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# Triple Conformal Geometric Algebra for Cubic Plane Curves 

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#### Abstract

Summary The Triple Conformal Geometric Algebra (TCGA) for the Euclidean $\mathbb{R}^{2}$-plane extends CGA as the product of three orthogonal CGAs, and thereby the representation of geometric entities to general cubic plane curves and certain cyclidic (or roulette) quartic, quintic, and sextic plane curves. The plane curve entities are 3-vectors that linearize the representation of non-linear curves, and the entities are inner product null spaces (IPNS) with respect to all points on the represented curves. Each IPNS entity also has a dual geometric outer product null space (OPNS) form. Orthogonal or conformal (angle-preserving) operations (as versors) are valid on all TCGA entities for inversions in circles, reflections in lines, and, by compositions thereof, isotropic dilations from a given center point, translations, and rotations around arbitrary points in the plane. A further dimensional extension of TCGA, also provides a method for anisotropic dilations. Intersections of any TCGA entity with a point, point pair, line or circle are possible. TCGA defines commutator-based differential operators in the coordinate directions that can be combined to yield a general $\mathbf{n}$-directional derivative.


## KEYWORDS:

Clifford algebra, cubic plane curve, conformal geometric algebra, 3 -vector entity

## 1 | INTRODUCTION

This paper ${ }^{1}$ assumes familiarity with Clifford's geometric algebras (GA) $\mathcal{G}_{p, q}[14]$ over an $(n=p+q)$-dimensional ( $n$ - D ) pseudoEuclidean ${ }^{2}$ vector space $\mathbb{R}^{p, q}$ over real numbers $\mathbb{R}$, and with the general theory of the corresponding Conformal Geometric Algebra (CGA) [23] $\mathcal{G}_{p+1, q+1}$, especially for an $(n=p)$-D Euclidean vector space $\mathbb{R}^{n}$. References [3] and [25] are books that include full discussions of CGA $\mathcal{C}_{4,1}$ for modeling three-dimensional Euclidean space $\mathbb{R}^{3}$. The paper [19] shows different aspects of $\mathcal{G}_{4,1}$, its quaternionic subalgebra structure, its relationship with Minkowski space algebra, and a self-contained Java implementation. Moreover, [21] provides a short self-contained general purpose tutorial for Clifford's geometric algebra (including CGA). See also [20], with an emphasis on common geometric features of the basic flat and round entities in CGA.

As further background, it is also helpful to have familiarity with the recently introduced Double Conformal Geometric Algebra $(\mathrm{DCGA}) \mathcal{C}_{2(3+1), 2(0+1)}=\mathcal{C}_{8,2}$ [5] for $\mathbb{R}^{3}$, and with the Double Conformal Space-Time Algebra $(\mathrm{DCSTA}) \mathcal{C}_{2(1+1), 2(3+1)}=\mathcal{C}_{4,8}$ for the Minkowski pseudo-Euclidean vector space $\mathbb{R}^{1,3}$ [9][6].

The CGA $\mathcal{G}_{p+1, q+1}$ non-linearly embeds the vector space $\mathbb{R}^{p, q}$ for which geometric entities are modeled conformally.

[^0]The use of CGA can be motivated ${ }^{3}$ by a consideration of stereographic embedding: In [25], the conformal embedding $\mathcal{C}(\mathbf{t})$ is defined as a stereographic embedding followed by a Minkowski homogenization. The stereographic embedding of $\mathbf{t} \in \mathbb{R}^{2}$, denoted $S(\mathbf{t})$, is the point on the unit 2 -sphere centered on the origin in $\mathbb{R}^{3}\left(\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{+}\right)$where the line $\mathbf{e}_{+}+t \hat{\mathbf{d}}$, with $\mathbf{d}=\mathbf{t}-\mathbf{e}_{+}$, intersects the unit circle around the origin in the $\hat{\mathbf{t}}_{+}-$plane. $\mathcal{S}(\mathbf{t})$ can be solved by similar triangles as

$$
\begin{equation*}
S(\mathbf{t})=\frac{2\|\mathbf{t}\|}{\|\mathbf{t}\|^{2}+1} \hat{\mathbf{t}}+\frac{\|\mathbf{t}\|^{2}-1}{\|\mathbf{t}\|^{2}+1} \mathbf{e}_{+} \tag{1}
\end{equation*}
$$

The Minkowski homogenization is $\mathcal{H}_{\mathcal{M}}(\mathbf{t})=\mathcal{S}(\mathbf{t})+\mathbf{e}_{-} \in \mathbb{R}^{3,1}$. Since $\mathcal{H}_{\mathcal{M}}(\mathbf{t})$ is homogeneous, it can be scaled by an arbitrary non-zero scalar without affecting the point that is represented. The common practice is to define the final conformal embedding as

$$
\begin{align*}
\mathcal{C}(\mathbf{t}) & =\frac{\|\mathbf{t}\|^{2}+1}{2}\left(S(\mathbf{t})+\mathbf{e}_{-}\right)  \tag{2}\\
& =\frac{\|\mathbf{t}\|^{2}+1}{2}\left(\frac{2\|\mathbf{t}\|}{\|\mathbf{t}\|^{2}+1} \hat{\mathbf{t}}+\frac{\|\mathbf{t}\|^{2}-1}{\|\mathbf{t}\|^{2}+1} \mathbf{e}_{+}+\mathbf{e}_{-}\right)  \tag{3}\\
& =\mathbf{t}+\frac{\mathbf{t}^{2}}{2}\left(\mathbf{e}_{-}+\mathbf{e}_{+}\right)+\frac{1}{2}\left(\mathbf{e}_{-}-\mathbf{e}_{+}\right)  \tag{4}\\
& =\mathbf{t}+\frac{\mathbf{t}^{2}}{2} \mathbf{e}_{\infty 1}+\mathbf{e}_{o 1} . \tag{5}
\end{align*}
$$

For $\|\mathbf{t}\|=0$, then $\mathcal{C}(\mathbf{t}) \rightarrow \mathbf{e}_{o 1}$, where

$$
\begin{equation*}
\mathbf{e}_{o 1}=\frac{1}{2}\left(\mathbf{e}_{-}-\mathbf{e}_{+}\right) \tag{6}
\end{equation*}
$$

is the entity for the point at the origin. In the limit as $\|\mathbf{t}\| \rightarrow \infty$, then $2 C(\mathbf{t}) /\|\mathbf{t}\|^{2} \rightarrow \mathbf{e}_{\infty 1}$, where

$$
\begin{equation*}
\mathbf{e}_{\infty 1}=\mathbf{e}_{-}+\mathbf{e}_{+} \tag{7}
\end{equation*}
$$

is the entity for the point at infinity. In CGA2 and CGA3, the corresponding points are named $\left\{\mathbf{e}_{\infty 2}, \mathbf{e}_{o 2}\right\}$ and $\left\{\mathbf{e}_{\infty 3}, \mathbf{e}_{o 3}\right\}$, respectively. Since $\mathcal{S}(\mathbf{t})$ is on the unit circle, then $\|\mathcal{S}(\mathbf{t})\|=1$ and $\left(\mathcal{S}(\mathbf{t})+\mathbf{e}_{-}\right)^{2}=0$ such that $\mathcal{C}(\mathbf{t})$ is a null vector.

The CGA $\mathcal{C}_{p+1, q+1}$ can also be extended by multiplicity $k$ to a $k$-CGA $\mathcal{C}_{k(p+1), k(q+1)}$ by multiplying $k$ orthogonal copies of CGA, each with the same scalar coefficients on corresponding CGA basis blades. The main motivation for this approach is, that in $k$-CGA one can formulate $k$-vectors by linear combination of $k$-vector extraction operators for coordinate polynomials of degree $k$ and (selectively) up to degree $2 k$. These $k$-vectors represent algebraic curves of degree $k$ to $2 k$, which is generally not possible in standard CGA. The advantages of this formulation are the complete freedom to use linear versor operators for translation, rotation, dilation, reflection, inversion and transversion, and the freedom of intersection with standard CGA entities is preserved.

The paper is organized as follows. Section 2 introduces the (simple) CGA model for the Euclidean plane $\mathbb{R}^{2}$, and provides some motivation for the non-linear embedding of points in CGA. Next, Section 3 describes a triple version of the CGA model for the Euclidean plane $\mathbb{R}^{2}$, called TCGA. It describes how plane curves of degree three to six can be encoded as 3 -vectors, how these curves can be transformed, intersected with simple CGA entities, and differentiated. It also includes selected examples of these curves, their TCGA expressions and graphs. Finally, Section 4 explains several aspects of computer algebra implementation of triple CGA. We conclude with Section 5. Appendix A explains computational details of efficient versor operations.

## 2 | CGA $\mathcal{G}_{2+1,1}$ OF THE EUCLIDEAN PLANE $\mathbb{R}^{2}$

The CGA $\mathcal{G}_{p+1, q+1}$ non-linearly embeds the $(n=p+q)$-dimensional vector space $\mathbb{R}^{p, q}$ for which geometric entities are modeled conformally ${ }^{4}$. These geometric entities in CGA represent points, point pairs, circles, spheres, hyperplanes, hy$\operatorname{per}\left(\mathrm{pseudo}\right.$ )spheres, and their intersections (by wedge products) in $\mathbb{R}^{p, q}$ (see [18]).

The Conformal Geometric Algebra (CGA) $\mathcal{G}_{3,1}$ of the Euclidean plane $\mathbb{R}^{2}$, also called the Compass Ruler Algebra [16], provides entities for points, lines, circles, and their intersections. The basis of $\mathbb{R}^{3,1}$ is given by four orthonormal vectors

[^1]$\left\{\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}=\mathbf{e}_{+}, \mathbf{e}_{4}=\mathbf{e}_{-}\right\}$, where the first three square to +1 , and the last to -1 . The unit pseudoscalar of the Euclidean plane $\mathbb{R}^{2}$ is $\mathbf{I}_{\mathcal{E}}=\mathbf{e}_{1} \mathbf{e}_{2}$. The unit pseudoscalar of $\mathcal{G}_{3,1}$ is $\mathbf{I}_{C}=\mathbf{e}_{1} \mathbf{e}_{2} \mathbf{e}_{+} \mathbf{e}_{-}$.

Notation 1. The symbolic test vector in the Euclidean plane is denoted by $\mathbf{t}=\mathbf{t}_{\mathcal{E}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}$, and a specific point by $\mathbf{p}=\mathbf{p}_{\mathcal{E}}=$ $p_{x} \mathbf{e}_{1}+p_{y} \mathbf{e}_{2}$.
Notation 2. The reverse $A^{\sim}$ of a multivector $A$ reverses the order of all vector products in $A$ (e.g., $\mathbf{I}_{\mathcal{E}}^{\sim}=\mathbf{e}_{2} \mathbf{e}_{1}$; in this case, $\mathbf{I}_{\mathcal{E}}^{\sim}=\mathbf{I}_{\mathcal{E}}^{-1}=-\mathbf{I}_{\mathcal{E}}$.

In the following subsections, we will introduce relevant multivector entities in CGA for geometric objects and for transformation operators.

## 2.1 | 2-D CGA null point entity

The CGA null point entity $\mathbf{T}_{C}=\mathcal{C}(\mathbf{t})$ is the conformal embedding

$$
\begin{equation*}
C(\mathbf{t})=\mathbf{t}+\frac{\mathbf{t}^{2}}{2} \mathbf{e}_{\infty 1}+\mathbf{e}_{o 1} \tag{8}
\end{equation*}
$$

of the $\mathcal{C}_{2}^{1}$ Euclidean vector ${ }^{5} \mathbf{t} \in \mathbb{R}^{2}$. The Minkowskian pair of null vectors representing origin and infinity is given by ${ }^{6}$

$$
\begin{equation*}
\mathbf{e}_{o}=\mathbf{e}_{o 1}=\frac{1}{2}\left(\mathbf{e}_{-}-\mathbf{e}_{+}\right), \quad \mathbf{e}_{\infty}=\mathbf{e}_{\infty 1}=\mathbf{e}_{-}+\mathbf{e}_{+} \tag{9}
\end{equation*}
$$

This implies that a conformally embedded point $\mathcal{C}(\mathbf{t})$ squares itself to zero, $\mathcal{C}(\mathbf{t})^{2}=0$ (i.e., it is also a null-vector). Furthermore, every finite conformal point is embedded, such that it is on the 3-D hyperplane in $\mathbb{R}^{3,1}$ defined by

$$
\begin{equation*}
C(\mathbf{t}) \cdot \mathbf{e}_{\infty 1}=-1, \tag{10}
\end{equation*}
$$

which includes the origin point as well, since by construction $\mathbf{e}_{o 1} \cdot \mathbf{e}_{\infty 1}=-1$.
A point $\mathbf{P}_{C}$ can be transformed to the normalized point $\hat{\mathbf{P}}_{C}$ that has unit scale on the homogeneous component $\mathbf{e}_{o 1}$ as

$$
\begin{equation*}
\hat{\mathbf{P}}_{C}=\frac{\mathbf{P}_{C}}{-\mathbf{P}_{C} \cdot \mathbf{e}_{\infty 1}} \tag{11}
\end{equation*}
$$

Points are assumed to be initially normalized as the embedding $\hat{\mathbf{P}}_{C}=\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$. After performing operations on a point (or other entity), it may no longer be normalized. Some operations do not preserve the homogeneous scale.

The projection (inverse embedding) of a CGA point $\mathbf{P}_{C}$ back to a vector $\mathbf{p}_{\mathcal{E}} \in \mathbb{R}^{2}$ is

$$
\begin{equation*}
\mathbf{p}_{\mathcal{E}}=\left(\hat{\mathbf{P}}_{C} \cdot \mathbf{I}_{\mathcal{E}}\right) \mathbf{I}_{\mathcal{E}}^{-1} \tag{12}
\end{equation*}
$$

The Euclidean distance $d(\mathbf{p}, \mathbf{q})$ between two finite CGA points $\mathbf{P}_{C}=\mathcal{C}(\mathbf{p})$ and $\mathbf{Q}_{C}=\mathcal{C}(\mathbf{q})$ is

$$
\begin{equation*}
d\left(\mathbf{P}_{C}, \mathbf{Q}_{C}\right)=\sqrt{-2 \hat{\mathbf{P}}_{C} \cdot \hat{\mathbf{Q}}_{C}}=\sqrt{(\mathbf{p}-\mathbf{q})^{2}} \tag{13}
\end{equation*}
$$

Remark 1. In the CGA $\mathcal{G}_{1,2+1}$ of the anti-Euclidean plane $\mathbb{R}^{0,2}$, we still have $\mathbf{e}_{\infty 1} \cdot \mathbf{e}_{o 1}=-1$, and the squared distance becomes $d^{2}=2 \hat{\mathbf{P}}_{C} \cdot \hat{\mathbf{Q}}_{C}$. The $\mathcal{G}_{1,3+1}$ CGA of the anti-Euclidean 3 -space $\mathbb{R}^{0,3}$, called Conformal Space Algebra (CSA), is used in DCSTA [9][6] and its subalgebras CSTA and CSA. Using $\mathbb{R}^{0,2}$ is possible by just changing the signs on some results ${ }^{7}$, like in the distance definition (13).

