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2	REACTION CENTRE
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ULTRAFAST STRUCTURAL CHANGES WITHIN A PHOTOSYNTHETIC

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Photosynthetic reaction centres harvest the energy content of sunlight by transporting electrons across an energy transducing biological membrane. We use time-resolved serial femtosecond crystallography¹ at an X-ray free electron laser² to observe light-induced structural changes in the photosynthetic reaction centre of *Blastochloris viridis* on a time-scale of picoseconds. Structural perturbations are first centred upon the protein's special pair of chlorophyll molecules that are photo-oxidized by light. Electron transfer to the menaquinone acceptor on the opposite side of the membrane induces a movement of this cofactor in concert with lower amplitude protein rearrangements. These observations reveal how proteins utilize conformational dynamics to stabilize the charge separation steps of electron transfer reactions.

Our biosphere depends upon the electron transfer reactions of photosynthesis as a primary source of energy. Photosystems and photosynthetic reaction centres form a family of integral membrane protein complexes found in plants, algae, cyanobacteria and photosynthetic bacteria that convert the energy of a captured photon into a charge separated state. The photosynthetic reaction centre of the purple non-sulphur bacterium *Blastochloris viridis* (RC_{vir}) contains three transmembrane subunits called H, L and M and a periplasmic subunit C. These subunits support four bacteriochlorophyll molecules (BCh), two bacteriopheophytin molecules (BPh), a tightly bound menaquinone (Q_A), a mobile ubiquinone (Q_B), a single non-haem iron and four haem co-factors (**Fig. 1**). Electron transfer reactions originate at a special pair (SP) of strongly interacting bacteriochlorophylls which in *Bl. viridis* have an absorption maximum at 960 nm. Photo-oxidation of the SP liberates an electron which is transferred to the active branch BPh_L within a few picoseconds, is transferred to the tightly bound menaquinone (Q_A) in less than a nanosecond, and is transferred to the mobile ubiquinone (Q_B) in microseconds. SP⁺ is reduced by subunit C and a second photo-oxidation event transfers a

second electron to Q_B^- , which is protonated from the cytoplasm and released into the membrane as ubiquinol (H_2Q). Other proteins participate in a cyclic flow that returns electrons to subunit C and the net effect is that two protons are transported across an energy-transducing membrane for every photon absorbed.

Electrons may tunnel between cofactors when they are separated by approximately 10 Å or less.³ The primary electron transfer step from SP to BPh_L occurs in 2.8 ± 0.2 ps⁴ over a distance of 10 Å by means of a two-step hopping mechanism *via* the monomeric BCh_L⁵ and is more rapid than conventional Marcus theory. By contrast, the 9 Å electron transfer step from BPh_L to Q_A has a single exponential decay time of 230 ± 30 ps⁶ which is consistent with conventional Marcus theory. Coherent nuclear motions⁷ and protein structural changes⁸ have been suggested to influence the initial charge-transfer reactions of photosynthesis yet the specific nature of these putative protein motions is unknown. Flash-freeze crystallographic trapping studies,⁹ time-resolved Laue diffraction¹⁰ and time-resolved serial femtosecond crystallography¹¹⁻¹⁴ (TR-SFX) have revealed structural changes in bacterial photosynthetic reaction centres^{9,10} and cyanobacterial photosystem II¹¹⁻¹⁴ that occur on the late microsecond to millisecond time-scale, yet no time-resolved crystallographic studies on the time-scale of the primary charge separation reactions of photosynthesis have been reported.

Here we apply time-resolved serial femtosecond crystallography¹ at an X-ray free electron laser (XFEL) to probe the ultrafast structural response of RC_{vir} to light. We photo-excited the special pair with 150 fs pulses centred at 960 nm (Extended Data Fig. 1). X-ray pulses 40 fs in duration were generated at the LCLS² and were used to record diffraction patterns from tens of thousands of microcrystals for the time points $\Delta t = 1$ ps, 5 ps (two repeats), 20 ps, 300 ps (two repeats) and 8 µs after photoexcitation (Extended Data Table 1). The time point $\Delta t = 1$ ps

populates the photo-excited charge transfer state of the SP in which charge rearrangements have occurred within the bacteriochlorophyll dimer but are prior to the primary electron transfer step; $\Delta t = 5$ and 20 ps are after the initial charge-transfer step and SP is oxidized and BPh_L is reduced; $\Delta t = 300$ ps is longer than the time constant for electron transfer to Q_A and menaquinone is reduced; and $\Delta t = 8$ µs corresponds to a meta-stable charge separated state.

Extended Data Fig. 2 presents overviews of the $|F_{obs}|^{light}$ - $|F_{obs}|^{dark}$ isomorphous difference Fourier electron density maps ("light" corresponds to data collected from photo-activated microcrystals whereas "dark" corresponds to data collected from microcrystals that were not photo-activated) for all time points. Difference electron density features are visible above 4.0 σ (σ is the root mean square electron density of the map) near SP for all time points and strong features associated with Q_A are visible for $\Delta t \geq 300$ ps (Extended Data Table 2). In contradistinction with ultrafast TR-SFX studies of bacteriorhodopsin, 15 photoactive yellow protein, 16 rsEGFP17 and bacterial phytochromes18 in which ultrafast structural changes are driven by the movements of atoms due to a photo-isomerization event, TR-SFX measurements of RC $_{vir}$ reveal a knock-on effect on protein structure due to the light-induced redistribution of charge. Electric-field induced conformational changes have been observed when fields of the order 10^8 V/m are applied across a protein-crystal 19 and this is the same order of magnitude as electric field perturbations due to the movement of an electron within RC $_{vir}$.

Recurring changes of electron density are visible as positive difference electron density in the region of overlap between the two bacteriochlorophylls SP_L and SP_M of the special pair, and complementary negative difference electron density features are visible primarily associated with SP_M (Fig. 2, Extended Data Fig. 3, Extended Data Table 2, Supplementary Video 1).

Singular value decomposition (SVD) of all seven difference Fourier electron density maps (Fig. 2e) reveals that the strongest positive and several of the strongest negative difference electron density features of the principal SVD component are associated with the SP (Extended Data Table 2). Quantification of electron density changes²⁰ within the RC_{vir} cofactors (Fig. 2f) and statistical tests against control difference Fourier electron density maps (Methods) provides a very high-level of confidence (p-value ≤ 0.001 , Extended data Table 3) that these recurring difference electron density features do not arise by chance. Thus photoexcitation causes the bacteriochlorophylls of SP to move closer together and the bending (an out of plane distortion) of SP_M could explain these observations. An out of plane distortion was used to model difference electron density features observed as carbon monoxide was photo-dissociated from the haem of myoglobin²¹ (Extended Data Fig. 4). Nonplanar distortions of chlorin and bacteriochlorin rings are observed in photosystem II and RC_{sph}, due to interactions with the surrounding protein²² and nonplanar porphyrins are also more easily oxidized than planar porphyrins.^{23,24} This suggests that the distortion of SP in advance of the primary charge separation event (Fig. 2a) could enhance the yield of the primary charge-transfer reaction, which has been optimized by evolution to achieve almost perfect quantum efficiency.²⁵

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When the C-subunit is fully reduced, an electron is transferred from haem₃ to SP⁺ in less than a microsecond.²⁶ The above reasoning implies that SP⁺ may be more easily reduced should SP_M return to a planar geometry before this electron transfer occurs. This hypothesis is consistent with our experimental observations since the amplitude of the positive difference electron density feature between the SP bacteriochlorophylls increases from $\Delta t = 1$ ps to 20 ps, decreases for $\Delta t = 300$ ps and is insignificant for $\Delta t = 8$ µs (Fig. 2, Extended Data Fig. 3, Extended Data Table 2). Moreover, neither TR-SFX studies of the S2 \rightarrow S3 transition of

cyanobacteria PSII¹³ ($\Delta t = 150 \,\mu s$ and 400 μs) nor TR-Laue diffraction studies of RC_{vir}¹⁰ ($\Delta t = 3 \,m s$) report a positive difference electron density feature in the region of overlap between the special pair of (bacterio)chlorophylls, implying that this feature has decayed.

Charge rearrangements cause SP^+ to move up to 0.3 Å towards the M-subunit by $\Delta t = 300$ ps and the side-chains of both $His173_L$ and $His200_M$ adjust to preserve their ligating interactions with the magnesium ions of SP^+ , as do $His168_L$ and $Tyr195_M$ adjust conformation in order to maintain their H-bond interactions to SP_L and SP_M respectively. These structural perturbations are revealed by paired negative and positive difference electron density features on the side-chain of $His173_L$ in the principal SVD components calculated from both the early (1ps, 2 × 5ps, 20 ps) and late (2 × 300 ps, 8 μ s) sub-sets of TR-SFX data, whereas positive difference electron density features associated with the side-chains of $His200_M$ and $Tyr195_M$ become noticeably stronger for the later sub-set of data (Extended Data Fig. 3i,j; Extended Data Table 2; Supplementary Video 1). These observations suggest that SP_L moves towards subunit M slightly in advance of SP_M , which may be due to dielectric asymmetry within photosynthetic reaction centres. Dielectric asymmetry is believed to underpin the phenomenon that electron transfer occurs only along the A-branch²⁷ in purple bacteria RCs and PSII (Fig. 1).

