# VARIATIONAL PROBLEMS WITH TWO PHASES AND THEIR FREE BOUNDARIES 

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ABSTRACT. The problem of minimizing $\int\left[\left.\nabla v\right|^{2}+q^{2}(x) \lambda^{2}(v)\right] d x$ in an appropriate class of functions $v$ is considered. Here $q(x) \neq 0$ and $\lambda^{2}(v)=\lambda_{1}^{2}$ if $v<0,=\lambda_{2}^{2}$ if $v>0$. Any minimizer $u$ is harmonic in $\{u \neq 0\}$ and $|\nabla u|^{2}$ has a jump

$$
q^{2}(x)\left(\lambda_{1}^{2}-\lambda_{2}^{2}\right)
$$

across the free boundary $\{u \neq 0\}$. Regularity and various properties are established for the minimizer $u$ and for the free boundary.

Introduction. In this paper we consider the problem of minimizing

$$
J(v)=\int_{\Omega}\left[|\nabla v|^{2}+q^{2}(x) \lambda^{2}(v)\right] d x, \quad v \in K
$$

where $q^{2}(x) \neq 0$,

$$
\lambda^{2}(v)= \begin{cases}\lambda_{1}^{2} & \text { if } v<0 \\ \lambda_{2}^{2} & \text { if } v>0\end{cases}
$$

and $\lambda^{2}(v)$ is lower semicontinuous at $v=0$; it is assumed that $\lambda_{i}^{2}>0$ and $\Lambda=\lambda_{1}^{2}-$ $\lambda_{2}^{2} \neq 0$. The class $K$ consists of all functions $v$ in $L_{\mathrm{loc}}^{1}(\Omega)$, with $\nabla v \in L^{2}(\Omega)$ such that $v=u^{0}$ on a given open subset $S$ of $\partial \Omega$, and $\Omega$ is a domain in $R^{n}$.

The analogous problem for functions in $K^{+}=\{v \in K, v \geqslant 0$ a.e. $\}$ was studied in [1]; in that paper it was proved that any (local) minimizer of $J(v)$ in $K^{+}$is Lipschitz continuous and, if $n=2$, the free boundary $\partial\{u>0\}$ is analytic if $q(x)$ is analytic.

The present variational problem is motivated by applications to the flow of two liquids in models of jets and cavities; these applications will be studied in other forthcoming papers [5,6]. The present work is aimed at extending results of [1]. In particular, we shall establish nondegeneracy theorems, the Lipschitz continuity of the solution, and some properties of the free boundary; for $n=2$ the free boundary is proved to be continuously differentiable.

A new and rather powerful tool introduced in this paper is the monotonicity formula (Lemma 5.1) asserting that, for a minimizer $u$, if $u\left(x_{0}\right)=0$ then

$$
r^{-4} \int_{B_{r}\left(x_{0}\right)} \rho^{2-n}\left|\nabla u^{+}\right|^{2} d x \cdot \int_{B_{r}\left(x_{0}\right)} \rho^{2-n}\left|\nabla u^{-}\right|^{2} d x \nearrow \quad \text { if } r \nearrow
$$

[^0]This is used in establishing Lipschitz continuity and in identifying blow-up limits. The differentiability of the free boundary for $n=2$ also involves a new set of ideas, exploiting among other things, the monotonicity formula.

1. Existence. Let $\Omega$ be a domain in $R^{n}$ with boundary $\partial \Omega$ which is locally a Lipschitz graph. Let $S$ be a nonempty open subset of $\partial \Omega$ and let $u^{0}$ be a given function in $L_{\mathrm{loc}}^{1}(\Omega)$ with $\nabla u^{0} \in L^{2}(\Omega)$. Let $q(x)$ be a strictly positive uniformly Lipschitz continuous function in compact subsets of $\bar{\Omega}$, and let $\lambda(u)$ be the function

$$
\lambda(u)= \begin{cases}\lambda_{1} & \text { if } u<0  \tag{1.1}\\ \lambda_{2} & \text { if } u>0\end{cases}
$$

where $\lambda_{1}, \lambda_{2} \geqslant 0$, and define $\lambda(0)$ such that

$$
\begin{equation*}
0 \leqslant \lambda(0) \leqslant \min \left\{\lambda_{1}, \lambda_{2}\right\} \tag{1.2}
\end{equation*}
$$

We assume that

$$
\begin{equation*}
\Lambda=\lambda_{1}^{2}-\lambda_{2}^{2} \neq 0 \tag{1.3}
\end{equation*}
$$

Finally, set $Q(u, x)=q(x) \lambda(u)$.
We introduce the convex set

$$
K=\left\{v \in L_{\mathrm{loc}}^{1}(\Omega), \nabla v \in L^{2}(\Omega), v=u^{0} \text { on } S\right\}
$$

and the functional

$$
J(v)=\int_{\Omega}\left(|\nabla v|^{2}+Q^{2}(v, x)\right) d x, \quad v \in K
$$

Problem (J). Find $u \in K$ such that $J(u)=\min _{v \in K} J(v)$.
Theorem 1.1. If $J\left(u^{0}\right)<\infty$ then there exists a solution of Problem (J).
Proof. Take a minimizing sequence $u_{k}$. Then the $\nabla u_{k}$ are uniformly bounded in $L^{2}(\Omega)$. Since $u_{k}-u^{0}=0$ on $S, S$ open and nonempty, we can estimate $u_{k}-u^{0}$ in $L^{2}\left(\Omega \cap B_{R}\right)$ for any ball $B_{R}=\{|x|<R\}$ and deduce that, for a subsequence,

$$
\begin{aligned}
\nabla u_{k} & \rightarrow \nabla u \quad \text { weakly in } L_{\mathrm{loc}}^{2}(\Omega), \\
u_{k} & \rightarrow u \quad \text { a.e. in } \Omega \\
Q^{2}\left(u_{k}, x\right) & \rightarrow \gamma \quad \text { weakly star in } L_{\mathrm{loc}}^{\infty}(\Omega),
\end{aligned}
$$

where $\gamma=Q^{2}(u, x)$ if $u \neq 0$, and $\gamma \geqslant Q^{2}(u, x)$ if $u=0$ (by (1.2)). Hence,

$$
\begin{aligned}
\int_{\Omega \cap B_{R}}\left(|\nabla u|^{2}+Q^{2}(u, x)\right) & \leqslant \liminf _{k \rightarrow \infty} \int_{\Omega \cap B_{R}}\left|\nabla u_{k}\right|^{2}+\lim _{k \rightarrow \infty} \int_{\Omega \cap B_{R}} Q^{2}\left(u_{k}, x\right) \\
& \leqslant \liminf _{k \rightarrow \infty} J\left(u_{k}\right) .
\end{aligned}
$$

Letting $R \rightarrow \infty$ we see that $u$ is an absolute minimum for $J$.
2. Continuity, subharmonicity and the free boundary condition. We denote a solution of Problem (J) by $u$.

Thegrem 2.1. For any compact subset $D$ of $\Omega$ there exists a constant $C$ such that

$$
|u(x)-u(y)| \leqslant C|x-y| \log (1 /|x-y|)
$$

if $x, y \in D .|x-y|<\frac{1}{2}$.
Proof. Let $B_{r}$ be any ball of radius $r$ in $D$ and denote by $v_{r}$ the solution of

$$
\begin{equation*}
\nabla v_{r}=0 \quad \text { in } B_{r}, \quad v_{r}=u \quad \text { on } \partial B_{r} . \tag{2.1}
\end{equation*}
$$

Then, by the minimality of $u$,

$$
\int_{B_{r}}\left(|\nabla u|^{2}+Q^{2}(u, x)\right) \leqslant \int_{B_{r}}\left(\left|\nabla v_{r}\right|^{2}+Q^{2}\left(v_{r}, x\right)\right) .
$$

It follows that $\int_{B_{r}}\left(|\nabla u|^{2}-\left|\nabla v_{r}\right|^{2}\right) \leqslant C r^{n}$. But the left-hand side is equal to

$$
\int_{B_{r}} \nabla\left(u-v_{r}\right) \cdot \nabla\left(u+v_{r}\right)=\int_{B_{r}}\left|\nabla\left(u-v_{r}\right)\right|^{2}+2 \int_{B_{r}} \nabla\left(u-v_{r}\right) \cdot \nabla v_{r}
$$

and the last integral vanishes, by (2.1). Consequently, $\int_{B,}\left|\nabla\left(u-v_{r}\right)\right|^{2} \leqslant C r^{n}$.
Proceeding as in [11, Theorem, 5.3.6], one can establish that

$$
\int_{B_{r}}\left|\nabla\left(u-v_{r}\right)\right|^{2} \leqslant C(R) r^{n}(\log R / r+1) \quad \text { if } 0<r<R
$$

so that

$$
\int_{B_{r}}|\nabla u|^{2} \leqslant C(R) r^{n}\left(\log \frac{R}{r}+1\right)
$$

from which the assertion follows as in [11, Theorem 3.5.2].
Theorem 2.2. The function $u$ is harmonic in $\{u \neq 0\}$.
Proof. For any $\zeta \in C_{0}^{1}(\Omega \backslash\{u=0\}), u \pm \varepsilon \zeta$ is in $K$ for any $\varepsilon>0$. Hence,

$$
0=\lim _{\varepsilon \downharpoonright 0} \frac{1}{2 \varepsilon}(J(u+\varepsilon \zeta)-J(\zeta))=\int_{\Omega} \nabla \zeta \cdot \nabla u .
$$

Theorem 2.3. If $\lambda(0)=\lambda_{1}$ and $\Lambda<0\left(\lambda(0)=\lambda_{2}\right.$ and $\left.\Lambda>0\right)$ then $u$ is subharmonic ( superharmonic) in $\Omega$.

Proof. Defining $v$ by (2.1), $B_{r} \subset \Omega$, we have $J(u) \leqslant J(\min (u, v)$ ), which gives, if $\lambda(0)=\lambda_{1}$.

$$
\begin{aligned}
I & \equiv \int_{B_{r}}\left[|\nabla u|^{2}-|\nabla \min (u, v)|^{2}\right] \leqslant \int_{B_{r}}\left[Q^{2}(\min (u, v), x)-Q^{2}(u, x)\right] \\
& =\int_{B_{r} \cap\{u>v\}}\left[Q^{2}(v, x)-Q^{2}(u, x)\right]=\int_{B_{r} \cap\{u \geqslant 0>v\}} \Lambda q^{2}(x) .
\end{aligned}
$$

But

$$
\begin{aligned}
I & =\int_{B_{r}} \nabla \max (u-v, 0) \cdot \nabla(u+v) \\
& =\int_{B_{r}} \nabla \max (u-v, 0) \cdot \nabla(u-v)+2 \int_{B_{r}} \nabla \max (u-v, 0) \cdot \nabla v \\
& =\int_{B_{r}}|\nabla \max (u-v, 0)|^{2} .
\end{aligned}
$$

Hence, $\Lambda<0$ implies $u \leqslant v$, i.e., $u$ is subharmonic. Similarly, if $\lambda(0)=\lambda_{2}$ and $\Lambda>0$, then $u$ is superharmonic.

Definition 2.1. The set $\Gamma=\partial\{u>0\} \cup \partial\{u<0\}$ is called the free boundary.
The next theorem shows that $u$ satisfies, in a generalized sense, the equation

$$
|\nabla u|^{2}-\left|\nabla u^{+}\right|^{2}=\Lambda q^{2}(x) \quad \text { on } \Gamma,
$$

provided the set $\{u=0\}$ has zero measure.
Theorem 2.4. Suppose meas $\{u=0\}=0$. Then, for any $\eta \in C_{0}^{1}\left(\Omega, R^{n}\right)$,

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \int_{\partial\{u<-\varepsilon\}}\left(|\nabla u|^{2}-\lambda_{1}^{2} q^{2}(x)\right) \eta \cdot \nu+\lim _{\delta \downarrow 0} \int_{\partial(u>\delta\}}\left(|\nabla u|^{2}-\lambda_{2}^{2} q^{2}(x)\right) \eta \cdot \nu=0 \tag{2.2}
\end{equation*}
$$

where $\nu$ is the outward normal.
Proof. Let $\tau_{\varepsilon}(x)=x+\varepsilon \eta(x), \varepsilon \neq 0$, and define $u_{\varepsilon} \in K$ by $u_{\varepsilon}\left(\tau_{\varepsilon} x\right)=u(x)$. Then

$$
\begin{aligned}
0 \leqslant & J\left(u_{\varepsilon}\right)-J(u) \\
= & \int_{\Omega}\left\{\left[\left|\nabla u\left(D \tau_{\varepsilon}\right)^{-1}\right|^{2}+Q^{2}\left(u, \tau_{\varepsilon}(x)\right)\right] \operatorname{det}\left(D \tau_{\varepsilon}\right)-\left(|\nabla u|^{2}+Q^{2}(u, x)\right)\right\} \\
= & \varepsilon \int_{\Omega}\left[|\nabla u|^{2}+Q^{2}(u, x)\right] \nabla \cdot \eta \\
& +\varepsilon \int_{\Omega}\left[-2 \nabla u D \eta \nabla u+\nabla_{x} Q^{2}(u, x) \cdot \eta\right]+O(\varepsilon)
\end{aligned}
$$

The linear term in $\varepsilon$ must vanish, giving (since $\Delta u=0$ in $\{u \neq 0\}$ )

$$
\begin{aligned}
0= & \lim _{\varepsilon \downarrow 0, \delta \downarrow 0} \int_{\Omega \backslash\{-\varepsilon<u<\delta\}} \nabla \cdot\left[\left(|\nabla u|^{2}+Q^{2}(u, x)\right) \eta-2 \eta \cdot \nabla u \nabla u\right] \\
= & \lim _{\varepsilon \downarrow 0} \int_{\partial\{u<-\varepsilon\}}\left[\left(|\nabla u|^{2}+Q^{2}(u, x)\right) \eta-2 \eta \cdot \nabla u \nabla u\right] \cdot \nu \\
& +\lim _{\delta \downarrow 0} \int_{\partial\{u>\delta\}}\left[\left(|\nabla u|^{2}+Q^{2}(u, x)\right) \eta-2 \eta \cdot \nabla u \nabla u\right] \cdot \nu \\
= & \lim _{\varepsilon \downarrow 0} \int_{\partial\{u<-\varepsilon\}}\left[\lambda_{1}^{2} q^{2}(x)-|\nabla u|^{2}\right] \eta \cdot \nu \\
& +\lim _{\delta \downharpoonright 0} \int_{\partial\{u>\delta\}}\left[\lambda_{2}^{2} q^{2}(x)-|\nabla u|^{2}\right] \eta \cdot \nu .
\end{aligned}
$$

Remark 2.1. If meas $\{u=0\}>0$ and if $\{u=0\}$ is a limit of increasing open sets $D_{\rho}(\rho \downarrow 0)$, then on the left-hand side of (2.2) there appears the additional term

$$
\lim _{\rho!0} \int_{\partial D_{\rho}}\left(|\nabla u|^{2}-\lambda^{2}(0) q^{2}(x)\right) \eta \cdot \nu
$$

3. Nondegeneracy. For any function $v$ and a ball $B_{r}=B_{r}\left(x^{0}\right)$ with center $x^{0}$ and radius $r$, we set

$$
f_{\partial B_{r}} v=\frac{1}{\left|\partial B_{r}\right|} \int_{\partial B_{r}} v, \quad\left|\partial B_{r}\right|=\text { surface area of } \partial B_{r} .
$$

Let

$$
\begin{equation*}
0 \leqslant q_{1} \leqslant q(x) \leqslant q_{2}<\infty, \quad|\Lambda| \geqslant I_{0}>0 . \tag{3.1}
\end{equation*}
$$

Theorem 3.1. Suppose $\Lambda<0$. For any $0<\kappa<1$ there is a positive constant $c$ depending only on $\kappa$ and $q_{1}^{2} l_{0}$ such that if $B_{r} \subset \Omega$ and $\frac{1}{r} f_{\partial B_{r}} u^{+}<c$, then $u^{+}=0$ in $B_{\kappa r}$.

