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Water oxidation and electron extraction kinetics in nanostructured tungsten trioxide photoanodes

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Abstract

A thorough understanding of the kinetic competition between desired water oxidation/electron extraction processes and any detrimental surface recombination is required to achieve high water oxidation efficiencies in transition metal oxide systems. The kinetics of these processes in high Faradaic efficiency tungsten trioxide (WO₃) photoanodes (>85%) are monitored herein by transient diffuse reflectance spectroscopy and correlated with transient photocurrent data for electron extraction. Under anodic bias, efficient hole transfer to the aqueous electrolyte is observed within a millisecond. In contrast, electron extraction is found to be comparatively slow (~ 10 ms), increasing in duration with nanoneedle length. The relative rates of these water oxidation and electron extraction kinetics are shown to be reversed in comparison to other commonly examined metal oxides (*e.g.* TiO₂, α -Fe₂O₃ and BiVO₄). Studies conducted as a function of applied bias and film processing to modulate oxygen vacancy density indicate that slow electron extraction kinetics result from electron trapping in shallow WO₃ trap states associated with oxygen vacancies. Despite these slow electron extraction kinetics, charge recombination losses on the microsecond to second timescales are observed to be modest compared to other oxides studied. We propose that the relative absence of such recombination losses, and the observation of a photocurrent onset potential close to flat-band, result directly from the faster water oxidation kinetics of WO₃. We attribute these fast water oxidation kinetics to the highly oxidising valence band position of WO₃, thus highlighting the potential importance of thermodynamic driving force for catalysis in outcompeting detrimental surface recombination processes.

Introduction

Photoelectrochemical water splitting is a promising strategy for sustainable fuel generation, but the limiting kinetics of the four-hole water oxidation reaction remain a key challenge to unlocking higher efficiencies.¹ Since the pioneering study on TiO₂ by Fukushima and Honda in the 1970s,² transition metal oxides have remained amongst the most widely studied photoanodes for water oxidation due to their earth abundance, good aqueous stability and facile synthesis.³⁻⁷ However, most such oxides, including TiO₂, α -Fe₂O₃, and BiVO₄, demonstrate sluggish water oxidation kinetics, allowing competing recombination processes to drastically lower performance.⁸⁻¹⁰ Heterojunctions, catalysts or overlayers are then required to accelerate water oxidation kinetics and/or slow recombination, increasing cost and complexity. With a narrower band gap than TiO₂, tungsten trioxide (WO₃) can absorb a larger proportion of the solar spectrum, while maintaining a deep valence band energy to provide a large thermodynamic driving force for water oxidation.^{11,12} WO₃ demonstrates good photocurrent onset and is also stable at acidic pH which allows preferential pairing with acid-stable photocathodes.^{11,13,14} While there exists extensive research on the semiconducting properties of WO₃, particularly with respect to photo- and electrochromic phenomena, the kinetics of water oxidation on this metal oxide remain, to the best of our knowledge, largely unexplored.¹⁵ In this study, we measure the kinetics of water oxidation and electron extraction in nanostructured WO₃ photoanodes, and compare these with other metal oxide photoanode materials. We reveal how these complementary processes of charge removal from the system compete with recombination on the microsecond to second timescale and enable the achievement of photocurrent onset potentials close to flat-band without added co-catalysts.

WO₃ is often reported to have high electrical conductivity,^{5,14,16} leading to its frequent implementation as an electron transport layer in various heterojunction photoanodes, including junctions formed with α -Fe₂O₃ and BiVO₄, in which conductivity is significantly poorer.^{11,12,17-21} It is proposed that high electrical conductivity improves water splitting performance by aiding the separation of photogenerated charges such that competing recombination processes are retarded.²² High conductivity in WO₃ is considered to be a result of a large density of charge carriers, caused by intrinsic oxygen vacancies, which act as n-type dopants.²³⁻²⁵ Carrier concentrations in the range of 10¹⁷ to 10²² cm⁻³ have been reported, depending on stoichiometry (WO_{3-x}).^{11,26,27} Similarly, changes in the stoichiometry of single crystals have been shown to modulate electrical conductivity between 10 – 10⁻⁴ S/cm.²³ However, such deviations from stoichiometry have also been suggested to introduce chemical defects that can result in increased charge trapping which, rather than boost performance, can often introduce additional recombination pathways.²⁸⁻³² For example, the presence of defect states in the band gap of α -Fe₂O₃ has been shown to enhance electron-hole recombination kinetics that can outcompete water oxidation.³³ In addition, high carrier densities can also

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5 result in shorter minority carrier lifetimes and thus also accelerated recombination losses.^{34,35} Water
6 oxidation is a kinetically slow process, requiring the accumulation of multiple holes. As such, kinetic
7 competition between water oxidation/electron extraction and surface charge recombination (sometimes
8 called back electron-hole recombination) can severely limit the efficiency of photoelectrochemical water
9 oxidation.³⁶⁻³⁹ Using sacrificial reagents, Amano et al found that the size of suspended WO₃ particles
10 affected the frequency of surface recombination with concurrent effects on oxygen evolution rate.⁴⁰
11 However, this competition has received limited attention to date for WO₃ photoanodes under
12 photoelectrochemical conditions.
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17 Reported herein is an analysis the dynamics of photogenerated charges in tungsten oxide photoanodes
18 synthesised by chemical vapour deposition (CVD). The synthesis method generates highly oxygen deficient
19 monoclinic WO_{3-x} nanoneedle structured films that are intensely blue in colour, which are then annealed in
20 air at elevated temperature to remove most of the oxygen vacancies. These photoanodes exhibit an early
21 photocurrent onset potential (*ca* 0.5 V vs RHE) and a high Faradaic efficiency for water oxidation (>85%).
22 The rate of water oxidation is analysed by transient diffuse reflectance spectroscopy and transient
23 photocurrent measurements and is found to be faster than other commonly used transition metal oxides,
24 and much faster than electron extraction. The effect of interband defects on electron transport is then
25 investigated. We find that electron trapping in the WO₃ needles is significant and propose a trap-mediated
26 mechanism of electron transport to the back contact (FTO). To examine how the density of traps may affect
27 this electron transport, we alter the annealing time and the needle length, and find that this has profound
28 effects on extraction time. We thus provide unique kinetic insight on the charge carriers of WO₃, a material
29 growing in popularity for solar fuel related applications,⁴¹⁻⁴⁵ and highlight the direct impact these kinetics
30 have on photoelectrochemical performance.
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Results

