# **ε-ENTROPY OF THE BROWNIAN MOTION WITH THE MULTI-DIMENSIONAL SPHERICAL PARAMETER**

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## §1. Introduction

M.S. Pinsker [3] has given a general method of calculating the  $\varepsilon$ -entropy of a Gaussian process and obtained, for example, an exact proof of the estimate for the  $\varepsilon$ -entropy of the ordinary Brownian motion B(t),  $0 \le t \le 1$ , which was presented without proof by A.N. Kolmogorov [1].

In this article, we estimate the  $\varepsilon$ -entropy of the Brownian motion with the multidimensional spherical parameter, by using the expansion of the Brownian motion with a multidimensional parameter by H.P. McKean [4] and by generalizing the Pinsker's method of calculating the  $\varepsilon$ -entropy.

Let  $X(A, \omega)$ ,  $A \in E^{d}$  (*d*-dimensional Euclidean space),  $\omega \in \mathcal{Q}(P)$ , be a Brownian motion with a parameter space  $E^{d}$ , that is,  $\{X(A), A \in E^{d}\}$ forms a Gaussian system and

1) E[X(A)] = 0 for every A,

2) X(O) = 0, where O is the origin of  $E^{d}$ ,

3)  $E[(X(A) - X(B))^2] = \operatorname{dis}(A, B)$ , where E(X) and  $\operatorname{dis}(A, B)$  denote the expectation of a random variable X and the Euclidean distance between A and B, respectively.

We shall call X(A) when the parameter A is restricted to the unit sphere<sup>1</sup>)  $S^{d-1}$  in  $E^d$  the Brownian motion with the d-dimensional spherical parameter and denote it, as in the preceding case, by X(A),  $A \in S^{d-1}$ .

The  $\varepsilon$ -entropy  $H_{\varepsilon}(X)$  of the process X(A) is defined as follows: Let  $\varepsilon > 0$  be arbitrarily fixed, and consider an approximating process X'(A) for the process X(A) on  $S^{d-1}$  satisfying the condition of reproducing accuracy,

(1) 
$$\int_{S^{d-1}} E[(X'(A) - X(A))^2] \, d\sigma(A) \leq \varepsilon^2$$

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<sup>1)</sup> Without loss of generality we may consider the unit sphere only.

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where  $d\sigma$  is the uniform probability measure on  $S^{d-1}$ . Then, the  $\varepsilon$ -entropy of the process X(A) is defined as

(2) 
$$H_{\varepsilon}(X) = \inf I(X', X),$$

where I(X', X) is the amount of information contained in a process X' with respect to the process X and the infimum is taken for all processes X' satisfying the condition (1).

Our aim is to prove that the  $\varepsilon$ -entropy of the Brownian motion on  $S^{d-1}$  is of order  $\varepsilon^{-2(d-1)}$  (Theorem 2);

(3) 
$$H_{\varepsilon}(X) = O(\varepsilon^{-2(d-1)}).$$

It seems to be interesting to note that the  $\varepsilon$ -entropy (in Kolmogorov-Tihomirov's sense, cf. Kolmogorov-Tihomirov [2]) of the space of  $\frac{1}{2}$ -Hölder continuous functions of (d-1)-variables with the sup-norm has the same order  $O(\varepsilon^{-2(d-1)})$ .

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## §2. The generalization of Pinsker's method

Pinsker's method of calculating the  $\varepsilon$ -entropy of a Gaussian process with one dimensional parameter is as follows: Let X(t),  $0 \le t \le T$ , be a Gaussian process with mean 0 whose covariance function r(s,t) = E[X(s)X(t)]is continuous in (s,t). Then the  $\varepsilon$ -entropy  $H_{\varepsilon}(X)$  of the process X(t) is given by the formula

(4) 
$$H_{\varepsilon}(X) = \frac{1}{2} \sum_{\lambda_i > \theta^2} \log \frac{\lambda_i}{\theta^2} ,$$

where  $\lambda_i$   $(i = 1, 2, \dots)$  are the eigen-values of the integral operator with the kernel r(s, t) in  $L^2[0, T]$ ,  $\lambda_1 \ge \lambda_2 \ge \dots \ge 0$ , and  $\theta$  is determined (uniquely) by the equation

(5) 
$$\sum_{i=1}^{\infty} \min \left( \theta^2, \lambda_i \right) = \varepsilon^2 \cdot \varepsilon^2$$

