

## Wing Transmission for a Micromechanical Flying Insect \*

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### Abstract

*Flapping wings provide unmatched maneuverability for flying micro-robots. Recent advances in modelling insect aerodynamics show that adequate wing rotation at the end of the stroke is essential for generating adequate flight forces. We developed a thorax structure using four bar frames combined with an extensible fan-fold wing to provide adequate wing stroke and rotation. Flow measurements on a scale model of the beating wing show promising aerodynamics. Calculations using a simple resonant mechanical circuit model show that piezoelectric actuators can generate sufficient power, force and stroke to drive the wings at 150 Hz.*

### 1 Introduction

Flapping flight for micro-robots is not only an intriguing mode of locomotion, but provides maneuverability not obtainable with fixed or rotary wing aircraft. Insects can fly with a payload equal to their body mass, and have peak accelerations approaching  $10m/s^2$  [May 1991]. Although they require relatively still air, flying micro-robots can fly over terrain which would be impassable for a legged micro-robot. Following the initial vision of Flynn [1987], pioneering work in micro-robotic flight was started by Shimoyama [Shimoyama et al 1993; Kubo et al 1994; Miki and Shimoyama 1998] and more recently, milli-robotic flapping flight by [Cox, Garcia, and Goldfarb, 1998].

This paper considers the kinematic and power requirements for a micro-robotic flying device using beating wings, and presents an initial design of a thorax for the device. As shown in Figure 1, we will be using flexural 4 bar elements to provide sufficient wing stroke, and a compliant wing which can deform to provide rotation. A summary of the MFI component dimensions is given in Table 1. Further work is

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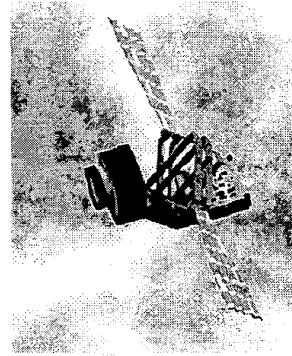


Figure 1: Conceptual drawing of micromechanical flying insect.

MFI component	size	total mass
4-bar frames (2 per wing)	links 5,5,4, 0.7 mm	20 mg
1 mm box beam base frame (mm)	10 × 4 × 1 mm	8 mg
piezo actuator (2 per frame)	0.25 × 5 × 0.2 mm (total-8 actuators)	15 mg
wings (polyester)	5 × 10 × .01 mm	3 mg
total structure		43 mg

Table 1: Components of final size MFI. Structure will be made from  $10\mu m$  thick 302 stainless.

required with models and prototypes before solving the critical issues of control, sensing, or power supply.

As a design target for the micromechanical flying insect (MFI), we are using the blowfly *Calliphora*, which has a mass of 100 mg, wing length of 11 mm, wing beat frequency of 150 Hz, and actuator power of about 8 mW. At this size scale, the current best understanding of non-steady state aerodynamics comes from experimental observations of real insects and kinematically similar mockups [Ellington et al 1996; Dickinson and Götz 1996].

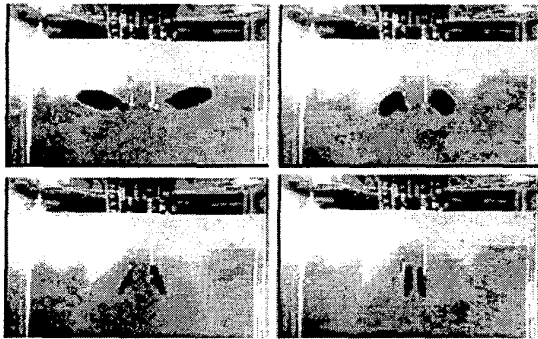


Figure 2: Robofly apparatus in mineral oil tank.