Sample Characterisation

Tungsten oxide samples were prepared by aerosol assisted chemical vapour deposition (AA-CVD) according to a previously established synthesis,⁴⁶ as detailed in the Supporting Information (SI, Methods). Figure 1a shows the cross-sectional SEM image of an annealed WO₃ sample, indicating the 4-micron long needles (yellow line) on top of the dense seed layer (400 nm, red line). X-ray diffraction data confirms the crystalline monoclinic nature of these WO₃ films, with the (002) peak most prominent due to the preferential growth along this axis (Supporting information, S1).

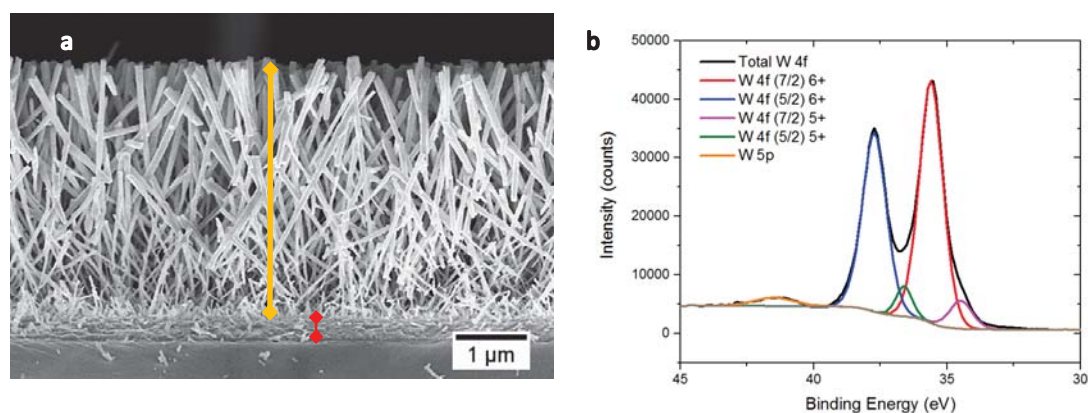


Figure 1: (a) Cross-sectional SEM image of annealed WO₃ needles (3.6 μm, yellow line) on 400 nm dense WO₃ seed layer (red line) deposited on FTO coated glass, 53 K magnification. (b) High resolution W 4f XPS scan indicating a small proportion of W⁵⁺ states.

Samples were also characterised by X-ray photoelectron spectroscopy (XPS). The XPS spectra of WO₃ samples before and after annealing indicate the films are of similar purity, with adventitious carbon the only detectable impurity (Figure S2b). The W 4f peaks in the un-annealed, blue sample show significant shoulders at lower binding energy as a result of a large number of W⁵⁺ species present, which can be directly correlated with oxygen vacancies (Figure S2c). The proportion of W⁵⁺ to W⁶⁺ analysed with this surface sensitive technique is approximately 30%. Figure 1b shows that after annealing in air at 500 degrees for 2 hours, the number of these reduced states is dramatically decreased to approximately 10% of the tungsten atoms probed. Photographs of the annealed and un-annealed samples are given in Figure S2a.

Photoelectrochemical water splitting performance

The current-voltage characteristics of annealed WO_3 were examined by linear sweep voltammetry under chopped simulated 1 sun irradiation from the front side (EE), as shown in Figure 2. As highlighted in Figure S3, the dark current was negligible over the voltage range examined, as was the photocurrent for the unannealed sample. As shown in Figure 2, the photocurrent reaches a plateau around $1.23 \text{ V}_{\text{RHE}}$ ($\sim 0.4 \text{ mA/cm}^2$), and the onset potential (*ca* $0.5 \text{ V}_{\text{RHE}}$) is only approximately 100 mV above the calculated flat-band for such nanostructured WO_3 .⁴⁷ Unlike photoanodes of several other transition metal oxides, there is an absence of transient spikes when the light is turned on (positive spikes) or off (negative spikes). Such photocurrent transients are indicative of current loss from recombination processes at the surface, typically the recombination of surface accumulated holes with bulk electrons, which compete with hole transfer to the electrolyte (e.g. by water oxidation).^{36,48} These transients are readily observed in $\alpha\text{-Fe}_2\text{O}_3$ and BiVO_4 , for which this recombination can be a significant loss pathway.^{49,50} The early photocurrent onset observed for the WO_3 needles combined with the lack of transient current spikes suggests that there is very little such surface recombination occurring, which may be a result of slower surface recombination or a faster reaction of accumulated holes, as we discuss further below. Oxygen evolution was measured with a calibrated Clark electrode as detailed in the SI. Figure S4 demonstrates the high Faradaic efficiencies obtained which were consistently between 85-90%. We note this to be higher than a large proportion of the Faradaic efficiency values reported in the literature for this material.^{14,25,39} This high Faradaic efficiency may be related to the high crystallinity of the WO_3 needles and/or preferential crystal orientation, reported by others to be significant.^{40,51} Further analysis on this aspect is, however, beyond the scope of this kinetic study.

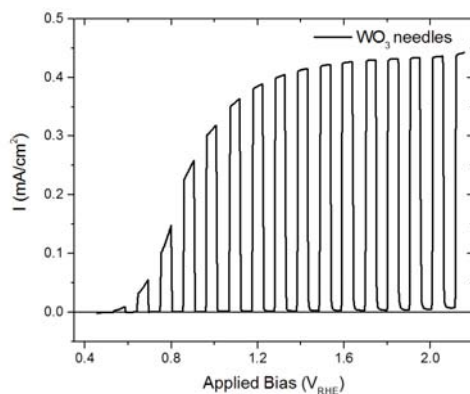


Figure 2: Chopped light linear sweep voltammetry, in $0.5 \text{ M H}_2\text{SO}_4$, of WO_3 needles under 1 sun conditions, indicating early photocurrent onset and reasonable maximum. Details of the 3-electrode set up can be found in the SI.