2) By Mercer's theorem

$$\sum_{i=1}^{\infty} \lambda_i = \sum_{i=1}^{\infty} \lambda_i \int_0^T [\varphi_i(t)]^2 dt = \int_0^T \sum_{i=1}^{\infty} \lambda_i [\varphi_i(t)]^2 dt = \int_0^T r(t,t) dt < \infty.$$

The right-hand side of the relation (4) also equals to the  $\varepsilon$ -entropy of the infinite dimensional Gaussian random variable  $X^* = (X_1^*, X_2^*, \cdots)^{3}$ :

(6) 
$$X_i^* = \int_0^T \varphi_i(t) X(t) dt^{(4)} \qquad (i = 1, 2, \cdots)$$

where  $\varphi_i(t)$  is the eigen-function of the integral operator corresponding to the eigenvalue  $\lambda_i$  and  $E[X_i^* X_j^*] = \lambda_i \delta_{ij}$ .

As an example, if in particular the sequence  $\lambda_1 \ge \lambda_2 \ge \cdots \ge 0$  of the eigen-values of the integral operator with the kernel corresponding to a Gaussian process takes the form:  $\lambda_k = ck^{-s}(s > 1; k = 1, 2, \cdots)$ , then, the  $\varepsilon$ -entropy of the process is

(7) 
$$H_{\varepsilon}(X) = O(\varepsilon^{-\frac{2}{s-1}}).$$

Now, we proceed to a Gaussian process X(A),  $A \in S^{d-1}$ , with mean 0. Assume the continuity of the covariance function r(A, B) = E[X(A) X(B)] in  $S^{d-1} \times S^{d-1}$ , so  $\sum_{i=1}^{\infty} \lambda_i$  is finite (see the discussion in the footnote 2)) where  $\lambda_i$ ,  $i = 1, 2, \cdots$ , are the eigenvalues of the integral operator with the kernel r(A, B) in  $L^2(S^{d-1}, d\sigma)$ . Then, the following entirely analogous result holds, and we state it as a theorem.

THEOREM 1. The  $\varepsilon$ -entropy  $H_{\varepsilon}(X)$  of the above Gaussian process X(A),  $A \in S^{d-1}$  is

(4') 
$$H_{\varepsilon}(X) = \frac{1}{2} \sum_{\lambda_i > \theta^2} \log \frac{\lambda_i}{\theta^2}$$

where  $\lambda_i (i = 1, 2, \dots)$  with  $\lambda_1 \ge \lambda_2 \ge \dots \ge 0$  are eigen-values of the integral operator and  $\theta$  is determined by the equation (5). The right-hand side of the relation (4') equals also to the  $\varepsilon$ -entropy of the infinite dimensional Gaussian random variable  $X^* = (X_1^*, X_2^*, \dots)$ :

(6') 
$$X_i^* = \int_{S^{d-1}} \varphi_i(A) X(A) d\sigma(A) \qquad (i = 1, 2, \cdots)$$

<sup>3)</sup> The  $\varepsilon$ -entropy of  $X^*$  is defined as  $H_{\varepsilon}(X^*) = \inf I(\widetilde{X}^*, X^*)$  where the infimum is taken for all infinite dimensional approximating random variables  $\widetilde{X}^* = (\widetilde{X}_1^*, \widetilde{X}_2^*, \cdots)$  satisfying the condition:  $\sum_{i=1}^{\infty} E[(\widetilde{X}_i^* - X_i^*)^2] \leq \varepsilon^2$ .

<sup>4)</sup> This (Bochner) integral is determined as an element of  $L^2(\Omega)$ .

where  $\varphi_i(A)$  is the eigen-function of the integral operator corresponding to the eigenvalue  $\lambda_i$ , and  $E[X_i^* X_j^*] = \lambda_i \delta_{ij}$ .

*Proof.* The proof is quite similar to the proof for one dimensional parameter case dealt by M.S. Pinsker [3], except for the construction of the process  $\dot{\xi}$  ([3], formula (132)). The proof, however, can be carried out by using the extension theorem of Urysohn, so that we shall not continue the proof further.

## §3. The main result

We are now in a position to prove our main result.

THEOREM 2. The  $\varepsilon$ -entropy of the Brownian motion with the d-dimensional spherical parameter is of order  $\varepsilon^{-2(d-1)}$ ;

(8) 
$$H_{\varepsilon}(X) = O(\varepsilon^{-2(d-1)}).$$

*Proof.* According to H.P. McKean [4] the Brownian motion with the *d*-dimensional parameter can be expanded as a sum of mutually independent Gaussian processes associated with spherical harmonics. We state this expansion and some related results with the Gaussian process X(A),  $A \in S^{d-1}$ .

