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Mills, B orcid.org/0000-0002-9141-0931, Watson, AJ, Goldblatt, C et al. (2 more authors) (2011) Timing of Neoproterozoic glaciations linked to transport-limited global weathering. Nature Geoscience, 4 (12). pp. 861-864. ISSN 1752-0894

https://doi.org/10.1038/ngeo1305

Timing of Neoproterozoic glaciations explained by transport-limited global weathering

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Snowball Earth glaciations are thought to be terminated by massive concentrations of CO_2, accumulated from millions of years of volcanic degassing [1, 2]. After deglaciation, high temperature would increase the rate of silicate weathering, which increases the removal rate of CO_2 until a steady state is achieved [3]. It has recently been shown that the expected weathering rate on snowball exit could be around 10 times modern day, giving a timescale on the order of 10^6 years to restore equilibrium [4]. However, over long timescales the supply of fresh rock becomes a limiting factor for silicate weathering [5], constraining the maximum sustainable weathering flux. Here we show that this limitation could explain the pattern of greenhouse – icehouse oscillations during the Neoproterozoic, with long gaps between extreme glaciations. Using a modified version of the COPSE biogeochemical model [6], our estimated limit gives a stabilisation time of $>10^7$ years. With a simple ice-albedo feedback included, the model readily produces glacial-interglacial oscillations on this timescale, compatible with those observed in the geological record. Sustained above-average interglacial nutrient levels imply prolonged positive excursions in carbonate $\delta^{13}C$ [7], as well as an increased flux of oxygen to the atmosphere [8], broadly consistent with data.

The Neoproterozoic era (1000-542Ma) is punctuated by at least three glaciations [9], the severe low-latitude Sturtian and Marinoan episodes being proposed as examples of `Snowball Earth' events [1, 10]. Figure 1 displays Neoproterozoic carbonate carbon isotope data [7], which shows a quasi-periodic pattern. Negative excursions associated with glaciation appear at ~50 Myr intervals between long periods of positive fractionation. The long interval between glaciations poses a puzzle given the standard model of a snowball Earth being terminated by very high CO_2 and temperature. The time taken to restore equilibrium after such a perturbation depends on the rate of CO_2 drawdown via silicate weathering, a process that would be greatly enhanced in the aftermath of snowball Earth. Highly weatherable rock flour produced by glacial grinding would likely cover a large surface area, and increased temperature and runoff should allow for an elevated weathering flux. Linked GCM and kinetic weathering models have determined the maximum weathering rate in this climate to be on the order of 10 times the modern day flux, implying a timescale of around 10^6 years to reduce atmospheric CO_2 to pre-glacial levels [4]. Based on these results, we would expect the system to establish equilibrium in a time far shorter than the interglacial periods following the Sturtian and Marinoan glaciations.

Here we propose that the timescale for CO_2 drawdown following a snowball glaciation should be extended due to transport limitation of the silicate weathering process. In a transport limited regime, silicate cations are completely leached from fresh regolith and therefore the rate of chemical weathering depends only on the physical erosion rate [11]. Modern continental cratons are transport limited, as seen by plotting the rate of denudation of silicate cations against total denudation rate [5]. In such a regime, increasing temperature or runoff does not increase the rate of CO_2 drawdown, because all the available silicate cations are already being processed. As global temperature and humidity rises, we would expect more weathering zones to fall into transport limitation, implying a theoretical maximum silicate weathering rate, where every available cation is leached.

Over the Phanerozoic, the mean continental erosion rate is estimated to be ~16m Myr^{-1} [12]. Using the average density and area of the present day continents (area=1.5×10^{14} m^2, density = 2.5×10^3 kg m^{-3}) yields a total mass of 6×10^{12} kg yr^{-1}. Assuming a cation weight fraction of 0.08 [5], we estimate a global silicate weathering rate maximum for the Phanerozoic of around 4.8×10^{11} kg yr^{-1}. This maximum transport limited rate is about 2.4 times greater than present day weathering rate [13].

