

# What does CIV $\lambda$ 1549 tell us about the physical driver of the Eigenvector quasar sequence?★,★★

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Received 21 December 2016 / Accepted 7 August 2017

## ABSTRACT

Broad emission lines in quasars enable us to “resolve” structure and kinematics of the broad-line emitting region (BLR) thought to involve an accretion disk feeding a supermassive black hole. Interpretation of broad line measures within the 4DE1 formalism simplifies the apparent confusion among such data by contrasting and unifying properties of so-called high and low accreting Population A and B sources.  $H\beta$  serves as an estimator of black hole mass, Eddington ratio and source rest frame; the latter being a valuable input for CIV $\lambda$ 1549 studies which allow us to isolate the blueshifted wind component. Optical and HST-UV spectra yield  $H\beta$  and CIV $\lambda$ 1549 spectra for low-luminosity sources while VLT-ISAAC and FORS and TNG-LRS provide spectra for high-luminosity sources. New high-S/N data for CIV in high-luminosity quasars are presented here for comparison with the other previously published data. Comparison of  $H\beta$  and CIV $\lambda$ 1549 profile widths/shifts indicates that much of the emission from the two lines arise in regions with different structure and kinematics. Covering a wide range of luminosity and redshift shows evidence for a correlation between CIV $\lambda$ 1549 blueshift and source Eddington ratio, with a weaker trend with source luminosity (similar amplitude outflows are seen over four of the five dex luminosity ranges in our combined samples). At low luminosity ( $z \lesssim 0.7$ ) only Population A sources show evidence for a significant outflow while at high luminosity the outflow signature begins to appear in Population B quasars as well.

**Key words.** quasars: general – quasars: emission lines – quasars: supermassive black holes

## 1. Introduction

Just over fifty years after their discovery, we are beginning to see progress in both defining and contextualizing the properties of type-1 active galactic nuclei (AGN)/quasars which can be argued to be the “parent population” of the AGN phenomenon. We now know that type-1 sources show considerable diversity in most observed properties. They all show broad optical/UV emission lines from ionic species covering a wide range of ionization potential. Balmer lines (especially  $H\beta$ ) and CIV $\lambda$ 1549 Å (hereafter CIV) have been used as representatives of low- and high-ionization lines, respectively. They are almost always accompanied by broad permitted FeII emission extending from the far ultraviolet (FUV) to the near infrared (NIR). The Balmer lines have played a prominent role in characterizing the broad-line region (BLR) in lower-redshift ( $z < 0.8$ ) quasars, with

$H\beta$  providing information for the largest number of sources (Osterbrock & Shuder 1982; Wills et al. 1985; Sulentic 1989; Marziani et al. 2003a; Zamfir et al. 2010; Hu et al. 2012; Shen 2013). Most quasars lie beyond  $z \sim 0.8$  where we lose  $H\beta$  as a characteristic low-ionization broad line in the absence of expensive IR spectra. Other higher-ionization broad lines appear in their optical spectra. Can any of these lines be used as useful substitute of  $H\beta$ ? The higher ionization CIV line has long been considered the best candidate for a BLR diagnostic beyond  $z \sim 1.4$ .

CIV and other high-ionization lines – unlike  $H\beta$  and other low-ionization lines – frequently show blue-shifted profiles over a broad range in redshift ( $z$ ) and luminosity (Gaskell 1982; Espey et al. 1989; Tytler & Fan 1992; Marziani et al. 1996; Richards et al. 2002; Bachev et al. 2004; Baskin & Laor 2005; Sulentic et al. 2007; Netzer et al. 2007; Richards et al. 2011; Coatman et al. 2016). The CIV emission line has always played a major role in the study of quasars at  $z \gtrsim 1.4$ . Over the years, we have come to consider outflows as almost ubiquitous in quasars precisely because of the blueshifts of CIV. Early reports at low  $z$  encompassed International Ultraviolet Explorer (IUE) observations of individual Seyfert 1 galaxies (e.g., Wamsteker et al. 1990; Zheng et al. 1995), or of high- $z$  quasars with CIV redshifted into the optical band (e.g., Osmer & Smith 1976). Since the 1990s the mainstream interpretation of the

\* Based on observations made with ESO telescopes at the La Silla Paranal Observatory under programme ID082.B-0572(A) and with the Italian Telescopio Nazionale Galileo (TNG) operated by the Fundación Galileo Galilei (INAF) at the Observatorio del Roque de los Muchachos.

\*\* The reduced spectra are only available at the CDS via anonymous ftp to [cdsarc.u-strasbg.fr](ftp://cdsarc.u-strasbg.fr) (130.79.128.5) or via <http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/608/A122>

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blueshift has involved emission by outflowing gas, possibly in a wind (Collin-Souffrin et al. 1988; Emmering et al. 1992; Murray et al. 1995) but alternate interpretations are still being discussed (Gaskell & Goosmann 2013). A less incomplete view came from the first analysis of *Hubble* Space Telescope (HST) observations (Wills et al. 1993; Corbin & Boroson 1996; Marziani et al. 1996). For the first time, it was possible to take advantage of moderate resolution in the UV domain, and a reliable rest frame set by optical narrow lines. For intermediate to high- $z$  quasars, a turning point came with the first data releases of the Sloan Digital Sky Survey (SDSS; Richards et al. 2002), where shifts were found to be common, especially in the amplitude range from a few hundred  $\text{km s}^{-1}$  to  $\sim 1000 \text{ km s}^{-1}$ . It was realized in the 1990s and later confirmed that the behavior of CIV in powerful radio-loud (RL) sources is different from the behavior in radio-quiet (RQ) sources, both in terms of shift amplitude and dependence on luminosity (Marziani et al. 1996; Richards et al. 2011; Marziani et al. 2016a). The previous works suggest that  $H\beta$  and CIV are not surrogates of one another – both lines arising, for example, in the same stratified medium (Marziani et al. 2010) and that RQ and RL sources should be considered separately when dealing with a high-ionization line such as CIV. The next major development came with the availability of IR spectrometers that allow the observations of narrow emission lines redshifted into the NIR and therefore to set CIV shifts measurements on a firm quantitative basis. Shifts larger than  $\sim 1000 \text{ km s}^{-1}$  were found to be frequent at high luminosity (e.g., Coatman et al. 2016; Marziani et al. 2016a; Vietri 2017b; Bisogni et al. 2017).

Current interpretation of the BLR in quasars sees the broad lines arising in a region that is physically and dynamically composite (Collin-Souffrin et al. 1988; Elvis 2000; Ferland et al. 2009; Marziani et al. 2010; Kollatschny & Zetzl 2013; Grier et al. 2013; Du et al. 2016). It appears that we need both CIV and  $H\beta$  to effectively characterize the low- and high-ionization gas in the BLR of quasars. For the past 20+ years our group has focused on spectroscopy of type-1 AGN/quasars, with one of our goals being the comparison of  $H\beta$  and CIV properties in the same sources. The work has involved two main steps.

1) LOWZ: using a large low- $z$  sample ( $\sim 160+$  sources) we introduced the 4DE1 contextualization to distinguish and unify type-1 quasar diversity. Measures of the  $H\beta$  profile width and CIV profile shift were adopted as two of the principal 4DE1 parameters in that formalism (Sulentic et al. 2000a,b, 2002). Comparison of  $H\beta$  measures with all high S/N CIV HST-UV spectra (130 sources by 2007) completed the low-redshift study (Sulentic et al. 2007). The 4DE1 formalism serves as a multidimensional set of HR diagrams for type-1 sources (Sulentic et al. 2000a). Broad-line measures, FWHM  $H\beta_{BC}$  and FeII strength, as parameterized by the intensity ratio involving the FeII blue blend at  $4570 \text{ \AA}$  and broad  $H\beta$  (i.e.,  $R_{\text{FeII}} = I(\text{FeII}\lambda 4570)/I(H\beta)$ ), coupled with measures of the CIV profile velocity displacement, provide three dimensions of the 4DE1 parameter space that are observationally independent and measure different physical aspects (see Marziani et al. 2010, and references therein). Earlier evidence for a correlation between the strength of a soft X-ray excess and some of the above measures (Wang et al. 1996; Boller et al. 1996) also motivated the inclusion  $\Gamma_{\text{soft}}$  as a 4DE1 parameter.

This first phase yielded a systematic view of the main observational properties along the 4DE1 sequence, which included parameters ranging from FIR colors to prevalence of radio-loud (RL) sources, to  $[\text{OIII}]\lambda 4959$ , 5007 blueshifts. Low- $z$  SDSS samples (Zamfir et al. 2010) corroborates a 4DE1 main

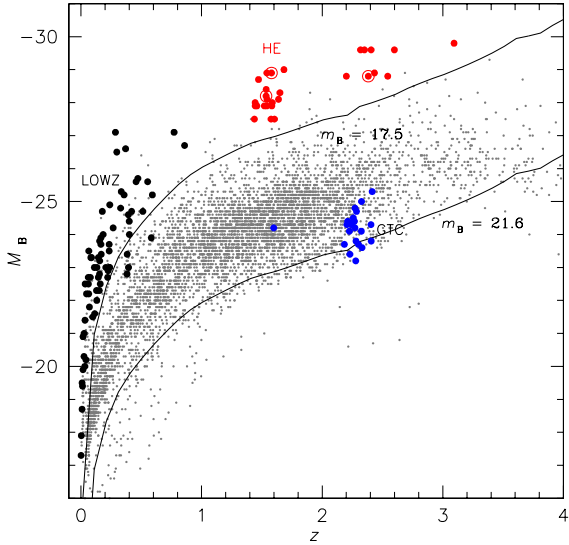
sequence extending from Population A quasars with  $FWHM H\beta < 4000 \text{ km s}^{-1}$ , strong optical FeII emission, a frequent CIV blueshift and soft X-ray excess and, at the other extreme, Population B sources usually showing broader lines, weaker optical FeII and absence of a CIV blueshift or soft X-ray excess (Sulentic et al. 2011; Fraix-Burnet et al. 2017, summarize several parameters that change among the 4DE1 sequence and are therefore systematically different between Pops. A and B).

We suggested that the distinction between NLSy1 and the rest of type-1 AGN was not meaningful, and that, at low  $z$ , a break in properties occurs at a larger  $FWHM(H\beta)$ ,  $\approx 4000 \text{ km s}^{-1}$  distinguishing the two populations, A and B, as narrower and broader (or high and low Eddington ratio) sources, respectively (Sulentic et al. 2011). Eventually, the break was associated with a critical Eddington ratio value (Marziani et al. 2003b), and possibly with a change in accretion disk structure (Sulentic et al. 2014b, and references therein). Eddington ratio ( $L/L_{\text{Edd}}$ ) was suggested by several authors as being the key parameter governing the 4DE1 sequence (e.g., Boroson & Green 1992; Marziani et al. 2001; Kuraszkiwicz et al. 2000; Dong et al. 2009; Ferland et al. 2009; Sun & Shen 2015). The roles of orientation, black hole mass, spin, and line-emitting-gas chemical composition in shaping the observed spectra remain unconvincingly understood to this date (Sulentic et al. 2012, and references therein), especially for RQ sources. Nonetheless, the 4DE1 has offered a consistent organization of quasars' observational properties whose physical interpretation is still lacunose.

2) High-luminosity, HE sample: the advent of the ESO VLT-ISAAC NIR spectrograph opened up the possibility for obtaining high-S/N spectra for  $H\beta$  in much-higher-luminosity quasars at intermediate-to-high redshift. Our extreme  $L$  quasar sample was chosen from the Hamburg ESO Survey (Wisotzki et al. 2000) in the range  $z = 0.9-3.1$  and bolometric luminosity  $\log L = 47.4-48.4$ . We obtained spectra of the  $H\beta$  region in 53 sources (Sulentic et al. 2004, 2006; Marziani et al. 2009). Some properties of the emission lines in the  $H\beta$  region at high luminosity  $\log L \gtrsim 47$  [ $\text{erg s}^{-1}$ ] are relevant to the interpretation of the results of the present paper. The  $H\beta$  profile still supported the dichotomy between narrower and broader sources, with the broader sources showing a redward asymmetry, even more prominent than at low- $L$ . The precise definition of Pop. A and Pop. B is however somewhat blurred at high  $L$ , since the minimum full width at half maximum (FWHM) of  $H\beta$  significantly increases with source bolometric luminosity (Marziani et al. 2009, hereafter M09).

In this paper we add very high S/N VLT-FORS and TNG-LRS spectral data for CIV of 28 high-luminosity quasars. Earlier high-S/N CIV spectra existed for other quasar samples, but no matter how good the CIV data, measures of CIV profiles remained of uncertain interpretation without a precise determination of the quasar rest frame that we have obtained in our sample from the  $H\beta$  spectral region. Accurate  $z$  measurement is not easy to obtain from UV broad lines, and redshift determinations from optical survey data, even at  $z \gtrsim 1$ , suffer systematic biases as large as several hundred  $\text{km s}^{-1}$  (Hewett & Wild 2010; Shen et al. 2016).

This paper presents new observations and data analysis along with detailed profile decompositions (reported in Sect. 2.1) that become possible with high-S/N spectra. Results from CIV data complete step 2 of our  $H\beta$  – CIV study and also allow us a comparison with previous low-redshift work. Our comparisons with low- $z$  sources is unprecedented because we introduce a sample of extreme-luminosity sources with accurate rest-frame estimates as well as spectra with high enough S/N to enable



**Fig. 1.** Location of the HE sample in the plane  $M_B$  vs. redshift (red spots), superimposed to a random subsample of the Véron-Cetty & Véron (2010) catalog (gray dots). Data points representing our low- $z$  RQ control sample and GTC sample of faint quasars at  $z \sim 2.3$  are also shown as black and blue filled circles respectively. The  $M_B$  associated with two apparent magnitude limited sample (17.5 appropriate for the HE survey, and 21.6 for the SDSS) are shown as a function of redshift.

decomposition of CIV into the blueshifted component and the unshifted component CIV counterpart to broad component H $\beta$ . Such comparisons are simply the best clues yet about the relative behavior of Pop. A and B quasars at low and high  $z$  and low and high  $L$ . Measurements and empirical models of spectra are reported in Sect. 3; they are discussed in Sect. 4 with a focus on the correlations between CIV shifts and luminosity and Eddington ratio (Sect. 4.3), which constrain the physical driver of the outflow (Sect. 4.4). A follow-up paper presents an analysis of possible virial broadening estimators based on rest-frame UV observations.

## 2. Samples, observations, and data analysis

### 2.1. Contextualizing our old and new Civ samples

Figure 1 shows the locations in a redshift – absolute magnitude plane of our extreme-luminosity Hamburg-ESO (HE) sample and the low- $z$  comparison sample of 70 radio-quiet quasars (RQ) used in this study. The HE sample, and quasars in the low- $z$  sample above  $z = 0.2$  involve the bright end of the optical luminosity function at their respective redshifts. Quasars with luminosities similar to those in the HE sample do not exist at low redshift. We also included in Fig. 1 a sample of 22 high-redshift ( $z \sim 2.3$ ) quasars (GTC sample) with luminosities similar to local quasars. Spectra for these sources were obtained with the 10 m GTC telescope (Sulentic et al. 2014a). We were motivated to see if low-luminosity quasars similar to our low- $z$  sample (luminosity analogs with  $\log L = 46.1 \pm 0.4$ ) existed at  $z \sim 2.3$  and whether or not they would show differences from their low- $z$  counterparts indicating a role for redshift evolution. The lack of H $\beta$  spectra for this sample required rest frame estimation from UV lines.