(9) 
$$X(A) = \sum_{n \ge 0} \sum_{l=1}^{D(n)} x_n^l(1) h_n^l(A), \ A \in S^{d-1}$$

where  $h_n^l(A)$  is a spherical harmonics of degree *n* satisfying

(10) 
$$\int_{S^{d-1}} h_n^l(A) h_m^k(A) d\sigma(A) = \begin{cases} 1, & \text{if } l = k, n = m \\ 0, & \text{otherwise,} \end{cases}$$

D(n) is the dimension of the vector space spanned by all the spherical harmonics of degree n,

(11) 
$$D(n) = (2n-2+d) \frac{(n-3+d)!}{(d-2)! n!} \quad (d \ge 2, n \ge 0)^{5}$$

and  $x_n^l(1)$   $(n \ge 0, 1 \le l \le D(n))$  are mutually independent Gaussian random variables which can be expressed in the form

(12) 
$$x_n^{l}(1) \equiv x_n^{l} = C(d) \int_0^1 C_n(u) dB_n^{l}(u) \, .$$

<sup>5)</sup> For d=2 and n=0, D(n)=1.

The processes  $B_n^l(u)$   $(n \ge 0, 1 \le l \le D(n))$  appeared in the above expression are mutually independent standard Brownian motions and

(13) 
$$C_n(u) = \frac{\int_0^{\cos^{-1}u} p_n(\cos\theta) \sin^{d-2}\theta d\theta}{\int_0^{\pi} \sin^{d-2}\theta d\theta} , \quad n \ge 0$$

with  $p_n(\cos\theta) = C_n^{\frac{d-2}{2}}(\cos\theta) / C_n^{\frac{d-2}{2}}(1)$ , where  $C_n^{\nu}(\cdot)$  is the Gegenbauer polynomial and C(d) is a constant depending only on d.

By the expansion (9) and by the independence of the random variables  $x_n^i$  with  $E[x_n^i] = 0$   $(n \ge 0, 1 \le l \le D(n))$  we easily see that the covariance function of the process X(A) is expressed in the form

(14) 
$$r(A, B) = \sum_{n \ge 0} \sum_{l=1}^{D(n)} E[(x_n^l)^2] h_n^l(A) h_n^l(B).$$

Using this, Mercer's expansion theorem shows us that the eigen-values  $\lambda_n^l (n \ge 0, 1 \le l \le D(n))$  of the integral operator with the kernel r(A, B) are equal to  $E[(x_n^l)^2]$ . Therefore, if we know the amount  $E[(x_n^l)^2]$  we can obtain the  $\varepsilon$ -entropy of the Brownian motion with the parameter space  $S^{d-1}$  by the formula (4'). In fact, we can prove in the following that for large n,  $E[(x_n^l)^2] = O(n^{-d})$ ,  $1 \le l \le D(n)$ , holds. Once the result is shown, then just by renumbering the double sequence of random variables  $x_0^l$ ,  $x_1^1$ ,  $x_1^2$ ,  $\cdots$ ,  $x_1^{D(1)}$ ,  $x_2^1$ ,  $\cdots$  into the ordinary sequence  $x_1'$ ,  $x_2'$ ,  $\cdots$ , while keeping the original order, we can easily apply Theorem 1 in §2. If  $x_k'$ , for large k, corresponds to the original random variable  $x_N^{M}(1 \le M \le D(N))$ , then by the relation  $\sum_{n=0}^{N} n^{d-2} = O(N^{d-1})$  (this nearly equals to k) and by the formula (11)  $(D(n) = O(n^{d-2})$  for large n), we obtain  $N = O(k^{\frac{1}{d-1}})$ , so that  $E[(x_k')^2] = O\left((k^{\frac{1}{d-1}})^{-d}\right) = O(k^{-\frac{d}{d-1}})$ . Then, by this and the formula (7), follows the desired result  $H_{\varepsilon}(X) = O(\varepsilon^{-\frac{2}{d-1}}) = O(\varepsilon^{-2(d-1)})$ .

Therefore, in the following, we are to prove that

(15) 
$$E[(x_n^i)^2] = O(n^{-d}), \quad 1 \le l \le D(n)$$

holds for large n.