Determining the global erosion rate in the Neoproterozoic is difficult, because it depends on the continental area and rate of uplift. Current estimates for Neoproterozoic uplift rates are close to present day values [14], and the majority of studies agree that the total continental area was probably less than it is now. Proxies for global denudation show very low values (<10m Myr^{-1}) for the early Phanerozoic, but are likely to be affected by sampling artefacts [15]. The rate of volcanic degassing in the Neoproterozoic is also important, as it is the balance between CO_2 degassing and its maximum consumption rate via weathering and burial that dictates the system response time to large perturbations. In carbon cycle models, degassing is usually assumed to be proportional to the seafloor spreading rate. Accounting for different continental growth models, the Neoproterozoic outgassing rate was probably between 1 and 5 times the present day rate [14, 16]. But smaller crustal reservoirs of carbon in the Neoproterozoic [17] may have decreased the CO_2 content of volcanic gas by up to 20%.

In figure 2, we use a modified version of the COPSE biogeochemical model [6] (see supplementary information) to investigate the effect of a weathering rate cap on the time taken to return to steady state after the suggested snowball exit concentration of 0.3 atm CO_2 [1, 2] is imposed. Silicate weathering rate is fixed at a prescribed maximum value, W_{max}, which is defined relative to the present day rate. We find that choice of W_{max} has a strong effect on the system: Assuming the Phanerozoic average erosion rate (W_{max}=2.4) yields a stabilisation time of ~10^7 years, even for conservative estimates of the CO_2 degassing rate *D*. A lower erosion rate, and/or a higher degassing rate increases the stabilisation time greatly. For further model runs we let global weathering follow a simple kinetic equation as described by Berner [18], but asymptote to W_{max} as the kinetic weathering rate approaches the maximum transport limited value, placing a cap on global weathering rates. The choice of kinetic weathering function, and the nature of the transition to W_{max} has negligible effect on results as the rate remains at W_{max} until CO_2 is very close to the stable level.

An important consideration for this work is weathering of rock flour left on the surface after a snowball glaciation, which would be expected to increase weathering kinetics as in the quaternary glacial cycle [19]. Global weathering fluxes would not become limited by transport of fresh rock until the flour produced during the glaciation had been completely leached. Le Hir et al [4] assume a thin soil profile in their model of the snowball earth aftermath, due to evidence of persistent weathering during glaciation [20]. Following their estimate of a 25cm reactive upper layer, we derive a weatherable equivalent of ~10^{17} moles C (see supplementary information).

Figure 3 shows model sensitivity to the initial quantity of rock flour. Here we allow a global weathering rate of 10 times present day when rock flour is present [4], switching to the transport limited equation once a specified amount of carbon has been buried, analogous to the abundance of glacial flour. We find that a weatherable equivalent on the order of 10^{20} moles C is required to significantly affect stabilisation time; we use a increased reactive layer depth of 2.5m (10^{18} mol C equiv.) for future model runs, due to uncertainty in estimation.

Our results indicate that the sequence of deep glaciations in the Neoproterozoic could be the result of a change of state in the long-term carbon-climate system to a regime which exhibits self-sustaining oscillations. If there was a long period in which global steady state temperature remained below the value required to trigger a snowball glaciation, this would be manifest as an oscillatory regime, with snowball glaciations alternating with warm phases. Such a temperature forcing may well be attributed to the continental configuration at this time. It has been shown that the position of the continents at low latitudes at 750Ma, along with the prevalence of basaltic lithologies, could provide the necessary cooling to trigger the first snowball event [21]. It is thought that the continents would have remained near low latitudes until 600Ma [22], after which they begin to drift to higher latitudes, relaxing the forcing.

To investigate this possible mechanism we parameterise a runaway ice-albedo feedback in our model by imposing a change in albedo when temperature falls below a given value $T_{crit,}$. Assuming the classic snowball scenario [23], we choose $T_{crit,}$ = 283K and allow deglaciation at 263K. Because deglaciation begins in the tropics, it is assumed to occur at lower temperature than is required for the ice sheets to initially advance. Throughout this work we assume a solar constant for 650Ma (1298 Wm^{-2} [24]), broadly representing the timeframe of interest. This allows glaciation at ~150ppm CO_2, close to other estimates [21].

We impose the described cooling scenario in the model, adding a parameter ρ to represent enhancement of kinetic weathering. This follows the treatment of vascular plant colonisation in the Phanerozoic COPSE model runs [6], acting as a multiplier on the kinetic weathering rate equation. To trigger oscillations we increase ρ by a factor of three for a period of 150Myrs. The magnitude of this enhancement is roughly analogous to the increase in basaltic surface area and mid-latitude runoff calculated in ref [21]. For present day CO_2 degassing rate (D = 1), we require W_{max} = 1.4 to produce a rough analogue of the Neoproterozoic record. This parameter choice is shown in figure 4. Assuming a higher CO_2 degassing rate shortens glacial duration and allows for larger values of W_{max} to produce the observed timing, in line with figure 2.