**Table 1.** Basic sample properties.

Name	$m_B$	$z$	Ref. <sup>a</sup>	$M_B^b$	$\log L$	Pop.	Notes
HE0035-2853	17.03	1.6377	1	-28.1	47.7	B	
HE0043-2300	17.06	1.5402	1	-27.9	47.6	A	RL
HE0058-3231	17.14	1.5821	1	-27.9	47.7	B	
HE0109-3518	16.44	2.4057	1	-29.6	48.3	A	
HE0122-3759	16.94	2.2004	2	-28.8	48.0	A	
HE0203-4627	17.34	1.4381	1	-27.5	47.5	B	RL
HE0205-3756	17.17	2.4335	1	-28.9	48.0	A	
HE0248-3628	16.58	1.5355	1	-28.4	47.9	A	
HE0251-5550	16.59	2.3505	1	-29.6	48.3	A	
HE0349-5249	16.13	1.5409	1	-28.9	48.1	B	
HE0359-3959	17.09	1.5209	1	-27.9	47.6	A	
HE0436-3709	16.84	1.4447	1	-27.9	47.7	B	
HE0507-3236	17.36	1.5770	3	-27.6	47.5	A	
HE0512-3329	17.03	1.5873	3	-28.0	47.7	A	
HE0926-0201	16.23	1.6824	3	-29.0	48.1	B	
HE0940-1050	16.70	3.0932	1	-30.1	48.4	A	
HE1039-0724	16.90	1.4584	1	-28.1	47.6	B	
HE1104-1805	16.45	2.3192	3	-29.6	48.3	A	
HE1120+0154	16.31	1.4720	1	-28.7	48.0	B	
HE1347-2457	16.83	2.5986	1	-29.6	48.3	A	
HE1349+0007	16.83	1.4442	1	-28.0	47.7	B	
HE1409+0101	16.92	1.6497	1	-28.3	47.8	B	
HE2147-3212	16.84	1.5432	1	-28.1	47.7	A	
HE2156-4020	17.39	2.5431	1	-28.8	48.0	B	
HE2202-2557	16.71	1.5347	2	-28.2	47.8	B	
HE2349-3800	17.46	1.6040	2	-27.5	47.5	B	RL
HE2352-4010	16.05	1.5799	2	-28.9	48.1	A	
HE2355-4621	17.13	2.3825	2	-28.8	48.0	B	

**Notes.** <sup>(a)</sup> References for redshift: 1: Marziani et al. (2009); 2: Sulentic et al. (2004); 3: Sulentic et al. (2006). <sup>(b)</sup> Absolute  $B$  magnitude, obtained from references in column Ref., corrected of galactic absorption and  $K$ -correction, and computed for  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_M = 0.3$ ,  $\Omega_\Lambda = 0.7$ .

### 2.2. The Hamburg-ESO (HE) sample

#### 2.2.1. General properties

The quasars considered in the present study are selected from ISAAC H $\beta$  observations of 52 quasars identified in the Hamburg ESO survey (Wisotzki et al. 2000), in the redshift range  $0.9 \lesssim z \lesssim 3$ . Previous observations of H $\beta$  were discussed in a series of papers (Sulentic et al. 2002, 2006; Marziani et al. 2009, hereafter M09). Out of the 52 H $\beta$  sources, 32 have  $z \gtrsim 1.4$  that allows for coverage of redshifted CIV, and of them, 28 were actually observed.

Table 1 summarizes the main properties of the sources including: identification in the HE catalog, apparent  $B$  magnitude taken from Wisotzki et al. (2000), redshift, reference to redshift, absolute  $K$ -corrected  $B$  magnitude also corrected for Galactic extinction, logarithm of the bolometric luminosity estimated from the tabulated  $m_B$ , assuming  $10\lambda L_\lambda(5100) \text{ \AA} = L$ , and population assignment (A or B, discussed in Sect. 2.7) in the 4DE1 scheme. In the last column, we identify RL sources. The redshifts were estimated from the narrow component of H $\beta$  and from [OIII] $\lambda\lambda$ 4959, 5007 (only if consistent with H $\beta_{\text{NC}}$ ) and are as reported in M09 and in the previous papers of the series dealing with H $\beta$  VLT/ISAAC observations. They are adopted as systemic redshifts for both the H $\beta$  and CIV spectral range.

#### 2.2.2. Objects belonging to particular classes

HE0512-3329 and HE1104-1805 are known to be gravitationally lensed quasars. They were discussed with the presentation of the H $\beta$  ISAAC spectra (Sulentic et al. 2006). Here we simply note that the FORS1 slit, PA was  $-112$  deg for HE1104, and that the PA values exclude contamination by the fainter lensed image. The HE0512 PA  $\approx -41$  deg implies an angle with

**Table 2.** Log of observations.

Name	Coordinates		Date	Telescope/ instrument	Exp. time (s)	Slit (arcsec)
	$\alpha$ (2000)	$\delta$ (2000)	observation			
HE0035-2853	00 38 06.5	-28 36 49	7-Oct.-08	VLT-U2/FORS	1980	0.7
HE0043-2300	00 45 39.5	-22 43 56	30-Oct.-09	VLT-U2/FORS	2340	1.0
HE0058-3231	01 00 39.2	-32 14 57	30-Oct.-09	VLT-U2/FORS	2160	1.0
HE0109-3518	01 11 43.5	-35 03 01	12-Oct.-08	VLT-U2/FORS	1800	1.0
HE0122-3759	01 24 17.4	-37 44 23	7-Oct.-08	VLT-U2/FORS	1980	0.7
HE0203-4627	02 05 52.4	-46 13 30	11-Oct.-08	VLT-U2/FORS	2520	1.0
HE0205-3756	02 07 27.2	-37 41 57	7-Oct.-08	VLT-U2/FORS	1980	0.7
HE0248-3628	02 50 55.3	-36 16 35	12-Oct.-08	VLT-U2/FORS	1800	1.0
HE0251-5550	02 52 40.1	-55 38 32	7-Oct.-08	VLT-U2/FORS	1800	0.7
HE0349-5249	03 50 59.3	-52 40 35	7-Oct.-08	VLT-U2/FORS	1800	1.0
HE0359-3959	04 01 14.0	-39 51 33	11-Oct.-08	VLT-U2/FORS	1440	1.0
HE0436-3709	04 38 37.3	-37 03 41	11-Oct.-08	VLT-U2/FORS	1440	1.0
HE0507-3236	05 09 17.8	-32 32 45	11-Oct.-08	VLT-U2/FORS	1440	1.0
HE0512-3329	05 14 10.8	-33 26 23	11-Oct.-08	VLT-U2/FORS	2340	1.0
HE0926-0201	09 29 13.5	-02 14 47	4-Mar.-08	TNG/LRS	2400	1.0
HE0940-1050	09 42 53.5	-11 04 27	27-Jan.-09	VLT-U2/FORS	1980	0.7
HE1039-0724	10 42 19.3	-07 40 37	6-Feb.-09	VLT-U2/FORS	2340	1.0
HE1104-1805	11 06 33.5	-18 21 25	6-Feb.-09	VLT-U2/FORS	1980	1.0
HE1120+0154	11 23 20.7	01 37 48	6-Feb.-09	VLT-U2/FORS	1980	1.0
HE1347-2457	13 50 38.8	-25 12 16	4-May-08	TNG/LRS	2400	1.0
HE1349+0007	13 51 50.5	00 07 39	1-Apr.-08	TNG/LRS	2400	1.0
HE1409+0101	14 12 21.7	00 47 19	3-May-08	TNG/LRS	2400	1.0
HE2147-3212	21 50 52.3	-31 58 26	21-Oct.-09	VLT-U2/FORS	1980	1.0
HE2156-4020	21 59 54.7	-40 05 50	30-Oct.-09	VLT-U2/FORS	2340	0.7
HE2202-2557	22 05 29.8	-25 42 23	2-Oct.-08	VLT-U2/FORS	1980	1.0
HE2349-3800	23 52 10.7	-37 43 22	30-Sep.-08	VLT-U2/FORS	2700	0.7
HE2352-4010	23 55 34.5	-39 53 54	2-Oct.-08	VLT-U2/FORS	1980	1.0
HE2355-4621	23 58 09.2	-46 05 00	30-Sep.-08	VLT-U2/FORS	2520	0.7

the lens axis of  $\approx 58$  deg on average. The small separation between the two lensed images  $\approx 0.64$  arcsec (which are of comparable brightness in the  $B$ -band) implies a separation within the slit of 0.5 arcsec, and that the two images should be well within the slit if the slit was centered on their midpoint. This is however not to be taken for granted since the  $A$  component is  $\approx 0.44$  mag brighter than  $B$  in the  $R$  band (Gregg et al. 2000). The possibility of lensing or of the occurrence of a close quasar pair was also raised for HE1120+0154 ( $\equiv$ UM 425). In this case, the separation of the FORS observations ( $\approx -171$  deg) excludes contamination by fainter lensed quasar images. In addition, the second-brightest component of the lens is very faint,  $V \approx 20$  (Meylan & Djorgovski 1989).

An interesting aspect of HE1120 is the presence of variable CIV absorption lines (Small et al. 1997). Given the limited radial velocity extent of the absorptions, the source would qualify as a mini-BAL, rather than a BAL QSO since the BALnicity index (Weymann et al. 1991) estimated from our spectrum is  $0 \text{ km s}^{-1}$ .

HE0359-3959 and HE2352-4010 belong to the class of weak-lined quasars (WLQs, Diamond-Stanic et al. 2009; Shemmer et al. 2010, HE 2352 as a borderline source), for which  $W(\text{CIV}) \lesssim 10 \text{ \AA}$ . WLQs in the HE sample are discussed in Sect. 4.5.

### 2.3. Observations and data reduction

Most (24) of the 28 HE quasar observations were obtained with FORS on the 8 m VLT-U2 telescope (ESO, Paranal Observatory, Chile) with four additional observations obtained with the Low Resolution Spectrograph (LRS) at the 3.6 m TNG telescope (Observatorio del Roque de los Muchachos, La Palma, Spain). Table 2 presents an observational log with: equatorial coordinates at J2000, date of observation, telescope and spectrograph, exposure time and slit width in arcsec. Long-slit observations with VLT/FORS1 were obtained with the 600B grism that provided a reciprocal dispersion of  $1.5 \text{ \AA/pixel}$ . The LRS

was equipped with the LR-B grism which yielded a dispersion of  $2.7 \text{ \AA/px}$ .

Data reduction was carried out using standard IRAF procedures. Spectra were first trimmed and overscan corrected. Bias subtraction and flat-fielding were performed each night. One dimensional (1D) wavelength calibration using arc lamps were obtained with the same configuration and slit width. The rms of the wavelength calibration was better than  $0.07 \text{ \AA}$  in all cases. The wavelength scales of all individual exposures were realigned with sky lines present in the spectrum before object extraction and background subtraction, to avoid significant zero point errors. The scatter of sky line peak wavelengths after reassignment  $\approx 0.5 \text{ \AA}$  or  $30 \text{ km s}^{-1}$  at  $5000 \text{ \AA}$  provides a realistic estimate of the uncertainty in the wavelength scale. The spectral resolution estimated from the FWHM of skylines is  $\sim 200\text{--}300 \text{ km s}^{-1}$  for VLT/FORS data. The S/N is  $\sim 100\text{--}200$  in the continuum of the CIV spectral region. The TNG/LRS observations yielded a spectral resolution of  $\approx 600 \text{ km s}^{-1}$  with a typical  $S/N \gtrsim 50$ , and  $\sim 200$  in the case of HE0926-0201 and HE1409+0101. Flux calibration was obtained each night from spectrophotometric standard stars observed with the same configuration. Due to the narrow slit width, and the wavelength coverage extending to the NUV ( $\approx 3300 \text{ \AA}$ ), an additional correction due to differential atmospheric refraction following Filippenko (1982), as well as an absolute flux recalibration, proved to be necessary.

We derived synthetic photometry on the spectra (computed using the sband task of IRAF). Specific fluxes in physical units ( $\text{erg s}^{-1} \text{ cm}^{-2} \text{ \AA}^{-1}$ ) were rescaled to the synoptic Catalina observations (Drake et al. 2009) for the wide majority of the objects, and to the SDSS (HE1120) and the NOMAD catalogs (HE0248; Zacharias et al. 2012; HE0058 and HE0203; Zacharias et al. 2005). Similarly, the original  $H\beta$  spectra synthetic magnitudes have been rescaled to the 2MASS  $J$ ,  $H$ ,  $K$  magnitudes (Cutri et al. 2003). The rescaling allowed for a validation of the bolometric luminosity computed from the  $1450 \text{ \AA}$  and  $5100 \text{ \AA}$  fluxes, assuming bolometric correction factors of 3.5 and 10, respectively. Following the rescaling, the bolometric luminosities derived from the optical and UV are in good agreement, with average  $\delta \log L = \log(3.5\lambda L_\lambda)_{1450} - \log(10\lambda L_\lambda)_{5100} \approx -0.072 \pm 0.310$  ( $\log L \gtrsim 43.5$  [ $\text{erg s}^{-1}$ ]) for the combined LOWZ+HE sample, and  $\delta \log L \approx -0.086 \pm 0.264$  for the HE sample. Only four sources of the HE sample (HE0205, HE0248, HE0436, and HE0512) have  $|\delta \log L| \gtrsim 0.5$ . The gravitationally lensed source, HE 0512, might have been affected by a different magnification in the UV and optical. If they are excluded, the HE average becomes  $\delta \log L \approx -0.05 \pm 0.17$ .

Systematic losses in emission line fluxes will cancel out when computing intensity ratios, provided that flux losses are the same for both lines. Nonetheless, great care should be used if computing intensity ratios of individual sources between the CIV and  $H\beta$  spectral range. Observations are not simultaneous and separation in time  $\Delta t$  between optical and IR span the range  $\Delta \tau \approx \Delta t/(1+z) \lesssim 1\text{--}3 \text{ yr}$  in the rest frame. Major optical/UV variability of highly luminous quasars is expected to occur on long timescales set by the expected emissivity-weighted radius of the BLR,  $c\Delta \tau \propto L(5100)_{48}^{1/2} \approx 4000 \text{ ld}$  (if we can extrapolate the radius-luminosity relation of Bentz et al. 2009). Previous works have also shown that, for a given time interval, higher-luminosity sources show lower variability levels (Vanden Berk et al. 2004; Simm et al. 2016; Caplar et al. 2017). However, moderate short-term variations ( $\approx 20\%$  on  $\Delta \tau \approx 15 \text{ d}$  on the rest frame) are possible (Woo et al. 2013; Punsly et al. 2016) and, along with intrinsic spectral energy distribution differences affecting the

bolometric correction, may be responsible for part of the scatter between rest frame optical- and UV-based  $L$  estimates.

#### 2.4. LOWZ: FOS-based comparison sample

A low- $L$  comparison sample was selected from the 71 RQ sources of Sulentic et al. (2007)<sup>1</sup>. Here we consider 70 RQ FOS sources (one source at  $\log L < 42 \text{ erg s}^{-1}$ , with  $L$  estimated from the H $\beta$  spectral range, was excluded due to its outlying low luminosity) as a complementary low- $L$ , FOS sample for a comparative analysis of CIV line parameters (LOWZ). Measurements on the H $\beta$  line profiles that include centroids, widths, and asymmetry are available for 49 of them from Marziani et al. (2003a). The comparison is restricted to this sample if H $\beta$  parameters beyond FWHM are involved. RL sources were excluded from the LOWZ sample because they are overrepresented with respect to their prevalence in optically selected samples ( $\approx 10\%$ , Zamfir et al. 2008) and especially because radio loudness is affecting the high-ionization outflows traced by CIV (Marziani et al. 1996; Punsly 2010; Richards et al. 2011; Marziani et al. 2016a).

The LOWZ sample cannot be considered complete, but: (1) previous work has shown that it samples well the 4DE1 optical main sequence; and (2) its bolometric luminosity  $L$  distribution is relatively uniform over the range  $44 \lesssim \log L \lesssim 47 \text{ [erg s}^{-1}\text{]}$ .

#### 2.5. Data analysis: CIV $\lambda$ 1549 and H $\beta$ profile interpretation

The broad profile of the H $\beta$  and CIV profiles in each quasar spectrum can be modeled by changing the relative intensity of three main components.

