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ZK-Burgers equation for three-dimensional Rossby solitary waves and its solutions as well as chirp effect

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Abstract

Two-dimensional Rossby solitary waves propagating in a line have attracted much attention in the past decade, whereas there is few research on three-dimensional Rossby solitary waves. But as is well known, three-dimensional Rossby solitary waves are more suitable for real ocean and atmosphere conditions. In this paper, using multiscale and perturbation expansion method, a new Zakharov-Kuznetsov (ZK)-Burgers equation is derived to describe three-dimensional Rossby solitary waves that propagate in a plane. By analyzing the equation we obtain the conservation laws of three-dimensional Rossby solitary waves. Based on the sine-cosine method, we give the classical solitary wave solutions of the ZK equation; on the other hand, by the Hirota method we also obtain the rational solutions, which are similar to the solutions of the Benjamin-Ono (BO) equation, the solutions of which can describe the algebraic solitary waves. The rational solutions of the ZK equations are worth of attention. Finally, with the help of the classical solitary wave solutions, similar to the fiber soliton communication, we discuss the dissipation and chirp effect of three-dimensional Rossby solitary waves.

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1 Introduction

When nonlinearity and dispersion are exactly matched in a dispersive wave system, a kind of the stable waves of permanent form, or the solitary waves, may arise. Russell observed a solitary wave firstly in 1834. Since then, the research on solitary waves has become an advanced subject in the fields of mechanics, physics, applied mathematics, and atmospheric and oceanic sciences and attracted more and more attention [1–12]. Rossby waves as the main research object of the present paper are an important branch of solitary waves.

Rossby waves exist widely in the real world, for example, the great red spot in the atmosphere of Jupiter, eddy motion of ocean current in the Gulf of Mexico, blocking highs of the atmosphere in the earth, and so on. Rossby waves have an important effect on the propagation of energy and form and keep the western boundary current in the ocean, such as Kuroshio, East Australia current. Meanwhile, they related to many natural phenomena

such as Nino phenomenon and north Atlantic interannual oscillation. The research on Rossby waves has great theoretical meaning and application value.

As we know, solitary waves such as Rossby solitary waves can be modeled by nonlinear partial differential equations. The Korteweg-de Vries (KdV) equation

$$u_t + \mu u u_x + \delta u_{xxx} = 0 \tag{1}$$

as the pioneer relevant equation describes the generation and evolution of Rossby solitary waves [13] and admits the two-dimensional solitary wave solution

$$u = a \operatorname{sech}^2 k(x - vt). \tag{2}$$

With the development of Rossby solitary wave theories, aiming at the stratified field, Wadati [14] derived the modified KdV equation

$$u_t + \mu u^2 u_x + \delta u_{xxx} = 0 \tag{3}$$

to describe the amplitude of Rossby solitary waves. This modified KdV equation [15] possesses a two-dimensional solitary wave solution of a similar ‘sech²’ form, but its nonlinearity is stronger than that of the KdV equation and can be used to reflect the dynamics of Rossby solitary wave when the initial disturbance is stronger. Besides of the two kinds of equations, for other initial disturbance, employing a different time and space stretching transform, Meng and Lv [16] and the authors of the paper [17] also derived the Boussinesq equation

$$u_{tt} + a_1 u_{xx} + a_2 u_{xx}^2 + a_3 u_{xxxx} = 0 \tag{4}$$

to show Rossby solitary waves.

Looking at the researches mentioned, we can find that these researches were finished by some two-dimensional solitary wave equations. These two-dimensional solitary wave equations can describe the propagation of Rossby solitary waves in a line. However, for the general theory and real problems, such as the propagation of Rossby solitary waves on the real ridges in the ocean and the real mountains in the atmosphere, two-dimensional models are not enough to explain the dynamics of Rossby solitary waves, and three-dimensional models must be considered. In fact, although three-dimensional equations such as the Kadomtsev-Petviashvili (KP) equation [18] have already been derived to describe the interaction of internal solitary waves, for three-dimensional Rossby solitary waves, commonly referred to as ‘lumps’, have received far less attention than the two-dimensional Rossby solitary waves. In this paper, the three-dimensional Rossby solitary waves fall into our research thesis.

On the other hand, the construction of the solutions of solitary wave models has been extensively investigated in the past decade, and many methods have been proposed to solve the soliton equation [19–24]. For the above-mentioned model, similar solutions of the ‘sech²’ form are presented, and the solitary waves owning this type solutions are called classical solitary waves. With the development of solitary wave theory, the integro-differential equation

$$u_t + a_1 u u_x + a_2 H(u_{xx}) = 0$$

is also derived to describe Rossby solitary waves and is called the BO equation, where H denotes the well-known Hilbert transformation. In the equation, the integral term $H(u_{xx})$ expresses the dispersive effect. The solution of the BO equation is the rational function

$$u(x, t) = \frac{\delta}{1 + \delta^2(x - \delta t)^2},$$

and therefore the solitary waves expressed by the BO model are called algebraic solitary waves [25, 26]. Recently, professor Ma and his cooperates investigated the rational solutions of the soliton equations by the Hirota method and Maple software in [27, 28], which motivated us to use the Hirota method for rational solutions of our three-dimension solitary wave model.

The concept of ‘chirp effect’ [29] is put forward to reflect frequency modulation effect in the fiber soliton communication. Because of the interaction between nonlinearity and dispersive, the excursion of the center wave happens, which can result in the occurrence of chirp effect. For the fiber soliton, the different parts of the pulse produce different frequencies, so the chirp effect occurs during the propagation. Motivated by the chirp effect in fiber soliton communication, we discuss the chirp effect that happens during the propagation of Rossby solitary waves in the atmosphere and ocean. This paper is little involved in the past research. Particularly, we note that dissipation effect plays an important role during the propagation of Rossby solitary waves in the real ocean and atmosphere, so that the dissipation effect on the three-dimensional Rossby solitary waves also falls into our research scope.

The main purpose of the present work is to establish a model of three-dimensional Rossby solitary waves and to reveal the effect of dissipation and chirp on the three-dimensional Rossby solitary waves during propagation. The plan of the paper is as follows. In Section 2, by adopting new three-dimensional time and space stretching transform, we derive a new ZK-Burgers equation to describe the evolution of three-dimensional Rossby solitary waves from the quasi-geostrophic vorticity equation. Based on the ZK-Burgers equation, the conservation laws of three-dimensional Rossby solitary waves are discussed and given in Section 3. In Section 4, we give the classical solitary wave solutions of the ZK-Burgers equation by using the sine-cosine method and the rational solutions by using the Hirota method. The rational solutions of ZK equation are obtained firstly in this paper. Based on these analytical solutions, the dissipation and chirp effect is discussed. Finally, some conclusions are given in Section 5.

2 ZK-Burgers equation

The theoretical basis of this study is the nondimensional barotropic and quasi-geostrophic potential vorticity equation including turbulent dissipation on a β -plane channel, written as

$$\left(\frac{\partial}{\partial t} + \frac{\partial \psi}{\partial x} \frac{\partial}{\partial y} - \frac{\partial \psi}{\partial y} \frac{\partial}{\partial x} \right) (\nabla^2 \psi + \beta y) = -\mu_0 \nabla^2 \psi + Q, \tag{5}$$

where ∇^2 is the Laplace operator, ψ is a geostrophic stream function, $\nabla^2 \psi$ denotes the vorticity dissipation caused by the Ekman boundary layer, $0 \leq \mu_0 \ll 1$ expresses the dissipative coefficient, and Q denotes the heating function to be given later.

Lateral boundary conditions is a nondimensional rigid wall boundary condition are taken as

$$\frac{\partial \psi}{\partial x} = 0, \quad y = 0, 1. \tag{6}$$

Here $y = 0$ and $y = 1$ denote the south and north boundaries of zonal flow, respectively.

We assume that the stream function is the sum of zonal flow and disturbance stream function of the form

$$\psi = - \int_0^y [u(s) - c + \epsilon^2 \alpha] ds + \epsilon^2 \psi', \tag{7}$$

where the small parameter ϵ is much less than 1, α is called the detuning parameter and reflects the proximity of the system to a resonate state, and the Rossby waves phase speed c is taken as a constant.

Substituting (7) into (5) and (6) and dropping the apostrophe for simplicity, we obtain the following nonlinear equation and boundary conditions for the disturbance stream function:

$$\frac{\partial}{\partial t} \nabla^2 \psi + (u(y) - c + \epsilon^2 \alpha) \frac{\partial}{\partial x} \nabla^2 \psi + [\beta - u''(y)] \frac{\partial \psi}{\partial x} + \epsilon^2 J(\psi, \nabla^2 \psi) = -\mu_0 \nabla^2 \psi, \tag{8}$$

$$\frac{\partial \psi}{\partial x} = 0, \quad y = 0, 1. \tag{9}$$

We introduce the time and space stretching transform as follows:

$$X = \epsilon x, \quad Y = \epsilon y, \quad T = \epsilon^3 t. \tag{10}$$

We note that this transform is different from the two-dimensional transform in [13]. Meanwhile, in order to balance the dissipation and nonlinearity, we assume that