We use output from the CO2SYS model [25] to approximate the atmospheric fraction of total ocean and atmosphere CO_2, assuming that there is gas exchange between atmosphere and ocean during glaciation [1]. The total solubility of CO_2 is higher in cold water than warm water, therefore deglaciation causes a large transfer of CO_2 from ocean to atmosphere.

The solid line for $\delta^{13}C$ shows the isotopic fractionation of marine carbonates, assuming the fractionation effect on burial takes into account the equilibrium fractionation between oceanic and atmospheric carbon, and a dependence on temperature, as in the full COPSE model [6]. The dashed line shows an alternative solution where fractionation effects are constant. Both treatments yield a continued positive fractionation during the interglacial period due to elevated burial of light organic carbon, due in turn to sustained above-average nutrient fluxes from weathering. Higher assumption of W_{max} increases nutrient delivery and therefore also increases fractionation. Low productivity during glaciations causes a negative excursion. We do not expect a simple model such as this to replicate exactly the isotope record. Negative excursions preceding glaciation are not reproduced by our model, and may be due to direct temperature effects on productivity, which are not included. Our aim is to demonstrate that the extended period of system disequilibrium following a snowball glaciation should contribute to prolonged positive excursions in $\delta^{13}C$, more complex analysis is required to fully understand the Neoproterozoic carbon cycle.

With the imposition of a suitable long term maximum weathering rate, oscillations in this simple carbon – climate model can provide a qualitative fit to the sequence of glaciations and carbon isotope variations in the Neoproterozoic. The globally transport limited scenario presents a prolonged period of elevated primary productivity, which would support suggested increases in oxygen concentration and phosphorous deposition over this time [7, 8, 26-29]. There is evidence for phosphorous deposition after the Marinoan glaciation but not after the Sturtian.

It is important to note that the mechanism we describe relies on a particular interpretation of the Neoproterozoic period, namely the Snowball Earth hypothesis [1, 10]. It is possible that the Neoproterozoic actually contained more frequent smaller glaciations, which would not terminate via a high CO_2 'super greenhouse'. Due to our extreme timeframe for CO_2 drawdown, our prediction is highly testable, with for example one recent study proposing a rapid decline in CO_2 following the

Marinoan glaciation [30]. Further work to establish the duration of any post-glacial greenhouse may thus enable validation or falsification of mechanisms to explain these fascinating events.

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Acknowledgements:

BM is funded by a UEA Dean's studentship. AJW's contribution was supported by a Royal Society Research Professorship. Part of CG's contribution was supported by a NASA Postdoctoral Fellowship at Ames Research Center and by the NASA Astrobiology Institute Virtual Planetary Laboratory lead team. AJW, RAB and TML are supported by the NERC project (NE/I005978/1) 'Reinventing the planet: The Neoproterozoic revolution in oxygenation, biogeochemistry and biological complexity'.

Author Contributions:

BM and AJW suggested the study. BM wrote the model, results were analysed by BM and AJW. Discussion with CG helped improve the method. RB and TML contributed to the manuscript.

Competing financial interests statement

The authors declare that they have no competing financial interests.

