

UNIFORM SPLINE INTERPOLATION OPERATORS IN L_2

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1. Let $m \geq 1$ and $n \geq 2$ be positive integers. Define the class \mathcal{S}^{2m} of cardinal splines of degree $2m - 1$ to be those functions S satisfying

(i) S is a polynomial of degree at most $2m - 1$ on each of the intervals

$$(1) \quad [i, i + 1], \quad i = 0, \pm 1, \pm 2, \dots$$

(ii) $S \in C^{2m-2}(-\infty, \infty)$

If in addition

$$(2) \quad S(x + n) = S(x) \quad \text{all } x,$$

we say S is a periodic spline and denote this class by \mathcal{S}_n^{2m} . Let $l_p(n)$ be the space of real n -tuples possessing the norm

$$\begin{aligned} \|y\|_p &= \left(\sum_{i=1}^n |y_i|^p\right)^{1/p}, & 1 \leq p < \infty \\ &= \max_{1 \leq i \leq n} |y_i|, & p = \infty. \end{aligned}$$

Then the periodic spline interpolation operator $\mathcal{L}_n^{2m} : l_p(n) \rightarrow L_p[0, n]$ is defined by letting $\mathcal{L}_n^{2m}y$ be that unique element of \mathcal{S}_n^{2m} satisfying

$$\mathcal{L}_n^{2m}y(i) = y_i, \quad i = 1, 2, \dots, n.$$

Similarly if $(y_i)_{i=-\infty}^{\infty}$ is in l_p , the class of doubly infinite real p -summable sequences, then the cardinal spline interpolation operator

$$\mathcal{L}^{2m} : l_p \rightarrow L_p(-\infty, \infty)$$

is defined by letting $\mathcal{L}^{2m}y$ be that unique element of $\mathcal{S}^{2m} \cap L_p(-\infty, \infty)$ which satisfies

$$\mathcal{L}^{2m}y(i) = y_i, \quad i = 0, \pm 1, \pm 2, \dots$$

The problem of calculating the norms of these operators for $p = \infty$ was first posed by Schurer and Cheney [6], who obtained

THEOREM 1 (Schurer and Cheney). *Let $\beta = 2 + \sqrt{3}$. Then*

$$\begin{aligned} (3) \quad \|\mathcal{L}_n^4\|_\infty &= 1 + \frac{3}{2}(\beta^k - \beta)(\beta^k + 1)^{-1}(\beta - 1)^{-1}, & n = 2k \\ &= 1 + \frac{3}{2}(\beta^k - \beta)(\beta^k + \beta)(\beta^{2k} + \beta)^{-1}(\beta - 1)^{-1}, & n = 2k - 1, \\ &\|\mathcal{L}^4\|_\infty = (1 + 3\sqrt{3})/4. \end{aligned}$$

Solutions were later obtained by Schurer [5] for $m = 3$ and Richards [1] for

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arbitrary m . The case $p = 1$ was investigated by the author [2], who established the following:

THEOREM 2. *Let $\alpha = 2 - \sqrt{3}$. Then*

$$\begin{aligned}
 (4) \quad \| \mathcal{L}_n^4 \|_1 &= \frac{3 + \sqrt{3}}{6} [\sqrt{3} - (9\sqrt{3} - 15)\alpha^{k-1} \\
 &\quad + (7\sqrt{3} - 12)\alpha^{2k-1}](1 - \alpha^{2k})^{-1} \quad n = 2k, \\
 &= \frac{3 + \sqrt{3}}{3} [\sqrt{3} - (7\sqrt{3} - 12)\alpha^{2k}] \\
 &\quad \cdot (1 + \alpha^{2k+1})^{-1}, \quad n = 2k + 1, \\
 \| \mathcal{L}^4 \|_1 &= (1 + \sqrt{3})/2.
 \end{aligned}$$

The case $p = 2$ will be solved in this paper. The contrast with the previous results is striking.

THEOREM 3.

$$(5) \quad \| \mathcal{L}_n^{2m} \|_2 = 1.$$

More precisely

$$(6) \quad \| \mathcal{L}_n^{2m} y \|_2 \leq \| y \|_2, \quad y \in l_2(n)$$

and equality holds in (6) if and only if $y_i \equiv \text{constant}$.

