# Towards a complete stellar mass function of the Hyades^ 

# I. Pan-STARRS1 optical observations of the low-mass stellar content 

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#### Abstract

Aims. The Hyades cluster is an ideal target to study the dynamical evolution of a star cluster over the entire mass range due to its intermediate age and proximity to the Sun. Methods. We wanted to extend the Hyades mass function towards lower masses down to $0.1 M_{\odot}$ and to use the full three-dimensional spatial information to characterize the dynamical evolution of the cluster. Results. We performed a kinematic and photometric selection using the PPMXL and Pan-STARRS1 sky surveys, to search for cluster members up to 30 pc from the cluster centre. We determined our detection efficiency and field star contamination rate to derive the cluster luminosity and mass functions down to masses of $0.1 M_{\odot}$. The thorough astrometric and photometric constraints minimized the contamination. A minimum spanning tree algorithm was used to quantify the mass segregation. Conclusions. We discovered 43 new Hyades member candidates with velocity perpendicular to the Hyades motion up to $2 \mathrm{~km} \mathrm{~s}^{-1}$. They have mass estimates between 0.43 and $0.09 M_{\odot}$, for a total mass of $10 M_{\odot}$. This doubles the number of Hyades candidates with masses smaller than $0.15 M_{\odot}$. We provide an additional list of 11 possible candidates with velocity perpendicular to the Hyades motion up to $4 \mathrm{~km} \mathrm{~s}^{-1}$. The cluster is significantly mass segregated. The extension of the mass function towards lower masses provided an even clearer signature than estimated in the past. We also identified as likely Hyades member an L0 dwarf previously assumed to be a field dwarf. Finally we question the membership of a number of previously published candidates, including a L2.5-type dwarf.


Key words. open clusters and associations: individual: Hyades - stars: luminosity function, mass function - stars: low-mass stars: individual: 2MASSI J0230155+270406 - stars: individual: 2MASS J05233822-1403022

## 1. Introduction

The determination of the (present day) mass functions of stellar clusters shed light on the formation of stars and brown dwarfs and of the production of the field population out of the stellar nurseries. The evolution of the cluster mass function, estimated from the study of various clusters of different ages, may be interpreted as the consequence of dynamical interactions of the cluster members and of the cluster with the Galactic gravitational potential. Studies such as Bouvier et al. (2008), Wang et al. (2011), Boudreault et al. (2012) of the intermediate-age clusters Pleiades, Praesepe and Hyades, have revealed the role of evaporation of low-mass members, as well as possibly different stages in the evolution at a given age.

[^0]Because of its proximity to the Sun (about 47 pc , for the most recent determinations see, e.g. van Leeuwen 2009; McArthur et al. 2011) and intermediate age (about 650 Myr , see e.g. Perryman et al. 1998; DeGennaro et al. 2009), the Hyades cluster is a valuable target to study the evolution of stellar clusters.

It is desirable to extend the Hyades census to lower masses to provide new benchmark objects at the end of the main sequence, with known age and metallicity. They can further constrain evolutionary models of the stellar core, of the stellar rotation (Delorme et al. 2011), and other stellar properties.

The Hyades cluster has a relatively large space velocity of $(U, V, W)=(-41,-19,-1) \mathrm{km} \mathrm{s}^{-1}$. This allows us to efficiently remove contaminants based on the proper motion. Röser et al. ( 2011, R11 thereafter) searched the vicinity of the Hyades cluster up to 30 pc from the centre, using the convergent point method and the wide-area surveys PPMXL (Röser et al. 2010) and CMC14 (Copenhagen University Obs. 2006).

The kinematic selection needs to be complemented by a photometric selection in order to improve the purity of the candidate sample. The Panoramic Survey Telescope And Rapid Response System (Pan-STARRS) will eventually comprise multiple telescopes and cameras, and currently consists of a single telescope and camera known as Pan-STARRS1 (or PS1, Kaiser et al. 2002). $\mathrm{PS}^{1}{ }^{1}$ is a new optical/red survey instrument with a large figure of merit thanks to its $7-\mathrm{deg}^{2}$ field of view, that provides new multi-epoch, multi-band observations of all the sky north of $-30^{\circ}$. It uses a dedicated $1.8-\mathrm{m}$ telescope located in Haleakalā on Maui. The pixel scale is 0.25 arcsec . It offers the perfect data set to study nearby clusters and extend their membership census to large cluster radii. PS1 has accumulated more than two years of survey data, with multiple grizy coverage of the Hyades. This data set allows us to push the study of R11 to greater depths and lower masses, and refine the member selection using multi-band photometry and astrometry unavailable until now.

The membership selection is a three step process. First, we select kinematic members by the convergent point method. Second, we further restrict the sample of candidates by photometric selection. Finally, we verify the PPMXL proper motions for all the candidates that survived the kinematic and photometric selections.

The paper is organized in the following way: we describe our data set in Sect. 2. Section 3 presents our kinematic selection while Sect. 4 presents our photometric selection. Candidates of particular interested are discussed in Sect. 5. The corrections to derive the luminosity function are described in Sect. 6, whereas the mass function is presented in Sect. 7. Finally we study the spatial distribution of our candidates in Sect. 8.

Through-out the article we conform to the usual practice and report optical photometry in the AB system (PS1, SDSS surveys) and near- and mid-infrared photometry in the Vega system (2MASS, UKIDSS, WISE surveys).

## 2. Observations

### 2.1. Pan-STARRS1 observations of the Hyades

A fraction of $56 \%$ of the observing time of PS1 is dedicated to the so-called $3 \pi$ survey, which monitors the sky north of $-30^{\circ}$. This survey covers the area in five filters, $g_{\mathrm{P} 1}, r_{\mathrm{P} 1}, i_{\mathrm{P} 1}, z_{\mathrm{P} 1}$ and $y_{\mathrm{P} 1}$, with repeated observations over 3 years. Details about the $3 \pi$ survey are in Chambers et al. (in prep.). The profiles of the blue filters, $g_{\mathrm{P} 1}, r_{\mathrm{P} 1}, i_{\mathrm{P} 1}$, are close to those of the SDSS filters (Stubbs et al. 2010; Tonry et al. 2012). In each filter, two visits of two exposures each are scheduled every year. The two exposures are separated by $30-90 \mathrm{~min}$, in order to detect fast-moving objects such as near-Earth objects. The telescope is pointed to the same direction within a few arcseconds, while the two visits are separated by a few days in the optical (gri), up to several months in the red bands $(z y)$ to improve the parallax factor range (Beaumont \& Magnier 2010). The static fill factor of the camera, taking into accounts the gaps between and within the chips, as well as areas of poor sensitivity, is about $75 \%$. As the survey progresses, the camera pointing and orientation is shifted to cover the areas missed in the previous seasons, and optimally improve the photometric and astrometric calibrations.

Between February 1st, 2010 and August 31, 2012, which form the photometric data we use in this article, Pan-STARRS1 has observed most of the area of the $3 \pi$ survey. With two observing seasons, most of the Hyades cluster (within $45^{\circ}$ of the

[^1]centre) has been observed 4 to 6 times in each filter under photometric or nearly-photometric conditions, with typical depths of 22 mag ( $\mathrm{AB}, 50 \%$ completeness limit, for individual exposures in gri, see Sect. 2.3). Extended periods of poor weather between mid-October and mid-January have prevented us from completing the programme in some limited areas (see Sect. 2.2). The image quality varies from $F W H M=1.4^{\prime \prime}$ in the $g$ band to $1.2^{\prime \prime}$ in the $z$ band, averaged over the whole Hyades area.

The images are processed on a nightly basis by the Image Processing Pipeline (IPP, Magnier 2006). The pipeline detrends and debiases the images, searches for sources $5 \sigma$ or more above background, and measures PSF photometry and various parameters. We select stars with good quality flags ${ }^{2}$.

The data are photometrically calibrated according to the algorithm of Schlafly et al. (2012), who use repeated observations of the same stars to constrain a model for the PS1 system throughput. The technique is an extension of that of Padmanabhan et al. (2008), used to calibrate the SDSS. Internal comparisons and comparison with the SDSS indicate that the photometric calibration is accurate at the $<1 \%$ level.

We transform the SDSS photometry into the PS1 photometric system using linear colour transformations (Finkbeiner, priv. comm.). We use the PS1 photometric system for candidate selection and in figures, but provide the original SDSS photometry in tables.

### 2.2. Sky coverage

Both PS1 and SDSS have incomplete sky coverage of the Hyades, the former because it is on-going and because of adverse weather pattern in the past two winter seasons; the latter because the cluster falls outside of the contiguous SDSS Northern Cap coverage.

To perform a spatial distribution analysis, it is necessary to create maps of the actual coverage. For PS1, the coverage variations have high spatial frequency, while SDSS is mostly contiguous over isolated stripes.

Because the incomplete sky coverage has its origin in the geometry of the focal plane (PS1) or of the survey (SDSS), we consider that the coverage fraction is independent of the candidate brightness.

About $25 \%$ of each exposure footprint is not observed because of the gaps between the sensitive devices and the masked areas due to bright stars, trails, etc. Therefore, a fraction of the kinematic candidates is not detected even if the whole area were covered a few times. The PS1 observations are conducted in pairs, called TTI pairs, with an identical pointing within a few arcseconds. The masks of the two images of the pair are thus largely overlapping. In our static mask, chip gaps appear every 20 arcmin, cell gaps every 2.5 arcmin. To produce a map with sufficient spatial resolution, we search the PS1 catalogue for matches for all 2MASS stars over the area of interest. We then calculate over $(12 \mathrm{arcmin})^{2}$ areas the fraction of 2MASS stars recovered in $g, r$ and/or $i$ (see Fig. 1).

This approach could lead to biases and underestimate the sky coverage in regions where 2MASS is deeper than PS1, i.e. in high extinction regions. We limit the 2MASS star faintness to $J<15$ mag to minimize this effect.

This calculation indicates a coverage in any of the $g, r$ or $i$ bands of $94 \%$, versus only $44 \%$ in all three bands. For this

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Fig. 1. Sky coverage of our PS1 data combined with SDSS-DR8 around the Hyades cluster in either the $g, r$ or $i$ filters. The colour scale indicates the fraction of 2MASS stars recovered in our catalogue (see text). The PS1 3-deg-diameter footprint is visible in a few locations. The white star indicates the cluster centre and the cross, the convergent point.
reason, we chose to select candidates in any band to improve coverage, at the cost of purity.

We perform a similar analysis for the SDSS and PPMXL surveys. The SDSS sky coverage is combined with that of PS1. For PPMXL, the average detected fraction of 2MASS stars with magnitudes between 9 and 13 is $98 \%$ (over ( 30 arcmin) ${ }^{2}$ areas). The lacking areas in PPMXL cover $1.2 \mathrm{deg}^{2}$ and are close to very bright stars and M77. We do not correct our coverage maps for PPMXL heterogeneities.

### 2.3. Depth

We determine the detection efficiency as a function of magnitudes and sky location of the surveys we use, namely PPMXL and Pan-STARRS1, by comparing them with SDSS DR8. For each SDSS star (type=6), we search within $1^{\prime \prime}$ of its position for a PPMXL or a PS1 $g, r$ or $i$ detection. We report the recovery rate as a function of SDSS $g, r$ or $i$ magnitudes. While the SDSS DR8 catalogue covers only a fraction of the area of interest, this procedure still gives a good estimate of the typical depth of both surveys over the survey region.

This method is not affected by the fact that SDSS is not much deeper than PS1. What matters is the purity of the SDSS pointsource catalogue, so that spurious SDSS detections are not interpreted as objects missed by PS1.

We determine the $50 \%$ completeness limit, (after removing the partial coverage incompleteness), over $\Delta \mathrm{RA} \times \Delta \mathrm{Dec}=1 \times$ $1 \mathrm{deg}^{2}$ areas (see Fig. 2).

In the case of PS1, we consider the $50 \%$ completeness relative to the completeness over the $16-18-\mathrm{mag}$ range. Because of the sharp decline in the PS1 recovery rate, the details of the normalization procedure does not affect much the measured depths.