An electron moves from SP to BPh_L in 2.8 ± 0.2 ps⁴ and from BPh_L to Q_A in 230 ± 30 ps.⁶ The tightly bound menaquinone is therefore neutral for $\Delta t = 1$ ps, 5 ps and 20 ps; three-quarters of the photo-activated population are reduced to semiquinone by $\Delta t = 300$ ps; and essentially all photo-activated molecules have Q_A reduced at $\Delta t = 8$ µs. Our difference Fourier electron density maps confirm these expectations since the few difference electron density features visible within the Q_A binding pocket for $\Delta t \leq 20$ ps are isolated whereas more continuous paired positive and negative difference electron density features are visible for Δt

≥ 300 ps (Fig. 3, Extended Data Fig. 5). These recurring features of the later sub-set of TR-SFX data (2 × 300 ps, 8 µs) produce strong difference electron density features in the principal SVD component that are associated with QA and its H-bond interaction with His217_M (Extended Data Table 2, Fig. 3d, Supplementary Video 2) and statistical tests establish that these recurring changes cannot be ascribed to noise (p-value ≤ 0.0125 , Extended Data Table 3). Structural refinement models these observations as due to a twist and translation of the semiquinone that brings the negatively charged head-group approximately 0.2 Å closer to the positive charge of the non-haem Fe²⁺ (Fig. 3f) and thereby stabilizes the reduced form of this cofactor. This interpretation receives support from QM/MM calculations that predict that the Q_A to His217_M H-bond is shortened by 0.17 Å when Q_A is reduced (Extended Data Fig. 6f) and suggest that semiquinone binding is stabilized by approximately 36 kJ.mol⁻¹ due to structural changes (Extended Data Fig. 6g,h, Methods), which is a sizeable fraction of the energy (125 kJ.mol⁻¹) of a 960 nm photon. Similar conclusions were drawn from an earlier analysis using a Density Functional Theory (DFT) formalism.²⁹ Light-induced electron density changes were visible for Q_A in TR-SFX studies of the S2→S3 transition of cyanobacteria PSII¹³ for the time points 150 µs and 400 µs, light-induced movements of the mobile quinone Q_B were also observed in $PSII^{11-14}$ for delays of hundreds of milliseconds, and larger light-induced motions of Q_B were reported in freeze-trapping studies of the Rhodobacter sphaeroides photosynthetic reaction centre.⁹

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For $\Delta t = 300$ ps, paired negative and positive difference electron density features are associated with the cytoplasmic portions of TM helices D_M and E_M (Fig. 3e) and indicate that RC_{vir} adjusts its structure in response to the movement of the semiquinone within the Q_A binding pocket (Fig. 3f). A more quantitative analysis (Methods, Extended Data Fig. 7, Supplementary Video 3) suggests that low-amplitude protein motions begin to arise already

by $\Delta t = 1$ ps (Fig. 4a) as observed in TR-SFX studies of bR^{15,20} and Mb:CO²¹; the amplitude of these motions increase with time and by $\Delta t = 5$ ps larger displacements are observed near the SP⁺ and BPh_L⁻ cofactors (Fig. 4b); and for $\Delta t = 300$ ps protein conformational changes extend throughout the A-branch of the electron transfer pathway from SP⁺ to Q_A⁻ (Fig. 4c). When the same representation is used to depict protein conformational changes predicted from QM/MM calculations (Supplementary Video 4) almost no structural changes are expected for the photo-excited charge transfer state (Fig. 4d); protein movements arise near the charged cofactors in the SP⁺:BPh_L⁻ charge-separated state (Fig. 4e); and structural changes extend throughout the A-branch in the SP⁺:Q_A⁻ charge-separated state (Fig. 4f). These findings demonstrate that RC_{vir} is not a passive scaffold but rather low-amplitude protein motions engage in a choreographed dance with electron movements taking the lead and protein conformational changes following. Conversely, as the protein's structure adjusts to stabilize these charge rearrangements the energetic barriers hindering the reverse electron transfer reaction increase, thereby extending the lifetime of the change separated species and enhancing the overall efficiency of photosynthesis.

In Marcus theory the total potential energy of an electron donor and its surroundings must be equal to that of the electron acceptor and its surroundings if an electron is to tunnel from donor to acceptor.³ Fluctuations in the organizational energy around protein cofactors are therefore essential to facilitate electron transfer reactions. Efforts aimed at understanding how protein conformational dynamics control the rates of electron transfer between cofactors^{8,30} have been hampered by a lack of experimental tools that characterize protein structural changes on the relevant time-scales. Our observations provide an experimental framework for extending the standard description of electron transfer reactions in photosynthesis³ to explicitly incorporate protein structural changes. Electron transfer reactions are ubiquitous in

- 217 nature and therefore a more nuanced understanding of the interplay between protein structural
- dynamics and the movement of electrons has far-reaching biochemical significance.

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Extended Data is available for this manuscript.

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Supplementary Information is available for this manuscript.

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Author Information

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Contributions

- 248 R.N. conceived the experiment, which was designed with input from G.B., R.D., J.D., S.B.,
- 249 M.L., S.C., D.M., M.S.H., G.J.W., J.D., D.P.D. and A.B. Samples were prepared by R.D.,
- 250 P.Bå., D.A. and R.B. Crystallization was supervised by G.B.. Time-resolved serial
- femtosecond crystallography experiments were performed at the LCLS by R.D., P.Bå., D.A.,
- 252 R.B. K.R.B., G.N., M.L., D.M., J.R., R.H., P.Be., E.M., L.J., R.A., S.C., E.C., C.E.C., P.D.,
- 253 G.H., M.S.H., C.L., S.L. C.S., A.S., G.J.W., C.W., J.D., D.P.D., A.B., G.B. and R.N. The CXI
- instrument was set up and run by M.L., M.S.H., G.J.W. and S.B.. The photoexcitation laser of
- 255 960 nm was operated and aligned by D.M., J.R., S.C. and J.D.. Sample delivery was
- performed by K.R.B., G.N., R.H., P.Be., P.D. and D.P.D. Data were analysed by R.D., P.Bå.,
- A.R., O.Y., A.B., G.B. and R.N. Structural refinement was performed by R.D. and P.Bå.
- Resampling, full occupancy structural refinement and analysis were performed by V.A.G,
- 259 G.K. and A.V. Integration within a sphere and statistical tests were performed by C.W., R.N.
- and P.Bå.. SVD analysis was performed by A.V. Quantum Mechanics/Molecular Mechanics
- analysis were performed by D.M., H.L.L and G.G.. Time-resolved IR spectroscopy
- measurements were performed by J.K, M.M and S.W. The manuscript was prepared by R.N.,
- P.Bå., R.D. and G.B. with additional input from all authors.

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Code Availability

- Software used for SVD analysis is available at https://github.com/Neutze-lab/SVD. Code
- 270 written in MATLAB to analyse difference electron density amplitudes is available at

https://github.com/Neutze-lab/maptool. Software associated with the resampling of X-ray diffraction data is available at https://github.com/Katona-lab/CFEL tools. **Data Availability** Atomic coordinates and structure factors have been deposited in the Protein Data Bank, www.pdb.org. PDB ID codes are the following: 5O4C the dark-conformation (a); 5NJ4 the dark conformation (b); 6ZHW the time-point $\Delta t = 1$ ps; 6ZI4 the time-point $\Delta t = 5$ ps (a); 6ZID the time-point $\Delta t = 5$ ps (b); 6ZI6 the time-point $\Delta t = 20$ ps; 6ZI5 the time-point $\Delta t = 10$ 300 ps (a); 6ZI9 the time-point $\Delta t = 300$ ps (b); 6ZIA the time-point $\Delta t = 8$ µs. Difference Fourier electron density maps and stream files containing X-ray diffraction intensities are deposited at the CXI database (http://www.cxidb.org/) with identification number ID 161. **Competing Interests** The authors declare no competing interests.

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FIGURE LEGENDS

Figure 1: Electron transfer steps of the photosynthetic reaction centre of Bl. viridis.

Cartoon representation of the H, L, M and C subunits. Cofactors are shown in black including the special pair of bacteriochlorophylls (SP), two monomeric bacteriochlorophylls (BCh), two bacteriopheophytins (BPh), a tightly bound menaquinone (Q_A), a mobile ubiquinone (Q_B), a non-haem iron (Fe^{2+}) and four haems. The approximate boundaries of the membrane are suggested in blue. The electron transfer pathway: $SP \rightarrow BPh_L \rightarrow Q_A$ is referred to as the A-branch. Approximate timescales for the first two electron transfer events, from SP to BPh_L and from BPh_L to Q_A , are depicted.