## 2.2 | 2-D CGA OPNS entities

A 2-D CGA point $\mathbf{T}_{C}=\mathcal{C}(\mathbf{t})$ is on a CGA geometric outer product null space (OPNS) [25] plane curve entity $\mathbf{X}_{C}^{*}$ if $\mathbf{T}_{C} \wedge \mathbf{X}_{C}^{*}=0$, where

$$
\begin{equation*}
\mathbf{X}_{C}^{*}=\mathbf{X}_{\mathcal{C}} \mathbf{I}_{\mathcal{C}}^{-1} \tag{14}
\end{equation*}
$$

[^2]is the CGA dual of the CGA IPNS entity $\mathbf{X}_{C}$, treated in Section 2.3. An entity and its dual entity represent the same plane curve. A CGA OPNS entity can be directly formed as the wedge of up to four CGA points as
\[

$$
\begin{equation*}
\mathbf{X}_{C}^{*}=\bigwedge \mathbf{P}_{C_{j}}, \quad \text { for } 1 \leq j \leq 4 \tag{15}
\end{equation*}
$$

\]

where the $\mathbf{P}_{C_{j}}$ are points on the plane curve that span the plane curve. The entire plane, as a surface, can be represented by the 4-blade entity $\mathbf{I}_{C}$.

The CGA OPNS 3-blade line $\mathbf{L}_{C}^{*}$ is the wedge of two CGA points $\mathbf{P}_{C_{j}}$ on the line and the point $\mathbf{e}_{\infty 1}$

$$
\begin{equation*}
\mathbf{L}_{C}^{*}=\mathbf{P}_{C_{1}} \wedge \mathbf{P}_{C_{2}} \wedge \mathbf{e}_{\infty 1}=\mathbf{L}_{C} / \mathbf{I}_{C} \tag{16}
\end{equation*}
$$

and is the CGA dual of the CGA IPNS 1-blade line $\mathbf{L}_{C}$.
The CGA OPNS 3-blade circle $\mathbf{C}_{C}^{*}$ is the wedge of three CGA points $\mathbf{P}_{C_{j}}$ on the circle

$$
\begin{equation*}
\mathbf{C}_{C}^{*}=\mathbf{P}_{C_{1}} \wedge \mathbf{P}_{C_{2}} \wedge \mathbf{P}_{C_{3}}=\mathbf{C}_{C} / \mathbf{I}_{C} \tag{17}
\end{equation*}
$$

and is the CGA dual of the CGA IPNS 1-blade circle $\mathbf{C}_{C}$.
The CGA OPNS 2-blade point pair $\mathbf{2}_{C}^{*}$ is the wedge of two finite CGA points $\mathbf{P}_{C_{j}}$

$$
\begin{equation*}
\mathbf{2}_{c}^{*}=\mathbf{P}_{C_{1}} \wedge \mathbf{P}_{C_{2}}=\mathbf{2}_{C} / \mathbf{I}_{C} \tag{18}
\end{equation*}
$$

and is the CGA dual ${ }^{8}$ of the CGA IPNS 2-blade point pair $\mathbf{2}_{C}$. If one of the points is $\mathbf{e}_{\infty 1}$, then it is a CGA OPNS 2-blade flat point $\mathbb{P}_{C}^{*}$.

The point pair decomposition [3]

$$
\begin{equation*}
\hat{\mathbf{P}}_{C_{ \pm}}=\frac{\mathbf{2}_{c}^{*} \mp \sqrt{\mathbf{2}_{c}^{*} \cdot \mathbf{2}_{c}^{*}}}{-\mathbf{e}_{\infty 1} \cdot \mathbf{2}_{c}^{*}} \tag{19}
\end{equation*}
$$

gives the two normalized points of the point pair $\mathbf{2}_{c}^{*}=\mathbf{P}_{C_{+}} \wedge \mathbf{P}_{C_{-}}$.
The CGA OPNS 2-blade flat point $\mathbb{P}_{C}^{*}$ is the wedge of one finite CGA point $\mathbf{P}_{C}$ and $\mathbf{e}_{\infty 1}$

$$
\begin{equation*}
\mathbb{P}_{c}^{*}=\mathbf{P}_{c} \wedge \mathbf{e}_{\infty 1}=\mathbb{P}_{c} / \mathbf{I}_{C} \tag{20}
\end{equation*}
$$

and is the CGA dual of the CGA IPNS 2-blade flat point $\mathbb{P}_{C}$. A unit scale (normalized) flat point is $\hat{\mathbb{P}}_{C}^{*}=\hat{\mathbf{P}}_{c} \wedge \mathbf{e}_{\infty 1}$. As explained in [3], a CGA IPNS 2-blade flat point $\mathbb{P}_{C}$ can represent the intersection of two non-parallel CGA IPNS 1-blade lines $\mathbf{L}_{C_{j}}$ in the plane as

$$
\begin{equation*}
\mathbb{P}_{c}^{*}=\mathbb{P}_{C} / \mathbf{I}_{C}=\left(\mathbf{L}_{C_{1}} \wedge \mathbf{L}_{C_{2}}\right) \mathbf{I}_{C}^{-1} \tag{21}
\end{equation*}
$$

The CGA point $\mathbf{P}_{C}$ of CGA OPNS 2-blade flat point $\mathbb{P}_{C}^{*}$ is projected $[3]$ to a vector as

$$
\begin{equation*}
\mathbf{p}_{\mathcal{E}}=\mathcal{C}^{-1}\left(\mathbf{P}_{C}\right)=\frac{\left(\mathbf{e}_{o 1} \wedge \mathbf{e}_{\infty 1}\right) \cdot\left(\mathbf{e}_{o 1} \wedge \mathbb{P}_{C}^{*}\right)}{-\left(\mathbf{e}_{o 1} \wedge \mathbf{e}_{\infty 1}\right) \cdot \mathbb{P}_{C}^{*}}=\frac{-\mathbb{P}_{C}^{*}}{\left(\mathbf{e}_{o 1} \wedge \mathbf{e}_{\infty 1}\right) \cdot \mathbb{P}_{C}^{*}} \cdot \mathbf{e}_{o 1}-\mathbf{e}_{o 1} . \tag{22}
\end{equation*}
$$

The null point embedding $\mathbf{P}_{C}=\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$, with the property $\mathbf{P}_{C}^{2}=\mathbf{P}_{C} \cdot \mathbf{P}_{C}+\mathbf{P}_{C} \wedge \mathbf{P}_{C}=0$, is both a CGA IPNS and OPNS point entity. There also exists the interpretation of undual IPNS 3-blade point $\mathbf{P}_{C} \mathbf{I}_{C}$.

## 2.3 | 2-D CGA IPNS entities

A CGA point $\mathbf{T}_{C}=\mathcal{C}(\mathbf{t})$ is on the CGA geometric inner product null space (IPNS) [25] plane curve entity $\mathbf{X}_{C}$ if $\mathbf{T}_{\mathcal{C}} \cdot \mathbf{X}_{\mathcal{C}}=0$. In the two dimensions of the plane $\mathbb{R}^{2}$, as considered in this paper, ${ }^{9}$ geometric entities represent plane curves and some of their intersections.

The CGA IPNS 1-blade circle ${ }^{10} \mathbf{C}_{C}$, centered at CGA point $\mathbf{P}_{C}=\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$ with radius $r$ or with finite surface point $\mathbf{Q}_{C}$, is defined as

$$
\begin{equation*}
\mathbf{C}_{C}=\mathbf{P}_{C}-\frac{1}{2} r^{2} \mathbf{e}_{\infty 1}=\mathbf{P}_{C}+\left(\mathbf{P}_{C} \cdot \hat{\mathbf{Q}}_{C}\right) \mathbf{e}_{\infty 1} \tag{23}
\end{equation*}
$$

[^3]A normalized circle $\hat{\mathbf{C}}_{c}$ has a unit scale (normalized) center point $\mathbf{P}_{C}=\hat{\mathbf{P}}_{c}$ and $\hat{\mathbf{C}}_{c}^{2}=r^{2}$. The circle entity $\mathbf{C}_{c}$ is derived, from the constant distance $r=d\left(\mathbf{P}_{C}, \mathbf{Q}_{C}\right)$ between the finite center point $\mathbf{P}_{C}$ and any point $\mathbf{Q}_{C}$ of the circle, by the inner product test with $\mathbf{T}_{\mathcal{C}}=\mathcal{C}\left(\mathbf{t}_{\mathcal{E}}\right)$ as

$$
\begin{equation*}
d^{2}\left(\mathbf{P}_{c}, \mathbf{Q}_{c}\right)=-2 \hat{\mathbf{P}}_{c} \cdot \hat{\mathbf{Q}}_{c}=-2 \hat{\mathbf{P}}_{c} \cdot \mathbf{T}_{c}=r^{2} \tag{24}
\end{equation*}
$$

which can be rewritten as

$$
\begin{equation*}
\mathbf{T}_{C} \cdot \hat{\mathbf{P}}_{C}+\frac{1}{2} r^{2}=\mathbf{T}_{C} \cdot\left(\hat{\mathbf{P}}_{C}-\frac{1}{2} r^{2} \mathbf{e}_{\infty 1}\right)=\mathbf{T}_{C} \cdot \mathbf{C}_{C}=0 \tag{25}
\end{equation*}
$$

The CGA IPNS 1-blade $\operatorname{lin} e^{11} \mathbf{L}_{C}$, normal to unit vector $\hat{\mathbf{n}}_{\mathcal{E}}$ at distance $d$ from the origin, or through point $\mathbf{p}_{\mathcal{E}}$, is defined as

$$
\begin{equation*}
\mathbf{L}_{\mathcal{C}}=\hat{\mathbf{n}}_{\mathcal{E}}+d \mathbf{e}_{\infty 1}=\hat{\mathbf{n}}_{\mathcal{E}}+\left(\mathbf{p}_{\mathcal{E}} \cdot \hat{\mathbf{n}}_{\mathcal{E}}\right) \mathbf{e}_{\infty 1} \tag{26}
\end{equation*}
$$

In terms of the line direction $\hat{\mathbf{d}}_{\mathcal{E}}$ through point $\mathbf{p}_{\mathcal{E}}$, where $\hat{\mathbf{d}}_{\mathcal{E}}^{*}=\hat{\mathbf{n}}_{\mathcal{E}}=\hat{\mathbf{d}}_{\mathcal{E}} \mathbf{I}_{\mathcal{E}}^{-1}$, the line entity can be written as

$$
\begin{equation*}
\mathbf{L}_{C}=\hat{\mathbf{d}}_{\mathcal{E}}^{*}+\left(\mathbf{p}_{\mathcal{E}} \cdot \hat{\mathbf{d}}_{\mathcal{E}}^{*}\right) \mathbf{e}_{\infty 1} \tag{27}
\end{equation*}
$$

The line $\mathbf{L}_{C}$ with unit direction $\hat{\mathbf{d}}_{\mathcal{E}}$ is also unit scale (normalized), where $\mathbf{L}_{c}=\hat{\mathbf{L}}_{c}$ and $\hat{\mathbf{L}}_{c}^{2}=1$.
In general, the wedge of $2 \leq k \leq 4$ CGA IPNS 1-blade entities forms a CGA IPNS $k$-blade entity that represents the intersection of the $k$ 1-blade entities.

## 2.4 | 2-D CGA versor operations

It can be shown that reflection of point $\mathbf{P}_{C}$ in circle $\mathbf{C}_{C}$ as the versor (sandwich product) operation

$$
\begin{equation*}
\mathbf{P}_{c}^{\prime}=\mathbf{C}_{C} \mathbf{P}_{C} \mathbf{C}_{c}^{-1} \tag{28}
\end{equation*}
$$

produces $\mathbf{P}_{c}^{\prime}$ as the inversion of $\mathbf{P}_{C}$ in the circle $\mathbf{C}_{C}$. The center point of circle $\mathbf{C}_{C}$ is $\mathbf{C}_{C} \mathbf{e}_{\infty 1} \mathbf{C}_{C}$. Successive inversions in two concentric circles, $\mathbf{C}_{C_{1}}$ of radius $r_{1}$ followed by $\mathbf{C}_{C_{2}}$ of radius $r_{2}$, produces isotropic dilation (uniform scaling) relative to the circle center by the factor $d=r_{2}^{2} / r_{1}^{2}$. It can also be shown that reflection of a point $\mathbf{P}_{C}$ in a line $\mathbf{L}_{C}$ as the versor operation

$$
\begin{equation*}
\mathbf{P}_{c}^{\prime}=\mathbf{L}_{c} \mathbf{P}_{c} \mathbf{L}_{c}^{-1} \tag{29}
\end{equation*}
$$

produces $\mathbf{P}_{C}^{\prime}$ as the reflection of $\mathbf{P}_{C}$ in the line $\mathbf{L}_{C}$. Successive reflections in two parallel lines that are separated by a displacement $\mathbf{d}_{\mathcal{E}} / 2$ from the first line toward the second line produces translation by displacement $\mathbf{d}_{\mathcal{E}}$. Successive reflections in two nonparallel lines subtending an angle $\theta / 2$ at their intersection point from the first line toward the second line produces rotation by angle $\theta$ around the intersection point in the direction of the first line toward the second line.

The line entity $\mathbf{L}_{\mathcal{C}}$ will also be called the 1-versor reflection operator (reflector), and the circle entity $\mathbf{C}_{\mathcal{C}}$ will also be called the 1 -versor inversion operator (inversor), since they are geometric entities and also version operators (versors).

The translator (translation versor, translation operator) $T=\mathbf{L}_{\mathcal{C}_{2}} \mathbf{L}_{\mathcal{C}_{1}}=\mathbf{L}_{\mathcal{C}_{2}} \cdot \mathbf{L}_{\mathcal{C}_{1}}+\mathbf{L}_{\mathcal{C}_{2}} \wedge \mathbf{L}_{\mathcal{C}_{1}}$ is defined as successive reflections in two parallel lines separated by half the translation displacement $\mathbf{d}_{\mathcal{E}} / 2=\left(\hat{\mathbf{L}}_{\mathcal{C}_{2}} \wedge \hat{\mathbf{L}}_{\mathcal{C}_{1}}\right) \cdot \mathbf{e}_{o 1}$, or $T^{2}=\mathbf{L}_{\mathcal{C}_{2}} \mathbf{L}_{\mathcal{C}_{1}}$ if by $\mathbf{d}_{\mathcal{E}}$. The rotor $R=\mathbf{L}_{C_{2}} \mathbf{L}_{C_{1}}$ is defined as successive reflections in two non-parallel lines subtending half the rotation angle $\theta / 2=$ $\operatorname{acos}\left(\hat{\mathbf{L}}_{\mathcal{C}_{2}} \cdot \hat{\mathbf{L}}_{\mathcal{C}_{1}}\right)$, or $R^{2}=\mathbf{L}_{C_{2}} \mathbf{L}_{\mathcal{C}_{1}}$ if subtending $\theta$. The isotropic dilator $D=\mathbf{C}_{C_{2}} \mathbf{C}_{C_{1}}$ is defined as successive inversions in two concentric circles with relative dilation factor $d=r_{2}^{2} / r_{1}^{2}$, or $D^{2}=\mathbf{C}_{C_{2}} \mathbf{C}_{C_{1}}$ with $d^{2}=r_{2}^{2} / r_{1}^{2} . T, R$, and $D$ are called 2 -versors $V \in\{T, R, D\}$. The product of $k$ vectors, with the product having an inverse, is called a $k$-versor $V_{k}$ [14][13]. Their operation on a CGA entity $\mathbf{X}$ has the form $\mathbf{X}^{\prime}=V \mathbf{X} V^{-1}$, called a versor "sandwich" operation. Each even CGA 2-versor $V$ has an exponential form $V=\exp (\log V)$.

