclear Fusion

Claudio Paoloni, *Senior Member, IEEE*, Diana Gamzina, Logan Himes, Branko Popovic, Robert Barchfeld, Lingna Yue, Yuan Zheng, Xiaopin Tang, Ye Tang, Pan Pan, Hanyan Li, Rosa Letizia, *Member, IEEE*, Mauro Mineo, Jinjun Feng, *Senior Member, IEEE*, and Neville C. Luhmann, Jr., *Fellow, IEEE*

**Abstract**—Understanding of the anomalous transport attributed to short-scale length microturbulence through collective scattering diagnostics is key to the development of nuclear fusion energy. Signals in the subterahertz (THz) range (0.1–0.8 THz) with adequate power are required to map wider wavenumber regions. The progress of a joint international effort devoted to the design and realization of novel backward-wave oscillators at 0.346 THz and above with output power in the 1 W range is reported herein. The novel sources possess desirable characteristics to replace the bulky, high maintenance, optically pumped far-infrared lasers so far utilized in this plasma collective scattering diagnostic. The formidable fabrication challenges are described. The future availability of the THz source here reported will have a significant impact in the field of THz applications both for scientific and industrial applications, to provide the output power at THz so far not available.

**Index Terms**—Backward-wave oscillator (BWO), double-corrugated waveguide (DCW), double-staggered grating (DSG), plasma diagnostic, terahertz (THz).

## I. INTRODUCTION

TERAHERTZ (THz) vacuum electron devices are gaining significant consideration when the generation of relatively high power in the frequency range below 1 THz is

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needed [1]–[9]. In particular, the backward-wave oscillator (BWO) is an effective solution to produce relatively high power and stable monochromatic THz signals. BWOs can be electronically tuned over a wide frequency range around the operating frequency and have high stability in frequency (up to  $10^{-7}$  to  $10^{-8}$  by phase locking). However, the only available BWOs are based on old technologies. No compact, affordable, and long-life THz BWOs is currently available.

Recently, the introduction of new high aspect-ratio fabrication processes derived from the MEMS technologies as the lithography, electroplating, and molding (LIGA) [5], [7] and advanced mechanical microfabrication as nano-CNC milling [10] provides an accuracy at the submicrometer level, which satisfies the demanding specifications of interaction structures or slow-wave structures (SWSs) to support the THz operation frequencies. Furthermore, the progress on simulation tools based on accurate 3-D electromagnetic and particle-in-cell (PIC) solvers permits now a reliable prediction of the THz vacuum source performance, to ease the fabrication phase. The development of innovative cathode materials [10] is leading to a novel generation of electron guns, with high current density and long lifetime, fundamental for the overall device mean time between failures and high-frequency stability. Nevertheless, the main challenges are to achieve a surface roughness below the skin depth (66 nm at 1 THz) for minimizing the ohmic losses and the assembly and alignment of the beam with tolerance in the order of tens of micrometers.

Nuclear fusion is one of the fields, where the availability of THz BWOs will have a relevant impact on improving the understanding of a critical phenomenon as the anomalous transport of the plasma. It still remains a fundamental area of investigation, which is essential for the development of fusion energy. The measurement technique is based on the collective Thomson scattering at THz frequency [12]–[15]. The plasma is illuminated by a THz beam that is scattered by the charged particles. The scattered beams due to the electron density fluctuations are detected by a receiver array and thereby map out the location, wavenumber spectrum, and strength of the turbulence.

In the recent upgrade of the high- $k$  scattering system at NSTX experiment at Princeton, the wavenumber  $k_{\theta}$  coverages, and it has been increased to target electron temperature gradient modes by increasing the probe frequency from 0.280 to 0.693 THz. The availability of solid-state sources is limited to about 0.3 THz and 30 mW, while about 100 mW are needed at 0.693 THz to excite a detectable scattering.

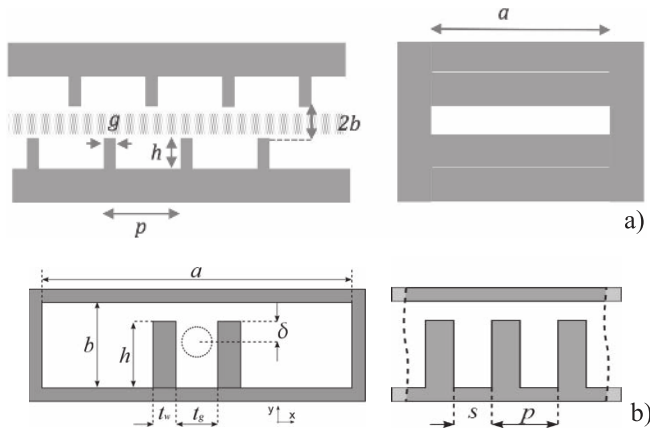


Fig. 1. SWS for 0.346-THz operation. (a) DSG. (b) DCW. The gray area in (a) and the dashed circle in (b) represent the sheet and cylindrical beam for the DSG and the DCW, respectively.

proposed for operation at THz frequencies and overcome the fabrication issues of the conventional structures.