Proof. Set $\gamma=\frac{1}{r} f_{\partial B_{r}} u^{+}$. The idea of this proof is to replace $u$ in $B_{r}$ by a function $v$ satisfying

$$
\begin{aligned}
v=0 & \text { on } \partial B_{r}, \\
v=u & \text { in } B_{r} \cap\{u \leqslant 0\}, \\
v=0 & \text { in } B_{\kappa r} \cap\{u>0\}, \\
\Delta v=0 & \text { in }\left(B_{r} \backslash B_{\kappa r}\right) \cap\{u>0\}
\end{aligned}
$$

and show that $J(v)<J(u)$ if $\gamma$ is sufficiently small.
For almost any $\varepsilon>0$ the surface $\{u=\varepsilon\}$ is smooth. Choose any such small $\varepsilon$ and consider the function $v_{\varepsilon}$ satisfying

$$
\begin{aligned}
& v_{\varepsilon}=u \quad \text { on } \partial B_{r}, \\
& v_{\varepsilon}=u \text { in } B_{r} \cap\{u<\varepsilon\}, \\
& v_{\varepsilon}=\varepsilon \text { in } B_{\kappa r} \cap\{u>\varepsilon\}, \\
& \Delta v_{\varepsilon}=0 \quad \text { in } D_{\varepsilon}^{+} \equiv\left(B_{r} \backslash B_{\kappa r}\right) \cap\{u>\varepsilon\} .
\end{aligned}
$$

The function $v_{\varepsilon}$ can be obtained by minimizing the Dirichlet integral over $B_{r}$ subject to the above constraints. Also $v_{\varepsilon}$ is continuous at $\{u=\varepsilon\} \cap\left(B_{r} \backslash \bar{B}_{\kappa r}\right)$ and $\min (u, 0)$ $\leqslant v_{\varepsilon} \leqslant u$. Since $\nabla v_{\varepsilon}$ is bounded in $L^{2}\left(B_{r}\right)$, the limit $v=\lim _{\varepsilon \rightarrow 0} v_{\varepsilon}$ exists and $\min (u, 0) \leqslant v \leqslant u$; hence $v$ is continuous in $B_{r}$ and has the desired properties.

We obtain

$$
\begin{aligned}
\int_{B_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right) & \leqslant \int_{B_{r}} q^{2}\left(\lambda^{2}(v)-\lambda^{2}(u)\right) \\
& \leqslant \int_{B_{k r} \cap\{u>0\}} \wedge q^{2} .
\end{aligned}
$$

Hence, setting $D^{+}=\left(B_{r} \backslash B_{\kappa r}\right) \cap\{u>0\}$,

$$
\begin{aligned}
& \int_{B_{\mathrm{xx}} \cap\{u>0\}}\left(|\nabla u|^{2}-\Lambda q^{2}\right) \leqslant \int_{D^{+}}\left(|\nabla v|^{2}-|\nabla u|^{2}\right) \\
&= \int_{D^{+}} \nabla(v-u) \cdot \nabla(u-v+2 v) \\
& \leqslant 2 \int_{D^{+}} \nabla v \cdot \nabla(v-u) \\
& \quad \leqslant \liminf _{F \rightarrow 0} 2 \int_{D^{+}} \nabla v_{\varepsilon} \cdot \nabla\left(v_{\varepsilon}-u\right) \\
&=\liminf _{\varepsilon \rightarrow 0} 2 \int_{\partial B_{k r} \cap\{u>\varepsilon\}}(u-\varepsilon)\left|\nabla v_{\varepsilon}\right| \equiv M
\end{aligned}
$$

where in the last formula we have used the integration by parts

$$
\begin{equation*}
\int_{D_{\varepsilon}^{+}} \nabla v_{\varepsilon} \cdot \nabla\left(v_{\varepsilon}-u\right)=\int_{\partial B_{k r}}(u-\varepsilon)\left|\frac{\partial v_{\varepsilon}}{\partial \nu}\right| ; \tag{3.2}
\end{equation*}
$$

notice that $\partial v_{\epsilon} / \partial \nu \leqslant 0$ on $\partial B_{\kappa r}$. Since $\partial B_{\kappa r}$ and $\partial\{u>\varepsilon\}$ form a corner at their intersection, one has to justify (3.2) by approximation. We shall do this later.

To estimate $M$ we introduce the function $w$ :

$$
\begin{aligned}
\Delta w & =0 & & \text { in } B_{r} \backslash B_{\kappa r}, \\
w & =u & & \text { on } \partial B_{r} \cap\{u>\varepsilon\}, \\
w & =\varepsilon & & \text { elsewhere on } \partial\left(B_{r} \backslash B_{\kappa r}\right) .
\end{aligned}
$$

Clearly $w \geqslant v_{f}$ and thus $|\nabla w| \geqslant\left|\nabla v_{\varepsilon}\right|$ on $\partial B_{\kappa r} \cap\{u>\varepsilon\}$. Since

$$
|\nabla w| \leqslant \frac{C}{r} f_{\partial B_{r}}(u-\varepsilon)^{+} \leqslant C \gamma \quad \text { on } \partial B_{\kappa r},
$$

we get

$$
\begin{equation*}
\left|\nabla v_{\varepsilon}\right| \leqslant C \gamma \quad \text { on } \partial B_{\kappa r} \cap\{u>\varepsilon\} . \tag{3.3}
\end{equation*}
$$

Hence

$$
\begin{aligned}
M \leqslant & C \gamma \int_{\partial B_{k r}} u^{+} \leqslant C \gamma\left(\int_{B_{k r}}\left|\nabla u^{+}\right|+\frac{1}{r} \int_{B_{k r}} u^{+}\right) \\
\leqslant & \frac{C \gamma}{|\Lambda|^{1 / 2} q_{1}} \int_{B_{k r}}\left(\left|\nabla u^{+}\right|^{2}+|\Lambda| q_{1}^{2} I_{\left\{u^{\prime}-0 \mid\right.}\right) \\
& +\frac{C \gamma}{|\Lambda| q_{1}^{2} r}\left(\sup _{B_{k r}} u^{+}\right) \int_{B_{k r}}|\Lambda| q_{1}^{2} I_{\left\{u^{+}>0\right\}} .
\end{aligned}
$$

Since $u$ is harmonic in $\{u>0\}, u^{+}$is subharmonic in $\Omega$; therefore $\sup _{B_{x r}} u^{+} \leqslant C \gamma r$. Hence

$$
\begin{aligned}
\int_{B_{x r} \cap\{u>0\}}\left(|\nabla u|^{2}-\Lambda q_{1}^{2}\right) \leqslant & \frac{C \gamma}{|\Lambda|^{1 / 2} q_{1}}\left(1+\frac{\gamma}{|\Lambda|^{1 / 2} q_{1}}\right) \\
& \cdot \int_{B_{\mathrm{kr}} \cap\{u>0\}}\left(|\nabla u|^{2}-\Lambda q_{1}^{2}\right) .
\end{aligned}
$$

Hence if $\gamma /\left(|\Lambda|^{1 / 2} q_{1}\right)$ is small enough then $u \leqslant 0$ in $B_{\kappa r}$.
It remains to justify (3.2). Approximate $D_{\varepsilon}^{+}$by domains $D_{m}$ by changing $D_{\varepsilon}^{+}$near $\partial B_{k r} \cap \partial\{u>\varepsilon\}$ so as to form a smooth boundary there. Denote the corresponding $v_{\varepsilon}$ by $v_{\varepsilon m}\left(v_{\varepsilon m}=\varepsilon\right.$ on the modified boundary $\partial D_{m}$ near $\left.\partial B_{\kappa r} \cap \partial\{u>\varepsilon\}\right)$. Then,

$$
\begin{align*}
& D v_{\varepsilon m} \rightarrow D v_{\varepsilon} \quad \text { on } \partial B_{\kappa r}, \text { away from } \partial\{u>\varepsilon\}, \\
& \left|D v_{\varepsilon m}\right| \leqslant C \quad \text { on } \partial D_{m}, \text { away from } \partial B_{r} \tag{3.4}
\end{align*}
$$

(by (3.3)). Since (3.2) holds for $v_{\varepsilon}=v_{\varepsilon m}$, taking $m \rightarrow \infty$ and using (3.4), the assertion (3.2) for $v_{F}$ follows.

Theorem 3.1 may be considered as a nondegeneracy theorem. It implies
Corollary 3.2. Suppose $\Lambda<0$. If $B_{r} \subset \Omega$ with center in the free boundary $\partial\left\{u>0\right.$; then $\frac{1}{r} f_{\partial B_{r}} u^{+} \geqslant c(c>0) ; c$ depends only on $q_{1}^{2} l_{0}$.

The analog of Theorem 3.1 and its corollary to the case $\Lambda>0$ are obvious.
Remark 3.1. If $\lambda^{2}(0)<\min \left\{\lambda_{1}^{2}, \lambda_{2}^{2}\right\}$ then the proof of Theorem 3.1 applies to both $u^{+}$and $u^{-}$. Consequently, if $B_{r} \subset \Omega$ with center in the free boundary $\partial\{u>0\}$ $(\partial\{u<0\})$, then

$$
\frac{1}{r} f_{\partial B_{r}} u^{+} \geqslant c \quad\left(\frac{1}{r} f_{\partial B_{r}} u^{-} \geqslant c\right)
$$

where $c$ is a positive constant depending only on $q_{1}^{2}\left\{\min \left(\lambda_{1}^{2}, \lambda_{2}^{2}\right)-\lambda^{2}(0)\right\}$.
4. Upper estimates on the averages. Let

$$
\begin{equation*}
\max \left\{\lambda_{1}^{2}, \lambda_{2}^{2}\right\} \leqslant l_{1} . \tag{4.1}
\end{equation*}
$$

Theorem 4.1. Assume that $\lambda(0)=\min \left(\lambda_{1}, \lambda_{2}\right)$. There exists a positive constant $C$ depending only on $q_{2}\left(\right.$ in (3.1)) and $l_{1}$ such that, if $B_{r} \subset \Omega$ with center in $\{u=0\}$, then

$$
\begin{equation*}
\frac{1}{r}\left|f_{\partial B_{r}} u\right| \leqslant C \tag{4.2}
\end{equation*}
$$

We shall prove the theorem in case $\Lambda<0$; the proof in case $\Lambda>0$ is similar. Since $\Lambda<0, \Delta u$ is a (positive) measure (by Theorem 2.3). In order to prove the theorem we first estimate the measure $\Delta u$.

Lemma 4.2. If $\Lambda<0$ and $B_{r} \subset \Omega$ then

$$
\begin{equation*}
\Delta u\left(B_{r / 2}\right) \leqslant C r^{n-1} . \tag{4.3}
\end{equation*}
$$

Proof. Defining $v$ as in (2.1) we have

$$
\int_{B_{r}}|\nabla u|^{2}-\int_{B_{r}}|\nabla v|^{2} \leqslant \int_{B_{r}}\left(Q^{2}(v, x)-Q^{2}(u, x)\right) \leqslant C r^{n} .
$$

The left-hand side is equal to

$$
\int_{B_{r}} \nabla(u-v) \cdot \nabla(u+v)=\int_{B_{r}} \nabla(u-v) \cdot \nabla u=\int_{B_{r}}(v-u) \Delta u
$$

where $\Delta u$ is a measure supported on $\{u=0\}$ (the continuity of $v-u$ is used in making sense out of the last integral); the last integral is equal to $\int_{B_{r}} v \Delta u$. Since $v \geqslant u$, also $v \Delta u \geqslant u \Delta u=0$, and thus

$$
\begin{equation*}
\int_{B_{r / 2} \cap\{u=0\}} v \Delta u \leqslant C r^{n} . \tag{4.4}
\end{equation*}
$$

We shall use the representation

$$
\begin{equation*}
u\left(x^{0}\right)=\int_{\partial B_{r}} P_{x^{v}}(y) u(y)-\int_{B_{r}} G_{x^{v}}(y) \Delta u(y) \tag{4.5}
\end{equation*}
$$

where $P_{x^{\prime \prime}}$ and $G_{x^{\prime \prime}}$ are Poisson's kernel and Green's function (in $B_{r}$ ), respectively. This formula can be justified by approximating $u$ by mollifiers $J_{\varepsilon} u$, applying the formula to $J_{\epsilon} u$ at $x^{0}$ and taking $\varepsilon \rightarrow 0$. If $x^{0} \in\{u=0\}$ then we obtain, from (4.5),

$$
\begin{equation*}
\int_{B_{r}} G_{x^{0}}(y) \Delta u(y)=\int_{\partial B_{r}} P_{x^{n}}(y) u(y) \tag{4.6}
\end{equation*}
$$

and the right-hand side is precisely $v\left(x^{0}\right)$. Thus we can rewrite (4.4) in the form

$$
\int_{B_{r / 2}}\left(\int_{B_{r}} G_{x}(y) \Delta u(y)\right) \Delta u(x) \leqslant C r^{n} .
$$

Noting that $G_{x}(y) \geqslant c r^{2-n}$ if $x, y \in B_{r / 2}(c>0)$ we obtain $c r^{2-n}\left(\Delta u\left(B_{r / 2}\right)\right)^{2} \leqslant C r^{n}$, and the assertion (4.3) follows.

Proof of Theorem 4.1. As before we take $\Lambda<0$. We may assume that the center of $B_{r}$ is in the origin. By (4.6),

$$
\begin{equation*}
\int_{\partial B_{r}} P_{0} u=\int_{B_{r}} G_{0}(y) \Delta u(y) . \tag{4.7}
\end{equation*}
$$

Suppose first that $\Delta u$ is smooth. Then

$$
I \equiv \int_{B_{r}} G_{0}(y) \Delta u(y)=\int_{0}^{r} G(s) h(s) d s
$$

with suitable functions $G$ and $h ; h(r)=r^{n-1} \int_{\partial B_{1}} \Delta u(r \xi) d H^{n-1}(\xi)$. By Lemma 4.2,

$$
\begin{equation*}
\int_{0}^{s} h(\tau) d \tau \leqslant C s^{n-1} \tag{4.8}
\end{equation*}
$$

Hence,

$$
I=\int_{0}^{r} G(s) \frac{d}{d s}\left(\int_{0}^{s} h(\tau) d \tau\right) d s=\left[G(s) \int_{0}^{s} h(\tau) d \tau\right]_{0}^{r}-\int_{0}^{r} G^{\prime}(s) \int_{0}^{s} h(\tau) d \tau d s .
$$

The expression in brackets vanishes at $s=r$ (since $G_{0}=0$ on $\partial B_{r}$ ) and at $s=0$ (by (4.8) and $G(s) \leqslant C s^{2-n}$ ). Hence,

$$
\begin{equation*}
\int_{B_{r}} G_{0}(y) \Delta u(y) \leqslant \int_{0}^{r} \frac{C}{s^{n-1}} C s^{n-1} d s \leqslant C r \tag{4.9}
\end{equation*}
$$

By using mollifiers $u_{\varepsilon}=u * \psi_{\varepsilon}$ we can establish the same estimate for the measure $\Delta u$. Here we use the estimate

$$
\begin{aligned}
\int_{\left.B_{r}\left(x_{0}\right)\right\}} \Delta u_{f}(x) & =\int_{B_{r}\left(x_{0}\right)} \int_{\left\{\left|\left.\right|_{1}\right|<f\right\}} \Delta u(x-y) \psi_{f}(y) d y \\
& =\int_{\{|| |<f\}} \int_{B_{r}\left(x_{0}\right)} \Delta u(x-y) \psi_{f}(y) d y \\
& =\int_{\{|| |<f\}} \Delta u\left(B_{r}\left(x_{0}-y\right)\right) \psi_{f}(y) d y \leqslant C r^{n-1} .
\end{aligned}
$$

From (4.7) and (4.9) we see that $\frac{1}{1} f_{\partial B} u \leqslant C$. Since $u(0)=0$ and $u$ is subharmonic, the last integral is actually positive and therefore (4.2) follows.