A translator is a type of versor. The CGA 2-versor translator $T$, for a translation by the displacement vector $\mathbf{d}_{\mathcal{E}}=d_{x} \mathbf{e}_{1}+d_{y} \mathbf{e}_{2}$, is defined by reflections in parallel lines as

$$
\begin{equation*}
T=1+\frac{1}{2} \mathbf{e}_{\infty 1} \mathbf{d}_{\mathcal{E}}=\exp \left(\frac{1}{2} \mathbf{e}_{\infty 1} \mathbf{d}_{\mathcal{E}}\right) \tag{30}
\end{equation*}
$$

where $1=\hat{\mathbf{L}}_{\mathcal{C}_{2}} \cdot \hat{\mathbf{L}}_{\mathcal{C}_{1}}=\cos (0)$, and $\mathbf{e}_{\infty 1} \mathbf{d}_{\mathcal{E}} / 2=\hat{\mathbf{L}}_{\mathcal{C}_{2}} \wedge \hat{\mathbf{L}}_{C_{1}}$ represents the intersection of the two parallel lines in infinity as a free vector [3] of the half displacement $\mathbf{d}_{\mathcal{E}} / 2$.

[^4]A rotor is a rotation operator. The CGA 2 -versor rotor $R$, for rotation around the point $\hat{\mathbf{P}}_{C}=\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$ by $\theta$ radians counter-clockwise in the right-handed $x y$-plane, is defined by translation of the origin rotor $\exp \left(\theta \mathbf{e}_{2} \mathbf{e}_{1} / 2\right)$ by $\mathbf{d}_{\mathcal{E}}=\mathbf{p}_{\mathcal{E}}$ (i.e., $\left.T \exp \left(\theta \mathbf{e}_{2} \mathbf{e}_{1} / 2\right) T^{-1}\right)$ as

$$
\begin{equation*}
R=\cos \left(\frac{\theta}{2}\right)+\sin \left(\frac{\theta}{2}\right)\left(\mathbf{e}_{2} \mathbf{e}_{1}-\left(\mathbf{p}_{\mathcal{E}} \cdot\left(\mathbf{e}_{2} \mathbf{e}_{1}\right)\right) \mathbf{e}_{\infty 1}\right)=\cos \left(\frac{\theta}{2}\right)+\sin \left(\frac{\theta}{2}\right)\left(\hat{\mathbf{P}}_{C} \wedge \mathbf{e}_{\infty 1}\right) \mathbf{I}_{C}=\exp \left(\frac{\theta}{2} \hat{\mathbb{P}}_{C}^{*} \mathbf{I}_{C}\right)=\exp \left(\frac{\theta}{2} \hat{\mathbb{P}}_{C}\right) \tag{31}
\end{equation*}
$$

where $\sin \left(\frac{\theta}{2}\right) \hat{\mathbb{P}}_{C}=\mathbf{L}_{\mathcal{C}^{2}} \wedge \mathbf{L}_{\mathcal{C}^{1}}$ is the IPNS flat point representing the intersection of the lines.
A dilator is a dilation operator. The CGA 2-versor isotropic dilator $D$, by dilation factor $d>0$ relative to the point $\hat{\mathbf{P}}_{C}=$ $\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$ in the plane, is defined by inversions in circles (centered on $\mathbf{p}_{\mathcal{E}}$ with $r_{1}=1$ and $r_{2}=\sqrt{d}$ ) as

$$
\begin{equation*}
D=\frac{d+1}{2}+\frac{d-1}{2} \hat{\mathbb{P}}_{c}^{*}=\sqrt{d} \hat{D}=\sqrt{d} \exp \left(\operatorname{atanh}\left(\frac{d-1}{d+1}\right) \hat{\mathbb{P}}_{c}^{*}\right)=\sqrt{d} \exp \left(\frac{\ln (d)}{2} \hat{\mathbb{P}}_{c}^{*}\right) \tag{32}
\end{equation*}
$$

where the exponential $\hat{D}=\exp (A)$ is a unimodular versor (n.b., $\hat{D}^{\sim}=\hat{D}^{-1}$ ) and the constant modulus (radius or invariant interval) is $\sqrt{d}=\sqrt{(d+1)^{2} / 4-(d-1)^{2} / 4}$ for hyperbolic rotation by angle $\varphi=\ln (d)$ through point $((d+1) / 2,(d-1) / 2)$ of the hyperbolic plane. If $\hat{D}$ is used as the dilator, then the homogeneous scale on $\mathbf{e}_{o}$ is not preserved.

For dilation around $\mathbf{e}_{o}$, then $\hat{\mathbb{P}}_{C}^{*}=\mathbf{e}_{o} \wedge \mathbf{e}_{\infty}=\mathbf{e}_{4} \mathbf{e}_{3}$ and the hyperbolic rotation is in the hyperbolic Minkowski $\mathbf{e}_{3} \mathbf{e}_{4}$-plane. Then, for this rotation to form the dilation of the point $\mathcal{C}\left(\mathbf{t}_{\mathcal{E}}\right)$, we must have the squared versor $D^{2}=\mathbf{C}_{C_{2}} \mathbf{C}_{C_{1}}$ (for $r_{1}=1, r_{2}=d$, and center $\mathbf{e}_{o}$ ) as the ratio

$$
\begin{align*}
D^{2} & =\left(d^{2} \frac{\mathbf{t}_{\mathcal{\varepsilon}}^{2}}{2} \mathbf{e}_{\infty 1}+\mathbf{e}_{o 1}\right)\left(\frac{\mathbf{t}_{\mathcal{E}}^{2}}{2} \mathbf{e}_{\infty 1}+\mathbf{e}_{o 1}\right)^{-1} \\
& =\left(\frac{d^{2} \mathbf{t}_{\mathcal{E}}^{2}-1}{2} \mathbf{e}_{3}+\frac{d^{2} \mathbf{t}_{\mathcal{E}}^{2}+1}{2} \mathbf{e}_{4}\right)\left(\frac{\mathbf{t}_{\mathcal{E}}^{2}-1}{2} \mathbf{e}_{3}+\frac{\mathbf{t}_{\mathcal{E}}^{2}+1}{2} \mathbf{e}_{4}\right)^{-1} \\
& =\left(\alpha^{\prime} \mathbf{e}_{3}+\beta^{\prime} \mathbf{e}_{4}\right)\left(\alpha \mathbf{e}_{3}+\beta \mathbf{e}_{4}\right)^{-1}=\frac{d^{2}+1}{2}+\frac{d^{2}-1}{2} \mathbf{e}_{4} \mathbf{e}_{3} \tag{33}
\end{align*}
$$

which can be verified after some algebra. For $r_{2}=\sqrt{d}$, we get the square root $\sqrt{D^{2}}=D$. It can also be verified that

$$
\begin{align*}
D\left(x \mathbf{e}_{1}+y \mathbf{e}_{2}\right) D^{\sim} & =\sqrt{d} \hat{D} \sqrt{d} \hat{D}^{\sim}\left(x \mathbf{e}_{1}+y \mathbf{e}_{2}\right)=d x \mathbf{e}_{1}+d y \mathbf{e}_{2}  \tag{34}\\
D^{2}\left(\alpha \mathbf{e}_{3}+\beta \mathbf{e}_{4}\right) & =D\left(\alpha \mathbf{e}_{3}+\beta \mathbf{e}_{4}\right) D^{\sim}=\alpha^{\prime} \mathbf{e}_{3}+\beta^{\prime} \mathbf{e}_{4}  \tag{35}\\
D C\left(\mathbf{t}_{\mathcal{E}}\right) D^{\sim} & =C\left(d \mathbf{t}_{\mathcal{E}}\right) \tag{36}
\end{align*}
$$

which preserves the homogeneous scale on the term $\mathbf{e}_{o}$. Dilation centered on $\mathbf{e}_{o}$ is generalized to dilation centered on a point $\mathbf{P}_{C}=\mathcal{C}\left(\mathbf{p}_{\mathcal{E}}\right)$ with the translation by $\mathbf{d}_{\mathcal{E}}=\mathbf{p}_{\mathcal{E}}, D^{\prime}=T D T^{-1}$.

By outermorphism, all of the CGA versors operate correctly on all of the CGA OPNS entities formed as wedges of points. By dualization, the versors also work correctly on all of the CGA IPNS entities. Consider the test point ${ }^{12} \mathbf{T}_{\mathcal{C}}=\mathcal{C}\left(\mathbf{t}_{\mathcal{E}}\right)$, an OPNS entity $\mathbf{X}_{C}^{*}$, and the outermorphism of their surface point test (their wedge) by a versor $V$ as

$$
\begin{equation*}
V\left(\mathbf{T}_{C} \wedge \mathbf{X}_{C}^{*}\right) V^{-1}=\left(V \mathbf{T}_{C} V^{-1}\right) \wedge\left(V \mathbf{X}_{c}^{*} V^{-1}\right)=\mathbf{T}_{c}^{\prime} \wedge\left(\bigwedge V \mathbf{P}_{C_{j}} V^{-1}\right)=\mathbf{T}_{c}^{\prime} \wedge \mathbf{X}_{c}^{* \prime} \tag{37}
\end{equation*}
$$

Since the transformations of points by the versor operations (compositions of inversions in general circles $\mathbf{C}_{C}$, including reflections in lines $\mathbf{L}_{C}$ ) are correct, then the transformation of $\mathbf{X}_{C}^{*}$ into the entity $\mathbf{X}_{C}^{* \prime}$ is also correct, and the surface point test is zero only for correctly transformed surface points $\mathbf{T}_{C}^{\prime}$. Considering the IPNS entity $\mathbf{X}_{C}=\mathbf{X}_{C}^{*} \mathbf{I}_{C}=\mathbf{X}_{C}^{*} \cdot \mathbf{I}_{C}$, then the test is

$$
\begin{equation*}
V\left(\mathbf{T}_{c} \cdot\left(\mathbf{X}_{C}^{*} \cdot \mathbf{I}_{C}\right)\right) V^{-1}=V\left(\left(\mathbf{T}_{\mathcal{C}} \wedge \mathbf{X}_{c}^{*}\right) \cdot \mathbf{I}_{C}\right) V^{-1}=V\left(\mathbf{T}_{c} \wedge \mathbf{X}_{C}^{*}\right) \mathbf{I}_{C} V^{-1}=V\left(\mathbf{T}_{c} \wedge \mathbf{X}_{C}^{*}\right) V^{-1} \mathbf{I}_{C} \tag{38}
\end{equation*}
$$

[^5]
## 3 | TRIPLE 2-D CGA $\mathcal{G}_{3(2+1), 3}$ OF THE EUCLIDEAN PLANE $\mathbb{R}^{2}$

The Triple Conformal Geometric Algebra (TCGA) $\mathcal{T}=\mathcal{G}_{3(2+1), 3}$ for the 2-D Euclidean plane $\mathbb{R}^{2}$ is a high-dimensional Geometric Algebra of the 12-D pseudo-Euclidean vector space $\mathbb{R}^{9,3}$. With $2^{12}$ basis blades, $\mathcal{T}$ is currently considered to be a high-dimensional Clifford algebra. TCGA is a staightforward extension of the concepts introduced with Double Conformal "Darboux Cyclide" Geometric Algebra (DCGA) $\mathcal{C}_{8,2}$ [10][8][7][5][4] and Double Conformal Space-Time Algebra (DCSTA) $\mathcal{C}_{4,8}$ [9][6]. Theoretically, there is no problem to similarly establish TCGA $\mathcal{G}_{12,3}$ for the 3-D Euclidean space $\mathbb{R}^{3}$. However, current computer algebra implementations cannot yet easily cope with elaborate computations in a high-dimensional Clifford algebra $\mathcal{G}_{12,3}$ of $2^{15}$ basis blades. Therefore, we pragmatically use the TCGA $\mathcal{T}$ of the Euclidean plane to explain the principles of this approach, where we only have to deal with currently feasible computations in a Clifford algebra of $2^{12}$ basis blades.

Triple conformal geometric algebra of the Euclidean plane $\mathcal{G}_{9,3}$ includes three copies of $\mathcal{G}_{3,1}$ CGA, which are named CGA1 $\mathcal{C}^{1}$, CGA2 $\mathcal{C}^{2}$, and CGA3 $\mathcal{C}^{3}$. Each $\operatorname{CGA} i \mathcal{C}^{i}$ has a Euclidean plane $\mathbb{R}^{2}$ subalgebra $\mathcal{C}_{2}$, denoted $\mathcal{E}^{i}$.

The metric for $\mathcal{C}^{1}$ is $m_{\mathcal{C}^{1}}=\operatorname{diag}(1,1,1,-1)=\left[\mathbf{e}_{i} \cdot \mathbf{e}_{j}\right]: i, j \in\{1,2,3,4\}$. The metric for $\mathcal{C}^{2}$ is $m_{\mathcal{C}^{2}}=\operatorname{diag}(1,1,1,-1)=\left[\mathbf{e}_{i} \cdot \mathbf{e}_{j}\right]$ $: i, j \in\{5,6,7,8\}$. The metric for $\mathcal{C}^{3}$ is $m_{\mathcal{C}^{3}}=\operatorname{diag}(1,1,1,-1)=\left[\mathbf{e}_{i} \cdot \mathbf{e}_{j}\right]: i, j \in\{9,10,11,12\}$. The metric for TCGA $\mathcal{T}$ combines the metrics for $\mathcal{C}^{1}, \mathcal{C}^{2}$ and $\mathcal{C}^{3}$ as $m_{\mathcal{T}}=\operatorname{diag}(1,1,1,-1,1,1,1,-1,1,1,1,-1)=\left[\mathbf{e}_{i} \cdot \mathbf{e}_{j}\right]: i, j \in\{1 \ldots 12\}$, because together the basis vectors of $\mathcal{C}^{1}, \mathcal{C}^{2}$ and $\mathcal{C}^{3}$ span $\mathbb{R}^{9,3}$. This can be summarized as metric $m_{\mathcal{J}}$ for the 1 -vectors of TCGA $\mathcal{G}_{9,3}^{1}$ with

$$
\begin{align*}
m_{\mathcal{T}}=\left[m_{i j}\right] & =\operatorname{diag}(1,1,1,-1,1,1,1,-1,1,1,1,-1) \\
m_{i j} & =\mathbf{e}_{i} \cdot \mathbf{e}_{j}: i, j \in\{1,2, \ldots, 12\} \tag{39}
\end{align*}
$$

The unit pseudoscalars for $\mathcal{E}^{1}, \mathcal{E}^{2}, \mathcal{E}^{3}$ are $\mathbf{I}_{\mathcal{E}^{1}}=\mathbf{e}_{1} \mathbf{e}_{2}, \mathbf{I}_{\mathcal{E}^{2}}=\mathbf{e}_{5} \mathbf{e}_{6}, \mathbf{I}_{\mathcal{E}^{3}}=\mathbf{e}_{9} \mathbf{e}_{10}$, respectively, and are the dualization operators on elements $A_{\mathcal{E}^{i}} \in \mathcal{E}^{i}$ by the division $A_{\mathcal{E}^{i}}^{*}=A_{\mathcal{E}^{i}} / \mathbf{I}_{\mathcal{E}^{i}}$. Similarly, the unit pseudoscalars for $\mathcal{C}^{1}, \mathcal{C}^{2}, \mathcal{C}^{3}$ are $\mathbf{I}_{\mathcal{C}^{1}}=\mathbf{e}_{1} \mathbf{e}_{2} \mathbf{e}_{3} \mathbf{e}_{4}, \mathbf{I}_{\mathcal{C}^{2}}=$ $\mathbf{e}_{5} \mathbf{e}_{6} \mathbf{e}_{7} \mathbf{e}_{8}, \mathbf{I}_{\mathcal{C}^{3}}=\mathbf{e}_{9} \mathbf{e}_{10} \mathbf{e}_{11} \mathbf{e}_{12}$, respectively, and are the dualization operators on elements $A_{\mathcal{C}^{i}} \in \mathcal{C}^{i}$ by the division $A_{\mathcal{C}^{i}}^{*}=A_{\mathcal{C}^{i}} / \mathbf{I}_{\mathcal{C}^{i}}$. The TCGA 12-blade unit pseudoscalar is $\mathbf{I}_{\mathcal{T}}=\mathbf{I}_{\mathcal{C}^{1}} \mathbf{I}_{\mathcal{C}^{2}} \mathbf{I}_{\mathcal{C}^{3}}=\mathbf{e}_{1} \mathbf{e}_{2} \ldots \mathbf{e}_{12}$.