Dynamics of electrons and holes

Transient diffuse reflectance spectroscopy (TDRS) measurements were performed on the microsecond to second timescale to monitor the kinetics of photogenerated charges under water oxidation conditions (see Supporting Information for details). Previous studies have shown that WO_3 exhibits two main spectral features: a broad transient signal centred at around 800 nm which, by use of chemical scavengers, has been assigned to photoexcited electrons around the conduction band; and a hole signal, which may be probed below 550 nm.⁵² The transient diffuse reflectance spectra for the WO_3 needles in this work, Figure 3a, are in agreement with these previous assignments. The stronger electron absorbance signal spans most of the visible spectrum down to approximately 550 nm. Monitoring the electrons at 800 nm, we observe a long lived transient signal that decays by approximately 1 second (black trace, Figure 3b). The transient photocurrent (TPC) was measured simultaneously and integrated to give the extracted electron density as a function of time delay (dashed blue trace, Figure 3b). By comparing the TDRS electron signal at 800 nm with the integrated TPC signal, it is clear that the integrated TPC signal rises as the 800 nm signal decays, supporting the assignment of the 800 nm TDRS signal to WO_3 electrons. On the other hand, the hole signal, monitored at 500 nm, presents a much faster decay, implying faster kinetics for the hole reaction with the electrolyte. Given the high Faradaic efficiency exhibited by this photoanode, this hole reaction can be assigned primarily to water oxidation. If we take approximate half-times for these decays with respect to the signal amplitude at 10 μs ; $t_{1/2}(\text{electrons})$ is ~ 10 ms and $t_{1/2}(\text{holes})$ is ~ 0.5 ms. These values indicate that holes react an order of magnitude faster than electrons are extracted by the external circuit. As will be highlighted in more detail in the following sections, this timescale of water oxidation is much faster than those observed for other commonly studied transition metal oxides.^{8,10,53}

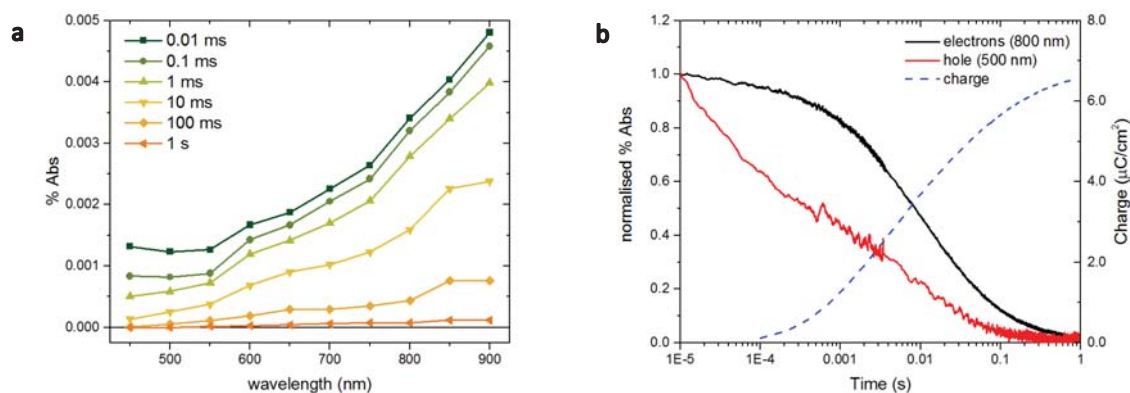


Figure 3: Transient diffuse reflectance data of WO_3 photoanode at 1.45 V_{RHE} in pH 1 electrolyte, 355 nm excitation at 0.3 mJ/cm^2 and 0.4 Hz. (a) transient spectra for time delays after photoexcitation from 0.01

ms to 1 s, and (b) kinetic traces at 800 nm (black) and at 500 nm (red), assigned to electron and hole signals, respectively, alongside integrated transient photocurrent data (dashed blue) corresponding to the kinetics of electron extraction to the external circuit.

Hole kinetics in water and propylene carbonate

To further examine the rapid water oxidation kinetics evident from the TDRS decay at 500 nm in Figure 2, the same kinetic measurement was repeated in inert organic solvent (propylene carbonate). As seen in Figure 4 (red trace), when there is no water present, a longer lived transient signal is observed at 500 nm. This trace still decays within the timescale of the measurement due to some oxidation of the solvent and/or other recombination processes. When a small percentage of water (pH 1 electrolyte) is added to the electrolyte, the rate of signal decay is increased, thus confirming that the fast decay observed at 500 nm is due to surface holes reacting with water. Although the poor miscibility of water in propylene carbonate limits our comparisons to low water concentrations, a clear trend of increasing rate with increasing water content is observed, with the 100% aqueous electrolyte displaying the fastest oxidation kinetics.

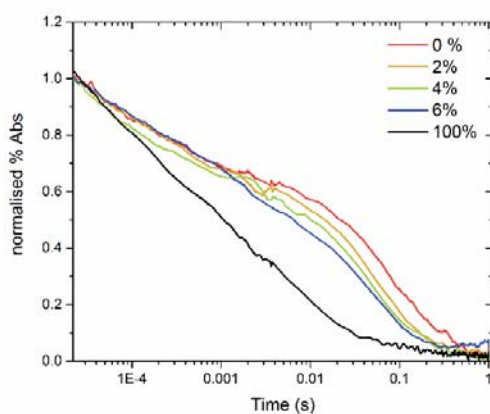


Figure 4: TDR decay traces for WO_3 needles at $1.45 V_{\text{RHE}}$ in propylene carbonate TBA SO_4 , with increasing aqueous electrolyte volume fraction ($0.1 \text{ M H}_2\text{SO}_4$, %), probed at 500 nm.