The DSG is conceived to support an elliptical sheet electron beam. The advantage of the sheet electron beam is the large cross section that permits to deliver a high beam current using a relatively beam current density. The sheet beam requires a careful design of the magnetic focusing system, but it is very promising to realize high-power vacuum electron devices by using low cathode loading guns.

The DCW is conceived to support a cylindrical electron beam. The DCW is of easy fabrication and assembly. The advantage of a cylindrical electron beam is that it is generated by the well-established Pierce gun technology and focused by the use of a conventional magnetic focusing system.

Both the SWSs are very promising to realize THz BWOs with a wide range of characteristics in terms of fabrication, output power, and cost.

### III. BWO DESIGN

The approach of using two different SWSs, the DSG and the DCW, to design a family of THz BWO is a breakthrough for tailored power generation at THz frequencies. Two BWOs based on the DSG and the DCW will be the first two devices for a new family of THz sources to cover a wide range of applications.

The main design targets are given in the following:

- 1) low cost;
- 2) easy assembly for high yield;
- 3) compact dimensions (200–300 cm<sup>3</sup>);
- 4) wide range of performance to potentially cover the sub-THz spectrum (0.1–1 THz);
- 5) tunable (at least 5%);
- 6) low beam voltage and compact power supply.

A 0.346-THz operating frequency is considered in the following for application in the plasma diagnostic in nuclear fusion, as described in Section I. The first design parameter defined was the beam voltage that determines the length of the period of the SWSs. A low beam voltage in the range of 12–18 kV favors to use of a compact and low-cost power supply. The resulting period was estimated to be in the range of the fabrication process.

Due to the different structures, different beam voltages were adopted. The DSG was designed to operate with 17-kV beam voltage, while the DCW was to support about 13 kV. The dimensions of the two SWSs are shown in Table I. It is notable that a period shorter than 200  $\mu\text{m}$  is required. In Fig. 2, the dispersion curve of the DCW with superimposed beam line at 12.8 kV is shown. The interaction impedance is typically low in the backward-wave mode.

1) *Couplers*: A detailed study based on 3-D electromagnetic simulation [18] on the coupler to transform the mode in the SWS in the TE<sub>10</sub> at the flange at the output port was carried out to maximize power transfer. The conductivity of copper considered in simulation is  $\sigma_{\text{cu}} = 3.9 \times 10^7$  S/m [7]. The coupler is a three-port network; one port is connected to the SWS, a second port is the beam tunnel to connect the gun, and the third port is the output port connected to the flange.

70 Presently, the only THz source to deliver  $\sim 100$  mW at  
71 0.693 THz needed to provide the minimum scattered signal  
72 level to be detected by the receiver array is a bulky optically  
73 pumped far infrared. A second high-frequency source is  
74 needed to provide local oscillator (LO) power for the  
75 receiver array. The use of a second laser is not feasible.  
76 It has been chosen to use an array based on sensitive room  
77 temperature subharmonic mixers working roughly at half of  
78 the illumination source frequency, namely, 0.346 THz. It is  
79 required 3 to 5 mW of LO power per mixer.

80 This paper describes an international joint effort of three  
81 leading institutions in China, the U.K., and the U.S. to design  
82 and construct a novel family of BWOs, operating at the  
83 frequency of 0.346 THz, to satisfy the quest for LO power for  
84 the matrix array for the NSTX-U plasma diagnostic and other  
85 future plasma diagnostic systems [16]. The design target is to  
86 achieve an output power in the range of hundreds of milliwatt  
87 by lightweight, compact, affordable, low-operating cost BWOs  
88 to enable a wide matrix of receivers.

89 This paper is organized as follows. The properties of two  
90 different SWSs suitable for THz BWO fabrication are reported  
91 in Section II. Section III describes the design aspects and  
92 the cold parameters. Section IV details the hot simulations  
93 and performance of the two BWOs. Challenges involved  
94 with microfabrication technologies of the proposed BWOs are  
95 discussed in Section V. Section VI reports on the gun and the  
96 window.