Figure 2: Light-induced electron density changes in RC_{vir} at the site of photo-oxidation.

a, Experimental $F_{obs}^{light} - F_{obs}^{dark}$ difference Fourier electron density map for $\Delta t = 1$ ps. b, Difference Fourier electron density map for $\Delta t = 5$ ps (data set a). c, Difference Fourier electron density map for $\Delta t = 20$ ps. d, Difference Fourier electron density map for $\Delta t = 300$ ps (data set a). e, Principal component from SVD analysis of all seven experimental difference Fourier electron density maps. All maps are contoured at $\pm 3.2 \, \sigma$ (blue, positive difference electron density; gold, negative difference electron density; σ is the root mean square electron density of the map). f, Relative amplitudes of difference electron density features integrated within a 4.5 Å sphere²⁰ centred upon the RC_{vir} co-factors (Extended Data Fig. 3j). The colour bars represent: cyan, $\Delta t = 1$ ps; blue, $\Delta t = 5$ ps a and b (in that order); purple, $\Delta t = 20$ ps; red, $\Delta t = 300$ ps b and a (in that order); mustard, $\Delta t = 8$ µs.

Figure 3: Light-induced electron density changes in RC_{vir} within the menaquinone binding pocket. a, Experimental $F_{obs}^{light} - F_{obs}^{dark}$ difference Fourier electron density map for $\Delta t = 5$ ps (data set a). b, Difference Fourier electron density map for $\Delta t = 300$ ps (data set a). c, Principal component from SVD analysis of the first four experimental difference Fourier electron density maps ($\Delta t = 1$ ps, 5 ps (a), 5 ps (b), 20 ps). d, Principal component from SVD analysis of the final three experimental difference Fourier electron density maps ($\Delta t = 300$ ps (a), 300 ps (b), 8 µs). All maps are contoured at ± 3.0 σ (blue, positive difference electron density; gold, negative difference electron density). e, Difference Fourier electron density map for $\Delta t = 300$ ps (data set a) showing the protein immediately surrounding Q_A and contoured at ± 3.5 σ . f, Superposition of the refined structures for the dark structure (yellow, Q_A in black) and $\Delta t = 300$ ps (purple structure).

Figure 4: Structural response of RC_{vir} to electron transfer events. a, Recurring movements of Cα atoms for $\Delta t = 1$ ps quantified by full occupancy structural refinement against 100 randomly resampled TR-SFX data sets. b, Recurring movements of Cα atoms for $\Delta t = 5$ ps (a) using the same representation. c, Recurring movements of Cα atoms for $\Delta t = 300$ ps (a) using the same representation. Recurring movements are represented as error weighted mean ratios relative to 100 control structural refinements (Methods) coloured from grey (< 80 % of the maximum error weighted mean ratio) to red (≥ 95 % of the maximum error weighted mean ratio). An identical representation is given for all time points in Extended Data Fig. 7. d, Movements of Cα atoms estimated from QM/MM energy minimization calculations associated with the SP photo-excited and all other cofactors in resting state: SP*:BPh_L⁰:Q_A⁰ (Methods). e, Movements of Cα atoms estimated from QM/MM energy minimization calculations associated with the SP photo-oxidized and BPh_L reduced: SP⁺¹:BPh_L⁻¹:Q_A⁰. f,

Movements of $C\alpha$ atoms estimated from QM/MM energy minimization calculations associated with SP photo-oxidized and Q_A reduced: SP^{+1} : BPh_L^0 : Q_A^{-1} . Movements are coloured from grey (no movements) to red (maximum $C\alpha$ motions). Transmembrane helices are drawn as rods.

METHODS

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419 Protein production and purification

The expression and purification of photosynthetic reaction centre from Bl. viridis cells was adapted from Wöhri et al.³¹. Cells were disrupted by three rounds of sonication followed by centrifugation in a JA20 rotor at 15000 rpm for 20 minutes to recover the membrane suspensions. Membranes were then purified by ultracentrifugation at 45000 rpm for 45 min in a Ti45 rotor. Membranes were homogenized in 20 mM Tris-HCl, pH 8.5 and diluted to $OD_{1012} = 10$. Membranes were then solubilized in 4 % lauryldimentylamine-N-oxide (LDAO) for 3 hours at room temperature. Unsolubilized membranes were removed by ultracentrifugation at 45000 rpm for 75 min in a Ti70 rotor. RCvir protein was purified by loading the supernatant onto a 250 ml POROS 50-µm HQ ion-exchange medium equilibrated with wash buffer (20 mM Tris-HCl, pH 8.5, 1% LDAO). The column was washed with 21 of wash buffer with 5 % elution buffer (20 mM Tris-HCL, pH 8.5, 1 M NaCl, 1 % LDAO) and eluted with an increasing concentration of elution buffer over 20 column volumes. Fractions with an $A_{280}/A_{830} < 3.5$ were pooled and concentrated in 100 kDa MW cut off concentration tubes (Vivaspin) to a volume of 10 ml. This was loaded in 5ml batches onto a HiPrep 26/60 Sephacryl S-300 column (GE) equilibrated with SE buffer (20 mM Tris-HCl, pH 8.5, 100 mM NaCl, 0.1 % LDAO) and eluted into 1.8 ml fractions. Fractions with an $A_{280}/A_{830} < 2.6$ were pooled and concentrated, followed by a 20-fold dilution in final protein buffer (20mM NaH_2PO_4/Na_2HPO_4 , pH 6.8, 0.1 % LDAO, 10 μM EDTA) and then concentrated again to 20 mg.ml⁻¹. Samples were flash-frozen in liquid nitrogen and stored at -80 °C.

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Protein crystallization

20 μl sitting drops were set up with a 1:1 ratio of protein solution (10 mg.ml⁻¹) and precipitant solution (3.6 M ammonium sulphate, 6 % heptane-1,2,3-triol, 20 mM NaH₂PO₄/Na₂HPO₄, pH

6.8) set up against a 1 ml reservoir of 2 M ammonium sulphate. Large crystals grew at 4 °C in 3 days. Crystals were harvested by pipette and crushed mechanically to create a seed stock by vortexing with seed beads for approximately 20 min with occasional cooling on ice³². For the XFEL experiment in April 2015 (run a) new 18.5 µl sitting drop vapour diffusion crystallization drops were set up in order to yield large numbers of micro-crystals. In these experiments the protein concentration was 8.5 mg.ml⁻¹ and a protein:precipitant concentration of 10:7.5 was used in the drops. 1 µl of undiluted crystal seed stock was spiked into the drops for a final v/v concentration of 5.4 %. Crystallization drops were then mixed by pipette and covered with a glass cover slide. Rod-like crystals grew over 5 days at 4 $^{\circ}$ C and were 10-20µm in the longest dimension. Microcrystals for the experiment in June 2016 (run b) were prepared as above, but with an additional round of microseeding using crushed microcrystals to seed an additional round of microcrystal growth³². Micro-crystals were harvested by pipette and concentrated up to three-fold by centrifugation at 1000 g for 1 min followed by removal of supernatant. These crystals were somewhat thicker and, while diffracting to higherresolution, they highlighted the compromise inherent in TR-SFX since a lower excited-state occupancy was usually observed when working with crystals of higher optical density.

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Sample injection and data collection

Microcrystals were transferred from Eppendorf tubes to a sample reservoir using a syringe and passing the microcrystal slurries through a stainless steel 20 μ m filter (VICI AG International) or a 20 μ m nylon filter (Sysmex). The reservoir was loaded into a temperature controlled rocking chamber and injected into the XFEL through a GDVN³³ using an internal diameter of 75 μ m. The microjet used a microcrystal suspension flow rate of 20 μ l.min⁻¹ and was focused to a 10 μ m diameter using helium gas. The X-ray beam was aligned to interact

with the liquid jet as close to the tip of the GDVN as practical and before Rayleigh breakup of the microjet.

Diffraction data were collected at 293 K at the CXI beam line³⁴ of the LCLS XFEL during beamtime awarded in April 2015 (run a) and June 2016 (run b). Diffraction data were recorded on a Cornell-SLAC Pixel Array detector³⁵. The X-ray wavelengths and equivalent pulse energies were 1.89 Å (6.56 keV) in 2015 and 1.31 Å (9.49 keV) in 2016. An X-ray pulse duration of 36 fs was used in 2015 and 45 fs in 2016. The XFEL beam was focused to a 3 μ m² spot for both experiments. The detector was located 89 mm from the microjet in 2015 and 145 mm from the microjet in 2016. Diffraction data were collected at a repetition rate of 120 Hz from microcrystals that were not exposed to any optical laser pump (dark-state) and for five time points corresponding to $\Delta t = 1$ ps, 5 ps, 20 ps, 300 ps and 8 μ s after photoexcitation. The time points $\Delta t = 5$ ps and 300 ps were repeated in both 2015 and 2016 and are referred to as data sets a and b respectively.