THEOREM 4.

$$(7) \quad \| \mathcal{L}^{2m} \|_2 = 1,$$

and

$$(8) \quad \| \mathcal{L}^{2m} y \|_2 < \| y \|_2, \quad y \in l_2.$$

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2. Before proceeding with the proofs of the theorems, we must first discuss some preliminary results found in [4]. Let

$$\begin{aligned}
 M_1(x) &= 1, & |x| &\leq \frac{1}{2} \\
 &= 0, & |x| &> \frac{1}{2}
 \end{aligned}$$

and define the central B -spline of order k , $M_k(x)$, to be the k -fold convolution of $M_1(x)$ with itself:

$$(9) \quad M_k(x) = M_1 * M_1 * \dots * M_1(x) \quad (k \text{ times}).$$

$M_k(x)$ is a cardinal spline of degree $k - 1$ having support on $[-k/2, k/2]$ and is a symmetric function. In addition, the Fourier transform of $M_k(x)$ is

easily computed to be

$$(10) \quad \int_{-\infty}^{\infty} M_k(x) e^{itx} dx = \psi_k(t),$$

where

$$(11) \quad \psi_k(t) = ((2/t) \sin \frac{1}{2}t)^k$$

Proof of Theorem 3. Define

$$(12) \quad \bar{M}_{2m}(x) = \sum_{\nu=-\infty}^{\infty} M_{2m}(x + \nu n).$$

Schoenberg [3] has shown that the functions $\bar{M}_{2m}(x - 1), \bar{M}_{2m}(x - 2), \dots, \bar{M}_{2m}(x - n)$ form a basis for \mathfrak{S}_n^{2m} . Hence if $\mathfrak{L}_n^{2m} y = S$, there exist reals c_1, c_2, \dots, c_n such that

$$(13) \quad S(x) = \sum_{i=1}^n c_i \bar{M}_{2m}(x - i)$$

and thus

$$(14) \quad y_i = S(i) = \sum_{j=1}^n \bar{M}_{2m}(i - j) c_j, \quad i = 1, 2, \dots, n.$$

Inverting this non-singular system and using matrix notation we obtain

$$(15) \quad c = \Omega y \quad (\Omega^{-1} = \bar{M}_{2m}).$$

Upon squaring both sides of (13) and then integrating, we get

$$(16) \quad \|\mathfrak{L}_n^{2m} y\|_2^2 = \int_0^n (S(x))^2 dx = \sum_{i=1}^n \sum_{j=1}^n \Lambda_{ij} c_i c_j = (c, \Lambda c),$$

where

$$(17) \quad \Lambda_{ij} = \int_0^n \bar{M}_{2m}(x - i) \bar{M}_{2m}(x - j) dx, \quad i, j = 1, 2, \dots, n.$$

Since \bar{M}_{2m} is symmetric, (15) and (16) imply

$$(18) \quad \|\mathfrak{L}_n^{2m}\|_2^2 = \sup_{\|y\|_2=1} (y, \Omega \Lambda \Omega y).$$

But Λ , and therefore $\Omega \Lambda \Omega$, is also symmetric. Thus if $\rho(\Omega \Lambda \Omega)$ denotes the spectral radius of $\Omega \Lambda \Omega$, then

$$(19) \quad \|\mathfrak{L}_n^{2m}\|_2^2 = \rho(\Omega \Lambda \Omega).$$

We shall now compute the eigenvalues and corresponding eigenvectors of the relevant matrices.

Because $\bar{M}_{2m}(x)$ has period n , $\bar{M}_{2m}(i - j)_{i,j=1}^n$ is the circulant matrix

$$\begin{aligned} \bar{M}_{2m} = C(\bar{M}_{2m}(0), \bar{M}_{2m}(1), \dots, \bar{M}_{2m}(m - 1), \\ \dots, \bar{M}_{2m}(n - m + 1), \dots, \bar{M}_{2m}(n - 1)). \end{aligned}$$

For the time being let us assume that

$$(20) \quad n \geq 4m$$

and thus each of the functions $M_{2m}(x + \nu n)$, $\nu = 0, \pm 1, \pm 2, \dots$ has disjoint support. Then by recalling (12) and the fact that $M_{2m}(x)$ has support on $[-m, m]$, we see that

$$\begin{aligned} \bar{M}_{2m}(i) &= M_{2m}(i), & i &= 0, 1, \dots, m-1, \\ &= 0, & i &= m, \dots, n-m, \\ &= M_{2m}(i-n), & i &= n-m+1, \dots, n-1. \end{aligned}$$

Hence

$$(21) \quad \bar{M}_{2m} = C(M_{2m}(0), M_{2m}(1), \dots, M_{2m}(m-1), 0, \dots, 0, M_{2m}(-m+1), \dots, M_{2m}(-1))$$

Thus if we define the function φ_k by

$$(22) \quad \varphi_k(\theta) = \sum_{\nu=-m+1}^{m-1} M_k(\nu) e^{i\nu\theta}$$

and let

$$\theta_j = 2\pi j/n, \quad j = 0, 1, \dots, n-1,$$

then \bar{M}_{2m} has eigenvalues $\varphi_{2m}(\theta_j)$ and corresponding eigenvectors

$$(23) \quad v_j = (1, \varepsilon_j, \varepsilon_j^2, \dots, \varepsilon_j^{n-1}), \quad j = 0, 1, \dots, n-1,$$

where $\varepsilon_j = e^{i\theta_j}$ is an n -th root of unity. Therefore Ω has eigenvalues $(\varphi_{2m}(\theta_j))^{-1}$.