## II. TERAHERTZ SLOW-WAVE STRUCTURES

98 At microwave frequencies, helices are the most common  
99 SWSs, but as the frequency increases toward the millimeter-  
100 wave range, their dimensions are too small for fabrication, and  
101 new geometries must be adopted. The simple structure of the  
102 rectangular corrugated waveguide inspired different structures  
103 that can be realized by the available fabrication processes with  
104 the dimensions to support THz frequencies.

105 In particular, the double-staggered grating (DSG) [17]  
106 [Fig. 1(a)] and the double-corrugated waveguide (DCW) [18]  
107 [Fig. 1(b)] are the two SWSs that have been successfully



TABLE I

DSG		DCW	
Parameter	$\mu\text{m}$	Parameter	$\mu\text{m}$
$a$	483	$a$	1500
$b$	45	$b$	200
$h$	170	$h$	140
$p$	185	$p$	140
$g$	44	$s$	80
		$t_g$	120
		$t_w$	60

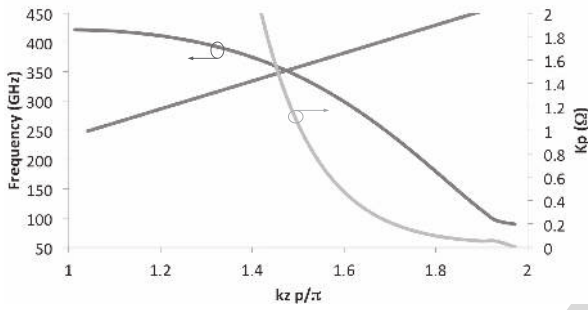


Fig. 2. DCW dispersion curve (brown curve), interaction impedance (orange curve), and beam line at 12.8 kV (blue line).

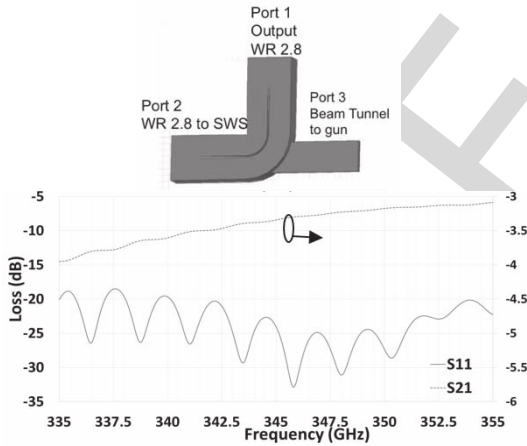


Fig. 3. DSG coupler S-parameters.

165 The coupler for the DSG is particularly challenging due to  
 166 the wide beam tunnel needed for the sheet beam. Having a  
 167 low cutoff frequency in the same range of the SWS, a ridge  
 168 is added to the bend (Fig. 3) to perturb the matching between  
 169 the SWS and the beam tunnel. The resulting  $S_{11}$  is better than  
 170  $-25$  dB over a wide frequency range around 0.346 THz.

171 A study of the coupler for the DCW was performed by  
 172 considering a back-to-back structure. First, a simple structure  
 173 including a tapered transition between a waveguide with  
 174 the same cross section of the DWG and the flanges is  
 175 designed, as shown in Fig. 4(a), to evaluate the effect of the