- The broad component (BC) has been referred to by various authors as the broad component, intermediate component, central broad component (e.g., Brotherton et al. 1994a; Popović et al. 2002). Here the BC is represented by a symmetric and unshifted profile (Lorentzian for Pop. A or Gaussian for Pop. B) and assumed to be associated with a virialized BLR subsystem.
- The blueshifted component (BLUE). A strong blue excess in Pop. A CIV profiles is not in doubt. In some CIV profiles, such as the extreme Population A prototype I Zw1, the blue excess is by far the dominant contributor to the total emission line flux (Marziani et al. 1996; Leighly & Moore 2004). BLUE is modeled by one or more skew-normal distributions (Azzalini & Regoli 2012). The “asymmetric Gaussian” use is, at present, motivated empirically by the often irregular appearance of the blueshifted excess in CIV and MgII $\lambda$ 2800.
- The very broad component (VBC). We expect a prominent CIV VBC associated with high ionization gas in the innermost BLR (Snedden & Gaskell 2007; Marziani et al. 2010; Wang & Li 2011; Goad & Korista 2014). The VBC was postulated because of typical H $\beta$  profile of Pop. B sources, that can be (empirically) modeled with amazing fidelity using the sum of two Gaussians, one narrower and almost unshifted (the BC) and one broader showing a significant redshift  $\sim 2000 \text{ km s}^{-1}$  (the VBC, Wills et al. 1993; Brotherton et al. 1994b; Zamfir et al. 2010). Past works provide evidence of a VBC in CIV and other lines (e.g., Marziani et al. 1996; Punsly 2010; Marziani et al. 2010).

In addition, superimposed to the broad profile there is narrower emission most likely associated with larger distances from the central continuum sources (i.e., with the narrow line region, NLR) and unresolved in the FORS spectra. The projected linear size at  $z \approx 1.5$  is 8.6 kpc/arcsec, meaning that within the ISAAC and FORS, slits were collecting light within a few kpc of the nucleus. This includes not only the innermost part of the NLR but also a significant contribution from the inner bulge of the host galaxy. A fourth component has been considered in the analysis of the integrated CIV and H $\beta$  profiles.

- The narrow component (NC). The H $\beta$ <sub>NC</sub> is not in question even in the HE sample, although the ratio H $\beta$ <sub>NC</sub> to H $\beta$ <sub>BC</sub> is on average lower at high  $L$ . Considering the large projected spatial apertures involved in our observations, the spatially-unresolved, emissivity-weighted H $\beta$ <sub>NC</sub> is expected to show a relatively narrow profile. The rationale for a significant CIV $\lambda$ 1549<sub>NC</sub> at low- $z$ , and moderate luminosity was provided by Sulentic & Marziani (1999). The lack of a well-defined critical density associated with CIV ([OIII] $\lambda$  $\lambda$ 4959, 5007 is instead suppressed at density  $n_{\text{H}} \gtrsim 10^6 \text{ cm}^{-3}$ ) and high ionization, favoring CIV emission, support the assumption that CIV NC could be up to three times broader than [OIII] $\lambda$  $\lambda$ 4959, 5007 (and H $\beta$ ) following the simple model of Netzer (1990) based on a density radial trend. Hence, the CIV $\lambda$ 1549<sub>NC</sub> profile may merge smoothly with CIV<sub>BC</sub>. Since CIV $\lambda$ 1549<sub>NC</sub> is mainly due to the inner NLR in this scenario, CIV $\lambda$ 1549<sub>NC</sub> blueshifts of a few hundred  $\text{km s}^{-1}$  per second are expected. Both CIV $\lambda$ 1549<sub>NC</sub> and H $\beta$ <sub>NC</sub> are modeled by a single symmetric Gaussian, but the H $\beta$ <sub>NC</sub> is assumed as reference for rest-frame, while CIV $\lambda$ 1549<sub>NC</sub> peak shift is allowed to vary in the CIV profile fits.

#### 2.6. Non-linear multicomponent fits and full-profile measurements of CIV $\lambda$ 1549 and H $\beta$

The data analysis of the present paper follows the same approach employed in several recent works, and consists in applying a non-linear multicomponent fit to continuum, a scalable FeII emission template, and emission line components (e.g., Negrete et al. 2013; Sulentic et al. 2014a). The minimum  $\chi^2_{\nu}$  multicomponent fit was carried out with the `specfit` routine within IRAF (Kriss 1994) yielding FWHM, peak wavelength, and intensity for the four line components BC, BLUE, VBC, and NC. An important caveat is that the line profile models are heuristic because the actual component shapes are unknown.

In addition to the CIV spectra, the H $\beta$  spectral range covered by the VLT-ISAAC was also measured using the IRAF task `specfit`. The `specfit` measurements were found to be generally consistent with the earlier ones reported in M09<sup>2</sup>. M09 FWHM,  $c(\frac{1}{2})$ , and EW of H $\beta$  are on average equal with a scatter  $\lesssim 20\%$ . Measures toward the line base are influenced by the line profile model; especially on the red side there can be differences associated with the correction by [OIII] $\lambda$  $\lambda$ 4959, 5007 emission. The correction of M09 was empirical, while in the present paper we assume an intrinsic H $\beta$  broad line shape (Sect. 2.5). This latter approach is preferable because of the strong and irregular [OIII] $\lambda$  $\lambda$ 4959, 5007 emission recently

<sup>2</sup> The measurement technique was different for the H $\beta$  observations in M09 and earlier papers by Sulentic et al. (2004) and Sulentic et al. (2006). In those papers FeII and continuum templates were sequentially subtracted before measuring H $\beta$  parameters, and not obtained through a multicomponent nonlinear fit.

<sup>1</sup> The FOS measurements used in this paper are available on CDS Vizier at <http://vizier.u-strasbg.fr>

detected in several luminous quasars (e.g., Cano-Díaz et al. 2012; Carniani et al. 2015; Marziani et al. 2016b, and references therein). The  $[\text{OIII}]\lambda\lambda 4959, 5007$  emission is modeled using a core component plus a semi-broad component, usually blueshifted (e.g., Zhang et al. 2011; Cracco et al. 2016; Marziani et al. 2016b).

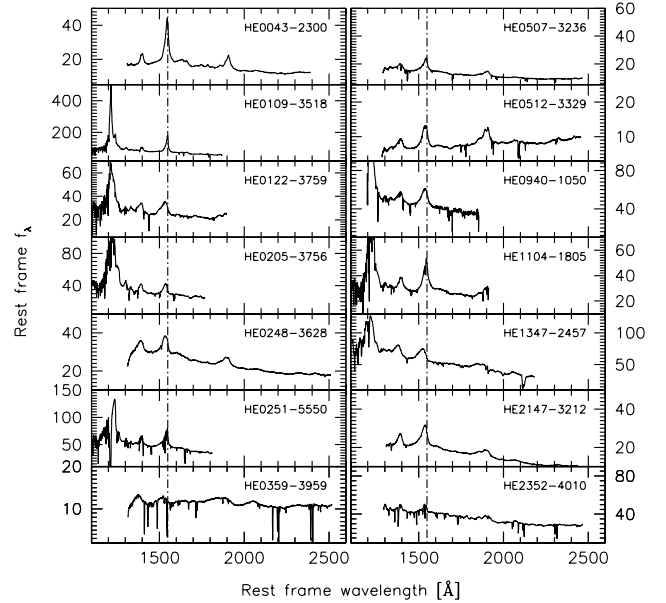
Line profile parameters (FWHM, centroids, asymmetry, and kurtosis defined by Zamfir et al. 2010) were measured on the full broad CIV and  $\text{H}\beta$  profiles (excluding the NC) to provide a qualitative description independent of the `specfit` modeling. Here we simply stress that line centroids at different fractional peak intensity ( $c(\frac{1}{4})$  for  $i = 1, 2, 3$  and  $c(0.9)$ ) are computed with respect to rest frame, while the asymmetry index (AI) at  $\frac{1}{4}$  peak line intensity, is computed with respect to the line peak. In addition  $c(\frac{1}{4})$  measures the deviations from rest frame close to the line base, while  $c(\frac{1}{2})$  measures the deviations from a symmetric and unshifted line at the fractional intensity where the broadening estimator (the FWHM) is measured. The  $c(0.9)$  is used as a proxy for the line peak radial velocity. The mock profiles of Fig. 2 in Marziani et al. (2016a) illustrate the quantitative centroid and width measurements extracted from the broad CIV profile (the narrow component is supposed to have been removed beforehand, if significant). The mock profiles show the main components expected in the CIV profile of Pop. A (BLUE and BC), and Pop. B sources (BLUE, BC, VBC).

## 2.7. Population classification

The A/B classification of the quasars in HE sample (Table 1) has been assigned by considering a limit in  $\text{FWHM}(\text{H}\beta)$  following the luminosity-dependent criterion of Sulentic et al. (2004):  $\text{FWHM}_{\text{AB}} \approx 3500 + 500 \times 10^{-0.06 \cdot (M_B + 20.24)}$  km s<sup>-1</sup>, where the  $M_B$  are reported in Table 1 (see Sulentic et al. 2004, for the derivation of the minimum FWHM as a function of the absolute magnitude).

Given the very weak dependence on  $L$ , the luminosity-dependent FWHM separation is consistent with an approximately constant limit at 4000 km s<sup>-1</sup> if  $\log L \lesssim 47$  [erg s<sup>-1</sup>]. The same population assignments are obtained if the luminosity-dependent limit  $\text{FWHM}_{\text{AB}}$  mentioned in the previous paragraph is applied to FOS sample in place of the constant limit at 4000 km s<sup>-1</sup>. If virial broadening dominates the width of  $\text{H}\beta$ , the line emitting region size scales roughly as  $r \propto L^{0.5}$ , and  $L/L_{\text{Edd}} \lesssim 1$ , then there is a minimum FWHM corresponding to the maximum  $L/L_{\text{Edd}}$ . At very high luminosity ( $\log L \gtrsim 47$  [erg s<sup>-1</sup>]), if all sources radiate sub-Eddington or close to the Eddington limit, we do not expect to observe sources with very narrow lines, such as the one falling in the domain of narrow-line Seyfert 1s (NLSy1s), as further discussed in Sect. 4.2. For instance, at  $\log L \gtrsim 47$  [erg s<sup>-1</sup>], sources with  $\text{FWHM} \lesssim 2500$  km s<sup>-1</sup> are not possible, if the conditions mentioned earlier in this section are satisfied. At the extreme  $L \approx 10^{15} L_{\odot}$ , no source may show  $\text{FWHM} \lesssim 4000$  km s<sup>-1</sup>. This underlines the inadequacy of assuming  $\text{FWHM} = 4000$  km s<sup>-1</sup> at the typical luminosity of the HE sources. The  $\text{FWHM}_{\text{AB}}$  limit is empirical (as discussed in Sect. 4.5.1 of Marziani et al. 2009). It is found a posteriori that Pop. A and B have systematically different  $L/L_{\text{Edd}}$  values (Marziani et al. 2003b; Sect. 3.3).

With respect to M09 the classification was changed for a few border-line sources close to the luminosity-dependent Pop. A/B boundary: HE1347-2457 and HE2147-3212  $\rightarrow$  Pop. A (they were previously assigned to spectral type B2,  $4000 < \text{FWHM}(\text{H}\beta) \leq 8000$  km s<sup>-1</sup>,  $1.0 \leq R_{\text{FeII}} \leq 1.5$  following



**Fig. 2.** Rest frame UV spectra of the 14 quasars belonging to Pop. A. Abscissae are rest frame wavelength in Å, ordinates are specific flux in units of  $10^{-15}$  erg s<sup>-1</sup> cm<sup>-2</sup> Å<sup>-1</sup>. The dot dashed line traces the rest-frame wavelength of the CIV line.

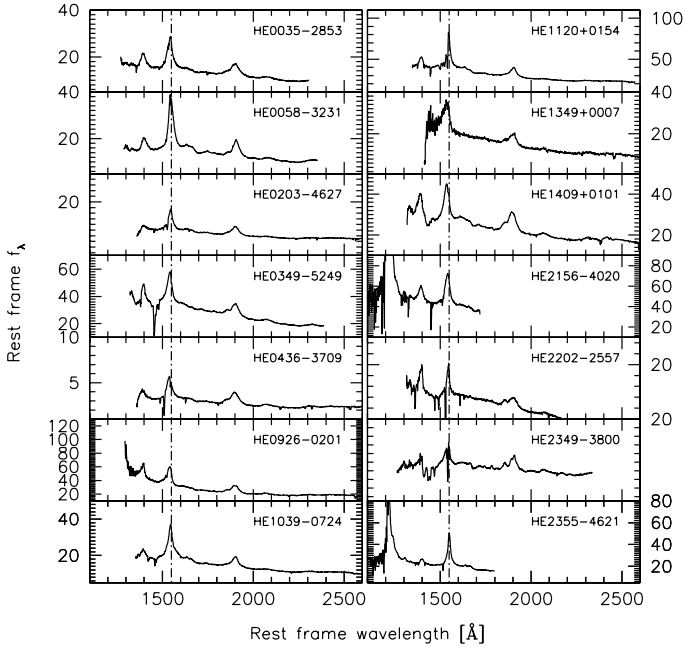
Sulentic et al. 2002), HE2156-4020  $\rightarrow$  Pop. B (was spectral type AM, which includes sources above the canonical limit at  $\text{FWHM} = 4000$  km s<sup>-1</sup>, and below the luminosity-dependent limit for Pop. A as defined in Sulentic et al. 2004).

## 3. Results

### 3.1. HE sample CIV $\lambda 1549$

We identified 14 Pop. A and 14 Pop. B sources in the HE sample. Their rest frame spectra are shown in Figs. 2 and 3, respectively. The first columns of Table 3 report the rest-frame-specific continuum flux  $f_{\lambda}$  at 1450 Å as well as the flux and equivalent width of CIV. The total flux of the 1400 feature (interpreted as a blend of  $[\text{OIV}]\lambda 1402$  and  $[\text{SiIV}]\lambda 1397$  emission) is reported in the last column; it will not be further considered in the present paper but will be used in a follow-up paper in preparation. Uncertainties on  $f_{\lambda}$  are derived from the magnitude uncertainties used in the rescaling, and are at the  $1\sigma$  confidence level. Measurements of the integrated line profile for CIV are reported in Table 4 which yields, in the following order: FWHM, asymmetry index, kurtosis, and centroids at  $\frac{1}{4}$  and  $\frac{1}{2}$  fractional intensity, along with the uncertainty at  $2\sigma$  confidence level. The results of the `specfit` analysis are shown in Figs. 4 and 5 for Pop. A and B sources, respectively. The  $[\text{OIII}]\lambda\lambda 4959, 5007$  emission is modeled with a core and a blueshifted component (Zhang et al. 2011; Marziani et al. 2016b). The  $\text{HeII}\lambda 1640$  emission on the red side of CIV has been modeled with BLUE and BC FWHM and shifts as for CIV, leaving only relative intensities as free parameters. The parameter values for the individual components are listed in Table 5. Flux normalized by the total line flux (including NC), FWHM, and centroid at  $\frac{9}{10}$  fractional intensity are reported for BLUE, BC, NC, and VBC (the VBC only in the case of Pop. B).

The amplitude of the CIV blueshifts is remarkable. Population A  $c(\frac{1}{2})$  blueshifts exceed 2000 km s<sup>-1</sup> in almost all cases, which is a shift amplitude extremely rare at  $\log L \lesssim 47$  [erg s<sup>-1</sup>] (Richards et al. 2011). In some extreme cases, the



**Fig. 3.** Rest frame UV spectra of the 14 Pop. B HE quasars presented in this paper. Abscissae are rest frame wavelength in Å, ordinates are specific flux in units of  $10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1}$ .

CIV profile appears to be fully blueshifted with respect to the rest frame. Shifts of comparable amplitude are observed only at high luminosity  $\log L \gtrsim 47$  [ $\text{erg s}^{-1}$ ] (e.g., Richards et al. 2011; Coatman et al. 2016, 2017, and references therein). Large shifts above  $1000 \text{ km s}^{-1}$  in amplitude are also observed in the CIV profile of most Pop. B sources.

### 3.2. HE sample H $\beta$ measurements paired to CIV $\lambda$ 1549

The last four columns of Table 3 report the rest-frame-specific continuum flux  $f_\lambda$  at  $5100 \text{ Å}$  as well as the flux and equivalent width of H $\beta$  and  $R_{\text{FeII}}$ . As for the UV range, uncertainties on  $f_\lambda$  are derived from the magnitude uncertainties used in the rescaling, and are at the  $1\sigma$  confidence level. The measurements on the H $\beta$  full line profile are reported in Table 6. We note in passing that there is only one source with an apparently discordant H $\beta$  FWHM with respect to its population assignment, HE2349-3800. In this case, the FWHM measurement is not stable, as the profile is the addition of a narrower BC and a much broader VBC (Fig. 5). The asymmetry index  $\text{AI} \approx 0.3$ , one of the largest in our sample, places this object among Pop. B sources. Table 7 reports the results of the *specfit* analysis for H $\beta$ , organized in the same way as for CIV. The new *specfit* H $\beta$  profile analysis yields results that are consistent with the previous work (M09) on the same objects. VBC shifts are  $\sim 2000 \text{ km s}^{-1}$  to the red, and the VBC contribution is large, above one half of the total line flux, consistent with the decomposition made by M09 on composite spectra. The large contribution of the VBC makes the BC FWHM especially uncertain.