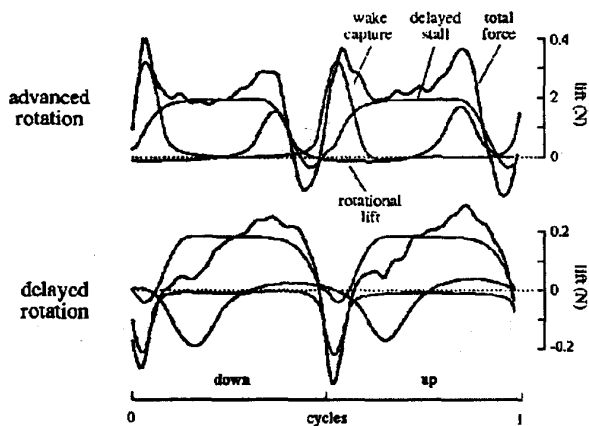


Figure 3: Experimental data from Robofly. The timing of wing rotation has a major effect on net lift. Advanced rotation, before the end of the stroke, significantly increases lift. (Adapted from Dickinson et al 1999).

## 2 Insect Unsteady Aerodynamics

The RoboFly apparatus (Figure 2) [Dickinson et al 1999], consists of a two-winged system driven by 3 stepping motors, which can closely mimic the stroke kinematics of a fruit fly, (or other arbitrary kinematics). Strain gauges are used to measure instantaneous wing forces, and the integral of forces around a closed wing beat cycle can be measured to determine net flight forces. Robofly running with wing beat of 1/6 Hz in oil has the same Reynold's number as a fruitfly with wing beat of 220 Hz in air. Flow was visualized using air bubbles in the oil tank and particle image velocimetry (PIV).

Results with the Robofly apparatus have identified the 3 key aerodynamic mechanisms used by insects: delayed stall, wake capture, and rotational circulation

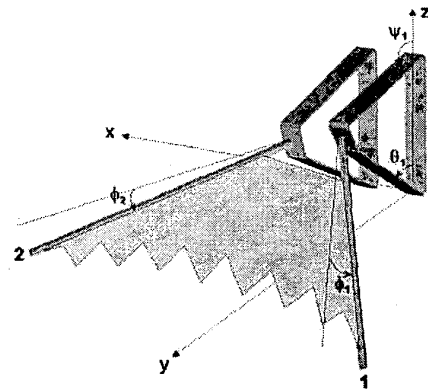


Figure 5: Wing kinematics parameters for two 4 bar frames and wing.

[Dickinson et al 1999]. Dickinson et al have found wing trajectories which generate peak lift forces of four times the equivalent insect weight. The timing of an equivalent of a back spin motion at the bottom of the wing stroke can change the net lift from positive to negative (Figure 3). The second key finding was the significant forces generated by wake capture at the top and bottom of the stroke. These results are directly relevant to the wing kinematics of the MFI, as it appears that a rapid wing rotation of 90 degrees needs to occur before the end of the down stroke to create adequate lift. This wing rotation and the resulting flow patterns can be seen in a sequence of images taken of *Drosophila* (Figure 4).

## 3 Thorax and Wing Design

We know that insect flight at the centimeter scale requires both large stroke amplitude and wing rotation [Dickinson et al 1999]. *Drosophila* uses a wing stroke of 160° combined with wing rotation of over 90°. Wing rotation is the challenging part of the design. The insect thorax uses a complicated arrangement of linkages and cams [Nachtigall, Wisser, and Eisinger, 1998] which is not yet fully understood, and is likely to be difficult to fabricate.

Our thorax design uses separate four bar frames to control the leading and trailing edges of a fanfold compliant wing, as shown in Figure 5. From simple geometric considerations, the total wing spar stroke angle can readily be shown to be  $\theta_w = 2\theta_{min}$  where  $\theta_{min}$  is the smaller base angle  $\theta_1$  when the 4 bar mechanism is in a triangle configuration. The motion magnification of the 4 bar will allow a piezoelectric unimorph with small angle deflection to be used as the actuator as described in Section 5.