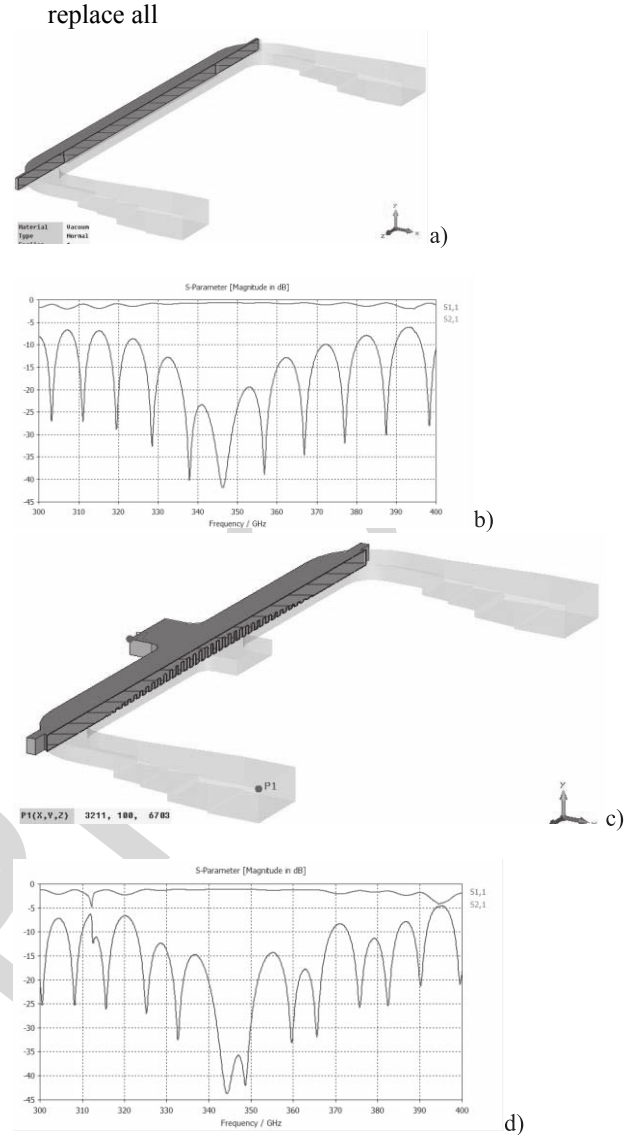


Fig. 4. DCW coupler. (a) Waveguide without DCW. (b) S-parameter structure. (c) Waveguide with DCW. (d) S-parameters.

176 waveguide tapering. Fig. 4(b) shows the obtained  $S_{11}$  better  
 177 than  $-25$  dB in the operation range. Next, a second structure  
 178 with similar topology, including three sections of pillars (two  
 179 tapered sections and one short section with a nominal height),  
 180 is designed for the best matching, as shown in Fig. 4(c).  
 181 Results show that  $S_{11}$  in this case is better than  $-35$  dB  
 182 in the region around the operating frequency [Fig. 4(d)].

183 Both the couplers' performance ensures the efficient prop-  
 184 agation of the RF signal from the interaction structure to the  
 185 flanges.

#### IV. LARGE SIGNAL SIMULATIONS

186 The design of the two BWOs is based on the definition  
 187 of the critical length for oscillations to set a proper number  
 188 of periods within the SWS. Results from this optimization  
 189 process are shown in Table II.

190 Next, the 3-D PIC simulations performed to evaluate the  
 191 electrical behavior of the BWOs. The DSG BWO supports an  
 192

TABLE II

BWO specifications	DSG	DCW
Beam Voltage	17.1 kV	12.8 kV
Beam current	14 mA	10 mA
Beam channel	483 x 90 $\mu\text{m}$	120 $\mu\text{m}$
Beam Aspect Ratio	5.4 : 1	Round
Beam Current Density	160 A/cm <sup>2</sup>	127 A/cm <sup>2</sup>
Magnetic field	0.35 T	0.5 T
No. Periods	65	116
Total length	~ 15 mm	~ 20 mm

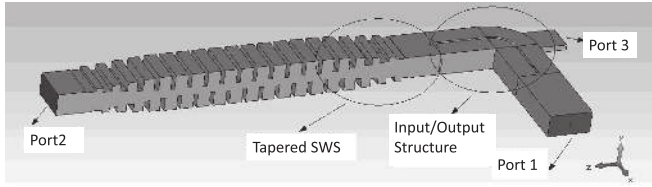


Fig. 5. PIC simulation setup for the DSG BWO.

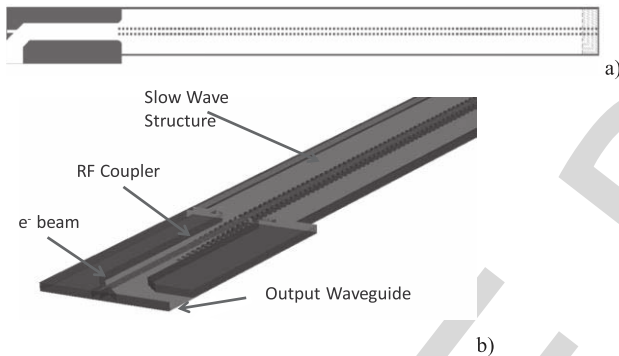


Fig. 6. PIC simulation setup for the DCW PIC simulations. (a) Top view. (b) Coupler detail.

193 elliptical electron beam with the cross section of  $400 \times 50 \mu\text{m}^2$   
 194 and the current of 14 mA (the aspect ratio of 5.4:1). The DCW  
 195 BWO supports a cylindrical electron beam of 50- $\mu\text{m}$  radius  
 196 and 10-mA current. Due to the different beam parameters  
 197 used, the overall device performance is different for the two  
 198 BWOs and should be evaluated in the context of the different  
 199 technological challenges required and the different application  
 200 needs. The DSG BWO is modeled by CST Particle Studio [19]  
 201 (simulation setup in Fig. 5) and the DCW BWO is modeled  
 202 by MAGIC3D [20] (simulation setup in Fig. 6).