$$\mu_0 = \epsilon^3 \mu, \quad Q = -\mu_0 \frac{du}{dy}. \tag{11}$$

Substituting (10) and (11) into (8) and (9), we have

$$\begin{aligned} & \left[\epsilon^3 \frac{\partial}{\partial T} + \epsilon(u - c + \epsilon^2 \alpha) \frac{\partial}{\partial X} + \epsilon^3 \mu \right] \left(\epsilon^2 \frac{\partial^2 \psi}{\partial X^2} + \epsilon^2 \frac{\partial^2 \psi}{\partial Y^2} + 2\epsilon \frac{\partial^2 \psi}{\partial Y \partial y} + \frac{\partial^2 \psi}{\partial y^2} \right) \\ & + \epsilon(\beta - u'') \frac{\partial \psi}{\partial X} + \epsilon^3 \frac{\partial \psi}{\partial X} \left(\epsilon^3 \frac{\partial^3 \psi}{\partial Y^3} + 3\epsilon \frac{\partial^3 \psi}{\partial y^2 \partial Y} + 3\epsilon^2 \frac{\partial^3 \psi}{\partial y \partial Y^2} + \frac{\partial^3 \psi}{\partial y^3} \right) \\ & - \epsilon^3 \left(\frac{\partial \psi}{\partial y} + \epsilon \frac{\partial \psi}{\partial Y} \right) \left(\epsilon^2 \frac{\partial^3 \psi}{\partial Y^2 \partial X} + \frac{\partial^3 \psi}{\partial y^2 \partial X} + 2\epsilon \frac{\partial^3 \psi}{\partial y \partial Y \partial X} \right) \\ & + \epsilon^5 \frac{\partial \psi}{\partial X} \frac{\partial^3 \psi}{\partial y \partial X^2} + \epsilon^6 \frac{\partial \psi}{\partial X} \frac{\partial^3 \psi}{\partial Y \partial X^2} - \epsilon^5 \frac{\partial \psi}{\partial y} \frac{\partial^3 \psi}{\partial X^3} - \epsilon^6 \frac{\partial \psi}{\partial Y} \frac{\partial^3 \psi}{\partial X^3} = 0, \end{aligned} \tag{12}$$

$$\frac{\partial \psi}{\partial X} = 0, \quad y = 0, 1. \tag{13}$$

Expanding the disturbance stream function in powers of ϵ according the WKB method

$$\psi = \psi_0 + \epsilon \psi_1 + \epsilon^2 \psi_2 + \dots \tag{14}$$

and then substituting (14) into (12) and (13), we obtain

$$\epsilon^1: \begin{cases} (u - c) \frac{\partial}{\partial X} \frac{\partial^2 \psi_0}{\partial y^2} + [\beta - u''(y)] \frac{\partial \psi_0}{\partial X} = 0, \\ \frac{\partial \psi_0}{\partial X} = 0, \quad y = 0, 1. \end{cases} \tag{15}$$

Assuming that ψ_0 has a separable solution of the form

$$\psi_0 = A(X, Y, T)\phi_0(y), \tag{16}$$

substituting (16) into (15), and assuming that $u \neq c$, we have

$$\begin{cases} \frac{d^2 \phi_0}{dy^2} + \frac{\beta - u''(y)}{u(y) - c} \phi_0 = 0, \\ \phi_0(y) = 0, \quad y = 0, 1. \end{cases} \tag{17}$$

In order to obtain the governing model of Rossby waves, we continue by solving the higher-order problem

$$\epsilon^2: \begin{cases} (u - c) \frac{\partial}{\partial X} \frac{\partial^2 \psi_1}{\partial y^2} + [\beta - u''(y)] \frac{\partial \psi_1}{\partial X} = -2(u - c) \frac{\partial^3 \psi_0}{\partial X \partial Y \partial y}, \\ \frac{\partial \psi_1}{\partial X} = 0, \quad y = 0, 1. \end{cases} \tag{18}$$

By analysis we find that it is reasonable to assume that the solution ψ_1 has the following form:

$$\psi_1 = A_Y(X, Y, T)\phi_1(y). \tag{19}$$

In a similar way, substituting (19) into (18) leads to

$$\begin{cases} \frac{d^2 \phi_1}{dy^2} + \frac{\beta - u''(y)}{u - c} \phi_1 = -2(u - c) \frac{\partial \phi_0}{\partial y}, \\ \phi_1(y) = 0, \quad y = 0, 1. \end{cases} \tag{20}$$

We find that the governing model of Rossby solitary waves cannot be derived in order to $O(\epsilon^2)$, so that we need to proceed to the higher-order equation

$$\epsilon^3: \begin{cases} (u - c) \frac{\partial}{\partial X} \frac{\partial^2 \psi_2}{\partial y^2} + [\beta - u''(y)] \frac{\partial \psi_2}{\partial X} = -2(u - c) \frac{\partial^3 \psi_1}{\partial X \partial Y \partial y} - \frac{\partial}{\partial T} \frac{\partial^2 \psi_0}{\partial y^2} - \mu \frac{\partial^2 \psi_0}{\partial y^2} \\ - \alpha \frac{\partial}{\partial X} \frac{\partial^2 \psi_0}{\partial y^2} - (u - c) \frac{\partial^3 \psi_0}{\partial X^3} - (u - c) \frac{\partial^3 \psi_0}{\partial X \partial Y^2} + \frac{\partial \psi_0}{\partial y} \frac{\partial^3 \psi_0}{\partial X \partial y^2} - \frac{\partial \psi_0}{\partial X} \frac{\partial^3 \psi_0}{\partial y^3}, \\ \frac{\partial \psi_2}{\partial X} = 0, \quad y = 0, 1. \end{cases} \tag{21}$$

Substituting (16) and (19) into (21), we have

$$\begin{aligned} \frac{\partial}{\partial X} \left(\frac{\partial^2 \psi_2}{\partial y^2} + \frac{\beta - u''}{u - c} \psi_2 \right) = & -2 \frac{\partial^3 A}{\partial X \partial Y^2} \phi_{1y} + \left(\frac{\partial A}{\partial T} + \alpha \frac{\partial A}{\partial X} + \mu A \right) \frac{\beta - u''}{(u - c)^2} \phi_0 - \frac{\partial^3 A}{\partial X^3} \phi_0 \\ & - \frac{\partial^3 A}{\partial X \partial Y^2} \phi_0 + \frac{1}{u - c} (\phi_{0y} \phi_{0yy} - \phi_0 \phi_{0yyy}) A \frac{\partial A}{\partial X}. \end{aligned} \tag{22}$$

Multiplying both sides of Eq. (22) by ϕ_0 and integrating it over y from 0 to 1, by virtue of (17) and (20) and the identical relation

$$\phi_0 \frac{\partial^2 \psi_2}{\partial y^2} = \frac{\partial}{\partial y} \left(\phi_0 \frac{\partial \psi_2}{\partial y} \right) - \frac{\partial}{\partial y} \left(\psi_2 \frac{d\phi_0}{dy} \right) + \psi_2 \frac{d^2 \phi_0}{dy^2}, \tag{23}$$

we get

$$A_T + \alpha A_X + a_0 A A_X + a_1 A_{XX} + \mu A + a_2 A_{XY} = 0, \tag{24}$$

where the coefficients satisfy

$$\begin{cases} a_0 = \int_0^1 \frac{\phi_0}{u-c} (\phi_{0y} \phi_{0yy} - \phi_0 \phi_{0yyy}) dy / \sigma, & a_1 = - \int_0^1 \phi_0^2 dy / \sigma, \\ a_2 = - \int_0^1 (\phi_0^2 + 2\phi_0 \phi_{1y}) dy / \sigma, & \sigma = \int_0^1 \frac{\beta-u''}{(u-c)^2} \phi_0^2 dy. \end{cases} \tag{25}$$

Equation (24) is three-dimensional. When $a_2 = 0$, Eq. (24) reduces to the KdV-Burgers equation, which has been derived to describe two-dimensional Rossby solitary waves; when the dissipation is absent, that is, $\mu = 0$, Eq. (24) becomes a high-dimensional KdV equation, the ZK equation, so we call it the ZK-Burgers equation.

3 Conservation laws of three-dimensional Rossby solitary waves

A conservation law [30] is an important concept to describe the conservation of fundamental physical quantity and plays an important role in analyzing unsteady problems of waves during propagation. In this section, we discuss the conservation laws of three-dimensional Rossby solitary waves.

In what follows, we make the following assumption:

$$A, A_X, A_{XX}, A_{XY} \rightarrow 0 \quad \text{as } |X|, |Y| \rightarrow \infty. \tag{26}$$

Equation (24) can be rewritten as

$$\frac{\partial}{\partial T} A + \frac{\partial}{\partial X} \left(\alpha A + \frac{a_0}{2} A^2 + a_1 A_{XX} + a_2 A_{XY} \right) + \mu A = 0. \tag{27}$$

First, integrating Eq. (27) with respect to X and Y from $-\infty$ to $+\infty$, by (26) we get

$$C_1 = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} A dX dY = \exp(-\mu T) \int_{-\infty}^{+\infty} A(X, Y, 0) dX dY. \tag{28}$$

It is easy to find that when dissipation is absent, that is, $\mu = 0$, C_1 is a time-invariant quantity and denotes the conservative mass of three-dimensional Rossby solitary waves. Furthermore, from (28) we can also find that the mass of the three-dimensional Rossby solitary waves decreases exponentially with the increasing of time T and the dissipative coefficient in the presence of dissipation.

Second, multiplying Eq. (24) by $2A$, we get

$$\begin{aligned} \frac{\partial}{\partial T} (A^2) + \frac{\partial}{\partial X} \left(\alpha A^2 + \frac{2a_0}{3} A^3 + 2a_1 A A_{XX} - a_1 A_X^2 + 2a_2 A A_{XY} + a_2 A_Y^2 \right) \\ - \frac{\partial}{\partial Y} (2a_2 A_X A_Y) + \mu A^2 = 0. \end{aligned} \tag{29}$$

In a similar way, integrating Eq. (29) with respect to X and Y from $-\infty$ to $+\infty$, by (26) we get

$$C_2 = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} A^2 dX dY = \exp(-2\mu T) \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} A^2(X, Y, 0) dX dY. \tag{30}$$

Here C_2 also is a time-invariant quantity and regarded as the momentum of three-dimensional Rossby solitary waves. Without dissipation, the momentum of three-dimensional Rossby solitary waves is conserved; with dissipation, it decreases exponentially with the increasing of time T and the dissipative coefficient; moreover, it decreases more quickly than the mass.

Third, multiplying both sides of Eq. (24) by $3A^2$, we obtain the first equation; taking the derivative of Eq. (24) with respect to X and then multiplying by $-6a_1A_X/a_0$ lead to the second equation; taking the derivative of Eq. (24) with respect to Y and then multiplying $-6a_2A_Y/a_0$ lead to the third equation. Summing the above three equations and neglecting dissipation and then integrating the obtained equation with respect to X and Y from $-\infty$ to $+\infty$, we find that

$$C_3 = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \left(A^3 - \frac{3a_1}{a_0} A_X^2 - \frac{3a_2}{a_0} A_Y^2 \right) dX dY \tag{31}$$

is also a time-invariant quantity and shows the conservative energy of three-dimensional Rossby solitary waves. After tedious calculation (which we omit here), we also find that the energy of three-dimensional Rossby solitary waves decreases with the increasing of time T and dissipative coefficient μ .