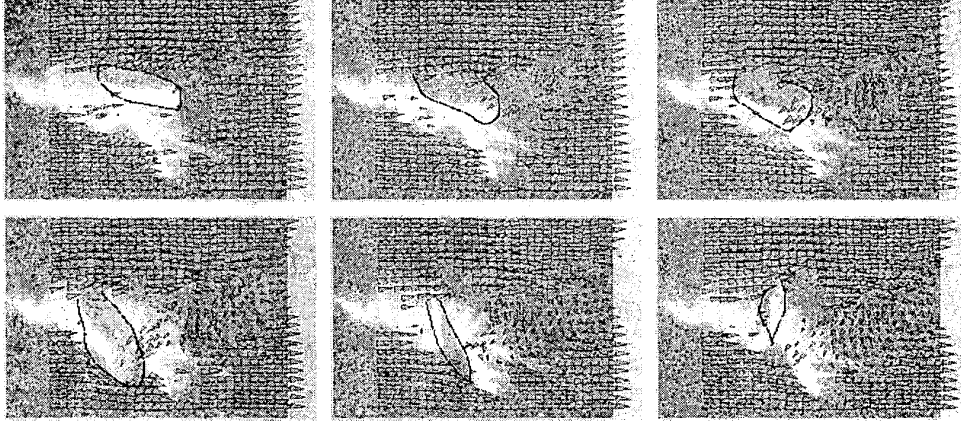


Figure 4: Flow for tethered *Drosophila*, with wing frequency 220 Hz. Images at  $200\mu\text{s}$  intervals (left to right, top to bottom) starting 1.4 ms from top of wing stroke. Note wing rotates  $\approx 90^\circ$  in 0.6 ms.

We analyze the kinematics of the thorax/wing structure to obtain desired flapping and rotational motions of the wing. The two wing spars are rotated by controlling the  $\theta_i$  crank angles ( $i = 1, 2$ ):

$$\begin{aligned}\psi_i(\theta_i) &= \tan^{-1}(B/A) \pm \cos^{-1}(C/\sqrt{A^2 + B^2}), \\ \phi_i(\theta_i) &= \tan^{-1}\left(\frac{b \sin \psi_i - a \sin \theta_i}{g + b \cos \psi_i - a \cos \theta_i}\right), \\ A &= 2ab \cos \theta_i - 2gb, \\ B &= 2abs \sin \theta_i, \\ C &= g^2 + b^2 + a^2 - h^2 - 2ag \cos \theta_i, \quad (1)\end{aligned}$$

where  $a, b, g, h$  are the 4 link lengths. A solution exists if and only if  $A^2 + B^2 - C^2 \geq 0$ .

### 3.1 Thorax and Wing Construction

Thorax construction materials should be chosen for high stress limits, high endurance, and low loss (Table 2). The insect thorax is constructed from cuticle with resilin connections [Jensen and Weis-Fogh, 1962], which compares very well in stress limits with stainless steel. Polysilicon is a good structural material, but is harder to handle and process than stainless. Note that peak stress must be limited to maintain life. This can be done with thin materials, and by keeping the overall system  $Q$  low since high  $Q$  increases dynamic stresses. While the overall  $Q$  of the actuator-wing system is predicted to be only  $\approx 2$  (Section 5.2) due to wing and actuator dissipation, for reasonable efficiency, the internal dissipation of the structural material [Hosaka, Ito, and Kuroda, 1994] should be kept low.

The four bar frames need to have a very high strength-to-weight ratio. As pointed out by Yeh et

al [1994], hollow beams can be several orders of magnitude stiffer than a solid beam of the same mass, and we use the same principle. Solid rectangular and triangular beams of width  $w$  and height  $h$  have respective cross-sectional moments of inertia [13]:

$$I_{rect} = wh^3/12 \quad \text{and} \quad I_{tri} = wh^3/36. \quad (2)$$

The stiffness of a cantilever beam of length  $l$  and modulus  $E$  in simple bending due to a concentrated force at the end is given by

$$k = 3EI/l^3. \quad (3)$$

A hollow beam has cross-sectional moment of inertia given by the difference in moment of inertia between a solid beam and a beam smaller by the wall thickness. Consider the links of the 4 bar structure shown in Figure 5, with a mass of  $0.24 \text{ mg/mm}$  length. For the same mass, we could use either a  $30\mu\text{m}$  thick cantilever, 1 mm wide, with cross-sectional moment of inertia  $2.3 \times 10^{-18} \text{ m}^4$ , or a 1 mm triangular beam constructed by folding a  $10 \mu\text{m}$  thick sheet, with cross-sectional moment of inertia  $700 \times 10^{-18} \text{ m}^4$ . Hence, the hollow triangular beam is approximately 300 times stiffer than the solid beam for the same mass. An 8 mm long box beam made from 302 stainless would have a stiffness of  $7000 \text{ N/m}$ . A peak load on the structure from the actuators of 100 mN (see Section 5) would cause a deflection of less than  $130\mu\text{m}$ , thus this beam is stiff enough.