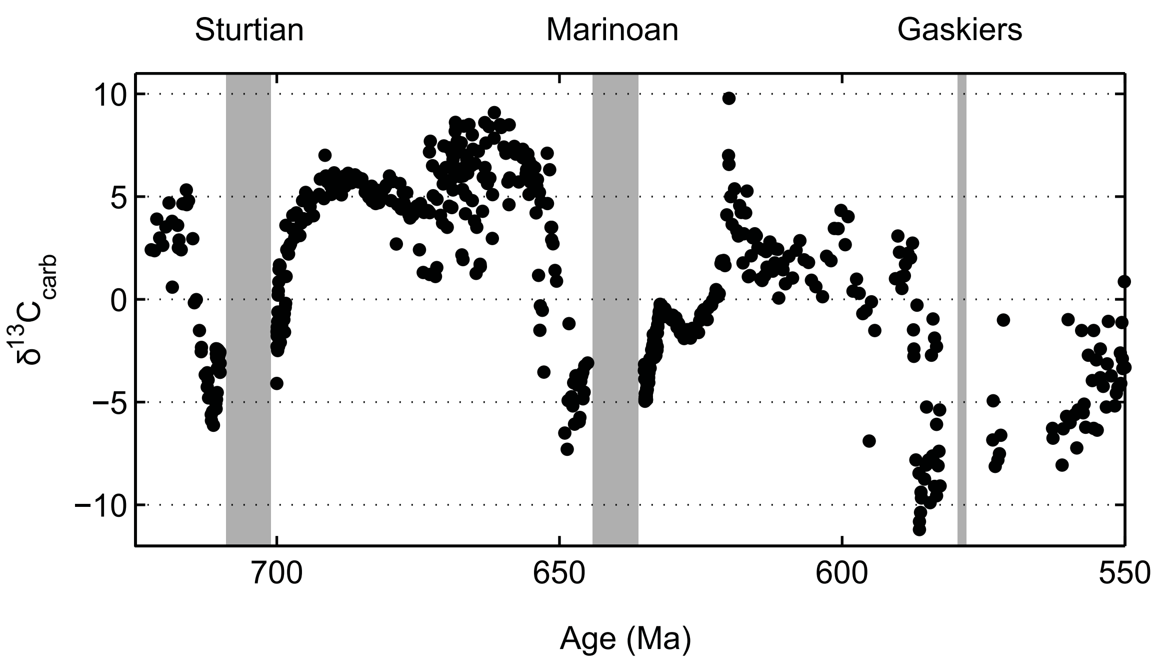
Figure captions

Figure 1: δ^{13}C record for the late Neoproterozoic. Isotopic composition of carbonates from ref. [7]. The vertical grey bars from left to right denote the Sturtian, Marinoan and Gaskiers glaciations.

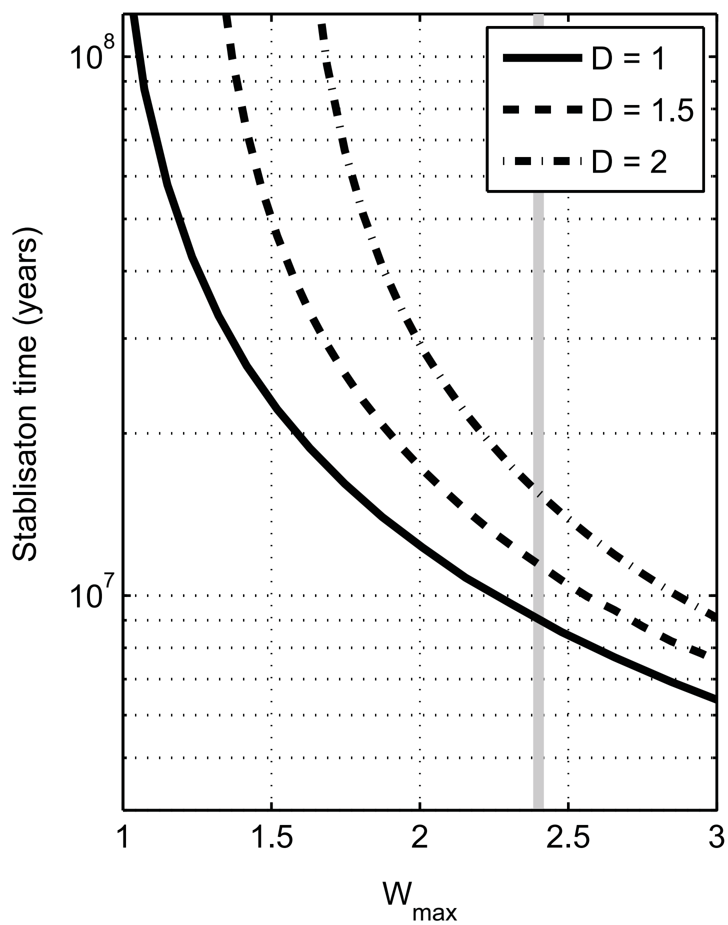
Figure 2: Phase portrait: stabilisation time versus maximum weathering rate, W_{max}. Here we assume an initial CO_2 concentration of 0.3 atmospheres, and fix the global weathering rate at W_{max}. The three lines show different choices of the relative CO_2 degassing rate, *D*. W_{max} is defined relative to present day silicate weathering rate, with the grey vertical line showing our estimate of W_{max}= 2.4 for the Phanerozoic. Increasing the weathering rate enhances nutrient delivery and therefore increases the organic burial fraction, allowing stability when W_{max} is somewhat smaller than D, providing W_{max} >1. See supplementary information for full model description.