We find that PS1 has a $50 \%$ completeness limit of $g=22.6$, $r=21.2$ and $i=21.8 \mathrm{mag}$, with respective widths of $0.44,0.48$ and 0.33 mag , over the region of interest (see Fig. 2). This is the depth of individual exposures, whose catalogues are combined (unique detections among multiple exposures are considered valid) on a filter per filter basis.


Fig. 2. Completeness of PS1 as a function of $g$ (blue), $r$ (green), and $i$ (red) SDSS magnitudes (solid curves, right to left), and histograms of the $50 \%$ completeness limit (dashed curves, normalized to unity). The similar completeness of CMC14 as a function of $r$ SDSS magnitude is shown (dotted curve, black). The maximum reflects the fraction of the area observed in each filter.

The limiting survey in this study is PPMXL (with a $50 \%$ completeness limit of $r=20.3 \mathrm{mag}$, see Sect. 6.1), except in some very restricted areas with poor PS1 observing conditions. Therefore we ignore the PS1 depth variations over the sky.

We perform the same analysis for CMC14, which is used to select low-mass Hyades candidates by R11. The result, shown in Fig. 2, is in agreement with the CMC14 documentation. The CMC14 and PS1 surveys complement each other: the PS1 saturation limit is brighter than CMC14's 95\% completeness limit, and the PS1 depth is about 4 mag fainter. In addition, CMC14 does not cover our survey area north of $+50^{\circ}$ and a small area at the edge of the cluster.

### 2.4. Other surveys

We complement the Pan-STARRS1 photometry and the 2MASS magnitudes from PPMXL with SDSS DR8 (Aihara et al. 2011) and UKIDSS DR8 (Lawrence et al. 2007; Hambly et al. 2008), searching within $1^{\prime \prime}$ of the predicted PPMXL position. We report the SDSS photometry (see Table 1) and use the optical data for the photometric selection.

The SDSS-DR8 coverage of the Hyades is limited to a few 2.5 -deg-wide stripes, making up less than $16 \%$ of the area. One of the stripes overlaps the centre. We use the SDSS photometry to complement the PS1 coverage, and to check the depth of the current PS1 catalogue.

UKIDSS, in its Galactic Cluster Survey component, also covers the Hyades cluster. Because our original kinematic selection requires 2MASS detection, UKIDSS data mostly allows us to refine the photometry of the faintest candidates (see Sect. 4.3).

Finally, WISE released its all-sky catalogue on March 14, 2012 (Wright et al. 2010). We marginally use the midinfrared photometry to select our candidates, and provide their photometry in the WISE filters (cf. Table 1).

Table 1. Basic description of the parameters of the R11 candidates and the 62 new candidates.

| Col. | Label | Unit | Description |
| :---: | :---: | :---: | :---: |
| 1 | ID |  | Running number (as in R11 if available) |
| 2 | RA | deg | PPMXL right ascension |
| 3 | DE | deg | PPMXL declination both J2000.0, epoch 2000.0 |
| 4 | RA_PS1 | deg | PS1 right ascension |
| 5 | DE_PS1 | deg | PS1 declination both J2000. 0 |
| 6 | MJDps | d | epoch of PS1 positions |
| 7 | PPMXL ID |  | PPMXL identifier |
| 8 | PPMXL flag |  | see Röser et al. (2010) |
| 9 | pmRA | mas $\mathrm{yr}^{-1}$ | Proper motion in RA $\cos (\mathrm{DE})$ |
| 11 | pmDE | mas yr ${ }^{-1}$ | Proper motion in DE |
| 13 | $\varpi$ | mas | Secular parallax |
| 15 | $r_{\text {c }}$ | pc | Distance from the cluster centre |
| 16 | $v_{r}$ | $\mathrm{km} \mathrm{s}^{-1}$ | Predicted radial velocity |
| 17 | $v_{\\|}$ | $\mathrm{km} \mathrm{s}^{-1}$ | Tangential velocity in direction to the convergent point |
| 18 | $v_{\perp}$ | $\mathrm{km} \mathrm{s}^{-1}$ | Tangential velocity perpendicular to the direction to the convergent point |
| 21 | $m$ | $M_{\odot}$ | mass of the star |
| 22 | $g_{\text {P1 }}$ | mag | PS1 g PSF median magnitude |
| 24 | $r_{\text {P1 }}$ | mag | PS1 $r$ PSF median magnitude |
| 26 | $i_{\text {P1 }}$ | mag | PS1 $i$ PSF median magnitude |
| 28 | $z_{\text {P1 }}$ | mag | $\mathrm{PS} 1 z \mathrm{PSF}$ median magnitude |
| 30 | $y_{\text {P1 }}$ | mag | PS1 y PSF median magnitude |
| 32 | $g^{\prime}$ | mag | SDSS DR8 $g$ magnitude |
| 34 | $r^{\prime}$ | mag | SDSS DR8 $r$ magnitude |
| 36 | $i^{\prime}$ | mag | SDSS DR8 $i$ magnitude |
| 38 | $z^{\prime}$ | mag | SDSS DR8 z magnitude |
| 40 | $J_{2 \mathrm{M}}$ | mag | 2MASS $J$ magnitude |
| 42 | $\mathrm{H}_{2 \mathrm{M}}$ | mag | 2MASS $H$ magnitude |
| 44 | $K_{s 2 \mathrm{M}}$ | mag | 2MASS $K_{\mathrm{s}}$ magnitude |
| 46 | W1 | mag | WISE $W 1$ magnitude |
| 48 | W2 | mag | WISE $W 2$ magnitude |
| 50 | W3 | mag | WISE $W 3$ magnitude |
| 52 | W4 | mag | WISE $W 4$ magnitude |
| 54 | Source flag |  | 0: R11; 1: this paper |

Notes. Errors are included for most parameters. See Vizier for a complete description and format of the file. (Table available at the CDS.)

## 3. Kinematic selection

### 3.1. The convergent point method

The convergent point method (CPM; for a detailed recent description, see van Leeuwen 2009) is an important tool to isolate (from the field) an ensemble of stars sharing the same 3D-space motion such as the stars in an open cluster. R11 have applied this method, using a subset of PPMXL having CMC14 $r$-band photometry, to find an empirical main sequence of present-day (and also of formerly bound) members of the Hyades down to an absolute magnitude of $M_{K_{\mathrm{S}}}$ of about 9 mag. We follow the approach of R11 (Sect. 4.1) and select only stars within a radius of 30 pc around the centre of the cluster, which we define at RA $=4$ h $28^{\prime} 25.7^{\prime \prime}$, Dec $=+16^{\circ} 42^{\prime} 45^{\prime \prime}(J 2000.0)$.

We allow a limit of $4 \mathrm{~km} \mathrm{~s}^{-1}$ for the velocity component $\left|v_{\perp}\right|$ perpendicular to the direction to the convergent point to select possible members. We also set an upper bound to the angle between the proper motion vector and the vector in the direction to the convergent point, of $9.5^{\circ}$, in order to limit the contamination for candidates closer to the convergent point. In order to increase the purity of our sample, we further require our candidates
to have $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$. We find that these parameters offer a good compromise of purity and completeness: the stronger constraint roughly decreases by a factor of two the field contamination (based on our Besançon simulation, see Sect. 6.2). On the other hand, a residual velocity of $2 \mathrm{~km} \mathrm{~s}^{-1}$ is still much larger than the velocity dispersion in the cluster, which is predicted to be smaller than $1 \mathrm{~km} \mathrm{~s}^{-1}$ (see R11; Ernst, priv. comm.).

However, the precision on our measurement of $\left|v_{\perp}\right|$ itself varies between 1 and $2 \mathrm{~km} \mathrm{~s}^{-1}$. We estimate the $\left|v_{\perp}\right|$ error by propagating the PPMXL proper motion measurement errors, and therefore it depends on target brightness and number of available epochs. For the faintest candidates close to the PPMXL detection limit, the incompleteness would be significant. For instance, the $15 \%$ of kinematic candidates with the largest errors $\left(\sigma\left(\left|v_{\perp}\right|\right)>1.5 \mathrm{~km} \mathrm{~s}^{-1}\right)$ have mostly a mass smaller than $0.1 M_{\odot}$, and are located 10 pc or more behind the cluster centre. The incompleteness for such members would raise to $30 \%$, but as we argue below (see Sect. 6.1), the present work concentrates on the mass range above this limit, for which the incompleteness is generally smaller than $10 \%$. Furthermore, Hyads in multiple systems have an orbital motion added to their secular motions. For wide binaries of the appropriate period, in the order of the PPMXL time baseline (a few decades) or longer, the proper motions of the components will differ from the proper motion of the centre of gravity, and the perpendicular velocity of each component may be larger than $2 \mathrm{~km} \mathrm{~s}^{-1}$ (see footnote 5 in Sect. 8 for an example; also Wielen et al. 1999; Frankowski et al. 2007).

We provide the list of candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ and those with $2<\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$ together in the Tables 5 and $6^{3}$

As R11 describe, the CPM is no final confirmation for a star to be co-moving with the bulk of the Hyades cluster, it "predicts" a secular parallax and a radial velocity for each candidate. Both quantities have to be verified to conclusively confirm membership. Using the predicted secular parallaxes one is able to construct colour-absolute magnitude diagrams (depending on the observed bands) to check if the candidates populate allowed loci in these diagrams.

The loci of M-stars in the $M_{K_{\mathrm{s}}}$ vs. $J-K_{\mathrm{s}}$ form an almost vertical line for $M_{K_{\mathrm{s}}} \geq 5.5 \mathrm{mag}$ with the consequence that the predicted $M_{K_{\mathrm{s}}}$ cannot be checked unless observations in a shorter wavelength band is available. Optical photometry from PPMXL itself is not useful to check the loci of kinematic candidates in a colour-magnitude diagram. It originates from photographic photometry from USNO-B based on POSS I and II, which is not accurate enough for this purpose. Hence, the R11 search for members turned out to be limited by the completeness limit ( $r \approx 17 \mathrm{mag}$ ) of CMC14.

The full PPMXL catalogue goes considerably deeper than CMC14 (Figs. 2 and 9), so including PS1 enables to check all the low-mass kinematical candidates for photometric membership. In this paper we selected 16543 such candidates, with partial overlap with the list published in R11. For the latter objects, the deeper PS1 photometry will refine the results based on the shallower CMC14 catalogue.

### 3.2. Confirmation of the PPMXL astrometry

The accuracy and the reliability of the proper motions are essential in the CPM. Röser et al. (2010) state that PPMXL contains a large number of stars with spuriously large proper motions,

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Fig. 3. Histograms of the offsets along RA and Dec between the expected and observed positions of the 247 candidates (thick line) and nearby field objects (normalized; thin line, with the fitted Gaussian as dashed line). The eight candidates with an offset larger than 300 mas are highlighted.
many of them are related to the (about $10 \%$ ) double entries in PPMXL. We can verify the PPMXL proper motions by the present-day PS1 astrometry itself. PPMXL uses USNO-B1.0 positions, complemented with 2MASS, with epochs spread over five decades for the objects detected in the first plate surveys, down to two decades for the objects only detected in the latest plate surveys. PS1 adds measurements 12 years on average after 2MASS's epoch.