Laser photoexcitation

An optical Ti:Sa pump laser 150 fs in duration was focused into a spot size of 190 μ m FWHM (323 μ m $^{1}/e^{2}$) and aligned to overlap with the LCLS X-ray pulse. The LCLS timing-tool³⁶ provided a timing accuracy of \pm 200 fs for the time point, Δt , between the arrival of the optical pump laser and the X-ray probe. A pump-laser wavelength of 960 nm was used to photo-excite RC_{vir} microcrystals, and this wavelength is at the absorption maximum of the special pair ($\epsilon_{960} \approx 100~000~M^{-1}.cm^{-1}$). The pump laser energy per pulse was 11.8 μ J in April 2015 and 11.0 μ J in June 2016. For an idealized Gaussian beam, 86.5 % of this light will pass through a spot with diameter $1/e^{2}$ and 50 % of this light will pass through a spot with diameter FWHM. Thus the average fluence within the FWHM spot can be estimated as 25 mJ/cm² and

23 mJ/cm² which equates to a pump-laser power-density of 138 GW/cm² for the 2015 experiment and 129 GW/cm² for the 2016 experiment. This calculation defines the units used throughout to specify the laser power-density. Both values are above 30 GW/cm² to 100 GW/cm² that has been recommended as an upper threshold to avoid nonlinear effects in bacteriorhodopsin. 37,38

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Extreme non-linear absorption was observed as ultrafast sample heating in time-resolved Xray scattering studies of RC_{vir} when pumped with 800 nm light.³⁹ When using 800 nm to photo-excite RC_{vir} it is the BCh cofactors (rather than the SP) which absorb light ($\epsilon_{800} \approx$ 180 000 M⁻¹.cm⁻¹). The pump-laser fluence used in that study³⁹ was 1560 GW/cm². Ultrafast sample heating within a GDVN liquid microjet has also been measured as a function of the 800 nm pump-laser fluence using time-resolved X-ray scattering (Fig. 28 of reference⁴⁰). These measurements show that the energy deposited into RC_{vir} samples is proportional to the pump-laser fluence (a linear response) up to 270 GW/cm² and that the measured heating then varies quadratically (a non-linear response) above a pump-laser fluence of 355 GW/cm². Thus either an idealized assumptions of a perfectly aligned Gaussian beam may not be realistic, and/or large losses occur as the incoming laser pulse is reflected from the surface of a GVDN liquid microjet, and/or a thresholds^{37,38} of 30 GW/cm² to 100 GW/cm² do not apply RC_{vir} when photo-excited at 800 nm. When 960 nm light is used to photo-excite the SP of RC_{vir} it is likely more difficult to induce non-linear effects because the photo-excited state SP* has an absorption maximum red-shifted 70 nm relative to the ground state⁴¹ and hole-burning⁴² has been observed in RCvir such that SP* is effectively transparent to the incoming light. Moreover, the absorbance of RC_{vir} at 960 nm is only 56 % of it absorbance at 800 nm and therefore non-linear effects are likely to arise at higher power densities when using 960 nm rather than 800 nm to photo-excite RC_{vir}. Non-linear ultrafast heating⁴⁰ is observed in RC_{vir}

delivered using a GDVN liquid microjet and photo-excited at 800 nm only above a power-density of 355 GW/cm². Therefore the 960 nm pump-laser power-densities of 138 GW/cm² and 129 GW/cm² used in this work are below where non-linear effects may reasonably be anticipated. These conclusions are supported by time-resolved IR spectroscopy measurements (Extended Data Fig. 1).

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Time-resolved infrared spectroscopy

Time-resolved vibrational spectroscopy measurements were performed with a near infrared (NIR) pump and mid-infrared (IR) probe setup using a regenerative amplifier (Spitfire Ace, Spectra Physics) to deliver pulses centred at 800 nm (1.2 mJ, 5 kHz). The amplifier output is used to pump a TOPAS-TWINS (Light Conversion) capable of generating tuneable femtosecond pulses at two different wavelengths. One path was used to generate mid-infrared probe light centred at 6000 nm via difference frequency generation whereas the other path generated 960 nm pump pulses via second harmonic generation of the idler beam. The 960 nm beam was chopped to 2.5 kHz and delayed in time relative to the probe pulses using an optical delay line. Two weak replicas derived from the midIR beam were used as probe and corresponding reference. Both probe and reference were dispersed in a Horiba spectrograph (grating with 75 gr/mm) and detected and integrated on a double-row MCT array with 64 pixels each on a shot-to-shot basis using a commercial detection system (Infrared Systems). Samples of RC_{vir} were prepared in a customized cell by enclosing ca. 15 uL of solution (RC_{vir} at ca. 0.4 mM in D₂O buffer) between two 2 mm thick CaF₂ windows separated by a 25 µm spacer. Probe and reference beams were focused at the sample position and collimated using 90° off-axis parabolic mirrors. The pump beam was focused using a 30-cm lens and overlapped with the probe beam at its focus. The sample cell was placed where pump and probe beams meet and translated continuously perpendicular to the beam direction during data acquisition. The focal spot size of the pump beam was determined using knife-edge scans and yielded perpendicular 1/e² radii of 57 μm and 56 μm. Different pump fluences were adjusted using reflective neutral density filters (Edmund Optics). For each fluence, twelve repeats over five time points (1000 pump shots per time point and repeat, at delays of -50, 1, 2, 5 and 300 ps) were recorded and less than 5% of shots were rejected during data treatment. Signals were calculated by subtracting consecutive pump-on from pump-off shots followed by application of the noise reduction algorithm.^{43,44} The spectral resolution is < 5 cm⁻¹. The results of these measurements are presented in Extended Data Fig. 1.

Data processing

Images containing more than 20 diffraction spots were identified as diffraction hits by Cheetah. Cheetah converted the raw detector data into the HDF5 format and data were then processed using the software suite CrystFEL version $0.6.2^{46,47}$ Crystals were indexed using a tetragonal unit cell (a = b = 226.4 Å, c = 113.7 Å, $\alpha = \beta = \gamma = 90^{\circ}$). Scaling and merging were performed using Monte Carlo methods using the same software. Data from the dark state and photo-excited states were scaled together using the custom dataset splitting option in the CrystFEL partialator module. Structure factors were calculated from merged intensities by the CCP4 module TRUNCATE and molecular replacement was performed using the CCP4 module Phaser using the ground-state RC_{vir} structures solved with XFEL radiation (PDB codes 5O4C and 5NJ4) as a search models. Statistics for data collection and refinement are detailed in Extended Data Table 1.

Electron density difference maps

Isomorphous $|F_{obs}|^{light}$ - $|F_{obs}|^{dark}$ difference Fourier electron density maps were calculated using the refined dark state structures for phases with the time-points $\Delta \underline{t} = 5$ ps (data set a) and 300

ps (data set a) calculated against data and coordinates using the pdb entry 5O4C whereas the time point with the time-points $\Delta t = 1$ ps, 5 ps (data set b), 20 ps, 300 ps (data set b) and 8 μ s were calculated against data and coordinates using the pdb entry 5NJ4. Thus all difference electron density map calculations used only data collected during the same experiment. Difference Fourier electron density maps represent measured changes in X-ray diffraction intensities as changes in electron density without bias towards the photo-activated state's structural model. The technique is extremely sensitive to small changes in electron density⁵⁰ and reveals more subtle features than are apparent from 2mFobs-DFcalc electron density maps alone (m is the figure of merit and D is estimated from coordinate errors). A Bayesian weighting calculation script⁵¹ using CNS software⁵² was also used to analyse the difference Fourier electron density maps. In this procedure structure factor amplitude differences were weighted by the product of the figure of merit of the ground state structure reflections and of a 14 of reference⁵³), weighting w (Equation which term, was calculated using Bayesian statistics developed to improve signal to noise. 53 For six of seven data-sets the recurring difference electron density features were slightly strengthened by this step. The exception was the time-point $\Delta t = 8$ µs which has difference electron density feature that are weaker than for the other maps (Fig. 2f) and appears due to a lower occupancy of the charge separated state in these microcrystals. It is possible that a fraction of the photo-oxidized SP population is reduced from the C-subunit by $\Delta t = 8 \mu s$, which is longer than the time-scale of this electron transfer step.²⁶ However, no efforts were made to reduce the C-subunit when preparing microcrystals and a similar occupancy (30% \pm 5%) is observed to persist in timeresolved spectroscopy measurements on crystals for up to millisecond delays. 10

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Singular value decomposition

SVD analysis of difference Fourier electron density maps was performed using an in-house code written in python that is based upon an approach previously described.⁵⁴ As has been discussed, 55 SVD may serve as a noise-filter to enhance the signal across a sequence of difference Fourier electron density maps. This step contains the assumption that the overall mechanism is linear and that changes in electron density are similar over the selected timewindows. When applying SVD we evaluate the expression $[U,\Sigma,V] = SVD(A)$, where A is a matrix of n difference Fourier electron density maps containing m elements; U is an n × n unitary matrix; Σ is an n × m rectangular matrix containing n diagonal elements (the singular values) arranged in decreasing order and all other matrix elements are zero; and the first right singular vector (the first row of the matrix V) is referred to as the principal component. Results from SVD analysis of all seven electron density maps are presented in Fig. 2e and Extended Data Fig. 31,m. Results from SVD analysis deriving from the first four time-points $(\Delta t = 1 \text{ ps}, 5 \text{ ps a and b}, 20 \text{ ps})$ and the last three time-points $(\Delta t = 300 \text{ ps a and b}, 8 \text{ µs})$ are shown in Fig. 3c,d; Extended Data Fig. 3h,i; Extended Data Fig. 5h,i; Supplementary Videos 1 and 2. This separation of the maps is motivated by the fact that photo-activated RC_{vir} molecules have menaquinone oxidized for the first sub-set of time-points yet most menaquinone molecules of photo-activated RCvir are reduced for the second sub-set of timepoints.