### 3.3. HE and LOWZ sample comparison

The meaning of the CIV $\lambda$ 1549 properties in the HE sample can be properly evaluated if compared to a low- $L$  sample. The main statistical results from the inter-sample LOWZ-HE comparison are summarized in Table 8 where parameter averages and

**Table 3.** Spectrophotometric measurements on CIV $\lambda$ 1549 and H $\beta$ .

Identification	CIV $\lambda$ 1549 spectral range				H $\beta$ spectral range			
	$f_{\lambda,1450}^a$	$f(\text{CIV})^b$	$W(\text{CIV})f(\text{SiIV+OIV})^b$	$f_{\lambda,5100}^a$	$f(\text{H}\beta)^b$	$W(\text{H}\beta)R_{\text{FeII}}$	$R_{\text{FeII}}$	
	[Å]	[Å]	[Å]	[Å]	[Å]	[Å]	[Å]	
Population A								
HE0043-2300	$17.6 \pm 1.3$	8.2	48	1.7	$3.2 \pm 0.2$	2.4	68	0.37
HE0109-3518	$80.4 \pm 7.8$	19.5	26	3.4	$3.4 \pm 0.3$	4.8	137	0.16
HE0122-3759	$24.8 \pm 2.7$	5.3	22	2.4	$2.2 \pm 0.3$	1.2	50	1.12
HE0205-3756	$24.5 \pm 3.2$	4.9	18	2.4	$9.7 \pm 0.6$	5.9	56	0.40
HE0248-3628	$29.9 \pm 0.9^c$	4.3	15	2.7	$0.8 \pm 0.0$	0.4	44	0.53
HE0251-5550	$46.7 \pm 4.7$	12.8	30	3.7	$5.9 \pm 0.3$	4.0	64	0.39
HE0359-3959	$11.2 \pm 1.3$	0.6	5.3	1.2	$1.8 \pm 0.1$	0.9	50	1.10
HE0507-3236	$15.1 \pm 2.3$	4.1	29	1.5	$2.1 \pm 0.2$	1.2	67	0.31
HE0512-3329	$6.4 \pm 1.2$	3.0	45	1.2	$2.7 \pm 0.1$	1.9	86	0.63
HE0940-1050	$45.3 \pm 3.6$	9.5	23	4.6	$5.2 \pm 0.4$	2.7	45	0.55
HE1104-1805	$27.8 \pm 3.3$	10.4	41	3.3	$3.0 \pm 0.3$	3.7	121	0.56
HE1347-2457	$59.9 \pm 4.8$	8.8	15	7.4	$3.9 \pm 0.2$	1.7	38	1.22
HE2147-3212	$20.7 \pm 2.7$	5.8	29	1.8	$1.7 \pm 0.1$	1.3	70	0.77
HE2352-4010	$42.1 \pm 3.4$	4.1	10.0	1.8	$6.3 \pm 0.2$	3.5	50	0.50
Population B								
HE0035-2853	$15.5 \pm 2.3$	5.6	39	1.8	$2.0 \pm 0.1$	1.5	77	0.72
HE0058-3231	$15.3^d$	9.9	72	1.9	$1.9 \pm 0.1$	1.9	92	0.42
HE0203-4627	$10.0^d$	2.7	27	0.4	$1.9 \pm 0.1$	0.8	60	0.36
HE0349-5249	$36.1 \pm 3.3$	9.7	29	4.1	$2.8 \pm 0.1$	1.8	64	0.74
HE0436-3709	$3.2 \pm 0.3$	1.1	36	0.3	$1.4 \pm 0.1$	1.2	80	0.54
HE0926-0201	$38.4 \pm 2.7$	10.6	32	5.5	$5.1 \pm 0.2$	3.0	70	0.40
HE1039-0724	$18.3 \pm 1.5$	7.9	47	1.3	$1.7 \pm 0.1$	1.5	82	0.24
HE1120+0154	$38.0 \pm 2.3$	13.8	37	3.2	$5.7 \pm 0.2$	3.9	65	0.18
HE1349+0007	$22.9 \pm 1.8$	6.1	28	...	$1.9 \pm 0.1$	1.6	73	0.34
HE1409+0101	$26.8 \pm 1.9$	7.7	29	3.8	$4.2 \pm 0.2$	4.1	91	0.49
HE2156-4020	$43.3 \pm 5.6$	13.0	32	4.0	$2.2 \pm 0.2$	2.3	99	0.43
HE2202-2557	$14.4 \pm 1.9$	1.7	12	1.2	$1.9 \pm 0.2$	1.0	51	0.40
HE2349-3800	$8.4 \pm 1.1$	2.4	27	0.4	$1.5 \pm 0.2$	0.9	57	0.47
HE2355-4621	$21.1 \pm 2.3$	7.7	40	0.9	$2.5 \pm 0.2$	2.6	99	0.18

**Notes.** (a) In units of  $10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1}$ . (b) In units of  $10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2}$ . (c) Rescaled to the  $V$  magnitude reported in the UCAC4 (Zacharias et al. 2012). (d) Rescaled to the  $V$  magnitude reported in the NOMAD1 catalog (Zacharias et al. 2005).

sample standard deviations are reported for Pops. A and B, along with the statistical significance of the Kolmogorov-Smirnov test,  $P_{\text{KS}}$ , of the null hypothesis that the HE and LOWZ samples are undistinguishable.

The first two rows list the results for  $z$  and UV luminosity  $\lambda L_\lambda$  at  $1450 \text{ Å}$ , which are obviously different at a high significance level. The third parameter listed is the Eddington ratio. The  $L/L_{\text{Edd}}$  has been computed from the  $5100 \text{ Å}$  luminosity, assuming a bolometric correction factor of 10 and  $L_{\text{Edd}} \approx 1.5 \times 10^{38} \cdot (M_{\text{BH}}/M_\odot) \text{ erg s}^{-1}$ . The black hole mass  $M_{\text{BH}}$  was estimated using the same  $5100 \text{ Å}$  luminosity, the FWHM H $\beta_{\text{BC}}$  with the scaling law of Vestergaard & Peterson (2006). Black hole masses in the HE sample are approximately several  $10^9 M_\odot$ , close to the largest  $M_{\text{BH}}$  values observed in active quasars (Sulentic et al. 2006; Marziani & Sulentic 2012, and references therein) and spent ones in the local Universe (such as Messier 87, Walsh et al. 2013). A  $L/L_{\text{Edd}}$  systematic difference between LOWZ and HE is present for both Pops. A and B. Indeed, several of the HE Pop. B sources are close to the boundary between Pops. A and B, set from the analysis of low- $z$  quasars. This is, at least in part, the effect of a redshift-dependent  $L/L_{\text{Edd}}$  bias in a flux limited survey (Sulentic et al. 2014a). In addition, a systematic increase of  $L/L_{\text{Edd}}$  with  $z$  is also expected (Trakhtenbrot & Netzer 2012). The sample difference in  $L/L_{\text{Edd}}$  is important for interpreting any luminosity difference, and makes it advisable to consider, in addition to Pops. A and B separately, also the full A+B sample (Sect. 4.3).

The second group of rows provides CIV line parameters: FWHM,  $c(\frac{1}{2})$ ,  $c(\frac{1}{4})$ , AI, kurtosis and rest frame EW. The third

**Table 4.** Measurements on the CIV $\lambda$ 1549 full line profile.

Identification	<i>FWHM</i> [km s <sup>-1</sup> ]	A.I.	Kurt.	$c(\frac{1}{4})$ [km s <sup>-1</sup> ]	$c(\frac{1}{2})$ [km s <sup>-1</sup> ]
Population A					
HE0043-2300	5140 ± 380	-0.04 ± 0.11	0.31 ± 0.06	-1110 ± 720	-980 ± 380
HE0109-3518	3440 ± 260	-0.16 ± 0.12	0.34 ± 0.08	-1050 ± 530	-820 ± 260
HE0122-3759	8510 ± 870	-0.30 ± 0.07	0.38 ± 0.04	-4800 ± 540	-3930 ± 870
HE0205-3756	4930 ± 330	-0.22 ± 0.18	0.36 ± 0.09	-3320 ± 1080	-2920 ± 330
HE0248-3628	7360 ± 400	-0.12 ± 0.10	0.38 ± 0.05	-3950 ± 810	-2790 ± 400
HE0251-5550	4220 ± 320	-0.30 ± 0.25	0.32 ± 0.13	-2170 ± 1860	-1640 ± 320
HE0359-3959	8590 ± 490	-0.13 ± 0.08	0.42 ± 0.04	-5990 ± 610	-5880 ± 490
HE0507-3236	5410 ± 430	-0.31 ± 0.10	0.27 ± 0.06	-3100 ± 970	-1670 ± 430
HE0512-3329	7130 ± 340	-0.10 ± 0.13	0.48 ± 0.07	-2960 ± 920	-2470 ± 340
HE0940-1050	8060 ± 600	-0.15 ± 0.08	0.41 ± 0.05	-3290 ± 660	-2800 ± 600
HE1104-1805	6710 ± 370	-0.07 ± 0.10	0.42 ± 0.07	-1630 ± 830	-1500 ± 370
HE1347-2457	8260 ± 440	-0.16 ± 0.09	0.45 ± 0.05	-5850 ± 660	-5490 ± 440
HE2147-3212	7770 ± 620	-0.27 ± 0.10	0.34 ± 0.05	-4440 ± 1090	-3110 ± 620
HE2352-4010	5180 ± 390	-0.31 ± 0.08	0.24 ± 0.04	-4370 ± 720	-2390 ± 390
Population B					
HE0035-2853	6730 ± 550	0.08 ± 0.10	0.32 ± 0.06	-1080 ± 970	-1530 ± 550
HE0058-3231	6290 ± 400	0.15 ± 0.13	0.37 ± 0.08	190 ± 1190	-290 ± 400
HE0203-4627	6850 ± 460	0.04 ± 0.10	0.34 ± 0.06	-1020 ± 880	-910 ± 460
HE0349-5249	6050 ± 460	-0.13 ± 0.11	0.33 ± 0.06	-1960 ± 900	-1750 ± 460
HE0436-3709	7040 ± 580	0.03 ± 0.10	0.30 ± 0.06	-1680 ± 1070	-1880 ± 580
HE0926-0201	5980 ± 410	-0.28 ± 0.12	0.36 ± 0.07	-3280 ± 1030	-2280 ± 410
HE1039-0724	7890 ± 740	-0.12 ± 0.10	0.28 ± 0.05	-490 ± 1250	-700 ± 740
HE1120+0154	5930 ± 420	-0.06 ± 0.17	0.31 ± 0.09	-1110 ± 1730	-670 ± 420
HE1349+0007	6630 ± 460	-0.22 ± 0.10	0.35 ± 0.05	-4090 ± 840	-3040 ± 460
HE1409+0101	6370 ± 420	-0.03 ± 0.11	0.36 ± 0.07	-2900 ± 1040	-2790 ± 420
HE2156-4020	6420 ± 370	0.01 ± 0.12	0.41 ± 0.07	-1870 ± 900	-1850 ± 370
HE2202-2557	5300 ± 420	0.08 ± 0.10	0.30 ± 0.06	-1280 ± 1000	-2080 ± 420
HE2349-3800	7140 ± 490	-0.20 ± 0.11	0.34 ± 0.07	-2810 ± 1120	-2070 ± 490
HE2355-4621	5740 ± 350	0.04 ± 0.18	0.37 ± 0.09	190 ± 1360	10 ± 350

groups provide first the 4DE1 parameters  $R_{\text{FeII}}$  and  $\text{FWHM } H\beta$ , and then the profile  $H\beta$  parameters reported as for CIV.

We note in passing that the CIV equivalent width is significantly lower for both Pop. A and B, reflecting the well-known anticorrelation between equivalent width of HILs and luminosity (the Baldwin effect e.g., Baldwin et al. 1978; Bian et al. 2012). The CIV EW is also strongly affected by  $L/L_{\text{Edd}}$  (Bachev et al. 2004; Baskin & Laor 2004; Shemmer & Lieber 2015); this issue will be discussed in a follow-up paper.

The 4DE1 parameter  $R_{\text{FeII}}$  (third row) is also significantly higher than at low  $z$ , but only for Pop. B sources, probably reflecting the increase of  $R_{\text{FeII}}$  as a function of Eddington ratio (Marziani et al. 2001; Sun & Shen 2015). A similar effect is expected to operate for  $W(H\beta)$ , which is systematically higher in Pop. B than in Pop. A at low  $L$ : Table 8 indicates a significant decrease in the HE sample,  $\approx 76 \text{ \AA}$  with respect to  $\approx 100$  at low  $L$ .

A main result concerning the CIV line profile is that the LOWZ and HE CIV centroids and  $\text{FWHM}$  are significantly different at a confidence level  $\geq 4\sigma$  for both Pop. A and B (4th to 6th rows in Table 8). While at low- $L$  a strong effect of the blueshifted excess is restricted mainly to extreme Pop. A sources, at high  $L$  the effect of the blueshifted excess is significant also for Pop. B sources, as shown by the profile decomposition of Fig. 5. In the case of the most extreme blueshifts, the kurtosis value can be  $\geq 0.4$  (Table 4), although sample differences between LOWZ and HE Pops. A and B are never significant (Table 8). On average, *profile shape* (AI and kurt.) parameters are preserved for both CIV and  $H\beta$ .

