

# POLAMI: Polarimetric Monitoring of AGN at Millimetre Wavelengths – I. The programme, calibration and calibrator data products

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## ABSTRACT

We describe the POLAMI (Polarimetric Monitoring of AGN at Millimetre Wavelengths) programme for the monitoring of all four Stokes parameters of a sample of bright radio-loud active galactic nuclei with the IRAM 30-m telescope at 3.5 and 1.3 mm. The programme started in 2006 October and accumulated, until 2014 August, 2300 observations at 3.5 mm, achieving a median time sampling interval of 22 d for the sample of 37 sources. This first paper explains the source selection, mostly blazars, the observing strategy and data calibration and gives the details of the instrumental polarization corrections. The sensitivity ( $1\sigma$ ) reached at 3.5 mm is 0.5 per cent (linear polarization degree),  $4^{\circ}.7$  (polarization angle), and 0.23 per cent (circular polarization), while the corresponding values at 1.3 mm are 1.7 per cent,  $9^{\circ}.9$  and 0.72 per cent, respectively. The data quality is demonstrated by the time sequences of our calibrators Mars and Uranus. For the quasar 3C 286, widely used as a linear polarization calibrator, we give improved estimates of its linear polarization, and show for the first time occasional detections of its weak circular polarization, which suggests a small level of variability of the source at millimeter wavelengths.

**Key words:** polarization – surveys – Galaxies: active – BL Lacertae objects: general – galaxies: jets – quasars: general.

## 1 INTRODUCTION

Radio-loud active galactic nuclei (AGN) are powerful emitters of radiation. Their relativistic jets radiate at all spectral bands from the longer radio wavelengths to the very high-energy  $\gamma$ -rays. This emission shows extremely fast and high-amplitude total flux density and polarization variability throughout the spectrum (e.g. Marscher et al. 2008, 2010; Abdo et al. 2010; Agudo et al. 2011b; Aleksić et al. 2015; Casadio et al. 2015).

Even after several decades of study since their discovery there are still relevant questions about some of the most fundamental physics driving these objects, especially in the innermost and most poorly explored regions. These questions include: Which models of jet formation, acceleration and collimation, compression of the plasma and magnetic field describe these systems best? What is

the mechanism at work for the high-energy emission? What is the composition of the plasma forming the relativistic jets? What is the ultimate origin of the variability in the putative accretion and jet ejection system? The linear and circular polarization spectra, and the strong polarization variability with doubling time-scales from hours to months (Saito et al. 2013; Hayashida et al. 2015; Marscher 2016) usually assumed to be produced by the dynamic evolution of inhomogeneities in the plasma forming the jets, are supposed to be intimately related to the answers to these questions.

Total flux density and polarimetric monitoring can therefore help shedding new light on these problems, especially in the short millimetre spectral band where AGN jets were not intensively explored before, and where they are believed to be optically thin (hence free of opacity effects). At radio wavelengths, AGN jets also tend to be affected by Faraday rotation. For rotation measures up to  $|RM| \approx 10^4 \text{ rad m}^{-2}$  sometimes encountered in AGN jets at centimetre wavelengths (Zavala & Taylor 2004; Hovatta et al. 2012), the observed polarization angle  $\chi_{\text{obs}} = \chi_{\text{int}} + RM\lambda^2$  can be rotated

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away from the intrinsic polarization angle  $\chi_{\text{int}}$  by several tens, even hundreds, of degrees. However, such rotation is usually negligible in the short millimetre range. This also makes millimetre observations less sensitive to Faraday de-polarization, hence reflecting better the intrinsic linear polarization properties, and therefore those of the magnetic field, of the system.

Previous programmes that focused on the millimetre emission of large samples of AGN have been dedicated to studies of the time dependence of the total flux density and linear polarization on typical time-scales much longer than a month. Little has been published on the study of circular polarization, or the evolution of relatively big AGN samples on shorter time-scales ( $\sim 2$  weeks). The only previous attempt to study the variability of the full-polarimetric (i.e. four Stokes parameters) properties of a large sample of AGN at millimetre wavelengths was presented in Agudo et al. (2014), where we investigated the variations of the 3.5 mm total flux density, linear and circular polarization parameters of a set of  $\sim 100$  AGN with the IRAM 30-m telescope in 2010 and in 2005 (see Agudo et al. 2010, 2014). The results from this variability study showed prominent total flux density and polarization changes in time-scales of a few years. Such changes reached median factors of  $\sim 1.5$  in total flux density and  $\sim 1.7$  in linear polarization degree (with maximum variations by factors up to 6.3 and  $\sim 5$ , respectively). We also noted drastic variations of the polarization angle, with random changes in 86 per cent of the sources with regard to the first survey in 2005.

In light of these large changes in most sources between our two surveys, we embarked on a higher cadence polarization monitoring programme of a selected group of AGN. The programme, dubbed POLAMI (Polarimetric Monitoring of AGN at Millimetre Wavelengths, see <http://polami.iaa.es>), started in 2006, and it measures all four Stokes parameters observed with the IRAM 30-m telescope at both 3.5 and 1.3 mm wavelengths. In this first paper of a series presenting the results of the first 8 yr of POLAMI observations, we provide detailed information about the observing programme (Section 2), the monitored source sample (Section 3), the data reduction and calibration (Section 4), and we demonstrate the quality of our data by showing the results obtained for our main calibrators (Section 5). The data obtained from our science targets, as well as an analysis of their circular polarization properties and of their total flux density and linear polarization variability are shown in the accompanying Paper II (Thum et al. 2018) and Paper III (Agudo et al. 2018), respectively. Detailed studies of specific sources and discussions of statistical aspects of our sample, like correlations with optical,  $\gamma$ -ray, and millimetre VLBI (very long baseline interferometry) data will be presented in separate publications, currently in preparation.

## 2 THE PROGRAMME

The POLAMI programme for the study of the polarimetric properties of AGN at short millimetre wavelengths evolved from technical work which monitored the calibration of the flux density scale and the polarization characteristics of the IRAM 30-m telescope. Observations of bright and point-like sources were recorded with complete polarization information, and it quickly became clear that this data set constituted a unique resource for AGN studies at short mm wavelengths.