The current PS1 astrometric registration is performed on a chip-per-chip basis ( 20 -arcmin on a side) against 2MASS (Magnier et al. 2008). Then the astrometric solutions are improved with PS1 detections alone, using overlapping PS1 exposures. The typical precision against 2MASS is 80 mas in each direction. However here we are interested in the PS1 precision relative to PPMXL (or vice versa). To estimate it we search the PS1 catalogue (all bands including $z_{\mathrm{P} 1}$ and $y_{\mathrm{P} 1}$ ) for PPMXL detections, up to 12 arcmin of each candidate. We only consider objects with proper motion smaller than $5 \mathrm{mas} \mathrm{yr}^{-1}$ in both directions, to avoid erroneously large PPMXL proper motions due to confusion. We measure the offset between the expected positions based on PPMXL astrometry (proper motion and positions at its published epoch 2000.0) and the PS1 observed positions on each exposure. Finally, we subtract from those offsets the mean offset of the objects in the vicinity of each candidate. This brings the PS1 astrometric system, which is relative to 2MASS and to the referential of the field bright stars and not to the inertial referential of extragalactic sources, into the IRCS. (Over 11 years, the correction is typically $\Delta \alpha=-23$ mas yr $^{-1}$ and $\Delta \delta=-47$ mas $\mathrm{yr}^{-1}$ with a dispersion of $35 \mathrm{mas} \mathrm{yr}^{-1}$.) We find that the core of the offset distributions along each direction, $\alpha$ and $\delta$, are well fitted by a Gaussian function of width 90 mas, which is therefore the dispersion of the PS1 astrometry relative to PPMXL (Fig. 3, blue curves).

Similarly, for each kinematic candidate, we extrapolate its PPMXL position using the PPMXL proper motion to the epoch of observations, and search for detections within 1 arcsec of that position. This search radius is much larger than the typical accuracy measured above, so that we do not need to use
an adaptive search radius as a function of the PPMXL position or proper motion uncertainties. On the other hand, erroneous large proper motions from PPMXL (wrong by more than $100 \mathrm{mas}_{\mathrm{yr}}{ }^{-1}$ ) would prevent us from recovering the (actually slow-moving) candidates. Those objects could not be Hyades members.

We find that most objects with an offset larger than 300 mas are likely slow-moving, distant sources. A few remaining candidates with erroneous proper motions due to confusion, partly identified by eye examination, can have small offsets if they have proper motions smaller than 60 mas $\mathrm{yr}^{-1}$. With these limitations in mind, we require that the offsets be smaller than 300 mas, measured with four or more PS1 measurements.

We do not use other surveys such as SDSS or UKIDSS to perform a similar verification, as the baseline between the 2MASS observations and SDSS or most of UKIDSS imaging is only a few years and do not set strong constraints.

## 4. Selection of Hyades member candidates

The kinematic sample is affected by background and foreground stars with the same proper motions mimicking the Hyades space velocities for their sky location. To remove these contaminants, we use optical photometry from Pan-STARRS1 and SDSS, combined with infrared photometry from 2MASS, and WISE.

### 4.1. Photometric selection

We again follow the steps of R11 and select our candidates in the absolute magnitude vs. colour diagrams, taking the secular parallax from the CPM.

We define our selection box empirically in colour-magnitude diagrams (see Fig. 4). Indeed, the available theoretical isochrones do not reproduce accurately the colours of the Hyades low-mass stars. As an example, we display the latest calculation (CIFIST2011) of the BT-Settl $600-\mathrm{Myr}$ isochrone (Allard, priv. comm.) ${ }^{4}$. On the other hand, each colour-magnitude diagram clearly exhibits a sequence above the cloud of field stars, particularly for $M_{g}<16 \mathrm{mag}$. In Fig. 4a we show the kinematic candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ within 9 pc of the cluster centre. Close to the cluster centre the field contamination is minimal and the over-density clearly visible. Most of the field stars do not share the cluster motion, so that their kinematic parallax is erroneous and their calculated absolute magnitude does not correspond to that of a Hyades-sequence star of their colour. In Fig. 4c, the locus of the R11 candidates (with CMC14 $r$-band photometry) overlaps with this sequence.

We note that the discrimination between the Hyades cluster sequence and the bulk of the background stars is the better the larger is the wavelength difference between the bands. In particular, $z$ and $y$ photometry alone combined with 2MASS photometry does not provide a good rejection (see Fig. 5). We note again that the isochrone does not reproduce the observations. We therefore concentrate on $g, r$ and $i$ bands in the following. For each CMD, we fit the Hyades stellar sequence with a second-order polynomial:
$M_{g}=9.628+0.8601 \times\left(g-K_{\mathrm{s}}\right)+0.04491 \times\left(g-K_{\mathrm{s}}\right)^{2}$
$M_{r}=9.834+1.1000 \times\left(r-K_{\mathrm{s}}\right)+0.05701 \times\left(r-K_{\mathrm{s}}\right)^{2}$
$M_{i}=10.312-1.4802 \times\left(i-K_{\mathrm{s}}\right)+0.07716 \times\left(i-K_{\mathrm{s}}\right)^{2}$.

[^4]

Fig. 4. a) Colour-magnitude diagram of $g_{\mathrm{P} 1}-K_{2 \mathrm{MASS}}$ for all the kinematic candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ within 9 pc of the cluster centre. b), c), d) Hess diagrams of $g_{\mathrm{P} 1}-K_{2 \mathrm{MASS}}$, of $r_{\mathrm{P} 1}-K_{2 \mathrm{MASS}}$, and of $i_{\mathrm{P} 1}-K_{2 \mathrm{MASS}}$, for the all kinematic candidates, i.e. with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{up}$ to 30 pc from the cluster centre. Filled (green) circles indicate candidates selected in $g, r$ and $i$; open (blue) circles: candidates selected in one or two filters; (red) crosses and pluses: candidates suspected to be giants (see Fig. 6) or to have bad proper motion measurements. Dots (magenta) represent R11 Hyades candidates with CMC14 r-band photometry in c) and PS1 and SDSS photometry in a), b), d). The dotted lines (blue) indicate our selection zone, while the thick (magenta) line represents the BT-Settl model 600-Myr isochrone. The grey area at the faint, red edge shows our incompleteness limit.

These constraints are only valid for stars with $M_{K_{\mathrm{s}}}>6$ mag. Candidates presumably brighter than this are better selected in R11 and are not considered in this study. To refine the selection box in order to minimize the contamination from field stars, we perform a numerical simulation of the field star contamination using the Besançon model (see Sect. 6.2). We adopt the following selection criteria: we select candidates redder than the fitted Hyades sequence shifted by -0.2 mag and bluer than the Hyades sequence of equal-luminosity binaries (see dotted lines in Fig. 4). This would allow the selection of equally-bright binaries.

Because of the partial sky coverage of those filters, we select candidates detected in any of those bands and satisfying our criterion. We reject a candidate if it fails our criterium in
any of the available CMDs. We select 273 PS1 candidates objects with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$. We perform the same analysis with SDSS-DR8 photometry, and select 30 SDSS candidates. There are 29 candidates selected with both PS1 and SDSS photometry. The SDSS DR8 catalogue contributes no candidate, and rejects no PS1 candidate, while PS1 rejects one SDSS candidate. The sample holds (again) 273 candidates.

In what precedes we have used the median PSF magnitudes of the good PS1 measurements. We can compare our PSF magnitudes with the aperture magnitude, to remove problematic measurements, such as poorly resolved visual binaries, or extended objects. The proper motions of such objects would likely be unreliable. We decide to remove the 13 candidates with PSF and aperture magnitudes discrepant by more than 0.5 mag. This cut


Fig. 5. Colour-magnitude diagram of $z_{\mathrm{P} 1}-K_{2 \mathrm{MASS}}$ for all the kinematic candidates, i.e. with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$, satisfying the photometric selections described in Sect. 4. Filled (green) circles indicate candidates selected in $g, r$ and $i$; open (blue) circles: candidates selected in one or two filters. Dots (magenta) represent R11 Hyades candidates (PS1 and SDSS photometry). The thick (magenta) line represents the BT-Settl model $600-\mathrm{Myr}$ isochrone.
only removes the few particularly poor measurements, and the dispersion of the magnitude differences of the remaining candidates is 0.05 mag . Most of the removed candidates were selected in a single band, have $\left|v_{\perp}\right|>2 \mathrm{~km} \mathrm{~s}^{-1}$, or are suspected to be giants (see below).

### 4.2. Rejection of background giant stars and erroneous proper motions

In Fig. 6 we present the traditional near-infrared colour-colour diagram where giants tend to populate the upper-left side of the diagram above the sequence of the dwarf stars. Again the isochrones do not describe accurately the observations. Figure 6 suggests a systematic offset of 0.04 mag in the $H$ band. We find 15 candidates that seem to display high extinction values. The excess maps of Schlegel et al. (1998) indicate large amounts of dust on their sky location (see Fig. 7), mostly at low Galactic latitude or towards the Taurus-Auriga star forming region. If truly reddened, these objects can only be located at large distances beyond the Hyades cluster. For instance, Tamburini et al. (2002), in a study of 1000 nearby stars, determines that the contribution to the polarization by the interstellar medium seems to become effective only after $\approx 70 \mathrm{pc}$, while our search excludes candidates beyond 77 pc . ( $90 \%$ of these objects have secular distances smaller than 70 pc .)

The space velocity of those objects whose near-infrared colours may indicate that they are giant stars, would be so large that they could be members of the Hyades; or, more likely, the PPMXL proper motion is erroneous.

We therefore reject candidates with:
$J-H>0.74$ mag and
$J-H>0.74+1.6 *\left(H-K_{\mathrm{s}}-0.34\right) \mathrm{mag}$.


Fig. 6. Near-infrared colour-colour diagram of our candidates, after the selection in the optical-near-infrared CMDs. The R11 candidates appear as magenta dots, the new candidates as green circles, and red crosses show kinematical candidates with suspiciously high reddening. Objects located above the dashed lines are rejected as background stars (red crosses). The BT-Settl model $600-\mathrm{Myr}$ isochrone down to $0.07 M_{\odot}$ is shown as the thick (magenta) line, and the reddening vector $\left(R_{V}=3.1\right)$ as the black arrow.

In addition, we inspect the neighbourhood of the candidates in VizieR. While PPMXL used the individual epochs from the sources of USNO-B1.0 (Röser et al. 2010) to derive proper motions, it relies, however, on the cross-matching automatically done in USNO-B1.0. So, for the small number of candidates that survived kinematic and photometric selection, we perform a visual inspection on VizieR and on the digitised sky survey charts from IRSA (NASA/IPAC Infrared Science Archive). Through our visual inspection and the analysis presented in Sect. 3.2, we remove 12 candidates that have a close neighbour or a likely erroneous proper motions, and keep 236 candidates.

The parameters available for each entry are described in Table 1. The full information, given in Table 6 is available at the CDS. Table 1 summarises its content: Cols. 1 and 2 list the J2000.0 positions as in PPMXL; Cols. 3 and 4, the PPMXL identifier and flag which describes the quality of the proper motion fit and the origin of the measurement; Cols. 5 to 8, the PPMXL proper motions and uncertainties; Cols. 9 and 10, the secular parallax and its uncertainty; Col. 11, the distance to the cluster centre; Col. 12, the predicted radial velocity; Cols. 13 to 16, the tangential velocity in the direction and perpendicular to the direction to the convergent point; Col. 17, the candidate mass; Cols. 18 to 49, the PS1, SDSS, 2MASS, and WISE magnitudes and their uncertainties; Col. 50, the source of the candidate, discovered in R11 or in this paper.

### 4.3. Further selection with UKIDSS and WISE photometry

We use additional surveys to refine the selection. UKIDSS observations are available for 92 candidates, mostly in the $K$ band. UKIDSS photometry confirms all those candidates except one, PPMXL J63.4757+15.8357, which has its $K$ magnitude


Fig. 7. Schlegel excess map of $E(B-V)$ (Schlegel et al. 1998), with the darkest areas having the largest excess. The Hyades candidates are super-imposed: green circles for the candidates detected in PS1, with low reddening; red crosses and pluses: candidates suspected to be giants (see Fig. 6) or to have bad proper motion measurements; magenta dots for R11. The blue diamond shows the convergent point.
discrepant with 2MASS's. It has two UKIDSS measurements, $K=12.02$ and 11.80 mag , both marked as close to saturated ( $\mathrm{k} \_1 \mathrm{ppErrBits}=0 \times 10$ ), while 2MASS has 11.35 mag , with a negligible colour term. There is no UKIDSS ZYJH photometry. We ignore the UKIDSS measurements and keep the candidate.