Kinetics of electron transport

To further investigate the relatively slow electron extraction kinetics observed from the 800 nm TDRS and TPC data in Figure 3, a series of transient diffuse reflectance decays were recorded at a range of potentials, along with integrated transient photocurrent signals, to determine the effect of applied bias (Figure 5a).

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With progressively larger anodic bias, the magnitude of the initial optical and final extracted charge signals increase simultaneously. The rise in these signal amplitudes indicates an increase in the density of photogenerated long-lived electrons that can be extracted to the external circuit. This can be rationalised as a consequence of greater band bending in the material under higher bias, creating a larger space charge layer to separate charges and thus a higher yield of electrons. In other words, the magnitude of the optical signal directly correlates to the number of electrons that remain in the electrode after any ultrafast bulk recombination occurs. In Figure 5b, we plot the correlation between our initial optical TDRS assay of long lived electrons (green triangles) and our integrated photocurrent assay of extracted electrons (blue circles) versus applied potential. The close agreement between these two assays indicates that, despite the slow electron extraction kinetics, the electrons observed on these timescales are still extracted reasonably efficiently by the external circuit, *i.e.* relatively few of the electrons monitored on the micro- to second timescales studies herein are recombining on this timescale. This is in agreement with the early photocurrent onset and lack of transient spikes observed in the chopped J-V curve (Figure 2a). If we consider the extracted charge density at both 50 ms (Figure 6b, pink squares) and 1 s (blue circles), we observe a small difference at low bias indicative of modest back electron-hole recombination losses. Back electron-hole (surface) recombination occurs over long timescales as a result of charge accumulation at the electrode-electrolyte interface. This build-up of charge draws electrons in the opposite direction of normal current flow, diminishing the total charge collected.³⁶ By 1.15 V_{RHE}, a large enough field is maintained to successfully prevent this slow timescale recombination process. This ‘surface recombination turnoff’ potential is lower than for α -Fe₂O₃ and BiVO₄ photoanodes,^{8,36,54,55} most probably due to the faster competing water oxidation reaction kinetics reported above. Applied potentials \geq 1.15 V_{RHE} result in increased charge extraction with time, as the majority of electrons are collected.

Figure 5c and d give, respectively, the normalised optical and photocurrent (without integration) transients for a series of applied biases. We note that increasing bias results in longer lived optical signals, and slower charge extraction kinetics. This could be a result of improved charge separation, although, as mentioned, there is little recombination at these timescales as most charges are extracted. Alternatively, and more likely, these slower kinetics under stronger anodic bias could result from increased trapping of charges with larger bias, resulting in hindered transport. It is widely acknowledged that WO₃ has a large propensity to form electronic defects within the band gap.^{24,29,56-58} With the conduction and valence band energies pinned at the electrolyte interface, increases in applied potential could empty, and thus make available for electron trapping, an increasing density of unoccupied shallow defect states within the space charge layer. This would slow extraction as electrons would encounter a greater number of progressively deeper states with increasing positive bias. This is analysed further in the discussion section below.

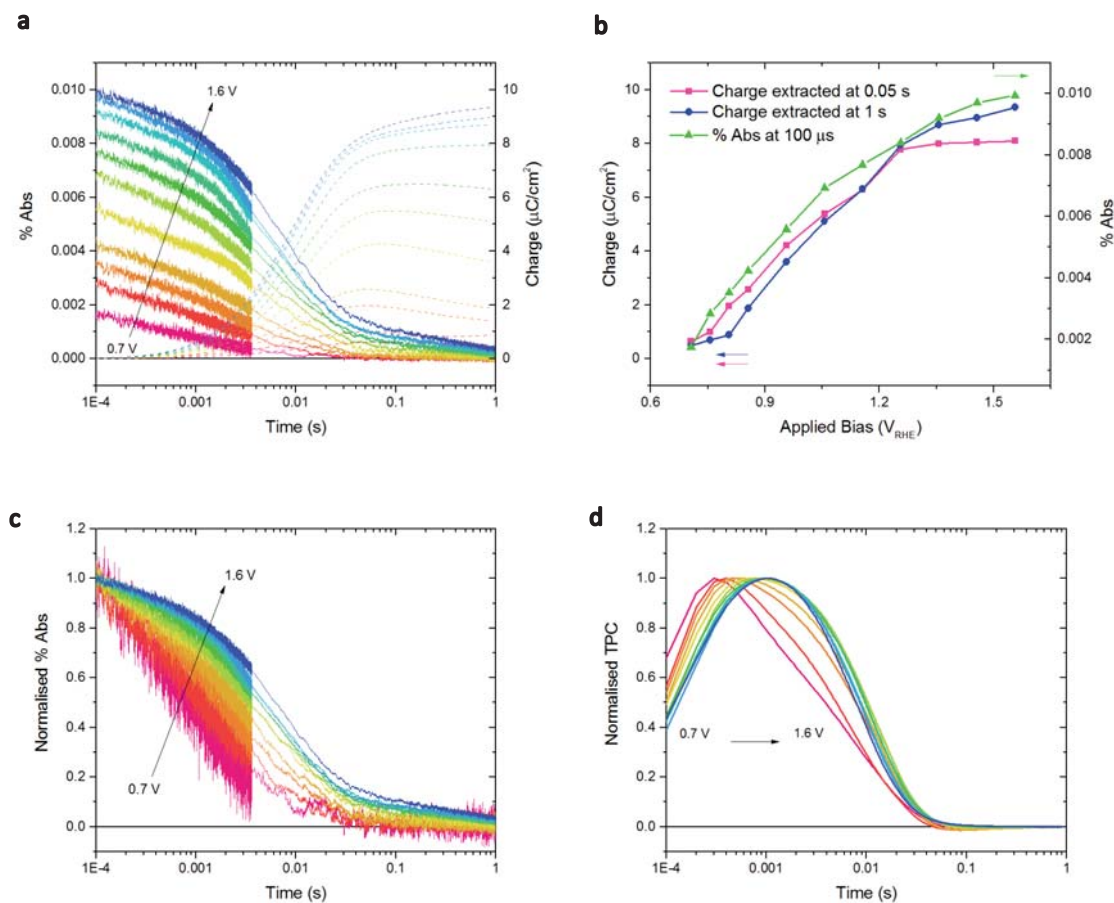


Figure 5: (a) TDR decays at 800 nm and the integrated transient photocurrent at a range of different applied potentials (0.7 – 1.6 V_{RHE}). (b) The initial magnitude of the optical signal (green triangles), and the extracted charge measured at 50 ms (pink squares) and at the end of the measurement, 1 s, (blue circles). (c) Normalised optical decays shown in (a). (d) Normalised transient photocurrents for the same applied potentials.