Besides to the above three conservation laws, let us finally construct the new quantity

$$\tilde{C}_4 = \frac{d}{dT} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} (X + Y)A dX dY, \tag{32}$$

which is related to the phase of the three-dimensional Rossby solitary waves. Neglecting dissipation and using (26), (28), and (30), we easily find that \tilde{C}_4 is a time-invariant quantity, that is, $\frac{d\tilde{C}_4}{dT} = 0$. We are more interested in the quantity

$$C_4 = \frac{\tilde{C}_4}{C_1}, \tag{33}$$

which can be used to express the velocity of the center of gravity for the ensemble of such three-dimensional Rossby solitary waves. Because C_1 and \tilde{C}_4 are both time-invariant quantities, we easily obtain that C_4 is also a time-invariant quantities; in other words, the velocity of the center of gravity of three-dimensional Rossby solitary waves is conserved.

4 The analytical solutions of ZK equation

As a well-known soliton equation, the KdV equation plays a key role in the development of soliton theory. However, the KdV equation is considered a $(1 + 1)$ -dimensional model (one spatial and one time dimensions). The best known two-dimensional generalizations of the KdV equation are the KP and ZK equations. For the ZK equation, there were many researches discussing the analytical solutions of the ZK equation in different conditions

with different methods [30–40], such as the ansatz method, the modified F -expansion method, and so on, some important results were obtained. Equation (24) we obtain in the paper is different from the common ZK equation due to the presence of dissipation effect. Because of the relation between KdV and ZK equation, it is reasonable to consider the classical solitary wave solutions of the ‘sech²’ form; on the other hand, motivated by the research about rational solutions of soliton equation by professor Ma, we have to think deeply whether rational solutions of the ZK equation exist. So this section is be divided into two subsections: one is devoted to seeking classical solitary wave solutions of the ZK equation by the sine-cosine method, and the other one is devoted to exploring the rational solutions of the ZK equation.

4.1 Classical solitary wave solutions

Assuming that

$$\xi = X + Y - (C + \alpha)T, \tag{34}$$

where C denotes the speed of wave, we have the following transform form:

$$\begin{aligned} \frac{\partial}{\partial T} &= -(C + \alpha)\frac{d}{d\xi}, & \frac{\partial}{\partial X} &= \frac{d}{d\xi}, & \frac{\partial}{\partial Y} &= \frac{d}{d\xi}, \\ \frac{\partial^2}{\partial X^2} &= \frac{d^2}{d\xi^2}, & \frac{\partial^2}{\partial Y^2} &= \frac{d^2}{d\xi^2}. \end{aligned} \tag{35}$$

Substituting (35) into Eq. (24) and neglecting dissipation (*i.e.*, $\mu = 0$), Eq. (24) can be rewritten as

$$-CA_\xi + \frac{a_0}{2}(A^2)_\xi + a_1A_{\xi\xi\xi} + a_2A_{\xi\xi\xi} = 0. \tag{36}$$

Integrating Eq. (39) with respect to ξ and taking the integrable constant as zero lead to

$$-CA + \frac{a_0}{2}A^2 + (a_1 + a_2)A_{\xi\xi} = 0. \tag{37}$$

Suppose that Eq. (24) has solutions of the following form:

$$A(X, Y, T) = [\lambda \cos^\tau(\gamma\xi)], \quad |\xi| \leq \frac{\pi}{2\gamma}, \tag{38}$$

or

$$A(X, Y, T) = [\lambda \sin^\tau(\gamma\xi)], \quad |\xi| \leq \frac{\pi}{2\gamma}, \tag{39}$$

where λ , τ , and γ are undetermined parameters; γ shows the wave number. Based on this assumption, we have

$$\begin{aligned} A(\xi) &= \lambda \cos^\tau(\gamma\xi), & A^2(\xi) &= \lambda^2 \cos^{2\tau}(\gamma\xi), \\ A_{\xi\xi} &= -\gamma^2\tau^2\lambda \cos^{\tau-2}(\gamma\xi) + \gamma^2\lambda\tau(\tau - 1) \cos^{\tau-2}(\gamma\xi) \end{aligned} \tag{40}$$

and

$$\begin{aligned}
 A(\xi) &= \lambda \sin^\tau(\gamma\xi), & A^2(\xi) &= \lambda^2 \sin^{2\tau}(\gamma\xi), \\
 A_{\xi\xi} &= -\gamma^2 \tau^2 \lambda \sin^\tau(\gamma\xi) + \gamma^2 \lambda \tau(\tau - 1) \sin^{\tau-2}(\gamma\xi).
 \end{aligned}
 \tag{41}$$

Substituting (40) and (41) into Eq. (37), we have

$$-C \cos^\tau(\gamma\xi) + \frac{a_0}{2} \lambda^2 \cos^{2\tau}(\gamma\xi) - (a_1 + a_2) \gamma^2 \tau \lambda \cos^{\tau-2}(\gamma\xi) (\tau \sin^2(\gamma\xi) - 1) = 0
 \tag{42}$$

and

$$-C \sin^\tau(\gamma\xi) + \frac{a_0}{2} \lambda^2 \sin^{2\tau}(\gamma\xi) - (a_1 + a_2) \gamma^2 \tau \lambda \sin^{\tau-2}(\gamma\xi) (\tau \cos^2(\gamma\xi) - 1) = 0.
 \tag{43}$$

In order to make these two equations hold, the following algebraic equations must be satisfied:

$$\begin{cases}
 \tau - 1 \neq 0, \\
 \tau - 2 = 2\tau, \\
 (a_1 + a_2) \gamma^2 \tau^2 = -C, \\
 (a_1 + a_2) \gamma^2 \lambda \tau (\tau - 1) = -\frac{a_0}{2} \lambda^2.
 \end{cases}
 \tag{44}$$

By a direct calculation we get

$$\tau = -2, \quad \gamma = \frac{1}{2} \sqrt{\frac{-C}{a_1 + a_2}}, \quad \lambda = \frac{3C}{a_0}.
 \tag{45}$$

If $\frac{C}{a_1 + a_2} < 0$, then it is easy to obtain periodic wave solutions of Eq. (24) of the following form:

$$A(X, Y, T) = \frac{3C}{a_0} \operatorname{csc}^2 \left(\frac{1}{2} \sqrt{\frac{-C}{a_1 + a_2}} [X + Y - (C + \alpha)T] \right), \quad \frac{C}{a_1 + a_2} < 0
 \tag{46}$$

and

$$A(X, Y, T) = \frac{3C}{a_0} \operatorname{sec}^2 \left(\frac{1}{2} \sqrt{\frac{-C}{a_1 + a_2}} [X + Y - (C + \alpha)T] \right), \quad \frac{C}{a_1 + a_2} < 0.
 \tag{47}$$

Meanwhile, if $\frac{C}{a_1 + a_2} > 0$, we can obtain solitary wave solutions of Eq. (24) of the following form:

$$A(X, Y, T) = \frac{-3C}{a_0} \operatorname{csch}^2 \left(\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} [X + Y - (C + \alpha)T] \right), \quad \frac{C}{a_1 + a_2} > 0
 \tag{48}$$

and

$$A(X, Y, T) = \frac{-3C}{a_0} \operatorname{sech}^2 \left(\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} [X + Y - (C + \alpha)T] \right), \quad \frac{C}{a_1 + a_2} > 0,
 \tag{49}$$

4.2 Rational solutions

In the same way, neglecting dissipation effect and assuming that $\alpha = 0$ and $\mu = 0$, Eq. (24) can be rewritten as

$$A_T + a_0AA_X + a_1A_{XXX} + a_2A_{XY} = 0. \tag{50}$$

For convenience, we normalize the coefficients in Eq. (50) and adopt the following transformation:

$$A = \frac{6\sqrt[3]{a_1}A'}{\sqrt[3]{3a_0}}, \quad X = \sqrt[3]{\frac{3}{a_1}}X', \quad Y = \frac{6\sqrt[3]{9a_1}}{\sqrt[3]{a_2}}Y'. \tag{51}$$

Omitting the apostrophes, Eq. (50) becomes

$$A_T + 6AA_X + 3A_{XXX} + 3A_{XY} = 0. \tag{52}$$

Introducing the variable substitution $A(X, Y, T) = v(\xi, t)$, $\xi = l_1X + l_2Y$, letting $6l_1 = \gamma_1$ and $3l_1^2 + 3l_2^2 = \gamma_2$, adopting the bilinear transform

$$v = \frac{12\gamma_2}{\gamma_1}(\ln \tau)_{\xi\xi}, \quad f = \frac{g}{\tau}, \tag{53}$$

and letting $v = w_\xi = \left(\frac{12\gamma_2 \tau_\xi}{\gamma_1 \tau}\right)_\xi$, after direct calculation, we have

$$\begin{aligned} &\frac{12\gamma_2}{\gamma_1} \left(\frac{\tau_{\xi t}}{\tau} - \frac{\tau_\xi \tau_t}{\tau^2} \right) + \frac{\gamma_1}{2} \frac{12\gamma_2}{\gamma_1} \left(\frac{\tau_{\xi\xi}^2}{\tau^2} - \frac{2\tau_{\xi\xi} \tau_\xi}{\tau^3} + \frac{\tau_\xi^4}{\tau^4} \right) \\ &+ \frac{12\gamma_2}{\gamma_1} \left(\frac{\tau_{\xi\xi\xi\xi}}{\tau} - \frac{4\tau_{\xi\xi\xi} \tau_\xi}{\tau^2} - \frac{3\tau_{\xi\xi}^2}{\tau^2} + \frac{12\tau_{\xi\xi} \tau_\xi^2}{\tau^3} - \frac{6\tau_\xi^4}{\tau^4} \right) = 0. \end{aligned} \tag{54}$$