## 5. Lipschitz continuity.

Lemma 5.1. Let $u$ be any function in $C^{0}\left(B_{R}\right) \cap H^{1.2}\left(B_{R}\right)$, where $B_{r}$ is a ball with radius $r$ and center $x^{0}, u\left(x^{0}\right)=0$, and $u$ is harmonic in $B_{R} \backslash\{u=0\}$. Set

$$
\phi(r)=\frac{1}{r^{2}} \int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2} d x \cdot \frac{1}{r^{2}} \int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2} d x
$$

where $\rho=\left|x-x^{0}\right|$. Then $\phi(r)<\infty$ and $\phi(r)$ is increasing in $r, r \in(0, R)$.
We shall refer to this result as the monotonicity lemma.
Proof. Set $S_{r}=\partial B_{r}$. We first assume that

$$
\begin{equation*}
\min _{S_{r}} u<0<\max _{S_{r}} u \quad \text { for all } r \in(0, R) \tag{5.1}
\end{equation*}
$$

Notice that the distribution $\Delta u^{+}$is a measure. Denote by $v_{m}$ mollifiers of $u^{+}$. Then $\Delta v_{m}^{2}=2\left|\nabla v_{m}\right|^{2}+2 v_{m} \Delta v_{m} \geqslant 2\left|\nabla v_{m}\right|^{2}$, so that

$$
2 \int_{B_{r} \backslash B_{r}}\left|\nabla v_{m}\right|^{2} \rho^{2-n} \leqslant \int_{B_{r} \backslash B_{r}} \Delta\left(v_{m}^{2}\right) \rho^{2-n}=r^{2-n} \int_{S_{r}} \frac{\partial}{\partial r} v_{m}^{2}+(n-2) r^{1-n} \int_{S_{r}} v_{m}^{2}-I_{F}
$$

where

$$
I_{\varepsilon}=\varepsilon^{2-n} \int_{S_{\varepsilon}} \frac{\partial}{\partial r} v_{m}^{2}+(n-2) \varepsilon^{1-n} \int_{S_{r}} v_{m}^{2} .
$$

Since $\left|D v_{m}\right|$ is bounded, $I_{\varepsilon} \rightarrow(n-2)\left|S_{1}\right| v_{m}^{2}(0)$ as $\varepsilon \rightarrow 0$. Hence,

$$
2 \int_{B_{r} \backslash B_{r}}\left|\nabla v_{m}\right|^{2} \rho^{2-n} \leqslant r^{2-n} \int_{S_{r}} \frac{\partial}{\partial r} v_{m}^{2}+(n-2) r^{1-n} \int_{S_{r}} v_{m}^{2} .
$$

Integrating with respect to $r, r_{0}<r<r_{0}+\delta$, and dividing by $\delta$, and then letting $m \rightarrow \infty$, we obtain

$$
\begin{aligned}
\frac{2}{\delta} \int_{r_{0}}^{r_{0}+\delta} d r \int_{B_{r} \backslash B_{r}}\left|\nabla u^{+}\right|^{2} \rho^{2-n} \leqslant & \frac{1}{\delta} \int_{r_{0}}^{r_{0}+\delta} r^{2-n} d r \int_{S_{r}} 2 u^{+} u_{r}^{+} \\
& +\frac{n-2}{\delta} \int_{r_{0}}^{r_{0}+\delta} r^{1-n} d r \int_{S_{r}}\left(u^{+}\right)^{2}
\end{aligned}
$$

Taking $\delta \rightarrow 0$ we obtain for a.a. $r_{0}$

$$
2 \int_{B_{r_{0}, ~} B} \mid \nabla u^{+} \Gamma^{2} \rho^{2-n} \leqslant r_{0}^{2-n} \int_{S_{r_{0}}} 2 u^{+} u_{r}^{+}+(n-2) r_{0}^{1-n} \int_{S_{r_{0}}}\left(u^{+}\right)^{2}
$$

Hence, for a.a. $r$.

$$
\begin{equation*}
2 \int_{B_{r}}\left|\nabla u^{+}\right|^{2} \rho^{2-n} \leqslant r^{2-n} \int_{S_{r}} 2 u^{+} u_{r}^{+}+(n-2) r^{1-n} \int_{S_{r}}\left(u^{+}\right)^{2} . \tag{5.2}
\end{equation*}
$$

Since a similar inequality holds for $u^{-}$, it follows that $\psi(r)$ is finite.
Since $r \rightarrow \int_{S_{r}}\left|\nabla u^{+}\right|^{2} \rho^{2-n}$ is in $L^{1}(0, R)$, we have

$$
\frac{d}{d r} \int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2}=\int_{S_{r}} r^{2-n}\left|\nabla u^{+}\right|^{2} \quad \text { a.e. }
$$

It follows that a.e.

$$
\begin{align*}
\phi^{\prime}(r)= & -\frac{4}{r^{5}} \int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2} \cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}+\frac{1}{r^{4}} \int_{S_{r}} r^{2-n}\left|\nabla u^{+}\right|^{2}  \tag{5.3}\\
& \left.\cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}+\frac{1}{r^{4}} \int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2} \cdot \int_{S_{r}} r^{2-n} \right\rvert\, \nabla u^{-2} .
\end{align*}
$$

We shall prove that $\phi^{\prime}(r) \geqslant 0$ a.e. in $(0, R)$. By scaling, we may assume that $r=1$.
Denote by $\nabla_{\theta} v$ the gradient of a function $v$ on $S_{1}$. Denote by $\Gamma_{1}$ the support of $u^{+}$ on $S_{1}$, and by $\Gamma_{2}$ the support of $u^{-}$on $S_{1}$. By assumption,

$$
\begin{equation*}
\operatorname{meas}\left(\Gamma_{i}\right) \neq 0 \quad \text { for } i=1,2, \ldots \tag{5.4}
\end{equation*}
$$

We introduce the constants

$$
\frac{1}{\alpha_{i}}=\inf _{v \in H_{l}^{\prime 2}, 2\left(\Gamma_{,}\right)} \frac{\int_{\Gamma_{,}}\left|\nabla_{\theta} v\right|^{2}}{\int_{\Gamma_{i}} v^{2}}
$$

For any $0<\beta_{1}<1$ we can write

$$
\begin{aligned}
\int_{S_{1}}\left(\left(u_{r}^{+}\right)^{2}+\beta_{1}^{2}\left|\nabla_{\theta} u^{+}\right|^{2}\right) & \geqslant 2\left\{\int_{S_{1}}\left(u_{r}^{+}\right)^{2} \cdot \int_{S_{1}} \beta_{1}^{2}\left|\nabla_{\theta} u^{+}\right|^{2}\right\}^{1 / 2} \\
& \geqslant 2 \frac{\beta_{1}}{\sqrt{\alpha_{1}}}\left\{\int_{S_{1}}\left(u_{r}^{+}\right)^{2} \cdot \int_{S_{1}}\left(u^{+}\right)^{2}\right\}^{1 / 2} \geqslant \frac{2 \beta_{1}}{\sqrt{\alpha_{1}}} \int_{S_{1}}\left|u^{+} u_{r}^{+}\right|
\end{aligned}
$$

and

$$
\int_{S_{1}}\left(1-\beta_{1}^{2}\right)\left|\nabla_{\theta} u^{+}\right|^{2} \geqslant \frac{1-\beta_{1}^{2}}{\alpha_{1}} \int_{S_{1}}\left(u^{+}\right)^{2}
$$

Choosing

$$
\begin{equation*}
\frac{1-\beta_{i}^{2}}{\alpha_{i}}=(n-2) \frac{\beta_{i}}{\sqrt{\alpha_{i}}} \tag{5.5}
\end{equation*}
$$

we find that

$$
\begin{equation*}
\int_{S_{1}}\left|\nabla u^{+}\right|^{2} \geqslant \frac{\beta_{1}}{\sqrt{\alpha_{1}}}\left\{\int_{S_{1}} 2\left|u^{+} u_{r}^{+}\right|+(n-2) \int_{S_{1}}\left(u^{+}\right)^{2}\right\} . \tag{5.6}
\end{equation*}
$$

The relations (5.2) and (5.6) hold also for $u^{-}$. Comparing with (5.3) we see that $\phi^{\prime}(r) \geqslant 0$ provided

$$
\begin{equation*}
\frac{\beta_{1}}{\sqrt{\alpha_{1}}}+\frac{\beta_{2}}{\sqrt{\alpha_{2}}} \geqslant 2 \tag{5.7}
\end{equation*}
$$

We easily compute that the $\beta_{i}$ satisfy (5.5) if

$$
\frac{\beta_{i}}{\sqrt{\alpha_{i}}}=\frac{1}{2}\left\{\left[(n-2)^{2}+\frac{4}{\alpha_{i}}\right]^{1 / 2}-(n-2)\right\} .
$$

If $\gamma_{i}$ is defined by

$$
\begin{equation*}
\gamma_{i}\left(\gamma_{i}+n-2\right)=1 / \alpha_{i}, \quad \gamma_{i}>0 \tag{5.8}
\end{equation*}
$$

then we obtain

$$
\begin{equation*}
\frac{\beta_{i}}{\sqrt{\alpha_{i}}}=\gamma_{i} \tag{5.9}
\end{equation*}
$$

The set function $\gamma_{1}$ as a function of $\Gamma_{1}$ was studied by Sperner [12] and by Friedland and Hayman [8]. In [12] it is proved that $\gamma_{1}(E) \geqslant \gamma_{1}\left(E^{*}\right)$ where $E$, $E^{*} \subset S_{1}$ provided $E^{*}$ is a spherical cap having the same ( $n-1$ )-dimensional Hausdorff measure as $E$. In [8] it is proved that $\gamma_{1}(E) \geqslant \psi(s)$ where $s=$ meas $(E) / \operatorname{meas}\left(S_{1}\right)$, and $\psi(s)$ is convex and decreasing:

$$
\psi(s)= \begin{cases}\frac{1}{2} \log \frac{1}{4 s}+\frac{3}{2} & \text { if } s<\frac{1}{4} \\ 2(1-s) & \text { if } \frac{1}{4}<s<1\end{cases}
$$

Setting $s_{1}=\operatorname{meas}\left(\Gamma_{i}\right) / \operatorname{meas}\left(S_{1}\right)$, we then have

$$
\gamma_{1}+\gamma_{2} \geqslant \psi\left(s_{1}\right)+\psi\left(s_{2}\right) \geqslant 2 \psi\left[\left(s_{1}+s_{2}\right) / 2\right] \geqslant 2 \psi(1 / 2)=2
$$

in view of (5.9), this completes the proof of (5.7), provided (5.1) is satisfied.
If (5.1) is not satisfied, let $R_{0}$ be the smallest value of $r$ for which at least one of the inequalities in (5.1) is invalid. Suppose for definiteness that $\min _{S_{R_{0}}} u \geqslant 0$. Then
$u^{-}$is harmonic in $D=B_{R_{0}} \cap\{u<0\}$, vanishing on $\partial D$; hence $u^{-}=0$ in $D$, which gives $\phi(r)=0$ if $0<r \leqslant R_{0}$. Since $\phi^{\prime}(r) \geqslant 0$ for a.a. $R_{0}<r<R$ (by the previous proof), the proof of the lemma is complete.

We shall now use Theorem 4.1 in order to establish Lipschitz continuity for any minimizer $u$.

Lemma 5.2. Assume that $\lambda(0)=\min \left(\lambda_{1}, \lambda_{2}\right)$. Then for any domain $D \Subset \Omega$ there exists a positive constant $C$ such that if $B_{r} \subset D$ with center in $\{u=0\}$ then

$$
\begin{equation*}
\frac{1}{r} f_{\partial B_{r}}|u| \leqslant C . \tag{5.10}
\end{equation*}
$$

Proof. By Green's formula $(0<\alpha<1)$

$$
\begin{align*}
f_{\partial B_{r}} r^{\alpha} u^{-} & =\int_{B_{r}} G_{0} \Delta\left(\rho^{\alpha} u^{-}\right)=-\int_{B_{r}} \nabla G_{0} \cdot \nabla\left(\rho^{\alpha} u^{-}\right)  \tag{5.11}\\
& =-\int_{B_{r}} \rho^{\alpha} \nabla G_{0} \cdot \nabla u^{-}+c_{n} \int_{B_{r}} \rho^{1-n} \alpha \rho^{\alpha-1} u^{-} \equiv J_{1}+J_{2},
\end{align*}
$$

and $G_{0}(\rho)=c \rho^{2-n}, c>0$ (we take for definiteness $n \geqslant 3$ ). Clearly,

$$
\begin{aligned}
\left|J_{1}\right| & \leqslant C\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}\right)^{1 / 2}\left(\int_{B_{r}} \rho^{2 \alpha-n}\right)^{1 / 2} \\
& \leqslant C r^{\alpha}\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}\right)^{1 / 2} .
\end{aligned}
$$

Introducing the function $\phi_{\varepsilon}(r)=\left(r^{\varepsilon} / r\right) \int_{\partial B_{r}} u^{-}(0<\varepsilon<\alpha)$ we also have

$$
J_{2} \leqslant c_{n} \alpha \int_{0}^{r} \rho^{\alpha-\varepsilon} \phi_{\varepsilon}(\rho) \leqslant \frac{c_{n} \alpha}{1+\alpha-\varepsilon} r^{1+\alpha-\varepsilon} \sup _{\rho \leqslant r} \phi_{\varepsilon}(\rho) ;
$$

notice that $\phi_{\varepsilon}(\rho)$ is bounded since $u^{-}$is Hölder continuous with any exponent $<1$. Dividing both sides of (5.11) by $r^{1+\alpha-\varepsilon}$ we then have

$$
\begin{equation*}
\phi_{\varepsilon}(r) \leqslant c_{n} \alpha \sup _{\rho \leqslant r} \phi_{\varepsilon}(\rho)+\frac{c r^{\epsilon}}{r}\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}\right)^{1 / 2} \tag{5.12}
\end{equation*}
$$

Similarly, if $\psi_{\varepsilon}(r)=\left(r^{\varepsilon} / r\right) f_{\partial B_{r}} u^{+}$then

$$
\psi_{\varepsilon}(r) \leqslant c_{n} \alpha \sup _{\rho \leqslant r} \psi_{\varepsilon}(\rho)+\frac{C r^{\varepsilon}}{r}\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2}\right)^{1 / 2}
$$

By Theorem $4.1 \phi_{\epsilon}(r)=\psi_{\varepsilon}(r)+O\left(r^{\varepsilon}\right)$. Hence,

$$
\begin{equation*}
\phi_{\varepsilon}(r) \leqslant c_{n} \alpha \sup _{\rho \leqslant r} \phi_{\varepsilon}(\rho)+\frac{C r^{\varepsilon}}{r}\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2}\right)^{1 / 2}+O\left(r^{\varepsilon}\right) . \tag{5.13}
\end{equation*}
$$

Taking the product of the left-hand sides of (5.12) and (5.13), we obtain

$$
\left(\phi_{\varepsilon}(r)\right)^{2} \leqslant C \alpha^{2}\left(\sup _{\rho \leqslant r} \phi_{\varepsilon}(\rho)\right)^{2}+C r^{2 \varepsilon}+C \frac{r^{2 \varepsilon}}{r^{2}}\left(\int_{B_{r}} \rho^{2-n}\left|\nabla u^{+}\right|^{2} \cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u^{-}\right|^{2}\right)^{1 / 2} .
$$

Using Lemma 5.1 and choosing $\alpha$ small enough, we obtain $\left(\phi_{\varepsilon}(r)\right)^{2} \leqslant C r^{2 \varepsilon}, C$ independent of $\varepsilon$. Hence, $\frac{1}{r} \jmath_{\partial B_{r}} u^{-} \leqslant C$ and therefore, also, $\frac{1}{r} f_{\partial B_{r}} u^{+} \leqslant C$.

Theorem 5.3. If $\lambda(0)=\min \left(\lambda_{1}, \lambda_{2}\right)$, then $u$ is Lipschitz continuous in any compact subset of $\Omega$.

Proof. Let $K$ be a compact subset of $\Omega$ and introduce $d=\operatorname{dist}(K, \partial \Omega)$. For any $x \in \Omega, x \notin\{u=0\}$, denote by $\rho=\rho(x)$ the distance from $x$ to $\{u=0\}$ and let $x^{0}$ be such that $\rho=\left|x-x^{0}\right|, u\left(x^{0}\right)=0$. If $\rho>d / 6$ then $u$ is harmonic in $B_{d / 6}(x)$, and thus $|D u(x)| \leqslant C / d$. Suppose next that $\rho(x)<d / 6$.