Since pin joints are difficult to place in a 1 mm structure without friction, wear, reinforcement, and alignment problems, we use flexural joints. As shown

property	music wire	poly silicon	302 stainless	polyester	cuticle	resilin
Elastic Modulus $E Nm^{-2}$	$200 \times 10^9$	$140 \times 10^9$	$181 \times 10^9$	$3.6 \times 10^9$	$10 \times 10^9$	$2 \times 10^6$
yield stress $Nm^{-2}$	$4 \times 10^9$	$0.7 \times 10^9$	$1 \times 10^9$	$55 \times 10^6$	$0.4 \times 10^9$	
ultimate extension	2%	.5%	0.5%	124%	2-3%	300%
endurance limit $Nm^{-2}$ ( $> 10^6$ cycles)	$10^9$ 0.5%	$0.7 \times 10^9$ 0.5%	$0.5 \times 10^9$ 0.3%	?	?	?
$Q = \frac{\pi}{\text{lossfactor}}$	200	$10^4$	?	?	30	$> 100$
density $mg/mm^3$	7.9	2.3	7.9	1.4	1.3	?
wave velocity $\sqrt{\frac{E}{\rho}} (ms^{-1})$	5000	7800	5000	1600	2800	?
reference	Oberg 88	Keller 98	matls.com	matls.com	Jensen 62	Jensen 62

Table 2: Material properties for thorax and wing.

in Figure 6, the 4 bar mechanism including box beam links and flexural joints can be fabricated by folding a sheet of photo-etched or laser cut stainless steel. The side and top links are triangular beams and the base is a box beam for easier attachment to the MFI frame.

We have used glass fiber reinforcing at the joints to increase lifetime with large bend angles of the  $5\times$  model, but a better long term approach is to use the “split tube flexure” of Goldfarb and Speich [1998] to obtain a more precise center of rotation for the joint and improved lifetime. Currently, the structures are folded manually using fixtures and bonded using cyanoacrylate adhesive. Automatic folding of microstructures is quite feasible, particularly using simple fixtures and a motion planning approach such as described by Lu and Akella [1999].

The wing for the  $5\times$  mockup is made from folded  $25\mu m$  thick polyester bonded to plastic wing spars as shown in Figure 7. The increased drag caused by the fan-folds is not thought to be significant at the low Reynolds number flow of the MFI.

#### 4 Preliminary Flow Results with $5\times$ Mockup

The  $5x$  wing mockup can be driven statically with a stroke of  $\pm 60^\circ$  or a rotation of  $\pm 45^\circ$ . Figure 8 shows PIV images of this wing being driven at  $f = 17Hz$  in a wind tunnel operating with a flow of  $0.14 m/s$  (to the right). Ellington [1984] approximates the mean Reynolds number during a wing stroke as

$$Re = \frac{Sf\Phi}{\nu} \quad (4)$$

where  $\nu = 1.6 \times 10^{-4} m^2/s$  is the kinematic viscosity of air,  $S = 32cm^2$  is the wing surface area for both wings, and  $\Phi = 1rad$  is the stroke amplitude. In this case,  $Re = 340$  which is comparable to the blowfly. Processing of the PIV images shows that induced velocities

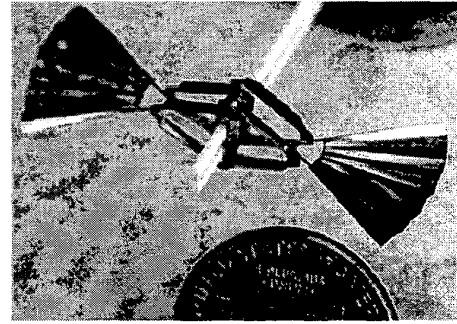
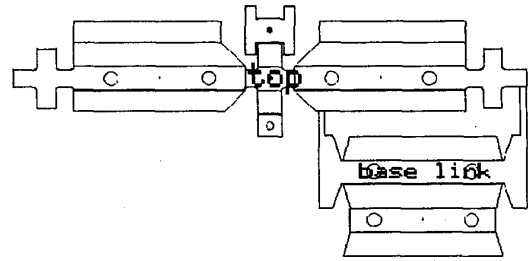


Figure 6: Flat pattern for one thorax 4 bar flexural frame, and assembled thorax at  $1.5\times$  scale.

as high as  $1.65m/s$  and  $0.79m/s$  were recorded during the upstroke and downstroke, respectively, with peak wing-tip speeds of  $2 m/s$ .