In Section 2 we already discussed 2-D CGA essentially in terms of CGA1 $\mathcal{C}^{1}=\mathcal{C}$ and its Euclidean subalgebra $\mathcal{E}^{1}=\mathcal{E}$ with unit vector elements $\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}, \mathbf{e}_{4}$, where $\mathbf{e}_{+}=\mathbf{e}_{3}$ and $\mathbf{e}_{-}=\mathbf{e}_{4}$. In the following subsections, we will introduce relevant multivector entities in TCGA for geometric objects and for transformation operators.

## 3.1 | TCGA IPNS 3-blade standard entities

The TCGA 3-blade point $\mathbf{P}_{\mathcal{J}}$ is the $k$-CGA of multiplicity $k=3$ (TCGA) embedding

$$
\begin{equation*}
\mathbf{P}_{\mathcal{T}}=\mathcal{T}\left(\mathbf{p}_{\mathcal{E}}\right)=\mathcal{C}^{1}\left(\mathbf{p}_{\mathcal{E}^{1}}\right) \mathcal{C}^{2}\left(\mathbf{p}_{\mathcal{E}^{2}}\right) \mathcal{C}^{3}\left(\mathbf{p}_{\mathcal{E}^{3}}\right)=\mathcal{C}^{1}\left(\mathbf{p}_{\mathcal{E}^{1}}\right) \wedge \mathcal{C}^{2}\left(\mathbf{p}_{\mathcal{E}^{2}}\right) \wedge \mathcal{C}^{3}\left(\mathbf{p}_{\mathcal{E}^{3}}\right) \tag{40}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathcal{C}^{1}\left(\mathbf{p}_{\mathcal{E}^{1}}\right)=\mathbf{p}_{\mathcal{E}^{1}}+\frac{1}{2} \mathbf{p}_{\mathcal{E}}^{2} \mathbf{e}_{\infty 1}+\mathbf{e}_{o 1}, \quad \mathcal{C}^{2}\left(\mathbf{p}_{\mathcal{E}^{2}}\right)=\mathbf{p}_{\mathcal{E}^{2}}+\frac{1}{2} \mathbf{p}_{\mathcal{E}}^{2} \mathbf{e}_{\infty 2}+\mathbf{e}_{o 2}, \quad \mathcal{C}^{3}\left(\mathbf{p}_{\mathcal{E}^{3}}\right)=\mathbf{p}_{\mathcal{E}^{3}}+\frac{1}{2} \mathbf{p}_{\mathcal{E}}^{2} \mathbf{e}_{\infty 3}+\mathbf{e}_{o 3}, \tag{41}
\end{equation*}
$$

and

$$
\begin{gather*}
\mathbf{p}_{\mathcal{E}^{1}}=p_{x} \mathbf{e}_{1}+p_{y} \mathbf{e}_{2}, \quad \mathbf{p}_{\mathcal{E}^{2}}=p_{x} \mathbf{e}_{5}+p_{y} \mathbf{e}_{6}, \quad \mathbf{p}_{\mathcal{E}^{3}}=p_{x} \mathbf{e}_{9}+p_{y} \mathbf{e}_{10}, \quad \mathbf{p}_{\mathcal{E}}^{2}=p_{x}^{2}+p_{y}^{2},  \tag{42}\\
\mathbf{e}_{\infty 1}=\mathbf{e}_{3}+\mathbf{e}_{4}, \quad \mathbf{e}_{\infty 2}=\mathbf{e}_{7}+\mathbf{e}_{8}, \quad \mathbf{e}_{\infty 3}=\mathbf{e}_{11}+\mathbf{e}_{12},  \tag{43}\\
\mathbf{e}_{o 1}=\frac{1}{2}\left(\mathbf{e}_{4}-\mathbf{e}_{3}\right), \quad \mathbf{e}_{o 2}=\frac{1}{2}\left(\mathbf{e}_{8}-\mathbf{e}_{7}\right), \quad \mathbf{e}_{o 3}=\frac{1}{2}\left(\mathbf{e}_{12}-\mathbf{e}_{11}\right) . \tag{44}
\end{gather*}
$$

The symbolic TCGA 3-blade test point $\mathbf{T}_{\mathcal{T}}=\mathcal{T}\left(\mathbf{t}_{\mathcal{E}}\right)$ is the product of the embeddings of the symbolic vector $\mathbf{t}_{\mathcal{E}^{1}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}$ in CGA1 and its copies in CGA2 and CGA3. The scalar coefficients on all corresponding CGA $i$ canonical basis blades are equal as copies.

Remark 2. Note that any CGAi IPNS $r$-blade entity $\mathbf{E}_{\mathcal{C}^{i}}$ has a representation as the TCGA IPNS $3 r$-blade entity $\mathbf{E}_{\mathcal{T}}=$ $\mathbf{E}_{\mathcal{C}^{1}} \mathbf{E}_{\mathcal{C}^{2}} \mathbf{E}_{\mathcal{C}^{3}}=\mathbf{E}_{\mathcal{C}^{1}} \wedge \mathbf{E}_{\mathcal{C}^{2}} \wedge \mathbf{E}_{\mathcal{C}^{3}}$. If the $\mathbf{E}_{\mathcal{C}^{i}}$ are IPNS 1-blade entities representing an implicit curve function $F(x, y)=\mathbf{T}_{C^{i}} \cdot \mathbf{E}_{\mathcal{C}^{i}}=0$, which is a circle, line, or point, then $\mathbf{E}_{\mathcal{J}}$ represents the function $F^{3}(x, y)=\mathbf{T}_{\mathcal{J}} \cdot \mathbf{E}_{\mathcal{J}}$ of multiplicity 3. Therefore, we also have the TCGA IPNS 3-blade circle $\mathbf{C}_{\mathcal{T}}$ and line $\mathbf{L}_{\mathcal{T}}$ entities, and also the IPNS 3r-blade entities representing their intersections.

## 3.2 | TCGA IPNS 3-vector extraction operators $T_{s}$

The test point $\mathbf{T}_{\mathcal{T}}$ is a sum of $4^{3}=64$ basis 3-blades, which is a subset of the $\binom{12}{3}=220$ canonical basis 3-blades in $\mathcal{C}_{9,3}$ TCGA. Each of the 64 3-blades has a scalar coefficient of a certain value in the variables $x$ and $y$. By using each of the 64 canonical basis 3-blade reciprocals ${ }^{13}$ as an extraction operator, each scalar coefficient can be extracted by inner product with $\mathbf{T}_{\mathcal{T}}$. Not every coefficient holds a unique value, but twenty unique values $s$ can be extracted by certain extraction operators $T_{s}$ (with index $s)$ as $s=\mathbf{T}_{\mathcal{T}} \cdot T_{s}$. Together, these twenty coefficients comprise all of the available distinct coordinate polynomial coefficients $\left(x, y, x^{2}, x y, y^{2}, \ldots, \mathbf{t}_{\mathcal{E}}^{6}\right)$ of cubic, quartic, quintic, and sextic degrees. The extraction operators $T_{s}$ are given in Tables 1,2 , and 3. For example, when expanding (40), the $x$-coordinate is found on three 3-blade terms: $x\left(\mathbf{e}_{1} \mathbf{e}_{o 2} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{5} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{o 2} \mathbf{e}_{9}\right)$. Furthermore, $-\mathbf{e}_{\infty 1},-\mathbf{e}_{\infty 2}$, and $-\mathbf{e}_{\infty 3}$ are reciprocal to $\mathbf{e}_{o 1}, \mathbf{e}_{o 2}$, and $\mathbf{e}_{o 3}$, respectively. Inspecting the expression for $T_{x}$ in Table 1, we see that it corresponds to reversing the product order in $\left(\mathbf{e}_{1} \mathbf{e}_{o 2} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{5} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{o 2} \mathbf{e}_{9}\right)$, and replacing $\mathbf{e}_{o i} \leftrightarrow \mathbf{e}_{\infty i}, i \in$ $\{1,2,3\}$, followed with division by 3 , since there are three terms with $x$ in (40). The other extraction operators in Tables 1,2 , and 3 are constructed by the same method.

Table 1. Extraction operators $T_{s}$ for general quadric plane curves.

$$
\begin{aligned}
& T_{x}=\frac{1}{3}\left(\mathbf{e}_{\infty 3} \mathbf{e}_{\infty 2} \mathbf{e}_{1}+\mathbf{e}_{\infty 3} \mathbf{e}_{5} \mathbf{e}_{\infty 1}+\mathbf{e}_{9} \mathbf{e}_{\infty 2} \mathbf{e}_{\infty 1}\right), T_{y}=\frac{1}{3}\left(\mathbf{e}_{\infty 3} \mathbf{e}_{\infty 2} \mathbf{e}_{2}+\mathbf{e}_{\infty 3} \mathbf{e}_{6} \mathbf{e}_{\infty 1}+\mathbf{e}_{10} \mathbf{e}_{\infty 2} \mathbf{e}_{\infty 1}\right) \\
& T_{x^{2}}=\frac{1}{3}\left(\mathbf{e}_{1} \mathbf{e}_{5} \mathbf{e}_{\infty 3}+\mathbf{e}_{1} \mathbf{e}_{\infty 2} \mathbf{e}_{9}+\mathbf{e}_{\infty 1} \mathbf{e}_{5} \mathbf{e}_{9}\right), T_{y^{2}}=\frac{1}{3}\left(\mathbf{e}_{2} \mathbf{e}_{6} \mathbf{e}_{\infty 3}+\mathbf{e}_{2} \mathbf{e}_{\infty 2} \mathbf{e}_{10}+\mathbf{e}_{\infty 1} \mathbf{e}_{6} \mathbf{e}_{10}\right) \\
& T_{x y}=\frac{1}{6}\left(\mathbf{e}_{1} \mathbf{e}_{6} \mathbf{e}_{\infty 3}+\mathbf{e}_{2} \mathbf{e}_{5} \mathbf{e}_{\infty 3}+\mathbf{e}_{1} \mathbf{e}_{\infty 2} \mathbf{e}_{10}+\mathbf{e}_{2} \mathbf{e}_{\infty 2} \mathbf{e}_{9}+\mathbf{e}_{\infty 1} \mathbf{e}_{5} \mathbf{e}_{10}+\mathbf{e}_{\infty 1} \mathbf{e}_{6} \mathbf{e}_{9}\right) \\
& T_{1}=\mathbf{e}_{\infty}=\mathbf{e}_{\infty 1} \mathbf{e}_{\infty 2} \mathbf{e}_{\infty 3}, T_{\mathbf{t}_{\varepsilon}^{2}}=\frac{2}{3}\left(\mathbf{e}_{o 1} \mathbf{e}_{\infty 2} \mathbf{e}_{\infty 3}+\mathbf{e}_{\infty 1} \mathbf{e}_{o 2} \mathbf{e}_{\infty 3}+\mathbf{e}_{\infty 1} \mathbf{e}_{\infty 2} \mathbf{e}_{o 3}\right) \\
& \hline
\end{aligned}
$$

Table 2. Extraction operators $T_{s}$ for general cubic plane curves.

$$
\begin{aligned}
& T_{x^{3}}=\mathbf{e}_{9} \mathbf{e}_{5} \mathbf{e}_{1}, T_{y^{3}}=\mathbf{e}_{10} \mathbf{e}_{6} \mathbf{e}_{2} \\
& T_{x y^{2}}=\frac{1}{3}\left(\mathbf{e}_{10} \mathbf{e}_{6} \mathbf{e}_{1}+\mathbf{e}_{10} \mathbf{e}_{5} \mathbf{e}_{2}+\mathbf{e}_{9} \mathbf{e}_{6} \mathbf{e}_{2}\right), T_{x^{2} y}=\frac{1}{3}\left(\mathbf{e}_{10} \mathbf{e}_{5} \mathbf{e}_{1}+\mathbf{e}_{9} \mathbf{e}_{6} \mathbf{e}_{1}+\mathbf{e}_{9} \mathbf{e}_{5} \mathbf{e}_{2}\right) \\
& T_{x t_{\varepsilon}^{2}}=\frac{1}{3}\left(\mathbf{e}_{\infty 3} \mathbf{e}_{5} \mathbf{e}_{o 1}+\mathbf{e}_{o 3} \mathbf{e}_{5} \mathbf{e}_{\infty 1}+\mathbf{e}_{9} \mathbf{e}_{\infty 2} \mathbf{e}_{o 1}+\mathbf{e}_{9} \mathbf{e}_{o 2} \mathbf{e}_{\infty 1}+\mathbf{e}_{\infty 3} \mathbf{e}_{o 2} \mathbf{e}_{1}+\mathbf{e}_{o 3} \mathbf{e}_{\infty 2} \mathbf{e}_{1}\right) \\
& T_{y t_{\varepsilon}^{2}}=\frac{1}{3}\left(\mathbf{e}_{\infty 3} \mathbf{e}_{6} \mathbf{e}_{o 1}+\mathbf{e}_{o 3} \mathbf{e}_{6} \mathbf{e}_{\infty 1}+\mathbf{e}_{10} \mathbf{e}_{\infty 2} \mathbf{e}_{o 1}+\mathbf{e}_{10} \mathbf{e}_{o 2} \mathbf{e}_{\infty 1}+\mathbf{e}_{\infty 3} \mathbf{e}_{o 2} \mathbf{e}_{2}+\mathbf{e}_{o 3} \mathbf{e}_{\infty 2} \mathbf{e}_{2}\right) \\
& \hline
\end{aligned}
$$

Table 3. Extraction operators $T_{s}$ for quartic, quintic, and sextic plane curves.

$$
\begin{aligned}
& T_{x^{2} t_{\varepsilon}^{2}}=\frac{2}{3}\left(\mathbf{e}_{1} \mathbf{e}_{5} \mathbf{e}_{o 3}+\mathbf{e}_{1} \mathbf{e}_{o 2} \mathbf{e}_{9}+\mathbf{e}_{o 1} \mathbf{e}_{5} \mathbf{e}_{9}\right), T_{y^{2} t_{\varepsilon}^{2}}=\frac{2}{3}\left(\mathbf{e}_{2} \mathbf{e}_{6} \mathbf{e}_{o 3}+\mathbf{e}_{2} \mathbf{e}_{o 2} \mathbf{e}_{10}+\mathbf{e}_{o 1} \mathbf{e}_{6} \mathbf{e}_{10}\right) \\
& T_{x y \mathbf{t}_{\varepsilon}^{2}}=\frac{1}{3}\left(\mathbf{e}_{1} \mathbf{e}_{6} \mathbf{e}_{o 3}+\mathbf{e}_{2} \mathbf{e}_{5} \mathbf{e}_{o 3}+\mathbf{e}_{1} \mathbf{e}_{o 2} \mathbf{e}_{10}+\mathbf{e}_{2} \mathbf{e}_{o 2} \mathbf{e}_{9}+\mathbf{e}_{o 1} \mathbf{e}_{5} \mathbf{e}_{10}+\mathbf{e}_{o 1} \mathbf{e}_{6} \mathbf{e}_{9}\right) \\
& T_{x t_{\varepsilon}^{4}}=\frac{4}{3}\left(\mathbf{e}_{o 3} \mathbf{e}_{o 2} \mathbf{e}_{1}+\mathbf{e}_{o 3} \mathbf{e}_{5} \mathbf{e}_{o 1}+\mathbf{e}_{9} \mathbf{e}_{o 2} \mathbf{e}_{o 1}\right), T_{y t_{\varepsilon}^{4}}=\frac{4}{3}\left(\mathbf{e}_{o 3} \mathbf{e}_{o 2} \mathbf{e}_{2}+\mathbf{e}_{o 3} \mathbf{e}_{6} \mathbf{e}_{o 1}+\mathbf{e}_{10} \mathbf{e}_{o 2} \mathbf{e}_{o 1}\right) \\
& T_{\mathbf{t}_{\varepsilon}^{4}}=\frac{4}{3}\left(\mathbf{e}_{\infty 1} \mathbf{e}_{o 2} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{\infty 2} \mathbf{e}_{o 3}+\mathbf{e}_{o 1} \mathbf{e}_{o 2} \mathbf{e}_{\infty 3}\right), T_{\mathbf{t}_{\varepsilon}^{6}}=8 \mathbf{e}_{o}=8 \mathbf{e}_{o 1} \mathbf{e}_{o 2} \mathbf{e}_{o 3} \\
& \hline
\end{aligned}
$$

[^6]
## 3.3 | TCGA IPNS 3-vector entities

The most general TCGA IPNS 3-vector plane curve entity $\boldsymbol{\Omega}$ is a linear combination of the twenty extraction operators $T_{s}$

$$
\begin{equation*}
\mathbf{\Omega}=\sum_{s} c_{s} T_{s} \tag{45}
\end{equation*}
$$

A linear combination of extraction operators from Tables 1 and 2 forms a TCGA IPNS 3-vector general quadratic or cubic plane curve entity. A linear combination that includes the extraction operators from Table 3 can form an entity for certain types of quartic (degree 4), quintic (degree 5), and sextic (degree 6) plane curves.

