

Time-Orthogonal Unitary Dilations and Noncommutative Feynman-Kac Formulae

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Abstract. An analysis of Feynman-Kac formulae reveals that, typically, the unperturbed semigroup is expressed as the expectation of a random unitary evolution and the perturbed semigroup is the expectation of a perturbation of this evolution in which the latter perturbation is effected by a cocycle with certain covariance properties with respect to the group of translations and reflections of the line. We consider generalisations of the classical commutative formalism in which the probabilistic properties are described in terms of non-commutative probability theory based on von Neumann algebras. Examples of this type are generated, by means of second quantisation, from a unitary dilation of a given self-adjoint contraction semigroup, called the time orthogonal unitary dilation, whose key feature is that the dilation operators corresponding to disjoint time intervals act nontrivially only in mutually orthogonal supplementary Hilbert spaces.

1. Introduction

For any positive-self-adjoint operator H in a Hilbert space \mathcal{h}_0 we construct a *time-orthogonal unitary dilation* $\{U_{s,v}, s \geq t\}$ of the contraction semigroup $\{e^{-tH}, t \geq 0\}$ in the Hilbert space $\mathcal{h}_0 \oplus L_2(\mathbb{R}, \mathcal{h}_0)$ which exhibits the following properties:

(1) the family $\{U_{s,v}, s \geq t\}$ is a strongly continuous unitary evolution in the sense that $U_{r,s}U_{s,t} = U_{r,t}$ for all $r \geq s \geq t$;

(2) it is time-orthogonal in the sense that $U_{s,t}$ acts nontrivially only in the subspace $\mathcal{h}_0 \oplus \mathcal{h}_{[s,t]}$ where $\mathcal{h}_{[s,t]} \subset L_2(\mathbb{R}, \mathcal{h}_0)$ is the subspace of functions with support contained in $]s, t]$;

(3) it is covariant under time translations and time reversal. This dilation is somewhat different from the classical dilations [14] of a one parameter con-

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traction semigroup into a one parameter group, examples of which may be found in [10, 4, 5].

Using the unitary evolution induced by second quantisation of the above dilation in the Fock space $\Gamma(\mathfrak{h}_0 \oplus L_2(\mathbb{R}, \mathfrak{h}_0))$ and contracting a suitable perturbation of this new evolution to the smaller Fock space $\Gamma(\mathfrak{h}_0)$ we obtain a noncommutative version of the Feynman-Kac formula for the semigroup $\exp\{-t(d\Gamma(H) + V)\}$ where V is a bounded operator in $\Gamma(\mathfrak{h}_0)$ and $\exp\{-td\Gamma(H)\} = \Gamma(\exp(-tH))$, Γ being the second quantisation map. For an alternative version of a Feynman-Kac formula based on classical dilations and completely positive maps in the theory of quantum stochastic processes the reader may refer to Accardi [1].

We use the following conventions. A product $\prod_{j=1}^n A_j$ of noncommuting elements is taken in the order $A_n A_{n-1} \dots A_1$. A *canonical pair* is a pair (p, q) of self-adjoint operators satisfying the Weyl relation

$$e^{ixp} e^{iyq} = e^{ixy} e^{iyq} e^{ixp}, \quad x, y \in \mathbb{R}.$$

We shall say that von Neumann algebras N_1, N_2, \dots, N_n in a Hilbert space \mathfrak{h} are *independent* (in the state Ω) if Ω is a unit vector in \mathfrak{h} for which

$$\left\langle \Omega, \prod_{j=1}^n A_j \Omega \right\rangle = \prod_{j=1}^n \langle \Omega, A_j \Omega \rangle, \quad A_j \in N_j, \quad j=1, \dots, n.$$

If \mathfrak{h}_0 is a second Hilbert space then the *conditional expectation* of $A \in \mathbf{B}(\mathfrak{h}_0 \otimes \mathfrak{h})$ given $\mathbf{B}(\mathfrak{h}_0)$ is the element $\mathbb{E}[A|\mathbf{B}(\mathfrak{h}_0)] \in \mathbf{B}(\mathfrak{h}_0)$ for which

$$\langle \psi, \mathbb{E}[A|\mathbf{B}(\mathfrak{h}_0)] \chi \rangle = \langle \psi \otimes \Omega, A \chi \otimes \Omega \rangle.$$

The following theorem is almost obvious [8].

Theorem 1.1. *Let N_1, N_2, \dots, N_n be independent von Neumann algebras and set*

$$N^{(j)} = \mathbf{B}(\mathfrak{h}_0) \otimes N_j, \quad j=1, \dots, n.$$

Then, for arbitrary $A^{(j)} \in N^{(j)}, j=1, \dots, n,$

$$\mathbb{E} \left[\prod_{j=1}^n A^{(j)} | \mathbf{B}(\mathfrak{h}_0) \right] = \prod_{j=1}^n \mathbb{E} \left[A^{(j)} | \mathbf{B}(\mathfrak{h}_0) \right]. \tag{1.1}$$

2. Evolutions and Perturbations

Let \mathfrak{h} be a separable Hilbert space and let $(U_{s,t}; s \geq t)$ be a family of bounded operators on \mathfrak{h} satisfying the conditions

- a) $U_{s,t}$ is strongly continuous in s for fixed t and in t for fixed s .
- b) $U_{s,s} = I$ for all $s \in \mathbb{R}$.
- c) $U_{r,s} U_{s,t} = U_{r,t}$ for $r \geq s \geq t$.

Such a family $(U_{s,t})$ is called an *evolution*. If each $U_{s,t}$ is unitary (respectively self-adjoint, a contraction) then the evolution is said to be *unitary* (respectively *self-adjoint*, *contractive*). If $U_{s,t} = U_{s-t}$ depends only on the difference $s - t$ then the

evolution is a *semigroup*. Every (strongly continuous) self-adjoint semigroup is of the form $U_{s-t} = \exp\{- (s-t)H\}$ where H is a not necessarily bounded self-adjoint operator called its *infinitesimal generator*.

A family $(M_{s,t} : s \geq t)$ of bounded operators in \mathcal{H} , strongly continuous in s and t separately and satisfying $M_{s,s} = I$ for all $s \in \mathbb{R}$ is said to be a *cocycle* for the evolution $(U_{s,t} : s \geq t)$ if, for $r \geq s \geq t$,

$$M_{r,t} = M_{r,s} U_{r,s} M_{s,t} U_{r,s}^{-1}. \tag{2.1}$$

It is easy to see that if $(M_{s,t})$ is a cocycle for the unitary evolution $(U_{s,t})$ and if, for $s \geq t$,

$$\tilde{U}_{s,t} = M_{s,t} U_{s,t},$$

then $(\tilde{U}_{s,t}, s \geq t)$ is also an evolution; indeed (2.1) is equivalent to the defining condition c) and continuity follows from the strong continuity of operator multiplication when one of the factors is uniformly bounded. We call the evolution $(\tilde{U}_{s,t})$ the *perturbation* of $(U_{s,t})$ by the cocycle $(M_{s,t})$.

The following theorem can be proved using the Picard method of successive approximations to approximate $M_{s,t}^B$ in operator-bound norm by $M_{s,t}^{(n)}$, defined inductively by $M_{s,t}^{(0)} = I$,

$$M_{s,t}^{(n)} = I - \int_t^s M_{s,r}^{(n-1)} U_{s,r} B U_{s,r}^{-1} dr. \tag{2.2}$$

See also [12].

Theorem 2.1. *Let $(U_{s,t})$ be an unitary evolution in the separable Hilbert space \mathcal{H} and let B be a bounded operator in \mathcal{H} . Then there exists a unique cocycle $(M_{s,t}^B : s \geq t)$ such that $M_{s,t}^B$ is strongly differentiable in t and*

$$\partial_t M_{s,t}^B = - M_{s,t}^B U_{s,t} B U_{s,t}^{-1}. \tag{2.3}$$

Moreover if the family $(U_{s,t} B U_{s,t}^{-1} : s \geq t)$ is commutative, then the strong operator valued integral $\int_t^s U_{s,\tau} B U_{s,\tau}^{-1} d\tau$ exists and, for $s \geq t$,

$$M_{s,t}^B = \exp \left\{ - \int_t^s U_{s,\tau} B U_{s,\tau}^{-1} d\tau \right\}. \tag{2.4}$$

We shall call $(M_{s,t}^B)$ the *cocycle determined by B* and the corresponding evolution $(U_{s,t}^B)$, with

$$U_{s,t}^B = M_{s,t}^B U_{s,t}$$

the *perturbed evolution determined by B* .