203 The results of output power for the two devices are shown  
 204 in Fig. 7. The DSG BWO [Fig. 7(a)] provides about 1 W  
 205 and the DCW BWO about 0.4 W [Fig. 7(b)]. The electron  
 206 energy distribution along the longitudinal direction for both  
 207 the BWOs is shown in Fig. 8. The spectral response at the  
 208 output port of the DCW BWO is shown in Fig. 9, showing the  
 209 highly monochromatic generation of signal at the frequency of  
 210 interest.

211 In order to demonstrate the tunability of the BWO designs,  
 212 the tuning range for the DSG and the DCW is shown  
 213 in Fig. 10(a) and (b), respectively. It can be noted that a

to replace

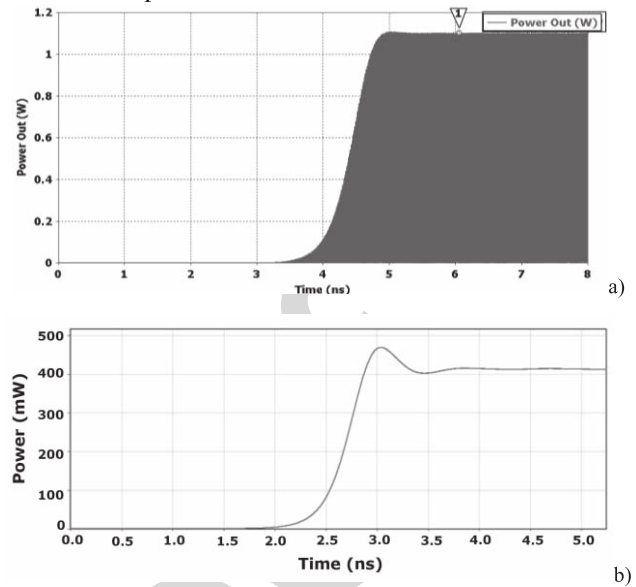


Fig. 7. Average output power for (a) DSG and (b) DCW BWOs.

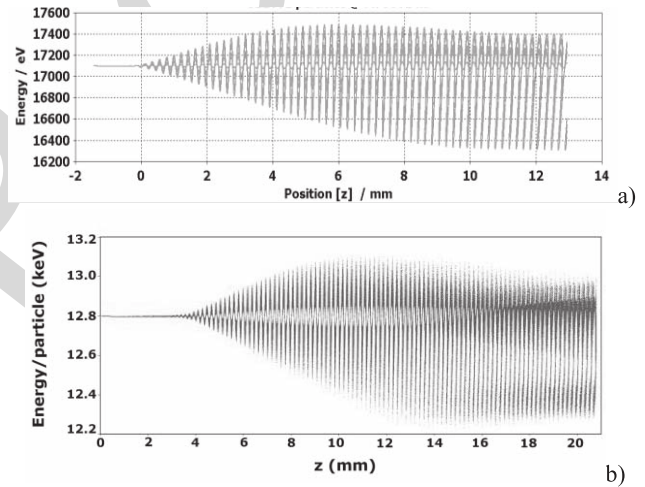


Fig. 8. Electron energy. (a) DSG BWO. (b) DCW BWO.

214 variation of beam voltage in the range of 15.8–17.8 kV permits  
 215 a variation in frequency of 12 GHz for the DSG and that the  
 216 same relative change in the nominal beam voltage allows a  
 217 tuning of 14 GHz for the DCW.

218 The BWO performances so far presented are at the state of  
 219 the art. The high power level and tuning features, not achiev-  
 220 able by any other technology today, represent a breakthrough  
 221 in the field.