Figure 3: Stabilisation time after 0.3 atm CO_2 perturbation for different initial abundances of rock flour. Here R_{max} denotes the maximum amount of carbon (in moles) that can be drawn down via weathering of glacial rock flour before it is depleted. This is set to 10^{18} (solid line), 10^{19} (dashed line) and 10^{20} (dash-dot line). This figure shows the situation where D = 1, W_{max} = 2.4. The grey vertical line shows the stabilisation time when no flour is present (as in fig2). The second drop in weathering rate here occurs as CO_2 returns to a stable concentration.

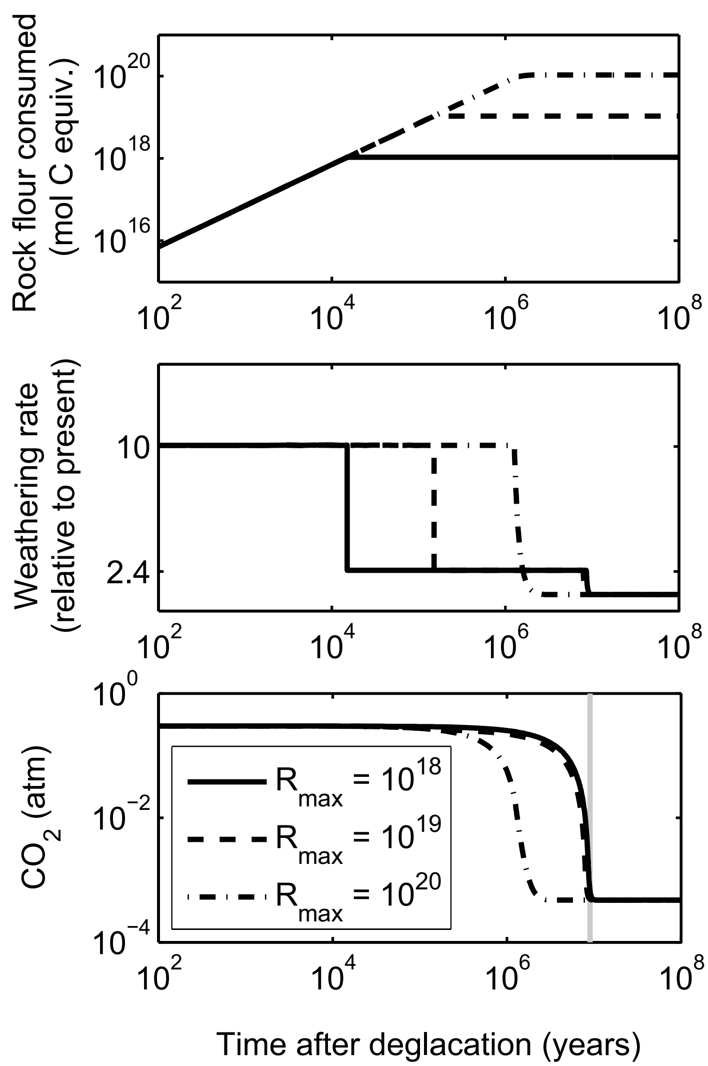
Figure 4: Cyclic solution when steady state temperature is forced below the ice-albedo runaway value for 150Myr. Here we let D = 1 and W_{max} = 1.4 to produce glacial timing on the order observed in the Neoproterozoic. The imposed kinetic weathering enhancement (ρ) is shown in grey in panel 1; in black is the weathering rate relative to present. Panel 2 shows total atmosphere/ocean carbon (grey), and atmospheric CO_2 (black). Panel 3 shows temperature alongside snowball entry/exit thresholds. Panel 4 shows $\delta^{13}C$. Solid line shows temperature/CO_2 dependent fractionation [6], dashed line shows solution when fractionation effects are constant.