3. Reduced Evolutions and Semigroups

Let \mathcal{H} be a separable Hilbert space. A *doubly filtered von Neumann algebra* in \mathcal{H} is a pair $(N, (N_{s,t} : s \geq t))$ comprising a von Neumann algebra N and a system $(N_{s,t} : s \geq t)$

of von Neumann subalgebras which generate N as a von Neumann algebra and which satisfy the condition

$$N_{s,t} \subset N_{a,b} \quad \text{whenever }]s, t] \subseteq]a, b].$$

A *covariant system* is a triple $(N, (N_{s,t} : s \geq t), (\gamma_t : t \in \mathbb{R}))$ consisting of a doubly filtered von Neumann algebra $(N, (N_{s,t} : s \geq t))$ and a one-parameter group (γ_t) of normal automorphisms of N satisfying

$$\gamma_r(N_{s,t}) \subseteq N_{r+s, r+t}.$$

A *reflective covariant system* is a quadruple $(N, (N_{s,t}), (\gamma_t), \varrho)$ consisting of a covariant system $(N, (N_{s,t}), (\gamma_t))$ together with a normal automorphism ϱ of N satisfying, for $r \in \mathbb{R}$ and $s \geq t$,

$$\begin{aligned} \varrho \gamma_r &= \gamma_{-r} \varrho, & \varrho &= \varrho^{-1} \\ \varrho N_{s,t} &\subseteq N_{-t, -s}. \end{aligned}$$

Let \mathcal{H}_0 be a second Hilbert space. A *reducing map* for a doubly filtered von Neumann algebra $(N, (N_{s,t}))$ is a strongly continuous linear map j from N to $B(\mathcal{H}_0)$ such that, for all $A \in N$,

$$j(A)^* = j(A^*)$$

and for $A \in N_{a,b}$, $B \in N_{c,d}$ with $]a, b] \cap]c, d] = \emptyset$,

$$j(AB) = j(A)j(B). \tag{3.1}$$

A *covariant reducing map* for a covariant system $(N, (N_{s,t}), (\gamma_t))$ is a reducing map j for $(N, (N_{s,t}))$ for which $j \circ \gamma_r = j$ for all $r \in \mathbb{R}$, and a *reflective covariant reducing map* for a reflective covariant system $(N, (N_{s,t}), (\gamma_t), \varrho)$ is a covariant reducing map for $(N, (N_{s,t}), (\gamma_t))$ for which $j \circ \varrho = j$.

An evolution $(U_{s,t})$ in \mathcal{H} is said to be *adapted* to the doubly filtered von Neumann algebra $(N, (N_{s,t}))$ if

$$U_{s,t} \in N_{s,t} \tag{3.2}$$

for all $s \geq t$, *covariantly* adapted for the covariant system $(N, (N_{s,t}), (\gamma_t))$ if, in addition to (3.2), for all $r \in \mathbb{R}$ and $s \geq t$

$$\gamma_r(U_{s,t}) = U_{r+s, r+t}, \tag{3.3}$$

and *reflectively* covariantly adapted for the reflective covariant system $(N, (N_{s,t}), (\gamma_t), \varrho)$ if, in addition to (3.2) and (3.3),

$$\varrho(U_{s,t}^*) = U_{-t, -s}. \tag{3.4}$$

It is clear from (3.1) that if $(U_{s,t})$ is an adapted evolution and j is a reducing map for $(N, (N_{st}))$

$$U_{s,t}^{(j)} = j(U_{s,t})$$

defines an evolution $(U_{s,t}^{(j)} : s \geq t)$ in \mathcal{H}_0 called the *reduced evolution* corresponding to j .

Theorem 3.1. (a) *The reduced evolution of a covariantly adapted evolution corresponding to a covariant reducing map is a semigroup.*

(b) *The reduced evolution of a reflectively covariantly adapted evolution corresponding to a reflective covariant reducing map is a self-adjoint semigroup.*

Proof. (a) If $(U_{s,t})$ is covariantly adapted and j is a covariant reducing map then, for $r \in \mathbb{R}, s \geq t$

$$\begin{aligned} U_{r+s,r+t}^{(j)} &= j(U_{r+s,r+t}) \\ &= j \circ \gamma_r(U_{s,t}) \\ &= j(U_{s,t}) \\ &= U_{s,t}^{(j)}, \end{aligned}$$

showing that $U_{s,t}^{(j)}$ depends only on the difference $s-t$ and hence that the reduced evolution is a semigroup.

(b) If in addition $(U_{s,t})$ is reflectively covariantly adapted and j is a reflective covariant reducing map then, for $s \geq t$,

$$\begin{aligned} [U_{s,t}^{(j)}]^* &= [j(U_{s,t})]^* \\ &= j(U_{s,t}^*) \\ &= j \circ \varrho(U_{-t,-s}) \\ &= j(U_{-t,-s}) \\ &= U_{-t,-s}^{(j)}. \end{aligned}$$

Thus, using the semigroup property, $U_{s,t}^{(j)}$ is self-adjoint. \square

Now we examine the reduction of perturbed evolutions.

Theorem 3.2. (a) *Let $(N, (N_{s,t}))$ be a doubly filtered von Neumann algebra and let B be an element of $\bigcap_{s \geq t} N_{s,t}$. Then, if $(U_{s,t})$ is an adapted evolution then the perturbed evolution $(U_{s,t}^B)$ is adapted.*

(b) *If in addition $(N, (N_{s,t}), (\gamma_t))$ is a covariant system, $(U_{s,t})$ is covariantly adapted and B is invariant under each γ_r , then the perturbed evolution $(U_{s,t}^B)$ is covariantly adapted.*

(c) *If in addition $(N, (N_{s,t}), (\gamma_t), \varrho)$ is a reflective covariant system, $(U_{s,t})$ is reflectively covariantly adapted, and B is self adjoint, then the semigroup reduced evolution obtained from $(U_{s,t}^B)$ from a reflective covariant reducing map j is self-adjoint with infinitesimal generator*

$$H^B = H + j(B),$$

where H is the infinitesimal generator of the reduced unperturbed evolution.

Proof. (a) To prove that $(U_{s,t}^B)$ is adapted we must show that, for $s \geq t, M_{s,t}^B \in N_{s,t}$. Recalling from Theorem 2.1 that $M_{s,t}^B$ was the uniform limit of operators $M_{s,t}^{(n)}$ defined by (2.2), it is sufficient to show that each $M_{s,t}^{(n)} \in N_{s,t}$. Since for $t \leq \tau \leq s, N_{s,\tau} \subseteq N_{s,t}$ and $U_{s,\tau}, B \in N_{s,\tau}$, the required conclusion follows by induction.

(b) To prove that $(U_{s,t}^B)$ is covariantly adapted we must show that, for $r \in \mathbb{R}$, $s \geq t$,

$$\gamma_r(M_{s,t}^B) = M_{r+s,r+t}^B.$$

It is sufficient to prove the corresponding property for the approximations $M_{s,t}^{(n)}$

$$\gamma_r(M_{s,t}^{(n)}) = M_{r+s,r+t}^{(n)}, \quad n=0, 1, 2, \dots, s \geq t. \tag{3.5}$$

The case $n=0$ is clear since each $M_{s,t}^{(0)} = I$; assuming (3.5) for $n=k-1$ we have, using (2.2),

$$\begin{aligned} \gamma_r(M_{s,t}^{(k)}) &= I - \gamma_r \left\{ \int_t^s M_{s,\tau}^{(k-1)} U_{s,\tau} B U_{s,\tau}^{-1} d\tau \right\} \\ &= I - \int_t^s \gamma_r(M_{s,\tau}^{(k-1)}) \gamma_r(U_{s,\tau}) B \gamma(U_{s,\tau})^{-1} d\tau \\ &= I - \int_t^s M_{r+s,r+\tau}^{(k-1)} U_{r+s,r+\tau} B U_{r+s,r+\tau}^{-1} d\tau \\ &= I - \int_{r+t}^{r+s} M_{r+s,\tau}^{(k-1)} U_{r+s,\tau} B U_{r+s,\tau}^{-1} d\tau \\ &= M_{r+s,r+t}^{(k)} \end{aligned}$$

and the result follows by induction.