The variability of AGN across the electromagnetic spectrum requires studies that monitor large samples over long time baselines. Unprecedented time sampling of the  $\gamma$ -ray and X-ray emission became available from space observatories, as well as exhaustive optical polarimetric monitoring and millimetre VLBI monitoring

of large sets of AGN. The combination of such comprehensive multispectral-range data sets has been crucial in understanding the innermost, and most poorly explored regions of jets in AGN, where the most energetic processes take place, (e.g. Marscher et al. 2010; Jorstad et al. 2010, 2013; Agudo et al. 2011a,b; Casadio et al. 2015). The POLAMI programme fits very well into this broad-band effort. In fact, it bridges a gap at short millimetre wavelengths where no systematic full-polarization AGN monitoring programmes exist.

The first POLAMI observing session took place on 2006 October 14. In this first publication, we present the data compiled until 2014 August 18, but the programme is still active and more recent data will be analysed on already planned publications. Observations were scheduled in 4–6-h sessions scattered very irregularly over this  $\sim 8$  yr period. The median time sampling of our programme  $t_S$  is 22 d, when all monitored AGN targets and all observations are considered over the monitoring period.  $t_S$  improved with time, dropping from  $\sim 42$  d during the first 4 yr to  $\sim 19$  d towards the end. Observations were interrupted during a brief period in 2009 from February to mid-June when a new generation of receivers (Section 4.1) was installed and commissioned.

All POLAMI observations are made at sky frequencies of 86.243 GHz (3.5 mm) and at 228.932 GHz (1.3 mm). At both frequencies the performance of the receivers is good (see Section 4.1), and the atmosphere is free of major absorption lines. Data are recorded throughout our programme using the XPOL procedure (Thum et al. 2008). All four Stokes parameters are obtained strictly simultaneously. In December 2009, the 3.5 and 1.3 mm bands (hereafter 3 and 1 mm, respectively) started to be observed in parallel, and XPOL then simultaneously recorded all Stokes parameters at both frequencies.

In a typical observing session, 15–20 sources are observed. Each polarimetric measurement of a target source is preceded by a pointing observation made on the same source. Integration times of the polarimetric measurement range from 4 min (sources of 10 Jy or brighter) to 15 min duration for the weaker sources in our sample (see Section 3) with  $\sim 1$  Jy. Whenever possible, observations of calibrators were included in each session. Mars and Uranus are used for the instrumental polarization (IP) and total flux density calibration, while 3C 286 and the Crab nebula allowed for cross-checks of the polarization calibration. IP parameters and total power calibration were sufficiently stable to allow interpolation for the few sessions without a calibrator measurement (see Section 4).

Until mid-August 2014, we accumulated 2700 polarimetric observations at 3 mm. After necessary data quality checks, we retain 2300 observations. All data, including pointing scans preceding every polarization integration, are compiled in a dedicated data base.<sup>1</sup> For each observation, the data base stores the values of the four Stokes parameters and their uncertainties, information about observing conditions, and a quality flag, for each frequency band. Polarization data are kept both in their original form before rotation to the celestial reference frame and before any IP is removed, as well as in their final form. Successive refinements of data analysis are then feasible without loss of information.

## 3 THE SOURCE SAMPLE

The entire POLAMI source sample is listed in Table 1 together with the most relevant properties of every source. The sample is designed

<sup>1</sup> The POLAMI data base is installed at IRAM–Granada’s computer system and is maintained by one of the authors (CT) supported by the Observatory’s computer group.

**Table 1.** List of sources in our POLAMI first priority sample and their most relevant properties.

Source name (1)	Common name (2)	Position reference (3)	Optical class (4)	SED class (5)	Redshift (6)	(S) (Jy)		N <sub>obs</sub>	
						3 MM (7)	1 MM (8)	3 MM (9)	1 MM (10)
0219+428	3C 66A	<i>a</i>	B	HSP	0.444	0.54	0.49	31	7
0235+164	AO 0235+164	<i>a</i>	B	LSP	0.94	1.61	0.91	50	15
0316+413	3C 84	<i>b</i>	G	LSP	0.0176	18.19	8.88	59	29
0336-019	CTA 26	<i>a</i>	Q	LSP	0.852	2.25	1.31	37	19
0355+508	NRAO 150	<i>b</i>	Q <sup>e</sup>	LSP	1.517 <sup>e</sup>	4.62	1.91	76	35
0415+379	3C 111	<i>d</i>	G	LSP	0.0491	3.70	2.44	91	47
0420-014	PKS 0420-01	<i>a</i>	Q	LSP	0.9161	4.21	2.28	46	29
0430+052	3C 120	<i>a</i>	G	LSP	0.033	2.21	1.37	26	14
0528+134	PKS 0528+134	<i>a</i>	Q	LSP	2.07	2.76	1.60	68	26
0716+714	S5 0716+71	<i>a</i>	B	ISP	0.127	3.78	2.87	108	64
0735+178	OI 158	<i>a</i>	B	LSP	0.45	0.87	0.61	52	31
0827+243	OJ 248	<i>a</i>	Q	LSP	0.942	1.52	1.93	73	42
0829+046	OJ 049	<i>a</i>	B	LSP	0.174	0.70	0.53	52	30
0836+710	4C 71.07	<i>a</i>	Q	LSP	2.218	1.48	0.65	91	48
0851+202	OJ 287	<i>a</i>	B	LSP	0.306	5.52	3.63	92	57
0954+658	S4 0954+65	<i>a</i>	B	LSP	≥0.45 <sup>f</sup>	1.36	0.97	79	42
1055+018	4C +01.28	<i>a</i>	Q	LSP	0.888	4.10	2.30	61	38
1101+384	Mrk 421	<i>d</i>	B	HSP	0.0308	0.50	0.32	25	18
1127-145	PKS 1127-14	<i>a</i>	Q	LSP	1.187	1.64	0.72	56	28
1156+295	4C 29.45	<i>a</i>	Q	LSP	0.725	1.37	0.82	59	32
1219+285	W Comae	<i>a</i>	B	ISP	0.103	0.44	0.34	33	17
1222+216	4C +21.35	<i>b</i>	Q	LSP	0.434	1.65	1.01	60	45
1226+023	3C 273	<i>a</i>	Q	LSP	0.1583	11.61	4.08	83	53
1253-055	3C 279	<i>a</i>	Q	LSP	0.536	20.71	12.35	77	37
1308+326	OP 313	<i>a</i>	Q	LSP	0.997	1.66	0.97	54	35
1328+307	3C 286	<i>c</i>	Q	LSP	0.846	0.90	0.31	57	26
1406-076	PKS 1406-076	<i>a</i>	Q	LSP	1.494	0.75	0.36	32	14
1510-089	PKS 1510-08	<i>a</i>	Q	LSP	0.36	2.75	1.89	53	27
1611+343	DA 406	<i>a</i>	Q	LSP	1.401	2.18	1.13	69	43
1633+382	4C 38.41	<i>a</i>	Q	LSP	1.813	4.11	2.85	76	48
1641+399	3C 345	<i>a</i>	Q	LSP	0.593	3.65	1.76	80	52
1730-130	NRAO 530	<i>a</i>	Q	LSP	0.902	3.13	1.50	48	26
1749+096	4C +09.57	<i>a</i>	B	LSP	0.322	3.22	1.70	62	44
2200+420	BL Lacertae	<i>a</i>	B	LSP	0.0686	6.98	5.59	70	44
2223-052	3C 446	<i>a</i>	Q	LSP	1.404	4.58	1.21	60	29
2230+114	CTA 102	<i>a</i>	Q	LSP	1.037	3.43	1.34	67	30
2251+158	3C 454.3	<i>a</i>	Q	LSP	0.859	24.20	24.22	80	39