Then based on Eq. (54), we obtain the bilinear form of Eq. (52):

$$D_\xi(D_t + \gamma_2 D_\xi^3)\tau \cdot \tau = 0. \tag{55}$$

The definition of the generalized bilinear D operator is

$$\begin{aligned} D_{p,\xi}^m D_{p,t}^n F \cdot F &= \left(\frac{\partial}{\partial \xi} + \alpha_p \frac{\partial}{\partial \xi'} \right)^m \left(\frac{\partial}{\partial t} + \alpha_p \frac{\partial}{\partial t'} \right)^n F(\xi, t) F(\xi', t') \Big|_{\xi'=\xi, t'=t} \\ &= \sum_{i=0}^m \sum_{j=0}^n \binom{m}{i} \binom{n}{j} \alpha_p^i \alpha_p^j \frac{\partial^{m-i}}{\partial \xi^{m-i}} \frac{\partial^i}{\partial \xi'^i} \frac{\partial^{n-j}}{\partial t^{n-j}} \frac{\partial^j}{\partial t'^j} F(\xi, t) F(\xi', t') \Big|_{\xi'=\xi, t'=t} \\ &= \sum_{i=0}^m \sum_{j=0}^n \binom{m}{i} \binom{n}{j} \alpha_p^i \alpha_p^j \frac{\partial^{m+n-i-j} F(\xi, t)}{\partial \xi^{m-i} \partial t^{n-j}} \frac{\partial^{i+j} F(\xi, t)}{\partial \xi^i \partial t^j}, \quad m, n \geq 0, \end{aligned}$$

where α_p^s is calculated as follows:

$$\alpha_p^s = (-1)^{r_p(s)}, \quad s = r_p(s) \bmod p,$$

$$D_{2,t}D_{2,\xi}F \cdot F = 2F_{\xi,t}F - 2F_tF_\xi, \quad D_{2,\xi}^4F \cdot F = 2F_{4\xi}F - 8F_{3\xi}F_\xi + 6F_{2\xi}^2.$$

Consequently, the linear combination of Eq. (55) yields the ZK equation

$$D_\xi(D_t + \beta D_\xi^3)F \cdot F = 2F_{\xi,t}F - 2F_tF_\xi + 2\beta F_{4\xi}F - 8\beta F_{3\xi}F_\xi + 6\beta F_{2\xi}^2. \tag{56}$$

We use the symbolic computation with Maple and obtain polynomial solutions of degrees of x and t less than eight to the generalized bilinear ZK equation,

$$F = \sum_{i=0}^8 \sum_{j=0}^8 c_{ij} \xi^i t^j, \tag{57}$$

where the c_{ij} are constants, and acquire 25 classes of polynomial solutions to the Eq. (56). Among the 25 classes of solutions, we enumerate four classes of solutions as follows:

$$\begin{aligned} & \frac{c_{2,7}c_{3,2}\xi^2t^2}{c_{3,7}} + \frac{c_{2,7}c_{3,1}\xi^2t}{c_{3,7}} + \frac{c_{2,7}c_{3,4}\xi^2t^4}{c_{3,7}} + \frac{c_{2,7}c_{3,6}\xi^2t^6}{c_{3,7}} + \frac{c_{2,7}c_{3,3}\xi^2t^3}{c_{3,7}} \\ & + \frac{c_{2,7}c_{3,5}\xi^2t^5}{c_{3,7}} + c_{0,7}t^7 + \frac{c_{2,7}^2c_{3,1}\xi t}{3c_{3,7}^2} + \frac{c_{2,7}^2c_{3,5}\xi t^5}{3c_{3,7}^2} + \frac{c_{2,7}^2c_{3,2}\xi t^2}{3c_{3,7}^2} \\ & + \frac{c_{2,7}^2c_{3,4}\xi t^4}{3c_{3,7}^2} + \frac{c_{2,7}^2c_{3,3}\xi t^3}{3c_{3,7}^2} + \frac{c_{2,7}^2c_{3,6}\xi t^6}{3c_{3,7}^2} + c_{3,0}\xi^3 + c_{3,1}\xi^3t \\ & + c_{3,4}\xi^3t^4 + c_{3,2}\xi^3t^2 + c_{2,7}\xi^2t^7 + c_{3,3}\xi^3t^3 + c_{3,6}\xi^3t^6 + c_{3,7}\xi^3t^7 + c_{3,5}\xi^3t^5 \\ & + \frac{(12\gamma_2c_{3,5}c_{3,7} - 12\gamma_2c_{3,6}^2 + c_{0,7}c_{3,6})t^6}{c_{3,7}} + \frac{(12\gamma_2c_{3,1}c_{3,7} - 12\gamma_2c_{3,2}c_{3,6} + c_{0,7}c_{3,2})t^2}{c_{3,7}} \\ & + \frac{(12\gamma_2c_{3,2}c_{3,7} - 12\gamma_2c_{3,3}c_{3,6} + c_{0,7}c_{3,3})t^3}{c_{3,7}} + \frac{(12\gamma_2c_{3,0}c_{3,7} - 12\gamma_2c_{3,1}c_{3,6} + c_{0,7}c_{3,1})t}{c_{3,7}} \\ & + \frac{(12\gamma_2c_{3,4}c_{3,7} - 12\gamma_2c_{3,5}c_{3,6} + c_{0,7}c_{3,5})t^5}{c_{3,7}} + \frac{(12\gamma_2c_{3,3}c_{3,7} - 12\gamma_2c_{3,4}c_{3,6} + c_{0,7}c_{3,4})t^4}{c_{3,7}} \\ & + 12\gamma_2c_{3,7}t^8 - \frac{c_{3,0}(12\gamma_2c_{3,6} - c_{0,7})}{c_{3,7}} + \frac{c_{2,7}c_{3,0}\xi^2}{c_{3,7}} + \frac{c_{2,7}^2\xi t^7}{3c_{3,7}} + \frac{c_{2,7}^2c_{3,0}\xi}{3c_{3,7}^2}, \tag{58} \\ & c_{3,0}\xi^3 + \frac{c_{2,6}^2c_{3,1}\xi t}{3c_{3,6}^2} + \frac{c_{2,6}^2c_{3,5}\xi t^5}{3c_{3,6}^2} + \frac{c_{2,6}^2c_{3,2}\xi t^2}{3c_{3,6}^2} + c_{3,1}\xi^3t + c_{3,4}\xi^3t^4 + c_{3,2}\xi^3t^2 \\ & + c_{3,3}\xi^3t^3 + c_{3,6}\xi^3t^6 + c_{3,5}\xi^3t^5 + \frac{c_{2,6}c_{3,2}\xi^2t^2}{c_{3,6}} + \frac{c_{2,6}c_{3,1}\xi^2t}{c_{3,6}} + \frac{c_{2,6}c_{3,4}\xi^2t^4}{c_{3,6}} \\ & + \frac{c_{2,6}c_{3,3}\xi^2t^3}{c_{3,6}} + \frac{c_{2,6}c_{3,5}\xi^2t^5}{c_{3,6}} + \frac{(12\gamma_2c_{3,4}c_{3,6} - 12\gamma_2c_{3,5}^2 + c_{0,6}c_{3,5})t^5}{c_{3,6}} \\ & + \frac{(12\gamma_2c_{3,3}c_{3,6} - 12\gamma_2c_{3,4}c_{3,5} + c_{0,6}c_{3,4})t^4}{c_{3,6}} + \frac{c_{2,6}^2c_{3,0}\xi}{3c_{3,6}^2} + \frac{c_{2,6}^2\xi t^6}{3c_{3,6}} \\ & - \frac{c_{3,0}(12\gamma_2c_{3,5} - c_{0,6})}{c_{3,6}} + \frac{(12\gamma_2c_{3,1}c_{3,6} - 12\gamma_2c_{3,2}c_{3,5} + c_{0,6}c_{3,2})t^2}{c_{3,6}} \\ & + \frac{(12\gamma_2c_{3,2}c_{3,6} - 12\gamma_2c_{3,3}c_{3,5} + c_{0,6}c_{3,3})t^3}{c_{3,6}} + \frac{(12\gamma_2c_{3,0}c_{3,6} - 12\gamma_2c_{3,1}c_{3,5} + c_{0,6}c_{3,1})t}{c_{3,6}} \end{aligned}$$