(c) By (b) and Theorem 3.1 (a) we know that the reduced evolution $((U_{s,t}^B)^j)$ is a semigroup so that we can write

$$(U_{s,t}^B)^j = T_{s-t}.$$

If f is in the domain $D(H)$ of H we have, for $t > 0$,

$$\begin{aligned} t^{-1}(T_t - I)f &= t^{-1}j(M_{0,-t}^{(B)} U_{0,-t} - I)f \\ &= j(t^{-1}(M_{0,-t}^{(B)} U_{0,-t} - U_{0,-t}))f + t^{-1}j(U_{0,-t} - I)f. \end{aligned}$$

Since j is strongly continuous and $f \in D(H)$ we can pass to the limit as $t \rightarrow 0$, using (2.3), to conclude that

$$\lim_{t \downarrow 0} t^{-1}(T_t - I)f = -j(B)f - Hf.$$

But this means that (T_t) is the one-parameter semigroup of which the self-adjoint operator $j(B) + H$ is the infinitesimal generator; in particular the reduced semigroup is self-adjoint. \square

We may summarise Theorem 3.2 (c) by the statement

$$e^{-(s-t)(H+j(B))} = j(M_{s,t}^B U_{s,t})$$

whenever $(U_{s,t})$ is a reflectively covariantly adapted evolution, j is a reflective covariant reducing map, and

$$e^{-(s-t)H} = j(U_{s,t}).$$

4. The Classical Case Generalised to Imprimitivity Systems

Let G be a complete, separable metric group acting on a Borel space (X, \mathcal{B}) and let \mathcal{H}_0 be a separable Hilbert space. Suppose that $g \rightarrow U_g$ is a strongly continuous unitary representation of G in \mathcal{H}_0 and $A \rightarrow E(A)$ is a spectral measure on \mathcal{B} relative to \mathcal{H}_0 satisfying

$$U_g E(A) U_g^{-1} = E(gA) \quad \text{for all } g \in G, \quad A \in \mathcal{B}. \tag{4.1}$$

Such a pair (U, E) is called an *imprimitivity system*.

Consider a strongly continuous, separable and measurable G -valued stochastic process $\{\omega(t), t \in \mathbb{R}\}$ defined on a probability space (Ω, F, μ) and satisfying the following conditions: (1) $\omega(0) = e$ with probability one; (2) $\omega(t)$ and $\omega(t)^{-1}$ have the same distribution for all t ; (3) for disjoint intervals $[t_j, s_j]$ the random variables $\omega(s_j)\omega(t_j)^{-1}$ are independent; (4) the distribution of $\omega(r+s)\omega(r+t)^{-1}$ does not depend on r whenever $s \geq t$ are fixed. Such a process is called a *symmetric G -valued process with stationary independent increments*. We denote by $F_{s,t}$ the smallest σ -algebra with respect to which all the random variables $\omega(a)\omega(b)^{-1}$, $t \leq a \leq b \leq s$ are measurable.

Regarding the existence of a large class of such processes with independent increments the reader may refer to [9, 11].

Let $\mathcal{H} = L_2(\mu, \mathcal{H}_0)$ be the Hilbert space of all \mathcal{H}_0 -valued μ -square integrable maps. We denote by N and $N_{s,t}$ respectively the von Neumann algebras of all bounded operators which are multiplications by F - and $F_{s,t}$ -measurable \mathcal{H}_0 -operator valued functions. For any $\omega \in \Omega$ and $t \in \mathbb{R}$ define $\tilde{\omega}$ and $\omega^t \in \Omega$ by putting $\tilde{\omega}(s) = \omega(s)^{-1}$; $\omega^t(s) = \omega(s+t)\omega(t)^{-1}$ for all $s \in \mathbb{R}$. Define the automorphisms γ_t , $t \in \mathbb{R}$ and ϱ of N by putting

$$(\gamma_t(A)f)(\omega) = A(\omega^t)f(\omega); \quad (\varrho(A)f)(\omega) = A(\tilde{\omega})f(\omega)$$

for every $A \in N$. Define the map $j: N \rightarrow \mathcal{B}(\mathcal{H}_0)$ by the Bochner integral

$$j(A)\xi = \int A(\omega)\xi d\mu(\omega), \quad \xi \in \mathcal{H}_0.$$

Let $\{U_{s,t}, s \geq t\}$ be the operators on \mathcal{H} defined by

$$(U_{s,t}f)(\omega) = U_{\omega(s)\omega(t)^{-1}}f(\omega), \quad f \in \mathcal{H}.$$

Then $(N, (N_{s,t}), \gamma_t, \varrho)$ is a reflective covariant system, j is a reflective covariant reducing map and $(U_{s,t})$ is a reflectively covariantly adapted evolution.

Suppose V is a bounded Borel function on (X, \mathcal{B}) and B is the operator of multiplication by the constant \mathcal{H}_0 -operator valued function $\int V(x)E(dx)$. Then an easy computation using (4.1) shows that

$$(U_{s,t} B U_{s,t}^{-1} f)(\omega) = \int V(\omega(t)\omega(s)^{-1}x)E(dx)f(\omega).$$

It follows from Theorem 2.1 that the cocycle determined by B is given by

$$(M_{s,t}^B f)(\omega) = \left[\exp \left\{ - \int_t^s V(\omega(\tau)\omega(s)^{-1}x) d\tau \right\} E(dx) \right] f(\omega).$$

Since $j(B) = \int V(x)E(dx)$, Theorem 3.2(c) implies that

$$e^{-(s-t)H_0}\xi = \int [U_{\omega(s)\omega(t)^{-1}}\xi]d\mu(\omega), \tag{4.2}$$

$$\begin{aligned} & e^{-(s-t)(H_0 + \int V(x)E(dx))}\xi \\ &= \int \left\{ \int \exp \left\{ - \int_t^s V(\omega(\tau)\omega(s)^{-1}x)d\tau \right\} E(dx) \right\} U_{\omega(s)\omega(t)^{-1}}\xi d\mu(\omega) \end{aligned} \tag{4.3}$$

for all $\xi \in \mathcal{H}_0$ where integration with respect to μ is in the Bochner sense. When $X = G = \mathbb{R}$, $\omega(t)$ is standard Brownian motion, $t \rightarrow U_t$ is the translation representation in $L_2(\mathbb{R})$ and $E(A)$ is multiplication by χ_A , Eq. (4.2) and (4.3) describe the classical Feynman-Kac formula with $H_0 = p^2/2$ and $\int V(x)E(dx) = V(q)$.

5. Time Orthogonal Unitary Dilations

To begin with we consider the case of a strongly continuous self-adjoint contraction semigroup $(A_t, t \geq 0)$ with a bounded positive self-adjoint infinitesimal generator H in a Hilbert space \mathcal{H}_0 . Let $\mathcal{H} = L_2(\mathbb{R}, \mathcal{H}_0)$ be the Hilbert space of all \mathcal{H}_0 -valued square integrable maps on \mathbb{R} and let the shifts $\{S_t, t \in \mathbb{R}\}$ and reflection S be defined on \mathcal{H} by

$$(S_t f)(s) = f(s-t); \quad (Sf)(s) = f(-s), \quad f \in \mathcal{H}, \quad t \in \mathbb{R}. \tag{5.1}$$

For any Borel set $A \subseteq \mathbb{R}$ we write \mathcal{H}_A for the subspace of \mathcal{H} comprising of functions with support in A . Consider the Hilbert space $\mathcal{H}_1 = \mathcal{H}_0 \oplus \mathcal{H}$ and denote by J the imbedding $\zeta \rightarrow \begin{pmatrix} \zeta \\ 0 \end{pmatrix}$ from \mathcal{H}_0 into \mathcal{H} . Define the operators $\{U_{s,t}, s \geq t\}$ in \mathcal{H}_1 by

$$U_{s,t} \begin{pmatrix} \zeta \\ f \end{pmatrix} = \begin{pmatrix} A_{s,t} & B_{s,t} \\ C_{s,t} & I + D_{s,t} \end{pmatrix} \begin{pmatrix} \zeta \\ f \end{pmatrix}, \tag{5.2}$$

where $A_{s,t}, B_{s,t}, C_{s,t}$, and $D_{s,t}$ are operators from \mathcal{H}_0 into $\mathcal{H}_0, \mathcal{H}$ into $\mathcal{H}_0, \mathcal{H}_0$ into \mathcal{H} and \mathcal{H} into \mathcal{H}_0 respectively determined by the equations

$$\left. \begin{aligned} A_{s,t} &= A_{s-t}; & B_{s,t}f &= -(2H)^{1/2} \int \chi_{[t,s]}(x)A_{s-x}f(x)dx; \\ C_{s,t}\zeta(x) &= \chi_{[t,s]}(x)(2H)^{1/2}A_{x-t}\zeta; \\ (D_{s,t}f)(x) &= -2 \int \chi_{[t,s]}(x)\chi_{[y,s]}(x)HA_{x-y}f(y)dy. \end{aligned} \right\} \tag{5.3}$$