We cross-matched our kinematic candidate catalogue with the WISE sky data release (Wright et al. 2010), within 5 arcsec of the J2011.0 PPMXL position. We recover all photometric candidates but a single $K=12.2-\mathrm{mag}$ candidate. We construct CMDs using optical and near-infrared (2MASS) photometry combined with WISE W1 and W2 bands (Fig. 8). Most candidates follow a sequence, again shifted from the theoretical isochrone, but only by +0.05 mag in colour. Only one candidate appears much bluer than the bulk of the candidate, PPMXL J101.3119+36.9271. This candidate, which has a double entry in PPMXL, is detected in PS1 $r$ band only, and is located 28 pc from the cluster centre. It is most probably a spurious candidate, and we remove it.

### 4.4. Status of the R11 candidates

We search the 724 candidates of R11 in our catalogues as we did for the kinematic sample as described above. Many R11 candidates are bright and possibly saturated in PS1 or SDSS, so that we only consider candidates with PS1 or SDSS $g>16, r>16$ or $i>16 \mathrm{mag}$, which is conservative. The CMC14 photometry of brighter objects is more reliable than the PS1 and SDSS photometry. We check the location of those fainter R11 candidates in the colour-magnitude diagrams, as described in Sect. 4.1, using SDSS and/or PS1 photometry. Seven candidates are rejected in one or more colour-magnitude diagrams (see Table 2) and we remove them from the subsequent analysis.

We also check the proper motions of the fainter R11 candidates following the analysis presented in Sect. 3.2. We find that six candidates have offsets in the PS1 catalogue of 300 mas or more (see Table 3), and we remove them from the sample.

### 4.5. Final sample

Our PS1+SDSS sample is made up of 236 candidates, with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$. Among those, 62 candidates were not included in the R11 candidates. We provide the PS1 photometry and a subset of relevant parameters of these new candidates in Table 5.

Among those 62 candidates, eight were previously identified as Hyades candidates (see Sect. 5). Turning now to the "best" candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$, we select 207 candidates, with 48 not included in the 620 remaining R11 candidates having $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$. We combine our two samples into a list of 668 cluster member candidates. The proper motions of R11 differs slightly from the PPMXL proper motions because the astrometry of CMC14 was combined with the PPMXL measurements. Whenever available we use the proper motions of R11.

## 5. Objects of special interest

We search our sample for known objects in Simbad and Vizier (all catalogues), within $5^{\prime \prime}$ of the epoch-2000.0 PPMXL position. We cannot discuss all observations of our candidates, so we note only those referring to cluster membership, spectral classification, age, or distance. The present study, based on the optical survey of USNO-B and the optical bands of PS1 is not very sensitive to brown dwarfs, except the closest ones. Indeed we find none of the candidates of Hogan et al. (2008) in our final list.

Eleven candidates are listed as field dwarfs in Slesnick et al. (2006, their Table A.1), identified from their Na-8189 index indicating an age $\gtrsim 100 \mathrm{Myr}$. Our candidates have an index between 0.78 and 0.82 typical of an intermediate age.

### 5.1. Hyades members from Bouvier et al. (2008)

We recover in our final list 13 of the candidates of Bouvier et al. (2008): CFHT-Hy-1-6, 8-10, 12, 13, 15, and 17. Eleven of those were already found in R11; here we add CFHT-Hy-12 and 13. We do not select the remaining candidates of Bouvier et al. (2008) for the following reasons: CFHT-Hy-11 is kinematically selected but photometrically rejected in both the $M_{K}$ vs. $g_{\mathrm{P} 1}-K_{\mathrm{s}}$ and $M_{K}$ vs. $i_{\mathrm{P} 1}-K_{\mathrm{s}}$ diagrams as too red by 0.04 and 0.09 mag respectively. With $\mathrm{v}_{\perp}=2.2 \mathrm{~km} \mathrm{~s}^{-1}$, it seems likely to be a contaminant. Three candidates are not in PPMXL (CFHT-Hy-16, 20 and 21), while the five remaining candidates were rejected on the basis of their kinematics, because of discrepant proper motions between Bouvier et al. (2008) and PPMXL. The proper motion of PPMXL for CFHT-Hy-12 and CFHT-Hy-13 differs from the proper motion of Bouvier et al. (2008) by between 10 and 24 mas $\mathrm{yr}^{-1}$, but the PPMXL baseline is much longer. We note that we do not select objects from their Table D. 1 (Hyades probable non-members).

### 5.2. Other known Hyades candidates

We identify six of the previously known Hyades candidates in our final list (ordered by RA):

- DENIS J0235495-071121 was classified as an M5.5 by Phan-Bao \& Bessell (2006), with a spectrophotometric distance of 25.5 pc , while we predict a secular distance of 28.3 pc. West et al. (2008) classify it as an M7 star based on SDSS spectroscopy. Gálvez-Ortiz et al. (2010) assigns it a distance of $28.2 \pm 3.4 \mathrm{pc}$ and a Hyades membership. We
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Fig. 8. Colour-magnitude diagrams for the kinematic candidates with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$ within 30 pc of the cluster centre: a) of $r_{\mathrm{P} 1}-K_{\mathrm{W} 1}$; b) of $J_{2 \text { MASS }}-\mathrm{W} 1$; c) of $K_{2 \mathrm{MASS}}-\mathrm{W} 1$; and $\mathbf{d}$ ) of $K_{2 \mathrm{MASS}}-\mathrm{W} 2$. Filled (green) circles indicate candidates selected in $g$, $r$ and $i$; open (blue) circles: candidates selected in one or two filters; (red) crosses and pluses: candidates suspected to be giants (see Fig. 6) or to have bad proper motion measurements. b), c), d)) Dots (magenta) represent R11 Hyades candidates, with CMC14 r-band photometry in c). The thick (magenta) line represents the BT-Settl model 600-Myr isochrone.
predict a somewhat large perpendicular velocity of (3.0 $\pm$ $0.8) \mathrm{km} \mathrm{s}^{-1}$.

- Cl* Melotte 25 VA 200 (LP 475-7) is listed in Giclas et al. (1962) as a Hyades member. We derive $v_{\perp}=(-3.8 \pm$ $0.8) \mathrm{km} \mathrm{s}^{-1}$.
- LP 415-19 (PPMXL J65.43489+20.40293) was identified as a possible Hyades member by Reid (1992, 1993). It is unresolved by Reid \& Mahoney (2000), but could be a double-lined system. The measured radial velocity, $36.82 \pm 2.10 \mathrm{~km} \mathrm{~s}^{-1}$, again agrees with our expected value of $37.5 \mathrm{~km} \mathrm{~s}^{-1}$.
- PPMXL J65.45657+19.48572 (LP 415-20) has been classified as M6.5e Hyades member by Bryja et al. (1994). It is reported as an M7+M9.5 binary with a separation of $0.1^{\prime \prime}$ by Siegler et al. (2003). Dupuy \& Liu (2011) predicts a period
of $14.4 \pm 0.4$ years, so it will be possible to measure its mass. It is detected by XMM (Watson et al. 2009).
- PPMXL J66.57934+17.05058 (LP 415-881) was identified as a possible Hyades member by Reid (1992). It is unresolved by Reid \& Mahoney (2000), using Keck I/HIRES échelle high-resolution spectroscopy to identify spectroscopic binaries. We derive $v_{\perp}=(-2.3 \pm$ 1.1) $\mathrm{km} \mathrm{s}^{-1}$.
- PPMXL J69.74372+15.66145 (LP 415-2128) was identified as a possible Hyades member by Bryja et al. (1992); Reid (1992). It is unresolved by Reid \& Mahoney (2000). The measured radial velocity, $39.5 \pm 0.5 \mathrm{~km} \mathrm{~s}^{-1}$, agrees well with our expected value of $39.7 \mathrm{~km} \mathrm{~s}^{-1}$.
A number of our member candidates have been identified as cool nearby stars, and have published spectral types obtained with

Table 2. Optical and 2MASS photometry of the R11 candidates rejected in our study because of their photometry.

| Name (deg, J2000.0) | $\begin{gathered} g_{\mathrm{P} 1} \text { or } g \\ (\mathrm{mag}) \\ \hline \end{gathered}$ | $r_{\mathrm{P} 1}$ or $r$ (mag) | $i_{\mathrm{P} 1}$ or $i$ <br> (mag) | $z_{\mathrm{pl}}$ or $z$ <br> (mag) | $\begin{gathered} y_{\mathrm{P} 1} \\ (\mathrm{mag}) \\ \hline \end{gathered}$ | surv. | $\begin{gathered} J \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} H \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} K \\ (\mathrm{mag}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $35.415+22.851$ | $18.06 \pm 0.01$ | $16.77 \pm 0.01$ | - | $14.27 \pm 0.00$ | $13.89 \pm 0.01$ | PS1 | $12.57 \pm 0.02$ | $11.96 \pm 0.02$ | $11.64 \pm 0.02$ |
|  | $18.14 \pm 0.01$ | $16.87 \pm 0.01$ | $15.10 \pm 0.00$ | $14.30 \pm 0.00$ | - | SDSS |  |  |  |
|  | - | $16.76 \pm 0.24$ | - | - | - | CMC |  |  |  |
| 60.394-0.726 | - | $16.84 \pm 0.01$ | - | $14.10 \pm 0.01$ | $13.69 \pm 0.01$ | PS1 | $12.35 \pm 0.02$ | $11.75 \pm 0.02$ | $11.48 \pm 0.03$ |
|  | - | $16.92 \pm 0.24$ | - | - | - | CMC |  |  |  |
| $65.320+23.198$ | $17.37 \pm 0.01$ | $16.08 \pm 0.01$ | $14.61 \pm 0.01$ | $13.71 \pm 0.01$ | $13.33 \pm 0.01$ | PS1 | $12.12 \pm 0.02$ | $11.51 \pm 0.02$ | $11.24 \pm 0.02$ |
|  | $17.27 \pm 0.01$ | $16.06 \pm 0.00$ | $14.47 \pm 0.00$ | $13.70 \pm 0.00$ | - | SDSS |  |  |  |
|  | - | $16.19 \pm 0.24$ | - | - | - | CMC |  |  |  |
| $48.638+19.720$ | $17.49 \pm 0.01$ | $16.17 \pm 0.01$ | $14.58 \pm 0.01$ | $13.82 \pm 0.01$ | $13.50 \pm 0.01$ | PS1 | $12.19 \pm 0.02$ | $11.64 \pm 0.03$ | $11.34 \pm 0.02$ |
|  | $17.51 \pm 0.01$ | $16.25 \pm 0.00$ | $14.61 \pm 0.00$ | $13.86 \pm 0.00$ | - | SDSS |  |  |  |
|  | - | $16.45 \pm 0.24$ | - | - | - | CMC |  |  |  |
| $73.164+47.698$ | - | $17.05 \pm 0.01$ | - | $14.54 \pm 0.01$ | $14.19 \pm 0.00$ | PS1 | $12.83 \pm 0.02$ | $12.20 \pm 0.02$ | $11.86 \pm 0.03$ |
|  | - | $17.44 \pm 0.24$ | - | - | - | CMC |  |  |  |
| $84.599+2.402$ | $18.10 \pm 0.01$ | $16.82 \pm 0.01$ | - | $14.33 \pm 0.01$ | $13.96 \pm 0.01$ | PS1 | $12.61 \pm 0.02$ | $12.05 \pm 0.03$ | $11.76 \pm 0.03$ |
|  | - | $16.84 \pm 0.24$ | - | - | - | CMC |  |  |  |
| $85.522+34.978$ | $17.80 \pm 0.01$ | $16.52 \pm 0.01$ | - | $14.39 \pm 0.01$ | $13.78 \pm 0.01$ | PS1 | $12.51 \pm 0.02$ | $11.93 \pm 0.02$ | $11.63 \pm 0.02$ |
|  | - | $16.80 \pm 0.24$ | - | - | - | CMC |  |  |  |

Table 3. PS1 photometry, PPMXL astrometry and PS1 offsets of the R11 candidates rejected in our study because of their astrometry.