The TCGA dual of any TCGA IPNS $k$-vector (or blade) entity $\mathbf{E}$ is its corresponding TCGA OPNS ( $12-k$ )-vector (or blade) entity

$$
\begin{equation*}
\mathbf{E}^{*}=\mathbf{E I}_{\mathcal{T}}^{-1} \tag{46}
\end{equation*}
$$

which represents the same geometric point set or plane curve, where $\mathbf{T}_{\mathcal{T}} \wedge \mathbf{E}^{*}=0$.

## 3.4 | TCGA IPNS intersection entities

The outer product of $2 \leq k \leq 4$ TCGA IPNS 3-blade entities forms the TCGA IPNS $3 k$-blade entity that represents their intersection, which corresponds to the CGA entities but has multiplicity 3.

The outer product of any one TCGA IPNS 3-vector entity $\boldsymbol{\Omega}$ and one or two ( $k \in\{1,2\}$ ) TCGA IPNS 3-blade entities forms the TCGA IPNS $3(k+1)$-vector entity that represents their intersection. The outer product of any two 3-vector entities (in the form of $\boldsymbol{\Omega}$ ) cannot be interpreted straightforward as intersection. Therefore, it is possible by the same method to intersect any TCGA entity with a 3-blade line or circle, but it is not possible to intersect the other general cubic and certain quadric, quintic, and sextic entities with each other. The reason for this is the general non-blade character ${ }^{14}$ of TCGA tri-vector entities for cubic curves, and is explained in more detail in the corresponding section of [10].

## 3.5 | TCGA versor operations

The CGAi 1-versors $\mathbf{C}_{C^{i}}, \mathbf{L}_{\mathcal{C}^{i}}$ and 2-versors $T_{\mathcal{C}^{i}}, R_{\mathcal{C}^{i}}$, and $D_{C^{i}}$ (collectively indicated by $V_{C^{i}}$ of grade $k \in\{1,2\}$ ) each have a representation as a TCGA $3 k$-versor $V_{\mathcal{T}}=V_{C^{1}} V_{C^{2}} V_{C^{3}}$ (the product of the same versor as represented in each CGA $i$ ). By outermorphism [25], a TCGA versor $V_{\mathcal{T}}$ operates correctly, as $\mathbf{E}^{\prime}=V_{\mathcal{J}} \mathbf{E} V_{\mathcal{T}}^{-1}$, on each CGA $i$ element or factor of any TCGA IPNS entity $\mathbf{E}$ (or its OPNS dual $\mathbf{E}^{*}$ ), and therefore a TCGA versor $V_{\mathcal{T}}$ operates correctly on any TCGA entity $\mathbf{E}$.

In particular, we can rotate (by rotor $R_{\mathcal{T}}=R_{\mathcal{C}^{1}} R_{\mathcal{C}^{2}} R_{\mathcal{C}^{3}}$ ), translate $\left(T_{\mathcal{T}}\right)$, and isotropically dilate ( $D_{\mathcal{T}}$ ) any TCGA geometric entity $\boldsymbol{\Omega}$ for general cubic plane curves and any of the types of quartic, quintic, and sextic plane curves that can be formed as linear combinations of extraction operators $T_{s}$. For any TCGA entity $\mathbf{E}$ (which may be a quadric, cubic, etc.) and any TCGA IPNS 3-blade circle $\mathbf{C}_{\mathcal{T}}$, then $\mathbf{E}^{\prime}=\mathbf{C}_{\mathcal{T}} \mathbf{E} \mathbf{C}_{\mathcal{T}}^{\sim}$ is the correct inversion of $\mathbf{E}$ in the circle $\mathbf{C}_{\mathcal{T}}$. Similarly, any TCGA entity can be reflected in the TCGA 3-blade line $\mathbf{L}_{\mathcal{T}}$.

We note, that only the use of three full copies of CGA in TCGA allows to use all conformal versor operators known from CGA. If instead only one or two copies of CGA would be used for $\mathbb{R}^{2}$, combined with two or one copies of GA for $\mathbb{R}^{2}$, then it would, e.g., not be possible to use the above translation operators for cubic plane curves, etc.

## 3.6 | TCGA differential operators

The extraction operators $T_{x^{3}}$ and $T_{y^{3}}$ are 3-blades that have inverses. This makes it possible to form the following two ratios of extraction operators

$$
\begin{equation*}
D_{x}=3 T_{x^{2}} T_{x^{3}}^{-1}, \quad D_{y}=3 T_{y^{2}} T_{y^{3}}^{-1} \tag{47}
\end{equation*}
$$

Clearly, $D_{x} T_{x^{3}}=3 T_{x^{2}}$ and $D_{y} T_{y^{3}}=3 T_{y^{2}}$, such that $D_{x}$ and $D_{y}$ act as differential operators on the extraction operators $T_{s}$. Using the commutator product $\times$, defined for any two multivectors $A$ and $B$ as

$$
\begin{equation*}
A \times B=\frac{1}{2}(A B-B A)=-B \times A \tag{48}
\end{equation*}
$$

[^7][^8]Table 4. Differential operations $D_{x} \times T_{s}, D_{y} \times T_{s}$ on quadric $T_{s}$ of Table 1 , and $\mathbf{t}_{\mathcal{E}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}$.

| $\times$ | $T_{x}$ | $T_{y}$ | $T_{x^{2}}$ | $T_{x y}$ | $T_{y^{2}}$ | $T_{1}$ | $T_{\mathbf{t}_{\varepsilon}^{2}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $D_{x}$ | $T_{1}$ | 0 | $2 T_{x}$ | $T_{y}$ | 0 | 0 | $2 T_{x}$ |
| $D_{y}$ | 0 | $T_{1}$ | 0 | $T_{x}$ | $2 T_{y}$ | 0 | $2 T_{y}$ |

Table 5. Differential operations $D_{x} \times T_{s}, D_{y} \times T_{s}$ on cubic $T_{s}$ of Table 2.

| $\times$ | $T_{x^{3}}$ | $T_{x y^{2}}$ | $T_{x^{2} y}$ | $T_{y^{3}}$ | $T_{x t_{\varepsilon}^{2}}$ | $T_{y t_{\varepsilon}^{2}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $D_{x}$ | $3 T_{x^{2}}$ | $T_{y^{2}}$ | $2 T_{x y}$ | 0 | $2 T_{x^{2}}+T_{\mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{x y}$ |
| $D_{y}$ | 0 | $2 T_{x y}$ | $T_{x^{2}}$ | $3 T_{y^{2}}$ | $2 T_{x y}$ | $2 T_{y^{2}}+T_{\mathbf{t}_{\varepsilon}^{2}}$ |

Table 6. Differential operations $D_{x} \times T_{s}, D_{y} \times T_{s}$ on quartic, quintic, and sextic $T_{s}$ of Table 3.

| $\times$ | $T_{\mathbf{t}_{\varepsilon}^{4}}$ | $T_{x^{2} \mathbf{t}_{\varepsilon}^{2}}$ | $T_{x y \mathbf{t}_{\varepsilon}^{2}}$ | $T_{y^{2} \mathbf{t}_{\varepsilon}^{2}}$ | $T_{x \mathbf{t}_{\varepsilon}^{4}}$ | $T_{y \mathbf{t}_{\varepsilon}^{4}}$ | $T_{\mathbf{t}_{\varepsilon}^{6}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $D_{x}$ | $4 T_{x \mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{x^{3}}+2 T_{x \mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{x^{2} y}+T_{y \mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{x y^{2}}$ | $4 T_{x^{2} \mathbf{t}_{\varepsilon}^{2}}+T_{\mathbf{t}_{\varepsilon}^{4}}$ | $4 T_{x y \mathbf{t}_{\varepsilon}^{2}}$ | $6 T_{x \mathbf{t}_{\varepsilon}^{4}}$ |
| $D_{y}$ | $4 T_{y \mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{x^{2} y}$ | $2 T_{x y^{2}}+T_{x \mathbf{t}_{\varepsilon}^{2}}$ | $2 T_{y^{3}}+2 T_{y \mathbf{t}_{\varepsilon}^{2}}$ | $4 T_{x y \mathbf{t}_{\varepsilon}^{2}}$ | $4 T_{y^{2} \mathbf{t}_{\varepsilon}^{2}}+T_{\mathbf{t}_{\varepsilon}^{4}}$ | $6 T_{y \mathbf{t}_{\varepsilon}^{4}}$ |

For a unit direction $\mathbf{n}=n_{x} \mathbf{e}_{1}+n_{y} \mathbf{e}_{2}$ in $\mathbb{R}^{2}, \mathbf{n}^{2}=1$, the $\mathbf{n}$-directional derivative operator is $D_{\mathbf{n}}=n_{x} D_{x}+n_{y} D_{y}$ and the $\mathbf{n}$-directional derivative of any TCGA IPNS 3-vector entity $\boldsymbol{\Omega}$ representing an implicit plane curve function $F(x, y)=\mathbf{T}_{\mathcal{J}} \cdot \boldsymbol{\Omega}$ is $\partial_{\mathbf{n}} \boldsymbol{\Omega}=D_{\mathbf{n}} \times \boldsymbol{\Omega}$, which represents the implicit plane curve function $\partial_{\mathbf{n}} F=(\boldsymbol{\nabla} \cdot \mathbf{n}) F=\mathbf{T}_{\mathcal{T}} \cdot\left(D_{\mathbf{n}} \times \boldsymbol{\Omega}\right)$, where the gradient operator is the geometric calculus symbolic vector $\boldsymbol{\nabla}=(\partial / \partial x) \mathbf{e}_{1}+(\partial / \partial y) \mathbf{e}_{2}$. For a TCGA IPNS 3-blade 3-CGA entity $\mathbf{C}_{\mathcal{T}}$, which represents the third power ${ }^{15}$ of a function $F^{3}(x, y)=F(x, y) F(x, y) F(x, y)$, the product $\mathbf{T}_{\mathcal{T}} \cdot\left(D_{x} \times \mathbf{C}_{\mathcal{T}}\right)=3 F^{2} \partial_{x} F$ etc., which is in agreement with the chain rule of differentiation. Mixed partial derivatives of higher orders are formed by nesting successive differential operations in any sequence (e.g., $D_{y} \times\left(D_{x} \times \boldsymbol{\Omega}\right)$ or $D_{x} \times\left(D_{y} \times \boldsymbol{\Omega}\right)$ ), etc.

## 3.7 | Examples of plane curve TCGA 3-vector entities

The TCGA plane curves are represented in algebraic geometry by implicit plane curve functions $F(x, y)$, with equation $F(x, y)=$ 0 , that have a straightforward representation as TCGA entities (i.e., as linear combinations of the extraction operators in Tables 1, 2, and 3), as will be shown in the following subsections. Points of a plane curve are solutions of the implicit equation $F(x, y)=0$, equivalent to the inner product null space of a TCGA IPNS 3-vector plane curve entity $\boldsymbol{\Omega}$ (i.e., points $\mathbf{P}_{\mathcal{T}}$ where $\mathbf{P}_{\mathcal{T}} \cdot \boldsymbol{\Omega}=0$ ). Many functions $F(x, y)$ include some real scalar parameters $a, b$, and $c$ that vary the curve.

An exhaustive description of all possible linear, quadric, cubic, quartic, quintic, and sextic plane curves that can be represented as TCGA IPNS 3-vector entities would be beyond the limits of this introductory paper. The current subsection instead presents some examples of TCGA IPNS 3-vector entities for some of the plane curves commonly discussed in algebraic geometry literature.

[^9]

Figure 1. Examples of cubics: right strophoid, nodal cubic, and folium of Descartes.

### 3.7.1 | General cubic plane curves

TCGA IPNS 3-vector entities can represent general cubic (degree 3) plane curves as linear combinations of the general quadric (degree 2) and general cubic extraction operators listed in Tables 1 and 2. In algebraic geometry, references [24] and [12] contain numerous implicit cubic plane curve equations $F(x, y)=0$, all of which can be represented as TCGA IPNS 3-vector entities. Examples of TCGA cubic plane curve entities are shown in Table 7.