Lemma 5.1. For $r \geq s \geq t$,

$$U_{r,s}U_{s,t} = U_{r,t}. \tag{5.4}$$

Proof. It is immediate from the definitions (5.3) that $A_{r,s}A_{s,t} = A_{r,t}$ and $B_{r,s}C_{s,t} = B_{r,t}$, $D_{s,t} = D_{r,t}$, $C_{s,t} = 0$. Since $H^{1/2}$ and $A_{s,t}$ commute for all $s \geq t$ we have from a routine computation

$$\begin{aligned} A_{r,s}B_{s,t} + B_{r,s} + B_{r,t}C_{s,t} &= B_{r,t}, \\ C_{r,s}A_{s,t} + C_{s,t} + D_{r,t}C_{s,t} &= C_{r,t}. \end{aligned}$$

We shall now complete the proof by showing that

$$C_{r,s}B_{s,t} + D_{r,s} + D_{s,t} + D_{r,s}D_{s,t} = D_{r,t}. \tag{5.5}$$

To this end we observe that the last term on the left hand side of the above equation is 0 and the kernels of the first three operator terms are given by

$$\begin{aligned} K_1(x, y) &= -2\chi_{[s,r]}(x)\chi_{[s,t]}(y)HA_{x-y}, \\ K_2(x, y) &= -2\chi_{[s,r]}(y)\chi_{[y,r]}(x)HA_{x-y}, \\ K_3(x, y) &= -2\chi_{[t,s]}(y)\chi_{[y,s]}(x)HA_{x-y}, \end{aligned}$$

respectively. A decomposition of the triangular region $\{(x, y) : t < y \leq x < r\}$ into a square and two triangular regions shows that

$$\chi_{[t,r]}(y)\chi_{[y,r]}(x) = \chi_{[s,r]}(x)\chi_{[s,t]}(y) + \chi_{[s,r]}(y)\chi_{[y,r]}(x) + \chi_{[t,s]}(y)\chi_{[y,s]}(x),$$

and hence

$$\sum_{i=1}^3 K_i(x, y) = -2\chi_{[t,r]}(y)\chi_{[y,r]}(x)HA_{x-y},$$

which is also the kernel of $D_{r,t}$. This proves (5.4). \square

Lemma 5.2. *For all $r \geq s \geq t$*

$$\begin{pmatrix} I & 0 \\ 0 & S_r \end{pmatrix} U_{s,t} \begin{pmatrix} I & 0 \\ 0 & S_r^{-1} \end{pmatrix} = U_{r+s,r+t}, \tag{5.6}$$

$$\begin{pmatrix} I & 0 \\ 0 & S \end{pmatrix} U_{s,t}^* \begin{pmatrix} I & 0 \\ 0 & S^{-1} \end{pmatrix} = U_{-t,-s}. \tag{5.7}$$

Proof. To prove (5.6) we first observe that $A_{s,t} = A_{s-t}$ and for any $f \in \mathcal{L}$ we have by a change of variables

$$\begin{aligned} B_{s,t}S_r^{-1}f &= -(2H)^{1/2} \int \chi_{[t,s]}(x)A_{s-x}f(x+r)dx \\ &= -(2H)^{1/2} \int \chi_{[r+t,r+s]}(x)A_{r+s-x}f(x)dx \\ &= B_{r+s,r+t}f. \end{aligned} \tag{5.8}$$

Similarly we have

$$S_rC_{s,t} = C_{r+s,r+t}. \tag{5.9}$$

Further for any $x \in \mathbb{R}$ we have

$$\begin{aligned} [S_rD_{s,t}S_r^{-1}f](x) &= (D_{s,t}S_r^{-1}f)(x-r) \\ &= -2 \int \chi_{[t,s]}(y)\chi_{[y,s]}(x-r)HA_{x-r-y}f(y+r)dy \\ &= -2 \int \chi_{[t,s]}(y-r)\chi_{[y-r,s]}(x-r)HA_{x-y}f(y)dy \\ &= -2 \int \chi_{[r+t,r+s]}(y)\chi_{[y,s+r]}(x)HA_{x-y}f(y)dy \\ &= [D_{r+s,r+t}f](x). \end{aligned} \tag{5.10}$$

Now (5.6) follows from (5.8)–(5.10).

To prove (5.7) we have to only observe that the adjoints of the operators $A_{s,t}$, $B_{s,t}$, $C_{s,t}$, and $D_{s,t}$ are given by

$$\left. \begin{aligned} A_{s,t}^* &= A_{s,t}; & (B_{s,t}^* \zeta)(x) &= -\chi_{[t,s]}(x)(2H)^{1/2} A_{s-x} \zeta; \\ C_{s,t}^* f &= (2H)^{1/2} \int \chi_{[t,s]}(x) A_{x-t} f(x) dx; \\ (D_{s,t}^* f)(x) &= -2 \int \chi_{[t,s]}(y) \chi_{[x,s]}(y) H A_{y-x} f(y) dy, \end{aligned} \right\} \quad (5.11)$$

respectively and follow a procedure similar to the proof of (5.6). \square

Lemma 5.3. $U_{s,t}$ is unitary.

Proof. Since $H^{1/2}$ and A_t commute we have for $\zeta \in \mathcal{H}_0$, using (5.11)

$$\begin{aligned} & (A_{s,t}^* A_{s,t} + C_{s,t}^* C_{s,t}) \zeta \\ &= A_{2(s-t)} \zeta + 2H \int \chi_{[t,s]}(x) A_{2(x-t)} \zeta dx \\ &= A_{2(s-t)} \zeta - \int \chi_{[t,s]}(x) \frac{d}{dx} A_{2(x-t)} \zeta dx = \zeta. \end{aligned} \quad (5.12)$$

Also, for any $f \in \mathcal{H}$,

$$\begin{aligned} & (A_{s,t}^* B_{s,t} + C_{s,t}^* + C_{s,t}^* D_{s,t}) f \\ &= -A_{(s-t)} (2H)^{1/2} \int \chi_{[t,s]}(x) A_{s-x} f(x) dx + (2H)^{1/2} \int \chi_{[t,s]}(x) A_{x-t} f(x) dx \\ &\quad - 2(2H)^{1/2} \int \chi_{[t,s]}(x) A_{x-t} \int \chi_{[t,s]}(y) \chi_{[y,s]}(x) H A_{x-y} f(y) dy dx \\ &= (2H)^{1/2} \{ \int \chi_{[t,s]}(x) [-A_{2s-t-x} + A_{x-t}] f(x) dx \\ &\quad - \int \chi_{[t,s]}(y) \{ \int \chi_{[y,s]}(x) 2H A_{2x-t-y} f(y) dx \} dy \}. \end{aligned}$$

The right hand side of the last equation is identically zero in view of the identity $2H A_{2x-t-y} = -\frac{d}{dx} A_{2x-t-y}$. Thus

$$A_{s,t}^* B_{s,t} + C_{s,t}^* + C_{s,t}^* D_{s,t} = 0. \quad (5.13)$$

Next we observe that $D_{s,t}^* D_{s,t}$ is an integral operator in $\mathcal{H} = L^2(\mathbb{R}, \mathcal{H}_0)$ with kernel

$$\begin{aligned} K(x, y) &= 4 \int \chi_{[t,s]}(x) \chi_{[x,s]}(z) H A_{z-x} \chi_{[t,s]}(y) \chi_{[y,s]}(z) H A_{t-y} dz \\ &= 2 \chi_{[t,s]}(x) \chi_{[t,s]}(y) H [A_{|x-y|} - A_{2s-x-y}]. \end{aligned} \quad (5.14)$$

The kernel of the integral operator $B_{s,t}^* B_{s,t}$ is

$$(2H)^{1/2} \chi_{[t,s]}(x) A_{s-x} (2H)^{1/2} \chi_{[t,s]}(y) A_{s-y} = 2 \chi_{[t,s]}(x) \chi_{[t,s]}(x) \chi_{[t,s]}(y) H A_{2s-x-y}. \quad (5.15)$$

The kernel of $D_{s,t}^* + D_{s,t}$ is

$$\begin{aligned} & 2 \chi_{[t,s]}(x) \chi_{[x,s]}(y) H A_{y-x} + 2 \chi_{[t,s]}(y) \chi_{[y,s]}(x) A_{x-y} \\ &= 2 \chi_{[t,s]}(x) \chi_{[t,s]}(y) H A_{|x-y|}. \end{aligned} \quad (5.16)$$

From (5.14)–(5.16) we conclude that

$$B_{s,t}^* B_{s,t} + D_{s,t}^* + D_{s,t} + D_{s,t}^* D_{s,t} = 0. \quad (5.17)$$