| Name <br> $(\mathrm{deg}, \mathrm{J} 2000.0)$ | $g_{\mathrm{Pl}}$ <br> $(\mathrm{mag})$ | $r_{\mathrm{P} 1}$ <br> $(\mathrm{mag})$ | $i_{\mathrm{P} 1}$ <br> $(\mathrm{mag})$ | $z_{\mathrm{Pl}}$ <br> $(\mathrm{mag})$ | $y_{\mathrm{P} 1}$ <br> $(\mathrm{mag})$ | $\mu_{\alpha} \cdot \cos (\delta)$ <br> $\left(\mathrm{mas} \mathrm{yr}^{-1}\right)$ | $\mu_{\delta}$ <br> $\left(\mathrm{mas} \mathrm{yr}^{-1}\right)$ | $\Delta \alpha$ <br> $(\mathrm{mas})$ | $\Delta \delta$ <br> $(\mathrm{mas})$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $64.045+11.534$ | $18.23 \pm 0.01$ | $16.87 \pm 0.01$ | $15.08 \pm 0.01$ | $14.24 \pm 0.01$ | $13.95 \pm 0.01$ | $+119.3 \pm 3.7$ | $-2.2 \pm 3.6$ | +211 | -292 |
| $74.096+14.101$ | $18.72 \pm 0.01$ | - | $15.71 \pm 0.00$ | $14.86 \pm 0.01$ | $14.44 \pm 0.00$ | $+75.3 \pm 4.0$ | $-11.3 \pm 3.8$ | -240 | -401 |
| $65.398+16.894$ | $16.11 \pm 0.01$ | $14.81 \pm 0.00$ | - | $12.45 \pm 0.01$ | $12.22 \pm 0.00$ | $+130.2 \pm 4.5$ | $-37.3 \pm 4.5$ | -584 | +267 |
| $75.715+17.451$ | $17.92 \pm 0.01$ | $16.61 \pm 0.00$ | $14.59 \pm 0.00$ | $13.62 \pm 0.00$ | $13.14 \pm 0.01$ | $+115.6 \pm 5.3$ | $-38.3 \pm 5.1$ | -487 | +98 |
| $71.754+23.263$ | $18.13 \pm 0.01$ | $16.82 \pm 0.01$ | $15.05 \pm 0.01$ | $15.01 \pm 0.01$ | $13.86 \pm 0.01$ | $+95.8 \pm 3.9$ | $-58.1 \pm 3.7$ | -97 | +290 |
| $77.825+39.908$ | $18.39 \pm 0.01$ | $17.14 \pm 0.01$ | $15.47 \pm 0.01$ | $14.77 \pm 0.01$ | $14.33 \pm 0.00$ | $+57.6 \pm 4.8$ | $-75.1 \pm 4.7$ | -235 | +469 |

Table 4. Published spectrophotometric and secular distances for a subset of the candidates with spectroscopic classification.

| Name | Sp.T. | References | Published | West et al. (2005) | Secular |
| :--- | :--- | :--- | :---: | :---: | :---: |
| 2M J0304+0045 | M5.5 | Bochanski et al. (2005) | 31 | 25.1 | 30.0 |
| DE J0304+0045 | M5.5 | Phan-Bao \& Bessell (2006) | 25.5 | 33.7 | 28.3 |
|  | M7 | West et al. (2008) | 25.5 | 20.2 | 28.3 |
| LP 415-20 | M7+M9.5 | Siegler et al. (2003) | $30 \pm 5$ | 27.2 | 40.2 |
| LP 413-53 | M9 | Gizis et al. (2000) | 19.9 | 20.3 | 33.7 |
| 2M J0230+2704 | L0: | Cruz et al. (2007) | $32.0 \pm 3.9$ | 28.7 | 44.4 |
| 2M J0523-1403 | L2.5 | Cruz et al. (2003) | $13.4 \pm 1.1$ | 13.3 | 23.6 |

spectroscopic observations. All those objects are classified as mid-M- to early-L-type low-mass stars, sometimes with hint of youth or binarity. These are described in the next two sections.

### 5.3. L-type dwarfs from Cruz et al. $(2003,2007)$

We identify in our candidate sample two objects published as field L-type dwarfs:

- 2MASSI J0230155+270406 is classified as a peculiar L0: dwarf, with a spectrophotometric distance of $32.0 \pm 3.9 \mathrm{pc}$ by Cruz et al. (2007). From the convergent-point analysis, we estimate its parallax to be $22.5 \pm 0.6$ mas, or $44.4 \pm 1.2 \mathrm{pc}$. This is only $3-\sigma$ away from the spectrophotometric distance, and we keep it as a candidate.
- 2MASS J05233822-1403022 is classified as a L2.5 dwarf in the optical (L5 in the near-infrared), with a spectrophotometric distance of $13.4 \pm 1.1 \mathrm{pc}$ by Wilson et al. (2003) and Cruz et al. (2003). We estimate its parallax to be $42.3 \pm 1.3$ mas, or $23.6 \pm 0.8 \mathrm{pc}$. Its spectrum shows signs of lithium (Reiners \& Basri 2008), and variable signs of X-ray,
radio and $\mathrm{H} \alpha$ emissions (Berger et al. 2010, and ref. therein). No companion was found with separation larger then $0.15^{\prime \prime}$ by Reid et al. (2008a).
Seifahrt et al. (2010) and Blake et al. (2010) measure a radial velocity of $+11.3 \pm 1.0 \mathrm{~km} \mathrm{~s}^{-1}$ and $+12.21 \pm$ $0.09 \mathrm{~km} \mathrm{~s}^{-1}$, while our kinematic analysis predicts a velocity of $+40.6 \mathrm{~km} \mathrm{~s}^{-1}$. Therefore its Hyades membership is questionable.


### 5.4. M-type dwarfs with spectroscopic classification

- LP 245-52 was first identified as a cool star by Cruz \& Reid (2002) and finally classified as a M6 dwarf by Cruz et al. (2003).
- 2MASS J03040207+0045512 (DENIS J030402.0+004551) was first identified as a late-type star by Tinney (1993), and classified as an M5.5 star by Bochanski et al. (2005). Their corresponding spectrophotometric distance is 31 pc , which agrees with our predicted distance of 30.0 pc .
- 2MASS J03300506+2405281 (LP 356-770) was classified as an M7 star by Gizis et al. (2000), with a spectrophotometric distance of 20 pc , smaller than our expected value of 37 pc . The radial velocity of $39.2 \pm 1.1 \mathrm{~km} \mathrm{~s}^{-1}$ (Reid et al. 2002) agrees marginally with our expected value of $31.1 \mathrm{~km} \mathrm{~s}^{-1}$.
- 2MASS J03540135+2316339 has been classified as M8 dwarf by Kirkpatrick et al. (1997). Cruz et al. (2003) estimates a spectrophotometric distance of $22.9 \pm 1.7 \mathrm{pc}$ for a spectral type of M8.5, again smaller than our expected value of 31.05 pc . It is not resolved into a binary (with separation larger then $0.10^{\prime \prime}$ ) by Close et al. (2003). It is also detected by XMM (Watson et al. 2009).
- PPMXL J57.73917+18.30189 (LP 413-53, 2MASS J03505737+1818069) was identified as an M9 star by Gizis et al. (2000). Various spectrophotometric distances have been proposed, close to 19.9 pc (Gizis et al. 2000), smaller (by about $3 \sigma$ ) than our prediction of 33.7 pc . Reid et al. (2002) observed the object when it was undergoing a strong flare. They measure and discuss its radial velocity: it is significant smaller (either $5.8 \pm 2.0 \mathrm{~km} \mathrm{~s}^{-1}$ or $-14.3 \pm 0.4 \mathrm{~km} \mathrm{~s}^{-1}$ depending on the method) than our prediction of $34.4 \mathrm{~km} \mathrm{~s}^{-1}$. However the unusual $\mathrm{H} \alpha$ profile, which may cause the discrepant measurements, is reminiscent of their observation of the Hyades SB2 binary.
- 2MASSI J0411063+124748 was identified as an M6 star by Cruz et al. (2007).
- 2MASS J04273708+2056389 is listed as a M5.75 field dwarf by Luhman (2006), based on youth indicators of the optical and IR spectra, meaning that it is older than a few Myr.
- 2MASS J04441479+0543573 is listed as a M8 dwarf by Reid et al. (2008b). Faherty et al. (2009) gives a proper motion of $\left.\mu_{( } \alpha, \delta\right)=(+95,-6) \pm 21$ mas yr ${ }^{-1}$ in agreement with the more accurate PPMXL measurement of $\left.\mu_{( } \alpha, \delta\right)=$ $(+77.01,+0.73) \pm 5.3 \mathrm{mas} \mathrm{yr}^{-1}$.


### 5.5. Spectrophotometric vs. secular parallaxes

For the candidates with spectroscopic classification, we can compare the spectrophotometric distance with our predicted secular distance. In Table 4, we list our secular distance, the distance we calculate based on the 2MASS $J$-band magnitude, using the $M_{J}$ magnitudes given by West et al. (2005), and the published distance from the literature. Deriving the distance from the spectral type and observed magnitudes, assuming a relationship based on old, field stars, is not an accurate process, given the spread of absolute magnitudes for a given spectral type. For most of the stars in Table 4, the distances differ, marginally for any object, but the spectrophotometric distance seems systematically underestimated.

Previous studies did not notice such large differences between the Hyades cluster members and the older field population: The latest BT-Settl models predict a radius of $R=0.116 R_{\odot}$ at 600 Myr for a mass of $m=0.09 M_{\odot}$, and $R=0.112 R_{\odot}$ for ages of 2 Gyr and older (Allard et al. 2012). Reid \& Mahoney (2000) compared their secular parallaxes with the spectrophotometric distances cited in Reid (1993), typically finding them 5\% larger. Siegler et al. (2003) did not notice any difference between the Hyades and field CMDs, within their uncertainties (their Sect. 4.2).

The differences in the distances may be caused by unresolved binaries. Alternatively, they could be field stars, although in this


Fig. 9. Completeness of PPMXL as a function of $g$ (blue), $r$ (green), and $i$ (red) SDSS magnitudes (solid lines, right to left), and histograms of the $50 \%$ completeness limit (dashed curves, normalized to unity).
case, we would also expect some candidates with an underestimated secular distance. We see none, but the sample size is small.

The low-mass stars of the Hyades may still be in the contraction phase, while the field population has settled on the main sequence. However this low-significance effect needs to be confirmed by our on-going parallax measurements with PS1 (see also Andrei et al. 2011).

## 6. Luminosity function

### 6.1. Detection efficiency

We performed an analysis similar to that of PS1 in Sect. 2.3 for the PPMXL catalogue.

We find that the PPMXL has a $50 \%$ completeness limit of $g=21.5, r=20.3$ and $i=19.5 \mathrm{mag}$, with respective widths of $0.24,0.22$ and 0.18 mag (see Fig. 9). Because PPMXL requires multiple detections in various filters, and also because the USNO-B1.0 photometric system differs from that of PS1 and SDSS, the detection efficiency is a function of the object colour. We find significant incompleteness for the faintest and the reddest objects. Therefore they are typically undetected on the blue plates, and detections on at least one red plate is required. Hence we calculate the completeness based on the $r$-band magnitude, as the $r$ filter profile is the closest to the red plate emulsion sensitivity profile. When that magnitude is unavailable, we extrapolate it from the $i$ or $g$ magnitudes, using the colours of our other Hyades candidates. In practice, the correction is smaller than a factor of 2 for all our candidates with a mass larger than $0.1 M_{\odot}$.

Because of our image examination by eye, because of the limit on the positional offsets (see Sect. 3.2), or because of the rejection of candidates with giant-like near-infrared colours and/or dust reddening, it is possible that a few genuine members are rejected. We do not try to estimate the possible incompleteness due to these selections. This would require a better understanding of the systematics in PPMXL, and an accurate calculation would require a detailed error model which is beyond the scope of this paper.