$$\begin{aligned}
 & + \frac{c_{2,6}^2 c_{3,3} \xi t^3}{3c_{3,6}^2} + \frac{c_{2,6}^2 c_{3,4} \xi t^4}{3c_{3,6}^2} + \frac{c_{2,6} c_{3,0} \xi^2}{c_{3,6}} + c_{0,6} t^6 + c_{2,6} \xi^2 t^6 + 12\gamma_2 c_{3,6} t^7, \tag{59} \\
 & \frac{(1,866,240\gamma_2^2 c_{6,0}^2 c_{6,1}^4 - 62,208\gamma_2 c_{3,1} c_{6,0} c_{6,1}^4 + 6,480\gamma_2 c_{5,1}^3 c_{6,0} c_{6,1}^2 - 432c_{1,1} c_{5,1} c_{6,1}^4)}{2,592c_{6,1}^6 c_{6,0}^{-1}} \\
 & - \frac{c_{6,0}(5,184c_{3,1}^2 c_{6,1}^4 - 720c_{3,1} c_{5,1}^3 c_{6,1}^2 + 25c_{5,1}^6)}{25,920c_{6,1}^6} + c_{1,1} x t + c_{3,1} x^3 t + 30\gamma_2 c_{5,1} \xi^2 t^2 \\
 & + \frac{c_{5,1}(72c_{3,1} c_{6,1}^2 - 5c_{5,1}^3) \xi^2 t}{144c_{6,1}^3} + \frac{c_{5,1} c_{6,0} \xi^5}{c_{6,1}} + \frac{\gamma_2(1,440\gamma_2 c_{6,0} c_{6,1}^2 - 48c_{3,1} c_{6,1}^2 + 5c_{5,1}^3) t^2}{2c_{6,1}^2} \\
 & - \frac{c_{5,1} c_{6,0}(4,320\gamma_2 c_{6,0} c_{6,1}^2 - 72c_{3,1} c_{6,1}^2 + 5c_{5,1}^3) \xi^2}{144c_{6,1}^4} + 60\gamma_2 c_{6,1} \xi^3 t^2 - 720\gamma_2^2 c_{6,1}^3 t^3 \\
 & + \frac{(18,662,400\gamma_2^2 c_{6,0}^2 c_{6,1}^4 + 4,320c_{1,1} c_{5,1} c_{6,1}^4 - 5,184c_{3,1}^2 c_{6,1}^4 + 720c_{3,1} c_{5,1}^3 c_{6,1}^2 - 25c_{5,1}^6) t}{25,920c_{6,1}^5} \\
 & + \frac{5\beta c_{5,1}^2 \xi t^2}{c_{6,1}} - \frac{c_{6,0}(60\gamma_2 c_{6,0} - c_{3,1}) \xi^3}{c_{6,1}} + \frac{5c_{5,1}^2 c_{6,0} \xi^4}{12c_{6,1}^2} - \frac{c_{6,0}(5\gamma_2 c_{5,1}^2 c_{6,0} - c_{1,1} c_{6,1}^2) \xi}{c_{6,1}^3} \\
 & + \frac{5c_{5,1}^2 \xi^4 t}{12c_{6,1}} + c_{6,1} \xi^6 t + c_{5,1} \xi^5 t + c_{6,0} \xi^6, \tag{60} \\
 & \frac{4,320c_{1,0} c_{5,0} c_{6,0}^4 - 5,184c_{3,0}^2 c_{6,0}^4 + 720c_{3,0} c_{5,0}^3 c_{6,0}^2 - 25c_{5,0}^6}{25,920c_{6,0}^5} + \frac{5\gamma_2 c_{5,0}^2 \xi t}{c_{6,0}} + c_{5,0} \xi^5 \\
 & - 720\gamma_2^2 c_{6,0} t^2 + \frac{c_{5,0}(72c_{3,0} c_{6,0}^2 - 5c_{5,0}^3) \xi^2}{144c_{6,0}^3} - \frac{\gamma_2(48c_{3,0} c_{6,0}^2 - 5c_{5,0}^3) t}{2c_{6,0}^2} \\
 & + c_{6,0} \xi^6 + c_{3,0} \xi^3 + \frac{5c_{5,0}^2 \xi^4}{12c_{6,0}} + c_{1,0} \xi + 60\gamma_2 c_{6,0} \xi^3 t + 30\gamma_2 c_{5,0} \xi^2 t, \tag{61}
 \end{aligned}$$

where the involved constants c_{ij} are arbitrary, provided that the solutions are meaningful. We can confirm that there are two distinct classes of rational solutions generated from (57) to the ZK equation by considering the transformation of the coefficient c_{ij} :

$$\begin{aligned}
 u & = (3xc_{3,7} + c_{2,7})(3x^3 c_{3,7}^2 + 36\beta t c_{3,7}^2 + 3x^2 c_{2,7} c_{3,7} - 36\beta c_{3,6} c_{3,7} + xc_{2,7}^2 + 3c_{0,7} c_{3,7})^{-2} \\
 & \times (-9x^3 c_{3,7}^3 + 216\beta t c_{3,7}^3 - 9x^2 c_{2,7} c_{3,7}^2 - 216\beta c_{3,6} c_{3,7}^2 \\
 & - 3xc_{2,7}^2 c_{3,7} + 18c_{0,7} c_{3,7}^2 - c_{2,7}^3) \tag{62}
 \end{aligned}$$

and

$$\begin{aligned}
 u & = -360 \frac{p}{q}, \tag{63} \\
 p & = 11,197,440x^{10} c_{6,0}^{10} + 18,662,400x^9 c_{5,0} c_{6,0}^9 + 60,466,176,000\beta^2 t^2 x^4 c_{6,0}^{10} \\
 & + 13,996,800x^8 c_{5,0}^2 c_{6,0}^8 + 483,729,408,000\beta^3 t^3 x c_{6,0}^{10} \\
 & + 40,310,784,000\beta^2 t^2 x^3 c_{5,0} c_{6,0}^9 + 6,220,800x^7 c_{5,0}^3 c_{6,0}^7 \\
 & + 80,621,568,000\beta^3 t^3 c_{5,0} c_{6,0}^9 + 10,077,696,000\beta^2 t^2 x^2 c_{5,0}^2 c_{6,0}^8 \\
 & + 2,015,539,200\beta t x^4 c_{3,0} c_{6,0}^9 - 186,624,000\beta t x^4 c_{5,0}^3 c_{6,0}^7 + 1,814,400x^6 c_{5,0}^4 c_{6,0}^6
 \end{aligned}$$

$$\begin{aligned}
 &+ 24,186,470,400\beta^2 t^2 x c_{3,0} c_{6,0}^9 - 1,119,744,000\beta^2 t^2 x c_{5,0}^3 c_{6,0}^7 \\
 &+ 1,343,692,800\beta t x^3 c_{3,0} c_{5,0} c_{6,0}^8 - 124,416,000\beta t x^3 c_{5,0}^4 c_{6,0}^6 \\
 &- 33,592,320x^5 c_{1,0} c_{6,0}^9 + 2,799,360x^5 c_{3,0} c_{5,0}^2 c_{6,0}^7 + 129,600x^5 c_{5,0}^5 c_{6,0}^5 \\
 &+ 4,031,078,400\beta^2 t^2 c_{3,0} c_{5,0} c_{6,0}^8 - 326,592,000\beta^2 t^2 c_{5,0}^4 c_{6,0}^6 \\
 &+ 335,923,200\beta t x^2 c_{3,0} c_{5,0}^2 c_{6,0}^7 - 31,104,000\beta t x^2 c_{5,0}^5 c_{6,0}^5 - 27,993,600x^4 c_{1,0} c_{5,0} c_{6,0}^8 \\
 &+ 16,796,160x^4 c_{3,0}^2 c_{6,0}^8 - 777,600x^4 c_{3,0} c_{5,0}^3 c_{6,0}^6 + 403,107,840\beta t x c_{3,0}^2 c_{6,0}^8 \\
 &- 37,324,800\beta t x c_{3,0} c_{5,0}^3 c_{6,0}^6 - 9,331,200x^3 c_{1,0} c_{5,0}^2 c_{6,0}^7 + 11,197,440x^3 c_{3,0}^2 c_{5,0} c_{6,0}^7 \\
 &- 1,296,000x^3 c_{3,0} c_{5,0}^4 c_{6,0}^5 + 36,000x^3 c_{5,0}^7 c_{6,0}^3 + 67,184,640\beta t c_{3,0}^2 c_{5,0} c_{6,0}^7 \\
 &- 10,886,400\beta t c_{3,0} c_{5,0}^4 c_{6,0}^5 + 432,000\beta t c_{5,0}^7 c_{6,0}^3 - 1,555,200x^2 c_{1,0} c_{5,0}^3 c_{6,0}^6 \\
 &+ 2,799,360x^2 c_{3,0}^2 c_{5,0}^2 c_{6,0}^6 - 388,800x^2 c_{3,0} c_{5,0}^5 c_{6,0}^4 + 13,500x^2 c_{5,0}^8 c_{6,0}^2 \\
 &- 129,600x c_{1,0} c_{5,0}^4 c_{6,0}^5 + 2,239,488x c_{3,0}^3 c_{6,0}^7 - 311,040x c_{3,0}^2 c_{5,0}^3 c_{6,0}^5 \\
 &+ 10,800x c_{3,0} c_{5,0}^6 c_{6,0}^3 + 1,866,240c_{1,0}^2 c_{6,0}^8 - 311,040c_{1,0} c_{3,0} c_{5,0}^2 c_{6,0}^6 \\
 &+ 21,600c_{1,0} c_{5,0}^5 c_{6,0}^4 + 373,248c_{3,0}^3 c_{5,0} c_{6,0}^6 - 77,760c_{3,0}^2 c_{5,0}^4 c_{6,0}^4 \\
 &+ 5,400c_{3,0} c_{5,0}^7 c_{6,0}^2 - 125c_{5,0}^{10} c_{6,0}^2, \\
 q = &(-25,920x^6 c_{6,0}^6 - 1,555,200\beta t x^3 c_{6,0}^6 - 25,920x^5 c_{5,0} c_{6,0}^5 + 18,662,400\beta^2 t^2 c_{6,0}^6 \\
 &- 777,600\beta t x^2 c_{5,0} c_{6,0}^5 - 10,800x^4 c_{5,0}^2 c_{6,0}^4 - 129,600\beta t x c_{5,0}^2 c_{6,0}^4 - 25,920x^3 c_{3,0} c_{6,0}^5 \\
 &+ 622,080\beta t c_{3,0} c_{5,0}^5 - 64,800\beta t c_{5,0}^3 c_{6,0}^3 - 12,960x^2 c_{3,0} c_{5,0} c_{6,0}^4 + 900x^2 c_{5,0}^4 c_{6,0}^2 \\
 &- 25,920x c_{1,0} c_{5,0}^5 - 4,320c_{1,0} c_{5,0} c_{6,0}^4 + 5,184c_{3,0}^2 c_{6,0}^4 - 720c_{3,0} c_{5,0}^3 c_{6,0}^2 + 25c_{5,0}^6)^2.
 \end{aligned}$$

In fact, the polynomial solutions in the first group of (58)-(59) and the second group of (60)-(61) change into (62) and (63), respectively.