Representing $u$ by Poisson's formula, $u(x)=\int_{\partial B_{\rho}(x)} P_{x}(y) u(y)$, we conclude that

$$
\begin{equation*}
|D u(x)| \leqslant \frac{C}{\rho} f_{\partial B_{\rho}(x)}|u(y)| . \tag{5.14}
\end{equation*}
$$

The function $|u|$ is subharmonic. Representing $|u(y)|\left(y \in \partial B_{\rho}(x)\right)$ by Green's function in $B_{\sigma}\left(x^{0}\right)$ we get

$$
|u(y)| \leqslant f_{\partial B_{0}\left(x^{0}\right)} \tilde{P}_{y}(z)|u(z)|, \quad 3 \rho<\sigma<5 \rho
$$

thus, $|u(y)| \leqslant C f_{\partial B_{o}\left(x^{0}\right)}|u(z)| \leqslant C \rho$ by Lemma 5.2. Substituting this into (5.14) we conclude that $|D u(x)| \leqslant C$ if $x \in K, u(x) \neq 0$. Since $u \in H_{\mathrm{loc}}^{1,2}(\Omega), D u=0$ a.e. on $\{u=0\}$, and thus $D u \in L^{\infty}(K)$.

Another Proof of Theorem 5.3. We shall give another proof, also based on Theorem 4.1 and Lemma 5.1.

Suppose $0 \in \Omega, u(0)=M>0$, and let $x^{0}$ be the nearest point to 0 on $\{u=0\}$. We assume first that $\left|x^{0}\right|=1$ and $B_{2} \subset \Omega$. By Harnack's inequality $u>c_{0} M$ in $B_{3 / 4}$ $\left(c_{0}>0\right)$ and therefore $f_{\partial B_{1}\left(x^{0}\right)} u^{+}>c M(c>0)$. From Theorem 4.1 it follows that

$$
\begin{equation*}
\int_{\partial B_{1}\left(x^{0}\right)} u^{-}>c M \tag{5.15}
\end{equation*}
$$

with another $c>0$, provided $M$ is large enough.
Let $y \in \partial B_{1 / 2}$ be a point on $\overrightarrow{0 x^{0}}$. Then $u>c_{0} M>0$ in $B_{1 / 4}(y)$. We shall use polar coordinates $(r, \omega)$ about $y$. Denote by $\Gamma$ the set of $\omega$ 's such that if $(r, \omega) \in$ $\partial B_{1}\left(x^{0}\right)$ then $u(r, \omega)<0$.

We integrate $u_{r}^{-}(r, \omega)$ over $(r, \omega) \in B_{1}\left(x^{0}\right), \omega \in \Gamma$. Using (5.15) and the fact that $u>0$ in $B_{1 / 4}(y)$ we obtain

$$
\begin{equation*}
c M \leqslant \int_{\partial B_{1}(y)} u^{-}=\int_{\Gamma} d \omega \int u_{r}^{-} \leqslant|\Gamma|^{1 / 2}\left\{\int_{B_{1}\left(x^{0}\right)}\left|\nabla u^{-}\right|^{2}\right\}^{1 / 2} . \tag{5.16}
\end{equation*}
$$

Next we integrate $u_{r}^{+}(r, \omega)$ and $(r, \omega) \in\left\{B_{1}\left(x^{0}\right) \backslash B_{1 / 4}(y)\right\}, \omega \in \Gamma$, and notice that $u^{+}(r, \omega) \geqslant c_{0} M$ in $B_{1 / 4}(y)$. We obtain

$$
\begin{equation*}
c M|\Gamma| \leqslant \int_{\Gamma} d \omega \int u_{r}^{+} \leqslant|\Gamma|^{1 / 2}\left\{\int_{B_{1}\left(x^{0}\right)}\left|\nabla u^{+}\right|^{2}\right\}^{1 / 2} \tag{5.17}
\end{equation*}
$$

Taking the product of both sides of the inequalities in (5.16) and (5.17), we get

$$
c M^{4} \leqslant \int_{B_{1}\left(x^{0}\right)}\left|\nabla u^{+}\right|^{2} \cdot \int_{B_{1}\left(x^{0}\right)}\left|\nabla u^{-}\right|^{2} .
$$

Using Lemma 5.1 we then obtain $M \leqslant C$. We have thus proved that

$$
\begin{equation*}
u(z) \leqslant C \rho(z) \quad(\rho(z)=\operatorname{dist}(z,\{u=0\})) \tag{5.18}
\end{equation*}
$$

if $z=0, \rho(z)=1, u(z)>0, B_{2 \rho(z)}(z) \subset \Omega$. The proof for general $z$ follows by considering $\tilde{u}(x) \equiv u(z+\rho(z) x) / \rho(z)$. From (5.18) we deduce that $|\nabla u(z)| \leqslant C$; the same estimate holds if $u(z)<0$. The proof that $u \in C_{\text {loc }}^{0.1}$ now readily follows.
6. Blow-up limits. The rest of this paper is devoted to the study of the free boundary. For definiteness we shall always assume that

$$
\begin{equation*}
\Lambda<0, \quad \lambda(0)=\lambda_{1} \tag{6.1}
\end{equation*}
$$

all the results obviously extend to the case $\Lambda>0, \lambda(0)=\lambda_{2}$. When (6.1) holds the free boundary coincides with

$$
\begin{equation*}
\Gamma^{+}=\partial\{u>0\} \tag{6.2}
\end{equation*}
$$

Indeed, for the remaining free boundary

$$
\begin{equation*}
\Gamma^{-}=\partial\{u<0\} \backslash \partial\{u>0\} \tag{6.3}
\end{equation*}
$$

we obviously have

$$
\begin{equation*}
u \leqslant 0 \quad \text { in a neighborhood } N \text { of } \Gamma^{-} . \tag{6.4}
\end{equation*}
$$

But since $\lambda(0)=\lambda_{1}$, the minimizer $u$ must be harmonic in $N$. Consequently, $\Gamma$ is empty.

Definition 6.1. A function $u$ is called a minimizer (of $J$ ) in $R^{n}$ if for any $B_{r} \subset R^{\prime \prime}$ and for any $v \in H^{1,2}\left(B_{r}\right), v=u$ on $\partial B_{r}$.

$$
J_{B}(u) \leqslant J_{B}(v)
$$

where $J_{B}(v)$ is the functional $J(v)$ with $\Omega$ replaced by $B_{r}$.
Suppose $u$ is a minimizer, $u\left(x_{k}\right)=0, x_{k} \rightarrow x_{0} \in \Omega, \rho_{k} \downarrow 0$, and set

$$
\begin{equation*}
u_{k}(x)=\frac{1}{\rho_{k}} u\left(x_{i}+\rho_{k} x\right) . \tag{6.5}
\end{equation*}
$$

We call $\left\{u_{k}\right\}$ a blow-up sequence with respect to $B_{\rho_{k}}\left(x_{k}\right)$. Since $\left|\nabla u_{k}(x)\right| \leqslant C$ in any bounded set and $u_{k}(0)=0$, we have, for a subsequence,

$$
\begin{align*}
u_{k}(x) & \rightarrow u_{0}(x) \quad \text { unformly in bounded sets, }  \tag{6.6}\\
\nabla u_{k} & \rightarrow \nabla u_{0} \quad \text { weakly in } L_{\text {loc }}^{\infty}\left(R^{\prime \prime}\right)
\end{align*}
$$

$u_{0}$ is called a blow-up limit.
Lemma 6.1. There holds

$$
\begin{gather*}
\partial\left\{u_{k}>0\right\} \rightarrow \partial\left\{u_{0}>0\right\} \quad \text { locally in the Hausdorff metric, }  \tag{6.7}\\
\nabla u_{k} \rightarrow \nabla u_{0} \quad \text { a.e. in } R^{n} . \tag{6.8}
\end{gather*}
$$

Proof. Suppose a ball $\bar{B}_{r}$ does not intersect $\partial\left\{u_{0}>0\right\}$. Then either $u_{0}>0$ in $\bar{B}_{r}$ or $u_{0} \leqslant 0$ in $\bar{B}_{r}$. In the first case $u_{k}>0$ in $\bar{B}_{r}$ if $k$ is large enough. In the second case $\frac{1}{r} f_{\partial B_{r}} u_{k}^{+}<\varepsilon$ for any $\varepsilon>0$ if $k$ is large enough, so that, by nondegeneracy, $u_{k} \leqslant 0$ in $B_{r / 2}$.

In both cases we conclude that $B_{r / 2}$ does not intersect $\partial\left\{u_{k}>0\right\}$ if $k$ is large enough.

Conversely, if $B_{r}$ does not intersect $\partial\left\{u_{k}>0\right\}$ for any large $k$ then either $u_{k}>0$ in $B_{r}$ or $u_{k} \leqslant 0$ in $B_{r}$. In the first case $u_{k}$ is harmonic in $B_{r}$ and then so is $u_{0}$; thus either $u_{0}>0$ in $B_{r}$ or $u_{0} \equiv 0$ in $B_{r}$, so that $B_{r}$ does not intersect $\partial\left\{u_{0}>0\right\}$. In the second case we have $u_{0} \leqslant 0$ in $B_{r}$ so that again $B_{r}$ does not intersect $\partial\left\{u_{0}>0\right\}$.

To prove (6.8) notice that, in every compact subset of $\left\{u_{0} \neq 0\right\}$, (6.8) is certainly valid. Next consider a density point $x^{0}$ of the set $\left\{u_{0}(x)=0\right\}$. By the Lipschitz continuity of $u_{0}$, we then deduce that $\left|u_{0}\right|=o(r)$ in $B_{r}$, and therefore, $\frac{1}{r} \int_{\partial B} u_{0}^{+}=o(1)$ as $r \rightarrow 0$.

Since $u_{k} \rightarrow u_{0}$ uniformly in $B_{1}$, we get $\frac{1}{r} \int_{\partial B_{r}} u_{k}^{+}<\varepsilon$ for any small $\varepsilon>0$, provided $k$ is large enough; hence by nondegeneracy, $u_{k} \leqslant 0$ in $B_{r}$. But then (since $\lambda(0)=\lambda_{1}$ ) $u_{k}$ is harmonic in $B_{r}$ and then so is $u_{0}$. Consequently, $\nabla u_{k} \rightarrow \nabla u_{0}$ uniformly in $B_{r / 2}$. We have thus proved that almost all the set $\left\{u_{0}=0\right\}$ can be covered by balls $B_{r}$, with suitable centers such that $\nabla u_{k} \rightarrow \nabla u_{0}$ in each $B_{r_{1}}$. It follows that $\nabla u_{k} \rightarrow \nabla u_{0}$ a.e. in the set $\left\{u_{0}=0\right\}$. This completes the proof of (6.8).

Lemma 6.2. $u_{0}$ is a minimizer in $R^{n}$ with respect to the function $Q_{0}(u, \lambda)=$ $q\left(x_{0}\right) \lambda(u)$.

Indeed, the proof is similar to the proof of Lemma 5.4 in [1]; that proof can be slightly simplified by using (6.8).

Theorem 6.3. Suppose $D \Subset \Omega, B_{r} \subset D$ with center $x^{0}$ in $\partial\{u>0\}$. Then

$$
\begin{equation*}
\frac{1}{r} f_{\partial B_{r}\left(x^{0}\right)} u \geqslant c, \quad c>0 . \tag{6.9}
\end{equation*}
$$

Strictly speaking, this result does not include Corollary 3.2 since the constant $c$ in (6.9) depends also on $D$ and on the Lipschitz coefficient of $u$.

Proof. Suppose the assertion is not true. Then there exist points $x_{m}^{0} \in D$ and $r_{m} \downarrow 0$ such that

$$
\frac{1}{r_{m}} f_{\partial B_{r}\left(x_{m}^{0}\right)} u \rightarrow 0, \quad x_{m}^{0} \in \partial\{u>0\}
$$

Setting $u_{m}(x)=u\left(x_{m}^{0}+r_{m} x\right) / r_{m}$ we may suppose that $x_{m}^{0} \rightarrow 0, u_{m} \rightarrow u_{0}$ uniformly in bounded sets. Then $u_{0}$ is subharmonic (since $u_{m}$ is subharmonic) and $f_{\partial B_{1}} u_{0}=0$ $=u_{0}(0)$. By the maximum principle it then follows that $u_{0}$ is harmonic in $B_{1}$.

Now $u_{0}$ is a local minimizer and $0 \in \partial\left\{u_{0}>0\right\}$, by (6.7). It follows that the free boundary $\partial\left\{u_{0}>0\right\}$ is nonempty; this set must be piecewise analytic since $u_{0}$ is harmonic. But then Theorem 2.4 shows that $\left|\nabla u_{0}\right|^{2}$ has jump $\Lambda q^{1}(0)$ across smooth parts of the free boundary. Since, however, $u_{0}$ is harmonic, $\left|\nabla u_{0}\right|^{2}$ cannot have a jump, i.e., $\Lambda q^{2}(0)=0$, a contradiction.

Consider a blow-up family

$$
u_{\varepsilon}(x)=\frac{1}{\varepsilon} u\left(x^{0}+\varepsilon x\right), \quad x^{0} \in \partial\{u>0\}, \varepsilon>0
$$

and let

$$
\begin{aligned}
I_{f}(r) & \equiv \frac{1}{r^{4}} \int_{B_{r}} \rho^{2-n}\left|\nabla u_{f}^{+}\right|^{2} \cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u_{\varepsilon}^{-}\right|^{2} \\
& =\frac{1}{(\varepsilon r)^{4}} \int_{B_{f r}\left(x^{0}\right)} \rho^{2-n}\left|\nabla u^{+}\right|^{2} \cdot \int_{B_{\varepsilon r}\left(x^{0}\right)} \rho^{2-n}\left|\nabla u^{-}\right|^{2} \equiv \tilde{I}_{\varepsilon r} .
\end{aligned}
$$

By Lemma 5.1, $\tilde{I}_{\rho}$ is an increasing function of $\rho$. Consequently there exists a nonnegative constant $\gamma$ such that

$$
\begin{equation*}
I_{\varepsilon}(r) \downarrow \gamma \quad \text { if } \varepsilon \downarrow 0 \tag{6.10}
\end{equation*}
$$

Now take a sequence $\varepsilon=\varepsilon_{k} \downarrow 0$ such that

$$
\begin{equation*}
u_{e_{\lambda}}(x) \rightarrow u_{0}(x) \quad \text { uniformly in bounded subsets of } R^{n} . \tag{6.11}
\end{equation*}
$$

Lemma 6.4. If (6.11) holds then, as $\varepsilon_{k} \downarrow 0$,

$$
\begin{equation*}
I_{\varepsilon_{r}}(r) \rightarrow \frac{1}{r^{4}} \int_{B_{r}} \rho^{2-n}\left|\nabla u_{0}^{+}\right|^{2} \cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u_{0}^{-}\right|^{2} . \tag{6.12}
\end{equation*}
$$

Proof. By the Lipschitz continuity of $u,\left|\nabla u_{\varepsilon}^{+}\right| \leqslant C$, and by Lemma 6.1, $\nabla u_{\varepsilon}^{+} \rightarrow$ $\nabla u_{0}^{ \pm}$a.e. Hence, (6.12) follows by the Lebesgue bounded convergence theorem.
Corollary 6.5. For any blow-up limit $u_{0}$ of $u_{\varepsilon}$ there holds

$$
\begin{equation*}
\frac{1}{r^{4}} \int_{B_{r}} \rho^{2-n}\left|\nabla u_{0}^{+}\right|^{2} \cdot \int_{B_{r}} \rho^{2-n}\left|\nabla u_{0}^{-}\right|^{2}=\gamma \tag{6.13}
\end{equation*}
$$

for all $r>0$.
Lemma 6.6. (i) If $\gamma=0$ then $u_{0} \geqslant 0$ in $R^{\prime \prime}$; (ii) if $\gamma>0$ and $n=2$ then $u_{0}(x)=$ $\mu_{2}(x \cdot e)^{+}-\mu_{1}(x \cdot e)^{-}$in $R^{n}$ where $e$ is a constant unit vector, $\mu_{i}$ are positive constants, and $\mu_{1}^{2}-\mu_{2}^{2}=\Lambda q^{2}\left(x^{0}\right)$.

The function $u_{0}$ in (ii) is called a 2-plane solution; if $\mu_{1}=0$ or $\mu_{2}=0$ then we call it a 1-plane solution.