Fig. 10. Colour-magnitude diagram of our "fake" kinematic candidates. As in Fig. 4b, filled circles indicate candidates selected in $g, r$ and $i$; open circles: candidates selected in one or two filters; triangles: objects selected in $g$ but rejected in $r$ or $i$. The statistics is threefold the actual statistics, as we plot three realizations of the Besançon model.

### 6.2. Contamination by field stars

We run the Besançon model simulation to create a catalogue of "mock" stars with kinematics and Johnson-Cousins BVRIK photometry (Robin et al. 2003). We run the model three times over the whole field covered by our search (Fig. 1). Despite this, the Poisson noise is large at the end of the main sequence at which a small number of fake candidates is selected. We randomize the positions, as the model is calculated over a coarse grid. We transform the Johnson-Cousins photometry into the gri SDSS system according to Jordi et al. (2006). Finally we use the PS1 recovery fraction determined in Sect. 2.2 and the PPMXL completeness (in SDSS $r$ band, see Sect. 6.1) to randomly decide whether a "fake" star has $g, r$ and $i$ detections or not.

We then run our kinematic and photometric selection to obtain a synthetic catalogue of field contaminants (Fig. 10). We find that almost all ( $92 \%$ ) contaminants are dwarfs, and the remaining, giants.

We find that the contamination by field stars is negligible ( $<10 \%$ ) in cluster shells up to 18 pc of the cluster centre (see Fig. 11). R11 considered actual observations of bright stars and the model of Kharchenko \& Schilbach (1996) and reached the same conclusion. At larger distances from the centre, the cluster member density decreases and the survey volume increases, and with it the field contamination. In these outer areas, for distances to the cluster centre between 18 and 30 pc , we find that the contamination increases up, for instance to about $17 \%$ for candidates with $7.5<M_{K}<9.5 \mathrm{mag}$. We stress that the precision of the simulation is rather poor, as the prediction of the Besançon model over such a large area and range of Galactic latitudes is inaccurate. The effect on the luminosity function, however, is small, especially at the very faint end, when Poisson noise and the effect of completeness corrections dominate the error budget.


Fig. 11. Luminosity function of the observed Hyades candidates (solid line) and of the simulated field contaminants (dotted line), for candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ and outside the cluster core, i.e. with a distance to the cluster centre between 3.1 and 30 pc . The light grey area for $M_{K}>9.5$ mag indicates the magnitudes for which our results are strongly affected by incompleteness.

### 6.3. Corrected luminosity function

We derive the cluster luminosity function over various cluster radii, to study its spatial dependence. The core of the cluster is considered to a radius of 3.1 pc . We then consider the volume up to the tidal radius $r_{\mathrm{t}}=9 \mathrm{pc}(\mathrm{R} 11)$. The halo is defined between the tidal radius and $2 r_{\mathrm{t}}=18 \mathrm{pc}$. Finally, the outer volume of our study covers a region between $2 r_{\mathrm{t}}$ and 30 pc .

We first correct the observed luminosity function for the detection efficiency: for each candidate, we obtain the PPMXL detection efficiency corresponding to its PS1 or SDSS $r$-band magnitude, as described above. Here we consider the detection efficiency, as determined in Sect. 6.1. For the 62 candidates found in PS1 or SDSS only, we also include the PS1+SDSS coverage probability. We then give the candidates a weight inverse of the combined detection efficiency.

For each cluster radius range, we then subtract the contamination derived from the simulation, corrected for incompleteness as the candidate sample. The correction of the PPMXL incompleteness is the largest effect at faint magnitudes, and our results for $M_{K}>9.5 \mathrm{mag}$ are very sensitive to the details of the procedure, and therefore not solid. The second largest effect is the correction of the field contamination, which mostly increases with the distance to the cluster centre, here up to $30 \%$ (for $M_{K}=9 \mathrm{mag}$ ).

The result is shown in Fig. 12 for four different shells around the cluster centre, as the system number density per magnitude bin. We emphasize that the luminosity function is our primary observational result, independent of evolutionary and atmospheric modelling. In order to compare with other studies, we discuss the properties of the mass function.
Table 5. The 62 Hyades candidates not included in R11. J2000.0 coordinates from PPMXL (rounded).

| RA (deg, | $\begin{array}{r} \text { Dec } \\ 000.0) \end{array}$ | $\begin{aligned} & \mu_{\alpha} \cdot \cos (\delta) \\ & \left(\text { mas }^{-1}\right) \\ & \hline \end{aligned}$ | $\begin{gathered} \mu_{\delta} \\ \left(\operatorname{mas~yr}^{-1}\right) \end{gathered}$ | $\begin{aligned} & \hline \text { kin. } \pi \\ & \text { (mas) } \end{aligned}$ | $\begin{gathered} g_{\mathrm{P} 1} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} r_{\mathrm{Pl}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} i_{\mathrm{P}_{1}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} K_{\mathrm{s}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} v_{\perp} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \\ \hline \end{gathered}$ | $\begin{gathered} r_{\mathrm{c}} \\ (\mathrm{pc}) \end{gathered}$ | Mass <br> $\left(M_{\odot}\right)$ | Ref. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 28.4133 | +32.3229 | $+214 \pm 4$ | $-24 \pm 4$ | $24.1 \pm 0.4$ | $19.05 \pm 0.01$ | $17.67 \pm 0.01$ | - | $11.84 \pm 0.02$ | -0.70 | 29.1 | 0.131 |  |
| 38.6231 | +36.2197 | $+224 \pm 4$ | $-47 \pm 4$ | $27.2 \pm 0.5$ | $18.86 \pm 0.01$ | $17.52 \pm 0.01$ | - | $11.47 \pm 0.02$ | +1.31 | 24.6 | 0.138 | 1 |
| 38.9566 | -7.1893 | $+284 \pm 5$ | $+82 \pm 5$ | $35.4 \pm 0.6$ | $19.48 \pm 0.01$ | $18.03 \pm 0.01$ | $15.68 \pm 0.01$ | $11.43 \pm 0.03$ | +3.00 | 29.1 | 0.111 | 2, 3, 4 |
| 41.5522 | +5.9406 | $+146 \pm 4$ | $+5 \pm 4$ | $18.4 \pm 0.5$ | $17.52 \pm 0.01$ | $16.25 \pm 0.00$ | $14.56 \pm 0.01$ | $11.14 \pm 0.02$ | -1.23 | 25.0 | 0.242 |  |
| 44.8713 | +8.4708 | $+164 \pm 5$ | $+13 \pm 5$ | $21.7 \pm 0.7$ | $19.73 \pm 0.01$ | $18.37 \pm 0.01$ | $16.09 \pm 0.01$ | $12.02 \pm 0.02$ | +1.66 | 18.5 | 0.134 |  |
| 46.0086 | +0.7642 | $+249 \pm 4$ | $+32 \pm 4$ | $33.3 \pm 0.5$ | $18.29 \pm 0.01$ | $16.91 \pm 0.01$ | $14.82 \pm 0.01$ | $10.88 \pm 0.02$ | -0.27 | 23.5 | 0.146 | 5, 6, 7 |
| 52.5212 | +24.0911 | $+180 \pm 4$ | $-53 \pm 4$ | $26.9 \pm 0.6$ | $19.41 \pm 0.03$ | $17.85 \pm 0.01$ | $15.59 \pm 0.01$ | $11.38 \pm 0.02$ | -1.26 | 14.5 | 0.144 | 8, 9, 10 |
| 55.1910 | +17.5799 | $+157 \pm 4$ | $-30 \pm 4$ | $24.4 \pm 0.6$ | $18.93 \pm 0.01$ | $17.49 \pm 0.01$ | $15.53 \pm 0.01$ | $11.73 \pm 0.02$ | -0.84 | 10.2 | 0.136 |  |
| 55.2289 | +19.4965 | $+163 \pm 5$ | $-32 \pm 5$ | $25.2 \pm 0.8$ | $20.15 \pm 0.04$ |  | $16.40 \pm 0.01$ | $12.24 \pm 0.02$ | +0.20 | 11.0 | 0.108 |  |
| 56.8816 | +12.0197 | $+141 \pm 4$ | $-14 \pm 4$ | $22.6 \pm 0.7$ | $19.44 \pm 0.03$ | $17.98 \pm 0.01$ | $15.75 \pm 0.01$ | $11.64 \pm 0.02$ | -0.89 | 8.9 | 0.151 |  |
| 57.7392 | +18.3019 | $+183 \pm 5$ | $-32 \pm 5$ | $29.7 \pm 0.8$ | $20.71 \pm 0.03$ | $19.04 \pm 0.01$ | $16.64 \pm 0.01$ | $11.78 \pm 0.02$ | +0.76 | 14.2 | 0.112 |  |
| 57.9864 | +41.2047 | $+145 \pm 4$ | -97 | $24.0 \pm 0.5$ |  | $16.88 \pm 0.01$ | $15.05 \pm 0.00$ | $11.54 \pm 0.02$ | -0.37 | 20.1 | 0.149 |  |
| 58.5058 | +23.2761 | $+192 \pm 5$ | $-66 \pm 5$ | $32.2 \pm 0.8$ | $21.00 \pm 0.04$ | $19.49 \pm 0.01$ | $16.83 \pm 0.01$ | $12.00 \pm 0.02$ | $-0.80$ | 16.7 | 0.098 | 11, 12, 13, 14 |
| 59.7924 | +9.2840 | $+143 \pm 4$ | $-1 \pm 4$ | $24.3 \pm 0.7$ | $18.41 \pm 0.01$ | $17.12 \pm 0.01$ | $15.19 \pm 0.01$ | $11.59 \pm 0.03$ | +0.31 | 9.4 | 0.145 |  |
| 60.8216 | -1.8217 | $+147 \pm 4$ | $+41 \pm 4$ | $25.9 \pm 0.6$ | $17.03 \pm 0.01$ | $15.74 \pm 0.01$ | $14.07 \pm 0.01$ | $10.75 \pm 0.02$ | +1.19 | 16.3 | 0.201 |  |
| 62.7767 | +12.7967 | $+125 \pm 5$ | $-18 \pm 5$ | $23.0 \pm 0.9$ | $19.24 \pm 0.02$ | $17.88 \pm 0.01$ | $15.76 \pm 0.01$ | $11.84 \pm 0.03$ | -0.53 | 5.4 | 0.137 | 15 |
| 64.5574 | +25.6610 | $+106 \pm 4$ | $-56 \pm 4$ | $20.9 \pm 0.7$ | $19.70 \pm 0.01$ | $18.27 \pm 0.01$ | $15.99 \pm 0.01$ | $11.86 \pm 0.02$ | -1.27 | 7.8 | 0.148 |  |
| 64.6898 | +21.0318 | $+108 \pm 4$ | $-53 \pm 4$ | $21.7 \pm 0.7$ | $19.71 \pm 0.02$ | $18.36 \pm 0.01$ | $16.12 \pm 0.01$ | $12.21 \pm 0.02$ | -2.97 | 3.9 | 0.124 | 10 |
| 64.7741 | +17.5732 | $+100 \pm 4$ | $-26 \pm 4$ | $19.3 \pm 0.7$ | $18.35 \pm 0.01$ | $17.02 \pm 0.01$ | $15.16 \pm 0.01$ | $11.51 \pm 0.02$ | +0.09 | 5.9 | 0.189 | 12 |
| 65.2000 | +15.2358 | $+161 \pm 5$ | $-59 \pm 5$ | $32.4 \pm 0.9$ | $14.44 \pm 0.01$ | $13.36 \pm 0.01$ |  | $8.63 \pm 0.02$ | -3.76 | 15.6 | 0.446 | 16 |
| 65.4349 | +20.4029 | $+101 \pm 4$ | $-39 \pm 4$ | $20.1 \pm 0.7$ | $20.10 \pm 0.02$ | $18.69 \pm 0.01$ | $16.58 \pm 0.01$ | $12.51 \pm 0.02$ | $-0.80$ | 4.8 | 0.118 | 17, 18 |
| 65.4412 | +20.3957 | $+95 \pm 5$ | $-33 \pm 5$ | $18.6 \pm 1.0$ | $21.21 \pm 0.06$ | $19.79 \pm 0.02$ | $17.30 \pm 0.01$ | $12.93 \pm 0.03$ | -0.07 | 8.1 | 0.107 | 10 |
| 65.4566 | +19.4857 | $+125 \pm 4$ | $-48 \pm 4$ | $24.9 \pm 0.7$ | $19.78 \pm 0.02$ | $18.35 \pm 0.01$ | $15.98 \pm 0.01$ | $11.67 \pm 0.02$ | -1.34 | 6.6 | 0.137 | 19, 20, 21, 22 |
| 66.5793 | +17.0506 | +97 | $-37 \pm 4$ | $20.3 \pm 0.8$ | $18.94 \pm 0.01$ | $17.59 \pm 0.01$ | $15.60 \pm 0.01$ | $11.91 \pm 0.03$ | -2.30 | 3.0 | 0.149 | 17, 18 |
| 66.9046 | +20.9441 | $+89 \pm 6$ | $-36 \pm 6$ | $18.3 \pm 1.1$ | $19.82 \pm 0.02$ | $18.47 \pm 0.01$ | $16.18 \pm 0.01$ | $12.07 \pm 0.02$ | -0.49 | 9.1 | 0.153 | 23, 10 |
| 67.8182 | +15.0034 | $+105 \pm 4$ | $-20 \pm 4$ | $21.9 \pm 0.9$ | $18.69 \pm 0.01$ | $17.33 \pm 0.01$ | $15.39 \pm 0.01$ | $11.71 \pm 0.02$ | +0.77 | 1.6 | 0.151 | 12 |
| 68.8065 | +20.1337 | $+99 \pm 6$ | $-33 \pm 6$ | $21.0 \pm 1.1$ | $21.06 \pm 0.04$ | $19.58 \pm 0.02$ | $17.12 \pm 0.01$ | $12.61 \pm 0.02$ | +1.39 | 3.4 | 0.109 |  |
| 68.9657 | +12.2556 | $+93 \pm 6$ | $-20 \pm 6$ | $20.5 \pm 1.2$ |  | $19.11 \pm 0.04$ | $16.81 \pm 0.01$ | $12.68 \pm 0.03$ | -1.43 | 4.6 | 0.109 |  |
| 69.2422 | +23.2774 | $+90 \pm 5$ | $-45 \pm 5$ | $19.8 \pm 1.0$ | $20.23 \pm 0.03$ | $18.72 \pm 0.01$ |  | $12.57 \pm 0.02$ | +0.11 | 7.1 | 0.117 |  |
| 69.7437 | +15.6614 | $+91 \pm 6$ | $-26 \pm 6$ | $20.4 \pm 1.3$ | $19.44 \pm 0.02$ | $18.12 \pm 0.01$ | $16.00 \pm 0.01$ | $12.20 \pm 0.03$ | $-0.40$ | 3.5 | 0.131 | 24,17,18 |
| 70.1052 | -4-6222 | $+63 \pm 5$ | $+32 \pm 5$ | $14.8 \pm 1.1$ | $20.04 \pm 0.02$ | $18.61 \pm 0.02$ | $16.59 \pm 0.01$ | $12.74 \pm 0.03$ | +2.10 | 30.0 | 0.140 |  |
| 71.0616 | +5.7326 | $+77 \pm 5$ | $+1 \pm 5$ | $17.9 \pm 1.2$ | $21.97 \pm 0.19$ | $20.07 \pm 0.03$ | $17.46 \pm 0.01$ | $12.52 \pm 0.03$ | -0.97 | 14.0 | 0.129 | 5,26 |
| 71.5113 | +19.4683 | $+76 \pm 6$ | $-41 \pm 6$ | $18.8 \pm 1.2$ | $19.52 \pm 0.02$ | $18.22 \pm 0.01$ | $16.09 \pm 0.01$ | $12.30 \pm 0.02$ | -2.08 | 8.2 | 0.135 |  |
| 71.5402 | +18.9578 | $+79 \pm 6$ | $-38 \pm 6$ | $19.3 \pm 1.3$ | $21.01 \pm 0.03$ | $19.41 \pm 0.01$ | $17.00 \pm 0.01$ | $12.72 \pm 0.03$ | -1.14 | 6.9 | 0.112 |  |
| 71.6875 | +24.6112 | $+90 \pm 6$ | $-61 \pm 6$ | $22.1 \pm 1.2$ | $19.73 \pm 0.02$ | $18.43 \pm 0.01$ | $16.14 \pm 0.01$ | $12.13 \pm 0.02$ | -1.12 | 7.2 | 0.126 | 10 |
| 71.9778 | -0.7934 | $+76 \pm 5$ | $+34 \pm 5$ | $19.0 \pm 1.2$ |  | $20.39 \pm 0.06$ | $17.68 \pm 0.01$ | $12.99 \pm 0.03$ | +2.80 | 16.7 | 0.103 |  |
| 72.1916 | +11.3857 | $+97 \pm 4$ | $-25 \pm 4$ | $24.0 \pm 1.0$ | $19.11 \pm 0.01$ | $17.73 \pm 0.01$ | $15.57 \pm 0.01$ | $11.53 \pm 0.03$ | -2.02 | 7.3 | 0.150 |  |
| 72.4634 | -16.5419 | $+90 \pm 4$ | $+74 \pm 4$ | $21.7 \pm 0.8$ |  |  | $13.92 \pm 0.01$ | $10.77 \pm 0.02$ | -0.52 | 26.7 | 0.243 |  |
| 73.4205 | -4.4760 | $+82 \pm 4$ | $+40 \pm 4$ | $21.1 \pm 1.0$ | $17.66 \pm 0.01$ | $16.33 \pm 0.01$ | $14.70 \pm 0.01$ | $11.35 \pm 0.02$ | +0.69 | 18.0 | 0.187 |  |
| 75.1623 | +24.3828 | $+66 \pm 4$ | $-52 \pm 4$ | $18.6 \pm 1.0$ | $19.11 \pm 0.01$ | $17.72 \pm 0.01$ | $15.82 \pm 0.01$ | $12.18 \pm 0.02$ | -1.14 | 11.9 | 0.144 | 10 |
| 75.1655 | +9.1261 | $+52 \pm 6$ | $-5 \pm 6$ | $14.2 \pm 1.5$ | $19.14 \pm 0.01$ | $17.79 \pm 0.01$ | $15.85 \pm 0.00$ | $12.19 \pm 0.03$ | -0.30 | 26.6 | 0.188 |  | References. (1) Cruz \& Reid (2002); (2) Phan-Bao \& Bessell (2006); (3) West et al. (2008); (4) Gálvez-Ortiz et al. (2010); (5) Tinney (1993); (6) Bochanski et al. (2005); (7) Deacon et al. (2009);