Equations (5.12), (5.13), and (5.17) imply that $U_{s,t}^*U_{s,t} = I$. We use (5.7) to observe that

$$\begin{aligned} U_{s,t}U_{s,t}^* &= \begin{pmatrix} I & 0 \\ 0 & S \end{pmatrix} U_{-t,-s}^* \begin{pmatrix} I & 0 \\ 0 & S^{-1} \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & S^{-1} \end{pmatrix} U_{-t,-s} \begin{pmatrix} I & 0 \\ 0 & S \end{pmatrix} \\ &= \begin{pmatrix} I & 0 \\ 0 & S \end{pmatrix} U_{-t,-s}^* U_{-t,-s} \begin{pmatrix} I & 0 \\ 0 & S \end{pmatrix} = I. \quad \square \end{aligned}$$

Lemma 5.4. *In the decomposition $\mathfrak{h}_1 = (\mathfrak{h}_0 \oplus \mathfrak{h}_{]t,s]}) \oplus \mathfrak{h}_{]t,s]}^\perp$ where $\mathfrak{h}_{]t,s]}^\perp$ is the orthogonal complement of $\mathfrak{h}_{]t,s]}$ in \mathfrak{h} , the operator $U_{s,t}$ assumes the form*

$$U_{s,t} = V_{s,t} \oplus I,$$

where $V_{s,t}$ is a unitary operator in $\mathfrak{h}_0 \oplus \mathfrak{h}_{]t,s]}$.

Proof. Let E be the projector onto $\mathfrak{h}_{]s,t]}$ in \mathfrak{h} . Since

$$B_{s,t}E = B_{s,t}; \quad EC_{s,t} = C_{s,t}; \quad ED_{s,t} = D_{s,t}E = D_{s,t},$$

we can write

$$U_{s,t} = \begin{pmatrix} A_{s,t} & B_{s,t} \\ C_{s,t} & I + D_{s,t} \end{pmatrix} = \begin{pmatrix} I & 0 \\ 0 & E \end{pmatrix} \begin{pmatrix} A_{s,t} & B_{s,t} \\ C_{s,t} & I + D_{s,t} \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & E \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & I - E \end{pmatrix}$$

and hence the lemma is proved. \square

Theorem 5.5. *Let H be a positive self-adjoint operator in a Hilbert space \mathfrak{h}_0 . Let $\mathfrak{h} = L_2(\mathbb{R}, \mathfrak{h}_0)$, $\mathfrak{h}_1 = \mathfrak{h}_0 \oplus \mathfrak{h}$ and let $J: \zeta \rightarrow \begin{pmatrix} \zeta \\ 0 \end{pmatrix}$ be the natural imbedding of \mathfrak{h}_0 in \mathfrak{h}_1 .*

*Then there exists a strongly continuous unitary evolution $\{U_{s,t}, s \geq t\}$ in \mathfrak{h} satisfying the following: (a) $J^*U_{s,t}J = e^{(t-s)H}$; (b) in the decomposition*

$$\mathfrak{h}_1 = (\mathfrak{h}_0 \oplus \mathfrak{h}_{]t,s]}) \oplus \mathfrak{h}_{]t,s]}^\perp$$

where $\mathfrak{h}_{]t,s]}$ is the subspace of \mathfrak{h} consisting of functions with support $]t, s]$ and $\mathfrak{h}_{]t,s]}^\perp$ is the orthogonal complement of $\mathfrak{h}_{]t,s]}$ in \mathfrak{h} , $U_{s,t} = V_{s,t} \oplus I$ where $V_{s,t}$ is a unitary operator in $\mathfrak{h}_0 \oplus \mathfrak{h}_{]t,s]}$; (c) $(I \oplus S_r)U_{s,t}(I \oplus S_r)^{-1} = U_{r+s, r+t}$ where $I \oplus S_r$ is the direct sum of the identity in \mathfrak{h}_0 and S_r is the shift defined by (5.1) in \mathfrak{h} ; (d) $(I \oplus S)U_{s,t}^(I \oplus S)^{-1} = U_{-t, -s}$ where $I \oplus S$ is the direct sum of the identity in \mathfrak{h}_0 and S , the reflection defined by (5.1) in \mathfrak{h} .*

Proof. First of all we observe that we may assume without loss of generality that the operator H is bounded. For if not, we may write \mathfrak{h}_0 as a direct sum $\bigoplus_{n=1}^\infty \mathfrak{h}^{(n)}$ such that H is bounded in each $\mathfrak{h}^{(n)}$, construct the evolutions $U_{s,t}^{(n)}$ in $\mathfrak{h}_0^{(n)} \oplus L_2(\mathbb{R}, \mathfrak{h}_0^{(n)})$ and put the pieces together without any difficulty. When H is bounded we construct $U_{s,t}$ by putting $A_{s,t} = e^{(t-s)H}$ and using (5.2) and (5.3). Then we conclude from Lemmas 5.1–5.4 that $U_{s,t}$ is a unitary evolution satisfying properties (1)–(4) of the theorem. Thus only strong continuity remains to be established. Because of unitarity this is equivalent to establishing the weak continuity of $U_{s,t}$ and hence

the weak continuity of the operators $A_{s,t}$, $B_{s,t}$, $C_{s,t}$ and $D_{s,t}$ defined by (5.3). By assumption $A_{s,t}$ is strongly, hence weakly continuous. If $\zeta \in \mathcal{H}_0$, $f, g \in \mathcal{H}$ we have

$$\begin{aligned} \langle \zeta, B_{s,t} f \rangle &= -2^{1/2} \int \chi_{[t,s]}(x) \langle \zeta, H^{1/2} A_{s-x} f(x) \rangle dx \\ &= -2^{1/2} \int_t^s \langle \zeta, H^{1/2} A_{s-x} f(x) \rangle dx. \end{aligned}$$

Since the right hand side is continuous in the variables t and s , the weak continuity of $B_{s,t}$ follows. A similar argument holds for the remaining operators. \square

In view of the properties (1)–(4) of the evolution $\{U_{s,t}, s \geq t\}$ we call it a *time-orthogonal unitary dilation* of the semigroup $\{e^{-tH}, t \geq 0\}$, and we note that it is reflective and covariant under time inversion and translations. It may be remarked that the dilation of Theorem 5.5 is different from the classical dilations of one parameter contraction semigroups of operators in a Hilbert space into one parameter groups of unitary operators in a larger Hilbert space (see [14, 4, 10]). The replacement of invariance under time translation by covariance and the time orthogonality property in Theorem 5.5(b) are the essential features of the new construction.

6. Second Quantisation

Let \mathcal{H} be a Hilbert space. The (*symmetric*) *Fock space* [6] over \mathcal{H} is the direct sum

$$\Gamma(\mathcal{H}) = \bigoplus_{n=0}^{\infty} \mathcal{H}^{(n)},$$

where $\mathcal{H}^{(0)} = \mathbb{C}$ and for $n \geq 1$ the *n-particle space* $\mathcal{H}^{(n)}$ is the symmetric component of the *n-fold tensor product* of \mathcal{H} with itself. The *exponential vectors*, [6]

$$e(f) = (1, f, (2!)^{-1/2} f \otimes f, (3!)^{-1/2} f \otimes f \otimes f, \dots), \quad f \in \mathcal{H}$$

form a total set in $\Gamma(\mathcal{H})$; the *vacuum*, denoted by Ω or $\Omega(\mathcal{H})$, is the vector

$$e(0) = (1, 0, 0, \dots).$$

If $\mathcal{H} = \mathcal{H}_1 \oplus \dots \oplus \mathcal{H}_n$ is a direct sum then there is a Hilbert space isomorphism, which we use to identify the two spaces, from $\Gamma(\mathcal{H})$ to $\Gamma(\mathcal{H}_1) \otimes \dots \otimes \Gamma(\mathcal{H}_n)$ under which each exponential vector $e(f_1, \dots, f_n)$ is mapped to the product vector $e(f_1) \otimes \dots \otimes e(f_n)$.