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Structural refinement of photo-excited states

Isomorphous $|F_{obs}|^{light}$ - $|F_{obs}|^{dark}$ difference Fourier electron density maps were inspected in COOT. Structural refinement was performed using Phenix.⁵⁶ A model was first placed within the unit cell using rigid body refinement followed by multiple rounds of partial-occupancy refinement where the SP, BCh_L, BPh_L, Q_A, portions of TM helices E_L, D_L, E_M and D_M, as well as connecting loops, and additional residues near cofactors (L153-178, L190, L230, L236-

248, M193-221, M232, M243-253, M257-266) were allowed to adopt a second conformation with 30 % occupancy and the dark-state structure (pdb entry 5O4C) was held fixed. The occupancy of 30 % was chosen by assessing the results from partial occupancy refinement when the occupancy was allowed to vary and was imposed for all structural refinements for consistency. Results from structural refinement were compared against the difference electron densities and some manual adjustments were made using COOT.⁵⁷ Refinement statistics are displayed in Extended Data Table 1. Validation of structure geometry was performed using MOLPROBITY⁵⁸ and PROCHECK.⁵⁹ Structural changes were also validated by calculating simulated difference Fourier electron density maps from the refined structures^{10,20} (Extended Data Figs. 2l and 4j).

Structural analysis of large-scale protein motions

The high multiplicity of SFX data was exploited for structural analysis by randomly selecting a sub-set of experimental observations from within each SFX data-set to create 100 separate (but not independent) serial crystallography data-sets for the two resting state data-sets and the seven photo-activated data-sets, amounting to 900 resampled data-sets in total. For each of these resampled data-sets the mean and uncertainty estimates (σ) for every unique Bragg reflection were determined. Structural refinement over a cycle of 100 rigid body and 100 isotropic restrained refinements with all atoms allowed to move and with every atom having 100 % occupancy were then performed against each of these 900 resampled data-sets using pdb entry 5NJ4 as a starting model. R_{free} values ranging from 22.1 % to 23.1% were recovered. Coordinate errors associated with each individual structural refinement are estimated 60 to be ≤ 0.2 Å.

The distances between the $C\alpha$ atoms of the photo-activated and resting RC_{vir} structures were compared pairwise using the miller package of CCTBX. A 100×100 Euclidian distance matrix was then calculated for every $C\alpha$ atom and every time point according to: $\Delta r_{ij}^{\Delta t, dark} = |\mathbf{r}_i^{\Delta t} - \mathbf{r}_j^{dark}|$, where i and j vary from 1 to100 and denote resampled dataset numbers, $\Delta r_{i,j}$ depicts the distance separating the $C\alpha$ coordinates of datasets i and j, and $\mathbf{r}_i^{\Delta t}$ and \mathbf{r}_j^{dark} are the refined coordinates obtained from the photo-activated or dark structures, respectively. A second order Taylor series expansion was then used to estimate the mean and error associated with the ratio $\Delta r_{ij}^{\Delta t, dark}/\Delta r_{ij}^{dark, dark}$ arising from coordinate variations within each set of 100 structural refinements. This expansion leads to the expression:

Error weighted mean ratio = $<\Delta r_{ij}^{state,dark}>/<\Delta r_{ij}^{dark,dark}>$

$$-var(\Delta r_{ij}^{~~dark,dark}) \times <\!\!\Delta r_{ij}^{~~state,dark}\!\!> /<\!\!\Delta r_{ij}^{~~dark,dark}\!\!>^3 + cov(\Delta r_{ij}^{~~state,dark},~\Delta r_{ij}^{~~dark,dark}) / <\!\!\Delta r_{ij}^{~~dark,dark}\!\!>^2$$

where <X> is the mean of the set X, var(X) is the variance of the set X, and cov(X,Y) is the covariance of two sets X and Y. The resulting error weighted mean ratios are represented in Fig. 4A to 4C and Extended Data Fig. 7 in which movements are coloured from grey (movements \le 80 % of the maximum ratio) to red (movements \ge 95 % of the maximum ratio). Full occupancy structural refinement avoided systematic bias in this analysis arising from partial occupancy structural refinement with a single dark-conformation held fixed, but at the

limitation, this analysis extracted recurring structural motions that evolve with time (Fig. 4a-c and Extended Data Fig. 7, Supplementary Video 3) and in a manner that is both consistent

cost of underestimating the magnitude of light-induced conformational changes. Despite this

with the known time-scales of the electron transfer reactions (Fig. 1) and theoretical

predictions (Fig. 4d-f, Supplementary Video 4).

Tests of the statistical significance of recurring difference electron density features

For each of the seven experimental difference Fourier electron density maps ($\Delta t = 1$ ps, 5 ps (data set a), 5 ps (data set b), 20 ps, 300 ps (data set b), 300 ps (data set a), 8 µs) a lower-pedestal of 3.0 σ was applied such that all electron density with an amplitude < 3.0 σ was set to zero. Both positive and negative difference electron densities were then integrated within a 4.5 Å radius sphere about a chosen coordinate (Extended Data Fig. 3j) as described for the analysis of TR-SFX data recorded from bacteriorhodopsin.²⁰ These positive (A+) and negative (A-) integrated difference electron density amplitudes were merged to yield a single amplitude according to: $A(\mathbf{r}) = ((A+)^2+(A-)^2)^{\frac{1}{2}}$ about the centre of integration \mathbf{r} . The results of this analysis are presented in Fig. 2f where six centres of integration, \mathbf{r} , are chosen as: the centre of the BPh_M ring; the magnesium atom of BCh_M; the mid-point between the two magnesium atoms of the two SP bacteriochlorophylls; the magnesium atom of BCh_L; the centre of the BPh_L ring; and the centre of the ketone containing six-carbon ring of menaquinone Q_A.

For tests of statistical significance (Extended Data Table 3), this set was complemented by the addition of amplitudes extracted by integration about the iron atoms of haem₁, hame₂, haem₃ and haem₄ to create a set of ten amplitudes for each of the seven time points:

[A(BPh_M), A(BCh_M), A(SP), A(BCh_L), A(BPh_L), A(Q_A), A(H₁), A(H₂), A(H₃), A(H₄)]_{Δt} arranged as a 10 × 7 element matrix. Control "noise only" $|F_{obs}|^{dark}$ - $|F_{obs}|^{dark}$ isomorphous difference Fourier electron density maps were calculated by first selecting sixteen resampled data-sets from the set of 100 generated from the 2015 RC_{vir} dark data, and sixteen resampled data-sets from the set of 100 generated from the 2016 RC_{vir} dark data. Eight $|F_{obs}|^{dark}$ - $|F_{obs}|^{dark}$ isomorphous difference Fourier electron density maps were then calculated by pair-wise comparisons between the sixteen resampled data-sets of the 2015 data, and another eight difference Fourier electron density maps were calculated by pairwise-comparisons of the

sixteen resampled data-sets of the 2016 data. Seven control difference Fourier electron density maps were then randomly selected from the set of sixteen "noise-only" maps, difference electron density values with an amplitude lower than 3σ were set to zero, and a set of $[A(\mathbf{r}, dark-dark)]$ were created by integrating the remaining difference electron density within a 4.5 Å radius sphere centred upon the RCvir cofactors as described above. A two-sample ttest was then performed in MATLAB to determine if the set of seven time-dependent amplitudes $[A(\mathbf{r},\Delta t)]$ and the set of seven "noise-only" amplitudes $[A(\mathbf{r},\mathrm{dark}-\mathrm{dark})]$ were indistinguishable from one another (the null hypothesis). The t-tests were then repeated 1000 times by randomly selecting a different combination of seven control amplitudes [A(r,darkdark)] from the sixteen "noise-only" difference Fourier electron density maps calculated above (of $16!/(9!\times7!) = 11440$ possible different combinations of the 16 control maps). The results of this analysis are summarized in Extended Data Table 3 and show that, when a threshold of $p \le 0.001$ is applied, the difference electron density amplitudes associated with the SP cannot be ascribed to noise. When a threshold of $p \le 0.0125$ is applied and the last three time-points ($\Delta t = 300 \text{ ps}$ (a), 300 ps (b), 8 µs) are examined as a set, the difference electron density amplitudes associated with the SP, BCh_L and Q_A cannot be ascribed to noise. Conversely, the set of difference electron density amplitudes associated with most other cofactors, as well as all sets of difference electron density amplitudes generated from noise-only maps, are indistinguishable from noise according to the results of this two sample t-test (Extended Data Table 3).