(8) Gizis et al. (2000); (9) Reid et al. (2002); (10) Slesnick et al. (2006); (11) Kirkpatrick et al. (1997); (12) Bouvier et al. (2008); (13) Cruz et al. (2003); (14) Close et al. (2003); (15) Cruz et al. 2007); (16) Giclas et al. (1962); (17) Reid (1992); (18) Reid \& Mahoney (2000); (19) Bryja et al. (1994); (20) Siegler et al. (2003); (21) Watson et al. (2009); (22) Dupuy \& Liu (2011); (23) Luhman (2006); (24) Bryja et al. (1992); (25) Reid et al. (2008b); (26) Faherty et al. (2009); (27) Reid et al. (2008a); (28) Seifahrt et al. (2010); (29) Blake et al. (2010).
Table 5. continued.


Fig. 12. Corrected luminosity function (density of candidates with $\left|v_{\perp}\right|<$ $2 \mathrm{~km} \mathrm{~s}^{-1}$ ) of the Hyades in the 2MASS $K_{\mathrm{s}}$ band, for regions of various cluster radii: 0 to $3.1 \mathrm{pc} ; 3.1$ to $9 \mathrm{pc} ; 9$ to 18 pc ; and 18 to 30 pc (top to bottom). The light grey area for $M_{K}>9.5$ mag indicates the magnitudes for which our results are strongly affected by incompleteness.

## 7. Mass function

We transform the luminosity of the new candidates using the mass-luminosity relation provided by Baraffe et al. (1998), used in the $K_{\mathrm{s}}$ band, as the models prove to better match the observations for that filter (R11). For the previously published candidates we use the masses reported in R11. We also try the new BT-Settl isochrones of Allard et al. (2012). We find only slightly larger masses, with the difference picking at less than $5 \%$ around $0.3 M_{\odot}$. We assume that the objects are individual stars and do not make correction for binarity, because of the lack of specific statistics describing the binarity as a function of spectral type.

We show the results in Fig. 13, again for the same four different cluster regions, and for the complete survey up to 30 pc . Here we consider the mass function as the number of systems between $\log m$ and $\log m+\mathrm{d} \log m$, as a function of the logarithm of their mass (Salpeter 1955):
$\xi(\log m)=\frac{\mathrm{d} N(\log m)}{\mathrm{d} \log m} \propto m^{-\Gamma}=m^{-(\alpha-1)}$.
For the volume within 30 pc , the mass function is well (minimum $\chi^{2}$ ) fitted for its high-mass ( $0.80-3.2 M_{\odot}$ ) part by a power law with (logarithmic) slope $\Gamma=2.53 \pm 0.30$ (or $\Gamma=2.21 \pm 0.37$ over $\left.0.80-2.0 M_{\odot}\right)$, steeper than the canonical Salpeter mass function (in this logarithmic representation the Salpeter powerlaw would have a slope of $\Gamma=1.35$, i.e. $\alpha=2.35$ ). In the intermediate part of the mass spectrum, 0.13-1.2 $M_{\odot}$, the mass function is nearly flat, with $\Gamma=0.15 \pm 0.06$. The error estimates on the slope are consistent with the more sophisticated analysis of Weisz et al. (2013).

Alternatively we can fit the range of 0.13-3.2 $M_{\odot}$ with a lognormal function (Chabrier 2003),
$\xi(m)=\mathrm{e}^{-\left(\log m-\log m_{c}\right)^{2} / 2 \sigma^{2}}$,
with parameters: $m_{\mathrm{c}}=(0.44 \pm 0.03) M_{\odot}$ and $\sigma=0.39 \pm 0.02$.
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Fig. 13. Corrected mass function (number of candidates with $\left|v_{\perp}\right|<$ $2 \mathrm{~km} \mathrm{~s}^{-1}$ ) of the Hyades in the 2MASS $K_{\mathrm{s}}$ band, for regions of various cluster radii: 0 to 3.1 pc (circle); 3.1 to 9 pc (filled square); 9 to 18 pc (triangle); and 18 to 30 pc (open square). The top histogram shows the complete mass function up to 30 pc , fitted with two power laws (thick lines), and a log-normal mass function (dashed line). The Salpeter mass function is shown for comparison (dotted line). The grey areas are the 1$\sigma$ locus of the simulation of Ernst et al. (2011) for the cluster centre ( $r_{\mathrm{c}}<3 \mathrm{pc}$, bottom) and the whole volume ( $r_{\mathrm{c}}<30 \mathrm{pc}$, top). The light grey area for $m<0.1 M_{\odot}$ indicates the masses for which our results are strongly affected by incompleteness.

Finally the last two mass bins with good completeness, $m=0.10-0.15 M_{\odot}$ show a deviation from the power-law mass function, and a significant deficiency in very-low-mass stars.