The first class of solutions in (62) reduces to

$$u = \frac{p}{q}, \tag{64}$$

$$\begin{aligned}
 p = &(-0.11979 \cdot 10^{-7} x^3 + 0.287496 \cdot 10^{-9} t - 0.1089 \cdot 10^{-7} x^2 + 0.1616264 \cdot 10^{-7} \\
 &- 0.33 \cdot 10^{-8} x)(0.33 \cdot 10 - 2x + 0.1 \cdot 10^{-2}),
 \end{aligned}$$

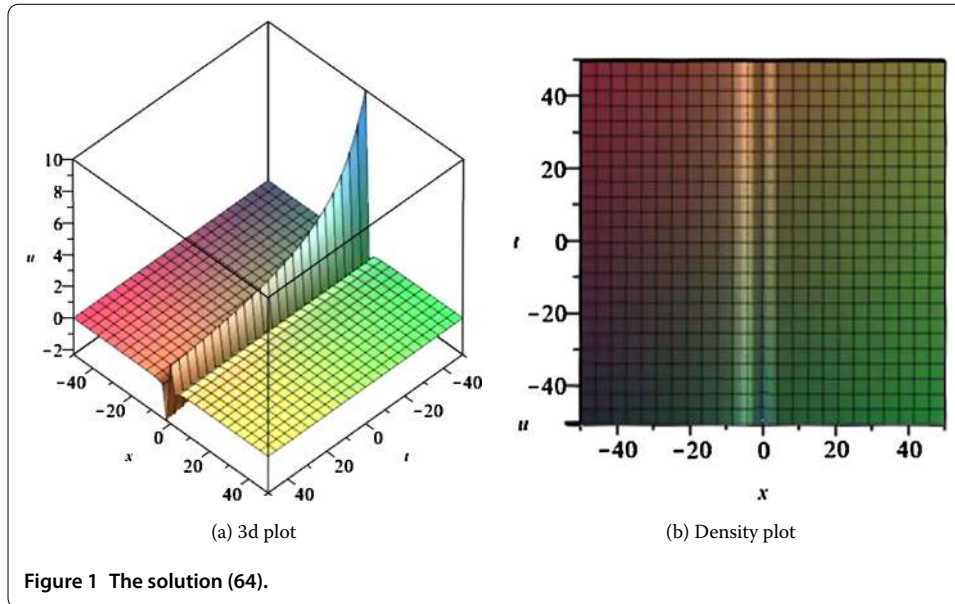
$$q = (0.363 \cdot 10^{-5} x^3 + 0.4356 \cdot 10^{-7} t + 0.33 \cdot 10^{-5} x^2 + 0.26004 \cdot 10^{-5} + 0.1 \cdot 10^{-5} x)^2,$$

when $c_{3,7} = 0.0011$, $c_{2,7} = 0.001$, $c_{3,6} = 0.001$, $c_{0,7} = 0.0008$, $\beta = 0.001$. The picture of the solution (64) is presented in Figure 1.

From the second class of solutions in (63) we obtain

$$u = \frac{p}{q}, \tag{65}$$

$$\begin{aligned}
 p = &-0.17640 \cdot 10^{-5} (0.4404601350 \cdot 10^{-38} t + 0.9501485144 \cdot 10^{-36} x \\
 &+ 0.5079242590 \cdot 10^{-40} t^2 + 0.1952022642 \cdot 10^{-42} t^3 + 0.3162999652 \cdot 10^{-34} x^{10} \\
 &+ 0.4518570932 \cdot 10^{-34} x^9 - 0.4663481371 \cdot 10^{-35} x^4 + 0.2904795599 \cdot 10^{-34} x^8
 \end{aligned}$$



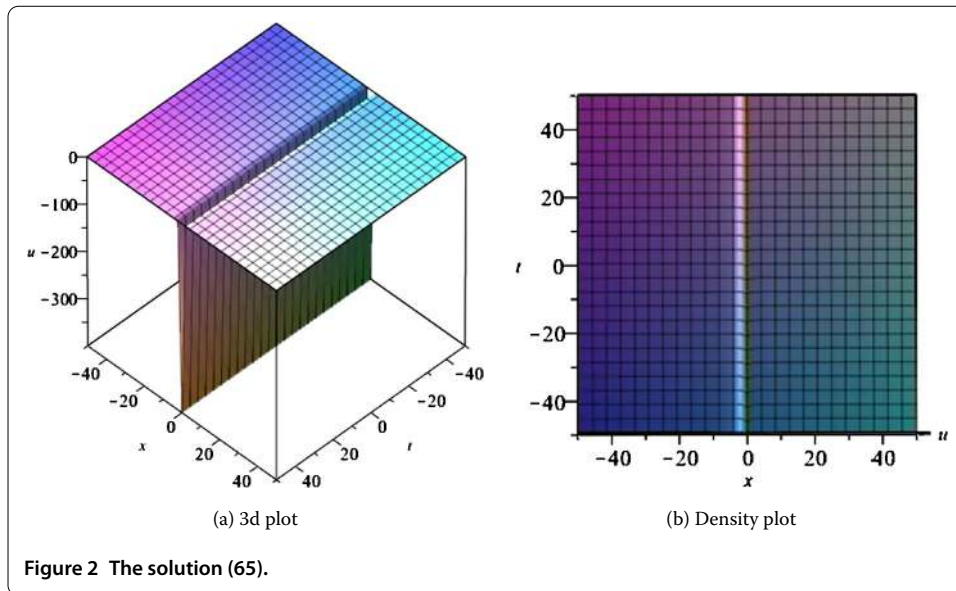
$$\begin{aligned}
 & + 0.1106588800 \cdot 10^{-34}x^7 + 0.2766471999 \cdot 10^{-35}x^6 - 0.2362228335 \cdot 10^{-34}x^5 \\
 & + 0.8273408826 \cdot 10^{-36}x^2 + 0.2225160652 \cdot 10^{-35}x^3 + 0.4658120333 \cdot 10^{-36} \\
 & + 0.1708019812 \cdot 10^{-38}t^2x^4 + 0.1366415850 \cdot 10^{-41}t^3x + 0.9760113212 \cdot 10^{-39}t^2x^3 \\
 & + 0.2091452831 \cdot 10^{-39}t^2x^2 + 0.2921394431 \cdot 10^{-36}tx^4 + 0.3704859301 \cdot 10^{-39}t^2x \\
 & + 0.1669368246 \cdot 10^{-36}tx^3 + 0.3577217671 \cdot 10^{-37}tx^2 + 0.3338736493 \cdot 10^{-37}tx), \\
 q = & (-0.3049462080 \cdot 10^{-20}x^6 - 0.1829677248 \cdot 10^{-22}tx^3 - 0.2613824640 \cdot 10^{-20}x^5 \\
 & + 0.2195612698 \cdot 10^{-25}t^2 - 0.7841473920 \cdot 10^{-23}tx^2 - 0.933508800 \cdot 10^{-21}x^4 \\
 & - 0.1120210560 \cdot 10^{-23}tx - 0.1742549760 \cdot 10^{-20}x^3 + 0.3702029184 \cdot 10^{-23}t \\
 & - 0.689653440 \cdot 10^{-21}x^2 - 0.871274880 \cdot 10^{-21}x + 0.45365184 \cdot 10^{-22})^2,
 \end{aligned}$$

when $c_{6,0} = 0.00007$, $c_{5,0} = 0.00006$, $c_{3,0} = 0.00004$, $c_{1,0} = 0.00002$, $\beta = 0.0001$. The picture of the solution (65) is presented in Figure 2.

Remarks As we know, the integro-differential equations (BO, ILW) which own the rational function solutions can be used to describe the algebraic solitary waves; these rational solutions can be used to describe the fission of solitary waves and to explain the possible physical mechanism of the squall line in the real atmosphere. On the other hand, the differential equations (KdV, mKdV) can describe the classic solitary waves. Now, through the above deduction, we can learn that the ZK-Burgers equation as a differential equation has rational function solutions. Can it also describe algebraic solitary waves? It is an open problem. We will discuss it in the future.

5 Dissipation and chirp effect of three-dimensional Rossby solitary waves

In Section 4, we have derived the analytical solutions of Eq. (24) without dissipation effect. In this section, based on the solitary wave solution (49), on the one hand, we will seek



an approximate analytical solution of Eq. (20) with dissipation and, furthermore, discuss the influence of dissipation on the three-dimensional Rossby solitary waves by using the obtained approximate analytical solution. On the other hand, we will study the chirp effect of three-dimensional Rossby solitary waves.

5.1 Dissipation effect of three-dimensional Rossby solitary waves

Because of $\mu \ll 1$ and $\mu \ll a_0 \sim a_1 + a_2$, we take the new space coordinate

$$\rho = X + Y - \int_0^T \alpha + \frac{a_0 A_0}{3} dT. \tag{66}$$

Assuming that $A_0 = A_0(\mu T)$ and substituting (66) into Eq. (24), we have

$$\frac{\partial A}{\partial T} - \frac{a_0 A_0}{3} \frac{\partial A}{\partial \rho} + a_0 A \frac{\partial A}{\partial \rho} + (a_1 + a_2) \frac{\partial^3 A}{\partial \rho^3} + \mu A = 0. \tag{67}$$

Introducing two time scales

$$\tau = T, \quad \eta = \mu T, \tag{68}$$

and assuming the solution as

$$A(\rho, T) = A_1(\rho, \tau, \eta) + \mu A_2(\rho, \tau) + \dots, \tag{69}$$

we obtain the following approximate equations:

$$\mu^0: \frac{\partial A_1}{\partial \tau} - \frac{a_0 A_0}{3} \frac{\partial A_1}{\partial \rho} + a_0 A_1 \frac{\partial A_1}{\partial \rho} + (a_1 + a_2) \frac{\partial^3 A_1}{\partial \rho^3} = 0, \tag{70}$$

$$\mu^1: \frac{\partial A_2}{\partial \tau} - \frac{a_0 A_0}{3} \frac{\partial A_2}{\partial \rho} + a_0 \left(A_1 \frac{\partial A_2}{\partial \rho} + A_2 \frac{\partial A_1}{\partial \rho} \right) + (a_1 + a_2) \frac{\partial^3 A_2}{\partial \rho^3} = -\frac{\partial A_1}{\partial \eta} - A_1. \tag{71}$$

Taking

$$\zeta = \rho + \frac{a_0 A_0}{3} \tau, \tag{72}$$

Eqs. (70) and (71) can be rewritten as

$$\mu^0: \frac{\partial A_1}{\partial \tau} + a_0 A_1 \frac{\partial A_1}{\partial \rho} + (a_1 + a_2) \frac{\partial^3 A_1}{\partial \rho^3} = 0, \tag{73}$$

$$\mu^1: \frac{\partial A_2}{\partial \tau} + a_0 \left(A_1 \frac{\partial A_2}{\partial \rho} + A_2 \frac{\partial A_1}{\partial \rho} \right) + (a_1 + a_2) \frac{\partial^3 A_2}{\partial \rho^3} = -\frac{\partial A_1}{\partial \eta} - A_1. \tag{74}$$