To handle Λ , we first note that the condition (20) ensures that any one ‘‘hump’’ of $\bar{M}_{2m}(x - i)$ will ‘‘hit’’ at most one other ‘‘hump’’ of $\bar{M}_{2m}(x - j)$. Then since (see [4, p. 177])

$$(24) \quad \int_{-\infty}^{\infty} M_{2m}(x - i) M_{2m}(x - j) = M_{4m}(i - j)$$

and using the periodicity of $\bar{M}_{2m}(x)$, it easily follows that

$$(25) \quad \Lambda_{ij} = \bar{M}_{4m}(i - j), \quad i, j = 1, \dots, n.$$

Thus Λ has eigenvalues $\varphi_{4m}(\theta_j)$.

These results show that $\Omega\Lambda\Omega$ has eigenvalues $\varphi_{4m}(\theta_j)/\varphi_{2m}^2(\theta_j)$ and corresponding eigenvectors v_j , $j = 0, 1, \dots, n-1$.

It will now be shown that

$$(26) \quad \varphi_{4m}(\theta)/\varphi_{2m}^2(\theta) \leq 1$$

with equality holding in (26) if and only if $\theta \equiv 0 \pmod{2\pi}$. This will

establish Theorem 3, as the supremum in (18) will be attained only when y is some multiple of $v_0 = (1, 1, \dots, 1)$.

Schoenberg [4] has shown that

$$(27) \quad \varphi_k(\theta) = \sum_{i=-\infty}^{\infty} \psi_k(\theta + 2\pi i).$$

Since $\psi_{2m}(\theta) \geq 0$, we have

$$(28) \quad \begin{aligned} \varphi_{2m}^2(\theta) &= \sum_{i=-\infty}^{\infty} \sum_{j=-\infty}^{\infty} \psi_{2m}(\theta + 2\pi i) \psi_{2m}(\theta + 2\pi j) \\ &\geq \sum_{i=-\infty}^{\infty} \psi_{2m}^2(\theta + 2\pi i) = \varphi_{4m}(\theta). \end{aligned}$$

This proves (26). Note that we get equality in (28) only if $\theta \equiv 0 \pmod{2\pi}$, since otherwise $\psi_{2m}(\theta + 2\pi i) > 0$ for all i .

We make the observation that the condition (20) may be discarded, since if $n < 4m$, the data $(y_i)_{i=1}^n$ and $\mathfrak{L}_n^{2m} y$ may be extended periodically to $[0, rn]$, where r is some integer such that $rn \geq 4m$. Theorem 3 may now be applied on $[0, rn]$. ■

Proof of Theorem 4. For $y \in l_2$, let $\mathfrak{L}^{2m} y = S$. Then there exist reals $c_i, i = 0, \pm 1, \pm 2, \dots$ such that

$$(29) \quad S(x) = \sum_{i=-\infty}^{\infty} c_i M_{2m}(x - i).$$

Proceeding as before we obtain

$$(30) \quad \|\mathfrak{L}^{2m} y\|_2^2 = \int_{-\infty}^{\infty} (S(x))^2 dx = \sum_{i=-\infty}^{\infty} \sum_{j=-\infty}^{\infty} \lambda_{ij} c_i c_j = (c, \lambda c),$$

where

$$(31) \quad \begin{aligned} \lambda_{ij} &= \int_{-\infty}^{\infty} M_{2m}(x - i) M_{2m}(x - j) dx = M_{4m}(i - j), \\ & \quad i, j = 0, \pm 1, \pm 2, \dots \end{aligned}$$

and

$$(32) \quad y_i = \sum_{j=-\infty}^{\infty} M_{2m}(i - j) c_j, \quad i = 0, \pm 1, \pm 2, \dots$$

Since $\varphi_{2m}(\theta) > 0$, we may invert the sequence convolution transformation (32) to get

$$(33) \quad c_i = \sum_j \omega(i - j) y_j,$$

where

$$(34) \quad \sum_{\nu=-\infty}^{\infty} \omega(\nu) e^{i\nu\theta} = (\varphi_{2m}(\theta))^{-1}$$

and by the Wiener-Lévy theorem

$$\sum_{\nu} |\omega(\nu)| < \infty.$$

Thus

$$(35) \quad \| \mathcal{L}^{2m} y \|_2^2 = (y, \omega \lambda \omega y).$$

But then using the correspondence

$$(y_\nu)_{\nu=-\infty}^{\infty} \leftrightarrow \tilde{y}(\theta) = \sum_{\nu=-\infty}^{\infty} y_\nu e^{i\nu\theta},$$

from Parseval's identity it follows that

$$(36) \quad \| \mathcal{L}^{2m} y \|_2^2 = \left(\tilde{y}, \frac{\varphi_{4m}}{\varphi_{2m}} \tilde{y} \right) = \frac{1}{2\pi} \int_0^{2\pi} \frac{\varphi_{4m}(\theta)}{\varphi_{2m}(\theta)} |\tilde{y}(\theta)|^2 d\theta.$$

Theorem 4 is an immediate consequence of (26) and (36).

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