Incident photon to current efficiency (IPCE) measurements were conducted at 1.23 V_{RHE} with front and back irradiation to probe whether slow electron extraction is limiting performance (SI, Figure S5a). With FTO absorption accounted for in both cases, higher efficiencies were obtained when the needles were irradiated from the back (36% vs 22% at 350 nm), with charges generated closer to the extraction point. This correlates with the slow extraction times observed by TDRS and suggests that the small amount of

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6 surface or back electron-hole recombination observed (Figure 5) has some effect on performance. The trend
7 in IPCE also agrees with the relatively large diffusion length of holes (150 nm) in WO₃, which has
8 previously been reported to be significantly larger than for other popular metal oxides (*e.g.* ~4 nm in α -
9 Fe₂O₃), sometimes larger than the depletion width.^{5,21,28} Since the WO₃ needles are approximately 100 nm
10 in diameter, holes may readily diffuse to the surface regardless of whether the electrode is irradiated from
11 the front or the back. IPCE data was also collected over longer wavelengths to confirm that any absorption
12 from oxygen vacancies does not contribute to enhanced visible light harvesting (Figure S5b).
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18 WO₃ modifications

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20 To aid confirmation of the cause of such slow electron extraction kinetics, various alterations were made to
21 the WO₃ needles. The first was to grow shorter needles (2 micrometres) through modification of the same
22 AACVD synthesis.⁴⁶ As is apparent in Figure 6a, the shorter WO₃ needles give a faster optical electron
23 signal decay (blue line) and earlier charge extraction kinetics (dashed lines), indicating that electron
24 extraction is kinetically limited by electron transport through the WO₃ needles. Front and back IPCE are
25 more comparable for these shorter needles as the electron transport distance has been reduced (Figure S6).
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30 The second modification was to anneal the as-synthesised blue WO_{3-x} needles in air for a longer period of
31 time. We have already confirmed through XPS that annealing in air reduces the number of oxygen vacancies
32 (Figure S2b versus Figure 1b). Similarly, observations in literature show that the converse approach of
33 annealing under oxygen-poor atmosphere (hydrogen or argon) increases the number of oxygen vacancies,¹³
34 which may act as electron trap sites.⁵⁹ Although the nature of any trap sites discussed herein have not been
35 probed directly, it is reasonable to assume that these would be associated with oxygen vacancies as
36 evidenced by our XPS measurements (Figure 1b) and the wealth of data on WO₃.^{56,57,60,61} Figure 6b
37 compares the electron decay of a long needle sample annealed in air at 500°C for 2 hours to one annealed
38 for 12 hours, and thus with a reduced population of oxygen vacancies. This subtle change in oxygen vacancy
39 density was determined by an increase in sub-bandgap polaron absorption (Figure S7), as previously
40 reported in the literature.^{26,62,63} As can be observed, the half-lifetimes of both the TDRS and photocurrent
41 decays are reduced for the sample that was annealed for a longer period of time, consistent with the idea of
42 trap states associated with oxygen vacancies slowing electron extraction.
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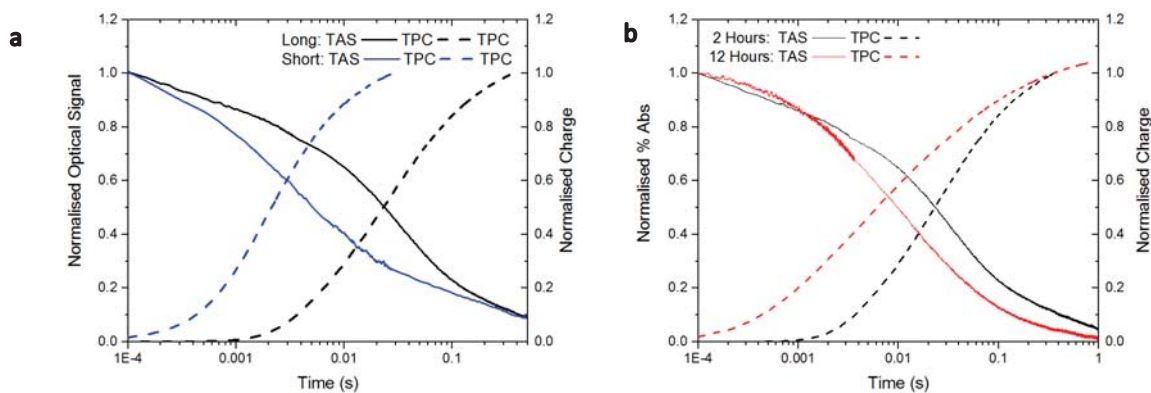


Figure 6: TDR decays at 800 nm (solid lines) and corresponding TPC data (dashed lines) at 1.45 V_{RHE} , pH 1 electrolyte. (a) Comparison of WO_3 needles with shorter WO_3 needles grown by the same method; (b) comparison of WO_3 needles to those annealed in air for a longer time to remove/fill more oxygen vacancies.

Comparisons with other transition metal oxides

We finally compare the kinetics of water oxidation and electron extraction reported herein for WO_3 with other transition metal oxides frequently employed as photoanodes for water splitting.