## V. BWO MICROFABRICATION

222 The main challenge in the THz frequency range is the fab-  
 223 rication of SWSs with the expected electromagnetic behavior  
 224 while establishing a reliable and repeatable process. For the  
 225 DSG circuit, vane height is the most sensitive dimension which  
 226 determines the bandwidth of the device. The period of the  
 227 structure affects the central operating frequency, whereas the  
 228 width of the DSG controls the dispersion curve. The DCW  
 229 structure is more sensitive to the  $h$  and  $p$  values driving the  
 230

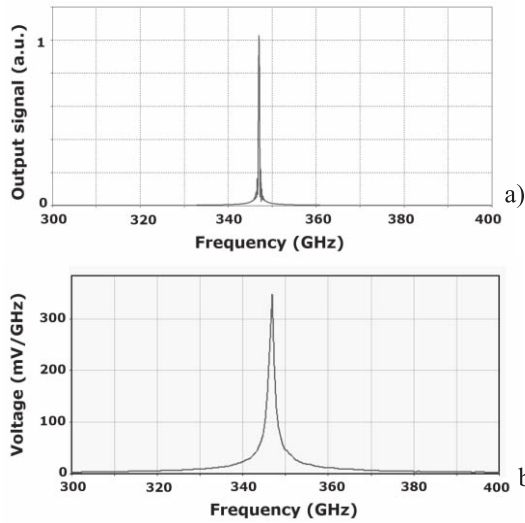


Fig. 9. Spectrum of (a) DSG and (b) DCW BWOs.

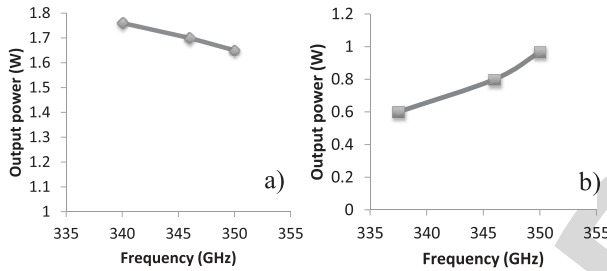


Fig. 10. Instantaneous power and tuning range of (a) DSG and (b) DCW BWOs for the beam voltage 15.8–17.8 and 11.75–13.3 kV, respectively.

231 dispersion curve. Machining tolerances are expected to be  
 232  $\pm 1 \mu\text{m}$ , which is sufficient to achieve the desired performance.

233 Photolithographic techniques, such as UV-LIGA, are  
 234 demonstrated suitable for the dimension accuracy required  
 235 for the two SWSs considered. However, especially for a  
 236 small number of pieces, the fabrication of the mold and  
 237 the electroforming process is not convenient. Furthermore,  
 238 a relevant effort to achieve a level of surface roughness better  
 239 than the skin depth (about 110 nm at 0.346 THz) to reduce  
 240 ohmic losses is required. CNC milling offers high flexibility  
 241 and possibility of patterning the third dimension. The state-  
 242 of-the-art prototype nano-CNC milling machine, developed  
 243 by DTL, a subsidiary of DMG-Mori-Seki, permits one to  
 244 achieve performance at the state of the art, with reduced  
 245 cost and high repeatability for dimensions suitable for THz  
 246 regime structures [10]. The high accuracy of the nanomilling  
 247 machine was proved to obtain levels of surface roughness  
 248 down to 40 nm, well below the skin depth at 0.346 THz. The  
 249 NN1000 nano/micromilling machine has a maximum spindle  
 250 speed of 50 000 r/min; the chip load is kept below 0.001-mm  
 251 feed per tool flute rpm. The proper setting of the machining  
 252 parameters is fundamental in achieving excellent surface finish  
 253 and tool lifetime.

254 In the case of the DSG and the DGW at 0.346 THz,  
 255 the dimensions shown in Table I represent a formidable  
 256 fabrication challenge. A test of feasibility for the fabrication  
 257 was performed realizing the DSG and the DCW in aluminum