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OM/MM Geometry optimizations

Initial coordinates were taken from PDB databank entry 5O4C and missing residues and cofactor segments were retrieved from PDB entry 1PRC.⁶² Protonation states of residues were chosen based on their reference pK_a values and structural criteria such as hydrogen bond interactions. After the addition of protons to the structure, a 200 step steepest descent geometry optimization was performed with Gromacs 4.5⁶³ to relax these coordinates. During this optimization the positions of the heavy atoms were constrained to their positions in the x-ray structure. As in previous work, the interactions were modelled with the Amber03 force-field.^{64,65} Non-bonded Coulomb and Lennard-Jones interactions were evaluated without periodic boundary conditions (PBC) and using infinite cut-offs.

After relaxing hydrogens with molecular mechanics (MM) optimization, we performed several Quantum Mechanics (QM)/MM geometry optimizations of all atoms in the reaction centre, using the interface between the TeraChem quantum chemistry package^{66,67} and Gromacs 4.5.⁶³ These optimizations were also performed without PBC and with infinite cutoffs for the Coulomb and Lennard-Jones interactions. The QM subsystems (Extended Data Fig. 6a-c) were modelled with unrestricted DFT. In these DFT calculation we used the PBE0 functional⁶⁸ in combination with the LANL2DZ basis set.⁶⁹ Empirical corrections to dispersion energies and interactions were introduced with Grimme's DFT-D3 model.⁷⁰ The remainder of the protein, including crystal water molecules, was modelled with the Amber03 force-field,^{64,65} in combination with the TIP3P water model.⁷¹ We searched for minimum-energy geometries in all relevant oxidation states of the system using the limited-memory Broyden-Fletcher-Goldfard-Shannon quasi-Newton optimization algorithm.

The goal of these optimization steps was to characterize the structural relaxation of the protein in response to changes in the electronic states of the cofactors along the A-branch of the photo-induced electron transfer process. We therefore examined the following electronic configurations:

i. All cofactors in their resting states: SP⁰, BPh_L⁰, Q_A⁰

- 739 ii. Special pair photo-excited, other cofactors in resting state: SP*, BPh_L⁰, Q_A⁰
- 740 iii. Special pair photo-oxidized, BPh_L reduced: SP⁺¹, BPh_L⁻¹, Q_A⁰
- iv. Special pair photo-oxidized, Q_A reduced: SP⁺¹, BPh_L⁰, Q_A⁻¹

Since including all co-factors into one large QM-region is computationally too demanding, we performed the optimizations with a different QM subsystem for each co-factor, including nearest residues, in all relevant electronic states. The structural response to the change in electronic state (Fig. 4d-f) was obtained by comparing the optimized geometries and potential

energies in the various oxidation states.

To quantify the effect of photo-absorption by the special pair (SP, i to ii) we first optimized the resting state with the SP plus nearby residues in the QM region (Extended Data Fig. 6a), described at the PBE0/LANL2DZ level of theory plus D3 dispersion corrections. This structure was used as a reference for the optimized structures in the excited state (SP*, ii) and after photo-oxidation (SP⁺¹, iii). Using the same QM/MM subdivision, we optimized the system in the first singlet excited state (S₁) by switching the QM description to the Time-dependent DFT within the Tamm-Dancoff approximation, 72 and in the photo-oxidized state by switching the spin state of the electronic wave function to the lowest energy doublet state (D₀). In the QM/MM optimization of the D₀ state of the SP, we modelled the BPh_L with point charges representing the reduced state of that co-factor. Only very modest protein structural changes were associated with the optimized geometries with the SP in the S₁ and D₀ relative to the reference structure in the resting state (S₀). Likewise, we also optimized the geometry of the protein with the BPh_L and nearby residues in the QM region (Extended Data Fig. 6d) in both the lowest energy singlet (S₀, i) and doublet (D₀, reduced, iii) states. In the optimization of the D₀ state of BPh_L, the partial charges on the SP were changed to reflect its photo-

oxidized (D_0, SP^{+1}) state. Again the structural response is rather minor, as the geometries are very similar (Extended Data Fig. 6e).

In the next step of the electron transfer process, the electron transfers from BPh_L to Q_A (iv). We optimized the protein with Q_A plus its immediate environment, including the non-heme Fe^{2+} site, in the QM region. The optimized structures in the resting and reduced states are compared in Extended Data Fig. 6f. Reduction of Q_A from menaquinone to (deprotonated) semiquinone induces significant structural changes in the Q_A binding pocket. In line with the difference densities observed at 300 ps after photo-excitation, the hydrogen bond between the Q_A carbonyl and $His217_M$ reduces by 0.17 Å. We suggest that the reduction of this hydrogen bond helps stabilizing the negative charge on the Q_A .

To quantify the overall structural response to the electron transfers, we computed the displacements of the atoms in the various states (ii-iv) with respect to the structure of resting state (i) and recorded these displacements as B-factors to the pdb coordinate file of the resting state. Because only one co-factor was included in the QM region of our QM/MM optimizations, we summed up the displacements of both QM/MM optimizations of each redox state. Fig. 4d-f represents these displacements as colours.

Stabilization energies

To estimate the energetic effects of the protein structural changes on the electron transfer process, we computed the adiabatic and vertical electron affinities for Q_A in isolation and in the optimized QM/MM protein models. These energies are shown schematically in Extended Data Fig. 6g,h. For the neutral Q_A in vacuum, the electron affinity *without* structural relaxation is 164.5 kJ.mol⁻¹ (vertical electron affinity, VEA). Structural relaxation in response

to adding the electron increases the affinity further by 24 kJ.mol⁻¹, so that the energy difference between the neutral reactant minimum on the one hand and the reduced product minimum is 188.5 kJ.mol⁻¹ (adiabatic electron affinity, AEA) on the other hand. The calculated AEA is in good agreement with results from previous computations, ²⁹ but is an overestimation with respect to the experimental value for the related 1.4-naphthoguinone (175 kJ.mol⁻¹).⁷³ Inside the protein environment, the VEA is much higher (258 kJ.mol⁻¹), part of which we attribute to the electrostatic interaction between the reduced QA co-factor with the positively charged Fe2+ ligand site. Structural relaxation of both the QA cofactor and the protein environment increases the electron affinity by 60 kJ.mol⁻¹ to yield an AEA of 318 kJ.mol⁻¹. Thus, the results of the computations suggest that the structural response of the protein adds another 36 kJ.mol⁻¹ to the intrinsic relaxation energy of Q_A (24 kJ.mol⁻¹ in vacuum) as concluded in earlier computations.²⁹ We note that in this analysis we focussed only on the effect of the structural response on the affinity of QA. To estimate the total reaction energy associated with the photo-induced electron transfer process from the SP to Q_A, we also need the absolute energies of the neutral, photo-excited and oxidized states of the SP as well as the neutral and reduced states of BPh_I. However, since these energies were not computed with identical QM/MM setups, we do not provide an accurate estimate here.