Proof. If $\gamma=0$ then either $u_{0}^{+}=0$ or $u_{0}^{-}=0$ in $R^{n}$. Since $u_{0}$ is subharmonic and $u_{0}(0)=0$, we conclude that $u_{0} \geqslant 0$. To prove (ii) we check the proof of Lemma 5.1 and find that equality can hold in (6.13) only if equality holds in the various Cauchy-Schwarz inequalities and $s_{1}=s_{2}=1 / 2$. Thus, with $S_{1}$ replaced by $S_{r}$,

$$
\begin{gathered}
\left|u_{r}^{+}\right|=C u^{+}, \quad u^{+} u_{r}^{+} \geqslant 0, C \text { constant }, \\
\int\left(u_{r}^{+}\right)^{2}=\beta_{1}^{2} \int\left|\nabla_{\theta} u^{+}\right|^{2}, \quad \int\left|\nabla_{\theta} u^{+}\right|^{2}=\frac{1}{\alpha_{1} r^{2}} \int\left(u^{+}\right)^{2} .
\end{gathered}
$$

It follows that $u_{r}^{+}=c u^{+} / r\left(c=c_{n}>0\right)$; a similar relation holds for $u^{-}$. Thus $u=r^{b} g(\theta)$ if $u \neq 0$. Since $u$ is bounded, $b \geqslant 0$. By nondegeneracy and Lipschitz
continuity, $b=1$. Thus $u=\operatorname{rg}(\theta)$ if $u \neq 0$ and then,

$$
\begin{equation*}
A g+(n-1) g=0 \quad \text { where } g(\theta) \neq 0 \tag{6.14}
\end{equation*}
$$

where $A$ is the Laplacian restricted to $\partial B_{1}$. If $n=2$ then $A g=g^{\prime \prime}$ and the assertions easily follow using Lemma 6.2 and Theorem 2.4.

Remark 6.1. We do not know whether the isoperimetric inequality $\gamma_{1}(E) \geqslant \gamma_{1}\left(E^{*}\right)$ used in the proof of Lemma 5.1 is a strict inequality whenever $E$ is not a spherical cap. If this is indeed the case then Lemma 6.6(ii) is valid for any $n \geqslant 2$. Indeed, from the proof of Lemma 5.1 we then conclude that, for any $r, u_{0}=l_{2}(x \cdot e)^{+}-l_{1}(x \cdot e)^{-}$ on $\partial B_{r}$ where $e=e(r), l_{i}=l_{i}(r)>0$. Setting

$$
\begin{aligned}
f(r) & =l_{1} e \\
& \text { if } u_{0}>0 \\
& =l_{2} e
\end{aligned} \quad \text { if } u_{0}<0, ~ \$
$$

we have $\Delta(x \cdot f(r))=0$ on $\left\{u_{0} \neq 0\right\}$, which gives

$$
\sum x_{i}\left(f_{i}^{\prime \prime}(r)+\frac{n+1}{r} f_{i}^{\prime}(r)\right)=0
$$

where $f=\left(f_{1}, \ldots, f_{n}\right)$. It follows that $f^{\prime \prime}+(n+1) f^{\prime} / r=0$, or $f(r)=C r^{-n-1}+c$ where $C, c$ are constant vectors in any component of $\left\{u_{0} \neq 0\right\}$. Since $u_{0}$ is bounded, $C=0$ and the assertions in (ii) easily follow.

## 7. Properties of the free boundary.

Theorem 7.1. There exists a positive constant $c \in(0,1)$ such that for any ball $B_{r} \subset \Omega$ with center in $\partial\{u>0\}$

$$
\begin{equation*}
c \leqslant \frac{巳^{n}\left(B_{r} \cap\{u>0\}\right)}{巳^{n}\left(B_{r}\right)} \leqslant 1-c . \tag{7.1}
\end{equation*}
$$

Proof. By nondegeneracy there exists a point $y \in \partial B_{r / 2}$ with $u(y) \geqslant c r$. Since $u$ is Lipschitz, $u(x)>0$ in $B_{\kappa r}(y)$, for some small enough $\kappa$. This establishes the left-hand side of (7.1). To obtain the second inequality, let

$$
\begin{array}{rlrl}
\Delta v & =0 & \text { in } B_{r}, \\
v & =u & & \text { on } \partial B_{r} .
\end{array}
$$

Then $v \geqslant u$ in $B_{r}$ and (cf. the proof of Theorem 2.1)

$$
\begin{equation*}
\int_{B_{r} \cap\{u \leqslant 0<v\}}|\Lambda| \geqslant \int_{B_{r}}|\nabla(u-v)|^{2} \geqslant \frac{c}{r^{2}} \int_{B_{r}}|u-v|^{2} \geqslant \frac{c}{r^{2}} f_{B_{k} r}|u-v|^{2} \tag{7.2}
\end{equation*}
$$

for any $0<\kappa<1$.
If $y \in B_{\kappa r}$ then (we take the center of $B_{r}$ to be at the origin)

$$
|v(y)-v(0)| \leqslant|y||\nabla v| \leqslant \kappa r \frac{C}{r} f_{\partial B_{r}}|u| \leqslant C \kappa r,
$$

and $v(0)=f_{\partial B_{r}} v=f_{\partial B_{r}} u,|u(y)| \leqslant C \kappa r$. It follows that $|v(y)-u(y)| \geqslant \int_{\partial B_{r}} u-C \kappa r$. Recalling Theorem 6.3 we obtain

$$
|v(y)-u(y)| \geqslant c r-C \kappa r \geqslant c r / 2
$$

if $\kappa$ is small enough. Using this estimate in (7.2) we find that

$$
\mathrm{E}^{\prime \prime}\left(B_{r} \cap\{u \leqslant 0\}\right) \geqslant \frac{C}{r^{2}} \int_{B_{k r}} c^{2} r^{2} \geqslant c r^{n} \quad(c>0) .
$$

Since $u^{\star}$ are continuous and subharmonic, the measures $d \lambda^{+}=\Delta u^{+}$and $d \lambda^{-}=$ $\Delta u^{-}$are Radon measures supported on $\Omega \cap \partial\{u>0\}$ and $\Omega \cap \partial\{u<0\}$, respectively.

Theorem 7.2. For any $D \Subset \Omega$ there exist positive constants $c, C$ such that, for any $B_{r} \subset D$ with center in $\partial\{u>0\}$,

$$
\begin{gather*}
c r^{n-1} \leqslant \int_{B_{r}} d \lambda^{+} \leqslant C r^{n-1},  \tag{7.3}\\
\int_{B_{r}} d \lambda^{-} \leqslant C r^{n-1} \tag{7.4}
\end{gather*}
$$

Proof. Let $x \in \partial\{u>0\}$. For almost all $r$ with $B_{r}(x) \subset \Omega$,

$$
\int_{B_{r}(x)} d \lambda^{+}=\int_{\partial B_{r}(x)} \nabla u^{+} \cdot \nu d H^{n-1} \leqslant C r^{n-1}
$$

since $u^{+}$is Lipschitz continuous. This proves the second inequality in (7.3). The proof of (7.4) is similar. The proof of the first inequality in (7.3) is similar to the proof given in [1, Theorem 4.3], with $u$ replaced by $u^{+}$.

Theorem 7.3 (Representative Theorem). (i) If $D 匹 \Omega$ then

$$
H^{n-1}(D \cap \partial\{u>0\})<\infty .
$$

(ii) There exist Borel functions $q_{u}^{ \pm}$such that

$$
\begin{equation*}
\Delta u^{ \pm}=q_{u}^{ \pm} H^{n-1} L \partial\{u>0\}, \tag{7.5}
\end{equation*}
$$

that is, for every $\zeta \in C_{0}^{\infty}(\Omega)$,

$$
\begin{equation*}
-\int_{\Omega} \nabla u^{ \pm} \cdot \nabla \zeta=\int_{\Omega \cap \partial\{u>0\}} \zeta q_{u}^{ \pm} d H^{n-1} \tag{7.6}
\end{equation*}
$$

(iii) For any $D \Subset \Omega$ there exist positive constants $c, C$ depending on $D, \Omega$, the constant $c$ in Corollary 3.2 and any bound on $|\nabla u|_{L^{x}(D)}$, such that for any ball $B_{r}(x) \subset D$ with $x \in \partial\{u>0\}$,

$$
\begin{align*}
c & \leqslant q_{u}^{+} \leqslant C  \tag{7.7}\\
c r^{n-1} & \leqslant H^{n-1}\left(B_{r}(x) \cap \partial\{u>0\}\right) \leqslant C r^{n-1},  \tag{7.8}\\
0 & \leqslant q_{u}^{-} \leqslant C . \tag{7.9}
\end{align*}
$$

Proof. For any compact set $E \subset D \cap \partial\{u>0\}$ and small $r$ choose a covering of $E$ with balls $B_{r}\left(y_{i}\right)$ such that $\sum I_{B_{2 r}\left(y_{i}\right)} \leqslant C$. Choosing $x_{i} \in B_{r}\left(y_{i}\right) \cap E$ we have, by Theorem 7.2,

$$
\sum_{i} r^{n-1} \leqslant C \sum_{i} \lambda^{+}\left(B_{r}\left(x_{i}\right)\right) \leqslant C \lambda^{+}\left(B_{4 r}(E)\right)
$$

which gives

$$
\begin{equation*}
H^{n-1}(E) \leqslant C \lambda^{+}(E) \tag{7.10}
\end{equation*}
$$

Thus (i) holds and $H^{n-1} L(D \cap \partial\{u>0\})$ is absolutely continuous with respect to $\lambda^{+}$.

Next, the support of $\lambda^{+}$is contained in $\partial\{u>0\}$ and, by Theorem 7.2,

$$
\begin{equation*}
\lambda^{+}\left(B_{r}\right) \leqslant C r^{n-1} \quad \text { for any ball } B_{r} \subset D \tag{7.11}
\end{equation*}
$$

from this it follows that $\lambda^{+}(E) \leqslant C H^{n-1}(E)$. We have thus shown that the Radon measure $\lambda^{+}$is absolutely continuous with respect to the Radon measure

$$
H^{n-1} L \partial\{u>0\}
$$

and vice versa. Setting $q_{u}^{+}=d \lambda^{+} / d H^{n-1} \mathscr{L}(\partial\{u>0\}$ ) we see that (7.5) holds (for $\Delta u^{+}$), and (7.10) and (7.11) establish (7.7) and (7.8).

Using the assertion (i) we can now proceed with proving (ii) and (iii) for $\lambda^{-}$by the same proof as for $\lambda^{+}$.

Since $\partial\{u>0\}$ has finite $H^{n-1}$ measure, the set $A=\Omega \cap\{u>0\}$ has finite perimeter locally in $\Omega$, that is, $\mu_{u} \equiv-\nabla I_{A}$ is a Borel measure and the total variation $\left|\mu_{u}\right|$ is a radon measure. We denote by $\partial_{\text {red }}\{u>0\}$ the reduced boundary of $\partial\{u>0\}$.

Theorem 7.4 (Identification Theorem). Let $x_{0} \in \partial_{\text {red }}\{u>0\}$ with

$$
\begin{align*}
& \quad \theta^{* n-1}\left(H^{n-1} L \partial\{u>0\}, x_{0}\right) \leqslant 1,  \tag{7.12}\\
& \int_{B_{r}\left(x_{0}\right) \cap \partial\{u>0\}}\left|q_{u}^{ \pm}-q_{u}^{ \pm}\left(x_{0}\right)\right|=o(1), \quad r \rightarrow 0 . \tag{7.13}
\end{align*}
$$

(i) If $\gamma>0$ (in Corollary 6.5) and $n=2$, then

$$
u\left(x_{0}+x\right)=\mu_{2}\left(x \cdot e\left(x_{0}\right)\right)^{+}-\mu_{1}\left(x \cdot e\left(x_{0}\right)\right)^{-}+o(|x|) \quad \text { as }|x| \rightarrow 0
$$

where $\mu_{i}>0, \mu_{1}^{2}-\mu_{2}^{2}=\Lambda q^{2}\left(x_{0}\right)$, and

$$
\left(\mu_{2}-\mu_{1}\right) e\left(x_{0}\right)=\left(q_{u}^{+}\left(x_{0}\right)-q_{u}^{-}\left(x_{0}\right)\right) \nu_{u}\left(x_{0}\right)
$$

(ii) If $\gamma=0$, then

$$
u\left(x_{0}+x\right)=q_{u}^{+}\left(x_{0}\right) \max \left\{-x \cdot \nu_{u}\left(x_{0}\right), 0\right\}+o(|x|)
$$

as $|x| \rightarrow 0$, and $\left(q_{u}^{+}\left(x_{0}\right)\right)^{2}=\left(\lambda_{2}^{2}-\lambda^{2}(0)\right) q^{2}\left(x_{0}\right)$.
Here $\nu_{u}\left(x_{0}\right)$ is the outward normal to $\partial\{u>0\}$ at $x_{0}$.
Proof. Take a blow-up sequence $u_{\varepsilon}(x)=u\left(x_{0}+\varepsilon x\right) / \varepsilon$ with $u_{\varepsilon} \rightarrow u_{0}$ uniformly in compact subsets. Then $\Delta u_{\varepsilon} \rightarrow \Delta u_{0}$ as distributions, and thus also as measures. From (7.5) we deduce that

$$
\Delta u_{\varepsilon}^{ \pm}=q_{u}^{ \pm}\left(x_{0}+\varepsilon x\right) H^{n-1} L \partial\left\{u^{\varepsilon}>0\right\} .
$$

If $\gamma>0$ and $n=2$ then, by Lemma 6.6,

$$
u_{0}=\mu_{2}(x \cdot e)^{+}-\mu_{1}(x \cdot e)^{-} \quad(e \text { constant vector })
$$

and therefore,

$$
\Delta u_{0}^{+}-\Delta u_{0}^{-}=\left(\mu_{1}-\mu_{2}\right) e d H^{n-1} L \Pi_{e}
$$

where $\Pi_{e}$ is the hyperplane orthogonal to $e$. We thus conclude that

$$
\left[q_{u}^{+}\left(x_{0}+\varepsilon x\right)-q_{u}^{-}\left(x_{0}+\varepsilon x\right)\right] d H^{n-1} L \partial\left\{u^{\varepsilon}>0\right\} \rightarrow\left(\mu_{1}-\mu_{2}\right) e d H^{n-1} L \Pi_{e} .
$$

Since $x_{0} \in \partial_{\text {red }}\{\mathbf{u}>0\}$ we have [10, Theorem 3.7]

$$
d H^{n-1} L \partial\left\{u^{\varepsilon}>0\right\} \rightarrow d H^{n-1} L \Pi_{0}
$$

where $\Pi_{0}=\left\{x ; \nu_{u}\left(x_{0}\right) \cdot x=0\right\}$. Recalling (7.12) we deduce that

$$
\left(q_{u}^{+}\left(x_{0}\right)-q_{u}^{-}\left(x_{0}\right)\right) d H^{n-1} L \Pi_{0}=\left(\mu_{2}-\mu_{1}\right) e d H^{n-1} L \Pi_{e}
$$

so that

$$
\left(\mu_{2}-\mu_{1}\right) e=\left(q_{u}^{+}\left(x_{0}\right)-q_{u}^{-}\left(x_{0}\right)\right) \nu_{u}\left(x_{0}\right) ;
$$

also $\mu_{1}^{2}-\mu_{2}^{2}=\Lambda q^{2}\left(x_{0}\right)$. Thus the $\mu_{i}$ and $e$ are uniquely determined, independently of the blow-up sequence, and assertion (i) follows.