First of all, we show the formula (15) in case the dimension d = 2 and 3, and then, generalizing it, we proceed to prove the formula (15) for  $d \ge 4$ , that is, (I) in case d is an even integer and (II) when d is odd.

In case d = 2,  $p_n(\cos \theta)$  in the expression (13) turns out to be  $\cos n\theta$ , so that  $C_n(u) = \frac{1}{n\pi} \sin(n \cos^{-1}u)$ . From this we have,

$$E[(x_n^l)^2] = \frac{1}{n^2\pi} \int_0^l \sin^2(n \cos^{-1}u) du$$
$$= \frac{1}{n^2\pi} \int_0^{\frac{\pi}{2}} \sin^2 n\theta \sin \theta \, d\theta$$
$$= O(n^{-2}).$$

While in case d = 3,  $p_n(\cos \theta) = P_n(\cos \theta)$ , hence we have  $C_n(u) = \frac{1}{2} - \frac{P_{n-1}(u) - P_{n+1}(u)}{2n+1}$  where  $P_n(\cdot)$  is the *n-th* Legendre polynomial. Then, by the orthogonality of the Legendre polynomials, we obtain

$$E[(x_n^l)^2] = \frac{1}{(2n+1)^2} \left\{ \int_0^1 (P_{n+1}(u))^2 \, du + \int_0^1 (P_{n-1}(u))^2 \, du \right\}$$
$$= O(n^{-3}).$$

In case  $d \ge 4$ , by the formula (12), we have

$$E[(x_n^l)^2] = (C(d))^2 \int_0^1 (C_n(u))^2 du$$
  
= (a constant depending on  $d$  only)  $\times \left\{ C_n^{\frac{d-2}{2}}(1) \right\}^{-2}$   
 $\times \int_0^1 \left\{ \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta \right\}^2 du$ 

and this expression becomes,

$$O(n^{-2d+6}) \times \int_0^1 \left\{ \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta \right\}^2 du$$

for large *n*, since  $C_n^{\frac{d-2}{2}}(1) = \frac{\Gamma(n+d-2)}{n! \Gamma(d-2)} = O(n^{d-3})$ .

To prove  $E[(x_n^{l})^2] = O(n^{-d})$ , we must show that the above integral (we denote it by  $I_a$ ) is of order  $O(n^{d-6})$ .

(I) The proof of the fact that  $I_d = O(n^{d-6})$  for d = 2p + 2  $(p \ge 1$ , integer).

First we estimate the integrand of the above integral. Let the following integral be denoted by  $I_p(u)$ ,

$$I_p(u) = \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta = \int_0^{\cos^{-1}u} C_n^p(\cos\theta) \sin^{2}\theta d\theta.$$

The integrand  $C_n^p(\cos \theta) \sin^{2p} \theta$  of the above integral becomes, by using the recurrence formula for the Gegenbauer polynomials

(16) 
$$\sin^2 \theta C_n^{\nu+1}(\cos \theta) = \frac{1}{2\nu} \left\{ (n+2\nu) C_n^{\nu}(\cos \theta) - (n+1) \cos \theta C_{n+1}^{\nu}(\cos \theta) \right\}$$

and the formula  $\sin \theta C_n^1(\cos \theta) = \sin (n+1)\theta$ ,

 $C_n^p(\cos\theta)\sin^{2p}\theta = \sin^2\theta C_n^p(\cos\theta)\sin^{2(p-1)}\theta$ 

$$= \frac{1}{2(p-1)} \left\{ (n+2(p-1))C_n^{p-1}(\cos\theta)\sin^{2(p-1)}\theta - (n+1)\cos\theta C_{n+1}^{p-1}(\cos\theta)\sin^{2(p-1)}\theta \right\}$$
  
$$= \frac{1}{2^{p-1}(p-1)!} \left\{ A_1^p(n)\sin\theta\sin(n+1)\theta + A_2^p(n)\cos\theta\sin\theta\sin(n+2)\theta + A_3^p(n)\cos^2\theta\sin\theta\sin(n+3)\theta + \cdots + A_p^p(n)\cos^{p-1}\theta\sin\theta\sin(n+p)\theta \right\}$$

where  $A_1^p(n), A_2^p(n), \dots, A_p^p(n)$  are polynomials of *n* of order (p-1). Noticing that  $\sin \theta \sin (n+1)\theta$ ,  $\cos \theta \sin \theta \sin (n+2)\theta$ ,  $\cdots$  and  $\cos^{p-1}\theta \sin \theta \sin (n+p)\theta$  are all expressed as the linear combinations of  $\cos n\theta$ ,  $\cos (n+2)\theta$ ,  $\cdots$ ,  $\cos (n+2p)\theta$ , we can show that the integral becomes