*Notes.* Columns are as follows: (1) IAU B1950 source name, (2) common (alternative) source name, (3) reference from where position was taken [see below references *a* to *d*], (4) optical classification into quasars (Q), BL Lac objects (B) and radio galaxies (G) after Véron-Cetty & Véron (2006), (5) classification of sources based on the broad-band spectral-energy distribution (low, intermediate, and high synchrotron peaked are LSP, ISP and HSP, respectively) following Ackermann et al. (2011, 2015) and the information provided in the MOJAVE data base Lister et al. (2009) (<http://www.physics.purdue.edu/astro/MOJAVE/>), (6) redshift from the MOJAVE data base, (7) average total flux density measured at 3 mm, (8) average total flux density measured at 1 mm, (9) number of valid observations at 3 mm and (10) number of valid observations at 1 mm. The number of 1 mm measurements is significantly smaller than the one at 3 mm. This is mainly because of the shorter time span of 1 mm monitoring observations (that started in 2009 December, see the text), but also because of the larger fraction of rejected data at 1 mm. At this wavelength, atmospheric instabilities are more pronounced, and sources are often considerably weaker than at 3 mm.

<sup>a</sup>Johnston et al. (1995).

<sup>b</sup>Beasley et al. (2002).

<sup>c</sup>Ma et al. (1998).

<sup>d</sup>Fey et al. (2004).

<sup>e</sup>From Agudo et al. (2007) and Acosta-Pulido et al. (2010).

<sup>f</sup>From Landoni et al. (2015).

as the intersection of two different larger samples. The first of them, i.e. the IRAM 30-m telescope's Pointing Source Catalog (Agudo et al. 2010), is a set of bright millimetre AGN with declination  $\gtrsim -25^\circ$ , which are easily visible to the 30-m telescope. The second one, is the source sample of the VLBA-BU-BLAZAR Monitoring

Program<sup>2</sup> (Jorstad & Marscher 2016; Jorstad et al. 2017) which compiles sequences of monthly 7 mm polarimetric VLBI images

<sup>2</sup> <http://www.bu.edu/blazars/VLBAproject.html>

since 2007. The intersection of the two sets includes 35 of the 37 sources currently monitored with VLBI by the VLBA–BU–BLAZAR Program, which happen to be among the brightest  $\gamma$ -ray AGN in the northern sky, and are therefore also intensively monitored by the *Fermi* Gamma Ray Space Observatory, the *Swift* X-ray satellite and a number of optical (Agudo et al. 2012a; Angelakis et al. 2016; Larionov et al. 2016; Itoh et al. 2016) and radio (Aller, Aller & Hughes 2003; Lister et al. 2009; Teräsranta et al. 2004; Max–Moerbeck et al. 2010) total flux density and polarimetric monitoring programmes. This helps maximizing the scientific output of the POLAMI programme. To these 35 sources, we added two more bright millimetre AGN, namely NRAO 150 and 3C 286, that we considered of interest to our astrophysical and astronomical objectives. Although NRAO 150 is close to the Galactic plane, and therefore its high Galactic absorption prevented its identification and classification until recently (Acosta-Pulido et al. 2010), this source is one of the brightest blazars in the northern sky (Agudo et al. 2007) and shows interesting polarimetric behaviour on its sub-milliarcsecond structure as revealed by 7 mm and 3 mm VLBI monitoring observations (Agudo et al. 2007; Molina et al. 2014). The quasar 3C 286 is a standard total flux density and linear polarization calibrator both at centimetre (Baars et al. 1977; Ott et al. 1994; Perley 1982) and millimetre wavelengths (Agudo et al. 2012b), and has therefore been an excellent control source for our programme.

All observations were subject of a detailed quality control as described in Section 4.4. The median number of accepted valid measurements per source at 3 mm is 61. The source with most measurements (108) is the BL Lac object 0716+714, and the source with the least number of observations (25) is 1101+384. The POLAMI source sample comprises 24 radio quasars, 11 BL Lacertae type objects<sup>3</sup> and 3 radio galaxies. Among the 11 BL Lac objects, most of them (7) are low-energy synchrotron peaked (LSP), 2 of them are intermediate-energy synchrotron peaked (ISP) and the 2 remaining ones have their synchrotron bump peaking at high energies (HSP). The redshift of the sources in the POLAMI sample ranges from  $z = 0.0176$  (for 3C 84) to  $z = 2.218$  (for 4C 71.07) and has a median of  $z = 0.815$ .