The CIV  $\text{FWHM}$  is correlated with CIV  $c(\frac{1}{2})$  in Pop. A sources (see Fig. 6; cf. Coatman et al. 2016):

$$c\left(\frac{1}{2}\right) \approx (-0.62 \pm 0.07) \cdot \text{FWHM} + (1630 \pm 380) \text{ km s}^{-1}, \quad (1)$$

with a significance for the null hypothesis of  $P \approx 10^{-8}$ . The correlation can be straightforwardly explained by the assumption

**Table 5.** Results of `specfit` analysis on CIV $\lambda$ 1549.

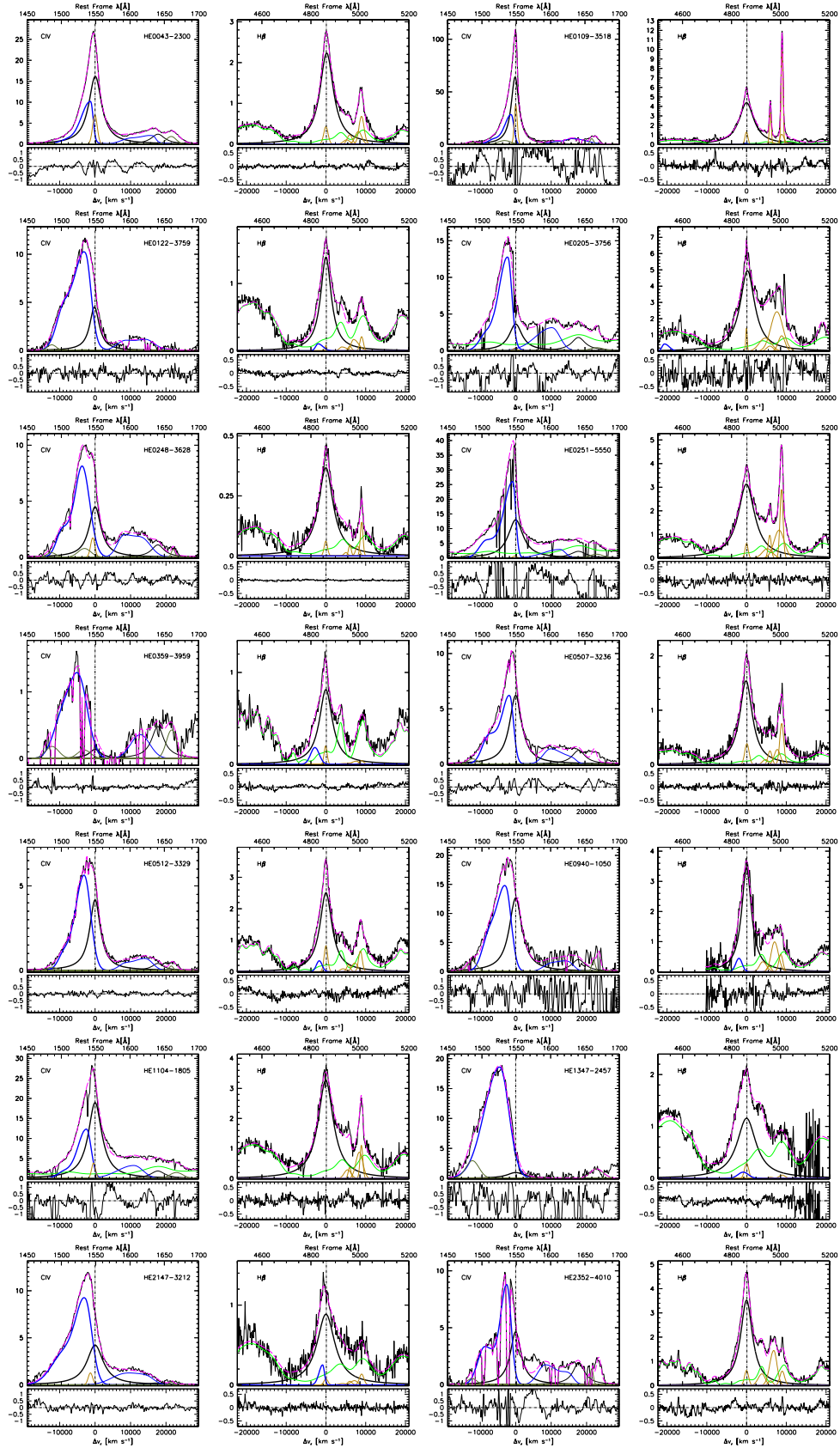
Identification	BLUE			BC			VBC			NC		
	$I/I_{\text{tot}}$	$\text{FWHM}$	Shift	$I/I_{\text{tot}}$	$\text{FWHM}$	Shift	$I/I_{\text{tot}}$	$\text{FWHM}$	Shift	$I/I_{\text{tot}}$	$\text{FWHM}$	Shift
Population A												
HE0043-2300	0.28	3840	-1600	0.65	4420	120	...	...	...	0.07	1560	30
HE0109-3518	0.27	3020	-1490	0.64	2660	-90	...	...	...	0.09	970	110
HE0122-3759	0.78	7520	-3350	0.22	3260	-90	...	...	...	0.00	...	...
HE0205-3756	0.73	4640	-2550	0.26	4430	80	...	...	...	0.00	1300	-60
HE0248-3628	0.63	5140	-3570	0.34	4070	80	...	...	...	0.03	1410	-600
HE0251-5550	0.58	4170	-1240	0.42	5050	80	...	...	...	0.00	1300	-160
HE0359-3959	0.93	7940	-5170	0.07	4050	80	...	...	...	0.00	...	...
HE0507-3236	0.48	4720	-2050	0.52	4210	80	...	...	...	0.00	...	...
HE0512-3329	0.58	5000	-3200	0.42	3700	80	...	...	...	0.00	...	...
HE0940-1050	0.53	6270	-3250	0.47	4360	80	...	...	...	0.00	...	...
HE1104-1805	0.31	4610	-2660	0.67	4650	80	...	...	...	0.02	1300	-350
HE1347-2457	0.95	7720	-4870	0.05	5880	80	...	...	...	0.00	...	...
HE2147-3212	0.67	6650	-3100	0.32	5550	80	...	...	...	0.02	1470	-1230
HE2352-4010	0.66	3760	-2680	0.34	3720	80	...	...	...	0.00	...	...
Population B												
HE0035-2853	0.43	5530	-2180	0.35	6630	40	0.17	11450	5480	0.04	1530	-260
HE0058-3231	0.18	5200	-2120	0.43	4960	-60	0.36	10620	3120	0.03	1130	-170
HE0203-4627	0.31	5520	-2210	0.43	5870	90	0.19	12390	2260	0.06	1730	70
HE0349-5249	0.52	5690	-2090	0.21	4560	-10	0.25	11360	2440	0.01	1050	10
HE0436-3709	0.48	5670	-2650	0.24	5430	80	0.28	11290	2620	0.00	...	...
HE0926-0201	0.82	5560	-2000	0.10	5090	80	0.07	13230	810	0.01	1030	50
HE1039-0724	0.29	4350	-1760	0.33	9190	80	0.30	12170	2500	0.07	2100	-250
HE1120+0154	0.39	4400	-1140	0.27	6200	360	0.25	14080	3020	0.09	1600	-180
HE1349+0007	0.87	5980	-3020	0.11	5440	80	0.00	13080	3190	0.02	2000	0
HE1409+0101	0.65	5390	-2920	0.13	7020	80	0.20	14630	2310	0.02	1400	-140
HE2156-4020	0.54	5000	-2520	0.17	4740	-30	0.28	13090	1420	0.01	860	-480
HE2202-2557	0.42	4490	-1940	0.27	7050	80	0.21	12420	3180	0.11	1210	-170
HE2349-3800	0.51	5710	-2760	0.19	3640	30	0.26	11160	1340	0.03	1090	80
HE2355-4621	0.04	5900	-3830	0.41	4660	0	0.42	14910	2150	0.13	1680	230

**Table 6.** Measurements on the  $H\beta$  full line profile.

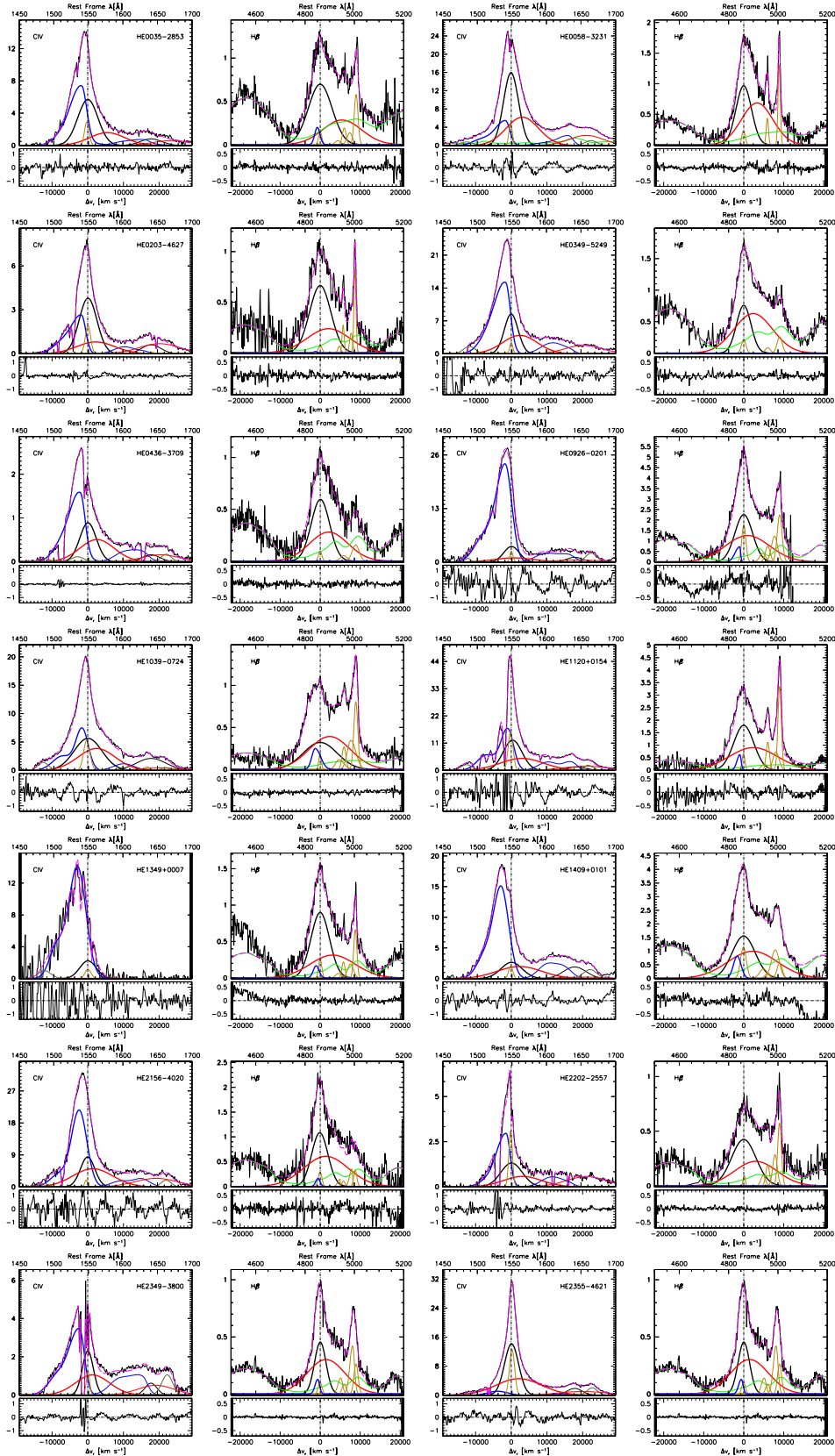
Identification	<i>FWHM</i> [km s <sup>-1</sup> ]	A.I.	Kurt.	$c(\frac{1}{4})$ [km s <sup>-1</sup> ]	$c(\frac{1}{2})$ [km s <sup>-1</sup> ]
Population A					
HE0043-2300	4150 ± 300	0.00 ± 0.10	0.33 ± 0.06	200 ± 700	200 ± 300
HE0109-3518	4450 ± 330	0.00 ± 0.10	0.32 ± 0.06	20 ± 760	20 ± 330
HE0122-3759	3440 ± 260	-0.01 ± 0.09	0.34 ± 0.05	-60 ± 500	-100 ± 260
HE0205-3756	4680 ± 350	-0.06 ± 0.10	0.33 ± 0.05	-70 ± 760	50 ± 350
HE0248-3628	4480 ± 320	0.00 ± 0.10	0.33 ± 0.06	-40 ± 760	-40 ± 320
HE0251-5550	5050 ± 360	0.00 ± 0.10	0.33 ± 0.06	-50 ± 850	-50 ± 360
HE0359-3959	5050 ± 350	-0.10 ± 0.09	0.33 ± 0.04	-500 ± 590	-610 ± 350
HE0507-3236	3880 ± 280	0.00 ± 0.10	0.33 ± 0.06	-100 ± 650	-100 ± 280
HE0512-3329	4080 ± 250	0.02 ± 0.09	0.38 ± 0.06	-20 ± 570	-180 ± 250
HE0940-1050	3420 ± 260	-0.06 ± 0.09	0.33 ± 0.05	-190 ± 470	-210 ± 260
HE1104-1805	4650 ± 330	0.00 ± 0.10	0.33 ± 0.06	-80 ± 780	-80 ± 330
HE1347-2457	5500 ± 390	0.03 ± 0.10	0.34 ± 0.06	10 ± 940	-70 ± 390
HE2147-3212	4680 ± 340	0.17 ± 0.12	0.31 ± 0.06	-30 ± 860	-160 ± 340
HE2352-4010	3640 ± 260	0.00 ± 0.10	0.33 ± 0.06	30 ± 610	30 ± 260
Population B					
HE0035-2853	7290 ± 590	0.36 ± 0.07	0.30 ± 0.05	1910 ± 1320	800 ± 590
HE0058-3231	6720 ± 640	0.31 ± 0.06	0.34 ± 0.04	2210 ± 770	1110 ± 640
HE0203-4627	7310 ± 420	0.09 ± 0.08	0.41 ± 0.05	620 ± 870	300 ± 420
HE0349-5249	6530 ± 650	0.28 ± 0.08	0.31 ± 0.04	1900 ± 970	710 ± 650
HE0436-3709	6640 ± 410	0.13 ± 0.09	0.39 ± 0.05	920 ± 950	420 ± 410
HE0926-0201	6250 ± 400	0.29 ± 0.10	0.37 ± 0.06	410 ± 1060	30 ± 400
HE1039-0724	9380 ± 690	0.26 ± 0.07	0.33 ± 0.04	1390 ± 960	1170 ± 690
HE1120+0154	6520 ± 480	0.34 ± 0.09	0.33 ± 0.06	950 ± 1330	330 ± 480
HE1349+0007	6090 ± 390	0.22 ± 0.10	0.38 ± 0.07	810 ± 1280	260 ± 390
HE1409+0101	7620 ± 530	0.34 ± 0.08	0.34 ± 0.05	1040 ± 1290	290 ± 530
HE2156-4020	5850 ± 430	0.18 ± 0.11	0.32 ± 0.06	740 ± 1290	140 ± 430
HE2202-2557	8770 ± 530	0.14 ± 0.08	0.41 ± 0.05	1330 ± 1010	720 ± 530
HE2349-3800	4390 ± 350	0.30 ± 0.12	0.28 ± 0.06	1110 ± 1220	200 ± 350
HE2355-4621	7100 ± 450	0.14 ± 0.11	0.37 ± 0.06	990 ± 1300	350 ± 450

of a blueshifted component of increasing prominence summed to an unshifted symmetric component represented through a scaled  $H\beta$  profile, as actually done with the `specfit` analysis (Sect. 3). The blueshifted excess resolved in radial velocity fully motivates the modelization of the CIV profile as due to the sum of BLUE and BC, and accounts for the CIV  $\text{FWHM}$  being larger than the  $\text{FWHM } H\beta$  by almost a factor  $\approx 2$  in Pop. A sources (Tables 4 and 6).





**Fig. 4.** Results of the `specfit` analysis on the spectra of Pop. A sources in the HE sample, in the CIV and H $\beta$  spectral range (adjacent left and right panels), after continuum subtraction. Horizontal scale is rest frame wavelength in  $\text{\AA}$  or radial velocity shift from CIV rest wavelength (*left*) or H $\beta$  (*right*) marked by dot dashed lines; vertical scale is specific flux in units of  $10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ \AA}^{-1}$ . The panels show the emission line components used in the fit: FeII emission (green), BC (thick black), and BLUE (thick blue), narrow lines (orange). The OIII] $\lambda$ 1663 line at  $\approx 1660 \text{ \AA}$  is shown in gray. Thin blue and black lines trace the BLUE and BC for HeII $\lambda$ 1640. The *lower panels* show (observed – `specfit` model) residuals.



**Fig. 5.** As in Fig. 4, for Pop. B sources, with the addition of the VBC (in red).

The large CIV blueshifts suggest that in Pop. A sources, the CIV emission is dominated by BLUE: BLUE is  $\lesssim \frac{1}{2}$  of the CIV total flux for 5 Pop. A objects and  $\geq 60\%$  for 9 of them. Conversely, 11 Pop. B sources show  $I(\text{BLUE}) \lesssim \frac{1}{2}$  the total CIV emission

$I(\text{CIV})$ , and only 3 Pop. B sources have  $I(\text{BLUE})/I(\text{CIV}) \geq 60\%$  (Table 5). The scheme in Marziani et al. (2016a) indicates that the lower shifts in Pop. B can be ascribed to the combined effect of BLUE plus a BC+VBC profile resembling the one of H $\beta$ .