(17) 
$$I_p(u) = \sum_{k=0}^p \frac{B_k^p(n)}{n+2k} \sin{(n+2k)\alpha}, \quad \alpha = \cos^{-1}u$$

where  $B_k^p$ ,  $k = 0, 1, \dots, p$ , are polynomials of *n* of order at most (p-1). Therefore, changing the variable of integration into  $\alpha$ , and making use of the fact

$$\int_{0}^{\frac{\pi}{2}} \sin(n+2k)\alpha \sin(n+2l)\alpha \sin \alpha d\alpha = \frac{1}{2} \left\{ \frac{1}{1-4(k-l)^{2}} + O(n^{-2}) \right\},$$

we have

$$I_{a} = \int_{0}^{1} \{I_{p}(u)\}^{2} du = \int_{0}^{\frac{\pi}{2}} \left\{ \sum_{k=0}^{p} \frac{B_{k}^{p}(n)}{n+2k} \sin(n+2k)\alpha \right\}^{2} \sin \alpha d\alpha$$

$$=\sum_{k,l=0}^{p} \frac{B_{k}^{p}(n)B_{l}^{p}(n)}{(n+2k)(n+2l)} \int_{0}^{\frac{\pi}{2}} \sin((n+2k)\alpha) \sin((n+2l)\alpha) \sin(\alpha) d\alpha$$
  
=  $O(n^{2p-4}) (=O(n^{d-6})).$ 

The last estimation is valid if the coefficient of the term  $n^{2p-4}$  never vanishes, that is, if at least one of the coefficients of the term  $n^{p-1}$  of the polynomials  $B_k^p(n)$   $(k = 0, 1, \dots, p)$  does not vanish. But this is true, for example,  $B_0^p(n)$  has non zero coefficient of  $n^{p-1}$ .

(II) The proof of the fact that  $I_d = O(n^{d-6})$  for d = 2p+3  $(p \ge 1$ , integer).

Similarly to (I), we denote the following integral by  $I_p(u)$ ,

$$I_{p}(u) = \int_{0}^{\cos^{-1}u} C_{n}^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta = \int_{0}^{\cos^{-1}u} C_{n}^{p+\frac{1}{2}}(\cos\theta) \sin^{2p+1}\theta d\theta$$

then, by the relation

(18) 
$$C_n^{p+\frac{1}{2}}(\cos\theta) = \frac{2^p p!}{(2p)!\sin^p\theta} P_{n+p}^p(\cos\theta)$$

for the half-integer Gegenbauer polynomial  $C_n^{p+\frac{1}{2}}$  and the associated Legendre polynomial  $P_{n+p}^p$ , we have

$$I_p(u) = c(d) \int_0^{\cos^{-1}u} P_{n+p}^p(\cos\theta) \sin^{p+1}\theta d\theta$$

where c(d) is a constant depending on d. By definition,

$$P_{n+p}^{p}(x) = (1-x^{2})^{\frac{p}{2}} \frac{d^{p}}{dx^{p}} P_{n+p}(x)$$

and by changing the variable of integration into  $x = \cos \theta$ , we get

$$\frac{1}{c(d)} I_p(u) = \int_u^1 \frac{d^p}{dx^p} P_{n+p}(x) (1-x^2)^p dx$$
$$= -(1-u^2)^p \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) + 2p \int_u^1 x(1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx$$

From this, the desired integral  $I_d$  is

$$[c(d)]^{2} \cdot I_{d} = [c(d)]^{2} \int_{0}^{1} \{I_{p}(u)\}^{2} du = \int_{0}^{1} \left\{ (1-u^{2})^{p} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^{2} du$$

$$(19) \qquad -4p \int_{0}^{1} (1-u^{2})^{p} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \left\{ \int_{u}^{1} x(1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\} du$$

$$+4p^{2}\int_{0}^{1}\left\{\int_{u}^{1}x(1-x^{2})^{p-1}\frac{d^{p-1}}{dx^{p-1}}P_{n+p}(x)dx\right\}^{2}du.$$