## 4 OBSERVATION AND CALIBRATION

### 4.1 Observation and antenna temperature calibration

The XPOL observing procedure (Thum et al. 2008) makes use of the Observatory’s horizontally (H) and vertically (V) polarized heterodyne receivers whose down-converted signals are detected into the backend VESPA after auto-correlations (Stokes  $I$  and  $Q$ ) and cross-correlation ( $U$  and  $V$ ). Stokes  $I$  and  $Q$  are calibrated like any spectroscopic observation using hot and cold loads. The system temperature  $T_{\text{sys}}$  of the  $U$  and  $V$  spectra are derived from those of the H and V spectra as  $T_{\text{sys}} = \sqrt{T_{\text{sys}}^{\text{H}} \cdot T_{\text{sys}}^{\text{V}}}$  (Thompson, Moran & Swenson 1991).

Each polarimetric observation of a source, typically of 3–15 min duration, was preceded by a pointing observation on the same source. Parameters of the pointing were recorded for an assessment of the quality of the polarimetric observation, that was based on the stability of the cross-scan profiles during the time spanned by the pointing measurements. The pointing accuracy of the telescope,

$<1$  arcsec for our observing strategy,<sup>4</sup> is negligible with regard to other effects affecting the quality of the observations, for example, the atmospheric stability.

During the course of our monitoring programme, in spring 2009, the receivers labelled ABCD (after the four different cryostats in which they were housed<sup>5</sup>) were replaced by a new receiver EMIR that consists of a single cryostat housing dual polarization (H and V) receivers for the four atmospheric windows at 3, 2, 1.3 and 0.9 mm (Carter et al. 2015). The warm optics of EMIR allowed simultaneous observations of the 3 and 1.3 mm bands, an option we used throughout our monitoring project as of 2009 December. Observing frequencies were 86.2 GHz (AB and EMIR) and 228.9 GHz (EMIR). The observing bandwidth, limited by the backend VESPA, was a nominal 640 MHz (AB and EMIR 3 mm) and 320 MHz (EMIR 1 mm). Receiver temperatures were in the range of 40–60 K (3 mm) and 80–100 K (1 mm). System temperatures ranged from 100 to 130 K at 3 mm and from 150 to 400 K at 1 mm. Beam widths are 28 arcsec (3 mm) and 12 arcsec (1.3 mm). All types of observing conditions were encountered during the observations, which were scheduled as a backup programme most of the time. The quality of the atmosphere was recorded in the observer’s logbook, but also evaluated from the pointing data (Section 4.4).

### 4.2 Total flux density calibration

After correcting for the dependence on elevation of the on-axis gain (Greve, Neri & Sievers 1998), we convert the antenna temperatures (corrected for atmospheric attenuation) into flux densities by applying a Kelvin-to-Jansky conversion factor  $f_{\text{Jy/K}}$ . We derive  $f_{\text{Jy/K}}$  at the two observing frequencies from observations of Mars and Uranus taken during the monitoring period and by comparing them with models of the emission spectrum of these two planets. The Mars model (Lellouch & Amri<sup>6</sup>) includes the disc-resolved thermal structure of the planet’s surface layers, which result in diurnal oscillations of its emitted flux density.

The brightness temperature of Uranus is strongly wavelength dependent (Griffin & Orton 1993). We applied the ESA4 model as used for SPIRE/*Herschel* photometry (Bendo et al. 2013),<sup>7</sup> 135.0 K (86 GHz) and 91.6 K (229 GHz). The long-term evolution of the planet’s 90 GHz brightness was investigated by Kramer, Moreno & Greve (2008), who found a seasonal modulation with an amplitude  $\leq 10$  per cent. The resulting brightness temperature extrapolated to our monitoring period is compatible with the brightness temperature adopted above.

Fig. 1 shows the flux densities observed at the two wavelengths and the predictions by the models after convolution with the telescope beam.<sup>8</sup> The final  $f_{\text{Jy/K}}$  factors that we applied to calibrate the total flux density, obtained from the comparison of the measurements in Fig. 1 and the models, are as follows:

$$\begin{aligned} f_{\text{Jy/K}}(3 \text{ mm}) &: (6.1 \pm 0.1) \text{ Jy K}^{-1}, \\ f_{\text{Jy/K}}(1 \text{ mm}) &: (8.6 \pm 0.4) \text{ Jy K}^{-1}. \end{aligned}$$

At 3 mm wavelength, the results from the two planets agree within measurement error, which, in turn, is comparable to the uncertainties

<sup>4</sup> See [http://www.iram.es/IRAMES/telescope/telescopeSummary/telescope\\_summary.html](http://www.iram.es/IRAMES/telescope/telescopeSummary/telescope_summary.html)