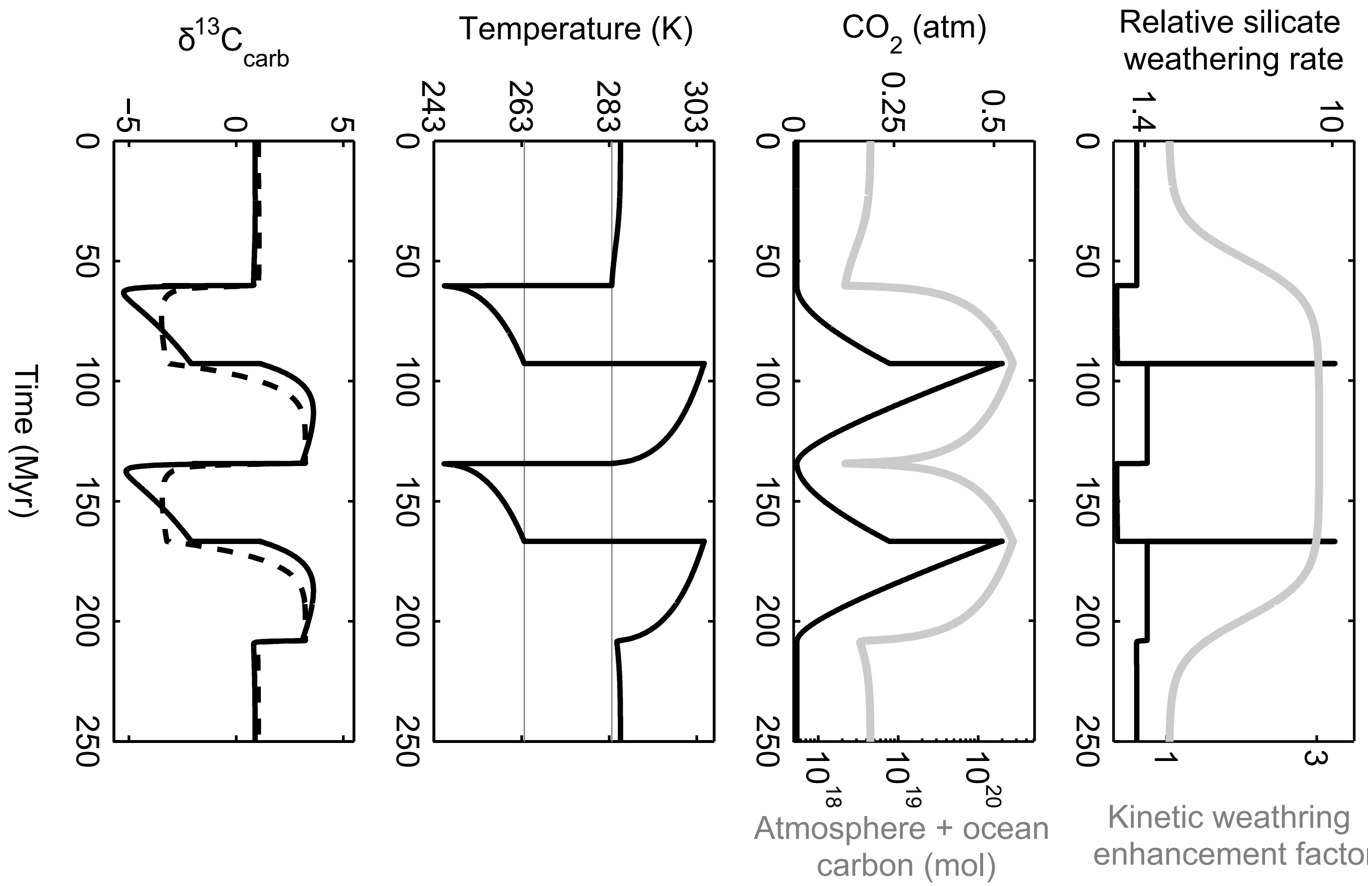
Sturtian
Marinoan
Gaskiers
10
5
0
−5
−10
700
650
600
550
$\delta^{13}C_{carb}$
Age (Ma)



D = 1
D = 1.5
D = 2
10^8
10^7
Stablisaton time (years)
1
1.5
2
2.5
3
W_{max}



Rock flour consumed (mol C equiv.)
10^{20}
10^{18}
10^{16}
10^2
10^4
10^6
10^8
Weathering rate (relative to present)
10
2.4
10^2
10^4
10^6
10^8
CO_2 (atm)
10^0
10^{-2}
10^{-4}
$R_{max} = 10^{18}$
$R_{max} = 10^{19}$
$R_{max} = 10^{20}$
10^2
10^4
10^6
10^8
Time after deglacation (years)



Relative silicate weathering rate
Kinetic weathring enhancement factor
CO2 (atm)
Atmosphere + ocean carbon (mol)
Temperature (K)
δ13Ccarb
Time (Myr)