Consider next the case $\gamma=0$. By Lemma 6.6 we then have $u_{0} \geqslant 0$ for any blow-up limit of the $u_{\mathrm{e}}$. We can then proceed as in Theorem 4.8 of [1]. Thus, taking $\nu_{u}\left(x_{0}\right)=e_{n}$, the proof that $u_{0}>0$ if $x_{n}<0, u_{0}=0$ if $x_{n}>0$ is the same as in [1]. Next, setting

$$
\begin{equation*}
\mu_{w}=-\nabla I_{(\Omega \cap\{w>0\})} \tag{7.14}
\end{equation*}
$$

for any function $w$ for which $\partial\{w>0\}$ has finite $H^{n-1}$ measure, we have, for every compact subset $E \subset B_{r}^{\prime}\left(B_{r}^{\prime}\right.$ is the ball in $\left.R^{n-1}\right)$,

$$
\begin{aligned}
H^{n-1}(E) & =\mu_{u^{+}}(E x(-1,1)) \cdot e_{n} \leqslant H^{n-1}\left(\partial\left\{u_{\varepsilon}^{+}>0\right\} \cap(E x(-1,1))\right) \\
& =H^{n-1}\left(\partial\left\{u_{\varepsilon}>0\right\} \cap(E x(-1,1))\right)
\end{aligned}
$$

and we can again proceed as in [1], thereby establishing that $u_{0}(x)=u_{0}^{+}(x)=$ $-q_{u}\left(x_{0}\right) x_{n}$ if $x_{n}<0$, and the proof of (ii) thereby follows; the last assertion in (ii) follows from Lemma 6.2 and Remark 2.1.

Remark 7.1. From Theorem 7.1 it follows (by [7, 4.5.6(3)]) that

$$
H^{n-1}\left(\partial\{u>0\} \backslash \partial_{\text {red }}\{u>0\}\right)=0 .
$$

From [7,4.5.6(2), 2.9.8 and 2.9.9] applied to $H^{n-1} L \partial\{u>0\}$ and the Vitali relation $\left\{\left(x, B_{r}(x)\right): x \in \partial\{u>0\}\right.$ and $\left.B_{r}(x) \subset \Omega\right\}$ it follows that for $H^{n-1}$ a.a. $x_{0} \in$ $\partial_{\text {red }}\{u>0\}$ the assumptions (7.12) and (7.13) are satisfied. Thus Theorem 7.4 shows that for $H^{n-1}$ a.a. $x \in \partial\{u>0\}$ the free boundary in a neighborhood of $x_{0}$ is approximately a hyperplane.

Remark 7.2. In special models arising in jet flows [5, 6] it has been shown that the free boundary is a continuous graph. In the next section we prove, more generally, that the free boundary is $C^{1}$ if $n=2$.
8. Differentiability of the free boundary ( $n=2$ ). In this section we prove that, in case $n=2$, the free boundary is continuously differentiable. The first lemma is valid for any $n \geqslant 2$. In proving it we shall use the fact that
(8.1) the sets $\{u>0\}$ and $\{u<0\}$ are connected to the boundary of $\Omega$.

To show this, suppose $K$ is a component of $\{u>0\}$ which is not connected to the boundary. Then, by replacing $u$ in $K$ by 0 we obtain a new function $\tilde{u}$ with smaller functional $J(\tilde{u})$, which is a contradiction.

Lemma 8.1. If $u$ and $\tilde{u}$ are both minimizers of $J$ in a bounded domain $D$, and if $\tilde{u}>u$ on $\partial D$, then $\tilde{u}>u$ in $\{u \neq 0\}$.

Proof. Set $v_{1}=\min \{u, \tilde{u}\}$ and $v_{2}=\max \{u, \tilde{u}\}$. Then $v_{1}=u$ on $\partial D$ and therefore, $J\left(v_{1}\right) \geqslant J(u)$. Similarly, $v_{2}=\tilde{u}$ on $\partial D$ and therefore, $J\left(v_{2}\right) \geqslant J(\tilde{u})$. However, $J\left(v_{1}\right)+J\left(v_{2}\right)=J(u)+J(\tilde{u})$ as seen by writing explicitly the terms in each $J$. It follows that $J\left(v_{1}\right)=J(u)$.

Suppose $u\left(x^{0}\right)=\tilde{u}\left(x^{0}\right) \neq 0$ and $u-\tilde{u}$ changes sign in any neighborhood of $x^{0}$. Then $v_{1}$ is not harmonic in any neighborhood of $x^{0}$. We introduce the function $w$ defined by

$$
\begin{aligned}
\Delta w & =0 \quad \text { in } B_{r}\left(x^{0}\right), \\
w & =v_{1} \quad \text { on } \partial B_{r}\left(x^{0}\right)
\end{aligned}
$$

for some small $r>0$, and $w=v_{1}$ in $D \backslash B_{r}\left(x^{0}\right)$. By the Dirichlet principle we find that $J(w)<J\left(v_{1}\right)=J(u)$, contradicting the minimality of $u$. Thus we conclude that either $\tilde{u} \geqslant u$ or $u \geqslant \tilde{u}$ in some neighborhood of $x^{0}$. Starting with $x^{0}$ near $\partial D$ and recalling (8.1), we deduce that $\tilde{u} \geqslant u$ on the set $\{u \neq 0\}$; furthermore, by the strong maximum principle, $\tilde{u}>u$ in this set.

From now on we make the assumptions

$$
\begin{equation*}
n=2, \quad q(x) \equiv 1 \tag{8.2}
\end{equation*}
$$

For definiteness we shall also assume that $\Lambda<0$. We denote points in $R^{2}$ by $X$ or $(x, y)$. Set $e_{1}=(1,0)$ and $e_{2}=(0,1)$.

Lemma 8.2. For any $\varepsilon_{0}>0, \eta>0$ there is a $\delta=\delta\left(\varepsilon_{0}, \eta\right)>0$ such that for any minimizer $u$ in the rectangle $I=\{-3<x<3,-1<y<1\}$ satisfying
(i) the free boundary contains $(0,0)$ and lies in the strip $\{|y|<\delta\}$,
(ii) $u(A)<-\eta$ where $A=\left(0,-\frac{1}{2}\right)$,
the free boundary in $I_{0}=\{-1<x<1,-1<y<1\}$ is a graph in any direction $\varepsilon e_{2} \pm e_{1}, \varepsilon \geqslant \varepsilon_{0}$.

Proof. Take a circle $K_{\mu}^{1}:(x+2)^{2}+(y-\mu)^{2}<\delta^{-3 / 2}$ with center $(-2, \mu)$ and radius $\delta^{-3 / 4}$ and increase $\mu$ from $-\infty$ until, at $\mu=\mu_{1}, \partial K_{\mu_{1}}^{1}$ touches the free boundary of $u$ for the first time. Since $\partial K_{\mu_{1}}^{1} \cap\{x=-2\}$ lies in $\{y<\delta\}$,

$$
\partial K_{\mu_{1}}^{1} \cap\left\{-3<x<-\frac{5}{2}\right\} \quad \text { and } \quad \partial K_{\mu_{1}}^{1} \cap\left\{-\frac{3}{2}<x<3\right\}
$$

both lie below $y=\delta-C \delta^{3 / 4}$ and thus, also below $y=-\delta$ if $\delta$ is small enough. Consequently, $\partial K_{\mu_{1}}^{1} \cap \partial\{u>0\}$ lies in $\left\{-\frac{5}{2}<\mathrm{x}<-\frac{3}{2}\right\}$ and contains a point $E_{1}=$ $\left(x_{1}, y_{1}\right)$ with $-\frac{5}{2}<x_{1}<-\frac{3}{2},-\delta<y_{1}<\delta$.

Similarly, we construct a circle $K_{\mu_{2}}^{2}$ whose closure intersects the free boundary only at points of $\partial K_{\mu_{2}}^{2}$ lying in $\left\{\frac{3}{2}<x<\frac{5}{2}\right\}$, and a point $E_{2}=\left(x_{2}, y_{2}\right)$ on $\partial K_{\mu_{2}}^{2} \cap$ $\partial\{u>0\}$, with $\frac{3}{2}<x_{2}<\frac{5}{2},-\delta<y_{2}<\delta$; further, $K_{\mu_{1}}^{1} \cap\{|y| \leqslant \delta\}$ and $K_{\mu_{2}}^{2} \cap\{|y| \leqslant$ $\delta\}$ are disjoint. We denote by $\sigma$ the curve consisting of (i) three line segments on
$y=-\delta$, from $(-3,-\delta)$ to the left endpoint of $\{y=-\delta\} \cap \partial K_{\mu_{1}}^{1}$, from the right endpoint of $\{y=-\delta\} \cap K_{\mu_{1}}^{1}$ to the left endpoint of $\{y=-\delta\} \cap \partial K_{\mu_{2}}^{2}$ and from the right endpoint of $(y=-\delta) \cap \partial K_{\mu_{2}}^{2}$ to ( $3,-\delta$ ), and (ii) the arcs of $\partial K_{\mu_{1}}^{i}$ lying in $\{|y|<\delta\}$.

Denote by $\Sigma_{-}$the part of $I$ lying below $\sigma$. Notice that $u<0$ in $\Sigma_{-}$.
From assumption (ii) and Harnack's inequality we get

$$
\begin{equation*}
u(X) \leqslant-c \eta \operatorname{dist}(X, \sigma) \quad \text { if } X \in \Sigma_{-}(c>0) \tag{8.3}
\end{equation*}
$$

We next claim
there exists a $C^{1}$ curve $\sigma_{i}: y=f_{i}(x)$ in $I$ such that $E_{i} \in \sigma_{i}$ and $u>0$ above $\sigma_{i}$ (in $I$ ), for $i=1,2$; furthermore, $f_{i}^{\prime}(x)-f_{i}^{\prime}\left(x_{i}\right)$ $\rightarrow 0$ as $x-x_{i} \rightarrow 0$, uniformly with respect to $u$.

Notice that $\sigma$ and $E_{i}$ depend on $u$ and so does $f_{i}$. To prove (8.4) suppose first that there exist sequences $E_{1}=E_{1}(m)=\left(x_{1}(m), y_{1}(m)\right), u=u_{m}$ and $Z_{m}=\left(\tilde{x}_{m}, \tilde{y}_{m}\right)$ with $u_{m}\left(Z_{m}\right) \leqslant 0$, such that $\left|Z_{m}-E_{1}(m)\right| \rightarrow 0$ and the angle between $\overrightarrow{E_{1}(m) Z_{m}}$ and the tangent to $\sigma$ at $E_{1}(m)$ does not converge to zero as $m \rightarrow \infty$. Set $r_{m}=\left|Z_{m}-E_{1}(m)\right|$ and consider a blow-up sequence of $u_{m}$ with respect to $B_{r_{m}}\left(E_{1}(m)\right)$. Let $w$ be a blow-up limit. We can rotate the coordinates in such a way that

$$
\begin{equation*}
w(x, y) \leqslant 0 \quad \text { if } y \leqslant 0 \tag{8.5}
\end{equation*}
$$

and then $w\left(x_{0}, y_{0}\right) \leqslant 0$ for some point $\left(x_{0}, y_{0}\right)$ with $y_{0}>0$. Consequently,

$$
\begin{equation*}
w \text { is not a 2-plane solution. } \tag{8.6}
\end{equation*}
$$

In view of (8.3) and the assumption $\Lambda<0, w$ does have two phases.
By Corollary 6.5 and Lemma 6.6,

$$
\begin{equation*}
w(X)=\alpha y+o(|X|) \quad \text { if } y<0,|X| \rightarrow 0 \tag{8.7}
\end{equation*}
$$

where $\alpha$ is determined by $\alpha^{2}\left(\alpha^{2}+|\Lambda|\right)=\gamma, \gamma=\lim _{r \rightarrow 0} \psi(r)$, where

$$
\psi(r)=\frac{1}{r^{4}} \int_{B_{r}}\left|\nabla w^{+}\right|^{2} \cdot \int_{B_{r}}\left|\nabla w^{-}\right|^{2} .
$$

Similarly, working with blow-up sequences $\frac{1}{m} w(m X)(m \rightarrow \infty)$ we find that

$$
\begin{equation*}
w(X)=\beta y+o(|X|) \text { if } y<0,|X| \rightarrow \infty \tag{8.8}
\end{equation*}
$$

where $\beta^{2}\left(\beta^{2}+|\Lambda|\right)=\gamma_{0}, \gamma_{0}=\lim _{r \rightarrow \infty} \psi(r)$. Since, by (8.6), $w$ is not a 2-plane solution, Lemma 6.6 shows that $\gamma_{0}>\gamma$ and, consequently,

$$
\begin{equation*}
\beta>\alpha \tag{8.9}
\end{equation*}
$$

Let $\Omega_{R}=\{w<0\} \cap B_{R}$. If we formally apply Green's formula to $w$ and $G=$ $y /\left(x^{2}+y^{2}\right)-y / R^{2}$ in $\Omega_{R} \backslash B_{\varepsilon}$, we obtain

$$
\begin{equation*}
\int_{\partial \Omega_{R \backslash B}}\left(G w_{\nu}-w G_{\nu}\right)+\int_{\partial B_{\varepsilon} \cap\{w<0\}}\left(G w_{\nu}-w G_{\nu}\right)=0 \tag{8.10}
\end{equation*}
$$

where $\nu$ is the inner normal. In order to justify (8.10) and make sense of the integrals over the free boundary we apply (7.6) with $u^{-}=w^{-}$and $\zeta=\eta G$ where $\eta=\eta(r)$ is given by

$$
\eta(r)= \begin{cases}1 & \text { if } r \leqslant R \\ 1-(r-R)^{2} / \delta^{2} & \text { if } R<r<R+\delta \\ 0 & \text { if } r>R+\delta\end{cases}
$$

and then let $\delta \rightarrow 0$. We then obtain (8.10) with $\int G_{\nu} w=0$ on the free boundary and $\int G w_{\nu}=-\int q_{\nu}^{-} G d H^{1}$ on the free boundary. By (8.5), $G \geqslant 0$ on the free boundary and therefore the last integral is nonnegative. We thus conclude from (8.10) that

$$
\begin{equation*}
\int_{\partial B_{r} \cap\{w<0\}} w G_{\nu} \leqslant \int_{\partial B_{f} \cap\{w<0\}}\left(G w_{\nu}-w G_{\nu}\right) . \tag{8.11}
\end{equation*}
$$

Using (8.8) we compute that

$$
\int_{\partial B_{R} \cap\{w<0\}} w G_{\nu}=\int 2 \frac{\sin \theta}{R^{2}}(\beta y+o(R))=2 \beta \int \sin ^{2} \theta d \theta+\eta(R)
$$

where $\eta(R) \rightarrow 0$ if $R \rightarrow \infty$. Similarly,

$$
-\int_{\partial B_{e} \cap\{w<0\}} w G_{\nu}=\alpha \int \sin ^{2} \theta d \theta+\eta_{0}(\varepsilon)
$$

where $\eta_{0}(\varepsilon) \rightarrow 0$ if $\varepsilon \rightarrow 0$. Finally, for a sequence $\varepsilon_{i} \rightarrow 0$ we have

$$
\int_{\partial B_{i} \cap\{w<0\}} G w_{\nu}=\alpha \int \sin ^{2} \theta d \theta+\eta_{1}\left(\varepsilon_{i}\right)
$$

with $\eta_{1}\left(\varepsilon_{i}\right) \rightarrow 0$ if $\varepsilon_{i} \rightarrow 0$. Indeed, this follows from

$$
\begin{aligned}
\frac{1}{\bar{\varepsilon}} \int_{0}^{\bar{\varepsilon}} d \varepsilon \int G w_{\nu} d s & =\frac{1}{\bar{\varepsilon}} \int \sin \theta[w]_{r=\bar{\varepsilon}} d \theta+\frac{o(\bar{\varepsilon})}{\bar{\varepsilon}} \\
& =\alpha \int \sin ^{2} \theta d \theta+\frac{o(\bar{\varepsilon})}{\bar{\varepsilon}}
\end{aligned}
$$

Using the preceding estimates in (8.11) we get

$$
2 \alpha \int_{\pi+o(1)}^{2 \pi+o(1)} \sin ^{2} \theta \geqslant 2 \beta \int_{\pi}^{2 \pi} \sin ^{2} \theta+o(1)
$$

where $o(1) \rightarrow 0$ if $\varepsilon \rightarrow 0, R \rightarrow \infty$; this contradicts (8.9).
We have thus proved that there cannot exist sequences $E_{1}(m), Z_{m}, u_{m}$ as above. It follows that, for each $u,\{u<0\} \cap\left\{x>x_{1}\right\}$ lies below a polygonal curve $\pi_{0}$ with sides $\overline{Z^{m} Z^{m+1}}$ having slope $\phi_{m}$ which decreases to the slope $\phi_{\infty}$ of $\sigma$ at $E_{1}$, uniformly with respect to $u$, as $\left|Z^{m}-E_{1}\right| \rightarrow 0$. We modify $\pi_{0}$ near its vertices so as to obtain a $C^{1}$ curve $y=f_{1}(x)$ lying above $\pi_{0}$ with slope converging to $\phi_{\infty}$ as $x \downarrow x_{1}$. Similarly, we can construct $y=f_{1}(x)$ for $x<x_{2}$, and this completes the construction of $\sigma_{1}$ as asserted in (8.4). $\sigma_{2}$ is constructed similarly.