<sup>5</sup> <http://www.iram.es/IRAMES/mainWiki/AbcdforAstronomers>

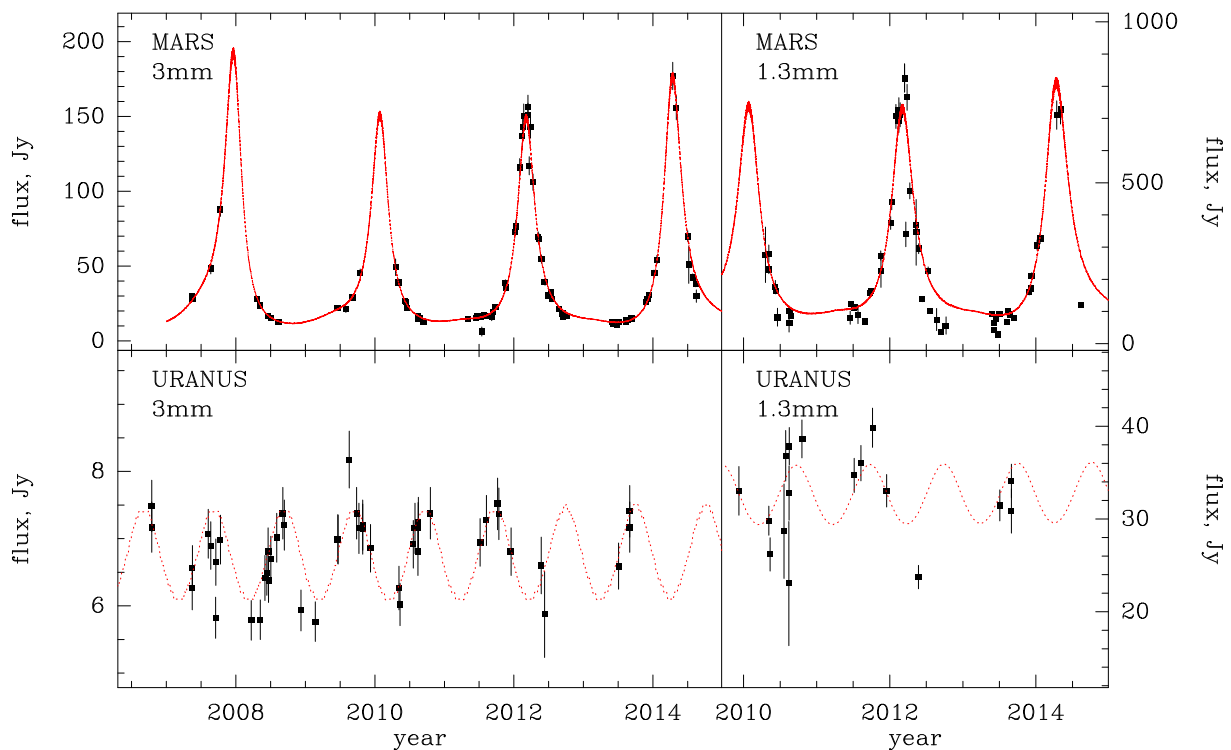
<sup>6</sup> <http://www.lesia.obspm.fr/perso/emmanuel-lellouch/mars/>

<sup>7</sup> <ftp://ftp.sciops.esa.int/pub/hsc-calibration/PlanetaryModels/ESA4/>

<sup>8</sup> For Uranus we used ASTRO which is part of the GILDAS software package. See <http://www.iram.fr/IRAMFR/GILDAS>

<sup>3</sup> BL Lac objects hereafter.





**Figure 1.** Light curves of Mars (top panel) and Uranus (bottom panel) at 3 mm (left-hand side) and 1 mm (right-hand side). Cyclic lines are light curves as predicted by the models after convolution with the planet’s disc (see the text).

estimated for the models ( $\lesssim 4$  per cent). The fact that virtually all 3 mm observations during the 8 yr of the monitoring period are well fitted by a single  $f_{Jy/K}(3\text{ mm})$  factor demonstrates the stability of the telescope and its calibration system. The increase of Uranus’s brightness temperature predicted by the seasonal effect (Kramer et al. 2008) during the monitoring period is very small ( $\lesssim 2$  per cent) and therefore not seen in our data.

At 1 mm, the telescope beam is often comparable or even smaller than the disc of Mars. The model prediction then depends on the exact beam shape used, including its near sidelobes, and is therefore less accurate than of the nearly point-like Uranus. For our estimate of  $f_{Jy/K}(1\text{ mm})$ , we therefore adopt the value derived from Uranus.

We note that the  $f_{Jy/K}$  factors derived here are slightly larger than those listed on the IRAM web page<sup>9</sup>, which were derived under best observing conditions and refer to an optimum state of the telescope. Our data taken under typical observing conditions depart by less than 10 per cent from the optimum values.

The formal errors on the determination of the  $f_{Jy/K}$  factors do not include the uncertainties in the models, for which we consider a conservative estimate of 5 per cent. The final uncertainty associated with each total flux density measurement ( $\sigma_S$ , see Table 4) include a quadratic sum of this 5 per cent (that typically dominates the 3 and 1 mm error), the error in  $f_{Jy/K}$ , and the statistical uncertainties of the measurement.

### 4.3 IP calibration

During the first phase of the monitoring programme, when the ABCD receivers were used, the main source of IP was the small ( $\leq 2$  arcsec), but unstable, misalignment between the horizontal and

**Table 2.** Evolution of the IP parameters of the XPOL polarimeter at the IRAM 30-m telescope.

Time period	$Q_i$ (per cent)	$U_i$ (per cent)	$V_i$ (per cent)	Receiver
3 mm				
2006-10-14	-1.13	-0.22	-0.11	ABCD
2006-10-15	-4.28	-0.04	0.10	ABCD
2007-05-09 → 2007-05-17	-2.79	0.46	-0.15	ABCD
2007-07-16 → 2007-07-17	-1.79	0.15	0.01	ABCD
2007-08-07 → 2007-10-11	-1.13	-0.22	-0.11	ABCD
2008-03-21 → 2008-05-07	-0.83	-0.18	-0.05	ABCD
2008-06-01	-1.70	0.11	0.12	ABCD
2008-06-11 → 2008-06-30	-2.34	-0.16	0.14	ABCD
2008-07-02 → 2008-09-02	-2.89	-0.09	0.07	ABCD
2008-09-10 → 2008-11-20	0.73	-0.18	0.18	ABCD
2008-12-07 → 2008-12-30	-1.60	0.05	0.44	ABCD
2009-02-16 → 2009-02-20	-0.04	-0.04	0.01	ABCD
2009-06-19	-2.68	-0.17	-0.44	EMIR
2009-06-20 → 2014-08-18	-0.70	-0.30	0.00	EMIR
1 mm				
2009-12-07 → 2011-10-26	-0.36	-1.07	2.57	EMIR
2011-11-16 → 2013-01-15	-0.36	-1.07	-0.42	EMIR
2013-01-29 → 2014-08-18	-0.36	-1.07	-0.92	EMIR

vertical receivers that were housed in separate dewars. Frequent calibration runs (Table 2) allowed us to track the evolution of the IP parameters  $Q_i$ ,  $U_i$ , and  $V_i$  measured in the Nasmyth coordinate system see (see Thum et al. 2008) where the receivers are located. Typically, 50 or more observations of unpolarized sources (mainly Mars and Uranus) and the Crab nebula (the source of well-known polarization properties, Aumont et al. 2010; Wiesemeyer et al. 2011)

<sup>9</sup> <http://www.iram.es/IRAMES/mainWiki/Iram30mEfficiencies>

**Table 3.** Systematic errors of Stokes parameters (measured from Mars and Uranus as standard variations in the Nasmyth system) due to uncontrolled uncertainties in IP.