**Table 7.** Results of specfit analysis on H $\beta$ .

Identification	BLUE			BC			VBC			NC		
	$I/I_{\text{tot}}$	$FWHM$	Shift [km s $^{-1}$ ]	$I/I_{\text{tot}}$	$FWHM$	Shift [km s $^{-1}$ ]	$I/I_{\text{tot}}$	$FWHM$	Shift [km s $^{-1}$ ]	$I/I_{\text{tot}}$	$FWHM$	Shift [km s $^{-1}$ ]
Population A												
HE0043-2300	0.00	...	...	0.96	4150	180	...	...	...	0.039	1290	-10
HE0109-3518	0.00	...	...	0.96	4570	0	...	...	...	0.037	860	-20
HE0122-3759	0.03	1450	-1980	0.96	3250	0	...	...	...	0.005	390	-20
HE0205-3756	0.05	...	-2640	0.93	4430	230	...	...	...	0.017	420	-20
HE0248-3628	0.00	...	...	0.98	4480	-20	...	...	...	0.025	1000	-30
HE0251-5550	0.00	...	...	0.98	5050	-40	...	...	...	0.023	890	10
HE0359-3959	0.09	2560	-2770	0.88	3960	-20	...	...	...	0.033	890	-80
HE0507-3236	0.00	...	...	0.95	3880	-130	...	...	...	0.054	1300	90
HE0512-3329	0.04	1730	-1680	0.90	3700	0	...	...	...	0.057	1060	-10
HE0940-1050	0.05	1970	-1900	0.95	2990	0	...	...	...	0.000	...	...
HE1104-1805	0.00	...	...	1.00	4650	0	...	...	...	0.000	...	...
HE1347-2457	0.03	3090	-840	0.93	5800	0	...	...	...	0.034	1240	10
HE2147-3212	0.05	2180	-830	0.94	5450	0	...	...	...	0.012	1200	0
HE2352-4010	0.08	3210	...	0.88	3640	0	...	...	...	0.035	1200	0
Population B												
HE0035-2853	0.04	1560	-710	0.57	6980	0	0.38	11240	5470	0.017	990	0
HE0058-3231	0.00	...	...	0.40	4580	0	0.58	9150	3350	0.021	1000	-60
HE0203-4627	0.00	...	...	0.58	6270	0	0.41	12390	2060	0.011	1070	0
HE0349-5249	0.00	...	...	0.30	4270	0	0.67	11250	2380	0.027	1350	0
HE0436-3709	0.00	...	...	0.50	5420	50	0.49	11190	2140	0.015	990	0
HE0926-0201	0.04	2680	-1020	0.37	5230	0	0.52	13010	950	0.067	1200	20
HE1039-0724	0.04	1420	-1170	0.54	8880	0	0.39	12170	2430	0.019	810	-100
HE1120+0154	0.03	1720	-960	0.43	5920	0	0.52	14040	2400	0.022	850	-60
HE1349+0007	0.03	1720	-1010	0.51	5450	0	0.43	13070	3110	0.025	910	0
HE1409+0101	0.06	2270	-1620	0.37	6810	-20	0.50	14560	2230	0.067	1690	-80
HE2156-4020	0.02	1340	-540	0.38	4760	-100	0.58	13050	1440	0.023	1030	-20
HE2202-2557	0.00	...	...	0.52	7170	0	0.47	12340	3000	0.010	630	120
HE2349-3800	0.03	1580	-620	0.30	3390	20	0.65	11090	1450	0.022	1110	-20
HE2355-4621	0.00	...	...	0.45	5660	40	0.52	13900	2150	0.027	1200	-80

**Table 8.** Two-sample comparison.

Parameter	A LOWZ		B LOWZ		A HE		B HE		A HE vs. LOWZ		B HE vs. LOWZ	
	Mean	$\pm\sigma$	Mean	$\pm\sigma$	Mean	$\pm\sigma$	Mean	$\pm\sigma$	$P_{\text{KS}}$		$P_{\text{KS}}$	
$z$	0.17	0.14	0.28	0.23	2.02	0.52	1.67	0.35	3.0E-10		1.2E-09	
$L$	45.42	0.79	45.76	1.00	47.69	0.36	47.48	0.22	9.3E-10		1.7E-08	
$L/L_{\text{Edd}}$	-0.56	0.43	-1.27	0.51	-0.02	0.22	-0.51	0.20	3.3E-05		5.8E-06	
$FWHM(\text{Civ})$	4321	1375	5266	1487	6479	1722	6454	661	2.8E-04		3.5E-04	
$c(\frac{1}{2})(\text{Civ})$	-1090	1245	-322	927	-3431	1586	-1656	1269	1.3E-04		2.5E-05	
$c(\frac{1}{2})(\text{Civ})$	-865	964	-250	557	-2742	1516	-1559	920	2.0E-04		5.0E-04	
$\text{AI}(\text{civ})$	-0.12	0.14	-0.04	0.14	-0.19	0.09	-0.03	0.13	3.3E-01		3.7E-01	
$\text{Kurt}(\text{Civ})$	0.36	0.08	0.33	0.05	0.37	0.07	0.34	0.03	7.8E-01		4.9E-01	
$W(\text{Civ})$	55	36	119	84	25	13	35	13	1.9E-03		7.3E-06	
$R_{\text{FeII}}$	0.66	0.48	0.27	0.21	0.62	0.32	0.42	0.17	4.2E-01		2.1E-02	
$FWHM(\text{H}\beta)$	2366	850	6390	2513	4368	630	6890	1221	6.8E-08		7.7E-02	
$c(\frac{1}{2})(\text{H}\beta)^a$	59	472	1225	1463	-63	153	1166	526	8.3E-02		2.1E-01	
$c(\frac{1}{2})(\text{H}\beta)^a$	1	309	497	886	-93	182	488	355	2.4E-01		6.1E-01	
$\text{AI}(\text{H}\beta)^a$	0.01	0.08	0.11	0.20	0.00	0.06	0.24	0.09	8.3E-02		3.0E-02	
$\text{Kurt}(\text{H}\beta)^a$	0.33	0.04	0.37	0.09	0.33	0.02	0.35	0.04	1.4E-01		1.1E-01	
$W(\text{H}\beta)^a$	88	34	108	33	68	29	76	16	1.4E-01		8.6E-04	

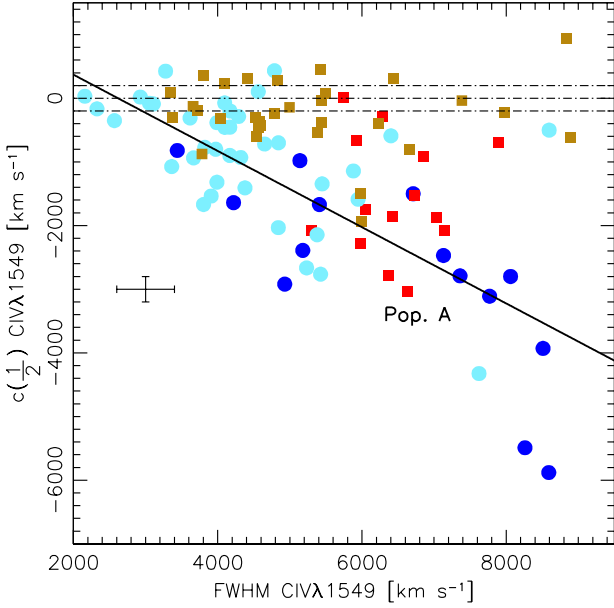
**Notes.** <sup>(a)</sup> Restricted to the sample for which H $\beta$  and Civ are both available (Sect. 2.4).

## 4. Discussion

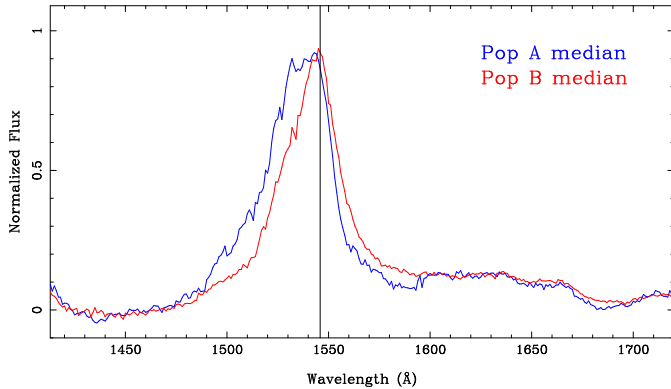
### 4.1. Civ in extreme quasars in 4DE1

The basic picture emerging from the HE sample involves a high prevalence of large Civ shifts, in several cases as large as several thousands km s $^{-1}$  with respect to the adopted rest frame, and in two cases exceeding 5000 km s $^{-1}$ . Twenty-one out of 24 HE sources show  $c(\frac{1}{2})$  Civ blueshifts larger than 1000 km s $^{-1}$  (see Table 4). This result is consistent with past finding (e.g., Richards et al. 2011) as well as with more recent works which show that quasars at the highest luminosities ( $\log L \gtrsim 47$  [erg s $^{-1}$ ]) display large Civ blueshifts (e.g., Coatman et al. 2016; Vietri 2017b).

A two-population distinction was introduced by several works, on the basis of the Civ.11549 blueshifts (Richards et al. 2011) and of the ratio FWHM to velocity dispersion (Collin et al. 2006; Kollatschny & Zetzl 2011). Population A and B quasars as defined in Sulentic et al. (2000a) provide an effective way to discriminate extremes in quasar properties. Figure 7 shows normalized median composite Civ profiles for the 14 Pop. A (blue) and 14 Pop. B (red) sources where the vertical line represents rest frame wavelength of Civ, derived from the matching H $\beta$  spectra. The median composite profiles presented in Fig. 7 show that Pop. A sources have a stronger signature of a blue component indicative of a wind or outflow. Median composite profiles are a good illustration of the individual blueshift measures presented in Table 4, where the



**Fig. 6.** Correlation between CIV FWHM and CIV  $c(\frac{1}{2})$ . Pop. A sources are indicated by filled circles and Pop. B sources by filled squares. HE sources are colored blue (Pop. A) and red (Pop. B), and LOWZ sources in pale blue (Pop. A) and gold (Pop. B). The LSQ best fit line (black) refers to Pop. A sources. The dot-dashed lines at  $\pm 200 \text{ km s}^{-1}$  mark the CIV  $c(\frac{1}{2})$  range consistent with no centroid shift. Error bars refer to  $1\sigma$  level uncertainty.

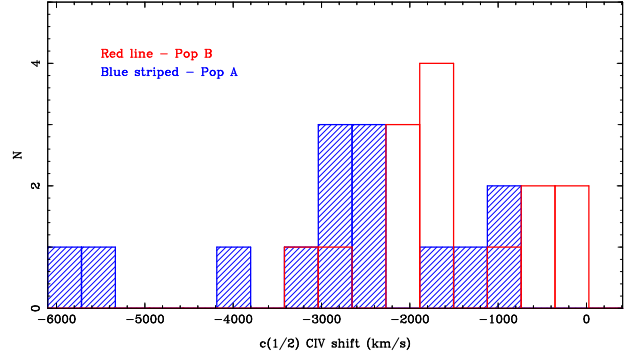


**Fig. 7.** Median composite CIV profiles for the 14 Pop. A (blue line) and 14 Pop. B (red) sources in the HE sample. Vertical line marks the source rest frame derived from the matching  $H\beta$  spectra.

high-luminosity Pop. A. profiles show blueshifts with a median of 80% of the flux on the blue side of the rest frame. The fraction of blue component flux among Pop. B. sources is 60%, also consistent with the one of the composite. In the comparison LOWZ sample only extreme Pop. A bins (A3, A4) show similar percentage of blueshift flux.

The distributions of the individual values as well as the averages and medians of profile measurements allow a first order comparison of Pop. A–B consistency in the three samples: one of lower  $L$ /higher  $z$  (GTC) and the other of higher  $L$ /higher  $z$  (HE), in addition to LOWZ.

Our LOWZ RQ sample showed Pop. A sources with a median CIV  $c(\frac{1}{2}) \approx -650 \text{ km s}^{-1}$  while Pop. B showed  $-230 \text{ km s}^{-1}$  (averages presented in Table 8 yield the same trend). A K-S test indicates a significant difference between the distribution of  $c(\frac{1}{2})$  values of LOWZ Pop. A and B at a  $4\sigma$  confidence level. The value for low- $z$  Pop. B is close to the typical uncertainty

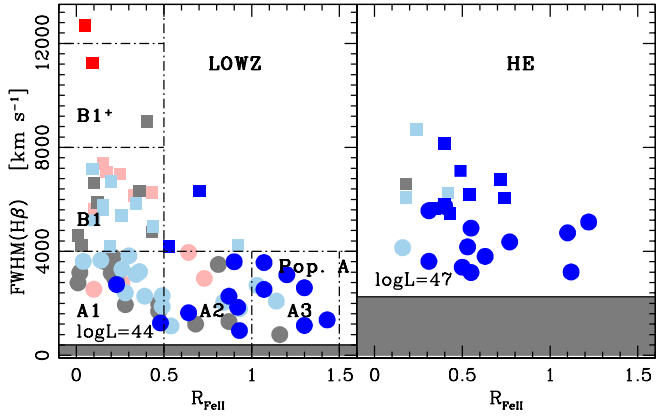


**Fig. 8.** Distribution of CIV profile shifts,  $c(\frac{1}{2})$ , for Pop. A (blue filled) and Pop. B (red line) in our extremely luminous HE quasars.

of individual measurements although a ranked-sign test (Dodge 2008) provides a Z-score  $\approx 2.3$ , for the null hypothesis that the value is not different from 0. In other words, Pop. B sources show modest blueshifts, often within the instrumental uncertainties. Hence our previous inferences that only higher  $L/L_{\text{Edd}}$  Pop. A quasars show a CIV blueshift (Sulentic et al. 2007). RL sources, which are predominantly Pop. B quasars, are overrepresented in the *Hubble* archive that provided low- $z$  CIV measures (Wills et al. 1993; Marziani et al. 1996). The median Pop. B CIV shift decreases to  $-40 \text{ km s}^{-1}$  if the RL quasars are included. Since the majority of HE (and GTC) quasars are RQ we exclude RLs from comparisons. The GTC sample showed median shifts of  $-400$  and  $-10 \text{ km s}^{-1}$  for Pops. A and B, respectively (significantly different at a  $\approx 4\sigma$  confidence level), consistent with LOWZ. In the LOWZ and GTC samples, Pop. B spectra show no significant blue flux excess with a large excess for extreme Pop. A sources (Sulentic et al. 2014a). The HE sample shows median CIV shifts of  $-2600$  and  $-1800 \text{ km s}^{-1}$  for Pops. A and B, respectively: CIV is dominated by blueshifted emission although a Pop. A–B difference persists in this high- $z$ , high- $L$  sample (Table 4). The difference between HE and LOWZ is highly significant for both Pops. A and B, as shown by K-S tests whose results are reported in Table 8: the  $P_{\text{KS}}$  corresponds to a significance level of  $\geq 3.5\sigma$ , for both CIV  $c(\frac{1}{2})$  and  $c(\frac{1}{4})$ .

Irrespective of whether one interprets the Pop. A and B designations as evidence for two distinct quasar populations (e.g., one capable of radio hyperactivity and one not) or simply a measure of extremes along a 4DE1 sequence, a search for CIV FWHM and profile shift sample differences in the Pop. A–B context gives us more sensitivity to changes due to redshift, luminosity, or Eddington ratio. Figure 8 shows a histogram of CIV shifts in the HE sample. We see that the majority of Pop. A HE sources show large blueshifts and a tendency towards smaller shifts in Pop. B thus preserving profile differences established at low  $z$ . However the difference between A and B appears to be decreasing in extreme sources. There may be a selection bias in the HE sample towards sources with higher accretion efficiency but a much larger sample will be needed to prove it (as discussed in Sect. 4.3). In the HE sample, we find only one Pop. B source without a blueshift and a few Pop. A sources with shifts similar to I Zw 1, a A3 prototype extreme blueshift source at low  $z$ , with  $c(\frac{1}{2}) = -1100 \text{ km s}^{-1}$ .