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EXTENDED DATA FIGURE LEGENDS

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Extended Data Figure 1: Dependence of transient infrared (IR) spectra on the pump laser fluence. a, Time-resolved IR difference spectra recorded from RC_{vir} in D₂O buffer for four time points and variable pump-laser fluence following 960 nm excitation. Spectral changes are consistent with earlier reports.^{74,75} Dominant time-dependent features are highlighted and include: (1) a negative band at 1687 cm⁻¹ (9-carbonyl stretch in SP_L and SP_M) and (2) a broad negative band centred at 1680 cm⁻¹ (9-carbonyl stretch in BCh_L). Transient changes from $\Delta t = 2$ ps to 5 ps correlate with the time-dependent photo-oxidation of SP. Grey columns indicate decreasing bands whereas cyan columns indicate increasing bands. b, Time resolve IR difference spectra normalized and superimposed. These spectra remain superimposable throughout the pump-laser fluence domain probed. c. Dependence of the magnitude of the difference IR signal (calculated as the sum of the absolute value over all pixels) on the pump laser fluence. Below 20 GW/cm² the absolute signal increases approximately linearly whereas above 60 GW/cm² the signal flattens out as a plateau is reached. This plateau is consistent with the complete bleaching of the special pair absorption at 960 nm in the photoexcited state and correlates with the phenomenon of hole-burning in RCvir. 42

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Extended Data Figure 2: Overview of the experimental $\mathbf{F}_{obs}^{light} - \mathbf{F}_{obs}^{dark}$ difference Fourier electron density maps calculated between the photo-excited and resting state data. a, Overview of the structure of RC_{vir} when viewed from the plane of the membrane. TM helices E_L and D_M are highlighted in red and gold respectively. b-h, Difference Fourier electron density maps for the time-points: b, $\Delta t = 1$ ps; c, $\Delta t = 5$ ps (a); d, $\Delta t = 5$ ps (b); e, $\Delta t = 20$ ps; f, $\Delta t = 300$ ps (b); g, $\Delta t = 300$ ps (a); h, $\Delta t = 8$ µs. All maps are contoured at ± 4.0 σ

(σ is the root mean square electron density of the map). Positive difference electron density is shown in blue and negative difference electron density is shown in gold.

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Extended Data Figure 3: Light-induced electron density changes in RCvir at the site of **photo-oxidation**. Experimental $F_{obs}^{light} - F_{obs}^{dark}$ isomorphous difference Fourier electron density maps for the time-points: \mathbf{a} , $\Delta t = 1$ ps. \mathbf{b} , $\Delta t = 5$ ps (data set a); \mathbf{c} , $\Delta t = 5$ ps (data set b). d, $\Delta t = 20$ ps; e, $\Delta t = 300$ ps (data set b); f, $\Delta t = 300$ ps (data set a); g, $\Delta t = 8$ µs. h-i, Principal component from SVD analysis of difference Fourier electron density maps for: h, the first four time-points, $\Delta t = 1$ ps, 5 ps (data set a), 5 ps (data set b) and 20 ps; i, the final three time-points, $\Delta t = 300$ ps (data set a), 300 ps (data set b) and 8 μ s. All maps are contoured at \pm 3.2 σ (blue, positive difference electron density; gold, negative difference electron density). Peak heights listed in Table 2 are marked with coloured circles in panels h and i. j, Representation of a sphere of radius 4.5 Å used to integrate difference electron density above a pedestal of 3.0 σ as described in reference²⁰. Positive difference density amplitudes (A+) and negative difference density amplitudes (A-) were merged according to $((A+)^2+(A-)^2)^{1/2}$ and are displayed in Fig. 2f. These integrated difference electron density values are also used in the statistical analyses presented in Extended Data Table 2. k, Simulated difference Fourier electron density maps near the SP calculated from the refined structure for $\Delta t = 20$ ps versus the refined resting state structure. Data are cut at 2.8 Å resolution for comparison with experimental data. The simulated map is contoured at \pm 12.0 σ. I, Relative magnitude of the singular values resulting from SVD analysis of the difference Fourier electron density maps for all seven time-points. m, First (blue) and second (red) columns of the unitary matrix U resulting from SVD analysis of all seven time-points weighted according to their corresponding singular values. The first right singular vector (the principal SVD component) makes a significant and positive contribution to all time-points. In contrast the difference electron density maps for $\Delta t = 5 \text{ps}$ (a) and 300 ps (a) contain strong negative contributions from the second right singular vector, which suggests that differences between the two experimental runs cause systematic differences in the difference Fourier electron density maps that are separated by SVD analysis.

Extended Data Figure 4: Electron density changes near the active site of myoglobin. The F_{obs}^{light} – F_{obs}^{dark} difference Fourier electron density map for $\Delta t = 10$ ps was calculated from data (pdb entries 5CNE minus 5CMV) recorded during TR-SFX studies of the photodissociation of carbon monoxide from the active site of myoglobin. The 21 a, Data from 15 Å to 1.8 Å were used to calculate the difference Fourier electron density map. b, Data from 15 Å to 3.0 Å were used to calculate the difference Fourier electron density map. Positive and negative difference electron density features associated with the heme group indicate slight motions of the protoporphyrin-IX. Blue represents positive difference density (contoured at 3.0 σ) and gold represents negative density (contoured at -3.0 σ). At 1.8 Å the maximum amplitude of the highlighted difference density features are: a+, 14.5 σ ; a-, 14.0 σ ; b+, 4.6 σ ; b- 4.0 σ ; c+, 4.1 σ , c- 4.0 σ ; d+, 3.8 σ . When the map is recalculated after data is cut to 3.0 Å resolution the corresponding values are: a+, 11.9 σ ; a-, 12.4 σ ; b+, 4.2 σ ; b- 4.1 σ ; c+, 3.7 σ , c- 2.9 σ ; d+, 3.2 σ .

Extended Data Figure 5: Light-induced electron density changes in RC_{vir} within the menaquinone binding pocket. Experimental $F_{obs}^{light} - F_{obs}^{dark}$ difference Fourier electron density maps for the time-points: \mathbf{a} , $\Delta t = 1$ ps. \mathbf{b} , $\Delta t = 5$ ps (data set a); \mathbf{c} , $\Delta t = 5$ ps (data set b). \mathbf{d} , $\Delta t = 20$ ps; \mathbf{e} , $\Delta t = 300$ ps (data set b); \mathbf{f} , $\Delta t = 300$ ps (data set a); \mathbf{g} , $\Delta t = 8$ μ s. All seven maps are contoured at ± 3.0 σ (blue, positive difference electron density; gold, negative difference electron density). \mathbf{h} - \mathbf{i} , Principal component from SVD analysis of difference

Fourier electron density maps for: \mathbf{h} , the first four time-points, $\Delta t = 1 \text{ps}$, 5 ps (data set a), 5 ps (data set b) and 20 ps; \mathbf{i} , the final three time-points, $\Delta t = 300 \text{ ps}$ (data set a), 300 ps (data set b) and 8 μ s. SVD difference Fourier maps are contoured at $\pm 3.2 \, \sigma$. Peak heights listed in Table 2 are indicated in panel i. \mathbf{j} , Simulated difference Fourier electron density maps near Q_A calculated from the refined structure for $\Delta t = 300 \, \text{ps}$ (data set a) versus the refined resting state structure. Data are cut at 2.8 Å resolution for comparison with experimental data. The simulated map is contoured at $\pm 12.0 \, \sigma$.

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Extended Data Figure 6: Results of QM/MM energy minimization calculations. a-c, QM subsystems used in the QM/MM optimizations of the co-factors. Atoms included in the QM region are shown in ball-and-stick representation, while the other atoms of the protein are shown as cartoons. Atoms belonging to the co-factor that are not included in the QM region are shown as sticks. Water molecules that are not part of the QM region are not shown. a, Atoms included in the QM region associated with the SP. b, Atoms included in the QM region associated with the BPh_L. c, Atoms included in the QM region associated with the Q_A. d, Structures of the energy minimized resting conformation (black: SP and BPh_L neutral) and that after the first electron transfer step (blue: SP⁺¹ and BPh_L⁻¹) shown near the special pair. e, Structures of the energy minimized resting conformation (black: SP and BPh_L neutral) and that after the first electron transfer step (blue: SP⁺¹ and BPh_L⁻¹) shown near BPh_L. **f**, Structures of the energy minimized resting conformation (black: SP and Q_A neutral) and that after the second electron transfer step (blue: SP^+ and Q_A^-) shown near Q_A . The H-bond between O_1 of the reduced semiquinone and $N_{\delta 1}$ of His217_M is predicted to be shortened by 0.17 Å when menaguinone is reduced. g-h, Schematic depiction of the potential energy surfaces for menaguinone in vacuum (g) and within the protein (h) in the resting (neutral) electronic state (black) and the reduced electronic state (red). Vertical electron affinities (VEA), adiabatic electron affinities (AEA) and relaxation energies E_r were computed at the PBE0-D3/LANL2DZ level of DFT for isolated Q_A and at the PBE0-D3/LANL2DZ/Amber03 QM/MM level for Q_A within the protein.

Extended Data Figure 7: Recurring movements of Ca atoms quantified by structural refinement. Recurring movements of Ca atoms quantified by full occupancy structural refinement against 100 randomly resampled TR-SFX data sets. Recurring movements are represented as error weighted mean ratios relative to 100 control structural refinements (Methods). Error weighted means ratios are coloured from grey (< 80 % of the maximum ratio) to red (\geq 95 % of the maximum ratio). **a,** Recurring movements of Ca atoms associated with refinements against data for $\Delta t = 1$ ps. **b,** Recurring movements of Ca atoms for $\Delta t = 5$ ps (data set a). **c,** Recurring movements of Ca atoms for $\Delta t = 5$ ps (data set b). **d,** Recurring movements of Ca atoms for $\Delta t = 300$ ps (data set b). **f,** Recurring movements of Ca atoms for $\Delta t = 300$ ps (data set a). **g,** Recurring movements of Ca atoms for $\Delta t = 8$ µs. Transmembrane helices are drawn as rods.