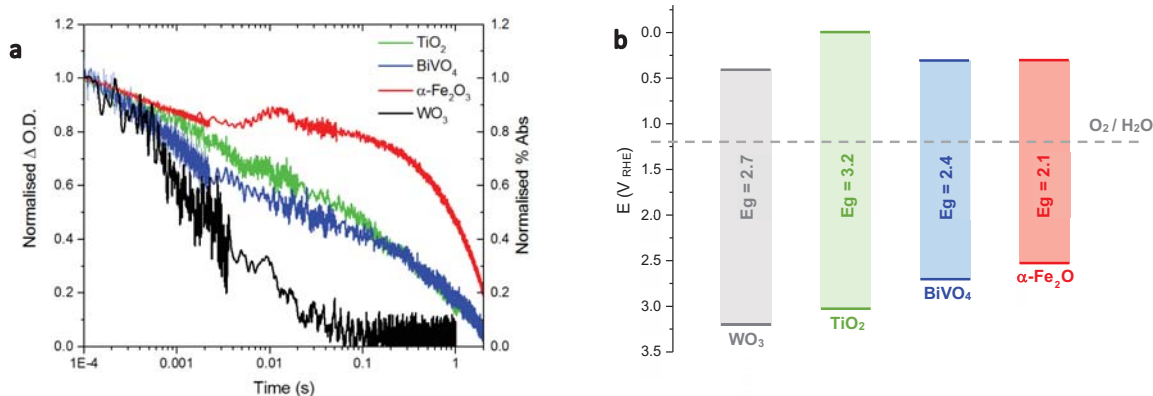


Figure 7: TDR decay traces for holes in the WO_3 needles (black), studied at 500 nm (pH 1, 1.45 V_{RHE}), compared to transient absorption measurements of hole decays in other transition metal oxides at similar applied bias, equating to the photocurrent plateau: BiVO_4 in blue (550 nm, pH 7, 1.5 V_{RHE}), TiO_2 in green (500 nm, pH 13, 1.15 V_{RHE}), and $\alpha\text{-Fe}_2\text{O}_3$ in red (650 nm, pH 13, 1.5 V_{RHE}). (b) Valence band alignments for the different metal oxides.^{17,19,64}

In Figure 7a, the transient optical signal for water oxidation by surface holes is compared between transition metal oxides commonly studied as photoanodes for water oxidation. Often considered the kinetic bottleneck of photoelectrochemical water splitting, the oxidation of water occurs on the timescale of 100 milliseconds to seconds for titania (TiO_2), hematite ($\alpha\text{-Fe}_2\text{O}_3$) and bismuth vanadate (BiVO_4) but is significantly faster for the WO_3 needle sample under similar applied bias. This trend can largely be correlated with the depth of the valence band (Figure 7b), with WO_3 displaying the most positive valence band energy and thus the largest thermodynamic driving force for water oxidation.

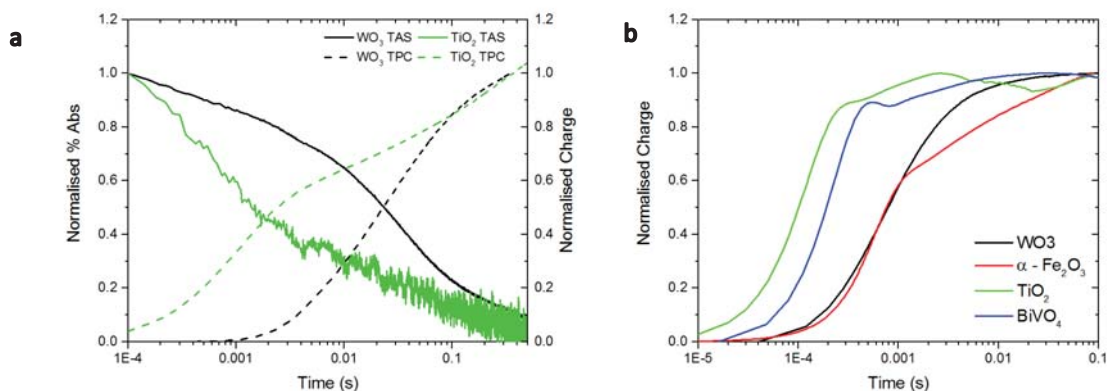


Figure 8: (a) Comparison of TDR decays at 800 nm (solid lines) and corresponding TPC data (dashed lines) at 1.45 V_{RHE} , pH 1 electrolyte for WO_3 needles (black) against TiO_2 needles (green) of similar length, also grown by CVD. (b) Comparison of flat, dense WO_3 (black trace, pH 1) with similar thicknesses of dense TiO_2 (green trace, pH 13), $\alpha\text{-Fe}_2\text{O}_3$ (red trace, pH 13) and BiVO_4 (blue trace, pH 7).

We also compared the electron extraction kinetics, observed both optically and by integrated photocurrent transients, of the WO_3 needles studied herein to TiO_2 needles of similar length (Figure 8a).⁶⁵ It is clear that the electron extraction is much faster for the TiO_2 sample, with a reduction in half-life of the decay from ~20 ms to ~1 ms. Using dense samples for ease of synthesis, Figure 8b gives a similar comparison of charge extraction for WO_3 versus dense TiO_2 , $\alpha\text{-Fe}_2\text{O}_3$ and BiVO_4 . Once again, we observe the slowest extraction times for the WO_3 sample, finding it similar to $\alpha\text{-Fe}_2\text{O}_3$ which is known to have low conductivity across such thicknesses when undoped (500 nm).^{21,66,67} This strongly suggests that slow electron transport is an intrinsic property of WO_3 . As suggested above, this is most likely related to its propensity to become oxygen deficient, forming vacancies that can trap charges.^{28,68}

Discussion

In the majority of photoelectrochemical systems examined for water oxidation, the water oxidation reaction itself is considered the kinetically limiting step whereas the extraction of the majority carrier (electrons) is comparatively fast.^{10,69,70} However, in this work on WO₃ photoanodes, we find the converse to be true. Direct analysis of transient photocurrent alongside the correlating optical decay of the electron signal has highlighted the extraction of electrons to be particularly slow, slower than the rate of water oxidation (Figure 3b). Rather surprisingly, we find evidence that holes in these nanostructured samples react faster at the electrolyte interface than for any other transition metal oxide we have previously investigated (Figure 7a), with an optical half-lifetime under bias of less than a millisecond ($t_{1/2}(h^+) \sim 0.5$ ms). Such rapid decay of hole signals in WO₃ has been observed previously by Cristino *et al.* but was assigned to non-faradaic oxidation processes rather than water oxidation.³⁹ Water oxidation was confirmed to be the dominant oxidation process at the surface of our WO₃ needles through oxygen evolution measurements yielding >85 % Faradic efficiency.