Full profile FWHM CIV measures for Pop. A and B sources yield medians of  $6900$  and  $6400 \text{ km s}^{-1}$ , respectively (averages are  $\approx 6500 \text{ km s}^{-1}$  for both populations, as reported in Table 8). These values are higher than LOWZ medians (with  $4130$



**Fig. 9.** Amplitude of CIV  $c(\frac{1}{2})$  in the 4DE1 optical plane (FWHM  $H\beta_{BC}$  vs.  $R_{FeII}$ ). Population A and B sources are marked as circles and squares, respectively. Colors represent shift amplitudes. Blue: large CIV blueshifts ( $c(\frac{1}{2}) > -1000$  km s $^{-1}$ ), aquamarine: significant blueshifts ( $c(\frac{1}{2}) = -200$  to  $-1000$  km s $^{-1}$ ); gray: no significant shift ( $|c(\frac{1}{2})| < 200$  km s $^{-1}$ ). Red and pink, large and significant  $c(\frac{1}{2})$  redshift, respectively. *Left:* 4DE1 plane for the RQ LOWZ sample. *Right:* 4DE1 plane for HE sample. The horizontal filled lines mark the minimum FWHM  $H\beta$  for virialized systems at  $\log L = 44$  and  $47$  (no sources may exist in the shaded area, Marziani et al. 2009).

and 4910 km s $^{-1}$  for Pop. A and B, respectively) and GTC ( $\approx 5200$  km s $^{-1}$  and 5800 km s $^{-1}$ ) estimates. The LOWZ and GTC FWHM CIV distributions differ from the HE sample with statistical significance  $\geq 4\sigma$ . The CIV FWHM in the HE sample is higher for both Pops. A and B due to the growing strength of the wind/outflow component shown in Sect. 3.3 – as CIV blueshift increases, FWHM CIV is likely to increase as well (Fig. 6; see also Fig. 7 in Coatman et al. 2016). This non-virial broadening associated with the wind (e.g., Richards et al. 2011) leads to overestimation of FWHM CIV-based black hole masses (e.g., Netzer et al. 2007; Sulentic et al. 2007; Park et al. 2013; Mejía-Restrepo et al. 2016; Coatman et al. 2017, and references therein). The amplitude of the effect depends on the location of 4DE1 sequence; being stronger for Pop. A (Sulentic et al. 2007; Brotherton et al. 2015).

#### 4.2. The Eigenvector 1 optical plane at low- and high- $L$

Figure 9 incorporates CIV into the 4DE1 optical plane using color to represent the strength of the CIV blueshift. The LOWZ sample (left panel) illustrates the strong preference for CIV blueshifts but only in Pop. A sources and especially sources with higher  $L/L_{Edd}$  in bins A2 and A3. Large blueshifts are rare in LOWZ Pop. B. The extreme HE sources presented in this paper (right panel Fig. 9) show a larger fraction of CIV blueshifts due to the blue excess in the CIV profiles of most sources. The FWHM( $H\beta$ ) of HE sample shift upwards as expected if  $H\beta$  arises in a structure where profile width is dominated by virial motions (Sect. 2.7; Marziani et al. 2009, a prime candidate is the accretion disk).

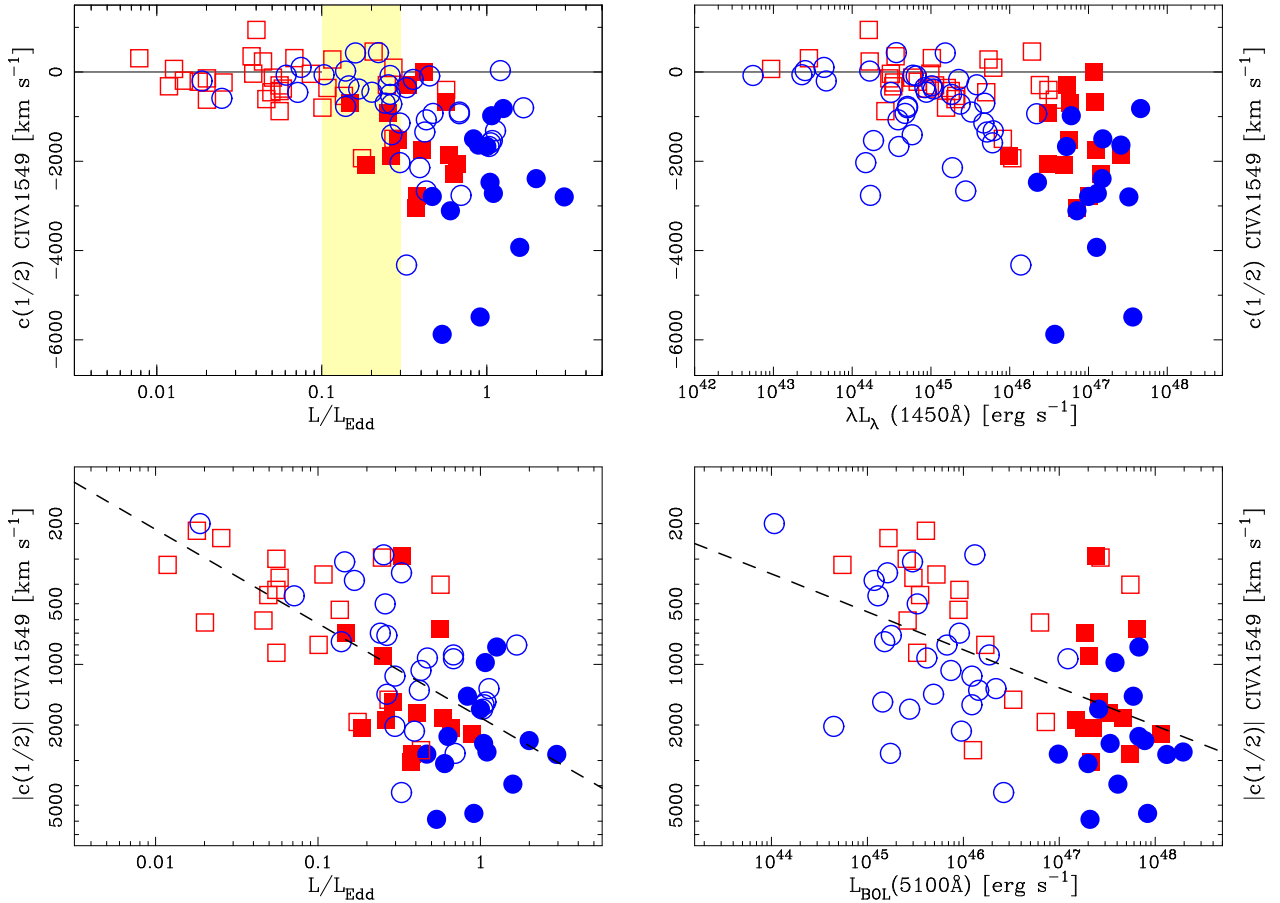
The advantage of the present work is to introduce accurate rest frame estimates and S/N high enough to decompose the CIV line profile in a symmetric virialized component and a blue shifted excess at high  $L$ . Clear trends then emerge by joining low- and high- $L$  sources which are not even guessed at in low- $z$  samples. The question is now whether the change in the HE sample is driven by  $L$  or  $L/L_{Edd}$ .

#### 4.3. Is it $L$ or $L/L_{Edd}$ that drives the CIV outflow?

Evidence at low  $z$  suggests that  $c(\frac{1}{2})$  CIV is an Eigenvector 1 parameter with  $L/L_{Edd}$  likely to be the driver of the wind/outflow. Source luminosity,  $L$ , has long been known to be an Eigenvector 2 parameter (Boroson & Green 1992) and therefore not strongly related to any of the 4DE1 parameters which, in low- $L$  samples have been shown to depend more strongly on  $L/L_{Edd}$  than on  $L$  (Marziani et al. 2003b; Sun & Shen 2015). In other words, luminosity effects are not strong in type-1, and, for a moderate luminosity range (1–2 dex), the diversity of quasars may be mainly accounted for by  $L/L_{Edd}$ , along with other parameters not immediately related to  $L$ , such as the viewing angle (e.g., Marziani et al. 2001). The addition of CIV measures for extreme HE sources, coupled with our low  $z$  sample, makes it possible to search for correlations over the entire range of  $L/L_{Edd}$  and over a range of more than 5 dex in quasar luminosity.

Figure 10 compares the correlations of  $c(\frac{1}{2})$  with  $L$  and  $L/L_{Edd}$  for 54 Pop A and 45 Pop B quasars in HE and LOWZ samples. We find weak correlation between  $L$  and  $c(\frac{1}{2})$  with sources showing large and similar CIV shifts over four of the five dex ranges in source luminosity. Excluding the four extreme blueshifts with  $\log L \geq 47$  [erg s $^{-1}$ ], blueshifts between 1000 and 4000 km s $^{-1}$  are observed over 4 dex in luminosity, as shown in Fig. 10 (upper right). The weak trend is consistent with previous results based on SDSS that found a dependence on luminosity with a large scatter (Richards et al. 2011; Shen et al. 2016). The left panels of Fig. 10 show a clearer trend between  $L/L_{Edd}$  and  $c(\frac{1}{2})$ . The largest blueshifts ( $>1000$  km s $^{-1}$ ) are observed for  $L/L_{Edd} \geq 0.2 \pm 0.1$ . In fact the observed distribution of points in the blueshift- $L/L_{Edd}$  diagram supports the suggestion of a discontinuity between  $L/L_{Edd}$  and  $c(\frac{1}{2})$  – a critical value of  $L/L_{Edd}$ , above which high amplitude winds are possible. If this interpretation is correct, then the Pop. A–B hypothesis would argue that the largest blueshifts are a Pop. A phenomenon and only arise in sources with  $L/L_{Edd} \geq 0.2$ . Figure 10 epitomizes a challenging aspect of the CIV blueshift interpretation: on the one hand, largest shifts are observed only at high  $L$ , at least in our samples; on the other hand,  $c(\frac{1}{2}) \lesssim -1000$  km s $^{-1}$  are not observed at low  $L/L_{Edd}$ . This is not to say that outflows do not occur at low  $L/L_{Edd}$ , but that the outflowing component BLUE is not dominating the line broadening. The threshold at  $L/L_{Edd} \approx 0.2 \pm 0.1$  would provide a necessary condition for such dominance to occur, with large negative values of  $c(\frac{1}{2})$  possible also at modest  $L$ . In other words, high  $L/L_{Edd}$  appears the condition sine qua non that makes possible outflows dominating the HIL profiles.

The lower panels of Fig. 10 center the attention on the 72 sources with larger blueshifts ( $>200$  km s $^{-1}$ ). The logarithm of  $|c(\frac{1}{2})|$  is shown as a function of  $L/L_{Edd}$  (left) and  $L$  (right) for 42 Pop. A and 30 Pop. B sources in the LOWZ and HE samples. The correlation results for the larger blueshift sources are reported in Table 9 where we list the correlation coefficients, computed following Pearson’s definition and the probability of a spurious correlation  $P$ . Linear fit parameters (intercept  $a$  and slope  $b$ ), obtained with unweighted least squares (LSQ) and with the orthogonal method (Press et al. 1992) are reported in the following columns. The orthogonal method is known to provide a robust estimator of fit parameters if they are not normally distributed. The values of the Pearson’s correlation coefficient  $r$  suggest significant correlations in all cases, at a confidence level  $>3\sigma$ . The values of  $r$  and fitting parameters are listed for Pops. A and B separately first. We then join Pops. A and B to improve the statistics since we do not notice



**Fig. 10.** Dependence of CIV  $c(\frac{1}{2})$  on  $L/L_{\text{Edd}}$  (left panels), on  $\lambda L_{\lambda}(1450 \text{ \AA})$  (right upper), and on  $\log L$  derived from the continuum at  $5100 \text{ \AA}$  (lower right). Population A sources are indicated by blue circles and Pop. B sources by red squares. Filled symbols represent HE sources and empty ones correspond to low- $z$  sources. The pale yellow band corresponds to  $L/L_{\text{Edd}} = 0.2 \pm 0.1$ . The dashed lines in bottom frames trace unweighted least squares.

**Table 9.** Results of correlation analysis  $c(\frac{1}{2})$  vs.  $\log L$  and  $L/L_{\text{Edd}}$ .

Sample	Var.	$r$	$P$	LSQ <sup>a</sup>		Orthogonal <sup>a</sup>		$r$ (partial) <sup>b</sup>
				$a$	$b$	$a$	$b$	
A	$L$	0.472	2.5E-03	$-3.867 \pm 2.095$	$0.151 \pm 0.045$	$-3.121 \pm 1.721$	$0.135 \pm 0.037$	$0.252 \pm 0.152$
B	$L$	0.617	7.2E-04	$-8.443 \pm 2.361$	$0.244 \pm 0.051$	$-9.569 \pm 2.468$	$0.268 \pm 0.053$	$0.268 \pm 0.179$
A+B	$L$	0.465	9.0E-05	$-4.818 \pm 1.870$	$0.169 \pm 0.040$	$-4.120 \pm 1.906$	$0.154 \pm 0.041$	$0.262 \pm 0.116$
A	$L/L_{\text{Edd}}$	0.537	5.9E-04	$3.269 \pm 0.049$	$0.527 \pm 0.102$	$3.346 \pm 0.078$	$0.749 \pm 0.208$	...
B	$L/L_{\text{Edd}}$	0.622	6.6E-04	$3.294 \pm 0.118$	$0.458 \pm 0.096$	$3.435 \pm 0.122$	$0.619 \pm 0.113$	...
A+B	$L/L_{\text{Edd}}$	0.625	1.4E-07	$3.265 \pm 0.045$	$0.451 \pm 0.048$	$3.334 \pm 0.055$	$0.584 \pm 0.084$	...

**Notes.**  $\log |c(\frac{1}{2})|$  in  $\text{km s}^{-1}$ , for  $c(\frac{1}{2}) \leq -200 \text{ km s}^{-1}$ , yielding 42 and 30 Pop. A and B sources.  $\log L = \log \lambda L_{\lambda}(1540 \text{ \AA}) + \text{BC}$ ,  $\text{BC} = \log(3.5)$ ,  $\log \lambda L_{\lambda}(1540 \text{ \AA})$  gives same slope  $b$  and intercept  $a' = a - 0.544$ . <sup>(a)</sup> Least square fit and orthogonal method parameters computed with SLOPES (Feigelson & Babu 1992). <sup>(b)</sup>  $L/L_{\text{Edd}}$  hidden variable.

relevant differences between the two populations. A significant correlation (slope  $b \approx 0.45$  or  $\approx 0.58$  for LSQ and orthogonal method, respectively) between blueshift CIV and  $L/L_{\text{Edd}}$  is found. Sources of the joined Pop. A and B sample (last row of Table 9) show an increase in blueshift amplitude with  $L/L_{\text{Edd}}$  and a Pearson's correlation coefficient of  $\approx 0.625$ , corresponding to a confidence level of  $\approx 5.2\sigma$  and a probability  $P \sim 10^{-7}$  for the null hypothesis of uncorrelated data. A weaker (slope  $b \approx 0.15$  for both LSQ and orthogonal method) but still significant dependence of  $c(\frac{1}{2})$  on bolometric  $L$  is also found (Fig. 10 bottom right). The luminosity trend becomes appreciable over a very wide range in luminosity (4 dex), when both the LOWZ and HE samples are put together, with Pearson's correlation coefficient of  $\approx 0.465$  implying a confidence level  $\approx 3.9\sigma$  (Table 9).

The higher frequency of large blueshifts at high luminosity may reflect a bias in  $L/L_{\text{Edd}}$  distribution of the observable quasars if large shifts are dependent on  $L/L_{\text{Edd}}$ . The partial correlation coefficient (e.g., Wall & Jenkins 2012) between  $c(\frac{1}{2})$  and  $L_{1450}$  with  $L/L_{\text{Edd}}$  as a hidden variable is presented in the last column of Table 9, along with its  $1\sigma$  uncertainty. The partial correlation reveals that, in our sample, the  $L/L_{\text{Edd}}$  bias is not strong enough to account for the  $c(\frac{1}{2}) - L$  correlation: the A+B partial correlation coefficient of  $c(\frac{1}{2}) - L$  is significant (at slightly more than  $2\sigma$  confidence level) if the  $L/L_{\text{Edd}}$  is considered as the third variable affecting  $c(\frac{1}{2})$  (Table 9). The residual luminosity dependence is rather evident from our data, as  $c(\frac{1}{2})$  blueshifts larger than  $2000 \text{ km s}^{-1}$  are not found in the LOWZ sample for

$\lambda L_{\lambda}(1450 \text{ \AA}) \lesssim 10^{44} \text{ erg s}^{-1}$ . It is unlikely that a large fraction of high-amplitude blueshifts would have gone undetected at  $10^{44} \text{ erg s}^{-1} \lesssim L \lesssim 10^{47} \text{ erg s}^{-1}$ .