Equation (73) is the common ZK equation; without loss of generality, based on (49), we take into account its solitary wave solution as

$$A_1(\zeta, \tau) = -A_0 \operatorname{sech}^2 \left\{ \left(\frac{A_0 a_0}{12(a_1 + a_2)} \right)^{\frac{1}{2}} \left(\zeta - \frac{a_0 A_0}{3} \tau \right) \right\}, \tag{75}$$

where $A_0 = \frac{3C}{a_0}$. Substituting (50) into (75), we have the following approximate solitary wave solution of Eq. (24):

$$A_1(X, T) = -A_0 \operatorname{sech}^2 \left\{ \left(\frac{A_0 a_0}{12(a_1 + a_2)} \right)^{\frac{1}{2}} \left(X + Y - \int_0^T (\alpha + C) dT \right) \right\} \frac{C}{a_1 + a_2} > 0. \tag{76}$$

Next, in order to determine the form of $A_0(\mu T)$, it is necessary to solve Eq. (74). Taking

$$A_2 = B(I), \quad I = \zeta - C\tau, \tag{77}$$

and substituting (77) into Eq. (74) lead to

$$-v \frac{dB}{dI} + a_0 \frac{\partial}{\partial I} (A_1 B) + (a_1 + a_2) \frac{\partial^3 B}{\partial I^3} = M(A_1), \tag{78}$$

where $M(A_1) = -\frac{\partial A_1}{\partial \eta} - A_1$. The solvability condition of Eq. (78) is

$$\int_{-\infty}^{+\infty} G(I) M(A_1) dI = 0, \tag{79}$$

where $G(I)$ satisfies

$$-v \frac{dG}{dI} + a_0 A_1 \frac{dG}{dI} + (a_1 + a_2) \frac{\partial^3 G}{\partial I^3} = 0. \tag{80}$$

Obviously, if $G(\pm\infty) = 0$ is satisfied, then the solution of Eq. (64) is

$$G = -A_0 \operatorname{sech}^2 \left\{ \left(\frac{A_0 a_0}{12(a_1 + a_2)} \right)^{\frac{1}{2}} (\zeta - vT) \right\}. \tag{81}$$

Substituting (81) into Eq. (79), we have

$$A_0 = \frac{3C}{a_0} e^{-\frac{4\mu T}{3}}. \tag{82}$$

So the approximate solitary wave solution of the ZK-Burgers equation is

$$A(X, T) = -A_1(T) \operatorname{sech}^2 \left\{ \left(\frac{A_1(T)a_0}{12(a_1 + a_2)} \right)^{\frac{1}{2}} \left(X + Y - \int_0^T \alpha + \frac{a_0 A_1(T)}{3} dT \right) \right\}, \quad (83)$$

where

$$A_1(T) = \frac{3C}{a_0} e^{-\frac{4\mu T}{3}}. \quad (84)$$

In the same way, we can also obtain other approximate periodic wave and solitary wave solutions. Based on (83), it is easy to find that the propagation speed and width of three-dimensional Rossby solitary waves are as follows, respectively:

$$c_s = \alpha + C e^{-\frac{4\mu T}{3}}, \quad (85)$$

$$W_s = 2 \sqrt{\frac{a_1 + a_2}{C}} e^{\frac{2}{3}\mu T}. \quad (86)$$

It is obvious that the propagation speed of three-dimensional Rossby solitary waves decreases exponentially with the increasing of μ and time T ; on the contrary, the width of three-dimensional Rossby solitary waves increases exponentially with the increasing of μ and time T . The figures of three-dimensional Rossby solitary waves at three different time points are as follows.

From Figures 3-5 we also find that, due to dissipation effect, the amplitude of three-dimensional Rossby solitary waves also decreases with time T .

Figure 3 Three-dimensional Rossby solitary wave with $\mu = 0.05$ at $T = 0$.

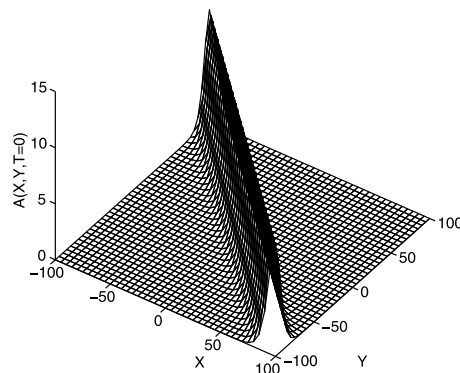
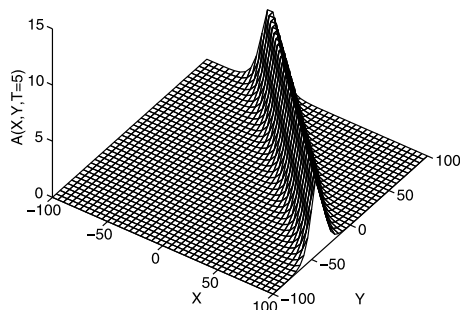
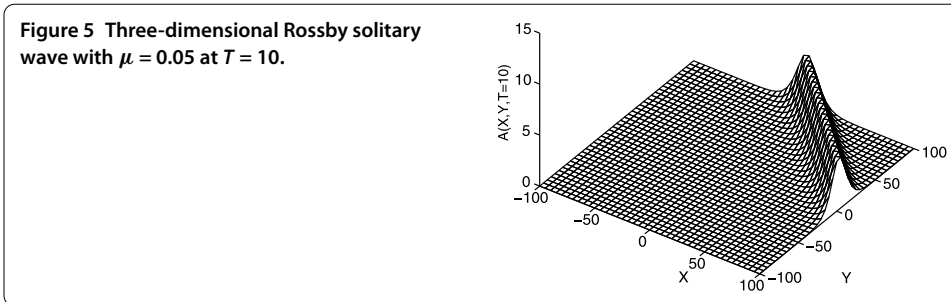


Figure 4 Three-dimensional Rossby solitary wave with $\mu = 0.05$ at $T = 5$.





5.2 Chirp effect of three-dimensional Rossby solitary waves

The Chirp effect is a concept from optical soliton communication area and is also called the frequency modulation effect. Because of the dispersion and nonlinear effect, the excursion phenomenon of the center wave occurs, and the chirp effect is caused. In this subsection, we introduce the chirp idea of optical soliton communication and analyze the dispersion and nonlinear effect during the propagation process of three-dimensional Rossby solitary waves. It is beneficial for studying the propagation feature of three-dimensional Rossby solitary waves.

For the ZK-Burgers equation (24), let us neglect dissipation effect ($\mu = 0$) and take the detuning parameter $\alpha = 0$, that is, the system is resonant. Then the equation can be rewritten as

$$A_T + a_0AA_X + a_1A_{XXX} + a_2A_{XYY} = 0, \tag{87}$$

where a_0 is the nonlinear coefficient, and a_1 and a_2 are dispersion coefficients. Here, based on the analytical solution of the ZK equation, we take the initial wave form of three-dimensional Rossby solitary waves as

$$A = \frac{-C}{a_0} \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right], \quad \frac{C}{a_1 + a_2} > 0. \tag{88}$$

First, we discuss the nonlinear effect of three-dimensional Rossby solitary waves alone. Then Eq. (87) becomes

$$A_T = -a_0AA_X. \tag{89}$$

Investigating the condition of time T from 0 to ΔT , where Δt is an infinitesimal variable, and introducing (88) into Eq. (89), we obtain the following approximate solution of Eq. (89):

$$A(\Delta T, X, Y) = \frac{C^2}{a_0^2} \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \exp \left\{ \frac{-2C^4 \Delta T}{a_0^3} \right. \\ \left. \times \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \tanh \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right\}. \tag{90}$$

Based on the approximate solution (90), it is easy to give the phase of the wave as follows:

$$\begin{aligned} \varphi_N &= \frac{-2C^4 \Delta T}{a_0^3} \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \\ &\quad \times \tanh \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right]. \end{aligned} \tag{91}$$

Then the chirp caused by nonlinearity is

$$\begin{aligned} \Delta v_N &= -\nabla \varphi_N \\ &= \frac{2C^4 \Delta T}{a_0^3} \sqrt{\frac{C}{a_1 + a_2}} \left\{ 2 \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right. \\ &\quad \left. \times \tanh^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] + \tanh^4 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right\}. \end{aligned} \tag{92}$$

Second, neglecting the nonlinearity effect, let us consider the dispersion effect alone. Then Eq. (87) becomes

$$A_T = -a_1 A_{XXX} - a_2 A_{XYT}. \tag{93}$$

In a similar way, we can get an approximate solution of Eq. (93),

$$\begin{aligned} A(\Delta T, X, Y) &= \frac{C^2}{a_0^2} \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \exp \left\{ -\frac{(a_1 + a_2) \Delta T}{16} \left(\frac{C}{a_1 + a_2} \right)^{\frac{3}{2}} \right. \\ &\quad \times \left[1 + 2 \operatorname{sech}^2 \left(\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right) \right. \\ &\quad \left. \left. \times \tanh \left(\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right) \right] \right\}, \end{aligned} \tag{94}$$

so that the phase of the wave meets

$$\begin{aligned} \varphi_D &= -\frac{C^{\frac{3}{2}} \Delta T}{16 \sqrt{a_1 + a_2}} \left\{ 1 + 2 \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right. \\ &\quad \left. \times \tanh \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right\}. \end{aligned} \tag{95}$$

Based on (95), we can obtain the chirp caused by the dispersion effect

$$\begin{aligned} \Delta v_D &= -\nabla \varphi_D \\ &= \frac{C^2 \Delta T}{a_1 + a_2} \left\{ 2 \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right. \\ &\quad \left. \times \tanh^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] + \tanh^4 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right\}. \end{aligned} \tag{96}$$

According to (92) and (96), the whole chirp is

$$\begin{aligned} \Delta v_S &= \Delta v_N + \Delta v_D \\ &= \left(\frac{2C^4}{a_0^3} + \frac{C^{\frac{3}{2}}}{8\sqrt{a_1 + a_2}} \right) \frac{C}{a_1 + a_2} \Delta T \left\{ 2 \operatorname{sech}^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right. \\ &\quad \left. \times \tanh^2 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] + \tanh^4 \left[\frac{1}{2} \sqrt{\frac{C}{a_1 + a_2}} (X + Y) \right] \right\}. \end{aligned} \tag{97}$$

By employing (97) we find that when the wave speed C satisfies the formula

$$C^5 = \frac{a_0^6}{2^8(a_1 + a_2)}, \tag{98}$$

the nonlinear effect balances to the dispersion effect, that is, $\Delta v_S = 0$; when the wave speed C satisfies

$$C^5 < \frac{a_0^6}{2^8(a_1 + a_2)}, \tag{99}$$

the dispersive effect is greater than the nonlinear effect, that is, $|\Delta v_D| > |\Delta v_N|$; on the contrary, when the wave speed C satisfies

$$C^5 > \frac{a_0^6}{2^8(a_1 + a_2)}, \tag{100}$$

the dispersion effect is less than the nonlinear effect, that is, $|\Delta v_D| < |\Delta v_N|$.