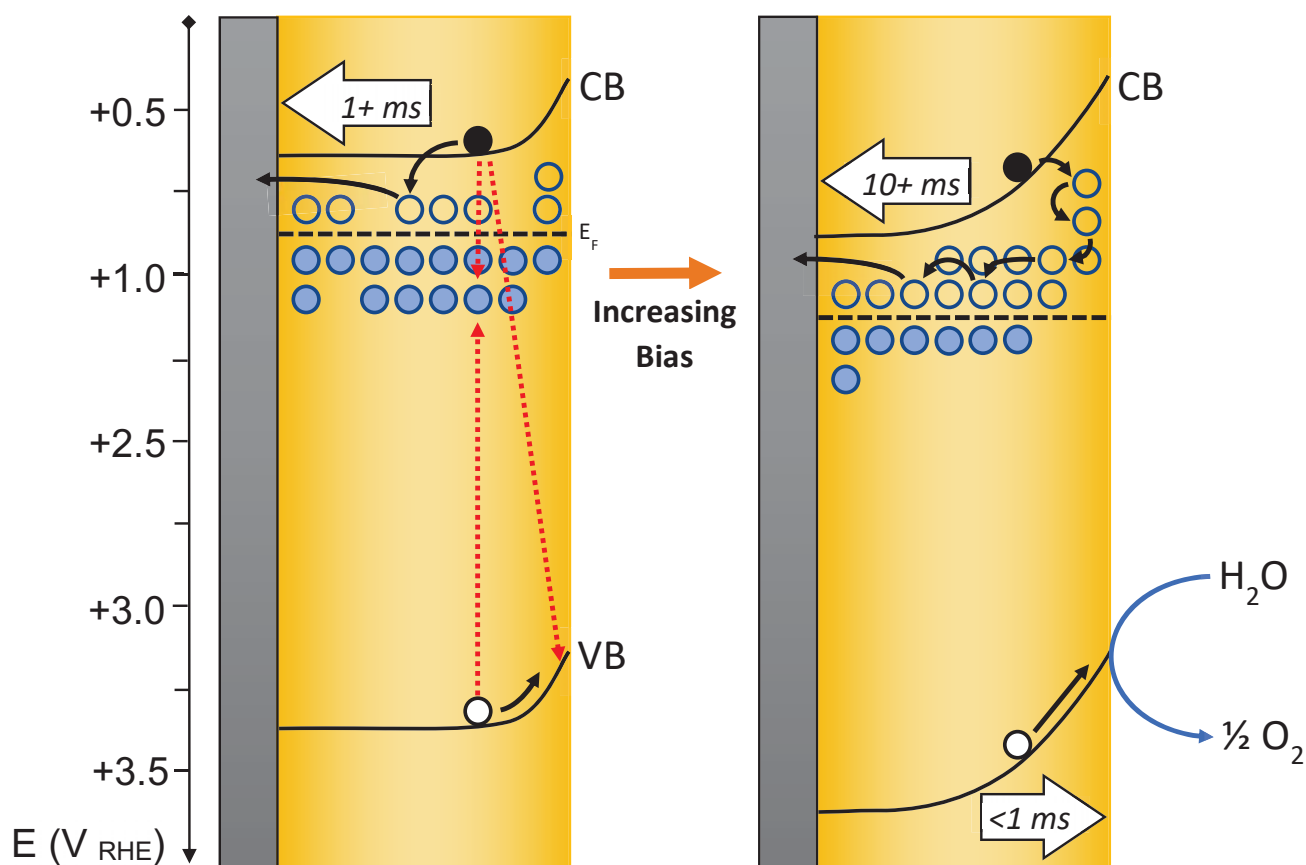
The importance of fast water oxidation kinetics is most apparent in the early onset of the photocurrent for this material (Figure 2), requiring a much smaller overpotential for water oxidation compared to α -Fe₂O₃ and BiVO₄.^{8,71} In the case of α -Fe₂O₃, it has been suggested that the requirement for overpotentials as large as 0.6 V above flat-band are the result of poor water oxidation kinetics,^{48,54,72} necessitating the use of co-catalysts to lower the onset potential.^{48,73,74} In this study, the earlier onset of WO₃, presented herein and observed by others,¹¹ correlates with faster water oxidation kinetics. Rapid reaction kinetics at the interface allow an increased proportion of holes that reach the surface to perform useful chemistry over competing back electron-hole recombination processes. This is true for our WO₃ nanoneedles without the need for co-catalysts or passivating overlayers. As shown in Figure 7b, the valence band edge is more positive in WO₃ than in BiVO₄, α -Fe₂O₃, and TiO₂,^{12,17,19,64,75} which may explain the origin of the faster water oxidation kinetics observed, given the excess thermodynamic driving force then available for the reaction to occur on WO₃. This relationship between valence band depth and kinetics has been proposed in other work,^{10,76} and would then suggest that there exists a valence band energy dependence on photoelectrochemical water oxidation that is analogous to the overpotential dependence often found in electrocatalytic water oxidation systems. The high crystallinity of the sample, the preferential growth orientation of the needles, or effects of pH on the water oxidation mechanism may also contribute to the faster kinetics observed for our WO₃, though further investigation on these aspects is beyond the scope of this study.