Visually, there are two interesting results that can be seen from the PIV data. Firstly, there is a significant momentum change in the particles directly under the wing from the end of the upstroke to the beginning of the downstroke and this would indicate a period of large lift forces. Secondly, vortex shedding at the trailing edge of the wing is readily apparent as the wing finishes the downstroke and begins the upstroke.

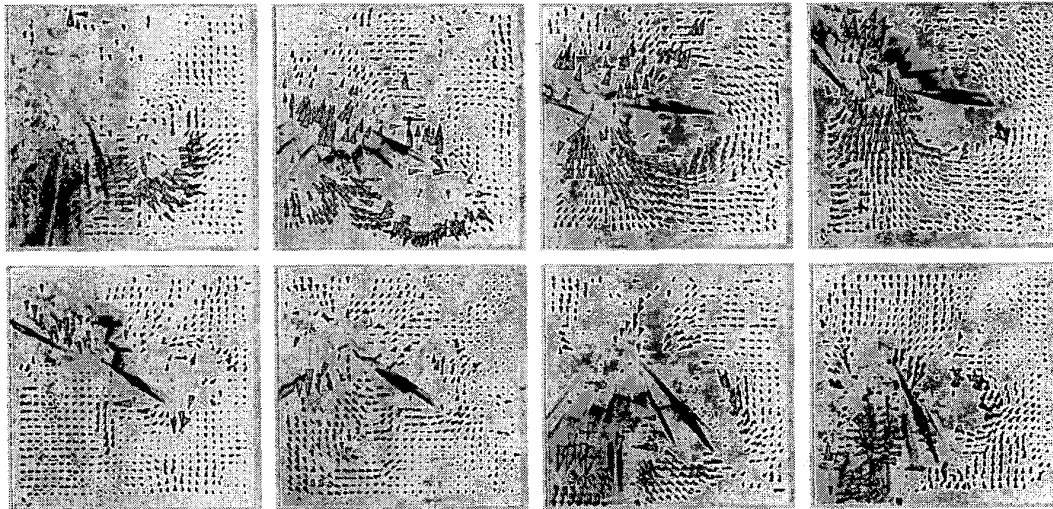


Figure 8: Flow structures measured using PIV for 50 mm wing beating at 17 Hz, showing complete wing stroke starting at bottom of stroke. Dark line is laser light sheet (negative image).

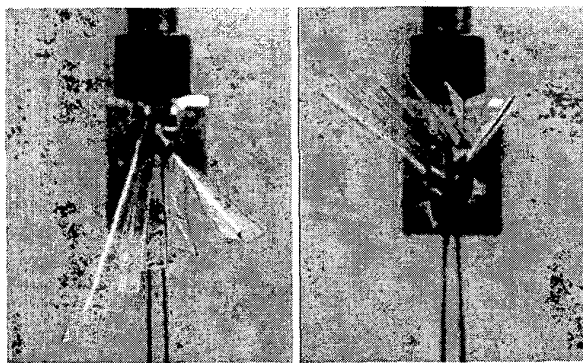


Figure 7: Static wing motion using voice coil actuator on 5× scale prototype.

## 5 Actuation, Transmission, and Power

The previous sections have discussed the design and fabrication of a prototype thorax for the MFI. The next critical issue is the delivery of sufficient mechanical power to the wings. Insect flight muscle delivers on the order of  $100Wkg^{-1}$  to the wings [Ellington 1985]. To avoid losses associated with accelerating inertial loads, real insects such as blowflies and fruitflies drive their wings at mechanical resonance [Sotavalta 1952]. We use the same principle in our thorax design.