6 Conclusion

In this paper, the ZK-Burgers equation is constructed to describe the three-dimensional Rossby solitary waves by employing the multiscale and perturbation expansion method. By analyzing the ZK-Burgers equation we obtain four conservation laws of three-dimensional Rossby solitary waves, that is, mass, momentum, energy, velocity of the center gravity and discuss their variation trends with dissipation effect. Moreover, analytical solutions of the ZK-Burgers equation are derived based on the sine-cosine and Hirota methods. With the help of analytical solutions, the dissipation and chirp effect are studied. The results show that, due to the dissipation effect, the propagation speed and amplitude of three-dimensional Rossby solitary waves decrease exponentially, and the width of three-dimensional Rossby solitary waves increases exponentially with the increasing of μ and time T ; on the other hand, the whole chirp effect is related to the zonal flow, β effect, and wave speed. When (82) is satisfied, the three-dimensional Rossby solitary waves can be spread steadily. The smaller the wave speed, the stronger the dispersive effect; the greater the wave speed, the stronger the nonlinear effect as the zonal flow and β effect is determined.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

The authors declare that the study was realized in collaboration with the same responsibility. All authors read and approved the final manuscript.

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References

- Gu, CH: Theory and Application of Solitary Wave. Zhejiang Science and Technology Press, Hangzhou (1990)
- Hasegawa, A: Optical Solitons in Fibers. Springer, Berlin (2003)
- Lou, SY, Tang, XY: Nonlinear Mathematical and Physics Methods. Science Press, Beijing (2006)
- Infeld, E, Rowlands, G: Nonlinear Waves, Solitons and Chaos. Cambridge University Press, Cambridge (2000)
- Xu, ZH, Yin, BS, Hou, YJ, Xu, YS: Variability of internal tides and near-inertial waves on the continental slope of the northwestern South China Sea. *J. Geophys. Res., Oceans* **118**, 197 (2013)
- Le, KC, Nguyen, LTK: Amplitude modulation of water waves governed by Boussinesq's equation. *Nonlinear Dyn.* **81**, 659 (2015)
- Gao, XY: Comment on 'Solitons, Bäcklund transformation, and Lax pair for the (2 + 1)-dimensional Boiti-Leon-Pempinelli equation for the water waves'. *J. Math. Phys.* **51**, 093519 (2015)
- Gao, XY: Bäcklund transformation and shock-wave-type solutions for a generalized (3 + 1)-dimensional variable-coefficient B-type Kadomtsev-Petviashvili equation in fluid mechanics. *Ocean Eng.* **96**, 245 (2015)
- Yang, HL, Song, JB, Yang, LG, Liu, YJ: A kind of extended Korteweg-de Vries equation and solitary wave solutions for interfacial waves in a two-fluid system. *Chin. Phys.* **16**, 3589 (2007)
- Gao, XY: Variety of the cosmic plasmas: general variable-coefficient Korteweg-de Vries-Burgers equation with experimental/observational support. *Europhys. Lett.* **110**, 15002 (2015)
- Sun, WR, Tian, B, Liu, DY, Xie, XY: Nonautonomous matter-wave solitons in a Bose-Einstein condensate with an external potential. *J. Phys. Soc. Jpn.* **84**, 074003 (2015)
- Xie, XY, Tian, B, Sun, WR, Wang, M, Wang, YP: Solitary wave and multi-front wave collisions for the Bogoyavlenskii-Kadomtsev-Petviashvili equation in physics, biology and electrical networks. *Mod. Phys. Lett. B* **29**, 1550192 (2015)
- Long, RR: Solitary waves in the westerlies. *J. Atmos. Sci.* **21**, 197 (1964)
- Wadati, M: The modified Korteweg-de Vries equation. *J. Phys. Soc. Jpn.* **34**, 34 (1973)
- Yang, HW, Yin, BS, Dong, HH, Ma, ZD: Generation of solitary Rossby waves by unstable topography. *Commun. Theor. Phys.* **57**, 473 (2012)
- Meng, L, Lv, KL: Nonlinear long-wave disturbances excited by localized forcing. *Chin. J. Comput. Phys.* **17**, 259 (2000)
- Yang, HW, Yin, BS, Shi, YL: Forced dissipative Boussinesq equation for solitary waves excited by unstable topography. *Nonlinear Dyn.* **70**, 1389 (2012)
- Grimshaw, RHJ, Zhu, Y: Oblique interactions between internal solitary waves. *Stud. Appl. Math.* **92**, 249 (1994)
- Ma, WX: Combined Wronskian solutions to the 2D Toda molecule equation. *Phys. Lett. A* **375**, 3931 (2011)
- Qiao, ZJ, Li, ST: A new integrable hierarchy, parametric solutions and traveling wave solutions. *Math. Phys. Anal. Geom.* **7**, 289 (2004)
- Abdou, MA: New solitons and periodic wave solutions for nonlinear physical models. *Nonlinear Dyn.* **52**, 129 (2008)
- Ma, WX, Zhang, Y, Tang, YN, Tu, JY: Hirota bilinear equations with linear subspaces of solutions. *Appl. Math. Comput.* **218**, 7174 (2012)
- Zhao, Q, Liu, SK: Application of Jacobi elliptic functions in the atmospheric and oceanic dynamics: studies on two-dimensional nonlinear Rossby waves. *Chin. J. Geophys.* **49**, 965 (2006)
- Zedan, HA, Aladrous, E, Shapll, S: Exact solutions for a perturbed nonlinear Schrödinger equation by using Bäcklund transformations. *Nonlinear Dyn.* **74**, 1145 (2013)
- Ono, H: Algebraic Rossby wave soliton. *J. Phys. Soc. Jpn.* **50**, 2757 (1981)
- Yang, HW, Wang, XR, Yin, BS: A kind of new algebraic Rossby solitary waves generated by periodic external source. *Nonlinear Dyn.* **74**, 1725 (2014)
- Zhang, YF, Ma, WX: A study on rational solutions to a KP-like equation. *Z. Naturforsch. A* **70**(4), 263 (2015)
- Shi, CG, Zhao, BZ, Ma, WX: Exact rational solutions to a Boussinesq-like equation in (1 + 1)-dimensions. *Appl. Math. Lett.* **48**, 170 (2015)
- Koch, TL: Optical Fiber Telecommunications IIIA. Academic Press, Pittsburgh (1997)
- Khalique, M, Magalakwe, G: Combined sinh-cosh-Gordon equation: symmetry reductions, exact solutions and conservation laws. *Quaest. Math.* **37**, 199 (2014)
- Biswas, A: 1-Soliton solution of the generalized Zakharov-Kuznetsov modified equal width equation. *Appl. Math. Lett.* **22**, 1775 (2009)
- Biswas, A, Zerrad, E: 1-Soliton solution of the Zakharov-Kuznetsov equation with dual-power law nonlinearity. *Commun. Nonlinear Sci. Numer. Simul.* **14**, 3574 (2009)
- Biswas, A: 1-Soliton solution of the generalized Zakharov-Kuznetsov equation with nonlinear dispersion and time-dependent coefficients. *Phys. Lett. A* **373**, 2931 (2009)

34. Biswas, A, Zerrad, E: Solitary wave solution of the Zakharov-Kuznetsov equation in plasmas with power law nonlinearity. *Nonlinear Anal., Real World Appl.* **11**, 3272 (2010)
35. Krishnan, EV, Biswas, A: Solutions to the Zakharov-Kuznetsov equation with higher order nonlinearity by mapping and ansatz methods. *Phys. Wave Phenom.* **18**, 256 (2010)
36. Suarez, P, Biswas, A: Exact 1-soliton solution of the Zakharov equation in plasmas with power law nonlinearity. *Appl. Math. Comput.* **217**, 7372 (2011)
37. Johnpillai, AG, Kara, AH, Biswas, A: Symmetry solutions and reductions of a class of generalized (2 + 1)-dimensional Zakharov-Kuznetsov equation. *Int. J. Nonlinear Sci. Numer. Simul.* **12**, 45 (2011)
38. Morris, R, Kara, AH, Biswas, A: Soliton solution and conservation laws of the Zakharov equation in plasmas with power law nonlinearity. *Nonlinear Anal., Model. Control* **18**, 153 (2013)
39. Wang, GW, Xu, TZ, Johnson, S, Biswas, A: Solitons and Lie group analysis to an extended quantum Zakharov-Kuznetsov equation. *Astrophys. Space Sci.* **349**, 317 (2014)
40. Güner, Ö, Bekir, A, Moraru, L, Biswas, A: Bright and dark soliton solutions of the generalized Zakharov-Kuznetsov-Benjamin-Bona-Mahoney nonlinear evolution equation. *Proc. Rom. Acad., Ser. A* **16**, 422 (2015)

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