In stark contrast to the hole kinetics, when a similar comparison is made between WO₃ electron dynamics and those of other transition metal oxides, we find that WO₃ displays slower electron extraction (Figure 8). This difference in extraction time is as large as an order of magnitude when comparing TiO₂ and WO₃

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6 needles and is still evident when comparing the dense films, demonstrating that slow electron extraction is
7 a property of WO_3 . Considering that good electron conductivity is often presented as a promising feature
8 of WO_3 ,^{5,14,16} this result was surprising. It suggests that the main benefit of adding WO_3 as an electron
9 transport layer in various heterojunctions is to aid charge separation and thus increase charge carrier
10 lifetime, rather than rapidly transport electrons to the contact. Irradiated from the front, we observe that
11 longer WO_3 needles have a slower charge extraction than shorter needles (Figure 6a), confirming that the
12 issue is related to transport through the material, rather than interfacial defects at the contact. We propose
13 that this slow electron extraction is due to the trapping of electrons below the conduction band and suggest
14 that the mobility of electrons is thereby reduced. The negative correlation of rate of extraction with
15 increasing anodic bias (Figure 5c,d) supports a trap-mediated transport mechanism. As the anodic bias is
16 increased, an increasing number of progressively deeper states below the conduction band are emptied and
17 thus become accessible to trap photoexcited electrons. With greater applied bias, both the density and the
18 depth of these available electron traps in the space charge layer increases. It follows that the electron
19 extraction is slowed because electrons trap and de-trap to reach the back contact, effectively hopping
20 through the needle. This model is consistent with the observation that shorter needles have a faster
21 extraction time, as the distance to the contact is reduced and so fewer trapping/de-trapping steps are
22 required. The fitting of the TDR electron decay to a stretched exponential gives further analytical grounding
23 that supports this transport mechanism, as it indicates the dispersive time range in which electrons are
24 extracted due to multiple trapping steps (Figure S8).⁷⁷ This trapping and de-trapping form of transport has
25 previously been explored with stretched exponential functions or power law decays in other metal oxide
26 based systems.^{78,79} It has also been tentatively proposed by Lindquist's group,³⁰ and is consistent with the
27 presence of interband states in WO_3 as confirmed in other studies.^{29,60} As previously acknowledged,
28 tungsten trioxide has a large propensity to form numerous oxygen vacancy defects.^{56,58,68,80} It is therefore
29 highly probable that these chemical defects form the interband states responsible for trapping electrons.
30 This is supported herein by the improvement in extraction time observed when the WO_3 is subjected to
31 prolonged annealing in air (Figure 6b): with increased opportunity provided for oxygen to permeate the
32 WO_3 structure and fill/remove such oxygen vacancies, longer annealing results in fewer trap sites to hinder
33 the transport of electrons to the contact. Furthermore, the lack of photocurrent in the un-annealed sample
34 (Figure S3b) suggests that excessive concentrations of oxygen vacancies trap electrons so strongly and
35 impede the charge carrier mobility to such an extent that most charges recombine. A model of electron
36 dynamics for WO_3 needles is summarised in Scheme 1.

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51 As evidenced in the bias dependence study (Figure 5a,b), we find, depending upon applied bias, that the
52 majority of electrons observed on these microsecond to second timescales are eventually extracted, negating
53 the possibility that a slow, trap-mediated recombination is predominantly being monitored. The lack of
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substantial recombination losses observed in the electron kinetics across these timescales (0.01 ms to 1 s) is supported by consideration of the internal quantum efficiencies obtained for the longer and shorter WO₃ needles (Figure S6) which remain similar under front illumination, despite the difference in electron extraction times shown in Figure 6a. Modest back electron-hole (surface) recombination is observed for anodic potentials close to flat-band, however this is readily removed by applied potentials > 1.15 V vs RHE. The lower back electron-hole recombination losses compared to other metal oxide photoanodes such as BiVO₄ and Fe₂O₃ is most likely due the faster water oxidation kinetics by WO₃ holes, as discussed above, which will result in less hole accumulation at the surface under conditions of steady-state photoelectrochemical water oxidation.^{8,36} On the other hand, the extensive electron trapping and associated slow electron transport kinetics observed for the WO₃ photoanodes studied herein can be expected to lead to internal resistance losses which will limit the fill factor of the *J/V* curve in a tandem cell device. Given the relatively efficient electron extraction observed for these WO₃ photoanodes, we can conclude that ultrafast bulk recombination is the most probable factor limiting device internal quantum efficiency.



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8 Scheme 1: Photogenerated electrons (black) have different extraction times depending on the applied
9 potential. This is due to the increased number of traps that are empty under bias and hence can accept
10 electrons (open circles). Electrons navigate through these states to the external circuit as indicated by
11 the black arrows. Little recombination, either from the surface or via the trap states (red arrows), is
12 observed on the timescales of the experiment as photogenerated holes (white) are rapid transported
13 to the surface to oxidise water within milliseconds. Left; $0.8 V_{\text{RHE}}$, Right; $1.45 V_{\text{RHE}}$.
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18 In summary, we have shown our WO_3 needles to be rapid water oxidation photoanodes, with high faradaic
19 efficiency and faster hole injection into the electrolyte than in other commonly studied transition metal
20 oxide photoanodes. These rapid water oxidation kinetics improve catalytic performance by suppressing
21 competing surface recombination and enabling photocurrent onset close to flat-band without added co-
22 catalysts; we propose that the kinetics correlate with the depth of the valence band. The extraction of
23 electrons, on the other hand, is observed to be particularly slow, which we assign to trapping of electrons
24 into vacant defects below the conduction band. We propose that these trap sites are vacant oxygen vacancies
25 which form the primary intrinsic dopants in the material. Annealing for longer in air therefore increased the
26 rate of charge extraction by reducing the number of oxygen vacancies. We thus conclude that the density
27 of intrinsic chemical defects, such oxygen vacancies, can have a significant effect on charge transport
28 through nanostructured WO_3 , with potential implications for other heavily doped materials.
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38
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