Ernst et al. (2011) modelled the R11 Hyades mass function using an $N$-body simulation of the cluster assuming an analytic Milky Way potential and the Kroupa (2001) initial mass function down to $0.08 M_{\odot}$. The agreement is good over the whole volume surveyed ( $r_{\mathrm{c}}<30 \mathrm{pc}$, see Fig. 13). However the observed knee at $1 M_{\odot}$ is smooth in the simulation, which may be due to unresolved binaries. In addition, in the mass bin $0.1-0.12 M_{\odot}$, we see a clear decrease in the mass function, while the model displays a small increase. This conclusion is only based on one mass bin in our own data, but seems consistent with the results of Bouvier et al. (2008). Similarly, for the cluster centre ( $r_{\mathrm{c}}<3 \mathrm{pc}$ ), the simulation predicts a flat mass function down to the end of the main sequence, while we clearly see a continuous decrease for lowmass stars. However the simulation by Ernst et al. (2011) did not take into account encounters with large molecular clouds, which may strip the cluster of its low-mass members, and this could the reason for the observed discrepancy. Alternatively, assuming that a fraction of stars are members of unresolved binaries also leads to a decrease of the (system) mass function at the low-mass end (Ernst, priv. comm.).

In Fig. 14, we compare our mass functions with published results of another nearby and well observed intermediate-age cluster, the Praesepe cluster (Boudreault et al. 2012), as well as previous mass functions of the Hyades (Bouvier et al. 2008; R11).


Fig. 14. Mass functions of the Hyades: this study (thick histogram, $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$ for consistency); R11 (thin dashed line); and Bouvier et al. (2008) (triangle); and of the Praesepe cluster: Boudreault et al. (2012) (square). The light grey area for $m<0.1 M_{\odot}$ indicates the masses for which our results are strongly affected by incompleteness.

Compared to R11, we determine the mass function for lower masses, with the consequence that the mass function plateau is extended towards lower masses, around $0.13 M_{\odot}$. Our result is similar to the Hyades mass function of Bouvier et al. (2008), but with better statistics thanks to the large volume probed. Therefore insignificant features such as the dip at $0.25 M_{\odot}$ is not seen in our mass function. In addition, our mass function is flatter because of the contribution of the Hyads at large cluster radii.

Comparing with the mass function of Boudreault et al. (2012), the Praesepe cluster seems to have retained its low-mass members at a higher rate, or has a steeper initial mass function, than the Hyades cluster. We point out that the different surveys use different selection methods and, of course, data sets. Thus they have different contamination rate and completeness. In that respect, our combined use of proper motion and multiband photometry for a nearby cluster, makes our study less prone to systematic errors than purely photometric studies.

## 8. Spatial structure

The total mass of the additional 62 candidates, within 30 pc , is $11 M_{\odot}$. By correcting for incompleteness as outlined above, we find about $14 M_{\odot}$. Within 9 pc of the centre, the 18 new candidates make up less than $3 M_{\odot}$.

In both cases, this is a small correction to the previous estimate by R11, i.e. $276 M_{\odot}$ within 9 pc . Therefore, we do not update the analysis presented in R11 as the results would be very similar.

In Fig. 15 we show the distribution of the candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ in cartesian Galactic coordinates, colour-coded by mass. It is clear that massive (blue) stars concentrate on the


Fig. 15. Projection of the candidates with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ on the Galactic coordinate planes. The colour scales with the mass (inner candidates within 9 pc of the centre are shown as smaller dots for clarity). The 9 -pc- and 18 -pc-radius circles are shown as dashed lines.
centre, while low-mass members are distributed over the whole area. The ellipticity of the cluster is clearly visible (Fig. 6 of Ernst et al. 2011).

As a proxy for mass segregation we use the minimum spanning tree (MST) method $\mathcal{M}_{\mathrm{MST}}^{\mathrm{P}}$ developed by Olczak et al. (2011). The MST is the graph connecting all $n$ vertices within a given sample $\zeta$ with the lowest possible sum of the $(n-1)$ edge lengths $e_{i, \zeta}$ excluding closed loops (Gower \& Ross 1969). The authors use the geometrical mean to define a measure of the MST of a sample $\zeta$,
$\gamma_{\mathrm{MST}}^{\zeta}=\exp \left[\frac{1}{n} \sum_{i=1}^{n} \ln e_{i, \zeta}\right]$,
with the corresponding statistics for a set of $m$ equal-sized samples $\zeta_{j}$,

$$
\begin{align*}
\left\langle\gamma_{\mathrm{MST}}^{\zeta_{j}}\right\rangle & =\frac{1}{m} \sum_{j=1}^{m} \gamma_{\mathrm{MST}}^{\zeta_{j}}, \\
\Delta \gamma_{\mathrm{MST}}^{\zeta_{j}} & =\sqrt{\frac{1}{m-1} \sum_{j=1}^{m}\left(\gamma_{\mathrm{MST}}^{\zeta_{j}}-\left\langle\gamma_{\mathrm{MST}}^{\zeta_{j}}\right\rangle\right)^{2}} . \tag{2}
\end{align*}
$$

Based on these quantities we define a normalized measure of mass segregation of a population of $N$ stars,

$$
\begin{align*}
\Gamma_{\mathrm{MST}} & =\frac{\left\langle\gamma_{\mathrm{MST}}^{\mathrm{ref}}\right\rangle}{\gamma_{\mathrm{MST}}^{\text {mass }}},  \tag{3}\\
\Delta \Gamma_{\mathrm{MST}} & =\frac{\Delta \gamma_{\mathrm{MST}}^{\text {ref }}}{\gamma_{\mathrm{MST}}^{\text {mass }}},
\end{align*}
$$

where the superscript "ref" refers to an ensemble of $m$ samples of $n \leq N$ random stars, while the superscript "mass" refers to a sample of the $n$ most massive stars. The number $m$ is chosen such that the ensemble of random stars is representative of the entire population. Hence, $\Gamma_{\mathrm{MST}}$ is a normalized measure of the concentration of the $n$ most massive stars relative to a representative ensemble of random stars. A value of one marks the unsegregated state. The larger $\Gamma_{\mathrm{MST}}$ the more concentrated are the massive stars compared to the reference ensemble. The associated standard deviation $\Delta \Gamma_{\text {MST }}$ quantifies the significance of the result.

Figure 16 shows the mass segregation analysis of two sets of our observational sample (points with error bars) compared to
the observational data of R11 as published in Ernst et al. (2011) (black line with shaded error region). The data in the upper panel are based on the full sample of 786 stars (with $\left|v_{\perp}\right|<4 \mathrm{~km} \mathrm{~s}^{-1}$ ), consisting of 724 candidates from R11 and our 62 new low-mass candidates, mostly located in the outer cluster parts. The effect of these new candidates is to increase the normalized mass segregation measure $\Gamma_{\text {MST }}$ of the more concentrated subsamples as is most evident for the five most massive stars (red). The result for the reduced sample of the "best" 669 candidates (with $\left|v_{\perp}\right|<2 \mathrm{~km} \mathrm{~s}^{-1}$ ) is presented in the lower panel ${ }^{5}$. We have further taken into account contamination by field stars in a statistical approach yet found no significant effect. The more conservative membership selection makes the shape of the distribution of the differential $\Gamma_{\text {MST }}$ much more regular. This characteristic signature of dynamically evolved mass segregation (see e.g. Ernst et al. 2011) implies a larger purity of the reduced sample and hence justifies the lower velocity threshold.

Our main findings from the mass segregation analysis of the revised Hyades candidates are the following:

1. The cumulative plot shows a significant degree of mass segregation for all samples. Comparison with the differential samples demonstrates the exceptionally compact configuration of the five most massive stars (red). The steady decline of $\Gamma_{\text {MST }}$ with stellar mass for the differential samples is a characteristic signature of dynamically evolved mass segregation (cf. Ernst et al. 2011).
2. The least massive Hyades candidates are clearly showing "inverse" mass segregation: $\Gamma_{\mathrm{MST}}<1$ for the last bin in the differential plot. This is another characteristic signature of dynamically evolved mass segregation that has removed the lowest mass members from the inner cluster parts.
3. The signature of mass segregation agrees fairly well with a moderately ( $S=0.3$ : cf. Subr et al. 2008) mass segregated star cluster model with 1000 members as shown in the middle panel of Fig. 4 in Olczak et al. (2011).
4. Our extended Hyades membership sample with 48 additional low-mass cluster candidates compared to the data of R11 demonstrates that the Hyades are stronger mass segregated than estimated before (Ernst et al. 2011).

[^5]

Fig. 16. Mass segregation analysis of our Hyades data and comparison with Ernst et al. (2011). Top: full sample (candidates with $\left|v_{\perp}\right|<$ $4 \mathrm{~km} \mathrm{~s}^{-1}$ ). Bottom: reduced sample with velocity cut as explained in the text and statistical removal of contaminants. The left hand-side shows the "cumulative" $\Gamma_{\text {MST }}$ (points) for the 5 (red), 10 (green), 20 (blue), 50 (magenta), 100 (cyan), 200 (brown), 500 (orange) most massive, and all stars (black). The right-hand side shows the "differential" $\Gamma_{\mathrm{MST}}$ (points) for the 5 (red), 6-10 (green), 11-20 (blue), 21-50 (magenta), 51-100 (cyan), 101-200 (brown), 201-500 (orange) most massive stars, and the remainder (black). The error bars mark the $1-\sigma$ uncertainties $\Delta \Gamma_{\text {MST }}$. The black solid line and the grey region represent $\Gamma_{\text {MST }}$ and $\Delta \Gamma_{\text {MST }}$ for the data of R11, respectively.

## 9. Conclusion

We have combined the PPMXL and Pan-STARRS1 catalogues to search for low-mass Hyades members up to 30 pc from the cluster centre. We select candidates based on the PPMXL kinematics and Pan-STARRS1 astrometry and photometry, combined with 2MASS, WISE and SDSS photometry, to produce a clean sample of candidates nearly complete down to $0.1 M_{\odot}$.

We discover 54 new candidates, with an average mass of $0.18 M_{\odot}$ and an average distance from the cluster centre of 18 pc . We select a purer sample of 43 new candidates with velocity perpendicular to the Hyades motion smaller than $2 \mathrm{~km} \mathrm{~s}^{-1}$. This doubles the number of Hyades candidates for masses smaller than $0.15 M_{\odot}$.

The combination of the astrometric selection of the fast Hyades members, relatively to the field stars, and multi-band photometric selection allows to increase the purity of our candidate sample to larger cluster radii, compared to others distant or slow clusters, and compared to studies using more limited data sets. Using the Besançon model to produce a fake catalogue of field stars, the estimated field dwarf contamination is
below $20 \%$ for the M-type dwarf in the outer shell, with distances between 18 and 30 pc , and well below $10 \%$ up to 18 pc in any mass bin. Our results are based on a detailed analysis of the sensitivity and contamination of our survey.

We find that the mass function is nearly flat for masses between 0.1 and $1 M_{\odot}$, with $-\alpha=0.15 \pm 0.06$. The cluster is clearly segregated in mass, with the lowest-mass members having been removed from the cluster centre.

We identify one known L-type dwarf, 2MASSI J023 0155+270406, as a likely Hyades member. This object, close enough to be accurately characterized, is a benchmark object at the frontier between low-mass stars and brown dwarfs, with a known age, metallicity and distance.

The Hyades cluster is close enough that PS1 will eventually measure the parallaxes of all low-mass star and L-type-dwarf cluster candidates within 30 pc of the centre. This will provide a strong confirmation for memberships, and allow a detailed analysis of the kinematics, binarity, activity (Seemann et al. 2012) of the cluster members. In the pre-Gaia area, the Hyades will therefore be a unique test of cluster evolution and of pre-main sequence evolution of low-mass stars.

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[^0]:    * Table of canditates is only available at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via
    http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/559/A43

[^1]:    1 http://www.ps1sc.org

[^2]:    ${ }^{2}$ Stars with good PSF fit, good background measurement, that are not saturated, not falling on chip defects, cosmic rays, diffraction spikes, electronic ghosts, or gaps.

[^3]:    3 Table 6 is only available in electronic form at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via
    http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/559/A43/

[^4]:    ${ }^{4}$ http://phoenix.ens-lyon.fr/Grids/BT-Settl/
    CIFIST2011/ISOCHRONES/

[^5]:    5 We recall that the massive binary member 77 Tau (PPMXL 2992334831999166121) is also part of this sample as its large $\left|v_{\perp}\right|=2.3 \mathrm{~km} \mathrm{~s}^{-1}$ is caused by the orbital motion.