Table 7. Examples of TCGA IPNS 3-vector cubic plane curve entities.

| Name | Cubic $\boldsymbol{F}(\boldsymbol{x}, \boldsymbol{y})$ | TCGA IPNS 3-vector entity |
| :--- | :--- | :--- |
| Apollonian cubic | $(x-a)\left(x^{2}+y^{2}\right)+b x+c y$ | $T_{x t_{\varepsilon}^{2}}-a T_{\mathbf{t}_{\varepsilon}^{2}}+b T_{x}+c T_{y}$ |
| cissiod of Diocles | $x^{3}+x y^{2}-y^{2}$ | $T_{x^{3}}+T_{x y^{2}}-T_{y^{2}}$ |
| cubic egg | $a x y^{2}+x^{2}+y^{2}-1$ | $a T_{x y^{2}}+T_{\mathbf{t}_{\varepsilon}^{2}}-T_{1}$ |
| egg of Newton | $y^{2}-x^{3}+a x^{2}+x-a$ | $T_{y^{2}}-T_{x^{3}}+a T_{x^{2}}+T_{x}-a T_{1}$ |
| folium of Descartes | $x^{3}+y^{3}-\sqrt{2} a x y$ | $T_{x^{3}}+T_{y^{3}}-\sqrt{2} a T_{x y}$ |
| nodal cubic | $y^{2}-x^{2}(x-a): a<0$ | $T_{y^{2}}-T_{x^{3}}+a T_{x^{2}}$ |
| serpentine | $x^{2} y+a^{2} y-a b x$ | $T_{x^{2} y}+a^{2} T_{y}-a b T_{x}$ |
| right strophoid | $y^{2}(x+a)+x^{2}(x-a)$ | $T_{x y^{2}}+a T_{y^{2}}+T_{x^{3}}-a T_{x^{2}}$ |
| trident of Newton | $x y-x^{3}-a x^{2}-b x-c$ | $T_{x y}-T_{x^{3}}-a T_{x^{2}}-b T_{x}-c T_{1}$ |
| witch of Agnesi | $x^{2} y+4 a^{2} y-8 a^{3}$ | $T_{x^{2} y}+4 a^{2} T_{y}-8 a^{3} T_{1}$ |

Figure 1 (A) shows the inversion of the quadric hyperbola $\boldsymbol{\alpha}$ in the circle $\mathbf{C}_{\mathcal{T}}$ as the cubic right strophoid $\boldsymbol{\beta}=\mathbf{C}_{\mathcal{T}} \boldsymbol{\alpha} \mathbf{C}_{\mathcal{J}}^{\sim}$ for $a=-4$. The hyperbola entity $\alpha=T_{\mathcal{J}}\left(T_{x^{2}}-T_{y^{2}}-r^{2} T_{1}\right) T_{\mathcal{J}}^{\sim}$ has radius $r=2$ and is translated by displacement $\mathbf{d}_{\mathcal{E}}=-r \mathbf{e}_{1}$ using the TCGA translator $T_{\mathcal{T}}$. The TCGA IPNS 3-blade circle $\mathbf{C}_{\mathcal{T}}$ has radius $r=4$ and center $\mathbf{e}_{o}$ (the origin). The right strophoid $\boldsymbol{\beta}$ is an entity that can also be represented in Double Conformal Geometric Algebra (DCGA) [5].

Figure 1 (B) shows the inversion of the nodal cubic $\delta$ for $a=-4$ in the circle $\mathbf{C}_{\mathcal{T}}$ (of radius $r=4$ centered on $\mathbf{e}_{o}$ ) as the quartic "defective hyperbola" $\gamma$ for $a=-4$ (see quartics Table 8).

Figure $1(\mathrm{C})$ shows a right strophoid $\beta$ and nodal cubic $\delta$, both for $a=-1$, which are both dilated by factor $d=4$ and rotated by angle $\theta=225^{\circ}$ using the TCGA dilator $D_{\mathcal{T}}$ and rotor $R_{\mathcal{T}}$ as $\boldsymbol{\beta}^{\prime}=R_{\mathcal{T}} D_{\mathcal{T}} \beta D_{\mathcal{T}}^{\sim} R_{\mathcal{T}}^{\sim}$ and $\delta^{\prime}=R_{\mathcal{J}} D_{\mathcal{T}} \delta D_{\mathcal{T}}^{\sim} R_{\mathcal{T}}^{\sim}$. The right
strophoid $\boldsymbol{\beta}^{\prime}$ and nodal cubic $\boldsymbol{\delta}^{\prime}$ both look similar to a folium of Descartes. However, the true folium of Descartes $\boldsymbol{\epsilon}$ for $a=4$ is also shown, and the shape differences are apparent.

By using TCGA entities and versor operations, it is easy to use TCGA as a tool to study the plane curves of Figure 1, and many others, and to learn how plane curves are related to each other through inversions in circles and by rotations, dilations, and translations.

### 3.7.2 | Various quartic plane curves

TCGA IPNS 3-vector entities can represent only the specific forms of quartic (degree 4) plane curves that can be written in terms of the general cubic extraction operators (all of Tables 1 and 2) plus the available quartic extraction operators of Table 3 , which can represent only the quartic terms $\mathbf{t}_{\mathcal{E}}^{4}, x^{2} \mathbf{t}_{\mathcal{E}}^{2}, x y \mathbf{t}_{\mathcal{E}}^{2}$, and $y^{2} \mathbf{t}_{\mathcal{E}}^{2}$, where $\mathbf{t}_{\mathcal{E}}^{2}=x^{2}+y^{2}$. TCGA cannot represent the individual general quartic terms $x^{4}, x^{2} y^{2}$, and $y^{4}$ with independent coefficients, because there are no extraction trivector operators to individually extract these coefficients from a general test point. Even with this limitation, there is a large variety of quartics in the algebraic geometry literature (e.g., in [24]) that can be written with this restricted choice of TCGA quartic terms. Some of the possible TCGA quartic plane curve entities are shown in Table 8.

Table 8. Examples of TCGA IPNS 3-vector quartic plane curve entities.

| Name | TCGA IPNS 3-vector quartic entity |
| :---: | :---: |
| cardioid | $T_{\mathbf{t}_{\varepsilon}^{4}}-4 a T_{\chi_{\text {xt }}{ }_{\varepsilon}^{2}}+4 a^{2} T_{x^{2}}-4 a^{2} T_{\mathbf{t}_{\varepsilon}^{2}}$ |
| Cartesian oval | $\begin{aligned} & b^{2}\left(T_{\mathbf{t}_{\varepsilon}^{4}}+2 a^{2} T_{\mathbf{t}_{\varepsilon}^{2}}+a^{4} T_{1}\right)-4 a b c\left(T_{x t_{\varepsilon}^{2}}+a^{2} T_{x}\right)+ \\ & 4 a^{2} c^{2} T_{x^{2}}-2 c\left(T_{\mathbf{t}_{\varepsilon}^{2}}+a^{2} T_{1}\right)+4 a b T_{x}+T_{1} \end{aligned}$ |
| Cassini oval | $T_{\mathbf{t}_{\varepsilon}^{4}}-2 a^{2}\left(T_{x^{2}}-T_{y^{2}}\right)+\left(a^{4}-b^{4}\right) T_{1}$ |
| conchoid of Durer | $\begin{aligned} & 2 T_{y^{2} \varepsilon_{\varepsilon}^{2}}-2 b\left(T_{x y^{2}}+T_{y^{3}}\right)+\left(b^{2}-3 a^{2}\right) T_{y^{2}}-a^{2} T_{x^{2}}+ \\ & 2 a^{2} b\left(T_{x}+T_{y}\right)+a^{2}\left(a^{2}-b^{2}\right) T_{1} \end{aligned}$ |
| conchoid of Nicomedes | $T_{x^{2} t_{\varepsilon}^{2}}-2 b T_{x t_{\varepsilon}^{2}}+b^{2} T_{t_{\varepsilon}^{2}}-a^{2} T_{x^{2}}$ |
| defective hyperbola | $a T_{x^{2} t_{\varepsilon}^{2}}+T_{y^{2} t_{\varepsilon}^{2}}-a^{2} T_{x^{3}}: a<0$ |
| deltoid (tricuspid) | $T_{\mathbf{t}_{\varepsilon}^{4}}-8 a\left(T_{x^{3}}-3 T_{x y^{2}}\right)+18 a^{2} T_{t_{\varepsilon}^{2}}-27 a^{4} T_{1}$ |
| hippopede | $T_{\mathbf{t}_{\varepsilon}^{4}}+4 b(b-a) T_{\mathbf{t}_{\varepsilon}^{2}}-4 b^{2} T_{x^{2}}$ |
| kappa | $T_{y^{2} t_{\varepsilon}^{2}}-a^{2} T_{x^{2}}$ |
| limacon | $T_{\mathbf{t}_{\varepsilon}^{4}}-4 a T_{\mathbf{x t}}{ }_{\underline{2}}+4 a^{2} T_{x^{2}}-b^{2} T_{\mathbf{t}_{\varepsilon}^{2}}$ |
| ovoid | $a T_{x^{2} t_{\varepsilon}^{2}}+T_{y^{2} t_{\varepsilon}^{2}}-a^{2} T_{x^{3}}: a>0$ |

A type of hippopede $\mathbf{H}$ can be formed as the inversion $\mathbf{H}=\mathbf{C}_{\mathcal{T}} \mathbf{E} \mathbf{C}_{\mathcal{T}}^{\sim}$ of an ellipse $\mathbf{E}=T_{x^{2}} / a^{2}+T_{y^{2}} / b^{2}-T_{1}, a>b$, in the circle $\mathbf{C}_{\mathcal{T}}$ of radius $r=b$ centered on $\mathbf{e}_{o}$. H has $y$-intercepts $\pm b$ and $x$-intercepts $\pm c=b^{2} / a$ (the inversion of $\pm a$ ). For fixed $b$, then $a=b^{2} / c$ gives $\mathbf{H}$ with $x, y$-intercepts $\pm c, \pm b$. As $a \rightarrow \infty$ for fixed $b, \mathbf{E}=T_{y^{2}} / b^{2}-T_{1}$ and $\mathbf{H}$ becomes a particular lemniscate of Booth (figure eight) that is the union of tangential spheres of radius $b$ through the origin. The implicit curve function $F(x, y)=\mathbf{T}_{\mathcal{T}} \cdot \mathbf{H}$ has a factor $\mathbf{t}_{\mathcal{E}}^{2}=x^{2}+y^{2}$ (making it appear sextic) representing the inversion $\mathbf{e}_{o}=\mathbf{C}_{\mathcal{J}} \mathbf{e}_{\infty} \mathbf{C}_{\mathcal{J}}^{\sim}$ since all entities, such as $\mathbf{E}$, not having the term $T_{\mathbf{t}_{\varepsilon}^{6}}=8 \mathbf{e}_{o}$ include $\mathbf{e}_{\infty}$ as a singular outlier handle point (exterior point). $F(x, y) / \mathbf{t}_{\mathcal{E}}^{2}$ is the quartic hippopede function.

Many of the quartics of Table 8, including $\mathbf{H}$ and hippopede, can also be formed in DCGA [5].
Figure 2 shows the hippopede $\mathbf{H}$ that has been discussed above and also a cardioid, Cartesian oval, and deltoid. The larger cardioid, call it $\mathbf{A}$, with $a=1$, was rotated $\theta=-90^{\circ}$, then dilated by $d=1 / 2$, and then translated by $\mathbf{d}_{\mathcal{E}}=2 \mathbf{e}_{1}$ into the smaller cardioid $\mathbf{A}^{\prime}=T_{\mathcal{J}} D_{\mathcal{T}} R_{\mathcal{J}} \mathbf{A} R_{\mathcal{J}}^{\sim} D_{\mathcal{T}}^{\sim} T_{\mathcal{J}}^{\sim}$. The Cartesian oval and the deltoid were both reflected in the circle (inversion in circle).


Figure 2. Examples of quartics: the hippopede, cardioid, Cartesian oval, and deltoid.

The reflected deltoid makes a shape that looks similar to an axe head, which is then rotated $180^{\circ}$ and then reflected again in the circle to make the larger B-like shape. The larger B-like shape is then dilated by $d=1 / 4$ (centered on $3 \mathbf{e}_{1}$ ) and translated by $\mathbf{d}_{\mathcal{E}}=4 \mathbf{e}_{1}$ to make the smaller B-like shape.

### 3.7.3 | Various quintic plane curves

In Table 3, we find that TCGA offers only two quintic (degree 5) extraction operators that can represent only the quintic terms $x \mathbf{t}_{\mathcal{E}}^{4}$ and $y \mathbf{t}_{\mathcal{E}}^{4}$, where $\mathbf{t}_{\mathcal{E}}^{2}=x^{2}+y^{2}$. This may seem limited, but in combination with all of the quartic and lesser degree terms available in TCGA, there is still a large variety of possible TCGA IPNS 3-vector quintic plane curve entities. The literature on quintic plane curves appears sparse, so we have only a couple of ad-hoc examples in Table 9.

Table 9. Examples of TCGA IPNS 3-vector quintic plane curve entities.

| Description | TCGA IPNS 3-vector quintic entity |
| :--- | :--- |
| modified Burnside curve | $T_{y^{2}}-T_{x t_{\varepsilon}^{4}}+T_{x}$ |
| serpentine-like | $T_{x \varepsilon_{\varepsilon}^{4}}-T_{x^{3}}-T_{y^{3}} / 125-\left(T_{x^{2}}+T_{y^{2}}\right) / 144+\left(T_{x}+T_{y}\right) / 100$ |

### 3.7.4 | Cranioid sextic plane curves

A sextic (degree 6) cranioid (head-like shape) plane curve can be formed by inversion in a circle of a cubic curve that is basically a straight line except for being curved as it passes through the origin (see Figure 3). Such a cubic curve is

$$
\begin{equation*}
y-b+\frac{b}{x^{2}+1}=0 \tag{49}
\end{equation*}
$$

As $x \rightarrow \infty$, it becomes the line $y=b$. As $x \rightarrow 0$, then $y \rightarrow 0$. The function is of even degree 2 in $x$ and symmetrical for $\pm x$. After multiplying by $x^{2}+1$, the TCGA IPNS 3-vector entity $\mathbf{B}_{0}$ for this cubic plane curve can be directly expressed in terms of the TCGA IPNS extraction operators as

$$
\begin{equation*}
\mathbf{B}_{0}=T_{x^{2} y}+T_{y}-b T_{x^{2}} . \tag{50}
\end{equation*}
$$

We will use $b=-1$, such that the curvature is upwards the $y$-axis. Instead of using $x^{2}$, we can use $\mathbf{t}_{\mathcal{E}}^{4}=\left(x \mathbf{e}_{1}+y \mathbf{e}_{2}\right)^{4}=$ $\left(x^{2}+y^{2}\right)^{2}$ for a sharper curvature. The TCGA IPNS 3-vector entity $\mathbf{B}_{1}$ for this quintic plane curve is

$$
\begin{equation*}
\mathbf{B}_{1}=T_{y_{\varepsilon}^{4}}+T_{y}-b T_{\mathbf{t}_{\varepsilon}^{4}} \tag{51}
\end{equation*}
$$



Figure 3. Cranioids $\mathfrak{C}_{0}, \mathfrak{C}_{0.5}$, and $\mathfrak{C}_{1}$.

The TCGA IPNS 3-vector entity $\mathbf{B}_{t}, 0 \leq t \leq 1$, that interpolates between $\mathbf{B}_{0}$ and $\mathbf{B}_{1}$ is

$$
\begin{equation*}
\mathbf{B}_{t}=(1-t) \mathbf{B}_{0}+t \mathbf{B}_{1} . \tag{52}
\end{equation*}
$$

Using the TCGA IPNS 3-blade circle $\mathbf{C}_{\mathcal{T}}=\mathbf{C}_{C^{1}} \mathbf{C}_{C^{2}} \mathbf{C}_{C^{3}}$ with center $2 \mathbf{e}_{2}$ and radius $r=3$, and $\mathbf{B}_{t}$ with parameters $0 \leq t \leq 1$ and $b=-1$, the TCGA IPNS 3-vector cranioid entity $\mathfrak{C}_{t}$ is the inversion

$$
\begin{equation*}
\mathfrak{C}_{t}=\mathbf{C}_{\mathcal{T}} \mathbf{B}_{t} \mathbf{C}_{\mathcal{T}}^{\sim} . \tag{53}
\end{equation*}
$$

For $t<0$ and $t>1$, the entity $\mathfrak{C}_{t}$ is generally no longer a cranioid but forms other interesting curves. Choosing a different $b$ for the curve $\mathbf{B}_{t}$ and different circles $\mathbf{C}_{\mathcal{T}}$ for the inversion could be the subject of further experimentation. Further research could also lead to various ways to parametrize the cranioid entity for certain shapes.