The exact MFI power requirement will depend on the non-linear kinematics of the two 4 bar structures and fan-fold wing, actuator compliance, and non-steady state aerodynamic loading and can not be

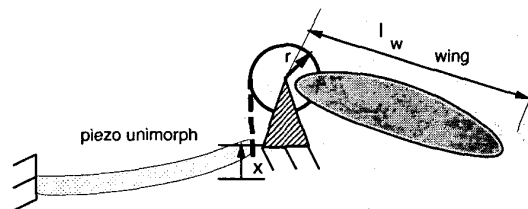


Figure 9: Mechanical schematic of actuator and wing connection. Virtual pulley represents thorax kinematics and motion amplification.

exactly modelled with current flow simulation tools. We thus examine actuating a single equivalent linear damping load through a simple pulley, which provides a rough bound on the power required to drive two wings generating sufficient lifting force. We consider a piezoelectric unimorph actuator driving a wing mounted on a pulley of radius  $r$  as shown in Figure 9. The equivalent wing damping is chosen such that the force at the end of the wing at maximum wing tip velocity is equal to twice the weight of the MFI:

$$B_L = 2 \frac{m_{MFI} g l_w}{r^2 \omega \theta} \quad (5)$$

where  $m_{MFI}$  is the mass,  $l_w$  is the length of the wing,  $\omega$  is the wing beat frequency, and  $\theta$  is the wing stroke amplitude (1/2 the total wing stroke).

For simplicity, we assume the inertia of the wing can be modelled as a point mass ( $m_w$ ) at  $r_w = 5.0mm$

parameter	value
wing pulley radius $r$	$0.5 \times 10^{-3} m$
wing point mass $m_w$	$0.5 \times 10^{-6} < m_w < 2.5 \times 10^{-6}$
wing inertia	$12 \times 10^{-12} < J < 60 \times 10^{-12} kgm^2$
wing beat $\omega$	$2\pi 150 sec^{-1}$
wing amplitude	$\theta = 1$ radian
load force $F_L$	$2.9 \times 10^{-2} N$
load damping $B_L$	$6.4 \times 10^{-2} Nsm^{-1}$
load mass $M_L$	$0.5 \times 10^{-4} < M_L < 2.5 \times 10^{-4} kg$
load stiffness $K_L$	chosen for resonance
required power	$7.0 \times 10^{-3} W$

Table 3: Model assumptions for MFI using linear damper model.

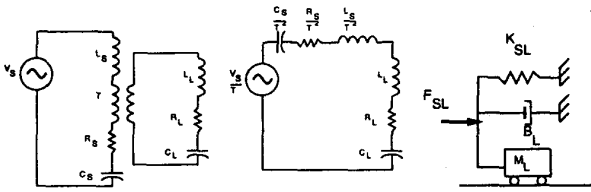


Figure 10: Electrical and mechanical equivalent resonant models for actuator, thorax, and wing.

from the joint. The thorax inertia may also be lumped into this parameter. Similarly to wing damping, wing inertia is mapped into a linear inertia  $M_L$  through a "wing pulley" of radius  $r$ .

$$M_L = \frac{m_w r_w^2}{r^2} \quad (6)$$

The load stiffness  $K_L$  can be chosen to tune the system to resonance. The resulting load parameters are summarized in Table 3.

### 5.1 Transmission

For many actuators, there is a large mismatch between source and load impedances. In addition, the large stiffness of the actuator would require too large a mass for resonance to be carried by the MFI. The leverage provided by the wing pulley (or equivalently the 4 bar thorax) can be modelled as a lossless transformer as shown in Figure 10. The actuator can be modelled as a voltage (force) source with equivalent damping ( $R_S$ ), stiffness ( $C_S^{-1}$ ), and inertia ( $L_S$ ). The thorax is modeled as a transformer with coupling ratio  $T$ , and the wing-hinge combination (the load) has damping ( $R_L$ ), stiffness ( $C_L^{-1}$ ), and inertia ( $L_L$ ). Typically, for unmatched piezo and load impedances, the stiffness of the actuator is much greater than the load stiffness, and the actuator damping is much less than the equivalent wing damping. Hence the efficiency is

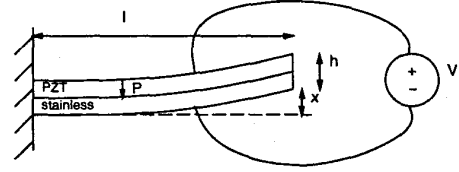


Figure 11: Piezoelectric unimorph actuator.

close to 100% but the power density is quite low compared to matched conditions.