Let C be a contraction from \mathcal{H} to a second Hilbert space \mathcal{H}' . The second quantisation $\Gamma(C)$ [5] is a contraction from $\Gamma(\mathcal{H})$ to $\Gamma(\mathcal{H}')$ whose action on exponential vectors is

$$\Gamma(C)e(f) = e(Cf). \tag{6.1}$$

The map $C \mapsto \Gamma(C)$ is continuous for the strong operator topologies in $\mathbf{B}(\mathcal{H}, \mathcal{H}')$ and $\mathbf{B}(\Gamma(\mathcal{H}), \Gamma(\mathcal{H}'))$. Second quantisation satisfies the functorial rules

$$\Gamma(I) = I, \tag{6.2}$$

$$\Gamma(C_1 C_2) = \Gamma(C_1) \Gamma(C_2), \tag{6.3}$$

$$\Gamma(C^*) = \Gamma(C)^*. \tag{6.4}$$

From (6.1) it is clear that, if $C = C_1 \oplus \dots \oplus C_n$ is a direct sum, then

$$\Gamma(C) = \Gamma(C_1) \otimes \dots \otimes \Gamma(C_n). \tag{6.5}$$

Let J be the inclusion injection from \mathfrak{h} into a larger space \mathfrak{h}' . Then, for $A \in \mathbf{B}(\Gamma(\mathfrak{h}'))$, $\Gamma(J)^* A \Gamma(J)$ is the vacuum conditional expectation of A given $\mathbf{B}(\Gamma(\mathfrak{h}))$, that is, the unique operator in $\mathbf{B}(\Gamma(\mathfrak{h}))$ such that, for arbitrary $\psi, \chi \in \Gamma(\mathfrak{h})$,

$$\langle \psi, \Gamma(J)^* A \Gamma(J) \chi \rangle = \langle \psi \otimes \Omega(\mathfrak{h}^\perp), A \chi \otimes \Omega(\mathfrak{h}^\perp) \rangle,$$

where \mathfrak{h}^\perp is the orthogonal complement of \mathfrak{h} in \mathfrak{h}' and we have used the identifications

$$\mathfrak{h}' = \mathfrak{h} \oplus \mathfrak{h}^\perp, \quad \Gamma(\mathfrak{h}') = \Gamma(\mathfrak{h}) \otimes \Gamma(\mathfrak{h}^\perp).$$

The Fock representation of the canonical commutation relations (RCCR) over \mathfrak{h} is the family of unitary Weyl operators $(W(f) : f \in \mathfrak{h})$ in $\Gamma(\mathfrak{h})$, whose actions on exponential vectors are

$$W(f)e(g) = \exp(-\frac{1}{4}\|f\|^2 + 2^{-1/2}i\langle f, g \rangle)e(g + 2^{-1/2}if).$$

The map $f \mapsto W(f)$ is continuous in the norm topology of \mathfrak{h} and the strong operator topology of $\mathbf{B}(\Gamma(\mathfrak{h}))$, the Weyl relation

$$W(f)W(g) = \exp\{-\frac{1}{2}i \operatorname{Im}\langle f, g \rangle\}W(f+g) \quad (f, g \in \mathfrak{h}) \tag{6.6}$$

holds and the vacuum expectation functional is

$$F(f) = \langle \Omega, W(f)\Omega \rangle = \exp(-\frac{1}{4}\|f\|^2).$$

The Fock representation is irreducible [6]. If U is a unitary operator in \mathfrak{h} then, for $f \in \mathfrak{h}$,

$$\Gamma(U)W(f)\Gamma(U)^{-1} = W(Uf). \tag{6.7}$$

If \mathfrak{h} is a direct sum, $\mathfrak{h} = \mathfrak{h}_1 \oplus \dots \oplus \mathfrak{h}_n$, then, for $f = (f_1, \dots, f_n) \in \mathfrak{h}$

$$W(f) = W(f_1) \otimes \dots \otimes W(f_n).$$

Now let \mathfrak{h}_0 be a Hilbert space, let $\mathfrak{h} = L^2(\mathbb{R}, \mathfrak{h}_0)$ be the Hilbert space of \mathfrak{h}_0 -valued square-integrable measurable vector-valued functions and let $\mathfrak{h}_1 = \mathfrak{h}_0 \oplus \mathfrak{h}$. Let N be the von Neumann algebra $\mathbf{B}(\Gamma(\mathfrak{h}_1))$ and, for $s \geq t$, let $N_{s,t}$ be the von Neumann algebra generated by the operators $W(\zeta, f)$, $\zeta \in \mathfrak{h}_0, f \in \mathfrak{h}_{[s,t]}$, where $\mathfrak{h}_{[s,t]}$ is the subspace of \mathfrak{h} comprising functions vanishing outside $]t, s]$.

Next let $S_t, t \geq 0$ and S be the shift operators and reflection in \mathfrak{h} and define automorphisms $\gamma_t, t \geq 0$, and ϱ of N by

$$\begin{aligned} \gamma_t(A) &= \Gamma(I \oplus S_t) A \Gamma(I \oplus S_t)^{-1} \\ \varrho(A) &= \Gamma(I \oplus S) A \Gamma(I \oplus S)^{-1}. \end{aligned}$$

Then using (6.5), (6.7) and the fact that $\mathbf{B}(\Gamma(\mathfrak{h}_1))$ is generated by the irreducible family of Weyl operators, it can be seen that $(N, (N_{s,t}), (\gamma_t), \varrho)$ is a reflective covariant system.

Next let J be the embedding $\zeta \rightarrow (\zeta, 0)$ of \mathfrak{h}_0 in \mathfrak{h}_1 and define $j : N \rightarrow \mathbf{B}(\Gamma(\mathfrak{h}_0))$ by

$$j(A) = \Gamma(J^*) A \Gamma(J),$$

thus $j(A)$ is the conditional expectation of A given $\mathbf{B}(\mathcal{H}_0)$. Finally let $(U_{s,t})$ be the time orthogonal unitary dilation of a strongly continuous self-adjoint contraction semigroup acting in \mathcal{H}_0 as in Theorem 5.5.

Theorem 6.1. *j is a reflective covariant reducing map and $(\Gamma(U_{s,t}), s \geq t)$ is a reflectively covariantly adapted evolution for $(N, (N_{s,t}), (\gamma_t), \varrho)$.*

Proof. That $j(A^*) = j(A)^*$ is clear from (6.4). To prove that j satisfies (3.1) we make the natural identifications

$$\begin{aligned} \mathcal{H} &= \mathcal{H}_{[a,b]} \oplus \mathcal{H}_{[c,d]} \oplus (\mathcal{H}_{[a,b]} \oplus \mathcal{H}_{[c,d]})^\perp \\ \mathcal{H}_1 &= \mathcal{H}_0 \oplus \mathcal{H}_{[a,b]} \oplus \mathcal{H}_{[c,d]} \oplus (\mathcal{H}_{[a,b]} \oplus \mathcal{H}_{[c,d]})^\perp \\ N_{a,b} &= \{W(\zeta) \otimes W(f_{[a,b]}) \otimes I \otimes I : \zeta \in \mathcal{H}_0, f \in \mathcal{H}_{[a,b]}\}'' \\ &= \mathbf{B}(\Gamma(\mathcal{H}_0) \otimes \Gamma(\mathcal{H}_{[a,b]})) \otimes I \otimes I \\ &= \mathbf{B}(\Gamma(\mathcal{H}_0)) \otimes \mathbf{B}(\Gamma(\mathcal{H}_{[a,b]})) \otimes I \otimes I, \end{aligned}$$

and similarly

$$N_{c,d} = \mathbf{B}(\Gamma(\mathcal{H}_0)) \otimes I \otimes \mathbf{B}(\Gamma(\mathcal{H}_{[c,d]})) \otimes I.$$

Since the von Neumann algebras $\mathbf{B}(\Gamma(\mathcal{H}_{[a,b]})) \otimes I \otimes I$ and $I \otimes \mathbf{B}(\Gamma(\mathcal{H}_{[c,d]})) \otimes I$ are independent in the vacuum state

$$\Omega(\mathcal{H}) = \Omega(\mathcal{H}_{[a,b]}) \otimes \Omega(\mathcal{H}_{[c,d]}) \otimes \Omega(\mathcal{H}_{[a,b]} \otimes \mathcal{H}_{[c,d]})^\perp,$$

(3.1) now follows from Theorem 1.1. Thus j is a reducing map. By second quantising the relations

$$(I \oplus S_t)^{-1} J = (I \oplus S)^{-1} J = J$$

and using (6.3), we have

$$\Gamma(I \oplus S_t)^{-1} \Gamma(J) = \Gamma(I \oplus S)^{-1} \Gamma(J) = \Gamma(J), \tag{6.8}$$

and hence, taking adjoints,

$$\Gamma(J)^* \Gamma(I \oplus S_t) = \Gamma(J)^* \Gamma(I \oplus S) = \Gamma(J)^*. \tag{6.9}$$

From (6.8) and (6.9) it is clear that $j \circ \gamma_r = j \circ \varrho = j$ so that j is reflective covariant.