to replace

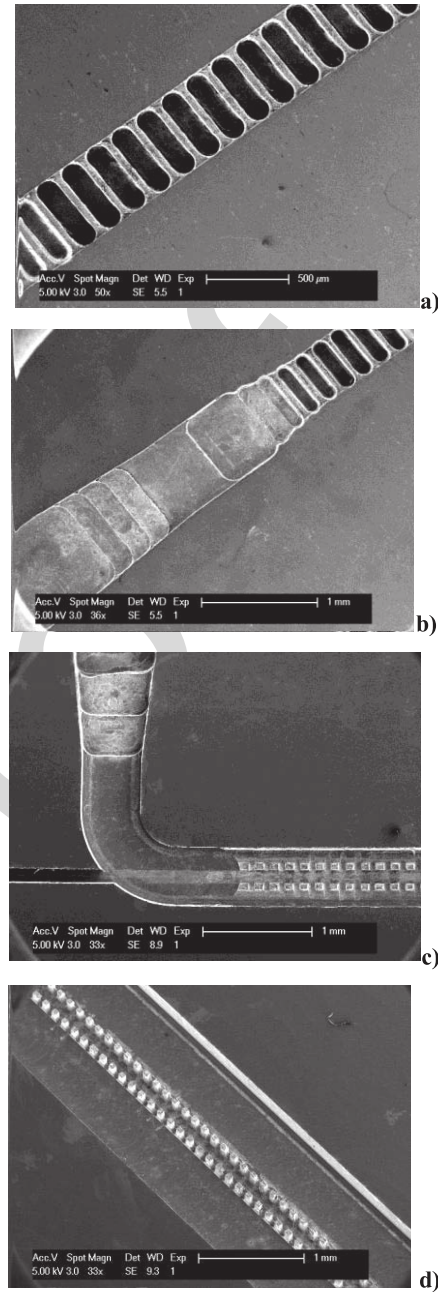


Fig. 11. SEM images of the SWSs fabricated by nano-CNC milling. (a) and (b) DSG. (c) and (d) DCW.

258 with the dimensions shown in Table I. Four different SEM  
 259 views of the DSG and the DCW realized by nano-CNC  
 260 milling are shown in Fig. 11. The high level of accuracy  
 261 for the very small dimensions is readily observed. The high  
 262 definition of the pillars is notable. Due to the characteristics  
 263 of aluminum, the surface roughness achieved was higher than  
 264 expected. The fabricated samples were primarily built to test  
 265 the microfabrication process in terms of dimensions achieved.  
 266 However, the measurements of the S-parameters were carried  
 267 out. The setup for the DSG measurement is shown in Fig. 12.  
 268 It consists of two halves assembled together by a system of  
 269 alignment pins. The setup for the DCW is similar to a lid  
 270 to close the waveguide that does not require a very accurate  
 271 alignment procedure.



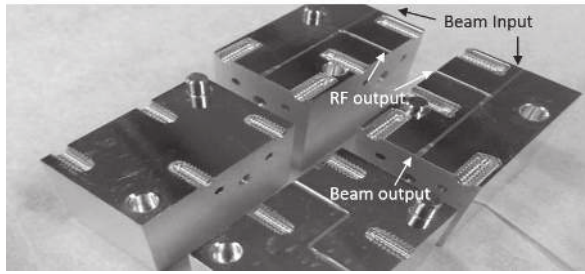


Fig. 12. Full DSG assembly with alignment pins and flanges.

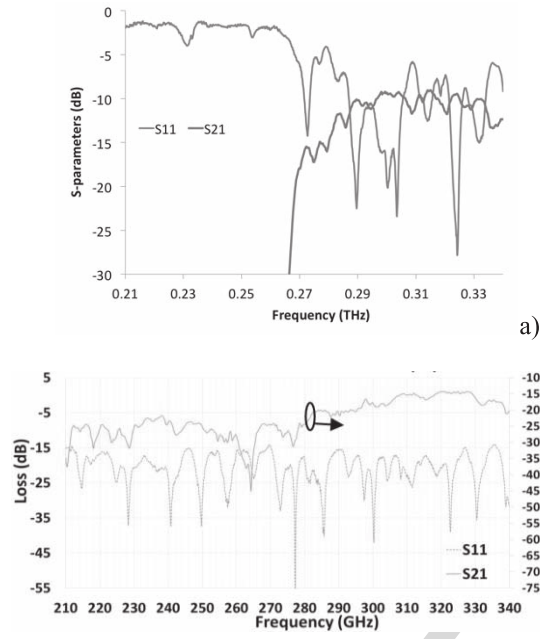


Fig. 13. Measurements of the S-parameters of (a) DCW and (b) DSG in the low range of the band.

272 The S-parameters of the fabricated structures were mea- 293  
 273 sured. The  $S_{11}$  and  $S_{21}$  for the fabricated DSG and DCW 294  
 274 are shown in Fig. 13(a) and (b), respectively. The mea- 295  
 275 surements are limited to the range of frequency below 0.34 THz due to 296  
 276 the available frequency range of the vector network analyzer. 297  
 277 The relatively high value of the transmission losses ( $S_{21}$ ) is 298  
 278 due to the difficulty to machining aluminum in this initial 299  
 279 fabrication test. An improved surface will be obtained by 300  
 280 the use of a different tooling and replacing aluminum with 301  
 281 copper. The transmission parameter of the DSG circuit is 302  
 282 lower than what was predicted in simulation models, and 303  
 283 this is due to poor surface roughness. The first DSG grating 304  
 284 circuit has surface roughness with  $R_a$  (arithmetic mean surface 305  
 285 roughness) of about 500 nm, and it is expected that this can 306  
 286 be improved to well below 100 nm by implementing diamond 307  
 287 tooling. However, the fabricated samples demonstrated the 308  
 288 CNC milling as a suitable process for SWS in the sub-THz 309  
 289 range. 310