The output force from the lever  $F_{SL}$  scales proportionally with input force  $F_S$ , while output position  $x_{SL}$  scales inversely with input position  $x_S$ . Hence, with transformer ratio  $T$ :  $K_{SL} = K_S/T^2$ ,  $B_{SL} = B_S/T^2$ ,  $M_{SL} = M_S/T^2$ , and  $F_{SL} = F_S/T$ . The attachment point of the actuator to the 4 bar structure, and the 4 bar link lengths determine  $T$ .

For resonance, the equivalent mass of the actuator  $M_{SL}$  combined with the wing inertia  $M_L$  must be chosen to have the same magnitude reactance at the resonant frequency  $\omega$  as  $K_{SL}$ . The equivalent mass of the piezo actuator is negligible compared to the wing inertia. Thus the maximum transformed actuator stiffness is  $K_{SLmax} = \omega^2 M_{Lmax}$ , where  $M_{Lmax}$  is the maximum translational equivalent wing inertia. With the assumed wing pulley radius  $r = 0.5mm$ , the maximum actuator stiffness is  $230Nm^{-1}$  at 150 Hz.

Using a figure of 7 mW for actuator power, basic AC circuit analysis shows the minimum actuator force  $F_{SL}$  for matched loads ( $B_{SL} \approx B_L$ ) and unmatched loads ( $B_{SL} \ll B_L$ ) is respectively found to be

$$F_{SLmin} = \sqrt{8B_L P_L} \text{ and } F_{SLmin} = \sqrt{2B_L P_L} \quad (7)$$

A low internal damping actuator needs only half as much force as a matched damping actuator and will run at higher efficiency, while requiring greater actuator mass. A piezoelectric actuator with low internal damping will need to provide 29 mN at 150 Hz to provide sufficient power.

### 5.2 Piezoelectric Actuation

Piezoelectric materials in  $d_{31}$  mode have too small a displacement and too high a stiffness to work well at wing beat frequencies of 150 Hz. A matching network could be used, but the high  $Q$  of the actuator system causes excessive stress in the piezo. A piezoelectric slab bonded to an elastic layer, forms a *uni-morph* which bends when an electric field is applied. A unimorph configuration as shown in Figure 11 gives a blocked force (following the bimorph analysis of Smits

et al [1991]):

$$F_S \approx \frac{3}{16} Y_{11} d_{31} \frac{hw}{l} E_3 h \quad (8)$$

where  $Y_{11}$  is stiffness,  $E_3$  is electric field,  $w$  is width, and  $d_{31}$  is the piezoelectric constant in  $mV^{-1}$ . The free displacement of the unimorph cantilever is given by [Smits et al 1991]:

$$x_S = \frac{3}{4} d_{31} \frac{l^2}{h^2} E_3 h. \quad (9)$$

Thus the actuator stiffness is just  $K_S = F_S/x_S$ .

Alternatively, a bimorph structure consisting of two layers of piezoelectric material would give twice the displacement and twice the force. However, the power limit for the actuator is typically determined by the dynamic stress limit [Berlincourt et al 1964]. A pre-stressed unimorph, such as the Face Int. THUNDER, can be safely operated with almost 10 times greater fields than a bimorph, giving a 25 times greater power density.

The resonant frequency of a homogeneous cantilever with density  $\rho$  is

$$\omega_o = \frac{h}{l^2} \sqrt{\frac{Y_{11}}{\rho}}. \quad (10)$$

Note that the cantilever resonant frequency is quite large compared to the wing resonance. Using eq. 10 and the relation between  $Q$  and damping for a series resonant circuit, the unimorph damping can be estimated as:

$$B_S = \frac{K_S}{Q\omega_o}, \quad (11)$$

where the elastic layer damping has been assumed small compared to the piezo damping. It is important to note that the quoted mechanical  $Q$  for piezoelectric materials is usually at low stress, as in a sensor application. At high stress the mechanical  $Q$  is about 10% of the unstressed  $Q$ . The MFI actuator will be running at high stress.