Figure 3 shows three cranioids $\mathfrak{C}_{0}$, $\mathfrak{C}_{0.5}$, and $\mathfrak{C}_{1}$, and it also shows a cranioid $\mathfrak{G}_{0.5}^{\prime}=T_{\mathcal{T}} R_{\mathcal{T}} \mathfrak{C}_{0.5} R_{\mathcal{J}}^{\sim} T_{\mathcal{T}}^{\sim}$ that has been rotated $35^{\circ}$ using a TCGA rotor $R_{\mathcal{T}}$ and then translated by $\mathbf{d}_{\mathcal{E}^{1}}=6 \mathbf{e}_{1}+3 \mathbf{e}_{2}$ using a TCGA translator $T_{\mathcal{T}}$.

Cayley's sextic, with equation $4\left(x^{2}+y^{2}-a x\right)^{3}=27 a^{2}\left(x^{2}+y^{2}\right)^{2}$, can also be represented as a TCGA entity.

## 3.8 | The anisotropic dilation operation

A TCGA 3-vector general cubic entity $\mathbf{X}_{\mathcal{T}}$ (a linear combination of extraction operators from Tables 1 and 2) can represent any cubic plane curve, including those that have been anisotropically dilated by a factor $d>0$ in an arbitrary direction $\hat{\mathbf{d}}_{\mathcal{E}}$ (directed scaling). The TCGA 6-versor dilator $D_{\mathcal{T}}$ is only for isotropic dilation (uniform scaling), and we have not been able to give an anisotropic dilation versor in $\mathcal{G}_{9,3}$ TCGA. However, the following describes a successful extension of TCGA in which we can obtain an anisotropic dilation versor $B_{\mathcal{T}}$ that is valid on any general TCGA cubic entity $\mathbf{X}_{\mathcal{T}}$.

TCGA can be generalized further by adding ${ }^{16}$ three new unit vector elements $\left[\mathbf{e}_{i} \cdot \mathbf{e}_{j}\right]=\left[m_{i j}\right]=\operatorname{diag}(-1,-1,-1): i, j \in$ $\{13,14,15\}$ as time-like dimensions, forming a triple of Minkowski space-times $\mathcal{M}^{i}: i \in\{1,2,3\}$ of signatures (2, 1), each having two space dimensions and a time-like dimension. In this new space-time, vectors are written ${ }^{17}$ as $\mathbf{t}_{\mathcal{M}^{1}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}+w \mathbf{e}_{13}$, and similarly for $\mathbf{t}_{\mathcal{M}^{2}}$ and $\mathbf{t}_{\mathcal{M}^{3}}$ using $\mathbf{e}_{14}$ and $\mathbf{e}_{15}$, respectively. We still have Euclidean spatial vectors $\mathbf{t}_{\mathcal{E}^{1}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}$, and similarly for $\mathbf{t}_{\mathcal{E}^{2}}$ and $\mathbf{t}_{\mathcal{E}^{3}}$, as previously defined. The new time-like coordinate $w$ can be written $w=c t$, with light speed $c=1$

[^10]and time $t$. In this larger $\mathcal{C}_{3(2+1), 3(1+1)} 15-\mathrm{D}$ algebra, which is less practicable ${ }^{18}$, it becomes now possible to define the unimodular hyperbolic rotor (boost ${ }^{19}$ versor) $B_{\mathcal{M}^{1}}=\exp \left(\varphi_{d} \hat{\mathbf{d}}_{\mathcal{E}^{1}} \mathbf{e}_{13} / 2\right)$ for anisotropic dilation in direction $\hat{\mathbf{d}}_{\mathcal{E}^{1}}$ by dilation factor $d>0$, with hyperbolic angle (rapidity) $\varphi_{d}=\operatorname{atanh}\left(-\beta_{d}\right)$ and natural speed $\beta_{d}=\sqrt{1-d^{2}}$, in accordance with the length $L_{0}$ contraction $L=d L_{0}=\gamma_{d}^{-1} L_{0}=\sqrt{1-\beta_{d}^{2}} L_{0}$ that is well-known in physics as part of the theory of special relativity. Note that, $\beta_{d}$ and $\varphi_{d}$ are imaginary numbers for dilation factor $d>1$, and then (just for this anisotropic dilation operation) we must use the triple conformal Clifford algebra $\mathcal{C} \ell(3(2+1), 3(1+1))$ for the complex vector space $\mathbb{C}^{2,1}$. Also note that, the rotation would be by the negative angle $-\varphi_{d}$ in the other Minkowski spacetimes with signatures $(1,2)$. The boost versor $B_{\mathcal{M}^{1}}$ corresponds to a space-time boost from the stationary observer space-time velocity vector $\mathbf{o}_{\mathcal{M}^{1}}=c \mathbf{e}_{13}$ into the space-time velocity vector $\gamma_{d} \mathbf{d}_{\mathcal{M}^{1}}=\gamma_{d}\left(\beta_{d} c \hat{\mathbf{d}_{\mathcal{E}^{1}}}+c \mathbf{e}_{13}\right)$, where
\[

$$
\begin{gather*}
\mathbf{d}_{\mathcal{M}^{1}}=\beta_{d} c \hat{\mathbf{d}}_{\mathcal{E}^{1}}+c \mathbf{e}_{13}, \quad \mathbf{o}_{\mathcal{M}^{1}}=c \mathbf{e}_{13}, \quad \mathbf{o}_{\mathcal{M}^{1}}^{-1}=-\mathbf{e}_{13} / c  \tag{54}\\
H=\mathbf{d}_{\mathcal{M}^{1}} / \mathbf{o}_{\mathcal{M}^{1}}=1-\beta_{d} \hat{\mathbf{d}}_{\mathcal{E}^{1}} \mathbf{e}_{13}, \quad|H|=\sqrt{1^{2}-\beta_{d}^{2}}=\gamma_{d}^{-1}, \quad \varphi_{d}=\operatorname{atanh}\left(-\beta_{d}\right),  \tag{55}\\
\hat{H}=\cosh \left(\varphi_{d}\right)+\sinh \left(\varphi_{d}\right) \hat{\mathbf{d}}_{\mathcal{E}^{1}} \mathbf{e}_{13}=\exp \left(\varphi_{d} \hat{\mathbf{d}}_{\mathcal{E}^{1}} \mathbf{e}_{13}\right)=|H|^{-1} H=\gamma_{d} H  \tag{56}\\
B_{\mathcal{M}^{1}}=\sqrt{\hat{H}}=\sqrt{\gamma_{d} H}=\sqrt{\gamma_{d} \mathbf{d}_{\mathcal{M}^{1}} \mathbf{o}_{\mathcal{M}^{1}}^{-1}}=\exp \left(\varphi_{d} \hat{\mathbf{d}}_{\mathcal{E}^{1}} \mathbf{e}_{13} / 2\right), \quad B_{\mathcal{M}^{1}}^{2}=\gamma_{d} H  \tag{57}\\
\gamma_{d} \mathbf{d}_{\mathcal{M}^{1}}=B_{\mathcal{M}^{1}}^{2} \mathbf{o}_{\mathcal{M}^{1}}=B_{\mathcal{M}^{1}} \mathbf{o}_{\mathcal{M}^{1}} B_{\mathcal{M}^{1}}^{\sim} \tag{58}
\end{gather*}
$$
\]

Similarly, we have the hyperbolic rotors in $\mathcal{M}^{2}$ and $\mathcal{M}^{3}$ as

$$
\begin{equation*}
B_{\mathcal{M}^{2}}=\exp \left(\varphi_{d} \hat{\mathbf{d}}_{\mathcal{E}^{2}} \mathbf{e}_{14} / 2\right), \quad B_{\mathcal{M}^{3}}=\exp \left(\varphi_{d} \hat{\mathbf{d}}_{\mathcal{E}^{3}} \mathbf{e}_{15} / 2\right) \tag{59}
\end{equation*}
$$

The triple hyperbolic rotor is $B_{\mathcal{T}}=\boldsymbol{B}_{\mathcal{M}^{1}} \boldsymbol{B}_{\mathcal{M}^{2}} \boldsymbol{B}_{\mathcal{M}^{3}}$. If we hold $w=0$, then the 15-D $\mathcal{G}_{9,6}$ algebra, denoted $\mathcal{T} \mathcal{M}$, reduces to the 12-D $\mathcal{G}_{9,3}$ TCGA $\mathcal{T}$. In effect, we can set $w=0$ in any entity $\mathbf{X}_{\mathcal{T} \mathcal{M}}$ in $\mathcal{T} \mathcal{M}$ by the projection

$$
\begin{equation*}
\mathbf{X}_{\mathcal{T}}^{\prime}=\left(\mathbf{X}_{\mathcal{T} \mathcal{M}} \cdot \mathbf{I}_{\mathcal{T}}\right) \mathbf{I}_{\mathcal{T}}^{-1}=\mathcal{P}_{\mathbf{I}_{\mathcal{T}}}\left(\mathbf{X}_{\mathcal{T} \mathcal{M}}\right) \tag{60}
\end{equation*}
$$

The anisotropic dilation, centered on (relative to) the origin $\mathbf{e}_{o}$ and in the direction $\hat{\mathbf{d}}_{\mathcal{E}}$ by dilation factor $d>0$, of any TCGA IPNS 3-vector cubic entity $\mathbf{X}_{\mathcal{T}}$ is defined as

$$
\begin{equation*}
\mathbf{X}_{\mathcal{T}}^{\prime}=\mathcal{P}_{\mathbf{I}_{\mathcal{T}}}\left(\boldsymbol{B}_{\mathcal{T}} \mathbf{X}_{\mathcal{T}} B_{\mathcal{T}}^{\sim}\right)=\mathcal{P}_{\mathbf{I}_{\mathcal{T}}}\left(\mathbf{X}_{\mathcal{T} \mathcal{M}}\right) \tag{61}
\end{equation*}
$$

A translated boost versor $B_{\mathcal{T}}^{\mathbf{p}}=T_{\mathcal{T}} B_{\mathcal{T}} T_{\mathcal{J}}^{\sim}$, with translator $T_{\mathcal{T}}$ for displacement $\mathbf{p}_{\mathcal{E}}$, creates a boost versor $B_{\mathcal{T}}^{\mathbf{p}}$ for anisotropic dilation centered on $\mathbf{p}_{\mathcal{E}}$. The boost $\mathbf{X}_{\mathcal{T} \mathcal{M}}=B_{\mathcal{T}}^{\mathbf{p}} \mathbf{X}_{\mathcal{T}}\left(B_{\mathcal{T}}^{\mathbf{p}}\right)^{\sim}$ of a cubic entity $\mathbf{X}_{\mathcal{T}}$ that is centered on $\mathbf{p}_{\mathcal{E}}$ is set into the constant velocity $\mathbf{v}_{\mathcal{E}}=\beta_{d} c \hat{\mathbf{d}}_{\mathcal{E}}$ of the boost, where $\mathbf{X}_{\mathcal{T} \mathcal{M}}$ is centered at $\mathbf{p}_{\mathcal{E}}+\mathbf{v}_{\mathcal{E}} t$ at time $t$, and the shape of $\mathbf{X}_{\mathcal{T} \mathcal{M}}$ and $\mathbf{X}_{\mathcal{T}}^{\prime}$ is contracted (or dilated) in direction $\hat{\mathbf{d}}_{\mathcal{E}}$ by the factor $d=\sqrt{1-\beta_{d}^{2}}$. The projection back to a general TCGA IPNS 3-vector cubic entity $\mathbf{X}_{\mathcal{T}}^{\prime}=\mathcal{P}_{\mathbf{I}_{\mathcal{T}}}\left(\mathbf{X}_{\mathcal{T} \mathcal{M}}\right)$ is at $w=c t=0$ and centered at its original center $\mathbf{p}_{\mathcal{E}}$ (or whatever it was before the boost), but its shape retains the $\hat{\mathbf{d}}_{\mathcal{E}}$-directional dilation by dilation factor $d$ away from the chosen center $\mathbf{p}_{\mathcal{E}}$.

Only the general cubic entities support a form, in terms of the TCGA extraction elements, that can be anisotropically dilated, and any attempt to anisotropically dilate other TCGA entities of degrees 4,5 , or 6 (quartic, quintic, sextic) will result in a different kind of transformation of the shape of the entity.

The anisotropic dilation operation only requires the inclusion of the vector units $\mathbf{e}_{13}, \mathbf{e}_{14}$, and $\mathbf{e}_{15}$, the ability of the algebra to use complex number scalars, and the versor $B_{\mathcal{T}}^{\mathbf{p}}$ and projection $\mathcal{P}_{\mathbf{I}_{\mathcal{T}}}\left(\mathbf{X}_{\mathcal{T} \mathcal{M}}\right)$. No other changes or additions are required to TCGA, including no changes are required to the TCGA point entity or other TCGA entities and versors.

Finally, Figure 4 shows an example for directed (anisotropic) scaling of the folium of Descartes, with $a=1$, dilated by $d=5$ in the diagonal direction $\mathbf{e}_{1}+\mathbf{e}_{2}$.

[^11]

Figure 4. Directed (anisotropic) scaling of the folium of Descartes $a=1$, dilated by $d=5$ in the diagonal direction $\mathbf{e}_{1}+\mathbf{e}_{2}$.

## 4 | IMPLEMENTATION OF TCGA

In symbolic calculations using a computer algebra software (e.g., SymPy [27] with the module $\mathcal{G \mathcal { A }}$ lgebra [1]), all implicit plane curve functions ${ }^{20} F(x, y)=0$ can be produced by the inner product $F(x, y)=\mathbf{T}_{\mathcal{T}} \cdot \boldsymbol{\Omega}=0$, where $\mathbf{T}_{\mathcal{T}}$ is the symbolic TCGA test point $\mathbf{T}_{\mathcal{T}}=\mathcal{T}\left(x \mathbf{e}_{1}+y \mathbf{e}_{2}\right)$ with $x$ and $y$ as symbols, and $\boldsymbol{\Omega}$ is any of the TCGA IPNS 3-vector entities of Section 3.7. An implicit plane curve equation $F(x, y)=0$ can be graphed or plotted using commonly available graphing, plotting, or scientific visualization software (e.g., MayaVi [26] using $x, y, z=\operatorname{mgrid}[-10: 10: 2000 j,-10: 10: 2000 j, 0: 1]$ for nearly 2-D graphing around the origin; MayaVi is primarily a 3-D visualization software).

[^12]For example, with SymPy [27] using the $\mathcal{G \mathcal { A }}$ lgebra [1] module, the sextic implicit plane curve equation for the cranioid $\mathfrak{C}_{t}$ of Section 3.7.4 is generated as

$$
\begin{equation*}
F(x, y)=\mathbf{T}_{\mathcal{J}} \cdot \mathfrak{G}_{t}=0 \tag{62}
\end{equation*}
$$

where the symbolic test point $\mathbf{T}_{\mathcal{T}}$ is the TCGA point embedding $\mathcal{T}\left(\mathbf{t}_{\mathcal{E}}\right)$ of $\mathbf{t}_{\mathcal{E}^{1}}=x \mathbf{e}_{1}+y \mathbf{e}_{2}$ as

$$
\begin{equation*}
\mathbf{T}_{\mathcal{T}}=\mathcal{T}\left(\mathbf{t}_{\mathcal{E}}\right)=\mathcal{C}^{1}\left(\mathbf{t}_{\mathcal{E}^{1}}\right) \mathcal{C}^{2}\left(\mathbf{t}_{\mathcal{E}^{2}}\right) \mathcal{C}^{3}\left(\mathbf{t}_{\mathcal{E}^{3}}\right) \tag{63}
\end{equation*}
$$

Figure 3 was produced by graphing $F(x, y)=0$ using the scientific visualization software MayaVi [26], and it was annotated with mathematical text using $\mathrm{T}_{\mathrm{E}} \mathrm{X}_{\mathrm{MACS}}$ [28].