Remark 8.1. The assertion (8.4) remains valid also if condition (ii) is dropped. Indeed, if in the previous proof $w$ is a 2-phase solution, then the proof is the same. If, on the other hand, $w$ is a 1-phase solution (and then $w \geqslant 0$ since $\Lambda<0$ ) then we
get a contradiction to Lemma 8.4 below; Lemma 8.4 is proved independently of Lemmas $8.1-8.3$. This remark will be used in proving Lemma 8.11 (which is an extension of Lemma 8.2 to the case where condition (ii) is dropped).

Now consider in the strip $I^{\delta}=I \cap\{|y|<\sqrt{\delta}\} \cap\left\{x_{1}<x<x_{2}\right\}$ the quotient difference $\Delta_{h . l} u$ of $u$ in the direction $l$ of $\varepsilon e_{2} \pm e_{1}$, with increment $h$, where $0<h<2 \delta$, i.e.,

$$
\Delta_{h, l} u(X)=(u(X+h l)-u(X)) / h .
$$

We claim that

$$
\begin{equation*}
\Delta_{h, l} u \geqslant c>0 \tag{8.12}
\end{equation*}
$$

in $I^{\delta}$. We first prove (8.12) on $y=\sqrt{\delta}$. If the assertion is not true then for sequences $u_{m}, X_{m}=\left(x_{m}, \delta_{m}^{1 / 2}\right)$ with $\delta_{m} \rightarrow 0$ there holds $\Delta_{h_{m}, I_{m}} u_{m}\left(X_{m}\right) \rightarrow 0$, with $0<h_{m}<2 \delta_{m}$ and $l_{m} \rightarrow l, l$ in direction $\varepsilon e_{2} \pm e_{1}, \varepsilon \geqslant \varepsilon_{0}$. Take a blow-up about free boundary points of $u_{m}$ on $\left\{x=x_{m}\right\}$ with radii $\leqslant 2 \delta_{m}^{1 / 2}$. Since the free boundary of $u_{m}$ lies in $\left\{|y|<\delta_{m}\right\}$, the blow-up limit $w$ is a 2-plane solution (we use here (8.3) and the assumption $\Lambda<0$ ) and its free boundary is the $x$-axis. Since

$$
\Delta_{h_{m}, I_{m}} u_{m}\left(X_{m}\right) \rightarrow \partial w\left(X_{0}\right) / \partial l
$$

where $X_{0}=(0,1)$ and $\varepsilon \neq 0$, we get a contradiction.
Similarly, we can establish (8.12) on $y=-\sqrt{\delta}$. Consider now the quotient difference on the vertical line $V_{1}$ of $\partial I^{\delta}$ passing through $E_{1}$. If (8.12) does not hold on $V_{1}$, say

$$
\Delta_{h_{m}, l_{m}} u_{m}\left(X_{m}\right) \rightarrow 0 \quad\left(X_{m} \in V_{1}\right)
$$

then we make a blow-up about $E_{1}$ with radii $r_{m}=\max \left\{h_{m},\left|X_{m}-E_{1}\right|\right\}$. Recalling that near $E_{1}$ the free boundary lies between $\sigma$ and $\sigma_{1}$ and using (8.4), we again deduce that the blow-up limit $w$ is a 2 -plane solution, with the $x$-axis as the free boundary; Further,

$$
\begin{array}{ll}
\frac{\partial w}{\partial l}\left(X_{0}\right)=0 & \text { if } h_{m}=o\left(\left|X_{m}-E_{1}\right|\right), \\
\Delta_{h_{0}, I_{w}}\left(X_{0}\right)=0 & \text { if }\left|h_{m}\right| \geqslant c_{0}\left|X_{m}-E_{1}\right|\left(c_{0}>0\right)
\end{array}
$$

for some $h_{0}, X_{0}$ and some $l$ in direction $\varepsilon e_{2} \pm e_{1}, \varepsilon \geqslant \varepsilon_{0}$. But this is impossible since $w$ is a function of $y$ only.

Having proved (8.12) on $\partial I^{\delta}$ we now translate $u$ in the direction $l$ by considering

$$
u_{\tau}(X)=u(X+\tau l), \quad \tau>0 .
$$

In view of (8.12), $u_{\tau}>u$ on $\partial I^{\delta}$ if $0<\tau<2 \delta$. Since $q(x) \equiv 1, u_{\tau}$ is a minimizer for the same functional $J$ as $u$. Appealing to Lemma 8.1 we conclude that $u_{\tau}>u$ in $\{u \neq 0\}$, from which the assertion follows.

Lemma 8.3. Any global minimizer $u$ with two phases must be a 2-plane solution.
Proof. For a sequence $m \rightarrow \infty$ we have $u_{m}(X) \equiv u(m X) / m \rightarrow v(X)$ where $v(X)$ is a 2-plane solution. Indeed,

$$
\psi(r) \equiv \frac{1}{r^{4}} \int_{B_{r}}\left|\nabla u^{-}\right|^{2} \int_{B_{r}}\left|\nabla u^{+}\right|^{2} \uparrow \gamma \quad \text { as } r \uparrow \infty,
$$

and since $\psi\left(r_{0}\right)>0$ for some $r_{0}>0$ (since $u$ has two phases) it follows that $\gamma>0$. On the other hand, $v$ satisfies (6.13) and, by Lemma 6.6(ii),

$$
v(x, y)= \begin{cases}\mu_{1} y & \text { if } y>0 \\ \mu_{2} y & \text { if } y<0\end{cases}
$$

where $\mu_{1}>0, \mu_{2}>0, \mu_{1}^{2} \mu_{2}^{2}=\gamma$.
Given $\varepsilon_{0}>0$ and $\eta=\mu_{2} / 2$, if $m$ is large enough then the $u_{m}$ restricted to $I$ satisfy the conditions of Lemma 8.2 (recall that the lemma is valid uniformly with respect to the class of all minimizers $u$ ). Hence, the free boundary $\partial\left\{u_{m}>0\right\}$ (for $m \geqslant m\left(\varepsilon_{0}\right)$ ) in $I_{0}$ is a graph in the direction $( \pm 1, \varepsilon)$ for any $\varepsilon \geqslant \varepsilon_{0}$. It follows that $\partial\{u>0\} \cap$ $\{|x|<m,|y|<m\}$ is a graph in any direction $( \pm 1, \varepsilon)$ where $\varepsilon \geqslant \varepsilon_{0}$. Since $\varepsilon_{0}$ can be chosen arbitrarily small (and $m \geqslant m\left(\varepsilon_{0}\right)$ ), $\partial\{u>0\}$ must coincide with the $x$-axis. By uniqueness for the Cauchy-Kowalewski theorem $u$ is thus linear in $y$ for $y>0$ and for $y<0$.

Lemma 8.4. Any global minimizer $u$ with one phase must be a 1-plane solution.
Naturally, to exclude a trivial case we assume that $u \geqslant 0$ in $R^{2}$ with (say) $0 \in \partial\{u>0\}$ and with $\lambda>0$, where $J(u)=\int\left(|\nabla u|^{2}+\lambda^{2} I_{\{u>0\}}\right)$.

Proof. The function $|\nabla u|$ is subharmonic and $|\nabla u|=\lambda$ on $\partial\{u>0\}$. Proceeding as in [2] (see also [9, p. 327]) we deduce that $|\nabla u|$ takes its maximum on the free boundary and, consequently, the free boundary is convex to $\{u>0\}$. If $\partial\{u>0\}$ is not a straight line then the blow-up limit of a subsequence of $u_{m}(X)=u(m X) / m$ converges to a minimizer $v$ whose free boundary includes two rays forming an angle $\neq \pi$ at the origin; this contradicts the Cauchy-Kowalewski theorem since $u=0$, $\partial u / \partial \nu=0$ on each of these rays.

Lemma 8.5. For any $\gamma>0$ and $C_{0}>0$ there is a $\delta=\delta\left(\gamma, C_{0}\right)$ such that if $u$ is a minimizer in $B_{1}$ with $|\nabla u| \leqslant C_{0}$ then for any ball $B_{\delta}\left(X^{0}\right) \subset B_{1 / 2}$ with center in the free boundary, the $\gamma$-flatness condition holds, i.e., the free boundary of $u$ in $B_{\delta}\left(X^{0}\right)$ lies in a strip with center $X^{0}$ and width $2 \gamma$.

Proof. If the assertion is not true then there is a sequence $B_{\delta_{m}}\left(X_{m}\right) \subset B_{1 / 2}$ with $\delta_{m} \rightarrow 0$ such that the flatness condition does not hold for some $u_{m} ; X_{m} \in \partial\left\{u_{m}>0\right\}$. A blow-up sequence with respect to $B_{\delta_{m}}\left(X_{m}\right)$ is convergent to a minimizer $v$ in $R^{2}$ and the free boundary of $v$ in $B_{1}(0)$ does not lie in a ( $2 \gamma$ )-strip with 0 in the centerline of the strip. If $v$ has two phases, this contradicts Lemma 8.3, whereas if $v$ has one phase, Lemma 8.4 is contradicted.

Lemma 8.6. If $u$ satisfies the $\gamma$-flatness condition in $B_{1}=B_{1}(0)$ in direction $(0,1)$ and if

$$
\begin{equation*}
u(A)>M u(P) \quad \text { where } A=\left(0, \frac{1}{2}\right), P \in\{u>0\} \cap B_{1 / 2}, \tag{8.13}
\end{equation*}
$$

then, for some absolute constant $C, \operatorname{dist}(P, \partial\{u>0\})<2 \gamma+C / M$.
Proof. By the flatness assumption $u>0$ in $B_{1} \cap\{y>\varepsilon / 2\}$ for any $\varepsilon>2 \gamma$. Suppose $\operatorname{dist}(P, \partial\{u>0\})>\varepsilon$; then also $\operatorname{dist}(P,\{y<\varepsilon / 2\})>\varepsilon / 2$. Applying Harnack's inequality in $B_{1} \cap\{y>\varepsilon / 2\}$ we get $u(P)>c \varepsilon u(A)$. Hence, by (8.13), $1 / M>c \varepsilon$, i.e., $\varepsilon<1 / c M$.

Lemma 8.7. For $\gamma$ sufficiently small let $\delta=\delta\left(\gamma, C_{0}\right)$ be as in Lemma 8.5, and let $B_{\delta}\left(X_{0}\right)$ be any ball in $B_{1 / 2}$ with $X_{0}$ for which the $\gamma$-flatness holds in the direction $(0,1)$, say, and $u(A)>0$ where $A=X_{0}+(0, \delta / 2)$. Then

$$
\begin{equation*}
u(A) \geqslant \gamma \sup _{B_{\delta / 2}\left(X_{0}\right)} u \tag{8.14}
\end{equation*}
$$

Proof. Take, for simplicity, $X_{0}=0$ and normalize by taking $\delta=1$. Set $A_{0}=A$. If the assertion (8.14) is not true then there exists a point $P_{0} \in B_{1 / 2} \cap\{u>0\}$ such that

$$
\begin{equation*}
u\left(P_{0}\right)>\frac{1}{\gamma} u\left(A_{0}\right) \tag{8.15}
\end{equation*}
$$

By Lemma 8.6

$$
\begin{equation*}
\operatorname{dist}\left(P_{0},\{y>\gamma\}\right)<C_{0} \gamma \tag{8.16}
\end{equation*}
$$

Let $E$ be a point on the free boundary with

$$
\begin{equation*}
\left|E-P_{0}\right|<\left(C_{0}+2\right) \gamma \tag{8.17}
\end{equation*}
$$

By the $\gamma$-flatness about $E$, the direction of flatness $\nu_{E}$ at $E$ differs from the direction $(0,1)$ by at most $C \gamma$.

We fix $\eta$ small, to be determined later (independently of $\gamma$ ) and take $\gamma \ll \eta$. By Harnack's inequality in $B_{1} \cap\{y>\eta / 2\}$ we have

$$
\begin{equation*}
\frac{1}{N} u\left(A_{0}\right) \leqslant u(X) \leqslant N u\left(A_{0}\right) \quad \text { if } X \in B_{1 / 2} \cap\{y>\eta\} \tag{8.18}
\end{equation*}
$$

where $N=N(\eta)$. Denoting by $G$ the Green function for $-\Delta$ in $\tilde{B} \equiv B_{1 / 4}(E) \cap\{y>$ $-2 \gamma\}$, we can represent the subharmonic function $u^{+}$at $P_{0}$ in the form

$$
u\left(P_{0}\right)=-\int_{\partial \tilde{B}} \frac{\partial G}{\partial \nu} u^{+}=-\int_{S}-\int_{T}
$$

where $S=\partial B_{1 / 4}(E) \cap\{y>\eta\}$ and $T=\partial B_{1 / 4}(E) \cap\{y<\eta\} \cap\left\{(X-E) \cdot \nu_{E}>\right.$ $\gamma\}$ (notice that $u^{+}=0$ on $\partial \tilde{B} \cap\left\{(X-E) \cdot \nu_{E} \leqslant \gamma\right\}$ and, in particular, on $\partial \tilde{B} \cap\{y$ $=-2 \gamma\}$ ). Setting $\sigma(\eta)=\operatorname{meas}(T)$, we have $\sigma(\eta) \rightarrow 0$ if $\eta \rightarrow 0$.

By (8.17), $-\partial G\left(P_{0}, X\right) / \partial \nu \leqslant C \gamma$ if $X \in S \cup T$. Consequently, $-\int_{T} \leqslant$ $C \gamma \delta(\eta) \sup _{T} u^{+}$and (using (8.18))

$$
-\int_{S} \leqslant N u\left(A_{0}\right)(1-\sigma(\eta)) C \gamma
$$

Recalling (8.15) we conclude that

$$
\frac{1}{\gamma} u\left(A_{0}\right) \leqslant C u\left(P_{0}\right) \leqslant N C \gamma u\left(A_{0}\right)+C \gamma \sigma(\eta) \sup _{T} u^{+}
$$

Choosing $\eta$ such that $2 C \sigma(\eta)<1$ we find that, provided $N C \gamma / 2 \gamma$, there holds $u\left(A_{0}\right) / \gamma^{2} \leqslant \sup _{T} u^{+}$. Thus, there is a point $P_{1} \in T$ such that

$$
\begin{equation*}
u\left(P_{1}\right)>\frac{1}{\gamma^{2}} u\left(A_{0}\right) \tag{8.19}
\end{equation*}
$$

Let $A_{1}$ be the point in $B_{1 / 2}(E)$ such that $\overrightarrow{A_{1} E}$ is in the direction $-\nu_{E}$, with $\left|A_{1}-E\right|=1 / 8$. Then, by Harnack's inequality,

$$
\begin{equation*}
u\left(A_{1}\right) \leqslant N u\left(A_{0}\right)<\frac{1}{\gamma} u\left(A_{0}\right) \tag{8.20}
\end{equation*}
$$

with the same $N$ as before (if $\eta$ is small enough). The previous setting for $A_{0}, P_{0}$ occurs also for $A_{1}, P_{1}$ since, by (8.19) and (8.20),

$$
u\left(P_{1}\right)>\frac{1}{\gamma^{2} N} u\left(A_{1}\right)>\frac{1}{\gamma} u\left(A_{1}\right) .
$$

We can now repeat the previous proof with $0, A_{0}, P_{0}$ replaced by $E, A_{1}, P_{1}$ and $B_{1 / 2}(0)$ replaced by $B_{1 / 4}(E)$. Thus there is a triple $E_{2}, A_{2}, P_{2}$ such that

$$
u\left(P_{2}\right)>\frac{1}{\gamma} \frac{1}{\gamma^{2} N} u\left(A_{1}\right)
$$

and $u\left(P_{2}\right)>u\left(A_{2}\right) / \gamma>0, E_{2} \in \partial\{u>0\}, \overrightarrow{E_{2} A_{2}}$ is in the direction $\nu_{E_{2}}$ of $\gamma$-flatness at $E_{2},\left|A_{2}-E_{2}\right|=\frac{1}{2} \cdot \frac{1}{4}$.