FOOTNOTES TO EXTENDED DATA TABLES 1044 1045 1046 Extended Data Table 1: Crystallographic data and refinement statistics Footnotes to Extended Data Table 1 1047 $^{\dagger}R_{split} = 1/\sqrt{2 \; \frac{\sum hkl \; |Ieven-Iodd|}{1/2 \sum hkl \; |Ieven+Iodd|}}$ 1048 [‡]Values in parenthesis is those of the highest resolution shell. 1049 1050 [£]Ratio of the number of indexed images to the total number of images. 1051 1052 1053 Extended Data Table 2: Difference Fourier electron density peak amplitudes Footnotes to Extended Data Table 2 1054 *Positive and negative difference electron density peaks are marked on Extended data Figs. 1055 3h, 3i and 5i. Difference electron density values were read out manually in COOT. 1056 Amplitudes $< 2.4 \sigma$ are not shown. 1057 1058 1059 **Extended Data Table 3:** Two sample t-tests of integrated electron density amplitudes 1060 Footnotes to Extended Data Table 3 1061 *Mean σ /sphere represents the root-mean-square of the difference electron density above and 1062 below a threshold of $\pm 3.0 \, \sigma$ within a sphere of radius 4.5 Å centred on the Mg²⁺ atoms of the 1063 BChs; the Fe atoms of the haems; and ring centres of the menaguinone and BPhe cofactors. 1064 All values are scaled relative to a mean value for SP = 1. 1065 [†]The number of elements in each sampled set of the two-sample t-test. 1066 [‡]Difference electron density associated with haem₃ gave the lowest mean σ /sphere (0.14) 1067 and an exceptionally low standard deviation (± 0.06). Randomly generated control data 1068 gave the corresponding values of 0.09 ± 0.22 . The coincidence of a set with a low standard 1069 deviation being compared against a control set with an exceptionally high standard 1070 1071 deviation frequently yielded low p-values. Nevertheless, the experimental difference 1072 features associated with haem₃ are weak and are therefore not physically meaningful. [£]Thirty two control data sets were generated by randomly selecting a sub-set of the dark 1073 1074 observations to generate a new data set. Sixteen control difference Fourier electron density maps were then calculated between two of the control data-sets to generate maps 1075 representing the noise inherent within the SFX experiment. Mean σ/sphere calculations and 1076 all other steps proceeded as with the light versus dark experimental difference Fourier 1077 1078 electron density maps. 1079 ^{ϵ}The reference set consisted of a set of seven maps (N_B = 7) randomly selected from sixteen control maps. Random selection of the control set and the t-test was repeated 1000 1080

times, from a possible set of $16!/(9!\times7!) = 11440$ different combinations of the 16 control maps. Percentage values give the fraction of occasions when $p \le threshold$ for the resulting two sample t-test.

\$\$The hypothesis that the two sets of the t-test are indistinguishable is either true (0) or false (1) at the specified confidence level.

\$\$Features identified as distinguishable from noise with a degree of confidence defined by the given p-value.

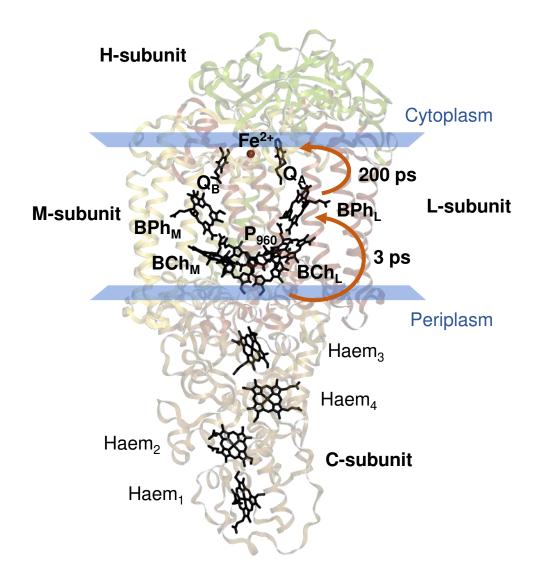


Figure 1: Dods et al.

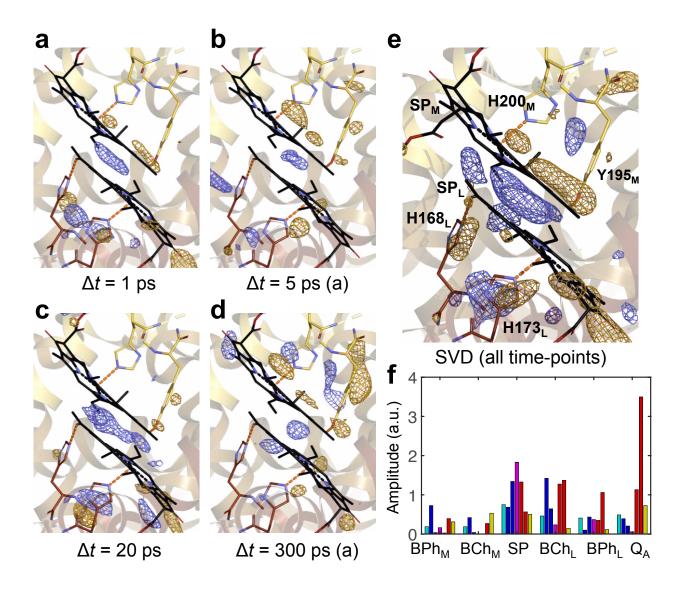


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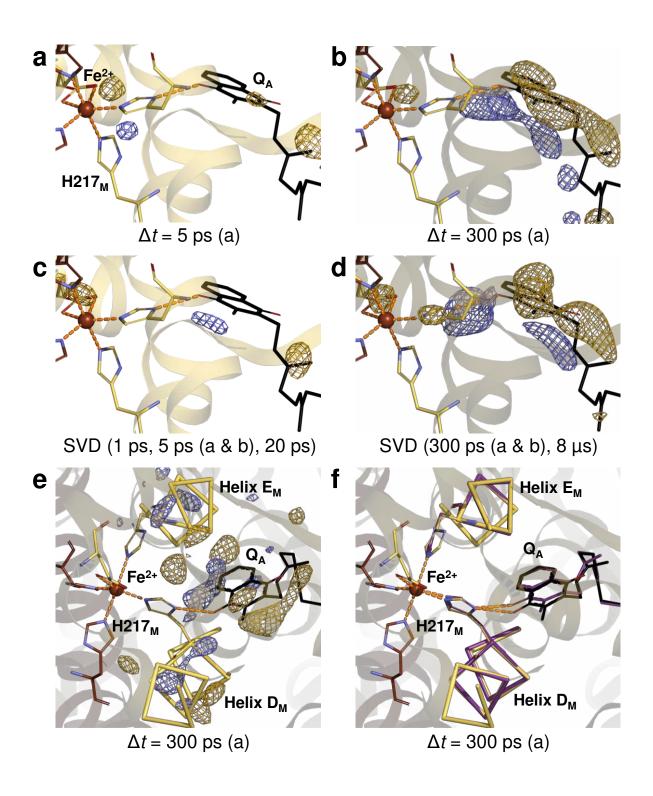


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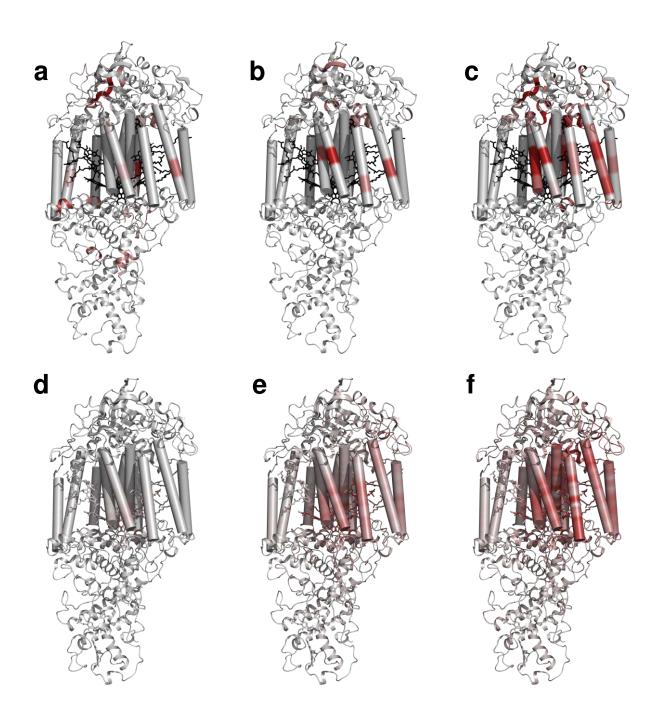


Figure 4: Dods et al.