	3 mm	1 mm
$\sigma_{Q_i}$	0.5 per cent	1.1 per cent
$\sigma_{U_i}$	0.4 per cent	1.7 per cent
$\sigma_{V_i}$	0.2 per cent	0.4 per cent

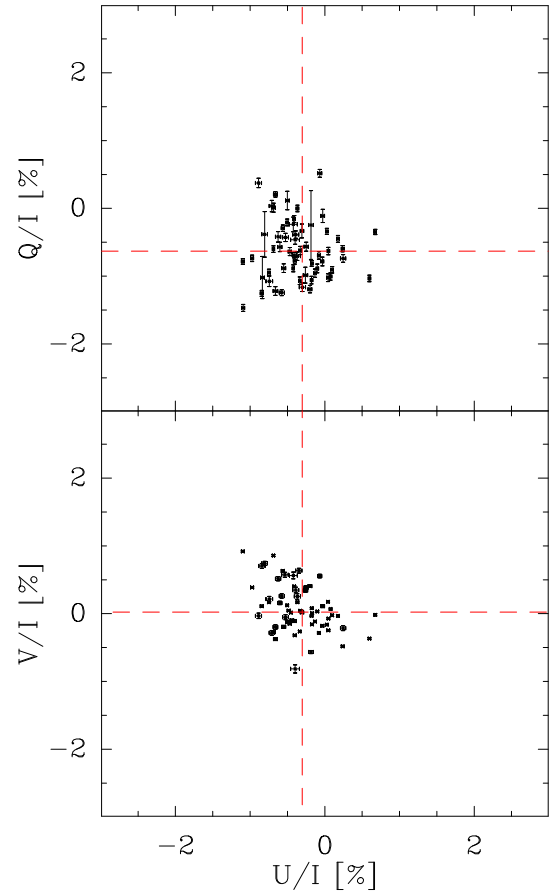
were made during the time range of our monitoring programme. After spring 2009, when EMIR was installed, the alignment between orthogonally polarized receivers became stable, since all receivers and their polarization splitting grids are housed inside a single dewar inaccessible to the outside. IP calibration runs had then to be made only after a technical warm up of the dewar.

In the data reduction, the IP parameters ( $Q_i$ ,  $U_i$  and  $V_i$ ) were subtracted from the Stokes parameters obtained in the Nasmyth system for a given session. For sessions in between IP calibration runs,  $Q_i$ ,  $U_i$  and  $V_i$  were interpolated. The resulting first stage IP-corrected calibrators were then collected in larger time bins with the aim to detect and correct for residual IP errors unnoticed in single calibration runs. Table 3 lists the first stage IP corrections adopted for the entire data set. Second stage IP corrections were then typically of the order of  $\sim 0.5$  per cent and  $\sim 1$  per cent for 3 and 1 mm, respectively, and were then subtracted from the target observations. Fig. 2 shows the results of a second stage IP measurement for the period of 2009 to 2014.

Another potential cause for IP, only affecting Stokes  $U$  and  $V$  is instrumental phase calibration errors. As outlined in Thum et al. (2008), the drift of the phase when using ABCD receivers is not larger than  $1^\circ$  per hour that can be easily calibrated. With EMIR, the phase drift is negligible, making  $V$  the potentially best determined Stokes parameter.

We see from Fig. 2 that the rms scatter of  $Q_i$ ,  $U_i$  and  $V_i$  is much larger than the individual errors of the calibrator measurements. The additional scatter results in a minor part from residual drifts of the IP, mainly due to uncontrolled small variations in the telescope optics and atmospheric fluctuations, both in brightness and in refractive index. Such imperfections, which could well be controlled or eliminated in specific short observing runs, are unavoidable in long-term monitoring programmes. We therefore take the rms scatter of the IP measurements of our calibrators (shown in Table 3) as a conservative estimate of the uncertainties of IP affecting our data. These values were added quadratically to the statistical errors of the Stokes parameters of every single measurement. The final error in almost all 3 mm observations is dominated by these systematic errors. At 1 mm, where the system temperature is much higher, the total errors are also dominated by the systematic error, except for sources weaker than about 2 Jy.

The last step of the polarization calibration is the rotation of the Stokes parameters  $Q_N$  and  $U_N$  as measured in the Nasmyth system to  $Q_{EQ}$  and  $U_{EQ}$  in the equatorial system. The former was stable enough during the course of our programme to guarantee that no additional calibration of the absolute linear polarization angle was needed. The two sets of Stokes parameters and their respective errors are related by a rotation matrix whose angle  $2\tau$  is given in Thum et al. (2008). We then derive the typical uncertainties of the linear polarization degree ( $m_L$ ), polarization angle ( $\chi$ ) and circular



**Figure 2.** IP measured at 3 mm during the time period between 2009 June 20 to 2014 August 18 from 55 independent measurements of Mars and Uranus. Dashed lines indicate weighted means of the measurements.