Finally, it is clear from (6.2), (6.3) and the continuity of second quantisation that $(\Gamma(U_{s,t}), s \geq t)$ is an evolution; second quantising (5.2) using (6.5) gives

$$\Gamma(U_{s,t}) = \Gamma(V_{s,t}) \otimes I \in \mathbf{B}(\Gamma(\mathcal{H}_0 \oplus \mathcal{H}_{[r,s]})) \otimes I = N_{s,t},$$

so the evolution is adapted. That the adaptation is reflective and covariant follows from the second quantisations of (5.3) and (5.4). \square

The reduced evolution corresponding to the reducing map j of the evolution $(\Gamma(U_{s,t}))$ is easily found by second quantising the relation $A_{s-t} = J^* U_{s,t} J$ to be the second quantisation $(\Gamma(A_t) : t \geq 0)$ of the original semigroup (A_t) . We write

$$\Gamma(A_t) = e^{-tH}.$$

Let us now consider perturbations.

It is now clear that elements B of $\bigcap_{s>t} N_{s,t}$ are precisely those of form $V \otimes I$ with $V \in \mathbf{B}(\mathscr{H}_0)$, and such elements are automatically invariant under the γ_t . From Theorem 3.2(c) we may thus write, for self-adjoint $V \in \mathbf{B}(\mathscr{H}_0)$,

$$\begin{aligned} e^{-(s-t)(H+V)} &= j(M_{s,t}^V \Gamma(U_{s,t})) \\ &= \mathbb{E}[M_{s,t}^V \Gamma(U_{s,t}) | \mathbf{B}(\Gamma(\mathscr{H}_0))], \end{aligned} \tag{6.10}$$

where $(M_{s,t}^V)$ is the cocycle for $\Gamma(U_{s,t})$ determined by $B = V \otimes I$.

We investigate conditions under which the family of operators

$$\Gamma(U_{s,t}) V \otimes I \Gamma(U_{s,t})^{-1}$$

is commutative, so that, according to Theorem 2.1, the cocycle can be expressed in integral form

$$M_{s,t}^V = \exp \left\{ - \int_t^s \Gamma(U_{s,\tau}) V \otimes I \Gamma(U_{s,\tau})^{-1} d\tau \right\}. \tag{6.11}$$

Let K_0 be a conjugation in \mathscr{H}_0 , that is a conjugate linear isometric map from \mathscr{H}_0 to itself whose square is the identity, such that

$$K_0 A_t = A_t K_0 \quad (t \geq 0).$$

Then if K is the conjugation in $\mathscr{H} = L^2(\mathbb{R}, \mathscr{H}_0)$ given by

$$(Kf)(x) = K_0(f(x))$$

and $K_1 = K_0 \oplus K$, we have

$$K_1 J = J K_0$$

and, for arbitrary $s \geq t$,

$$K_1 U_{s,t} = U_{s,t} K_1.$$

Let N_0, N_1 be the von Neumann algebras generated by the operators

$$\{W(\zeta) : \zeta \in \mathscr{H}_0, K_0 \zeta = \zeta\}, \quad \{W(f) : f \in \mathscr{H}_1, K_1 f = f\}$$

respectively. Then, using (6.6), it is easily seen that N_0 and N_1 are commutative. Also, since by (6.7),

$$\begin{aligned} \Gamma(U_{s,t}) W(\zeta) \otimes I \Gamma(U_{s,t})^{-1} &= \Gamma(U_{s,t}) W(J\zeta) \Gamma(U_{s,t})^{-1} \\ &= \Gamma(U_{s,t} J\zeta), \end{aligned}$$

and if $K_0 \zeta = \zeta$, then

$$K_1 U_{s,t} J\zeta = U_{s,t} K_1 J\zeta = U_{s,t} J K_0 \zeta = U_{s,t} J\zeta,$$

we see that conjugation by $\Gamma(U_{s,t})$ maps $N_0 \otimes I$ into the commutative algebra N_1 . It follows that, for V belonging to N_0 , the operators $\Gamma(U_{s,t}) V \otimes I \Gamma(U_{s,t})^{-1}$ commute so that by Theorem 2.1 the cocycle can be written in integral form so that (6.10) becomes

$$e^{-(s-t)(H+V)} = \mathbb{E} \left[\exp \left\{ - \int_t^s \Gamma(U_{s,\tau}) V \otimes I \Gamma(U_{s,\tau})^{-1} d\tau \right\} \Gamma(U_{s,t}) | \mathbf{B}(\Gamma(\mathscr{H})) \right]. \tag{6.12}$$

A conjugation K_0 in \mathcal{H}_0 effects a separation of the set of fields $r(\zeta)$, $\zeta \in \mathcal{H}_0$, that is the self-adjoint operators defined by

$$W(t\zeta) = e^{itr(\zeta)}, \quad \zeta \in \mathcal{H}_0, \quad t \in \mathbb{R}$$

into ‘‘coordinate’’ fields, for which $K_0\zeta = \zeta$ and ‘‘momentum’’ fields, for which $K_0\zeta = -\zeta$. Thus we may say that the integral form (6.12) of Feynman-Kac formula is valid provided the perturbation V is a function of the coordinate fields alone, in the sense of being an element of the von Neumann algebra generated by the spectral projectors of these fields alone.

7. An Example

We consider the case [8] of the semigroup $A_t = e^{-t}$ acting in the Hilbert space $\mathcal{H}_0 = \mathbb{C}$, with complex conjugation as the conjugation K_0 . The second quantisation $(\Gamma(A_t))$ is then the semigroup in $\Gamma(\mathbb{C}) = l^2$ whose infinitesimal generator is the number operator

$$a^\dagger a(z_0, z_1, z_2, \dots) = (0, z_1, 2z_2, \dots),$$

or equivalently the renormalised harmonic oscillator Hamiltonian $\frac{1}{2}(p^2 + q^2 - 1)$, where p and q are the canonical pair defined by

$$e^{ip} = W(it), \quad e^{iq} = W(t), \quad t \in \mathbb{R}$$

and

$$a = 2^{-1/2}(q - ip), \quad a^\dagger = 2^{-1/2}(q + ip).$$

We shall show that the operator $\Gamma(U_{s,t})$ can be expressed formally as a product integral

$$\Gamma(U_{s,t}) = \prod_t^s e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)}, \tag{7.1}$$

where (P, Q) is essentially the (Fock) canonical Wiener process of [3] of minimal variance $\sigma^2 = 1$. We recall that for $t \geq 0$ the operators $P(t), Q(t)$ of the process are defined in $\Gamma(L^2(\mathbb{R}))$ by

$$e^{ixP(t)} = W(ix\chi_{[0,t]}), \quad e^{ixQ(t)} = W(x\chi_{[0,t]}) \quad (x \in \mathbb{R});$$

for $t < 0$ we define $P(t), Q(t)$ by

$$e^{ixP(t)} = W(-ix\chi_{[t,0]}), \quad e^{ixQ(t)} = W(-x\chi_{[t,0]}),$$

so that, for arbitrary $s > t$,

$$P(s) - P(t) = (s - t)^{1/2} p_{s,t}, \quad Q(s) - Q(t) = (s - t)^{1/2} q_{s,t},$$

where $(p_{s,t}, q_{s,t})$ is the canonical pair defined by

$$e^{ixp_{s,t}} = W(ix(s - t)^{-1/2}\chi_{[t,s]}), \quad e^{ixq_{s,t}} = W(x(s - t)^{-1/2}\chi_{[t,s]}) \quad (x \in \mathbb{R}).$$

To give meaning to the product integral (7.1) we first consider the formal “Riemann product” for a partition $\mathcal{A} = \{t = t_0 < t_1 < \dots < t_n = s\}$,

$$\begin{aligned} T_{\mathcal{A}} &= \prod_{j=1}^n \exp \{i \sqrt{2} (p \otimes (Q(t_j) - Q(t_{j-1})) - q \otimes (P(t_j) - P(t_{j-1})))\} \\ &= \prod_{j=1}^n \exp \{i \theta_j (p \otimes q_j - q \otimes p_j)\}, \end{aligned}$$

where

$$\theta_j = (2(t_j - t_{j-1}))^{1/2}$$

and (p_j, q_j) is the canonical pair

$$p_j = p_{t_{j-1}, t_j}; \quad q_j = q_{t_{j-1}, t_j}.$$

Recognising $p \otimes q_j - q \otimes p_j$ as an angular momentum, we may define $\exp \{i \theta_j (p \otimes q_j - q \otimes p_j)\}$ rigorously as the second quantisation $\Gamma(R_j)$ of the rotation R_j which acts in

$$\mathbb{C} \oplus L^2(\mathbb{R}) = \mathbb{C} \oplus \mathbb{C} \chi_j \oplus (\mathbb{C} \chi_j)^\perp,$$

where

$$\chi_j = (t_j - t_{j-1})^{-1/2} \chi_{[t_{j-1}, t_j]},$$

as

$$R_j(z, f) = (z_0, f_0),$$

where

$$\begin{aligned} z_0 &= \cos \theta_j z - \sin \theta_j \langle \chi_j, f \rangle \\ f_0 &= \sin \theta_j z \chi_j + (\cos \theta_j - 1) \langle \chi_j, f \rangle \chi_j + f. \end{aligned}$$