Versors of high grade, operating on entities of high grade, incur high computational costs in terms of the total number of arithmetic operations involved. In general, a $k$-versor $V_{k}$ can be factored into a product of $k$ vectors as $V_{k}=\mathbf{v}_{1} \mathbf{v}_{2} \ldots \mathbf{v}_{k}$, and then the versor operation can be implemented as a succession of vector reflection operations in the $\mathbf{v}_{i}$, which may lower the computational costs. Intelligent usage of the associativity of the geometric product can reduce the dimensions and grades of products more quickly than naïve usage, and does generally improve the efficiency of computations with the geometric product, including versor operations. In our experience it has proved optimal to compute $k$-CGA $\{k, 2 k\}$-versor operations $A_{\mathcal{K}}^{\prime}=V_{\mathcal{K}} A_{\mathcal{K}} V_{\mathcal{K}}^{-1}$ as $k$ successions of the CGA1 $\ldots k\{1,2\}$-versor operations $A_{\mathcal{K}}^{\prime}=V_{C^{1}}\left(V_{\mathcal{C}^{2}} \ldots\left(V_{\mathcal{C}^{k}} A_{\mathcal{K}} V_{\mathcal{C}^{k}}^{-1}\right) \ldots V_{C^{2}}^{-1}\right) V_{\mathcal{C}^{1}}^{-1}$ using the associativity of the geometric product (see Appendix A).

For example, the TCGA 6-versor operations (combinations of translations, dilations and rotations) were tested as successions of the three CGAi 2 -versor operations for each operation, otherwise the symbolic test computations were very slow in SymPy with the $\mathcal{G} \mathcal{A}$ lgebra module (compare e.g., Fig. 2). When implementing and testing TCGA in SymPy [27] using the $\mathcal{G} \mathcal{A l g e b r a}$ [1] module, it is easy to create functions that return entities and versors according to parameters and to implement the successions of versor operations on the entities to transform them.

Furthermore, the anisotropic dilation operation of Section 3.8 has been successfully tested using SymPy [27] with the $\mathcal{G A} \mathcal{A}$ lgebra [1] module, and it is indeed practicable for some calculations if the three 2-versors $B_{\mathcal{M}^{i}}^{\mathbf{p}}=T_{\mathcal{C}^{i}} B_{\mathcal{M}^{i}}\left(T_{\mathcal{C}^{i}}\right)^{\sim}$ operate on a general cubic entity $\mathbf{X}_{\mathcal{T}}$ in succession instead of directly using the 6-versor operation $B_{\mathcal{T}}^{\mathbf{p}} \mathbf{X}_{\mathcal{T}}\left(\boldsymbol{B}_{\mathcal{T}}^{\mathbf{p}}\right)^{\sim}$ (see Appendix A).

## 5 | CONCLUSION

The Triple Conformal Geometric Algebra (TCGA or 3-CGA) $\mathcal{C}_{3(2+1), 3}$ for the Euclidean plane $\mathbb{R}^{2}$ has IPNS 3-blade entities representing points, point pairs, lines, and circles and IPNS 3-vector entities representing general cubic plane curves and their inversions in circles (and generalizing to reflections in lines). TCGA also has IPNS 3-vector entities for representing certain types of cyclidic (or roulette) quartic, quintic, and sextic plane curves. The circle and line entities are also the operators for inversion and reflection in them, which are implemented as vector reflections. Conformal operations (as versor operations with exponential forms) for dilation, translation, and rotation of any TCGA entity can be defined in terms of inversions in circles, reflections in parallel lines, and reflections in non-parallel lines, respectively. As explained at the end of Section 3.4, for any TCGA plane curve entity representing a general cubic or certain types of quartic, quintic, or sextic curve, it is possible to form an entity representing its intersection with a line or circle. TCGA includes commutator-based differential operators applicable to any TCGA IPNS 3-blade or 3-vector entity and that provide a standard method of differentiation (or differential calculus).

As already indicated in the introduction, further extensions generalizing TCGA to $k$-CGAs, will allow for the definition of IPNS $k$-vector entities representing geometric surfaces of general degree $k$ (and also Darboux $k$-cyclidic entities for surfaces of degrees $(k+1) \ldots 2 k)$ and their intersections with $k$-blade entities ( $k$-CGA hyperplanes and hyperspheres), and algebraic differential operators of the form $D_{x}=k T_{x^{k-1}} T_{x^{k}}^{-1}$ (as products of certain extraction operators $T_{s}$ ) using the commutator product $\partial_{x} \boldsymbol{\Omega}=D_{x} \times \boldsymbol{\Omega}$ on any $k$-vector entity $\boldsymbol{\Omega}$ [4]. However, as the dimension $n=p+q$ of the geometric vector space $\mathbb{R}^{p, q}$ and the multiplicity $k$ of the $k$-CGA increase, the dimension $k(n+2)$ of the $k$-CGA increases rapidly and the number of canonical algebraic basis blades $2^{k(n+2)}$ [25] increases exponentially, which may lead to impracticable computational costs unless the implementation is very efficient [11] and optimized [15][17]. Computational costs can also be reduced in the computations of versor operations by intelligent usage of the associativity of the geometric product and of the Cartan-Dieudonné theorem on orthogonal transformations to compute slow $\{k, 2 k\}$-versor operations as $k$ successions of much faster $\{1,2\}$-versor operations (see Appendix A).

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## APPENDIX A: ON EFFICIENT USAGE OF ASSOCIATIVITY IN VERSOR OPERATION COMPUTATIONS

An intelligent (efficient) usage of the associativity of the geometric product is to associate groups of factors in a chain of geometric products into parentheses, changing the order of multiplications, to reduce or minimize the total number of algebra operations or computational costs (in runtime) that are required as compared to the naïve (inefficient) usage of the associativity of the geometric product, which is the default (usually left to right) order of multiplication of the factors. Different choices of associative groupings in parentheses to order the products may be more or less efficient (or expensive) than each other according to computational costs
(usually measured by runtime). To determine the most intelligent (efficient) usage requires gathering intelligence in the form of estimates of computational costs for alternative usages and then intelligently choosing the usage that minimizes computational costs.

As an example of a naïve usage of the associativity of the geometric product, consider the versor "sandwich" operation $\boldsymbol{\Omega}^{\prime}=$ $\mathbf{C}_{\mathcal{T}} \boldsymbol{\Omega} \mathbf{C}_{\mathcal{J}}^{\sim}=\mathbf{C}_{\mathcal{C}^{1}} \mathbf{C}_{C^{2}} \mathbf{C}_{\mathcal{C}^{3}} \boldsymbol{\Omega} \mathbf{C}_{\mathcal{C}^{3}} \mathbf{C}_{C^{2}} \mathbf{C}_{\mathcal{C}^{1}}$ for the inversion of a TCGA IPNS 3-vector entity $\boldsymbol{\Omega}$ in a TCGA IPNS 3-blade standard circle (3-versor inversor) $\mathbf{C}_{\mathcal{T}}=\mathbf{C}_{\mathcal{C}^{1}} \mathbf{C}_{C^{2}} \mathbf{C}_{\mathcal{C}^{3}}$. The TCGA circle $\mathbf{C}_{\mathcal{T}}$ is factored into the product of the three corresponding CGA $i$ circles $\mathbf{C}_{C^{i}}$. Each CGAi IPNS 1-blade circle (1-versor inversor) $\mathbf{C}_{C^{i}}$ is a vector with up to four component terms. Without any parentheses to perform the three vector reflections in succession, the six geometric products in $\boldsymbol{\Omega}^{\prime}$ are evaluated in the default (left to right) order. Assuming that like terms are added during each of the six multiplications, this naïve usage requires up to $4^{2}+4^{3}+4^{3} \times 64+4 \times \sum_{i=0,2,4,6}\binom{12}{i}+4 \times \sum_{i=1,3,5,7}\binom{12}{i}+4 \times \sum_{i=0,2,4,6,8}\binom{12}{i}=16+64+4096+4 \times 1486+4 \times 1816+4 \times 1981=$ 25308 basis blade multiplications to compute $\boldsymbol{\Omega}^{\prime}$.

A more intelligent usage of the associativity of the geometric product is $\boldsymbol{\Omega}^{\prime}=\mathbf{C}_{C^{1}}\left(\mathbf{C}_{C^{2}}\left(\mathbf{C}_{\mathcal{C}^{3}} \boldsymbol{\Omega} \mathbf{C}_{C^{3}}\right) \mathbf{C}_{\mathcal{C}^{2}}\right) \mathbf{C}_{\mathcal{C}^{1}}$, with parentheses to associate each successive vector reflection as three nested vector reflections. This requires up to $3 \times 64 \times\left(4+4^{2}\right)=3840$ basis blade multiplications if each successive vector reflection in parentheses transforms up to 64 terms in $\boldsymbol{\Omega}$ into another 64 transformed terms (this depends on adding like terms during each successive vector reflection).

Another choice (usage) of associativity is $\boldsymbol{\Omega}^{\prime}=\left(\mathbf{C}_{C^{1}} \mathbf{C}_{C^{2}} \mathbf{C}_{\mathcal{C}^{3}}\right) \boldsymbol{\Omega}\left(\mathbf{C}_{C^{3}} \mathbf{C}_{C^{2}} \mathbf{C}_{\mathcal{C}^{1}}\right)=\mathbf{C}_{\mathcal{T}} \boldsymbol{\Omega} \mathbf{C}_{\mathcal{T}}^{\sim}$. With this choice, $\boldsymbol{\Omega}^{\prime}$ requires up to $2 \times 4^{3}+4^{3} \times 64+4^{3} \times \sum_{i=0,2,4,6}\binom{12}{i}=128+4096+95104=99328$ basis blade multiplications, which is more naïve (less efficient or intelligent) than the default (left to right) associativity. After the multiplications, like terms could be added, resulting in up to 64 transformed 3-blade terms. It is possible that this choice would be the first choice used by a novice, but it is, perhaps surprisingly, the most naïve (inefficient). There are still other possible choices (usages) of the associativity in computing $\boldsymbol{\Omega}^{\prime}$, which may each have different computational costs.

For the example choices discussed above, the successive (nested) vector reflections appear to be most efficient, requiring the least basis blade multiplications. In a parallel computing environment, a nested expression may not be the most efficient or intelligent choice for minimizing runtime. In applications that use precomputed results, the final results are the same for all choices of associativity after all like terms are added; however, the computational costs to precompute the results may depend on the choice of associativity during the precomputing, leading to faster or slower precomputation.

## AUTHOR BIOGRAPHY



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R.B. Easter prepared this revised TCGA preprint paper with corrections and improvements to the published version.


Eckhard Hitzer. Second author. E. Hitzer jointly prepared the revisions for publication with R.B. Easter, including corresponding with the publishing editors and responding to the peer reviewers. E. Hitzer also presented the TCGA paper (by preparing and presenting a 38 -slide PDF presentation based on the original pre-submission TCGA manuscript and other referenced works) at the Empowering Novel Geometric Algebra in Graphics and Engineering (ENGAGE) Workshop (Session: ENGAGE III, June 27 14:00-15:30) at the Computer Graphics International 2017 (CGI2017) conference (June 27-30, 2017 / Yokohama, Japan).

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    ${ }^{2}$ We use Euclidean for vector spaces with $n=p, q=0$, and pseudo-Euclidean for vector spaces with $p<n, q>0$.

[^1]:    ${ }^{3}$ For the notation used in this paragraph, please refer to Section 2.
    ${ }^{4}$ In the framework of $k$-CGA, standard (single) CGA could be labeled 1-CGA. To avoid cluttered notation we omit this. The current section briefly reviews standard CGA of the Euclidean plane.

[^2]:    ${ }^{5}$ In Clifford algebra, the notation $\mathcal{G}_{p, q}^{k}$ indicates the subset of $k$-vectors, i.e. all elements of grade $k, 0 \leq k \leq(n=p+q)$, in $\mathcal{G}_{p, q}$.
    ${ }^{6}$ In standard CGA of the Euclidean plane $\mathcal{G}_{3,1}$, usually the notation for origin null vector $\mathbf{e}_{o}$, and for infinity null vector $\mathbf{e}_{\infty}$, respectively, are used. But to prepare for the TCGA notation, and to avoid confusion, we already give the notation $\mathbf{e}_{o 1}$ and $\mathbf{e}_{\infty 1}$ as well.
    ${ }^{7}$ Note that in Clifford analysis [2] the $m$-dimensional Euclidean space is usually defined by $\mathbb{R}^{0, m}$, and the norm of a vector is then defined as $|\mathbf{p}|=\sqrt{-\mathbf{p}^{2}}$.

[^3]:    ${ }^{8}$ Note, that the two point pairs $\mathbf{2}_{C}$ and its dual $\mathbf{2}_{c}^{*}$ are orthogonal to each other. We thank the anonymous reviewers for pointing out this clarification.
    ${ }^{9}$ In three or more dimensions $\mathbb{R}^{n \geq 3}$, not considered in this paper, geometric entities more generally represent surfaces or hypersurfaces and some of their intersections.
    ${ }^{10}$ Sphere in $\mathbb{R}^{3}$.

[^4]:    ${ }^{11}$ Plane in $\mathbb{R}^{3}$.

[^5]:    ${ }^{12}$ A surface in CGA is described by an outer products of points on the surface in OPNS, or the dual of this outer product in IPNS. A surface in $k$-CGA is described by a $k$-vector in IPNS. We call a conformally embedded general point test point, its Euclidean position is the corresponding test vector. In IPNS, setting to zero the inner product of a test point with a multivector representing a surface, produces an equation for all points on the surface.

[^6]:    ${ }^{13}$ Reciprocals may also be called pseudoinverses (see [25] and [10]).

[^7]:    ${ }^{14}$ At the moment we see no possibility to model higher degree cubic, quadric, quintic and sextic curves and surfaces by blades in $k$-CGA, and preserve the expressions of geometric transformations by means of versors at the same time. See the discussion in [10] and on conics in [25].

[^8]:    $D_{x}$ and $D_{y}$ produce the correct and exact derivative extraction operators $D_{x} \times T_{s}$ and $D_{y} \times T_{s}$ as shown in Tables 4,5 , and 6 .

[^9]:    ${ }^{15}$ The third power arises, because we use three copies of CGA to produce TCGA.

[^10]:    ${ }^{16}$ Please note, that this increase in dimension may lead to a substantial increase in computation time, when compared to TCGA for the Euclidean plane in $\mathcal{G}_{9,3}$.
    ${ }^{17}$ Note, that the indexes 13,14 , and 15 do not indicate bivectors, but rather enumerate vectors.

[^11]:    ${ }^{18}$ The notion of less practicable only refers to the higher computational cost due to the increase in dimensions of the Clifford algebra thus generated.
    ${ }^{19}$ The term boost is taken over from special relativity, but the hyperbolic rotor we use currently only serves as anisotropic dilation operator for the purpose of shape manipulation.

[^12]:    ${ }^{20}$ Note, that for an implicit plane curve equation $F(x, y)=0$, a point $(x, y)$ is on the plane curve when $F(x, y)=0$. In the neighborhood of ( $x, y$ ), the function $F(x, y)=\mathbf{T}_{\mathcal{T}} \cdot \boldsymbol{\Omega}$ is positive on one side of the plane curve and negative on the other side of the plane curve. This can also be used to judge if a point is inside or outside a closed curve, like in Figure 2.