Using the resonant circuit model of Figure 10, we analyzed power requirements and wing stroke amplitude assuming a piezoelectric unimorph (Figure 9). We obtained design parameters as given in Table 4, using an actuator  $5 \times 0.2 \times 2mm^3$  in size. A 15 mg mass of PZT5H can generate the 7 mW of mechanical power required at a frequency of 150 Hz, using a driving voltage of 200 volts. Note that this is the total actuator volume required for both wings. Calculations show the mechanical efficiency is better than 95%. (Electrical efficiency can be readily calculated

parameter	value
frequency $\omega$	$2\pi 150 sec^{-1}$
high $E_3$ mechanical $Q$	$\approx 10$
material stiffness $Y_{11}$	$8.2 \times 10^{10} Nm^{-2}$
piezoelectric constant $d_{31}$	$-274 \times 10^{-12} m/V$
ideal actuator power density	$1.5 kW kg^{-1}$
actuator dimensions $lwh$	$5 \times .2 \times 2 \times 10^{-9} m^3$
actuator mass	$15 \times 10^{-6} kg$
applied field $E_3$	$2 \times 10^6 Vm^{-1}$
blocked force $F_S$	0.135 N
free displacement $x_{DC}$	$5.1 \times 10^{-5} m$
resonant freq. $\omega_o$	$4.4 \times 10^4 sec^{-1}$
$K_S = F_S/x_{DC}$	$2600 Nm^{-1}$
$M_S = K_S/\omega_o^2$	$1.4 \times 10^{-4} kg$
$B_S$	$.134 Nsm^{-1}$
transmission ratio $T$	4.5
$F_{SL}$	30 mN
$K_{SL}$	$130 Nm^{-1}$
$B_{SL}$	$2.9 \times 10^{-3} Nsm^{-1}$
$M_{SL}$	$6.7 \times 10^{-6} kg$
required $M_L$	$1.4 \times 10^{-4} kg$
wing stroke at 150 Hz	$110^\circ$
wing power at 150 Hz	6.5 mW
overall $Q$	2.1
efficiency $\eta$	95.5%
actuator power density	$400 W kg^{-1}$
peak beam stress	$1.2 \times 10^8 Nm^{-2}$
peak beam strain (compressive)	0.15%
peak beam deflection $x_S$	$1 \times 10^{-4} m$

Table 4: Model assumptions for PZT5 actuation of both wings using a pre-stressed unimorph.

considering the power supply source impedance as described in Stein et al [1994].) This actuator is limited by its strength to peak tensile strains less than 0.1%. The wing power and amplitude for this model are shown in Figure 12. Our analysis of alternative piezo materials such as single-crystal PMNPT shows even better performance, but single-crystal PMNPT is currently difficult to obtain.

## 6 Conclusion

At this stage in the MFI design, we have a good understanding of the necessary wing kinematics, forces, velocities and power from measurements on Roboffly and real insects. We have designed kinematic structures using folded stainless steel which give adequate wing motion when driven by external (high power) actuators. Using appropriate fixturing, folding stainless steel is feasible at the final size. Current flexure life time for the  $5 \times$  structure is on the order of  $10^4$  cycles; using improved flexural elements, we should be able to

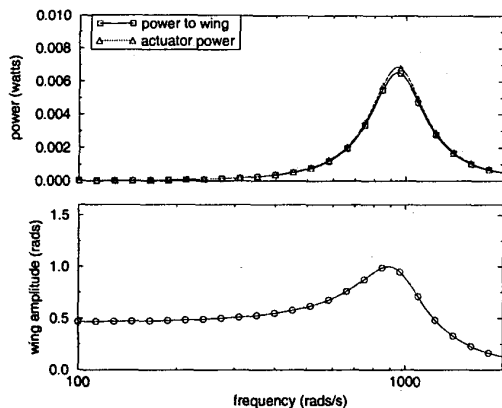


Figure 12: Wing power and wing amplitude for PZT-5H piezo unimorph. Predicted performance is adequate for flight.

achieve our goal of 1 million wing beats before failure.

In the next stage of the project, we will be quantifying flight forces on the 5× scale MFI mockup using force sensors mounted on the wing spars, in combination with PIV. Our plan is to develop closed-loop wing controllers which can react to wing forces and modify wing stroke patterns as needed to achieve stable flight.

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