## 290 VI. GUN AND WINDOW

291 The design of the electron gun for the cylindrical beam 312  
 292 is based on a conventional Pierce gun and does not present 313

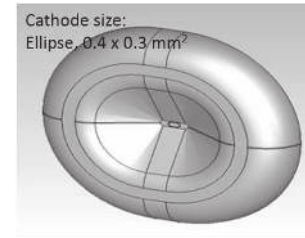


Fig. 14. Electron gun schematic.

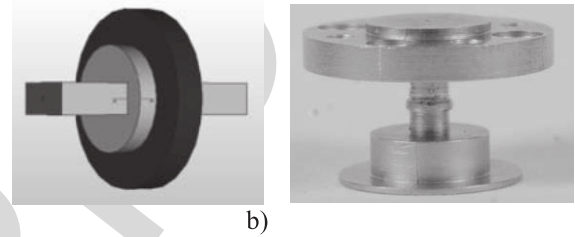
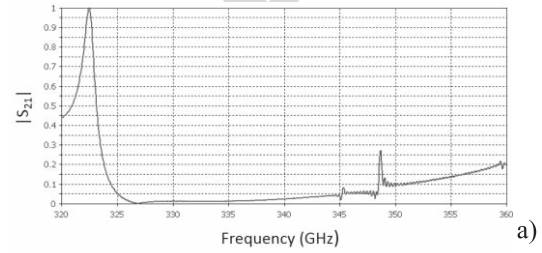


Fig. 15. 0.346-THz window. (a) Simulated S-parameter ( $S_{21}$ ). (b) Schematic. (c) Prototype.

specific novelty. On the contrary, the sheet beam requires an 293  
 accurate design of the gun and the magnetic focusing system. 294

A planar cathode is considered to generate the cylindrical 295  
 electron beam. A preliminary simulation and test was per- 296  
 formed, where a beam voltage of 17.4 kV and a current of 297  
 14 mA have been obtained. The elliptical electron beam has 298  
 a ratio 5.4:1 with a current density of 94 A/cm<sup>2</sup> and a 50% 299  
 fill factor [6]. The schematic for the gun is shown in Fig. 14. 300

Different solutions of magnetic focusing systems based 301  
 on solenoidal structures are under investigation to obtain up 302  
 to 1.3 T for the full length of the DSG BWO. Based on 303  
 PIC analysis performed in CST, 98.5% beam transmission 304  
 efficiency is expected. The solenoid magnet structure has a 305  
 radial component of magnetic field of 1.3 T and a longitudinal 306  
 component of 0.35 T. External dimensions of the magnetic 307  
 structure are  $62 \times 32 \times 35.4$  mm<sup>3</sup>. Engineering estimates 308  
 predict that the weight of the full system, including magnets, 309  
 will be under 10 pounds. 310

A window, suitable for both the DSG and DWG BWOs, 311  
 was designed and simulated using CST MWS [Fig. 15(a)] 312  
 and tested in the frequency range 327–347 GHz. The window 313  
 is a pillbox-type with MPCVD diamond as the disk. The 314  
 MCVDP diamond dielectric constant is 5.6 with the loss 315  
 tangent of 0.003 in the simulation. The thickness of the disk is 316  
 0.3 mm, the diameter is 2 mm, and the depth and the diameter 317  
 of the circular waveguides are 0.7 and 1.2 mm, respectively. 318  
 The flange connecting the internal SWS and the outer load 319  
 is WR2.8 rectangular waveguide. Further refinements are 320

in progress. Fig. 15(b) and (c) shows the schematic and a first prototype of the window, respectively.

## VII. CONCLUSION

An international collaboration of leading institutions in vacuum electronics in China, the U.K., and the U.S. is working on building a new family of THz vacuum electron devices for medium power generation. The availability of these sources will permit to enable a new high- $k$  plasma diagnostic to improve the understanding of plasma turbulence in nuclear fusion reactor and many other applications in the THz range. Two different topologies of SWSs have been adopted to design the 0.346-THz BWOs. The DSG and the DCW have been demonstrated to be suitable interaction structures to provide a wide range of performance, with a tailored design. The fabrication of the SWS is a formidable challenge. The samples realized by CNC milling have proved the high accuracy of the process.

The fabrication of all the parts for the final assembly of the BWOs is in progress.

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