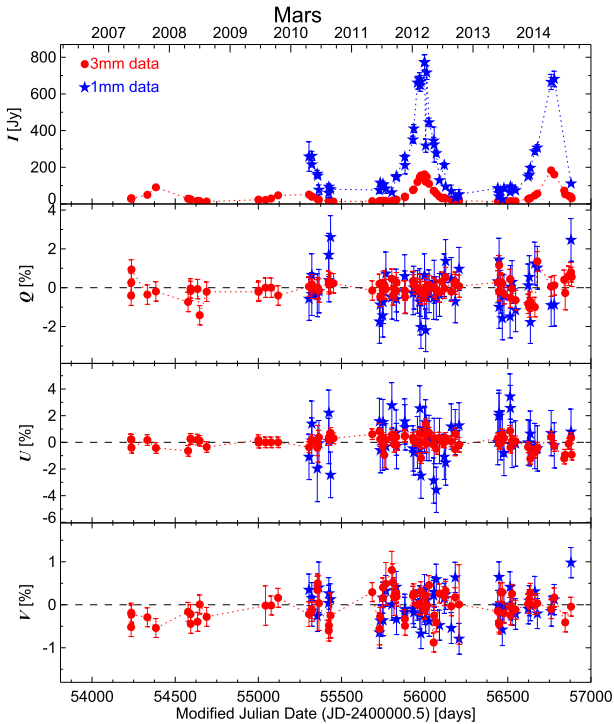
**Table 4.** Median errors of the total flux density and polarization data measured in the equatorial system from the entire source sample defined in Table 1.

	3 mm	1 mm
$\bar{\sigma}_S$	5.2 per cent	5.4 per cent
$\bar{\sigma}_{m_L}$	0.51 per cent	1.71 per cent
$\bar{\sigma}_\chi$	4:7	9:9
$\bar{\sigma}_{m_C}$	0.23 per cent	0.72 per cent

polarization ( $m_C$ ), estimated as medians over all measurements on our entire sample. These uncertainties are given in Table 4.

#### 4.4 Data quality control

An assessment of the quality of individual observations is based on the observer's logbook and an automated check that evaluates the quality of the pointing scan preceding each polarimetric observation. Four criteria were applied to every four cross-scan pointing observation: (1) the scatter of the positions derived from Gaussian fits, (2) the scatter of the width of the Gaussians with regard to the nominal beam width at our observing frequencies, (3) the receiver temperature as a diagnosis tool for technical problems, and (4) the system temperature for adverse weather conditions. Quality thresholds were set for each of the four criteria, and an observation was rejected whenever one of the criteria was passing its threshold.



**Figure 3.** Daily averages of total flux density, and Stokes  $Q$ ,  $U$  and  $V$  (Nasmyth coordinates) fully calibrated measurements of Mars made by our programme at 3 and 1 mm as a function of time. The clustering of  $Q$ ,  $U$  and  $V$  measurements of this unpolarized source around 0.0 per cent demonstrates the goodness of our polarization calibration. The dispersion on the  $Q$ ,  $U$  and  $V$  measurements reflect the accuracy on the IP determination discussed in Section 4.3. The dashed lines on the last three plots show the  $Q = 0$ ,  $U = 0$  and  $V = 0$  lines, respectively.

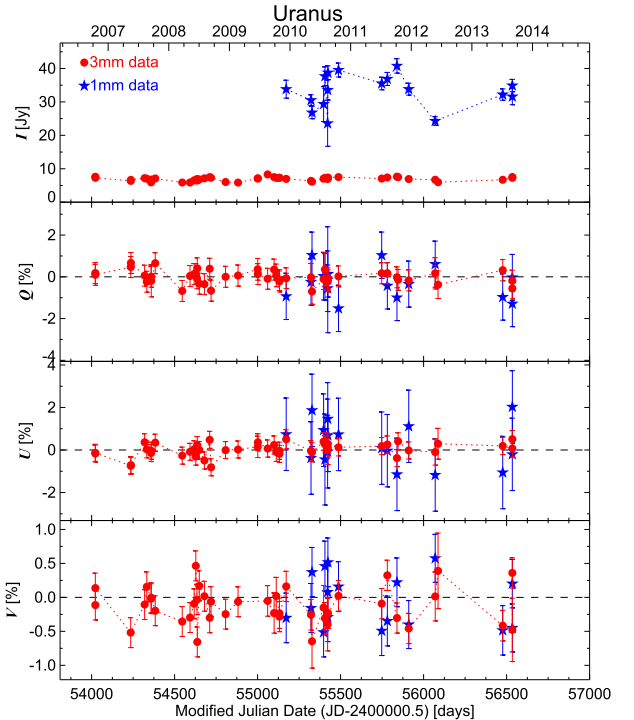
Additionally, if a 3 mm observation was rejected, the simultaneous 1 mm observation was also rejected.

A particular difficulty arose during periods of increased atmospheric instability when anomalous refraction occurred (Altenhoff et al. 1987). Under these conditions, the telescope beam can be displaced or distorted in ways which are not always caught by our criteria 1 and 2. Anomalous refraction events that preferentially occur during summer afternoons and affect observations more at 1 mm than at 3 mm explain most of the 1 mm measurements in Fig. 1, which fall below the model curves. The most severe departure occurs for Mars around 2012.6 where five consecutive observations are affected. Inspection of these sessions shows, however, that the values of  $Q$ ,  $U$  and  $V$  of the main unpolarized calibrators (Mars and Uranus) are still acceptable, i.e. consistent with zero within the errors (see next section). The data from these session albeit with somewhat larger errors, were consequently retained. Fortunately, the presence of strong anomalous refraction is easily recognized by an experienced observer on bright sources. Corresponding entries in the logbooks were used for rejecting individual observations and very occasionally whole observing sessions.