#### 4.4. Inferences on the outflowing component dynamical conditions

The main results of the present investigation involve (1) the detection of extreme CIV blueshifts, (2) the dependence of  $c(\frac{1}{2})$  on  $L/L_{\text{Edd}}$ , and (3) the weaker dependence on luminosity that becomes appreciable over a wide range in luminosity (i.e., large blueshifts in Pop. B at high  $L$  even if their  $L/L_{\text{Edd}}$  is  $\lesssim 0.3$ ). The threshold above which large CIV blueshifts ( $\gtrsim 1000 \text{ km s}^{-1}$ ) are observed may be associated with a change in accretion disk structure: at  $L/L_{\text{Edd}} \gtrsim 0.2$ , a geometrically and optically thick accretion (“slim”) disk may form in an advection-dominated accretion flow (ADAF; e.g., Abramowicz et al. 1988; Abramowicz & Straub 2014, and references therein). In this case, outflows may be confined at the edge of a funnel (Sądowski et al. 2014), and geometry is expected to favor large shift amplitudes. Estimates of the radiative efficiency suggest that “slim” disks may be very frequent at high-luminosity and high- $z$  where highly accreting quasars are preferentially selected (Netzer & Trakhtenbrot 2014; Trakhtenbrot et al. 2017). The slim configuration may develop in the innermost part of the disk, while further out the disk is expected to retain a radiatively-efficient, geometrically thin regime (e.g., Frank et al. 2002, and references therein). The interplay between wind and the BLR in general and the disk structure at high  $L/L_{\text{Edd}}$  is a complex problem that it is just beginning to be faced (e.g., Wang et al. 2014, and last paragraph of Sect. 4.5) and that goes beyond the scope of the present paper. In the following, we mainly restrict our discussion to basic considerations on the outflow acceleration mechanism.

Wind models suitable for accretion disks around supermassive black holes can be broadly grouped into two major classes: models based on magnetohydrodynamic centrifugal acceleration (Emmering et al. 1992; Everett 2005; Elitzur et al. 2014), and radiation-driven models (Murray et al. 1995; Elvis 2000; Everett 2005; Proga 2007; Risaliti & Elvis 2010). Both acceleration mechanisms predict that the outflow terminal velocity is strongly dependent on Eddington ratio: for magneto-centrifugal acceleration, the terminal velocity scales as  $v_{\text{t,mag}} \propto L^{\frac{1}{3}} L/L_{\text{Edd}}$ . This case implies the conservation of the angular momentum of the gas originally in the Keplerian accretion disk. Predicted line profiles are fairly symmetric and often show a prominent core. The second class can more easily account for fully blueshifted profiles as observed in Pop. A sources because radiation-driven winds are expected to give rise to a predominantly radial outflow. Results (1) and (2) also indicate a predominant role of radiative forces, and condition (3) is expected in the case of a radiation-driven outflow.

When radiation forces dominate (which may be the case of a radiation driven wind, or of discrete clouds), the outflow terminal velocity can be written as  $v_t \propto v_K(r) \sqrt{\mathcal{M} \frac{L}{L_{\text{Edd}}}} \propto M_{\text{BH}}^{\frac{1}{2}} L^{-\frac{\alpha}{2}} \sqrt{\mathcal{M} \frac{L}{L_{\text{Edd}}}}$ , if  $r \propto L^{\alpha}$  (e.g., Laor & Brandt 2002; Netzer & Marziani 2010).  $\mathcal{M}$  is the force multiplier, and  $v_K$  the Keplerian velocity at the wind launching radius. The slope of the  $L/L_{\text{Edd}}$  ratio dependence (Table 9) is in agreement with the predictions of a dominant role of radiation forces. A concomitant dependence on luminosity occurs for increasing black hole mass if the term  $v_K$  or  $v_K \sqrt{\mathcal{M}}$  depend on luminosity. For example,

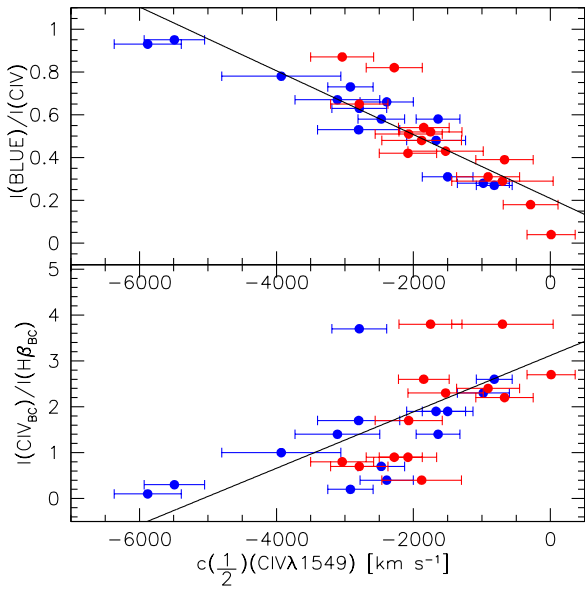
$v_K \sqrt{\mathcal{M}} \propto L^k$  with  $k \approx 0.15$  is possible if the force multiplier scales roughly with the Eddington ratio.

#### 4.5. Physical conditions of the line emitting gas

The increase in CIV FWHM along with CIV shift argues against a smaller value of the emissivity weighted distance for CIV in a virial velocity field as the dominant broadening effect for CIV; if this were the case, the profile would remain roughly similar to the one of  $H\beta$ .

The balance between radiative and gravitational forces is not the same for the whole BLR. Specifically, if we consider the high-ionization and low-ionization lines (HILs and LILs, respectively) emitting BLR as two distinct sub-regions, radiation forces may dominate in the high-ionization sub-region, while gravitational ones in the denser and thicker low-ionization sub-region (Ferland et al. 2009; Marziani et al. 2010; Negrete et al. 2012). The  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}})$  intensity ratio values (which can be deduced from the data reported in Tables 3, 5, and 7) imply low ionization for the gas emitting the virialized BC, and confirm much higher ionization for emission associated with BLUE. The median values for the BC intensity ratios  $I(\text{CIV})/I(\text{H}\beta)$  are  $\approx 1.4$  and  $\approx 2.2$ , and for BLUE  $I(\text{CIV BLUE})/I(\text{H}\beta \text{ BLUE}) \approx 34$  and 60 (Pops. A and B, respectively, Tables 5 and 7). The overall scenario envisaged at low  $z$  is basically confirmed for the high- $z$  quasars. The low BC  $I(\text{CIV})/I(\text{H}\beta)$  along with several other diagnostic flux ratios such as  $I(\text{Ly}\alpha)/I(\text{CIV})$ ,  $I(\text{AlIII}\lambda 1860)/I(\text{SiIII}\lambda 1892)$ ,  $I(\text{MgII}\lambda 2800)/I(\text{Ly}\alpha)$ ,  $I(\text{FeII})/I(\text{H}\beta)$ , and  $I(\text{Ly}\alpha)/I(\text{H}\beta)$  suggested that the BC is being emitted by a dense region of low-ionization ( $n_{\text{H}} \sim 10^{11} - 10^{12} \text{ cm}^{-3}$ , ionization parameter  $\sim 10^{-2} - 10^{-2.5}$ ) while BLUE requires lower density and higher ionization ( $n_{\text{H}} \lesssim 10^{10} \text{ cm}^{-3}$ , ionization parameter  $\sim 10^{-1}$ , Marziani et al. 2010). The  $I(\text{CaII})/I(\text{H}\beta_{\text{BC}})$  intensity ratio measured by Martínez-Aldama et al. (2015) for several sources of the HE sample also indicates the existence of a region of low ionization and high density as the emitter of the BC, and sets a robust lower limit on the density of the line emitting gas,  $\sim 10^{10.5} \text{ cm}^{-3}$  (as found by Matsuoka et al. 2007). The inferences on the physical conditions of the two emitting regions add further support to the idea that the emitting gas resolved in radial velocity – that is, emitting BC and BLUE – is also spatially segregated (Collin-Souffrin et al. 1988; Elvis 2000).

Figure 11 shows the relation between  $c(\frac{1}{2})$  and the prominence of BLUE with respect to total CIV emission (upper panel), and the  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}})$  intensity ratio (lower panel). The BLUE prominence and  $c(\frac{1}{2})$  are almost perfectly anti-correlated (Pearson’s  $r \approx -0.90$ ), with intensity ratio  $I(\text{CIV BLUE})/I(\text{CIV}) \approx (-0.15 \pm 0.014) c(\frac{1}{2}) + 0.21 \pm 0.04$ , which is not surprising given the definition of BLUE. More interestingly,  $c(\frac{1}{2})$  and the intensity ratio of the broad component of CIV and  $H\beta$  are also correlated ( $r \approx 0.65$ ):  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}}) \approx (0.615 \pm 0.14) c(\frac{1}{2}) + (3.12 \pm 0.36)$ . As the shifts grow larger, the BC of CIV diminishes dramatically. The equivalent width  $W(\text{CIV})$  is anti-correlated with  $c(\frac{1}{2})$  and the  $c(\frac{1}{2})$  amplitude is correlated with  $R_{\text{FeII}}$  (Marziani et al. 2016a): the most extreme shifts are associated with the weakest emission lines, and strongest FeII emitters. In other words, the largest blueshifts are associated with a very low ionization degree:  $c(\frac{1}{2})$  is highly correlated with  $R_{\text{FeII}}$ :  $r \approx -0.82$ , corresponding to a significance above  $4\sigma$  for 28 objects. In addition, minimum  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}})$  and the BLUE/CIV $\lambda$ 1549 are also associated with the strongest FeII emission: BLUE/CIV $\lambda$ 1549 is correlated with  $R_{\text{FeII}}$  with  $r \approx 0.63(3.9\sigma)$ ; and  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}})$



**Fig. 11.** Relation between  $c(\frac{1}{2})$  CIV and intensity ratio between the BLUE component and the full CIV line (*upper panel*), and intensity ratio of the broad components of CIV and H $\beta$  (*lower panel*) for the sources of the HE sample. The horizontal scale is in units of km s<sup>-1</sup>. Population A sources are indicated by blue symbols and Pop. B by red ones.

is anti-correlated with  $r \approx -0.36$  ( $\approx 2\sigma$ ). This is why it is so difficult to identify a correction to the FWHM of CIV that is directly connected to  $c(\frac{1}{2})$ : the  $c(\frac{1}{2})$  values are large when blue dominates the CIV emission, and the CIV BC is concomitantly very faint, especially in Pop. A. When the CIV shift is small, however, it might be possible to recover a CIV-based virial estimator. This issue will be the subject of an eventual paper.

Extreme Pop. A ( $R_{\text{FeII}} \gtrsim 1.0$ , in the HE sample: HE0122-3759, HE0359-3959, HE1347-2457) all have  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}}) \lesssim 1$ . The ratio  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}})$  for HE0359-3959 is just  $\approx 0.1$ , since the CIV is almost fully blue shifted. In general, density and ionization parameter appear to reach extreme values in correspondence of sources with  $R_{\text{FeII}} \gtrsim 1$ ,  $\log n_{\text{H}} \sim 10^{12} \text{ cm}^{-3}$ , and  $\log U \lesssim -2.5$  (Negrete et al. 2012). These sources (called xA by Marziani & Sulentic 2014) show very low equivalent width of emission lines along the 4DE1 main sequence. Several (but not all; in our sample HE0359-3959 and the borderline object HE2352-4010) meet the definition of weak-lined quasars (WLQs, Diamond-Stanic et al. 2009,  $W(\text{CIV}) \leq 10 \text{ \AA}$ ). The WLQs in our sample show high-amplitude blueshifts:  $-5900$  (HE0359) and  $-2390$  (HE2352) and also very low  $I(\text{CIV}_{\text{BC}})/I(\text{H}\beta_{\text{BC}}) \lesssim 0.5$ . WLQs as a class are also characterized by high-amplitude CIV $\lambda 1549$  blueshifts (Plotkin et al. 2015; Shemmer & Lieber 2015). WLQs properties are, in general, consistent with the properties of xA sources. It is interesting to note that WLQ are present also at low  $L$ . By joining the Sulentic et al. (2007) and the Plotkin et al. (2015) data, Marziani et al. (2016a) showed that 9 out of 11 sources are consistent with the xA location in the 4DE1 optical plane FWHM(H $\beta$ ) versus  $R_{\text{FeII}}$ . In turn, xA sources show low  $W \text{ CIV} \lesssim 30 \text{ \AA}$  could perhaps provide a more appropriate physical limit for isolating the weakest CIV $\lambda 1549$  emitters: WLQs are extreme cases of xA sources but apparently not fundamentally different from the other xA sources, as they show prominent HIL blueshifts along with strong FeII emission.

The results of the present paper yield some constraints on the structure of the BLUE and BC emitting regions in Pop. A. The LIL-BLR may be exposed to a low-level ionizing continuum, due to scatter or self-absorption by the high-ionization wind (not necessarily the same gas emitting the CIV line, e.g., Leighly et al. 2007). An alternative is that the LIL-BLR may be shadowed from the continuum by the optically and geometrically thick structure as provided by an ADAF. At the same time, the 1900 blend of HE sources (described in the follow-up paper) shows a rather high  $I(\text{AlIII}\lambda 1860)/I(\text{SiIII}\lambda 1892)$  ratio, implying high density following Negrete et al. (2012). This result suggests that LIL emission may occur from a remnant emitting regions made of the densest gas, while much-lower-density gas is being ablated away and dispersed into the quasar circumnuclear regions.

## 5. Conclusions

Comparison of low (H $\beta$ ) and high (CIV) excitation broad emission line profiles in high- and low-redshift quasars offers a method to search for quasar evolution in redshift, luminosity, and source Eddington ratio. However there are three desiderata in order to make such comparisons: 1) an accurate measure of the local rest frame in each source, 2) spectral S/N must be high enough to allow decomposition of both H $\beta$  and CIV, and 3) a context in which one can interpret quasar diversity (i.e., Pop. A and B discrimination); this explains why there have been few studies of this kind (e.g., Espey et al. 1989; Marziani et al. 1996; Sulentic et al. 2007; Coatman et al. 2016). Low redshift comparisons show that the two lines are very different in Pop. A and B quasars. This paper combines samples of low- $z$  (ground based/HST-FOS), high- $z$ /low- $L$  (GTC sample) and high- $z$ /high- $L$  (HE ISAAC/FORS and TNG-LRS sample) quasars which explore the full observed ranges of  $L$  and  $L/L_{\text{Edd}}$ . CIV measures of extreme quasars show blueshift values among the largest ever observed. Previously published H $\beta$  spectra provide accurate rest frame measures, Eddington ratios and Pop. A–B classifier while high S/N CIV measures, through accurate line profile decomposition, add insights into disk structure and especially into outflows or winds arising from the disk. Major results are:

1. CIV outflows become stronger and more frequent at higher  $z$  in sources with 2–4 dex higher luminosity than low- $z$  quasars. The emission component associated with the outflow becomes a dominant source of line broadening, and confirms that the full CIV profile is not a useful virial estimator of black hole mass for most sources.
2. We find that the 4D Eigenvector 1 sequence persists but is less strong among the extreme HE quasars as far as CIV blueshifts are concerned: many Pop B. sources show a significant CIV blueshift. In the optical plane of the 4DE1 parameter space, data points are displaced upwards toward larger FWHM(H $\beta$ ) values, as expected for the systematic increase of line width with luminosity in virialized systems.
3. We show evidence that the CIV outflows are driven by  $L/L_{\text{Edd}}$  rather than source luminosity. A clear dependence of CIV blueshifts on Eddington ratio is observed while a weaker trend is observed with source luminosity that becomes appreciable only over a wide luminosity range (4–5 dex). Large shifts ( $c(\frac{1}{2}) \lesssim -1000 \text{ km s}^{-1}$ ) are apparently possible above a threshold  $L/L_{\text{Edd}} \approx 0.2 \pm 0.1$ .
4. The relations between  $c(\frac{1}{2})$  and  $L/L_{\text{Edd}}$  and  $L$  are consistent with a radiation-driven outflow, and the largest shifts may



support the existence of an optically thick ADAF in at least some of the Pop. A sources.

**Acknowledgements.** A.d.O., J.W.S. and M.L.M.A. acknowledge financial support from the Spanish Ministry for Economy and Competitiveness through grants AYA2013-42227-P and AYA2016-76682-C3-1-P. J.P. acknowledges financial support from the Spanish Ministry for Economy and Competitiveness through grants AYA2013-40609-P and AYA2016-76682-C3-3-P. D.D. and A.N. acknowledge support from grants PAPIIT108716, UNAM, and CONACyT221398.

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