Hence

$$\prod_{j=1}^n R_j(z, f) = (z_1, f_1),$$

where

$$\begin{aligned} z_1 &= \left(\prod_{j=1}^n \cos \theta_j \right) z - \sum_{j=1}^n \langle \chi_j, f \rangle \sin \theta_j \prod_{k=j+1}^n (\cos \theta_k) \\ f_1 &= \sum_{j=1}^n \sin \theta_j \left(\prod_{k=1}^{j-1} f_k \sin \theta_k - \sum_{k=1}^{j-1} f_k \sin \theta_k \prod_{l=k+1}^{j-1} (\cos \theta_l) \right) \chi_j \\ &\quad + \sum_{j=1}^n (\cos \theta_j - 1) f_j \chi_j + f. \end{aligned}$$

From this it can be seen that, for sufficiently smooth f ,

$$\lim_{\mathcal{A}} \prod_{j=1}^n R_j(z, f) = (z', f'),$$

where

$$z' = e^{-(s-t)}z - \sqrt{2} \int \chi_{[t, s]}(x) e^{-(s-x)} f(x) dx,$$

$$f'(x) = \sqrt{2} e^{-(x-t)} \chi_{[t, s]}(x) z - 2 \int \chi_{[t, s]}(y) \chi_{[y, s]}(x) e^{-(x-y)} f(y) dy,$$

and the limit is as the mesh $|A| = \max(t_j - t_{j-1})$ decreases to zero. But this is precisely $U_{s,t}(z, f)$. Since $U_{s,t}$ is unitary the product $\prod_{j=1}^n R_j$ converges strongly to $U_{s,t}$. Using the strong continuity of second quantisation it follows that $\Gamma(U_{s,t})$ is indeed given by the limit of the Riemann products $T_A = \prod_{j=1}^n \Gamma(R_j)$.

The algebra N_0 in this case is clearly the von Neumann algebra generated by the operators e^{ixq} , $x \in \mathbb{R}$, thus a self-adjoint element of N_0 takes the form $V(q)$ for some real, bounded measurable function V . Since, by (6.7)

$$\begin{aligned} \Gamma(U_{s,t}) e^{ixq} \otimes I \Gamma(U_{s,t})^{-1} &= \Gamma(U_{s,t}) W(x, 0) \Gamma(U_{s,t})^{-1} \\ &= W(U_{s,t}(x, 0)) \\ &= W(e^{-(s-t)}x, xf), \end{aligned}$$

where

$$f(\tau) = \sqrt{2} \chi_{[t, s]}(\tau) e^{-(\tau-t)}, \tag{7.2}$$

we have

$$\Gamma(U_{s,t}) V(q) \otimes I \Gamma(U_{s,t})^{-1} = V(q_t^s), \tag{7.3}$$

where q_t^s is defined formally as

$$q_t^s = e^{-(s-t)} q \otimes I + \sqrt{2} I \otimes \int_t^s e^{-(s-\tau)} dQ(\tau), \tag{7.4}$$

using the stochastic integral notation [3] for field operators associated with the canonical Wiener process; q_t^s is rigorously defined by

$$e^{ixq_t^s} = W(e^{-(s-t)}x, xf),$$

with f given by (7.2). Thus, from (6.12) we obtain the formula [7]

$$\begin{aligned} &e^{-(s-t)(\frac{1}{2}(p^2 + q^2 - 1) + V(q))} \\ &= \mathbb{E} \left[\exp \left\{ - \int_t^s V(e^{-(s-\tau)}) q \otimes I + \sqrt{2} I \otimes \int_t^s e^{-(s-\tau)} dQ(\tau') \right\} \prod_t^s e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)} | \mathbf{B}(I_2) \right]. \end{aligned} \tag{7.5}$$

As was brought to our attention by a comment of L. van Hemmen, the Feynman-Kac formula (7.5) is essentially the ‘‘oscillator process’’ formula of [13, p. 52]. To see this, note first that, because of their relationship (7.4) to the q -component of the canonical Wiener process, in the state $\Omega_0 \otimes \Omega(L^2(\mathbb{R}))$ where Ω_0 is the oscillator ground state, the vacuum in $\Gamma(\mathbb{C})$, the commuting family of operators $q(t) = q_{t,0}$, $t \geq 0$, constitute a realisation of the Ornstein-Uhlenbeck velocity process.

Moreover, for vectors in $l^2 = \Gamma(\mathbb{C})$ expressed in the form $f(q)\Omega_0, g(q)\Omega_0$ we have, from (7.4) with $]t, s]$ replaced by $]0, t]$,

$$\begin{aligned} &\langle f(q)\Omega_0, \exp\{-t(\frac{1}{2}(p^2 + q^2 - 1) + V(q))\}g(q)\Omega_0\rangle \\ &= \left\langle (f(q)\Omega_0) \otimes \Omega, \exp\left\{-\int_0^t V(q(\tau))d\tau\right\} \prod_0^t e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)}(g(q)\Omega_0) \otimes \Omega \right\rangle \\ &= \left\langle f(q \otimes I)\Omega_0 \otimes \Omega, \exp\left\{-\int_0^t V(q(\tau))d\tau\right\} \prod_0^t e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)}g(q \otimes I)\Omega_0 \otimes \Omega \right\rangle \\ &= \left\langle f(q(0)\Omega_0) \otimes \Omega, \exp\left\{-\int_0^t V(q(\tau))d\tau\right\} g(q(t)) \prod_0^t e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)}\Omega_0 \otimes \Omega \right\rangle, \end{aligned}$$

using (7.3),

$$= \left\langle f(q(0))\Omega_0 \otimes \Omega, \exp\left\{-\int_s^t V(q(\tau))d\tau\right\} g(q(t))\Omega_0 \otimes \Omega \right\rangle,$$

using the invariance of the vacuum $\Omega_0 \otimes \Omega$ under the second quantisation

$$\begin{aligned} \prod_0^t e^{i\sqrt{2}(p \otimes dQ - q \otimes dP)} &= \left\langle \Omega_0 \otimes \Omega, \overline{f(q(0))}g(q(t)) \exp\left\{-\int_0^t V(q(\tau))d\tau\right\} \Omega_0 \otimes \Omega \right\rangle \\ &= \mathbb{E} \left[\overline{f(q(0))}g(q(t)) \exp\left\{-\int_0^t V(q(\tau))d\tau\right\} \right], \end{aligned}$$

where the expectation is over the Ornstein-Uhlenbeck velocity process $(q(t))$. Thus, as in [13],

$$\begin{aligned} &\langle f(q)\Omega_0, \exp\{-t(\frac{1}{2}(p^2 + q^2 - 1) + V(q))\}g(q)\Omega_0\rangle \\ &= \mathbb{E} \left[\overline{f(q(0))}g(q(t)) \exp\left\{-\int_s^t V(q(\tau))d\tau\right\} \right]. \end{aligned}$$

From the viewpoint of noncommutative probability theory, it is natural to regard the operators

$$q(t) = e^{-t}q \otimes I + \sqrt{2}I \otimes \int_0^t e^{-\tau}dQ(\tau), \quad t \geq 0$$

as the q -component of the canonical Ornstein-Uhlenbeck velocity process $(p(t), q(t), t \geq 0)$, where

$$p(t) = e^{-t}p \otimes I + \sqrt{2}I \otimes \int_0^t e^{-\tau}dP(\tau), \quad t \geq 0.$$

Because $p(t)$ and $q(t)$ constitute a canonical pair for all $t \geq 0$, the process $((p(t), q(t)))$ is essentially a noncommutative W^* -stochastic process in the sense of [2], insofar as the von Neumann uniqueness theorem gives rise to a family of von Neumann algebra isomorphisms $j(t), t \geq 0$ from the von Neumann algebra generated by the spectral projectors of $p(t)$ and $q(t)$ to that generated by those of the fixed canonical pair (p, q) .

The perturbation cocycle $(M_{t,0}^B)$ corresponding to an element B of $\mathbf{B}(l^2)$ not in N_0 can be expressed formally as the continuous product

$$M_{t,0}^B = \prod_0^t \exp \{ -B(p(\tau), q(\tau)) d\tau \};$$

here $B(p(\tau), q(\tau))$ is the pre-image of B under $j(\tau)$.

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