## 5 RESULTS

### 5.1 Unpolarized calibrators

Figs 3 and 4 show plots of fully calibrated  $I$ ,  $Q$ ,  $U$  and  $V$  measurements of the Stokes parameters of Mars and Uranus versus time, both for 3 and 1 mm. While the Stokes  $I$  shows considerable time



**Figure 4.** Same as Fig. 3 but for Uranus.

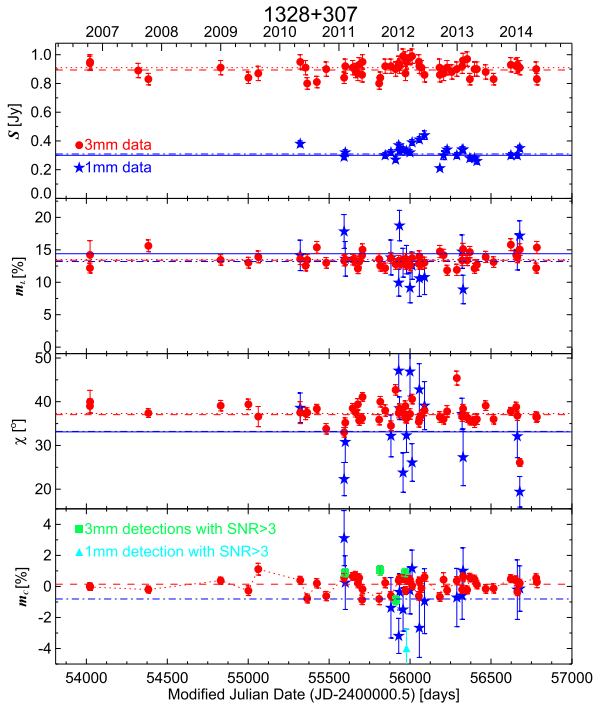
dependent variability as reproduced by the models shown in Fig. 1, the remaining Stokes parameters representing the linear and circular polarization remain rather constant and consistent with zero. This is the expected behaviour for any unpolarized source like Mars and Uranus for the IRAM 30-m Telescope, and constitutes a measure of the goodness of the polarization calibration, within the uncertainties discussed in Section 4.3.

### 5.2 Polarized calibrators

Because of the known total flux density and polarization stability of 3C 286 at centimetre wavelengths, also demonstrated in the short millimetre regime (Agudo et al. 2012b, and this paper), the target was used as a control source for our calibration. Fig. 5 shows the time series of 3 and 1 mm total flux density ( $S$ ), linear polarization degree ( $m_L$ ), linear polarization angle ( $\chi$ ) and circular polarization ( $m_C$ ) of fully calibrated measurements of 3C 286 taken during the course of our programme (see also Table 5 where we give their weighted averages and standard deviations). These measurements confirm the results shown in Agudo et al. (2012b) where we demonstrate the low level of variability of the source in  $S$ ,  $m_L$  and  $\chi$  at both 3 and 1 mm, and therefore its goodness as total flux density and polarization calibrator also at short millimetre wavelengths.

Closer inspection of the circular polarization data at 3 and 1 mm, presented here for 3C 286, shows that the observed scatter is nearly twice as high as the median values over the entire sample as given at the end of Section 4.3. Four 3 mm observations and one at 1 mm have signal-to-noise ratios  $\geq 3$  and therefore formally qualify as actual  $m_C$  detections. Contrary to the other Stokes parameters, circular polarization in this source does not appear to be stable. In Paper II (Thum et al. 2018), it is argued that  $V$  is the Stokes parameter with the highest variability, and 3C 286 may be an extreme example.

We also made measurements of the Crab nebula using the standard observing setup of our programme as a check on the overall



**Figure 5.** Daily averages of total flux density ( $S$ ), linear polarization degree ( $m_L$ ), linear polarization angle ( $\chi$ ), and circular polarization ( $m_C$ ) of fully calibrated measurements of 3C 286 (1328+307) made by our programme as a function of time. The error bars of the 1 mm total flux density measurements are typically smaller than the size of the corresponding symbols. The dashed lines indicate weighted averages of every 3 mm data train, whereas dot–dashed lines show the 1 mm weighted averages. Dotted lines symbolize the 3 mm weighted averages measured from a shorter data train, as given in Agudo et al. (2012b), and the continuous lines indicate 1 mm weighted averaged data reported in Agudo et al. (2012b). Squares and the triangle symbols on the bottom plot show 3 mm and 1 mm  $m_C$  detections, respectively, with signal-to-noise ratio larger than 3.

**Table 5.** Weighted averages and standard deviations of 3 and 1 mm measurements of  $S$ ,  $m_L$ ,  $\chi$  and  $m_C$  from 3C 286.

	3 mm	1 mm
$\langle S \rangle$	$(0.90 \pm 0.01)$ Jy	$(0.31 \pm 0.01)$ Jy
$\sigma_S$	0.05 Jy	0.06 Jy
$\langle m_L \rangle$	$(13.4 \pm 0.1)$ per cent	$(13.2 \pm 0.7)$ per cent
$\sigma_{m_L}$	1.0 per cent	2.9 per cent
$\langle \chi \rangle$	$(37.0 \pm 0.4)^\circ$	$(33.2 \pm 2.0)^\circ$
$\sigma_\chi$	2:7	8:0
$\langle m_C \rangle$	$(0.14 \pm 0.07)$ per cent	$(-0.81 \pm 0.44)$ per cent
$\sigma_{m_C}$	0.5 per cent	1.8 per cent

health of the acquisition system, in particular, the sign of Stokes  $U$  and the calibration of the phase. These measurements successfully reproduced the known linear and circular polarization properties of the source (Aumont et al. 2010; Wiesemeyer et al. 2011).

## 6 SUMMARY

In this first of a series of papers, we present the basis of the POLAMI programme. This is an ongoing full-polarization simultaneous 3 and 1 mm monitoring programme of a set of 37 bright AGN, mostly blazars. The observations were performed, from 2006

October to 2014 August, with the XPOL polarimeter on the IRAM 30-m Telescope with a median time sampling of 22 d. We presented here the motivation for this observing programme, its strategy, the source sample and the details regarding the calibration of the data. In particular, we discuss the total flux density and IP calibration of the IRAM 30-m telescope. We demonstrate the suitability of the IRAM 30-m telescope for full-polarization observations of blazars at short millimetre wavelengths. The polarization sensitivity is sufficiently high and stable for the successful development of long-term monitoring projects like POLAMI. We confirm the suitability of 3C286 as a linear polarization (degree and angle) calibrator, and we present evidence that the source is variable in circular polarization.

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