To estimate these integrals, we first express  $(1-u^2)^p \frac{d^p}{du^p} P_n(u)$  in terms of  $P_n(u)$  and  $P_{n-1}(u)$ . For this purpose, we make use of the recurrence formula of the Legendre polynomials  $(1-x^2)P'_n(x) = n(P_{n-1}(x) - xP_n(x))$  and the differential equation derived from the Legendre's differential equation

(20) 
$$(1-x^2) \frac{d^k}{dx^k} P_n(x) - 2(k-1)x \frac{d^{k-1}}{dx^{k-1}} P_n(x) + (n+(k-1))(n-(k-2)) \frac{d^{k-2}}{dx^{k-2}} P_n(x) = 0, \quad (k \ge 2).$$

For any  $p \ge 1$ , we have

(21) 
$$(1-u^2)^p - \frac{d^p}{du^p} P_n(u) = P_{n-1}(u)Q_{n-1,p}(u) + P_n(u)Q_{m,p}(u)$$

where  $Q_{n-1,p}(u)$  and  $Q_{n,p}(u)$  are polynomials of u of the form

(22) 
$$Q_{n-1,p}(u) = \sum_{k=0}^{p-1} C_k(n) u^k, \quad Q_{n,p}(u) = \sum_{k=0}^p D_k(n) u^k.$$

The coefficients  $C_0(n), C_1(n), \dots, C_{p-1}(n), D_0(n), D_1(n), \dots, D_p(n)$  have the following properties: (i)  $C_{p-1}(n) \neq 0$ ,  $D_p(n) \neq 0$  (ii) they are the polynomials of n with the order at most p (iii) if p is an even integer, then  $D_0(n)$  is the polynomial of order p and if p is odd,  $C_0(n)$  is the polynomial of order p. By these facts and by the property of the Legendre polynomial:  $\int_0^1 \{P_n(x)\}^2 dx = O(n^{-1})$  for large n, we can easily show that the first integral of the right-hand side of the equality (19) becomes,

$$\int_{0}^{1} \left\{ (1-u^{2})^{p} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^{2} du = \int_{0}^{1} (1-u^{2})^{2} \left\{ (1-u^{2})^{p-1} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^{2} du$$
$$= O(n^{2(p-1)}) \cdot O(n^{-1}) = O(n^{d-6}) .$$

For the second integral of the right-hand side of (19), we have

$$\left| \int_{0}^{1} (1-u^{2})^{p} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \left\{ \int_{u}^{1} x(1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\} du \right|$$

$$\leq \left\{ \int_{0}^{1} \left\{ (1-u^{2})^{p} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^{2} du \right\}^{1/2} \cdot \left\{ \int_{0}^{1} \left\{ \int_{u}^{1} x(1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^{2} du \right\}^{1/2}$$

The first term of the product on the right-hand side of the inequality, by the above result, has the order  $O(n^{\frac{d-6}{2}})$  and the integrand of the second term can be evaluated as follows:

$$\begin{split} \left\{ \int_{u}^{1} x (1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^{2} &\leq \int_{u}^{1} x^{2} dx \cdot \int_{u}^{1} \left\{ (1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) \right\}^{2} dx \\ &< \int_{0}^{1} x^{2} dx \cdot \int_{0}^{1} \left\{ (1-x^{2})^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) \right\}^{2} dx = O(n^{d-6}) \,. \end{split}$$

Hence the second integral is at most of order  $O(n^{d-\theta})$ . As for the last integral of the equality (19), by a similar approach, we estimate it to be at most of order  $O(n^{d-\theta})$ . This proves the desired result for  $d = 2p + 3(p \ge 1)$ , and thus we have proved the theorem completely.

### References

- [1] A.N. Kolmogorov: Theory of the transmission of information, Amer. Math. Soc. Translations Ser. 2, 33 (1963), 291-321.
- [2] A.N. Kolmogorov and V.M. Tihomirov: ε-entropy and ε-capacity of sets in functional spaces. Amer. Math. Soc. Translations Ser. 2, 17 (1961), 277-364.
- [3] М.С. Пинскер, Гауссовские источники, Проблемы Перебачи Информации, 14 (1963) 59–100.
- [4] H.P. McKean: Brownian motion with a several-dimensional time, Theory Prob. Appl., 8-4 (1963), 335-354.
- [5] A. Erdélyi and others: Higher transcendental functions I-III, McGraw-Hill Publ. New York (1953-1955).

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