Continuing in this way, step by step, we construct a sequence ( $E_{n}, A_{n}, P_{n}$ ) such that

$$
\begin{equation*}
u\left(P_{n}\right)>\frac{1}{\gamma^{n+1}} \frac{1}{N^{n-1}} u\left(A_{n-1}\right) \tag{8.21}
\end{equation*}
$$

and $u\left(P_{n}\right)>u\left(A_{n}\right) / \gamma>0, u\left(E_{n}\right)=0, \overrightarrow{E_{n} A_{n}}$ is in the direction $\nu_{E_{n}}$ of $\gamma$-flatness about $E_{n},\left|A_{n}-E_{n}\right|=\frac{1}{2} 2^{-n}$. Recall that, by Harnack's inequality, $u\left(A_{1}\right)>u\left(A_{0}\right) / N$. Since the configuration of each pair $A_{n}, A_{n-1}$, with respect to the free boundary, is similar (after scaling) to that of $A_{1}, A_{0}$ (using the $\gamma$-flatness in each ball $B_{2^{-n}}\left(E_{n}\right)$ and the fact that the directions $\nu_{E_{n}}, \nu_{E_{n-1}}$ differ by at most $C \gamma / 2^{n}$ ), we also have, by Harnack's inequality, $u\left(A_{n}\right)>u\left(A_{n-1}\right) / N$ (with $N$ independent of $n$ ). Recalling (8.21) we obtain

$$
u\left(P_{n}\right) \geqslant \frac{1}{N^{n-1}} \frac{1}{\gamma^{n+1}} \frac{1}{N^{n-1}} u\left(A_{0}\right)=\frac{1}{\gamma^{2}} \frac{1}{\left(\gamma N^{2}\right)^{n-1}} u(A)
$$

Choosing $\gamma<N^{2}$ we conclude that $u\left(P_{n}\right) \rightarrow \infty$ if $n \rightarrow \infty$, which is impossible. This completes the proof of (8.14).

Lemma 8.7 extends to $u^{-}$, that is, if $A_{*}=X_{0}-(0, \delta / 2)$ then

$$
\begin{equation*}
u\left(A_{*}\right)<0, \quad u^{-}\left(A_{*}\right)>\gamma \sup _{B_{\delta / 2}\left(X_{0}\right)} u^{-} . \tag{8.22}
\end{equation*}
$$

Corollary 8.8. If $\gamma$ is small enough, say $\gamma<\gamma_{0}$, then

$$
\begin{equation*}
\int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{+}\right|^{2} \leqslant C\left(u^{+}(A)\right)^{2}, \quad \int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{-}\right|^{2} \leqslant C\left(u^{-}\left(A_{*}\right)\right)^{2} \tag{8.23}
\end{equation*}
$$

where $C=C\left(\gamma_{0}\right)$ and $R=\delta\left(\gamma_{0}\right) / 4$.

Indeed, introducing $G(X)=\log 2 R /\left|X-X_{0}\right|$ in $B_{2 R}\left(X_{0}\right)$, we have, by Green's formula,

$$
\begin{aligned}
-\int_{\partial B_{2 R}\left(X_{0}\right)}\left(u^{ \pm}\right)^{2} \frac{\partial G}{\partial \nu} & =\iint_{B_{2 R}\left(X_{0}\right)}\left[\Delta\left(u^{ \pm}\right)^{2} G-\left(u^{ \pm}\right)^{2} \Delta G\right] \\
& =2 \iint_{B_{2 R}\left(X_{0}\right)}\left|\nabla u^{ \pm}\right|^{2} G \geqslant c \iint_{B_{R}\left(X_{0}\right)}\left|\nabla u^{ \pm}\right|^{2},
\end{aligned}
$$

and the left-hand side is estimated by (8.14) and (8.22).
Lemma 8.9. If $X_{0} \in \partial\{u>0\}$ and

$$
\limsup _{x \rightarrow X_{0}}\left|\nabla u^{-}(X)\right|=\alpha, \quad \underset{X \rightarrow x_{0}}{\limsup }\left|\nabla u^{+}(X)\right|=\beta,
$$

then

$$
\begin{align*}
& \alpha^{2}\left(\alpha^{2}+|\Lambda|\right) \leqslant \frac{1}{R^{4}} \int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{-}\right|^{2} \cdot \int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{+}\right|^{2},  \tag{8.24}\\
& \beta^{2}\left(\beta^{2}-|\Lambda|\right) \leqslant \frac{1}{R^{4}} \int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{-}\right|^{2} \cdot \int_{B_{R}\left(X_{0}\right)}\left|\nabla u^{+}\right|^{2} \tag{8.25}
\end{align*}
$$

here $u$ is any minimizer in $B_{R}\left(X_{0}\right)$.
Proof. It suffices to prove (8.24). We take $X_{0}=0$ and $X_{n} \rightarrow 0$ with $\left|\nabla u^{-}\left(X_{n}\right)\right| \rightarrow \alpha$; we may suppose that $\alpha>0$. Let $Y_{n}$ be the nearest point to $X_{n}$ on the free boundary. Consider a blow-up sequence with respect to $B_{r_{n}}\left(Y_{n}\right), r_{n}=\left|X_{n}-Y_{n}\right|$. Since $\alpha>0$ and $\Lambda<0$, the blow-up limit has two phases and, by Lemma 8.3, it is a 2-plane solution with slopes $\alpha$ and $\bar{\alpha}$ satisfying $\alpha^{2}-\bar{\alpha}^{2}=\Lambda$. It easily follows that, as $\varepsilon \rightarrow 0$,

$$
\frac{1}{\varepsilon^{2}} \int_{B_{\varepsilon}}\left|\nabla u^{-}\right|^{2} \rightarrow \alpha^{2}, \quad \frac{1}{\varepsilon^{2}} \int_{B_{\varepsilon}}\left|\nabla u^{+}\right|^{2} \rightarrow \bar{\alpha}^{2} .
$$

The assertion (8.24) now follows using the monotonicity lemma.
Lemma 8.10. Under the conditions of Corollary 8.8 (with $\gamma<\gamma_{0}$ ),

$$
\begin{equation*}
\left|\nabla u^{-}(X)\right| \leqslant C u^{-}\left(A_{*}\right)\left(u^{+}(A)+1\right) \quad \text { in } B_{R / 2}\left(X_{0}\right), \tag{8.26}
\end{equation*}
$$

where $C$ is a constant (depending on $\gamma_{0}$ ).
Proof. The function $w=\left|\nabla u^{-}\right|$is subharmonic. By Lemma 8.9 and Corollary 8.8,

$$
\limsup _{X \rightarrow X_{0}} w(X) \leqslant C\left[u^{+}(A) u^{-}\left(A_{*}\right)\right]^{1 / 2}
$$

where $A=A\left(X_{0}\right), A_{*}=A_{*}\left(X_{0}\right)$. If $\gamma$ is small enough then by Harnack's inequality

$$
\begin{equation*}
u^{+}\left(A\left(X_{0}\right)\right) \leqslant C u^{+}(A), \quad u^{-}\left(A_{*}\left(X_{0}\right)\right) \leqslant C u^{-}\left(A_{*}\right) \tag{8.27}
\end{equation*}
$$

where $A, A_{*}$ correspond to the free boundary point 0 and $X_{0} \in B_{R}$. Hence,

$$
\begin{equation*}
\limsup _{X \in B_{R}, \operatorname{dist}(X, \partial(u>0)) \rightarrow 0} w(X) \leqslant C\left[u^{+}(A) u^{-}\left(A_{*}\right)\right]^{1 / 2} \tag{8.28}
\end{equation*}
$$

On the other hand, by Corollary 8.8 and (8.27),

$$
\int_{B_{R} \cap\{u<0\}} w^{2} \leqslant C\left(u^{-}\left(A_{*}\right)\right)^{2}
$$

Set $W=\max \left\{w, C\left[u^{+}(A) u^{-}\left(A_{*}\right)\right]^{1 / 2}\right\}$ in $B_{R}$. By (8.28), $W$ is a continuous subharmonic function and, therefore, by elliptic estimates,

$$
W^{2}(X) \leqslant C \int_{B_{R}} W^{2} \leqslant C u^{+}(A) u^{-}\left(A_{*}\right)+C\left(u^{-}\left(A_{*}\right)\right)^{2}
$$

and (8.26) follows.
Lemma 8.11. Lemma 8.2 remains true without the assumption (ii).
Proof. It suffices to establish that

$$
\begin{equation*}
\Delta_{h, l} u>0 \quad \text { on } \partial I^{\delta} \tag{8.29}
\end{equation*}
$$

for all $h, l, u$. Suppose this is not true for a sequence $u_{m}$ with $X=X_{m}, h=h_{m}$, $l=l_{m}$. If the intervals $\tilde{l}_{m}:\left(X_{m}, X_{m}+h_{m} l_{m}\right)$ lie in $\left\{u_{m}>0\right\}$ then we can proceed as before. Indeed, the blow-up limit $w$ with respect to $B_{\delta_{\delta^{\prime 2}}}\left(X_{m}\right)$ (or $B_{r_{m}}\left(E_{i}\right), E_{i}$ depends on $m$ ) is either a 1-plane solution with $w \geqslant 0($ since $\Lambda<0)$ or a 2-plane solution and its free boundary is $\{y=0\}$ (here we use Remark 8.1); thus we get a contradiction as before.

If $\tilde{l}_{m}$ lies in $\left\{u_{m}<0\right\}$ and if a blow-up limit $w$ turns out to be a 1-plane solution with $w=0$ if $\{y<0\}$, we do not get a contradiction. In order to derive a contradiction we shall work with $U_{m}=u_{m} / u_{m}^{-}\left(A_{*}\right)$ instead of $u_{m}$, where $A_{*}$ is chosen as in Lemma 8.10 ( $A_{*}$ depends on $m$ ). Then $U_{m}\left(A_{*}\right)=-1$ and $U_{m}^{-}$is uniformly Lipschitz continuous (by Lemma (8.10)). Taking a blow-up limit $W$ of $U_{m}^{-}$ with respect to $B_{\delta_{m}^{\prime 2}}\left(X_{m}\right)$ (or $B_{r_{m}}\left(E_{i}\right)$ ) we find that the free boundary of $W$ is $\{y=0\}$; hence, by Liouville's theorem (reflecting first $W$ across $\{y=0\}$ ) $W \equiv c y$ if $y<0(c>0)$, and therefore, $\Delta_{h_{m}, I_{m}} U_{m} \geqslant c$ uniformly with respect to $\tilde{l}_{m}$ in $\left\{u_{m}<0\right\}$, that is,

$$
\begin{equation*}
\Delta_{h_{m}, l_{m}} u_{m}>c u_{m}\left(A_{*}\right)>0 \tag{8.30}
\end{equation*}
$$

uniformly with respect to $h_{m}, l_{m}, X_{m}$.
It remains to establish uniform positivity (in the sense of (8.30)) in case $\tilde{l}_{m}$ lies partially in $\left\{u_{m}>0\right\}$ and partially in $\left\{u_{m}<0\right\}$. In this case we can write it as a disjoint union of intervals $\tilde{l}_{m}=l_{m}^{1}+l_{m}^{2}+l_{m}^{3}$ where $l_{m}^{1} \subset\left\{u_{m}>0\right\}, l_{m}^{2} \subset\left\{u_{m}<0\right\}$ and $l_{m}^{3}$ is an interval with endpoints on $\sigma$ and $\sigma_{i}$. By Remark 8.1, meas $\left(l_{m}^{3}\right)=o\left(h_{m}\right)$ and thus either meas $\left(l_{m}^{1}\right)>c h_{m}$, or meas $\left(l_{m}^{2}\right)>c h$, or both inequalities hold. By the previous arguments for $\tilde{l}_{m}$ in $\left\{u_{m}>0\right\}$ and for $\tilde{l}_{m}$ in $\left\{u_{m}<0\right\}$ we deduce that the incremental quotients $\Delta_{l_{m}^{\prime}} u$ with respect to $l_{m}^{i}$ satisfy

$$
\begin{aligned}
& \Delta_{l_{m}^{\prime}} u \geqslant c \operatorname{meas}\left(l_{m}^{1}\right) / \operatorname{meas}\left(\tilde{l}_{m}\right), \\
& \Delta_{l_{m}^{2}} u \geqslant c u_{m}\left(A_{*}\right) \operatorname{meas}\left(l_{m}^{2}\right) / \operatorname{meas}\left(\tilde{l}_{m}\right) .
\end{aligned}
$$

Since also $\Delta_{l_{m}} u \geqslant 0$, the assertion (8.29) holds. We can now proceed as in Lemma 8.2 to complete the proof of Lemma 8.11.

Theorem 8.12. The free boundary $\partial\{u>0\} \cap \Omega$ is continuously differentiable.
Proof. By Lemma 8.5, for any small $\gamma>0$ there is a $\delta=\delta(\gamma)>0(\delta \downarrow 0$ if $\gamma \downarrow 0)$ such that the $\gamma$-flatness condition holds in every ball $B_{\delta}$ with center in the free boundary. Take such a ball $B_{\delta}$ and suppose for simplicity that its center is at the origin and that the flatness direction is $(0,1)$. By Lemma 8.11 the free boundary in $B_{\delta / 2}$ has the form $y=f(x)$ with $f(x)$ Lipschitz continuous.

Denote by $\gamma=\gamma(\delta)$ the inverse of the function $\delta=\delta(\gamma)$.
Take $x_{1}, x_{2}$ in $(-\delta / 4, \delta / 4)$ and set

$$
r=\left|x_{1}-x_{2}\right|, \quad X_{i}=\left(x_{i}, f\left(x_{i}\right)\right), \quad B_{i}=B_{2 r}\left(X_{i}\right) .
$$

Each $X_{i}$ must lie in the flatness strip of the disc $B_{j}(j \neq i)$. Therefore, the angles between the directions of flatness at $X_{1}$ and $X_{2}$ are bounded by $C \gamma(r)$. It follows that $\left|f^{\prime}\left(x_{1}\right)-f^{\prime}\left(x_{2}\right)\right| \leqslant C \gamma(r)$ for any two points $x_{1}, x_{2}$ where $f(x)$ is differentiable. Thus $f^{\prime}(x)$ has a continuous version.

The next result is concerned with the continuity of the normal derivative of $u$. Letting

$$
\gamma=\lim _{r+\infty} \frac{1}{r^{4}} \int_{B_{r}\left(X_{0}\right)}\left|\nabla u^{-}\right|^{2} \cdot \int_{B_{r( }\left(X_{0}\right)}\left|\nabla u^{+}\right|^{2}
$$

where $X_{0}$ is a free boundary point, we define $\beta=\beta(\gamma)>0$ by $\beta^{2}\left(\beta^{2}-|\Lambda|\right)=\gamma$ and denote by $\nu=\nu_{X_{0}}$ the normal to the free boundary at $X_{0}$ (pointing into \{ $u>0\}$ ).

Theorem 8.13. For any sector

$$
\Sigma_{c}=\left\{X ;\left(X-X_{0}\right) \cdot \nu>c\left|X-X_{0}\right|\right\}, \quad c>0
$$

there holds $u_{\nu}(X) \rightarrow \beta$ if $X \in \Sigma_{c}, X \rightarrow X_{0}$.
Proof. Let $X_{m} \in \Sigma_{c}, X_{m} \rightarrow X_{0}$ and take a blow-up sequence $u_{m}$ with respect to $B_{\left|X_{m}-X_{0}\right|}\left(X_{0}\right)$. Then $u_{m}(X) \rightarrow v(X)=\beta y(y>0)$ and $\partial u_{m}\left(X_{m}\right) / \partial \nu \rightarrow \partial v\left(Y_{0}\right) / \partial y$ since $Y_{0}$ lies in $\{y>0\}$. Since $\partial v\left(Y_{0}\right) / \partial y=\beta$, the assertion follows.

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